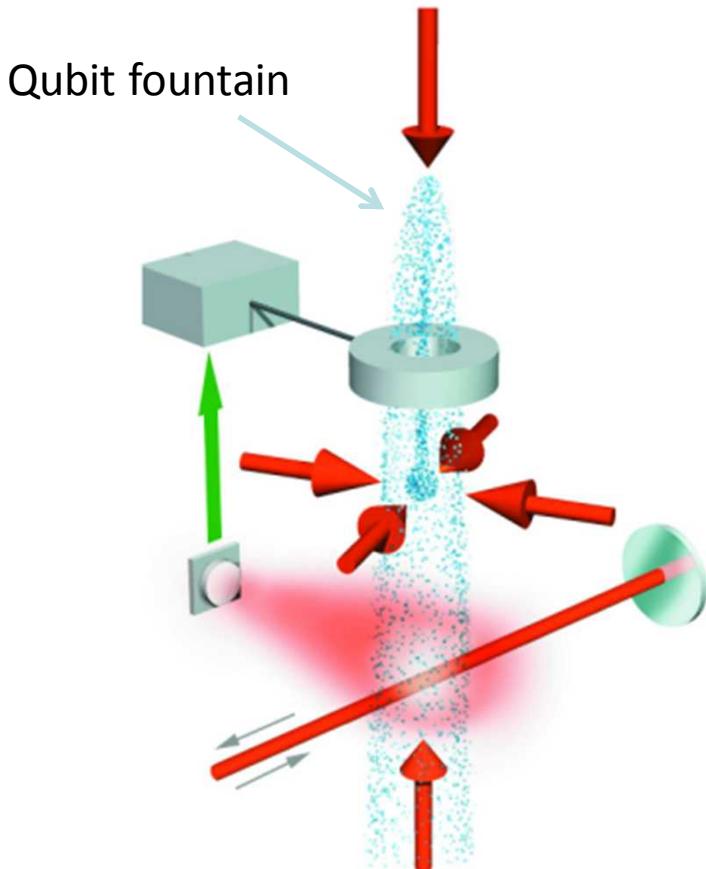


Quantum sensors

Grant Biedermann

Quantum-Coherence



Atomic fountain principle

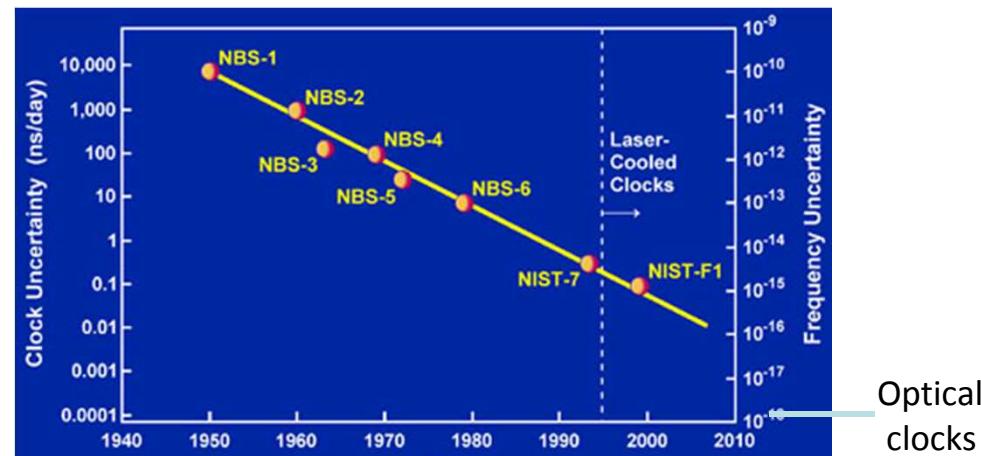
http://smsc.cnes.fr/PHARAO/GP_instrument.htm

Outstanding quantum coherence in neutral atoms enables precision metrology and quantum information

- Example: atomic clocks

$$\left| 6^2 S_{1/2}; F = 3, M_F = 0 \right\rangle \leftrightarrow \left| 6^2 S_{1/2}; F = 4, M_F = 0 \right\rangle$$

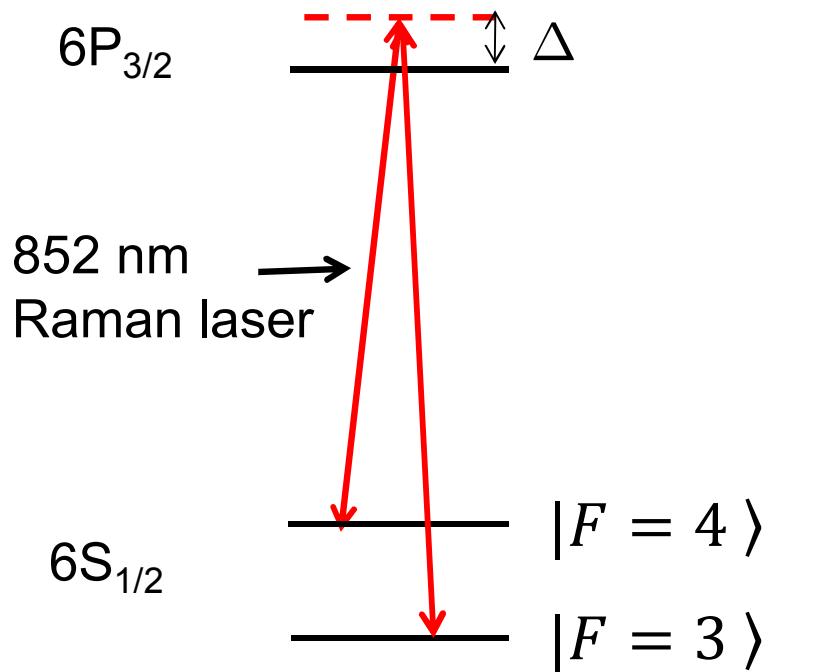
$$\left| 0 \right\rangle \leftrightarrow \left| 1 \right\rangle$$



<http://www.nist.gov>

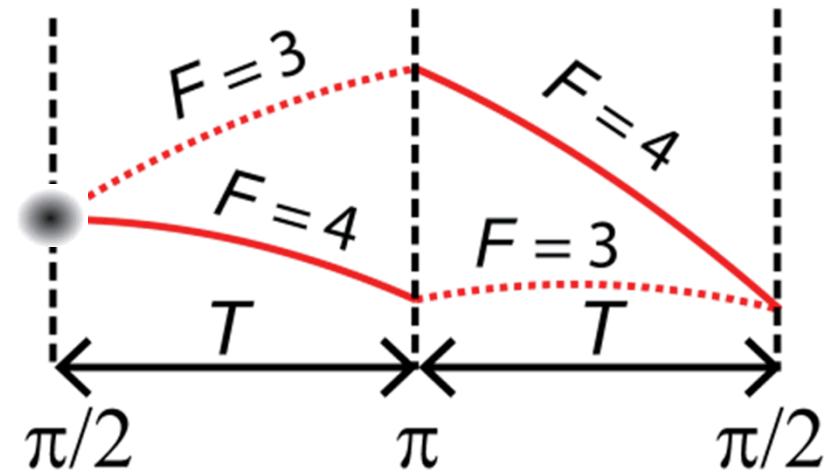
Typical accuracy now better than one part in 10^{15}

Light-pulse atom interferometry



stimulated Raman transition

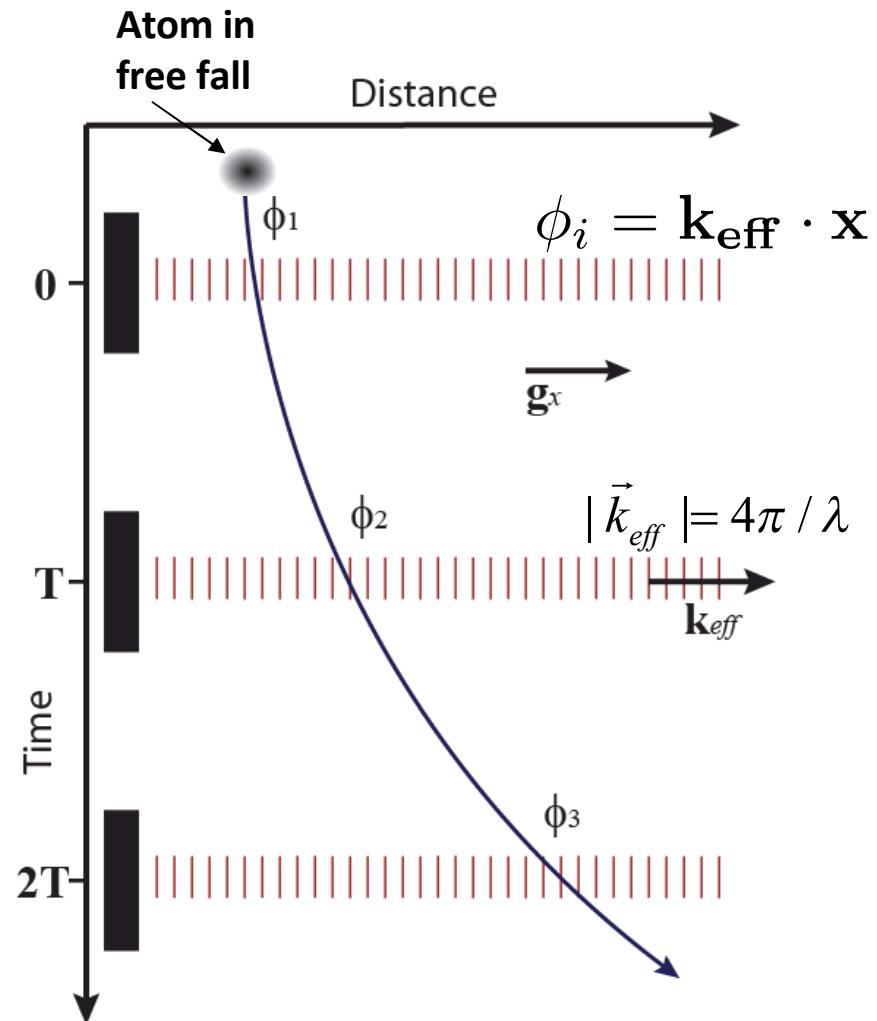
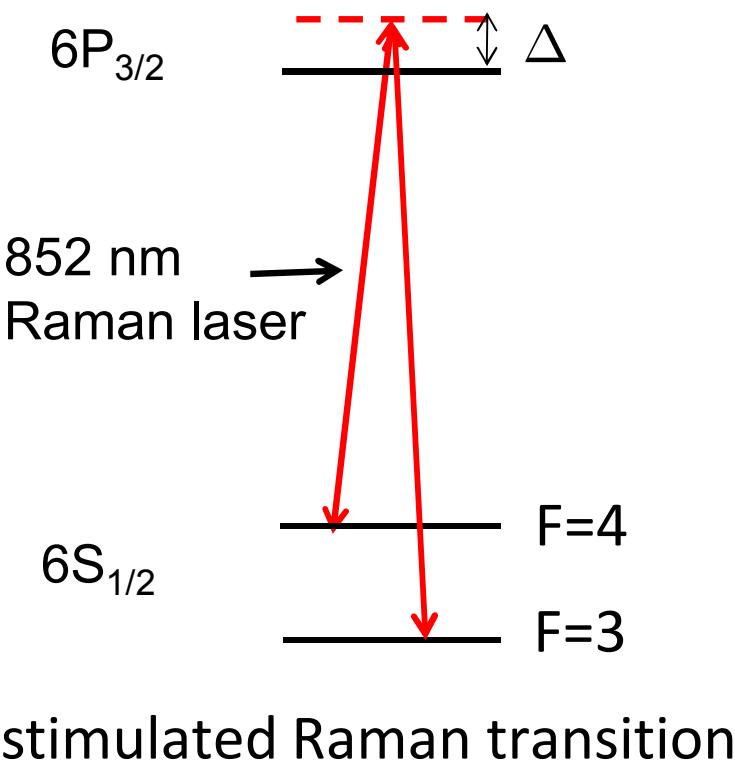
wavepacket trajectory



- Exceptional accelerometers and gyroscopes nrad/s/ $\sqrt{\text{Hz}}$, ng/ $\sqrt{\text{Hz}}$ to pg/ $\sqrt{\text{Hz}}$
- Growing commercial and govt. interest in fielding this technology

Kasevich, and Chu, Phys. Rev. Lett. 67, 181–184 (1991)

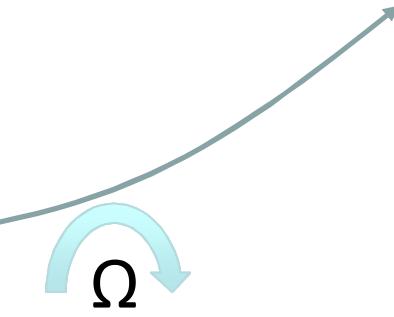
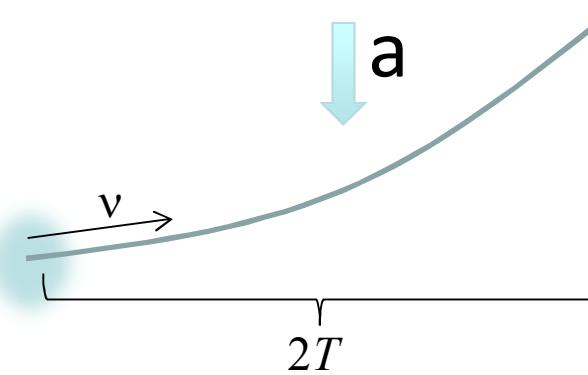
Light-pulse atom interferometry



For an atom starting in F=3: $P_{F=4} = \frac{1}{2}(1 - \cos \Delta\phi)$ & $\Delta\phi = \phi_1 - 2\phi_2 + \phi_3$

Measuring acceleration and rotation with a particle in free-fall

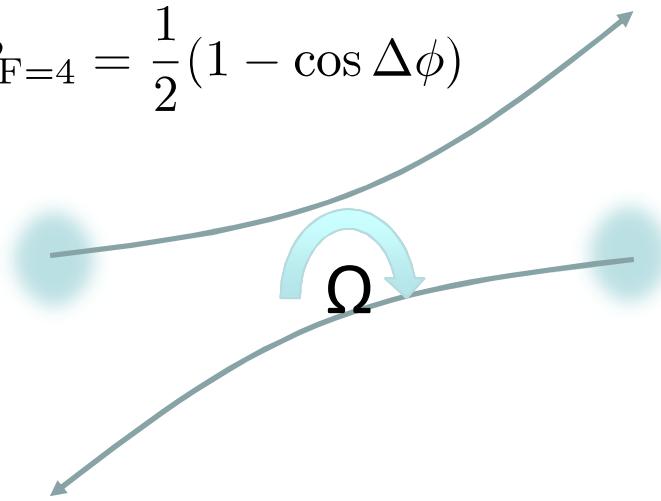
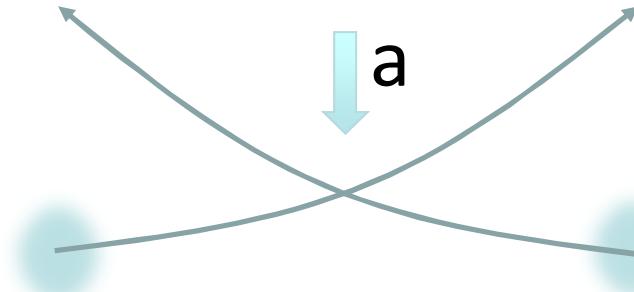
1-particle



Interferometer phase shift: $\Delta\phi = \vec{k}_{eff} \cdot \left(\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2 \right)$

$$P_{F=4} = \frac{1}{2}(1 - \cos \Delta\phi)$$

2-particles



Sensitivity increases with T^2

Launch and recapture

Steady state atom number:

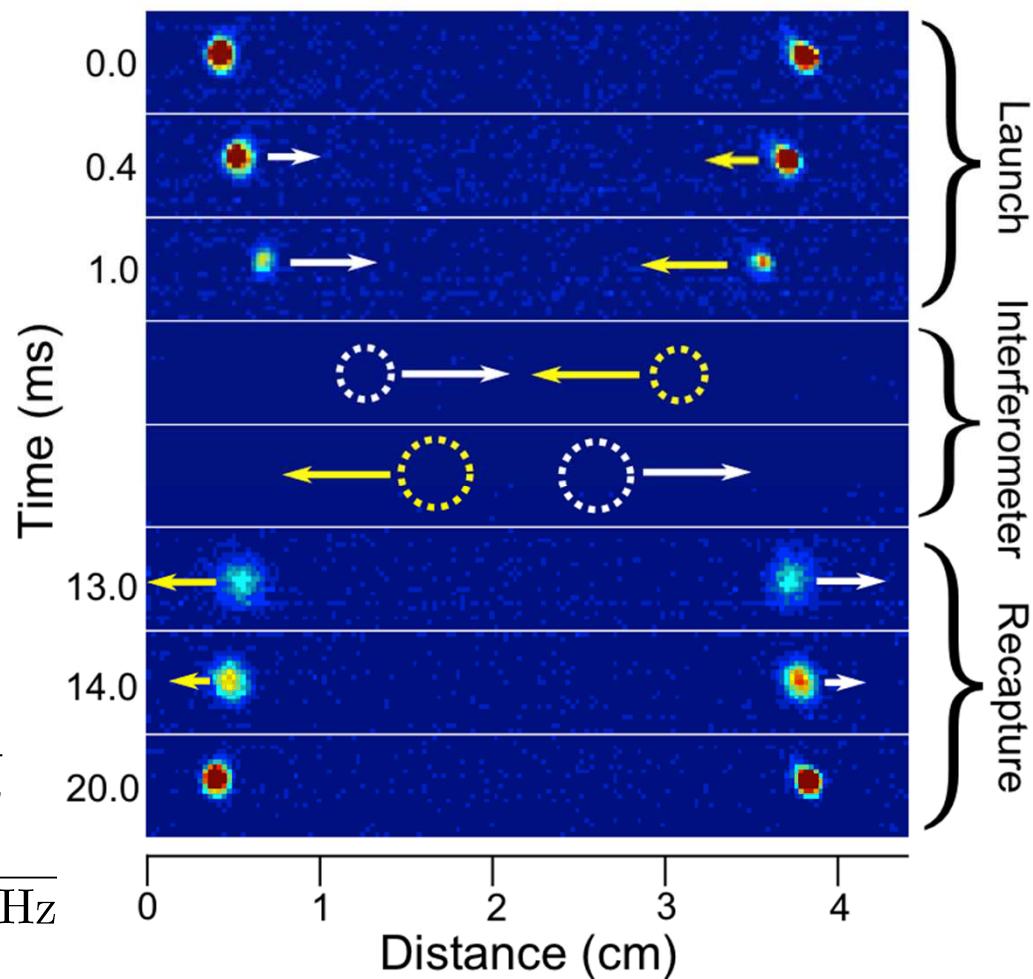
$$N_s = \frac{\alpha\eta T_c}{\beta T_c + (1 - r_0)}.$$

Base recapture efficiency $r_0 = 96\%$

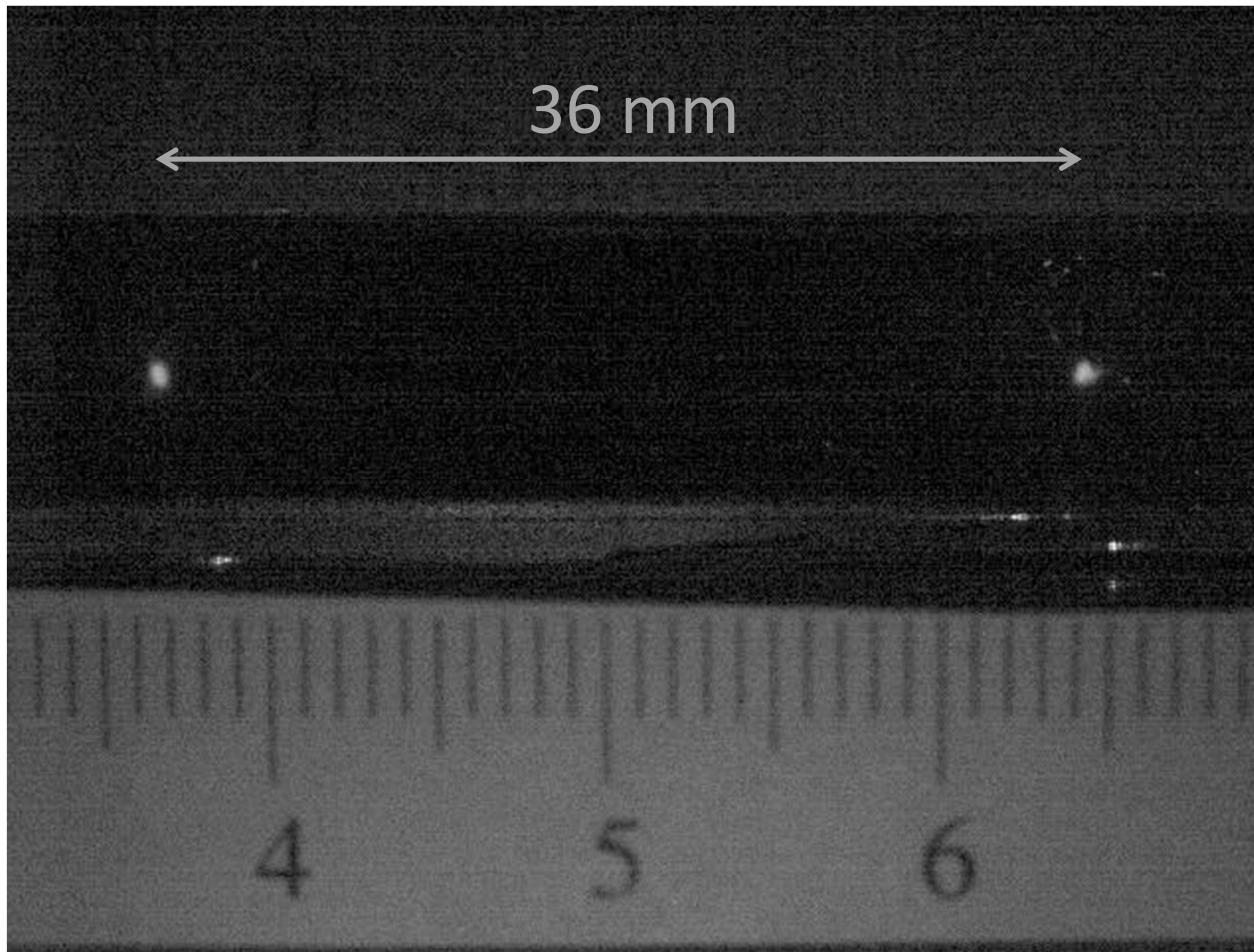
Benefits

- Increases signal by 10x
- Data rates > 50 Hz
- Minimizes cycle dead time
- Reduced complexity
- Sufficient for:
 $33 \text{ ng}/\sqrt{\text{Hz}} \text{ & } 70 \text{ nrad}/\text{s}/\sqrt{\text{Hz}}$
- Demonstrated:
 $900 \text{ ng}/\sqrt{\text{Hz}} \text{ & } 1100 \text{ nrad}/\text{s}/\sqrt{\text{Hz}}$

