

# FULLY-KINETIC PARTICLE-IN-CELL SIMULATIONS OF TRIGGERED THREE-ELECTRODE GAS SWITCHES\*

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\* Sandia National Laboratories is a multi-program laboratory managed and operated by Sandia Corporation, a wholly owned subsidiary of Lockheed Martin Corporation, for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000.

## Abstract

We present the results of 2D and 3D fully-kinetic electromagnetic particle-in-cell Monte Carlo (PICMC) simulations of triggered three-electrode gas switches using dry air as a gas (at pressures greater than 1 ATM). In such switches the AK gap voltage is set slightly below the breakdown threshold. A voltage pulse applied to a trigger needle placed in the AK gap allows breakdown to occur between, first, the trigger and anode, followed by the trigger and cathode. We demonstrate that a fully-kinetic PICMC approach can be used to follow the entire evolution of the switch, from the initial avalanche and streamer formation up to the fully conducting phase. We utilize an 18-species air chemistry model which is shown to agree with swarm parameters (breakdown threshold, drift velocity) obtained by experiment. Photon transport and photo-ionization are also included to permit the modeling of cathode directed streamers. This computational model will be used to help design closing switches for pulsed-power systems.

### (1) Introduction and Outline

We describe the results of 2D and 3D hybrid-PIC simulations with the LSP [1] code of a **triggered three-electrode railgap closing switch** being designed and fielded at Sandia National Laboratory (SNL) for pulsed power applications.

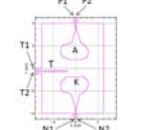
These are **fully-kinetic** electromagnetic simulations of the switch using air as the working gas at a pressure of many atmospheres.

A high density gas is required to hold off a 200 kV voltage (for an AK gap of ~ 1 cm) until a pulse applied to the trigger electrode initiates a multi-step closing process.

#### Outline:

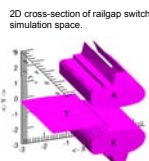
- Description of the operation of three-electrode railgap switch and introduction of switch geometry, driving circuits, and operating parameters.
- Description of 18-species air chemistry including photon generation and transport, and description of PICMC algorithms for kinetic particle-based chemistry modeling.
- 2D Cartesian railgap simulation results.
- 3D Cartesian railgap simulation results.
- Conclusions

### (3) Three Electrode "Railgap" Switch



2D LSP simulation geometry for railgap switch in xy plane

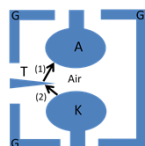
P1,P2, etc. are open simulation boundaries which are connected to the circuits shown in the adjacent panel (4).



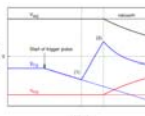
Full 3D geometry of railgap switch

**Full 3D railgap geometry**  
For the full switch the "rail" electrodes extend in the z-direction for 12 cm (6 cm length shown here).

### (2) Schematic Operation of Three-Electrode Closing Switch



Schematic drawing of cross-section of three-electrode triggered closing switch. AK gap ~ 1.5 cm



Schematic of electrode voltages as a function of time

The three electrodes are the cathode (K), anode (A), and trigger (T).

Prior to opening of switch A and K are held fixed potentials (e.g. +100 kV and 100 kV). The surrounding metal is held at ground (G) and T floats electrically.

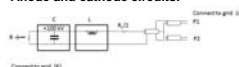
The initial AK gap voltage must be small enough to hold off breakdown of the air.

But a voltage pulse applied to the trigger allows breakdown of the TA gap (1).

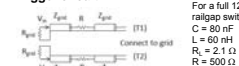
Due to TA gap current, the T voltage rises until the KT gap can break down. Resulting in a complete current path between A and K (2).

### (5) LSP simulation of Railgap Switch

#### Anode and cathode circuits:

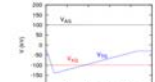


#### Trigger circuit



To avoid self-breakdown for  $V_{AK} = 200$  kV and 1.5 cm AK gap requires an air density of 6.5 ATM

#### Railgap voltages in vacuum



### (6) Air chemistry model

**18 Total Species:** e,  $N_2$ ,  $O_2$ , Ar,  $N_2^+$ ,  $O_2^+$ , Ar<sup>+</sup>, N,  $N_2^+$ ,  $O_2^-$ , O, O<sup>-</sup>, NO<sup>+</sup>, N<sup>+</sup>, O<sup>+</sup>,  $\gamma$ ,  $N_2^*$  where  $N_2^*$  is an excited state of  $N_2$  (13.0 eV) which can radiatively relax and act as a source of photons.

