



Dielectric Resonator Metasurfaces: Optical Magnetism, Emission and Optical Devices

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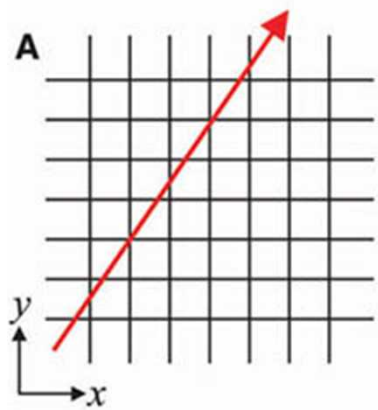


Outline

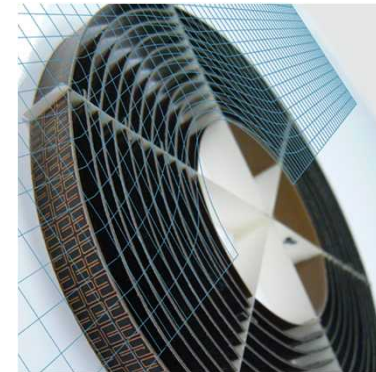
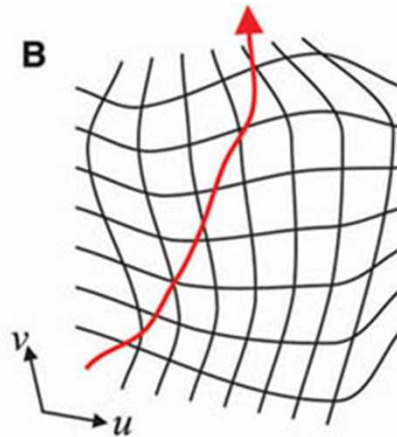
- **Dielectric Resonators as Metamaterials**
- **Optical Magnetism**
- **Directional Emission, Fano Resonances, and Third-Harmonic Generation**

The Promise of 3D Metamaterials

- Engineer ϵ and μ everywhere in space using deep subwavelength structures (usually resonators)
- Then light will do wonderful things

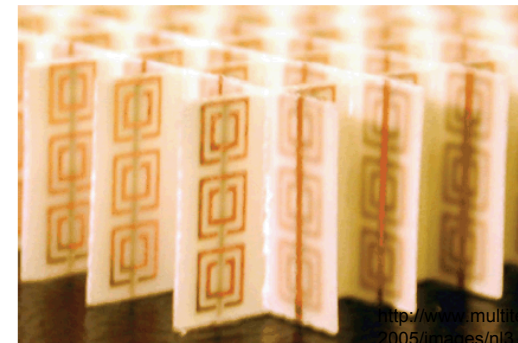


Sir John Pendry



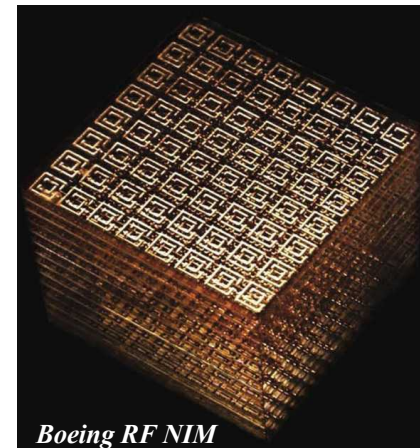
David Smith, Duke

- For RF ($\sim 3\text{GHz}$), $\lambda \sim \text{cm}$
- We need to create sub-wavelength “inductors” and “capacitors” $\sim \text{XXmm}$

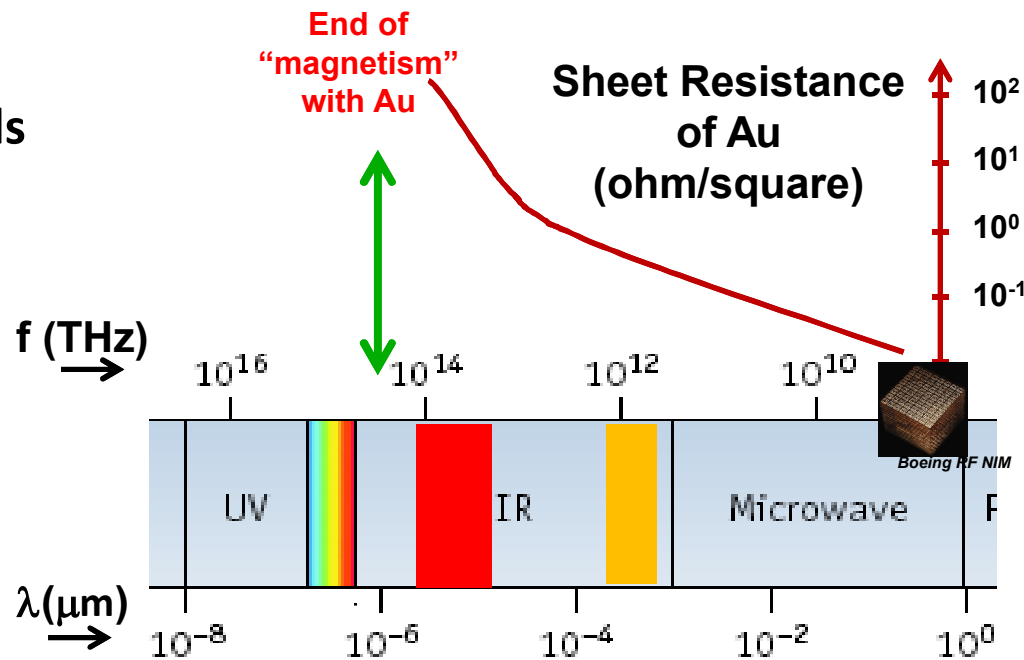


Metamaterials at Optical Frequencies?

Issue 1: Fabrication



Issue 2: Loss of metals

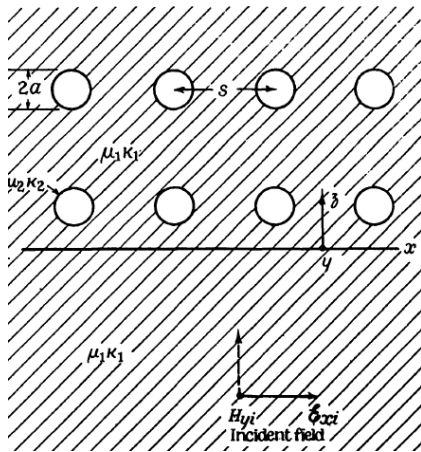


Dielectric Resonators for Metamaterials

THE ELECTRICAL CONSTANTS OF A MATERIAL LOADED WITH SPHERICAL PARTICLES*

By L. LEWIN.†

(The paper was first received 4th March, and in revised form 27th September, 1946.)



$$\mu'_{re} = \mu_{r1} \left(1 + \frac{3v_f}{\frac{F(\theta) + 2b_m}{F(\theta) - b_m} - v_f} \right)$$

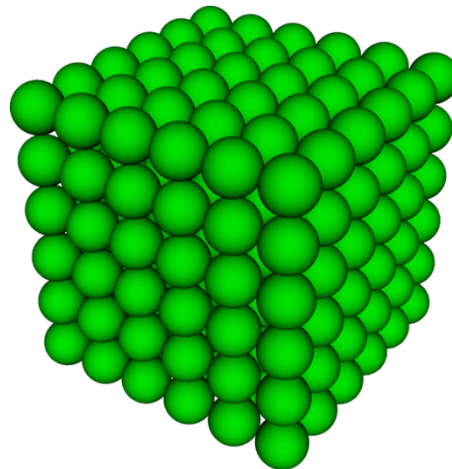
$$\epsilon'_{re} = \epsilon_{r1} \left(1 + \frac{3v_f}{\frac{F(\theta) + 2b_e}{F(\theta) - b_e} - v_f} \right)$$

$$F(\theta) = \frac{2(\sin \theta - \theta \cos \theta)}{(\theta^2 - 1) \sin \theta + \theta \cos \theta}$$

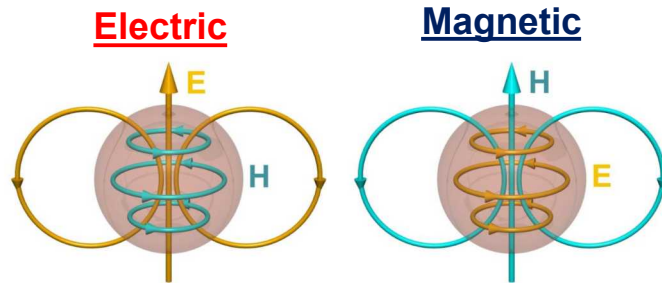
$$\theta = k_0 a \sqrt{\epsilon'_{r2} \mu'_{r2}}$$

$$b_e = \frac{\epsilon_1}{\epsilon_2},$$

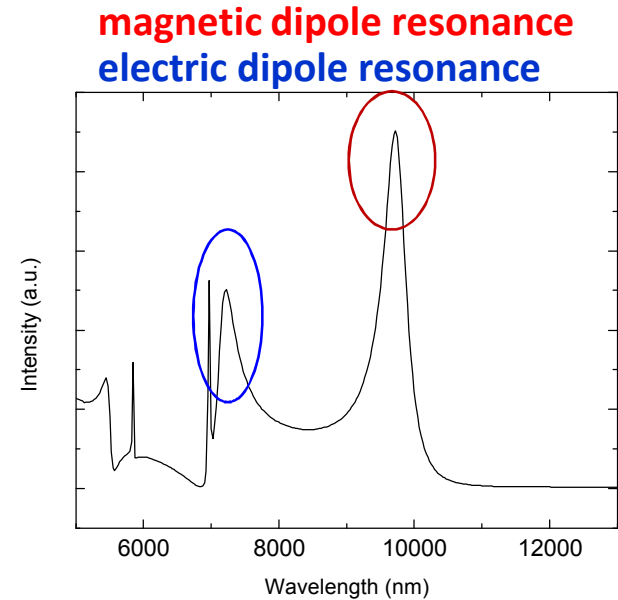
$$b_m = \frac{\mu_1}{\mu_2}$$



Dielectric Resonators

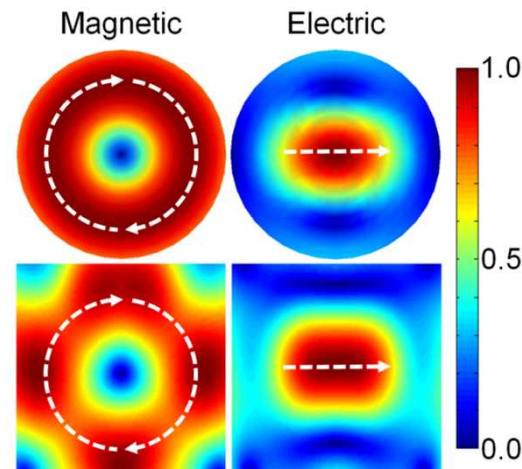


Images: A. Miroshnichenko



Magnetic dipole resonance: tailor μ
Electric dipole resonance: tailor ϵ

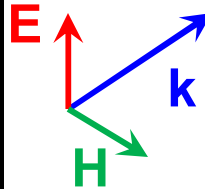
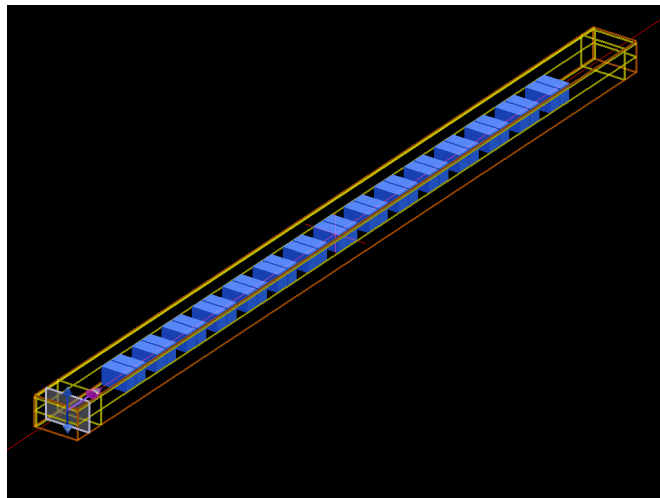
- Cubes work fine too
- Introducing “cuts” in cube can move relative positions of resonances