CCD images of ensemble exchange

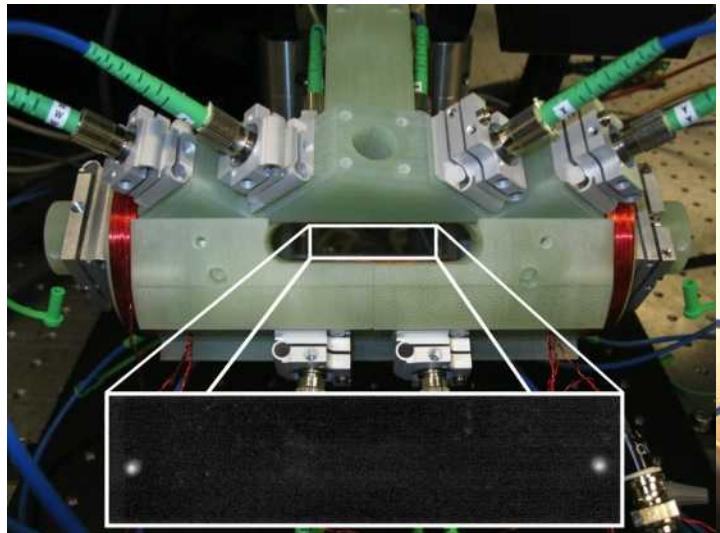


Launch and recapture

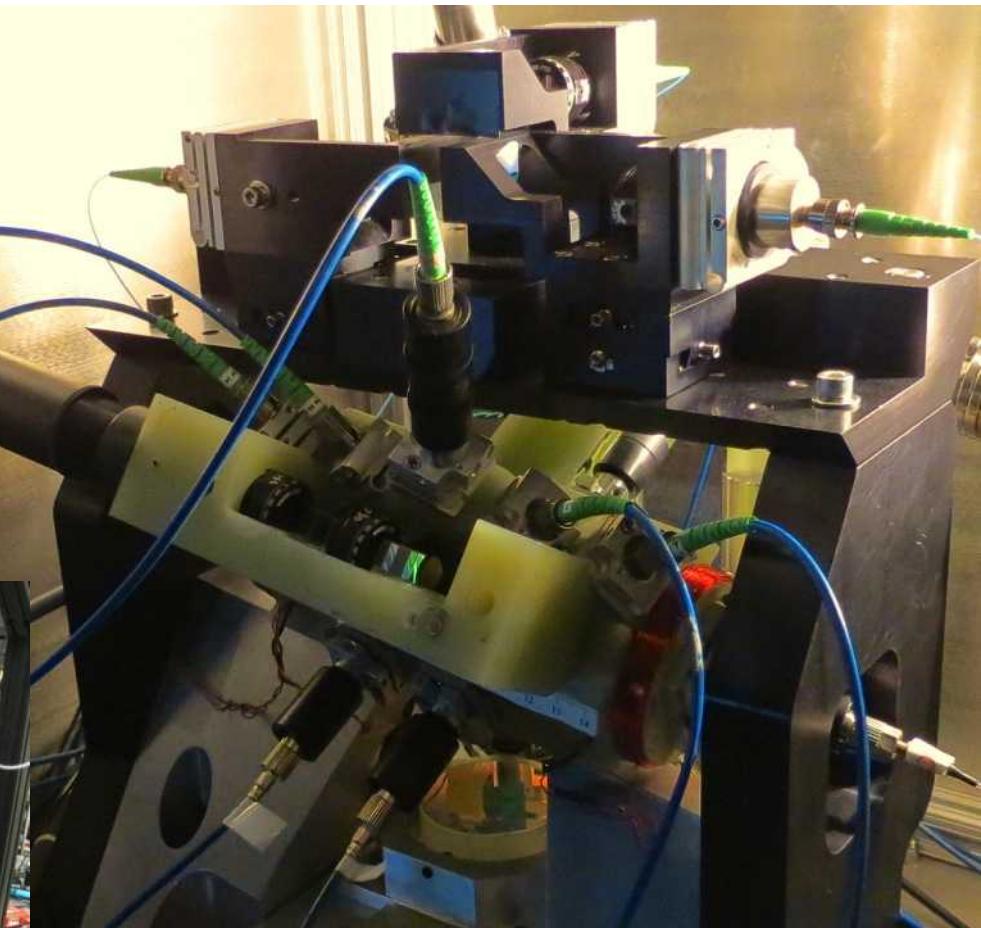


- Repeats at ≈ 60 measurements per second

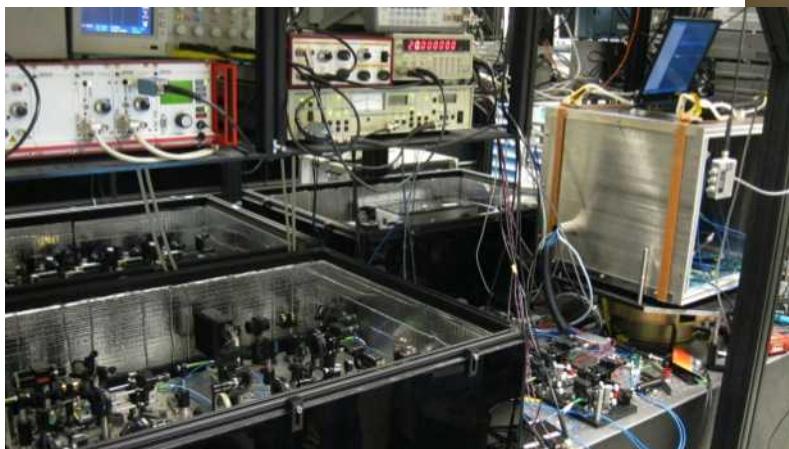
Experiment platform



Picture of interferometer sensor

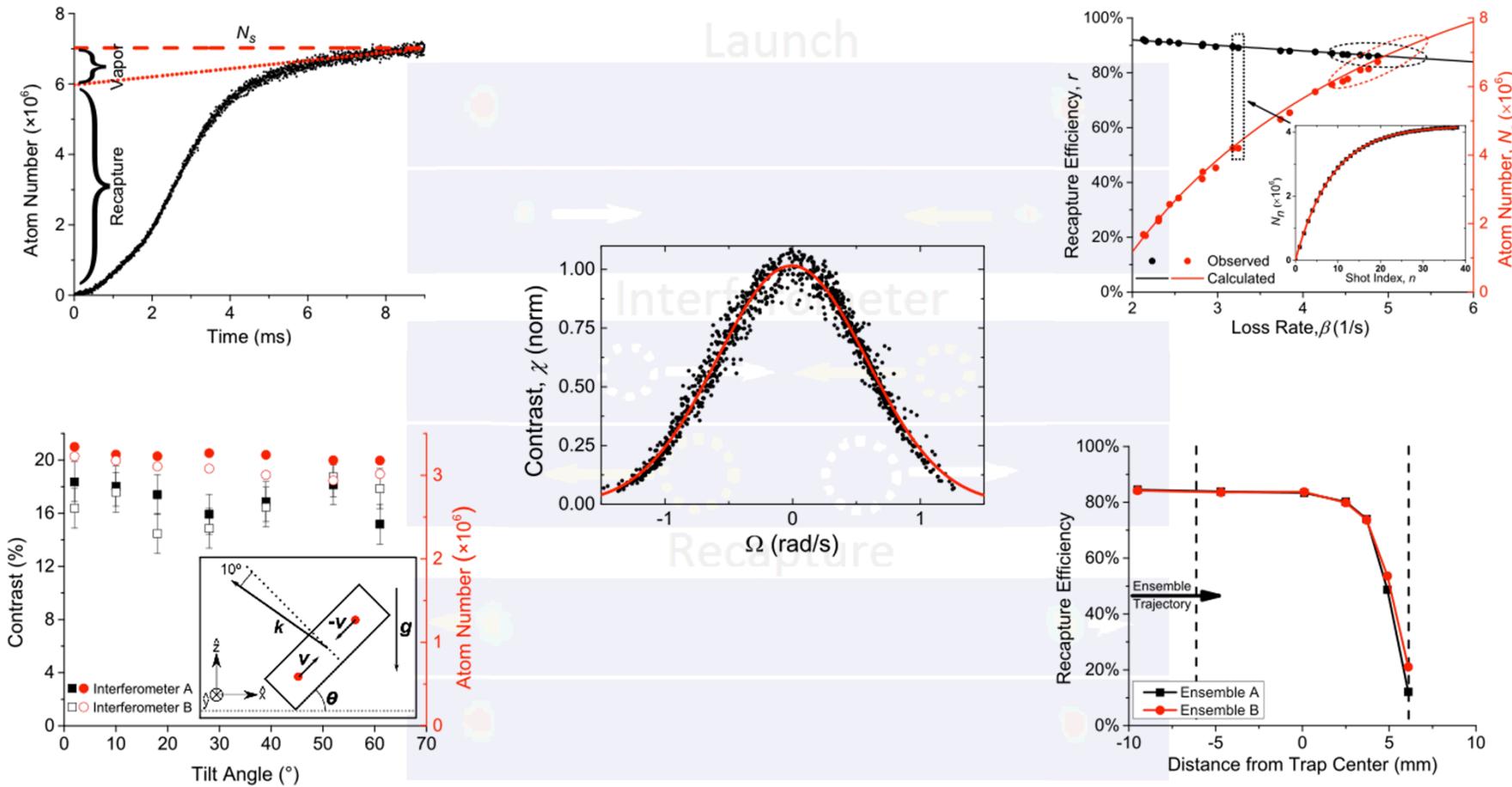


Picture of system hardware



Characterizing ensemble exchange

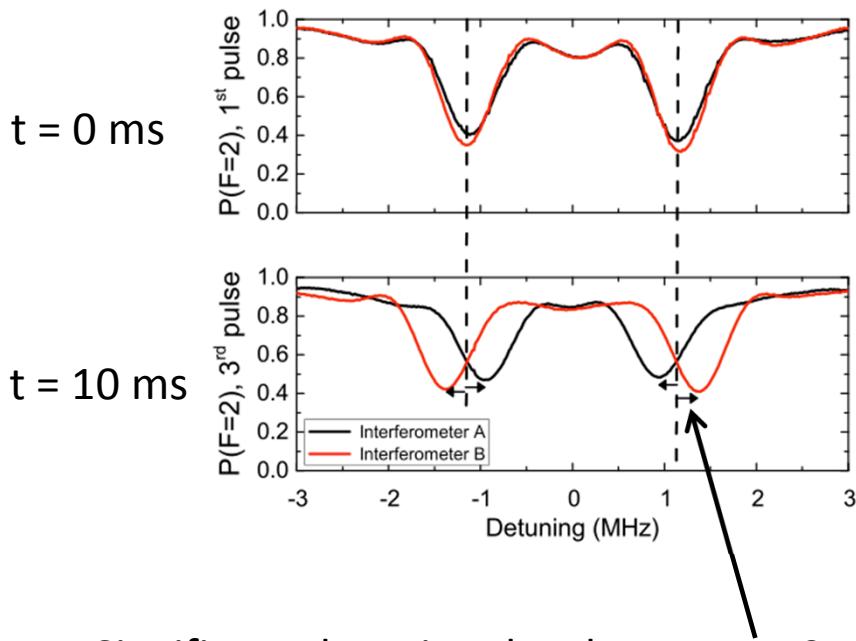
- Dynamic aspects of Ensemble Exchange characterized
- Robust to rotations, tilts and displacements



Controlling jerks

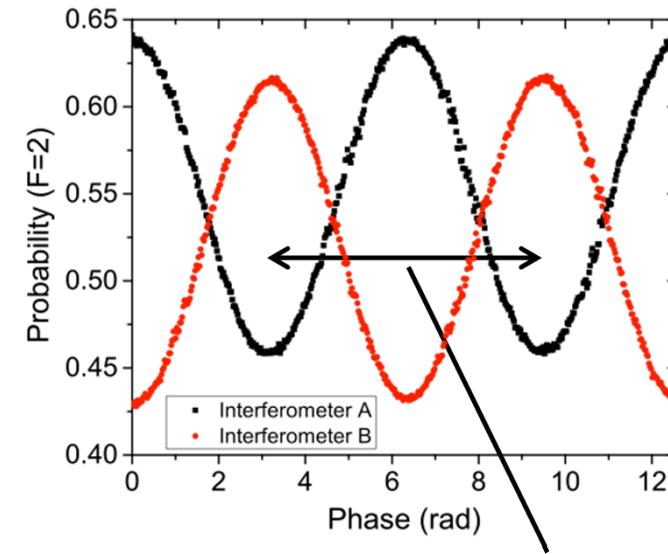
Controlling frequency and phase in a dynamic environment

Raman laser frequency scan



Significant detuning develops over 10 ms due to gravitationally-driven Doppler shift

Raman laser phase scan

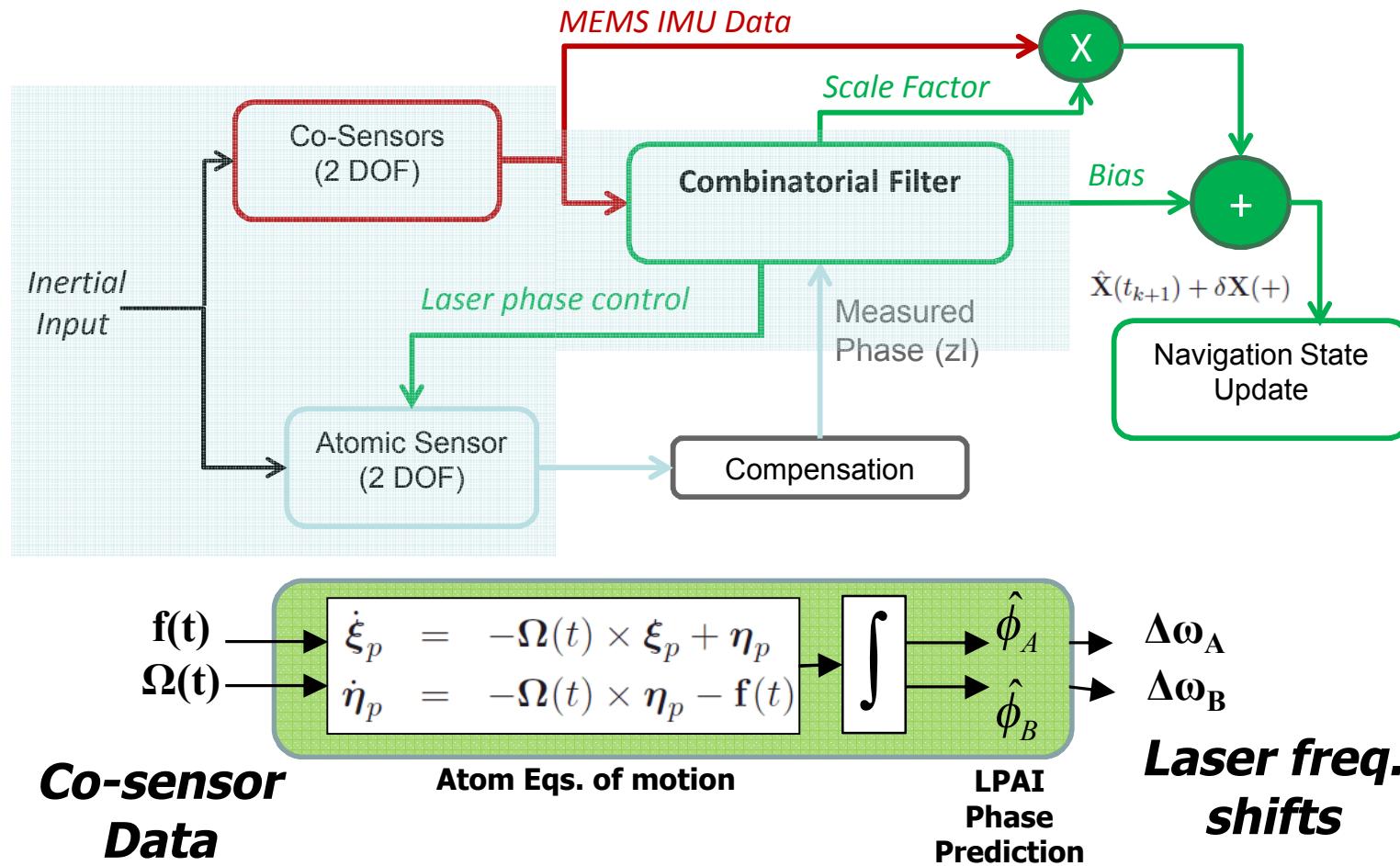


$\approx \text{mg OR } \approx \text{mrad/s (400 deg/hr)}$

Small changes in acceleration or rotation during one interferometer cycle cause phase slip

Technical challenge

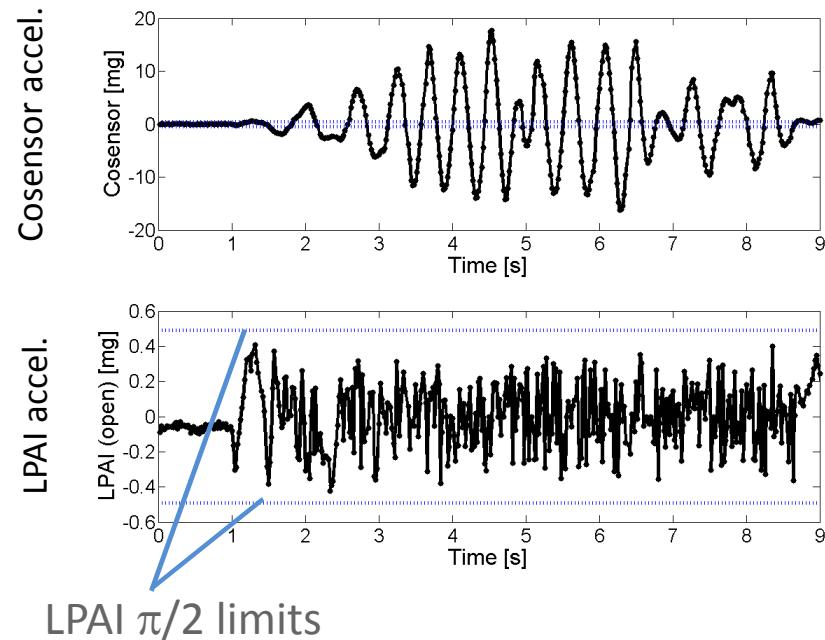
- Real-time cosensor read, Runge-Kutta propagator calculation, and phase/frequency write to interferometer laser in less than 0.5 ms (note: $T = 4.5$ ms)



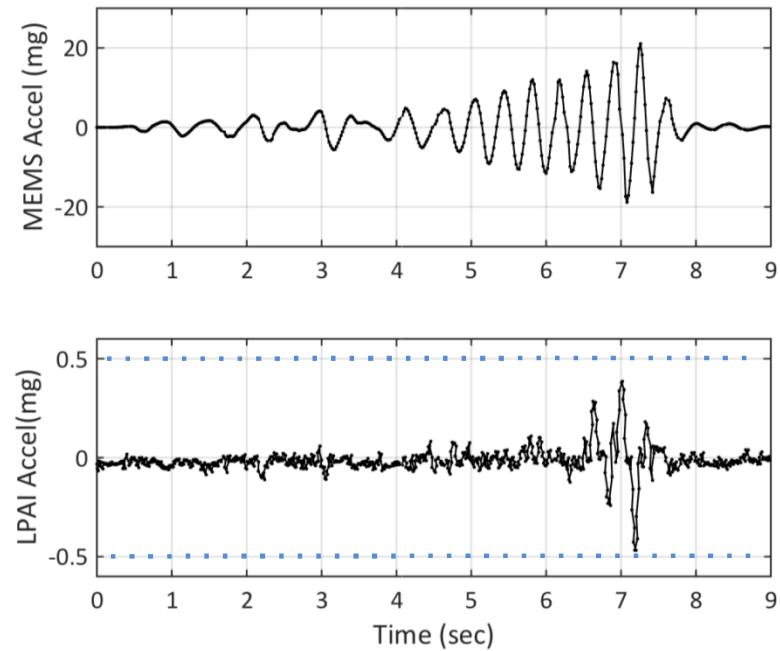
Phase feed-forward results



OPEN LOOP



FEED FORWARD

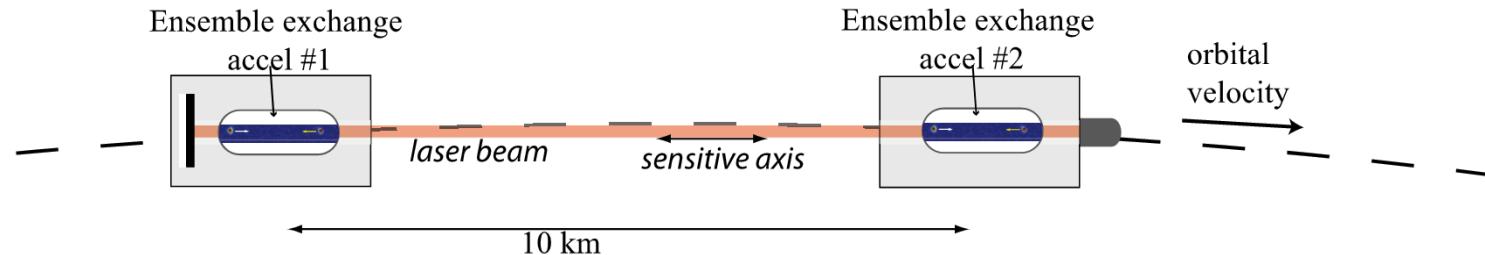


Initial results:

- Accelerometer only
- Feed forward successfully locks LPAI phase to within π radians
- Extends dynamic range to $\approx \pm 10$ mg

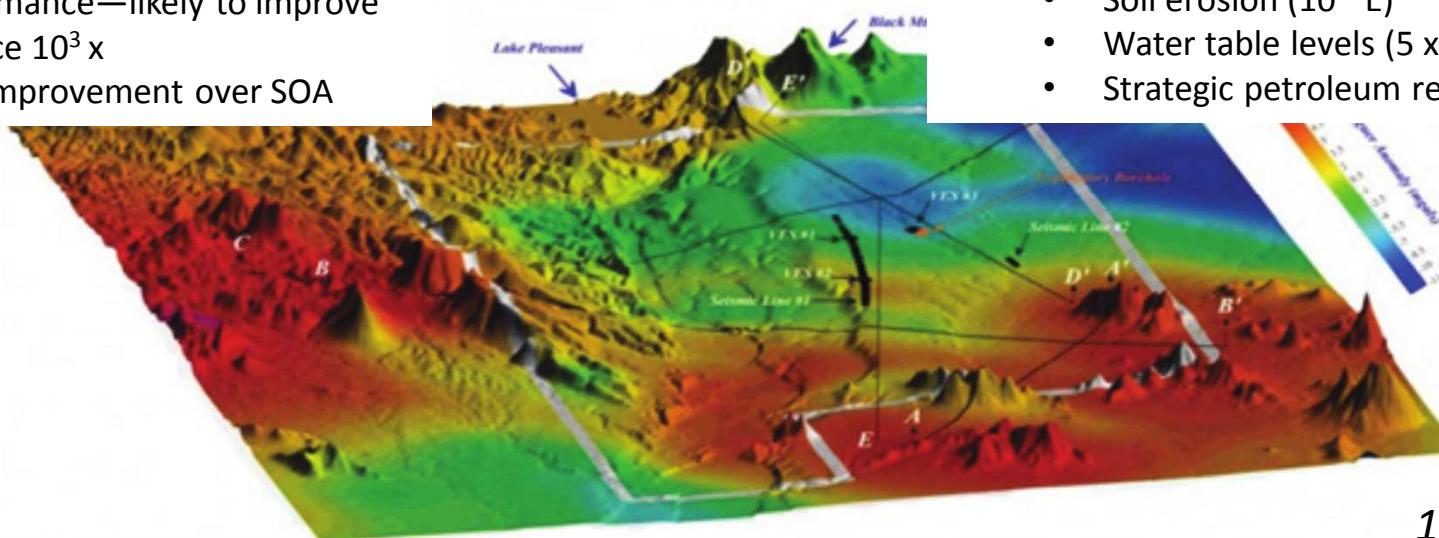
Gradiometer survey—path finder

Simultaneous opposed gradiometers—bias rejection



GOALS

- 10^4 s stability—multiple orbital passes ($SOA < 1$ s)
- 10^{-3} E per shot with ground based performance—likely to improve in space 10^3 x
- 10^4 x improvement over SOA

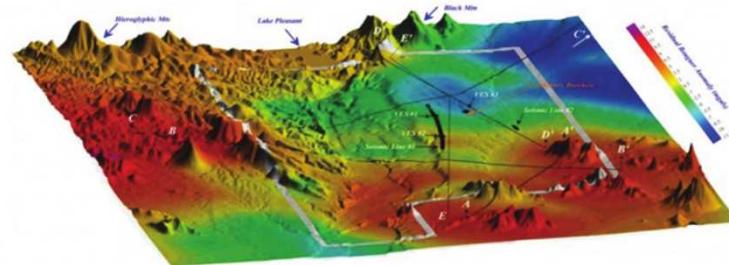
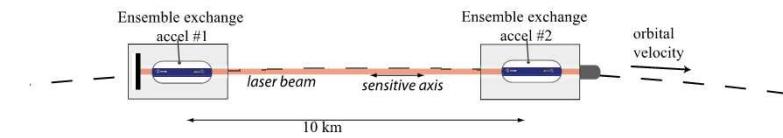
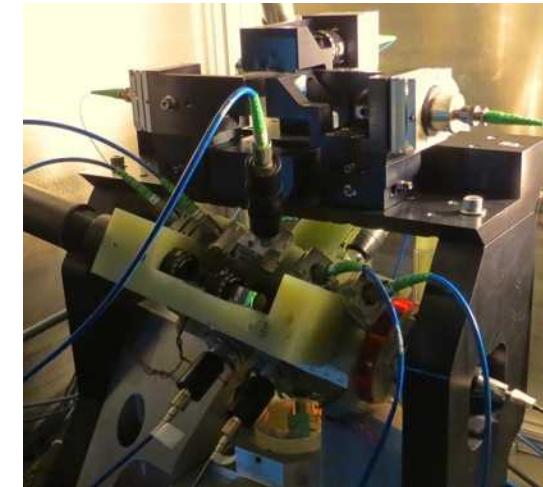


UTILITY

- Improved gravity maps in contested areas for GPS-denied navigation
- Other targets
 - Soil erosion (10^{-3} E)
 - Water table levels (5×10^{-3} E)
 - Strategic petroleum reserves (10^{-6} E)

