

# Predictivity in fracture modeling

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Science, and Engineering  
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*Exceptional  
service  
in the  
national  
interest*



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# Outline

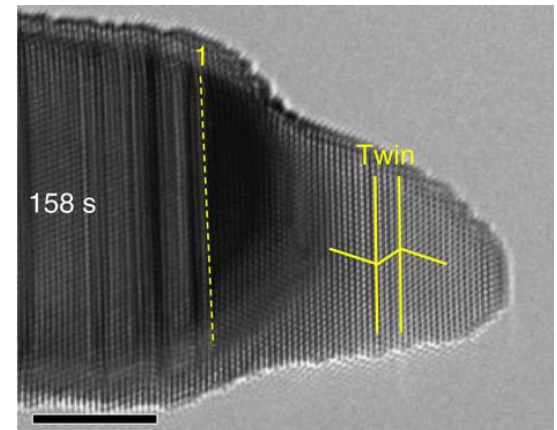
- Why is fracture a challenging problem?
- Treatment of fracture in the local theory and in peridynamics
- Reality and usefulness of nonlocality
- Material variability and the probabilistic nature of fracture

# Physical aspects

- It involves many length and time scales.
- It can involve multiple mechanisms that lead to fracture only in combination with each other.
- It can involve many interacting defects.



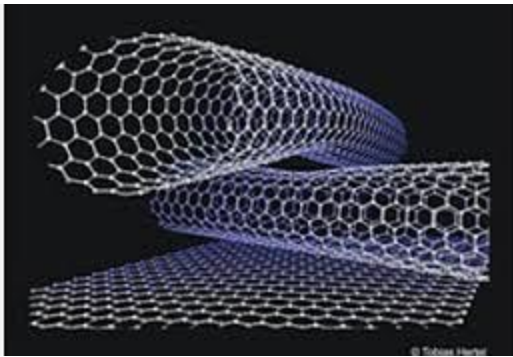
Fragmented glass  
(image: Washington Glass School)



Nanoscale twins in fractured gold crystal  
J. Wang et al, Nature Communications (2013)

# Mathematical aspects

- The nature of fracture as a discontinuity is incompatible with the local theory.
- The Cauchy (local) theory assumes a continuous deformation.
- This leads to need to hack in extraneous mathematical relations (LEFM).
- Fracture processes are often physically nonlocal at some level.



Carbon nanotubes (image: nsf.gov)

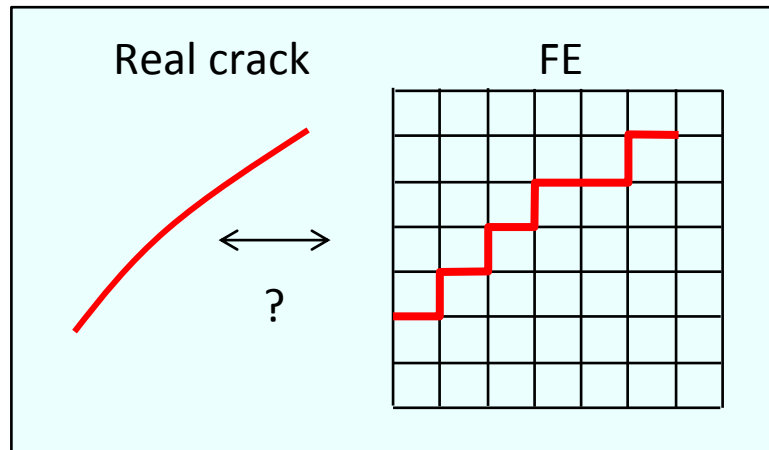
$$\nabla \cdot \sigma + b = 0$$



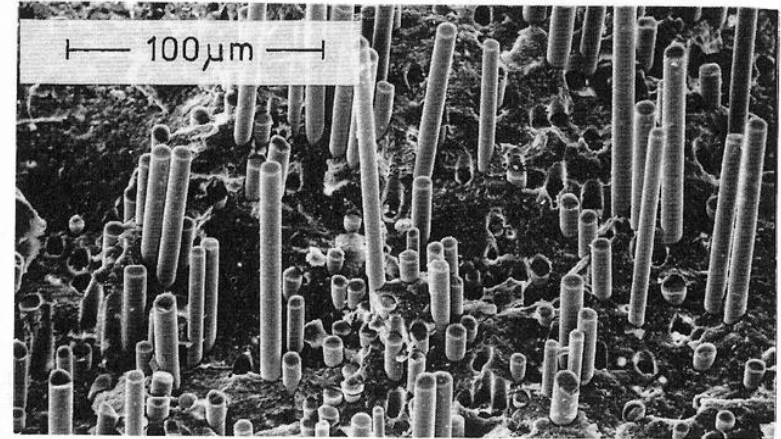
Augustin-Louis Cauchy, 1840  
(image: Library of Congress)

# Computational aspects

- Traditional FEM (e.g. displacement Galerkin) inherits the incompatibility of the local theory with fracture.
- This leads to the need to hack in a fix after discretization.
- MD is better but too slow.



(b) Complex crack path in a composite



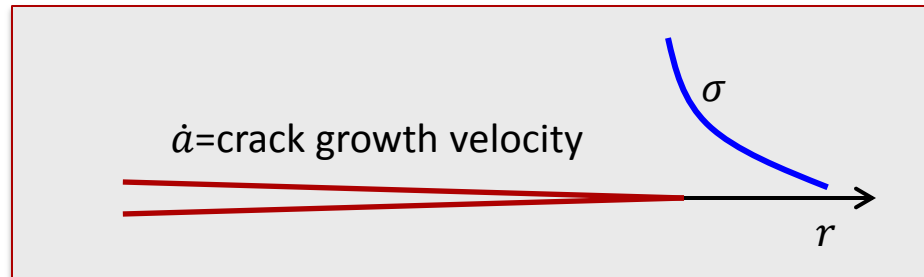
**Figure 11.20** Pull-out: (a) schematic diagram; (b) fracture surface of 'Silceram' glass-ceramic reinforced with SiC fibres. (Courtesy H. S. Kim, P. S. Rogers and R. D. Rawlings.)

# Fracture modeling in the local theory

- Smeared crack models, continuum damage mechanics
  - Attempt to adjust bulk material properties without addressing individual cracks.
  - Damage  $D$  at  $x$  evolves according to local conditions.
$$\dot{D} = \phi(F, \dot{F}, \theta, \dots)$$
where  $F$  = deformation gradient tensor,  $\theta$  = temperature.
  - Damage degrades the material properties. Examples:
    - $Y = (1 - D)Y_0$  ... flow stress (smeared crack model).
    - $E = (1 - D)E_0$  ... elastic modulus (CDM).
- Advantage: Easy to include in FEM.
- Disadvantage: Failure occurs in many elements rather than on 2D surfaces (discrete cracks).

# Linear elastic fracture mechanics (LEFM)

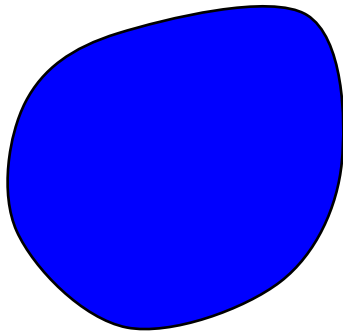
- Assumes that the classical asymptotic stress field ( $\sigma = K/\sqrt{r}$ ) holds within a small region near the crack tip.
- Crack growth occurs according to a given function  $\dot{a}(K)$  that is obtained empirically.



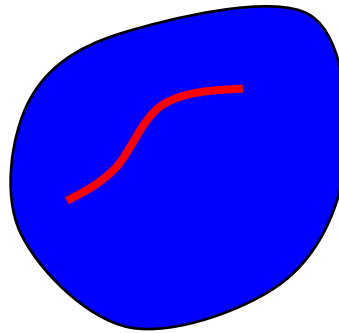
- Advantages:
  - Doesn't model the details of the process zone.
  - Compatible with energy balance concept of brittle fracture (Griffith).
- Disadvantages:
  - Doesn't model the details of the process zone.
  - The function  $\dot{a}(K)$  can be measured only in ideal cases.
  - Assumes a pre-existing crack.

# Purpose of peridynamics\*

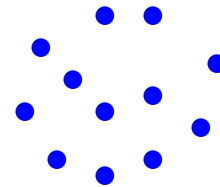
- To unify the mechanics of continuous and discontinuous media within a single, consistent set of equations.



Continuous body



Continuous body  
with a defect



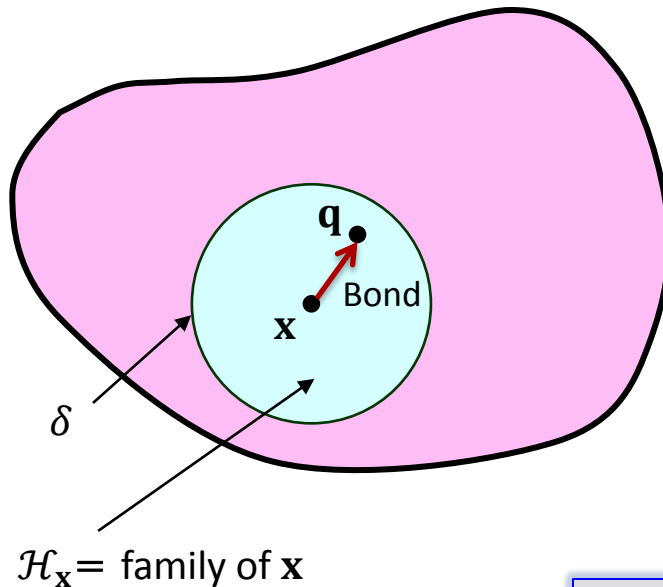
Discrete particles

- Why do this?
  - Avoid coupling dissimilar mathematical systems (A to C).
  - Model complex fracture patterns.
  - Communicate across length scales.

\* Peri (near) + dyn (force)

# Peridynamics: Horizon and family

- Any point  $\mathbf{x}$  interacts directly with other points within a distance  $\delta$  called the “horizon.”
- The material within a distance  $\delta$  of  $\mathbf{x}$  is called the “family” of  $\mathbf{x}$ ,  $\mathcal{H}_{\mathbf{x}}$ .



Equilibrium equation

$$\int_{\mathcal{H}_{\mathbf{x}}} \mathbf{f}(\mathbf{q}, \mathbf{x}) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x}) = 0$$

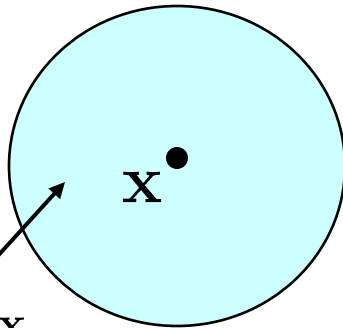
$\mathbf{f}$  = bond force density

General references

- SS, Journal of the Mechanics and Physics of Solids (2000)
- SS and R. Lehoucq, Advances in Applied Mechanics (2010)
- Madenci & Oterkus , *Peridynamic Theory & Its Applications* (2014)

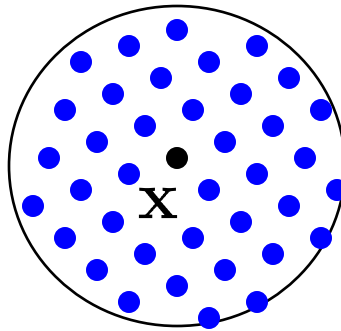
# Peridynamics: Strain energy at a point

Continuum

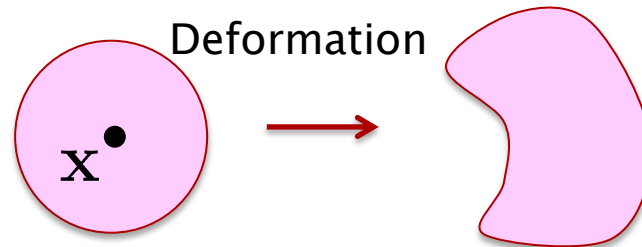
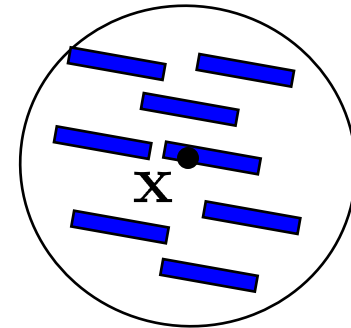


Family of  $\mathbf{x}$

Discrete particles



Discrete structures



- Key assumption: the strain energy density at  $\mathbf{x}$  is determined by the deformation of its family.