Full Wave Simulation of Propagation Through Split Cube Array

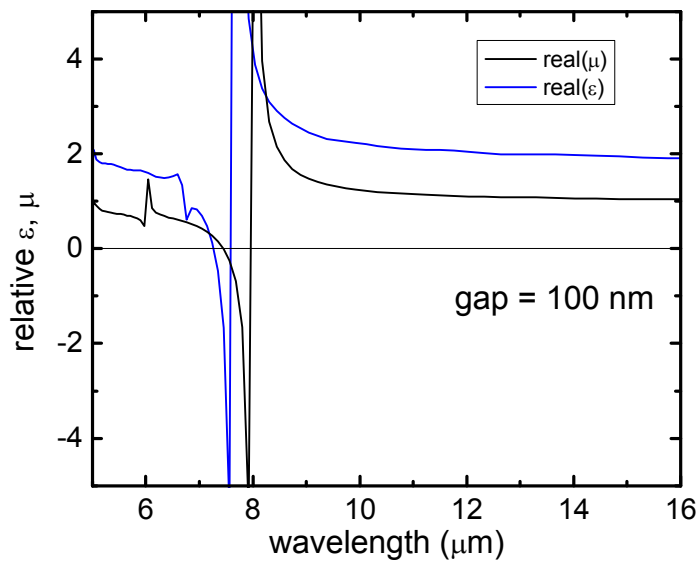
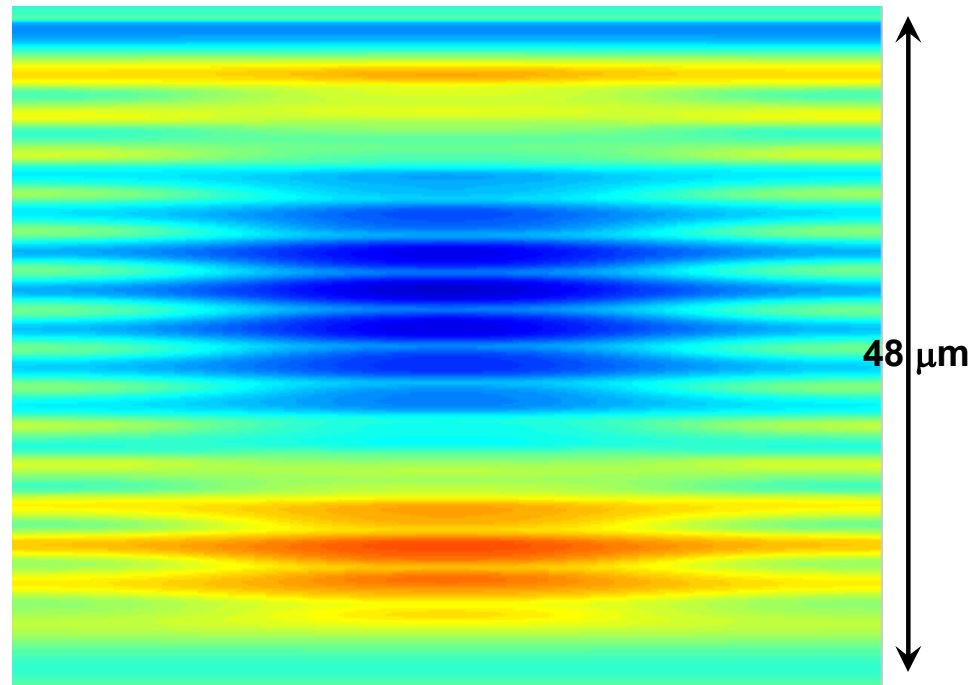
$\epsilon=32$, edge = $1.53 \mu\text{m}$, gap = 100 nm



H_x at top of unit cell ($z=1.3 \mu\text{m}$)
Incident waves



$2.6 \mu\text{m}$



What is a mirror?



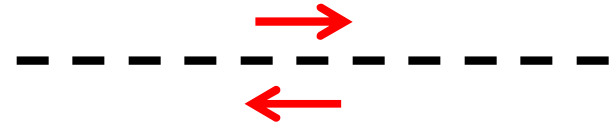
“Regular” mirrors invert the phase of the reflected wave

Dipoles Close to Surfaces

Electric dipole on top of a perfect **electric** conductor (PEC)



IMAGE THEORY

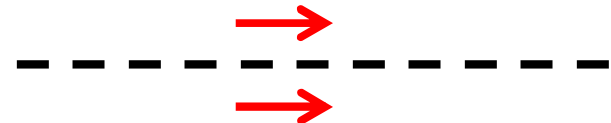


- Because of boundary condition of the PEC surface, the electric field at the dashed plane has to be zero
- This means that the radiation of an electric dipole close to PEC is quenched

Electric dipole on top of a perfect **magnetic** conductor (PMC)



IMAGE THEORY

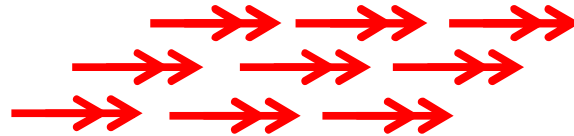


- This means that the radiation of an electric dipole on a PMC is enhanced

Magnetic Dipoles

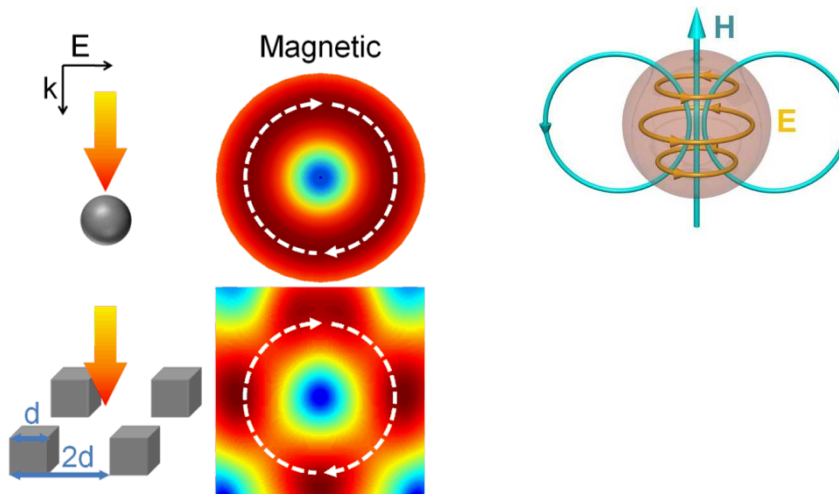
However, a perfect magnetic conductor does not exist in nature

Array of magnetic dipoles



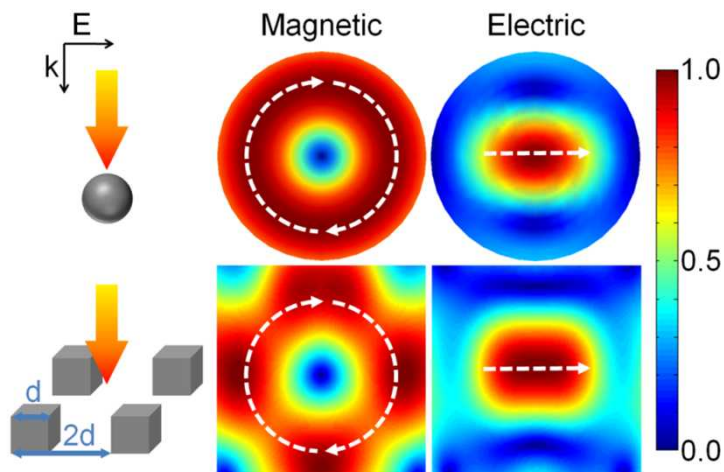
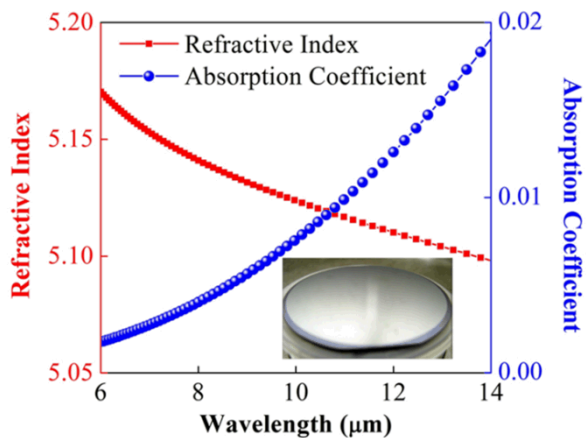
Because the magnetic dipole responds in phase with the electric field, this represents an **artificial magnetic conductor**

