

Pulsed Power Driven High Energy Density Plasmas

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*Exceptional
service
in the
national
interest*



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²*Laboratory of Plasma Studies, Cornell University, Ithaca, NY*

³*Lawrence Livermore National Laboratory, Livermore, CA*

⁴*General Atomics, San Diego, CA*

⁵*Raytheon Ktech, Albuquerque, NM*

⁶*Naval Research Laboratory, Washington, DC*

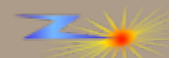


Synopsis

- Pulsed-power technology produces large currents (15-27 MA) in a short pulse (100-600 ns) on the Z machine
- Large currents generate large magnetic fields
= tremendous pressure
- Large pressures enable access to High Energy Density regimes ($> \sim 10^{11} \text{ J m}^{-3}$, or $> \sim 1 \text{ Mbar}$)
- Magnetically-driven implosions and explosions of many types provide many interesting applications
- These applications go well beyond traditional concepts known colloquially as “z-pinches”

Pressure is
equivalent to
energy density

1 Mbar =
1.01e6 atmospheres =
1.e11 Pascals =
100 GPa =
5 MGauss =
1.e11 J/m³



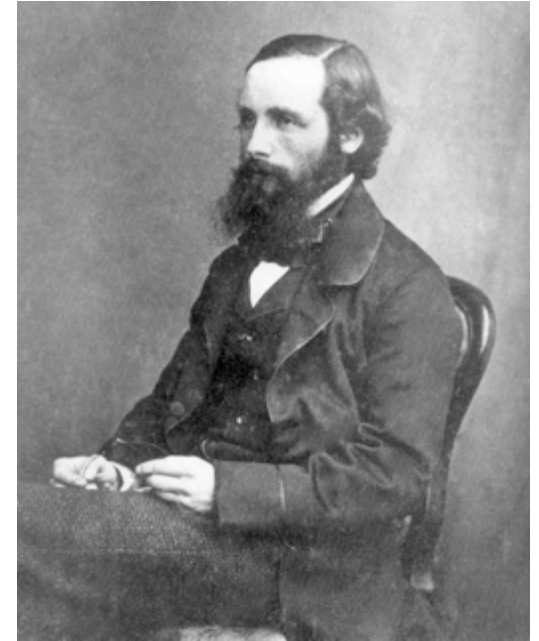
Pulsed power is an amazing technology

Pulsed-power accelerators:

- Serve as precision scientific instruments
- Deliver megajoules of energy to milligrams of matter on a time scale of nanoseconds
- Achieve extreme states of matter over macroscopic volumes
- Drive a wide variety of high-energy-density-physics experiments in support of the U.S. national-security mission



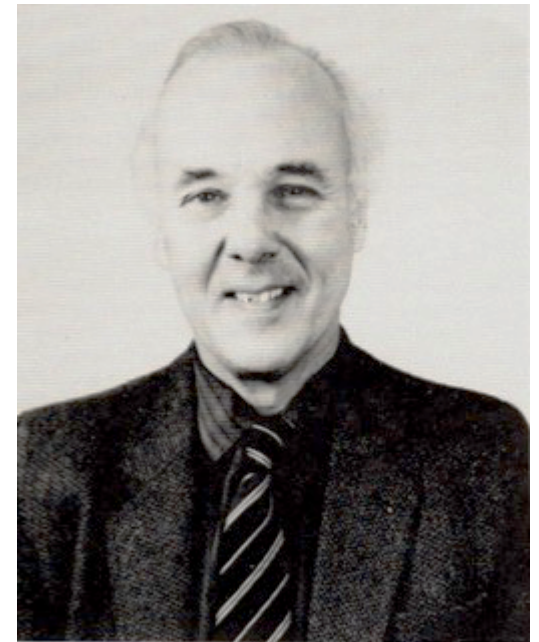
Michael Faraday



James Clerk Maxwell



Erwin Otto Marx

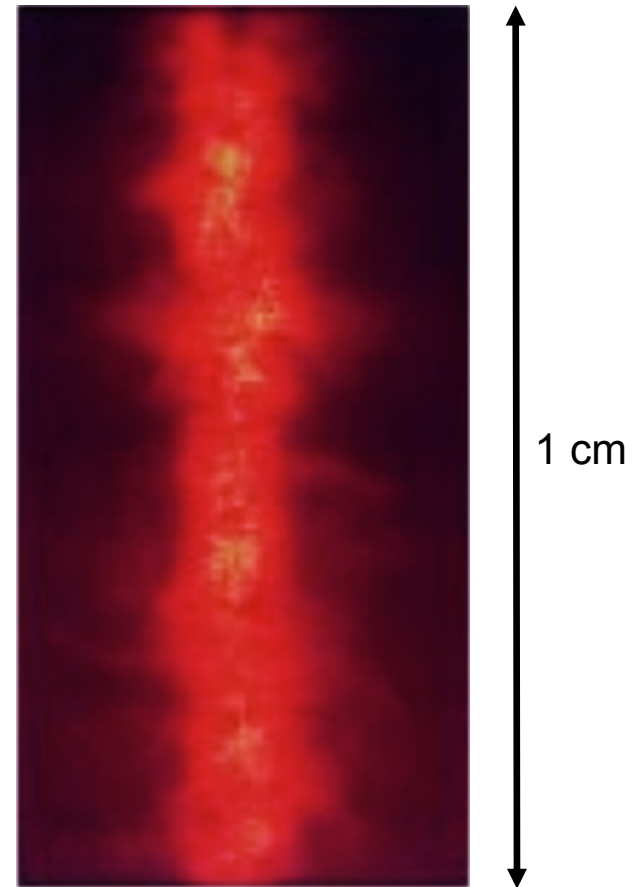


J. C. "Charlie" Martin

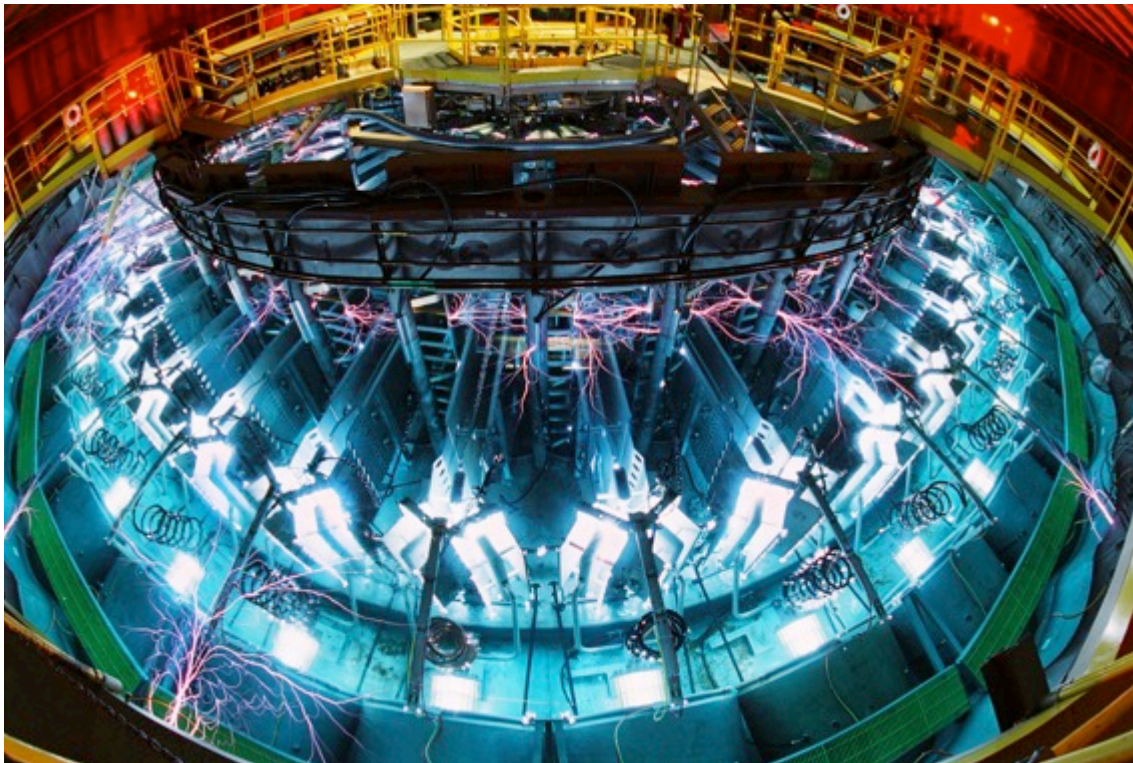
Pulsed-power experiments presently deliver the following:

- Kinetic energy per atom: 1 MeV
- Implosion velocities: 100 cm/ μ s
- Shock velocities: 30 km/s
- Temperatures: 5 keV
- Magnetic pressures: 5 Mbar
- Energy radiated in K-shell x rays: 400 kJ
- Energy radiated in thermal x rays: 2.2 MJ
- Power radiated in thermal x rays: 330 TW

100-ps x-ray
image of a
280-TW z pinch
at stagnation



(Deeney, Douglas, Spielman,
and colleagues, PRL, 1998.)



20-TW 33-m-diameter Saturn
accelerator

(Bloomquist, Corcoran,
Spielman, and colleagues.)

Sandia's Z accelerator is *presently* the world's largest and most powerful pulsed-power machine

$E_{\text{stored}} = 20 \text{ MJ}$

$V_{\text{stack}} = 4 \text{ MV}$

$I_{\text{load}} = 26 \text{ MA}$

$E_{\text{radiated}} = 2.2 \text{ MJ}$

$P_{\text{electrical}} = 80 \text{ TW}$

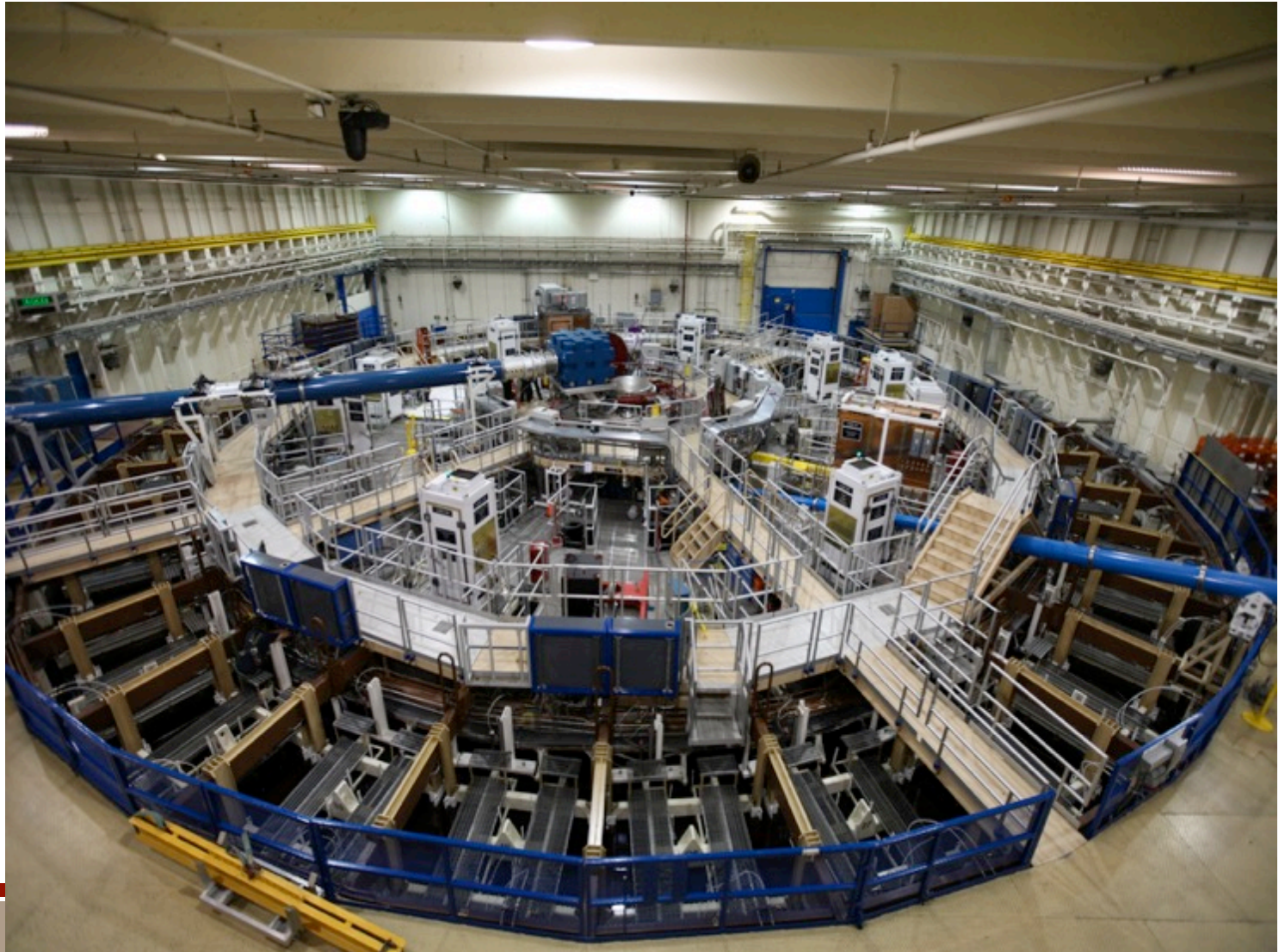
$L_{\text{vacuum}} = 12 \text{ nH}$

$\tau_{\text{implosion}} = 130 \text{ ns}$

diameter = 33 m

- Since 1997 we have conducted, on average, 160 Z shots each year.
- To date, 2800 Z shots have been conducted altogether.

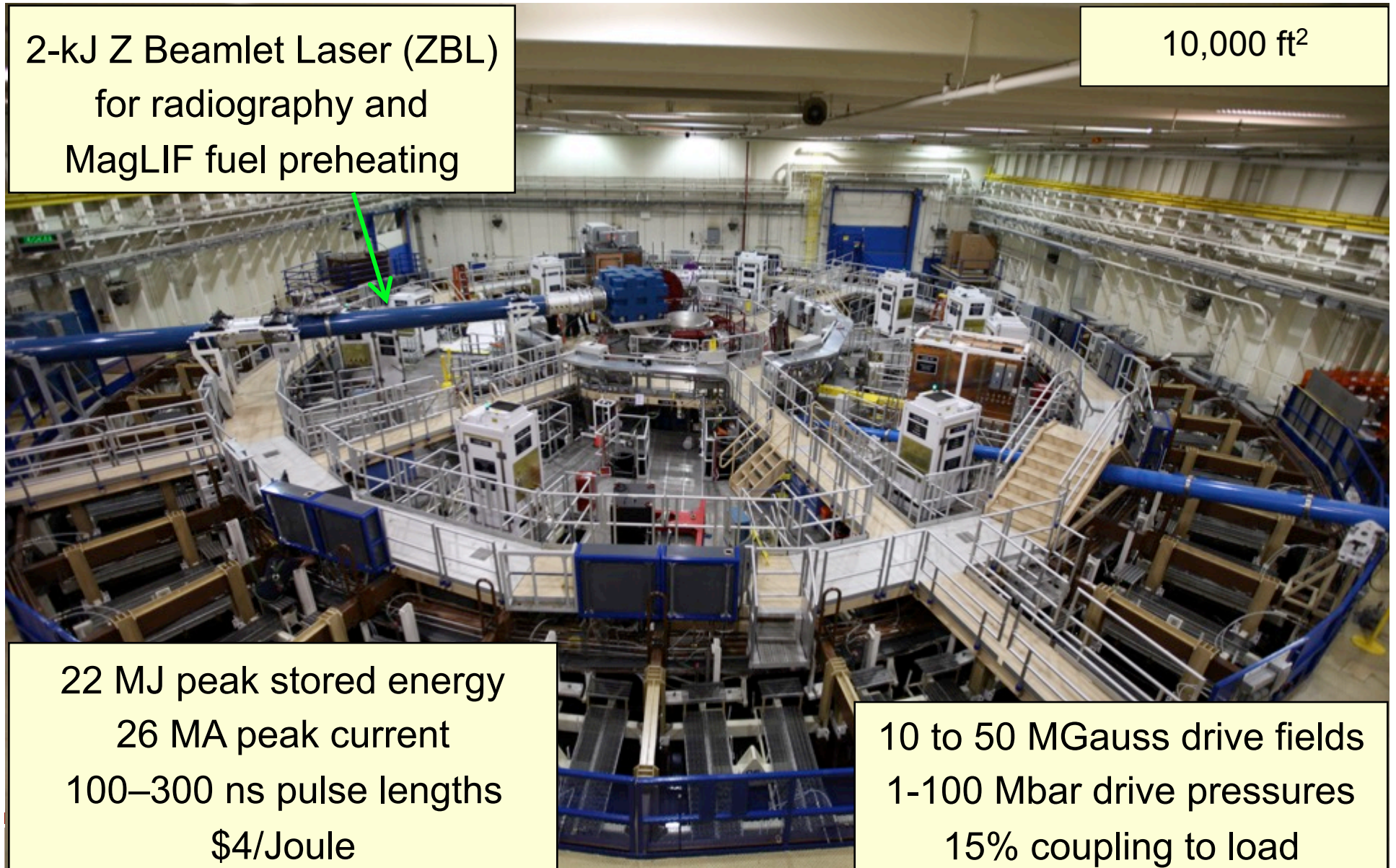
McDaniel,
Bloomquist, Ives,
Johnson, LeChien,
Lehr, Maenchen,
Savage, Spielman,
Stoltzfus, Weed, and
Weinbrecht.



The Z pulsed-power facility combines a compact MJ-class target physics platform (the Z accelerator) with a TW-class laser (ZBL)

2-kJ Z Beamlet Laser (ZBL)
for radiography and
MagLIF fuel preheating

10,000 ft²



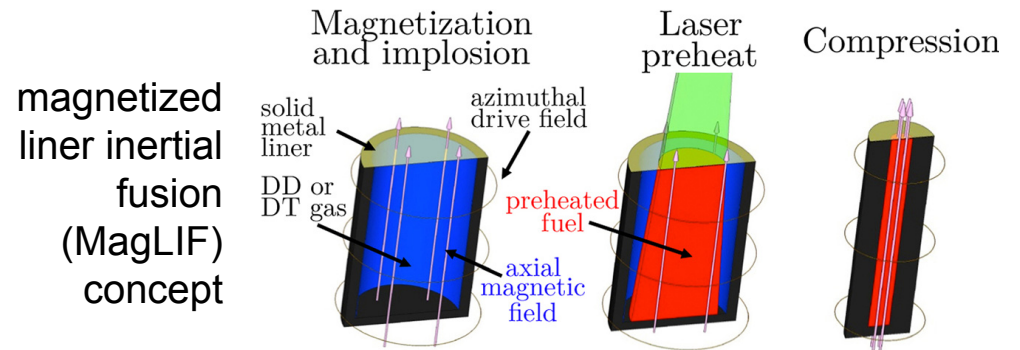
22 MJ peak stored energy
26 MA peak current
100–300 ns pulse lengths
\$4/Joule

10 to 50 MGauss drive fields
1-100 Mbar drive pressures
15% coupling to load

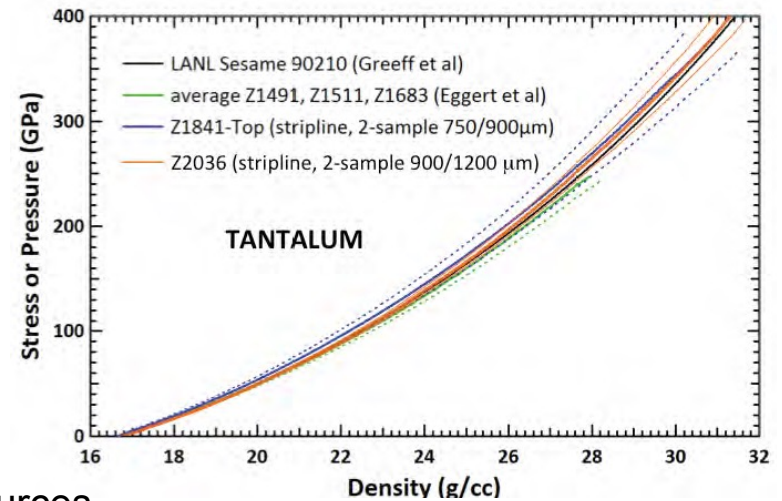
Z drives a wide variety of experiments in support of the U.S. national-security mission

- Inertial confinement fusion
- Radiation physics
- Material physics
- Laboratory astrophysics

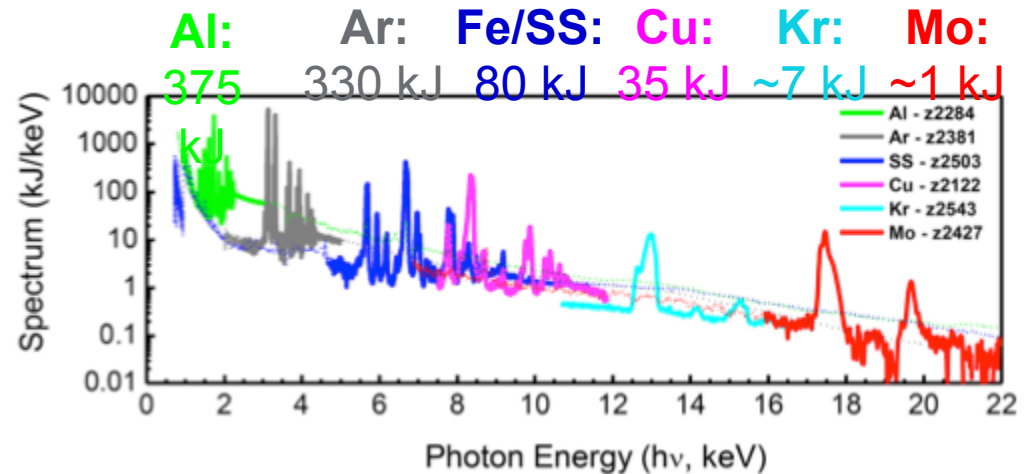
Results of experiments conducted on Z have been published in over 500 peer-reviewed journal articles.