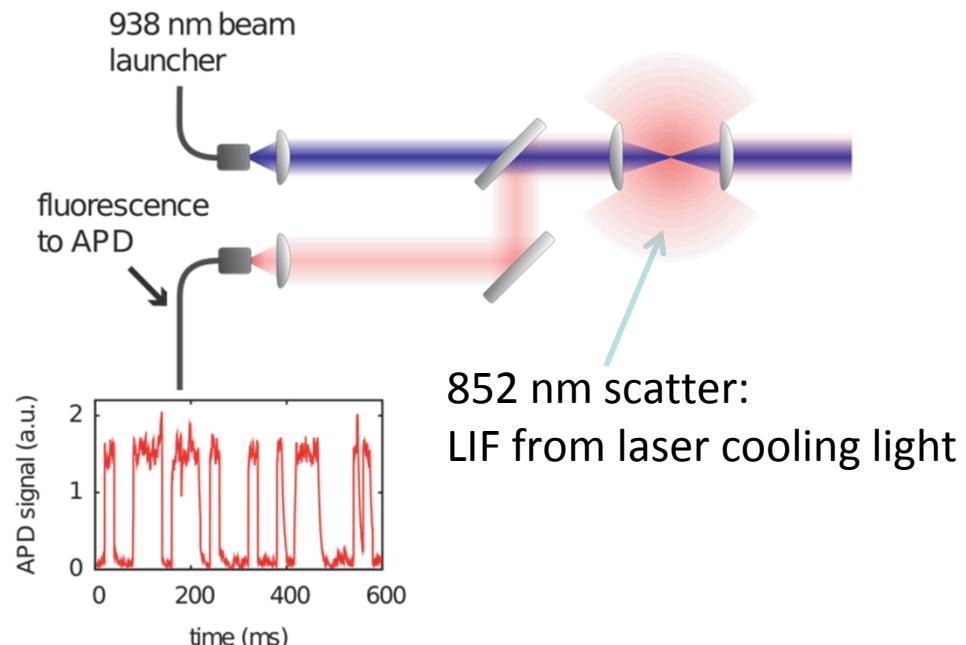
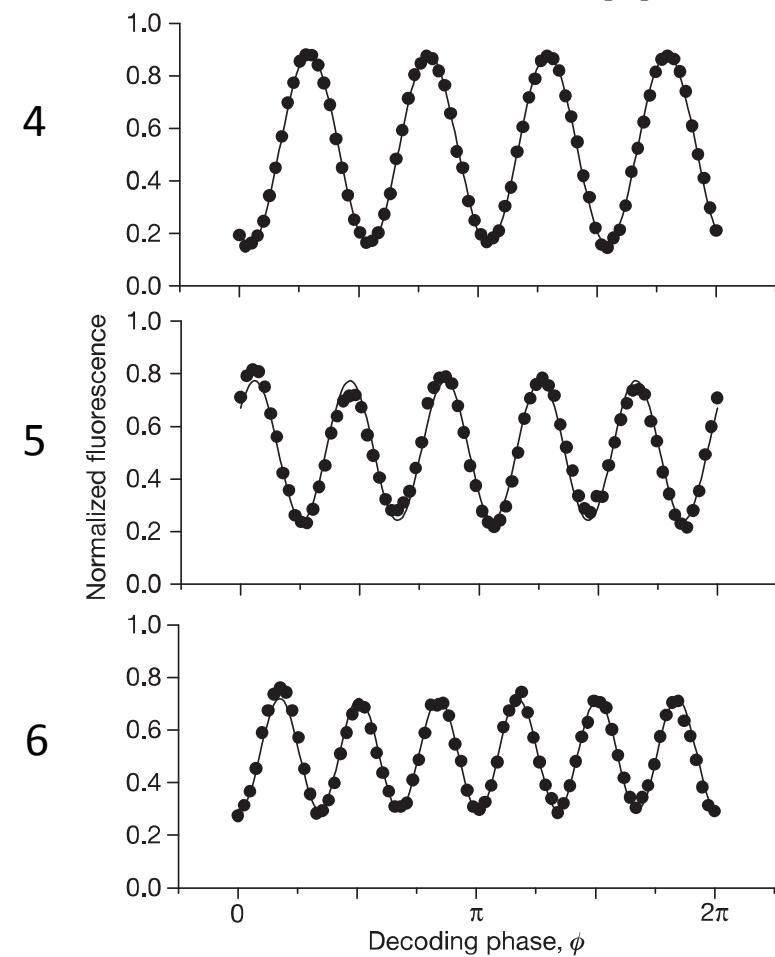
Challenges and Opportunities

- Exquisite precision demonstrated
- Challenging system hardware
- Foreign investments
 - Inertial measurement units
 - Space programs
- Natural technology for a calibrating reference
- Natural technology for large measurements systems
- Natural technology for entanglement-enhanced precision

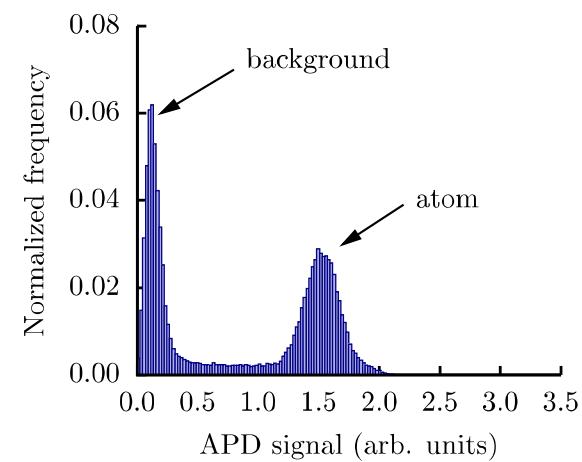


Entangled states for metrology

Cat states with ions [1]

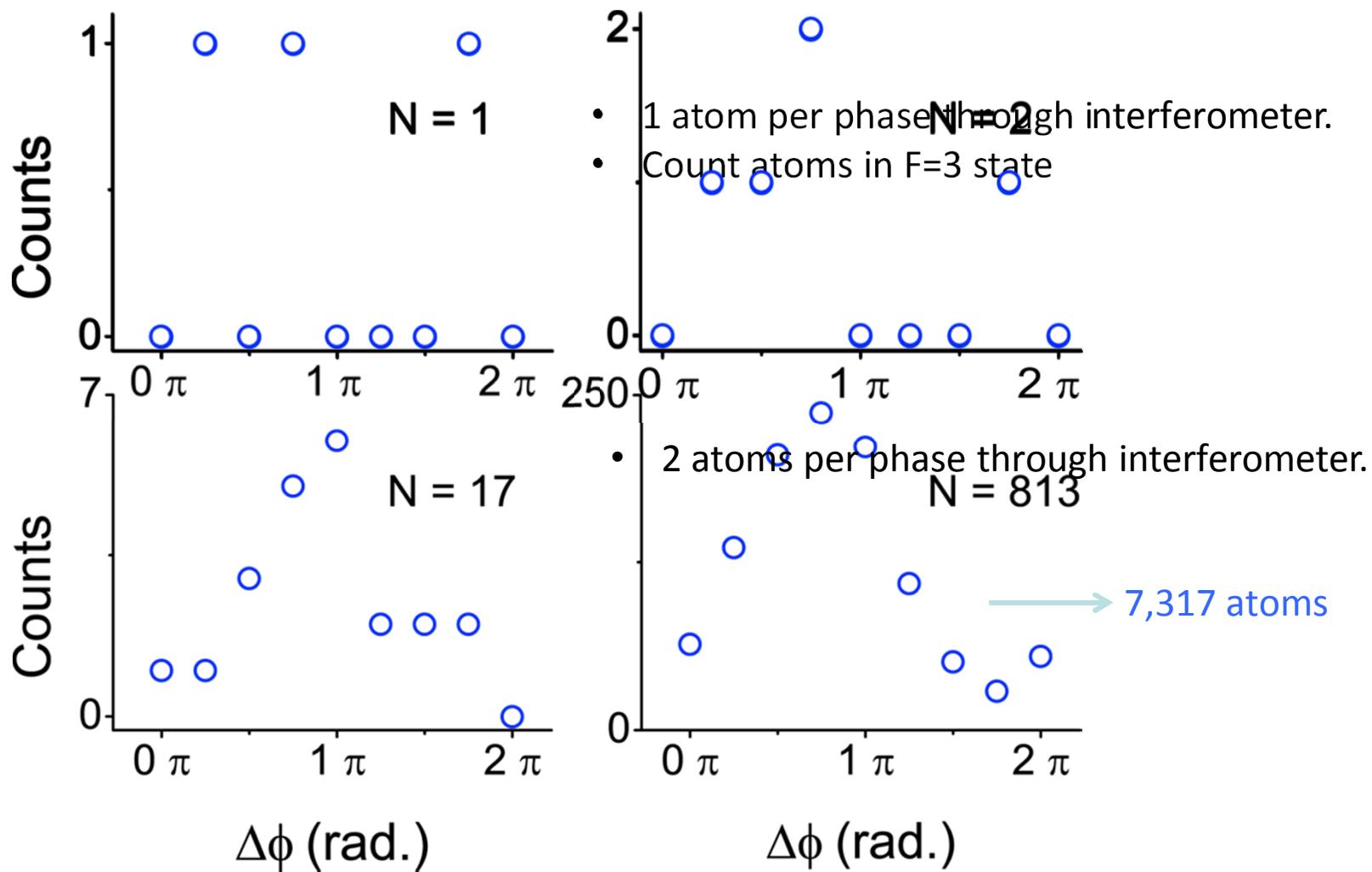


*Micron-scale:
inherently low
power*

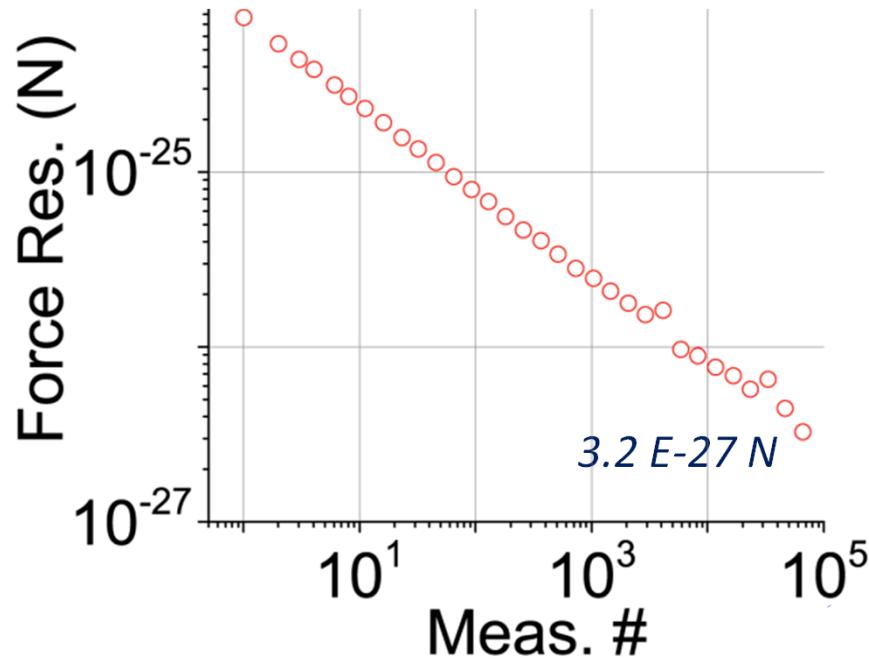


[1] Leibfried, et al., "Creation of a six-atom 'Schrödinger cat' state", *Nature* **438**, 639 (2005)

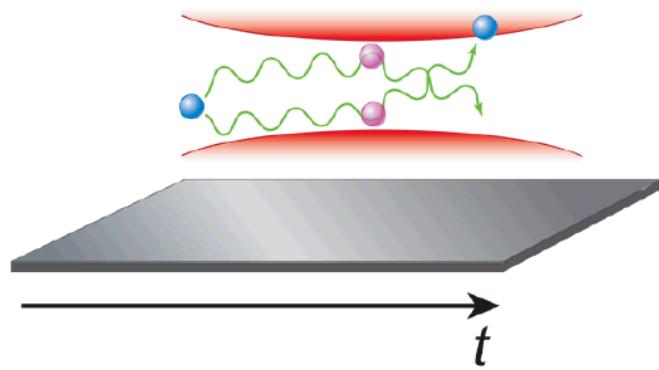
Building a fringe, one atom at a time



Single atom interferometry



Atom interferometer with single atoms

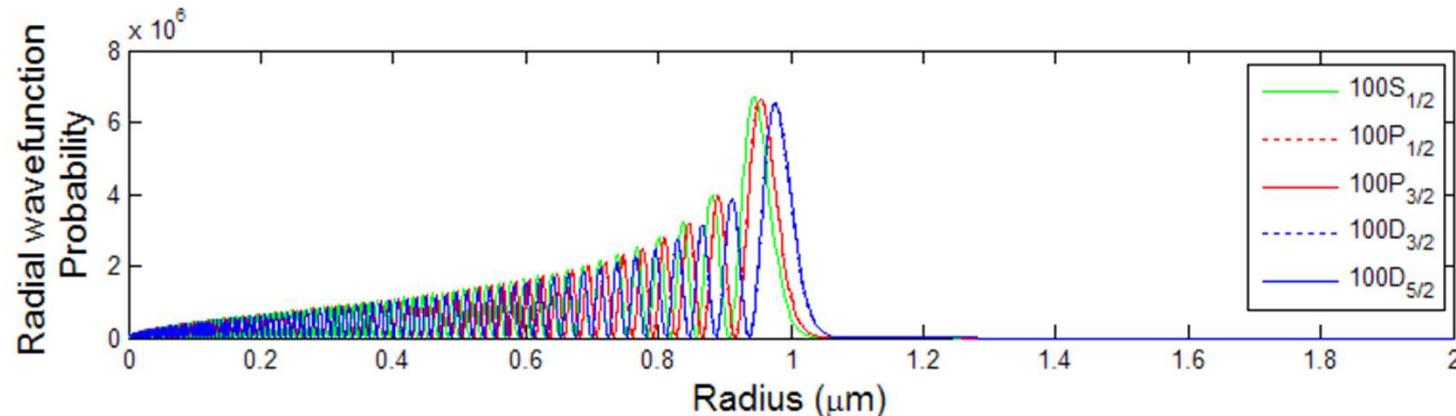


E. Rasel, Physics 5, 135 (2012)

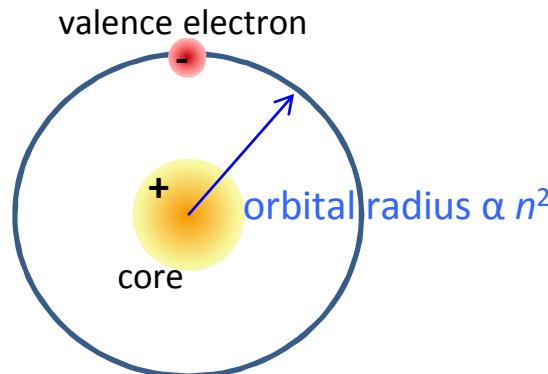
- We showed one can use single atoms
- Single atom control: gateway to harnessing quantum control in sensing
- $10^{-27} \text{ N} \approx mg$ for a cesium atom

Rydberg state mediated interaction

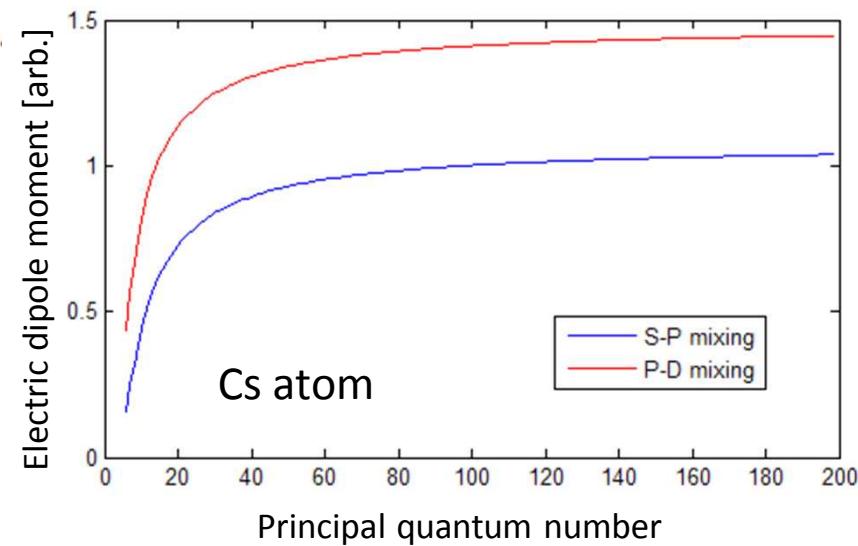
An example of the radial wavefunctions of a Cs atom at $n = 100$:



A Rydberg atom can have a strong electric dipole moment.

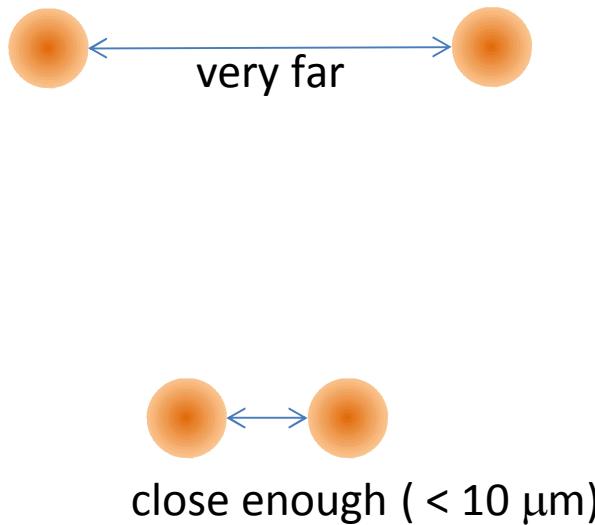


A classical picture of an atom

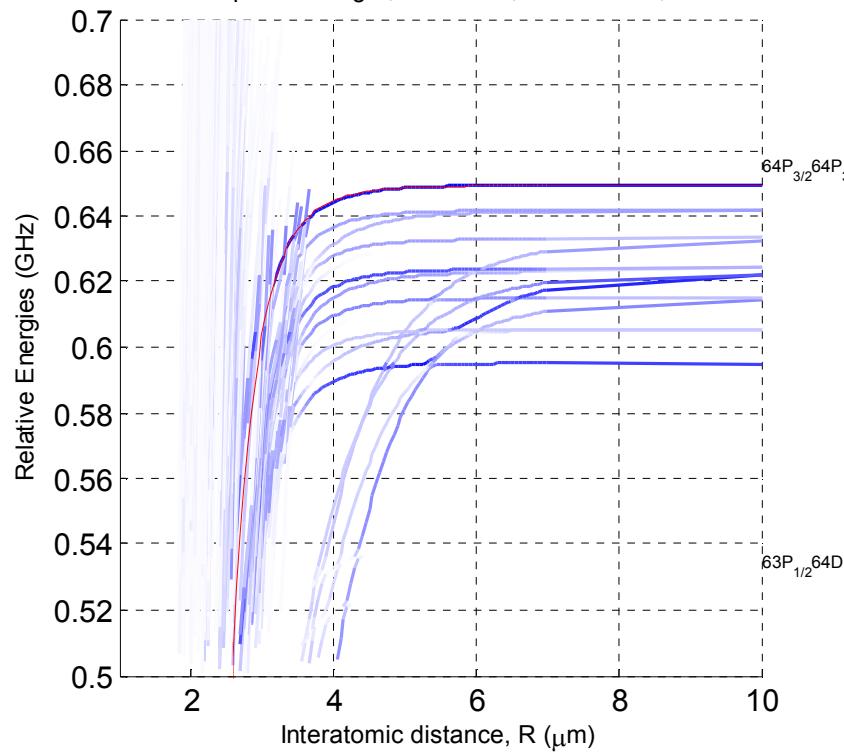


Blockade & electric dipole-dipole interaction

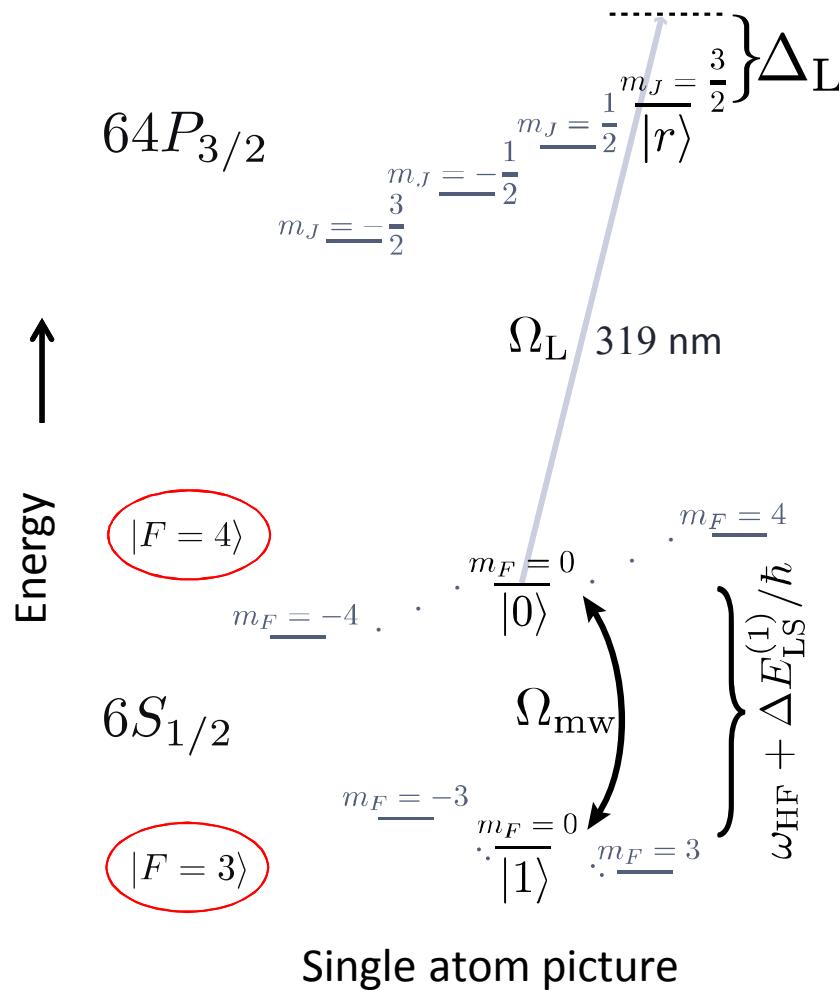
$$H_{\text{atoms}} = \sum_i H_0^{(i)} + \frac{1}{4\pi\epsilon_0 r^3} \sum_{i \neq j} (\mathbf{D}^{(i)} \cdot \mathbf{D}^{(j)} - 3\mathbf{D}^{(i)} \cdot \hat{\mathbf{r}}\hat{\mathbf{r}} \cdot \mathbf{D}^{(j)})$$



Weighted Rydberg Energy levels: Excitation from ground-state to $64\text{P}_{3/2}$
x-polarized light; $B = 4.8$ G; $E = 6.4$ V/m;