Dielectric Spheres

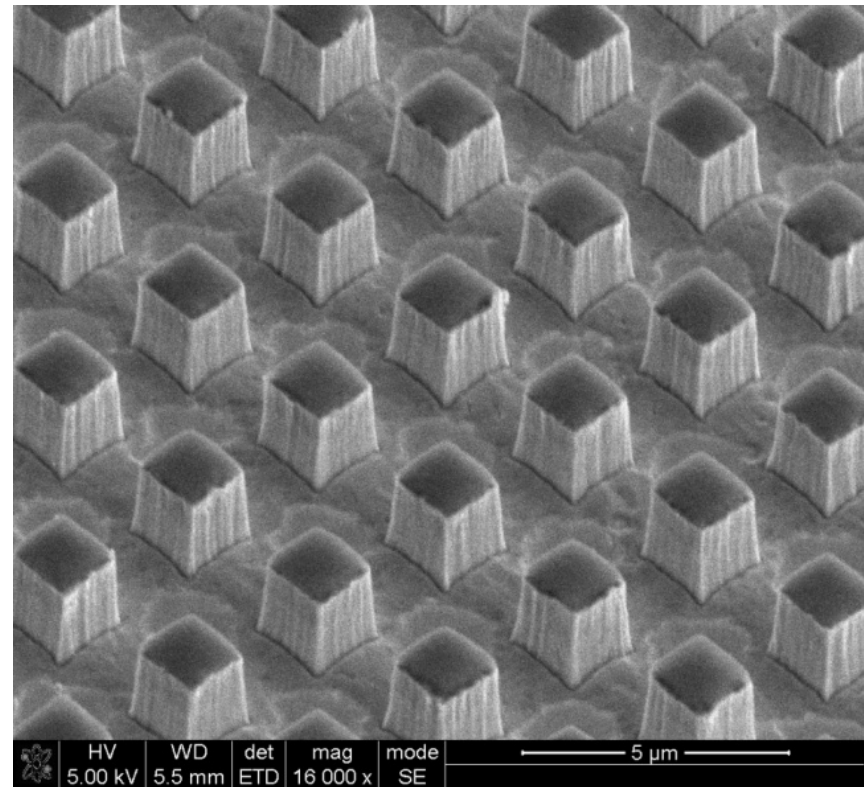




Dielectric Resonator IR Metasurface: Te/BaF_2



Mask + RIE etch

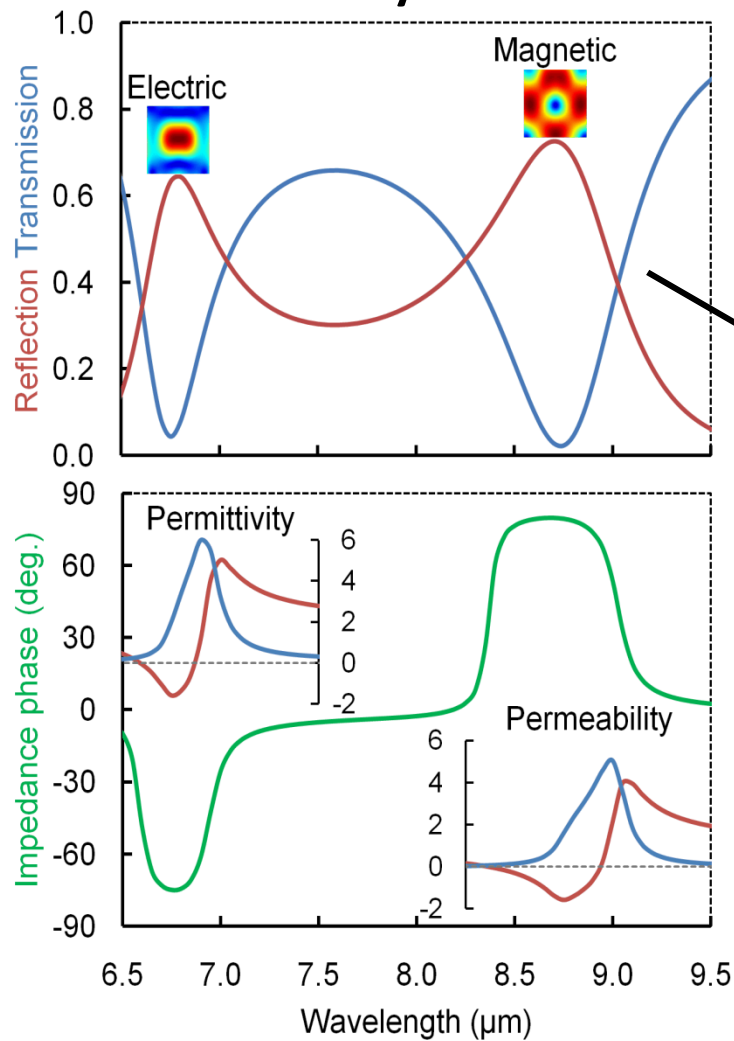


$1.53 \times 1.53 \times 1.7 \text{ mm}^3$
10 deg wall slope

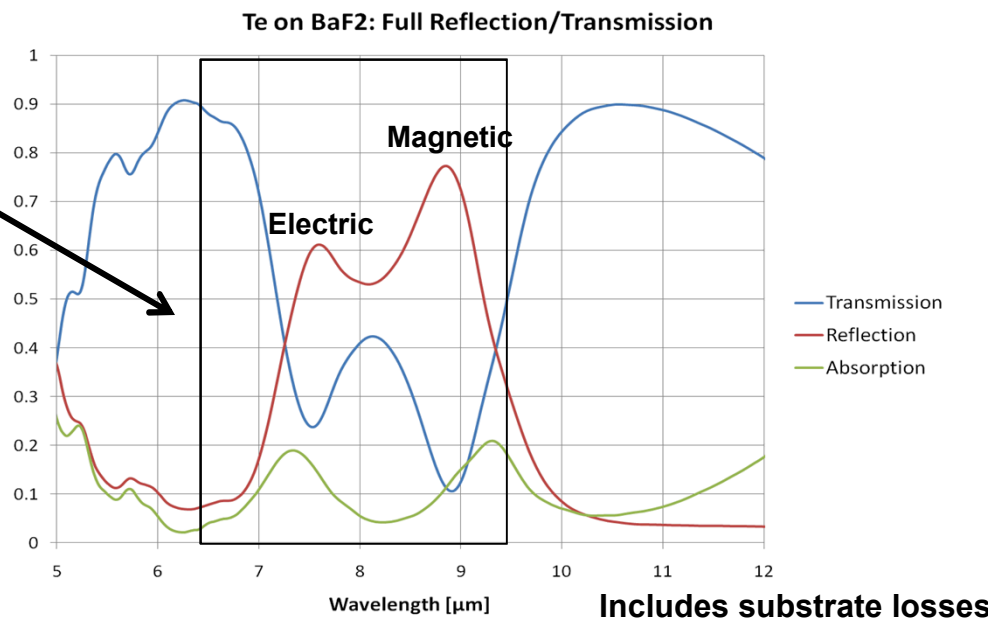


Dielectric Resonator IR Metasurface: Te/BaF_2

Theory

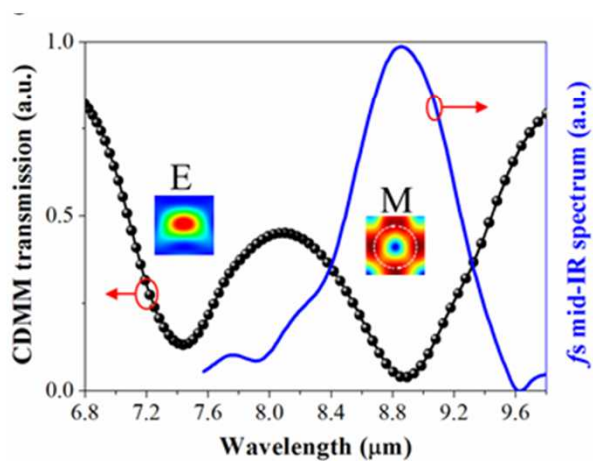
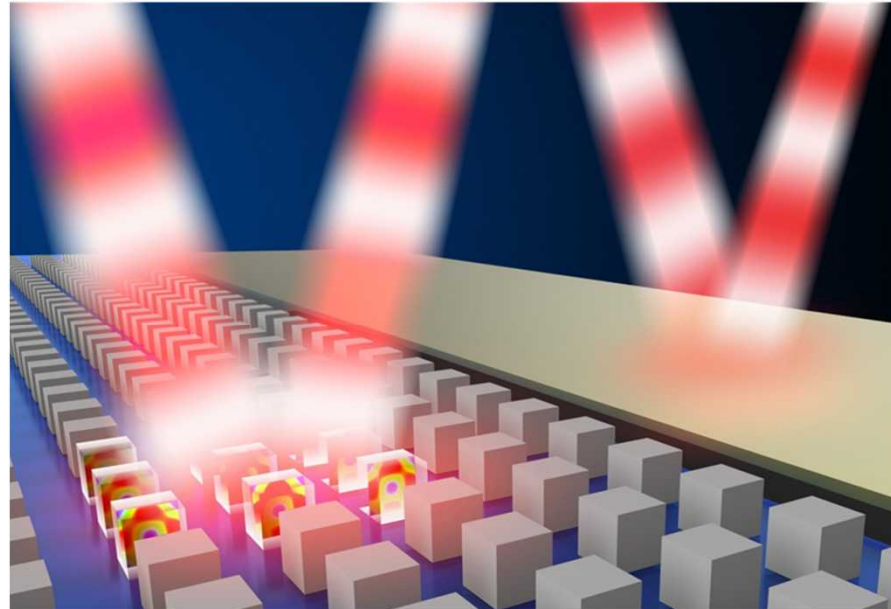


