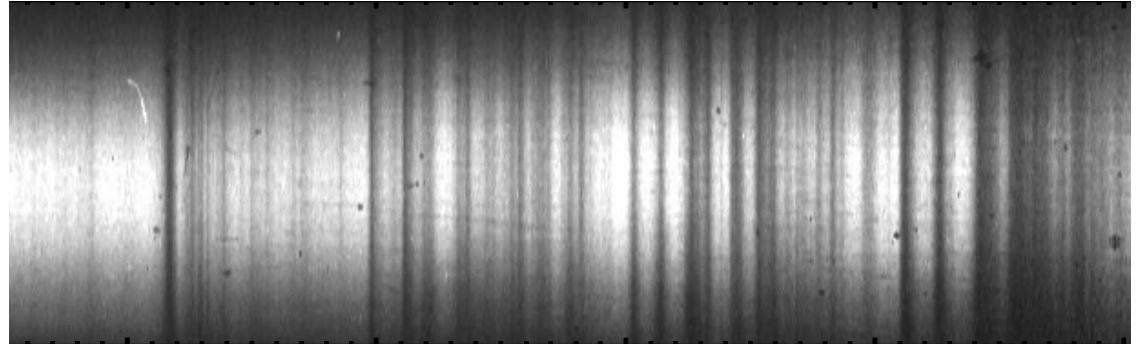
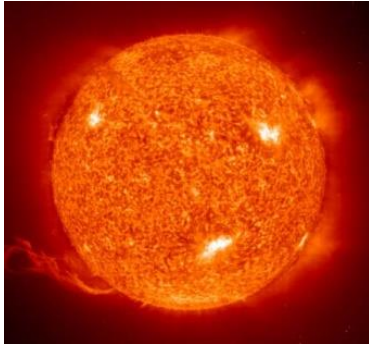


*Exceptional service in the national interest*



## Stellar interior opacity measurements

Jim Bailey

Sandia National Laboratories

# The stellar opacity collaboration involves universities, U.S. national labs, a private company, and the French CEA laboratory



J.E. Bailey, T. Nagayama, G.P. Loisel, G.A. Rochau, S.B. Hansen  
**Sandia National Laboratories, Albuquerque, NM, 87185-1196**



C. Blancard, Ph. Cosse, G. Faussurier, F. Gilleron, J.-C. Pain  
**CEA, France**



A.K. Pradhan, C. Orban, M. Pinsonneault, and S.N. Nahar  
**Ohio State University, Columbus, Ohio, 43210**



C.A. Iglesias and B. Wilson  
**Lawrence Livermore National Laboratory, Livermore, CA, 94550**

J. Colgan, C. Fontes, D. Kilcrease, and M. Sherrill  
**Los Alamos National Laboratory, Los Alamos, NM 87545**

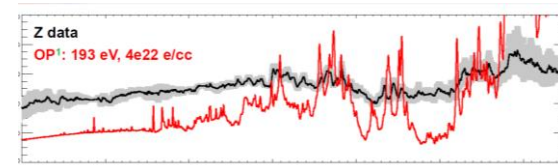
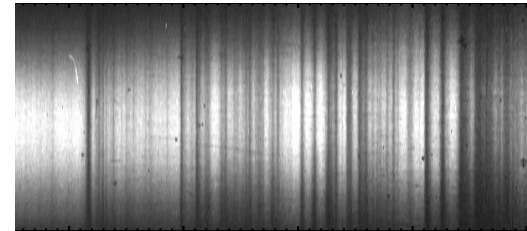
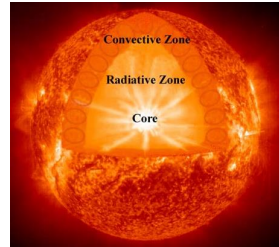
J.J. MacFarlane, I. Golovkin  
**Prism Computational Sciences, Madison, WI**

R.C. Mancini  
**University of Nevada, Reno, NV**



# Z iron opacity experiments refine our understanding of the sun and photon absorption in high energy density matter.

- Solar interior predictions don't match helioseismology
  - Arbitrary opacity increase of 10-20% would fix the problem, but is this the correct explanation?
- Z experiments have measured iron plasma opacity at nearly solar convection zone base conditions
  - Experiment temperature is the same as in sun, density within a factor of 2
- Opacity models disagree with measurements at near-solar-interior conditions
  - The solar Rosseland mean opacity is ~ 7% higher using Z iron data instead of OP calculations



The measurements imply photon absorption in high energy density matter is different than previously believed

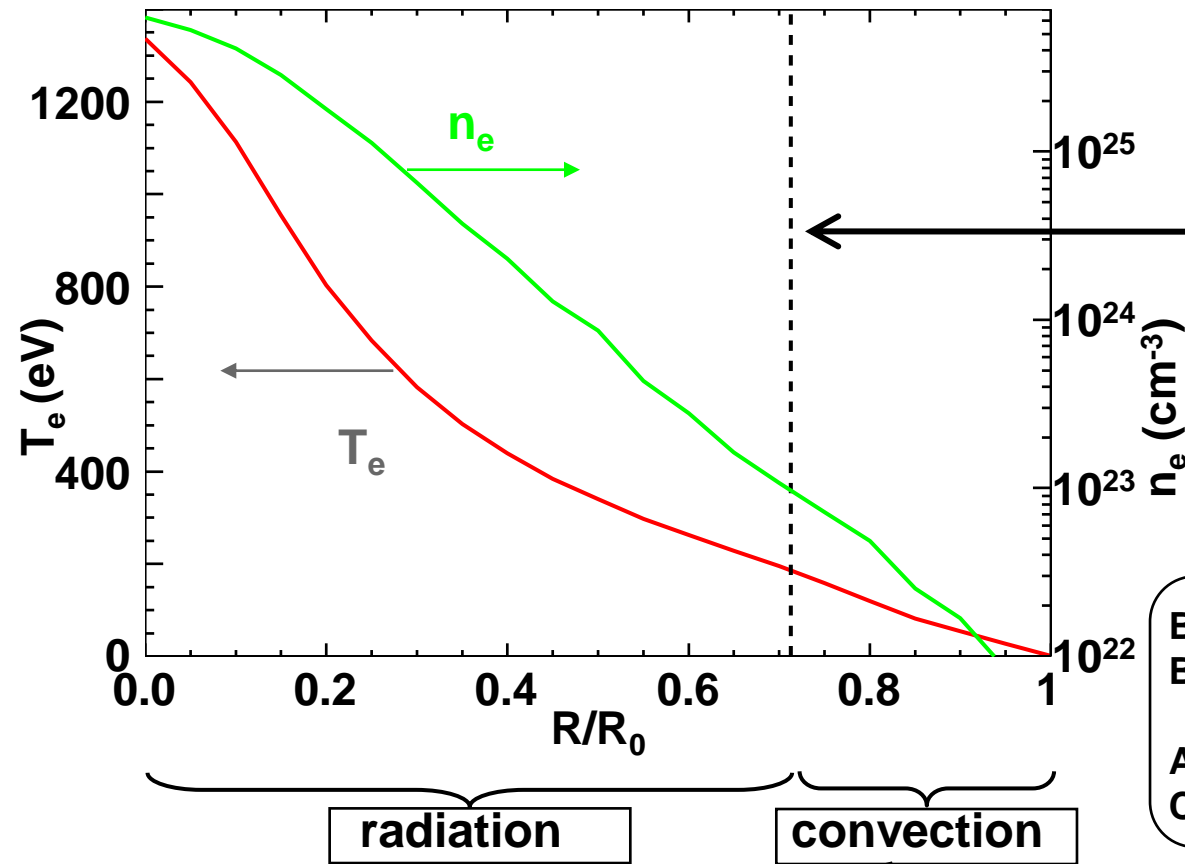
# If our opacity measurements are correct, we must revise our understanding for atoms in HED plasmas

- Measured opacities are generally higher than theory predicts – e.g., Rosseland mean is 30-60% higher
- The measurements alter our understanding of the sun
- Solar physics calibrates many other objects. Therefore the measurements alter our understanding of every main sequence star in the sky, including exoplanet host stars
- The measurements imply likely revisions for ICF capsule dopants

*These serious consequences mandate continued effort*

- We invested the last 2 years investigating possible errors and refining results
- The major conclusions survived this scrutiny
- Going forward, new experiments will further test hypotheses for possible theory deficiencies and possible experiment deficiencies
- Our job isn't finished until we resolve the model/data discrepancy

# Standard solar model predictions of the solar structure disagree with helioseismology



• predicted boundary location different from measured

{~ 10-20  $\sigma$  difference}

Density and sound speed structure are also very different

Bahcall et al, ApJ (2004)

Basu & Antia Physics Reports 2008

Asplund et al Ann Rev AA (2009)

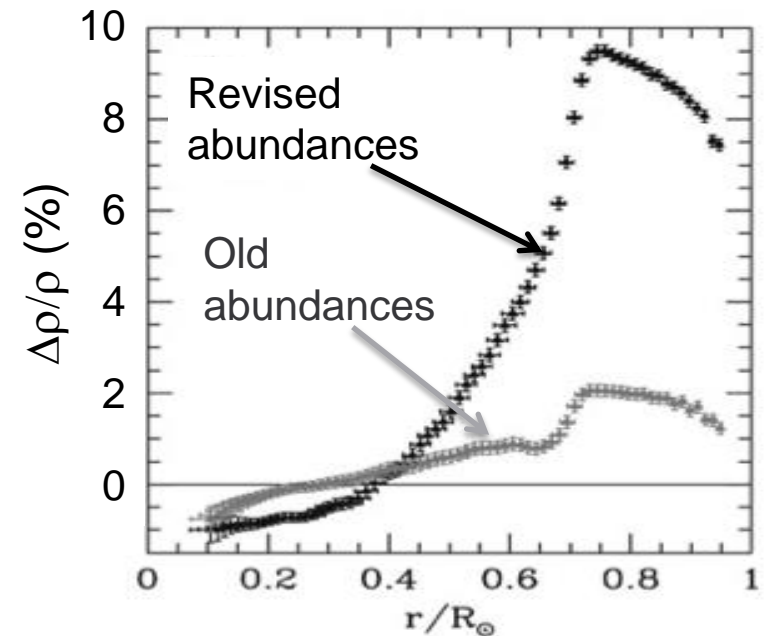
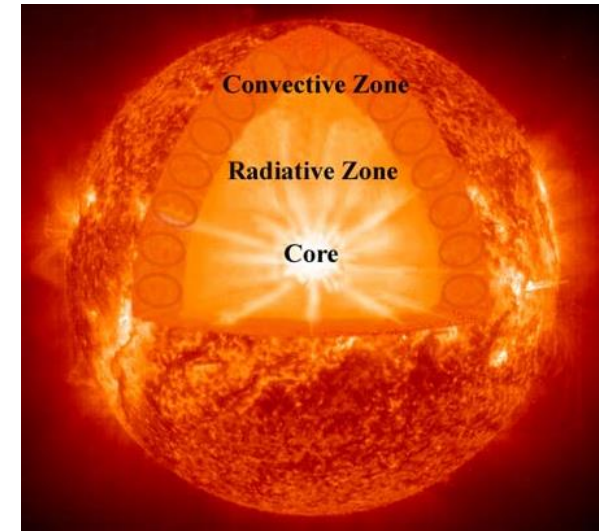
Christensen-Dalsgaard et al A&A (2009)

- Boundary location depends on radiation transport
- A 1% opacity change leads to observable changes.
- This accuracy is a challenge – experiments are needed to know if the solar problem arises in the opacities or elsewhere.

# The disagreement arose after the solar abundance revision that began in 2000

- Standard solar model (simulation)  
Inputs:
    - Abundance
    - EOS
    - Opacity
    - Etc.
  - Helioseismology (measurements)
- 
- Solar abundance revised in 2005
    - C, N, O, Ar, Ne  $\rightarrow$  lowered by 35-45 %
  - Now, standard solar model disagrees with helioseismic measurements

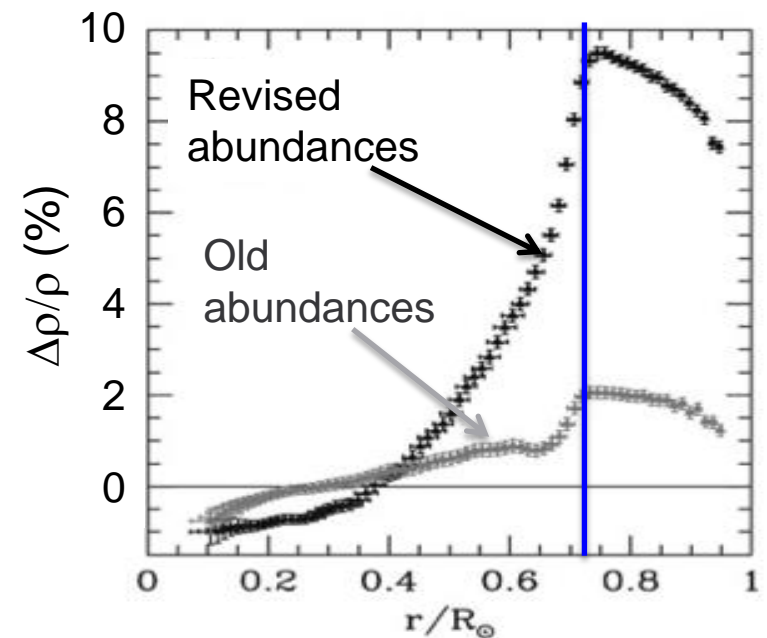
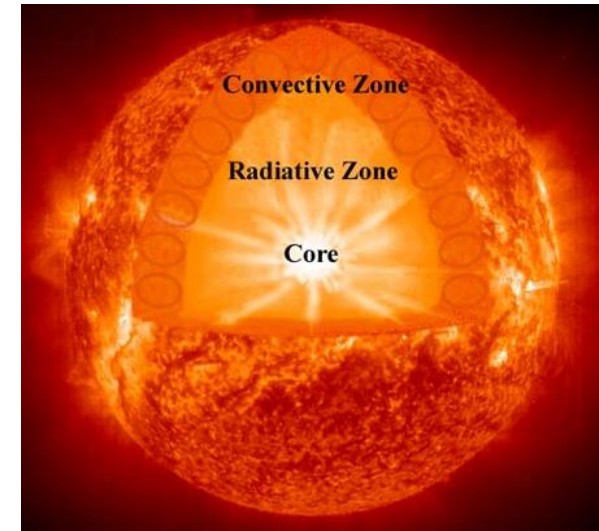
CZB location:  $1\sigma \rightarrow 13\text{-}30\sigma$



# The disagreement arose after the solar abundance revision that began in 2000

- Standard solar model (simulation)  
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CZB location:  $1\sigma \rightarrow 13\text{-}30\sigma$

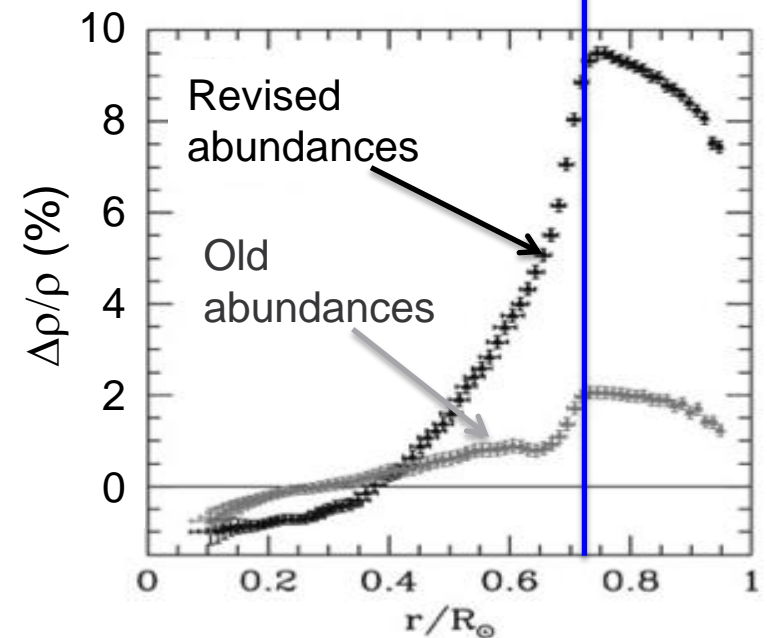
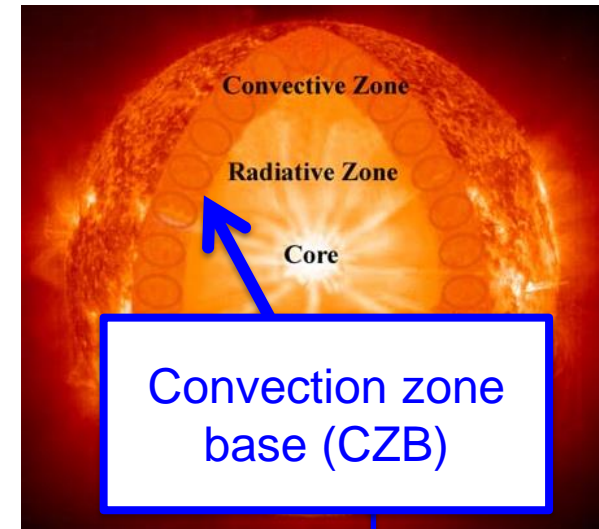




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Inputs:
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- 
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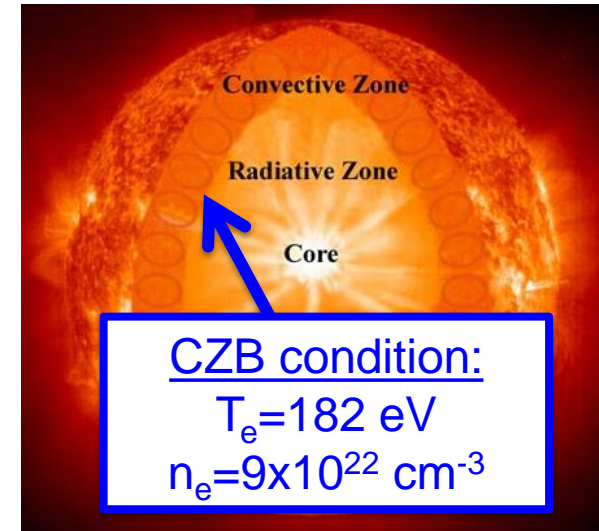
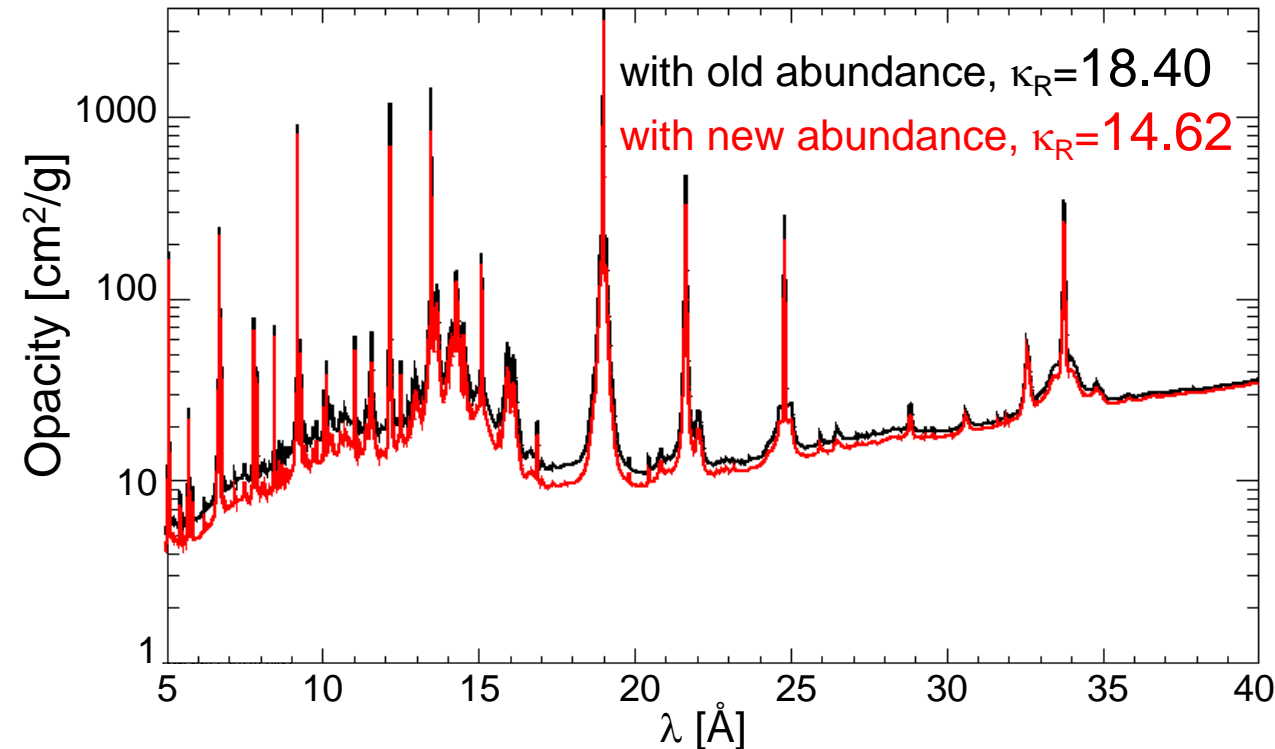
CZB location:  $1\sigma \rightarrow 13\text{-}30\sigma$





# Disagreement could be resolved if the true opacity is higher than predicted

Solar mixture opacity at **C**onvection **Z**one **B**ase (CZB)