# Peridynamic vs. local equations

- State notation: State<bond>=vector

Relation	Peridynamic theory	Standard theory
Kinematics	$\underline{\mathbf{Y}}\langle \mathbf{q} - \mathbf{x} \rangle = \mathbf{y}(\mathbf{q}) - \mathbf{y}(\mathbf{x})$	$\mathbf{F}(\mathbf{x}) = \frac{\partial \mathbf{y}}{\partial \mathbf{x}}(\mathbf{x})$
Linear momentum balance	$\rho \ddot{\mathbf{y}}(\mathbf{x}) = \int_{\mathcal{H}} \left( \mathbf{t}(\mathbf{q}, \mathbf{x}) - \mathbf{t}(\mathbf{x}, \mathbf{q}) \right) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x})$	$\rho \ddot{\mathbf{y}}(\mathbf{x}) = \nabla \cdot \boldsymbol{\sigma}(\mathbf{x}) + \mathbf{b}(\mathbf{x})$
Constitutive model	$\mathbf{t}(\mathbf{q}, \mathbf{x}) = \underline{\mathbf{T}}\langle \mathbf{q} - \mathbf{x} \rangle, \quad \underline{\mathbf{T}} = \hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}})$	$\boldsymbol{\sigma} = \hat{\boldsymbol{\sigma}}(\mathbf{F})$
Angular momentum balance	$\int_{\mathcal{H}} \underline{\mathbf{Y}}\langle \mathbf{q} - \mathbf{x} \rangle \times \underline{\mathbf{T}}\langle \mathbf{q} - \mathbf{x} \rangle dV_{\mathbf{q}} = \mathbf{0}$	$\boldsymbol{\sigma} = \boldsymbol{\sigma}^T$
Elasticity	$\underline{\mathbf{T}} = W_{\underline{\mathbf{Y}}} \text{ (Fréchet derivative)}$	$\boldsymbol{\sigma} = W_{\mathbf{F}} \text{ (tensor gradient)}$
First law	$\dot{\varepsilon} = \underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} + q + r$	$\dot{\varepsilon} = \boldsymbol{\sigma} \cdot \dot{\mathbf{F}} + q + r$

$$\underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} := \int_{\mathcal{H}} \underline{\mathbf{T}}\langle \boldsymbol{\xi} \rangle \cdot \dot{\underline{\mathbf{Y}}}\langle \boldsymbol{\xi} \rangle dV_{\boldsymbol{\xi}}$$

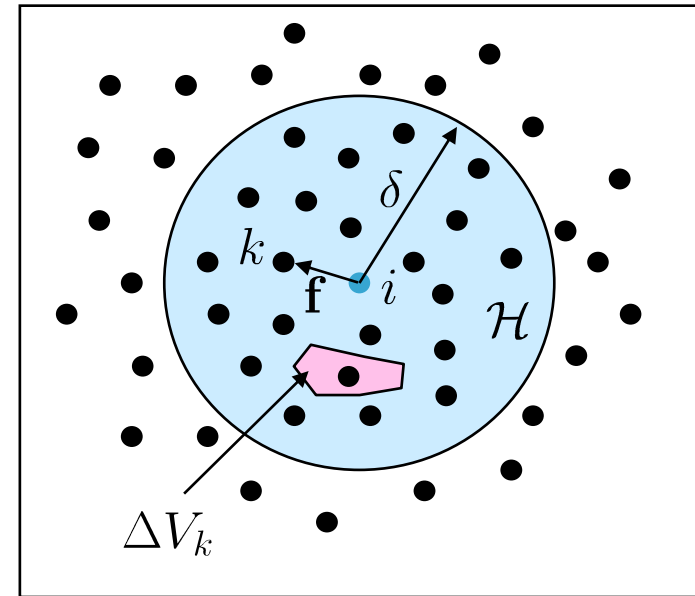
# Peridynamics: EMU numerical method

- Integral is replaced by a finite sum: resulting method is [meshless](#) and [Lagrangian](#).

$$\rho \ddot{\mathbf{y}}(\mathbf{x}, t) = \int_{\mathcal{H}} \mathbf{f}(\mathbf{x}', \mathbf{x}, t) dV_{\mathbf{x}'} + \mathbf{b}(\mathbf{x}, t) \quad \longrightarrow \quad \rho \ddot{\mathbf{y}}_i^n = \sum_{k \in \mathcal{H}} \mathbf{f}(\mathbf{x}_k, \mathbf{x}_i, t) \Delta V_k + \mathbf{b}_i^n$$

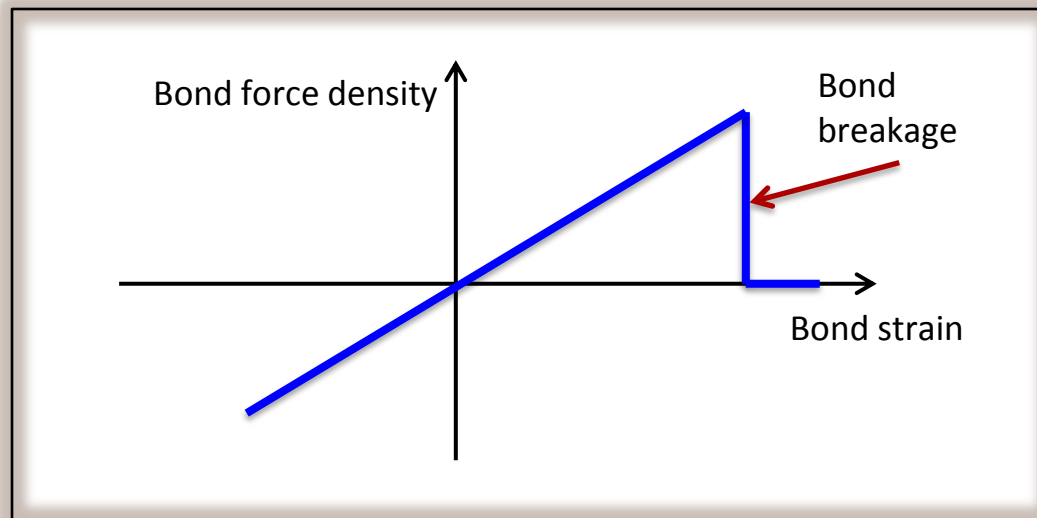
- Linearized model:

$$\rho \ddot{\mathbf{u}}_i = \sum_{k \in \mathcal{H}_i} \mathbf{C}_{ik} (\mathbf{u}_k - \mathbf{u}_i) \Delta V_k + \mathbf{b}_i$$

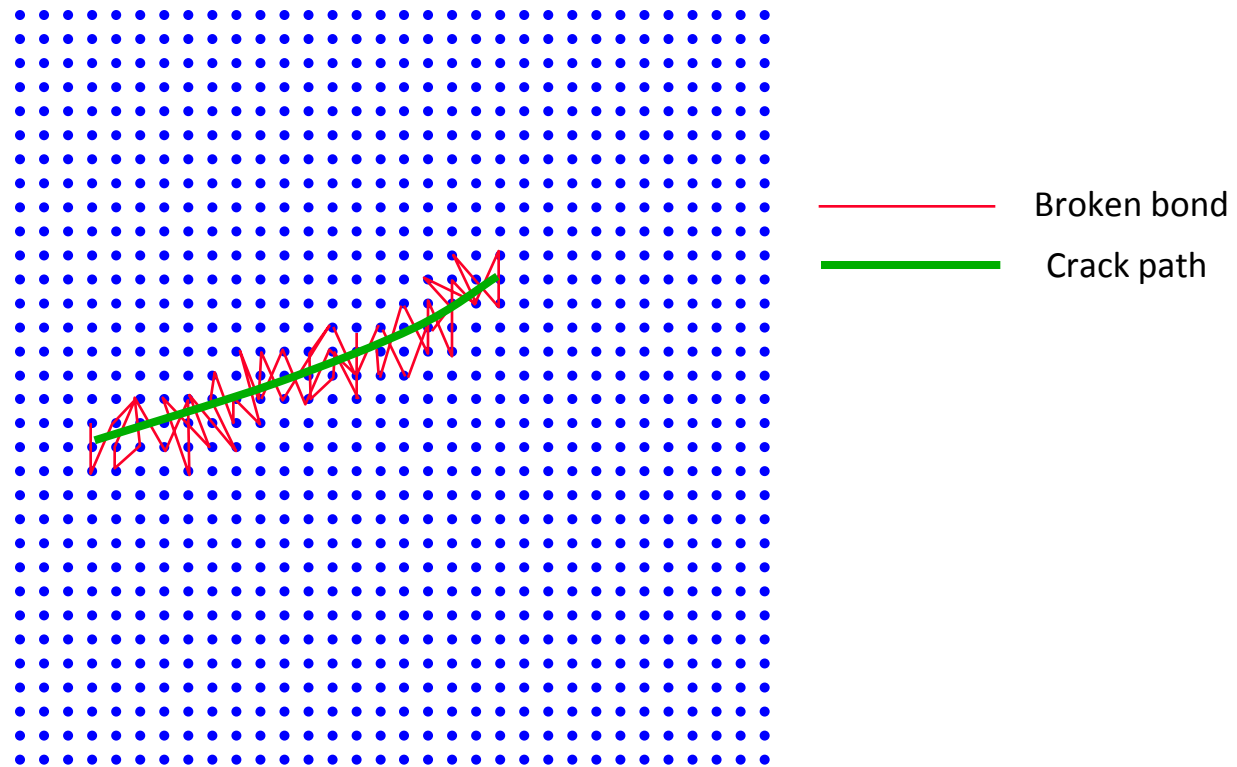


# Bond based material models

- If each bond response is independent of the others, the resulting material model is called bond-based.
- The material model is then simply a graph of bond force density vs. bond strain.
- Damage can be modeled through bond breakage.
- Bond response is calibrated to:
  - Bulk elastic properties.
  - Critical energy release rate.

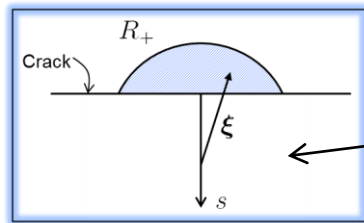
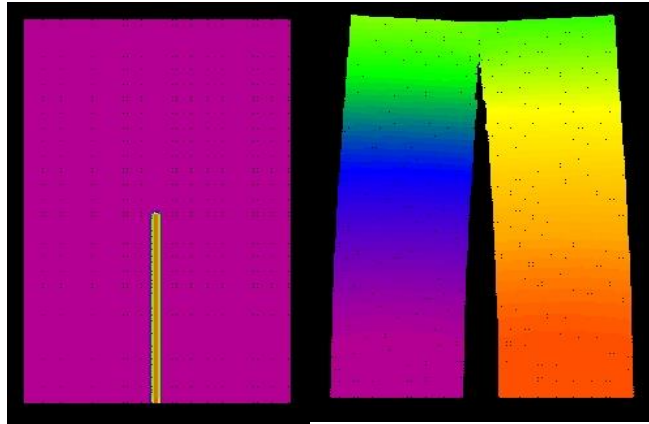


# Autonomous crack growth



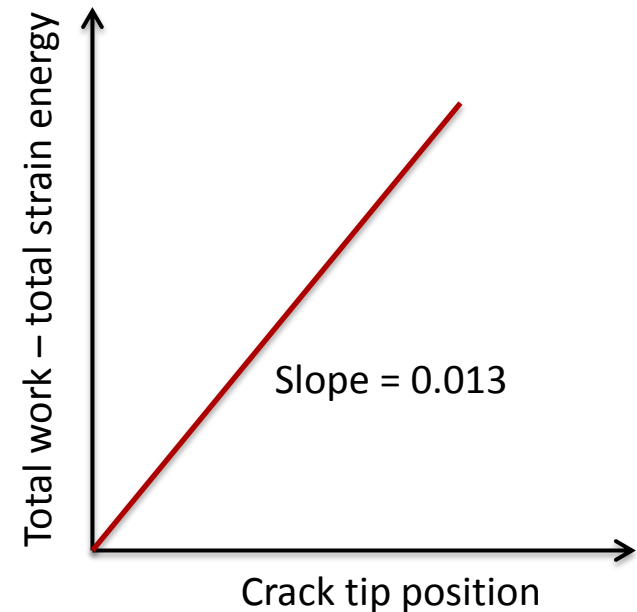
- When a bond breaks, its load is shifted to its neighbors, leading to progressive failure.