Experiments

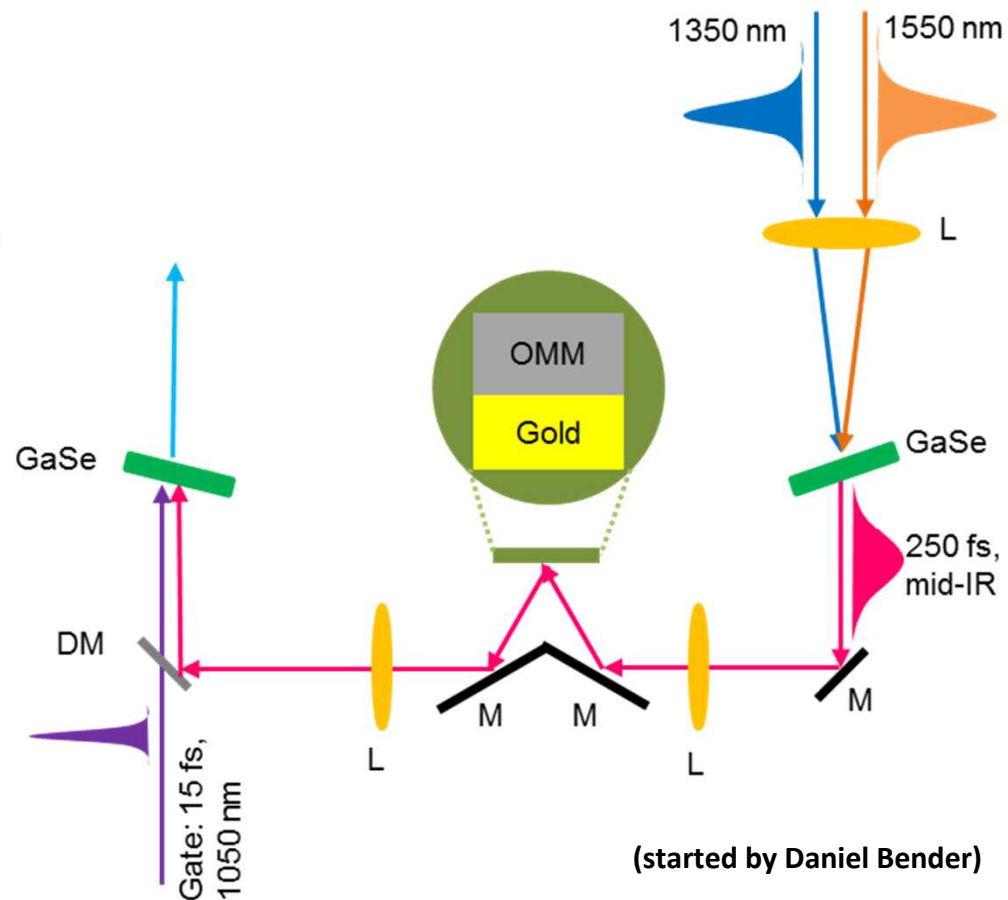


Proving Optical Magnetism: Measure Absolute Phase of Reflected Wave

a

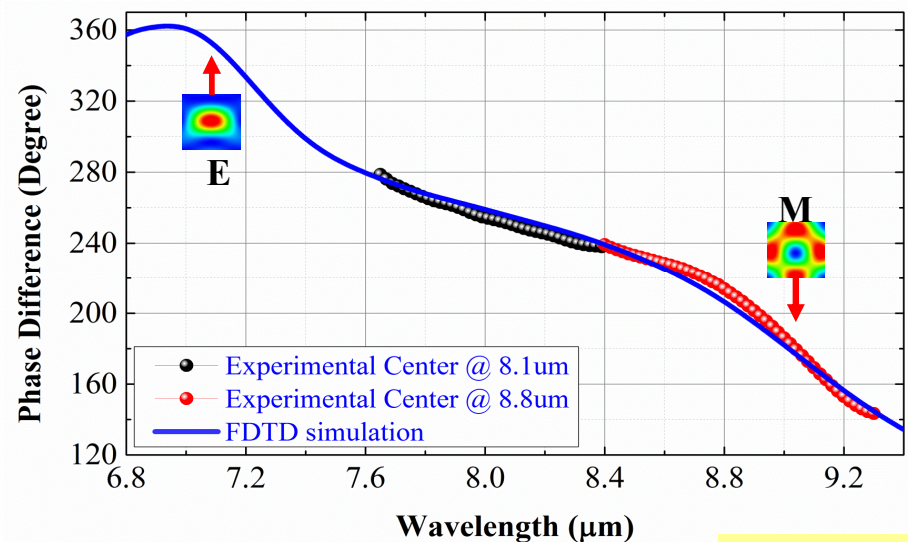
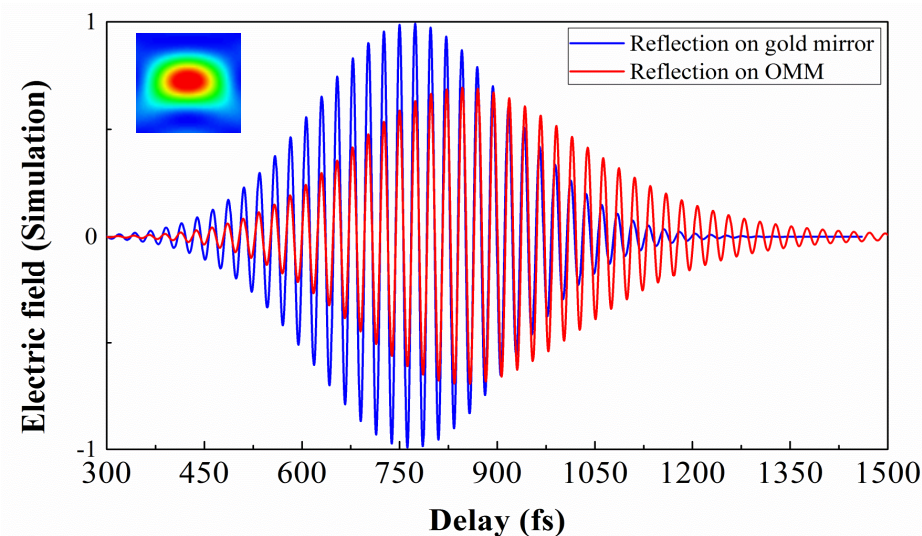
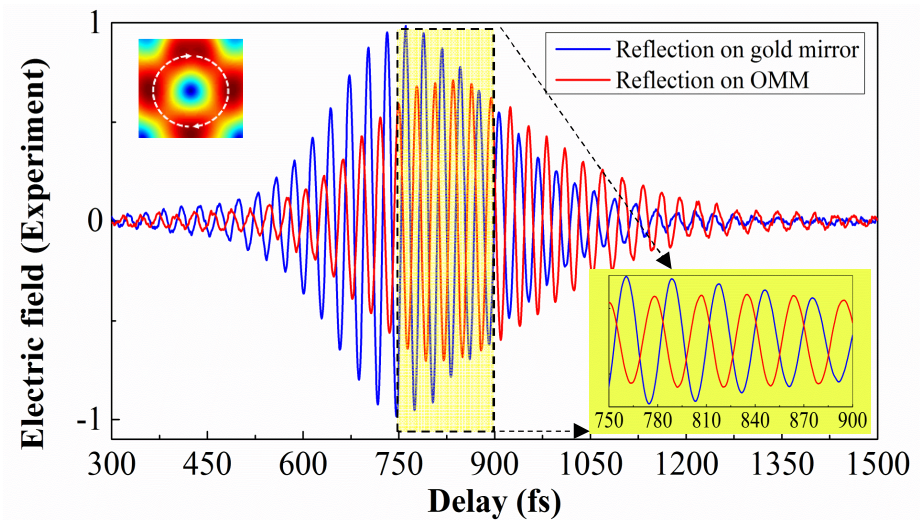
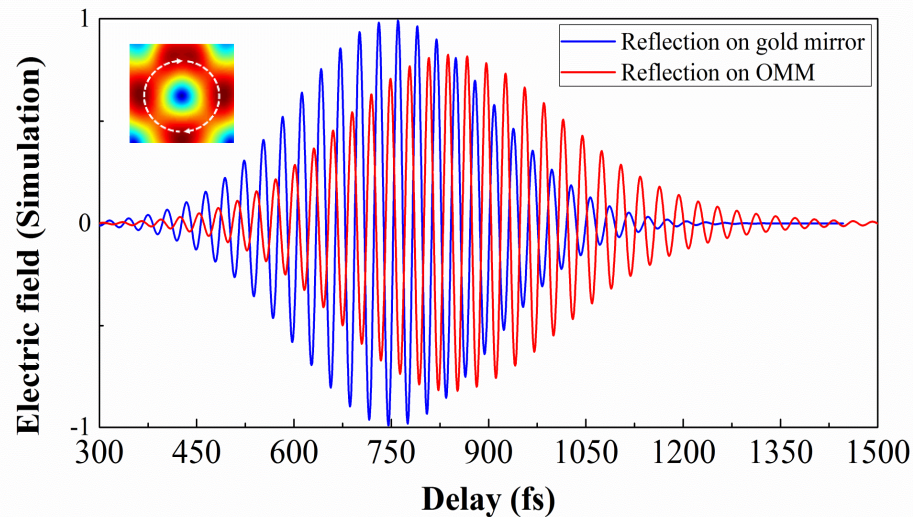


Phase-locked Time Domain Spectroscopy

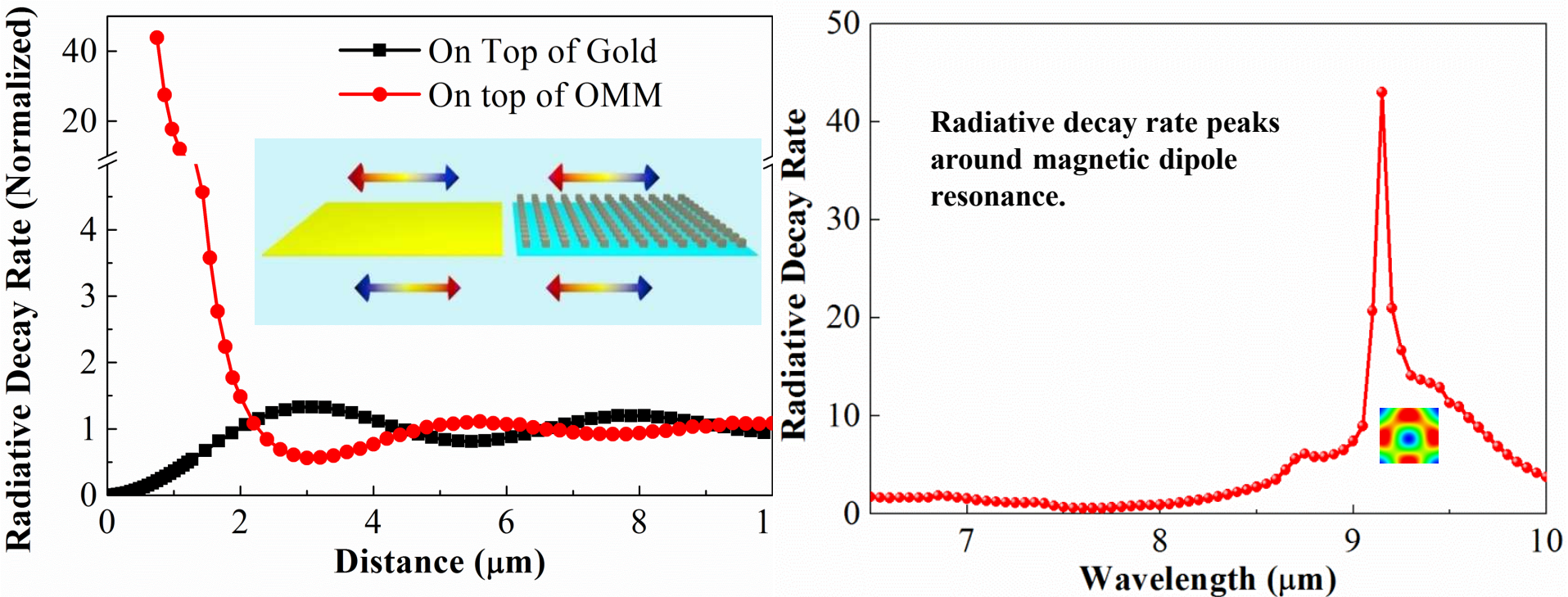




Experimental Demonstration of “Optical Magnetism”

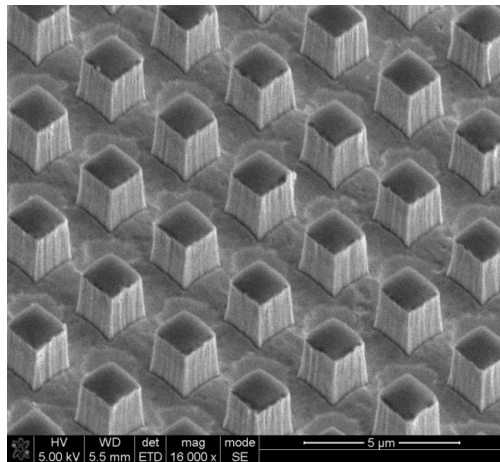


Radiative Decay Rate of a Transverse Electric Dipole Near Au and OMM Surfaces

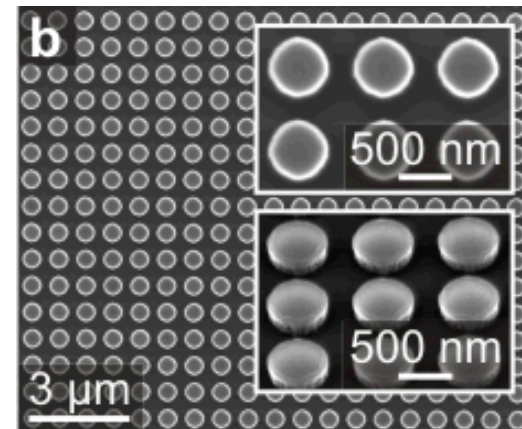


- Oscillatory dependence on distance is shifted by about half a period
- Dipole emission near the magnetic mirror is enhanced even for very small distances