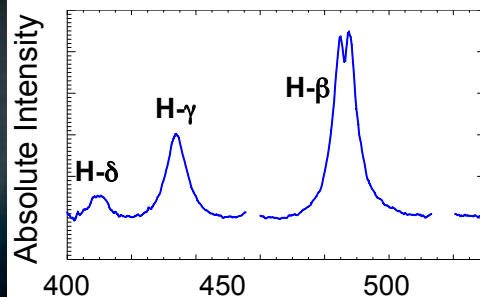
quasi-isentrope of tantalum to 4 Mbar



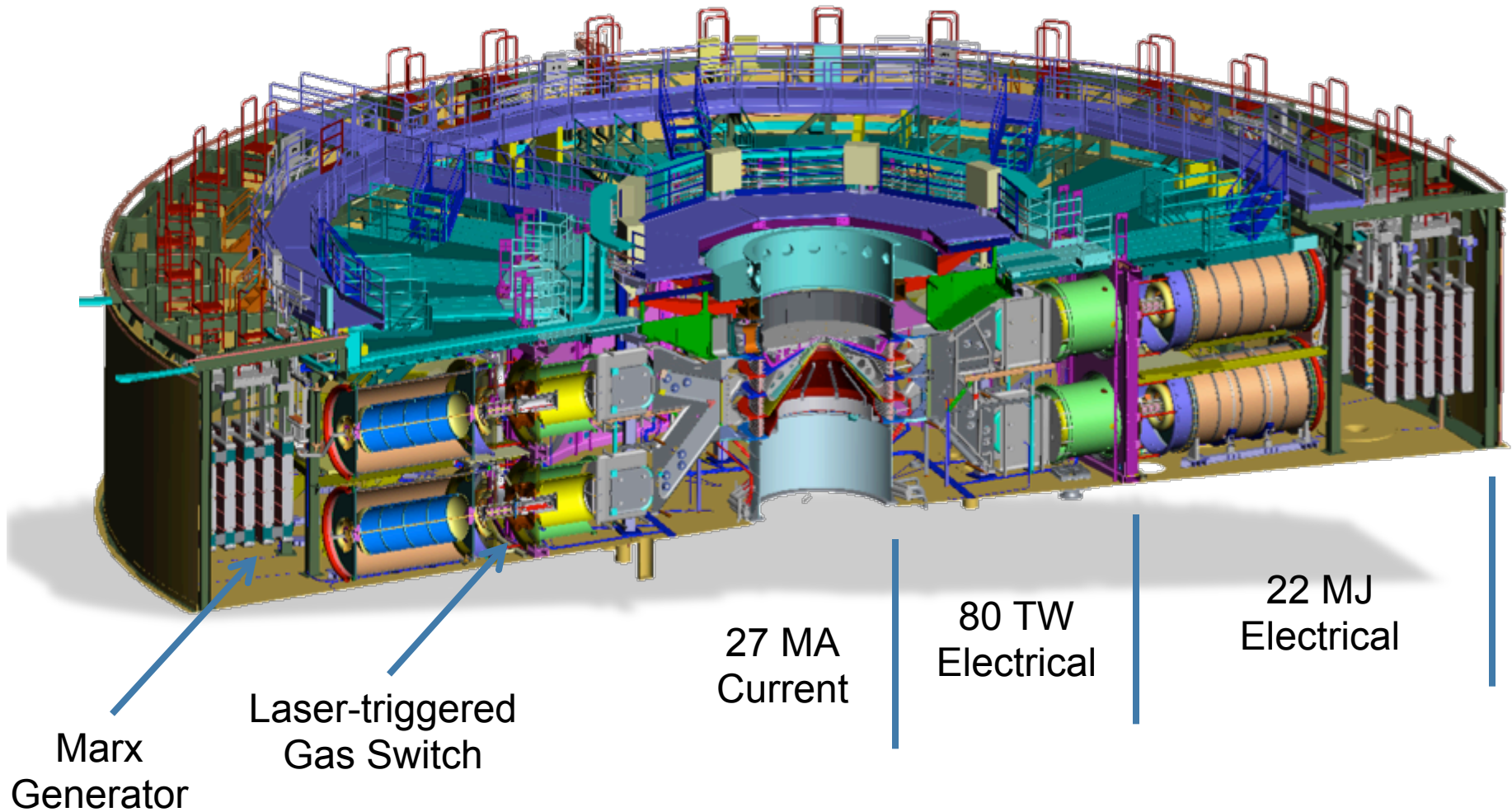
K-shell x-ray sources



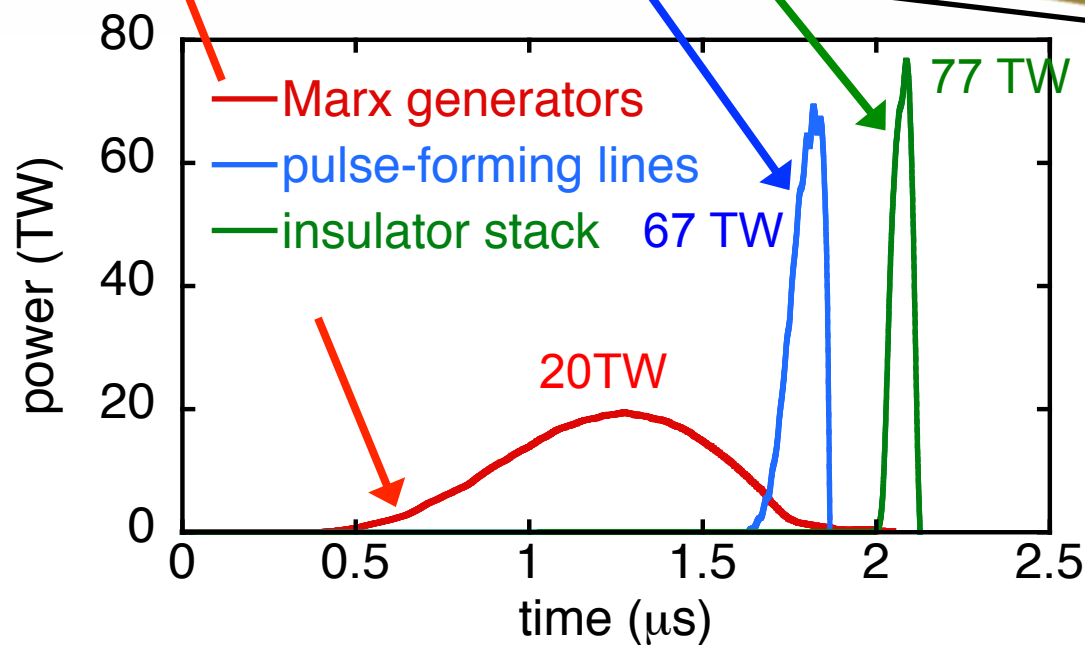
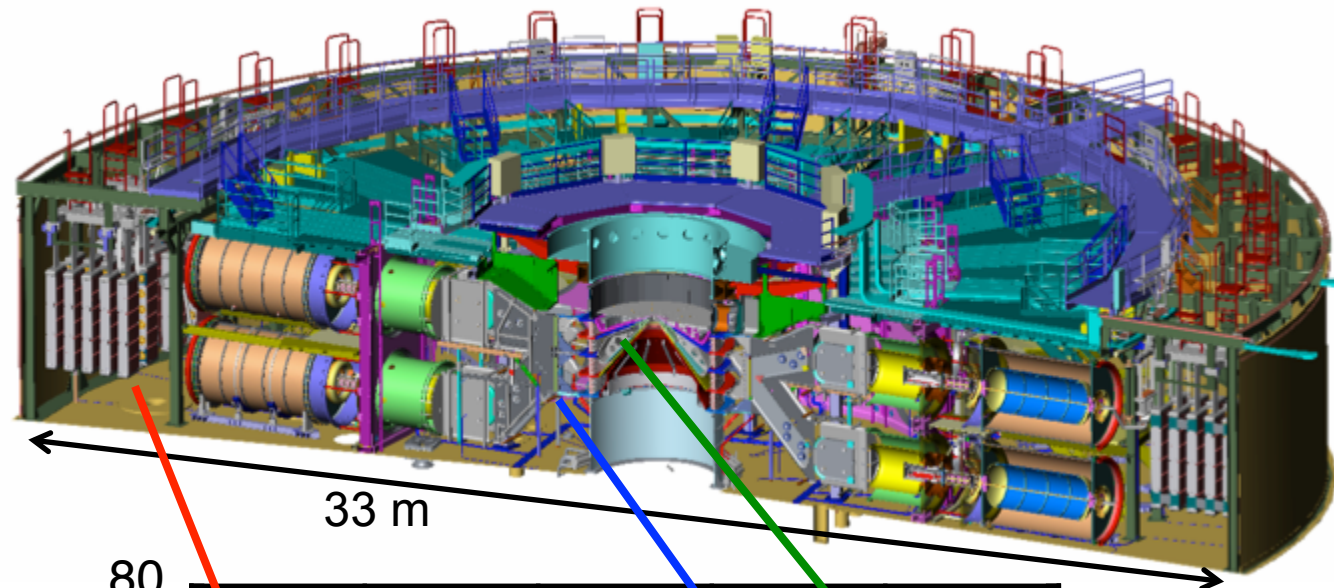
white-dwarf line shapes



Cross section of the Z facility at Sandia National Laboratories



Z works by compressing energy in space and time to reach high energy densities (pressures)

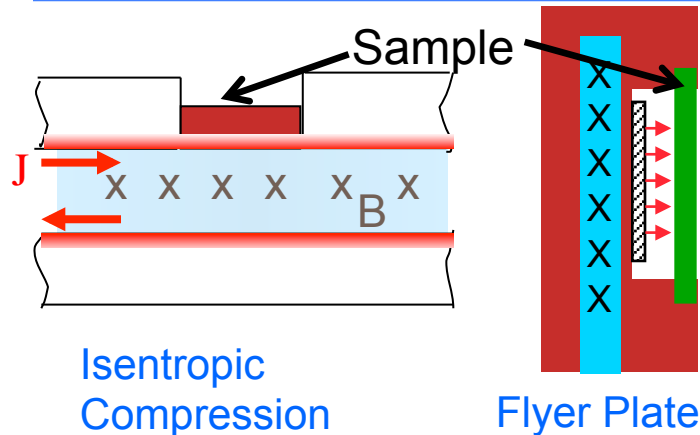


- High electrical power density achieved through multiple stages of spatial and temporal pulse compression

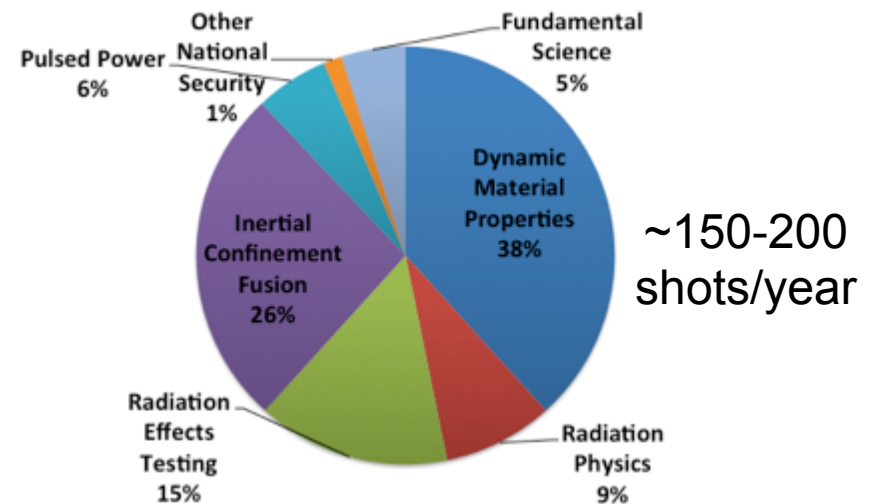
- Further increases of energy density are achieved through shock compression, implosion, stagnation

We use magnetic fields on Z in several ways to create High Energy Density matter for stockpile stewardship applications

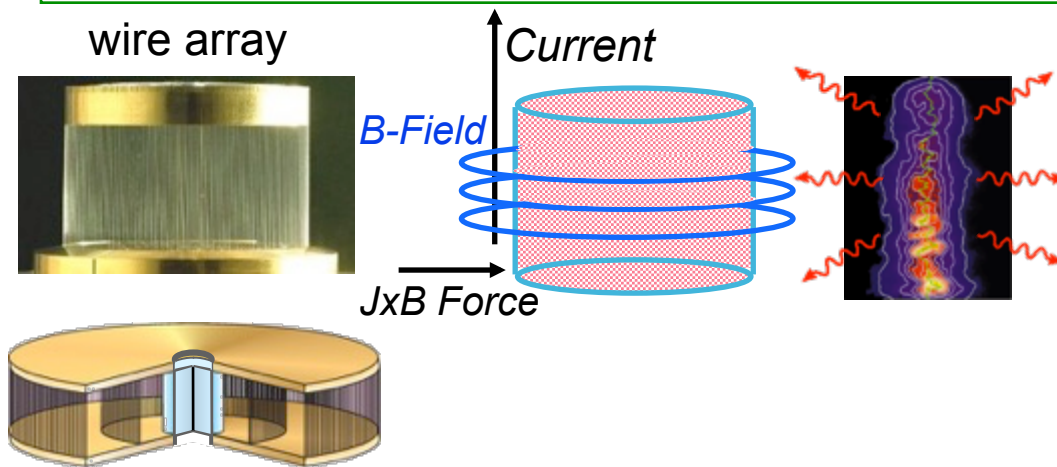
Dynamic Material Properties



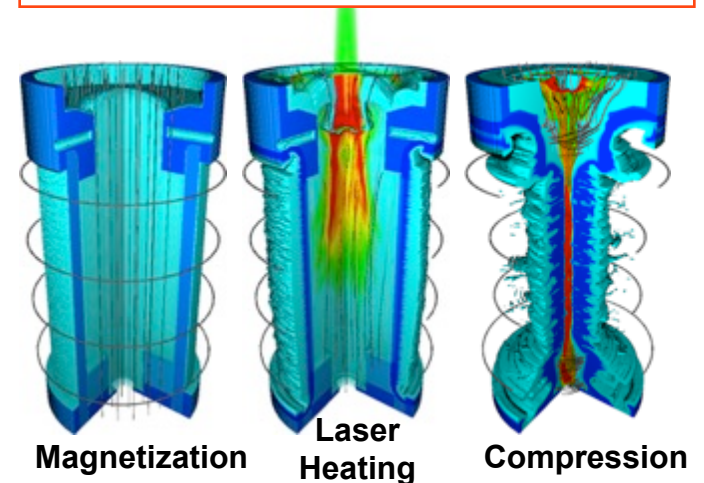
CY13 Z shot distribution



Z-Pinch X-ray Sources (RES, Rad. Physics)



Inertial Confinement Fusion

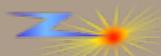
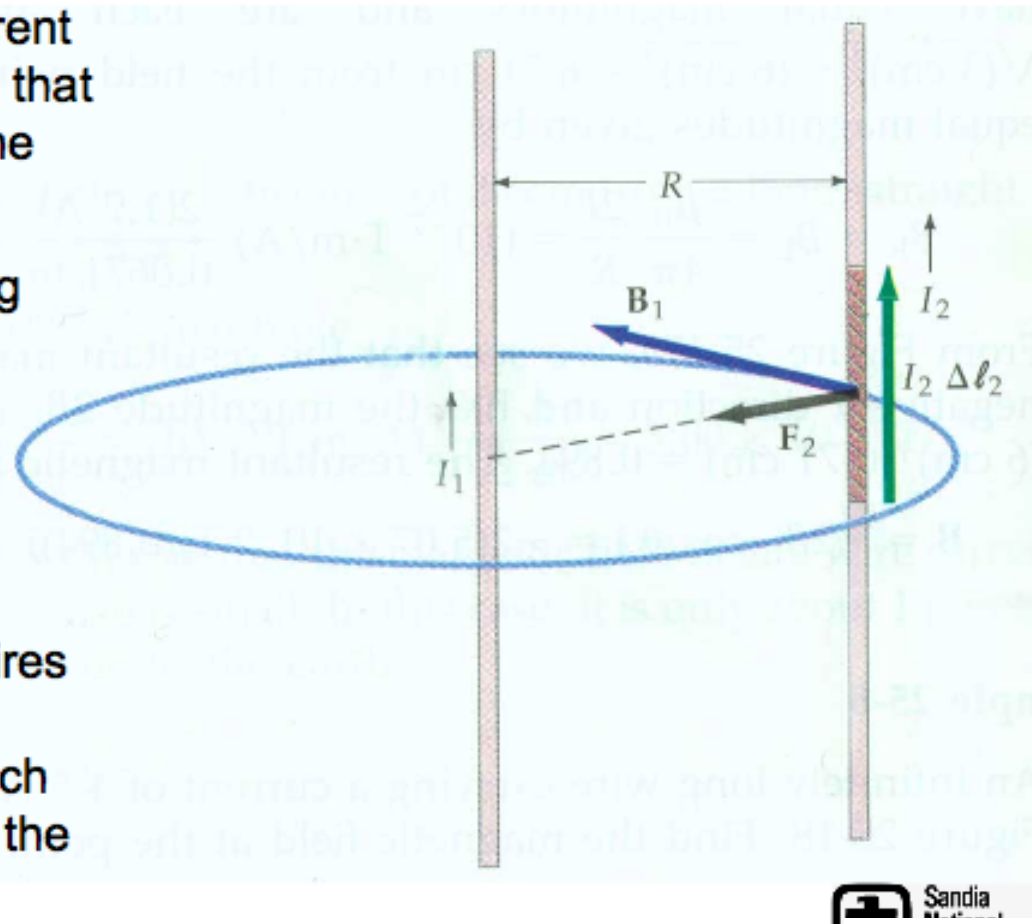


1st year physics refresher: the “pinch effect”

A single wire carrying current produces a magnetic field that encircles it according to the right-hand rule

Two parallel wires carrying current along the same direction will attract each other (Biot-Savart Law, “ $\mathbf{J} \times \mathbf{B}$ force”)

Definition of an Ampere:
If two very long parallel wires 1 m apart carry equal currents, the current in each is defined to be 1 A when the force/length is 2×10^{-7} N/m



Large currents can create large B fields!

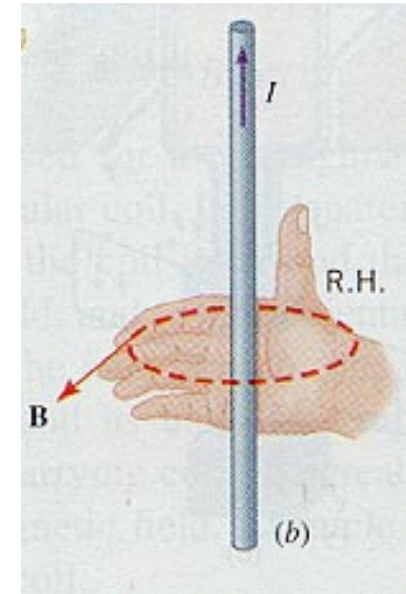
$$\nabla \times \mathbf{B} = \frac{4\pi \mathbf{J}}{c}$$

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \frac{4\pi}{c} \iint_S \mathbf{J} \cdot d\mathbf{S} \quad \text{Ampere's law}$$

For an axial current I :

$$2\pi r B_\theta = \frac{4\pi}{c} I \quad B_\theta = \frac{2}{c} \frac{I}{r} \quad (\text{cgs})$$

$$B_\theta (\text{G}) = \frac{I(\text{A})}{5 r(\text{cm})} \longrightarrow \mathbf{P_{mag} \sim B^2 \sim I^2 r^{-2}}$$



100 A at 2 mm radius is 100 G

1.0×10^7 A (**10 MA**) at **4 mm** radius is 5×10^6 G = **1 MBar** of pressure!

2.5×10^7 A (**25 MA**) at **1 mm** radius is 5×10^7 G = **100 MBar** of pressure!! **←Z Machine**
(~1000x more than high explosives)

LARGE CURRENTS → LARGE MAGNETIC FIELDS → LARGE PRESSURES!



How strong is this pressure?

$N = kg\ m\ s^{-2}$ (mks) versus $dyne = g\ cm\ s^{-2}$ (cgs)

So $1\ N = 10^5\ dyne$, and, in pressure units:

$$1\ N\ m^{-2}\ (Pa) = 10\ dyne\ cm^{-2} = 10^{-5}\ bar$$

$$P_m(dyne / cm^2) = \frac{B(G)^2}{8\pi}$$

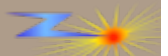
A typical refrigerator magnet is 100 gauss \sim 400 dyne/cm²

A 5000 G (0.5 T) magnetic field \sim 10⁶ dyne/cm² \sim 1 atmosphere \sim 1 Bar

A 5x10⁶ G (500 T) magnetic field \sim 1 Million atmospheres = 1 Megabar (MB)=
High energy density physics (“HEDP”)

A 5x10⁹ G (500 kT) magnetic field \sim 1 Trillion atmospheres = 1 Terabar (TB) >
pressure in the center of the sun

Note that high explosives have pressure \sim 100,000-300,000 atmospheres
 \sim 0.1-0.3 Mbar (not “HEDP”) \sim equivalent \sim 50-150 T or 5x10⁵-1.5x10⁶ G



We can incorporate the effect of magnetic fields into the plasma fluid equations as an effective pressure

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0$$

mass conservation

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla P$$

momentum conservation
(F=ma) cgs

For slowly varying fields we can approximate: $\nabla \times \mathbf{B} = \frac{4\pi \mathbf{J}}{c}$ (Ampere's law, ignoring displacement current)

We re-write JxB as: $\mathbf{J} \times \mathbf{B} = \frac{c}{4\pi} (\nabla \times \mathbf{B}) \times \mathbf{B} = -\frac{c}{4\pi} \mathbf{B} \times (\nabla \times \mathbf{B})$

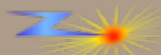
From vector identities: $\mathbf{B} \times (\nabla \times \mathbf{B}) = \frac{1}{2} \nabla (\mathbf{B} \cdot \mathbf{B}) - \mathbf{B} \cdot \nabla \mathbf{B} = \nabla \left(\frac{B^2}{2} \right) - \mathbf{B} \cdot \nabla \mathbf{B}$

So JxB becomes: $\mathbf{J} \times \mathbf{B} = \frac{c}{4\pi} \left(\mathbf{B} \cdot \nabla \mathbf{B} - \nabla \left(\frac{B^2}{2} \right) \right)$

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla P = \underbrace{\frac{1}{4\pi} \mathbf{B} \cdot \nabla \mathbf{B}}_{\text{magnetic tension}} - \nabla \left(\underbrace{P}_{\text{fluid pressure}} + \underbrace{\frac{B^2}{8\pi}}_{\text{magnetic pressure}} \right)$$

Plasma momentum is affected by **magnetic fields**

In the case of an axisymmetric z-directed current (B_θ field), the magnetic tension is zero



Large currents and the corresponding magnetic fields can be used to create and control high energy density matter

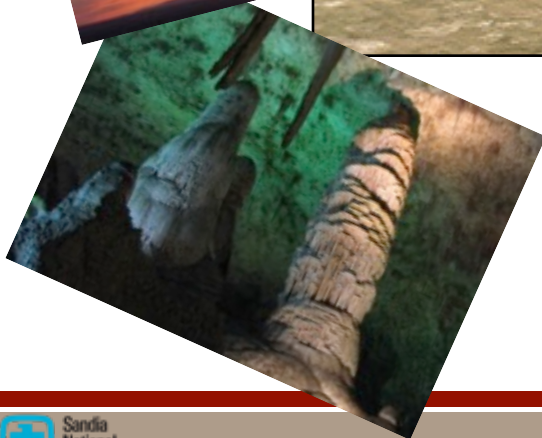
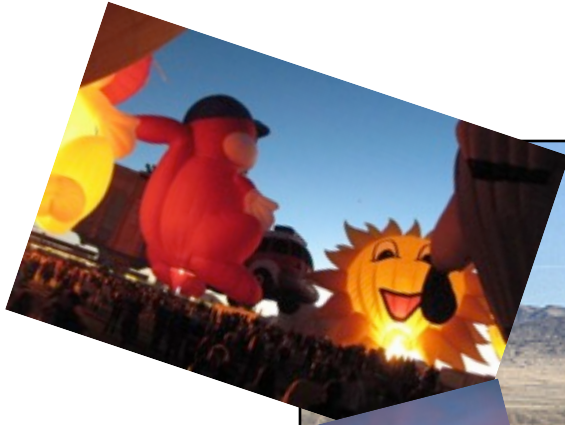
$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla P \approx \frac{1}{4\pi} \mathbf{B} \cdot \nabla \mathbf{B} - \nabla \left(P + \frac{B^2}{8\pi} \right)$$

- Magnetic pressure can be created efficiently over cm³ volumes, enabling large samples and energetic sources
- Magnetic pressure increases in a converging geometry
- High pressures can be created *without making material hot*
- Magnetic fields can be generated over long time scales with significant control over the time history (pulse shaping)
- Energy density is increased through planar compression or cylindrical implosion
- Magnetic fields change the way particles and energy are transported in a plasma

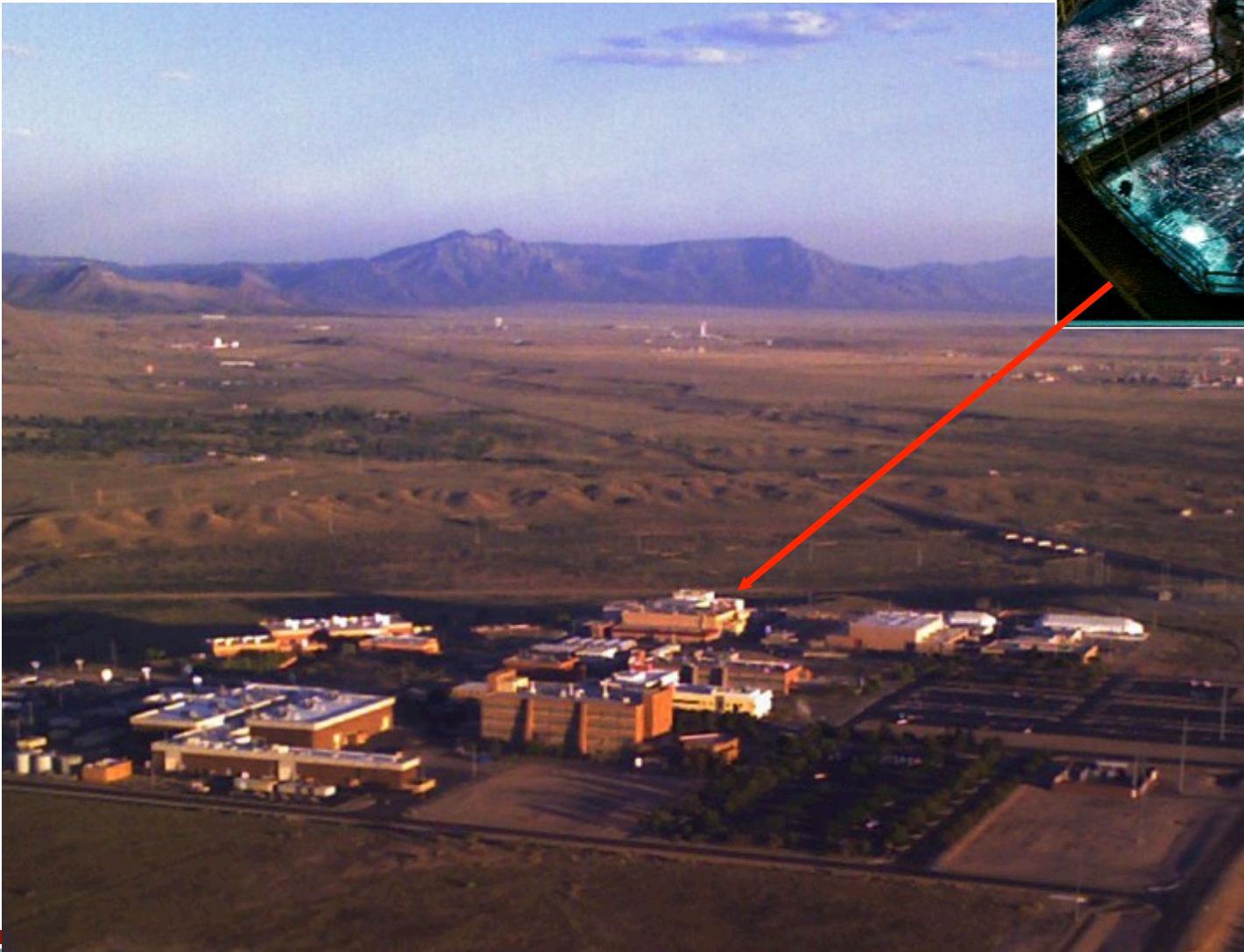


The “Z” pulsed-power facility is located at Sandia National Laboratories in Albuquerque, New Mexico