# PD reproduces the Griffith crack energy



From bond properties, energy release rate should be

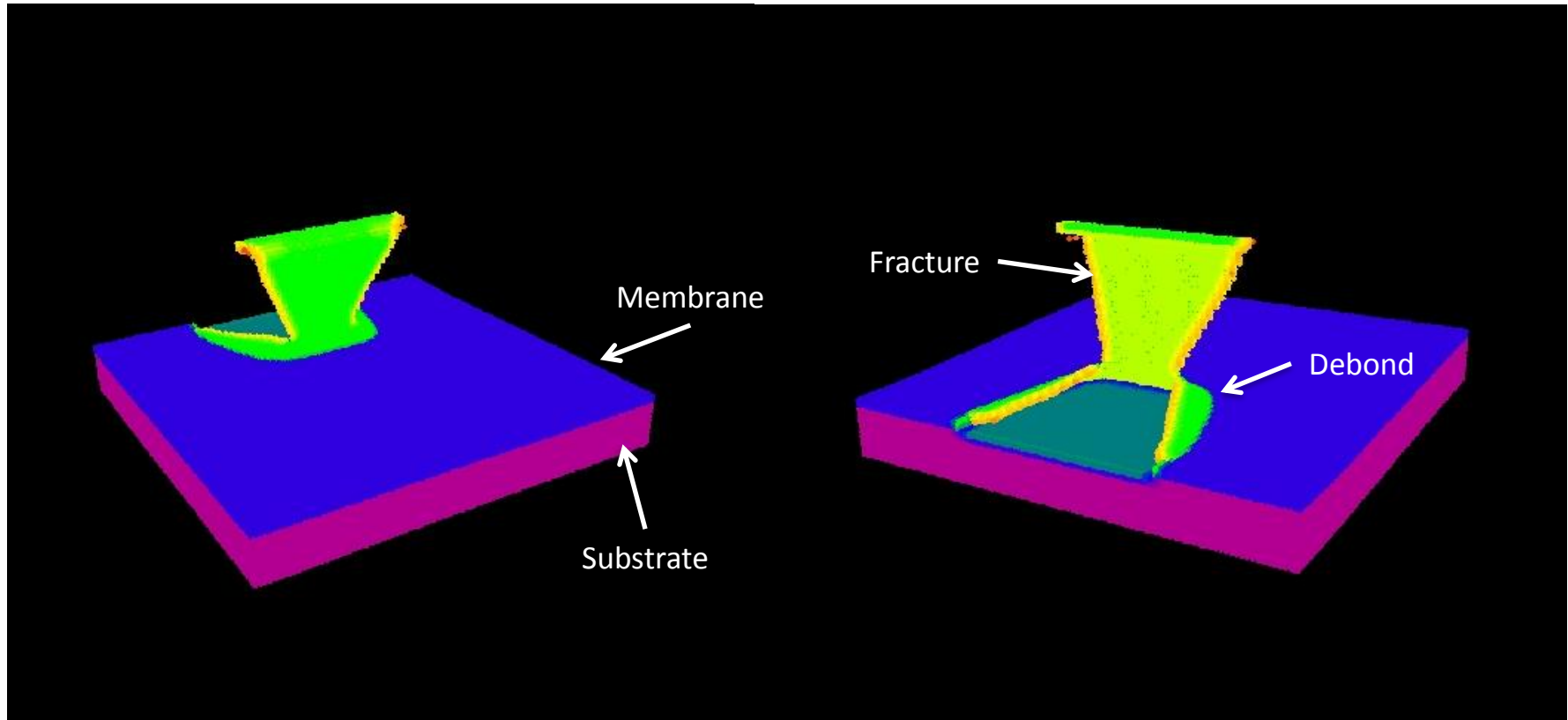
$$G = 0.013$$



- This confirms that the energy consumed per unit crack growth area equals the expected value from bond breakage properties.

# Fracture and debonding of membranes

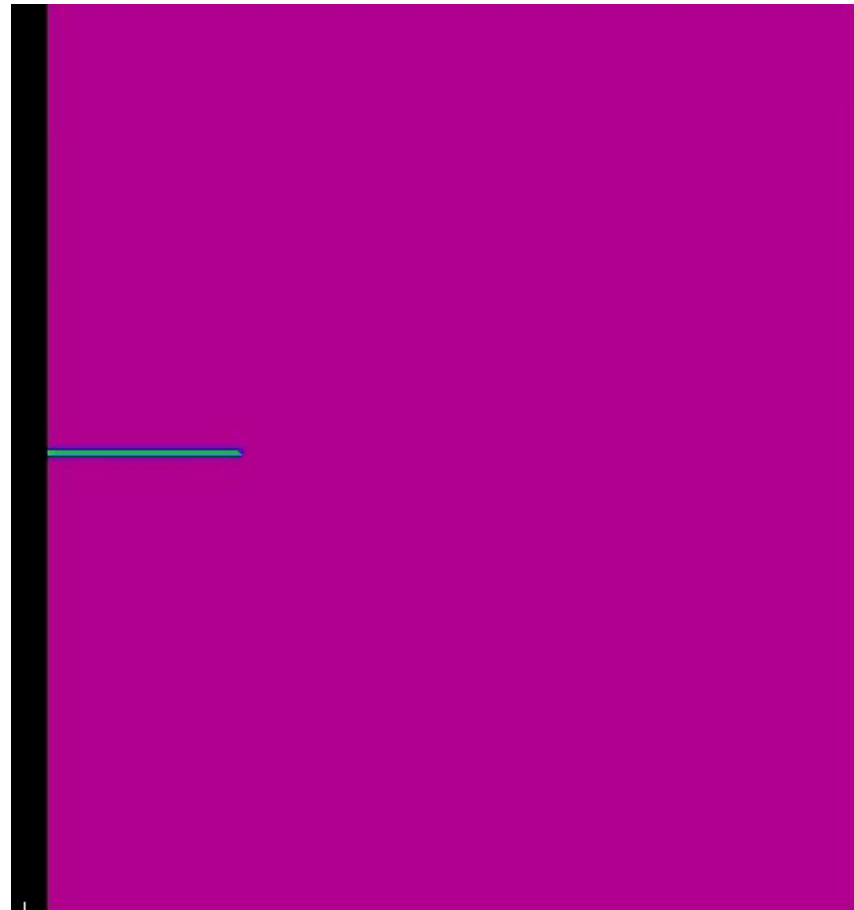
- Simulation of peeling illustrates interplay between fracture (tearing) and debonding (peeling).



# Dynamic crack branching

Video

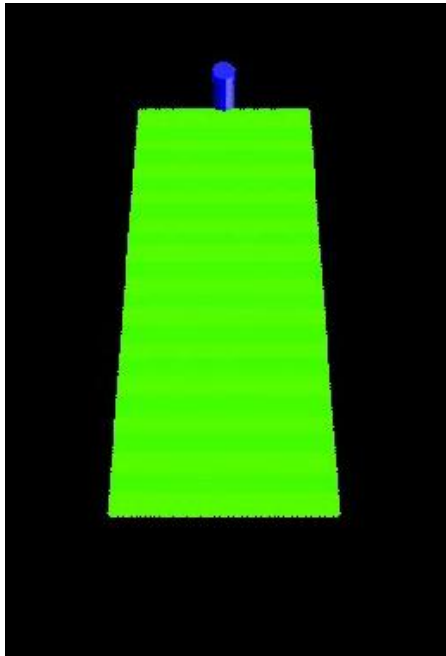
- Red indicates bonds currently undergoing damage.
- These appear ahead of the visible discontinuities.
- Blue/green indicate damage (broken bonds).
- More and more energy is being built up ahead of the crack – it can't keep up.
  - Leads to fragmentation.



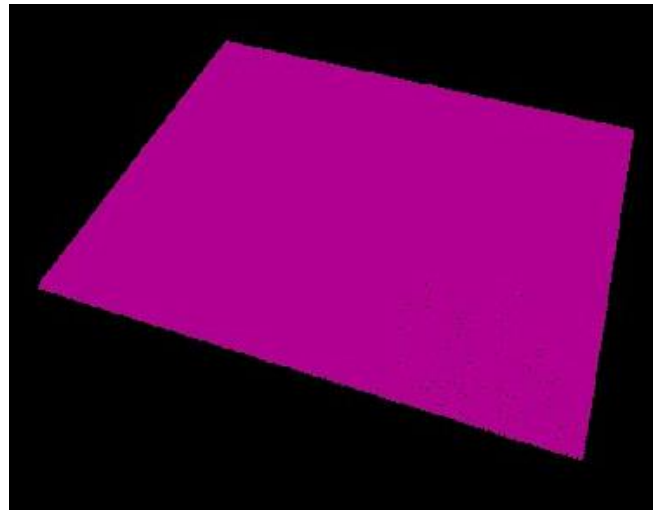
More on dynamic fracture: see Ha & Bobaru (2010, 2011)

# Membranes and thin films

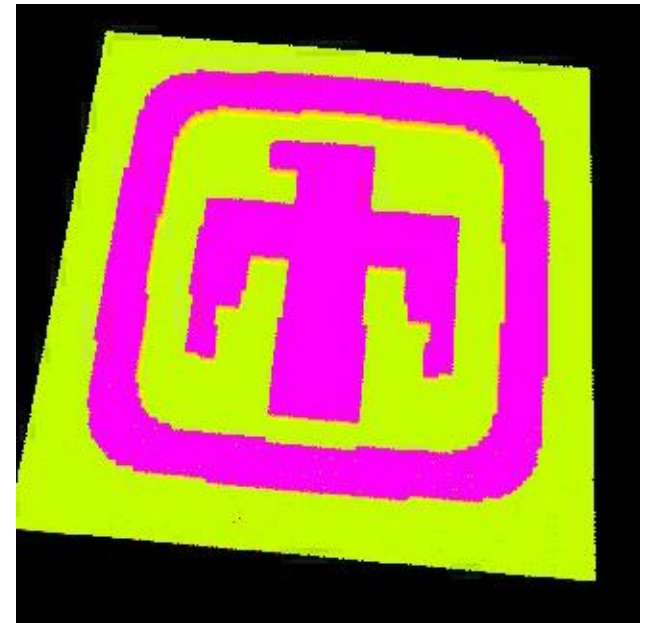
Videos



Oscillatory crack path

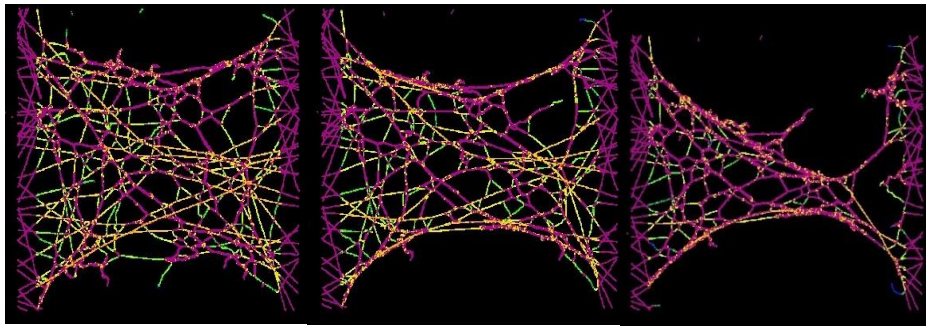


Crack interaction in a sheet

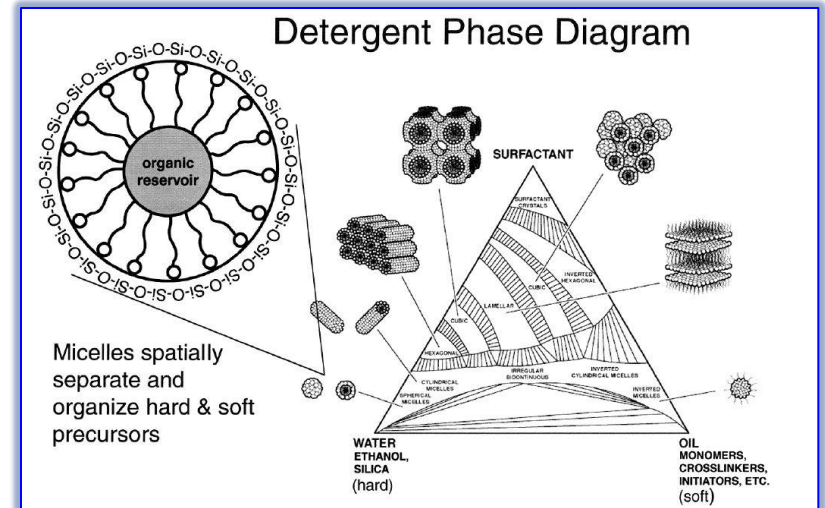


Aging of a film

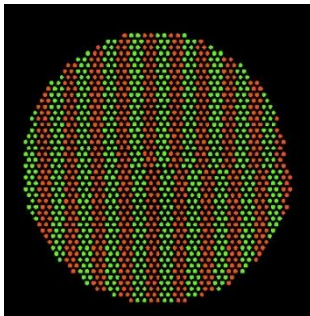
- Potential importance for self-assembled nanostructures.
- All forces are treated as long-range.