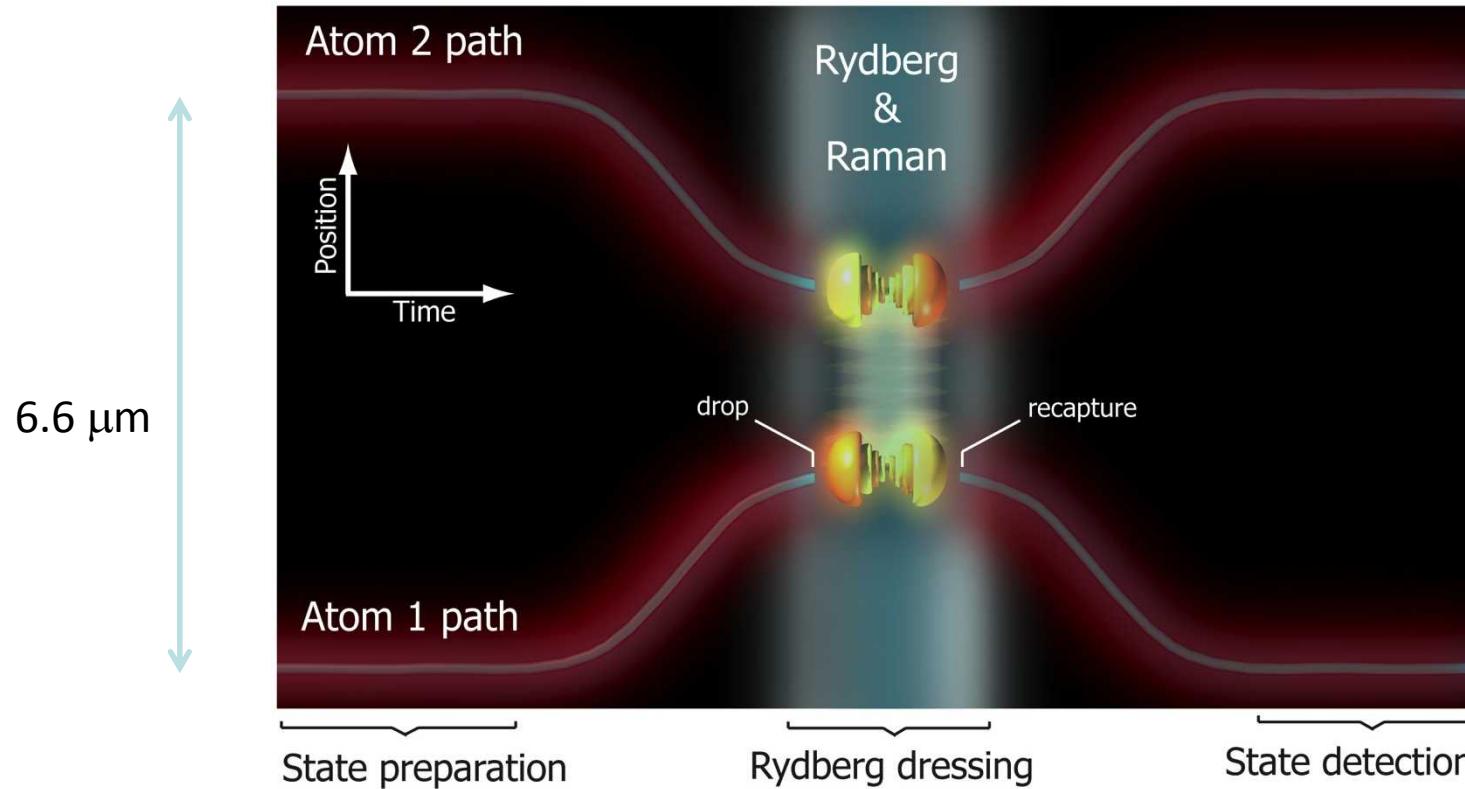


Rydberg-dressed ground state interaction

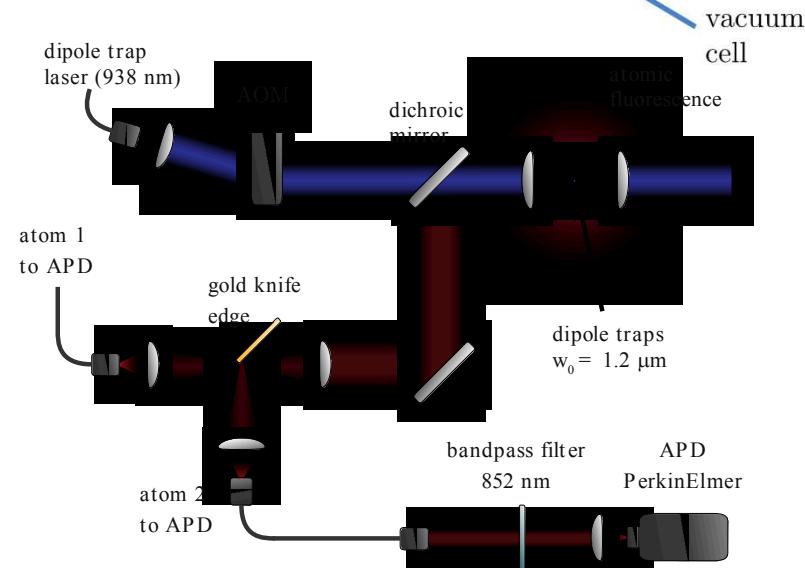
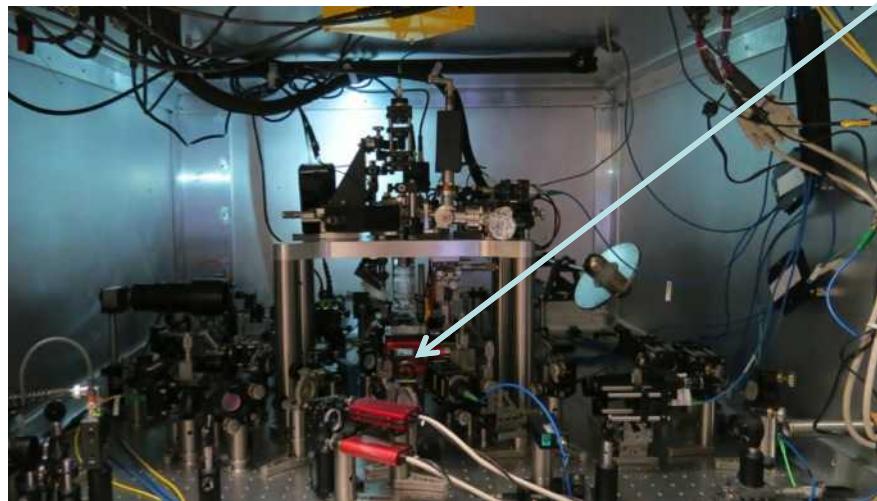
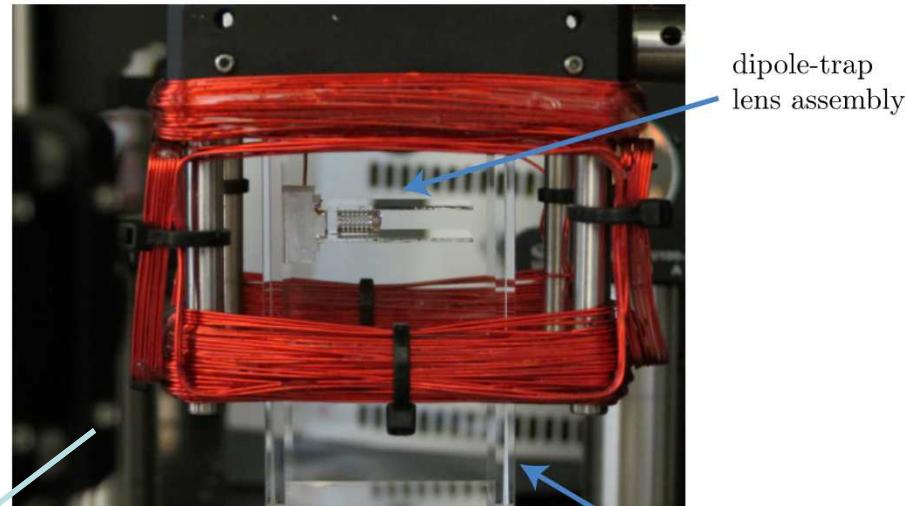
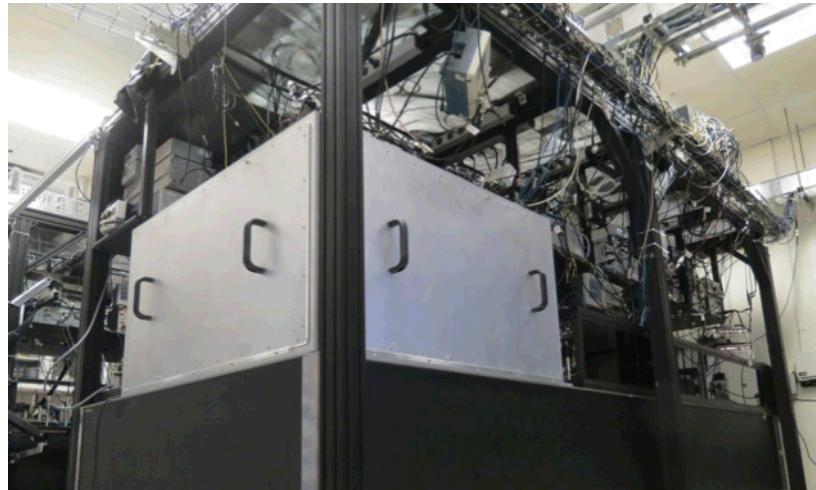


- Interaction range **increases** as principal quantum number n increases
- However, oscillator strength **decreases** as n increases—making Ω_L smaller and thus J
- Target smallest n that your optical resolution can accommodate
- Solution—*dynamic positioning*

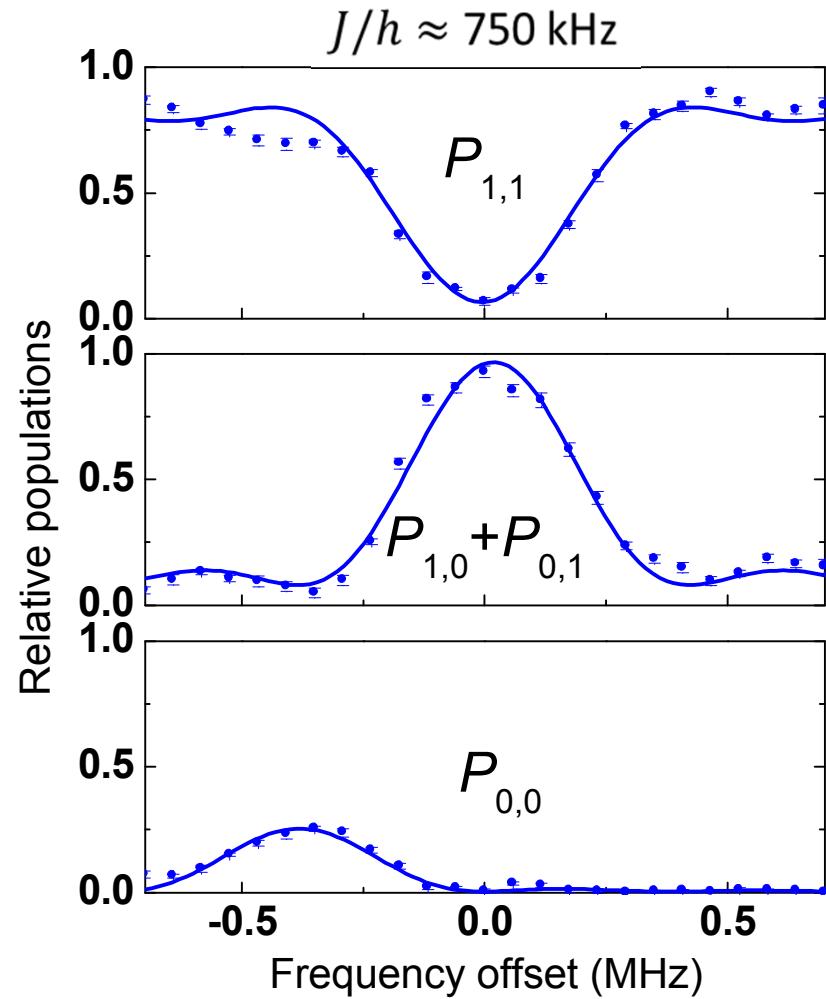
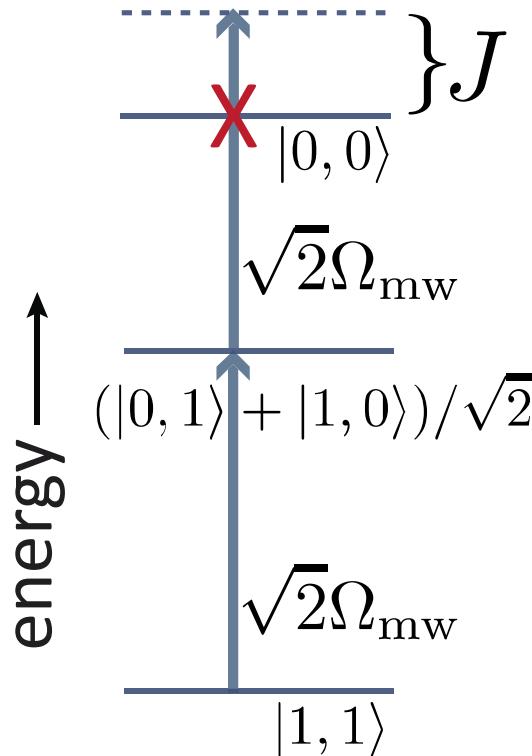
Dynamic atom positioning



Apparatus

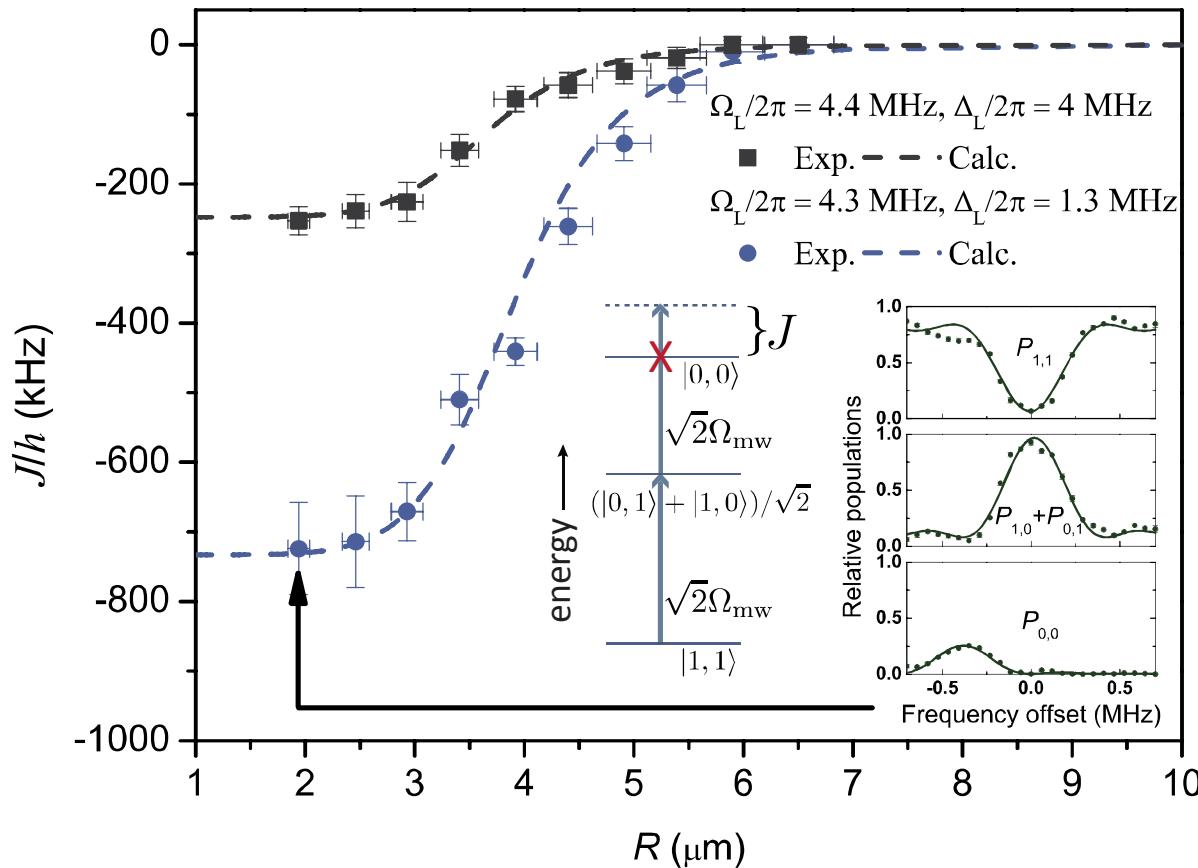


Two-qubit microwave resonances



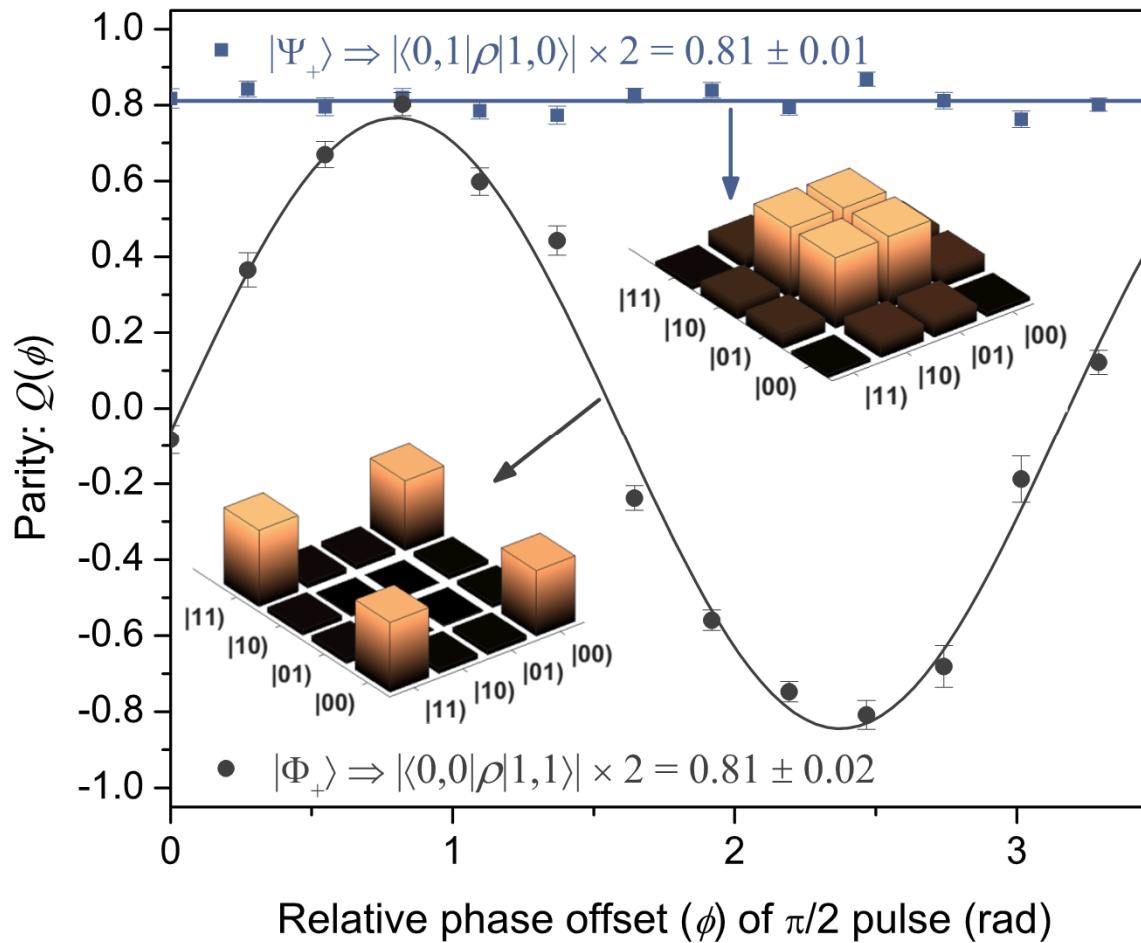
J vs. R , no longer elusive

Direct measurement of two-qubit interaction strength J as a function of two-atom separation with two conditions.



Entanglement Fidelity $\geq 81\%$

Verify the entanglement via parity measurements



Prepare two Cs atoms in Bell state $|\Psi_+\rangle$ or $|\Phi_+\rangle$

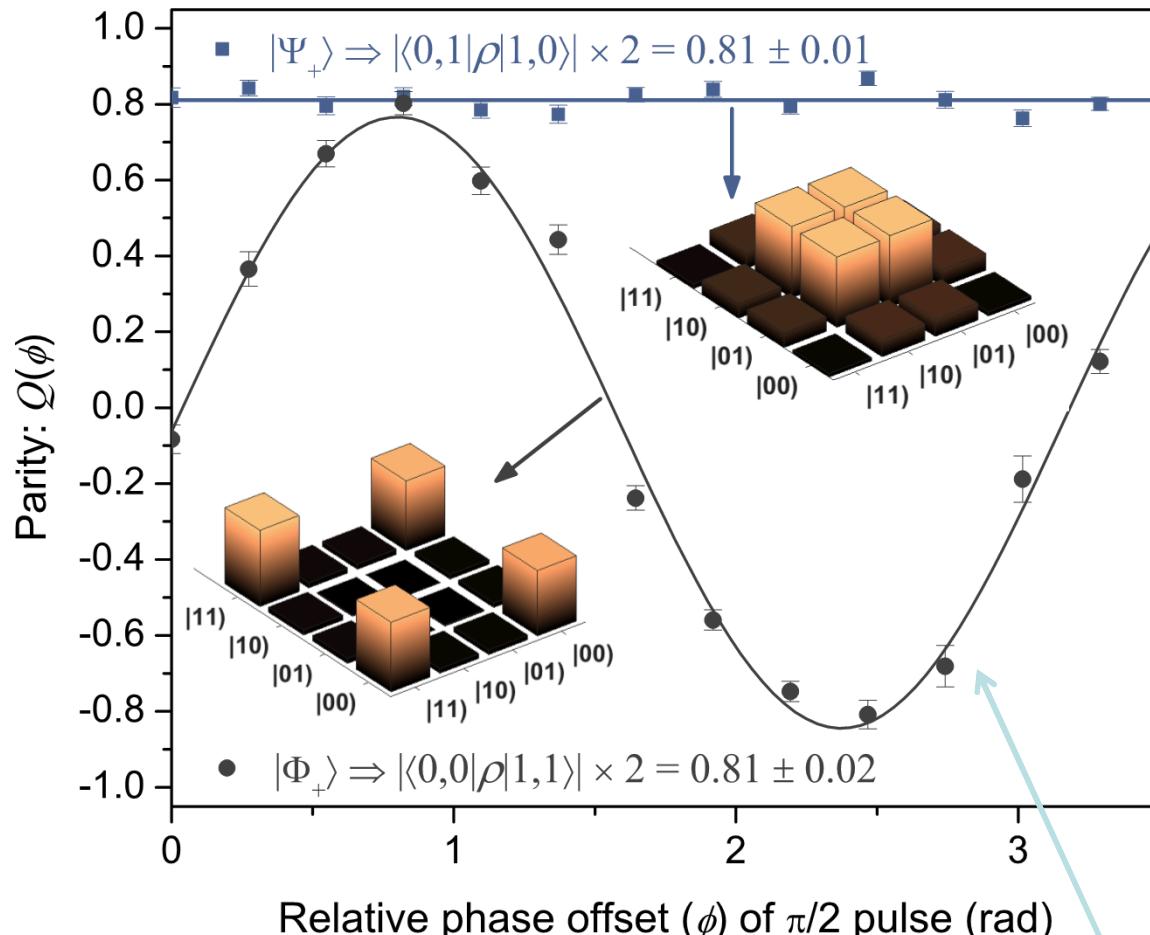
Apply a global $\pi/2$ rotation with a given phase

Perform parity measurement

$$Q = P_{11} + P_{00} - (P_{01} + P_{10})$$

Obtain the two-qubit entanglement fidelity F , where $Q \leq F \leq 1$.

