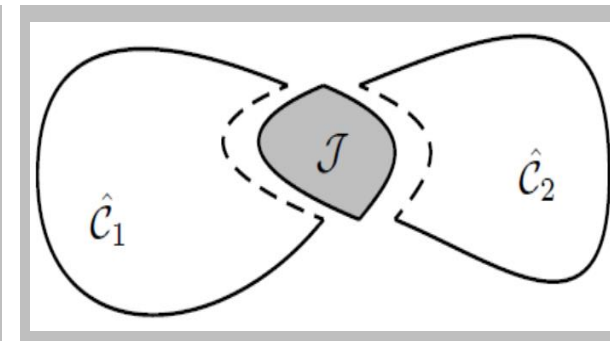
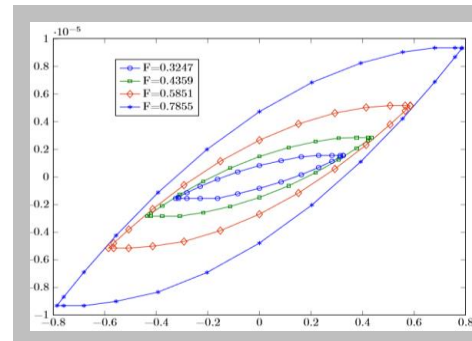
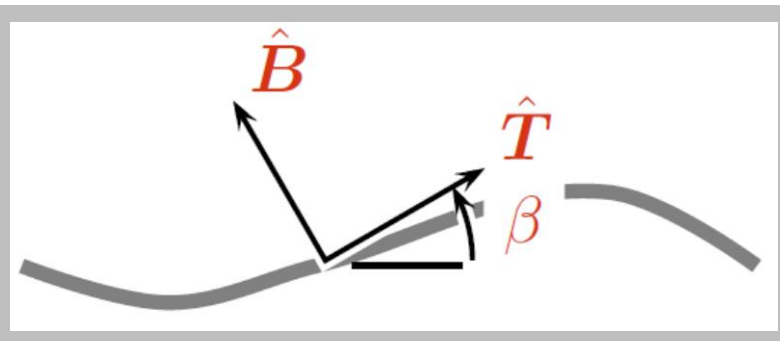


*Exceptional service in the national interest*



# Continuum Shell Models for Structural Damping

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Dan Segalman – Michigan State University

