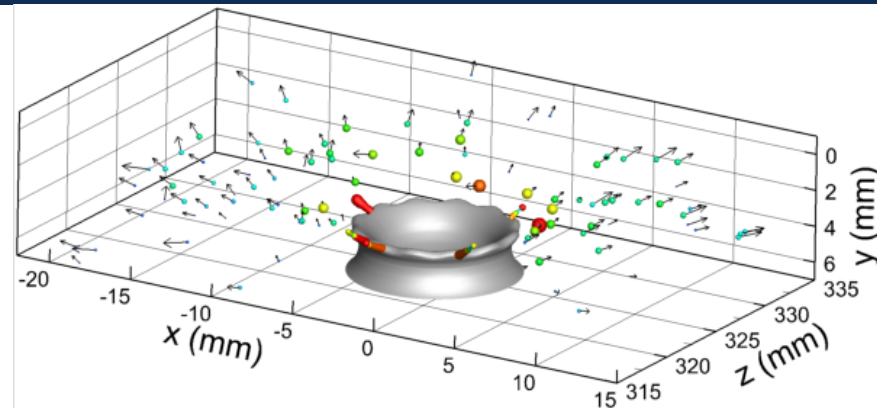
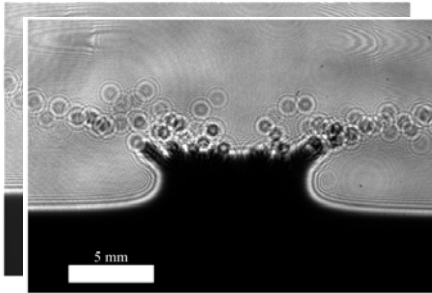


Exceptional service in the national interest



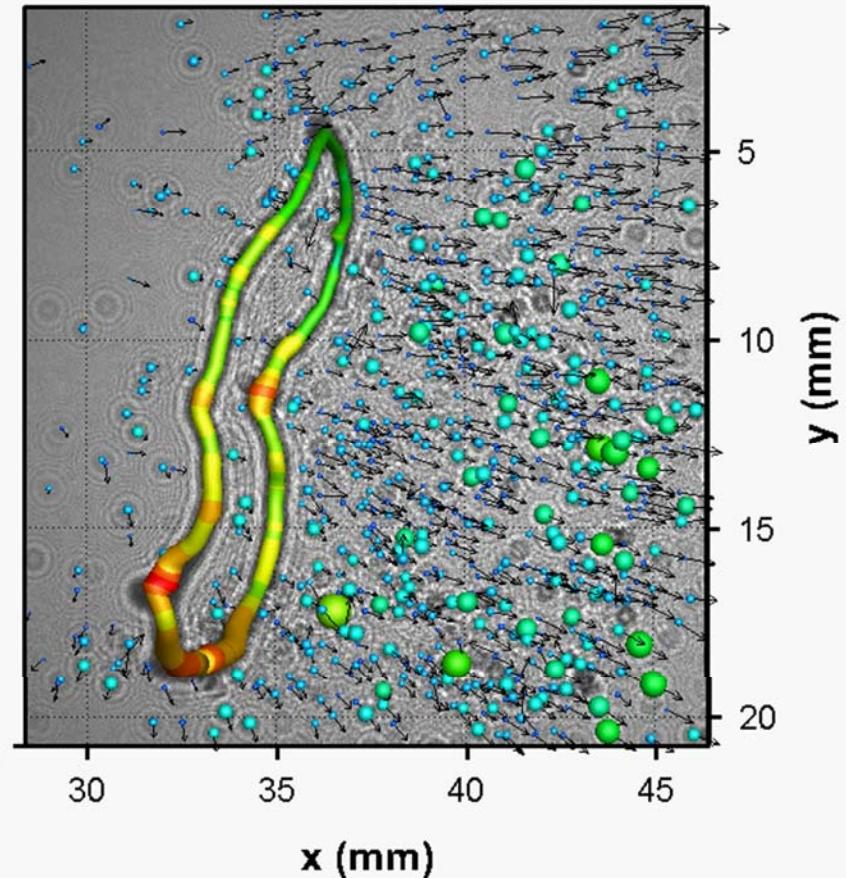
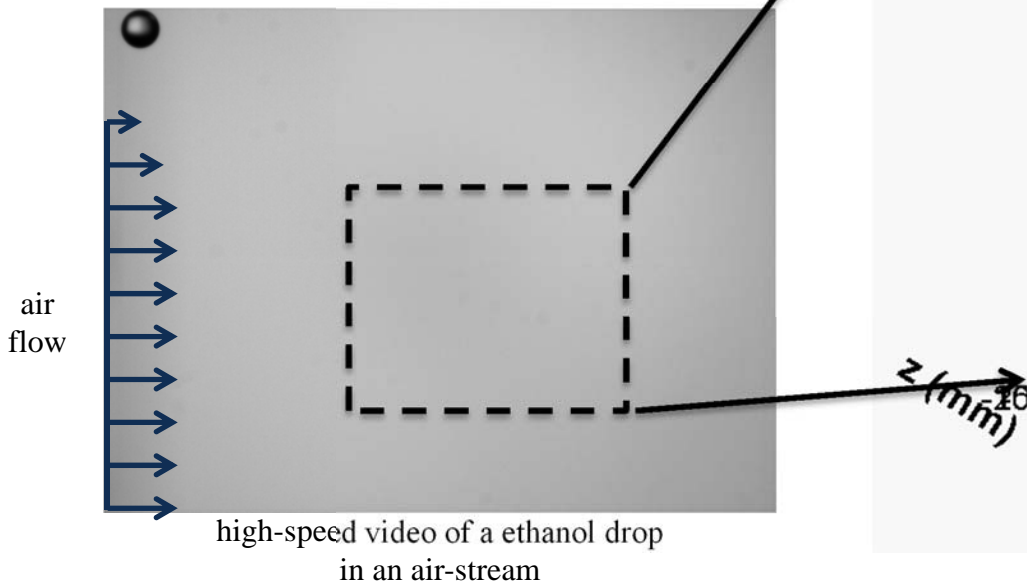
Digital In-line Holography (DIH) for 3D Quantification of Liquid Sprays

Daniel R. Guildenbecher
October 14, 2015

Motivation: 3D imaging for a 3D world

Widely available 2D imaging or point-wise measurement techniques are often insufficient to resolve 3D flow phenomena

- Repetition needed to capture spatial statistics

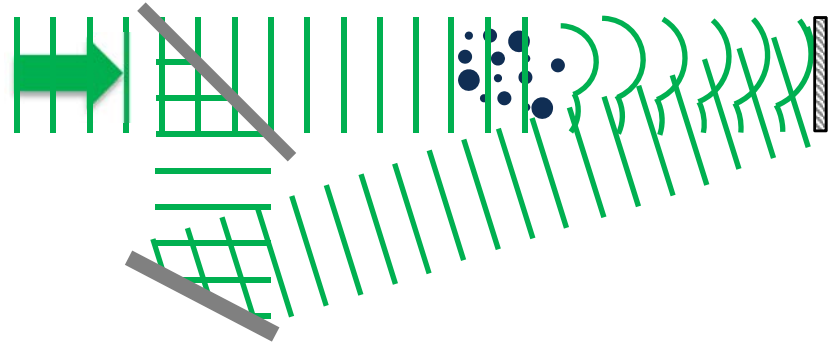


digital holographic measurement
(Gao, Guildenbecher et al, 2013, *Opt. Lett.*)

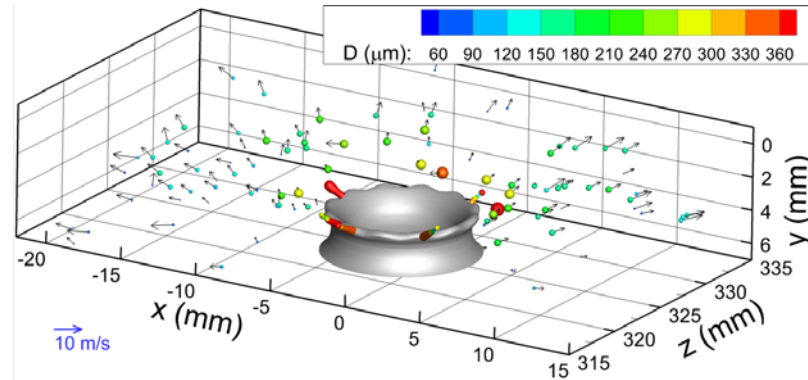
Holography is an optical technique to record and reconstruct a 3D light field

Outline for talk

Introduction to holography and
the “digital revolution”



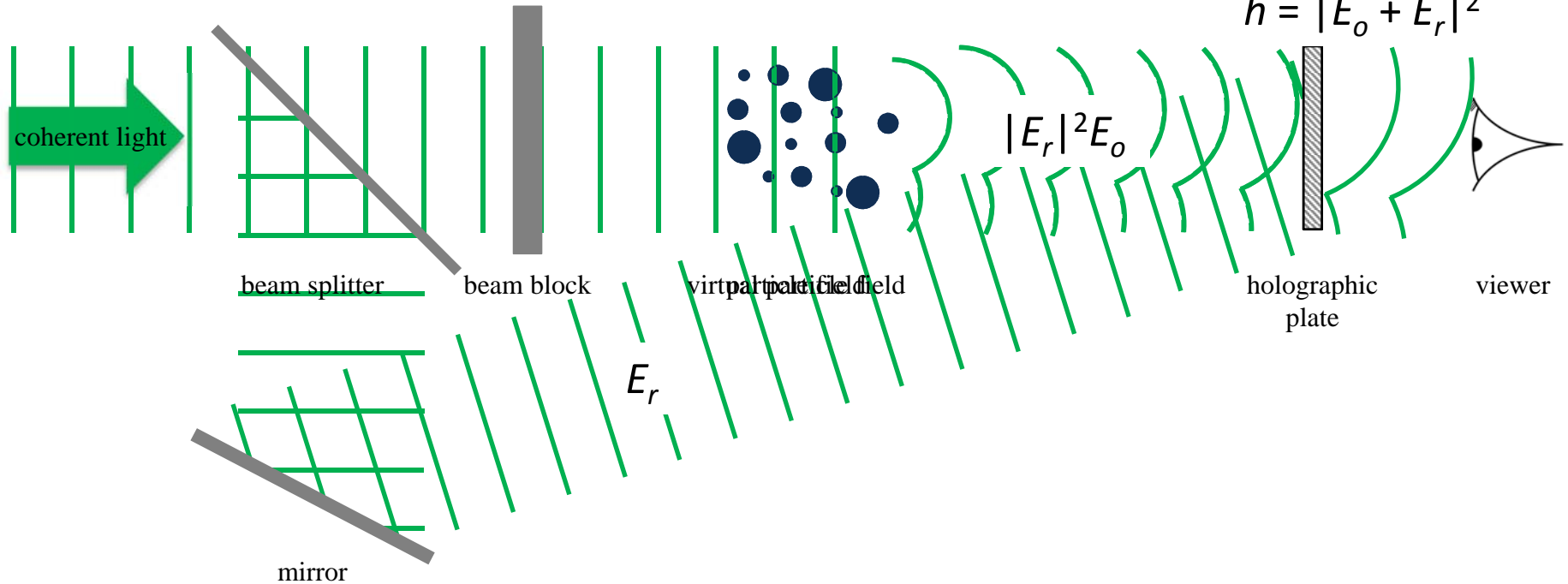
Application to liquid sprays



Propellant fire measurements



What is holography?



Optical method first proposed by Gabor in 1948

1. Coherent light diffracted by particle field forms the object wave, E_o
2. Interference with a reference wave, E_r , forms the hologram: $h = |E_o + E_r|^2$
3. Reconstruction with E_r forms virtual images at original particle locations

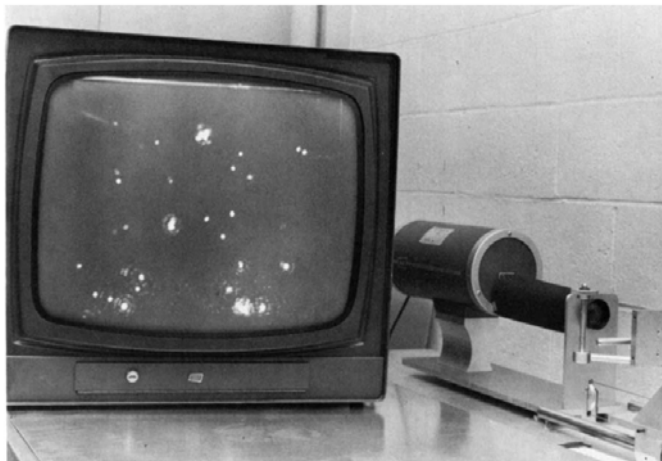
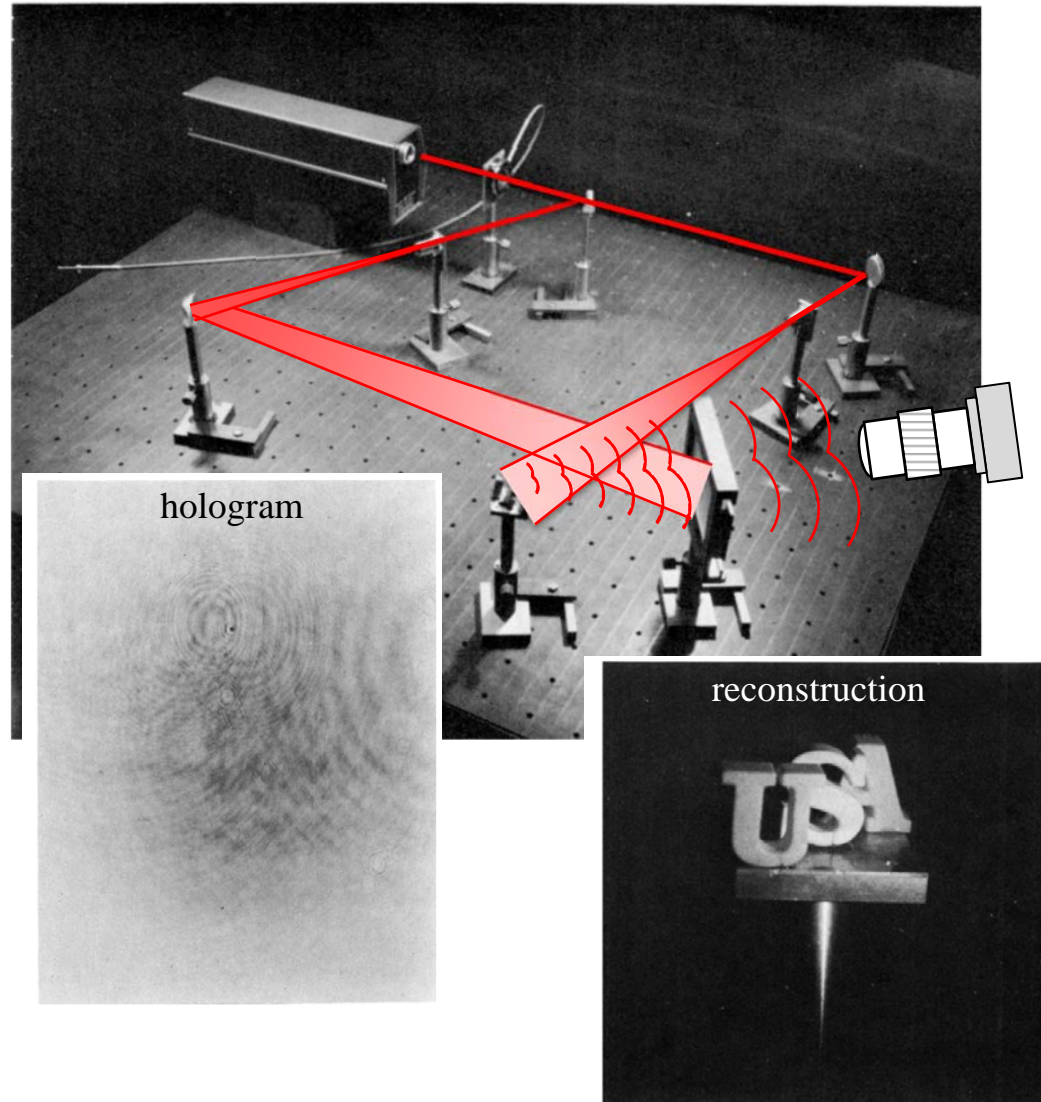
$$h \cdot E_r = \underbrace{(|E_o|^2 + |E_r|^2)E_r}_{\text{DC term}} + \underbrace{|E_r|^2 E_o}_{\text{virtual image}} + \underbrace{E_r^2 E_o^*}_{\text{real image}}$$

Analog holography

Applications of holography took off with invention of the laser in 1960

Challenges:

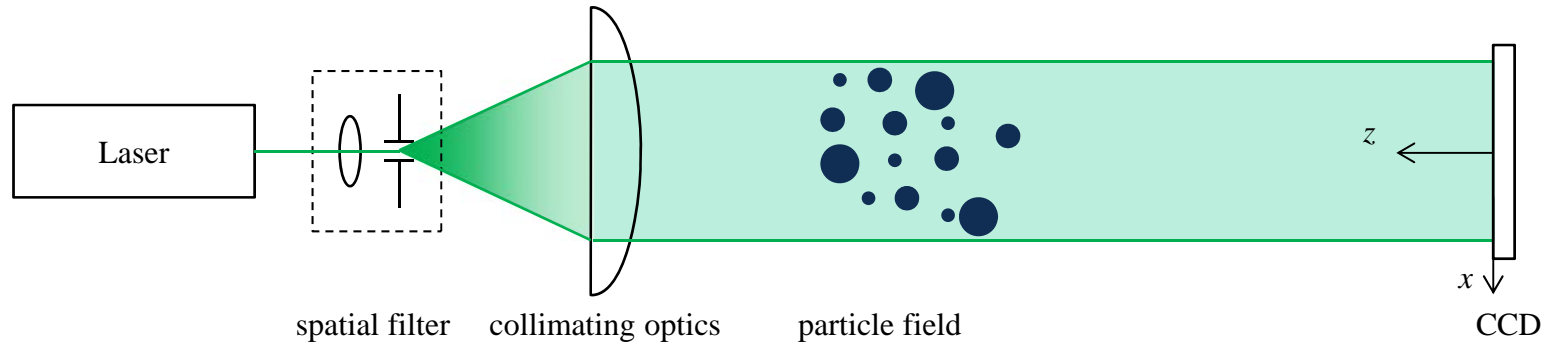
- Darkroom needed to process the hologram
- Limited temporal resolution
- Manual post processing



Thompson et al, 1967, *Appl. Opt.*

Collier et al, 1971, *Optical Holography*

Digital in-line holography (DIH)



Holographic plate and wet-chemical processing replaced with digital sensor

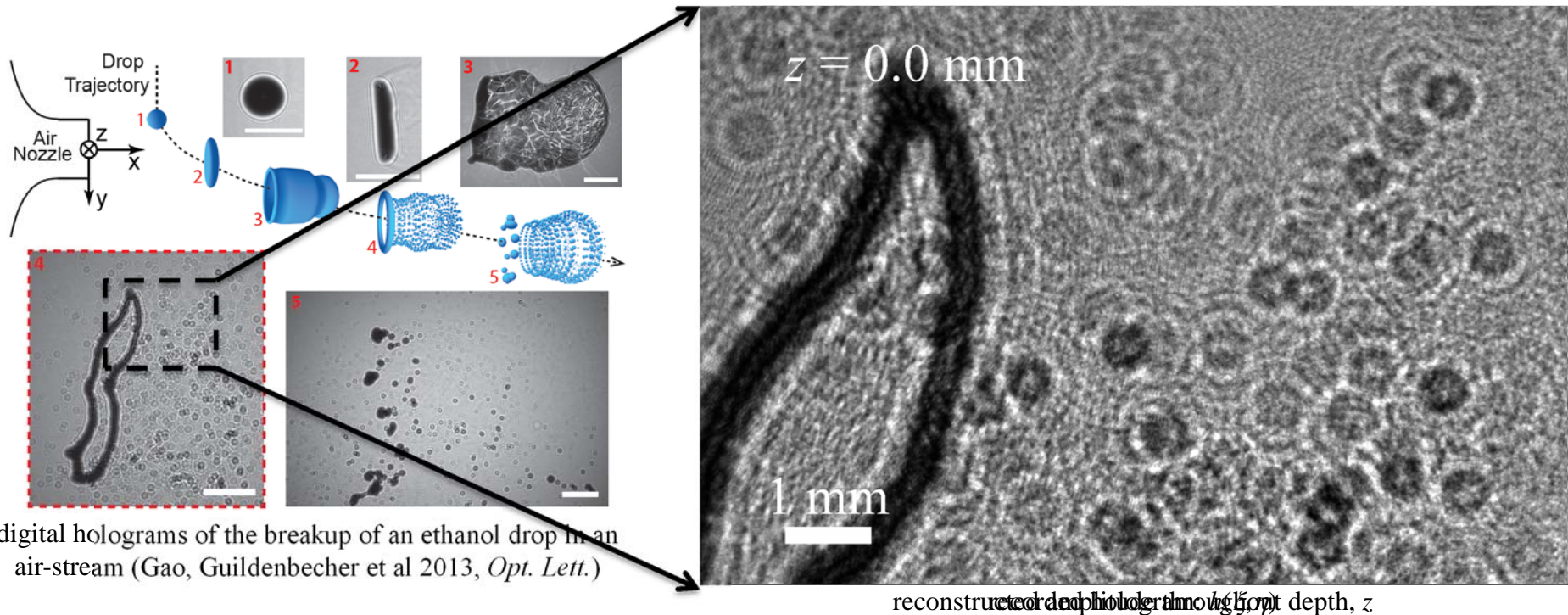
- First proposed by Schnars and Jüptner in '90s
- Advantages: (1) no darkroom, (2) temporal resolution is straight forward, (3) results can be numerically refocused and post-processed
- Challenge: Resolution of digital sensors (order 100 line pairs/mm) is much less than resolution of photographic emulsions (order 5,000 line pairs/mm)
 - For suitable off axis angles, θ , the fringe frequency, f , is typically too large to resolve with digital sensors ($f = 2\sin(\theta/2)/\lambda$)
 - Rather, the in-line configuration ($\theta = 0$) is typically utilized

Numerical refocusing

Light propagation in a non-absorbing, constant index of refraction medium is described by the diffraction integral equation:

$$E(x, y, z) = \frac{1}{\lambda} \iint E(\xi, \eta, z=0) \frac{e^{-jkr}}{r} d\xi d\eta \quad \text{where: } r = \sqrt{(\xi - x)^2 + (\eta - y)^2 + z^2}$$

- $E(\xi, \eta, 0) \equiv$ complex amplitude at hologram plane = $h(\xi, \eta) \cdot E_r^*$
- $E(x, y, z) \equiv$ refocused complex amplitude at optical depth z

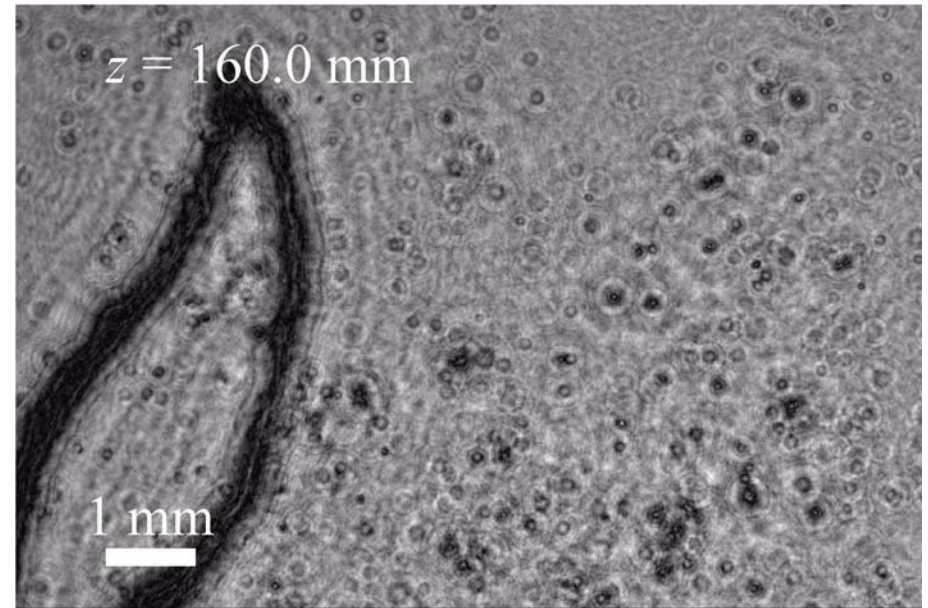


Data processing

Acquisition and refocusing of a digital hologram is relatively straightforward.

However...

For quantitative measurements, methods are required to locate and measure particles.