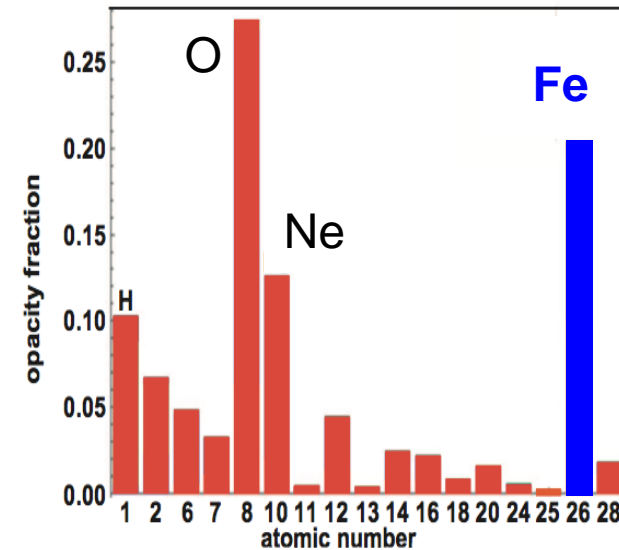
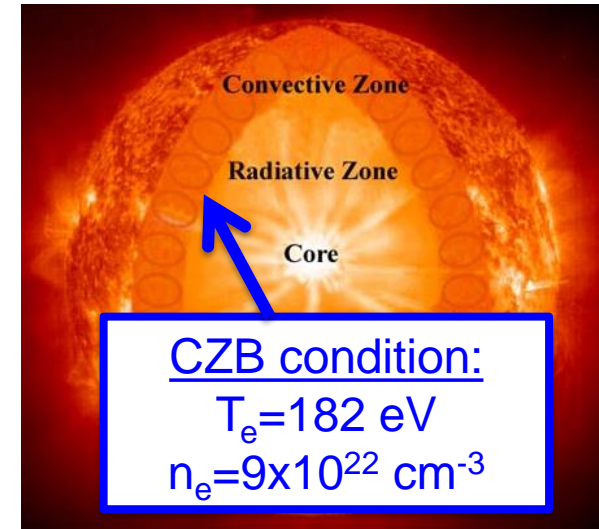
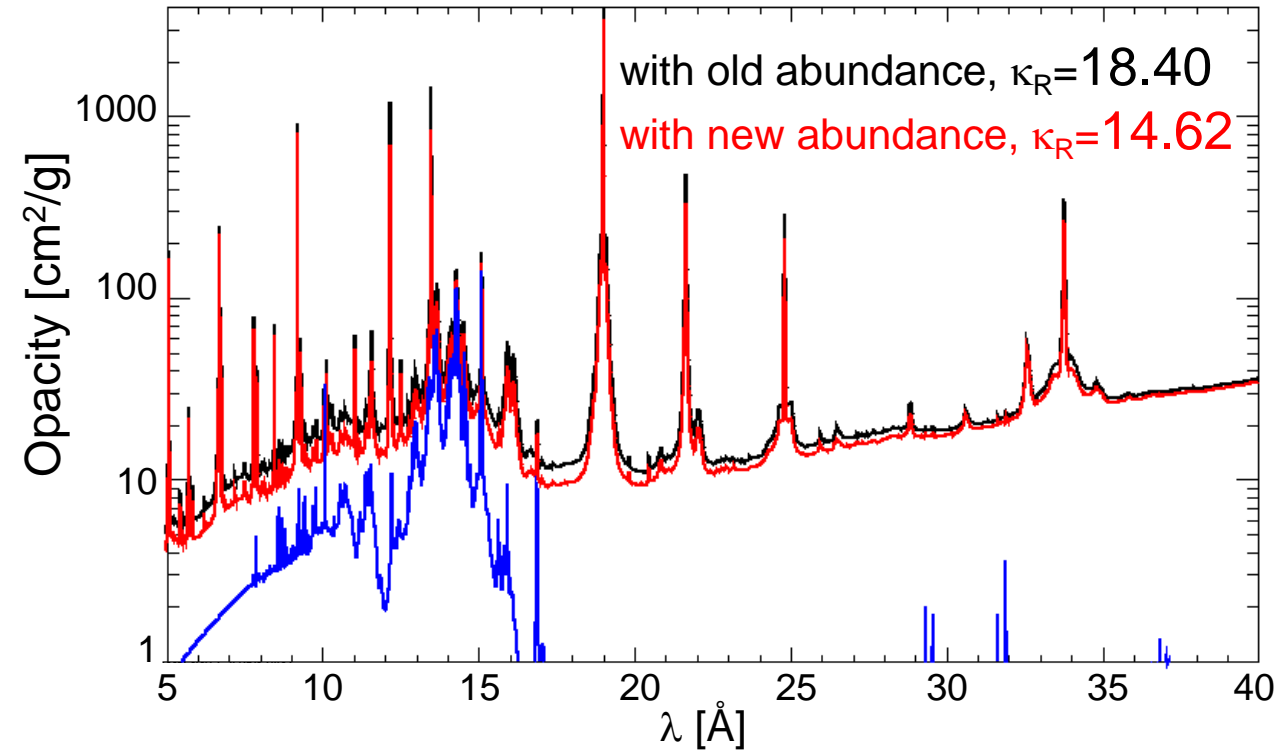
The disagreement is mostly resolved by ad-hoc opacity increases of 10-30%

Rosseland mean opacity  $\rightarrow$  heat transfer by radiation

$$\frac{1}{k_R} = \frac{\int \frac{1}{k_n} \frac{\mathcal{B}_n}{T} dn}{\int \frac{\mathcal{B}_n}{T} dn}$$

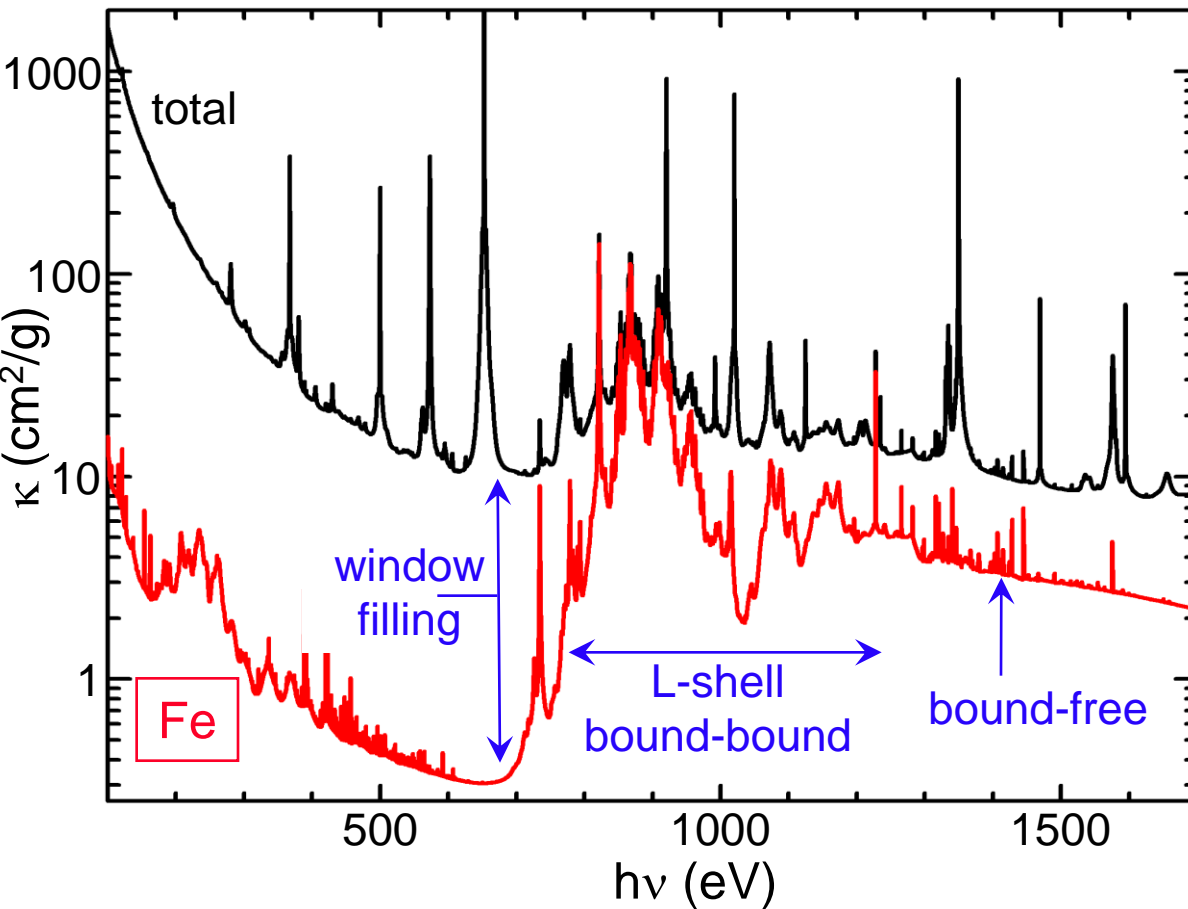
# Iron makes an important contribution to the solar opacity

Solar mixture opacity at **C**onvection **Z**one **B**ase (CZB)



CEA OPAS calculation

# At the convection zone base iron contributes opacity predominantly from L-shell transitions

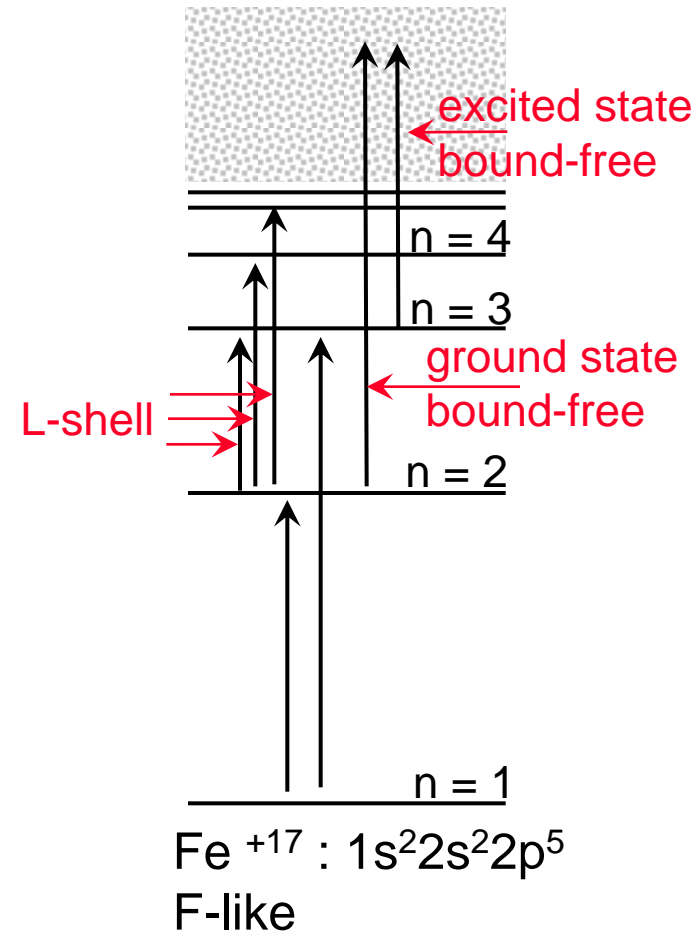


Opacity Project  
Badnell *et al.*,  
MNRAS 2005

$T_e$ ,  $n_e$  determine ionization  
{190 eV,  $9e22$  e/cc}  
Ionization determines which  
transitions are important  
At CZB, dominant charge  
states are  $Fe^{+16}$ ,  $Fe^{+17}$ ,  $Fe^{+18}$

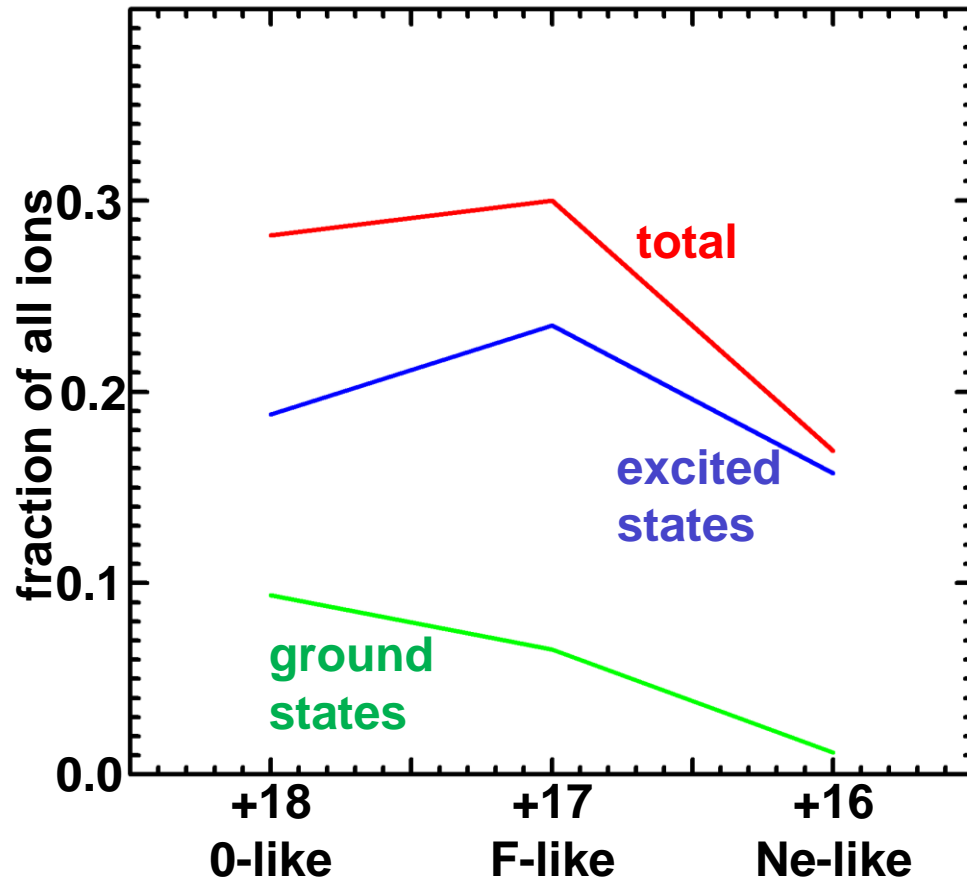
Photons are transported in opacity windows, but windows are filled by other elements

# Multiple entangled physical processes are a concern for opacity models



- Energy level structure and detail
- Multiply excited states
- Autoionizing levels
- Photoionization
- Line broadening
- Continuum lowering

# One challenge for iron opacity models is that significant population resides in myriad excited states



Iron at 195 eV,  $4e22$  electrons/cc  
SCRAM calculation

Accurate energy level description  
required for *all* excited states

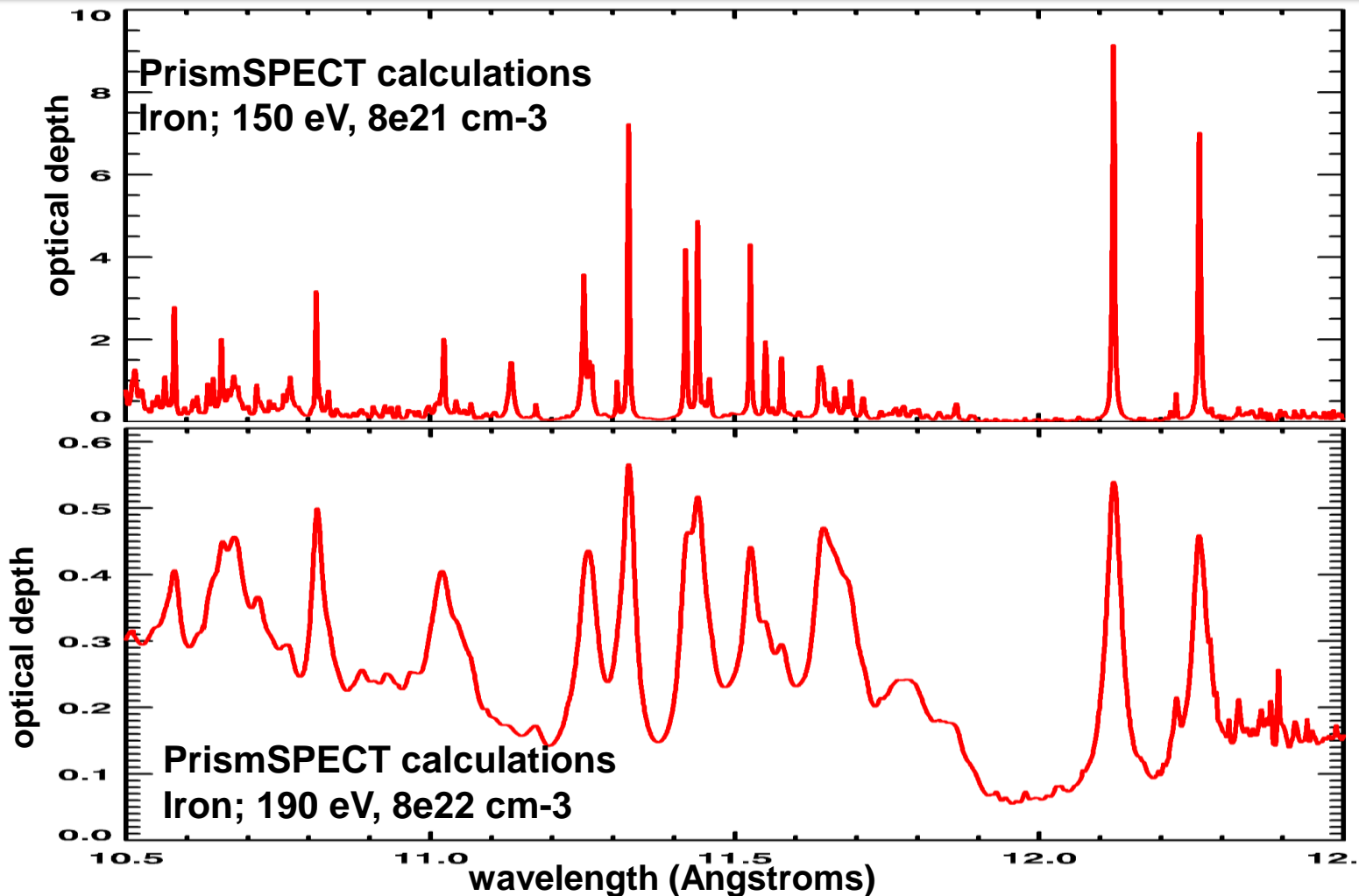
Plasma effects more easily modify  
excited states

Example: Ne-like iron  
excited state fraction = 93%  
ground state fraction = 7%

Implies a ~3% increase in excited state  
population causes ~40% decrease in  
ground state population

i.e., 40% decrease in lines originating  
from ground state

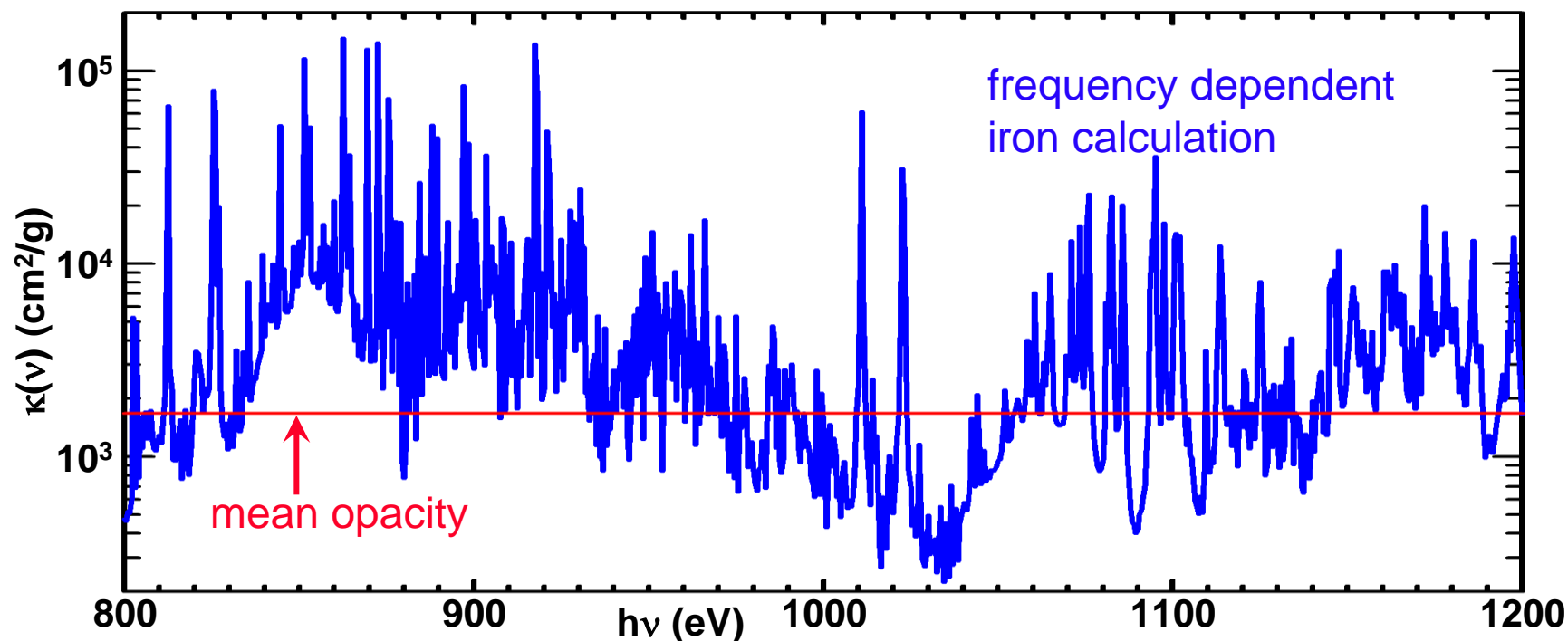
# Another challenge: Line broadening influences opacities but models for many-electron ions are untested



- Broadening tends to close the opacity windows between lines
- Modeling high-n and multiply-excited states is difficult



# Strategy: frequency-dependent transmission measurements test opacity model physics



Detailed information about the physical basis for opacity models is encoded in the frequency dependent transmission spectra.

# The importance of stellar opacity was recognized nearly a century ago, but no laboratory measurements have been done up to now. Why?