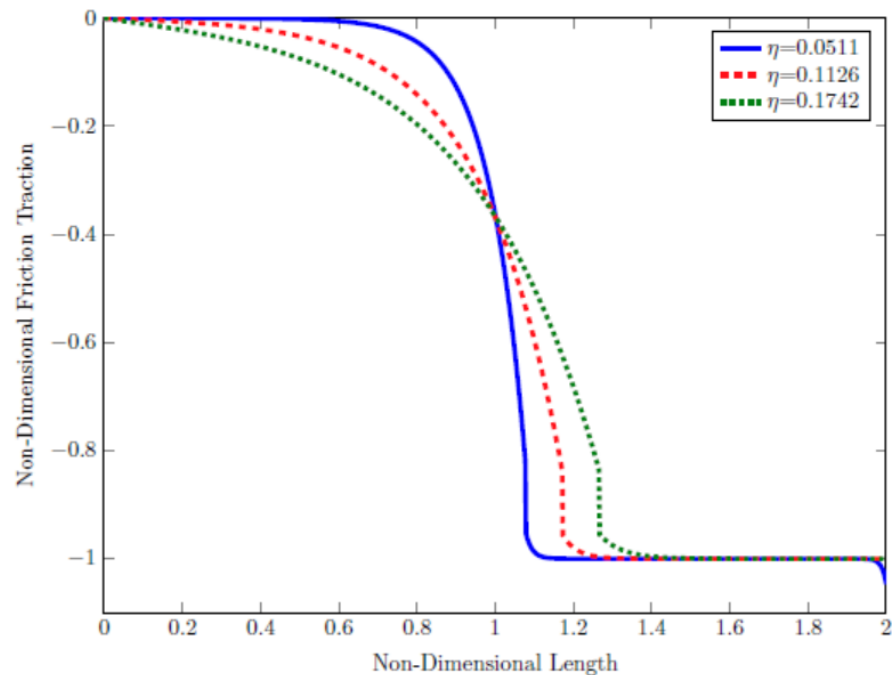
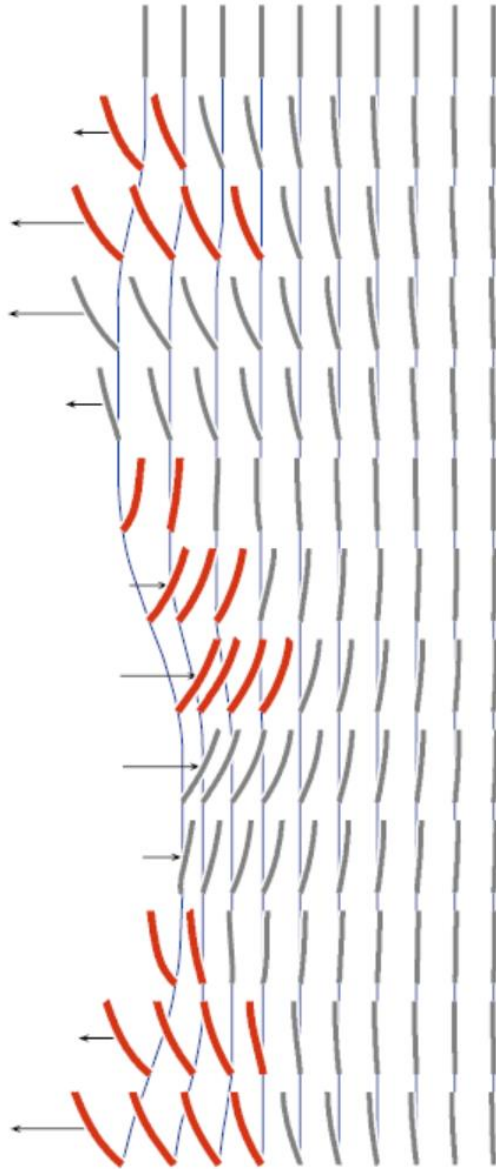


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# Research Motivation

- Develop shell models which capture structural damping and joint stiffness in a reduced manner.
  - Capture Micro- and Macro-Slip with physics incorporated directly into the shell.
  - Bar-like model developed by Quinn and Segalman.
  - Shearable beam-like model developed by Brink and Quinn.
- Develop a convenient framework to include the nonlinear joint into modal dynamics.
  - Quinn developed a method which incorporates joint forces into the modal equations of motion.