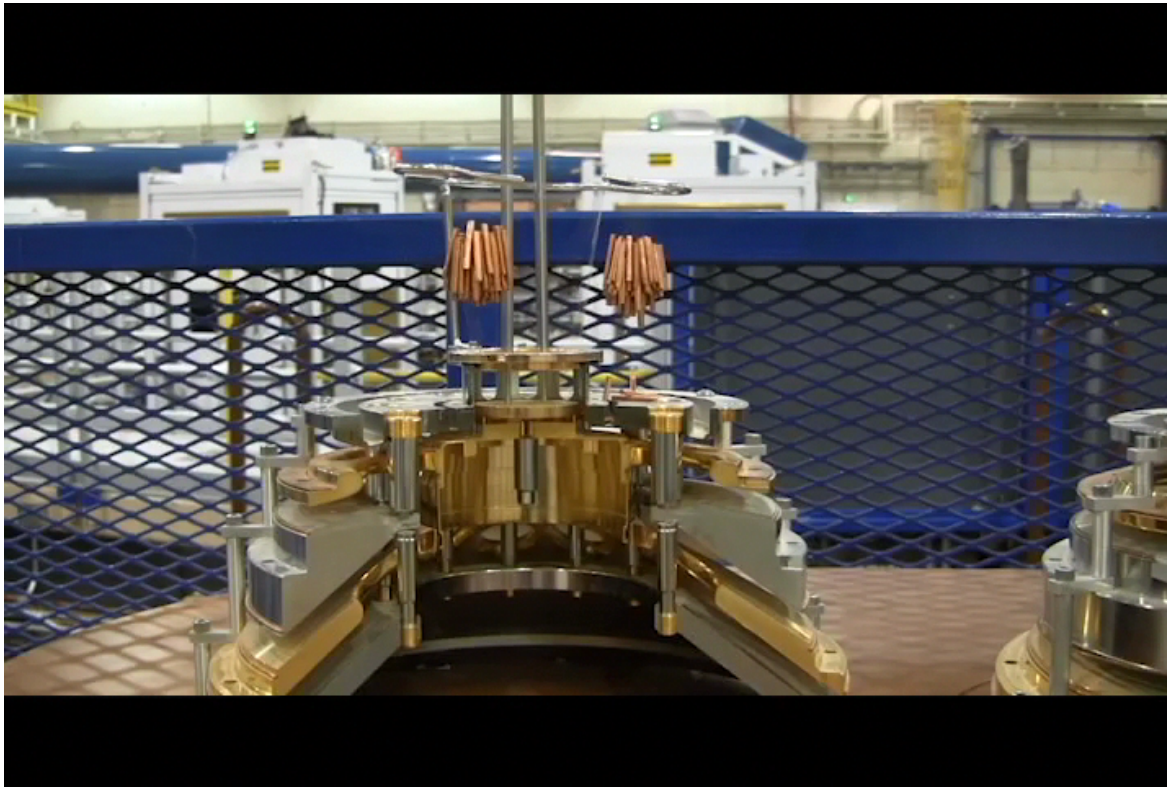
Youtube.com: search for the BBC TV show:
“Horizon: Can we make a star on earth?”



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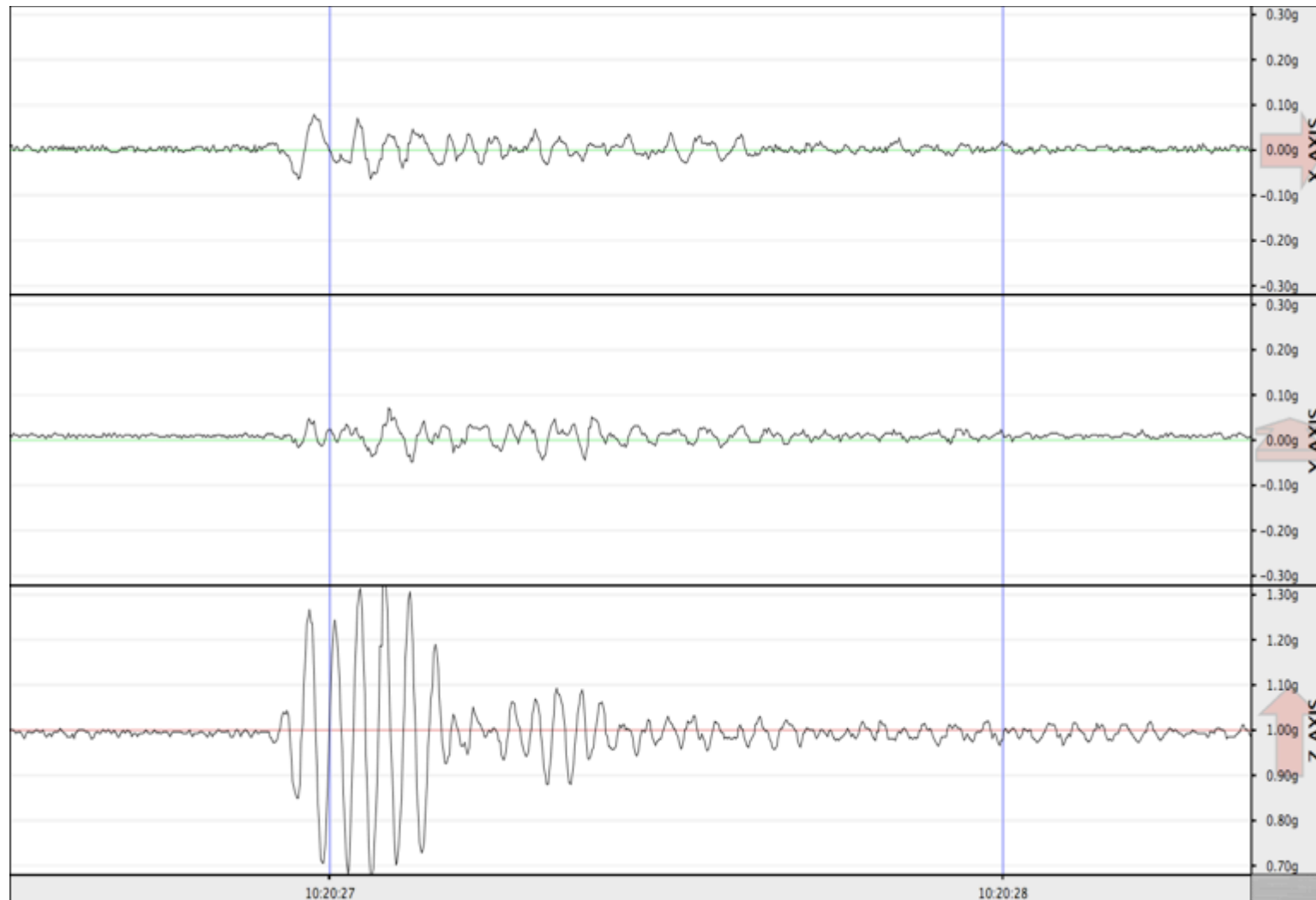


Z poses a challenging environment for experiments



- 30 – 100 g instantaneous shock at center

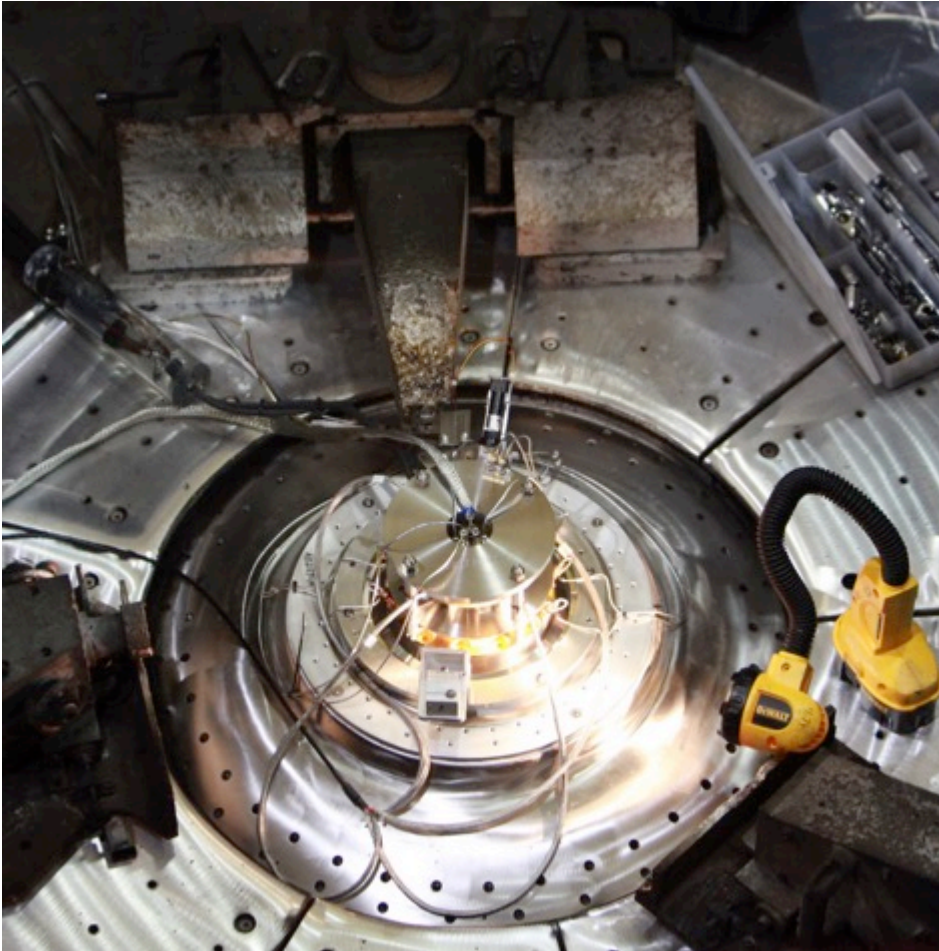
The ground shock is $\pm 1g$ of acceleration over 200 ms



Seismac program



Debris from all experiments must be carefully managed (several MJ energy release equivalent to few sticks of dynamite)



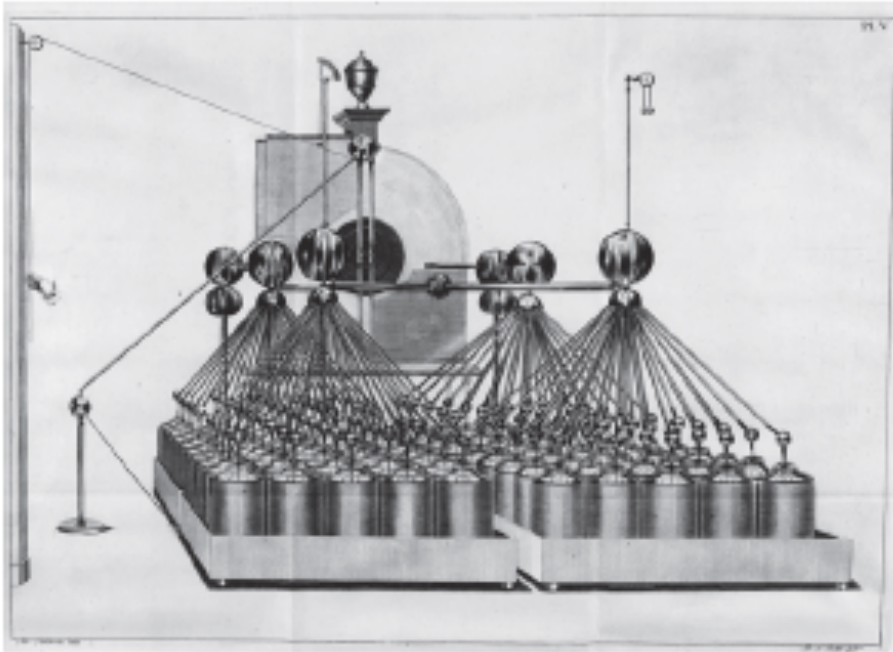
Pre-shot photo of coils & target hardware



Post-shot photo

MagLIF Experiment

The earliest z-pinch experiment on record



Martinus van Marum

Location: Amsterdam

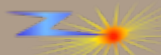
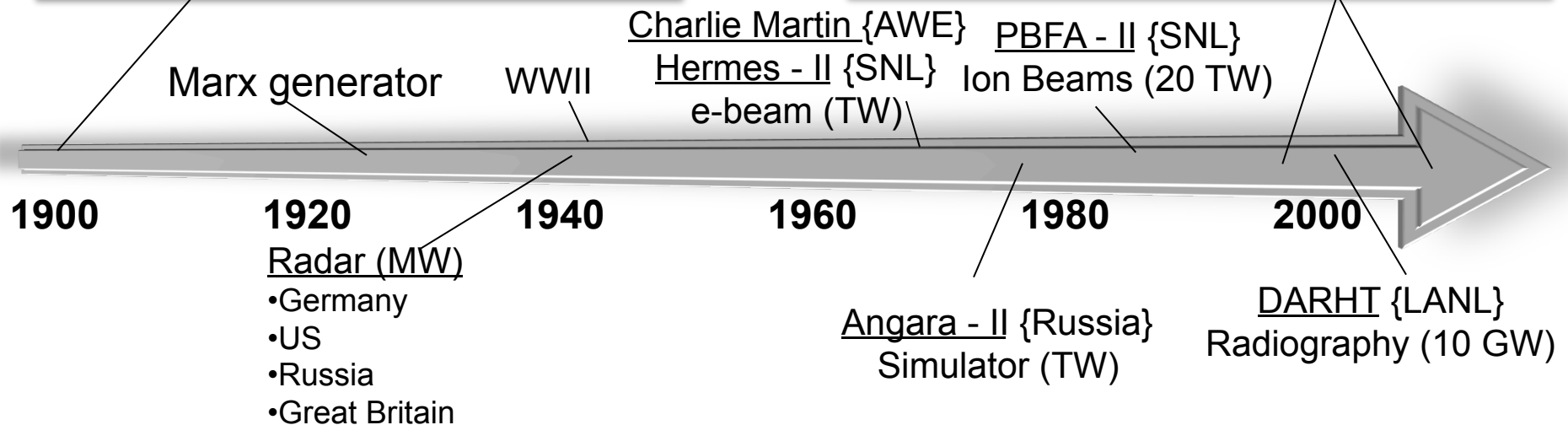
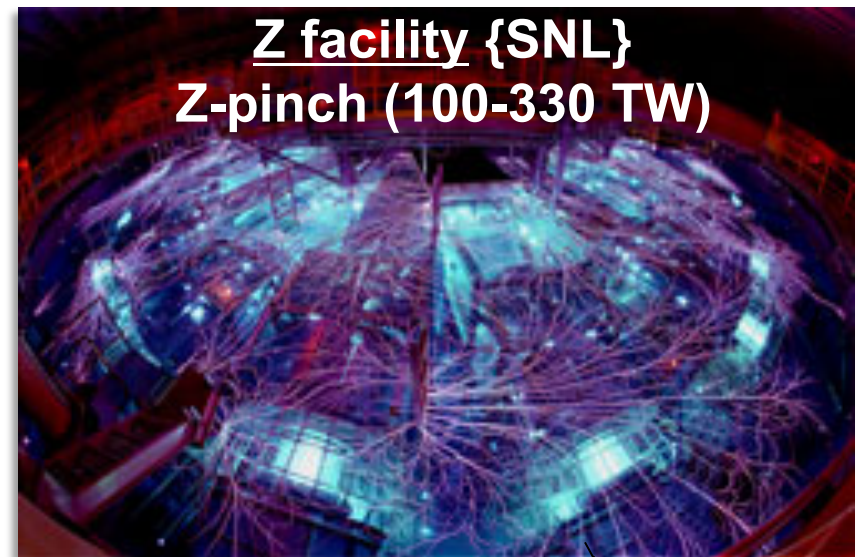
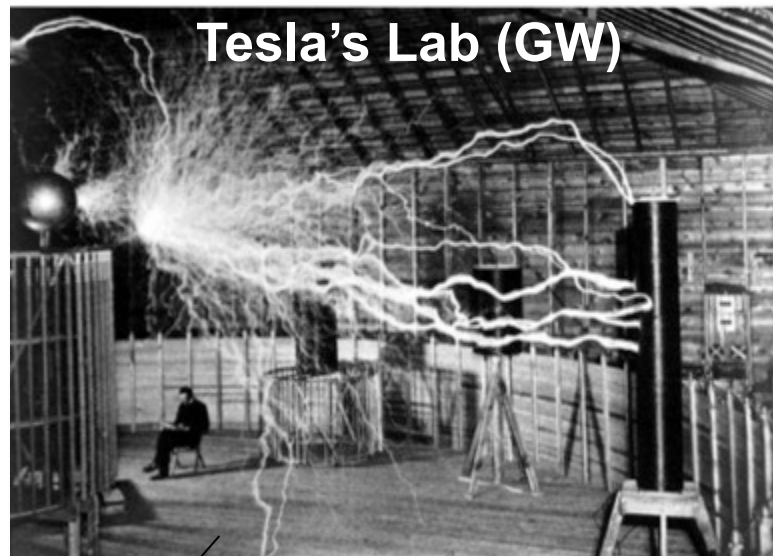
Date: 1790

Parameters:
1 kJ energy storage

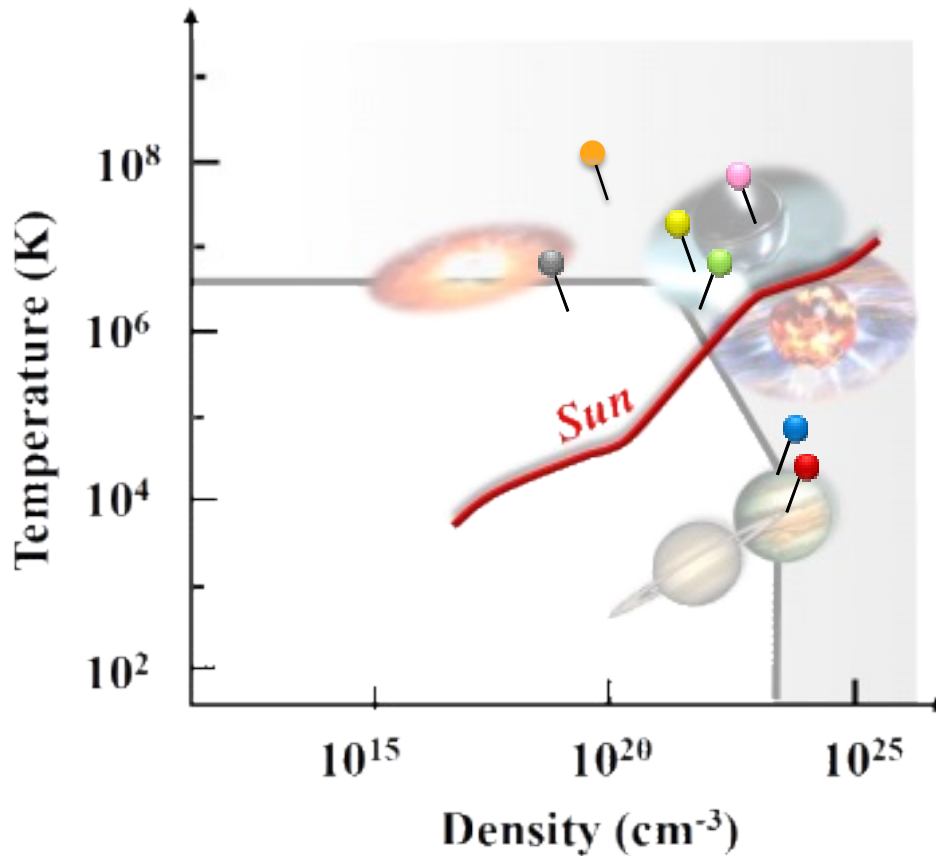
Technology:
100 Leyden jars

Application:
explosion and vaporization of 1 m
long wire

The accumulation and transmission of electromagnetic energy, called “pulsed power”, has been investigated for more than a century



HED experiments on Z address issues of fundamental importance



Diamond at 10 Mbar

Knudson et al., Science 322 (2008)

D₂ EOS at 1 Mbar

Knudson et al., PRL 87 (2001)

Photoionized Plasmas

Foord et al., PRL 94 (2004)

Radiating Shocks

Rochau et al., PRL 100 (2008)

Opacity at $T_e > 150$ eV

Bailey et al., PRL 99 (2007)

Fusing Plasmas

Bailey et al., PRL 93 (2004)

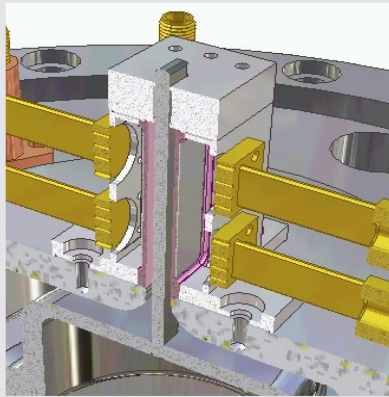
High-Z Ion Radiating Plasmas at 40 Mbar

Ampleford et al., submitted PoP (2013)

There are multiple ways to use the current on Z

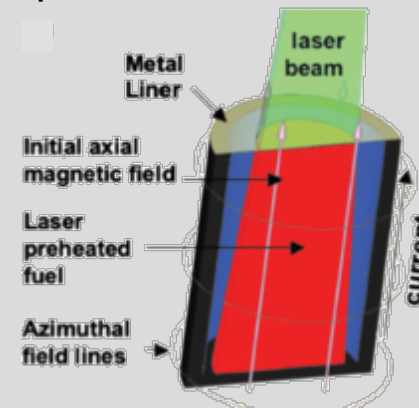
Shock Compression

Velocities > 40 km/s



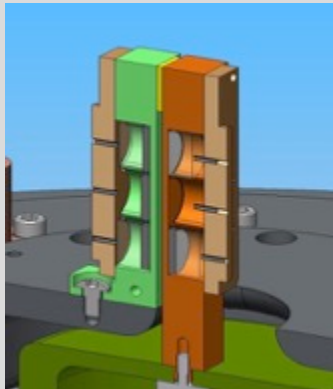
Cylindrical compression

- Fusion (P up to 1 Gbar)
- EOS up to 6.5 Mbar



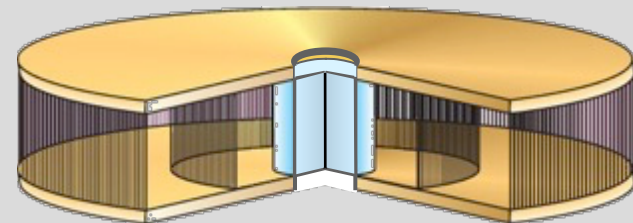
Planar ramp-compression

continuous compression up to ~4 Mbar

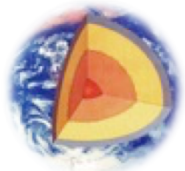


Z-pinch x-ray sources

X-ray power up to 330 TW, 2.2 MJ



The stockpile stewardship is divided into a broad portfolio of “science campaigns” with diverse requirements



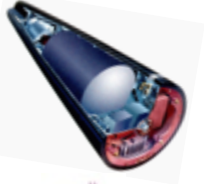
Dynamic Materials Properties

Materials of interest at pressure and temperatures of interest



Inertial Confinement Fusion

Researching into creating “self-heating or burning plasmas”