#### Included Reactions:

- Collisional ionization (including elastic and inelastic scattering)
  - $e + O_2 \rightarrow e + O_2^+$
  - $e + N_2 \rightarrow e + N_2^+$
  - $e + Ar \rightarrow e + Ar^+$
- Electron dissociation
  - $e + O_2 \rightarrow O + O$
  - $e + NO \rightarrow N + O$
- Electron attachment
  - $e + O_2 \rightarrow M + O_2^-$
  - $e + N_2 \rightarrow M + N_2^-$
  - $e + O + O_2 \rightarrow O^- + O_2$
- Charge exchange
  - $N_2^+ + O_2 \rightarrow O_2^+ + N_2$
- Formation and destruction of cluster ions
  - $O_2^+ + O_2 \rightarrow O_2^+ + O_2$
  - $N_2^+ + N_2 \rightarrow N_2^+ + N_2$
  - $O_2^+ + N_2 \rightarrow O_2^+ + N_2$
  - $e + N_2 \rightarrow N_2^-$
  - $e + O_2 \rightarrow O_2^-$
- Electron-ion recombination
  - $e + O_2^+ \rightarrow O + O$
  - $e + N_2^+ \rightarrow N + N$
- Detachment
  - $O_2^- + M \rightarrow e + O_2 + M$
  - $N_2^- + M \rightarrow e + N_2 + M$
  - $O^- + O_2 \rightarrow e + O_2$
- Ion-ion recombination
  - $N_2^+ + O_2^- \rightarrow N_2 + O_2$
  - $O_2^+ + O_2^- \rightarrow O_2 + O_2$
- Miscellaneous
  - $N^+ + O_2 \rightarrow NO^+ + O$
  - $O^+ + N_2 \rightarrow NO^+ + N$

### (7) Scattering and Chemistry Reactions

#### Algorithms and Cross-sections

Reactions (1a), (1b), and (1c) are all modeled by standard PICMC methods [8] [7].

All cross-sections for these reactions are obtained from the cross-section database of the Boltzmann code EEDF [8].

Reactions (2)-(9) are modeled by the binary scattering algorithms described by Narbut and Yonemura [9] [10].

Cross-sections for these reactions are obtained from the literature: (2a) [11], (2b), (3a), (3b), (8a), (8b), (9a), (9b) [12], (4), (5a), (5b), (6a), (6b), (7b) [13], (5a) [14], (5c), (7a), [15].

In many cases only the reaction rate,  $k(T)$ , is provided rather than the cross-section, which is needed for the kinetic PIC treatment. In these cases we obtain the cross section as follows [16]:

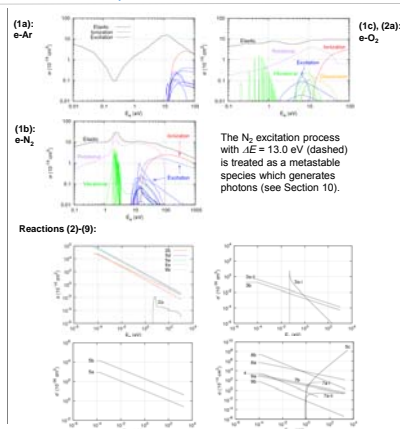
If the reaction rate is given by

$$k(T) = k_0 T^n \exp(-E/kT)$$

then the center-of-mass cross-section, for any  $n > -3/2$ , can be written

$$\sigma(E) = \sqrt{\pi/8} \frac{k_0}{\Gamma(n+3/2)} \frac{(E-E_0)^{n+1/2}}{E} \Theta(E-E_0)$$

### (8) Cross-section plots

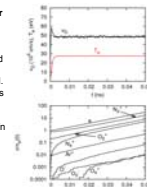


### (9) Swarm Simulation Test of Air Model

Swarm simulations to test the air chemistry model:

- Cell PIC simulation with fixed applied electric field.
- Particle momenta are advanced but not particle positions.
- Particle self-fields are neglected.
- A seed electron population starts the avalanche.

At steady-state we obtain swarm parameters (usually functions of  $E/n$  for weakly-ionized plasma).



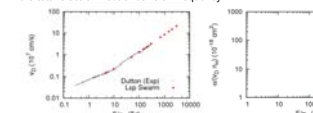
Sample swarm simulation: Air at STP (78%  $N_2$ , 21%  $O_2$ , 1% Ar)  $E/n = 500$  Td ( $T_d = 10^{-17}$  V-cm<sup>2</sup>). No photons.