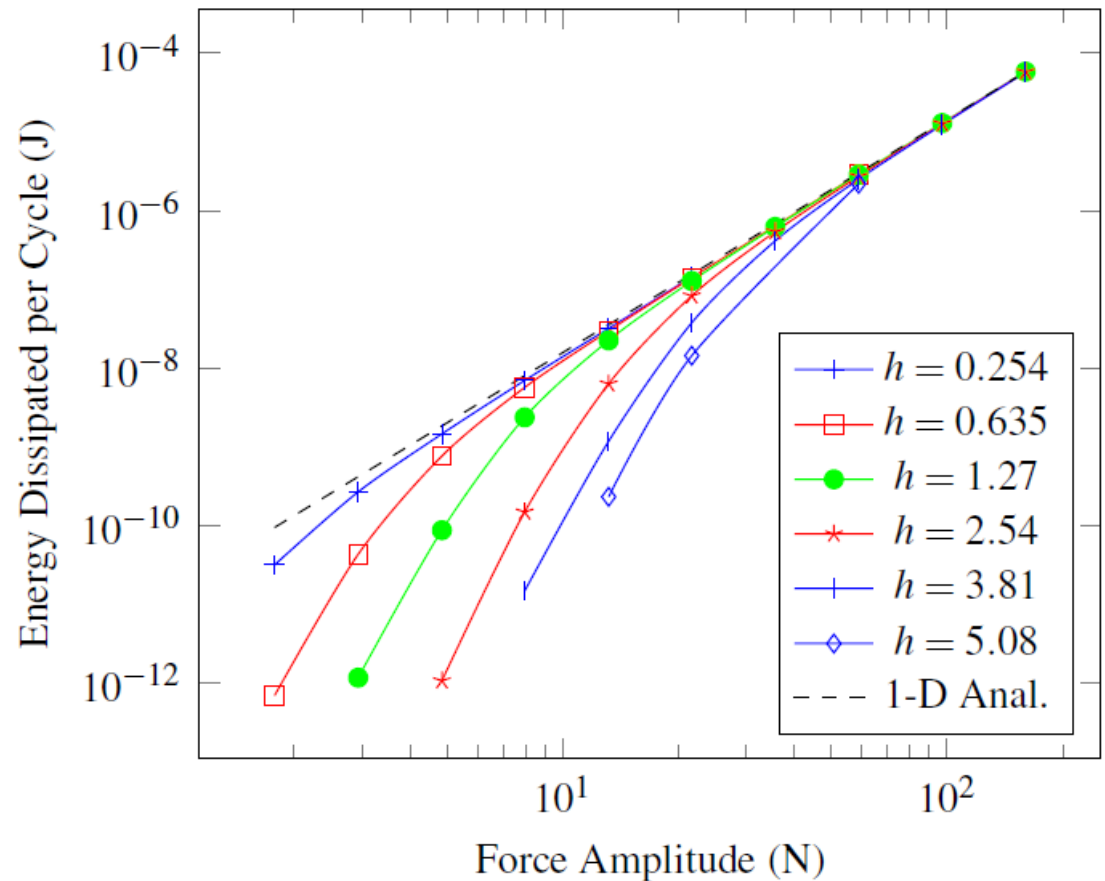
# Importance of Including Shear



- Shearing allows forces to transmit passed the slip initiation length.
- The shorter the cross-section, the more bar-like the behavior becomes

# Importance of Including Shear

- Shearing allows for a precipitous drop off from the cubic power law bar solution.



# Shearable Shell Derivation

We use a geometrically exact formulation for shell theory presented by Libai and Simmonds (restricted here to quasi-static behavior)

Strains and Internal Forces

$$\mathbf{y}' = (1 + e) \hat{\mathbf{T}} + g \hat{\mathbf{B}},$$

$$\beta' = k$$

$$\mathbf{m} \equiv \hat{\mathbf{k}} \times \mathbf{y}',$$

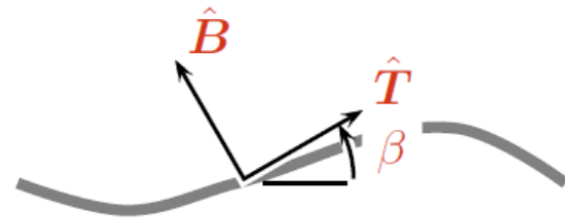
$$\mathbf{F} = N \hat{\mathbf{T}} + Q \hat{\mathbf{B}}$$

Linear and Angular Momentum Balance

$$\left( N \hat{\mathbf{T}} + Q \hat{\mathbf{B}} \right)' + \mathbf{p} = \mathbf{0}, \quad M' + \mathbf{m} \cdot \left( N \hat{\mathbf{T}} + Q \hat{\mathbf{B}} \right) + \ell = 0.$$

Constitutive Laws

$$N = EHc_1 e, \quad Q = EHc_2 g, \quad M = EH^3c_3 k$$



# Shearable Shell Derivation

External forces and moments

$$\mathbf{p} = F_t \hat{\mathbf{e}}_x + (R_n - F_n) \hat{\mathbf{e}}_y,$$

$$\ell = h (F_t \cos \beta + (R_n + F_n) \sin \beta).$$

The contact force is assumed to follow Coulomb's law:

Stick:  $u = v - h \sin \beta$

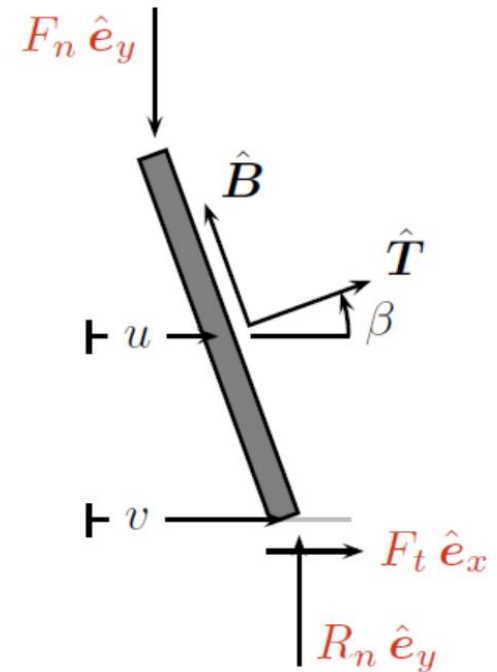
Slip:  $F_t = \mu R_n \operatorname{sgn}(\dot{v})$

The equations of motion reduce to

$$(N' - \beta' Q) + F_t \cos \beta - (R_n - F_n) \sin \beta = 0,$$

$$(Q' + \beta' N) + F_t \sin \beta + (R_n - F_n) \cos \beta = 0,$$

$$M' + (1 + e) N - g Q + h (F_t \cos \beta + (R_n + F_n) \sin \beta) = 0.$$