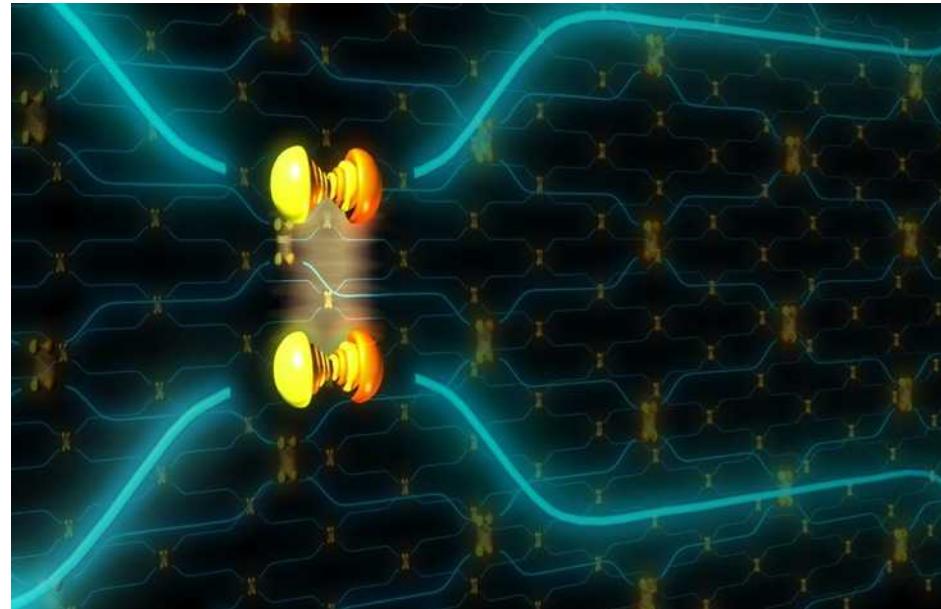
Application to metrology



Cat state 2x response to phase

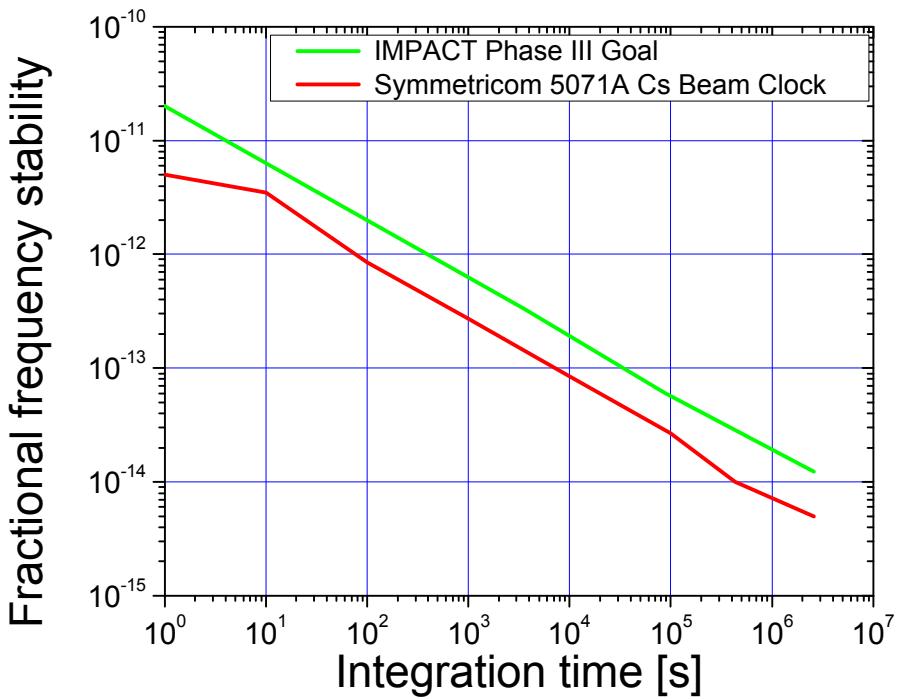
Challenges and Opportunities

- System hardware naturally low power and compact
- Rydberg transition wavelength laser diode challenging
- Potential applications
 - Quantum computation/simulation
 - Quantum-enhanced metrology
- Low TRL



IMPACT Project Goals

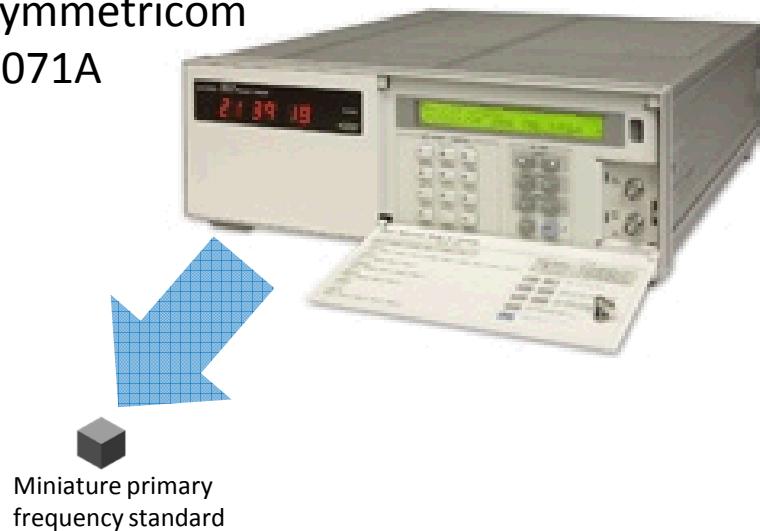
- Achieve Cs Beam Clock performance in a mass and power constrained package
- 5 cm³, 50 mW, 10⁻¹⁴ performance



Applications--Excellent timing for:

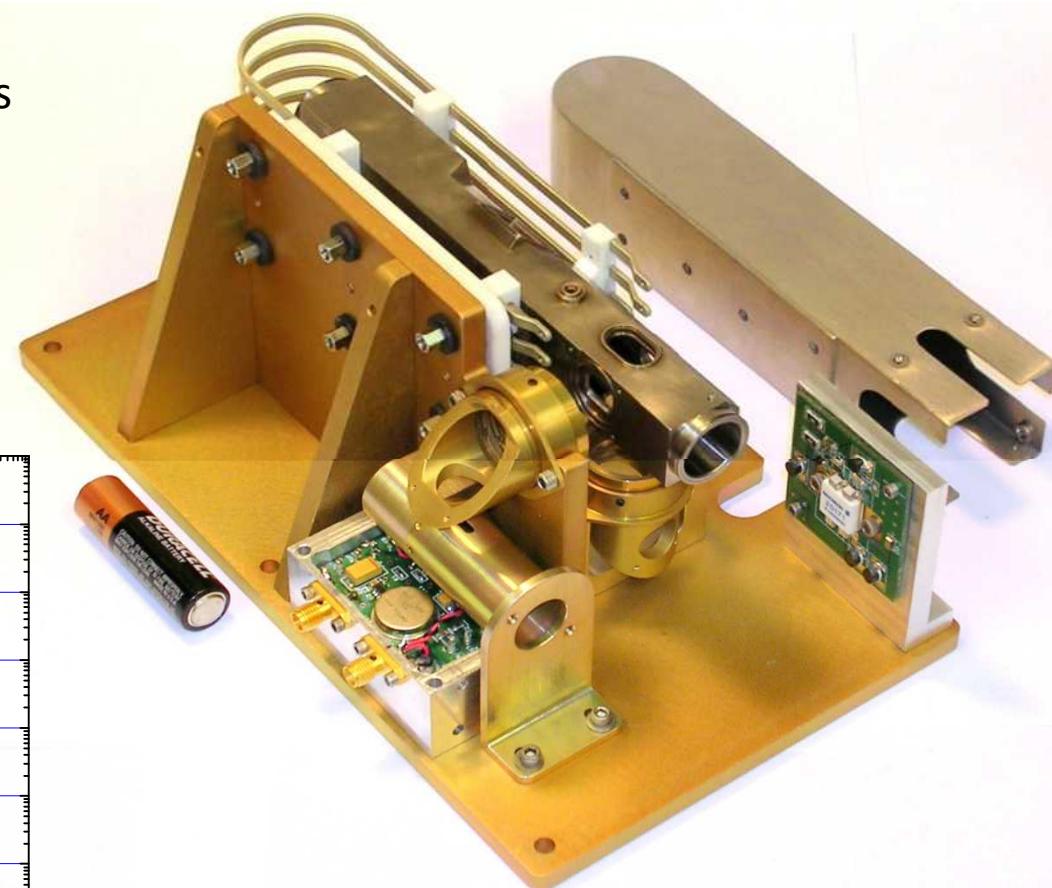
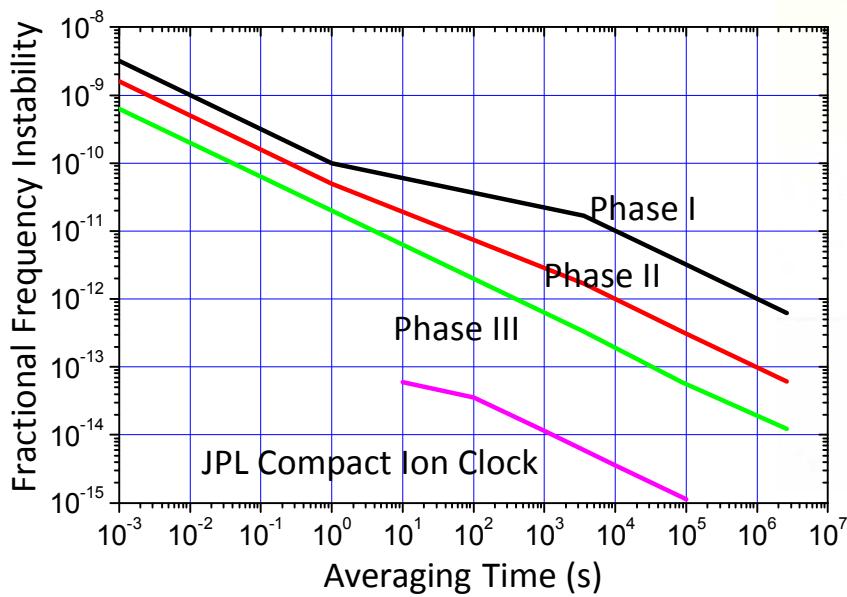
- Rapid GPS acquisition, and GPS denied navigation and timing
- Nano/pico (cube) satellites
- Pulsed radio and spread spectrum communications

Symmetricom
5071A



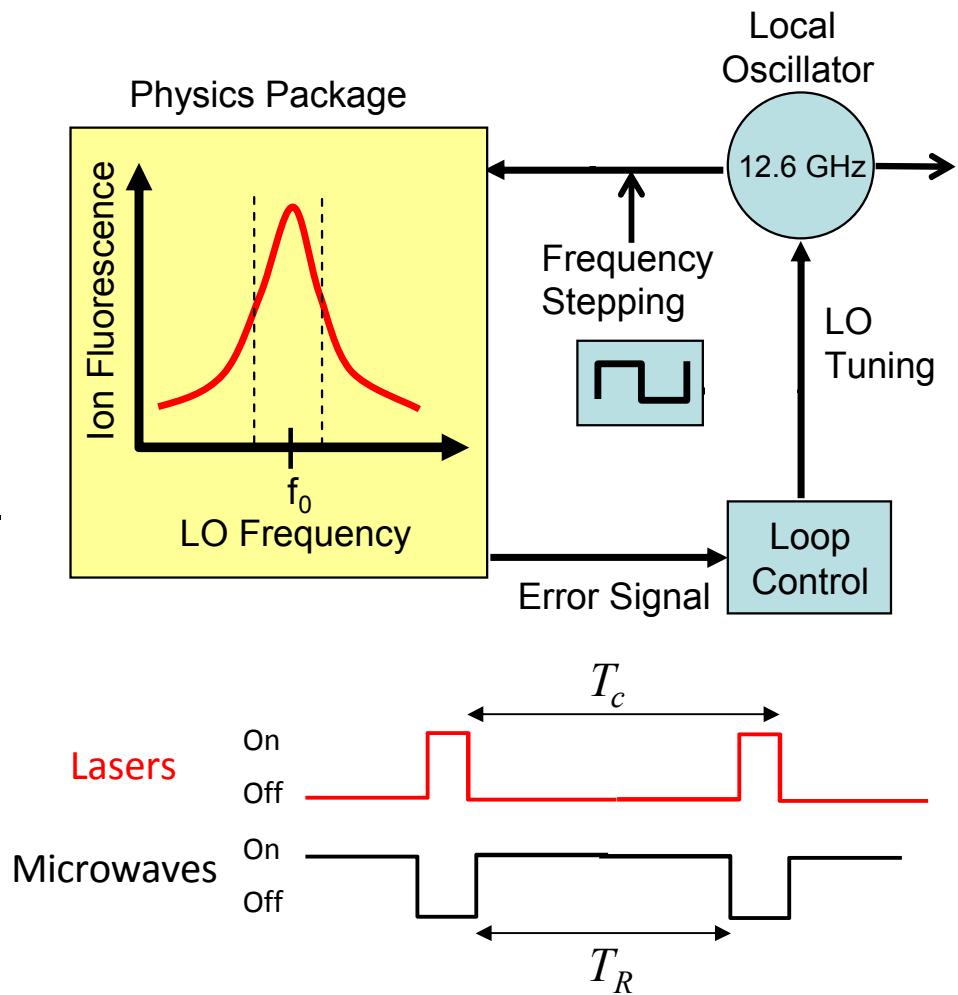
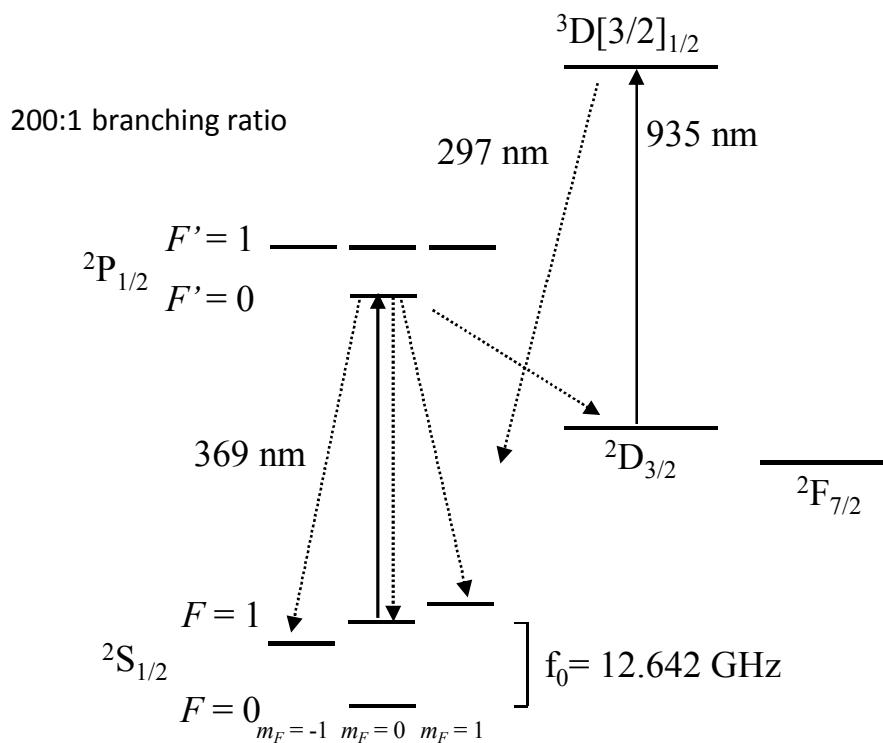
Trapped Ion Clock for Miniaturization

- Trapped ion clocks are already compact while delivering excellent performance.
- Low mass, size, power
- Trapped ion lifetime: up to 10,000 hrs
- Coherence time: > 100s
- Other approaches for IMPACT:
 - Miniature fountain clock
 - Locked optical frequency converted microwaves with WGM resonator



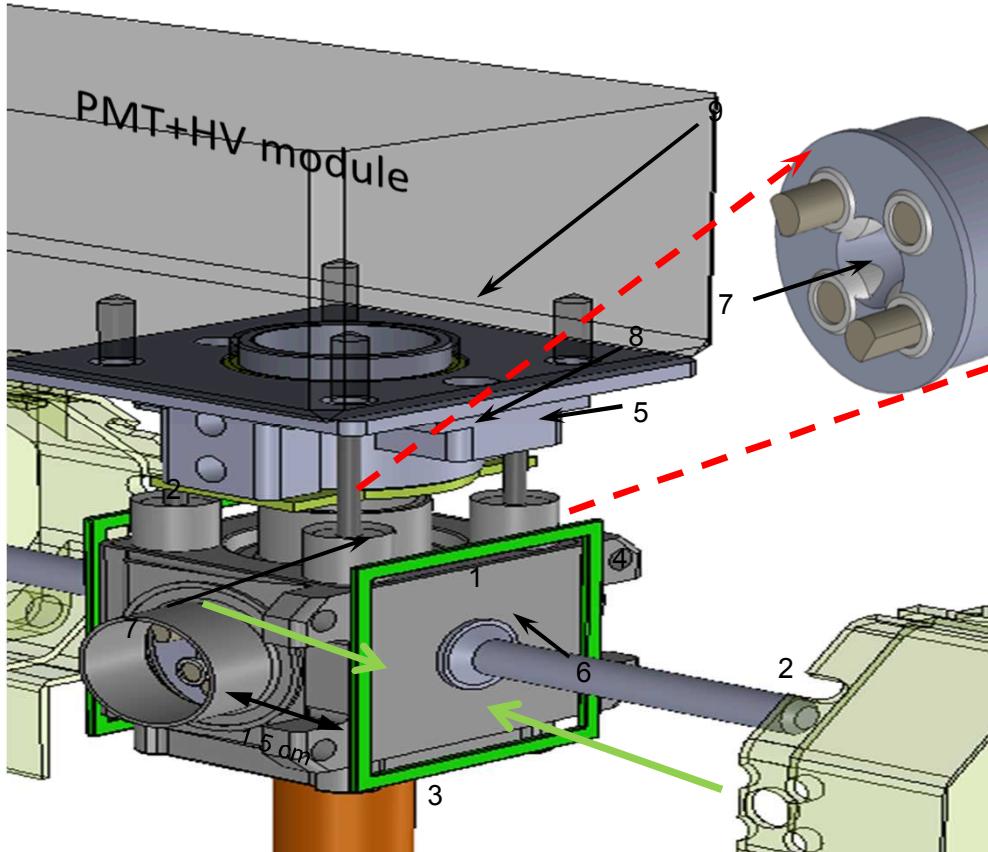
^{199}Hg Trapped Ion Clock from JPL

Atomic Frequency Reference with $^{171}\text{Yb}^+$

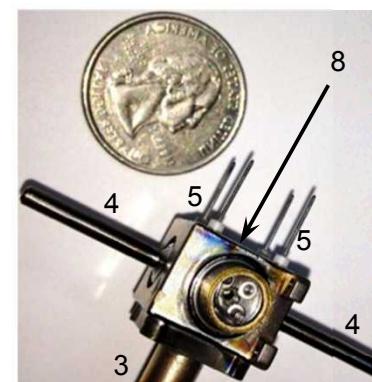
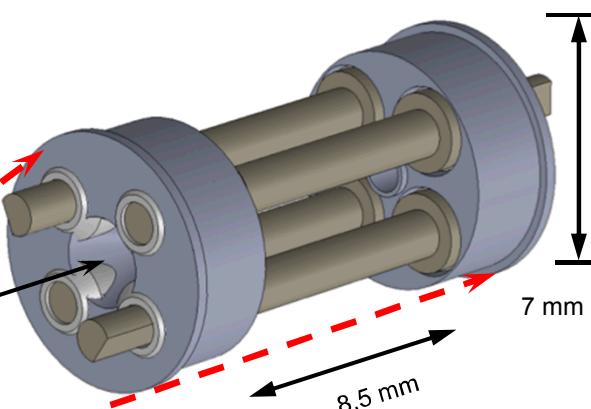


3 cm³ Vacuum Package and Ion Trap

Vacuum package w/ Detector



Ion-trap electrodes



1. Vacuum package
2. μ -metal shield
3. Copper pump-out tube
4. Yb oven appendage
5. Electrical feedthroughs
6. C-field coils
7. Laser port (sapphire)
8. Fluorescence collection window (sapphire)
9. Lens and filters tube

- Titanium body with sapphire windows.
- Linear Quadrupole RF Paul Trap
- Pinched off since April 25th, 2012

- Buffer gas cooling with He
- Getter Pumped.
- Trapped ion lifetime > 3 weeks.

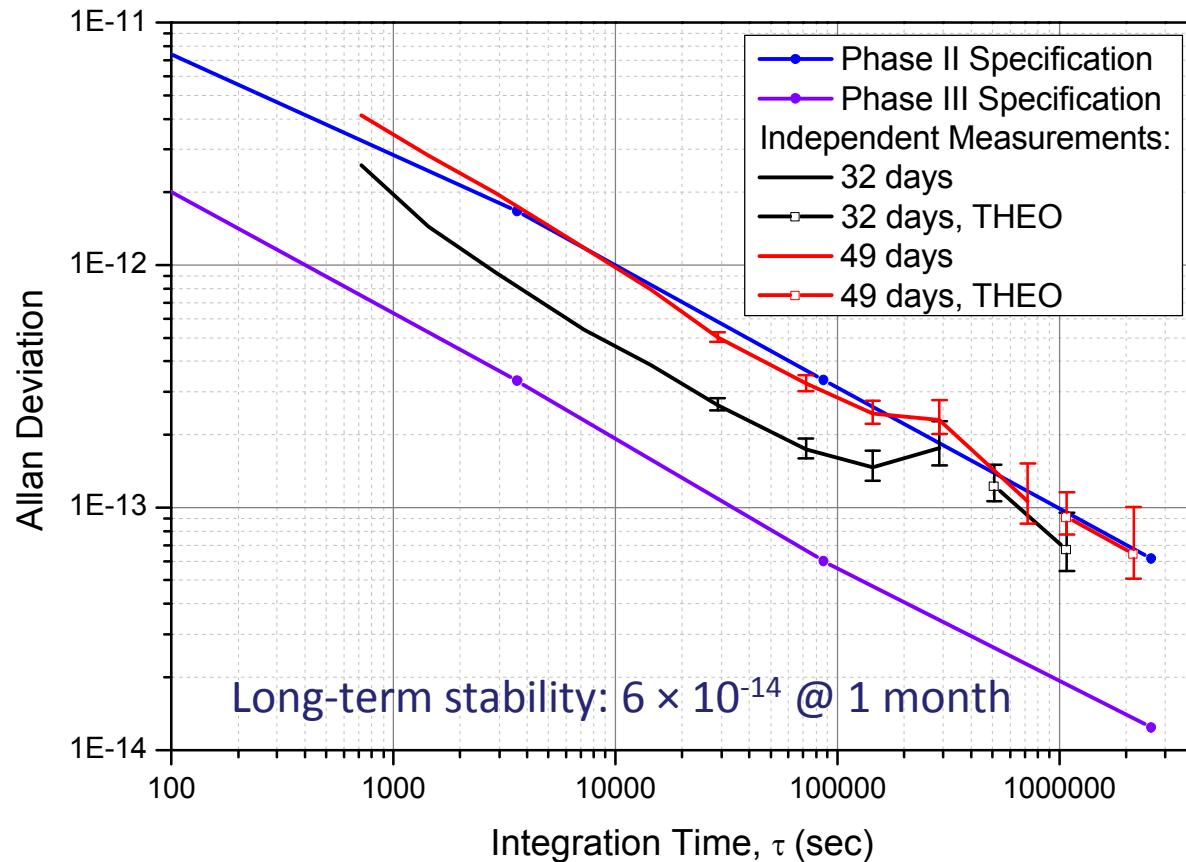
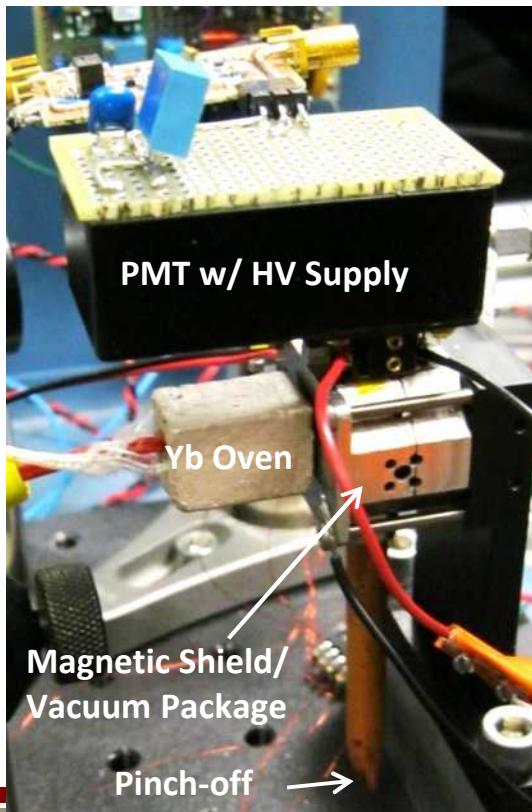
Independent Government Testing of Long-Term Stability

- Bread board clock with 935 nm VCSEL and MEMS shutter
- Large doubled 369 nm laser

- Blu-ray burner diode at 405 nm to create ions.