Challenge: depth-of-focus problem

The spatial extent of the diffraction pattern limits the angular aperture, Ω , from which a particle is effectively reconstructed (Meng et al, 2004, *Meas. Sci. Technol.*)

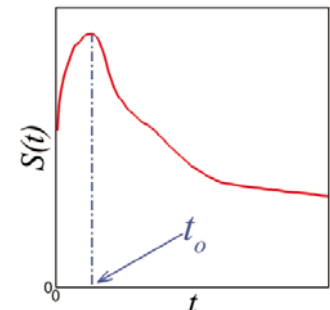
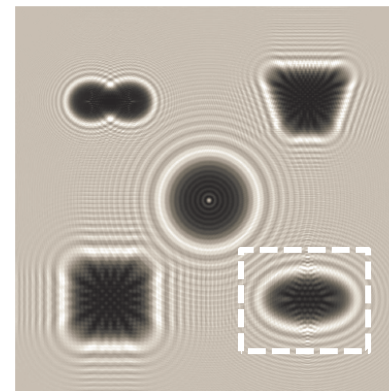
- From the central diffraction lobe $\rightarrow \Omega \approx 2\lambda/d$
- Using the traditional definition of depth-of-focus, δ , based on change of intensity within the particle center $\rightarrow \delta \approx 4\lambda/\Omega^2$
- Therefore: for in-line holography, $\delta \approx d^2/\lambda$
 - Example: $d = 300 \mu\text{m}$, $\lambda = 532 \text{ nm} \rightarrow \delta \approx 170 \text{ mm!}$

Literature contains two basic methods to find the focal plane:

1. Fit a model to the observed diffraction patterns (inverse method)
 - Generally accurate with small depth uncertainty
 - Limited to objects with known diffraction patterns (spheres)
2. Reconstruct the amplitude (or intensity) throughout depth and apply a focus metric to find “in-focus” objects
 - No *a-priori* knowledge of particle shape required
 - Accuracy is a strong function of the chosen focus metric

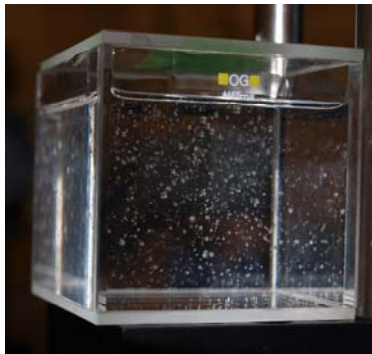
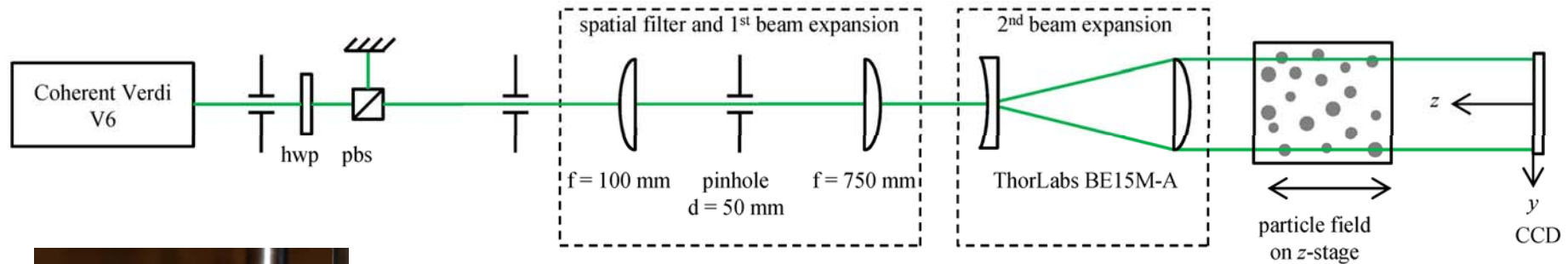
Hybrid method:

- Focus metric is a combination of amplitude minimization and edge sharpness maximization
 - Details in Guildenbecher et al 2013, *Appl. Opt.*; Gao et al 2013, *Opt. Express*; Gao et al 2014, *Appl. Opt.*



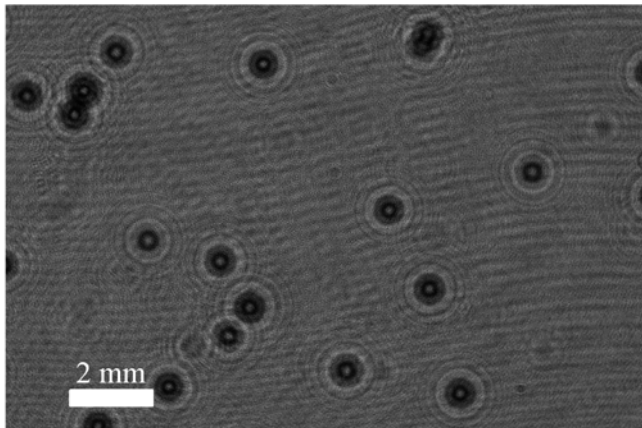
Gao et al 2014, *Appl Opt.*

Experimental validation

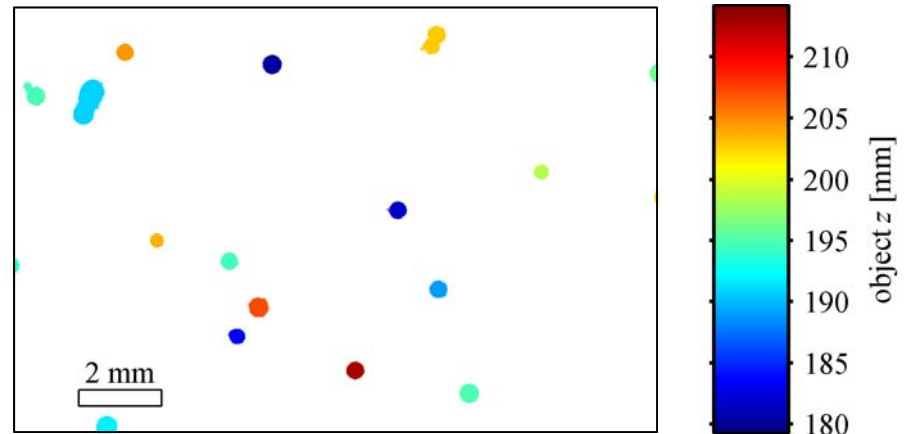


particle field

- Quasi-stationary particle field
 - Polystyrene beads ($\bar{d} \approx 465 \mu\text{m}$) in 10,000 cSt silicone oil
 - Settling velocity $\approx 0.8 \mu\text{m/s}$
- Multiple holograms recorded, displacing the particle field 2 mm in the z-direction between each acquisition

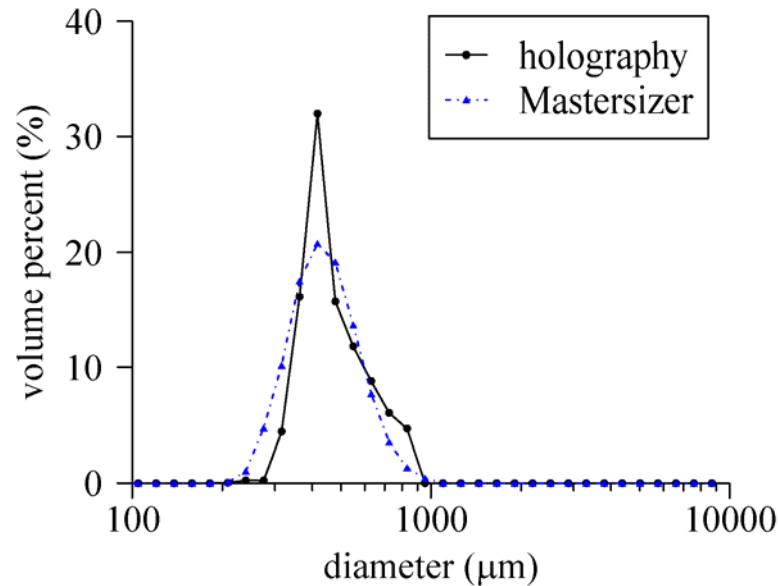


hologram



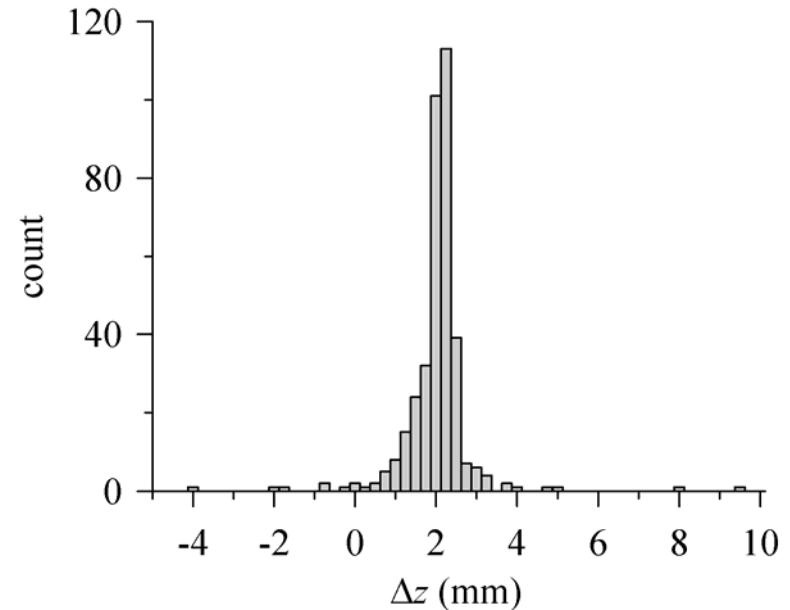
Detected objects colored by z-position

Experimental validation



Diameter measured from area of the detected 2D morphology

- Actual mass median diameter = 465 μm
- Measured mass median diameter = 474 μm
 - Error of 2.0% with respect to actual value



Displacement found by particle matching between successive holograms

- Actual displacement = 2.0 mm
- Mean detected displacement = 1.91 mm +/- 0.81 mm
 - Standard deviation of 1.74 times mean diameter

Application to Liquid Sprays

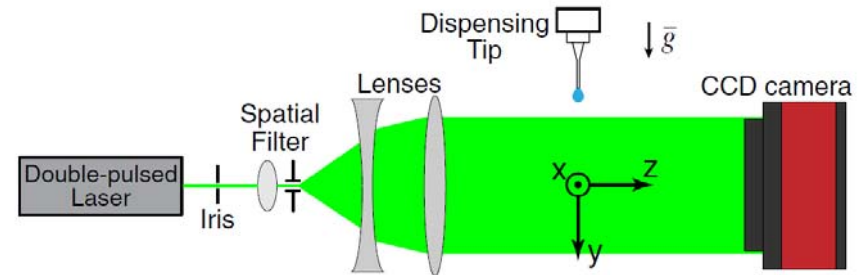
Aerodynamic drop fragmentation

Experimental configuration: Double-pulsed laser and imaging hardware as typically used in PIV

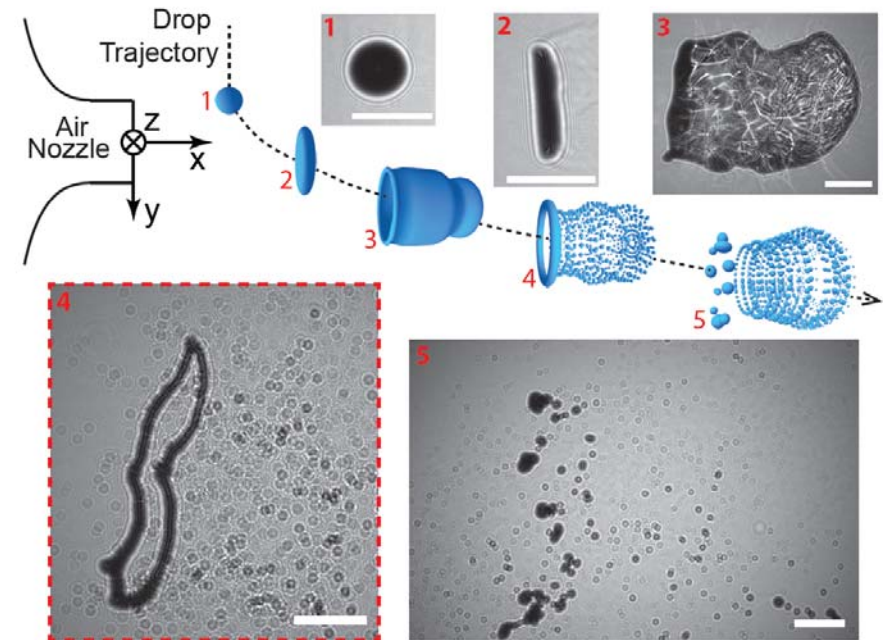
- $\lambda = 532 \text{ nm}$, 5 ns pulsewidth
- Interline transfer CCD (4008 \times 2672, 9 μm pixel pitch)
- Temporal separation, $\Delta t = 62 \mu\text{s}$, determined by laser timing

Note: without a separate reference wave, coherence length requirements in DIH are greatly relaxed.

- Expensive injection seeders are not always needed
- Faster lasers (ps or fs) can be used with some advantages (e.g. Nicolas et al 2007, *Opt. Express*)



Optical configuration (Gao, Guildenbecher et al 2013, *Opt. Lett.*)

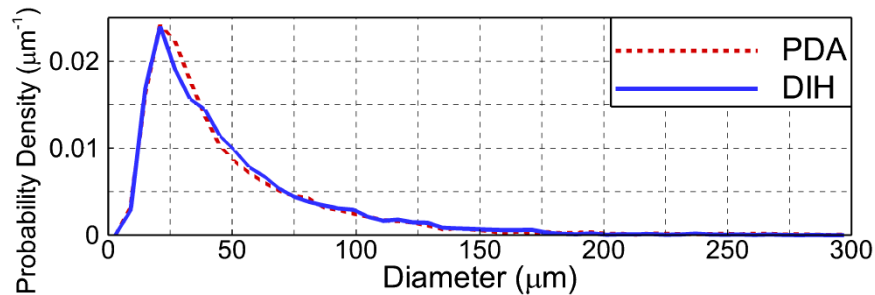


digital holograms of the breakup of an ethanol drop in an air-stream (Gao, Guildenbecher et al 2013, *Opt. Lett.*)

Aerodynamic drop fragmentation

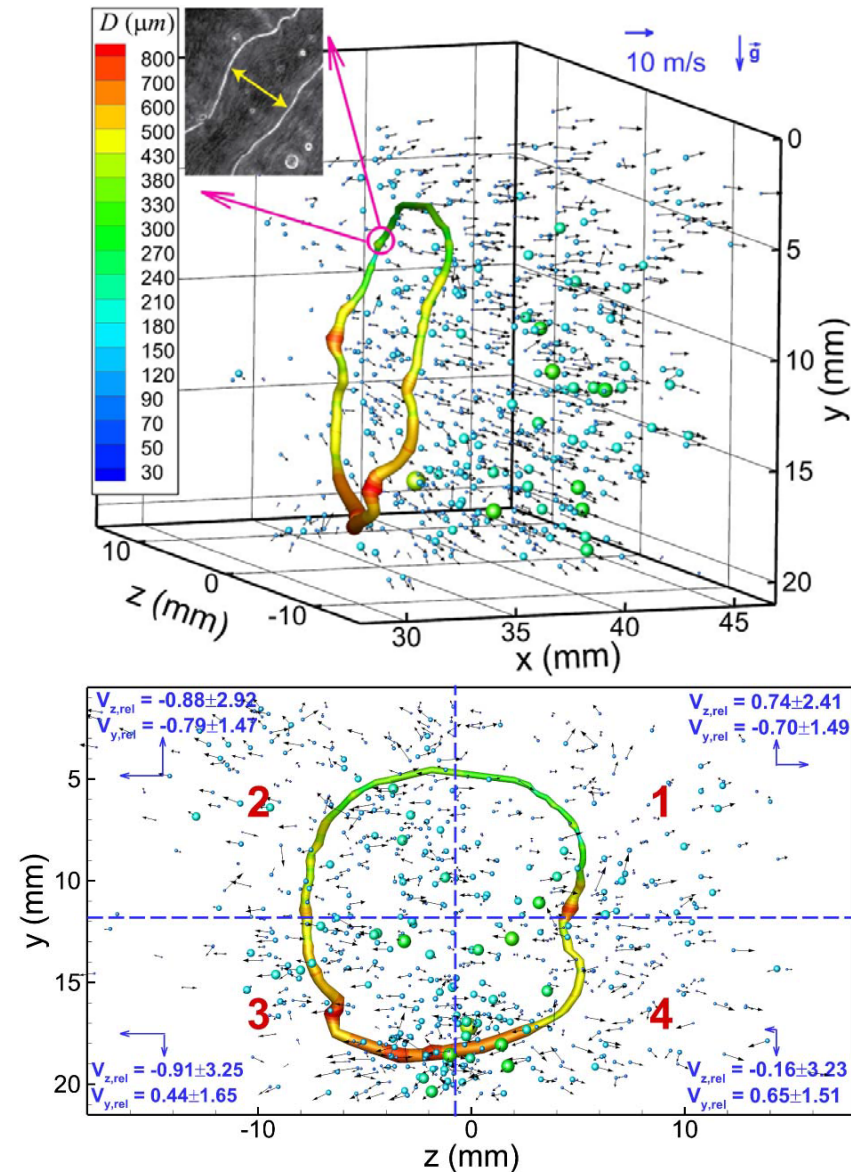
Secondary drop sizes/positions extracted by the hybrid method

- Comparison with phase Doppler anemometer (PDA) data confirms accuracy of measured sizes



Ring measured from z-location of maximum image gradient

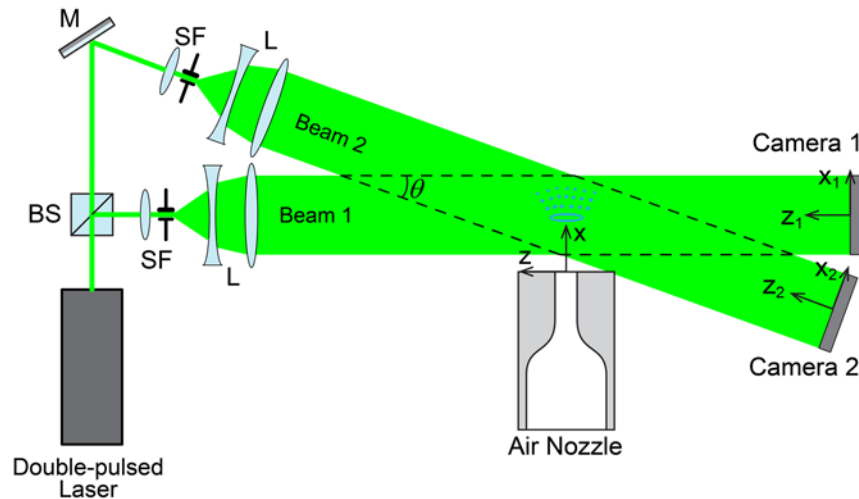
- Total volume of ring + secondary drops is within 2.2% of the initial volume



Aerodynamic drop fragmentation

Velocimetry suffers from uncertainty in the out-of-plane (z) position

- A stereo-view configuration is one solution

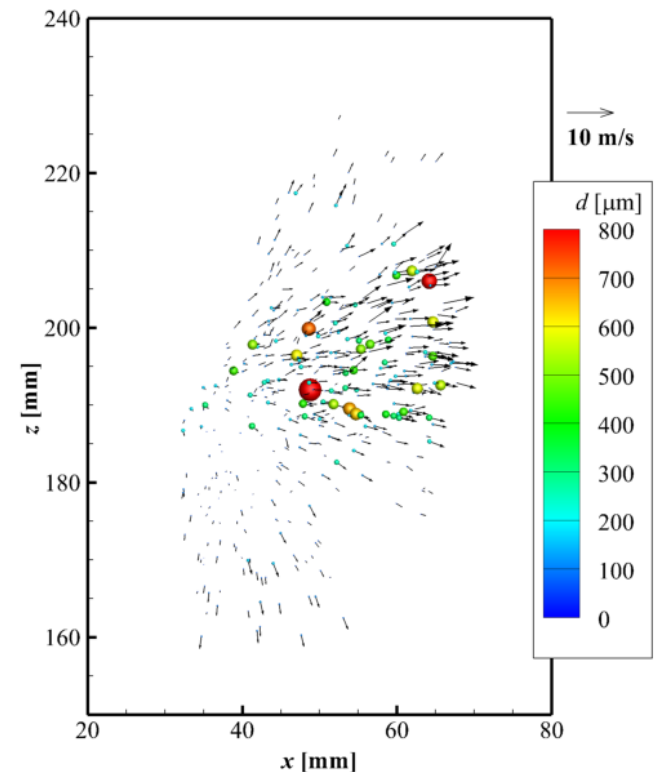
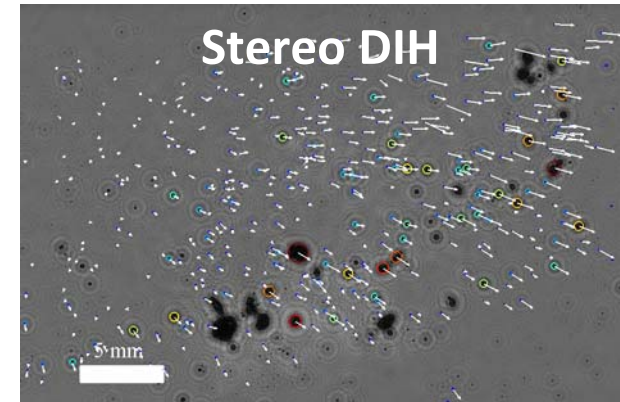


Advantages:

- Improved z -uncertainty
- Eliminates false particle size and position measurements

Challenges:

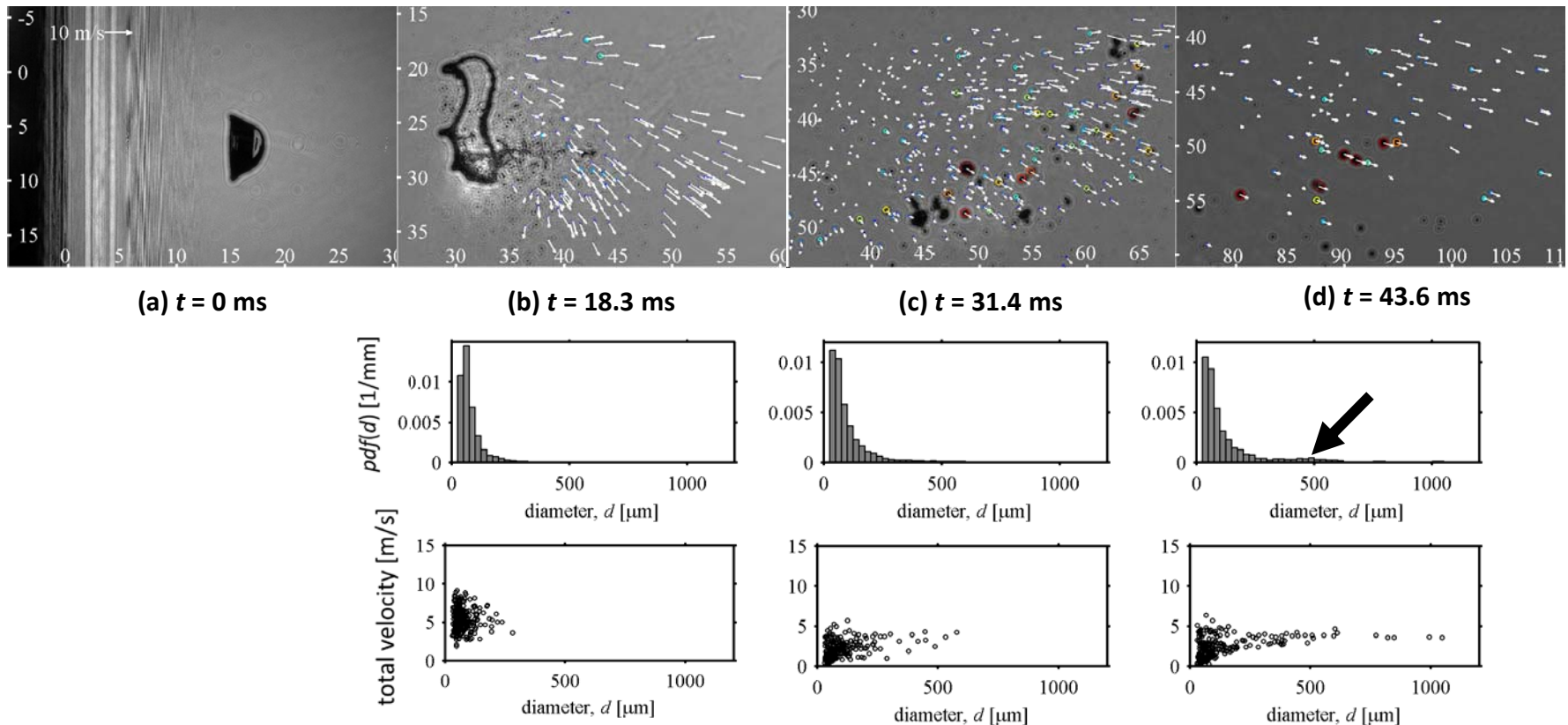
- Increased experimental complexity
- Careful calibration required



Aerodynamic drop fragmentation

Ensemble averaging of 44 realizations at each condition

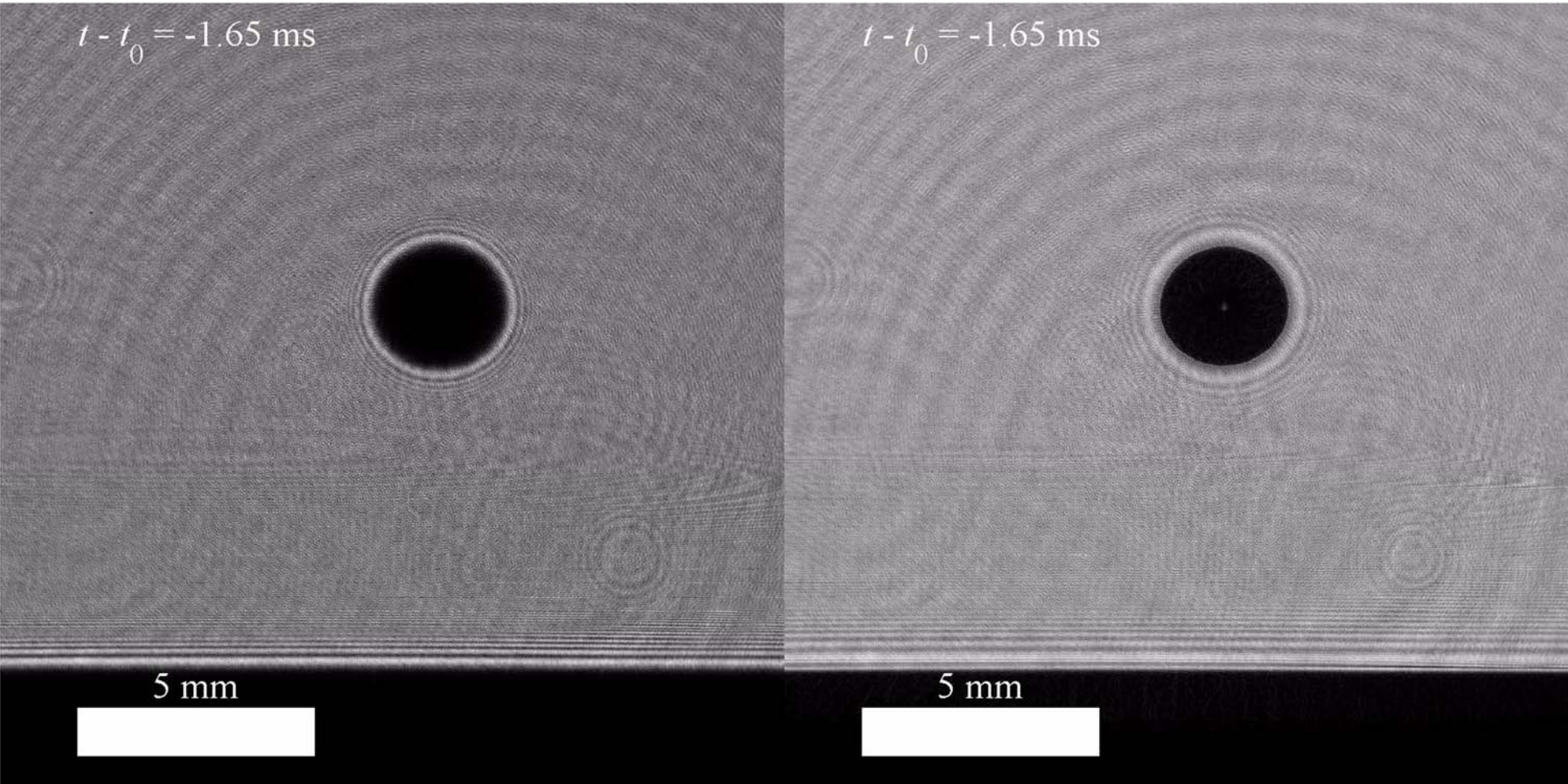
- Roughly 10,000 individual drops measured per condition