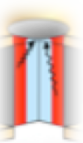
Radiation Effects Sciences

Study the impact of radiation on circuits and components



Primary Assessment

Simulate and study process of importance in NW primaries



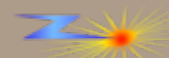
Secondary Assessment

Simulate and study process of importance in NW secondaries



Fundamental Science

Use the conditions to study basic science questions in astrophysics, geophysics, plasma, atomic physics



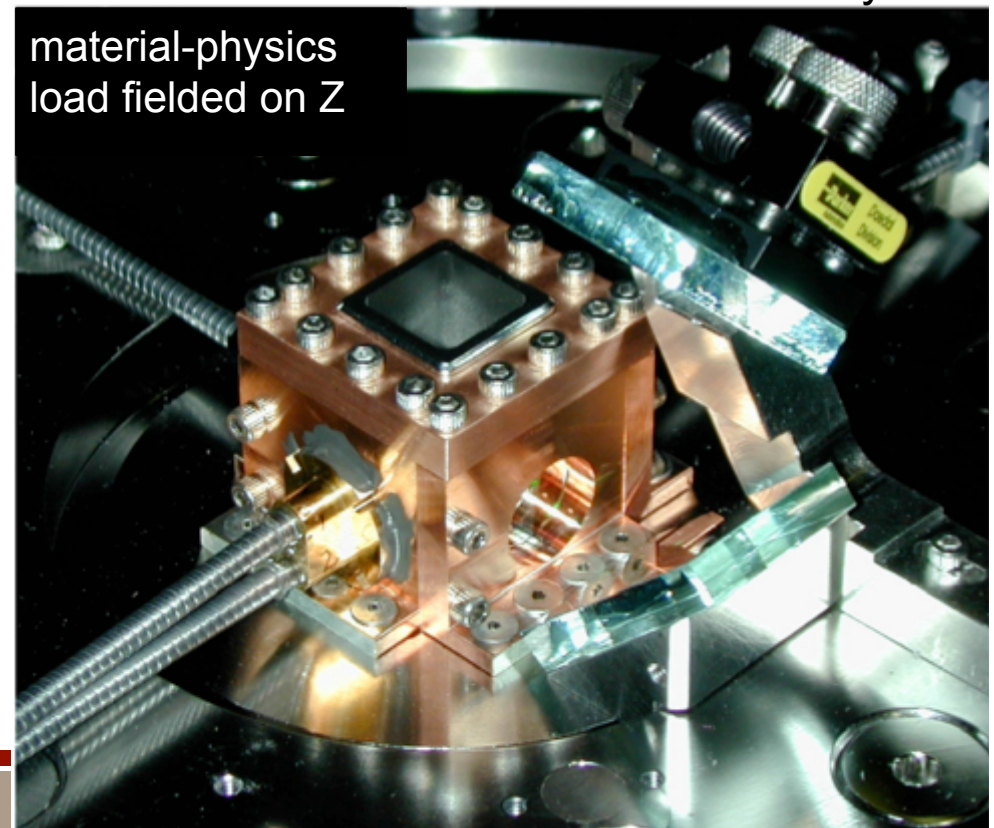
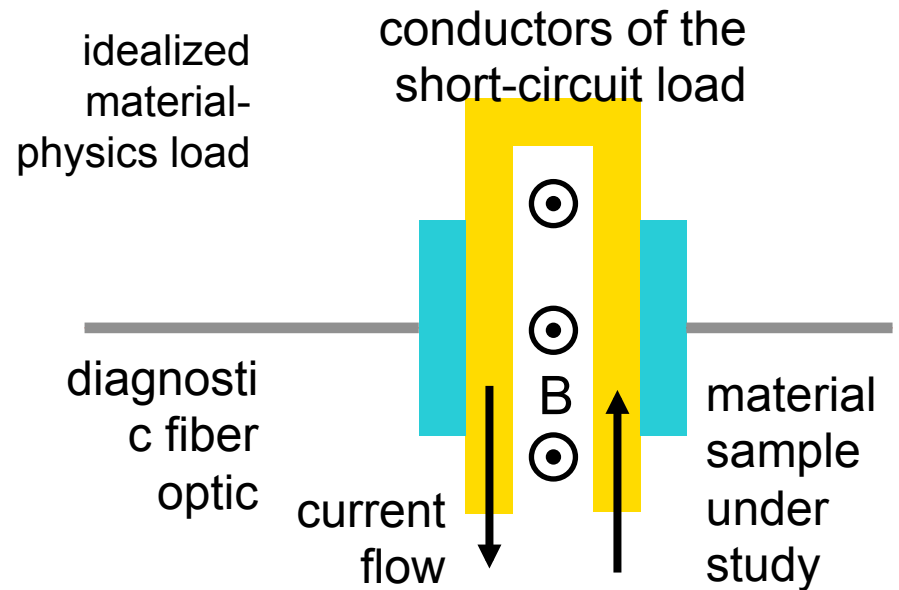
Z and other pulsed-power machines are used to drive material-physics experiments

- This application is outlined in seminal publications by Reisman and colleagues (JAP, 2001) and Hall and co-workers (RSI, 2001).
- The magnetic pressure generated within a short-circuit load drives the experiment.

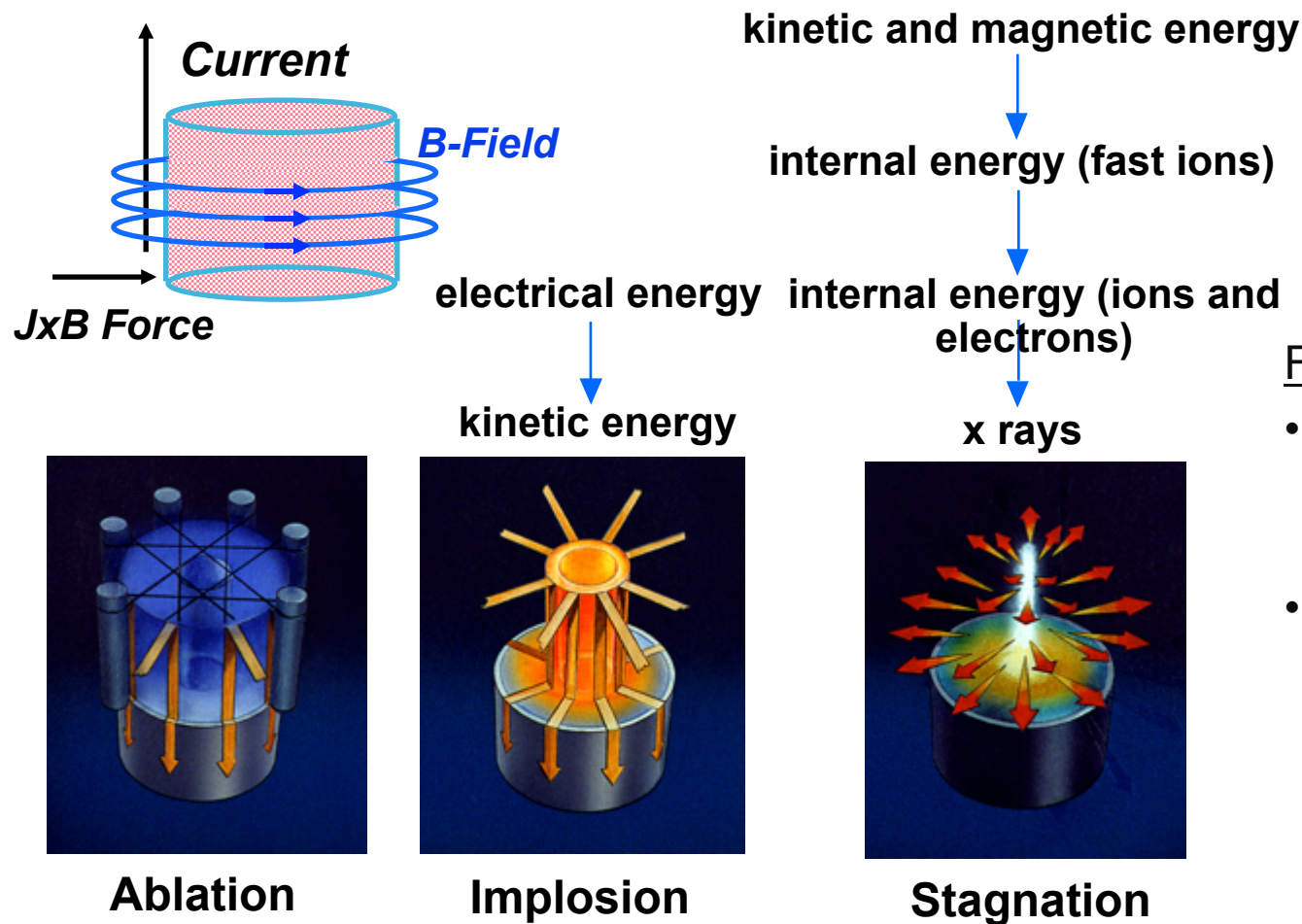
$$P_{\text{magnetic}} = \frac{\mu_0 I^2}{2 w^2}$$

I = current

w = width of the conductor



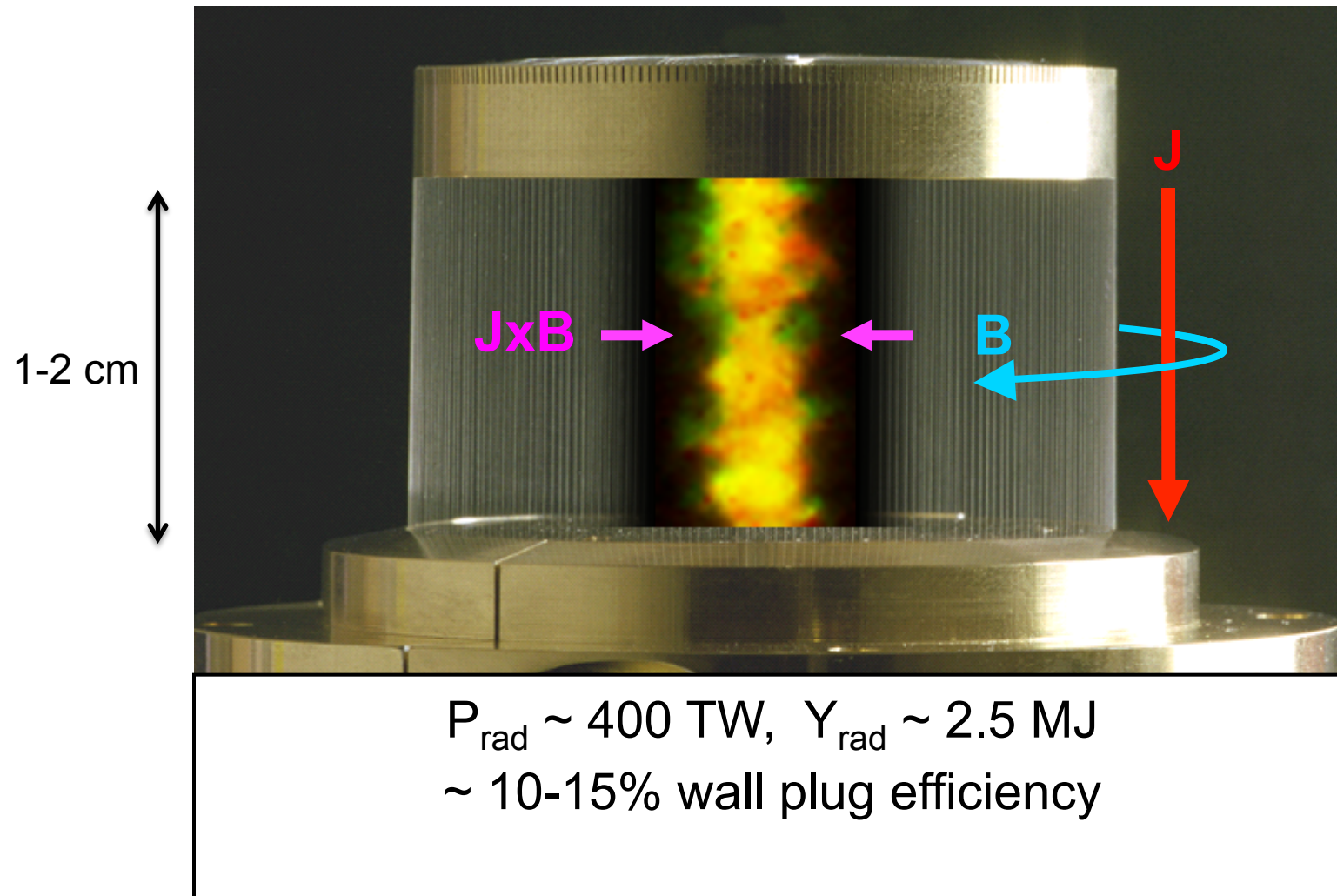
Magnetically-driven fast z-pinch implosions efficiently convert electrical energy into radiation



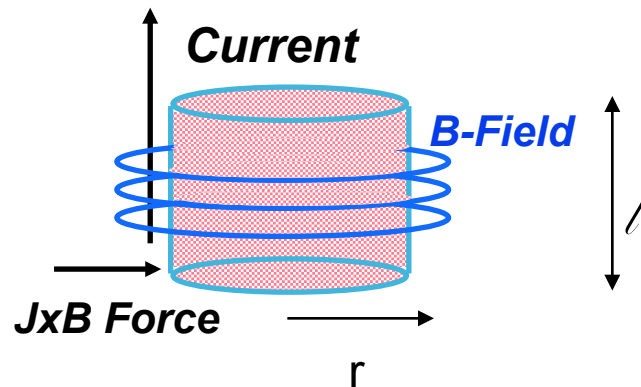
Fast wire z-pinch loads:

- Z-pinchs are imploded in 60-120 ns, and radiate x-rays in 5 ns
- Energy: x-ray ~ 15% of stored electrical

Magnetically driven implosions are efficient, powerful, x-ray sources from 0.1 to 10 keV



How much magnetic energy can we put into kinetic energy ?



According to the MHD equations we can treat the magnetic field working on the wires as a magnetic pressure $P_{\text{mag}} \sim B^2$

For a thin shell with all the current on the outside:

MKS units

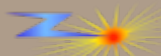
$$ma = F$$

$$m(d^2r/dt^2) = PA \quad A = 2\pi r\ell \quad P = \frac{B^2}{2\mu_0} \quad B = \frac{\mu_0 I}{2\pi r}$$

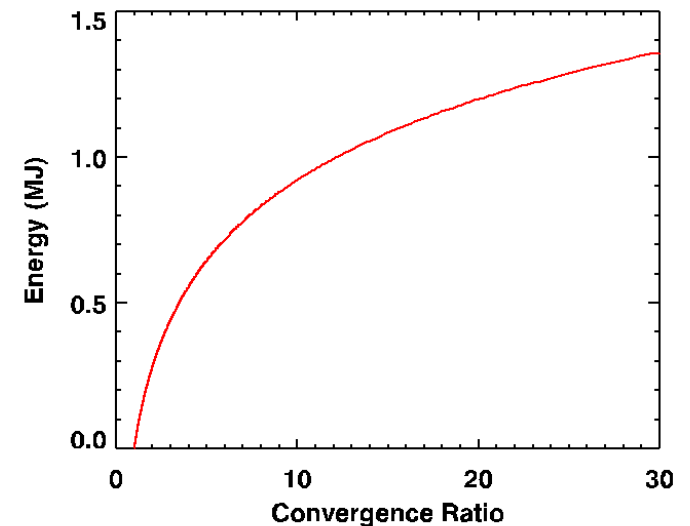
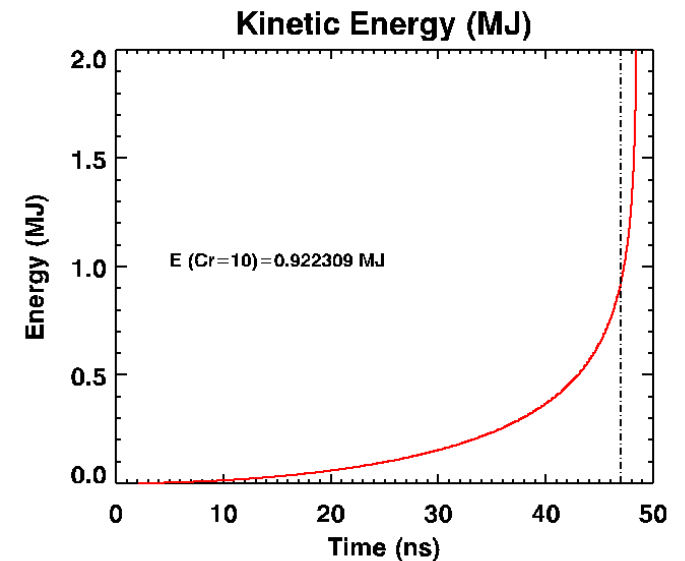
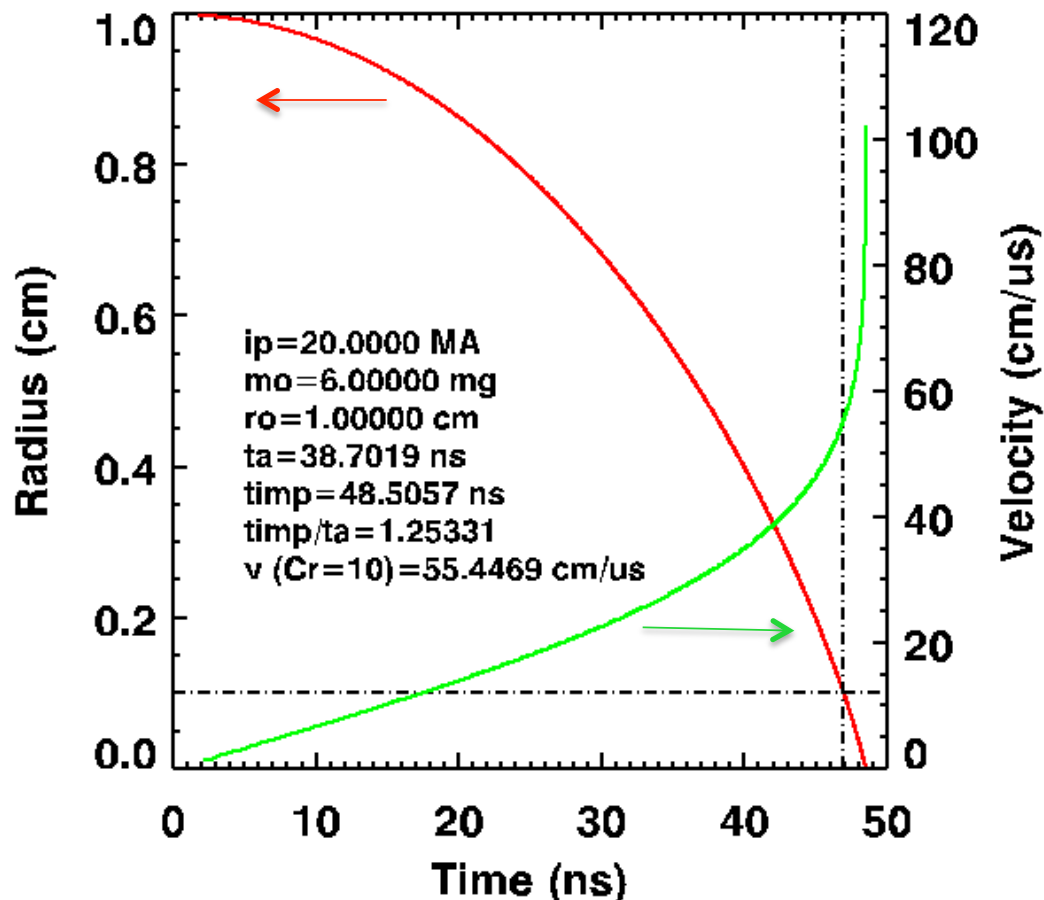
$$\frac{m}{\ell} \frac{d^2r}{dt^2} = -\frac{\mu_0}{4\pi} \frac{I^2}{r}$$

Acceleration increases with I^2 , but it also **increases** during implosion (r^{-1}) !

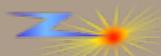
We can integrate this equation to get the kinetic energy given to the imploding shell



An analytic solution to the equation of motion with $I \sim \text{constant}$ shows rising magnetic pressure accelerates the pinch



- Velocity at convergence ratio of 10 ~ 50 cm/ μ s ~ 500 km/s



We employ energies of ~ 1 MJ in every day objects



An energy of 1 MJ:

- Kinetic energy of F150 at ~ 60 mph
- 0.48 x energy in a stick of dynamite
- 100 W incandescent light bulb uses 1 MJ in 4.3 hours

A velocity of 500 km/sec:

- $\sim 1,100,000$ miles per hour
- New York to LA in ~ 8 seconds
- $1/600$ speed of light

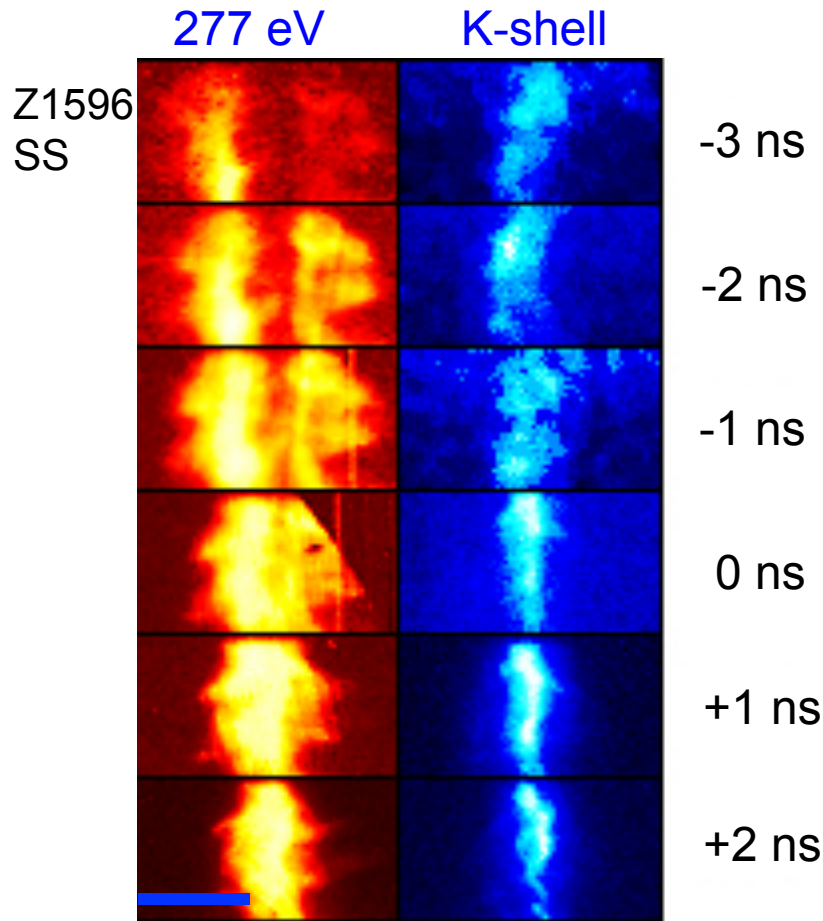
- $m_{\text{F150}} = 2950 \text{ kg}$
- $v_{\text{F150}} = 94 \text{ km/hour (58 mph)}$
- $E = 1 \text{ MJ}$
- In a typical z-pinch, this 1 MJ is released in 5 ns

Self-emission x-ray imaging shows that magnetically driven implosion generates large volumes (10 to 100 mm³) of hot plasma

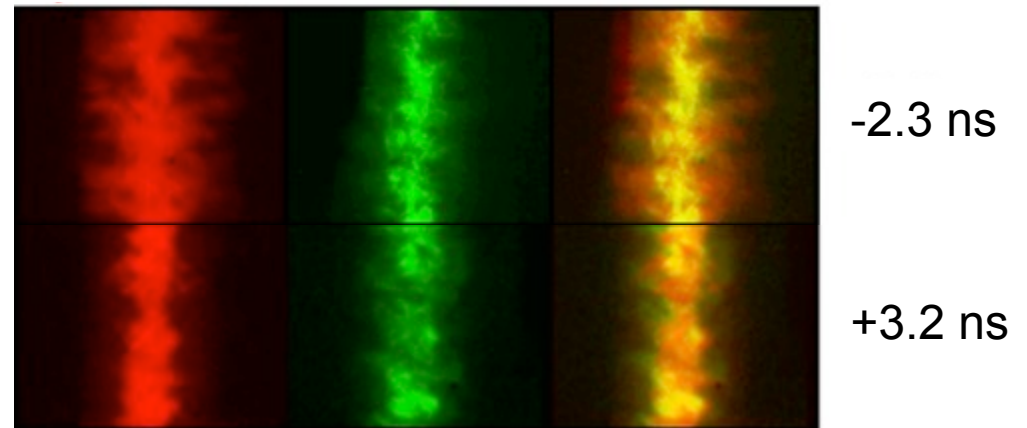
Z1519, Al

277 eV

1.7 keV Al
K-shell Overlay



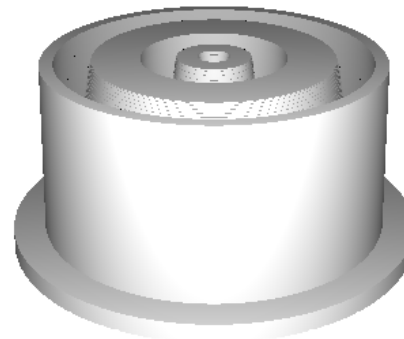
5 mm



- As noted previously,
 - Evidence of instability at stagnation
 - Observed structure varies with atomic number
- Hot, dense column emitting K-shell on axis with colder material still imploding