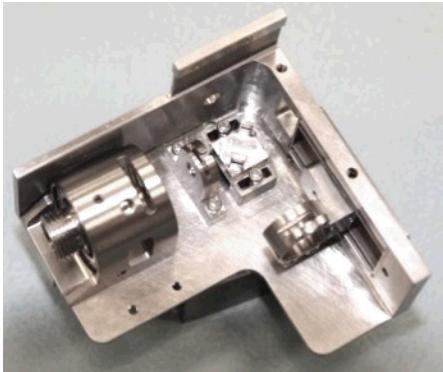
- Allan deviation derived from data sets of 6 days, 26 days, 3 days, 4 days, and 10 days.
- 49 days of data

Integrated Vacuum Package



Complete Physics Package

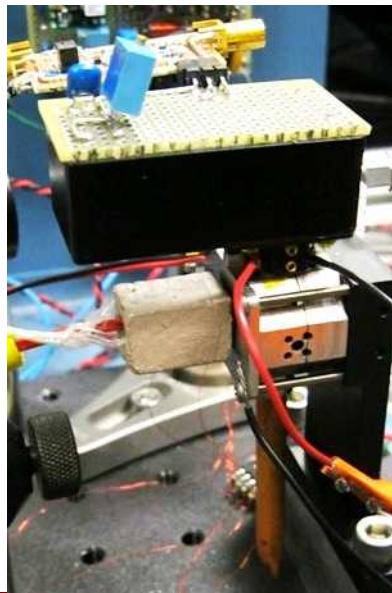
Integrated Optics



Physics Package Interface



Vacuum Package



MEMS
Shutter
Optics &
Lasers

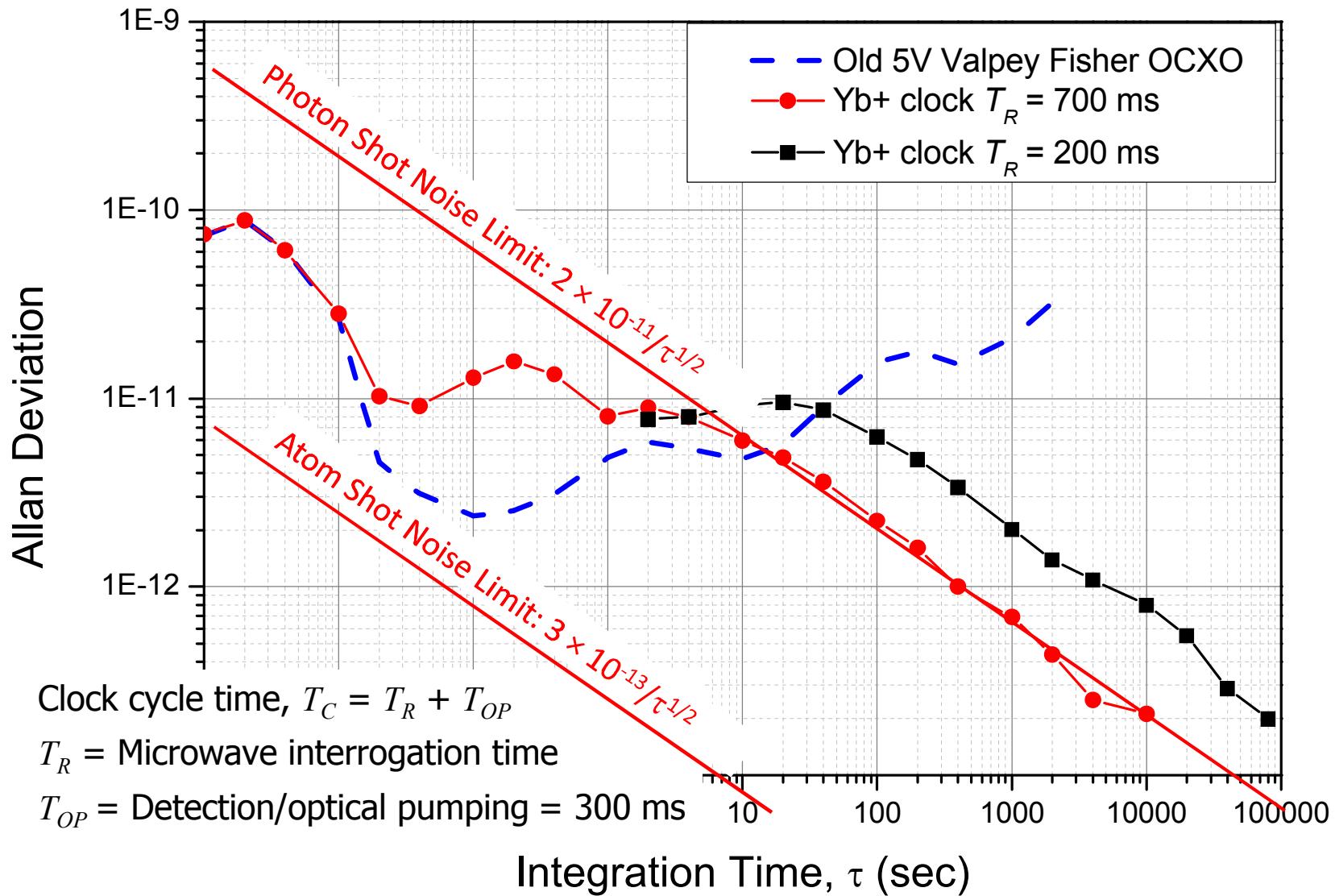
935 nm
VCSEL

Signal Detection
PMT

Physics Package
Interface PCB

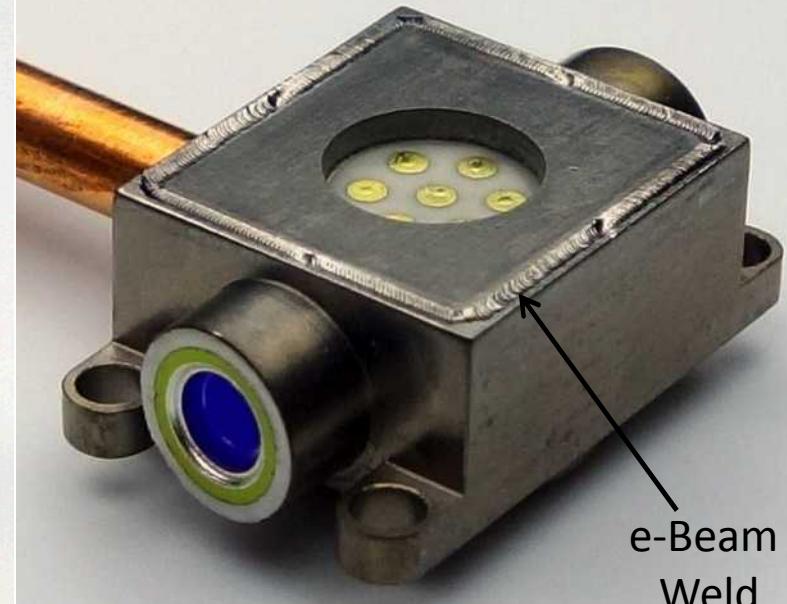
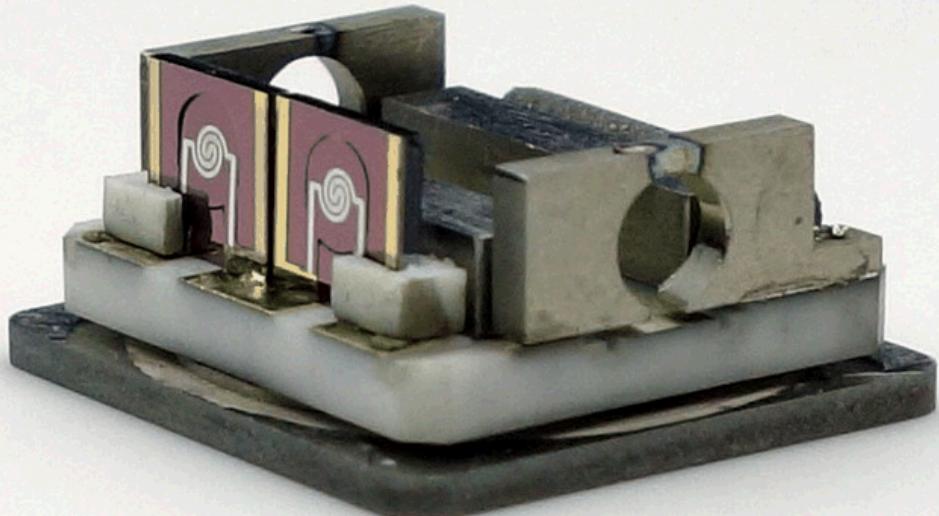
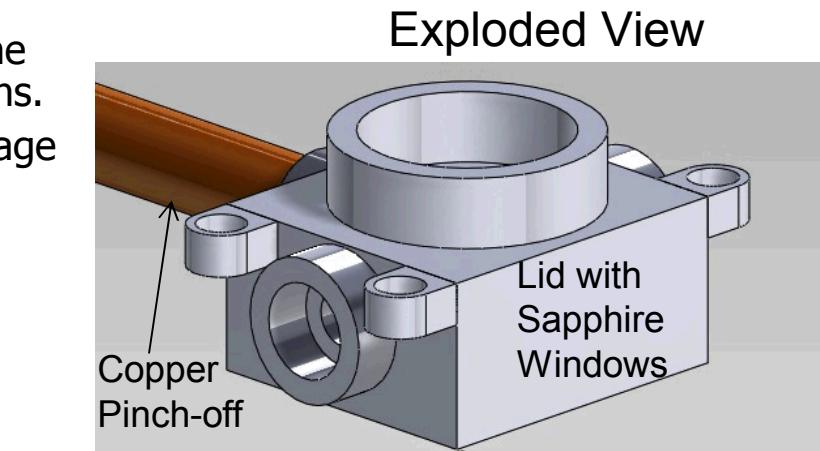
Vacuum Package

Integrated Clock Performance



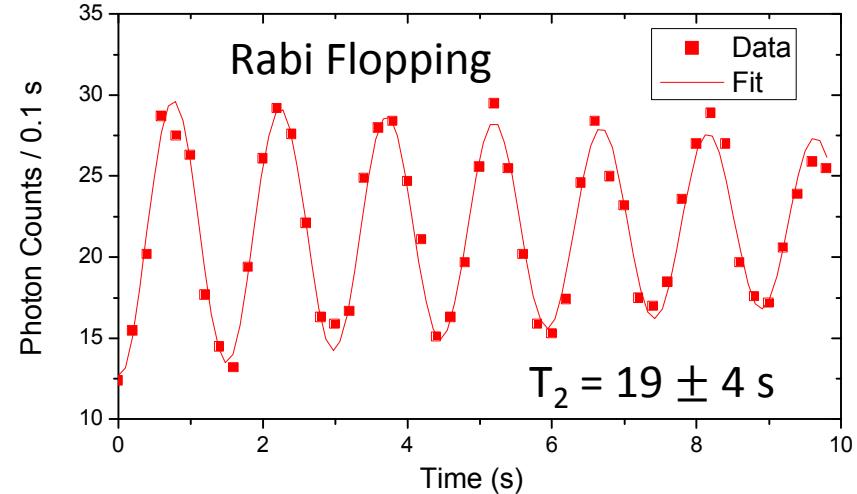
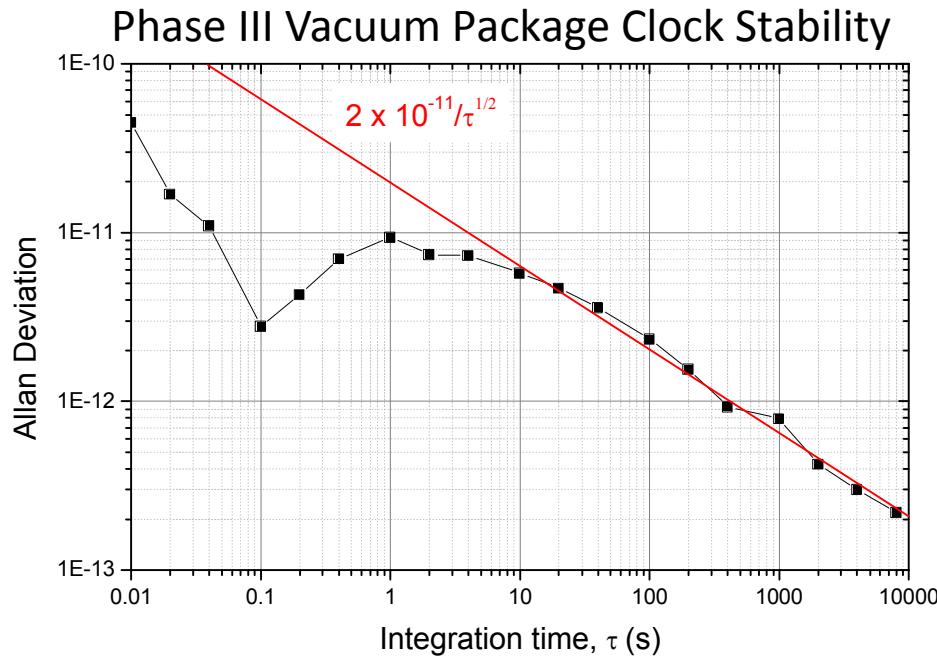
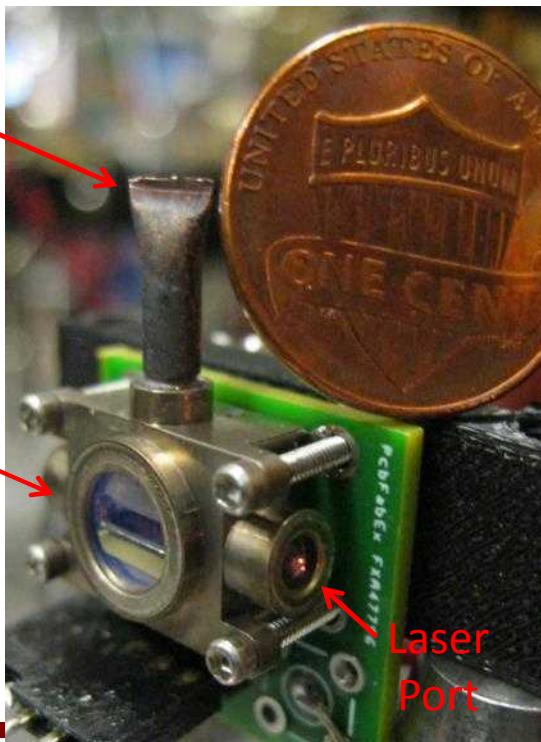
Co-Fired Hybrid Ceramic/Titanium Package

- High temperature co-fired ceramic (HTCC) simplifies the electrical vacuum feedthroughs and internal connections.
- Make the ion trap an integral part of the vacuum package for maximum miniaturization.
- Integrated Yb sources: Silicon micro hotplate
- CuAg braze the Ti parts to the HTCC
- AuGe solder Si hotplates to HTCC
- E-beam weld the base to the lid
- Size of package cube: 0.8 cm³



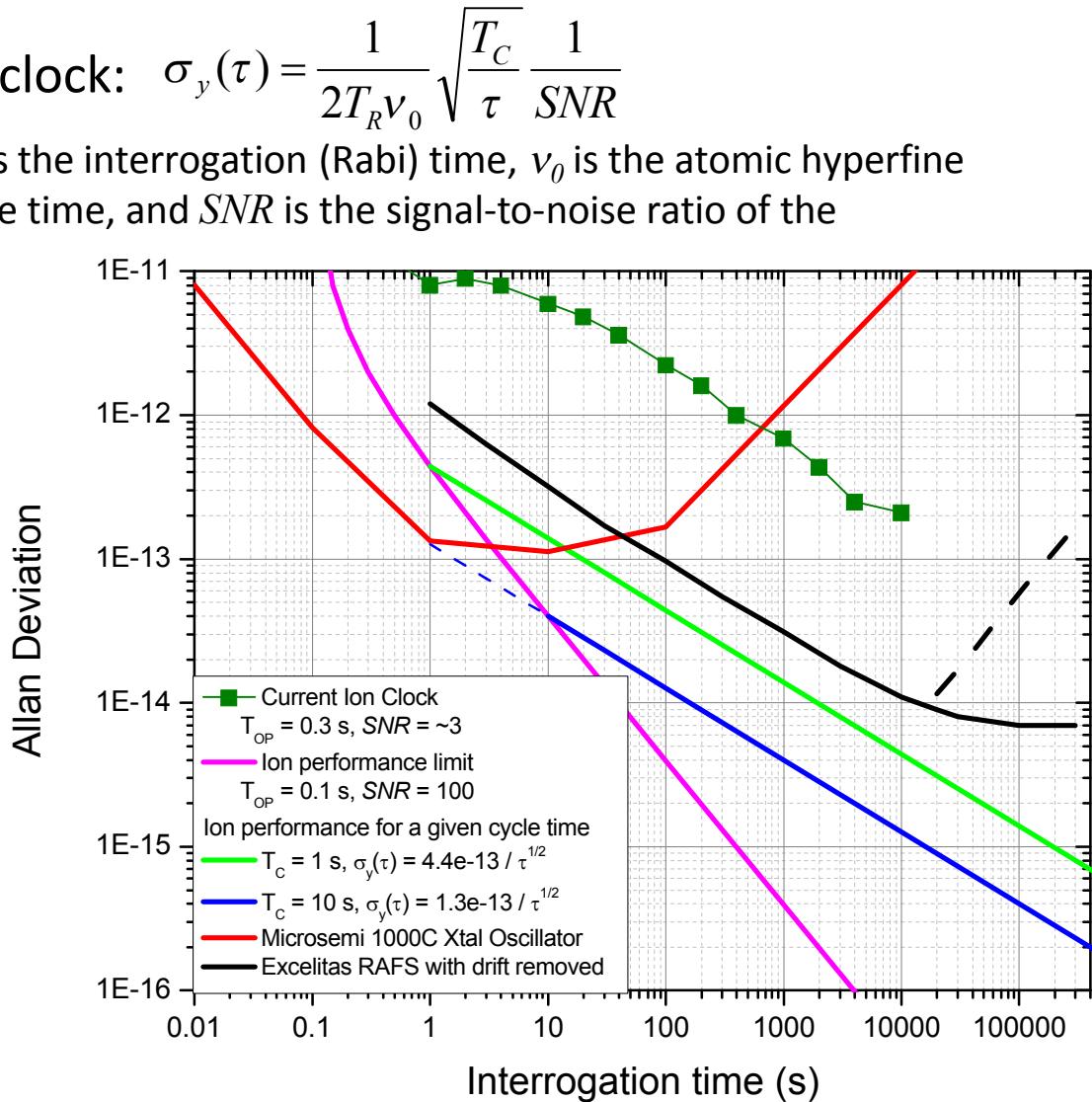
Performance of the Phase III Vacuum Package

- The vacuum package was pinched-off on Thursday, October 30th, 2014.
- Trapped ion lifetime is \sim 50 hours.
- Achieving Phase III performance
 - $T_{\text{microwave}} = 700 \text{ ms}$
 - $T_{\text{optical pumping}} = 300 \text{ ms}$
- Magnetic field correlations removed



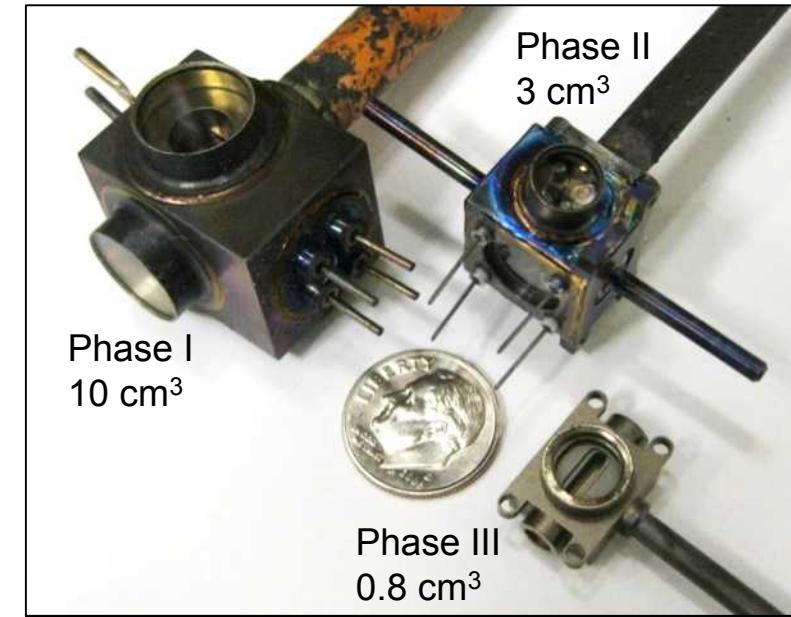
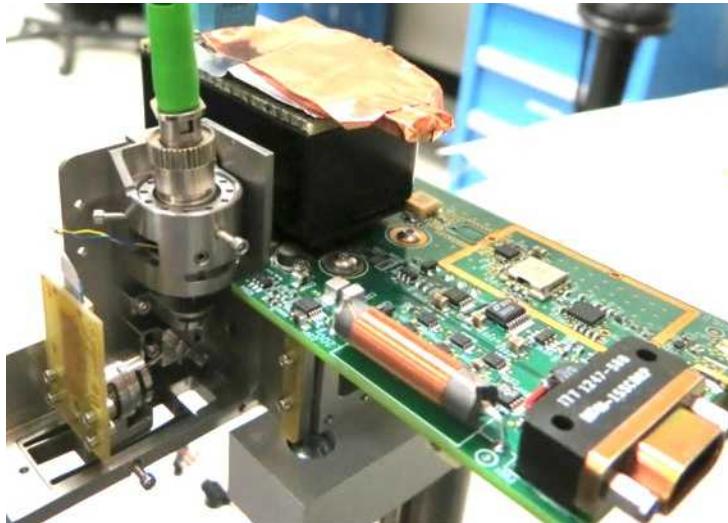
Improving the Yb Ion Clock

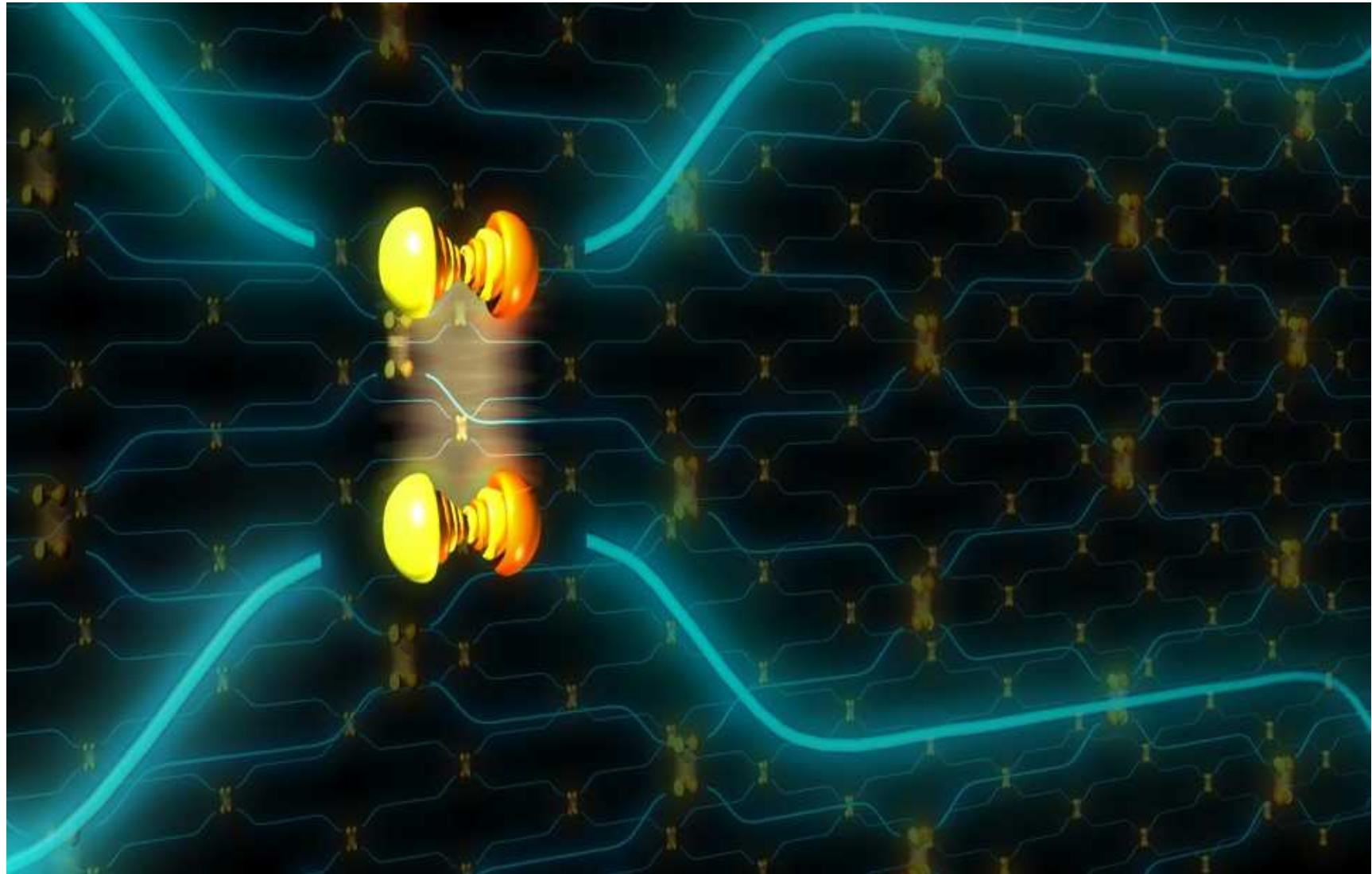
- Allan deviation for the ion clock: $\sigma_y(\tau) = \frac{1}{2T_R v_0} \sqrt{\frac{T_C}{\tau}} \frac{1}{SNR}$
- where τ is the averaging time, T_R is the interrogation (Rabi) time, v_0 is the atomic hyperfine resonance frequency, T_C is the cycle time, and SNR is the signal-to-noise ratio of the fluorescence detection.
- SNR improvement
 - Switching to 369 nm detection
 - Collect more light
 - Reducing the F-state trapping
 - Trapping more ions
- Faster optical pumping:
 - More laser power: 200-400 μ W
- Need a better local oscillator