Scaling Dielectric Resonators to the Near IR: Silicon Nanostructures

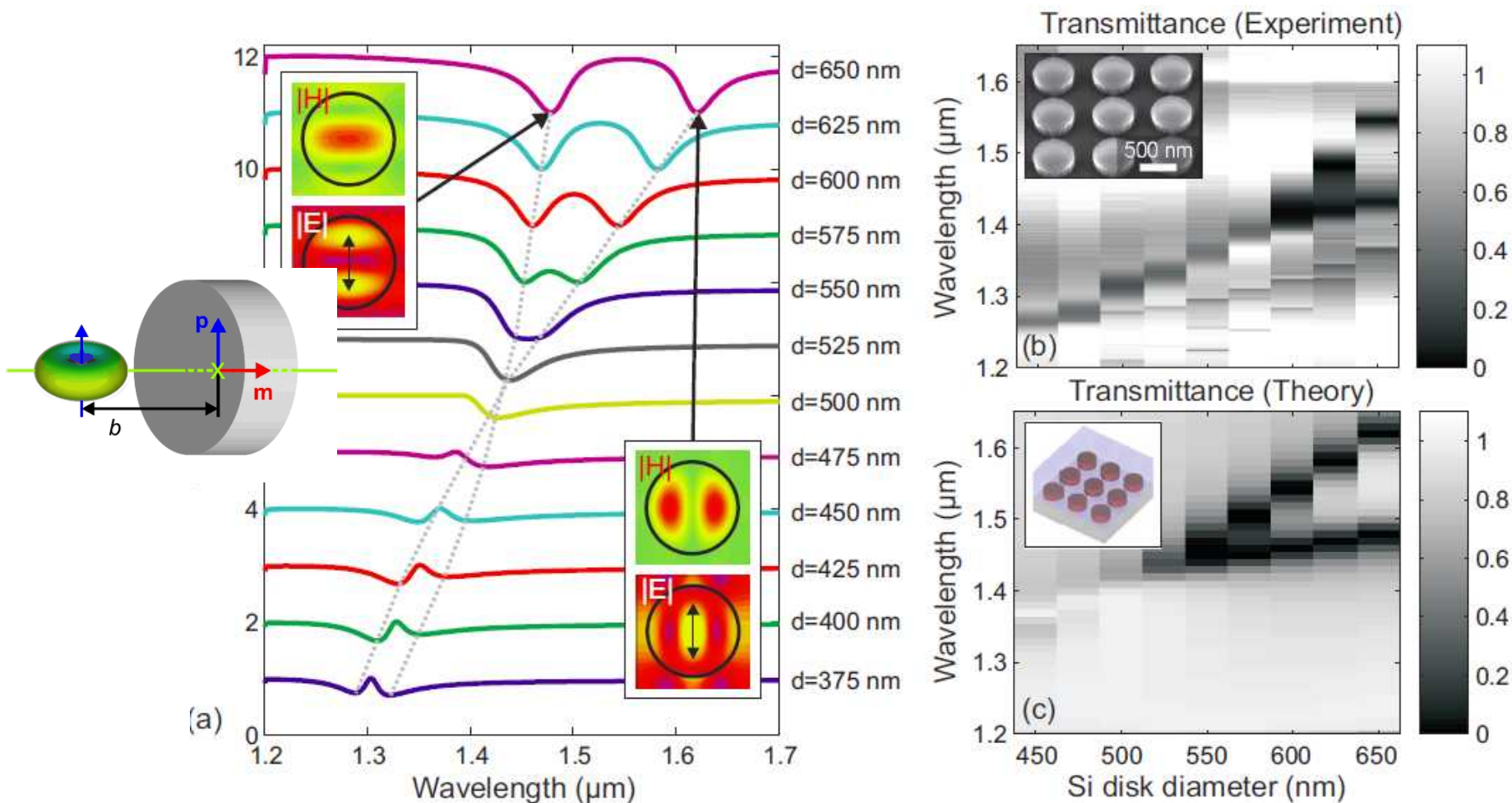


Tellurium: $n \sim 5$, Size $\sim 1.5 \mu\text{m}$
 $\lambda > 5 \mu\text{m}$



Silicon: $n \sim 3.5$, Size $< 200\text{-}500\text{nm}$
 $\lambda > 1 \mu\text{m}$

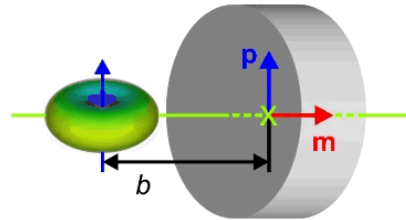
Scaling Dielectric Resonators to the Near IR: Silicon Cylinders (with I. Staude & Y. Kivshar, ANU)



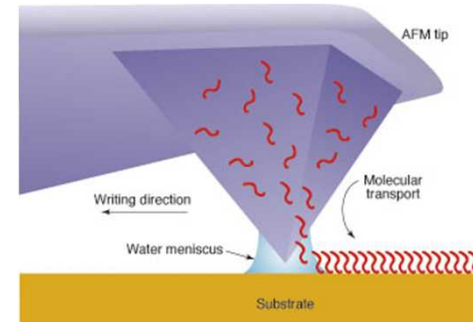
Changing the aspect ratio of the nanocylinder, changes relative position of E & H resonances

Integration with Quantum Dots

With J. Hollingsworth
and F. Darwood, LANL

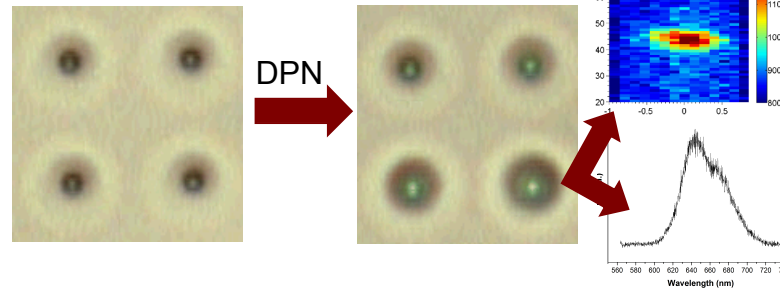


Dip-pen nanolithography



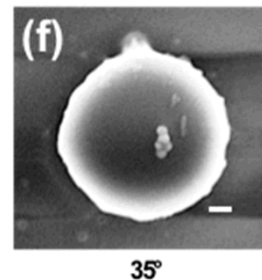
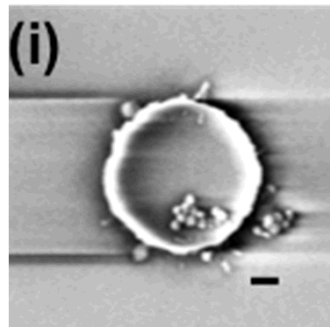
<http://str.llnl.gov>

Many QDs



QDs in protective polymer coating selectively deposited onto the tops of Si nanoresonators