# Shearable Shell Derivation

The equations of motion can be nondimensionalized and linearized in the deformation

$$\begin{array}{ll} \text{Linear} & \kappa u'' + F_t = 0, \quad -\kappa c_1 \beta' + (R_n - F_n) = 0, \\ \text{Angular} & \alpha \eta^2 \beta'' + \kappa c_2 \beta + \eta (F_t + (R_n + F_n) \beta) = 0. \end{array}$$

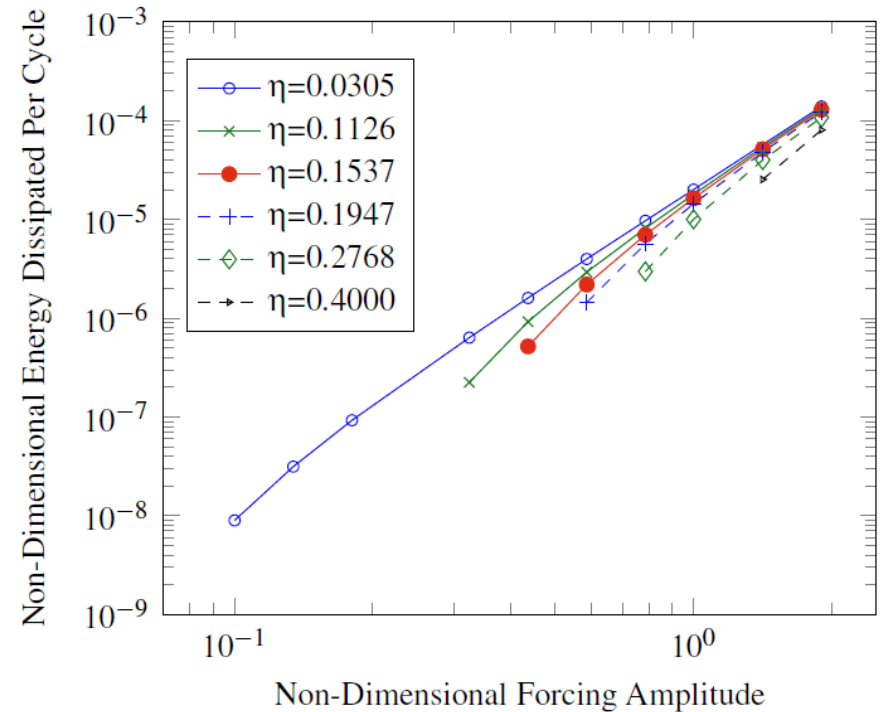
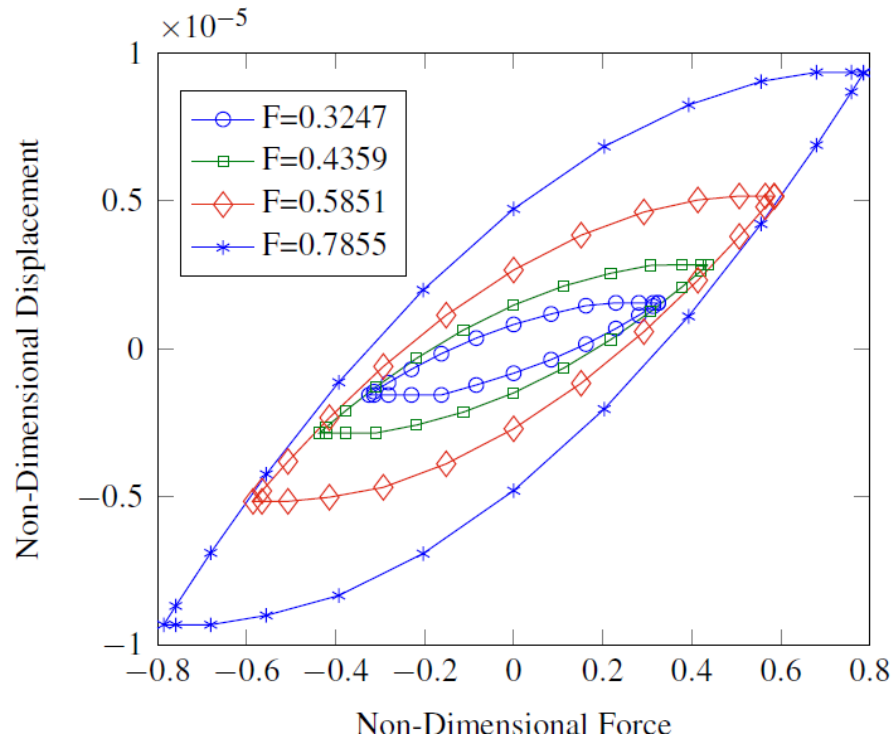
$\alpha$ : nondimensional bending stiffness