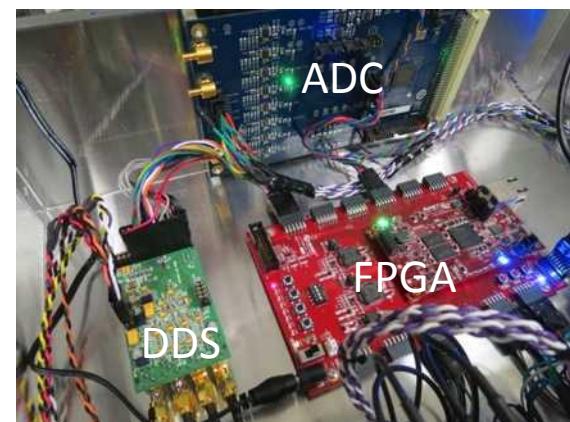
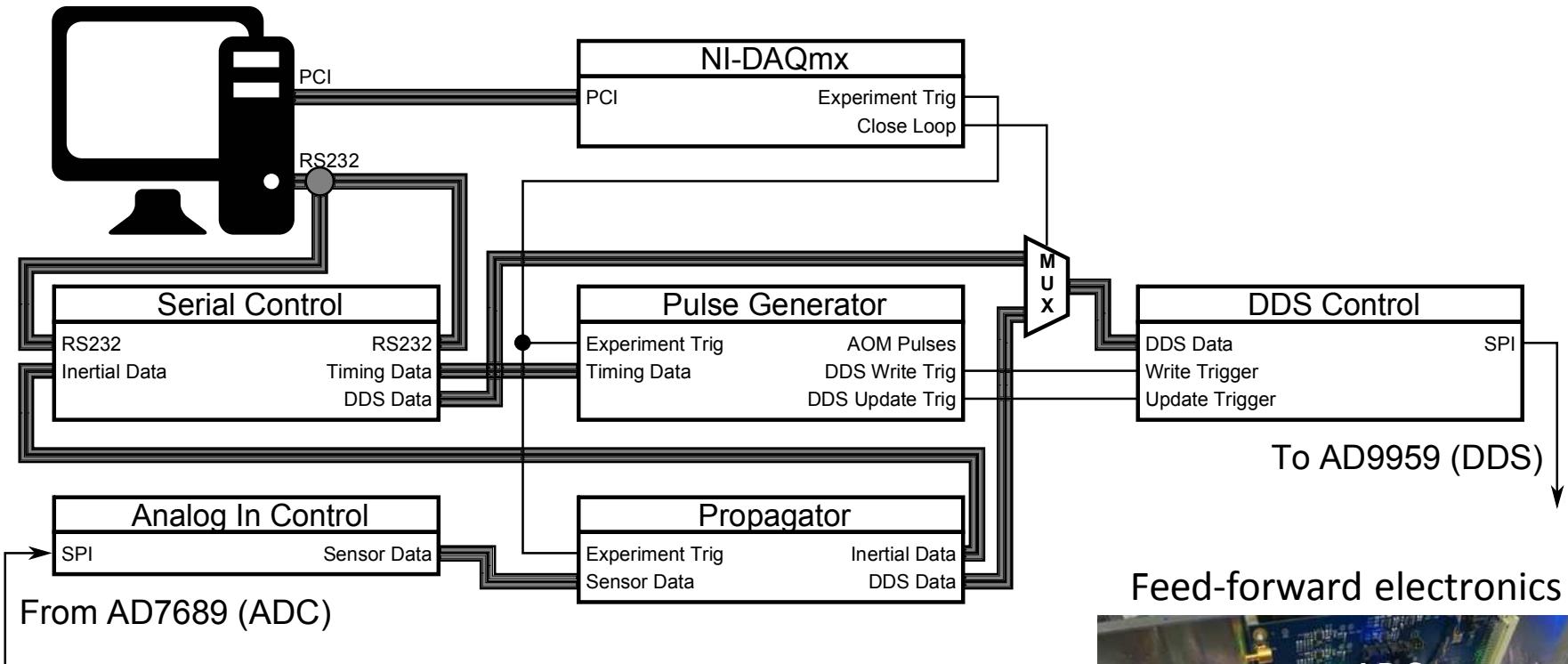
Challenges and Opportunities

- Tremendous technology development
- Need for technological improvements
 - 369 nm laser with low-power and long-term stability
 - Low power local oscillator
 - Improve signal-to-noise ratio
- Potential for a very high performance miniature clock
 - $4 \times 10^{-13} / \sqrt{\tau}$

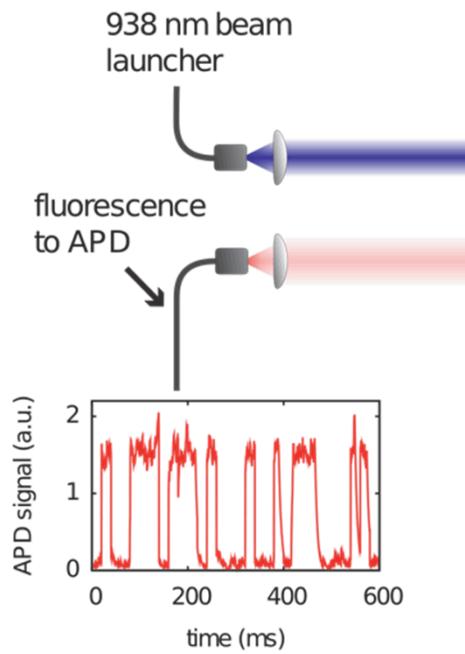




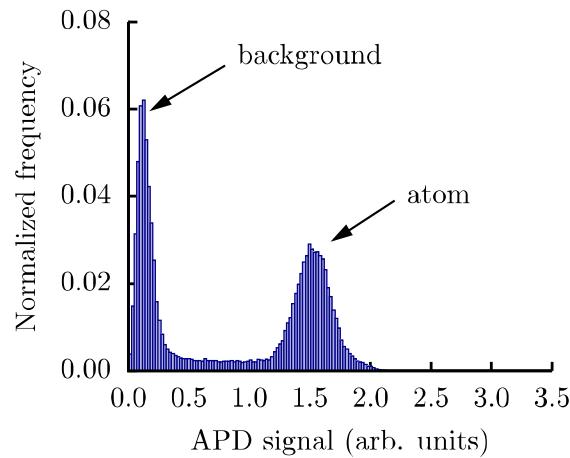
Control hardware



Single atom control

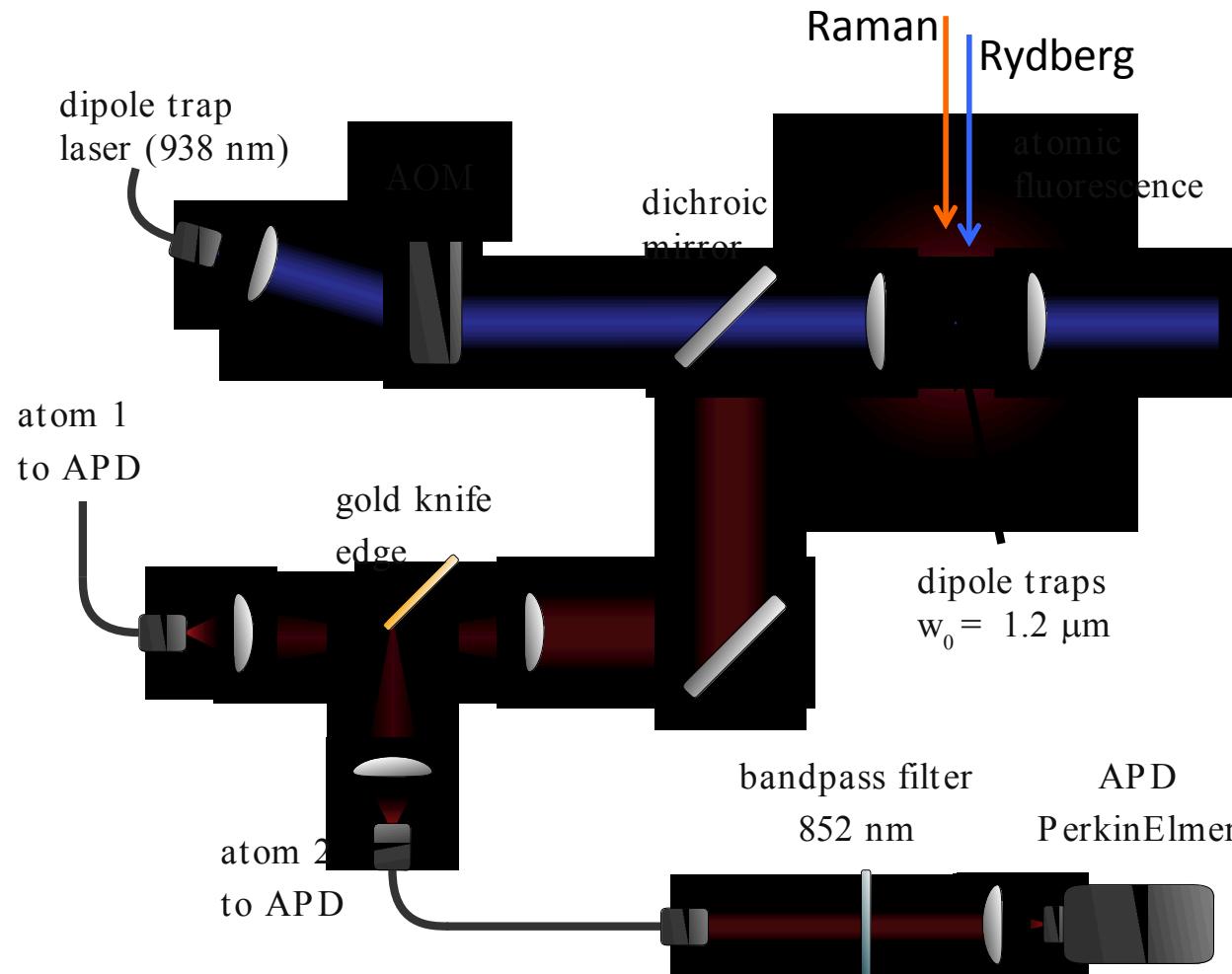


852 nm scatter:
LIF from laser cooling light



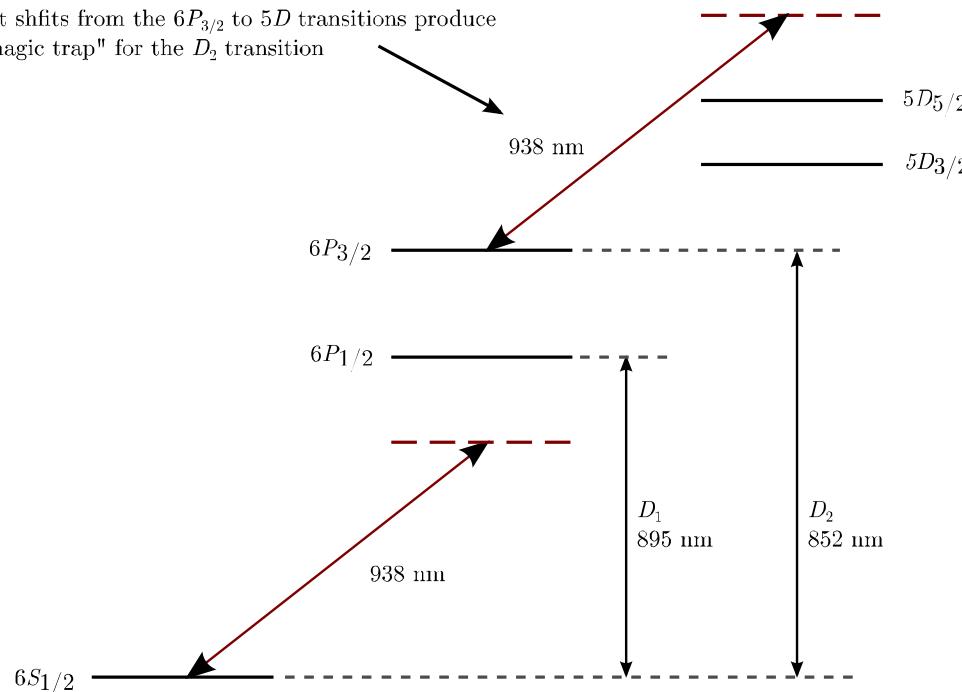
Spot size $\approx 1 \mu\text{m}$ —collisional blockade

Experiment schematic

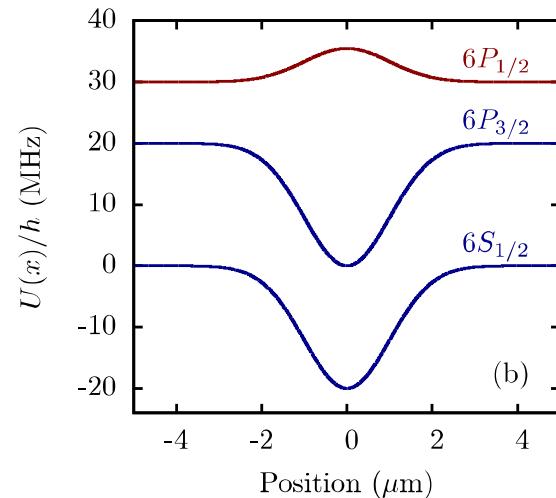


Single atom control

Light shifts from the $6P_{3/2}$ to $5D$ transitions produce a "magic trap" for the D_2 transition



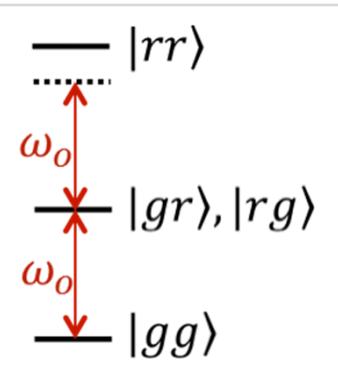
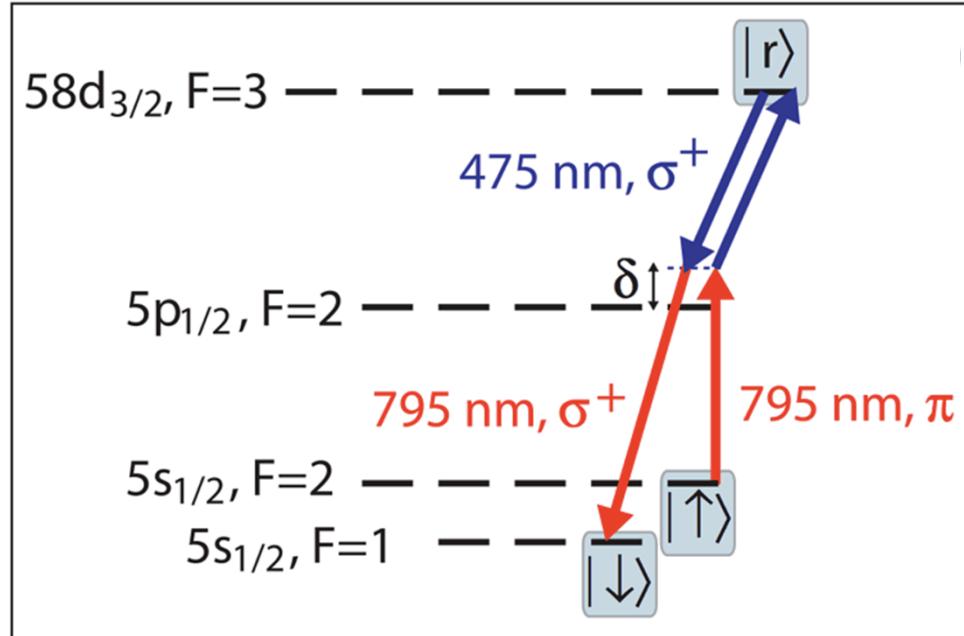
Why 938 nm? It's magic for the cooling transition.



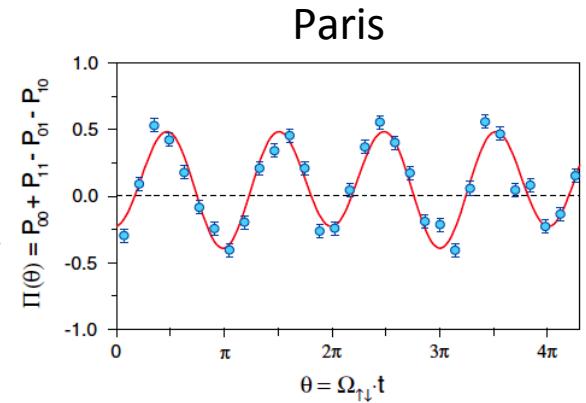
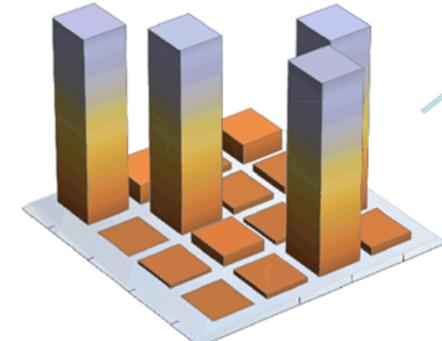
- ≈ 5 mW, 43 nm red
- focused to ≈ 1 μ m
- gives ≈ 20 MHz or ≈ 1 mK

On-demand interactions

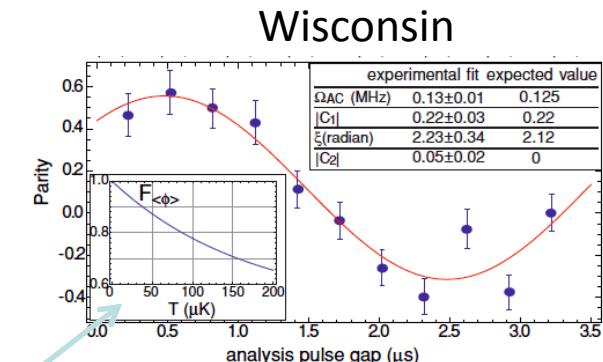
Two-atom entanglement using Rydberg blockade



CNOT gate



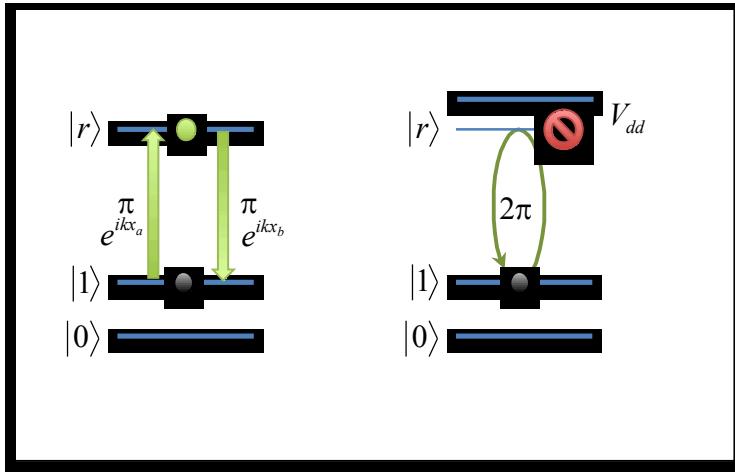
Wilk et al., PRL 104,
010502 (2010)



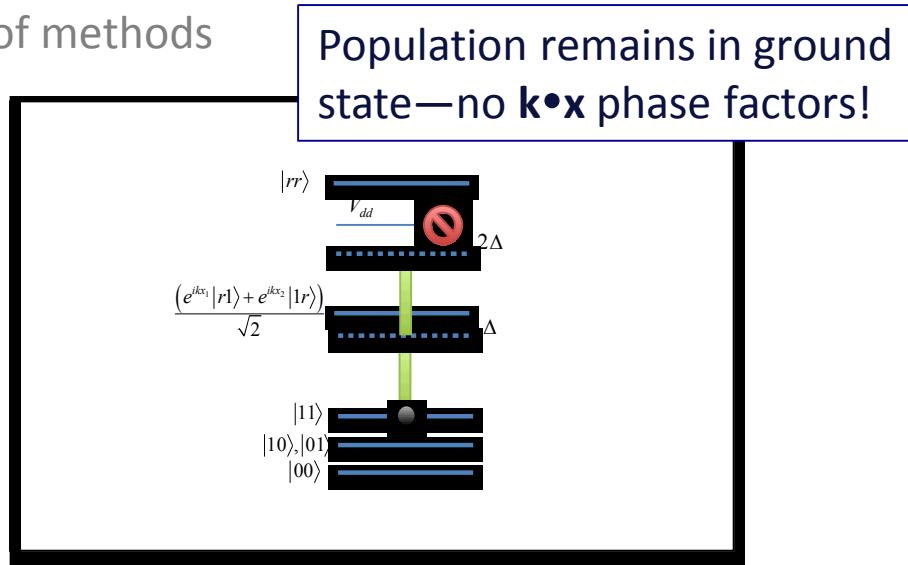
Zhang et al., PRA 82,
030306(R) (2010)

New options for Rydberg-state-mediated interactions

Comparison of methods



Direct excitation



Rydberg dressed

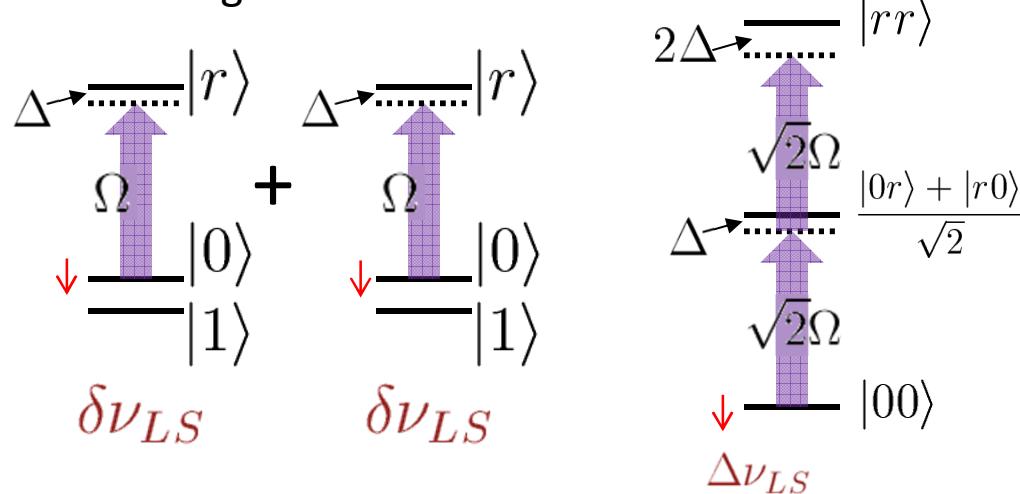
"Elaborate theoretical proposals for the realization of various complex phases and applications in quantum simulation exist. Also a simple model has been already developed that describes the basic idea of Rydberg dressing in a two-atom basis. However, an experimental realization has been **elusive so far**."

T. Pfau's group, Stuttgart, Germany

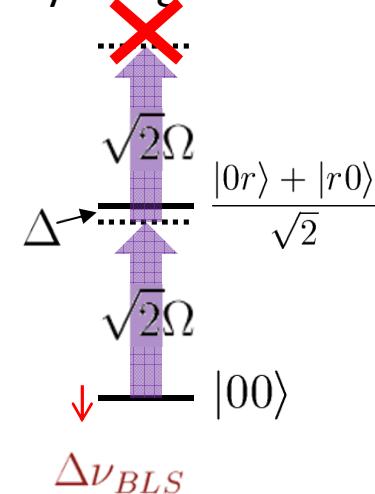
J. B. Balewski, *et al.*, *N. J. Phys.* 16, 063012 (2014)

Interaction between two Rydberg-dressed atoms

Normal light shift:



With Rydberg blockade:



$$H_{BLS} = H_{LS} + H_{int} = \begin{bmatrix} 2\delta\nu_{LS} & 0 & 0 & 0 \\ 0 & \delta\nu_{LS} & 0 & 0 \\ 0 & 0 & \delta\nu_{LS} & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} J & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \text{ for } \begin{bmatrix} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{bmatrix}$$

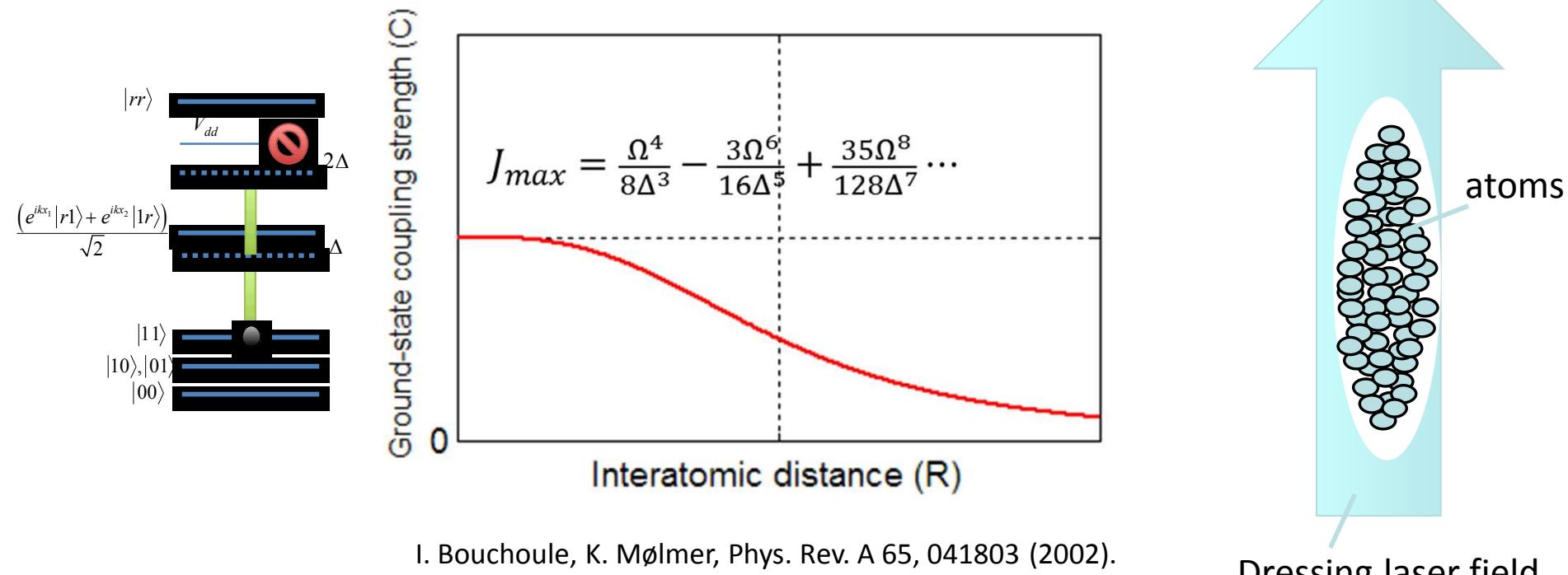
$$J = \Delta\nu_{BLS} - \Delta\nu_{LS} = \frac{\Omega^4}{8\Delta^3} - \frac{3\Omega^6}{16\Delta^5} + \frac{35\Omega^8}{128\Delta^7} - \dots$$

$$H_{int} = \frac{J}{4} (\sigma_z^{(1)} + 1)(\sigma_z^{(2)} + 1)$$

Rydberg-dressed interactions

Tunable interaction strength (J), low sensitivity to atom motion, and effectively strong ground-state interactions.