Eddington, “The Internal Constitution of the Stars”, 1926



High transmission accuracy is needed since  $\tau = -\ln(T)$

High accuracy requires:

Macroscopic samples uniformly heated to stellar interior conditions

Backlight bright enough to overcome emission at stellar interior temperatures

Stellar opacity measurements are possible for the first time:

MegaJoule class facilities like Z and NIF

3 decades of opacity research at smaller scale facilities to hone our approach

Advanced plasma diagnostic techniques

# Benchmark quality opacity experiment requirements have been developed over 30 years

Overarching requirements for each application:

Ideally: Reproduce the temperature, density, and radiation

Minimum: Reproduce the same charge states and measure the same transitions

Experiment requirements:

1. Accurate transmission measurements ( $\sim \pm 5\%$ )
2. Demonstrated uniformity
3. Reliable plasma diagnostics
4. Freedom from self emission
5. Freedom from background contamination
6. Multiple areal densities (for dynamic range and systematic error tests)
7. Thorough sample characterization
8. An evaluation of suitable the LTE approximation is
9. Multiple  $T_e$ ,  $n_e$  conditions, to aid disentangling physical effects
10. Multiple atomic number elements, to aid disentangling physical effects and help verify robustness against systematic errors
11. Multiple experiments of each type, to confirm reproducibility
12. Peer review and documentation

Example references:

Davidson *et al.* Appl. Phys. Lett. 1988

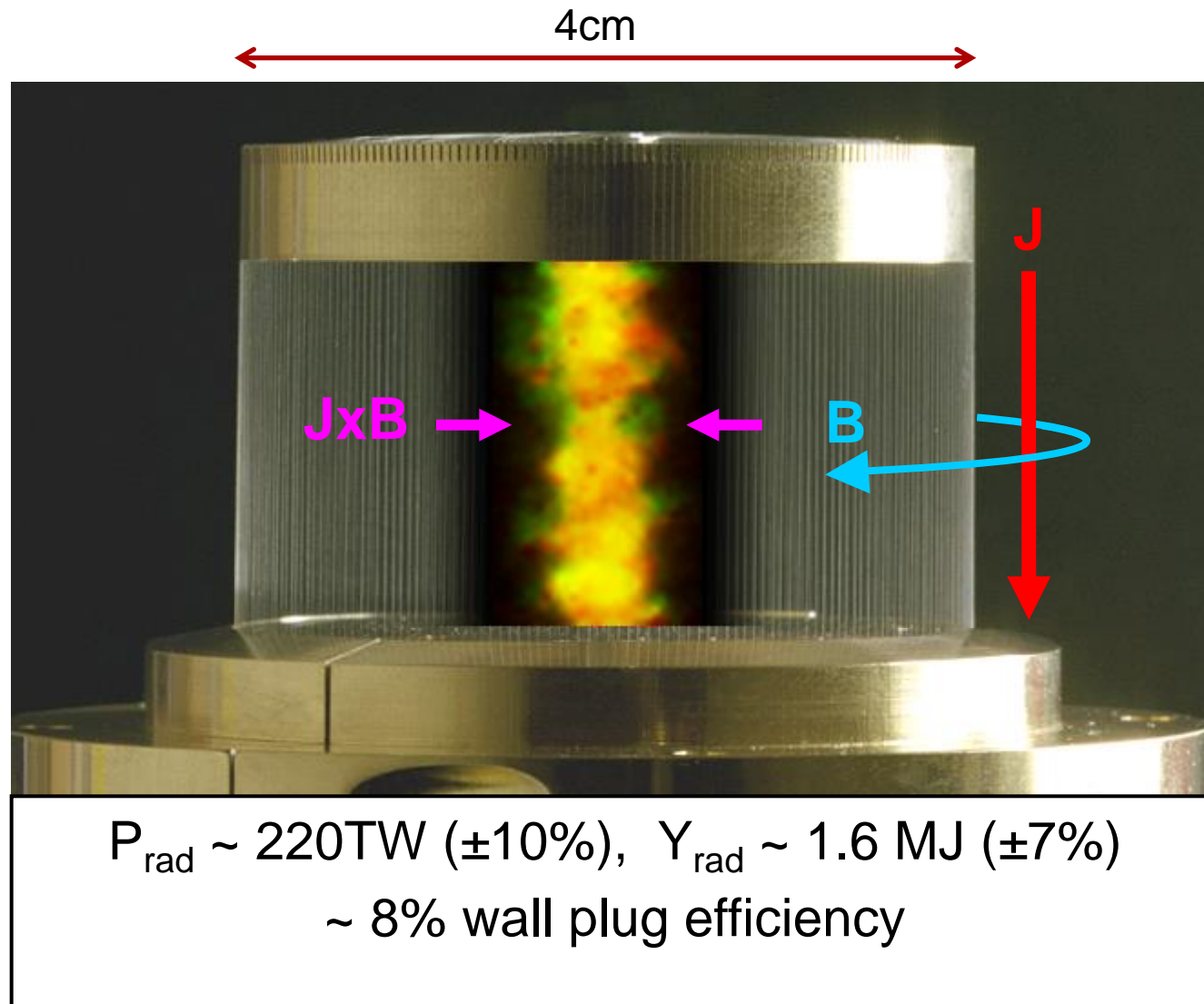
Perry *et al.* Phys. Rev. Lett 1991

Foster *et al.* Phys. Rev. Lett. 1991

Perry *et al.* Phys. Rev. E 1996

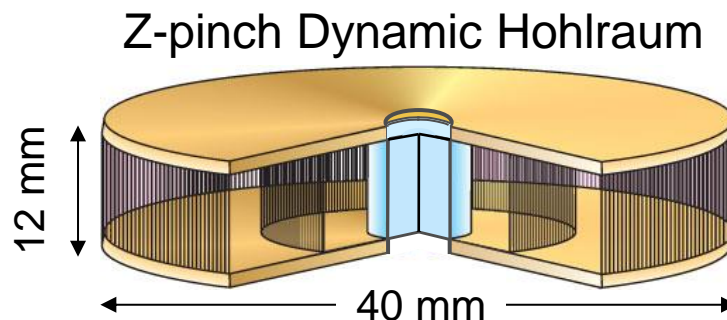
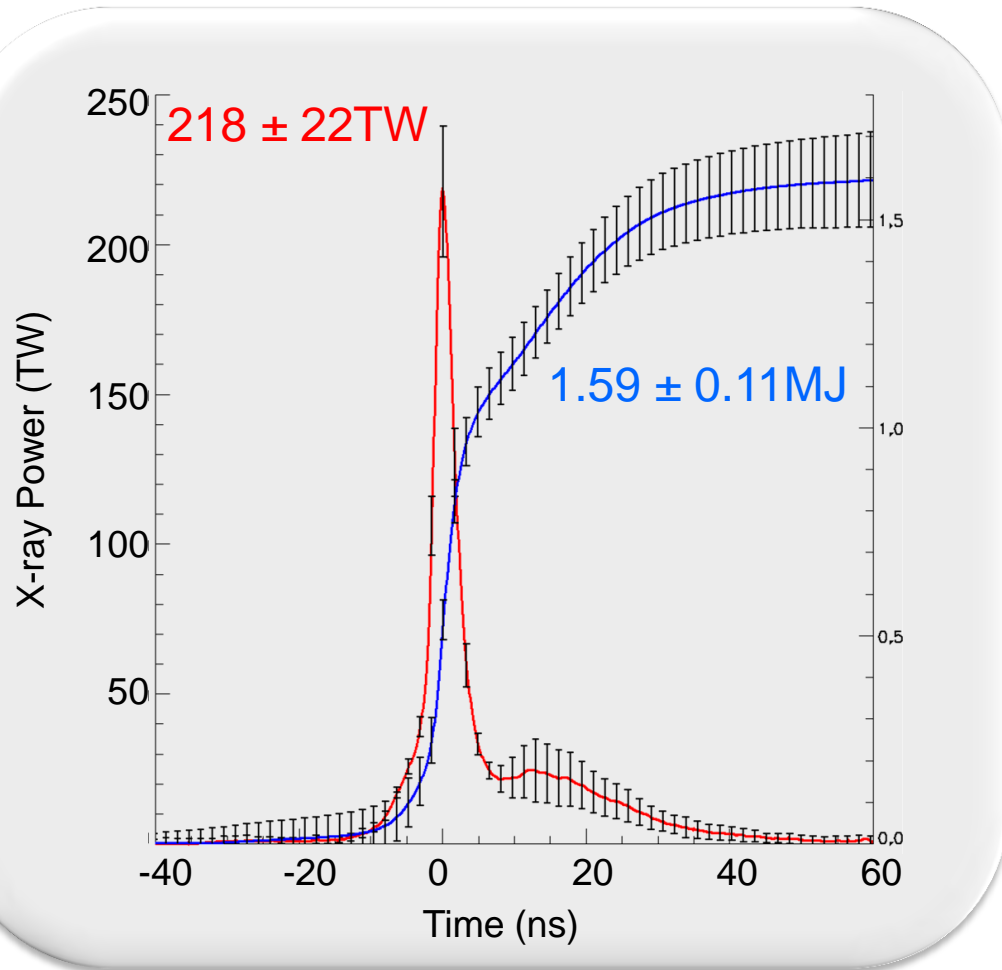
Springer *et al.* JQSRT 1997

# We use the Z machine to create energetic and powerful x-ray sources



# The ZPDH x-ray emission is reproducible to $\pm 10\%$ in peak power and $\pm 7\%$ in energy

## Radial X-ray Power and Energy (20 shot average)

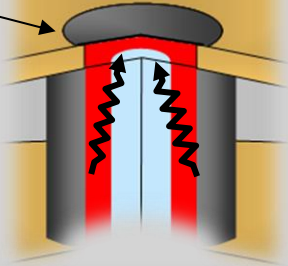


	ZR	Z
Marx Energy	21 MJ	12 MJ
I <sub>peak</sub>	25.8 MA	21.8 MA
Mass	8.5 mg	3.8 mg
Peak Power	220 TW (10%)	120 TW (14%)
Radiated Energy	1.6 MJ (7%)	0.82 MJ (17%)

# The ZPDH radiating shock is used to both heat and backlight samples to stellar interior conditions.

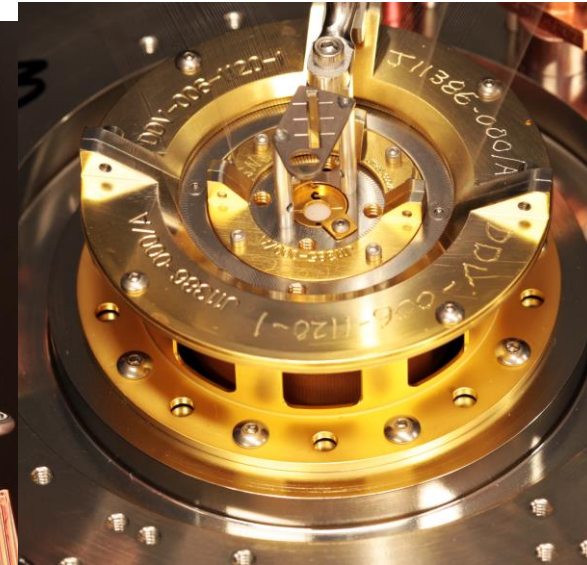
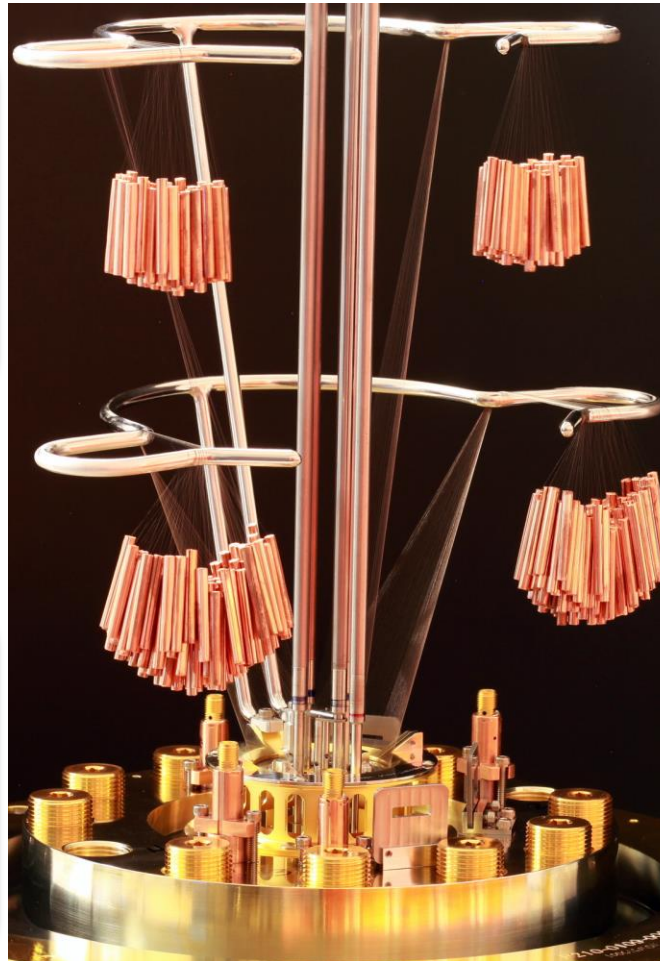
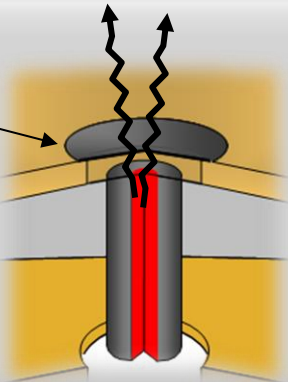
**Foil is heated during  
the ZPDH implosion**

Thin  
Foil



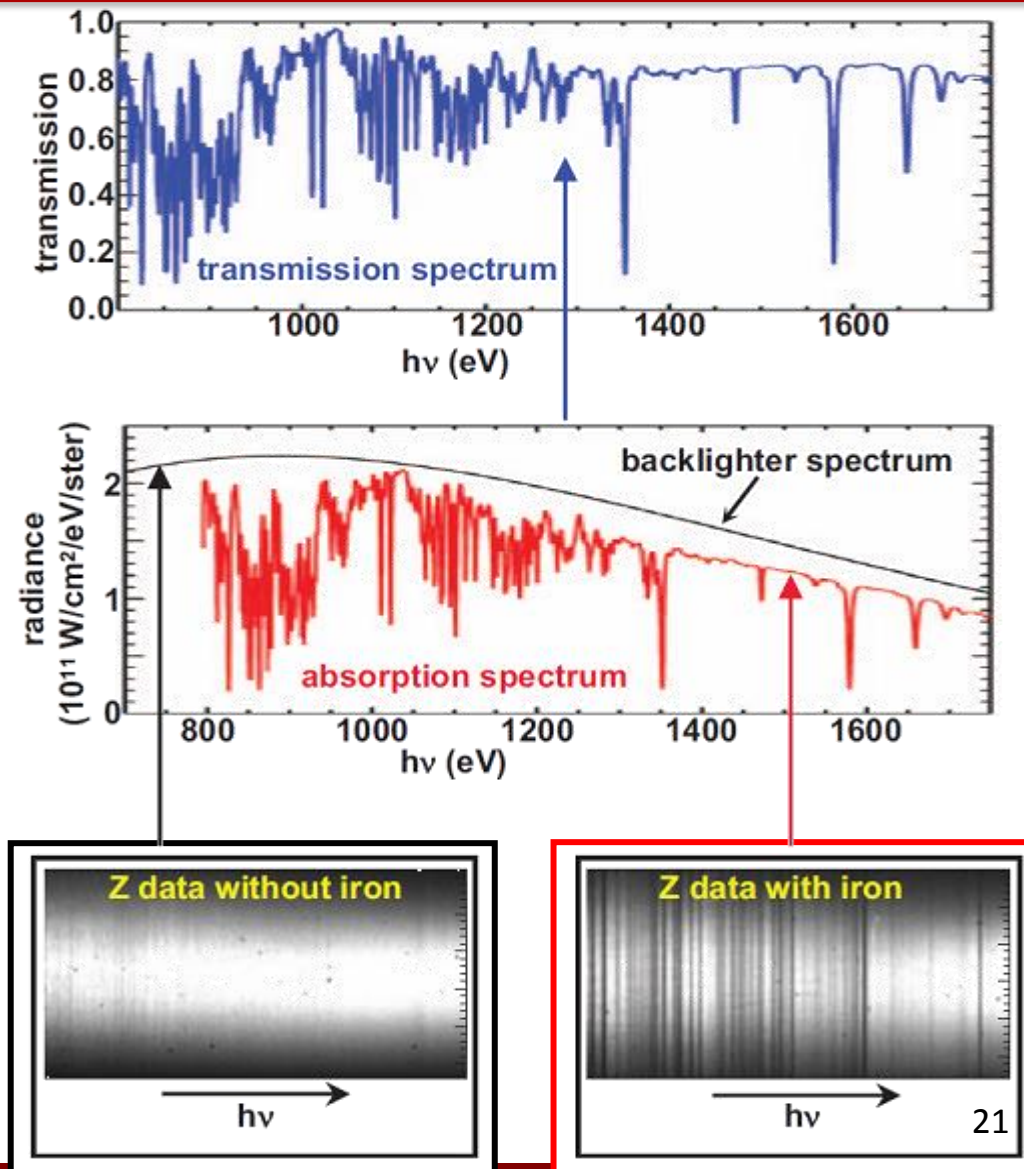
**Foil is backlit  
at shock stagnation**

Thin  
Foil

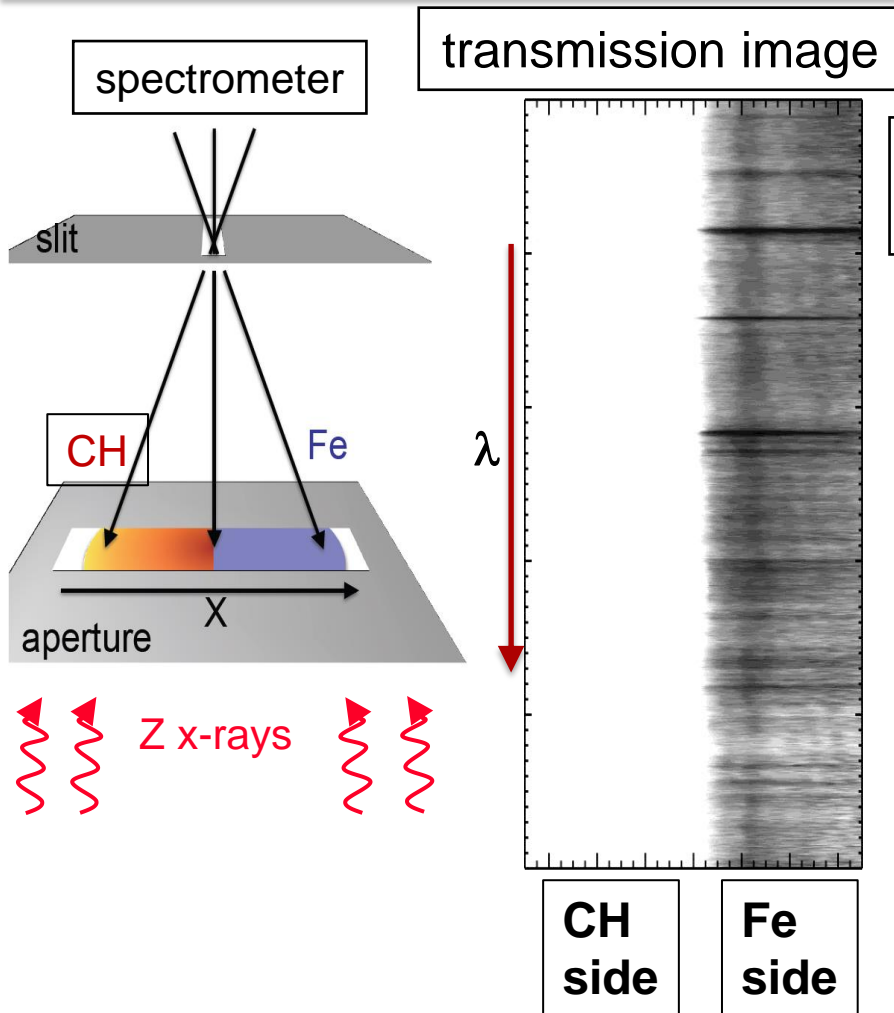




# Transmission is inferred by dividing the attenuated spectrum by the unattenuated spectrum.



# Measurements with half-moon shaped samples enable transmission determination from single experiments



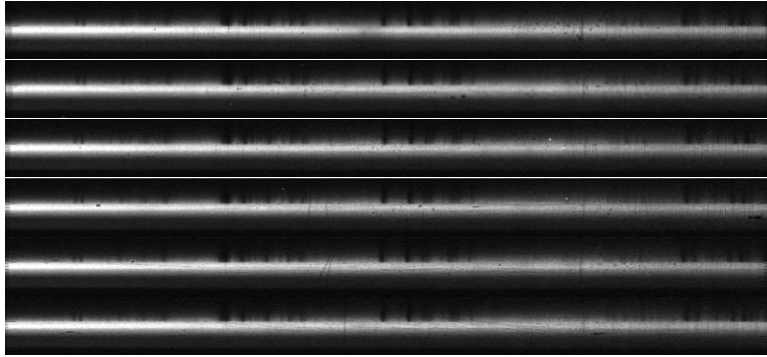
**Backlit spectra with and without sample determine transmission**

T.S. Perry *et.al.* Phys. Rev. E (1996)  
J.E. Bailey *et.al.* Phys. Plasmas (2009)

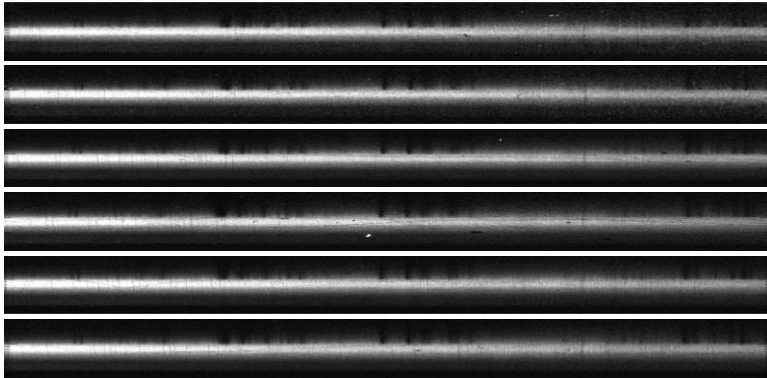
# Opacity data are recorded with an array of crystal spectrometers