Numerical modeling tools are mature and complement Z gas puff experiments

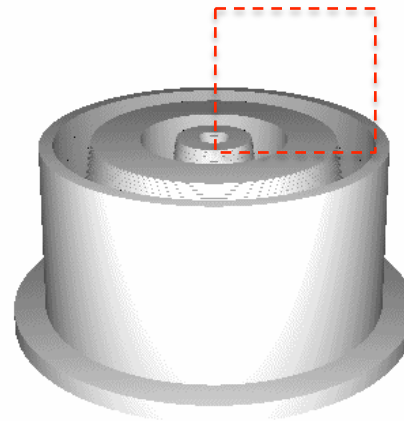
- Cold gas flow models may be validated using experimental interferometer data
- Benchmarked simulated profiles can be used to initiate MHD simulations
- Tabulated atomic data are used to estimate K-shell x-ray outputs
- Pre-shot NRL modeling [Thornhill *et al.*, HEDP **8**, 197 (2012)] was consistent with SNL Gorgon simulations (Jennings)



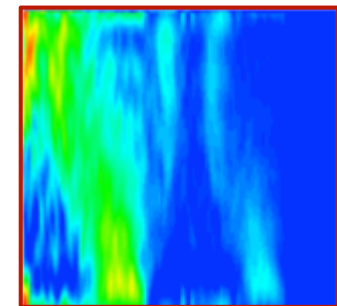
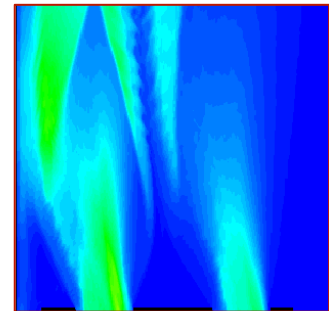
C.A. Jennings, SNL, GORGON

Numerical models are being used to design experiments and benchmarked post-shot to gain physics insight

- Cold gas flow models may be validated using experimental interferometer data
- Benchmarked simulated profiles can be used to initiate MHD simulations
- Tabulated atomic data are used to estimate K-shell x-ray outputs
- Pre-shot NRL modeling [Thornhill *et al.*, HEDP **8**, 197 (2012)] was consistent with SNL Gorgon simulations (Jennings)



Simulation

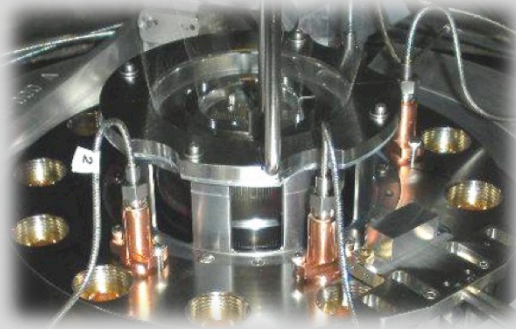


Interferometer

C.A. Jennings, SNL, GORGON

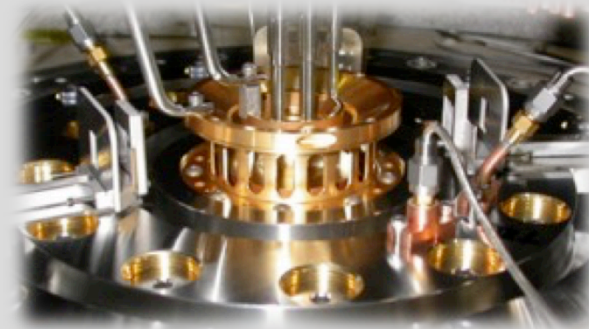
Many x-ray experiments use just one of two types of routine, established wire array z-pinch sources

Large-diameter K-shell sources



- Diameter: 4-8 cm
- Height: 2 cm
- Radial Beamlines: 9
- Axial Beamlines: 1
- Optimized for K-shell emission
- Spectrum/Power/Energy depends on wire material: Al, SS, Cu

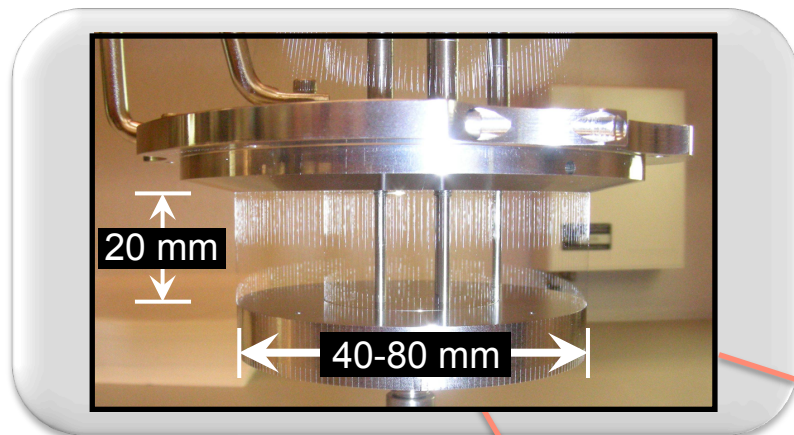
Z-pinch dynamic hohlraum



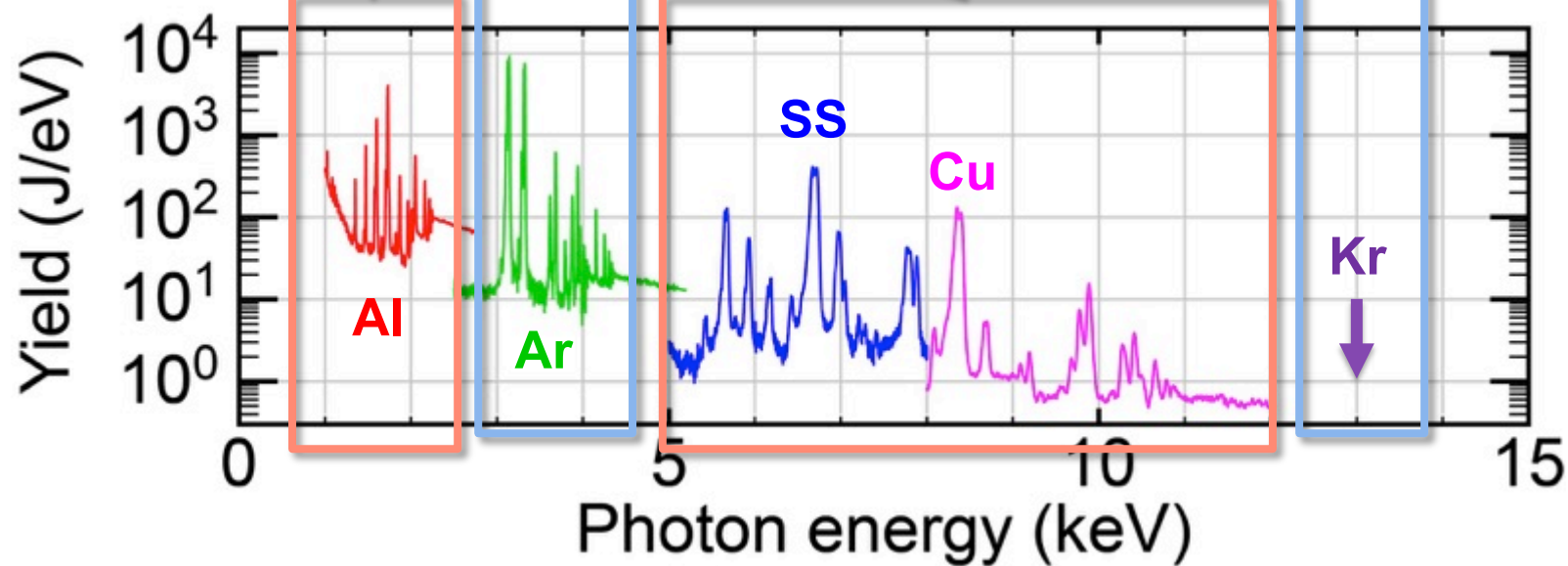
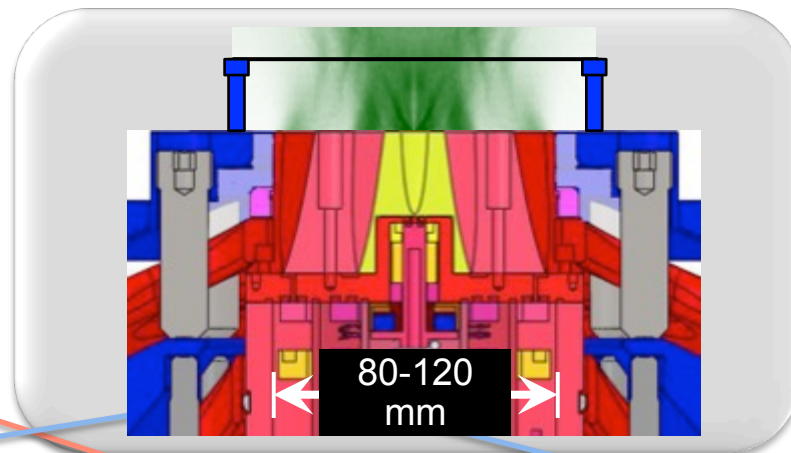
- Diameter: 4 cm
- Height: 1.2 cm
- Radial Beamlines: 9 or 18
- Axial Beamlines: 1
- Near-Planckian emission spectrum
- Broadband Energy Emission: 1 MJ
- Peak Power: 170 TW
- Power FWHM: 3 ns

Magnetically-driven implosions to $\sim 500\text{--}1000\text{ km/s}$ are used to access K-shell emissions from different materials

Nested Wire Arrays

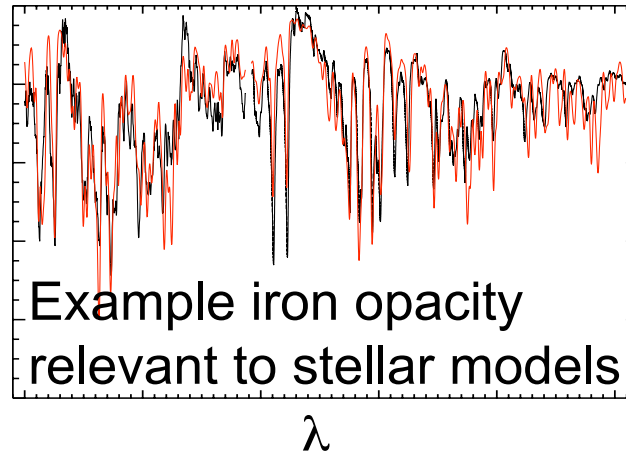


Structured Gas Puffs

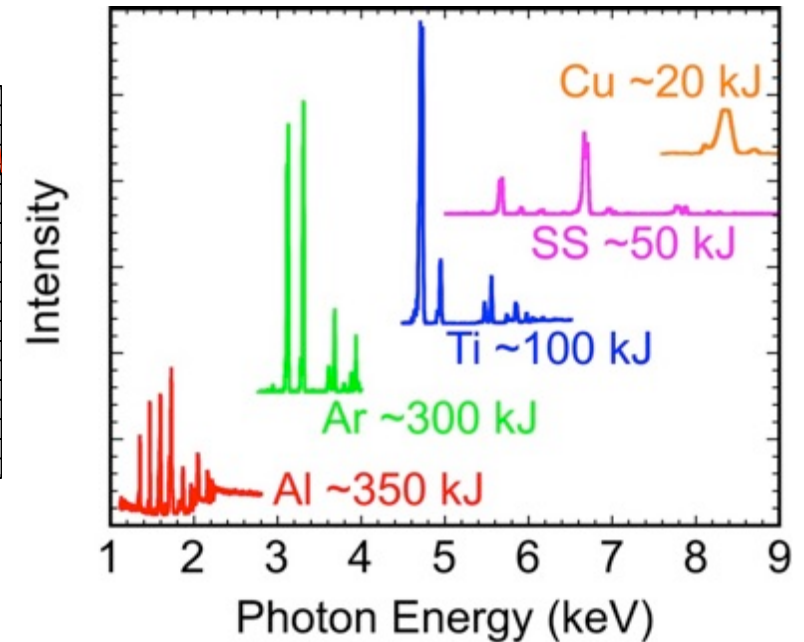


Magnetically-driven x-ray sources are being used for a variety of fundamental and applied science applications

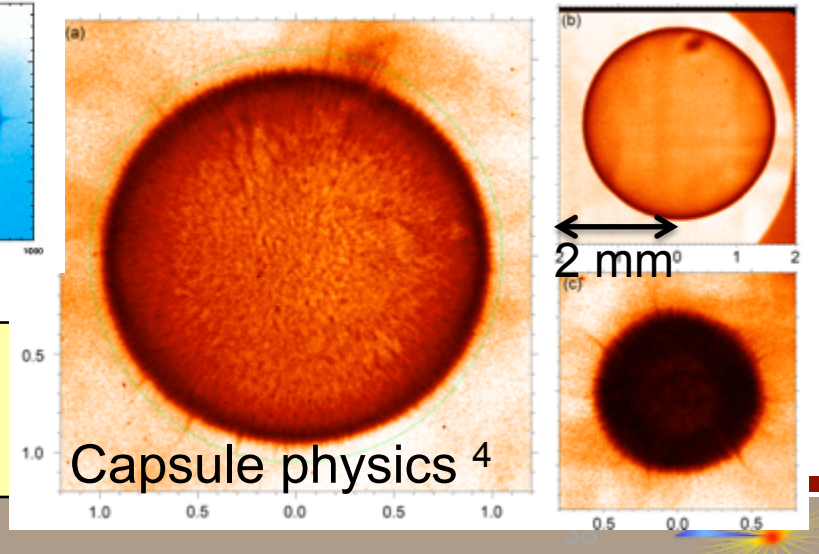
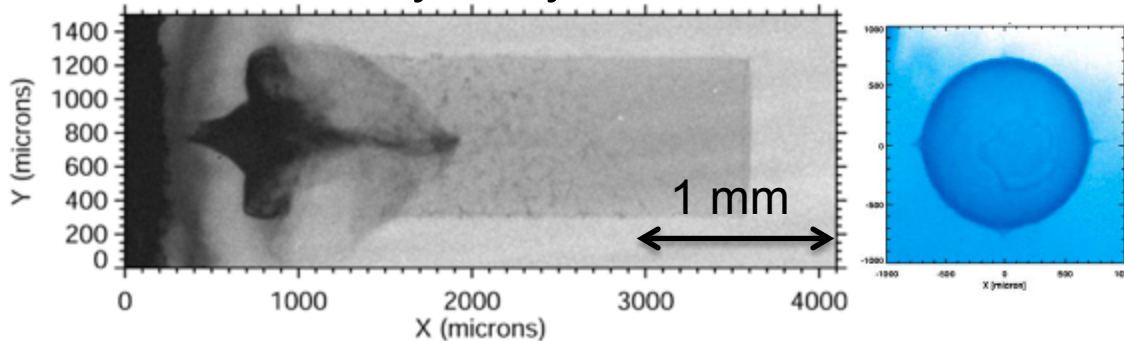
Opacity Measurements and Photo-ionized plasmas ¹



K-Shell Sources ³



Radiation Hydrodynamics ²

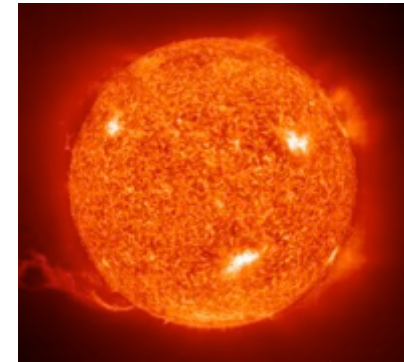
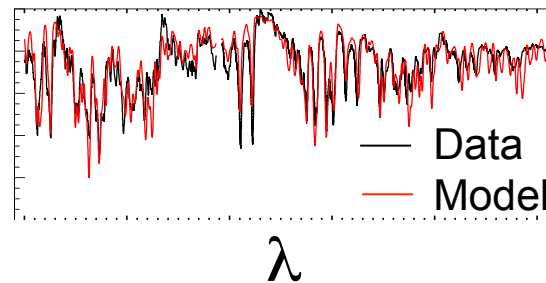


- ¹ J.E. Bailey, et al, Phys. Rev. Lett. 99, 265002 (2007).
- ² G.R. Bennett, et al., Phys. Rev. Lett. 205003 (2007).
- ³ C. A. Coverdale et al., IEEE T. Plas. Sci. 35, 582 (2007).
- ⁴ M.E. Cuneo et al., IEEE T. Plas. Sci. 40, 3222 (2012).

Magnetically-driven x-ray sources are being used for a variety of fundamental and applied science applications

Solar opacity at $T_e > 150\text{-}170$

Bailey et al., PRL 99 (2007)



White dwarf physics

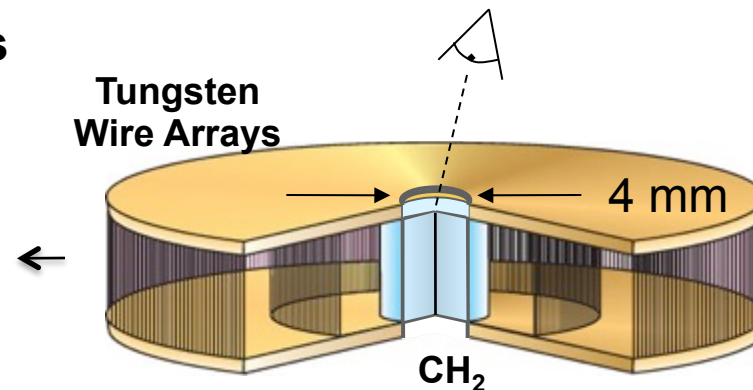
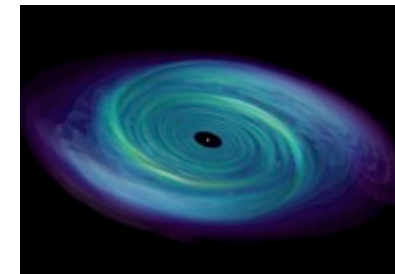
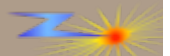


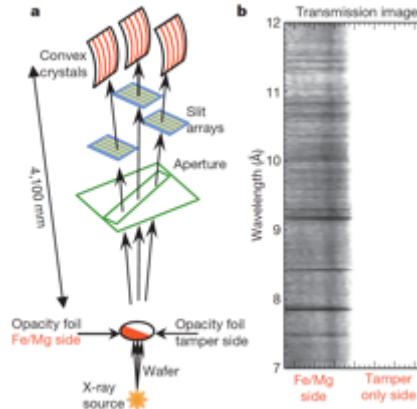
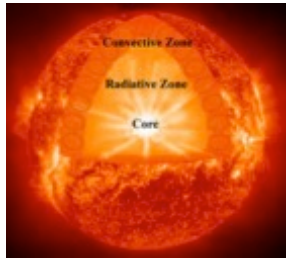
Photo-ionized plasmas



- 1 μg of stellar interior at $R \sim 0.7R_{\text{sol}}$
- 10^{-3} liters of accretion disk at $R \sim 100 - 1000$ km from black hole
- ~ 0.1 liters of white dwarf photosphere
- weapons science



We are making fundamental discoveries in Astrophysics and Planetary Science within the Z Fundamental Science Program



A higher-than-predicted measurement of iron opacity at solar interior temperatures

- Jim Bailey, et. al., Nature 517, 14048 (2015)
- Sandia, Ohio State University, University of Nevada Reno, PRISM, LANL, LLNL, and CEA
- High-precision measurements at the almost unreachable conditions of the radiation/convection boundary in the sun

A higher opacity of iron at solar conditions explains contradictions between helioseismic observations and established solar models

LETTER

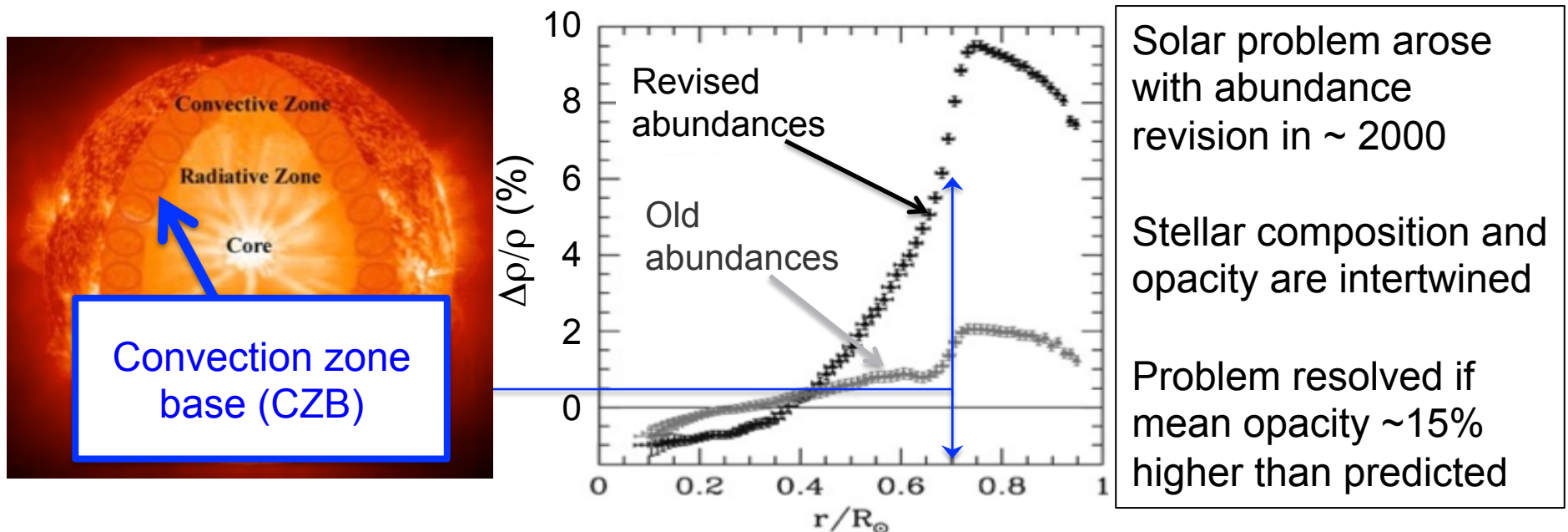
doi:10.1038/nature14048

A higher-than-predicted measurement of iron opacity at solar interior temperatures

J. E. Bailey¹, T. Nagayama¹, G. P. Loisel¹, G. A. Rochau¹, C. Blancard², J. Colgan³, Ph. Cosse², G. Faussurier², C. J. Fontes³, F. Gilleron², I. Golovkin⁴, S. B. Hansen¹, C. A. Iglesias⁵, D. P. Kilcrease³, J. J. MacFarlane⁴, R. C. Mancini⁶, S. N. Nahar⁷, C. Orban⁷, J.-C. Pain², A. K. Pradhan⁷, M. Sherrill³ & B. G. Wilson⁵

Z iron opacity experiments refine our understanding of the sun, stars, and laboratory HED systems

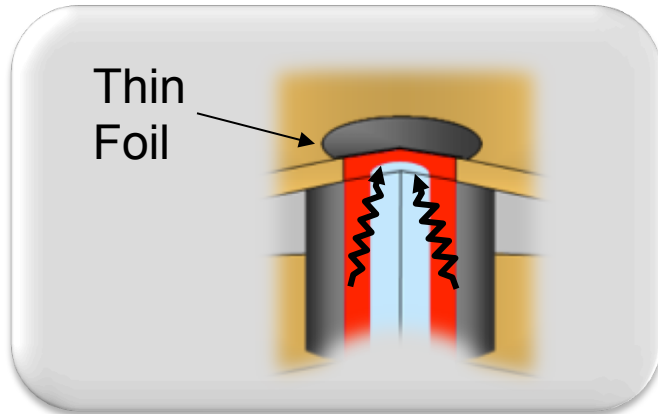
- Opacity is the fundamental property that controls photon absorption in matter
- Opacity model uncertainty has been proposed as an explanation for the decade old problem that solar models do not match helioseismic observations



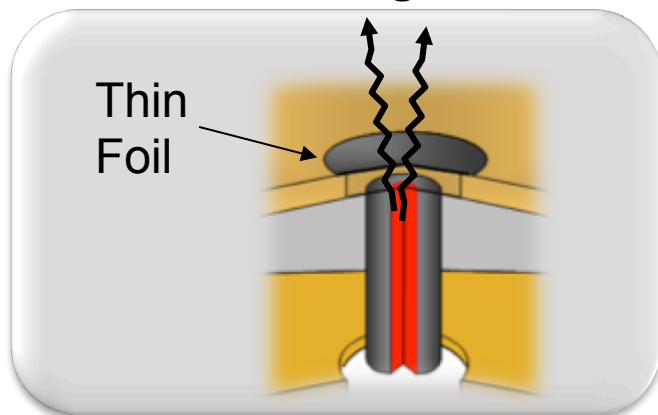
- **Conditions at the solar CZB : 182 eV, $9e22$ electrons/cc**
- **Similar conditions exist in many astrophysical and laboratory plasmas**
- **Thus, solar opacity model refinements have far-reaching implications**

The ZPDH radiating shock is used to both heat and backlight samples to stellar interior conditions.