$\kappa$ : nondimensional axial stiffness

$\eta$ : nondimensional height

The external loading is prescribed at the boundaries

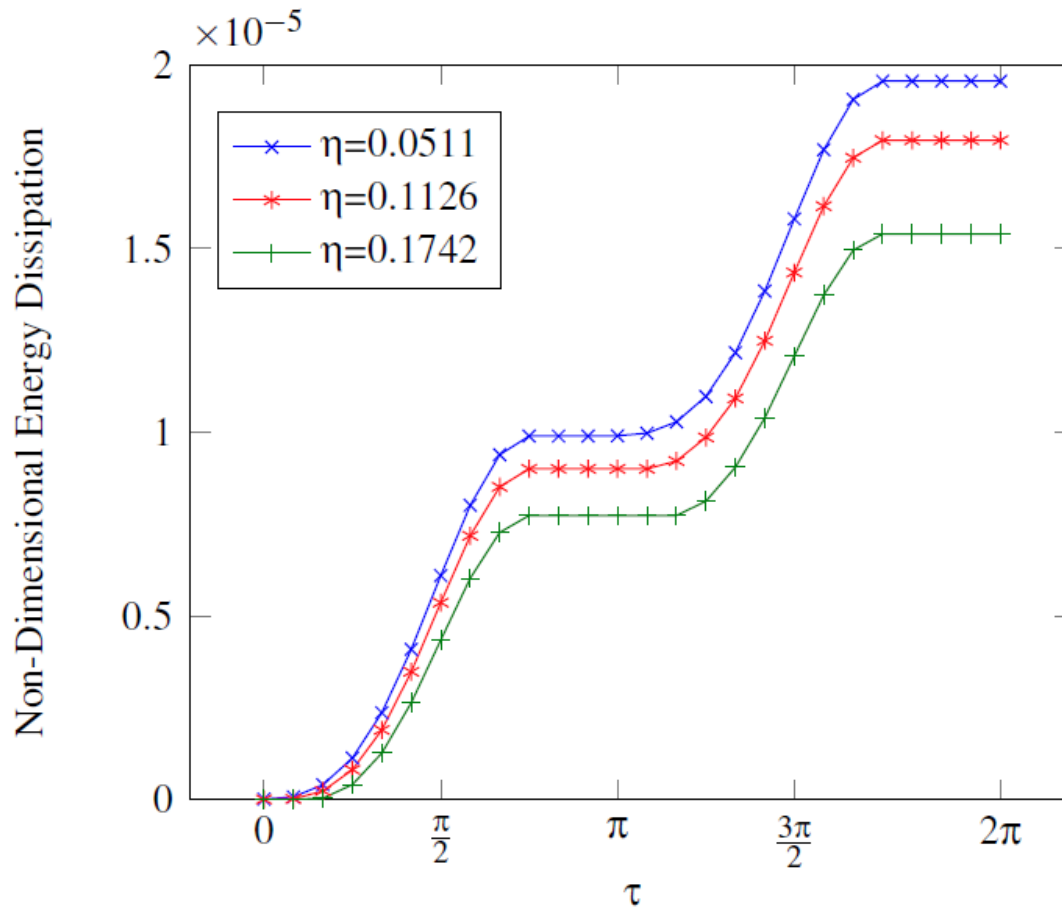
# Some Shell Results



Hysteresis curves and energy dissipated per cycle versus forcing amplitude are as expected.



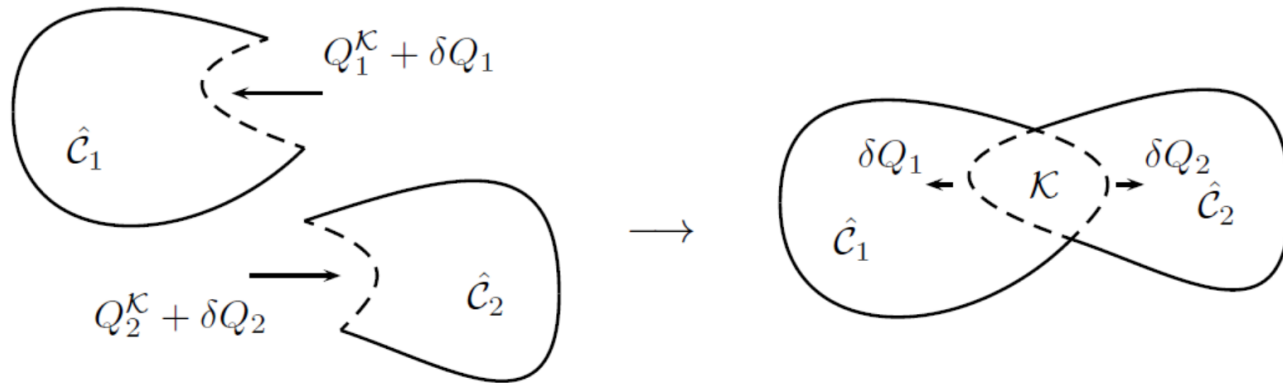
# More Shell Results



Energy dissipated over a cycle shows expected flat spotting.

# Modal Analysis

- Quinn solves modal equations of motion for a monolithic structure, then adds the effect of the joint back in.



$$\left[ \int_{\mathcal{M}} \phi_i(x) \rho_{\mathcal{M}}(x) \phi_i(x) dx \right] \ddot{A}_i(t) + \left[ \int_{\mathcal{M}} \frac{\partial \phi_i(x)}{\partial x} EA_{\mathcal{M}}(x) \frac{\partial \phi_i(x)}{\partial x} dx \right] A_i(t)$$

Monolithic Response

$$= -(\phi_i(s_1) \delta Q_1(t) + \phi_i(s_2) \delta Q_2(t)).$$

Forces Arising From Joint at Interfaces

# Modal Analysis – Elastic Bar

Consider the response of an elastic rod with simply supported boundaries and a distributed interface.

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} = 0, \quad x \in (0, 1),$$

The response of the monolithic structure can be expressed as

$$u(x, t) = \sum_{k=1}^{\infty} A_k(t) \phi_k(x), \quad \phi_k(x) = \sqrt{2} \sin(k \pi x).$$