DIH is particularly advantageous for rapid quantification of particle statistics

High-speed (kHz) DIH

Increased temporal resolution is possible using high-speed (kHz rate) cameras



Challenges: (1) higher readout noise, fewer pixels, larger pixel pitches
(2) very large data sets (10s of Gb)

High-speed (kHz) DIH

Processing of a single hologram can take roughly 30 min on a typical CPU

- Much of that time spent on numerical refocusing:

$$E(x, y, z) = FFT^{-1} \left[FFT[h(\xi, \eta)] \cdot G(f_x, f_y, z) \right]$$

- Refocusing to a single depth, z , requires:

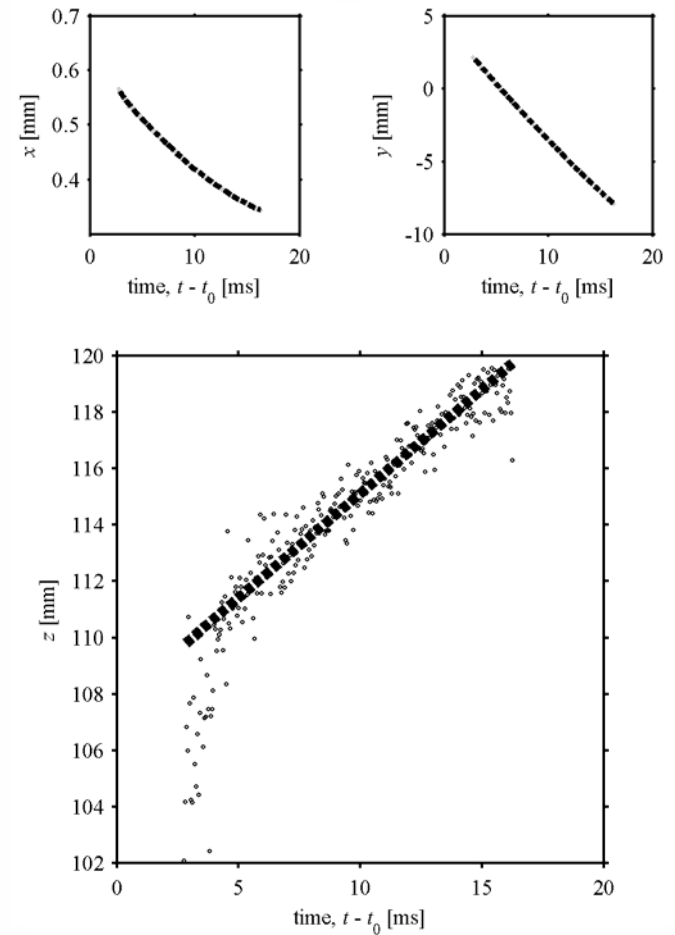
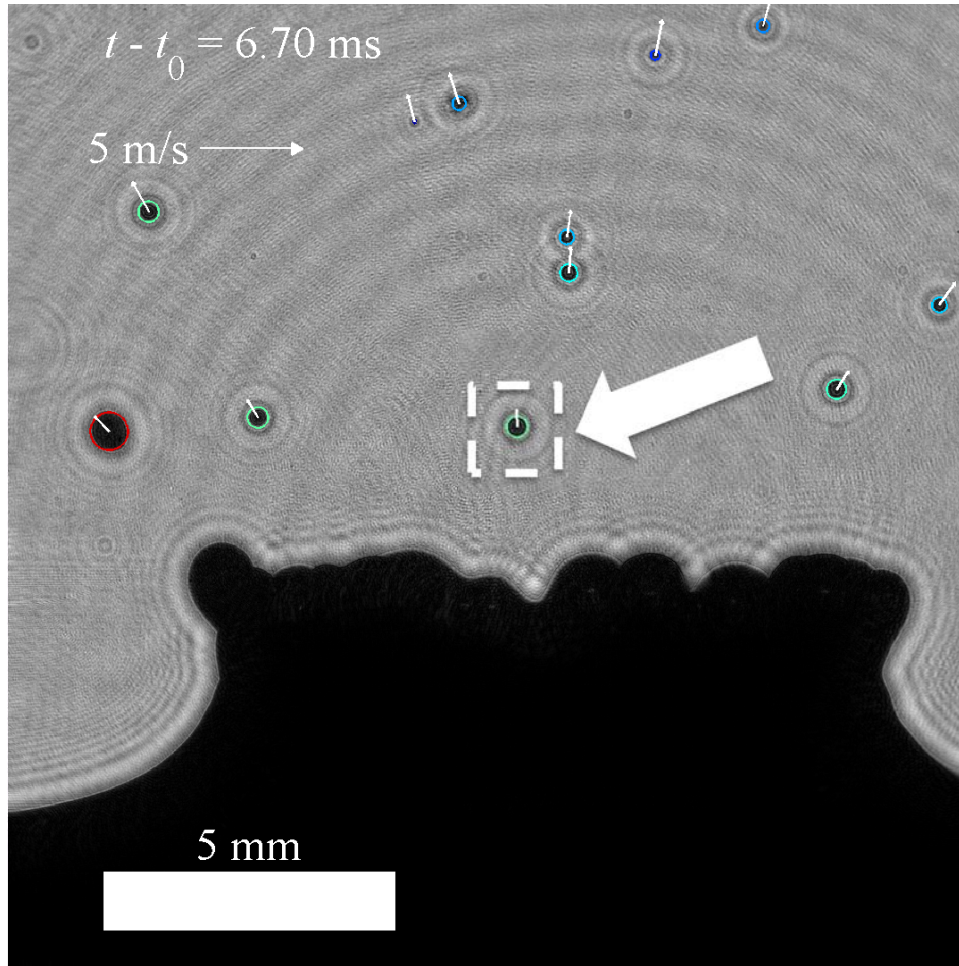
- (1) calculation of $G(f_x, f_y, z) = \exp \left[-jkz \sqrt{1 - \lambda^2 f_x^2 - \lambda^2 f_y^2} \right]$
- (2) multiplication of two large arrays, $FFT[h(\xi, \eta)] \cdot G(f_x, f_y, z)$
- (3) a two-dimensional inverse FFT

Graphical processing units (GPUs) are well suited to these tasks

- E.g. NVIDIA Tesla K40 GPU, Dual Xeon CPU, Matlab v2014a with parallel computing toolbox → per-frame processing time of ~7 seconds

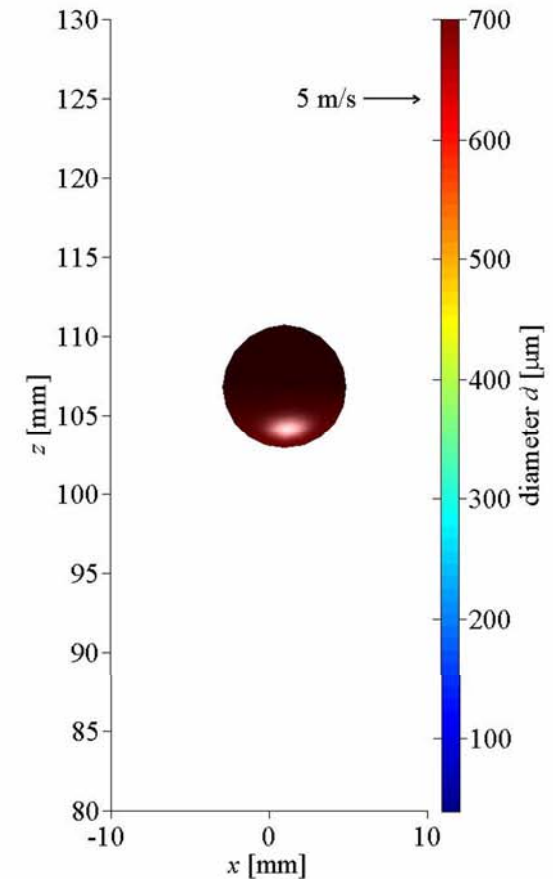
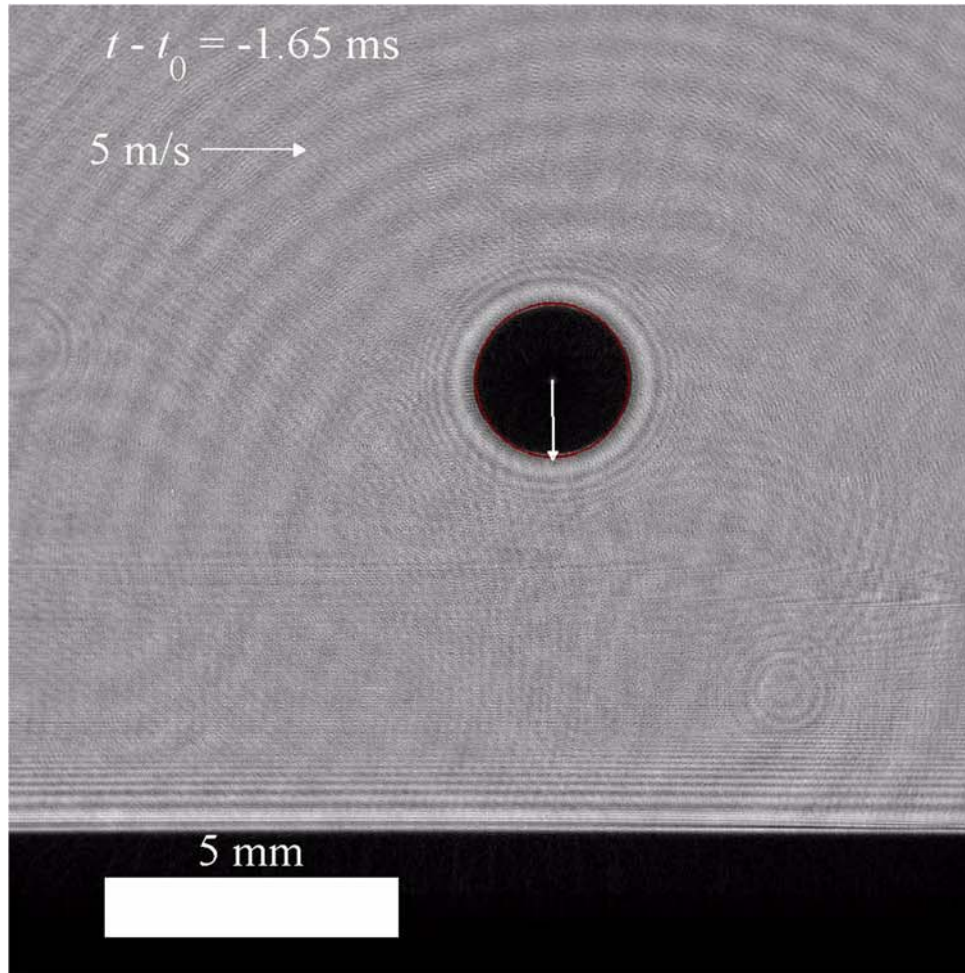


High-speed (kHz) DIH



- Frame-to-frame particle matching illustrates the depth-of-focus problem
- With sufficient temporal resolution, particles trajectories can be fit to temporal models

High-speed (kHz) DIH



- Multi-frame trajectory fitting leads to a 36X reduction in z -uncertainty

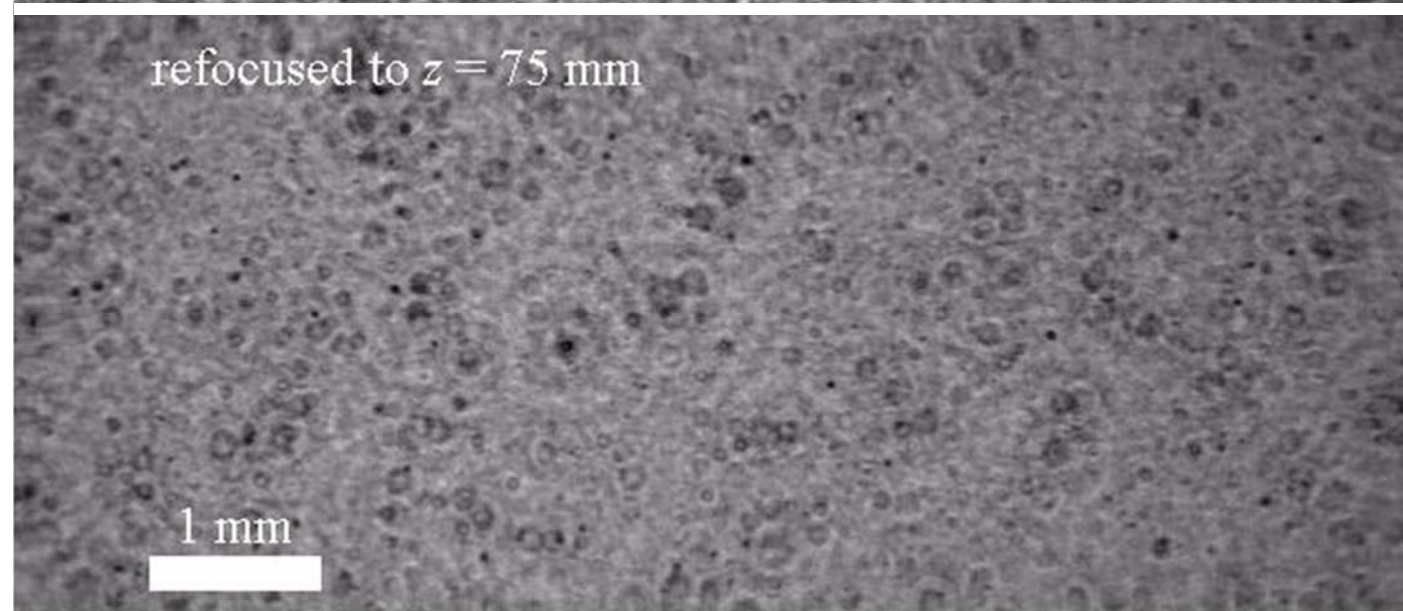
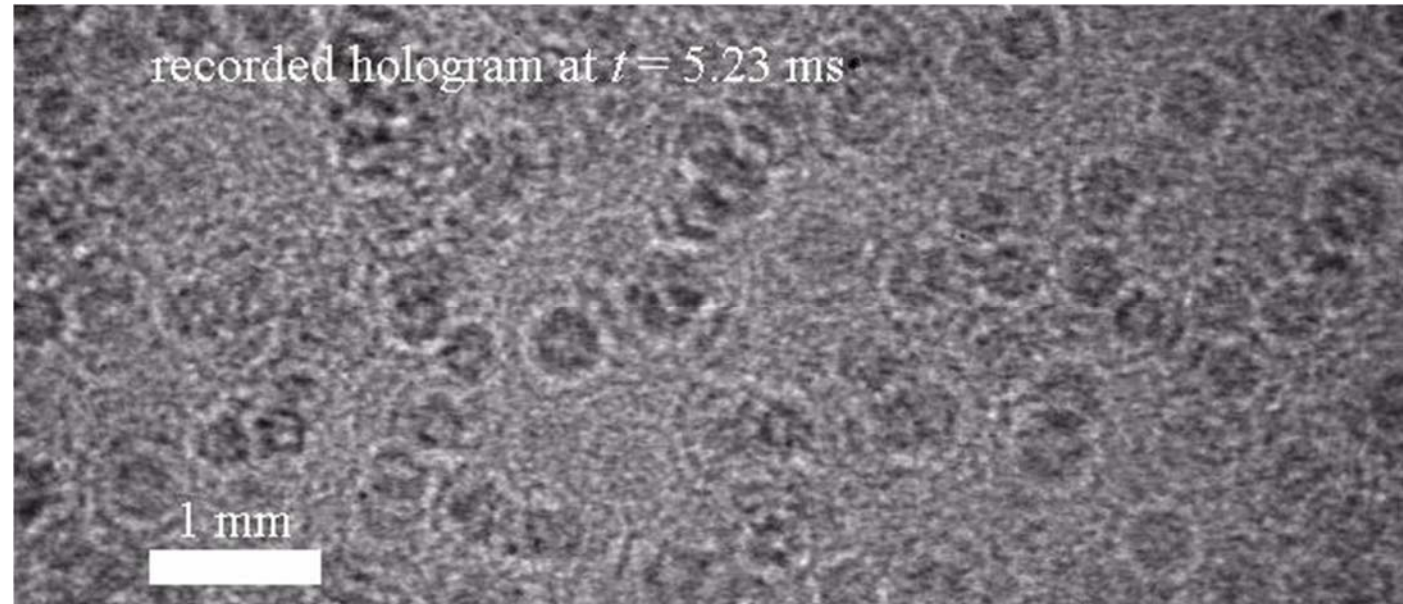
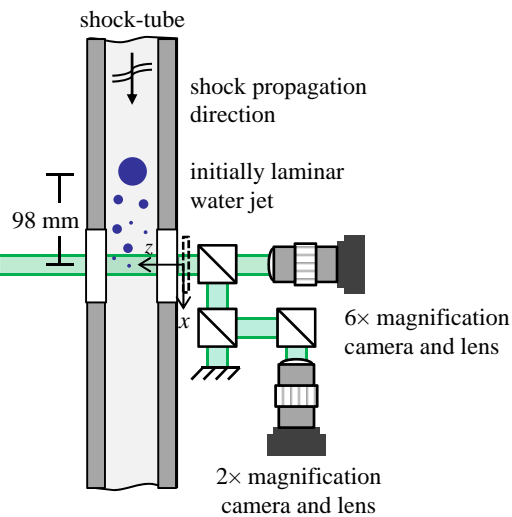
Breakup of a water jet in a shock-tube



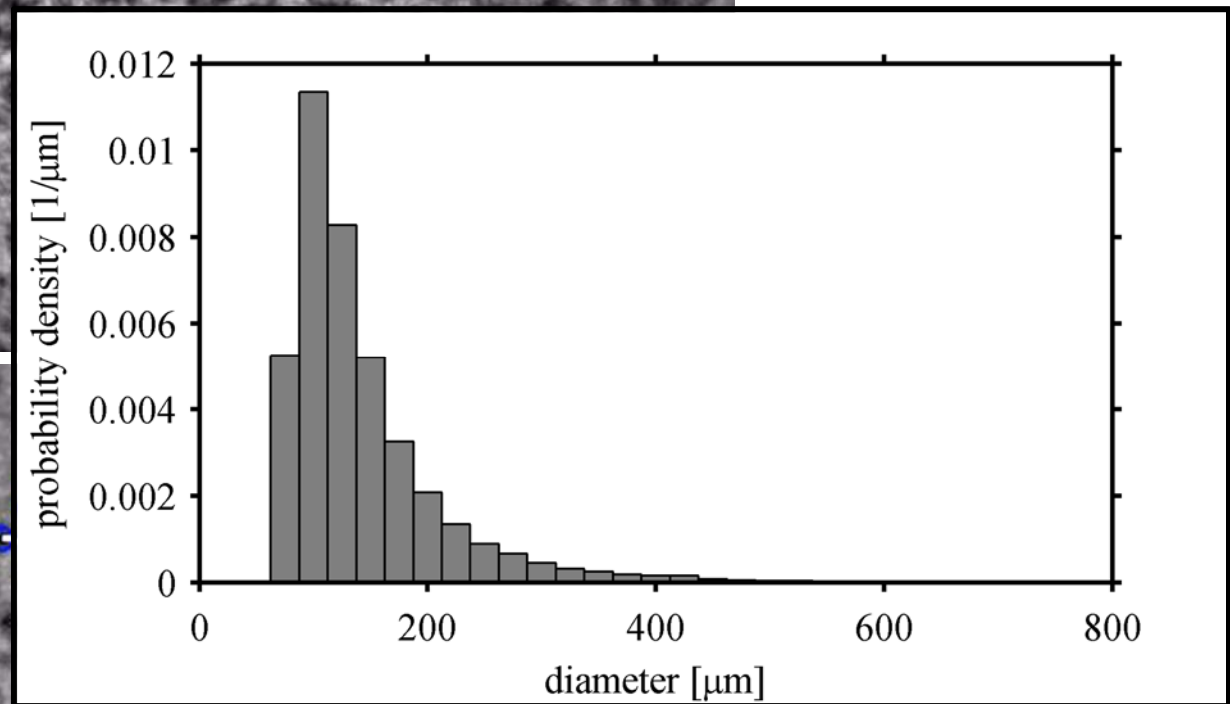
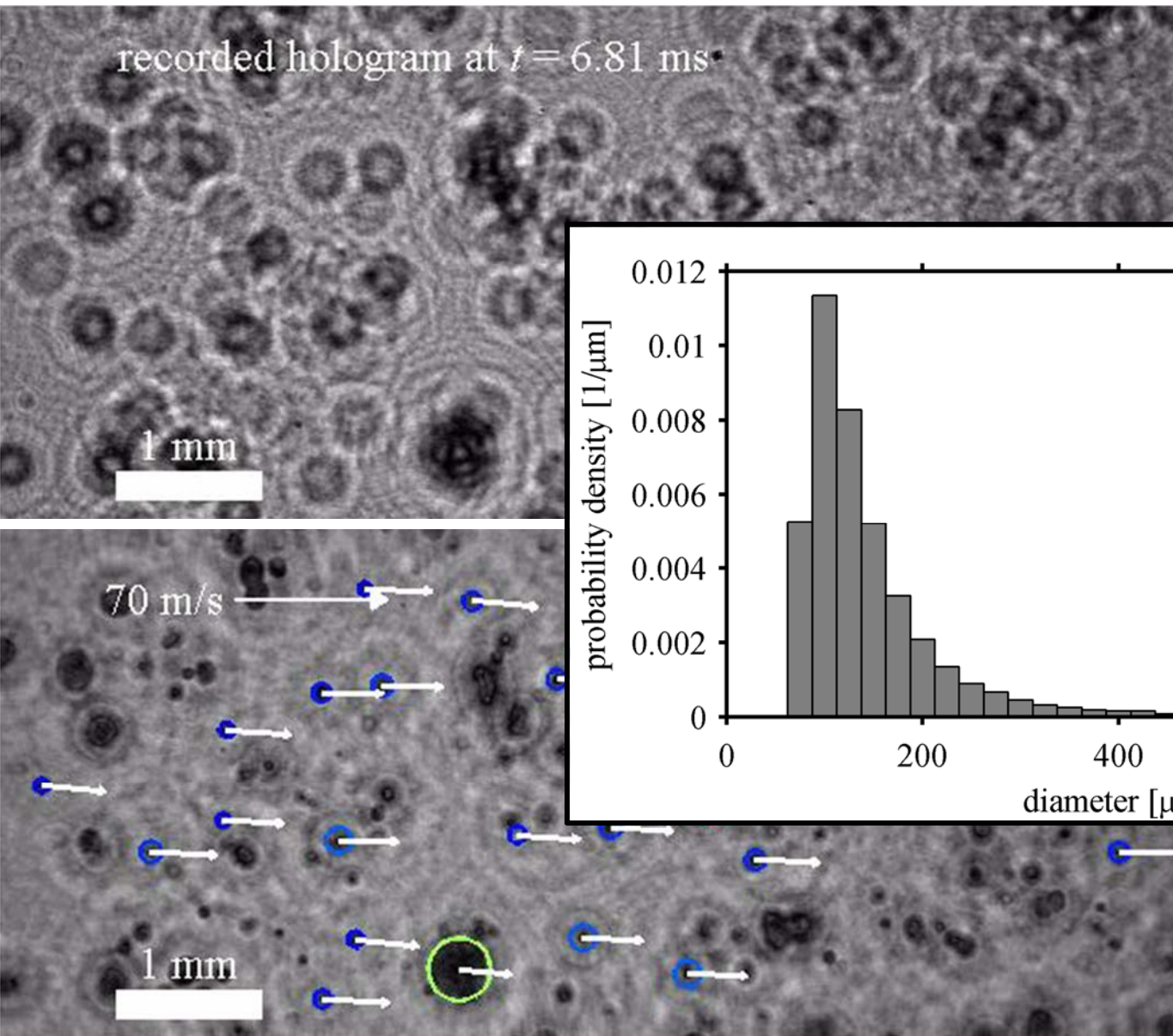
Goals:

1. Quantify the fragment sizes and velocities as a function of shock strength
2. Investigate the relation between surface instabilities and fragment properties