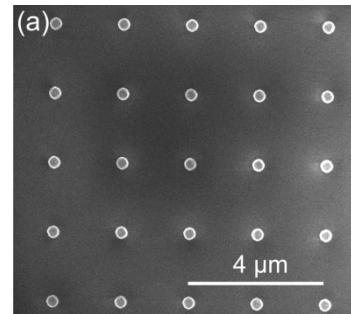
Towards Single QDs



(near-IR g-NQDs)
SEM: C. Sheehan, LANL

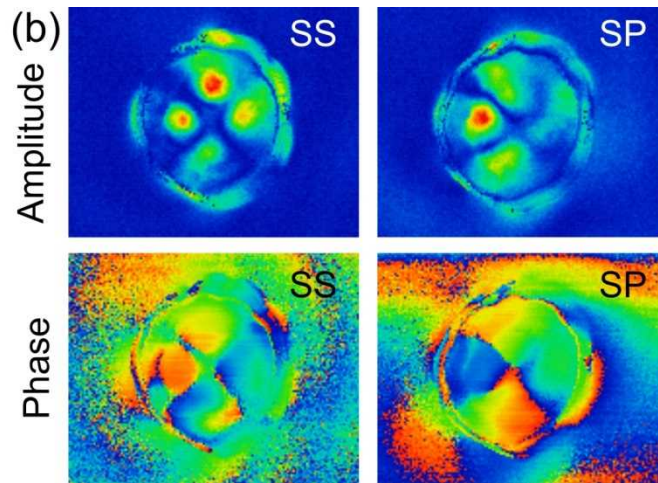
Near Field Imaging of Localized Modes

With Prof. Habteyes, UNM

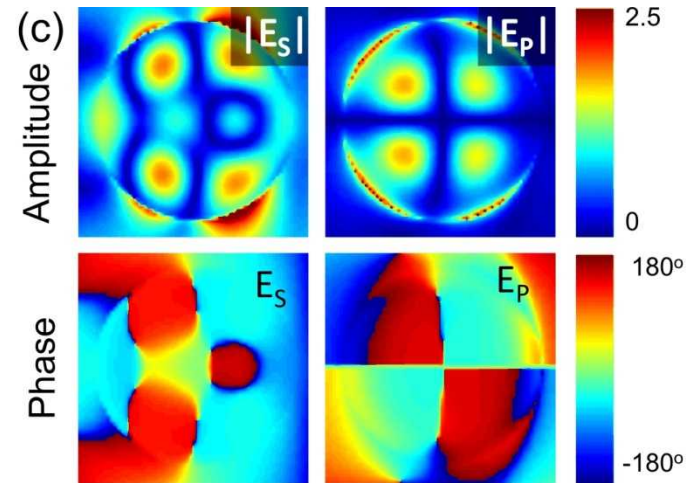


nanodisk diameter 412 nm,
nanodisk height 140 nm
 $\lambda=633\text{nm}$

Measurements



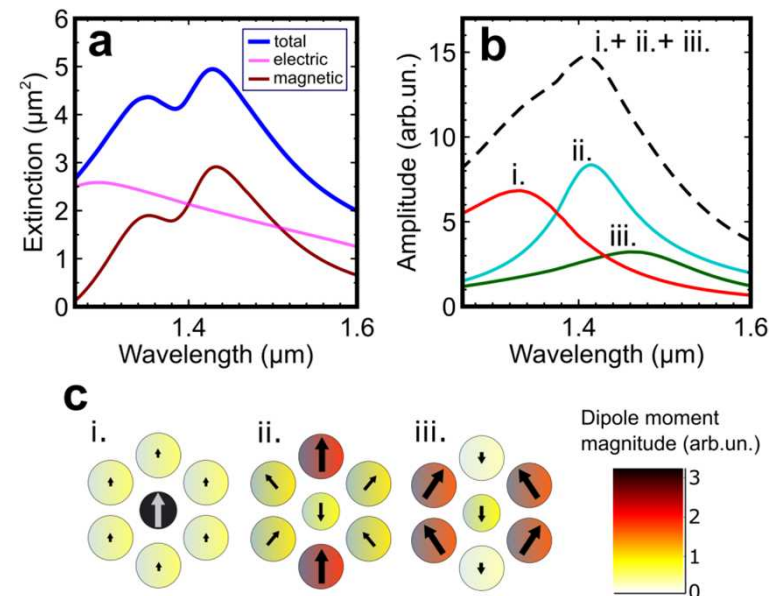
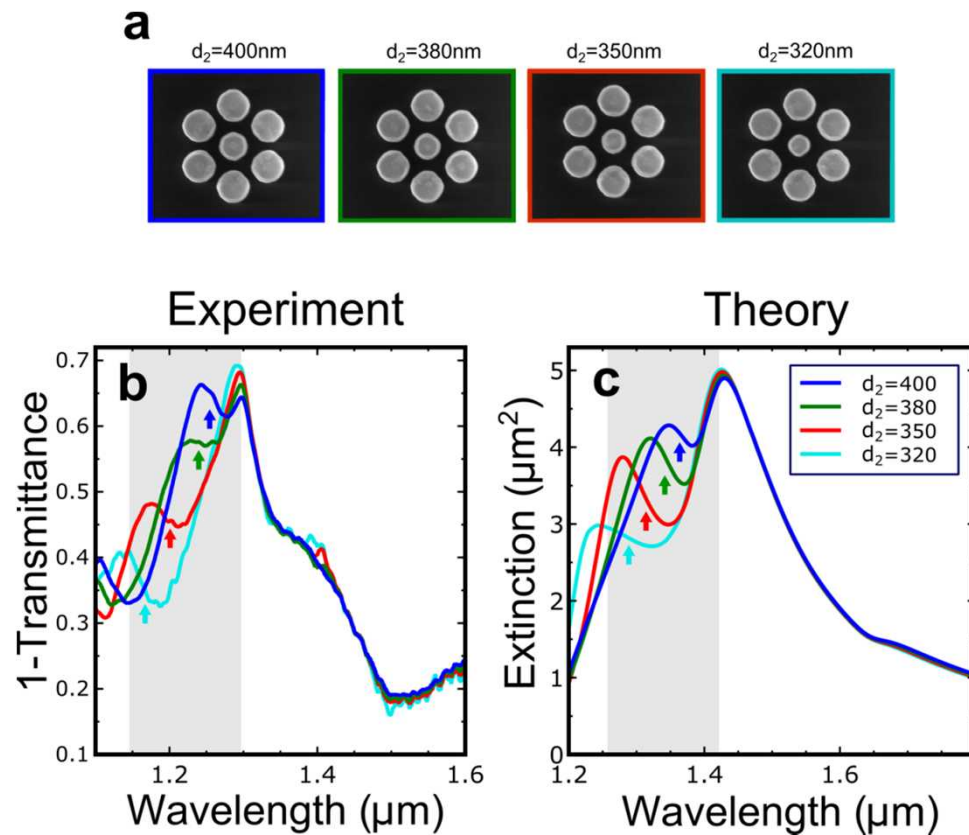
Calculated near-field amplitudes



Mostly electric quadrupole

Fano Resonances in All-Dielectric Nanoparticle Oligomers

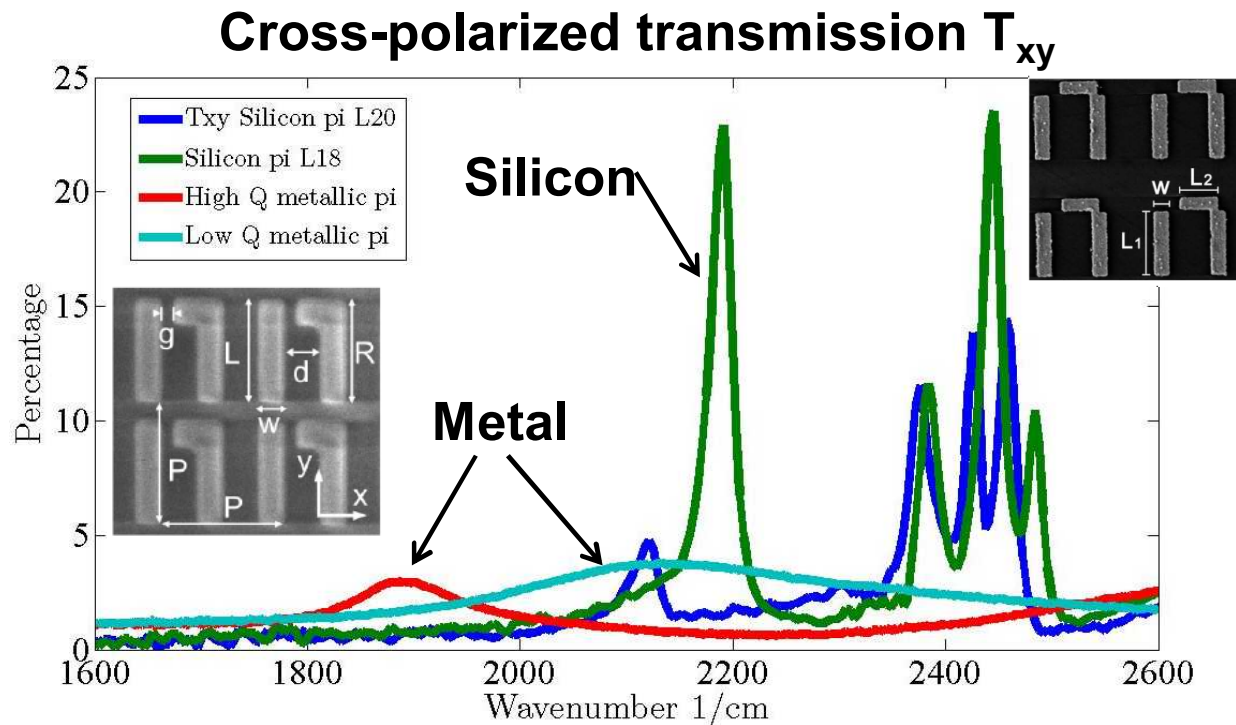
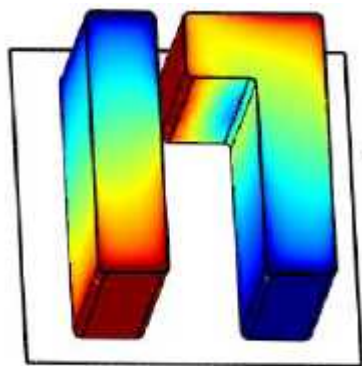
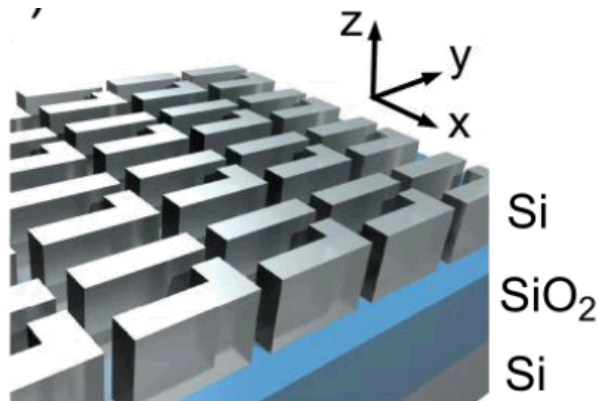
(with I. Staude & Y. Kivshar, ANU)



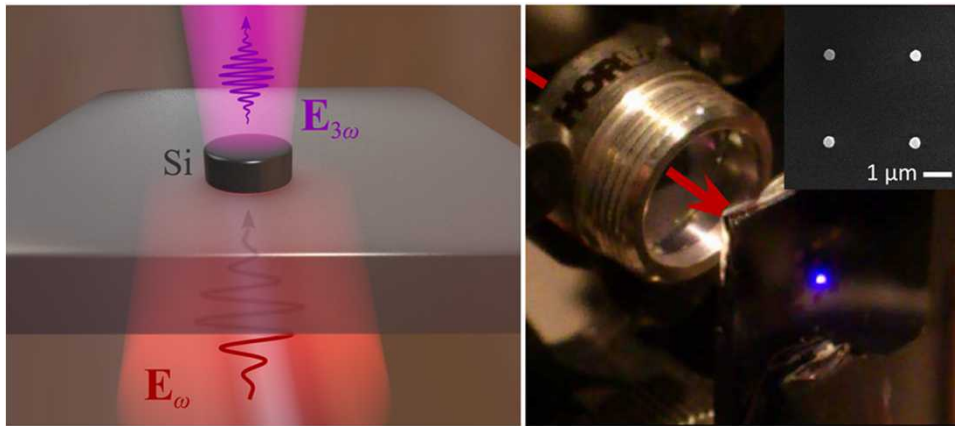


All-Dielectric Fano-Resonant 2D Chiral Metasurfaces:

Order of Magnitude Higher Q Than Similar Metallic Metasurfaces (with G. Shvets, UT Austin)

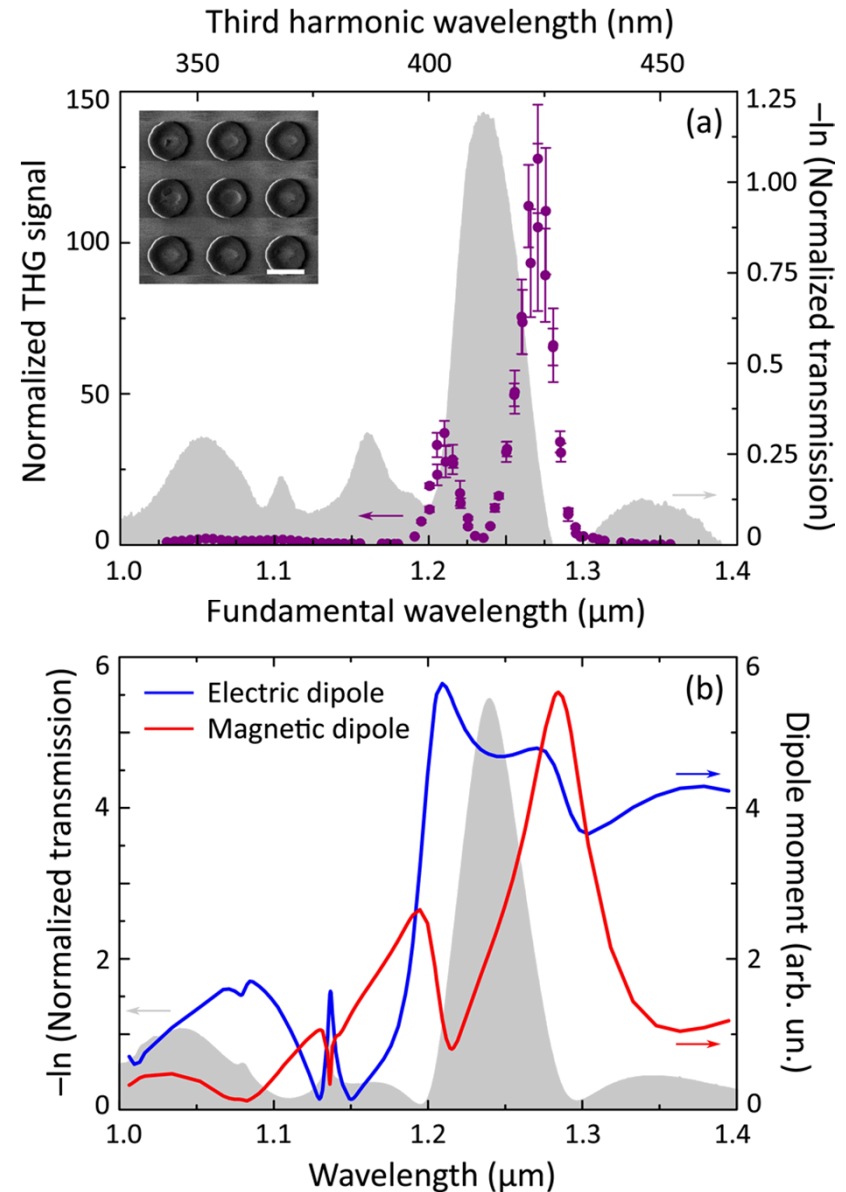


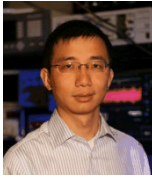
Third Harmonic Generation



(With ANU & Moscow State)