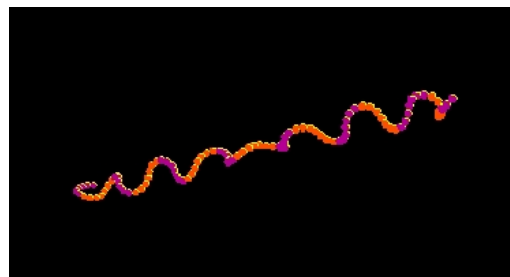
## Failure in a nanofiber membrane (F. Bobaru, Univ. of Nebraska)



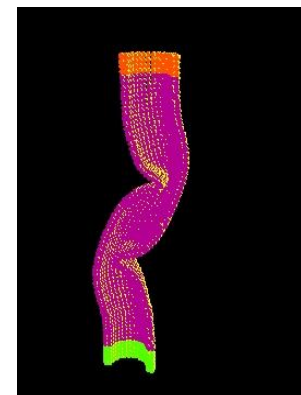
Self-assembly is driven by long-range forces  
Image: Brinker, Lu, & Sellinger, Advanced Materials (1999)



## Dislocation



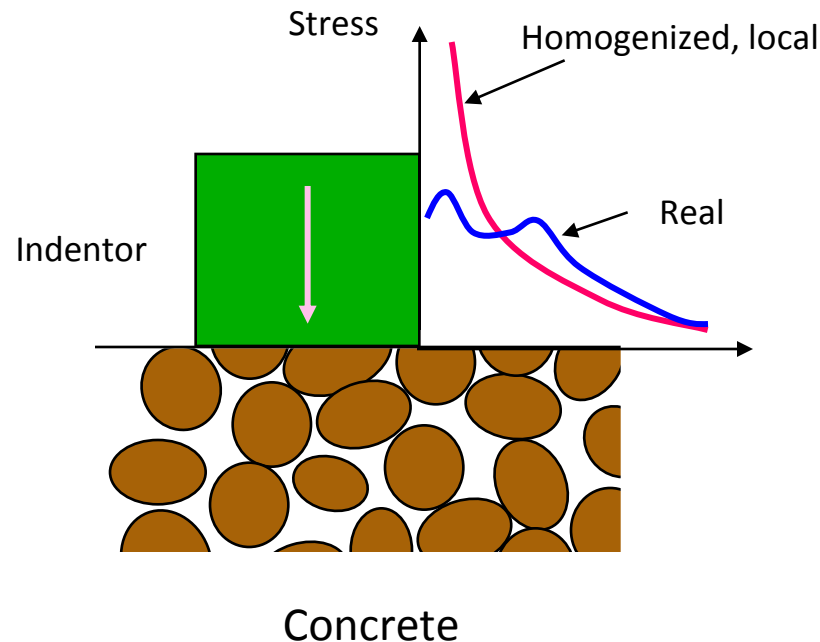
## Nanofiber self-shaping



## Carbon nanotube

# Nonlocality can arise from the way we choose to model things

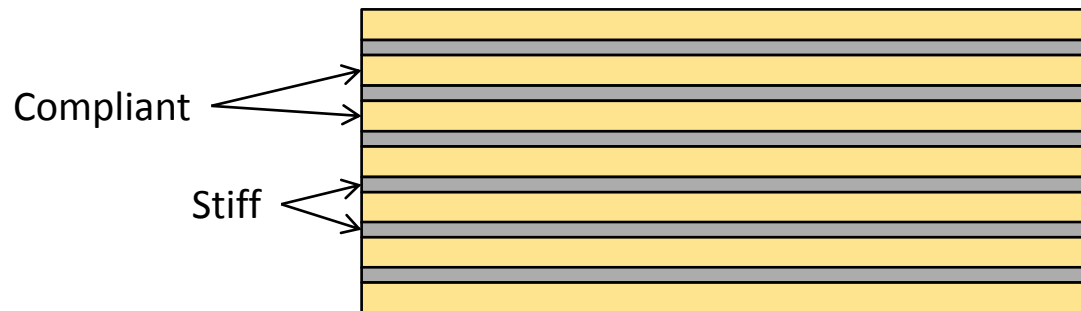
- Homogenization, neglecting the natural length scales of a system, often doesn't give good answers.



**Claim: Nonlocality is an essential feature of a realistic homogenized model of a heterogeneous material.**

# Nonlocality in a composite

- When we use a “smoothed out” displacement field, nonlocality appears in the equations.
- Example: alternate stiff and compliant layers.
- Prescribe the mean displacement (1D) to be a step function of position.



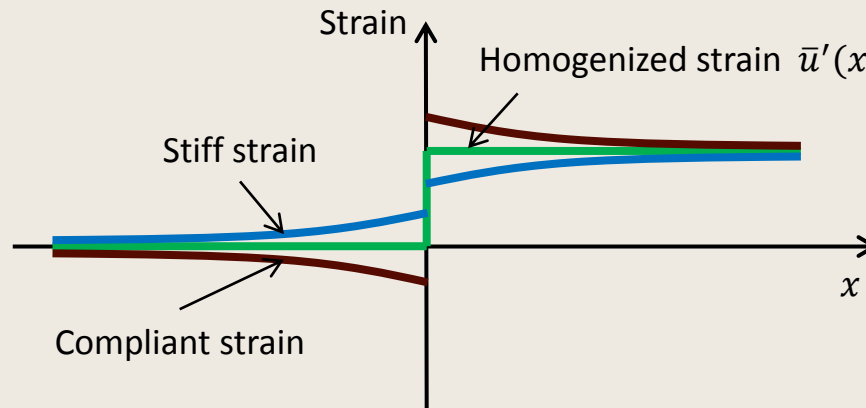
# Nonlocality in a composite, ctd.

- The evolution equation for the mean displacement field is nonlocal.

$$\rho \ddot{u}(x, t) = E_c \bar{u}''(x, t) + \gamma k \lambda^4 \int_{-\infty}^{\infty} (\bar{u}(p, t) - \bar{u}(x, t)) e^{-\lambda|x-p|} dp + b(x, t),$$

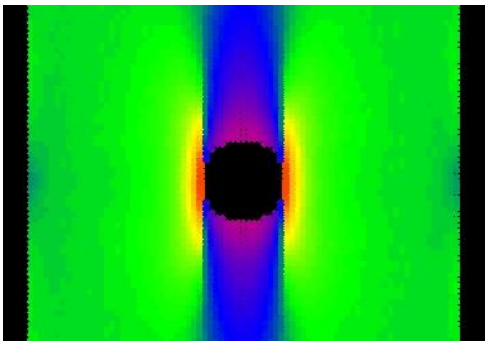
$$\frac{1}{\lambda} = \sqrt{\frac{E_s h_s h_c^2}{3\mu_c(h_s + h_c)}} = \text{length scale.}$$

Strain in each phase if the homogenized strain follows a step function

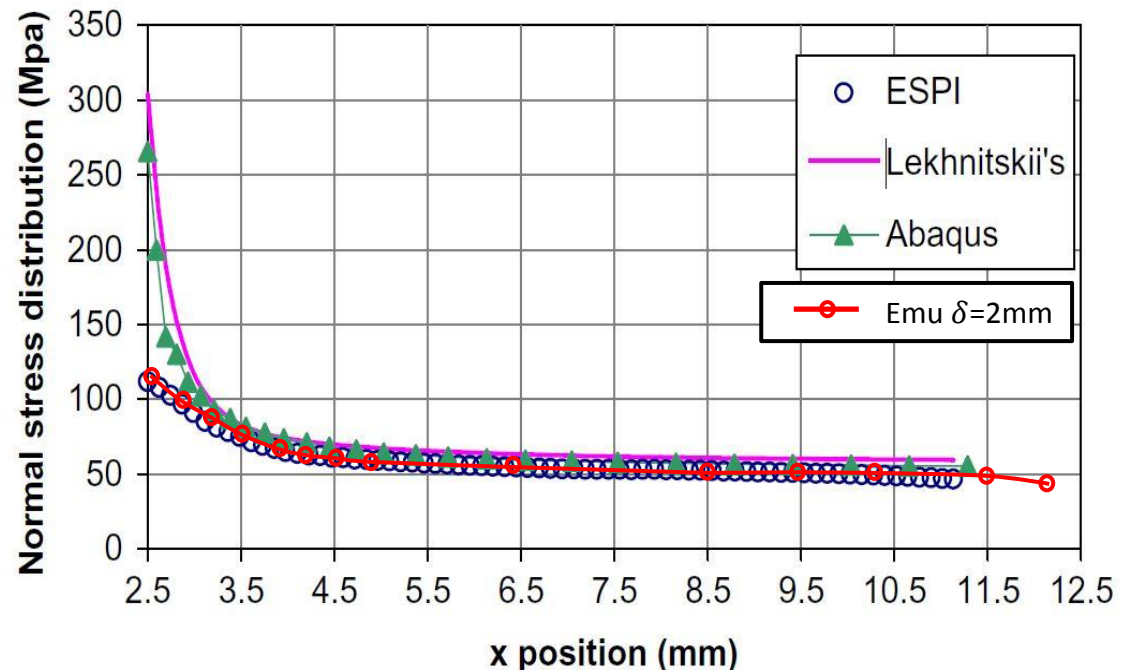


# Are composites nonlocal?

- Peridynamic model is more accurate than the local model for predicting stress concentration in a laminate.
- $h_s = h_c = 0.4\text{mm}$ ,  $E_s = 150\text{GPa}$ ,  $\mu_c = 4\text{GPa}$ .
- $\Rightarrow 1/\lambda = 1.41\text{mm}$ .



EMU: contours of longitudinal stress  
Horizon = 2mm



Data of Toubal, Karama, and Lorrain, Composite Structures 68 (2005) 31-36

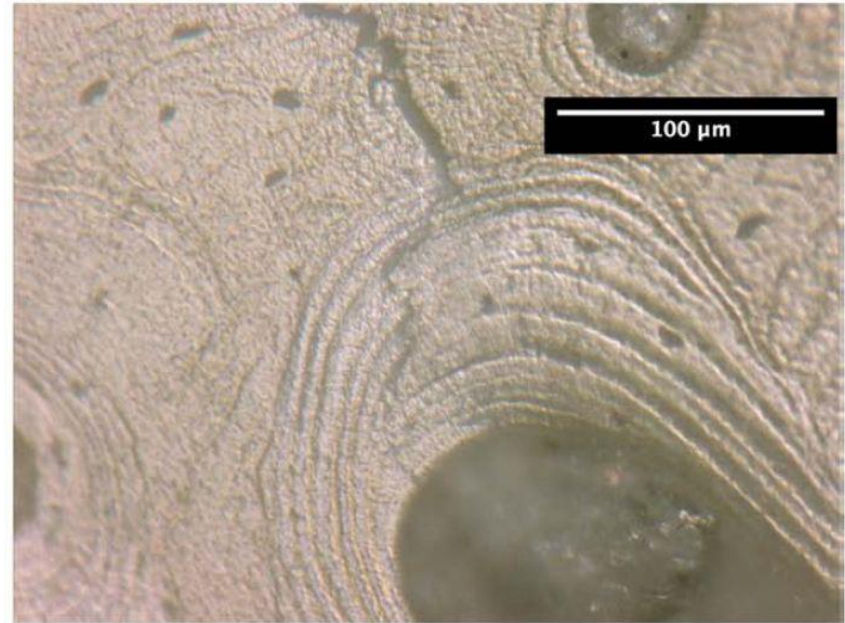
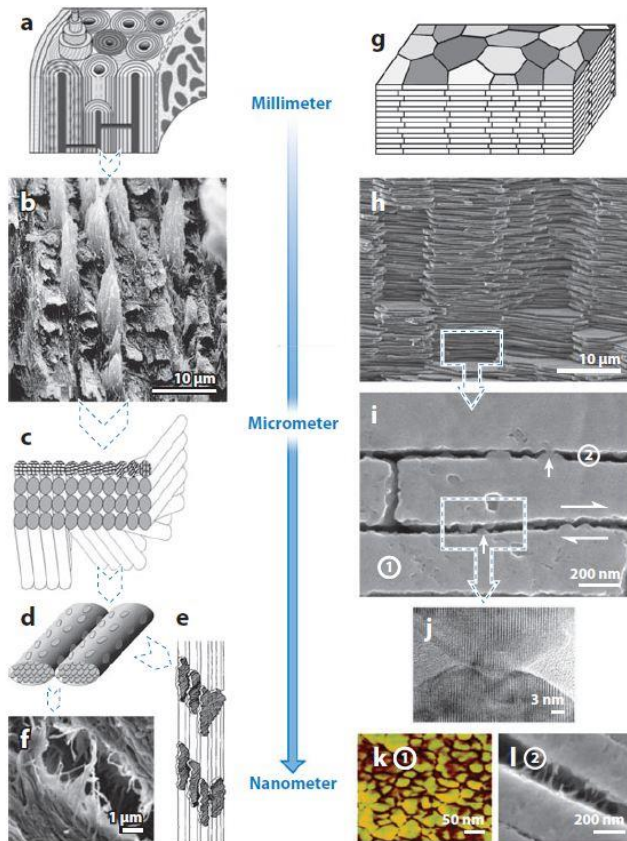
# Survey of some nonlocal theories