Breakup of a water jet in a shock-tube

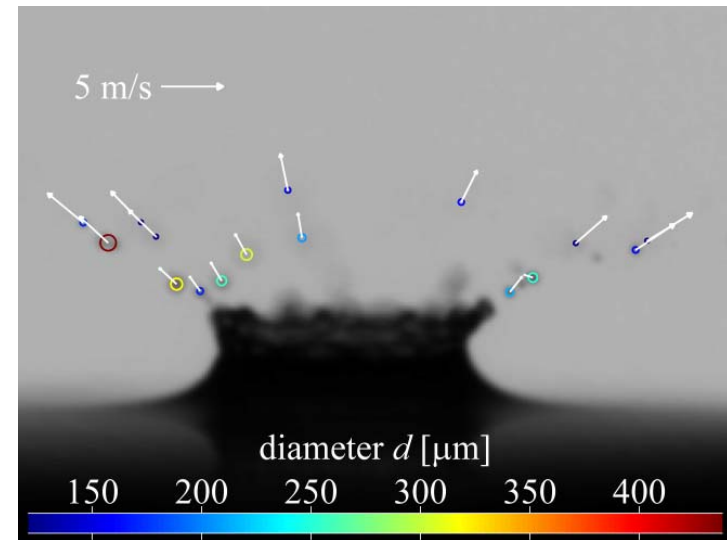
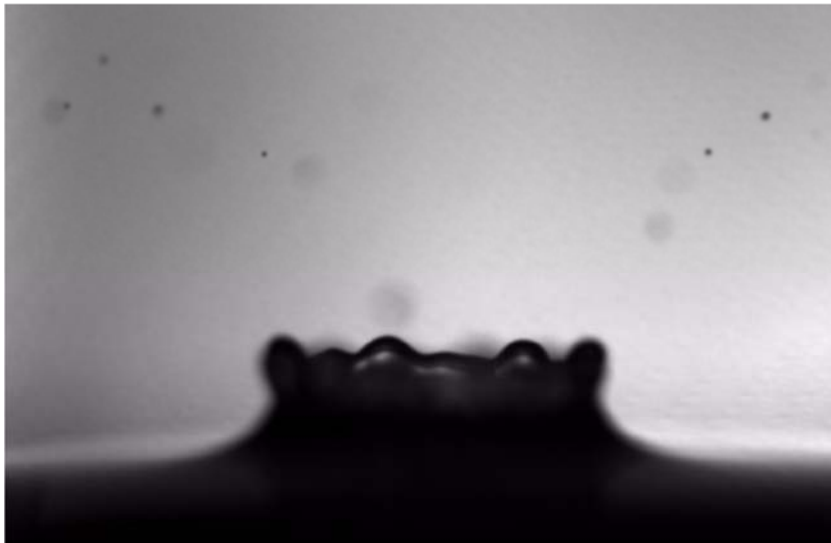
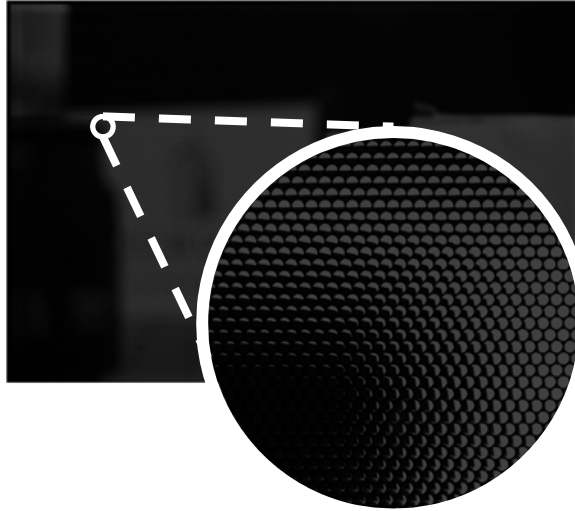


Breakup of a water jet in a shock-tube



Alternative 3D measurements

Plenoptic cameras use micro-lens arrays and white light to create a 3D image



Thoughts on applications to fuel sprays

DIH has many advantages:

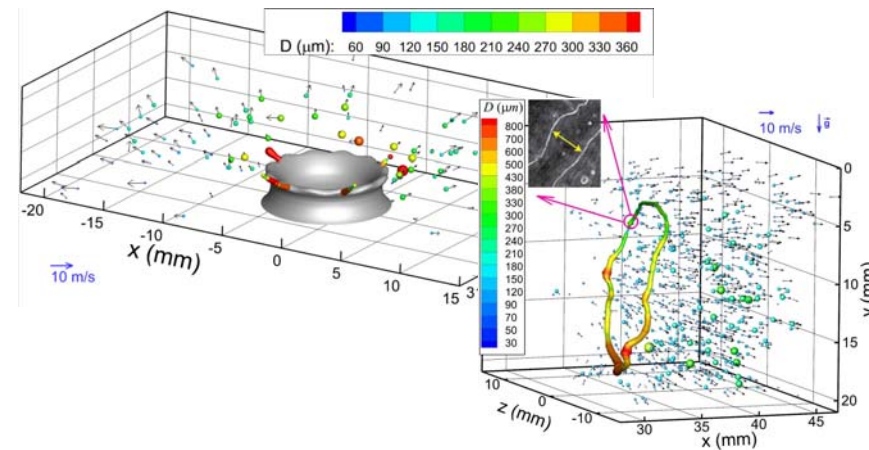
- Simple optical configuration
- 3D-3C measurement
- Rapid quantification of statistics
- Captures details of transient events

... and challenges:

- Depth-of-focus problem
- Data processing
- Small field of view
- Limited to dilute sprays

For modeling of liquid fuel sprays
DIH/plenoptic imaging could provide:

- Detailed particle statistics of laboratory scale problems which form the basis of fuel spray models (e.g. drop impact, aerodynamic breakup)
- Qualitative, 3D imaging of larger, more realistic phenomena
 - Quantitative imaging may be possible in sub-regions of the flow and/or downstream positions where particle density is reduced

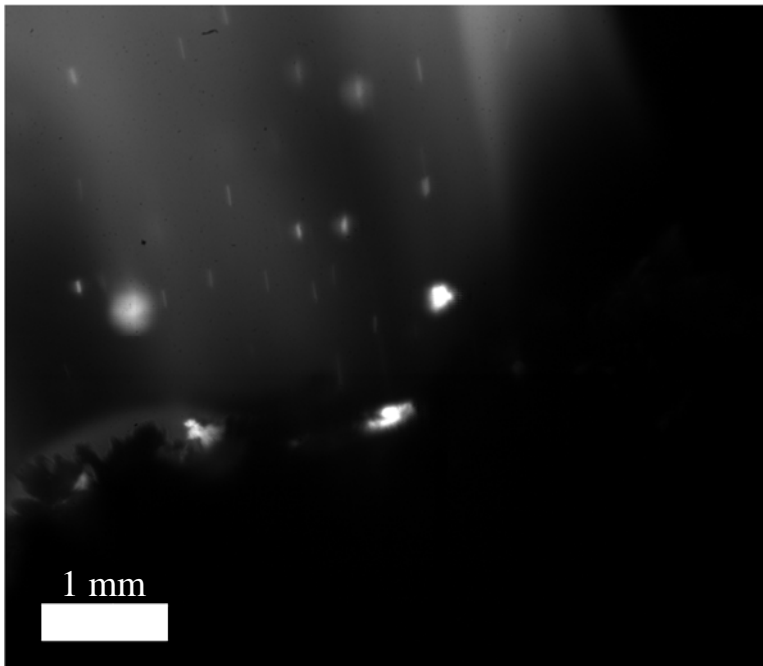


Propellant fire measurements

Aluminum drop combustion in propellants

Motivation: rocket failures can lead to propellant fires

- Sandia Laboratories is interested in predicting the response of objects in this environment



high-speed video of a burning propellant

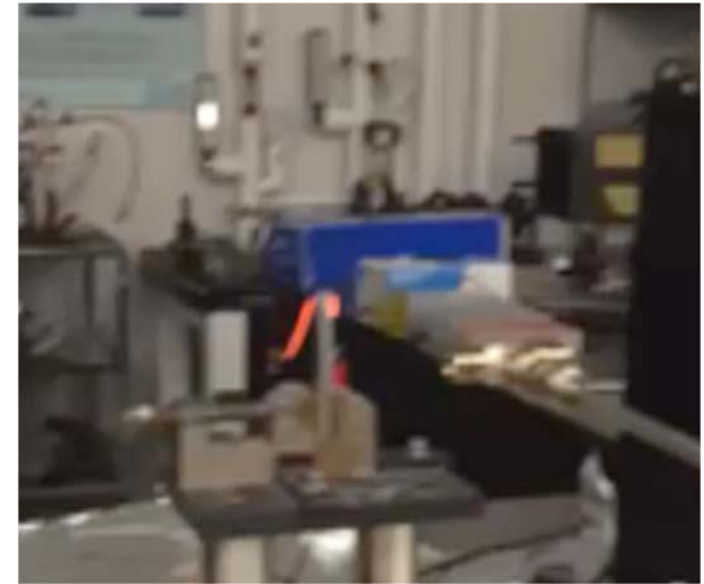
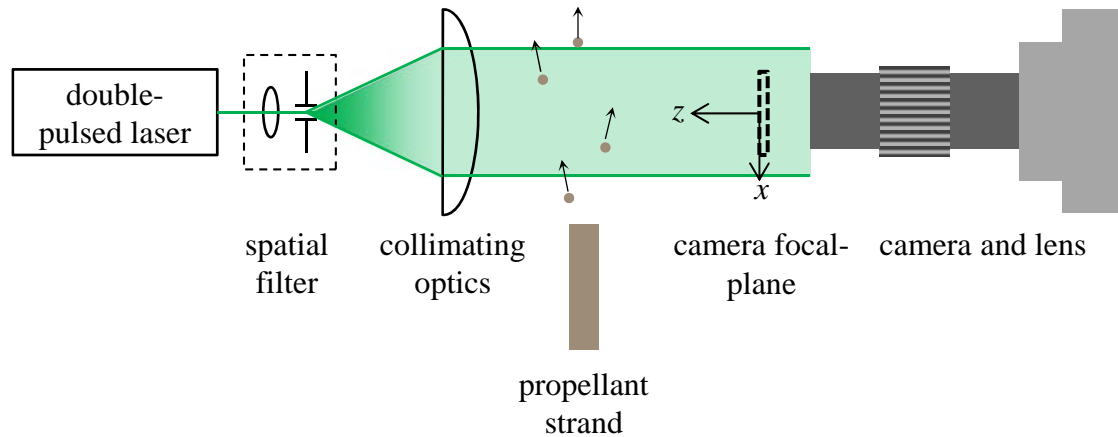


<http://www.cbsnews.com/news/rocket-crash-no-immediate-threat-to-station-but-cause-is-unknown/>

Aluminum agglomeration at the surface yields large reacting drops with high damage potential

- Prediction requires knowledge of particle *size, velocity, and temperature*

Aluminum drop combustion in propellants



propellant in the test fixture

Propellant: solid-rocket propellant pressed into a pencil size strand

- Combusts from the top surface down, ejecting molten aluminum particles traveling a few m/s

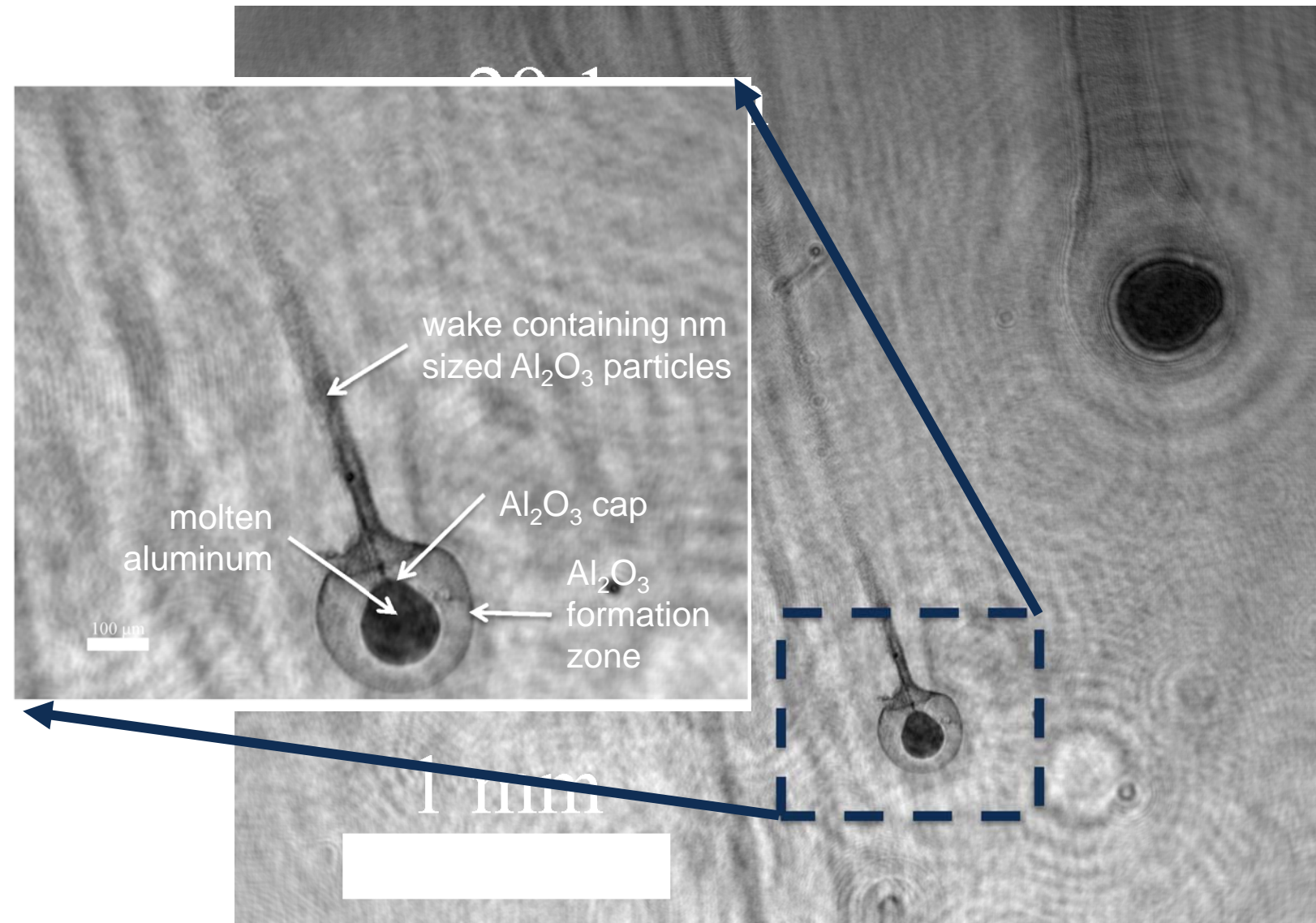
Laser: Continuum Minilite Nd:YAG, 532 nm wavelength, 5 ns pulse duration

Camera: sCMOS from LaVision at 15Hz

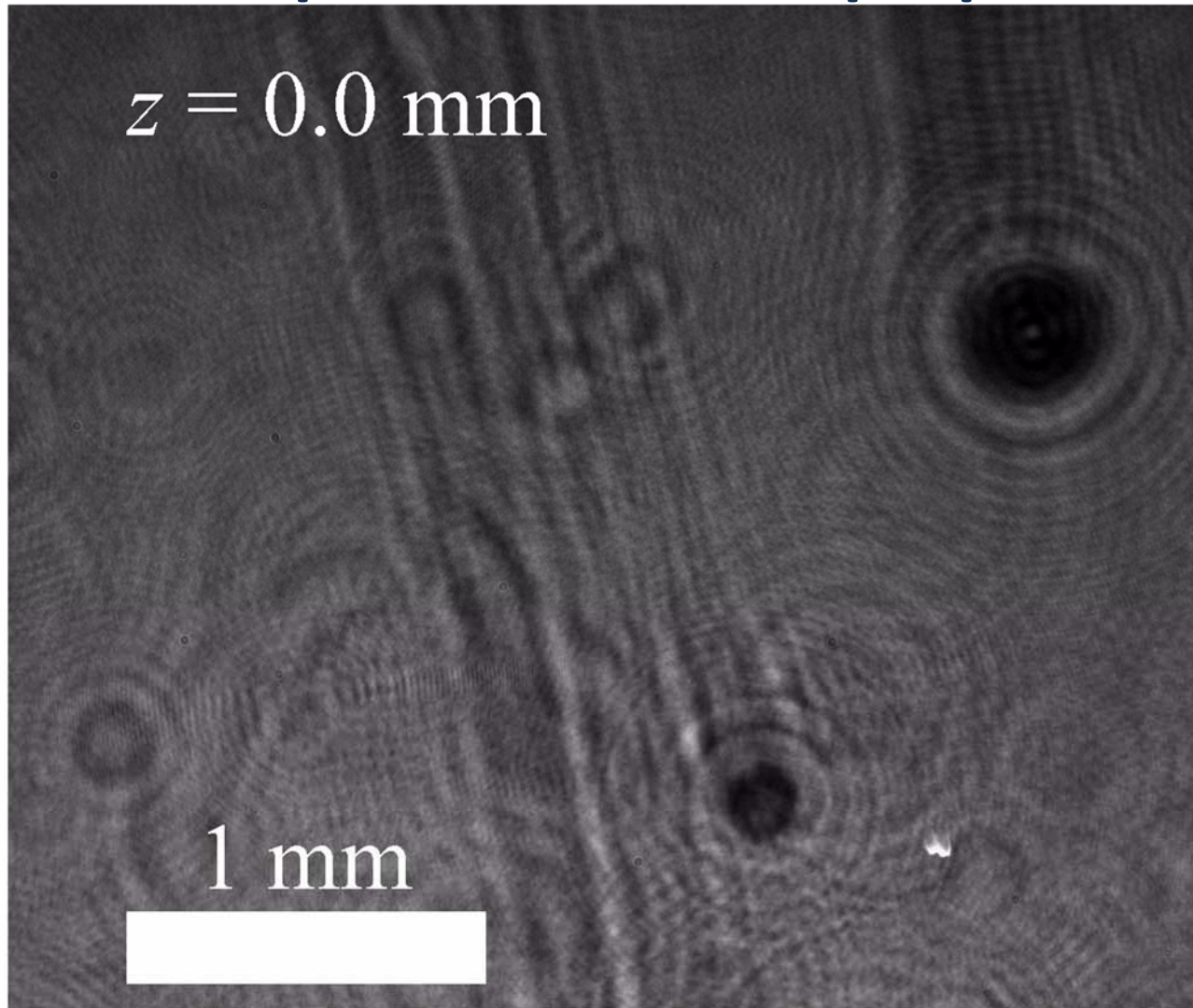
Lens: Infinity K2 long distance microscope with CF-4 objective

- ~ 6X magnification

Aluminum drop combustion in propellants



Aluminum drop combustion in propellants



Algorithms automatically measure unique features of burning aluminum

Aluminum drop combustion in propellants

Three strand burns → 5594 images
and 17496 measured drops

- Main peak due to agglomerated particulates
- Peak at 50 μm due to non-agglomerated particulate

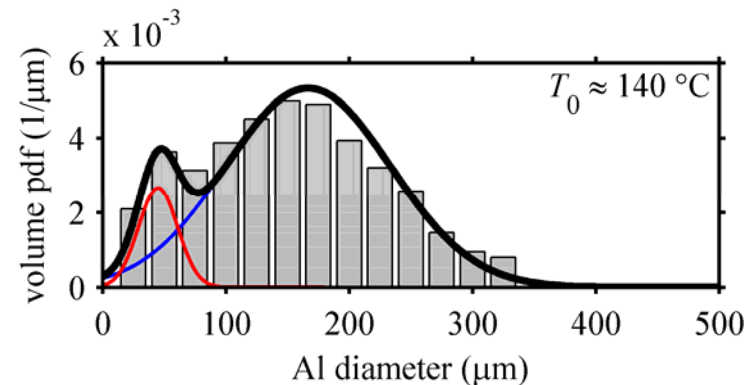
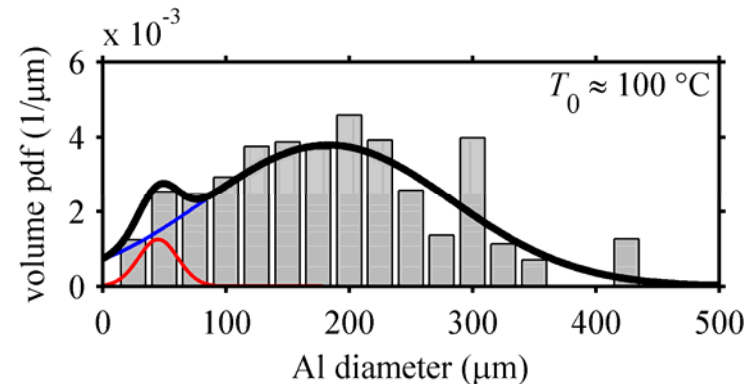
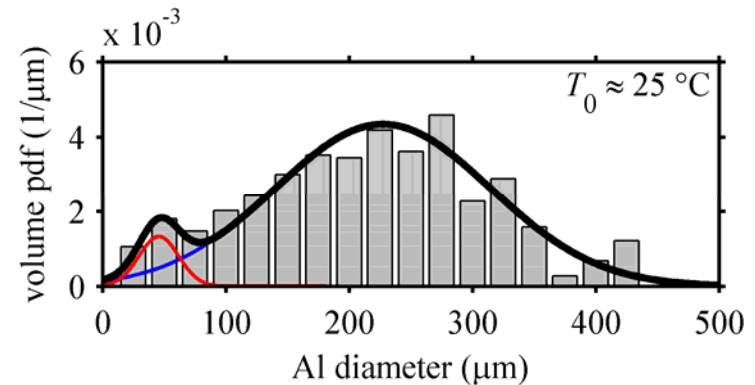
Experiments repeated at higher initial temperature (faster burn rate)

- Main peak is reduced due to decreased residence time for agglomeration

- Peak at 50 μm remains

Trend is consistent at still higher initial temperatures

- Main peak reduced further
- Peak at 50 μm remains

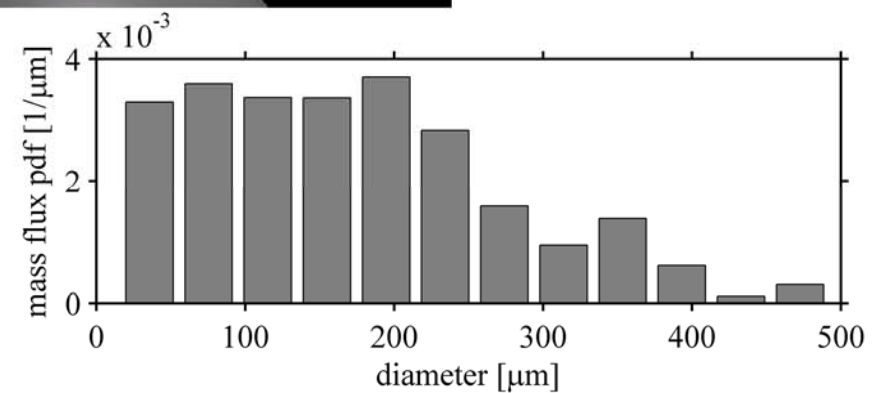
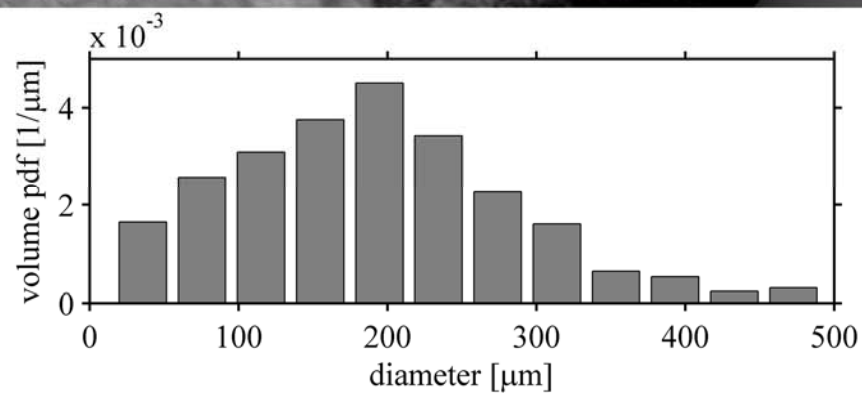
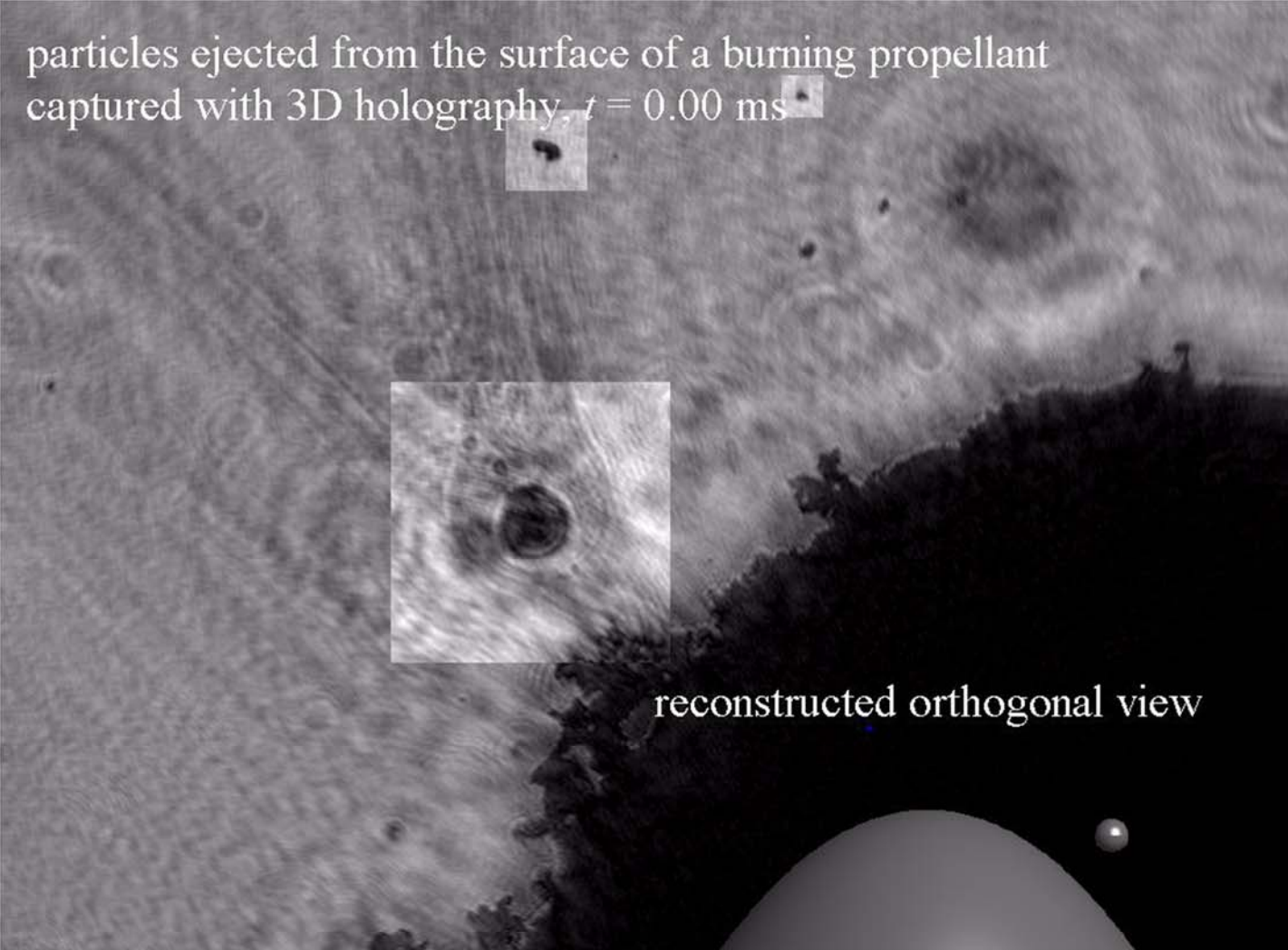


Recorded at
20,000 fps

Camera: Photron
SA-Z

Laser: Coherent
Verdi V6

43,684 frames →
15,991 measured
drops



Acknowledgements:

This work was supported by the Laboratory Directed Research and Development and the Weapons Systems Engineering Assessment Technology program at Sandia National Laboratories (SNL)

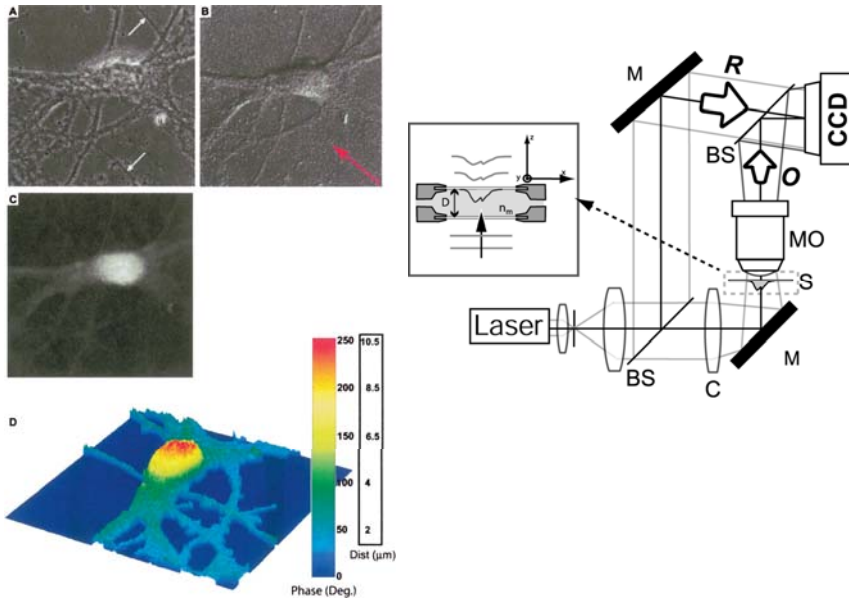
Many thanks to all of my excellent collaborators: *Jian Gao* (Johns Hopkins University), *Phillip L. Reu* (SNL), *Jun Chen* (Purdue University), *Sean P. Kearney* (SNL), *Kathryn G. Hoffmeister* (SNL), *Paul E. Sojka* (Purdue University), *Thomas W. Grasser* (SNL), *H. Lee Stauffacher* (SNL), *Marcia A. Cooper* (SNL), *Luke Engvall* (University of Colorado), *Justin L. Wager* (SNL), *Thomas A. Reichardt* (SNL), *Paul A. Farias* (SNL), and many others....