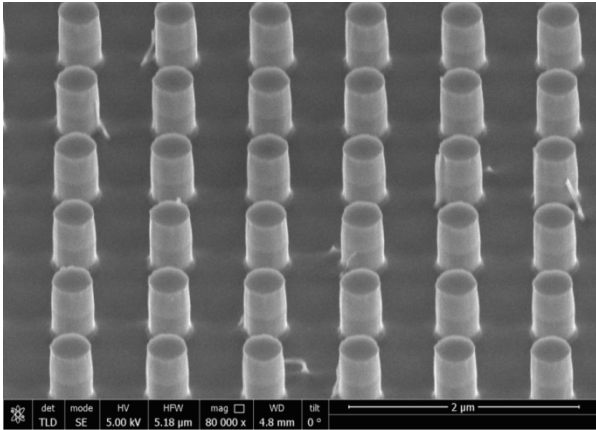
Nano Letters (2015)





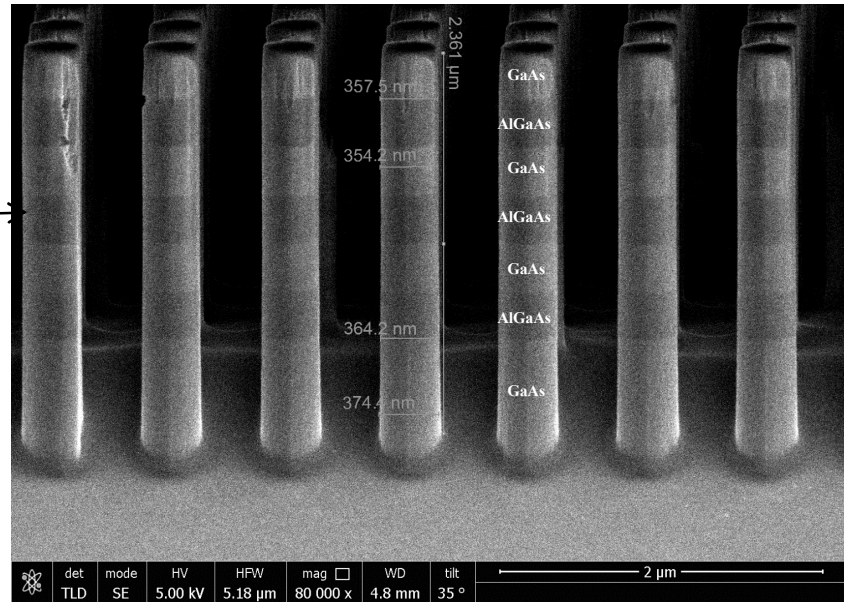
Sheng Liu

Dielectric Metamaterials with III-V's?

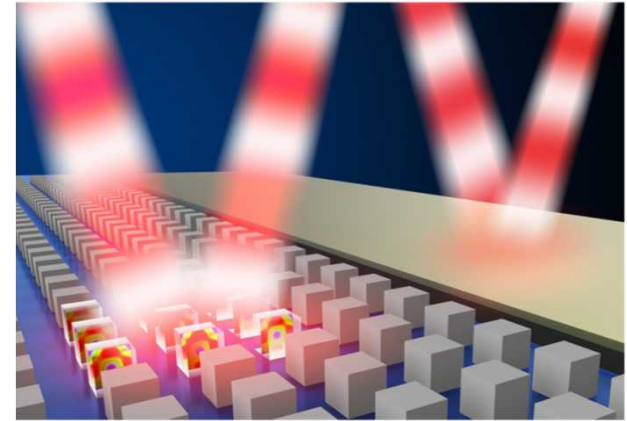


GaAs nanocylinders

Oxidized



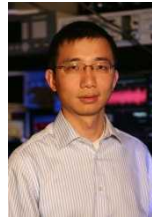
Summary



- Arrays of dielectric resonators were used to create an “optical magnetic mirror”.
- Proved magnetic mirror behavior using absolute optical phase measurements
- Through geometry and mode control, high index dielectric resonators offer a platform to create optical devices and metamaterial functionality.

ACKNOWLEDGMENTS:

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- **ANU:** Y. Kivshar, I. Staude, M. Decker, D. Neshev, ...
- **Lomonosov Moscow State:** M. Shcherbakov,...
- **UNM:** T. Habteyes
- **UT Austin:** G. Shvets



Center for Integrated Nanotechnologies: CINT



Center for Nanoscale
Materials
Argonne National Lab.

Molecular Foundry
Lawrence Berkeley National
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Center for Functional
Nanomaterials
Brookhaven National Lab.



Center for Integrated
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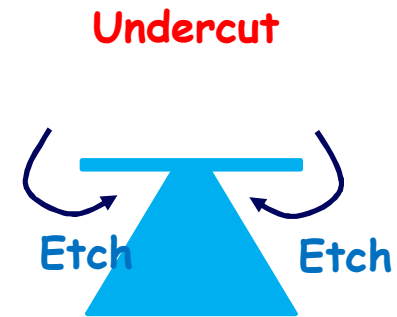
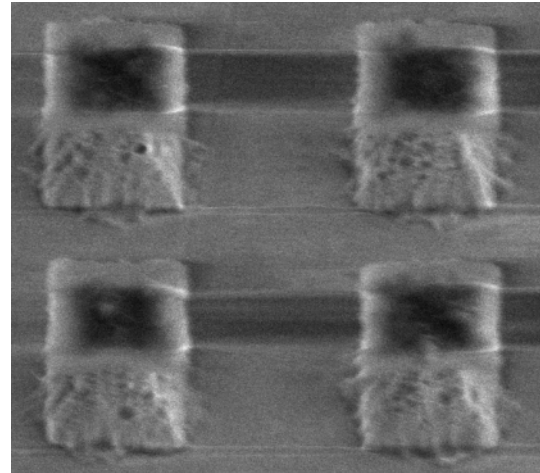
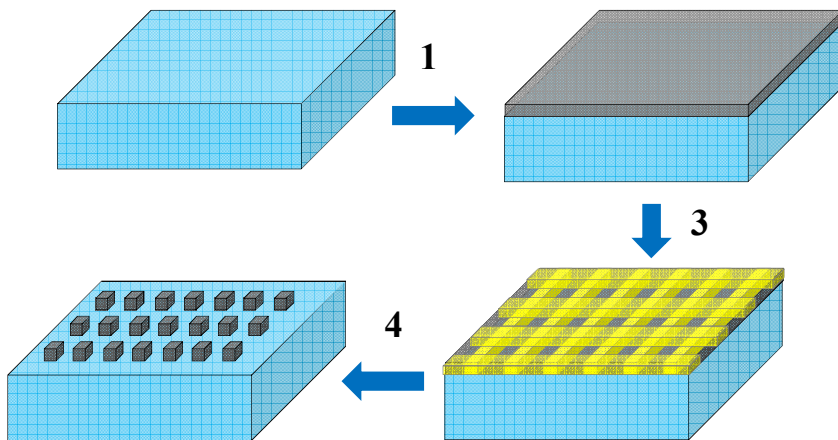
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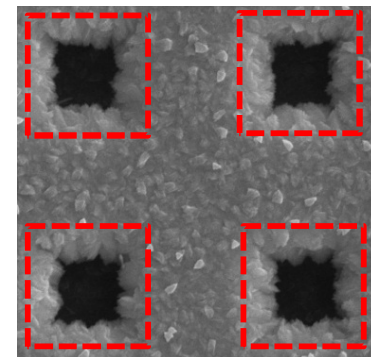
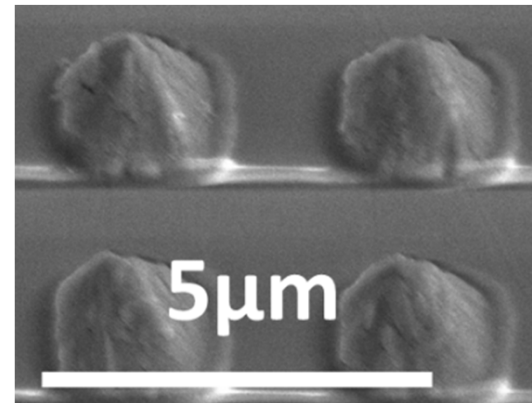
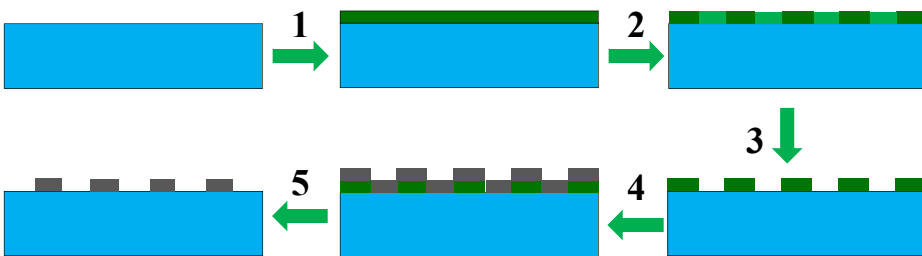
END