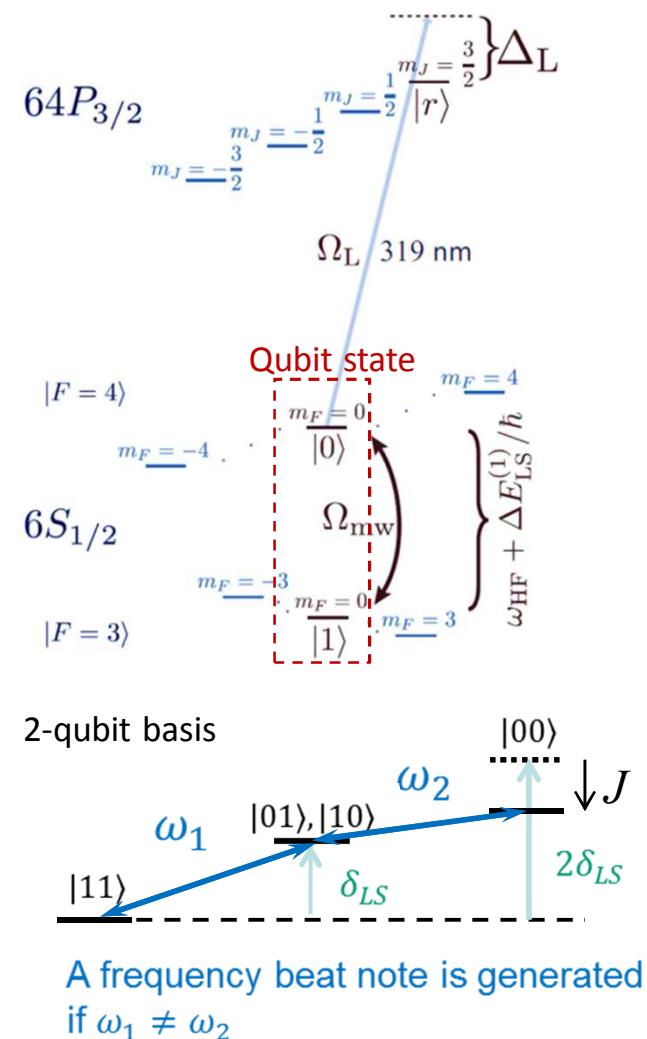
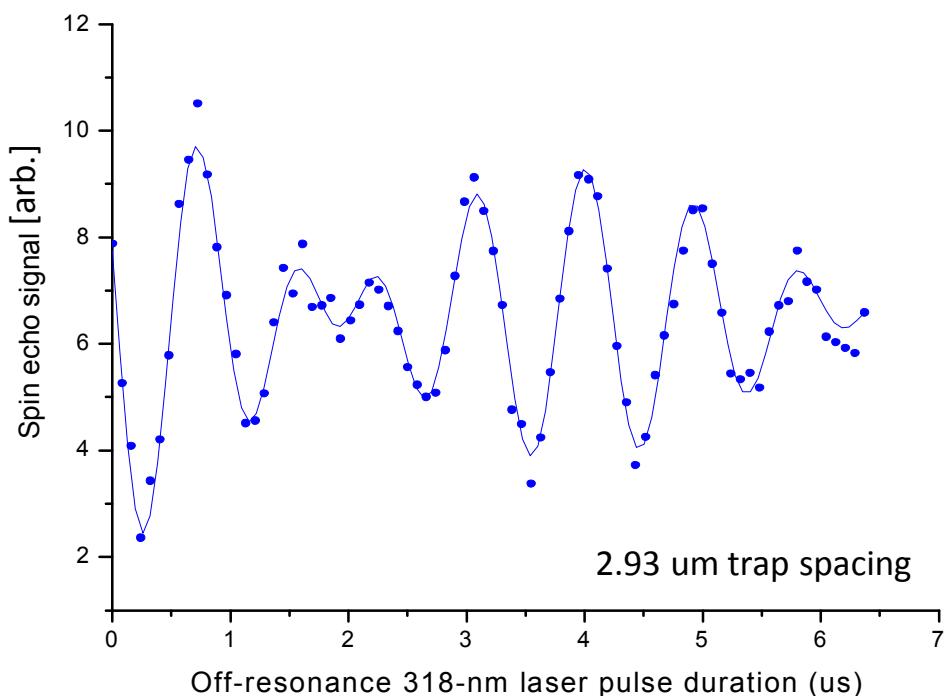
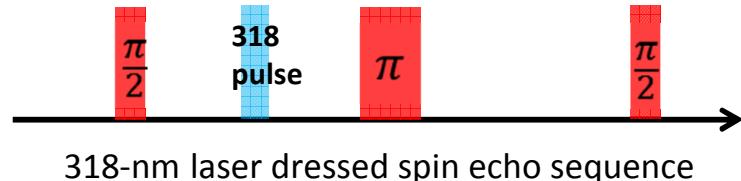
$$H_{int} = \sum_{ij} \frac{J_{ij}}{4} (\sigma_z^{(i)} + 1)(\sigma_z^{(j)} + 1)$$



I. Bouchoule, K. Mølmer, Phys. Rev. A 65, 041803 (2002).
 J. Johnson, S. Rolston, Phys. Rev. A 82, 033412 (2010).

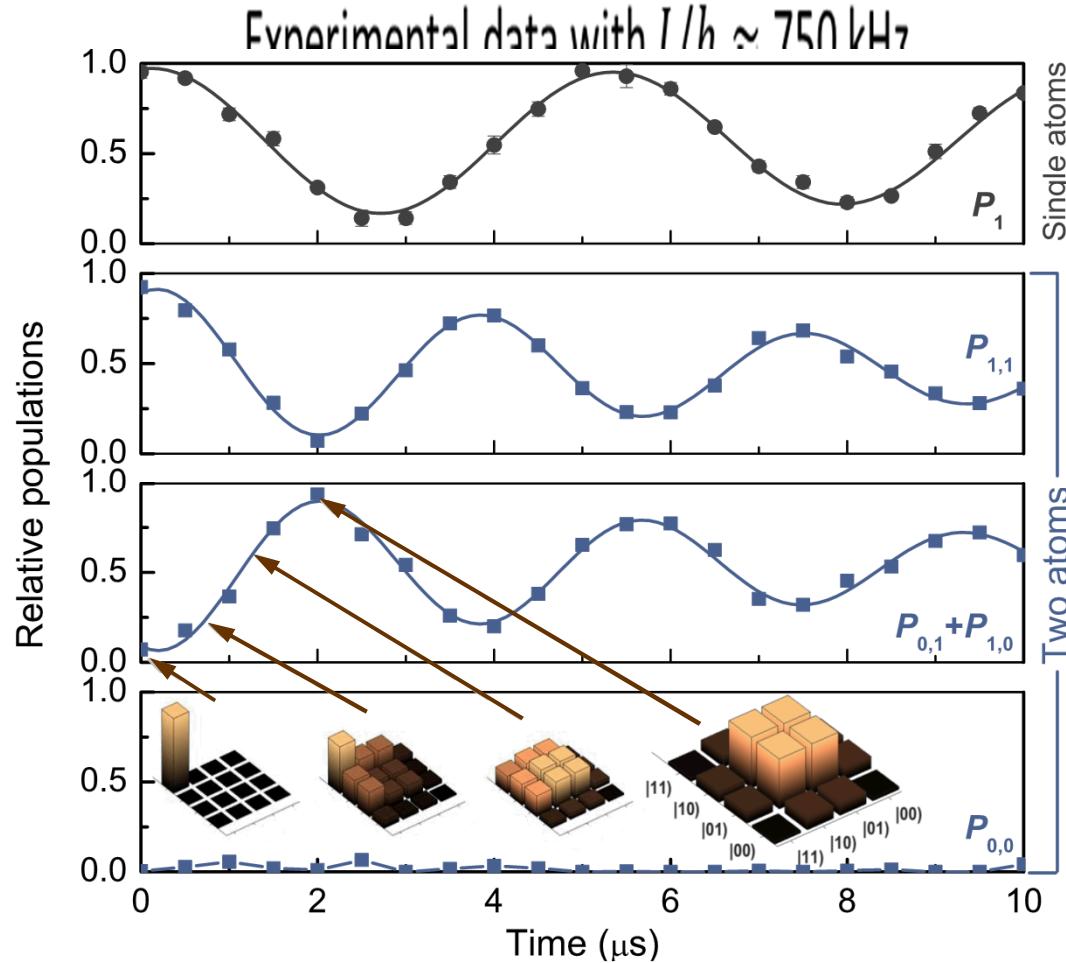
First evidence of Rydberg-dressed interaction

Microwave transition is via Raman laser



Producing Bell-state entanglement

Initial state is $|1\rangle|0\rangle|1\rangle|1\rangle$, then apply CNOT and Rabi lasers



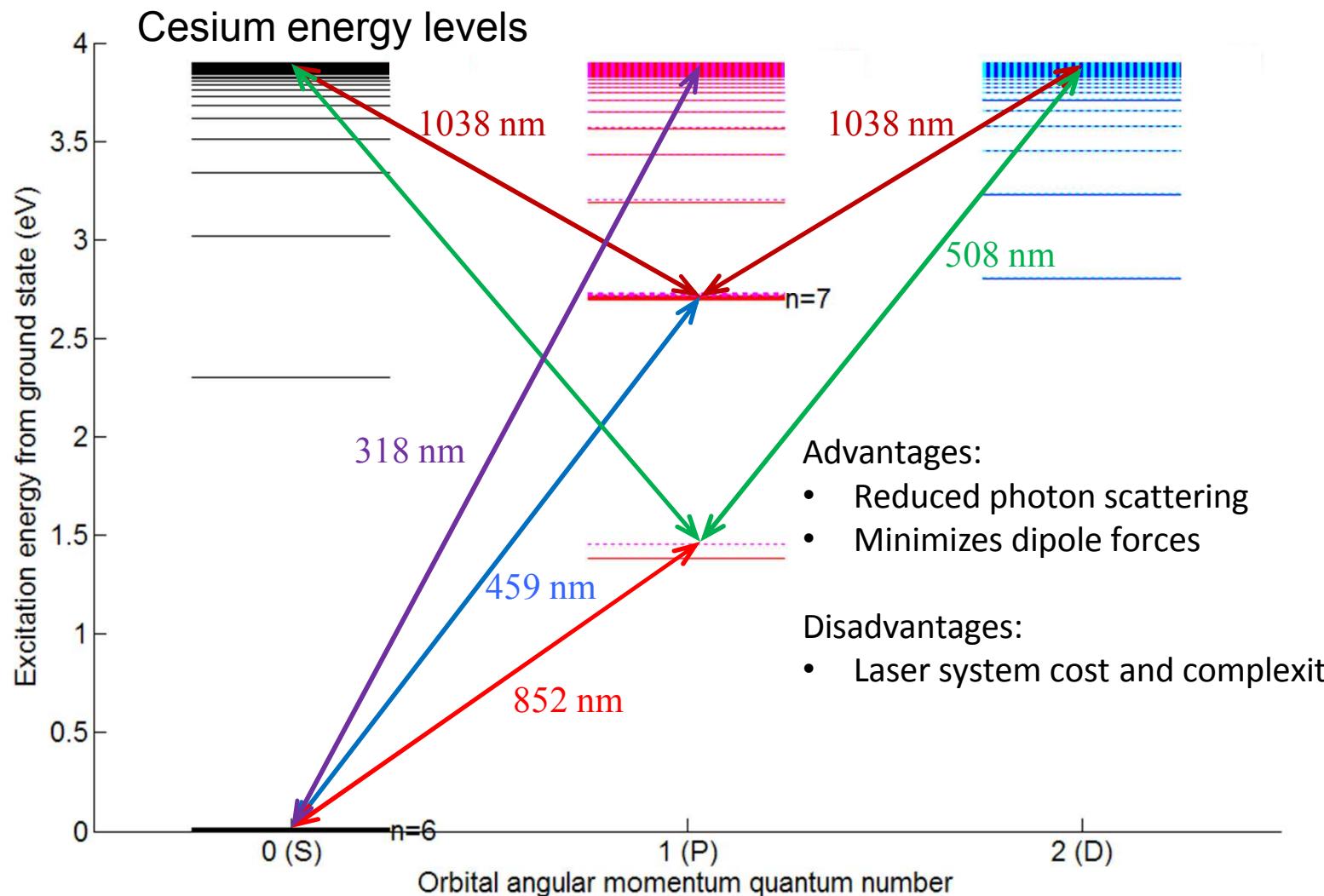
Single-atom Rabi

Two-atom Rabi oscillation:

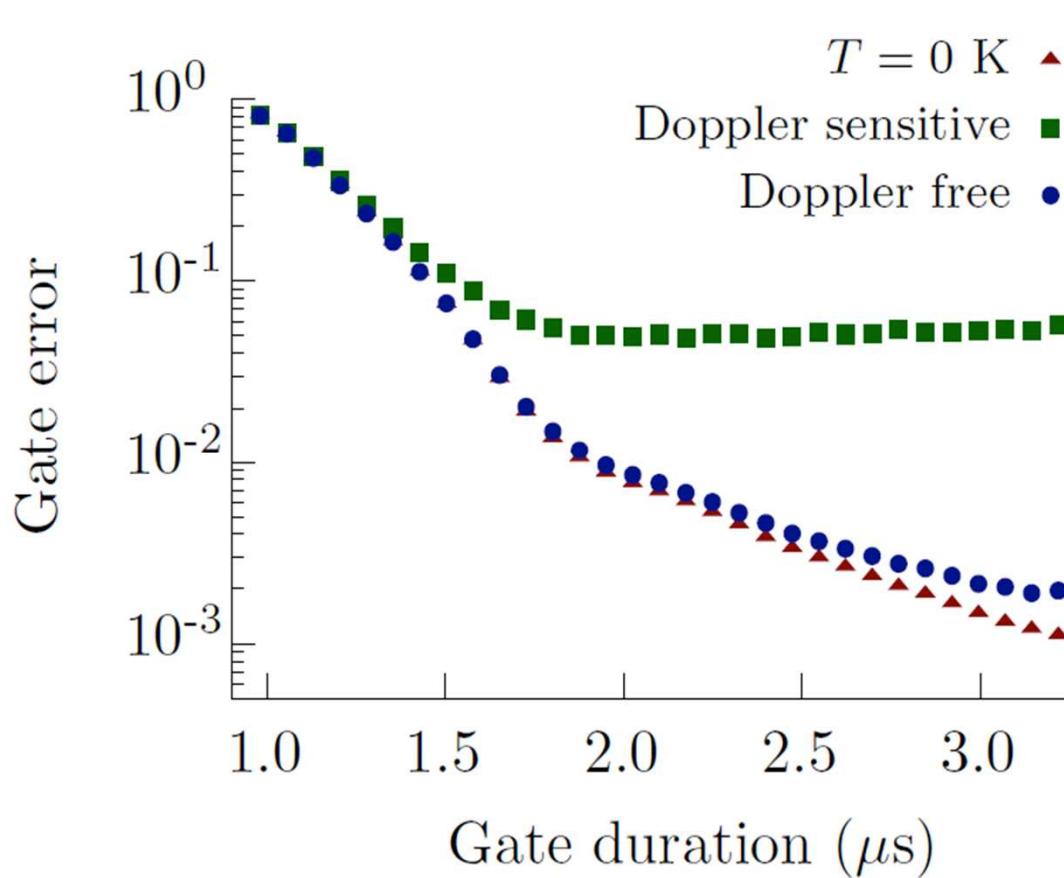
- $\sqrt{2}$ times faster
- No significant population being transferred to $|00\rangle$
- Bell state $|\Psi_+\rangle$ is produced at $t = \pi/\sqrt{2}\Omega_{mw}$

Process occurs entirely and directly in the ground state

Optimizing for long-term relationships



Simulated CPHASE gate fidelities



$\Omega = 0 \rightarrow 3 \text{ MHz}$
 $\Delta/2\pi = 6 \rightarrow 0 \text{ MHz}$
 $\Gamma = 3.7 \text{ kHz}$
 $T = 16 \mu\text{K}$

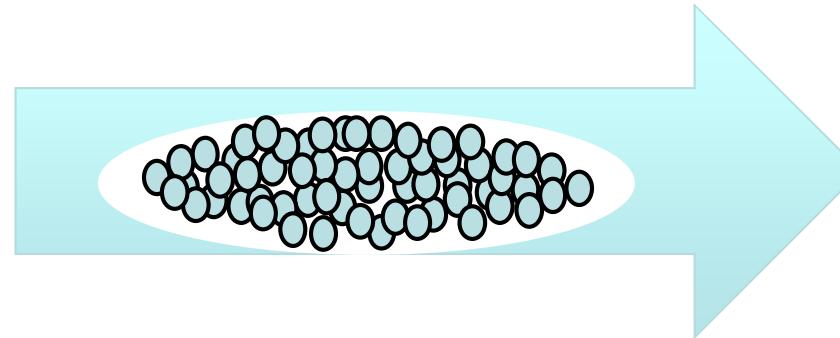
- Motional errors set a high floor on error for the original scheme.
- The Doppler-free scheme is limited by the much smaller photon scattering rate.
- Entanglement fidelity expected to be even larger

Published: Phys. Rev. A 91, 012337 (2015)

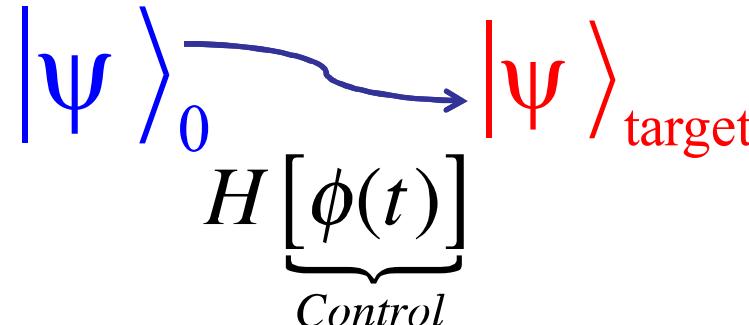


Quantum Control of Ensembles

Symmetrically couple ensemble of atoms
localized with Rydberg blockade radius



Optimal Control

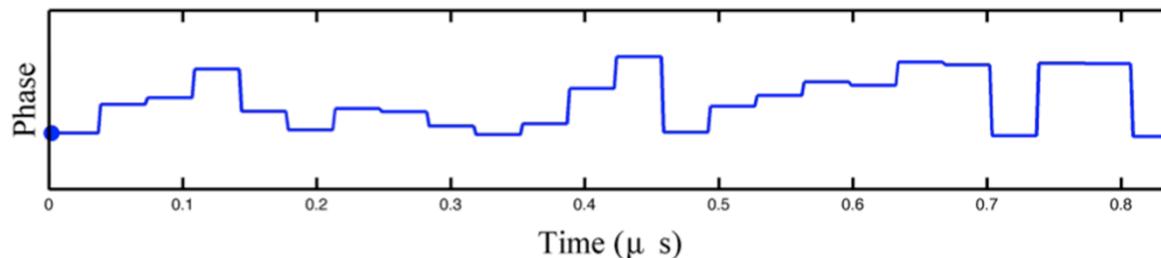


For n atoms

$$H = \sum_n \left[\frac{\Omega_{\mu w}}{2} \left(e^{i\phi(t)} |0\rangle\langle 1| + e^{-i\phi(t)} |1\rangle\langle 0| \right)^{(n)} + \frac{\Omega_R}{2} (|1\rangle\langle r| + |r\rangle\langle 1|)^{(n)} + \Delta |r\rangle\langle r|^{(n)} \right]$$

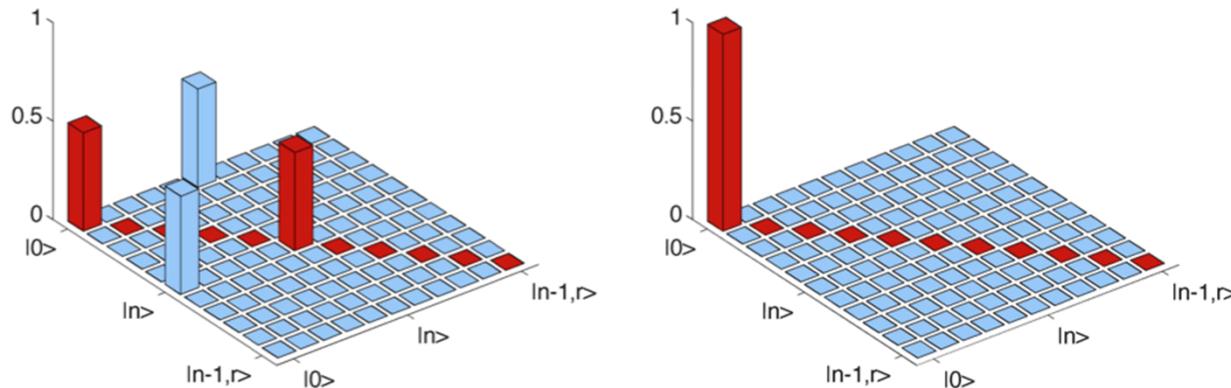
$$+ V_{dd} \sum_{n \neq m} |r\rangle\langle r|^{(n)} |r\rangle\langle r|^{(m)}$$

Example: A 5-atom “Cat State”



ρ target

ρ evolved



$$|\psi_{cat}\rangle = \frac{1}{\sqrt{2}}(|0\rangle|0\rangle|0\rangle|0\rangle|0\rangle + |1\rangle|1\rangle|1\rangle|1\rangle|1\rangle)$$

$$|\psi(t)\rangle$$

Miniature Laser Sources

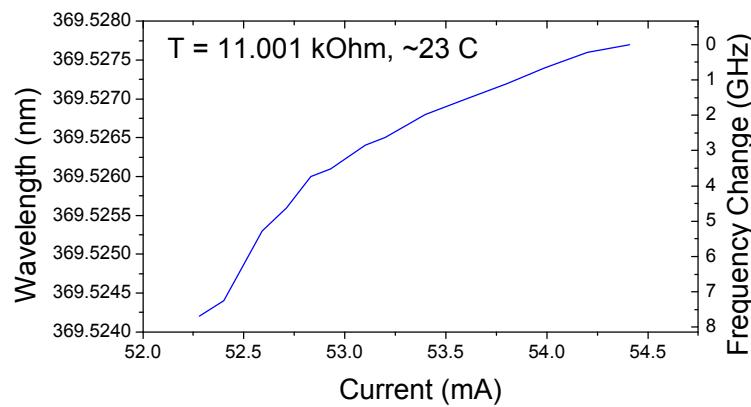
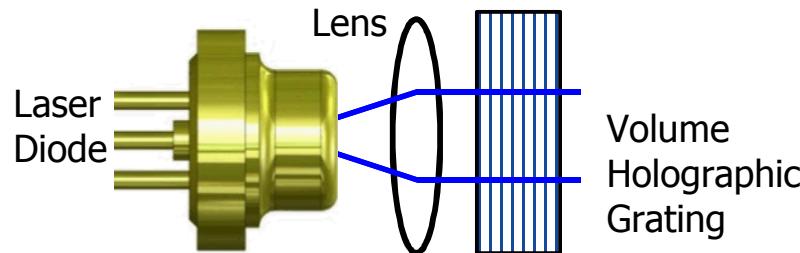
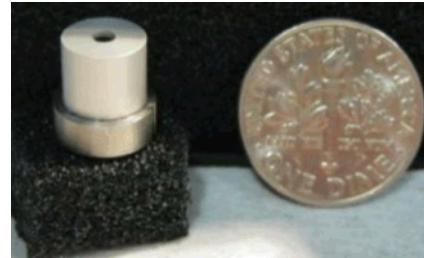
369 nm Direct Diode

Nichia diode packaged by Ondax. Size: 1 cm³

Mode-hop free tuning: 8 GHz

Power consumption: 260 mW (laser) + 100 mW (TEC)

Power output: 1-5 mW



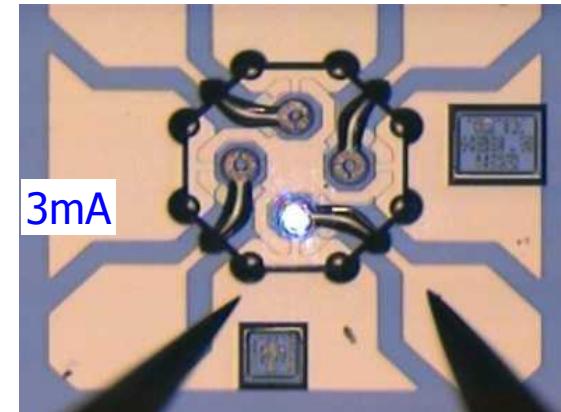
935 nm VCSEL

935 nm VCSEL fabricated at Sandia.

Single mode P > 1 mW

Power consumption: 2.5 mW (laser) + 18 mW (heater)

935nm VCSEL (V122571)



Dynamic atom positioning