Questions

Backup slides

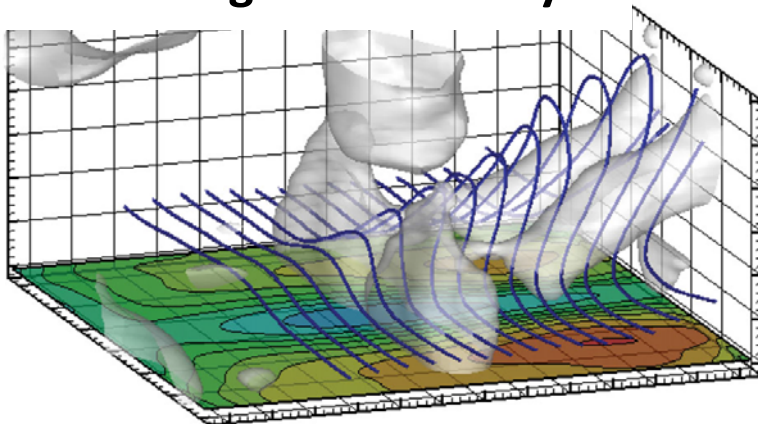
DIH in the literature

Microscopy



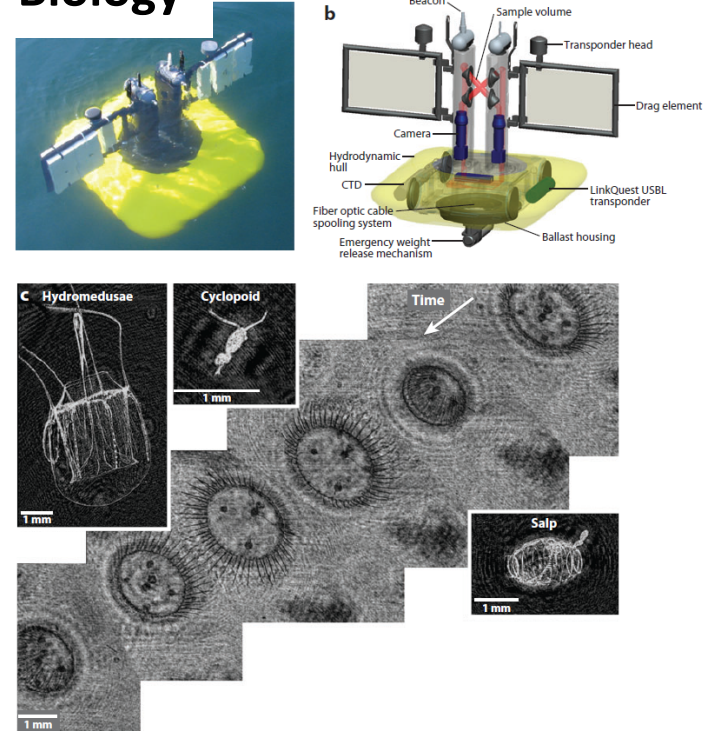
Marquet et al 2005, *Opt. Lett.*

Particle Image Velocimetry



Sheng et al 2009, *J. Fluid Mech.*

Biology

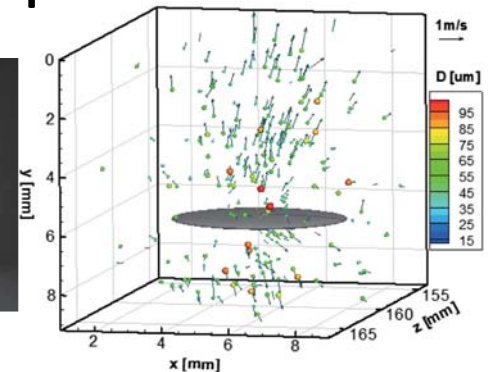


Katz and Sheng 2010, *Annu. Rev. Fluid Mech.*

Multiphase Flows



Yao et al 2015, *Appl Opt.*

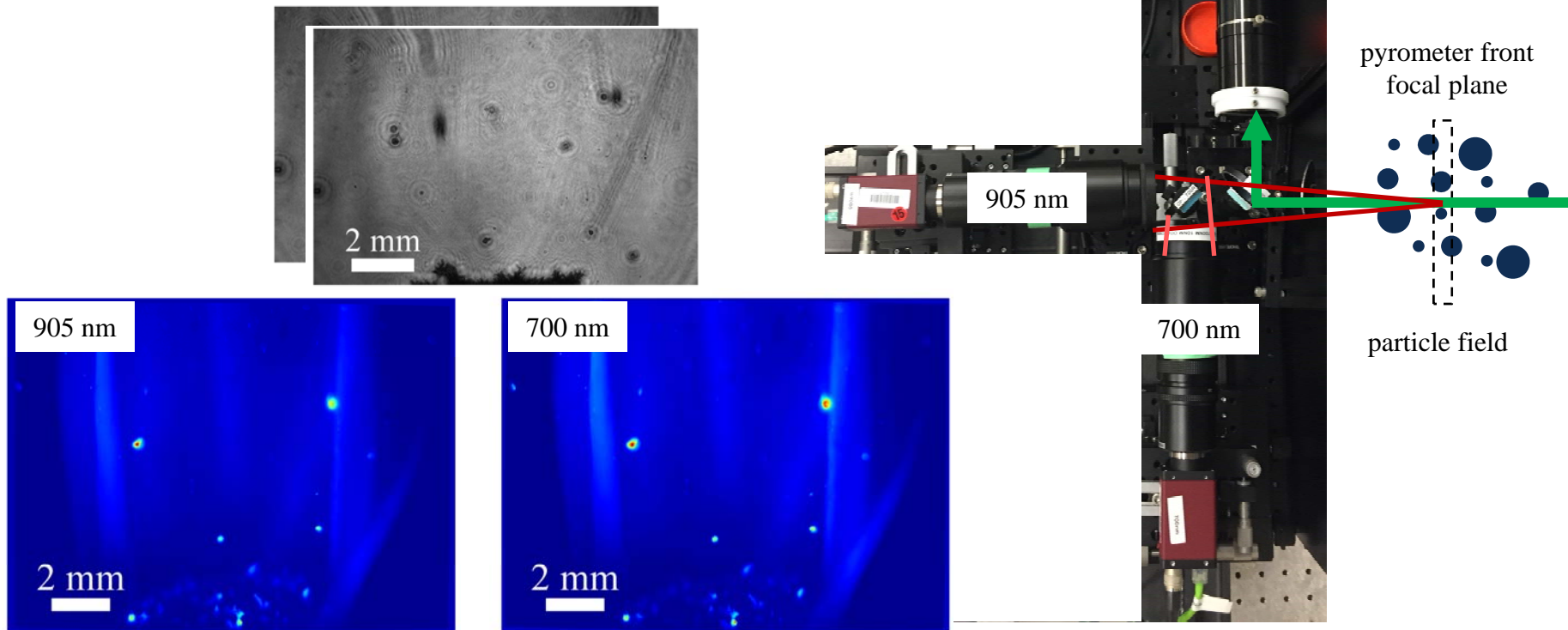


Aluminum drop combustion in propellants

DIH gives mass transfer (particle size + velocity)

We really need to quantify the heat transfer (particle and gas phase)

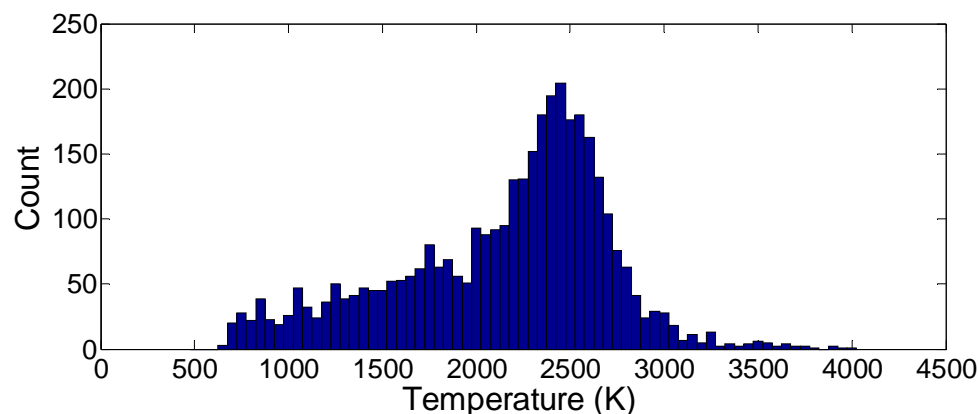
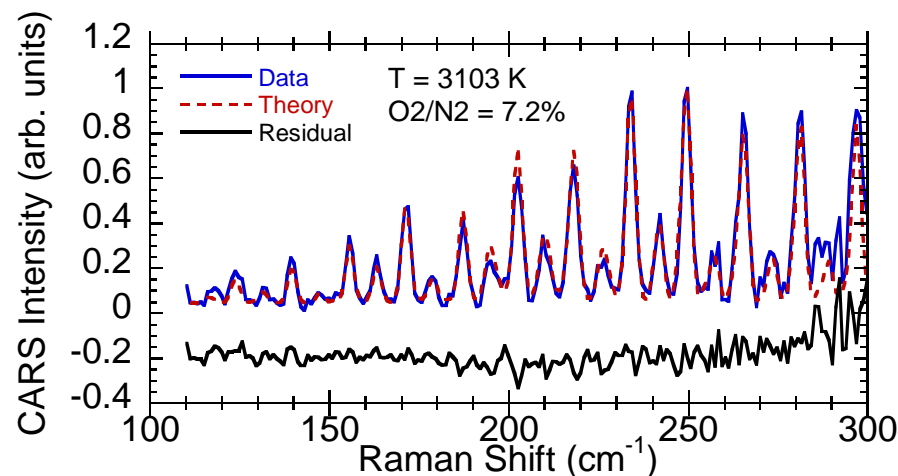
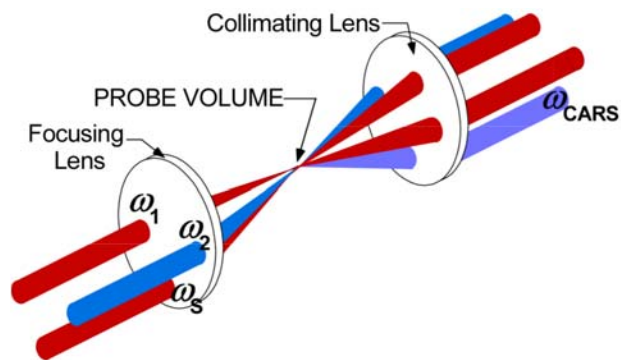
- Combination of DIH and two-color pyrometry → particle size + velocity + *temperature*



Aluminum drop combustion in propellants

Gas phase temperature can be measured using fs/ps CARS

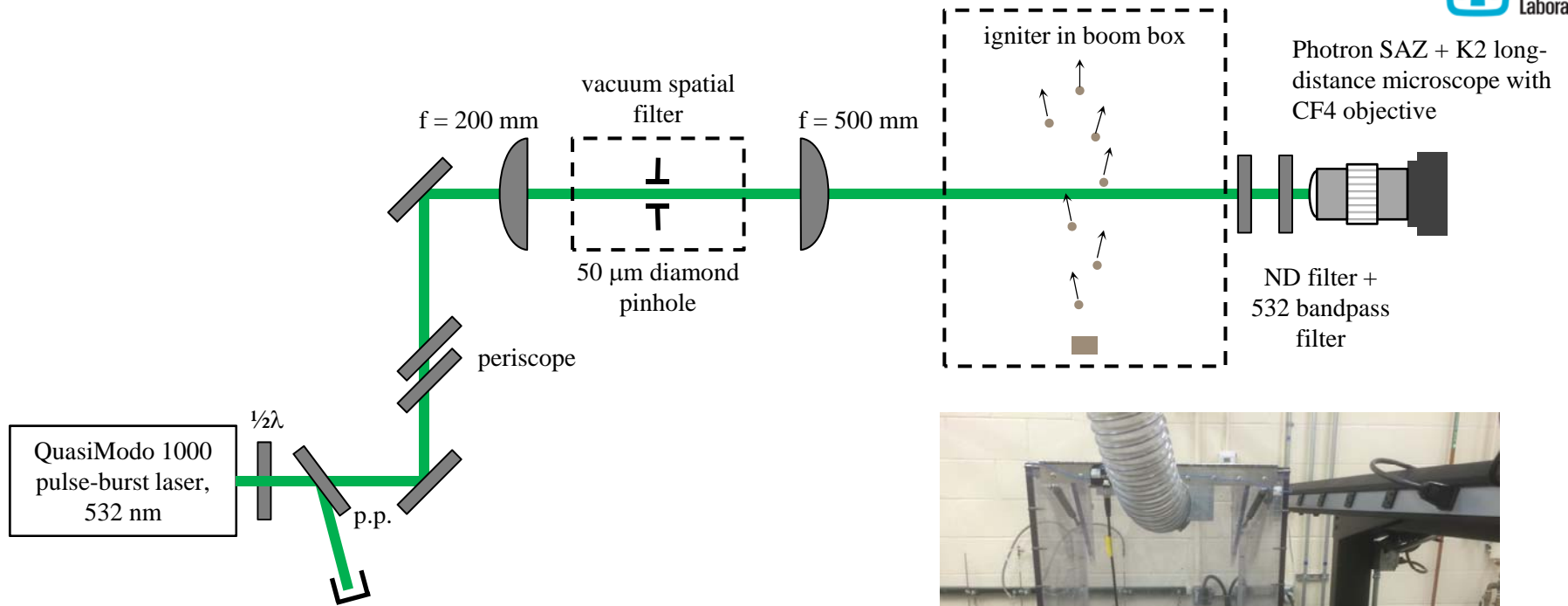
- Advantages compared to ns CARS:
 - Low (mJ) pulse energies → reduces dielectric breakdown
 - Time-delayed probe → eliminates background signal
 - Enhanced precision ~ 1%



See poster: *Hybrid fs/ps CARS for sooting and metalized flames* by Kathryn Hoffmeister, Sean Kearney, Daniel Guildenbecher, and Caroline Winters

New concepts and opportunities

Pulse-burst DIH

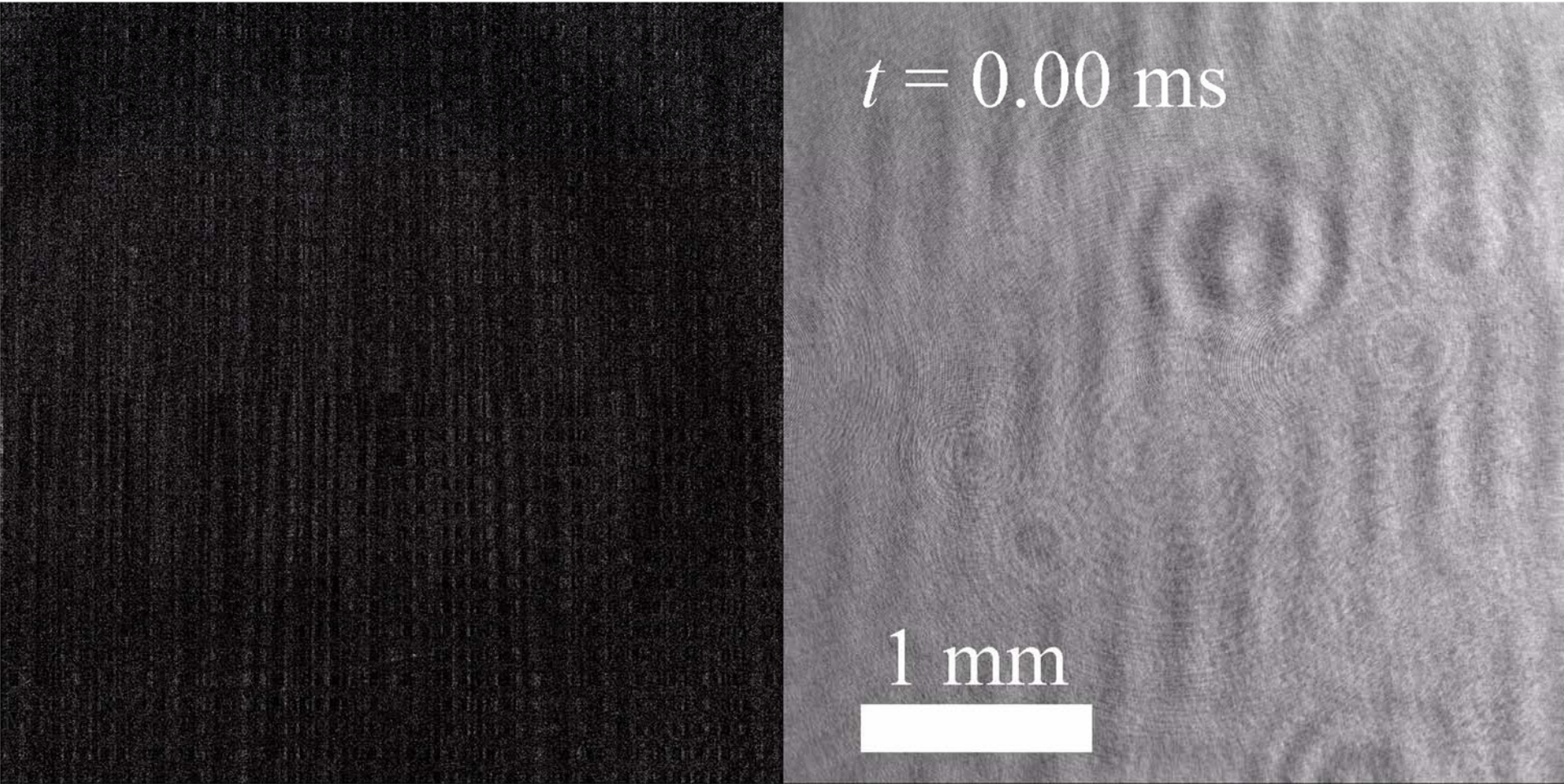


Spectral Energies pulse-burst laser



boom-box and high-speed DIH imager

Pulse-burst DIH

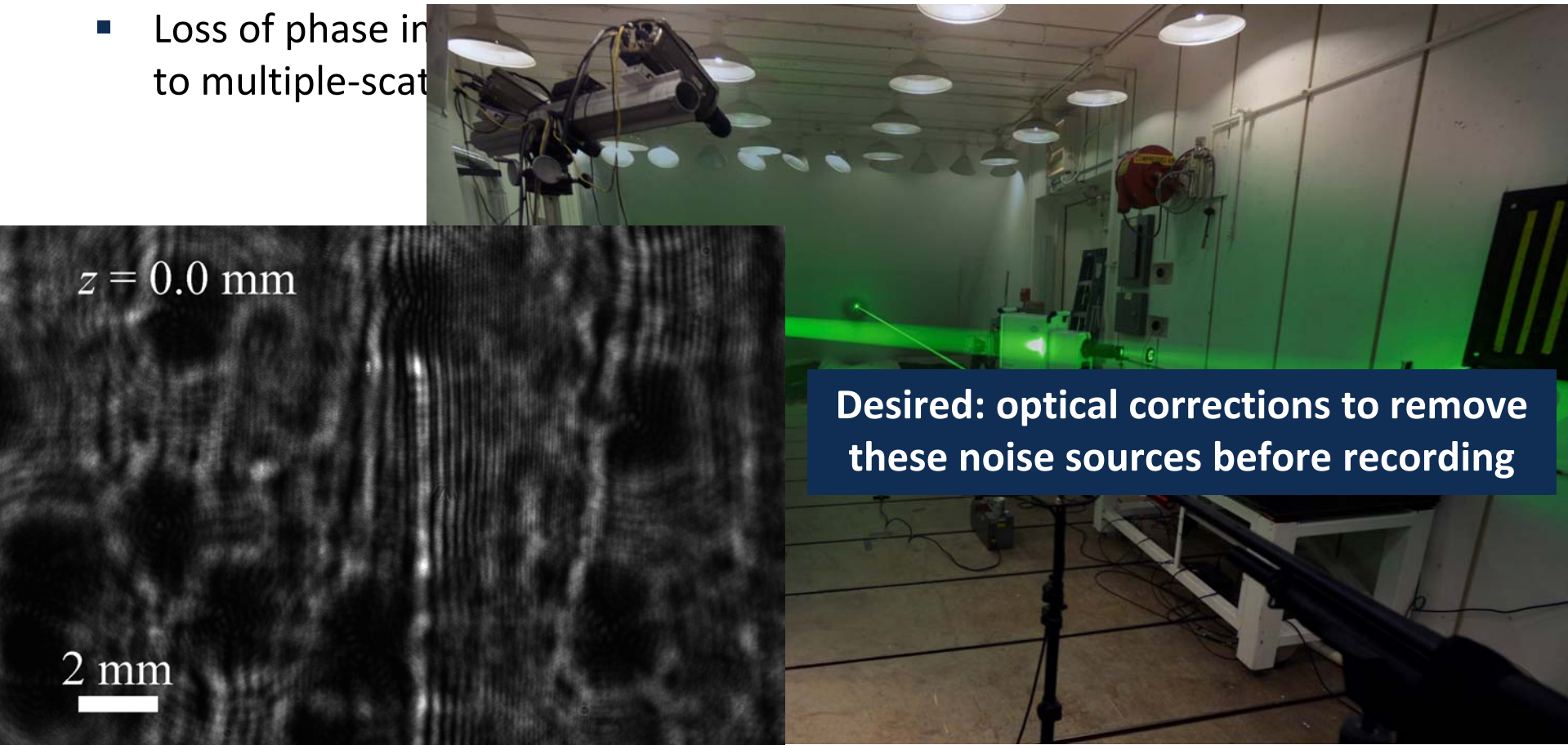


- Beam quality is sufficient for DIH
- Freezes high-speed particles and penetrates through flash and smoke
- Noise due to soot and index-of-refraction gradients

Optical challenges in DIH

Coherent imaging is susceptible to:

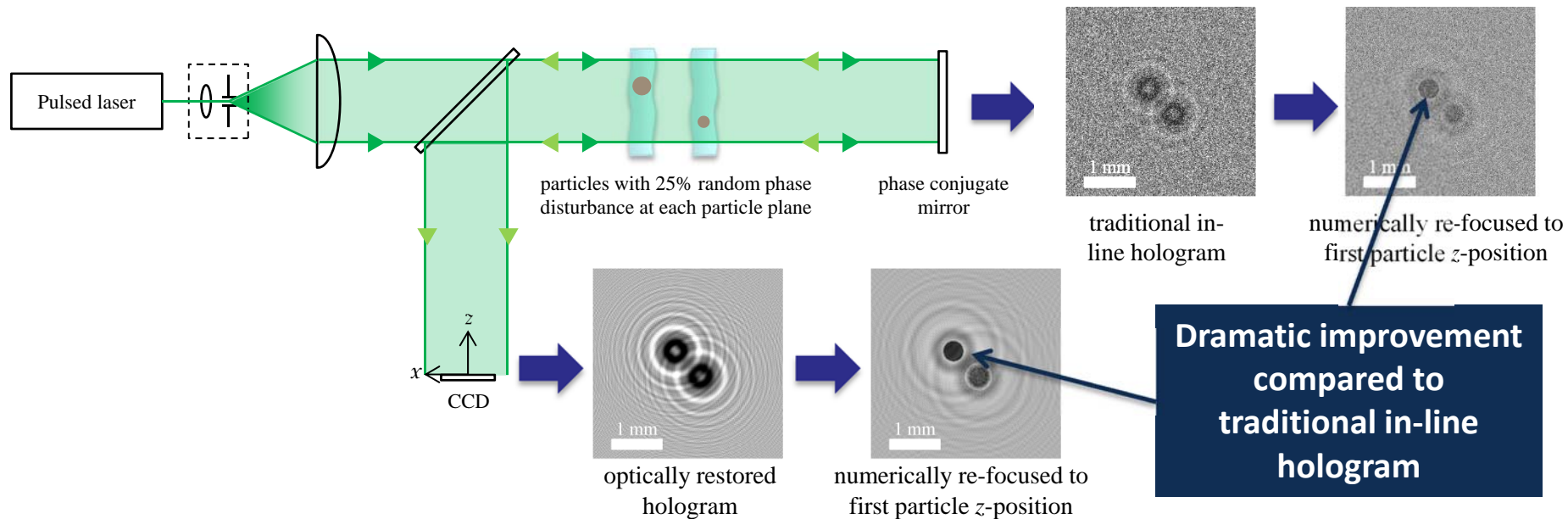
- Image distortion through index of refraction gradients
- Loss of phase information due to multiple-scattering



Reconstructed amplitude throughout depth, z

Holography configuration for shotgun investigations

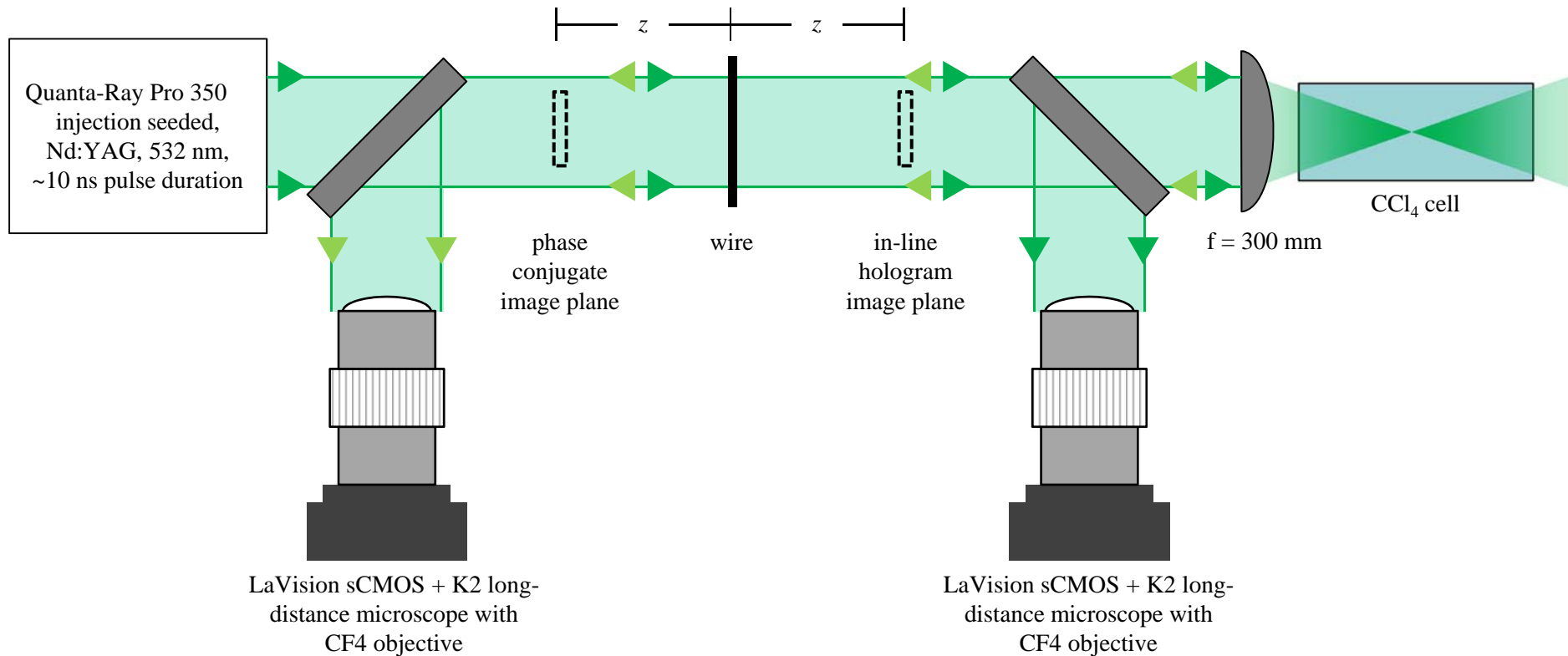
Phase-conjugate DIH theory



- Phase-conjugate mirror reflects the incoming wave with opposite phase
 - Non-linear optical effect achieved through passive means (stimulated Brillouin scattering) or active means (degenerate four-wave mixing)
- After double passing, the phase disturbance is canceled

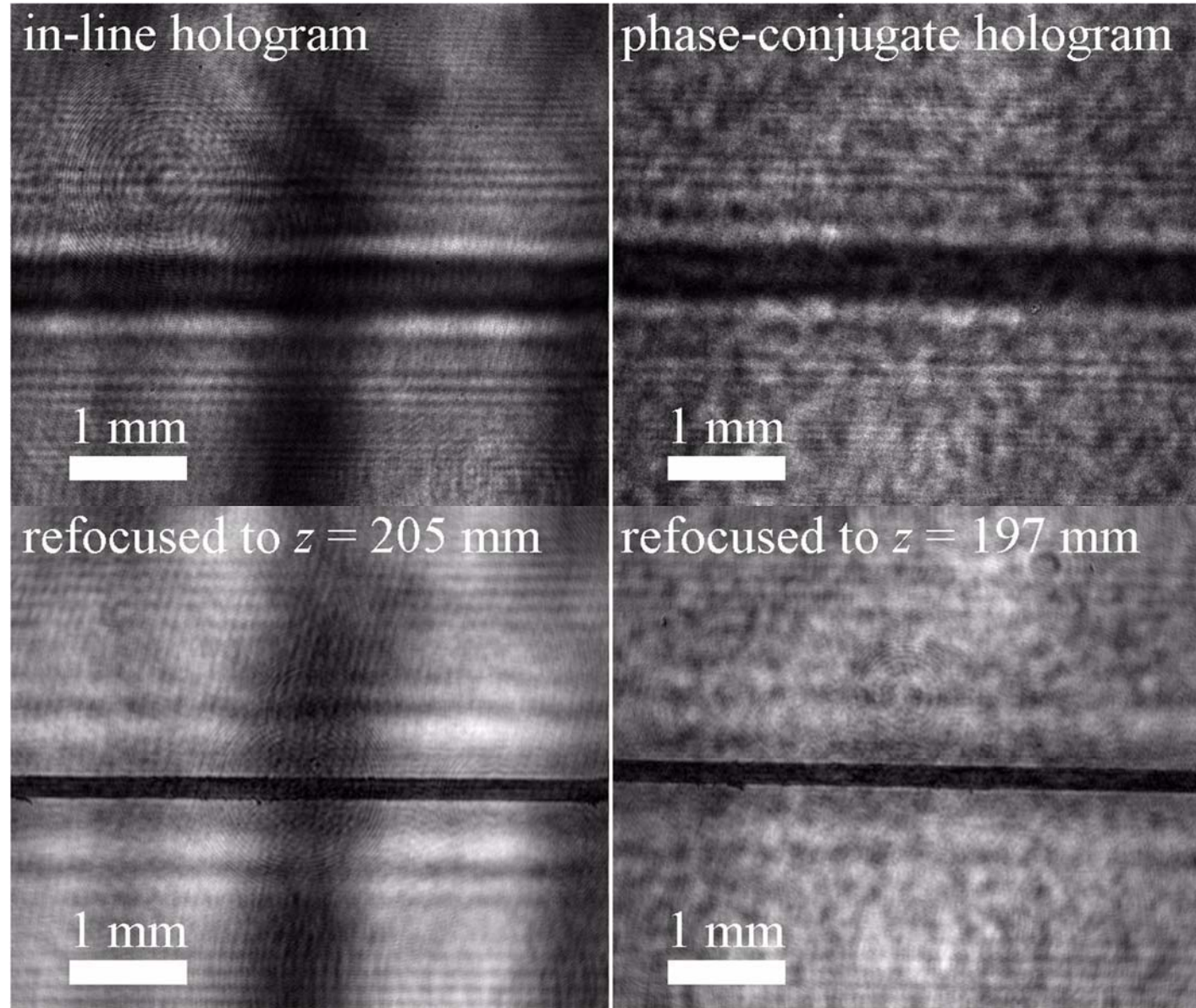
SBS phase-conjugate DIH

A focused beam in a non-linear medium induces phase conjugation via stimulated Brillouin scattering (SBS)



SBS phase-conjugate DIH

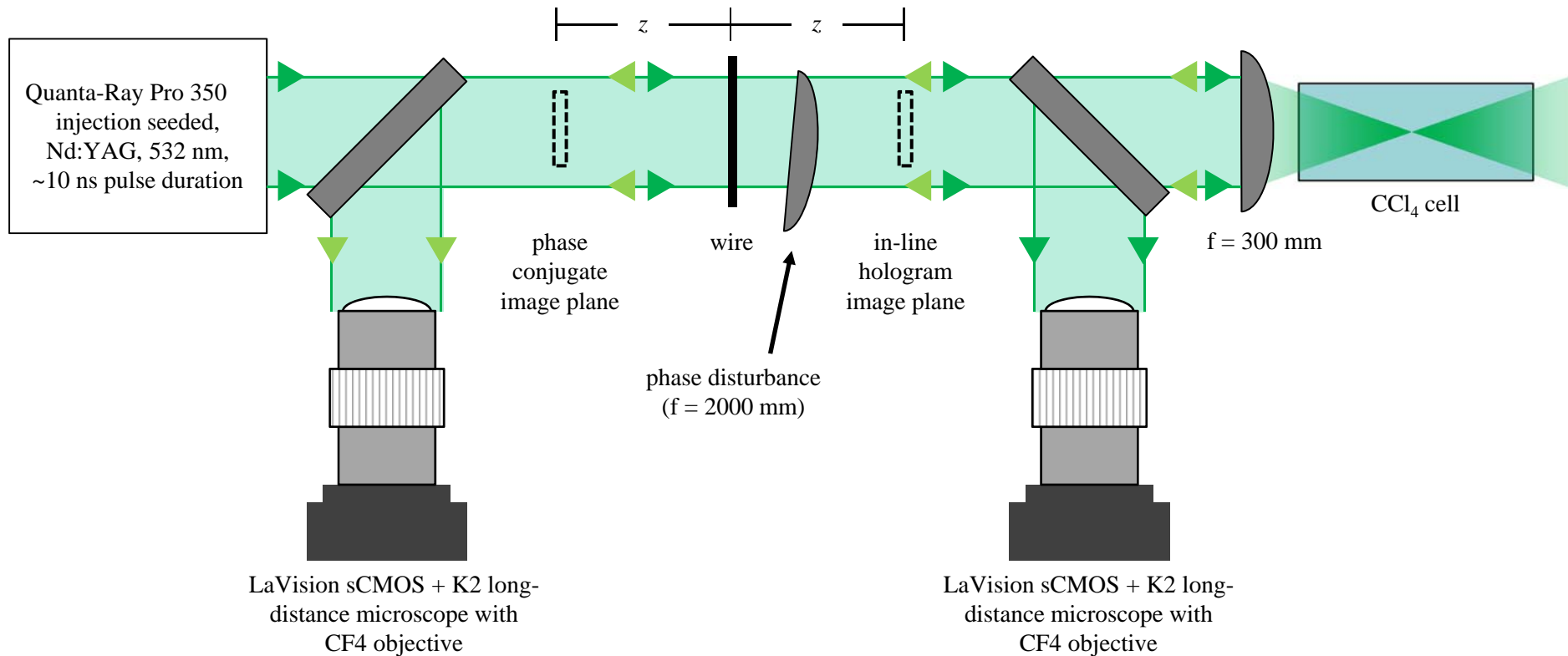
Without a
disturbance
both views give
similar results



SBS phase-conjugate DIH

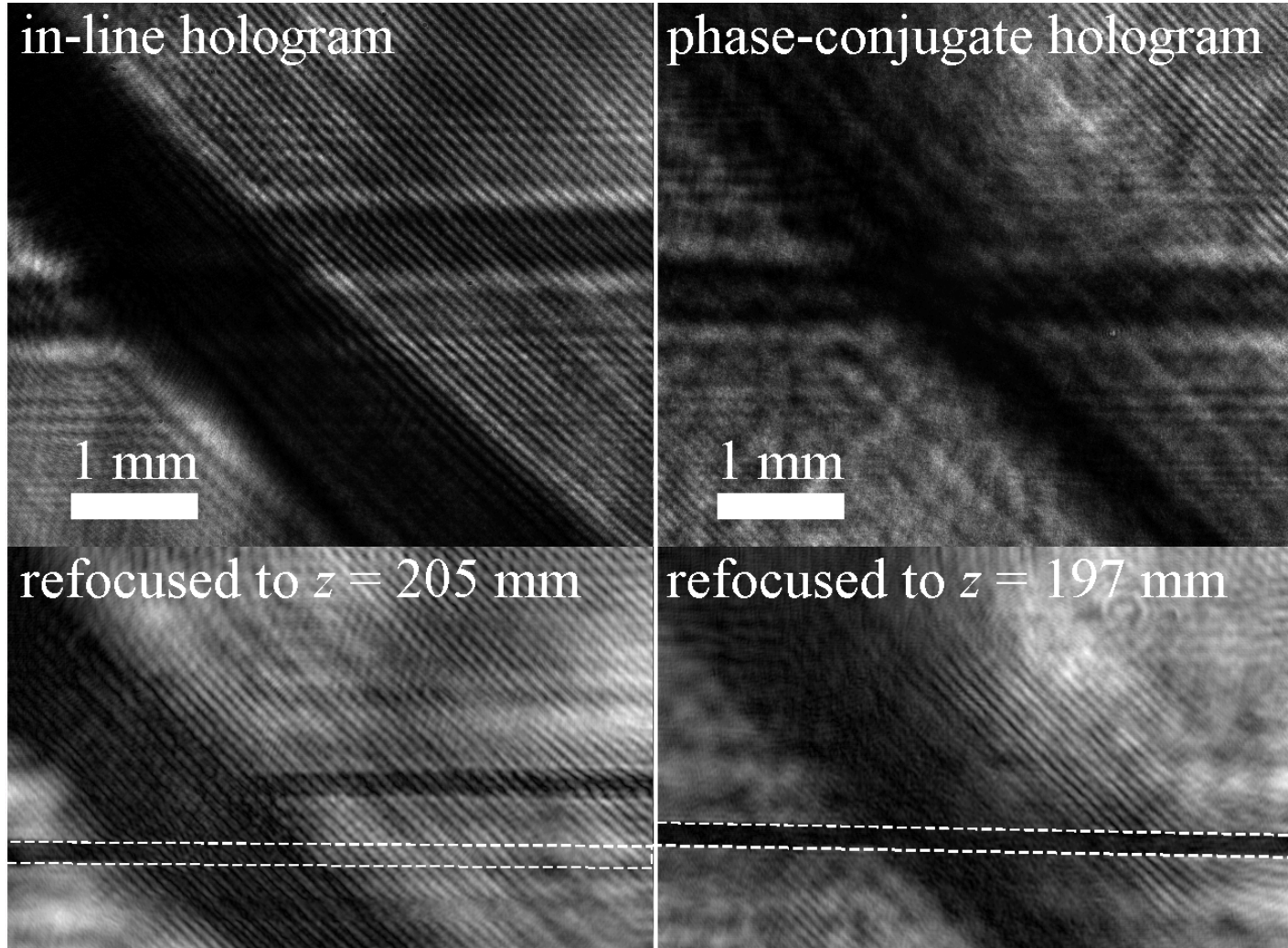
A focused beam in a non-linear medium induces phase conjugation via stimulated Brillouin scattering (SBS)

- A misaligned lens in the beam path causes a phase disturbance



SBS phase-conjugate DIH

Phase
conjugation
corrects image
distortion

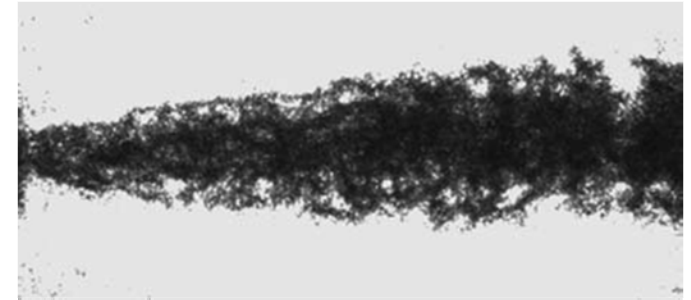


See poster: *Phase conjugate digital inline holography (PC-DIH)* by
Kathryn Hoffmeister, Sean Kearney, and Daniel Guildenbecher

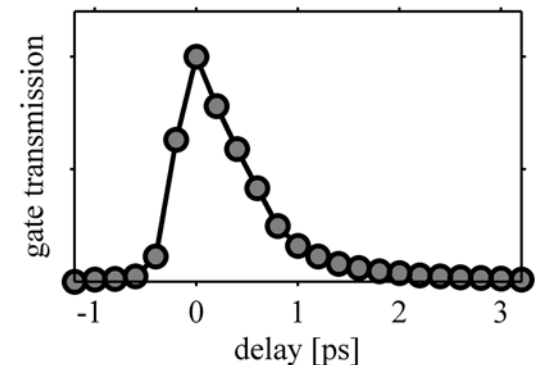
Ballistic DIH

Multiple scattering can be reduced through ps time gating

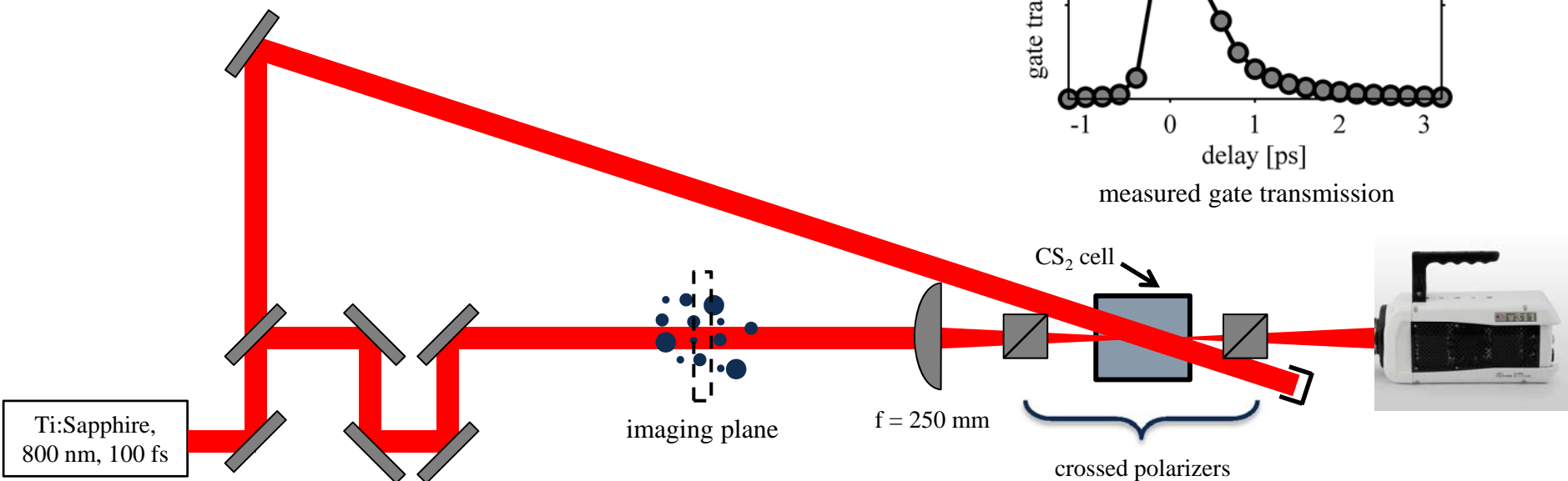
- Combination with DIH might enable scatter free 3D imaging through optically dense media
 - First proposed by: Trolinger et al 2011, *International Journal of Spray and Combustion Dynamics*



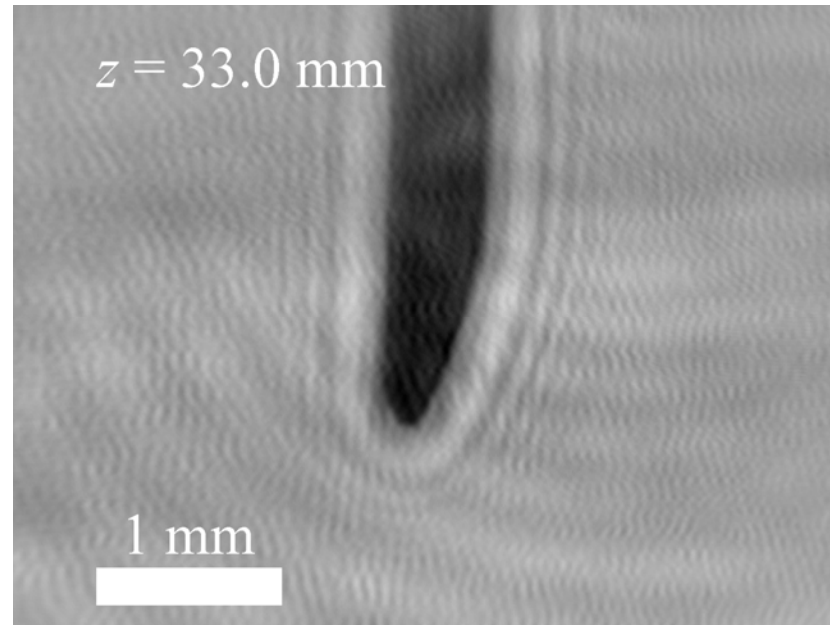
ballistic image of a diesel spray
(Linne et al 2006, *Exp. Fluids*)



measured gate transmission



DIH imaging through a Kerr gate (no scatter sources)



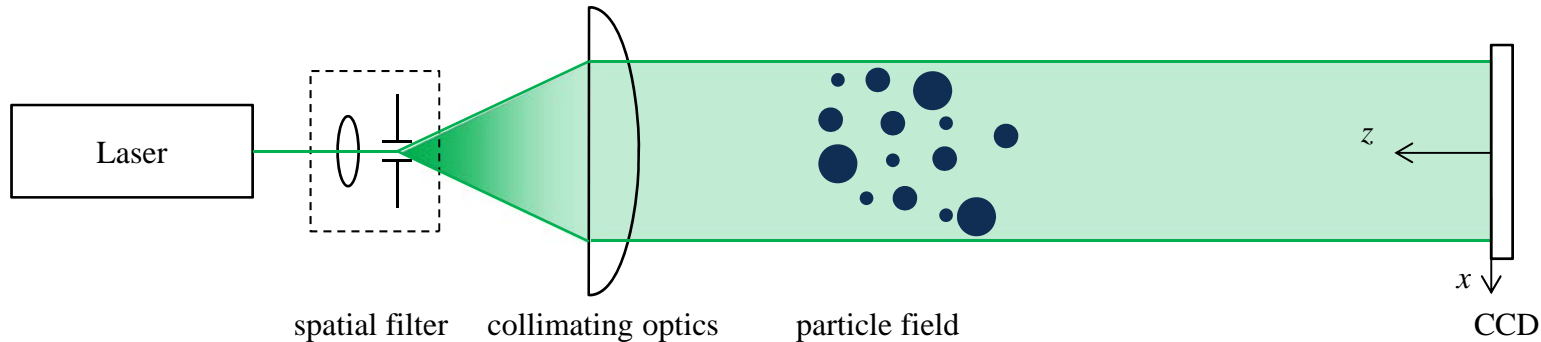
DIH image of a needle recorded with the ballistic configuration (1.6 ps switch delay)

**See poster: *Ballistic imaging holography* by
Derek Dunn-Rankin, Ali Ziaee, Jim Trolinger**

Next step: Explore ballistic DIH through dense scattering sources

- Challenge: *Can we retain sufficient image fidelity and coherence to resolve 3D phenomena?*

Where is the reference wave?