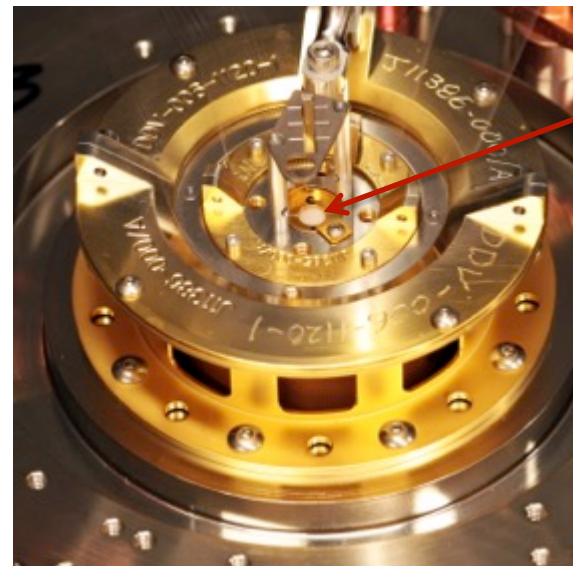
Foil is heated during the ZPDH implosion



Foil is backlit at shock stagnation

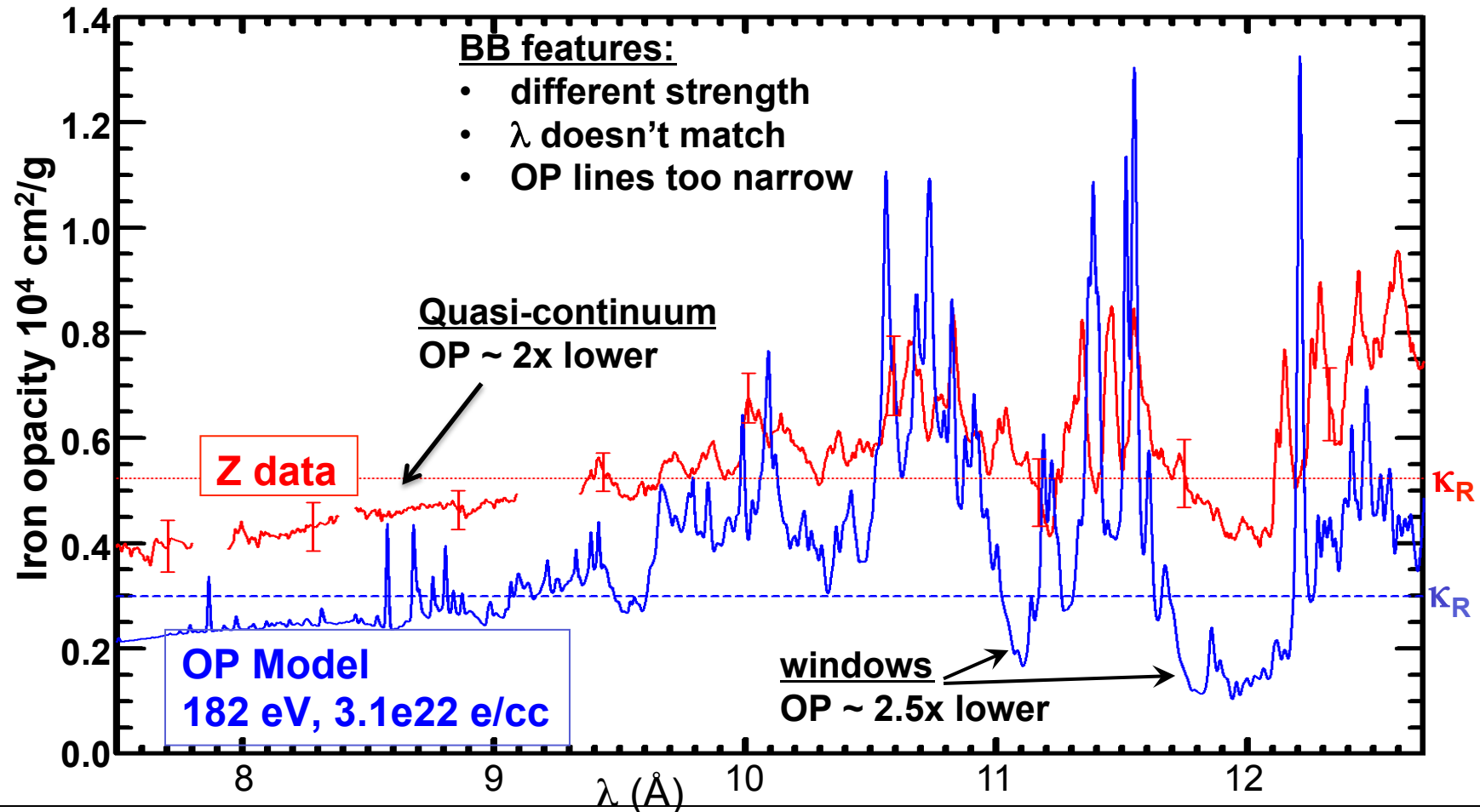


Bailey et al., *Nature* 2015



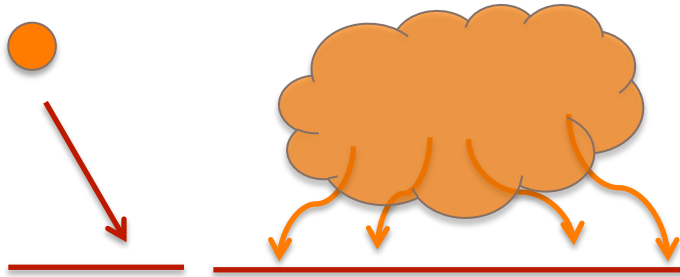
opacity sample

Opacity model discrepancy at solar-interior conditions implies photon absorption in HED matter is different than previously believed



- Calculated Rosseland mean for iron is 1.75x lower than measured
- This difference accounts for roughly half the opacity change needed to resolve the solar problem

We are making fundamental discoveries in Astrophysics and Planetary Science within the Z Fundamental Science Program



Iron rain following a meteor impact explains the iron-enriched mantle of the earth and a key earth/moon difference

Impact vaporization of planetesimal cores in the late stages of planet formation

- Richard D. Kraus, et. al., Nature Geoscience, DOI: 10.1038/NGEO2369 (2015)
- Sandia, Harvard, UC Davis, and LLNL
- Multi-Mbar dynamical material experiments to measure properties of vaporized iron at conditions of planetary impacts

nature
geoscience

LETTERS

PUBLISHED ONLINE: 2 MARCH 2015 | DOI: 10.1038/NGEO2369

Impact vaporization of planetesimal cores in the late stages of planet formation

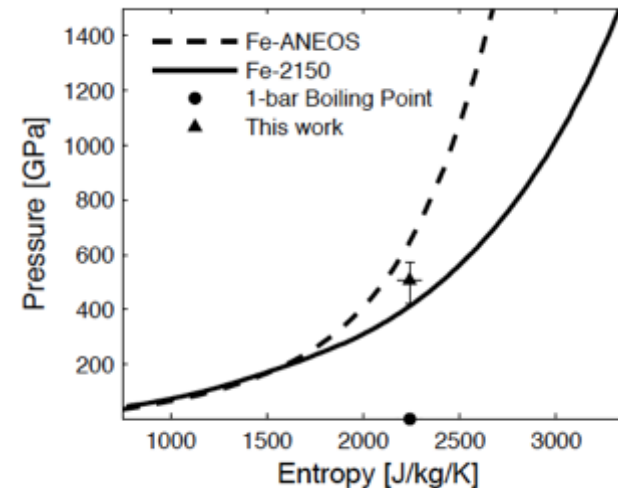
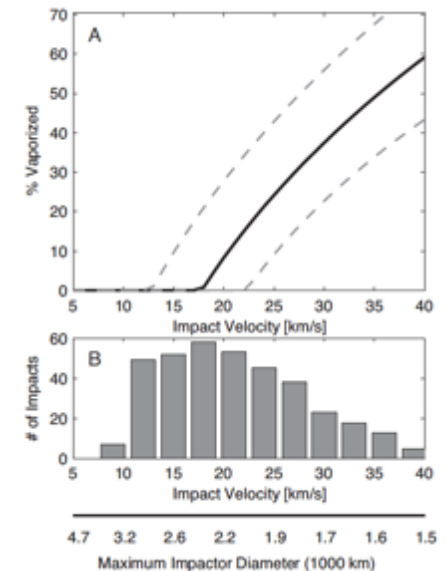
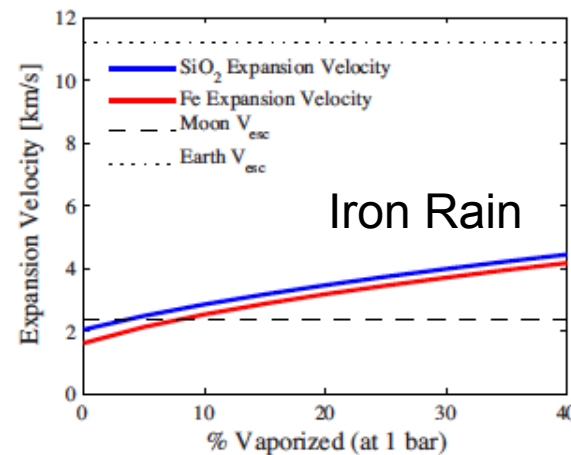
Richard G. Kraus^{1,2★}, Seth Root³, Raymond W. Lemke⁴, Sarah T. Stewart^{1,5}, Stein B. Jacobsen¹ and Thomas R. Mattsson⁴

Planetesimal impacts played an important role in shaping the solar system

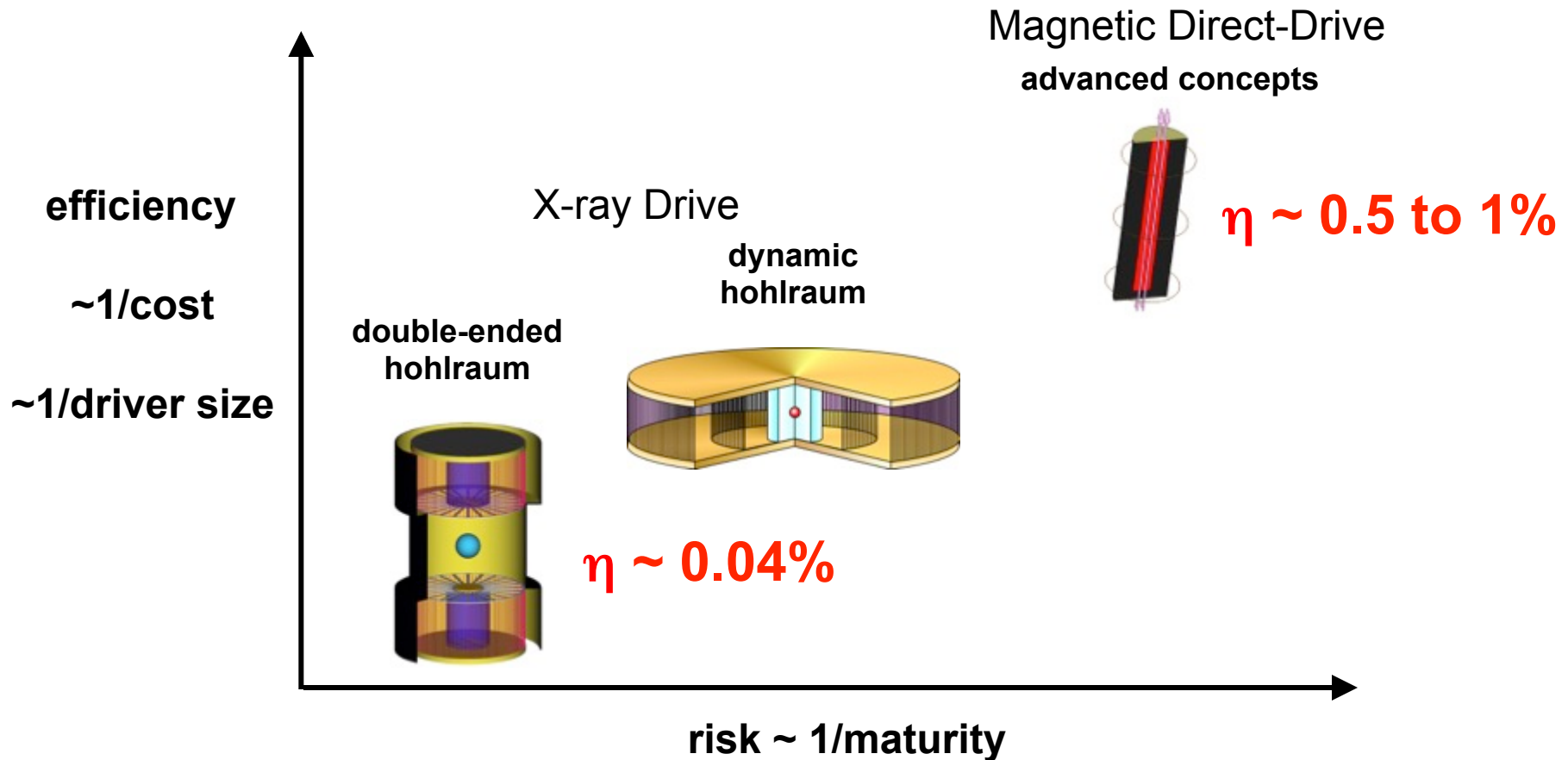


R. G. Kraus et al., Nature Geoscience
04/2015; 8(April):269-272

A significantly lower shock pressure (500 GPa) is required to vaporize iron upon release than previously used in planetary modeling (890 GPa)

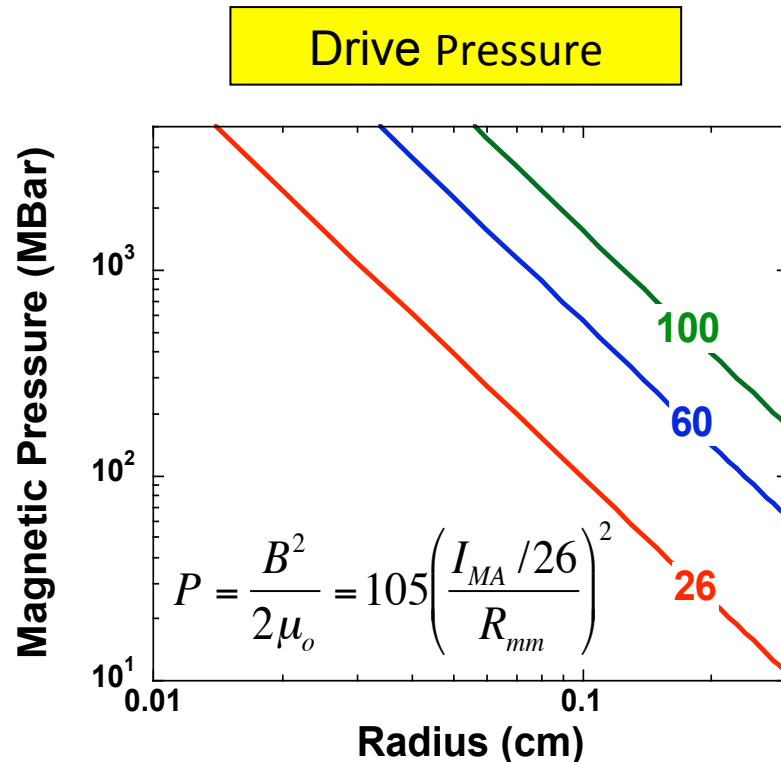


Are there more efficient pulsed power methods for heating and compressing fusion fuel?

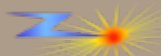


- Pulsed power can flexibly drive many target types
- Direct fuel compression and heating with the magnetic field could be up to 20X more efficient

Direct-magnetically-driven targets efficiently absorb large energies

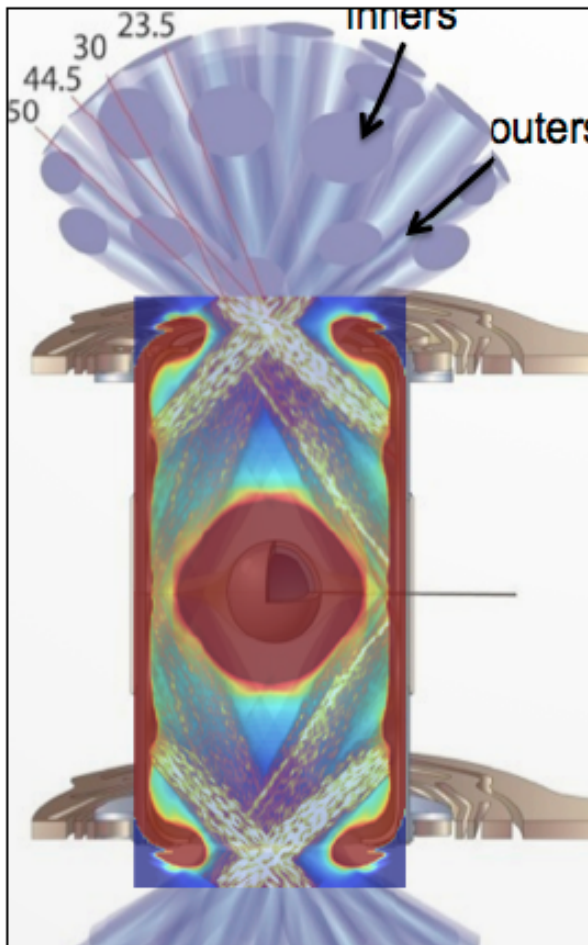


- Large drive pressures:
 - 100 MBar on Z at 26 MA
 - >500 MBar at > 60 MA
- Efficient. No energy is wasted on:
 - Heating target
 - Hohlraum heating
 - Conversion to x-rays
- Energy rich:
 - 100kJ into fusion fuel (~0.5%) on Z
 - 7 MJ into target (~5%) at 60 MA

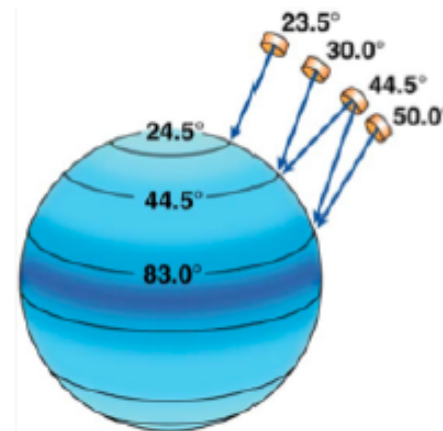
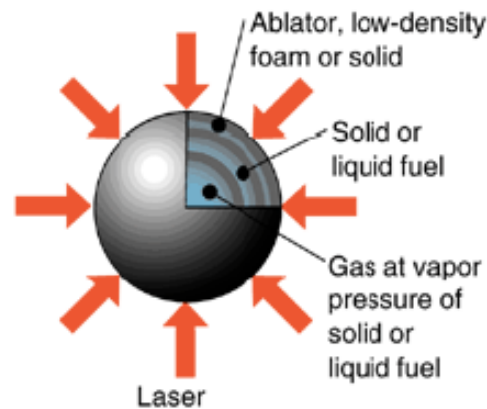


The United States ICF program has focused on three main approaches

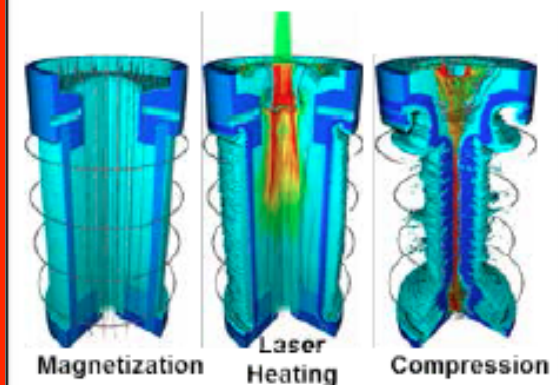
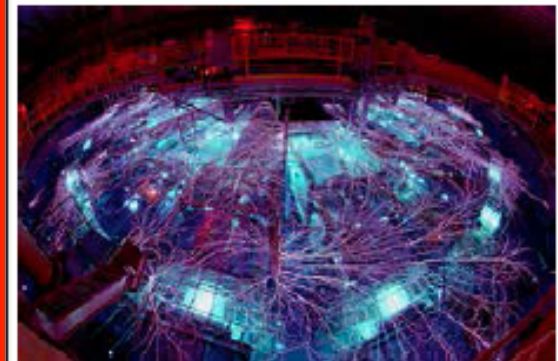
Radiation-driven implosions