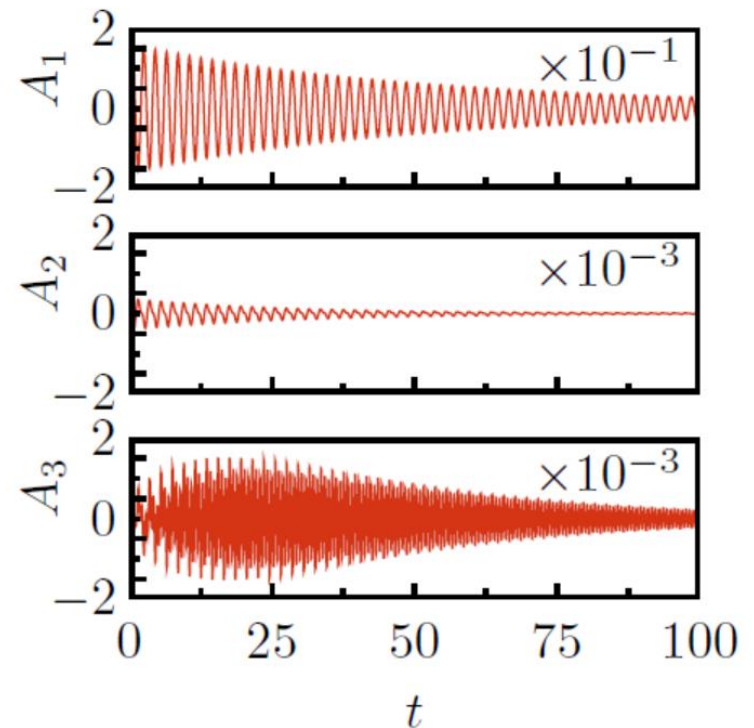
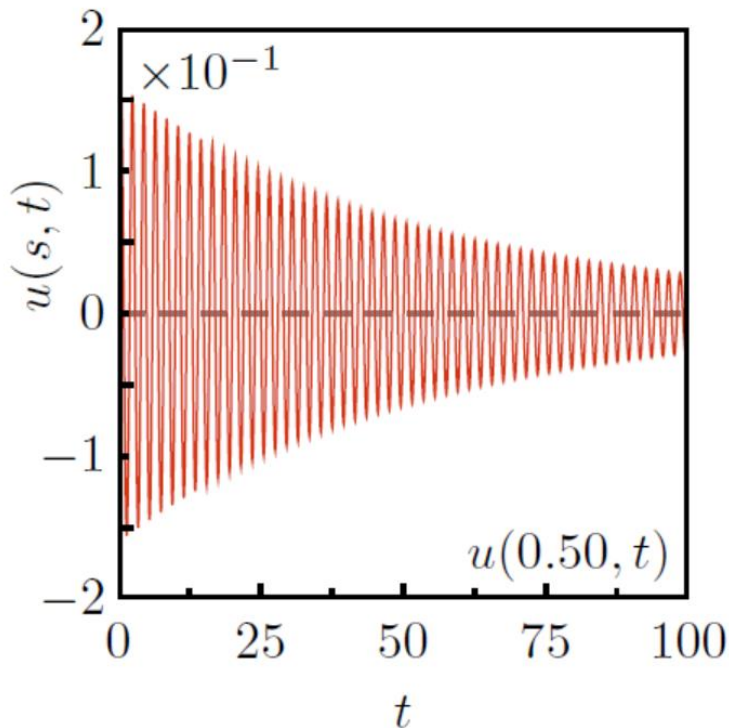
An  $N$  mode truncation

$$\ddot{A}_j + 2 \zeta(j \pi) \dot{A}_j + (j \pi)^2 A_j + (\phi_j(s_2) - \phi_j(s_1)) \delta Q \left( \sum_{k=1}^N A_k(t) (\phi_k(s_2) - \phi_k(s_1)) \right) = 0,$$

# Modal Analysis – Elastic Bar

The interface is located between  $s_1 = 0.20$  and  $s_2 = 0.30$  and the initial conditions only excite the fundamental mode of the rod.

$$N = 5, \quad \zeta = 0.005, \quad \beta = 1.00, \quad \ell_0 = 1.00, \quad F_0 = 0.10.$$



# Conclusions

- A nonlinear shearable elastic shell is developed which directly incorporates friction into its formulation.
- The shell is capable of modeling micro- and macro-slip phenomena.
- Quinn's modal analysis method is a concise way to introduce joint nonlinearities into modal framework.

# Future Work

- Compare the shearable elastic shell to the elastic bar using the modal analysis techniques.
- Introduce the reduced order shell models and Quinn's modal analysis techniques into a finite element framework.