Hologram is the combination of object and reference waves: $h = |E_o + E_r|^2$

- Reconstruction with E_r gives: $h \cdot E_r = \underbrace{(|E_o|^2 + |E_r|^2)E_r}_{\text{DC term}} + \underbrace{|E_r|^2 E_o}_{\text{virtual image}} + \underbrace{E_r^2 E_o^*}_{\text{real image}}$
 - In off-axis holography, these terms are spatially separated as we attempt to reconstruct the original object wave, E_o
 - In in-line holography, we actually want to reconstruct the combination of the reference wave and object wave, $E_o + E_r$
 - Rearranging: $h \cdot E_r = \underbrace{|E_o|^2 E_r}_{\text{DC term}} + \underbrace{|E_r|^2 (E_o + E_r)}_{\text{virtual image}} + \underbrace{E_r^2 E_o^*}_{\text{real image}}$

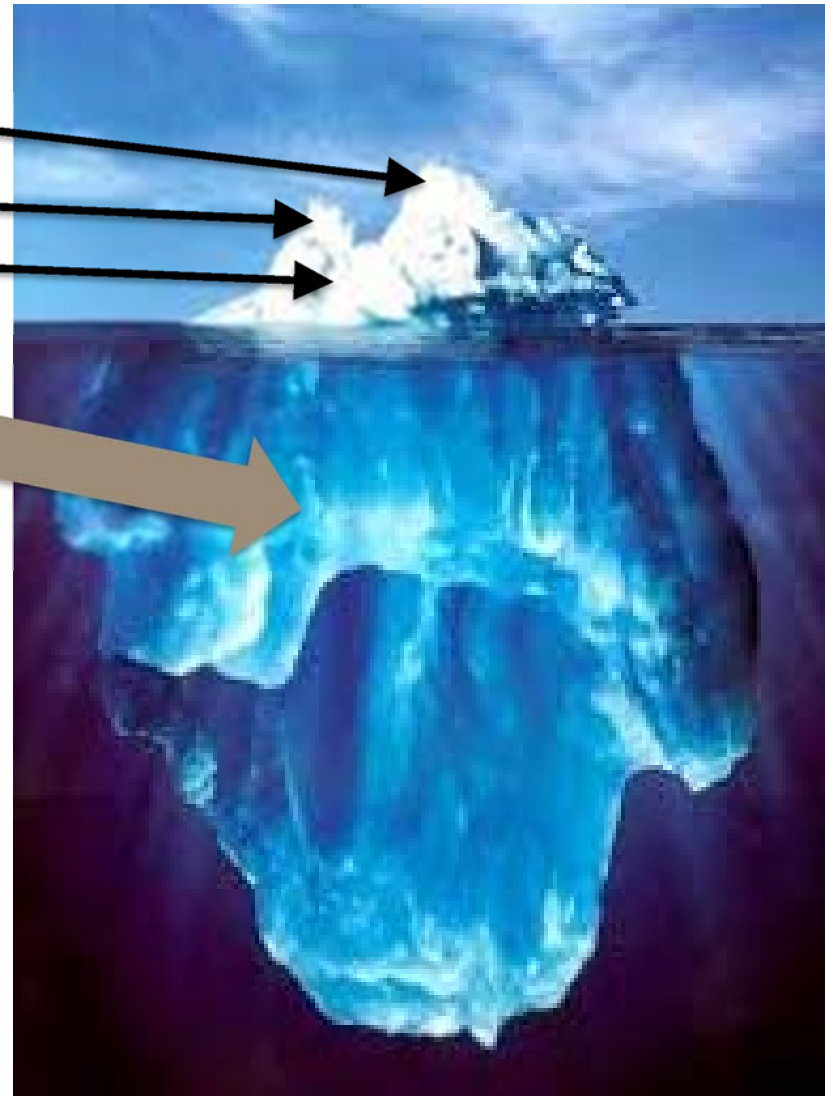
Data processing

The basic DIH system includes:

- Coherent light source (laser)
- Particle field
- Image recorder (digital camera)
- **DATA ANALYSIS SOFTWARE**

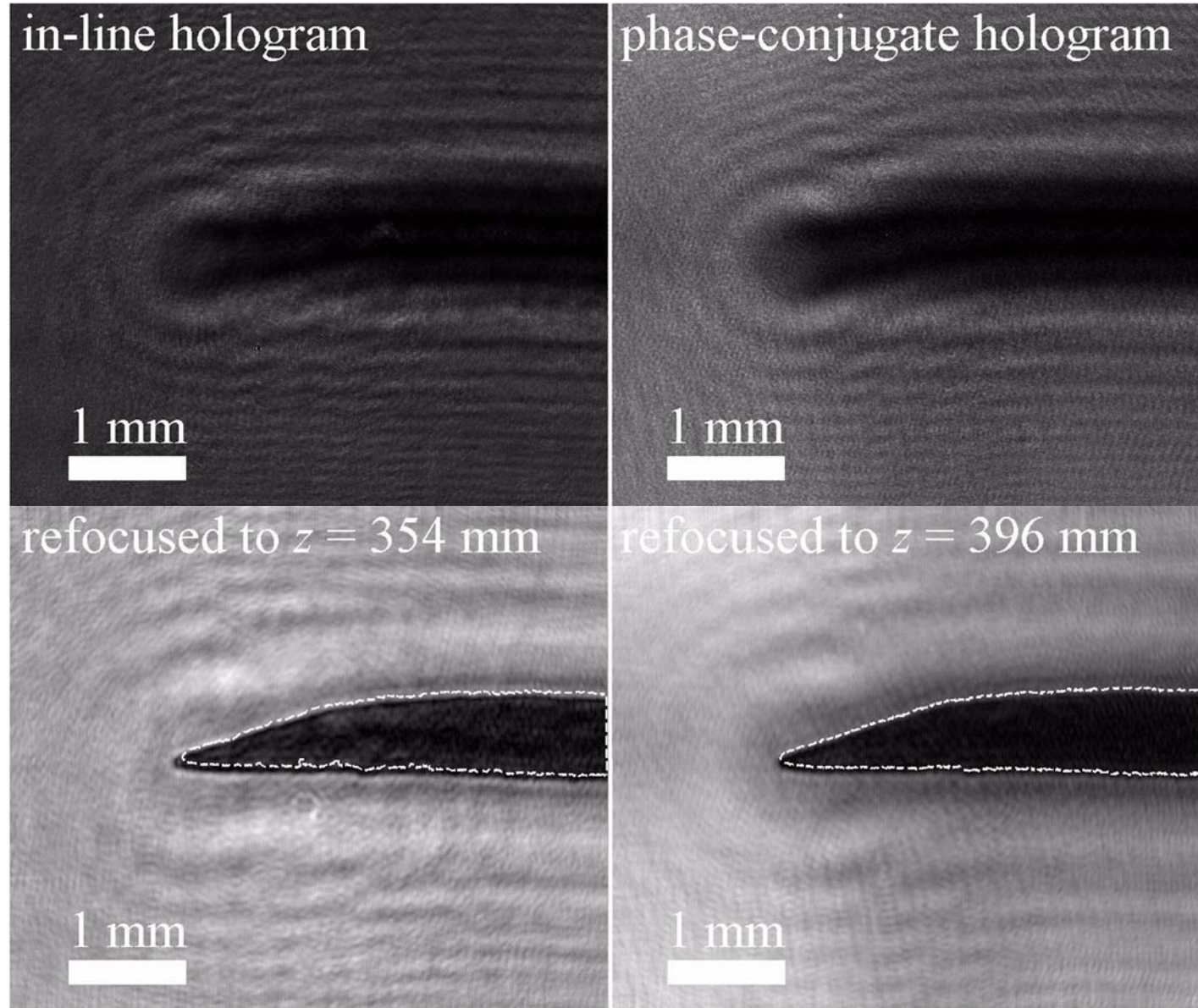
Currently each group has their own code:

- Hybrid : Guildenbecher, Gao, et al
- Laplacian: Choi and Lee
- Correlation coefficient: Yang et al
- Minimum edge intensity: Tian et al
- Variance: Palero et al.
- Etc...



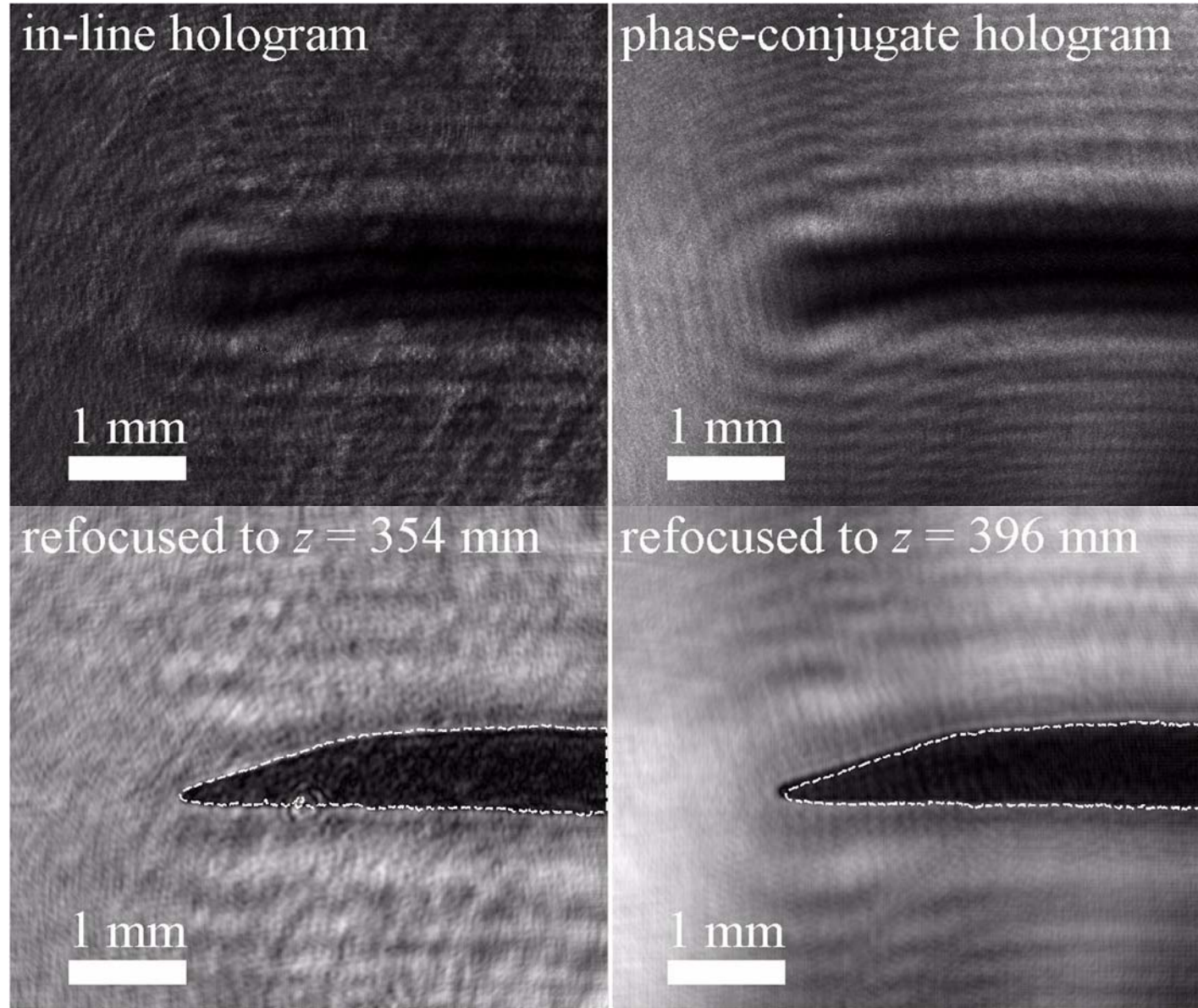
SBS phase-conjugate DIH

A hot plate
creates a phase
disturbance in
the air

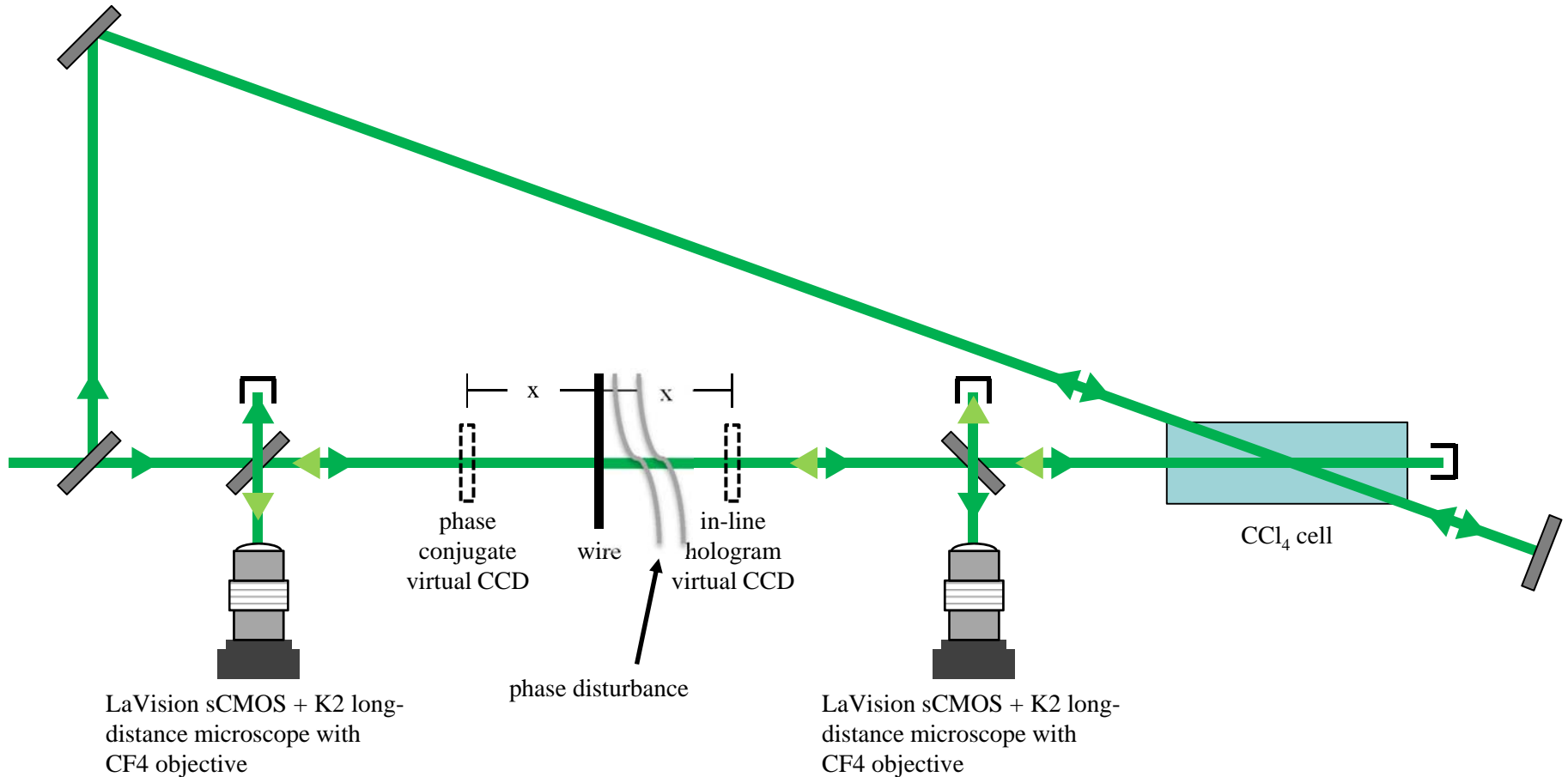


SBS phase-conjugate DIH

A butane
igniter creates
a more severe
phase
disturbance

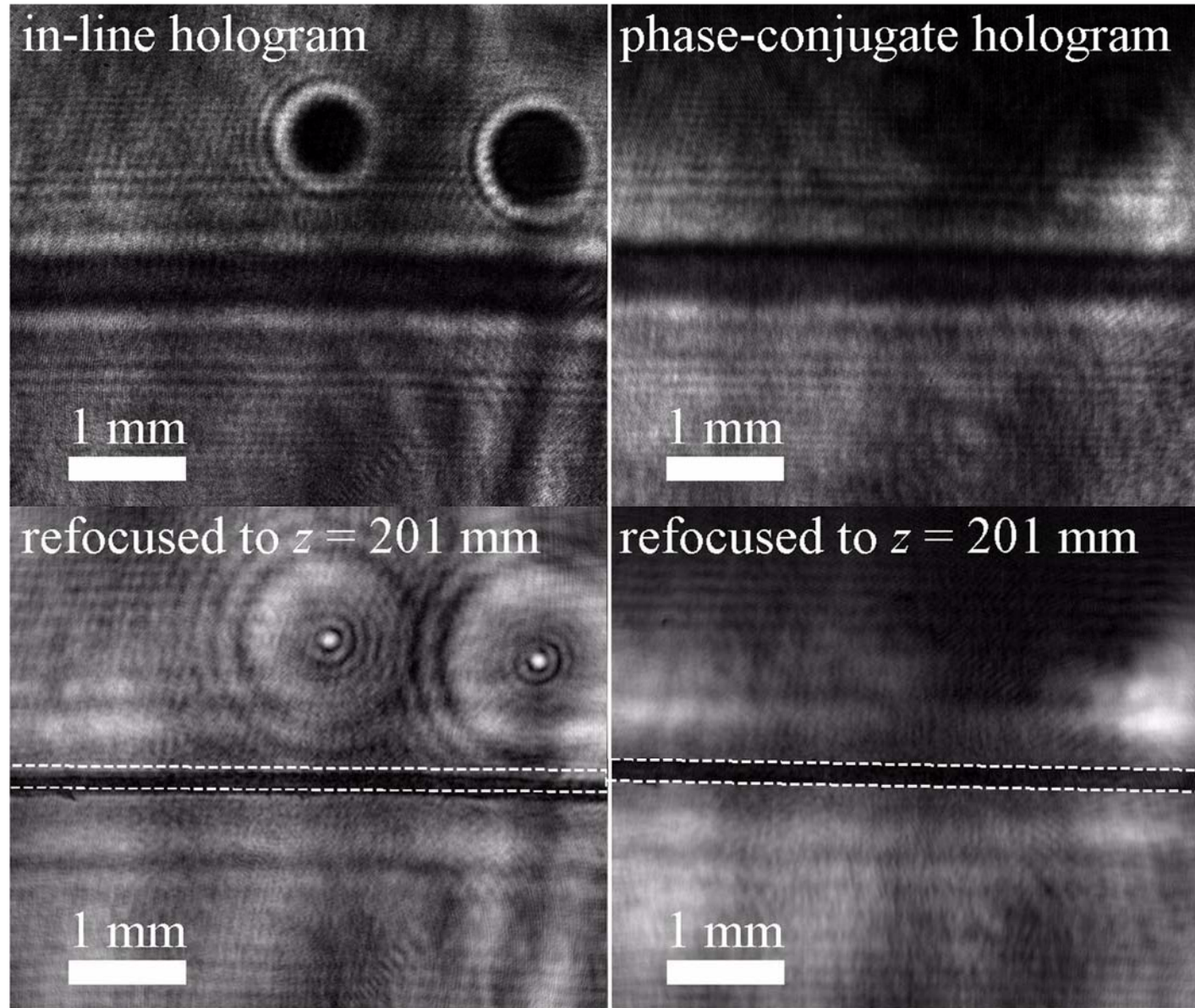


4-wave mixing phase-conjugate DIH



4-wave mixing phase-conjugate DIH

Glass with a
uneven layer of
optical glue
creates a severe
distortion



Drop impact on a thin film

Motivation: measurement of secondary droplets by other methods requires significant experimental repetition

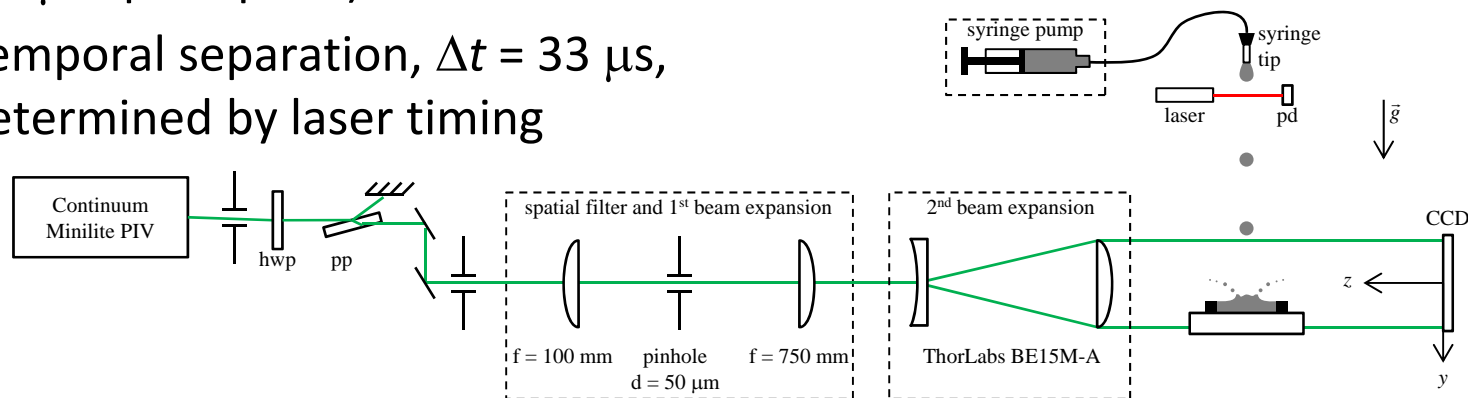
- Process symmetry provides opportunities to validate accuracy

Experimental configuration:

- Double pulsed laser ($\lambda = 532$ nm, 5 ns pulsewidth)
- Interline transfer CCD (4872×3248 , $7.4 \mu\text{m}$ pixel pitch)
- Temporal separation, $\Delta t = 33 \mu\text{s}$, determined by laser timing



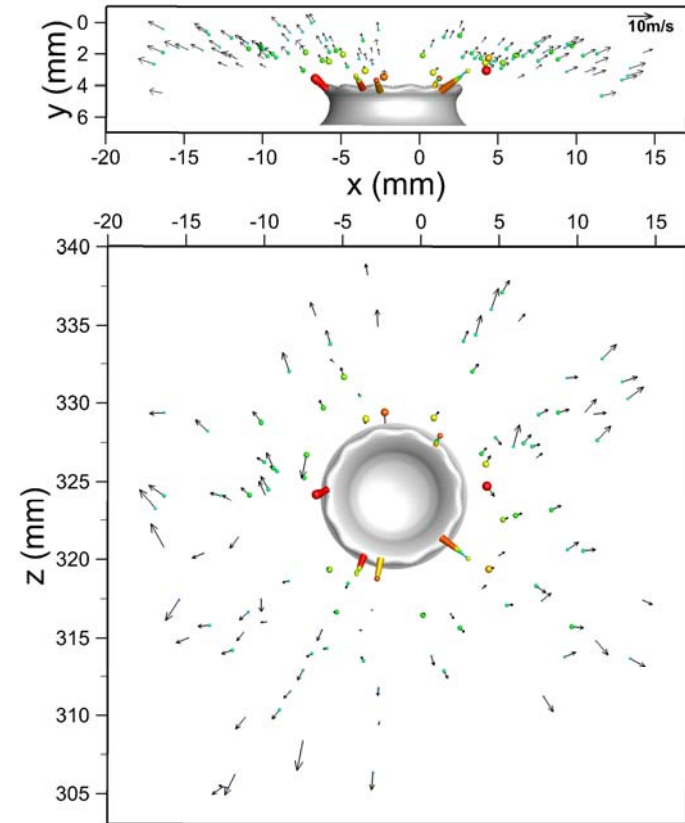
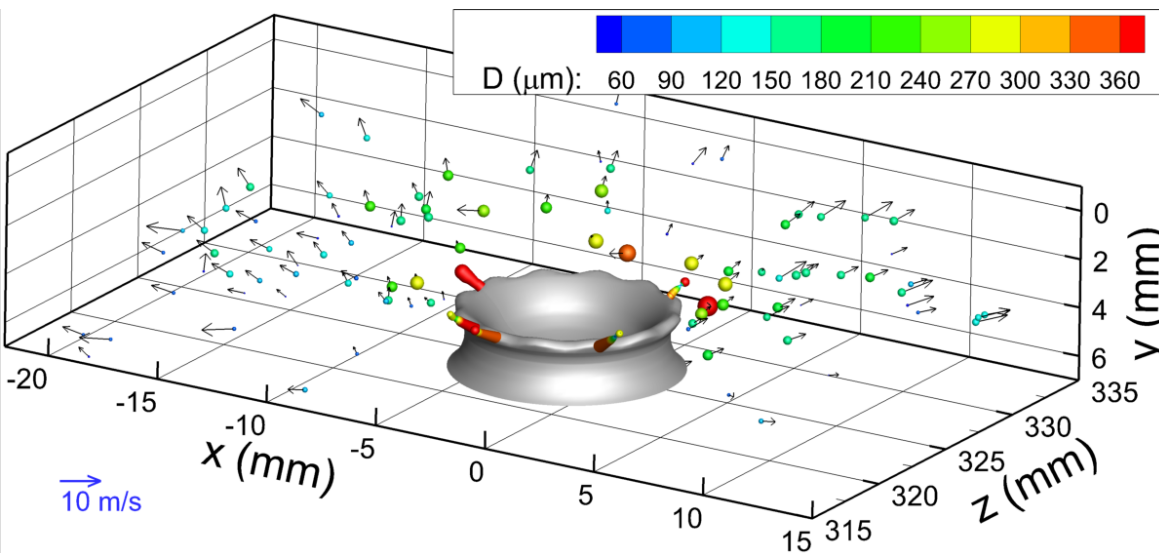
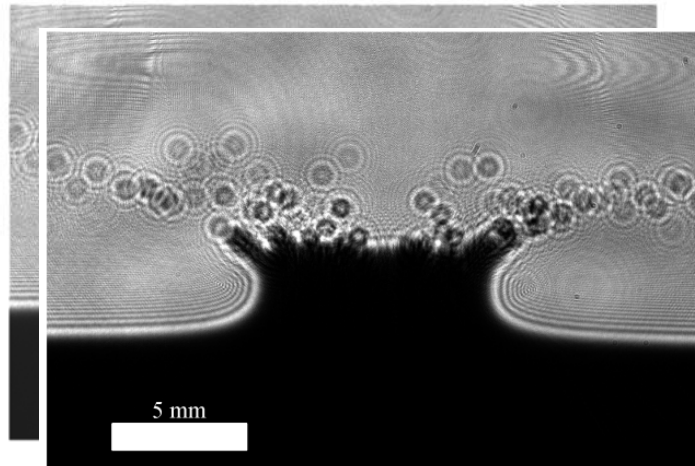
impact of a 3 mm water drop on a 2 mm water film
(Guildenbecher et al, 2013, *Exp. Fluids*.)



experimental configuration of holographic recording of drop impact on a thin film
(Guildenbecher et al, 2014, *Exp. Fluids*.)