# Hundreds of spectra were measured and analyzed to support the experiment reliability and reproducibility

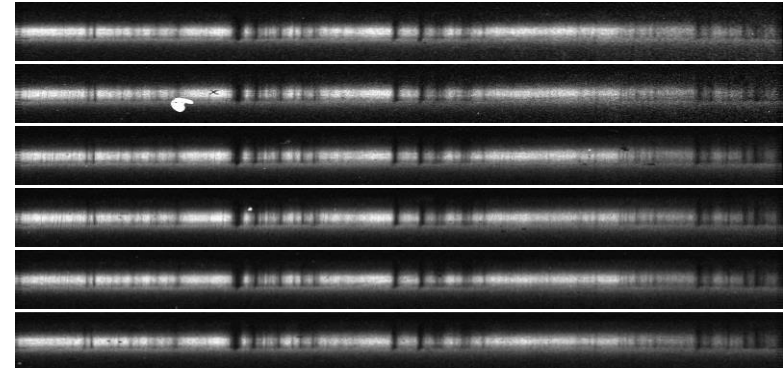
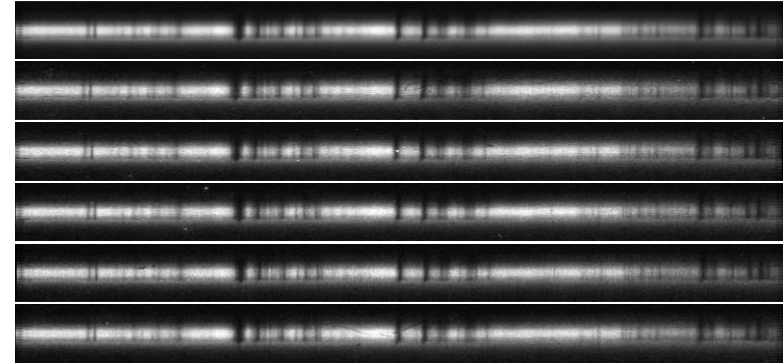


ccp4a



ccp4b

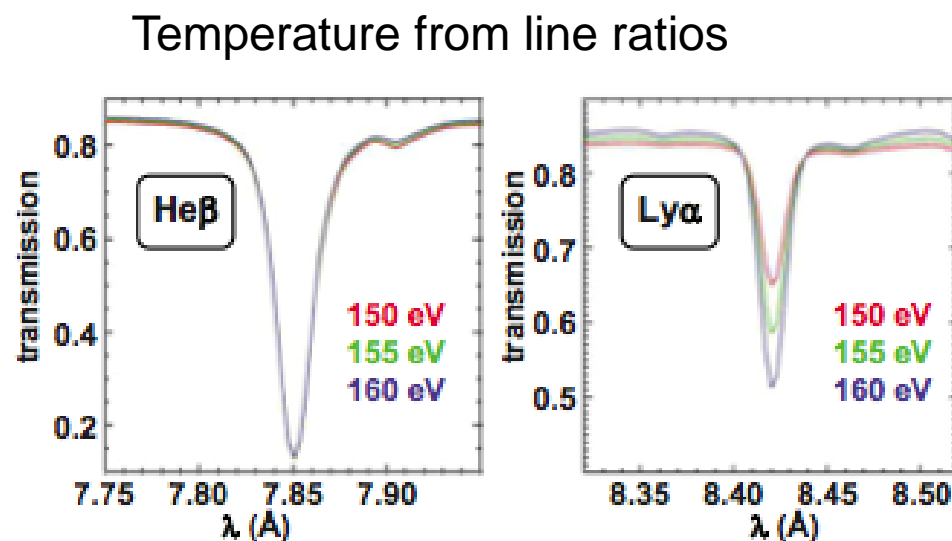
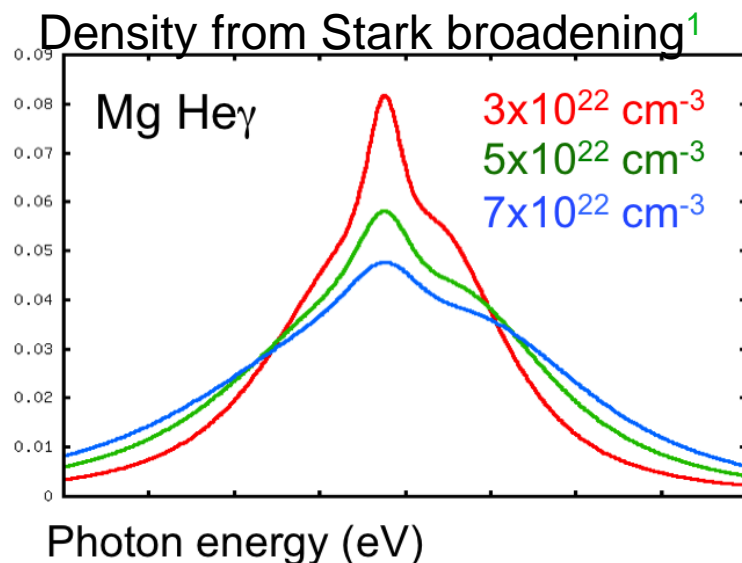
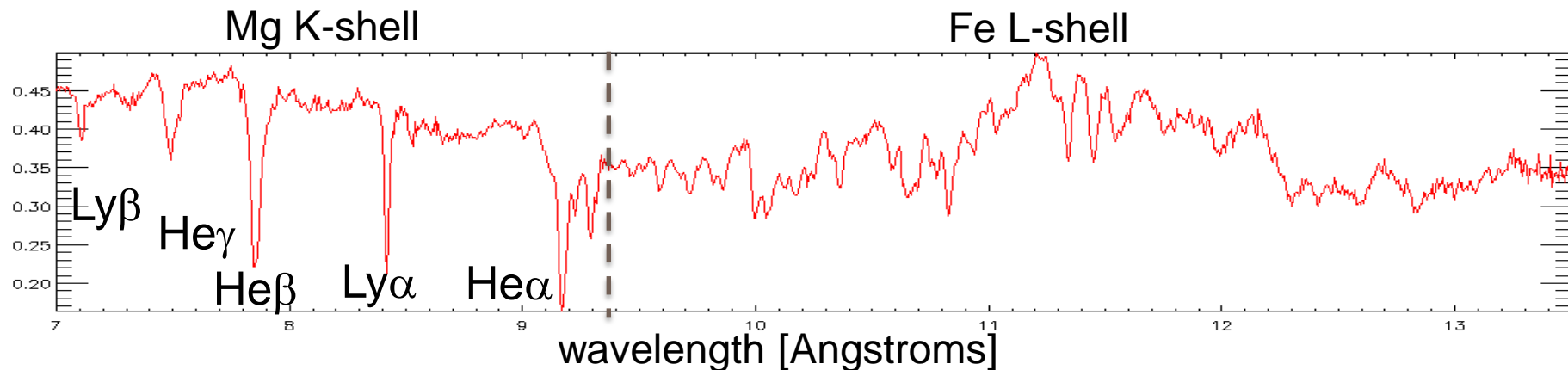
ccp10a



ccp10b

Data from z2762

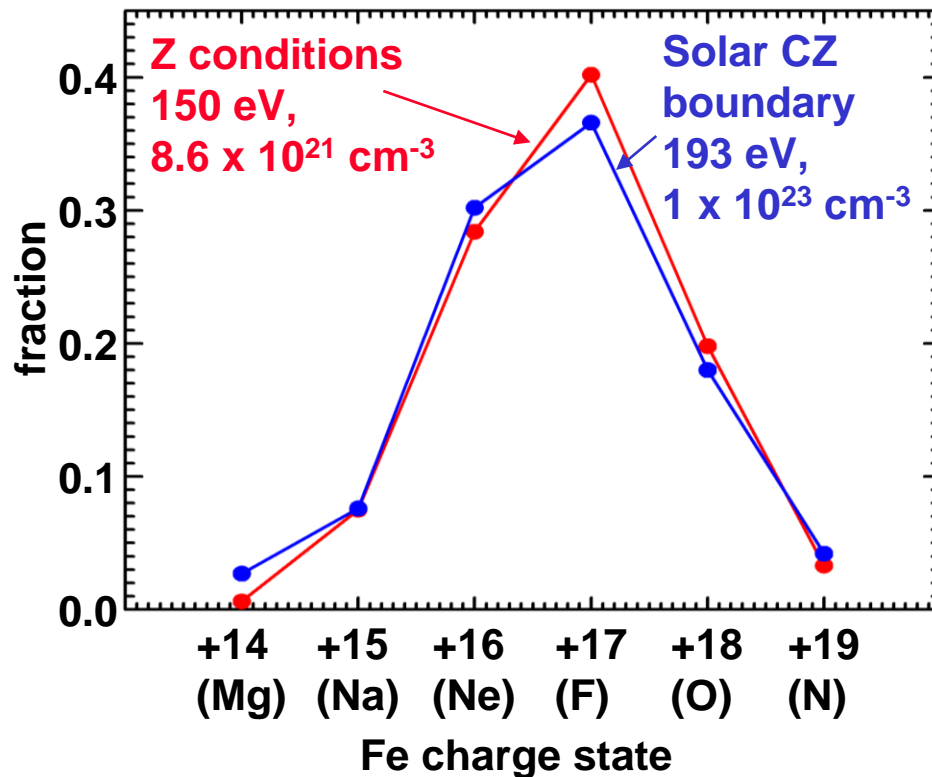
# Plasma conditions are inferred by mixing Mg with Fe and using K-shell line transmission spectroscopy



R. C. Mancini, comp. phys. commun. (1991)  
T.N. Nagayama et. al. RSI (2013)  
T.N. Nagayama et. al. POP (2014)



# In 2007, Z experiments produced the iron charge states that exist in the solar interior



Producing the correct charge states enables opacity model tests:

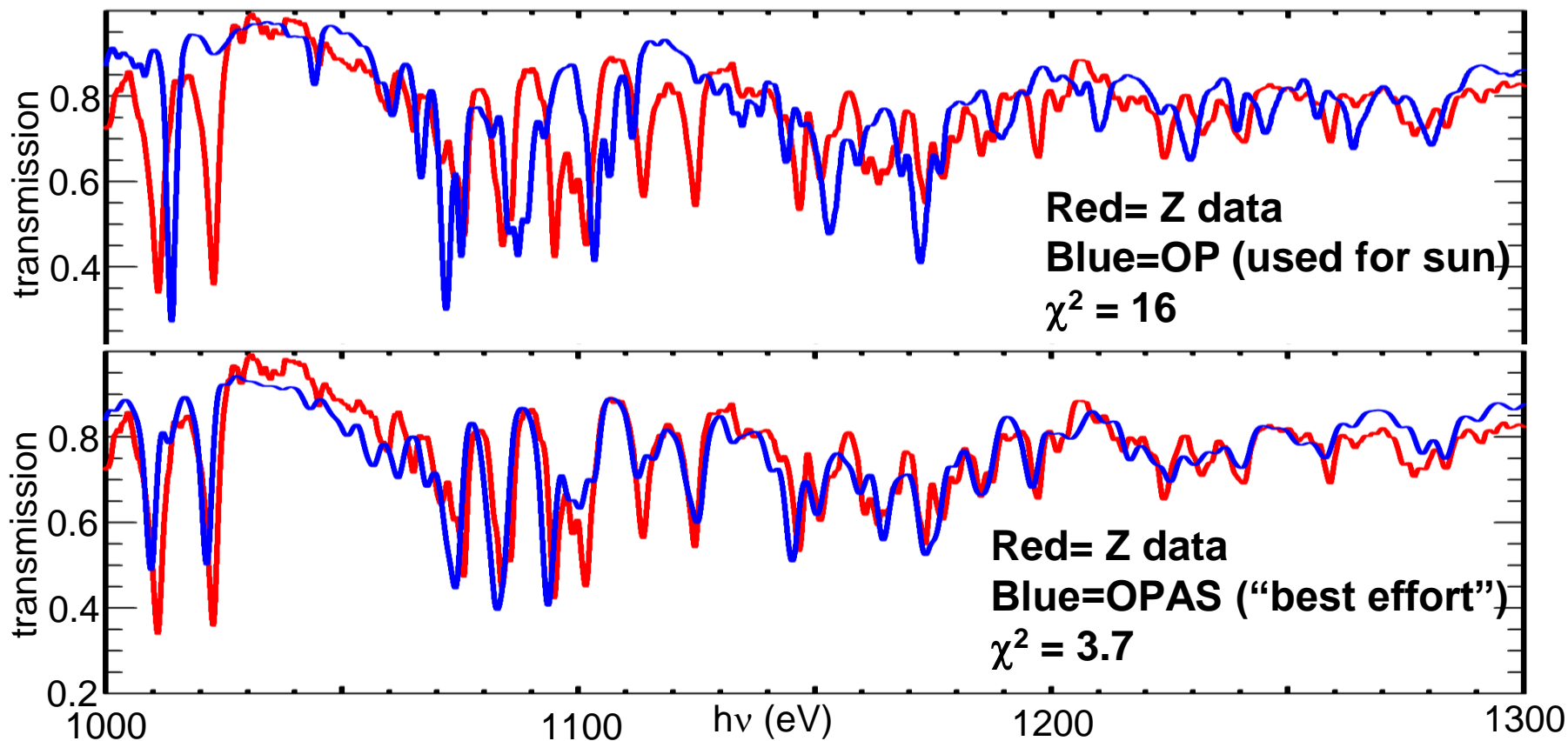
- 1) Charge state distribution
- 2) Energy level description

High density and high temperature studies required further progress

[Bailey *et al.*, PRL (2007)]



# The 2007 Z data was matched well by “best-effort” models, but not by a model used in solar research

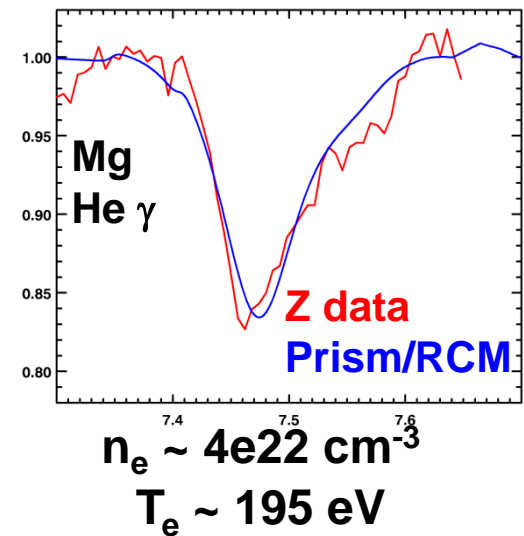
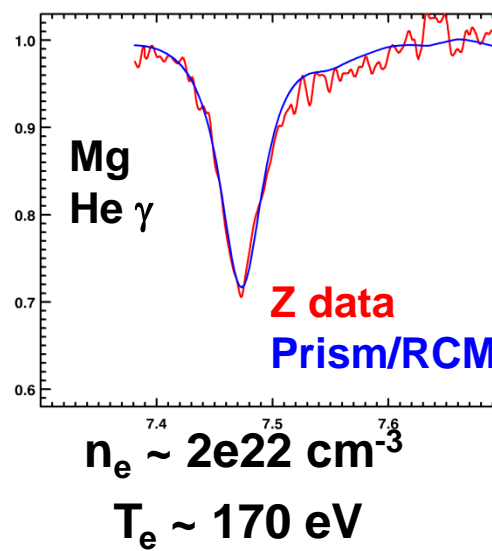
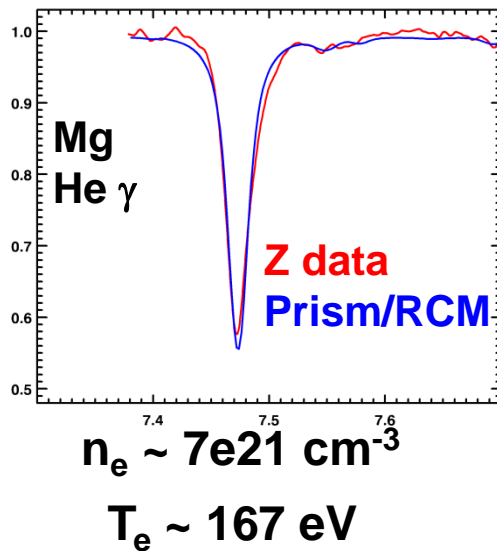
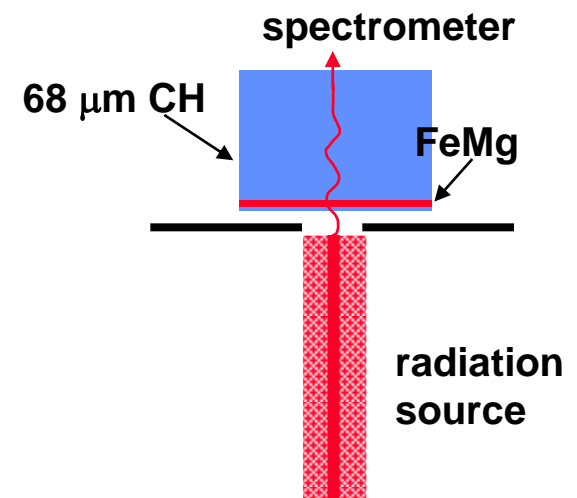
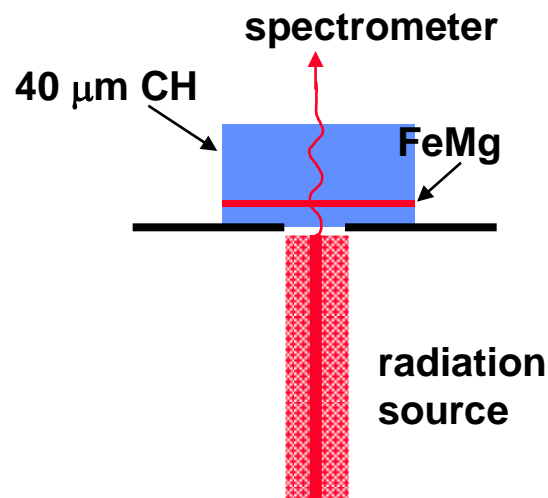
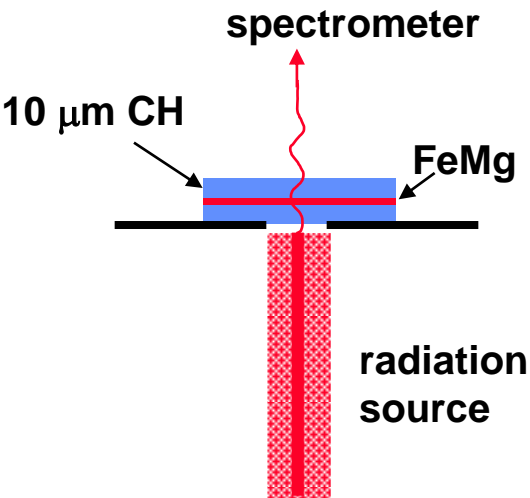


**OP Rosseland mean is ~ 1.5x lower than OPAS at Z conditions.**

**If this difference persisted at solar conditions, it would solve the CZ problem**

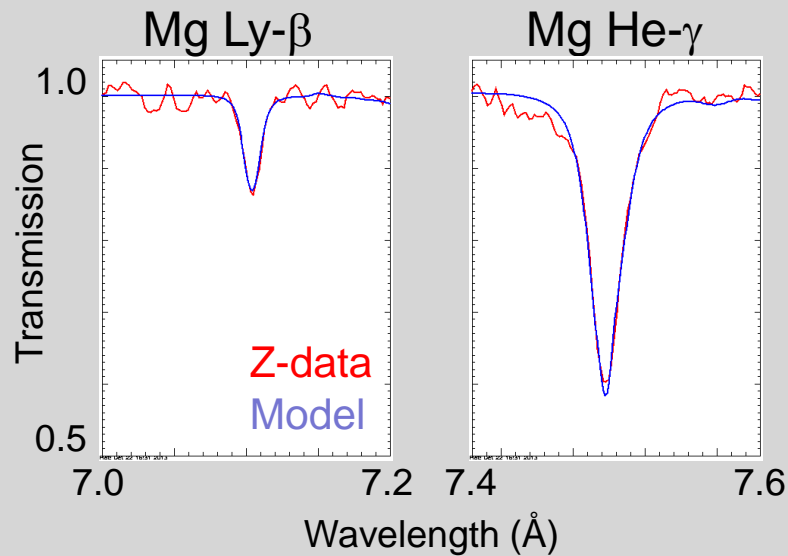
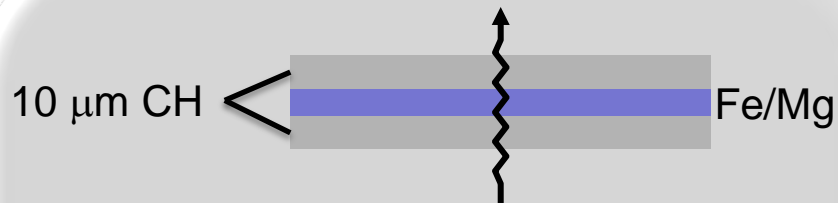
**Experiments at higher density needed**

# Adjusting the CH tamper thickness controls the opacity sample density and temperature



# The achieved temperature and density depend on the target design.

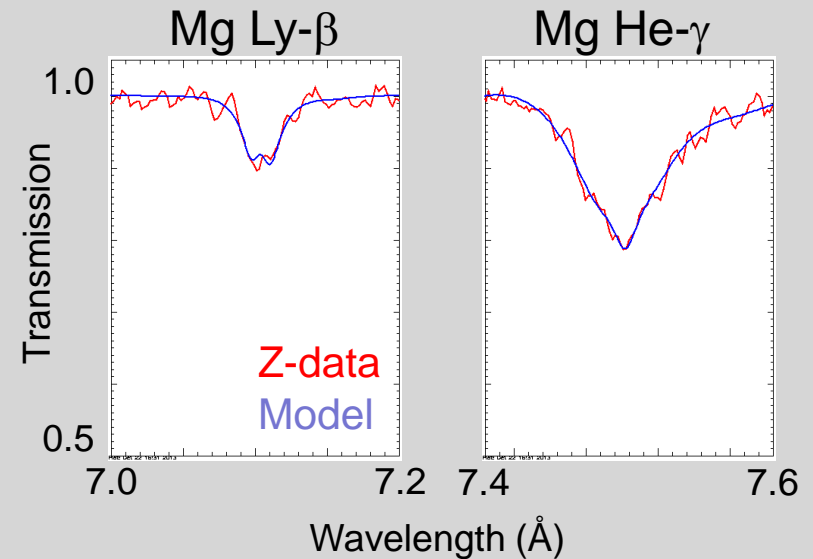
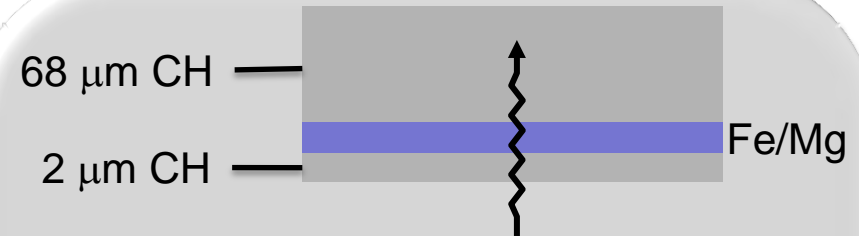
## Thin Tamper