- MD-like forces, Cauchy-Born kinematics (Navier).
- Add nonlocal terms to the standard local model (Kroner).
- Average the local stress and use this in the Cauchy equilibrium equation (Eringen).
- Average the damage and use this in a local material model (Belytschko, Bazant).
- Include rotational DOFs in a local model (Cosserat).
- Include second partial derivatives of displacement in a local model (Coleman).
- Strongly nonlocal, linear elastic, mechanical only (Kunin).

*All of the above are unsuitable for modeling fracture.*

- Strongly nonlocal, nonlinear, mechanical + consistent thermodynamics + diffusion + damage (peridynamics).

# Bone: A composite material with many length scales

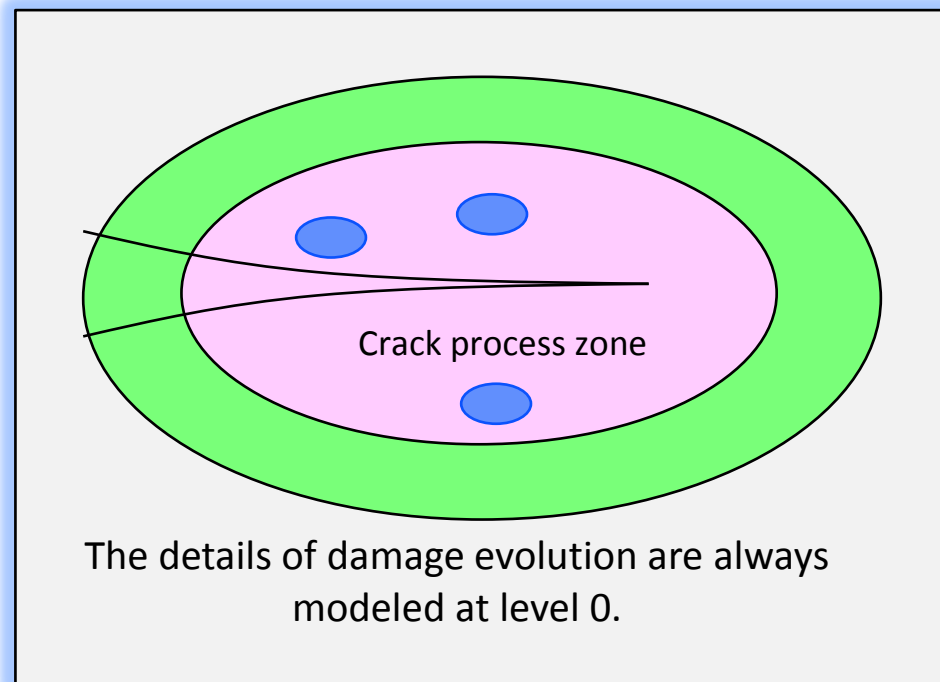


Bone structure helps delay, deflect crack growth. Image: Chan, Chan, and Nicolella, *Bone* 45 (2009) 427–434

Bone contains a hierarchy of structures at many length scales. Image: Wang and Gupta, *Ann. Rev. Mat. Sci.* 41 (2011) 41-73

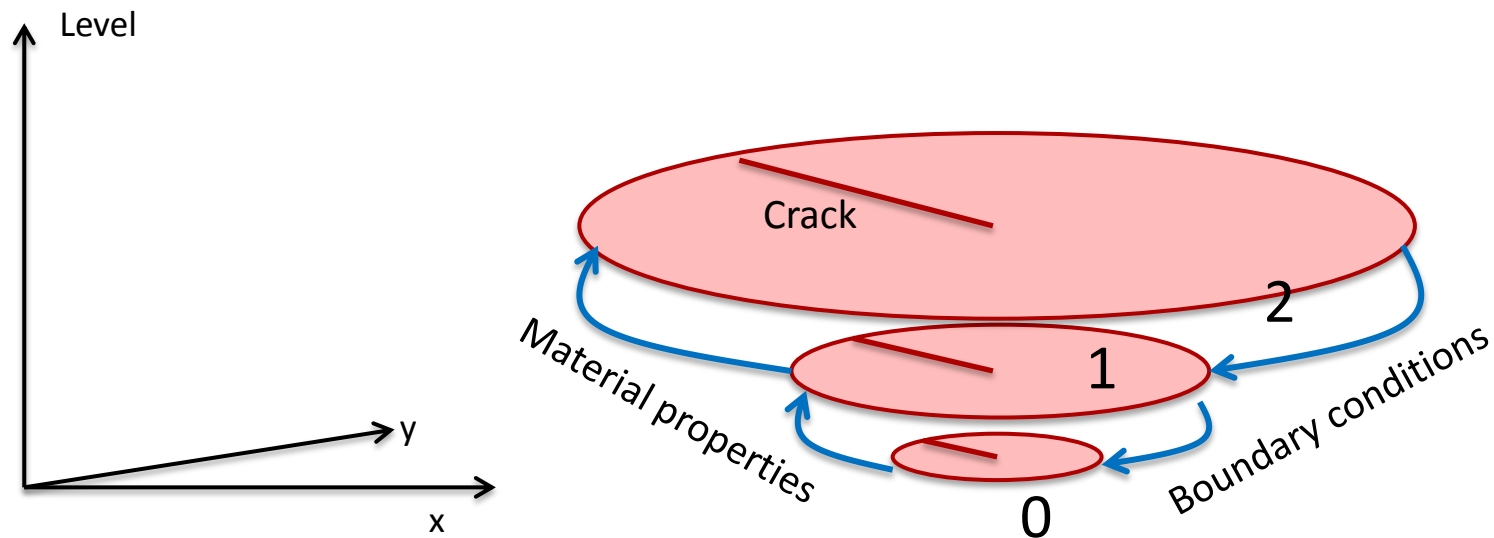
# Concurrent multiscale method for defects

- Apply the best practical physics at the smallest length scale (near a crack tip).
- Scale up hierarchically to larger length scales.
- Each level is related to the one below it by the same equations.
  - Any number of levels can be used.
- Adaptively follow the crack tip.



# Concurrent solution strategy

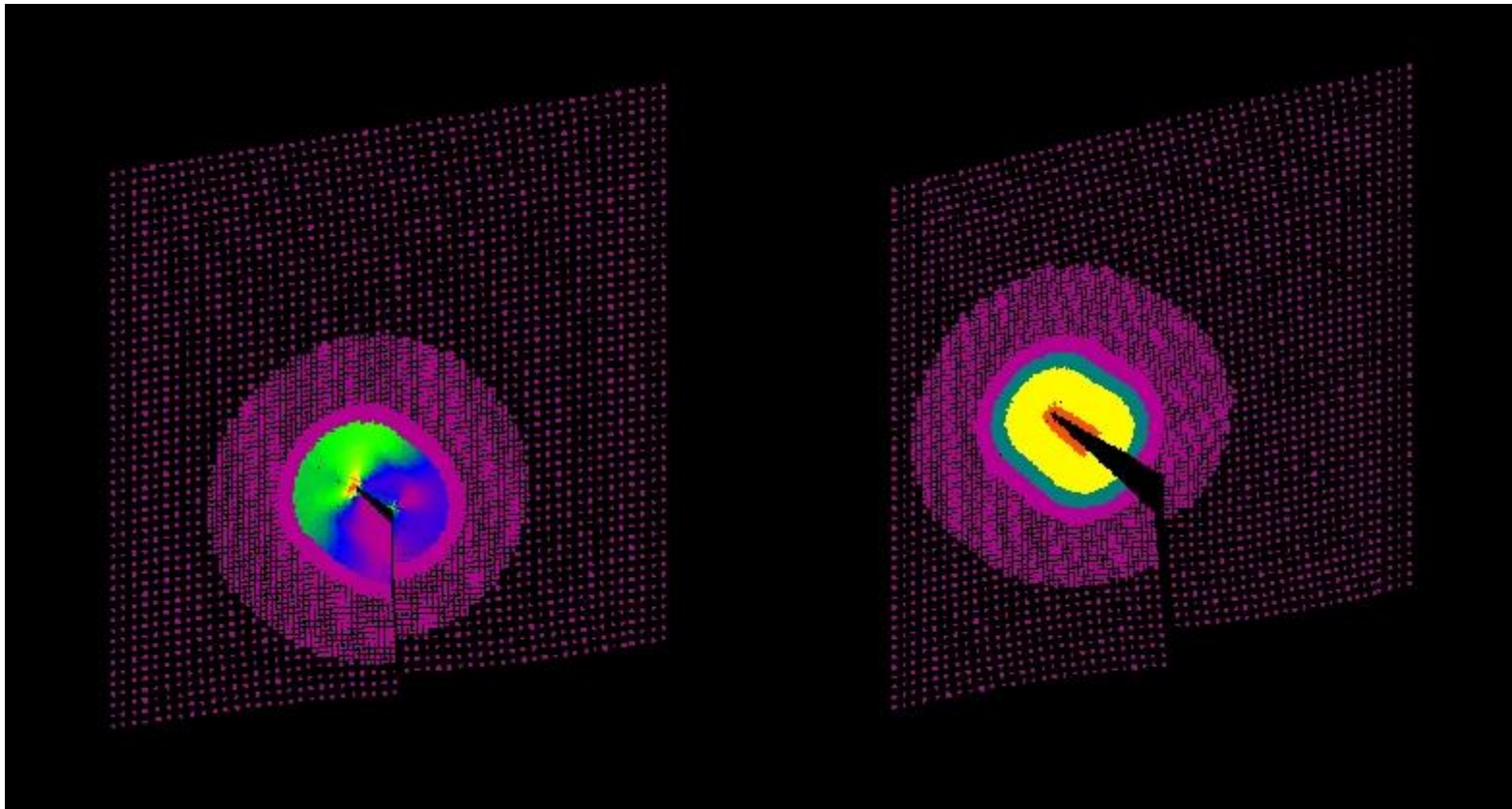
- The equation of motion is applied only within each level.
- Higher levels provide boundary conditions on lower levels.
- Lower levels provide coarsened material properties (including damage) to higher levels.
- In principle, a large number of levels can be used, all coupled in the same way: “scalable multiscale” method.



Schematic of communication between levels in a 2D body

# Concurrent multiscale example: shear loading of a crack

- Level 0 region adaptively follows the crack tip.