- Note:
- 1) There's no initial  $N^+$  or  $O^+$ . This implies that you never get any, and consequently never get  $NO^+$ .
  - 2) Density of  $O^-$  = density of  $O$ .
  - 3) No N density because it can only be made by 2<sup>nd</sup> order reaction  $e + N_2 \rightarrow N + N$ .

Electron drift velocity and temperature go to steady-state values, and there is a constant ionization rate.

### (10) Drift velocity and Ionization rate as a function of $E/n$ for air:

Vary field for air at 1 ATM. Timestep chosen so that  $v_{dt} \sim 0.1$  where  $v$  is the total electron-neutral collision frequency.



At 1 ATM for  $E/n < 120$  Td, air becomes attachment dominated. (The critical field does not scale exactly as  $E/n$ , as the attachment process is 3-body)

Good agreement with experimental swarm data [17].

### (12) 2D Railgap Simulations

- Initial AK voltage of 200 kV and 1.5 cm AK gap.
- Air at 6.5 ATM ( $n = 1.7 \times 10^{21}$  cm<sup>-3</sup>,  $T = 273$  K).
- Driving circuits and trigger pulse shown in section 4
- Circuit parameters are scaled to account for 1-cm length in virtual direction.
- Variable cell size, but  $\Delta x = \Delta y = 400$   $\mu$ m in AK gap.
- $\Delta t = 3 \times 10^{-15}$  ns (to resolve e-n collision frequency)
- $\sim 10^6$  total cells,  $> 10^7$  particles.
- Must be run for millions of timesteps,  $\sim 100$  processors. Long runtimes ( $\sim 1$  week)

### (13) 3D Railgap Simulations

- Problems with 2D results:
- Maximum electron density in 2D sim:  $10^{16}$  cm<sup>-3</sup>. Expect 1-2 orders of magnitude higher electron density.
  - Maximum switch resistance in 2D sim: 10  $\Omega$ . Expect 1-2 orders of magnitude lower.

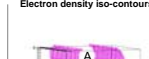
Motivation for 3D simulations: In full 3D railgap simulation we expect current filaments to form along the rails. Pinching of the filaments should produce higher current densities and more Joule heating which should lead to more ionization, and presumably higher electron densities and lower switch resistance.

#### 3D Cartesian railgap simulation:

- Simulate only a 1.5-cm segment in the z-direction, and assume periodic boundaries. Edge effects are neglected.
- $\sim 10^6$  total cells,  $> 10^7$  particles.

Must again be run for millions of timesteps. On 100s of processors very long runtimes ( $\sim 1$  month).

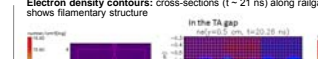
### Electron density iso-contours



$t = 16.2$  ns,  $n_e = 2.2 \times 10^{16}$  cm<sup>-3</sup>

Density (and current) filaments are formed from both anode and cathode to trigger.

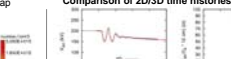
### Electron density contours: cross-sections ( $t \sim 21$ ns) along railgap shows filamentary structure



Squares show size of cells in xz plane (400x400  $\mu$ m<sup>2</sup>)

Filaments are only a few cells in width. But simulations at higher resolution give same filament size.

### Comparison of 2D/3D time histories



Despite presence of filamentary structure in 3D, the amount of charge and current per unit length are the same as 2D!

LSP simulations do not predict higher currents in 3D due to stronger pinching at higher dimensionality.

## Conclusions

- We have performed 2D and 3D fully-kinetic hybrid-PIC simulations of the three-electrode triggered railgap closing switch.
- An air density of 6.5 ATM is required to hold off the 200 kV across the 1.5-cm AK gap until the switch is closed by a pulse applied to the trigger electrode.
- 2D and 3D simulations predict a total current of  $\sim 1.5$  kA per unit length (along the rails). For a 12-cm length railgap this leads to a switch resistance on the order of 10  $\Omega$ . The electron number density in 2D and 3D reaches  $\sim 10^{16}$  cm<sup>-3</sup>.
- 3D simulations show the development of longitudinal current filaments (radius  $\sim 1$  mm, spacing  $\sim$  few mm). But there is no sign of enhanced current (or reduced switch resistance) due to three-dimensional pinching of filaments.
- These simulations qualitatively display closing switch behavior using a fully-kinetic PIC-based approach, but the electron densities are believed to be 1-2 orders of magnitude too small, and the switch resistance 1-2 orders of magnitude too large.
- We are currently investigating the possible causes of the low ionization state of the simulation results.

## References

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