$$T_e = 156 \pm 6 \text{ eV}$$

$$n_e = 6.9 \pm 1.7 \times 10^{21} \text{ cm}^{-3}$$

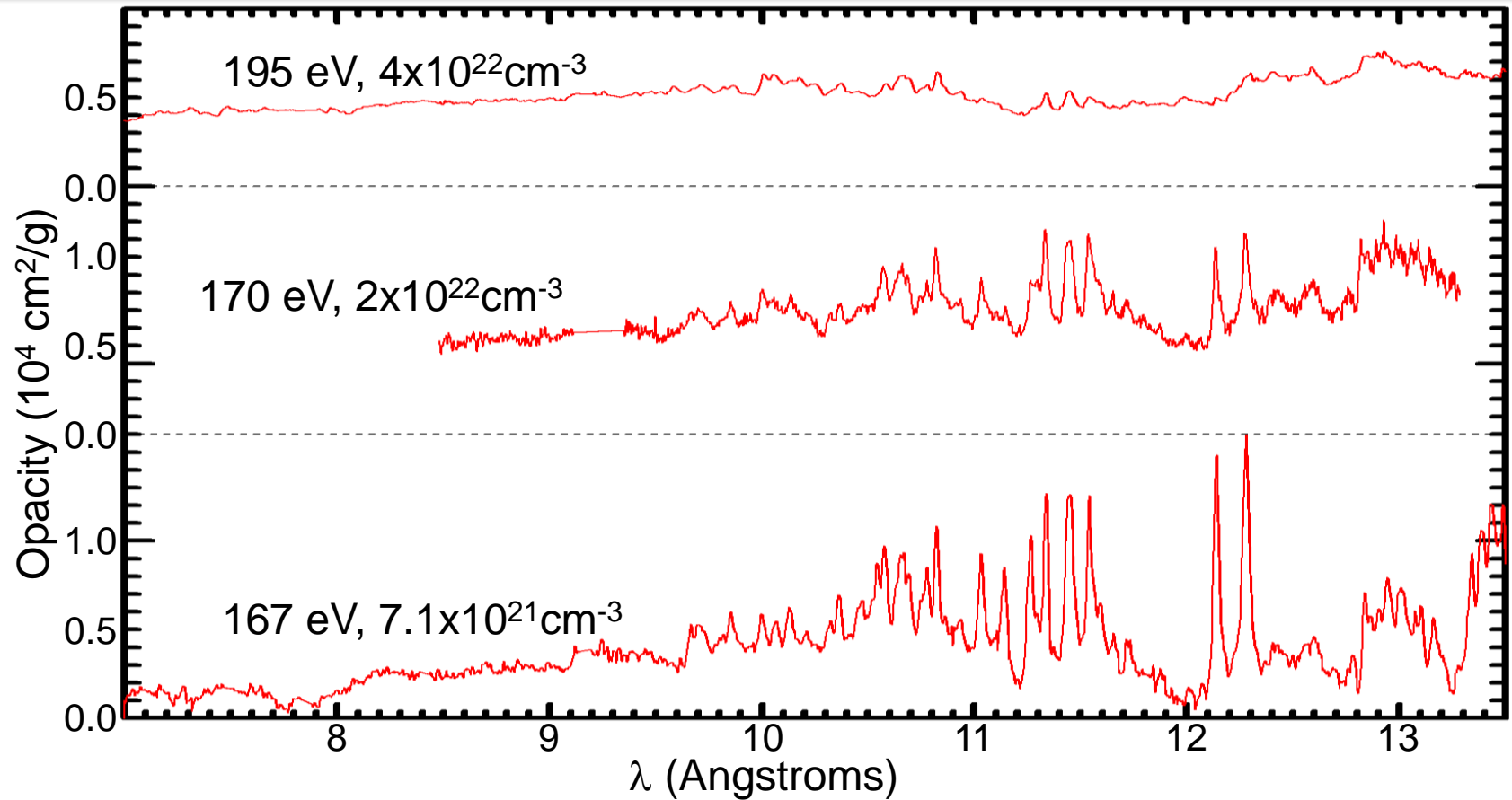
## Thick Tamper



$$T_e = 196 \pm 5 \text{ eV}$$

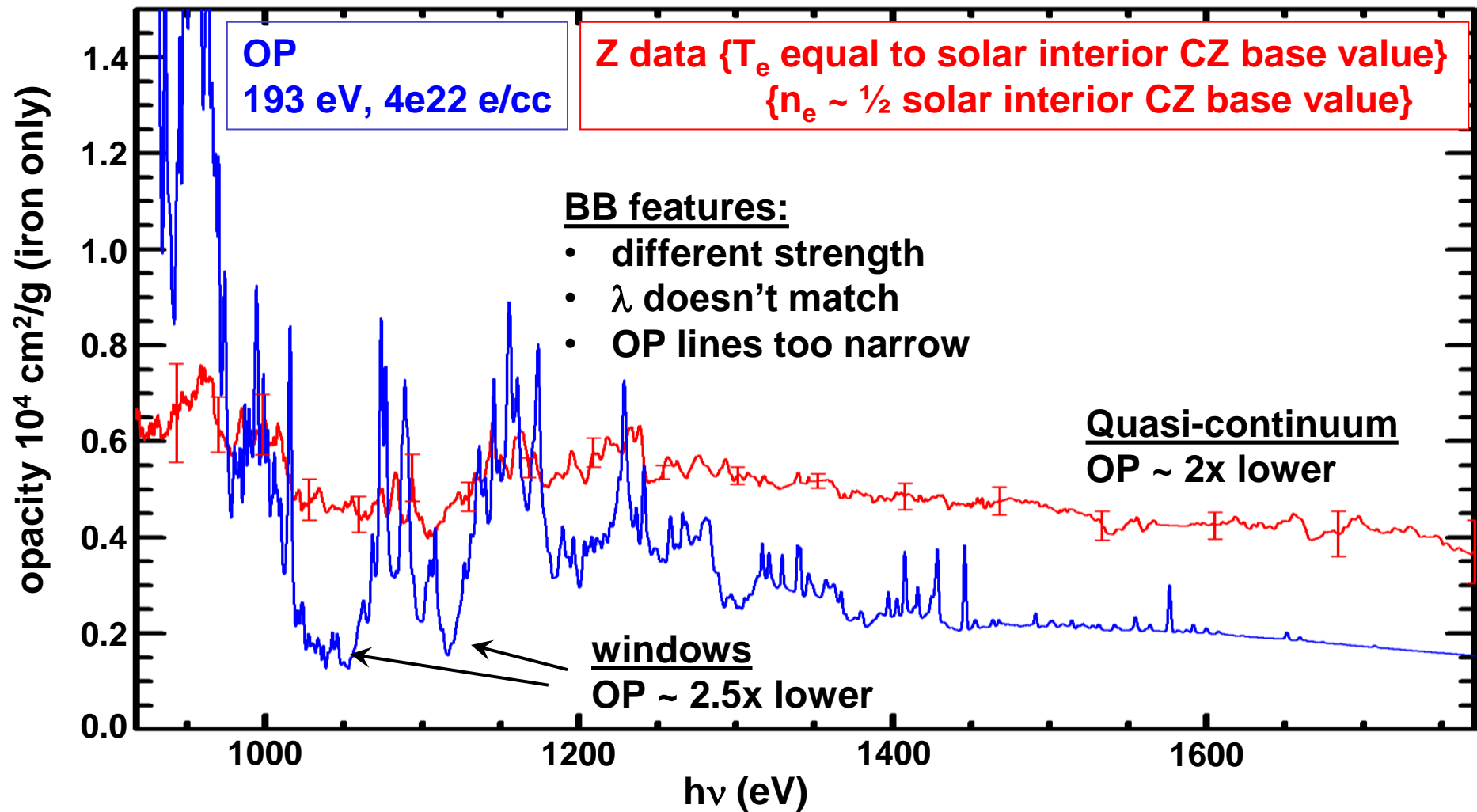
$$n_e = 39. \pm 9. \times 10^{21} \text{ cm}^{-3}$$

# Iron opacity spectra have been acquired at conditions approaching the solar convection zone base



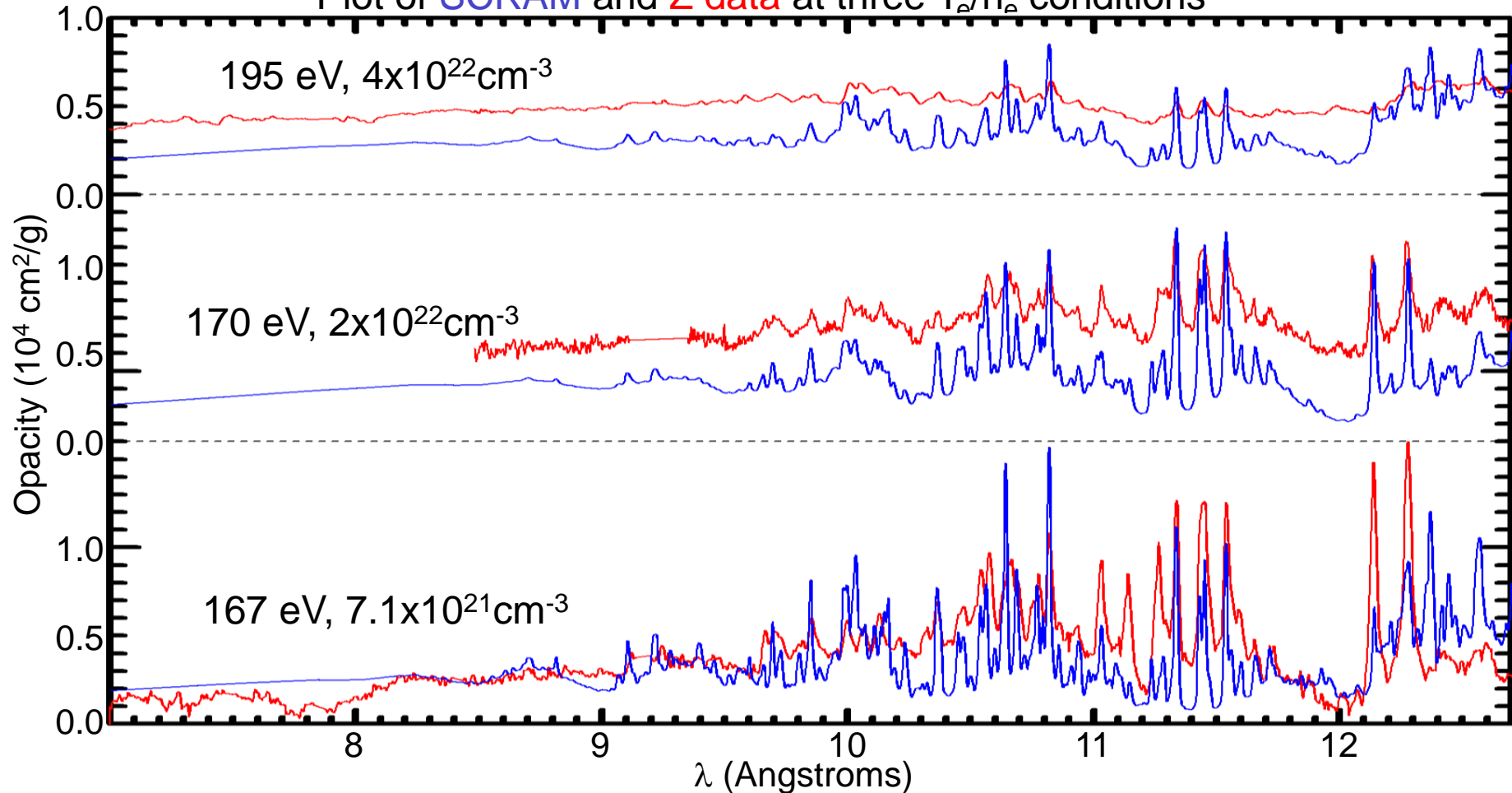
- Multiple conditions helps dis-entangle the different physical processes
- Some clear trends are observed as  $T_e$ ,  $n_e$  increase: shorter, fatter lines; windows fill in; quasi-continuum opacity increases

# The OP opacity model is used in solar models but it disagrees with Z iron plasma opacity measurements



# “Best Effort” opacity models “match” the data at lower $T_e/n_e$ conditions but not at conditions near the CZB

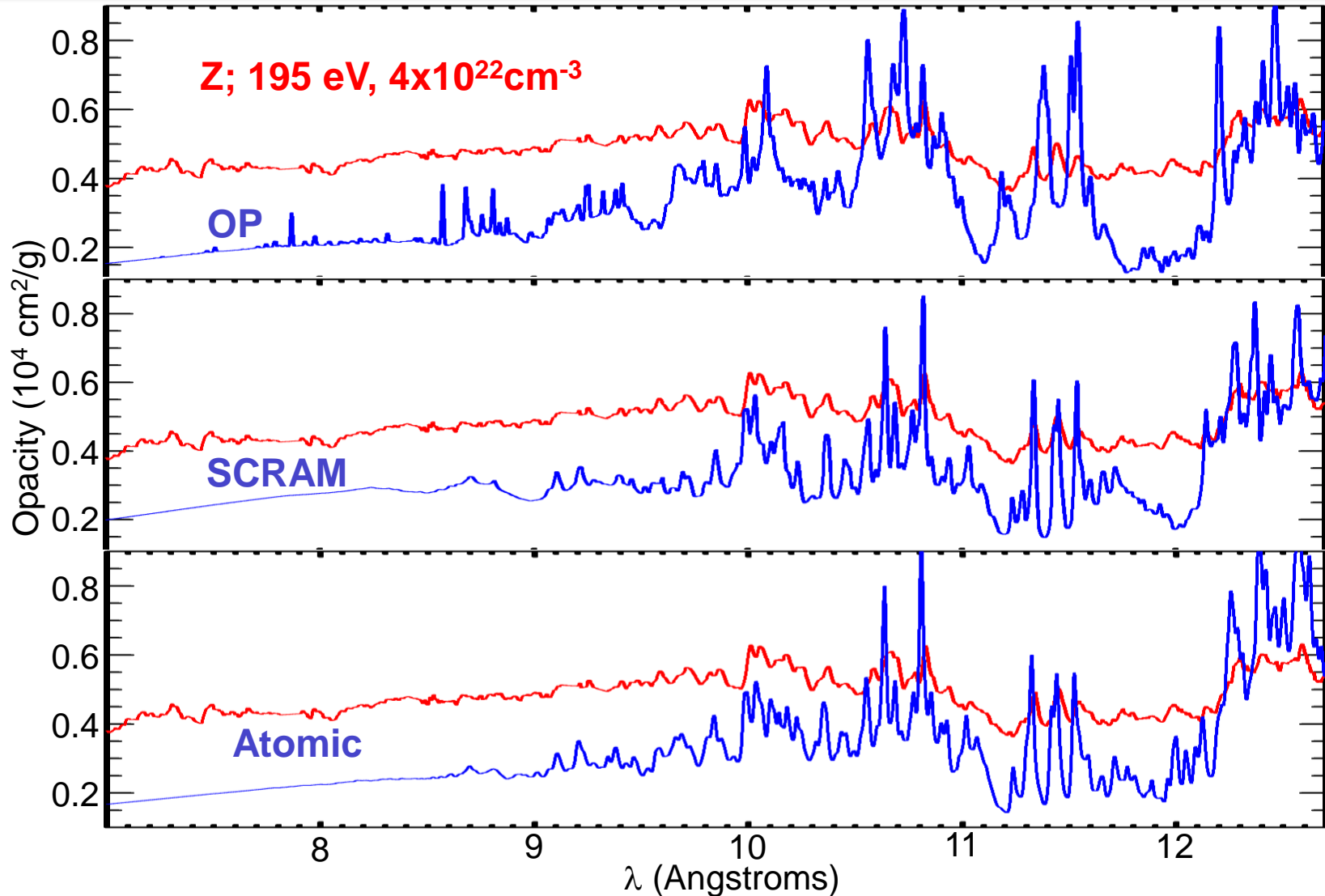
Plot of **SCRAM** and **Z data** at three  $T_e/n_e$  conditions



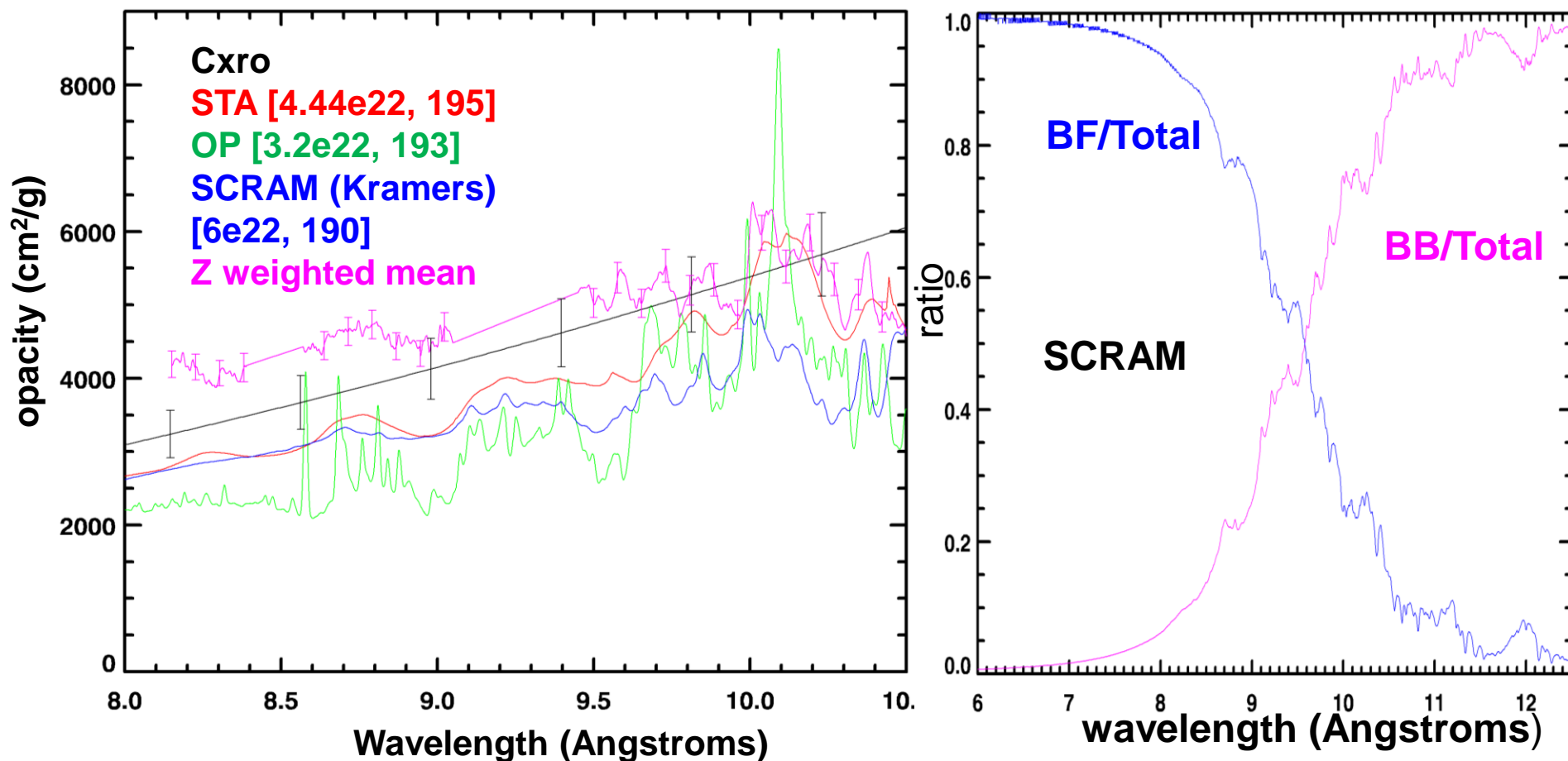
**At high temperature , density, calculations are generally lower than the data**



# No model examined up to now has satisfactory agreement with iron opacity measured at near-CZB conditions



# The measured hot iron plasma opacity exceeds past measurements of the cold iron opacity



This could imply errors in either our hot opacity measurement, the cold opacity measurement, photoionization calculations, or BB opacity calculations

# The serious implications of this research mandate scrutiny of random and systematic uncertainties

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Random error determination: average many spectra from multiple experiments

Systematic error determination:

Experiment tests of hypotheses

Simulations of sample conditions, postprocessed

Specific potential errors investigated include:

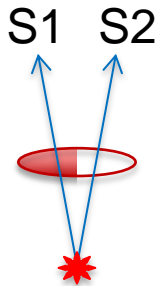
Transmission accuracy – tamper-only experiments, reproducibility, Beer's Law scaling

Sample uniformity – direct measurements, simulations

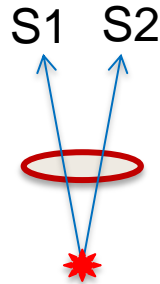
Self emission – direct observations, Beer's Law tests, post-processed simulations

Tamper attenuation – change tamper material, Beer's Law, post-process simulations

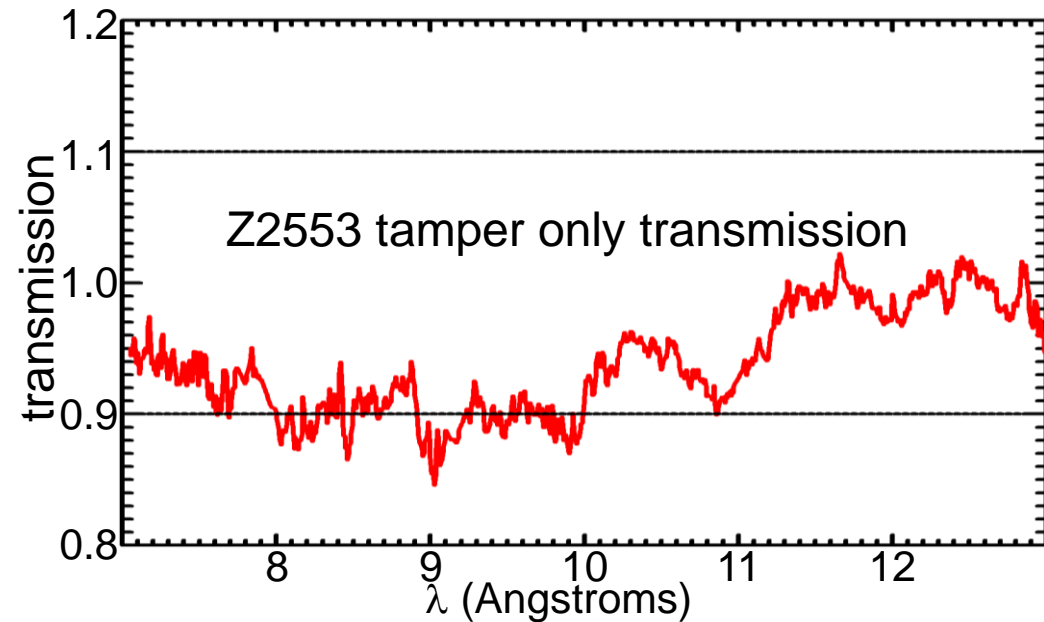
# Tamper-only experiments confirm transmission accuracy