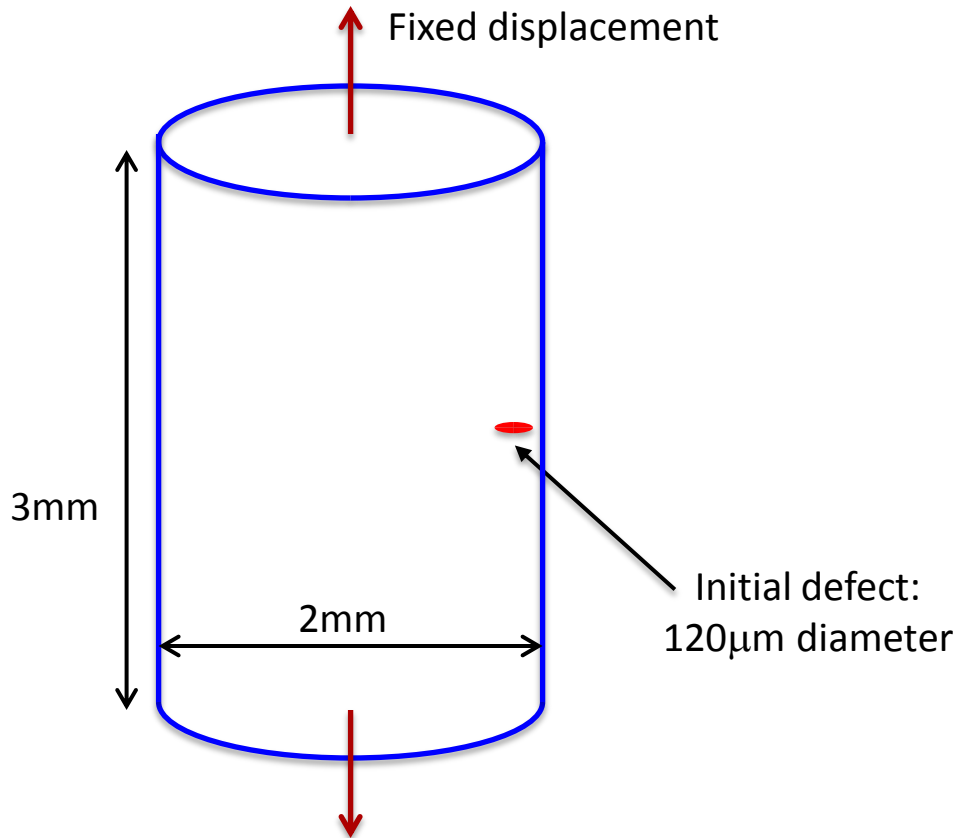


Bond strain

Damage process zone

# Failure of a glass rod in tension

- A classical test problem for fractography.
- We will try to reproduce key fractographic features.
- Multiscale approach allows us to make the horizon  $\ll$  geometric length scales.



$$\rho = 3000 \text{ kg/m}^3$$

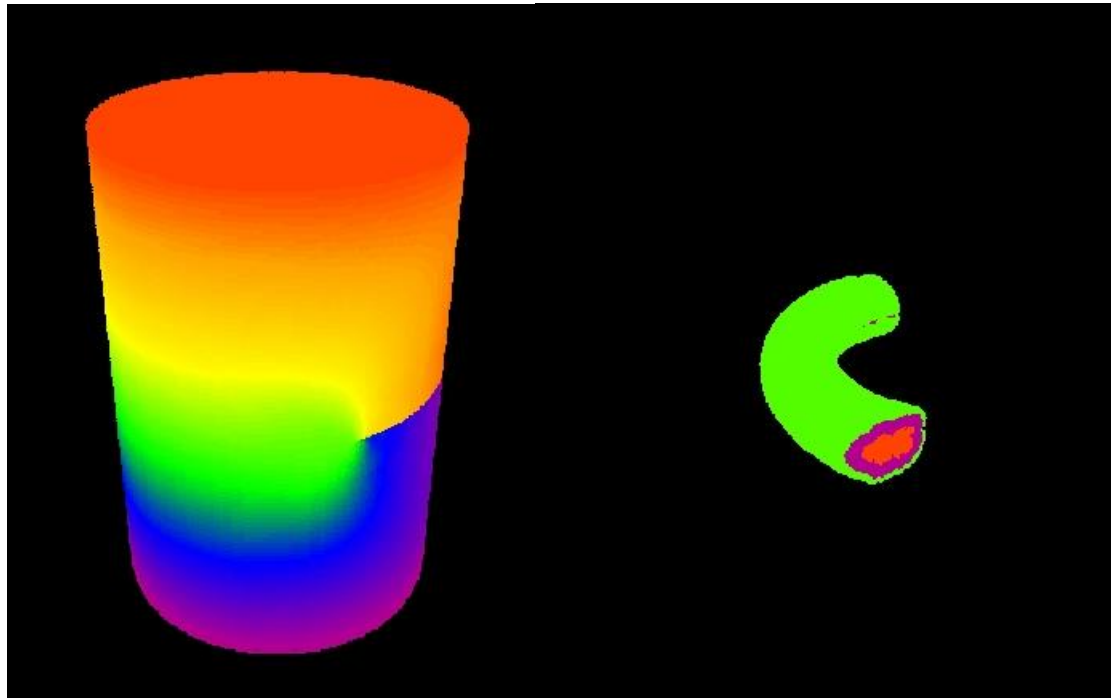
$$E = 70.5 \text{ GPa}$$

$$\nu = 0.25$$

$$G_{Ic} = 7.0 \text{ J/m}^2$$

$$\delta = 25 \mu\text{m}$$

# Failure of a glass rod in tension



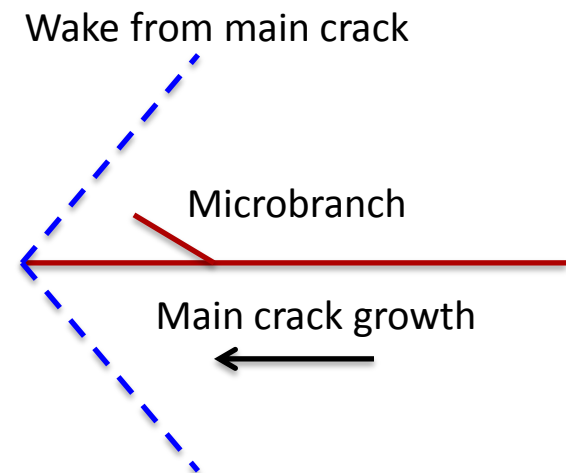
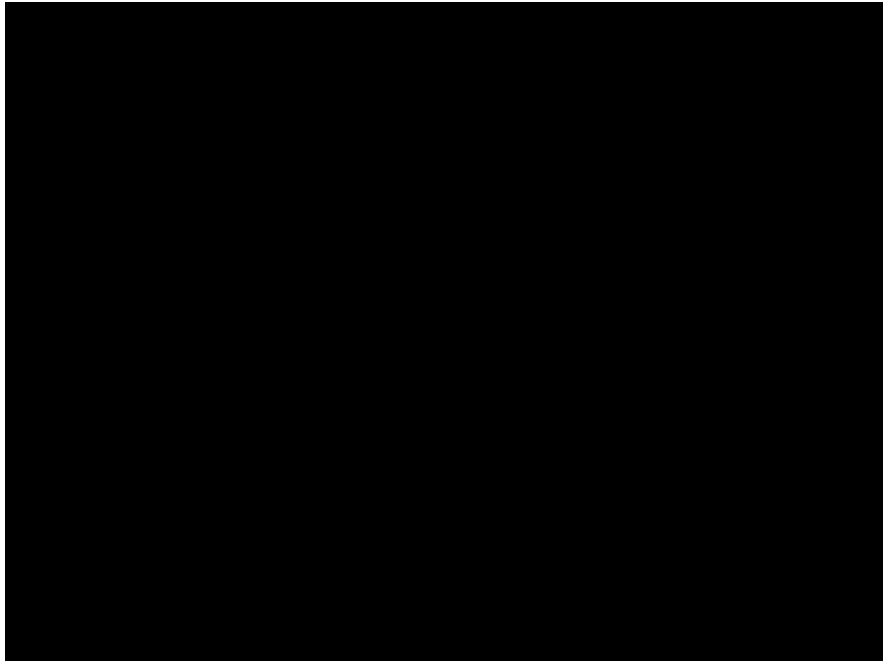
Level 1 displacement

Level 0 surrounds the crack front

- Level 1 multiscale.
- 20,000,000 level 0 sites (most are never used).
- Level 0 horizon is 25 $\mu$ m.

# Failure of a glass rod in tension (movie)

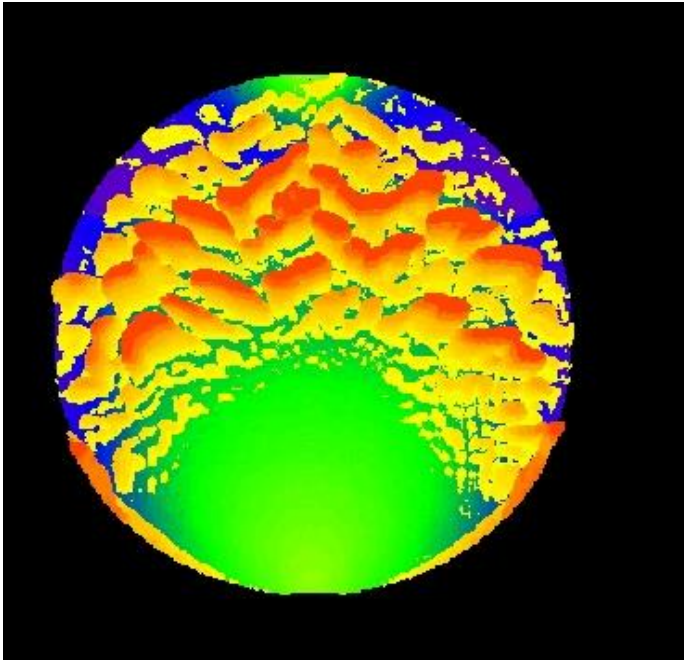
Evolution of surface roughness (movie)



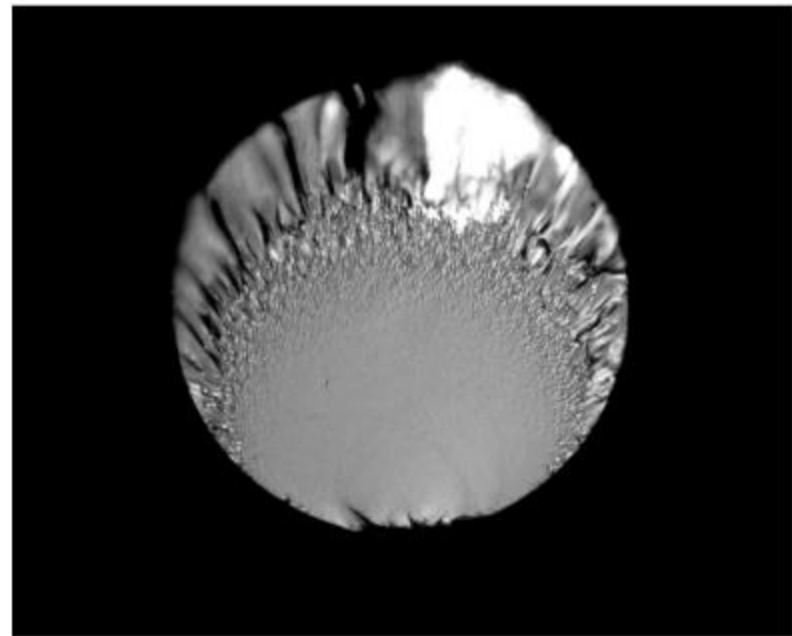
- Rough features branch off from the main crack.
- Each one grows slower than the main crack and eventually dies.

# Mirror-mist-hackle

- Model predicts roughness and microbranches that increase in size as the crack grows.
- Transition radius decreases as initial stress increases – trend agrees with experiments.