Drop impact on a thin film

Processed with the hybrid method



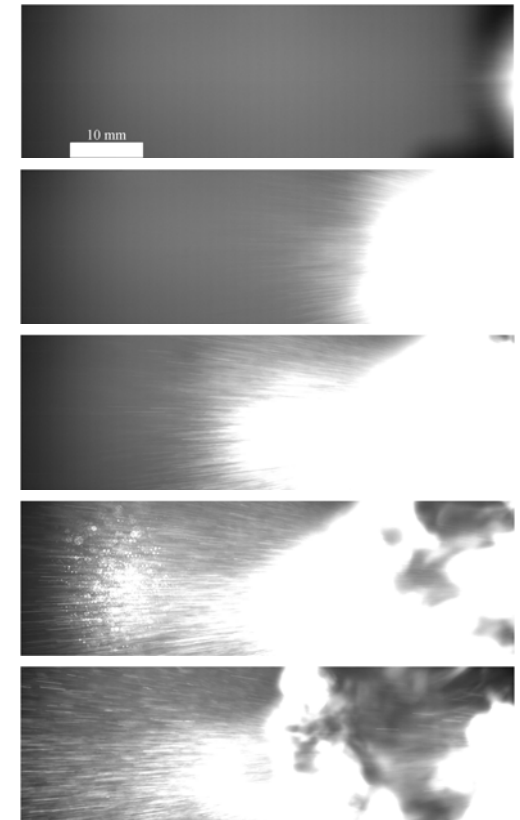
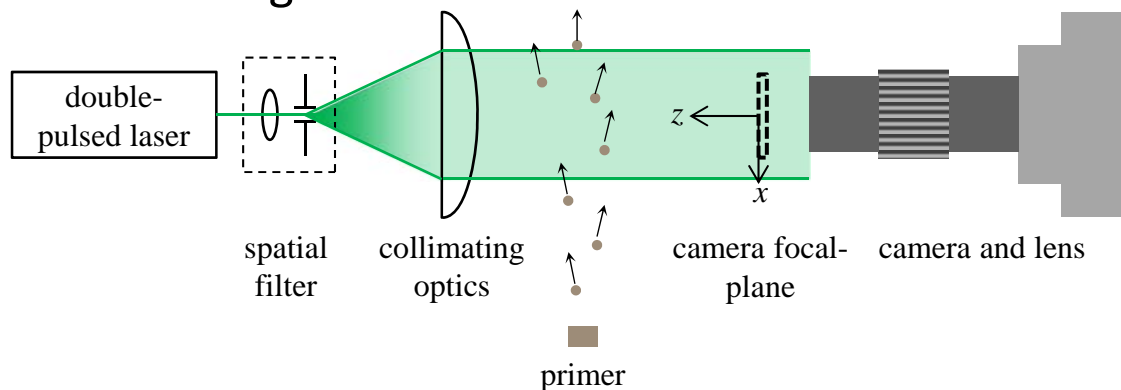
holographic reconstruction of
drop impact on a thin film
(Guildenbecher et al, 2014, *Exp. Fluids*.)

Percussion primers

Motivation: No viable technique currently exists to quantify the size and velocity distribution from the hot particles in percussion primers

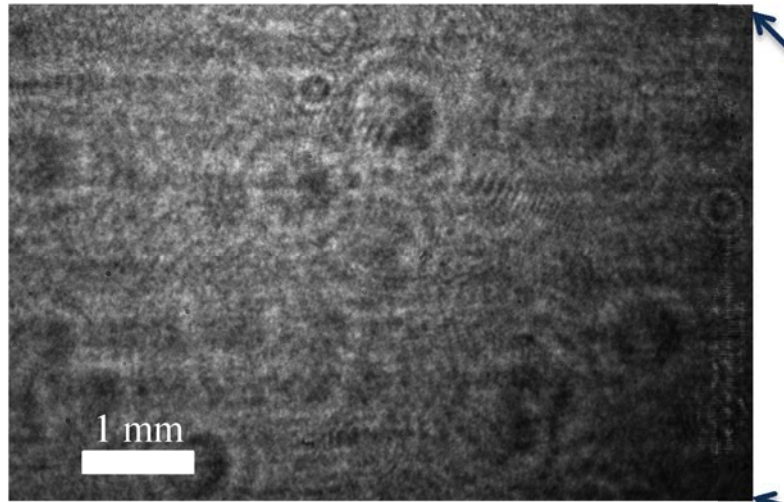
Experimental configuration:

- Double pulsed laser ($\lambda = 532$ nm, 5 ns pulsewidth)
- Interline transfer CCD (4872×3248 , $7.4 \mu\text{m}$ pixel pitch)
- $\sim 6\text{X}$ magnification achieved using Infinity K2 long distance microscope with CF-4 objective
- Temporal separation, $\Delta t = 2 \mu\text{s}$, determined by laser timing

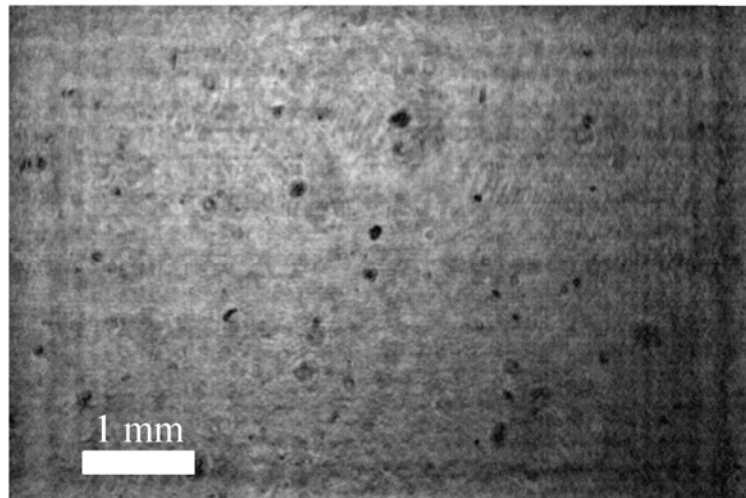


High-speed video of event

Percussion primers

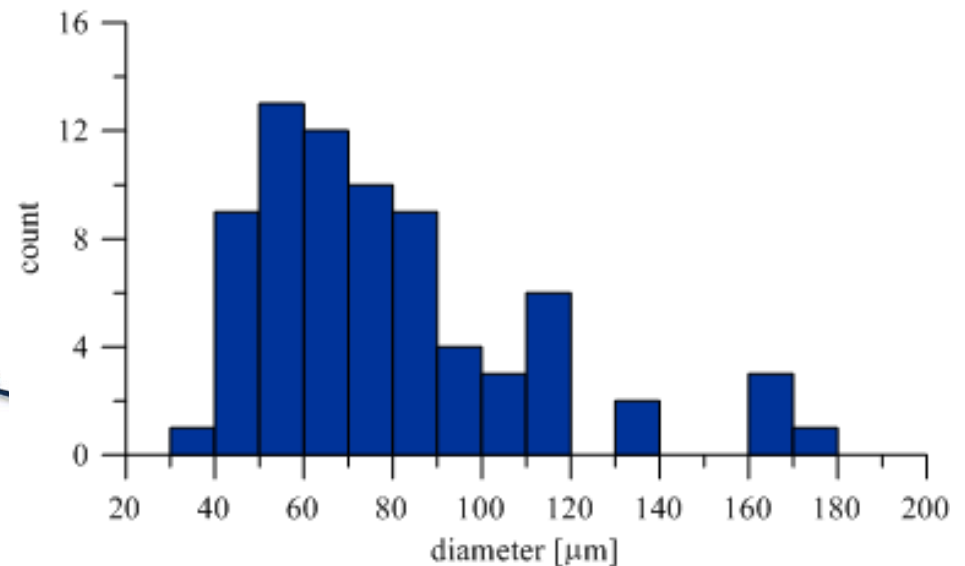


Recorded hologram



Numerically re-focused to $z = 200$ mm from the CCD

Five holograms recorded at these conditions



First known quantification of particle size

- Particle size distribution shows the expected behavior
- Probability goes to zero at large and small particle diameters

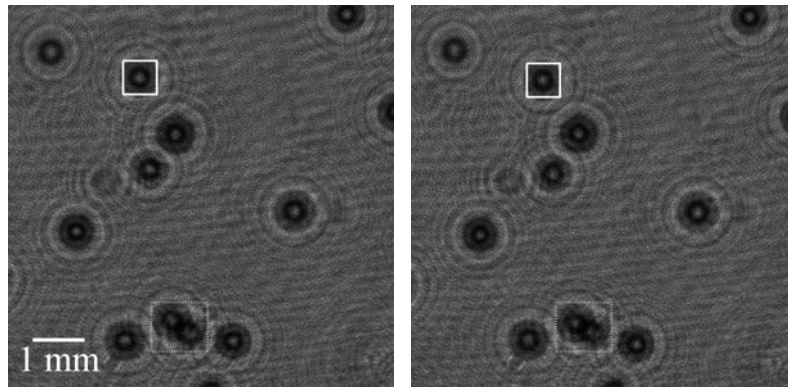
Cross-correlation method

Theory: in-focus particle images from two sequential holograms contain correlated information

- The maximum cross-correlation, c , gives the displacement $(\Delta x, \Delta y)$

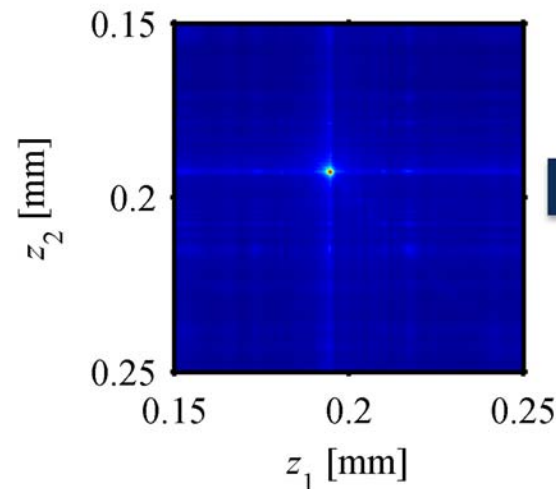
$$c = \max_{\Delta x, \Delta y} \left[\sum_m \sum_n \text{Img}_1(m, n) \text{Img}_2^*(m, n)(m - \Delta x, n - \Delta y) \right]$$

- Img_1 and Img_2 chosen as the edge sharpness images from the two frames
- z positions in each frame (z_1 and z_2) are found from the maximum value of c over all possible combinations of z_1 and z_2



hologram
(Guildenbecher et al,
2013, *Opt. Lett.*)

hologram after displacing
the particle field by 2 mm
(Guildenbecher et al,
2013, *Opt. Lett.*)

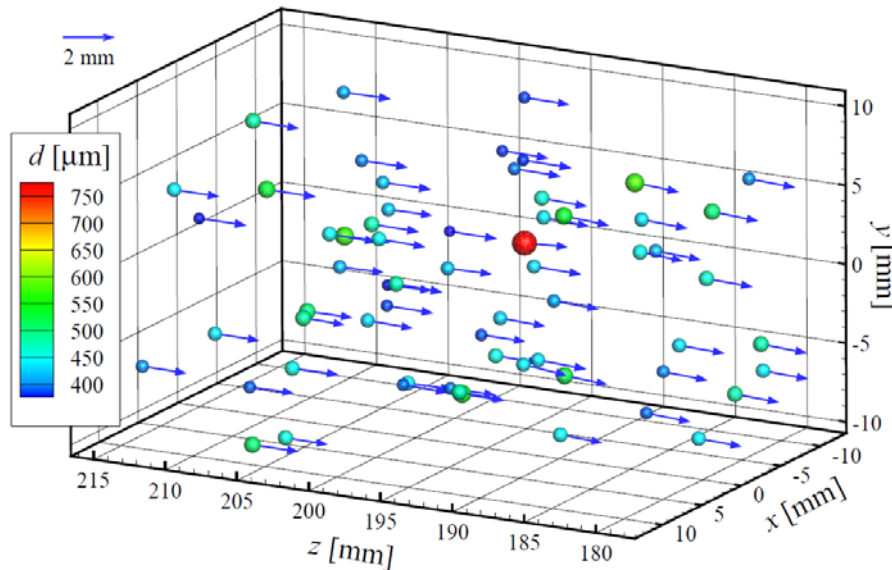


maximum value of c for the particle in the white
boxes (Guildenbecher et al, 2013, *Opt. Lett.*)

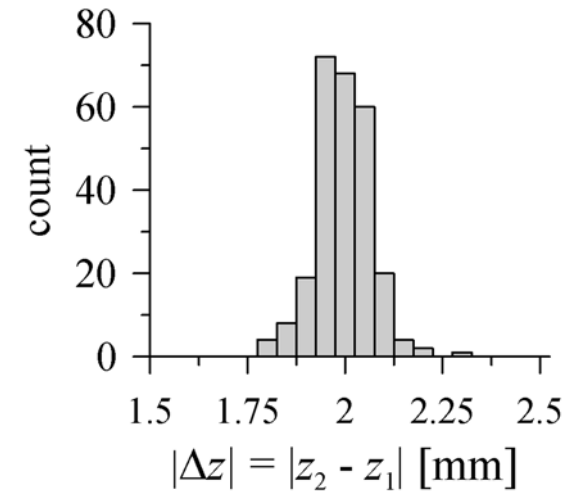
$z_1 = 194.72 \text{ mm},$
 $z_2 = 192.72 \text{ mm},$
 $\Delta z = 2.00 \text{ mm}$

Cross-correlation method

Again, experimentally validated with quasi-stationary particles in silicone oil



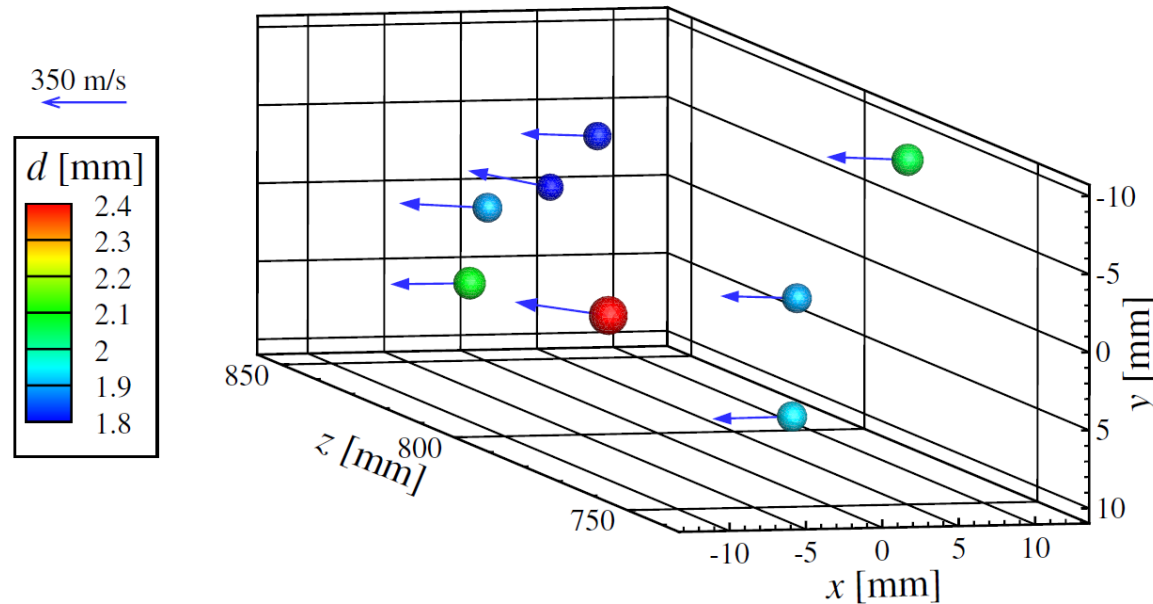
measured displacement field from one realization
(Guildenbecher et al, 2013, *Opt. Lett.*)



measured z -displacements from all realizations
(Guildenbecher et al, 2013, *Opt. Lett.*)

- Actual displacement = 2.0 mm
- Mean detected displacement = 1.996 mm +/- 0.072 mm
 - Standard deviation of 0.15 times mean diameter
 - Order of magnitude improvement compared to uncertainties in the literature

Sonic pellets from a shotgun

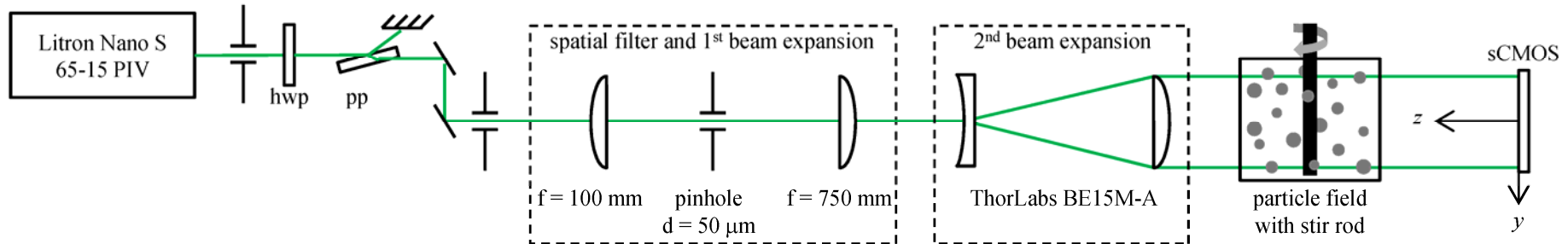


particle field from the shotgun measured with the cross-correlation method
(Guildenbecher et al, 2013, *Opt. Lett.*)

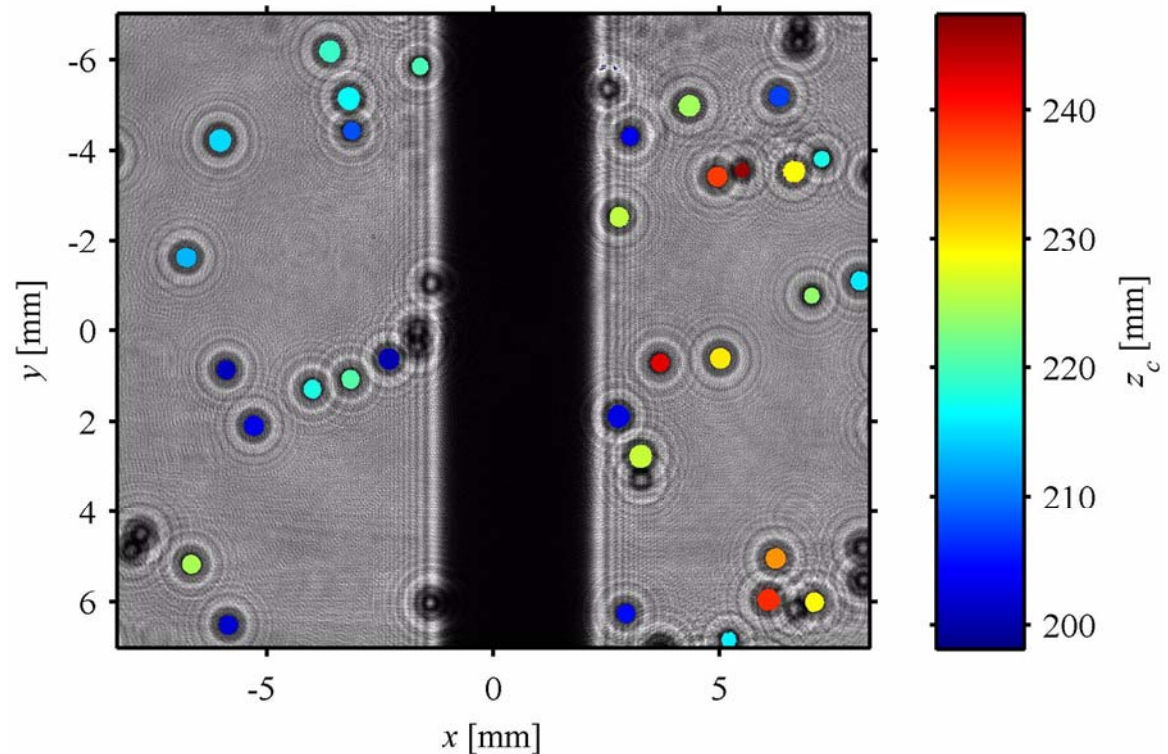
Results closely match the expected mean velocity (350 m/s) and diameter (2.0 mm)

- Uncertainty in Δz is on the order of 0.2 particle diameters

3D, 3C fluid velocity measurements?



- Particles stirred by a rotating rod ($r_0 = 1.58$ mm, $\omega_0 = 100$ rpm)
- Recorded at 15 Hz with a LaVision sCMOS camera (2560×2160 , 6.5 μm pixel pitch)



particles measured with the hybrid method, background shows the recorded holograms

3D, 3C fluid velocity measurements?

For all trajectories

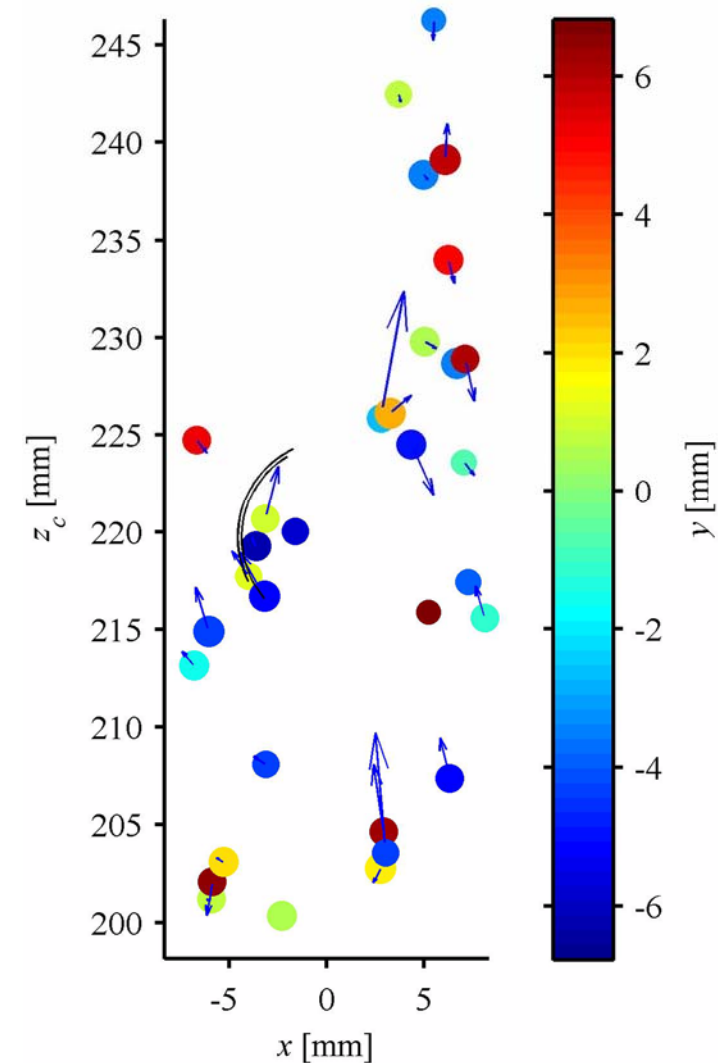
- Error in measured $z = -0.04 \pm 1.51$ mm
- Error in measured $\Delta z = -0.03 \pm 1.05$ mm
 - Standard deviation of $2.3 \cdot \bar{d}$

Experiments repeated with smaller particles
($\bar{d} = 118 \mu\text{m}$, see paper for details)

- Error in measured $z = -0.003 \pm 0.379$ mm
- Error in measured $\Delta z = -0.001 \pm 0.302$ mm
 - Standard deviation of $2.6 \cdot \bar{d}$

Next steps:

- Compare results with alternative particle detection methods
- Use results to quantify effects of particle overlap and other experimental noise sources



all measured x - z trajectories vs. predicted

3D, 3C fluid velocity measurements?

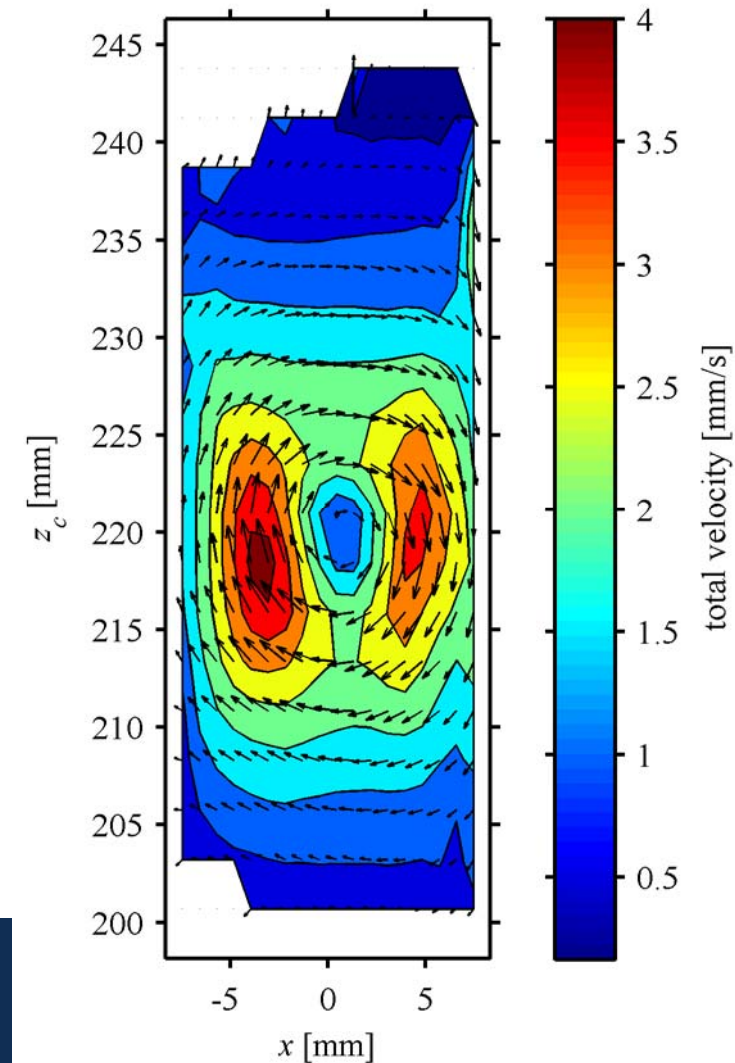
Advantages:

- Simple optical setup requiring only one line-of-sight view
- Large depth of field (hundreds of mm possible)
- Particle sizes can be measured (if desired)

Challenges:

- High uncertainty in the z-direction
- Particle field must be relatively sparse providing only limited vectors
- Vectors at random positions
- Methods not as mature as PIV or even tomographic-PIV

Note: the literature contains many works on holographic-PIV. My own work has not been focused on these applications



mean measured x-z velocities