Laser-driven implosions

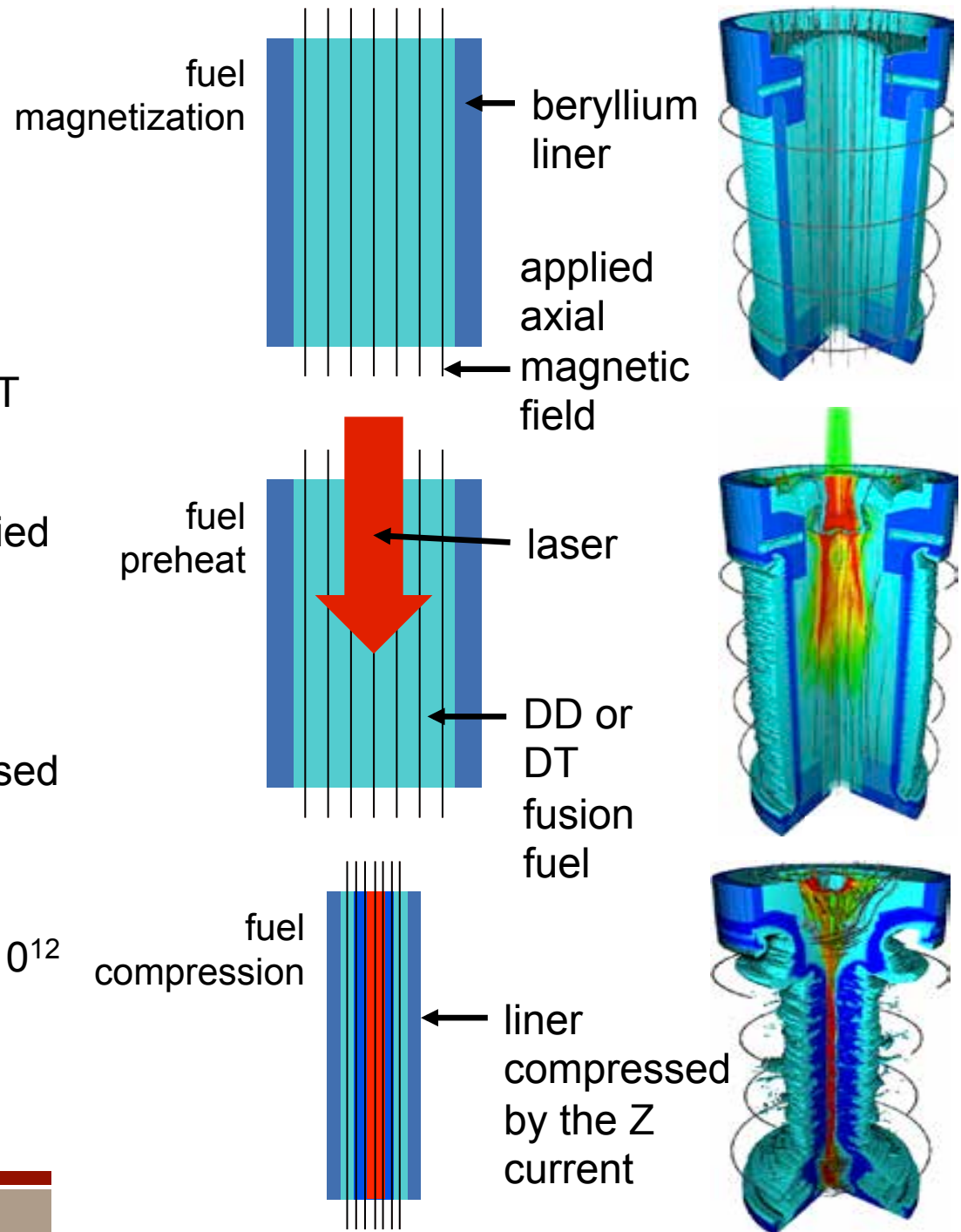


Magnetically-driven implosions

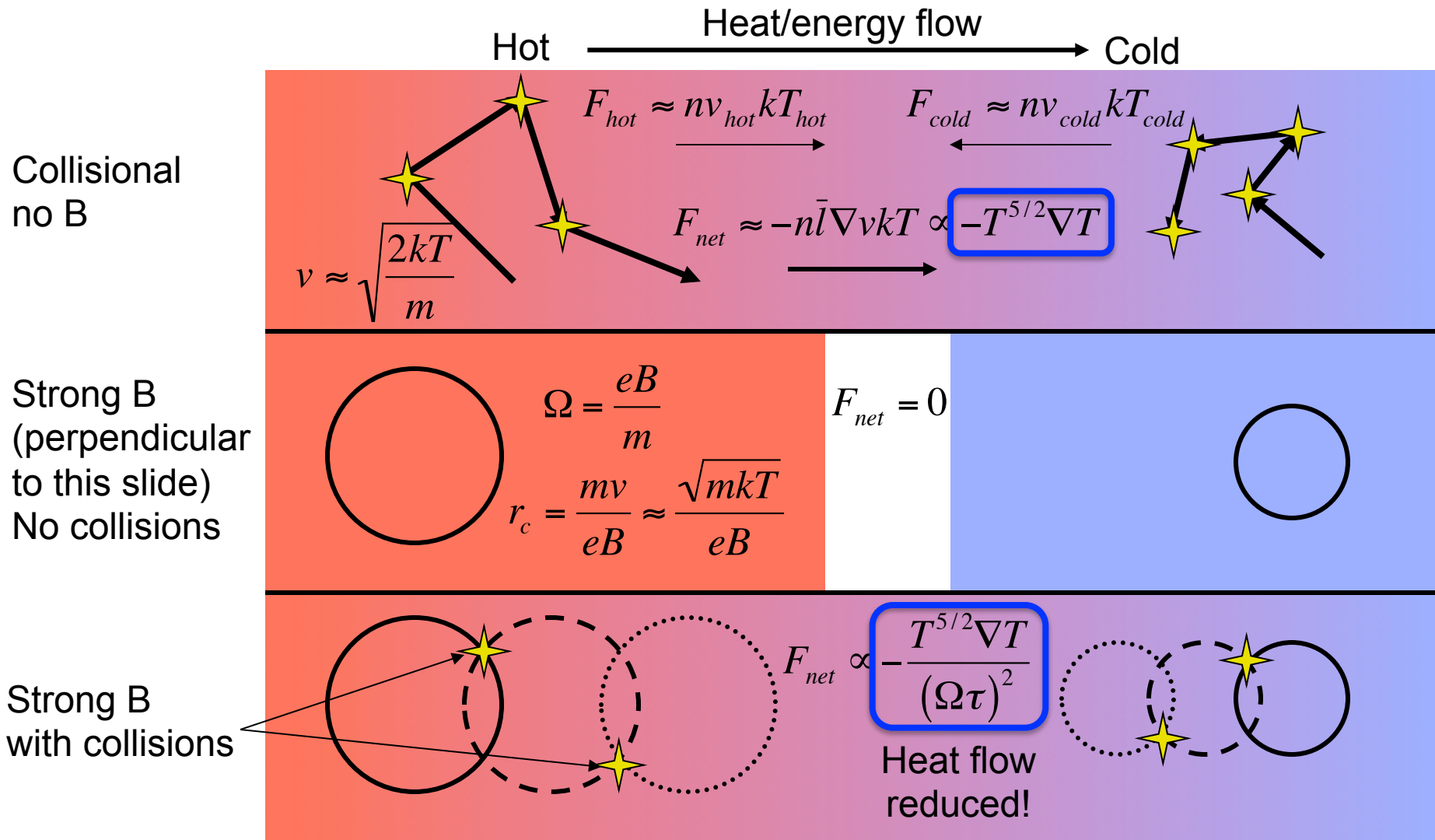


Sandia is conducting magnetized-liner inter-fusion (MagLIF) experiments on Z

- A beryllium liner contains DD or DT fusion fuel
- The fuel is magnetized by an applied axial magnetic field
- The fuel is preheated by a laser
- The fuel is subsequently compressed by the Z-accelerator current
- To date, the MagLIF concept has achieved DD neutron yields $\sim 2 \times 10^{12}$

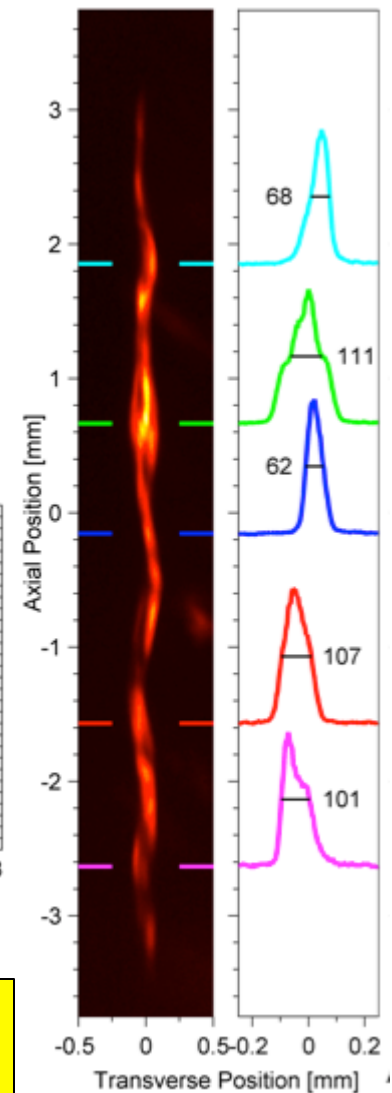
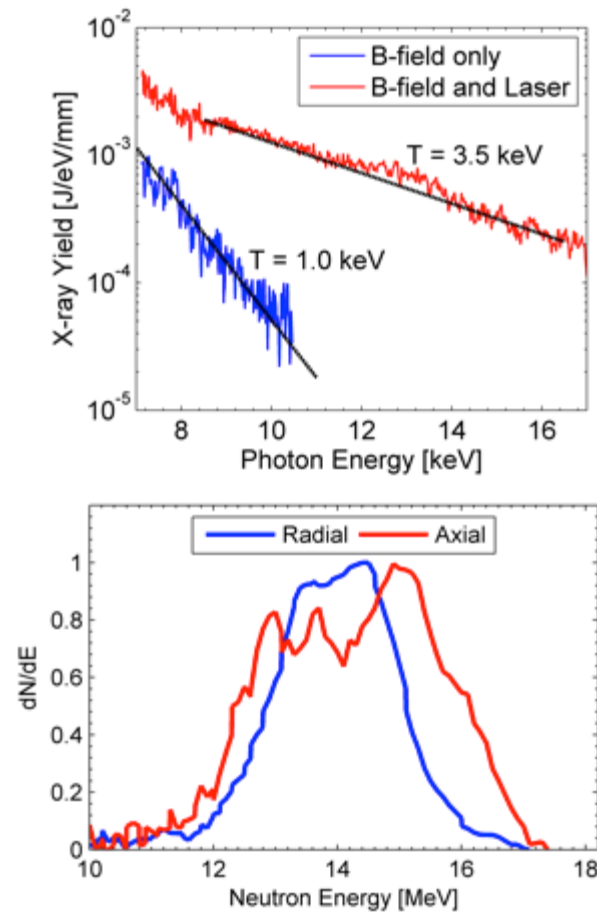
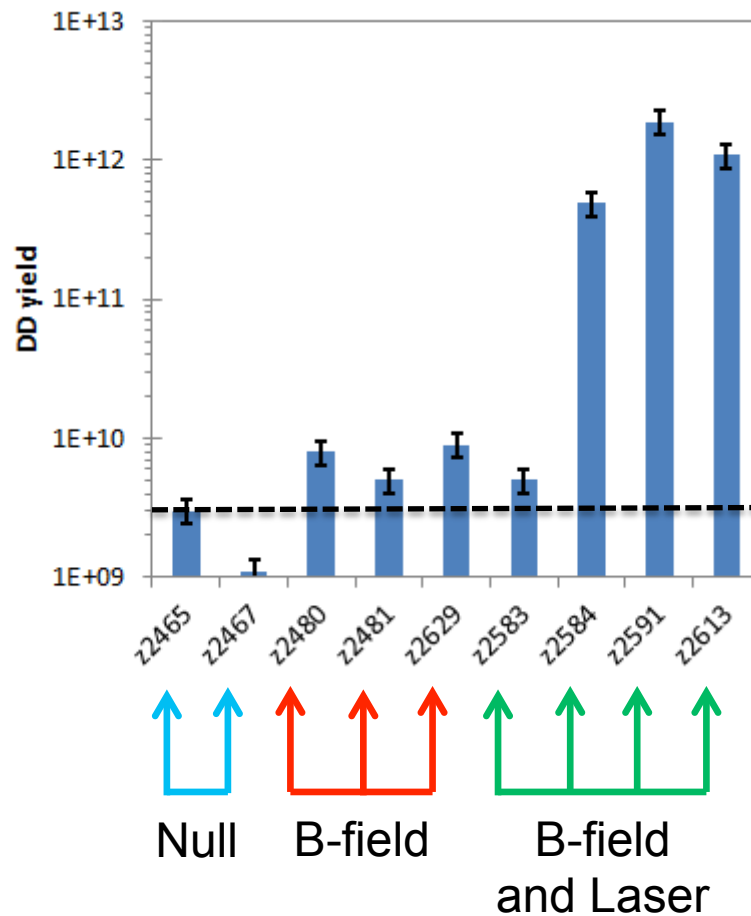


The presence of a magnetic field can strongly affect transport properties, e.g. electron heat conduction



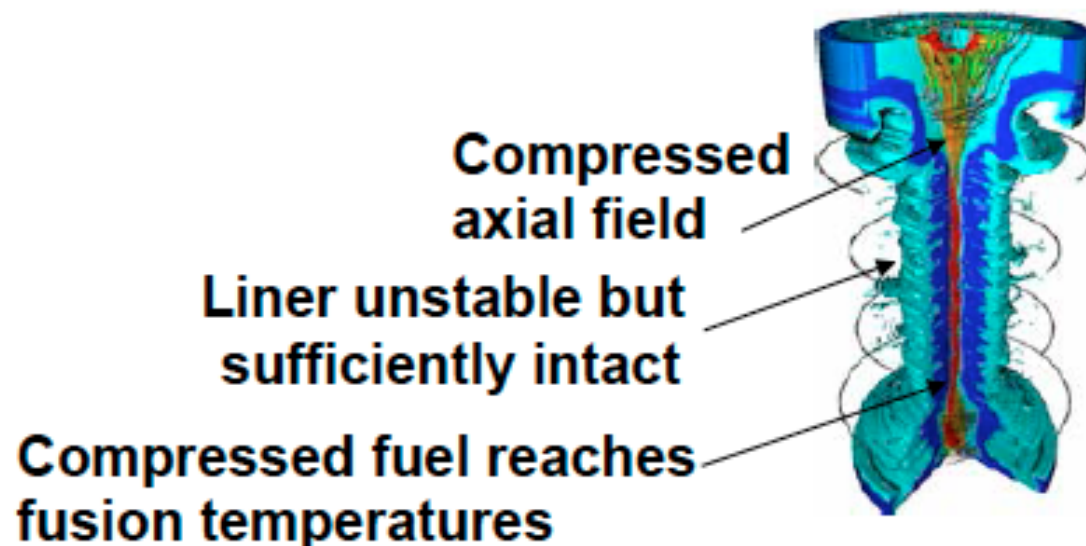
“Anomalous” heat transport can reduce the benefit of magnetic fields (e.g., in tokamaks) but there remains a significant benefit

We have obtained promising initial results with MagLIF



Sefkow, Slutz, et al., Phys. Plasmas (2014)
 Gomez et al. Phys. Rev. Lett. (2014), Phys. Plasmas (2015)
 Schmit et al. Phys. Rev. Lett. (2014)
 Knapp et al. Phys. Plasmas (2015)

MagLIF is a concept well suited to pulsed power and may lower requirements for plasma self-heating



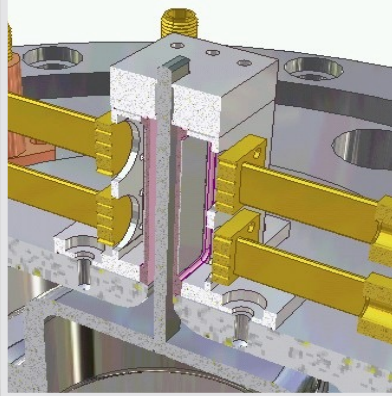
Vimp ~ 7 to 10 cm/μsec
Telectron ~ Tion ~ 3 keV
DD Neutron Yield ~ 2.e12
DT Neutron Yield ~ 5.e10
P ~ 1 GBar
BR ~ 0.4 MGauss-cm
 $\rho R_{\text{fuel}} \sim 1.5 \text{ mg/cm}^2$
 $\rho R_{\text{liner}} \sim 0.9 \text{ g/cm}^2$

See Ryan McBride's talk on Thursday

Magnetically-driven planar and cylindrical geometries are used for the dynamic materials program

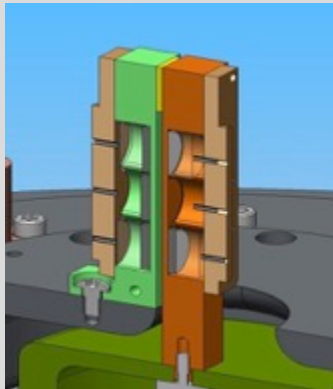
Shock Compression

Velocities > 40 km/s



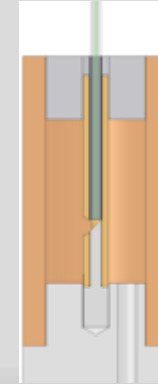
Planar Ramp-Compression

continuous compression up to ~4 Mbar



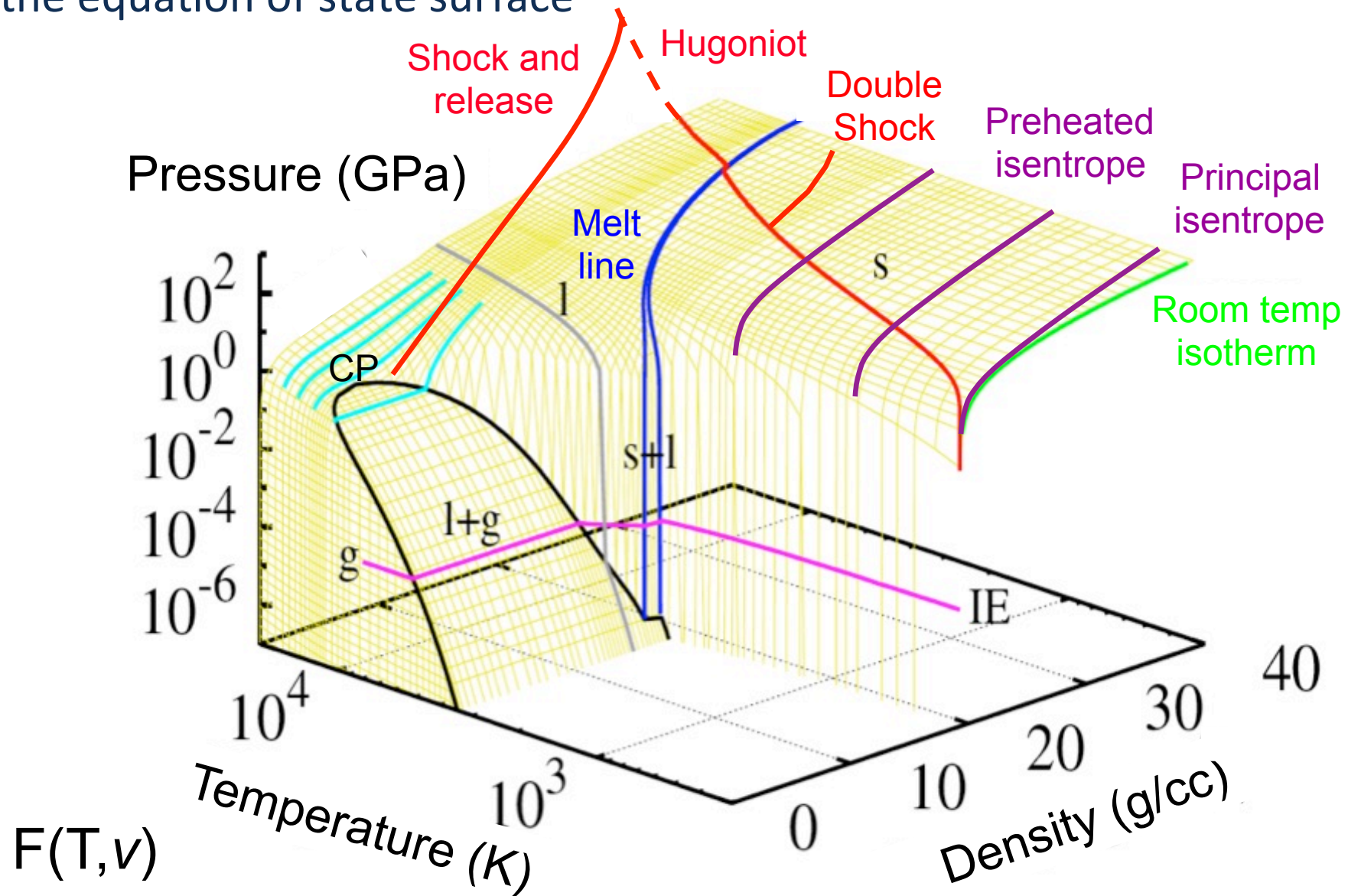
Cylindrical Ramp-Compression

Possibly continuous compression up to ~10-20 Mbar



- Shock or ramp compression of materials increases densities by 2 to **10:1**

Magnetic pressure on Z allows access to a large region of the equation of state surface



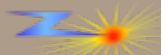
The dynamic materials program measures equations of state and other properties (strength) for elements and compounds of interest

<div>hydrogen 1 H 1.0079</div>		<div>D₂ LiD CO₂ H₂O</div>										<div>MgO CaF LiF SiO₄ (e.g. quartz)</div>										<div>helium 2 He 4.0026</div>	
<div>lithium 3 Li 6.941</div>		<div>beryllium 4 Be 9.0122</div>											<div>boron 5 B 10.811</div>		<div>carbon 6 C 12.011</div>	<div>nitrogen 7 N 14.007</div>	<div>oxygen 8 O 15.999</div>	<div>fluorine 9 F 18.998</div>	<div>neon 10 Ne 20.180</div>				
<div>sodium 11 Na 22.990</div>		<div>magnesium 12 Mg 24.305</div>											<div>aluminum 13 Al 26.982</div>	<div>silicon 14 Si 28.086</div>	<div>phosphorus 15 P 30.974</div>	<div>sulfur 16 S 32.065</div>	<div>chlorine 17 Cl 35.453</div>	<div>argon 18 Ar 39.948</div>					
<div>potassium 19 K 39.098</div>		<div>calcium 20 Ca 40.078</div>	<div>scandium 21 Sc 44.956</div>	<div>titanium 22 Ti 47.867</div>	<div>vanadium 23 V 50.942</div>	<div>chromium 24 Cr 51.996</div>	<div>manganese 25 Mn 54.938</div>	<div>iron 26 Fe 55.845</div>	<div>cobalt 27 Co 58.933</div>	<div>nickel 28 Ni 58.693</div>	<div>copper 29 Cu 63.546</div>	<div>zinc 30 Zn 65.39</div>	<div>gallium 31 Ga 69.723</div>	<div>germanium 32 Ge 72.61</div>	<div>arsenic 33 As 74.922</div>	<div>selenium 34 Se 78.96</div>	<div>bromine 35 Br 79.904</div>	<div>krypton 36 Kr 83.80</div>					
<div>rubidium 37 Rb 85.468</div>		<div>strontium 38 Sr 87.62</div>	<div>yttrium 39 Y 88.906</div>	<div>zirconium 40 Zr 91.224</div>	<div>niobium 41 Nb 92.906</div>	<div>molybdenum 42 Mo 95.94</div>	<div>technetium 43 Tc [98]</div>	<div>ruthenium 44 Ru 101.07</div>	<div>rhodium 45 Rh 102.91</div>	<div>palladium 46 Pd 106.42</div>	<div>silver 47 Ag 107.87</div>	<div>cadmium 48 Cd 112.41</div>	<div>indium 49 In 114.82</div>	<div>tin 50 Sn 118.71</div>	<div>antimony 51 Sb 121.76</div>	<div>tellurium 52 Te 127.60</div>	<div>iodine 53 I 126.90</div>	<div>xenon 54 Xe 131.29</div>					
<div>caesium 55 Cs 132.91</div>		<div>barium 56 Ba 137.33</div>	<div>57-70 ★</div>	<div>lanthanum 71 Lu 174.97</div>	<div>hafnium 72 Hf 178.49</div>	<div>tantalum 73 Ta 180.95</div>	<div>tungsten 74 W 183.84</div>	<div>rhenium 75 Re 186.21</div>	<div>osmium 76 Os 190.23</div>	<div>iridium 77 Ir 192.22</div>	<div>platinum 78 Pt 195.08</div>	<div>gold 79 Au 196.97</div>	<div>mercury 80 Hg 200.59</div>	<div>thallium 81 Tl 204.38</div>	<div>lead 82 Pb 207.2</div>	<div>bismuth 83 Bi 208.98</div>	<div>polonium 84 Po [209]</div>	<div>astatine 85 At [210]</div>	<div>radon 86 Rn [222]</div>				
<div>francium 87 Fr [223]</div>		<div>radium 88 Ra [226]</div>	<div>89-102 ★ ★</div>	<div>lawrencium 103 Lr [262]</div>	<div>rutherfordium 104 Rf [261]</div>	<div>dubnium 105 Db [262]</div>	<div>seaborgium 106 Sg [266]</div>	<div>bohrium 107 Bh [264]</div>	<div>hassium 108 Hs [269]</div>	<div>meitnerium 109 Mt [268]</div>	<div>unnilium 110 Uun [271]</div>	<div>ununium 111 Uuu [272]</div>	<div>unubium 112 Uub [277]</div>	<div>unusquadrium 114 Uuq [289]</div>									