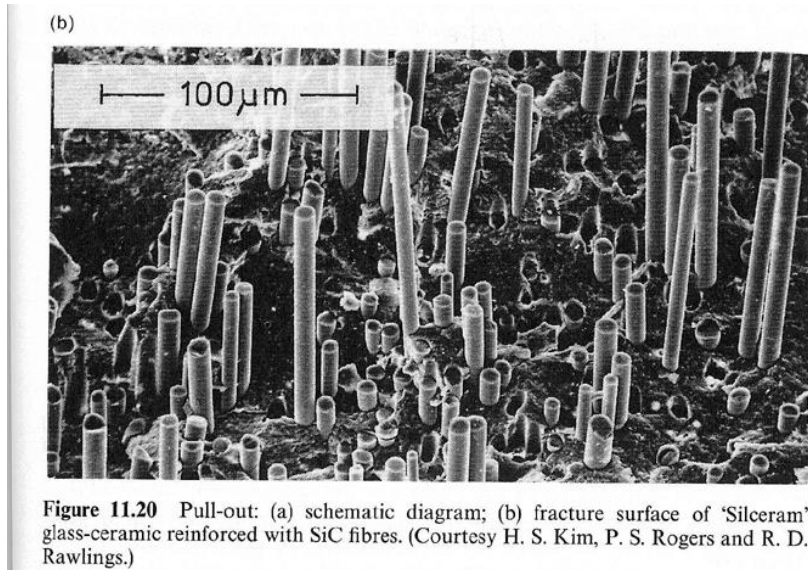


3D peridynamic model

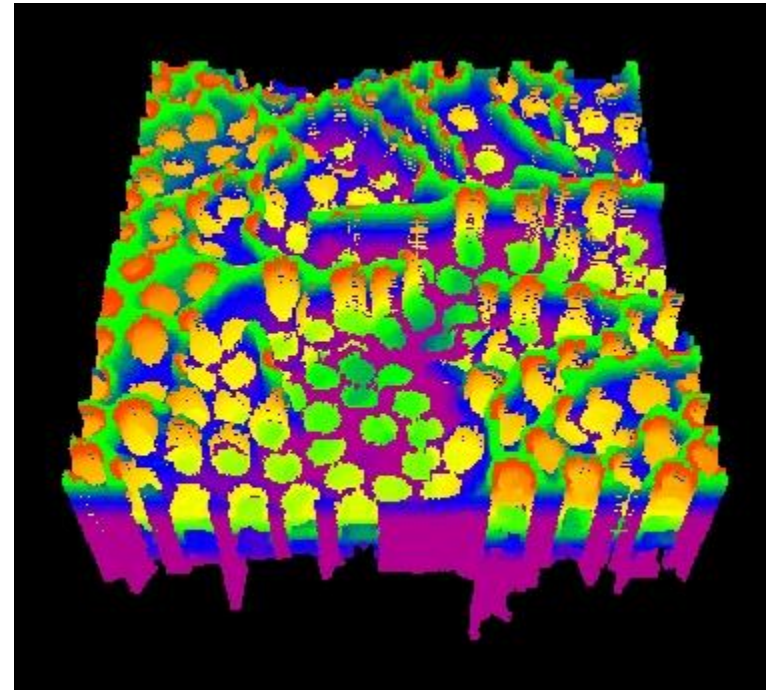


Fracture surface in a glass optical fiber  
(Castilone, Glaesemann & Hanson, Proc. SPIE (2002))

# Composite fracture features

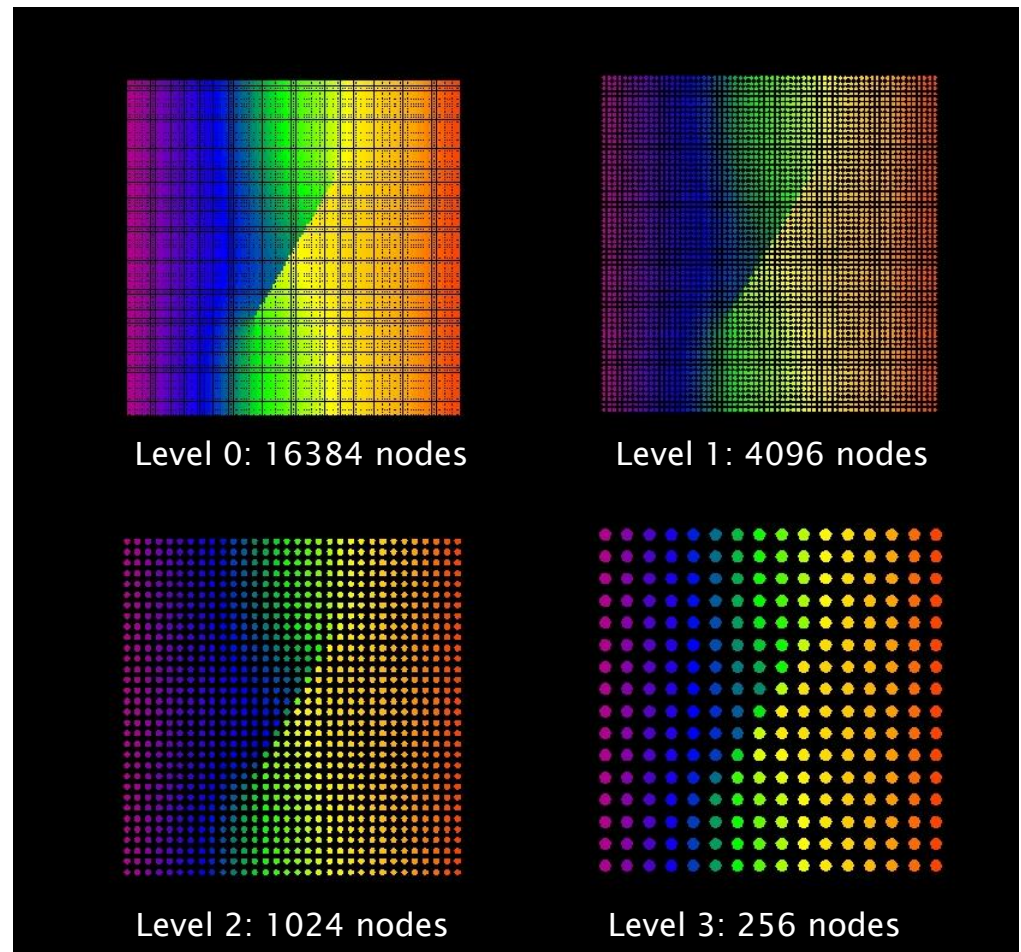


Complex crack path in a composite



# Multiscale verification: crack in a plate

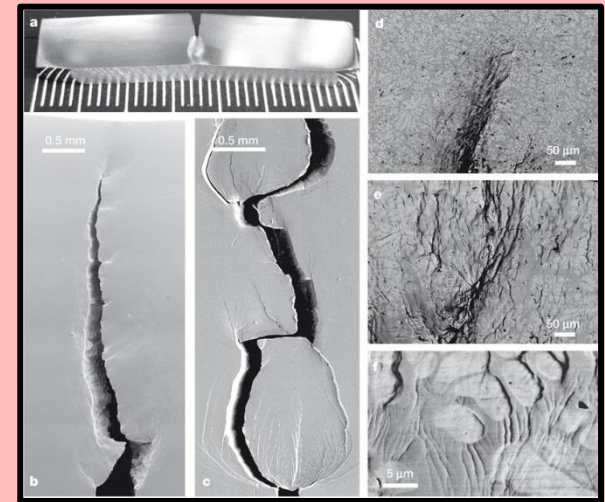
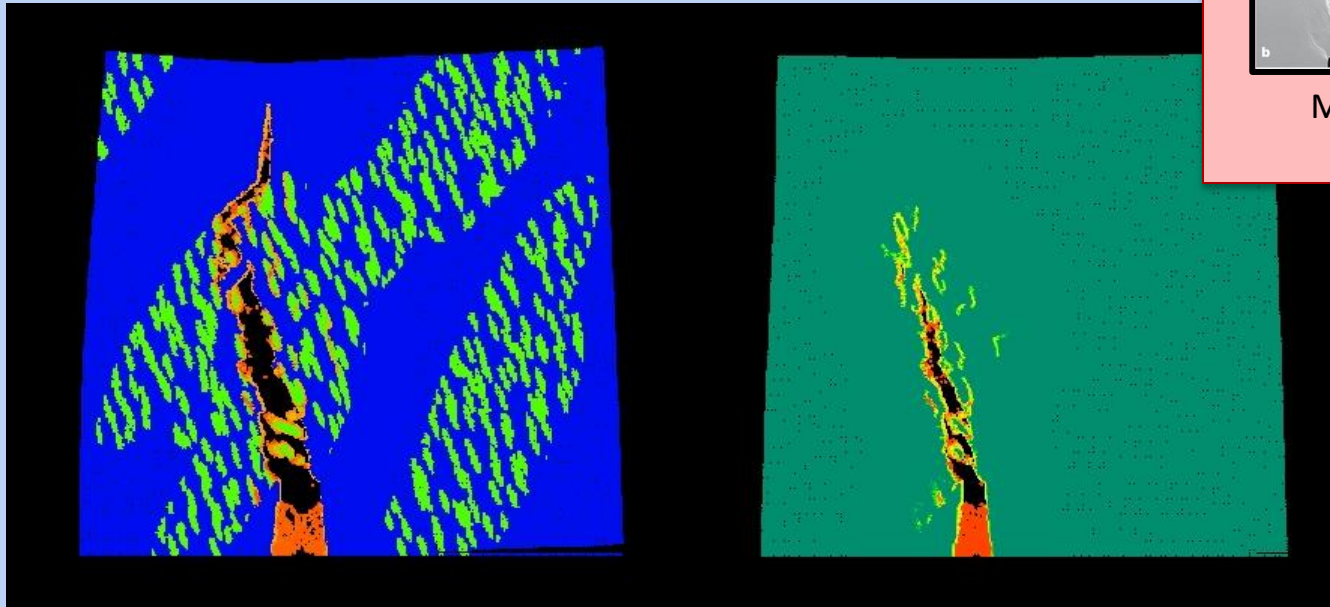
- Example: Solve the same problem in four different levels using the successively upscaled material properties – results are the same.



# Multiscale modeling reveals the structure of brittle cracks

- Material design requires understanding of how morphology at multiple length scales affects strength.
- This is a key to material reliability.

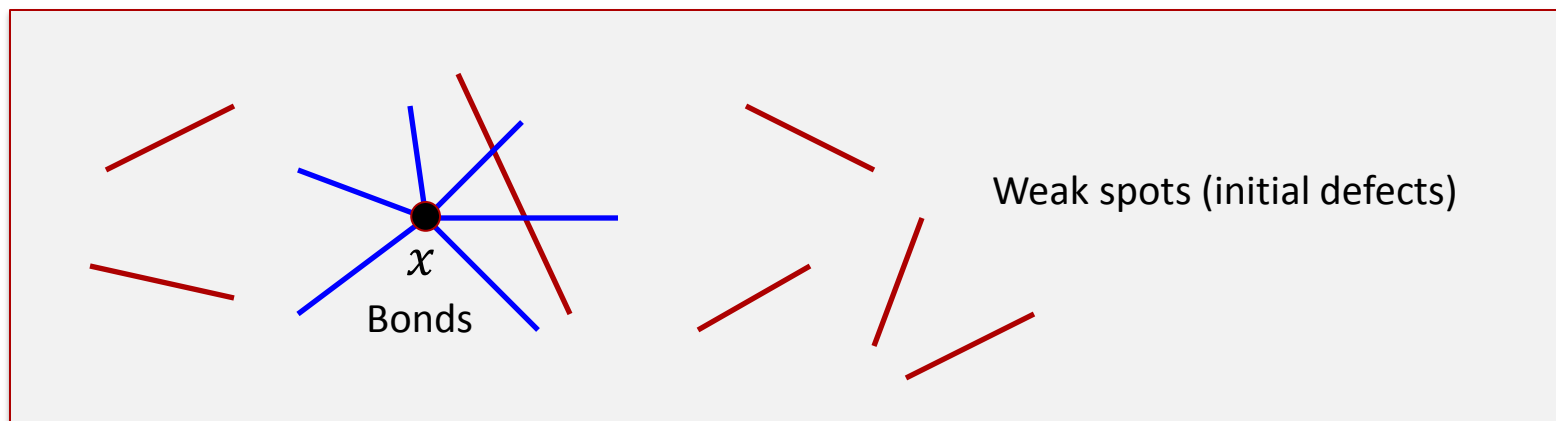
Multiscale model of crack growth through a brittle material with distributed defects



Metallic glass fracture (Hofmann et al, Nature 2008)

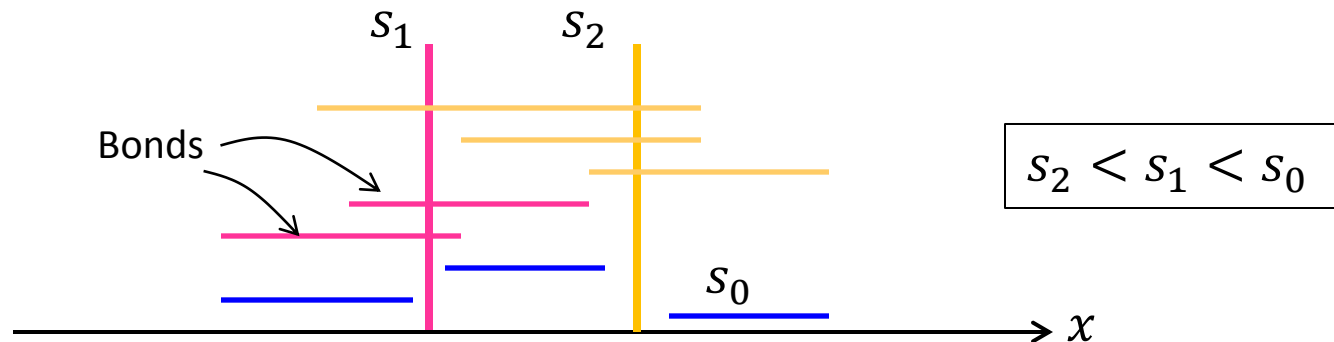
# How should we treat a random distribution of defects?