Half-moon sample:  
Transmission =  $S1/S2$

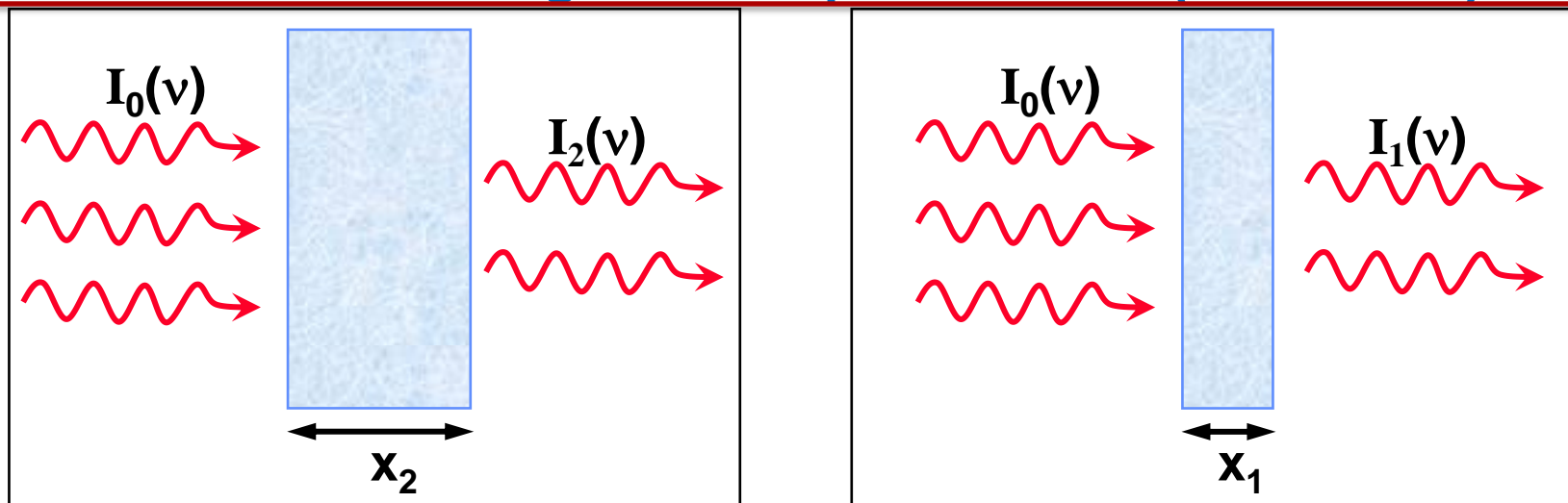


Tamper-only sample:  
Transmission =  $S1/S2 = 1.0$



- For this example the average absolute error is ~7%
- Errors are further reduced by averaging repeated experiments
- We repeat this test in every experiment series to avoid the possibility of anomalously large errors

# Possible experiment flaws can be evaluated from transmission scaling with sample thickness (Beer's Law)



Expected scaling with thickness :  $T_1 = T_2^{(x_1/x_2)}$

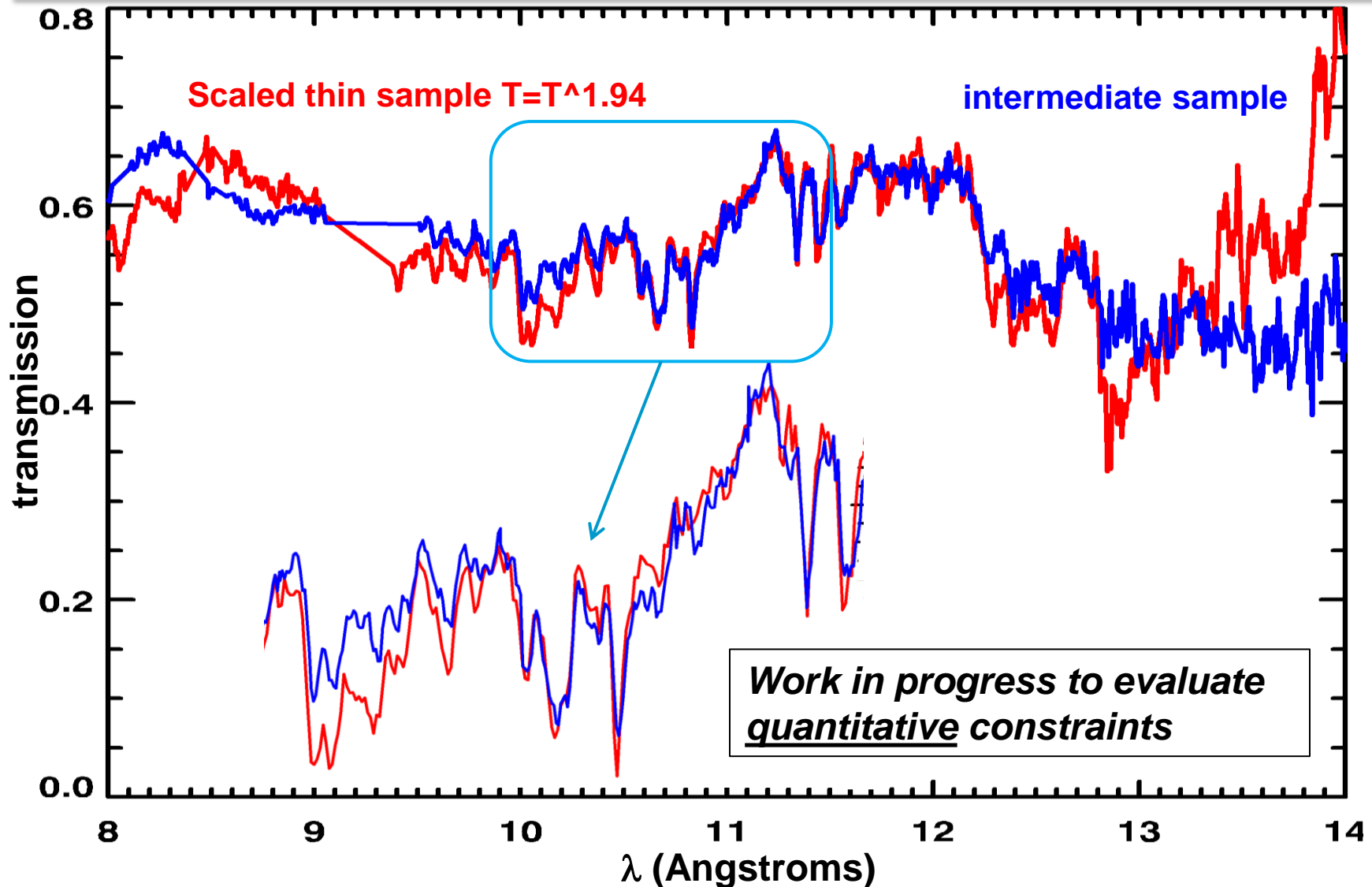
e.g., if  $x_2 = 2 * x_1$ , then  $T_2 = T_1 * T_1$

experiment problems cause transmission scaling to deviate:

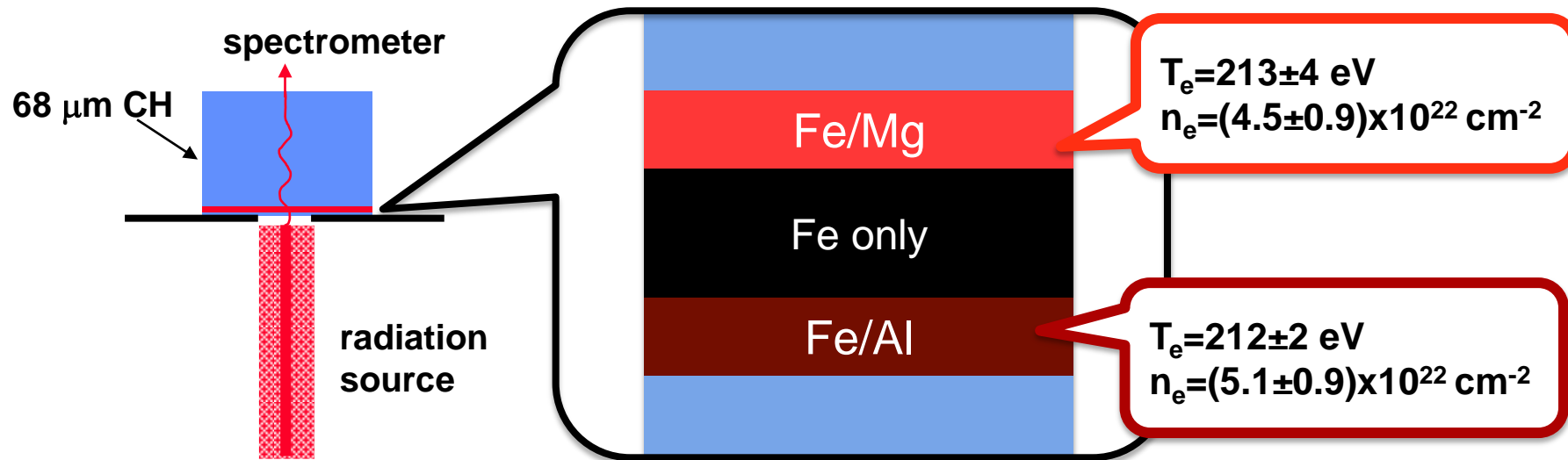
- Sample emission
- Background subtraction
- Crystal defects
- Gradients

***Most potential experiment problems cause the scaled thin sample transmission to be lower than the thick sample transmission***

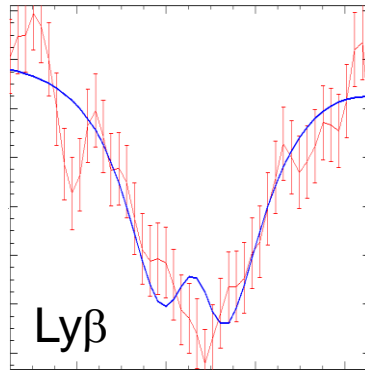
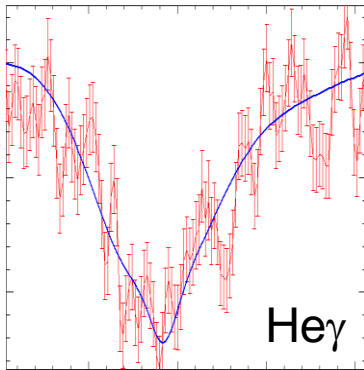
# Beers Law test confirms reliability of high Te/ne iron data in the 8-13 Angstrom range



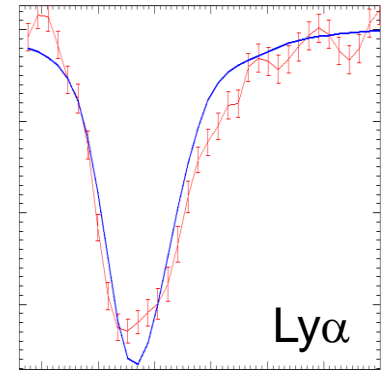
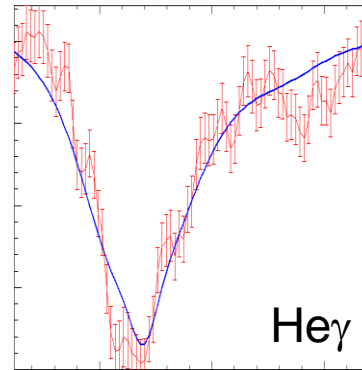
# Direct uniformity measurement confirmed that there is no significant spatial gradient in the sample



## Mg lines



## Al lines





# Simulations can provide significant insight and quantitative estimates for systematic error hypotheses

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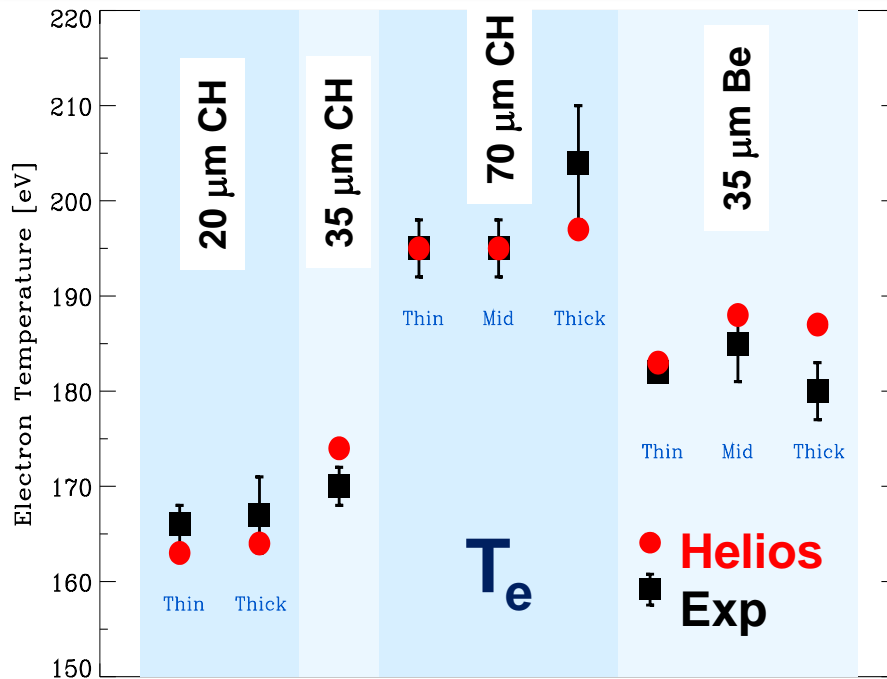
1. Why do  $T_e$  and  $n_e$  increase with tamper thickness?
2. Is sample shadowing predicted to be important? [Yes]
3. Are spatial gradients predicted to be important? [No]
4. Are temporal gradients predicted to be important? [No]
5. Is the time-integrated Fe/Mg self-emission negligible compared to the time-integrated backlighter? [No]
6. Is the time-integrated CH self-emission negligible compared to the time integrated CH attenuated backlighter? [Yes]

We use 1D HELIOS simulations to help answer these questions

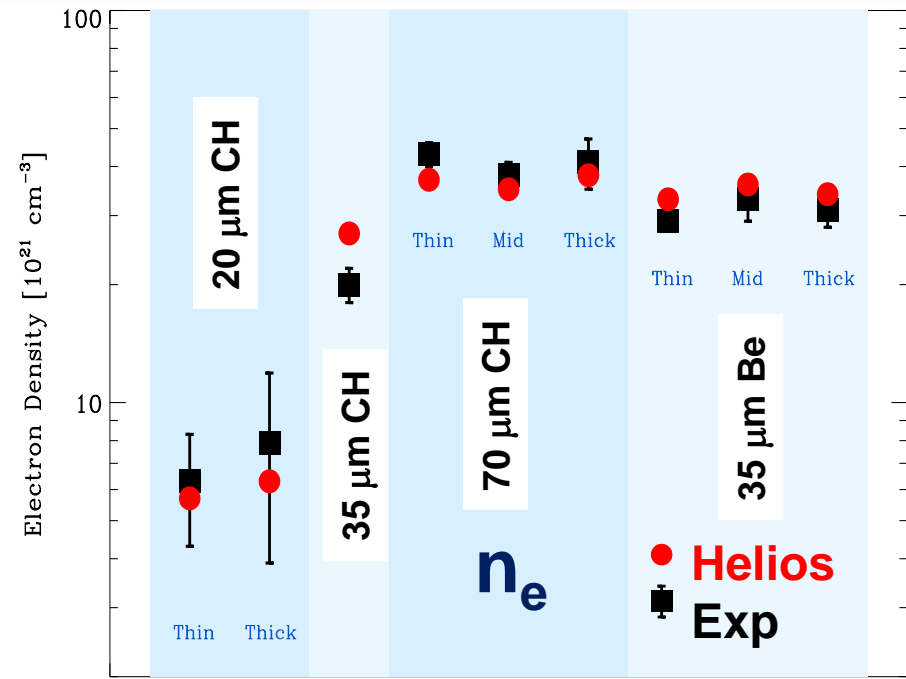
We use our best experimental knowledge of the radiation drive, but this knowledge is imperfect

Therefore a key question is simulation credibility

# Simulations reproduce sample temperature and density for a wide range of experiments



Target configurations

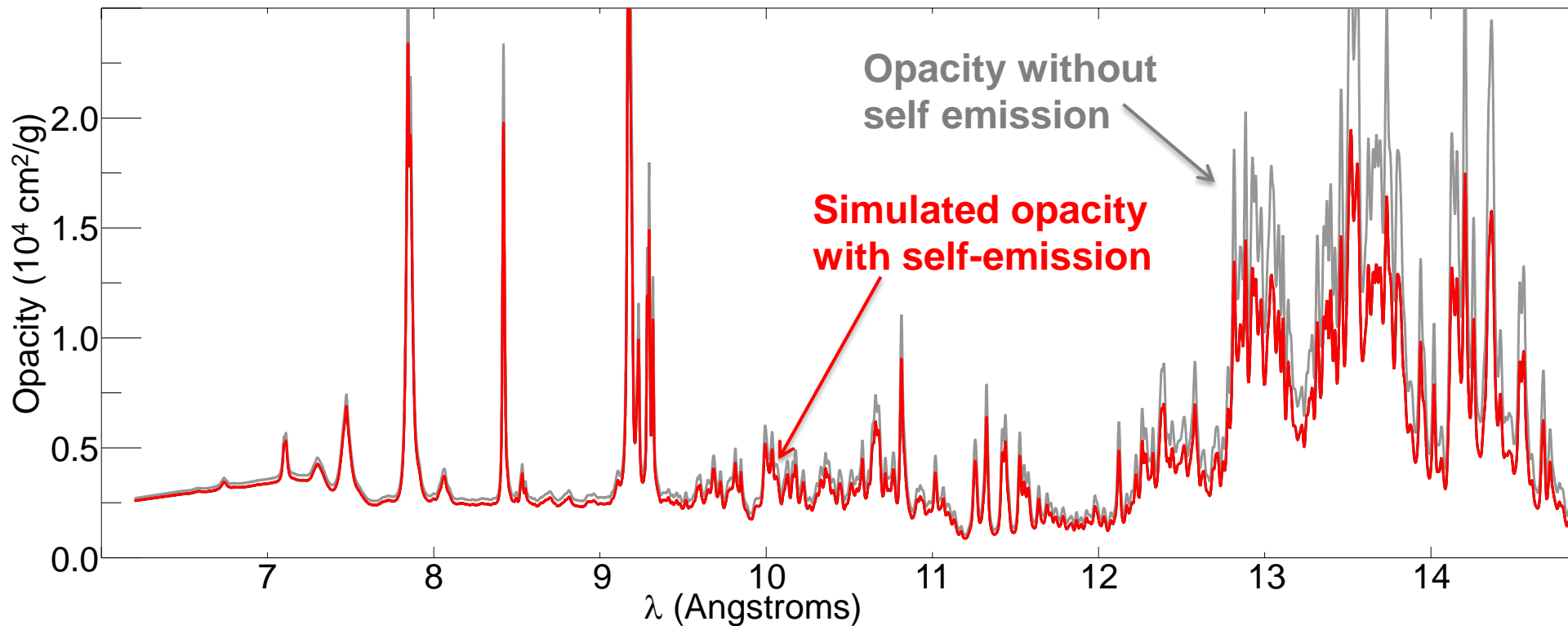


Target configurations

Credibility is supported by the fact that simulations reproduce multiple experiments using the same drive and backlight for all:

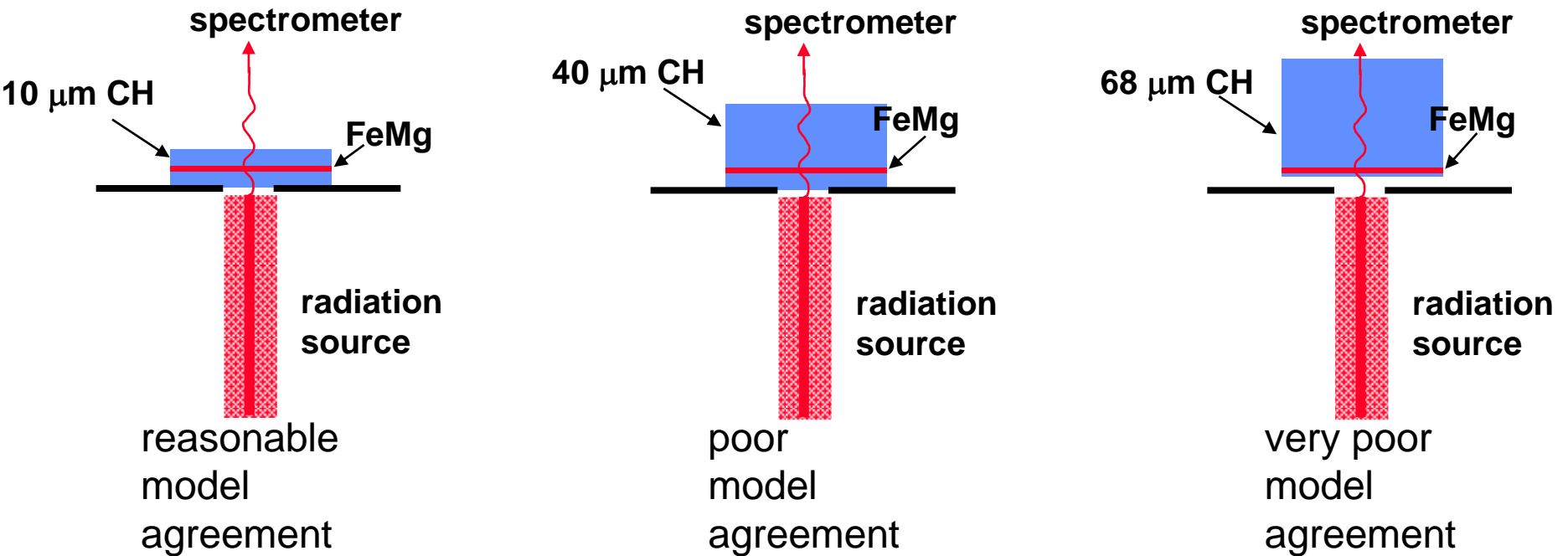
- Tamper thickness changes
- Sample thickness changes
- Tamper material changes

# Self emission influence on opacity is modest for wavelengths below ~12.5 Angstroms



- If present, self emission always reduces the inferred opacity
- Any self emission correction will increase model discrepancies for  $\lambda < 13 \text{ \AA}$
- We observe no self emission, but the quantitative constraint this provides is still under evaluation

# A valid question is whether the rear tamper thickness alters the inferred opacity

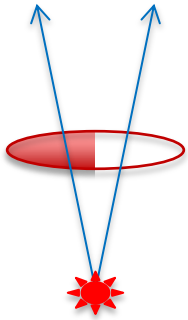


We use the same heating radiation, backlight, diagnostics, sample fabrication for all  
The only difference is the tamper thickness