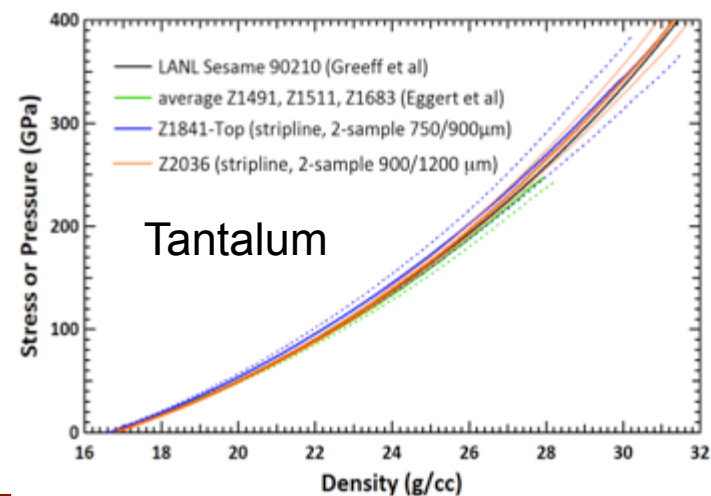
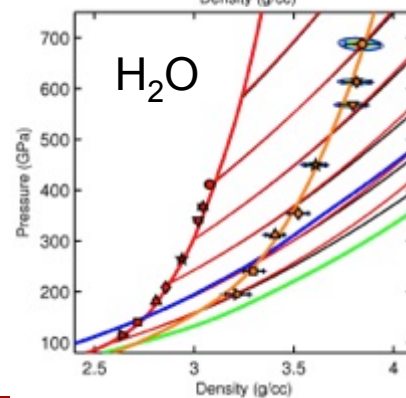
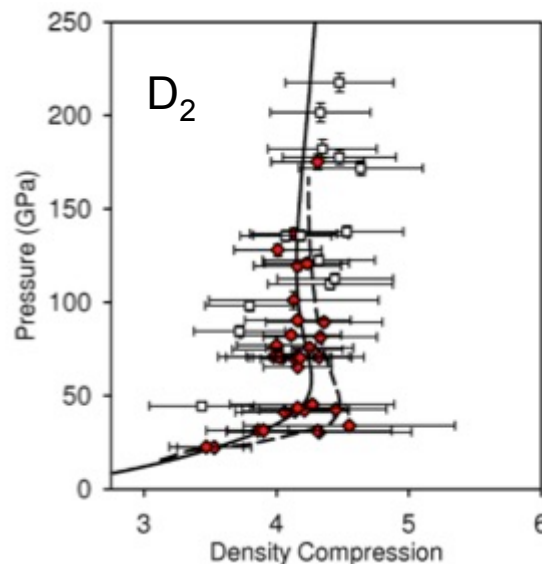
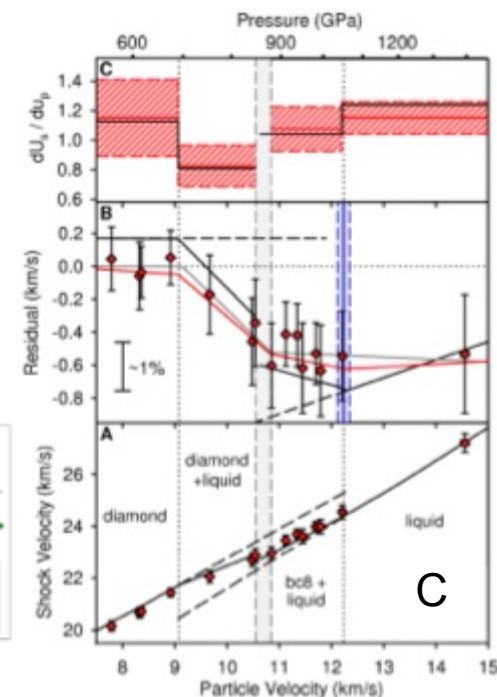
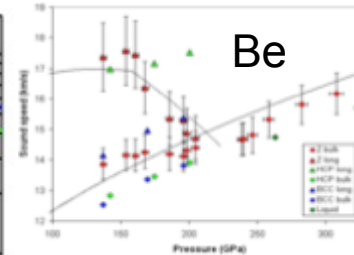
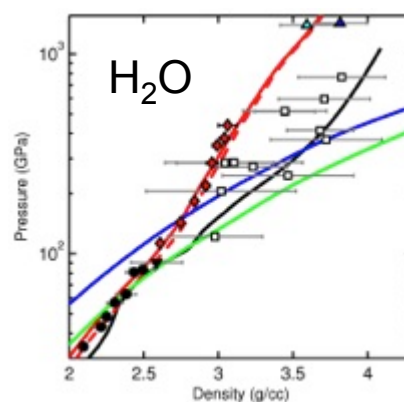
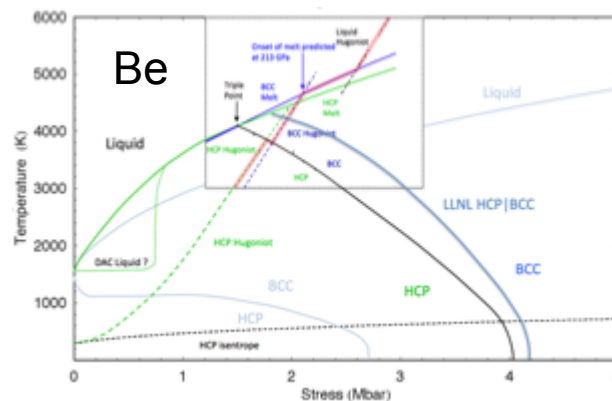
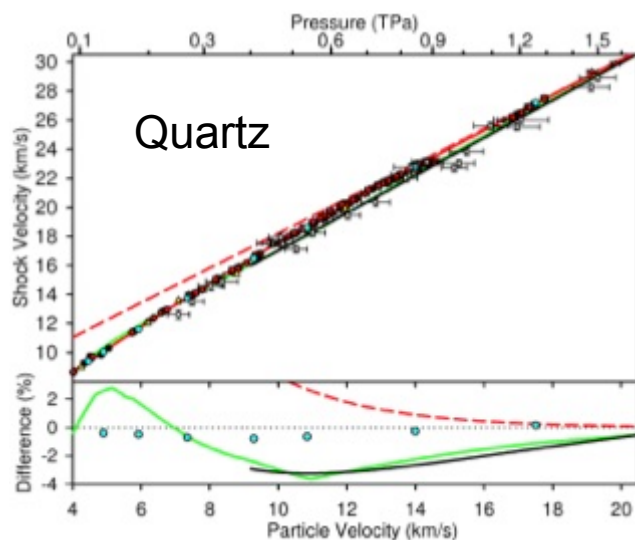
*Lanthanide series

* * Actinide series

lanthanum 57	cerium 58	praseodymium 59	neodymium 60	promethium 61	samarium 62	euporium 63	gadolinium 64	terbium 65	dysprosium 66	holmium 67	erbium 68	thulium 69	ytterbium 70
La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb
138.91	140.12	140.91	144.24	[145]	150.36	151.96	157.25	158.93	162.50	164.93	167.26	168.93	173.04
actinium 89	thorium 90	protactinium 91	uranium 92	neptunium 93	plutonium 94	americum 95	curium 96	berkelium 97	californium 98	einsteinium 99	fermium 100	mendelevium 101	nobelium 102
Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No
[227]	232.04	231.04	238.03	[237]	[244]	[243]	[247]	[247]	[251]	[252]	[257]	[258]	[259]



Z has been used to address several interesting problems in the multi-Mbar regime



There is a need for facilities that achieve thermonuclear ignition and high-yield fusion

- The U.S. conducted its last underground nuclear test 23 years ago (in 1992)
- Since then, the U.S. has not conducted thermonuclear-ignition or high-fusion-yield experiments
- There will soon be no full-time scientists with direct experimental ignition or high-yield experience

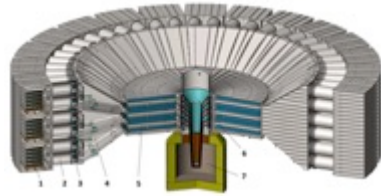
The second U.S. underground nuclear test was a 1-kiloton explosion (1955).

Preparation for an underground test at the Nevada Test Site in the 1990s.

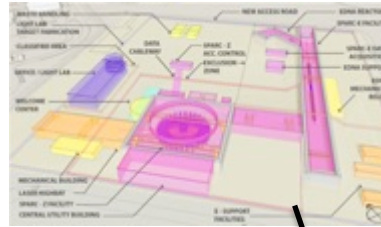
It's clear the U.S. needs a thermonuclear-burn facility



We have a long-term vision to achieve ignition and high-yield fusion in the future



2019: Baikal operations begin
(50 MA Russian fusion facility)



SPARC-Z operations begin



CD-0 for SPARC-Z High-Yield Facility

Z-300 demonstrates ignition
and high-gain scaling

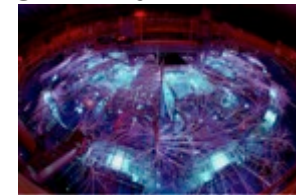
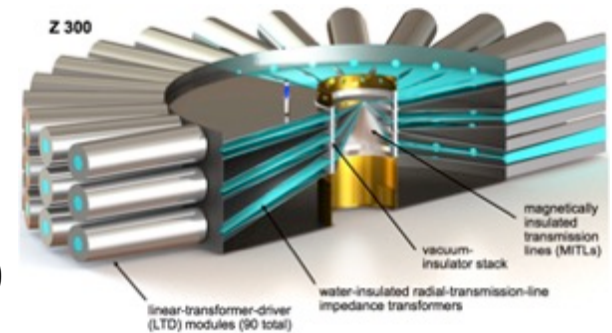
Z-300 operations begin

Review of Ignition on Z-300; CD-0 for Z300

Demonstrate LTD module prototypes (e.g., radiography)

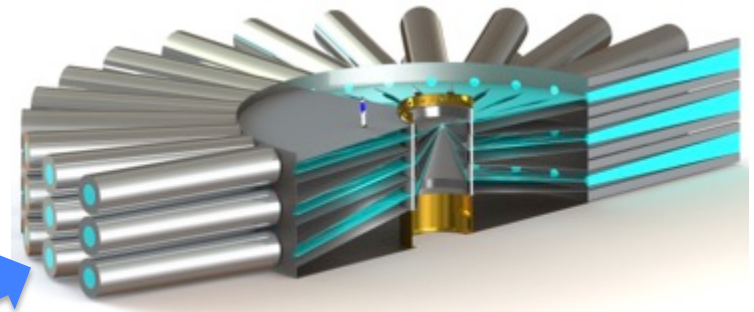
2015: National ICF program review

2013: First integrated tests of new MagLIF idea on Z



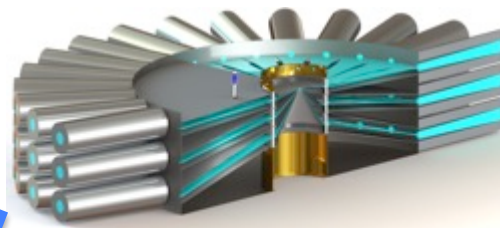
We are currently exploring target designs and pulsed power architectures that may be on the path to 0.5-1 GJ yields that also meet the needs of the science campaigns

Z-800 Goal
Yield = E_{machine}



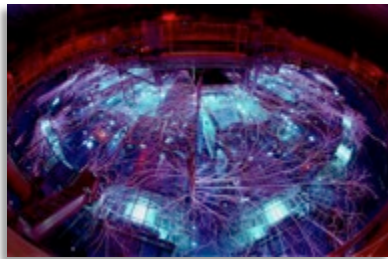
- 890 TW
- 52 m diameter
- 130 MJ stored energy
- 66 MA

Z-300 Goal
Yield = E_{target}



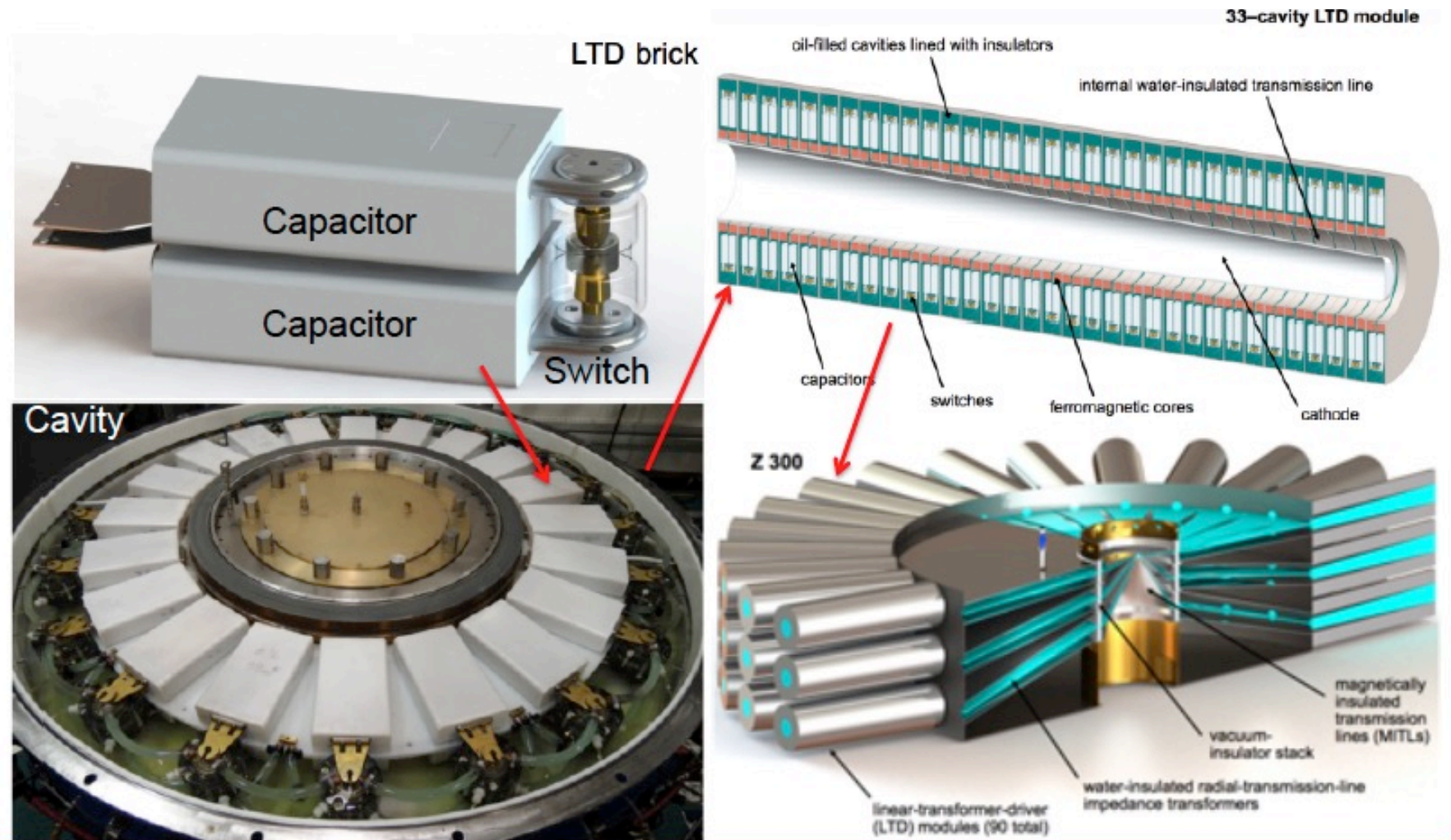
- 320 TW
- 35 m diameter
- 47 MJ stored energy
- 49 MA

Z Goal
Yield = E_{fuel}



- 80 TW
- 33 m diameter
- 22 MJ stored energy
- 26 MA peak current

We have developed a linear transformer driver architecture that can scale to 800-1000 TW and that is twice as efficient and twice as compact as Z



Z 300 will deliver 49 MA to a MagLIF load.
The goal: thermonuclear ignition (i.e., a liner gain of ~ 1)

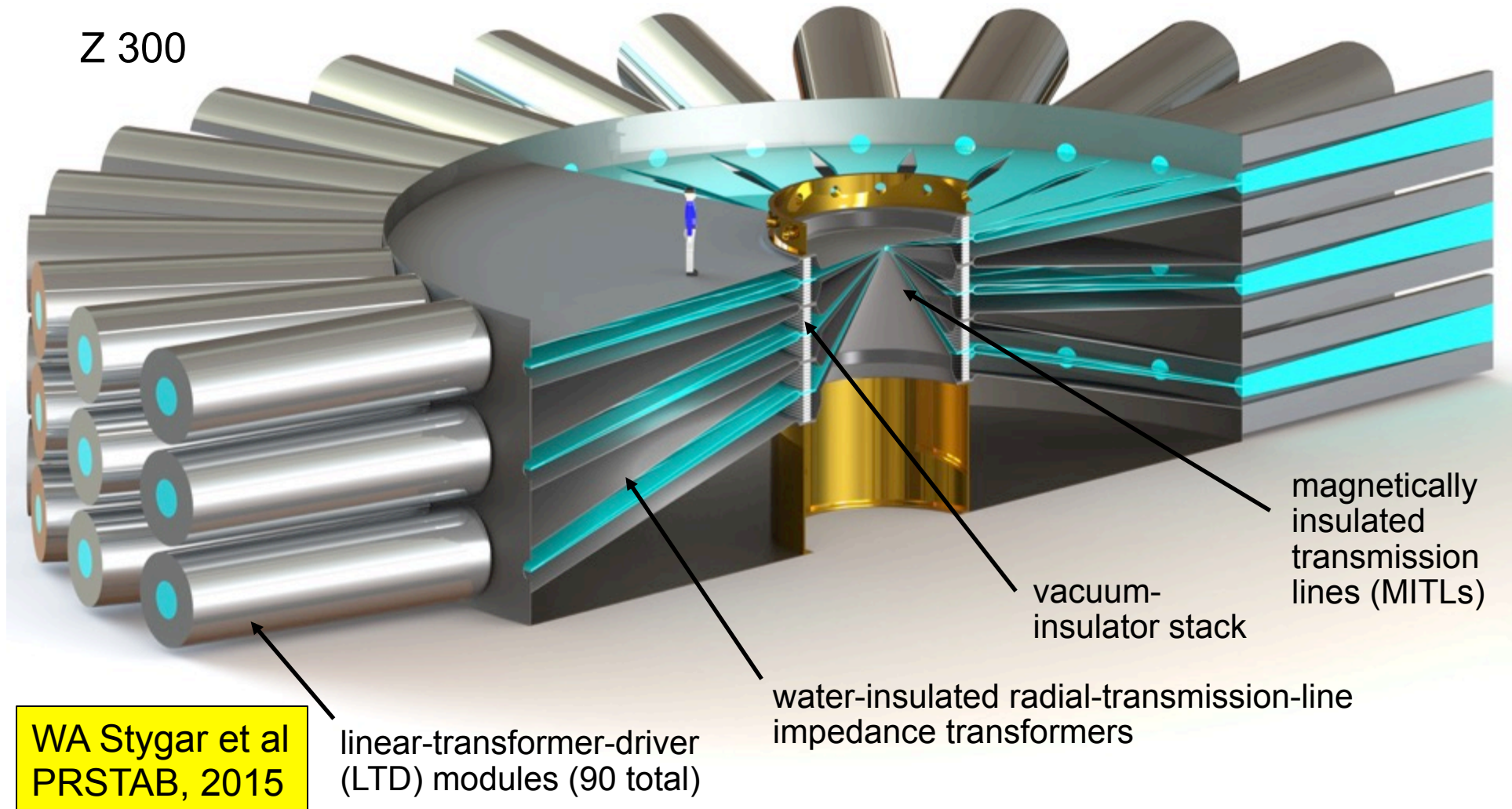
$$E_{\text{LTDs}} = 47 \text{ MJ}$$
$$P_{\text{LTDs}} = 320 \text{ TW}$$

$$V_{\text{stack}} = 7.7 \text{ MV}$$
$$L_{\text{vacuum}} = 15 \text{ nH}$$

$$I_{\text{load}} = 49 \text{ MA}$$
$$\tau_{\text{implosion}} = 150 \text{ ns}$$

$$\text{diameter} = 35 \text{ m}$$
$$\text{fusion yield} \sim 5 \text{ MJ}$$

Z 300



WA Stygar et al
PRSTAB, 2015

Z 800 will deliver 66 MA to a MagLIF load.
The goal: ~500 MJ of fusion yield (i.e., a machine gain of ~4)

$$E_{\text{LTDs}} = 130 \text{ MJ}$$

$$P_{\text{LTDs}} = 890 \text{ TW}$$

$$V_{\text{stack}} = 15 \text{ MV}$$

$$L_{\text{vacuum}} = 20 \text{ nH}$$

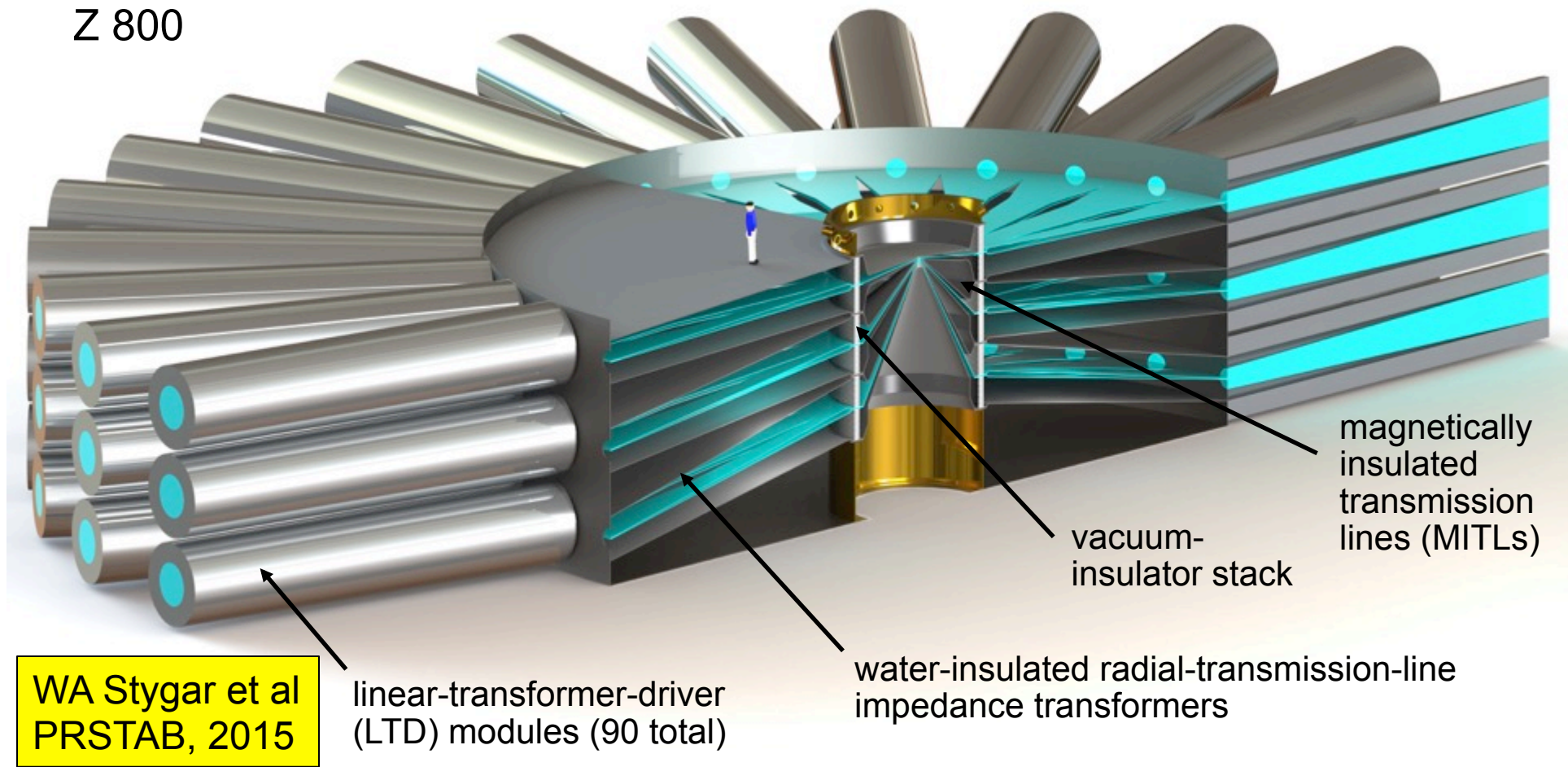
$$I_{\text{load}} = 66 \text{ MA}$$

$$\tau_{\text{implosion}} = 114 \text{ ns}$$

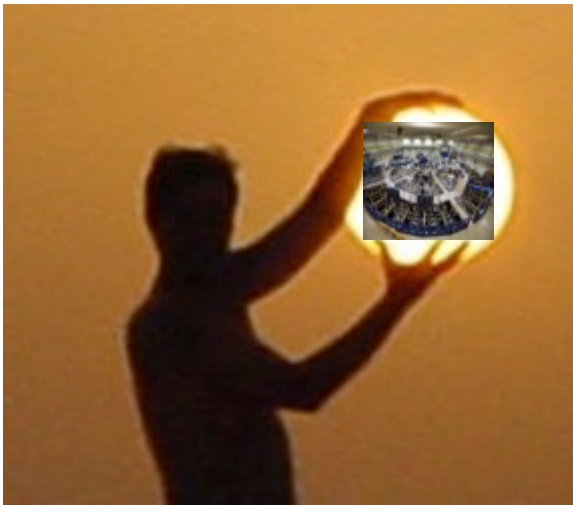
$$\text{diameter} = 52 \text{ m}$$

$$\text{fusion yield} \sim 500 \text{ MJ}$$

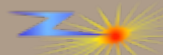
Z 800



Large HED facilities are as close as we can get to “holding the Sun” on Earth



- We make sufficient enough macroscopic quantities of high energy density, hot-dense matter to allow accurate study
- Pulsed power, in particular, allows large energy delivery to the largest samples or targets, generating large scale size, hot, dense plasmas with large volume-to-surface area ratios, and at low cost per shot
- We may someday be able to “create a star” in the laboratory
- When we make measurements in the HED field, ***quite often*** the measurements disagree with scientific predictions or expectations before the experiments → experiments are needed to validate our models



Summary

- Large currents create large magnetic fields, and large magnetic fields create large pressures, which are needed to access high energy density regimes
- The Z machine creates large currents, allowing us to address fundamental issues in HED science, dynamic materials, and inertial confinement fusion
- Pulsed power can inexpensively, efficiently, and flexibly drive many different kinds of applications at large currents and high voltages
- These applications go well beyond traditional concepts known colloquially as “z-pinches”
- Magnetic drive is a low cost and efficient way to generate high energy density conditions over large volumes
- The upper limits on magnetically-driven performance in achieving high energy densities are not known
- There is a lot of room for innovation!