- Can the random nature of fracture be reflected in continuum level properties?
  - Wrong ways:
    - Local: randomly assign element properties.
    - Peridynamic: randomly assign critical strains to bonds.
  - Right way:
    - Start with a realization of a random distribution of initial defects.
    - Assign critical bonds strains that reflect each realization.

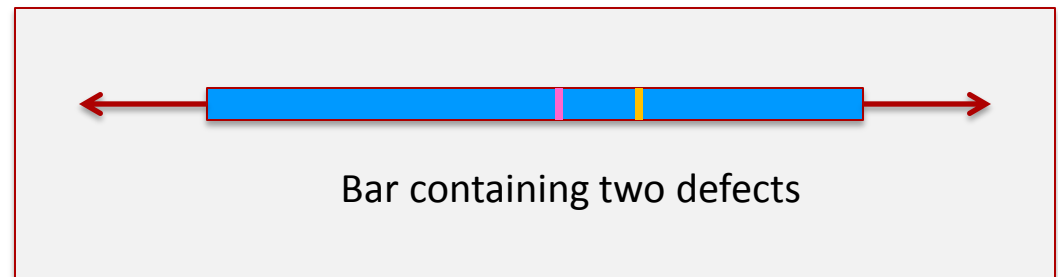


# Modeling randomness in fracture

- Assign a critical strain to each defect  $s_i$ .
- Set each bond critical strain to the minimum of those of the defects it crosses.

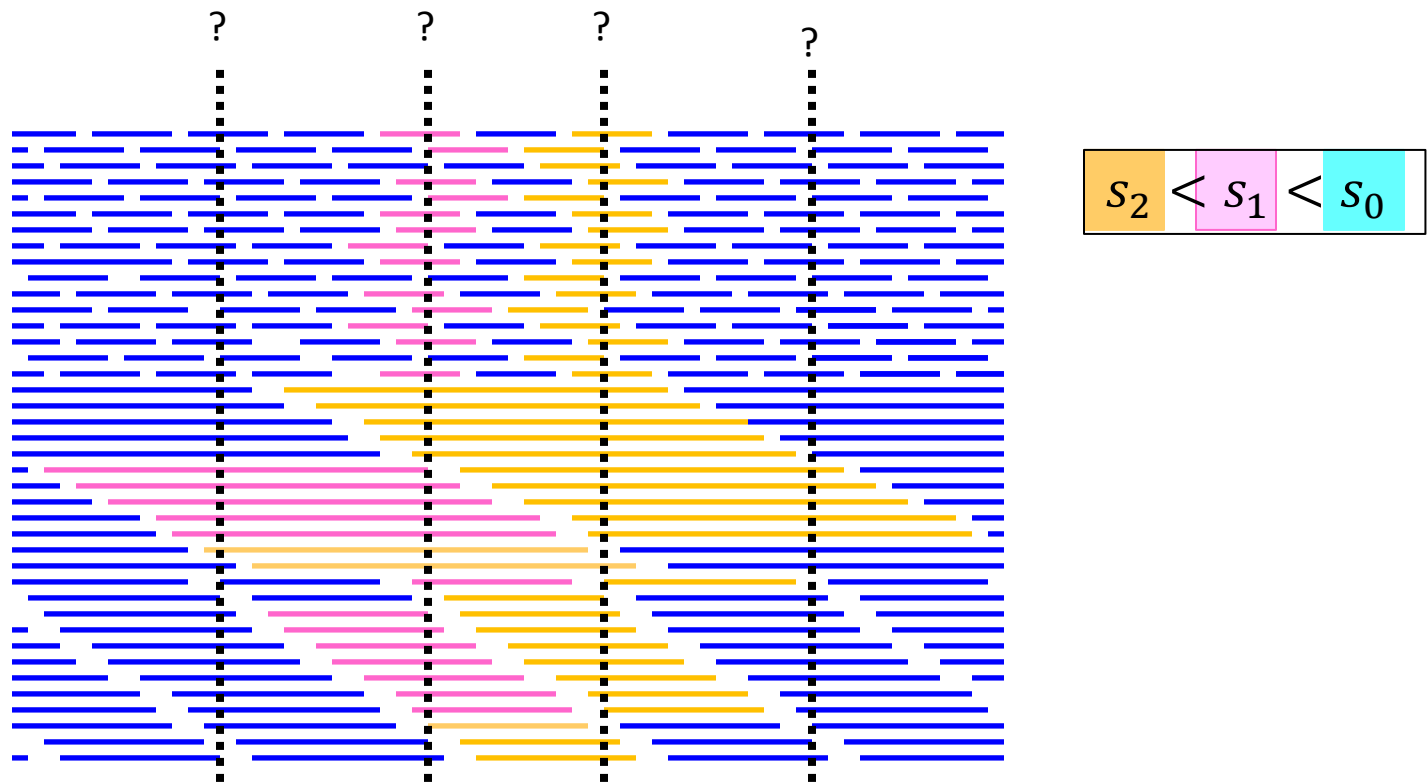


$s_0 \dots$  undamaged critical strain  
 $s_1 \dots$  defect 1 critical strain  
 $s_2 \dots$  defect 2 critical strain



# Where would fracture occur in this model?

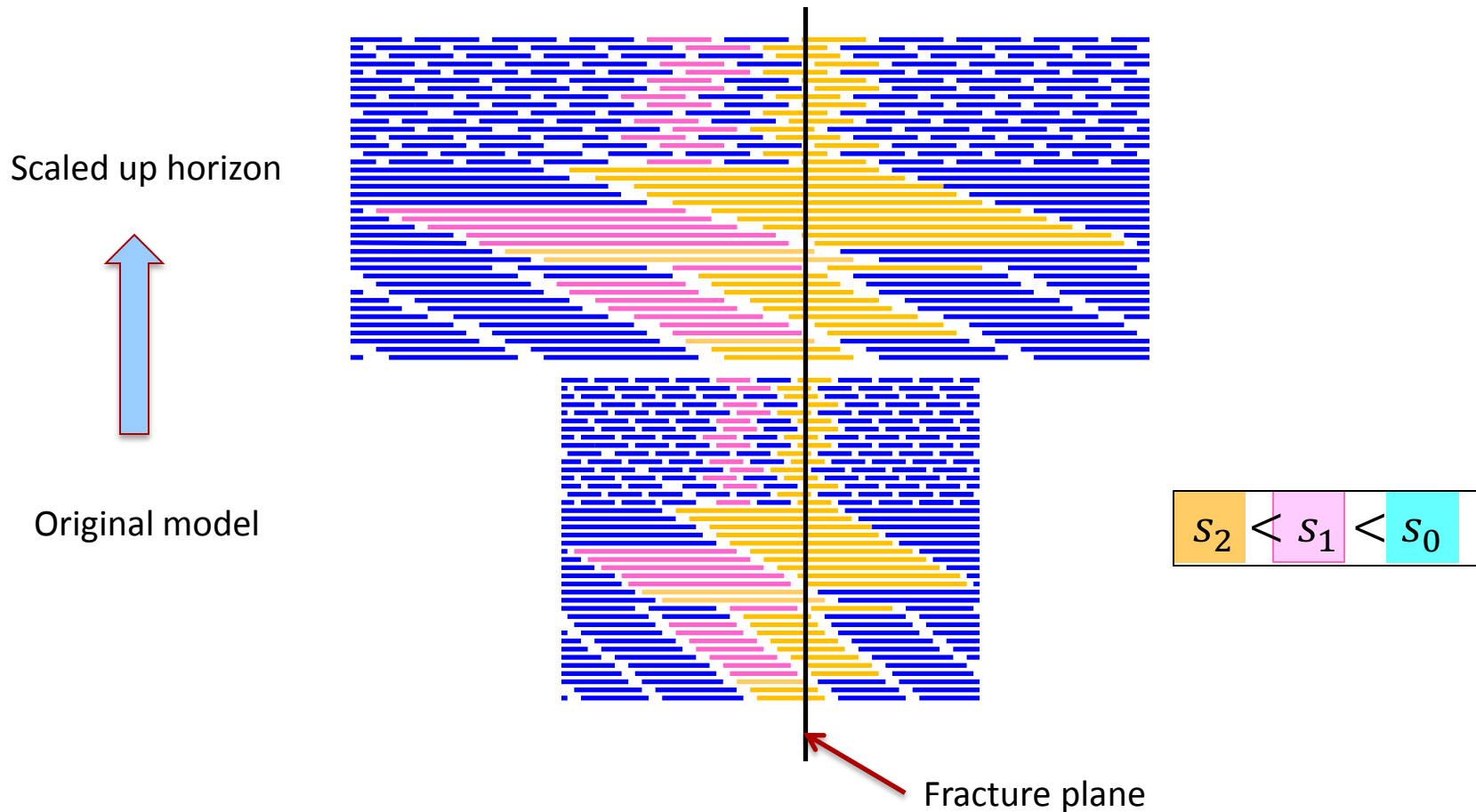
- Assign a critical strain to each defect  $s_i$ .
- Set each bond critical strain to the minimum of those of the defects it crosses.



Colors show critical strain in some bonds

# Rescaling the critical strains

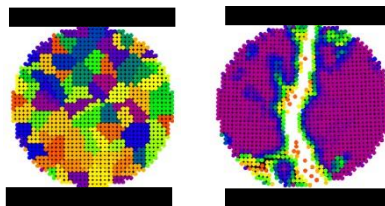
- Set each new bond critical strain to the minimum of those of the original bonds it intersects.
- The rescaled model still fractures at the weakest spot.



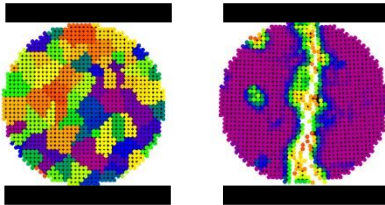
# Strategy for treating randomness in fracture due to material variability

- Define a set of realizations of randomly distributed initial defects.
- Assign level 0 bond strains.
- Scale up to any desired length scale The scaled up model contains the defects.

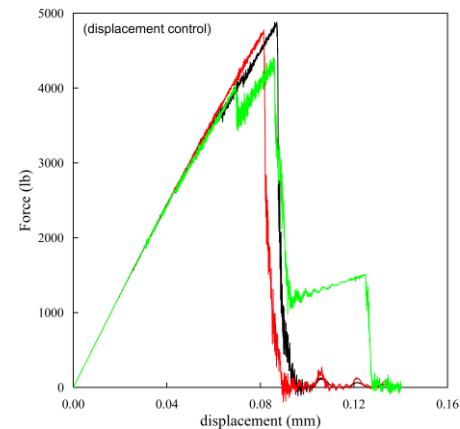
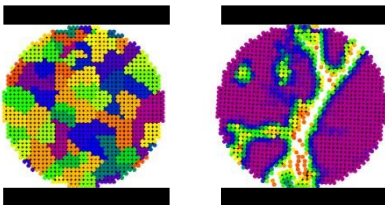
material realization 1



material realization 2



material realization 3



Brazilian test: J Bishop & J. Holland

# Discussion

- It's fairly clear that nonlocality is an essential feature of any mechanics theory that seeks to accomplish what peridynamics seeks to accomplish.
- Nonlocal interactions can arise due to the way we choose to model things.
  - Especially heterogeneous media.
- Nonlocality may also be an essential feature of a scalable multiscale approach to treating large numbers of small defects.

*Do we need a different word for nonlocal continuum mechanics?*