# Hypothesis: Does the FeMg sample “shadow” the CH tamper behind it enough to enhance inferred attenuation?

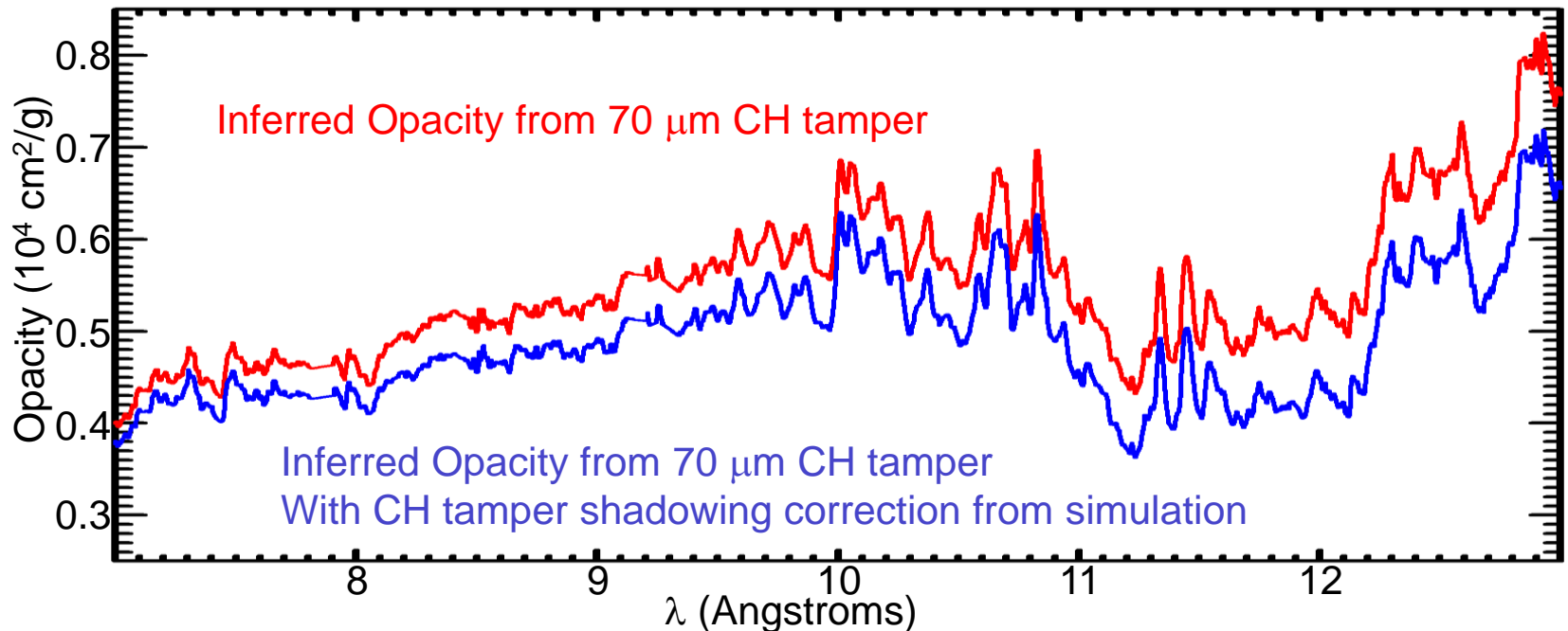
S1 S2

Half-moon sample: Transmission =  $S1/S2$

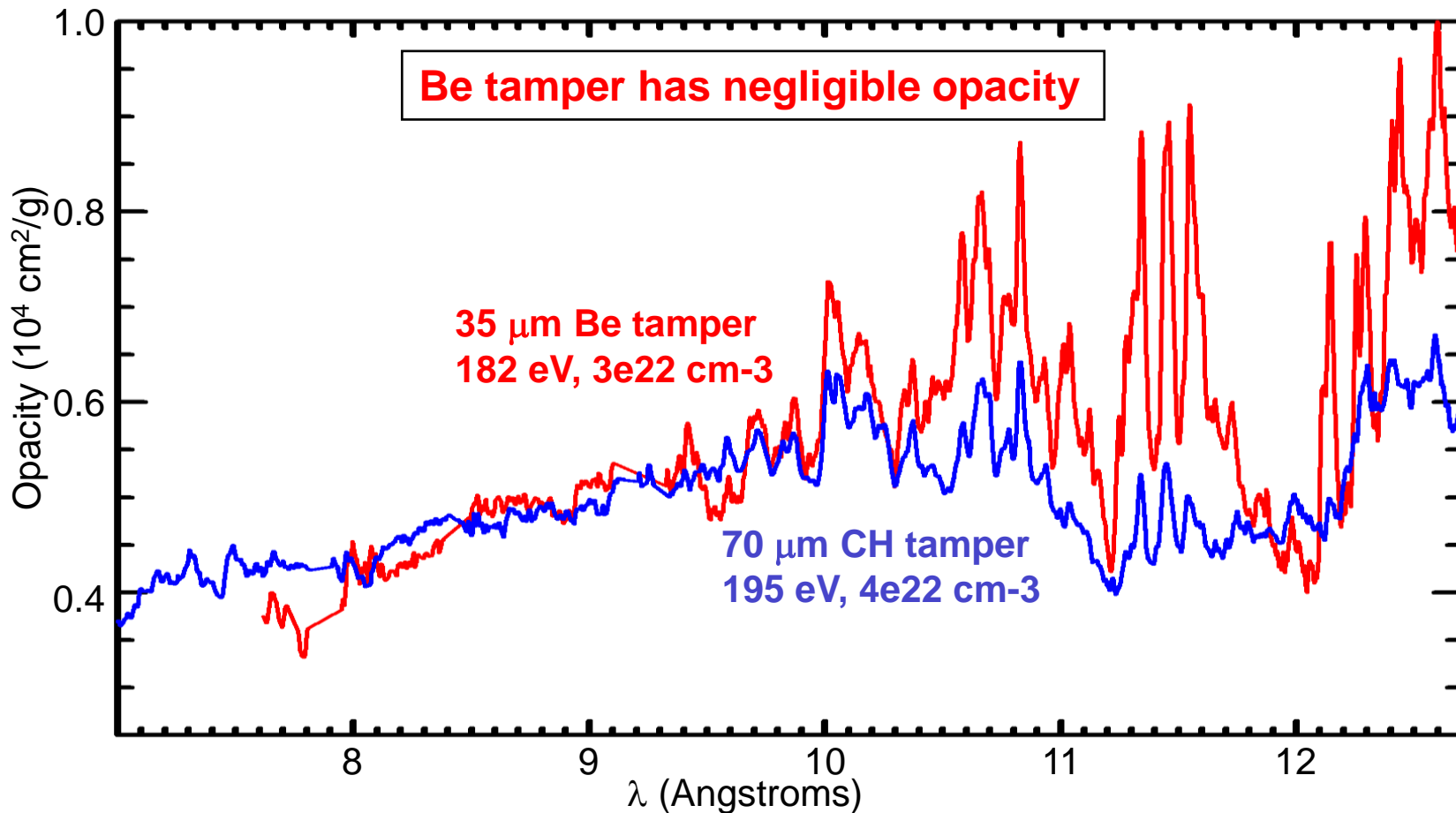


This assumes the CH tamper on the side with the FeMg is heated the same as the side without the FeMg

Simulations suggest the cooler CH behind the FeMg may alter the inferred opacity by up to 20% ; the correction increased with  $\lambda$

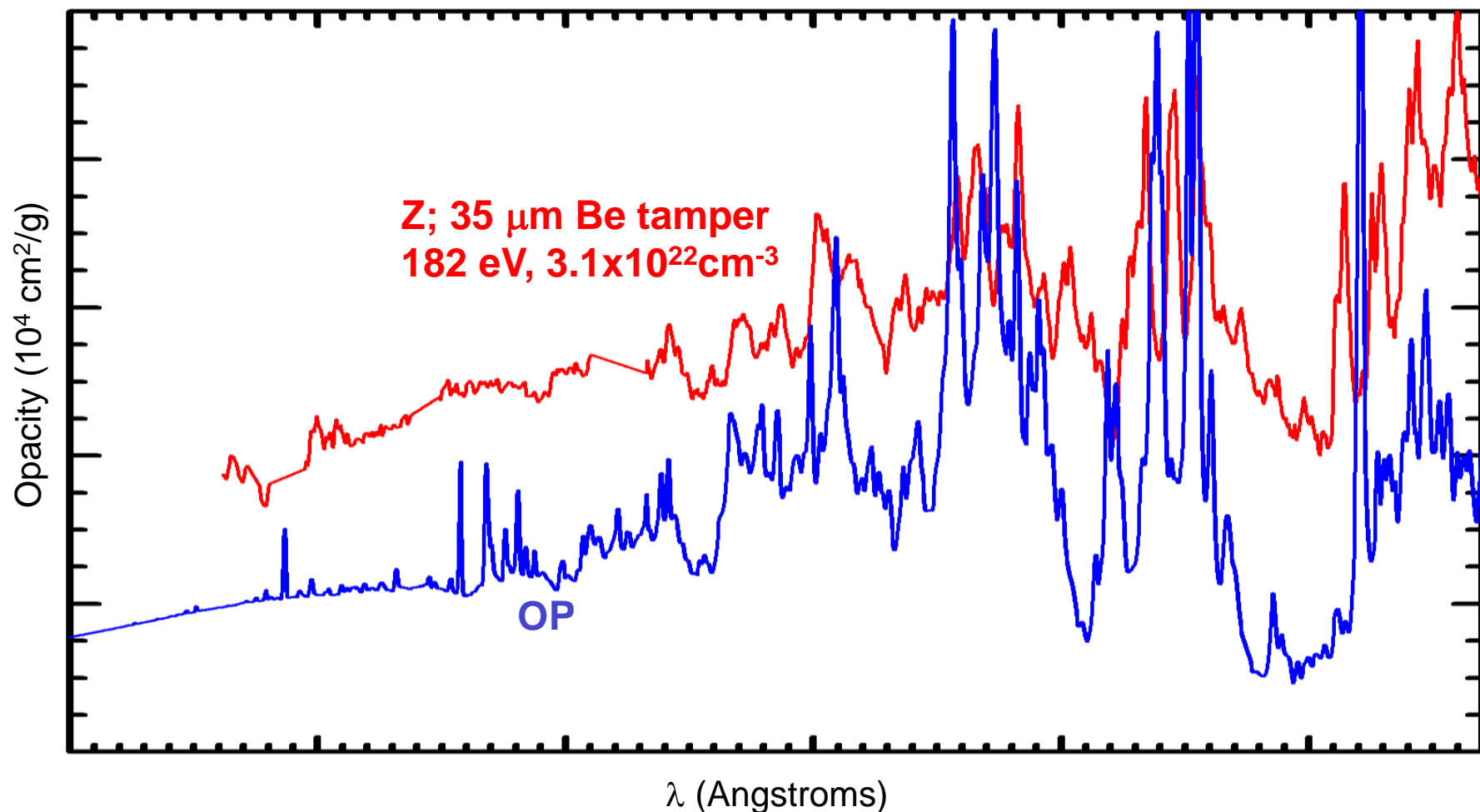


# Be tamper experiments eliminate possible tamper effects on inferred iron transmission / opacity



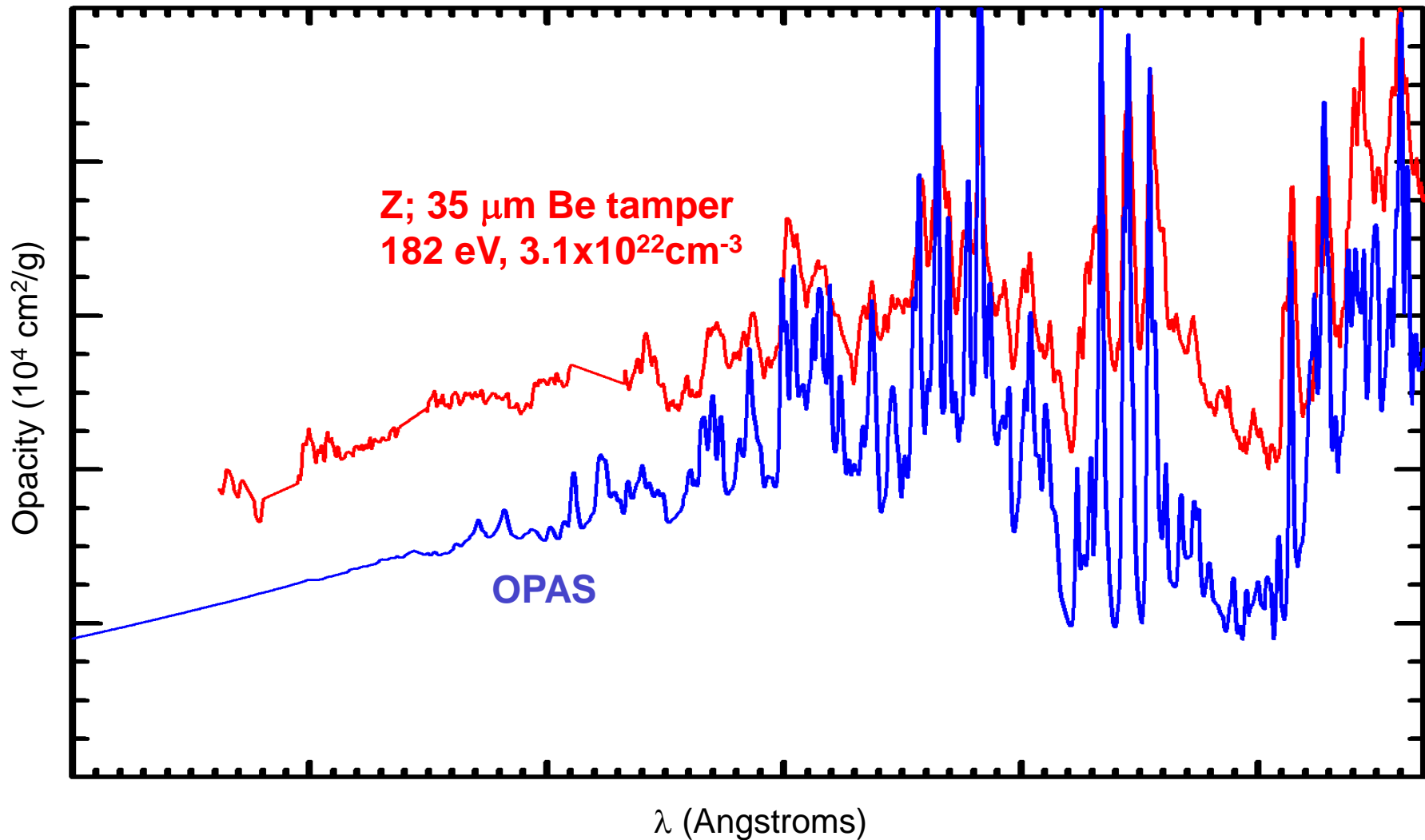
The Be tamper experiments confirm CH-tamped data accuracy at  $\lambda < 10.5$  Angstroms  
AND  
the need for a CH tamper correction at longer wavelengths

# The Be-tamped iron opacity measurement re-affirm the need for significant opacity model revisions

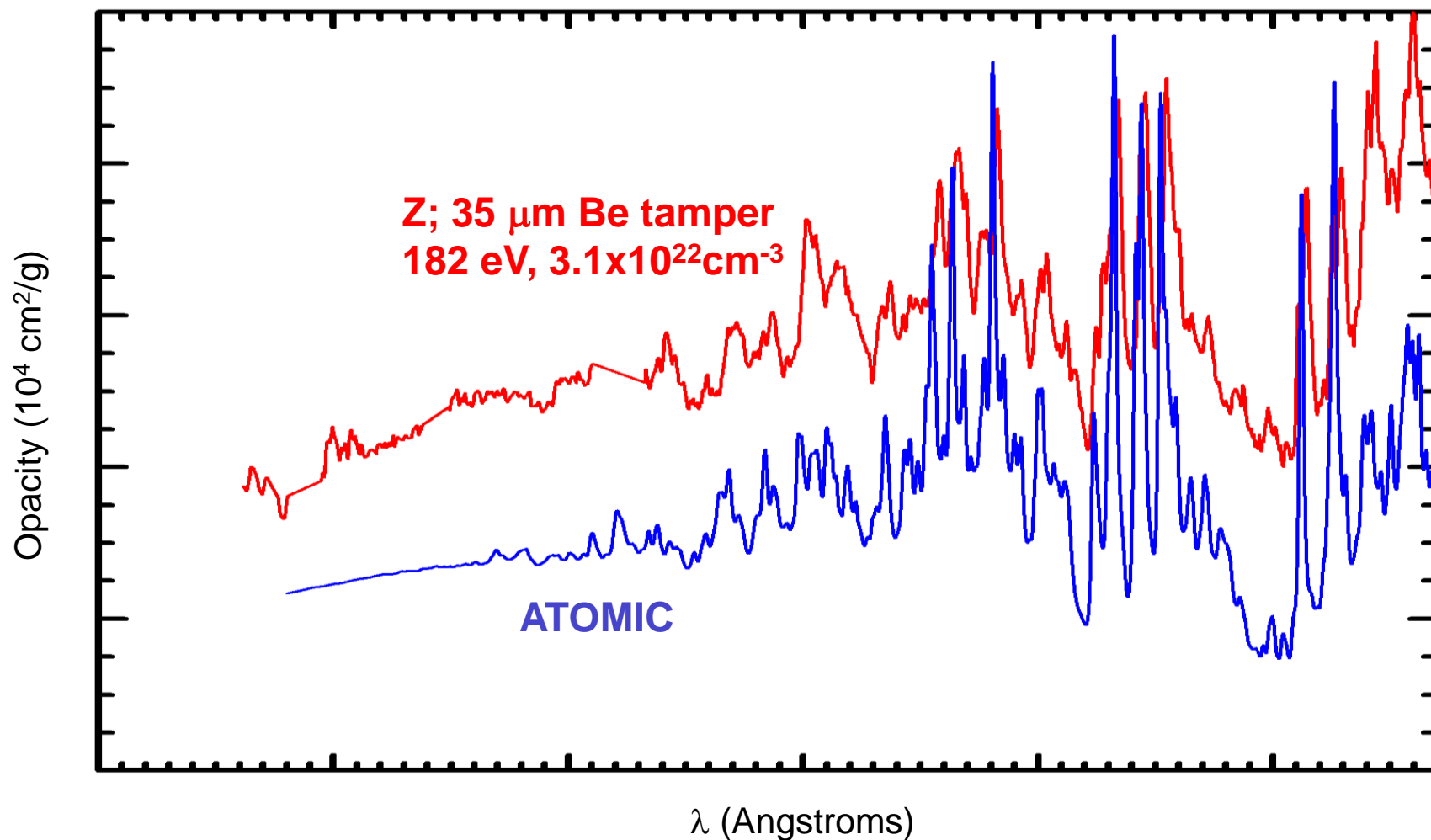




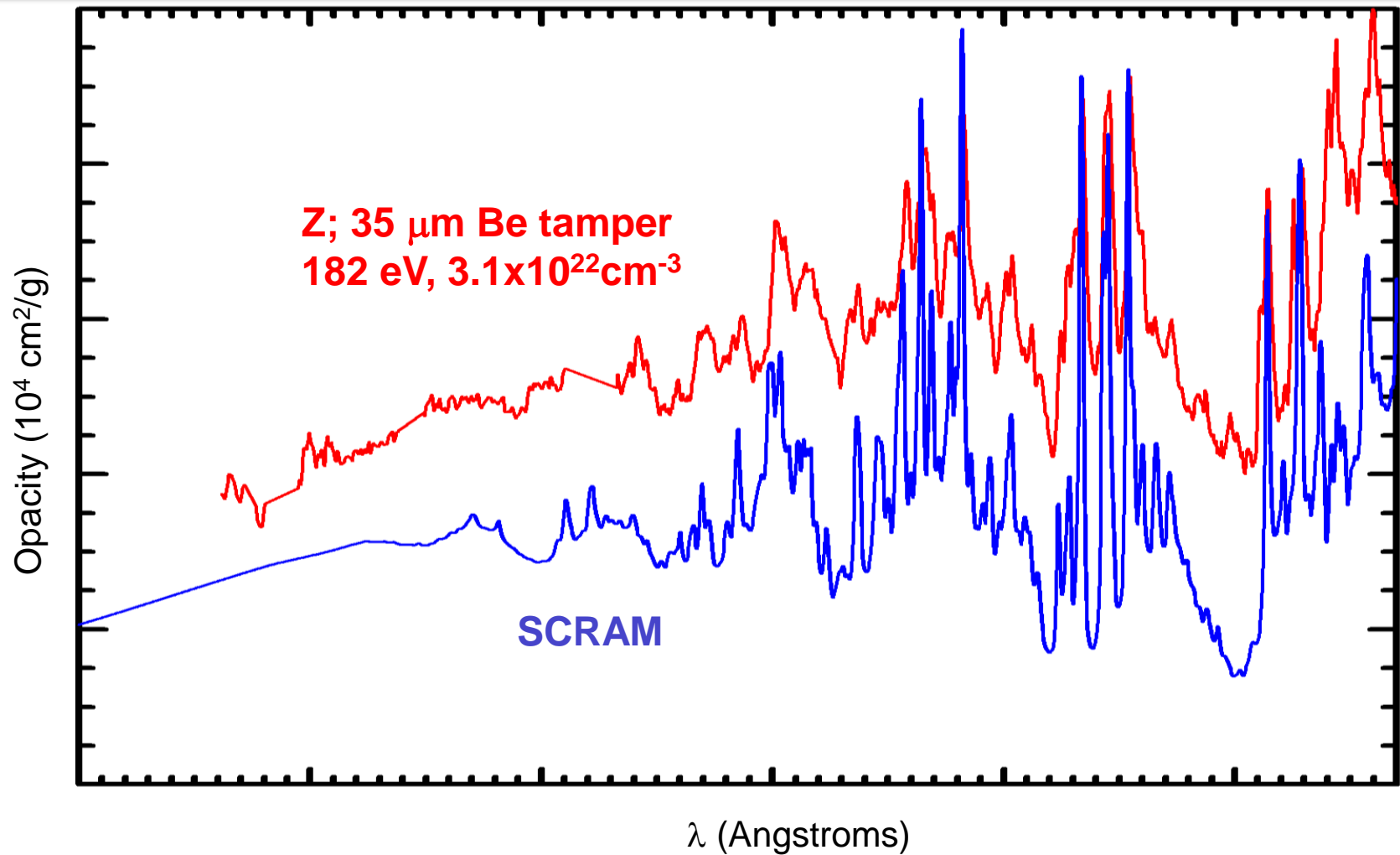
# The Be-tamped iron opacity measurement re-affirm the need for significant opacity model revisions



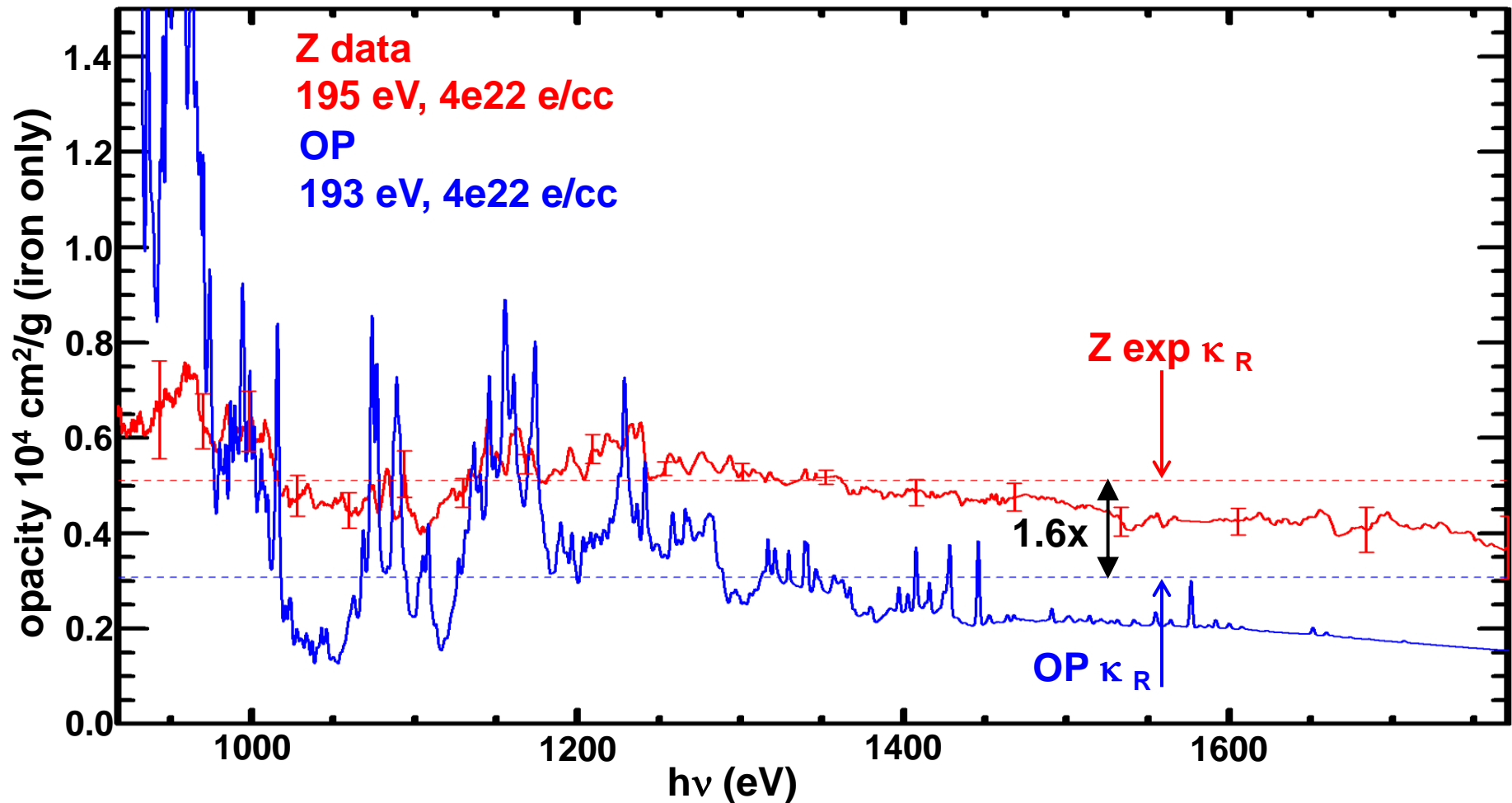
# The Be-tamped iron opacity measurement re-affirm the need for significant opacity model revisions



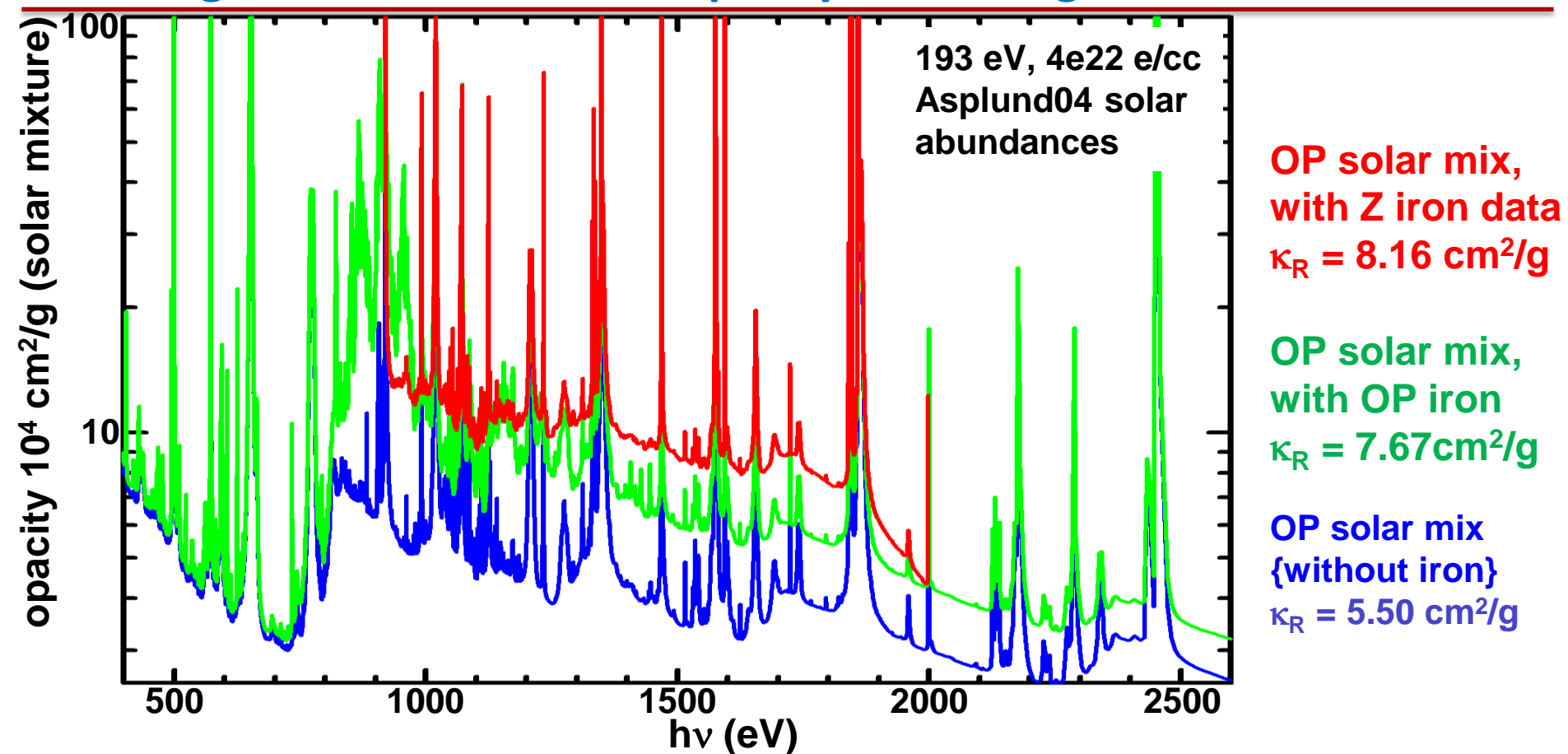
# The Be-tamped iron opacity measurement re-affirm the need for significant opacity model revisions



# The measured pure iron Rosseland mean opacity is ~1.6x higher than calculated with OP



# A solar mixture plasma using Z iron data has ~ 6% higher Rosseland mean opacity than using OP iron



- A 6% Rosseland increase partially resolves the solar problem, but the measured iron opacity by itself cannot account for the entire discrepancy
- Other elements and regions deeper in the sun could contribute
- As we refine the analysis the conclusion could still change

# Future opacity experiments will examine opacity model physics issues

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2-3 transition arrays - Ni

Bound free quasi continuum – Cr

Line broadening

Higher densities – thicker Be tamper

Plasma composition – Fe mixed with low Z

# What are the hypotheses for the discrepancy and how can we test them?

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## Hypotheses:

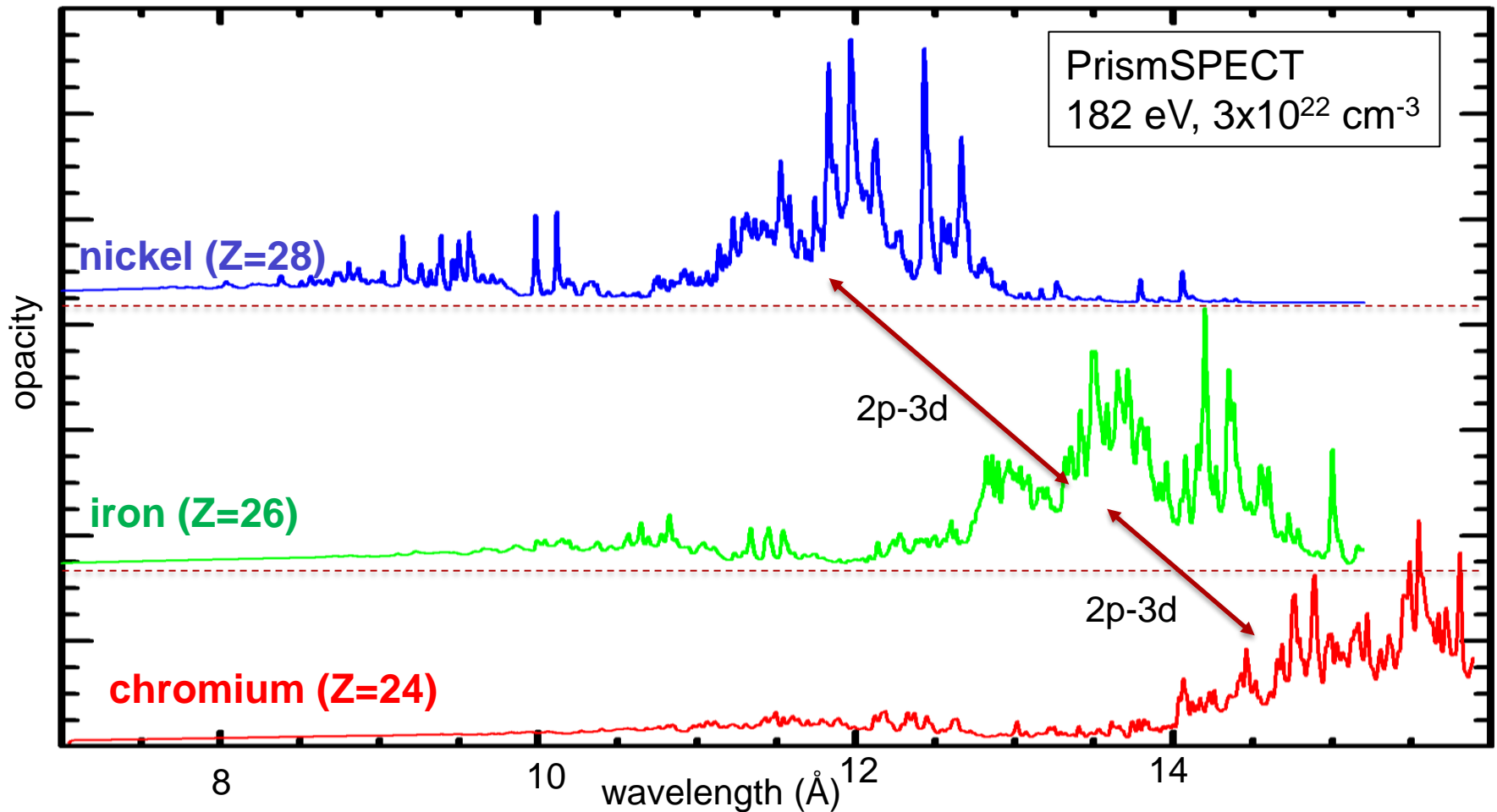
- 1) Despite all our effort, iron measurement is flawed somehow
- 2) Photon absorption is shifted from long  $\lambda$  to short  $\lambda$  by a process that is as yet undetermined
- 3) Models have difficulty predicting opacity for open L-shell configurations
- 4) Models have difficulty predicting highly excited configurations

## Tests:

- A) Z experiments with lower and higher atomic number elements
- B) Experiments on a different platform (NIF)

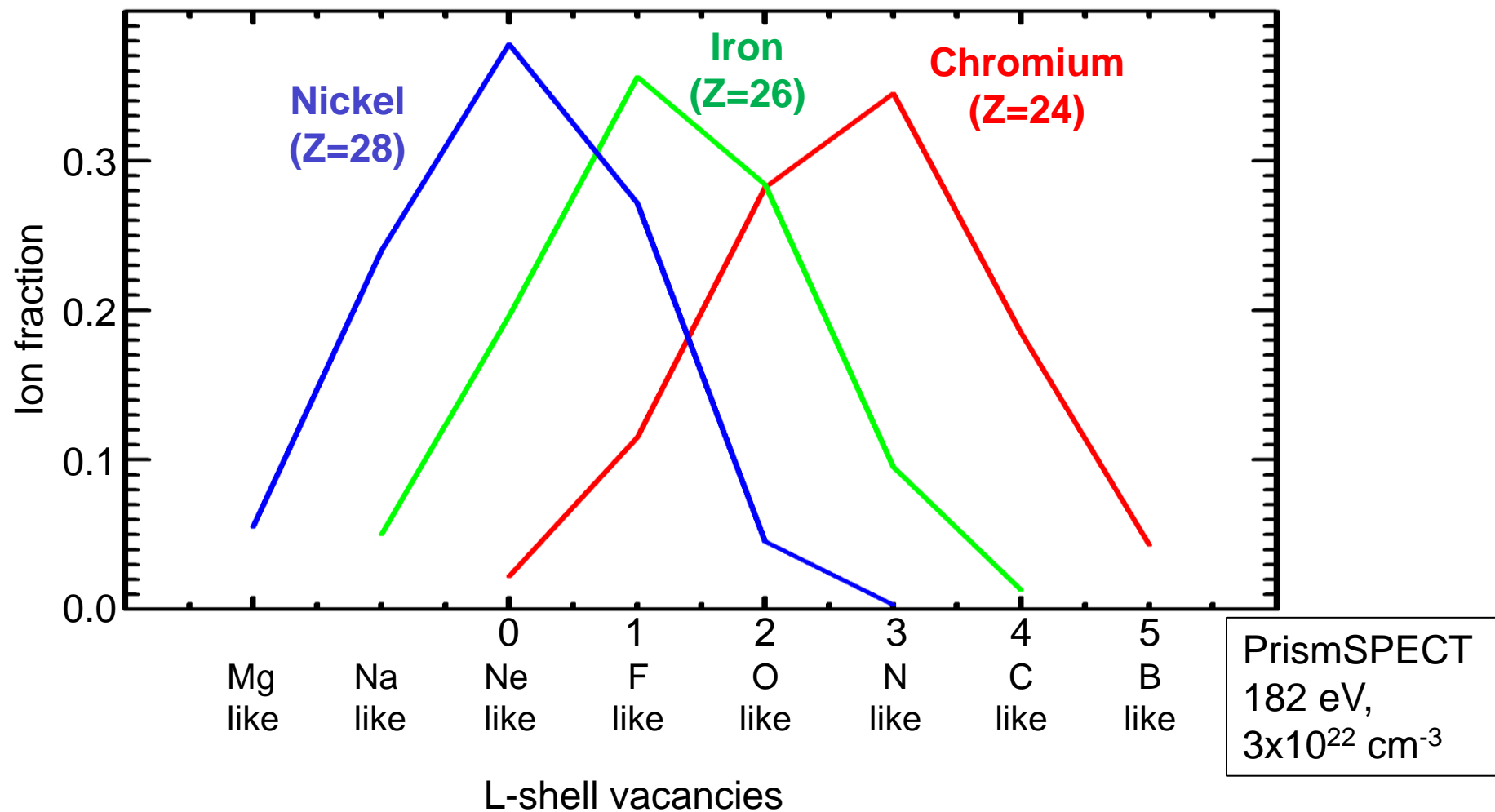


# Experiments with different elements shift different spectral regions into the highest accuracy experiment range



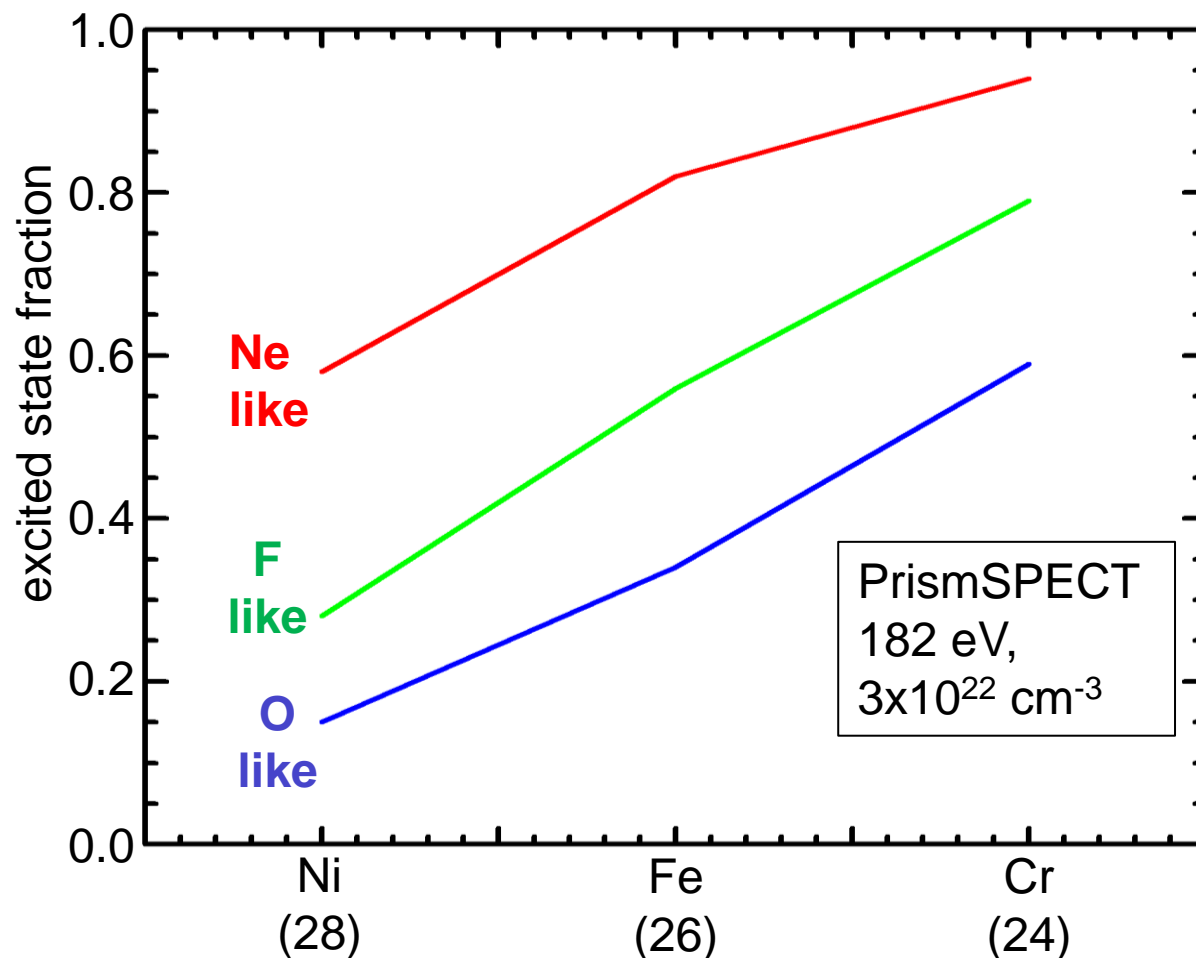
Experiments with different elements also can help identify possible experiment peculiarities with the iron measurements (e.g., unknown contaminants)

# The number of L shell vacancies changes with the sample element



Opacity from transitions with an open L-shell may be more complex to model

# The fractional excited state population increases as the atomic number decreases



Opacity from ions with high excited state populations may be more complex to model  
These difficulties increase as atomic number decreases

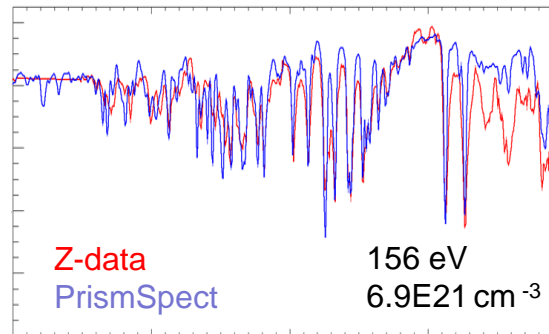
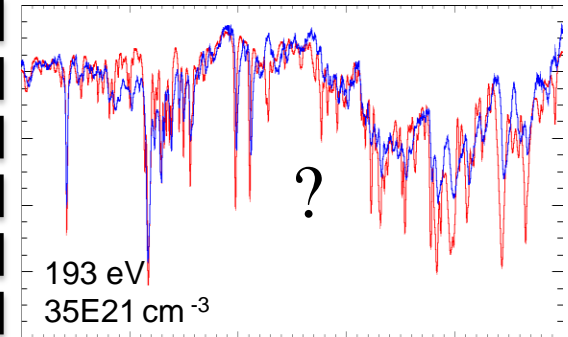
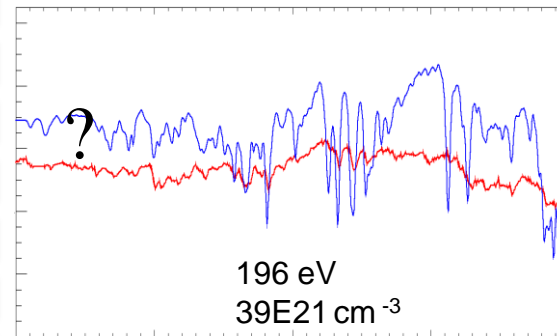
# We will untangle the complex opacity issues through precise measurements across a range of $T_e$ , $n_e$ , and $Z$ .

fewer L-shell vacancies, lower excited state populations

Chromium (open L-shell)

Nickel (closed L-shell)

Iron (open L-shell)

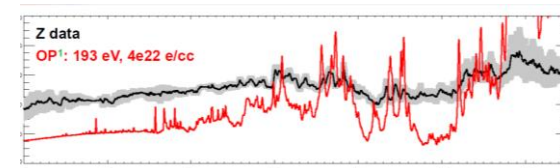
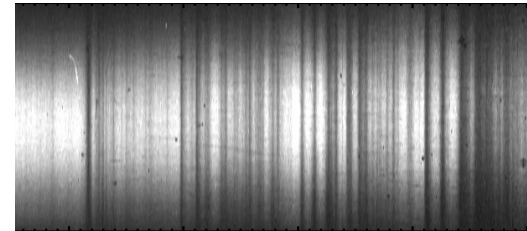
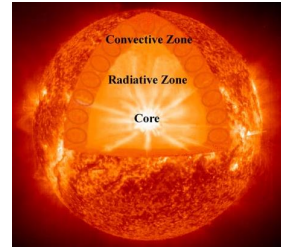


Increased Temp. and Density ↑

→ Increased Atomic Number

# Z iron plasma opacity experiments refine our understanding of solar interior structure

- Solar interior predictions don't match helioseismology
  - Arbitrary opacity increase of 10-30% would fix the problem, but is this the correct explanation?
- Z experiments have measured iron plasma opacity at nearly solar convection zone base conditions
  - Experiment temperature is the same as in sun, density within a factor of 2
- Opacity models disagree with measurements at near-solar-interior conditions



The Rosseland mean opacity for solar matter is ~ 6 % higher using Z iron data instead of OP model calculations