Prior (Difficult) Fabrication

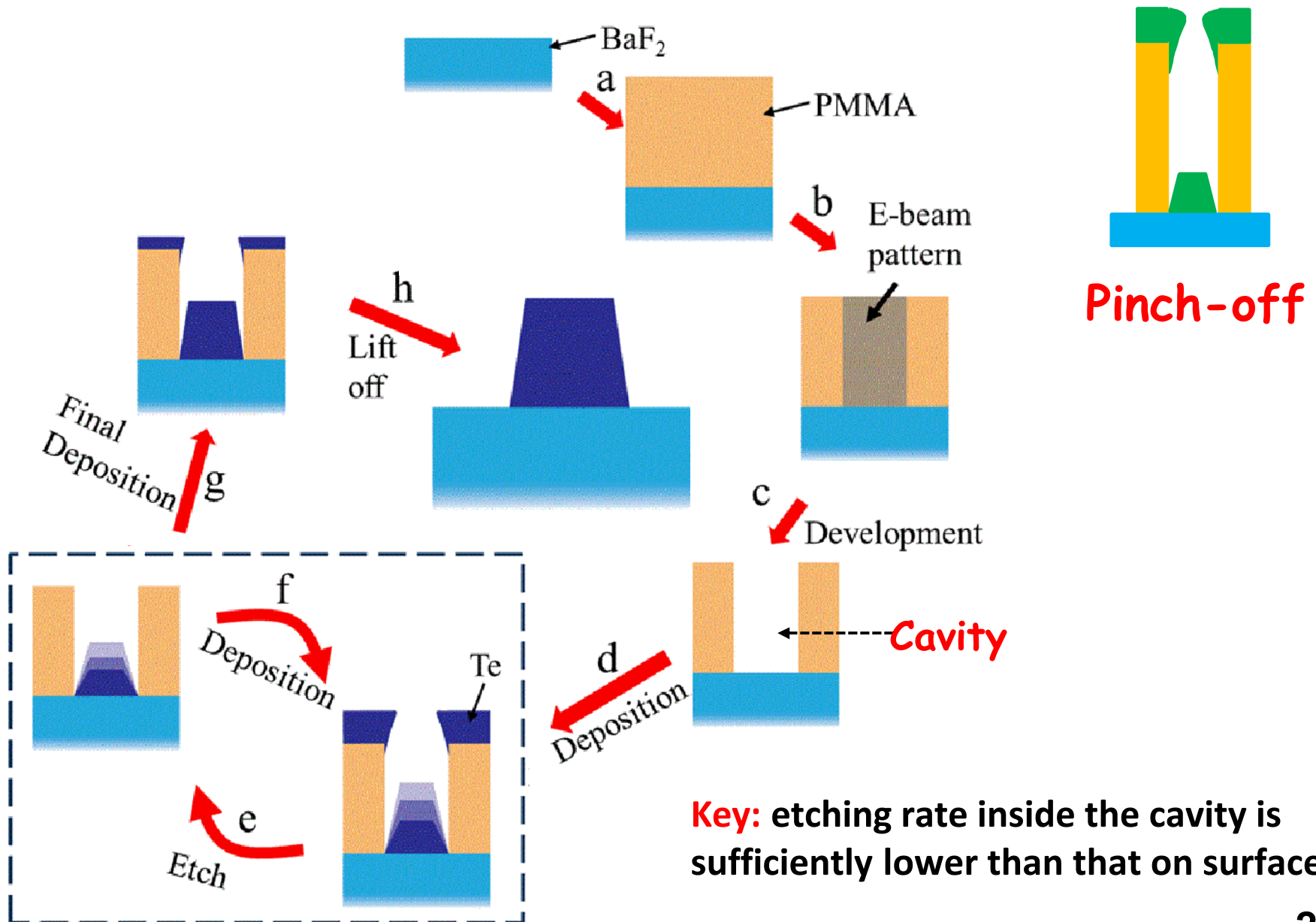
Deposition + Mask + Etching




Deposition + Liftoff

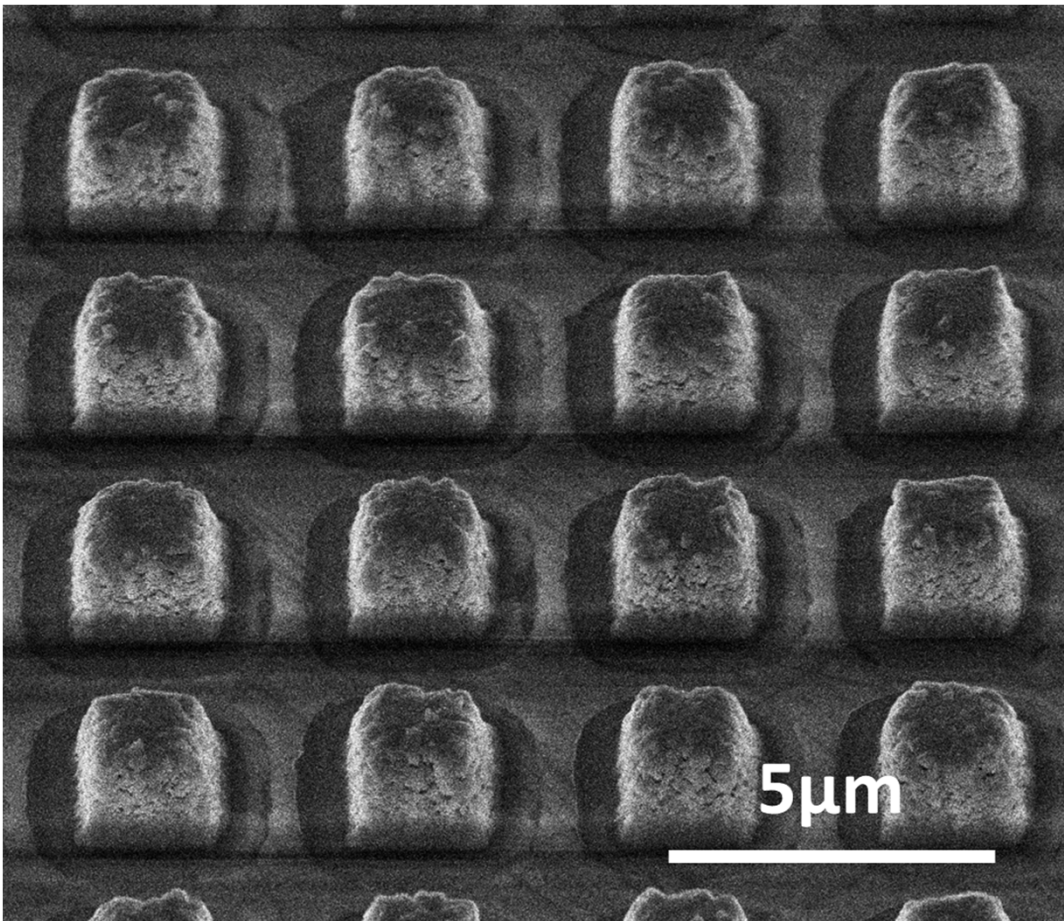


Multi-cycle Deposition-etch Process

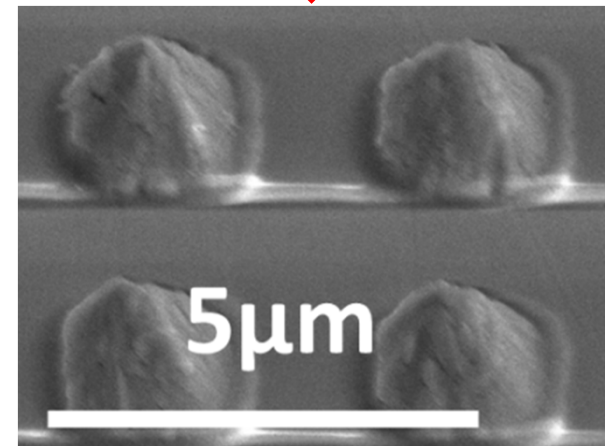


Multi-cycle Deposition-etch Process


2 Dep-etch
cycles

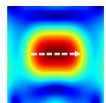
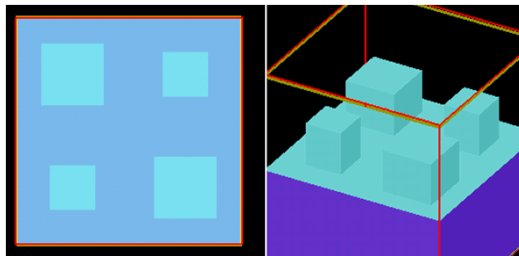


Single
deposition

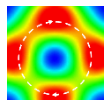


Transmission Spectra of Te metamaterial

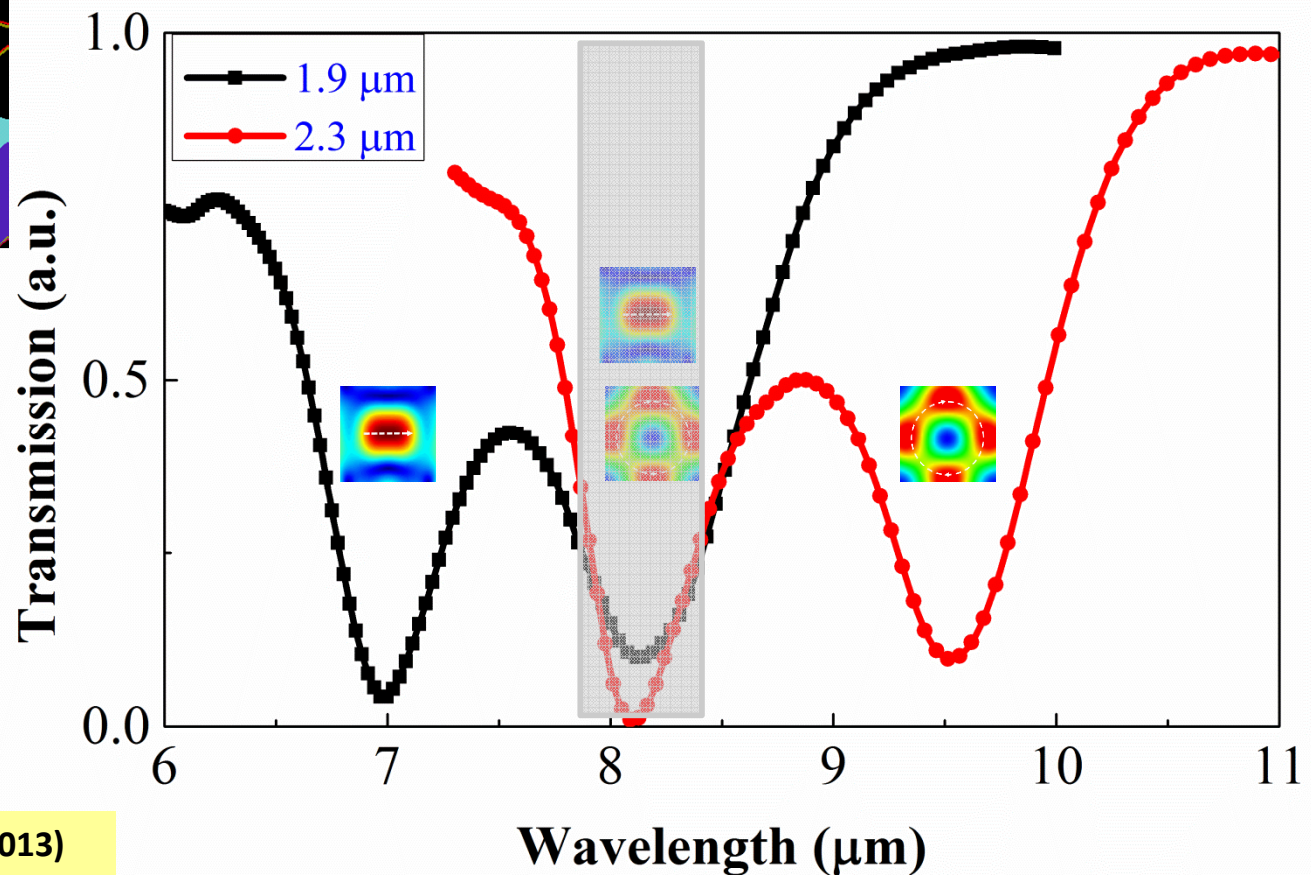
- Tuning of electric and magnetic resonances
- Overlap of electric and magnetic resonances (potential for **negative index material**)



$$\epsilon < 0$$

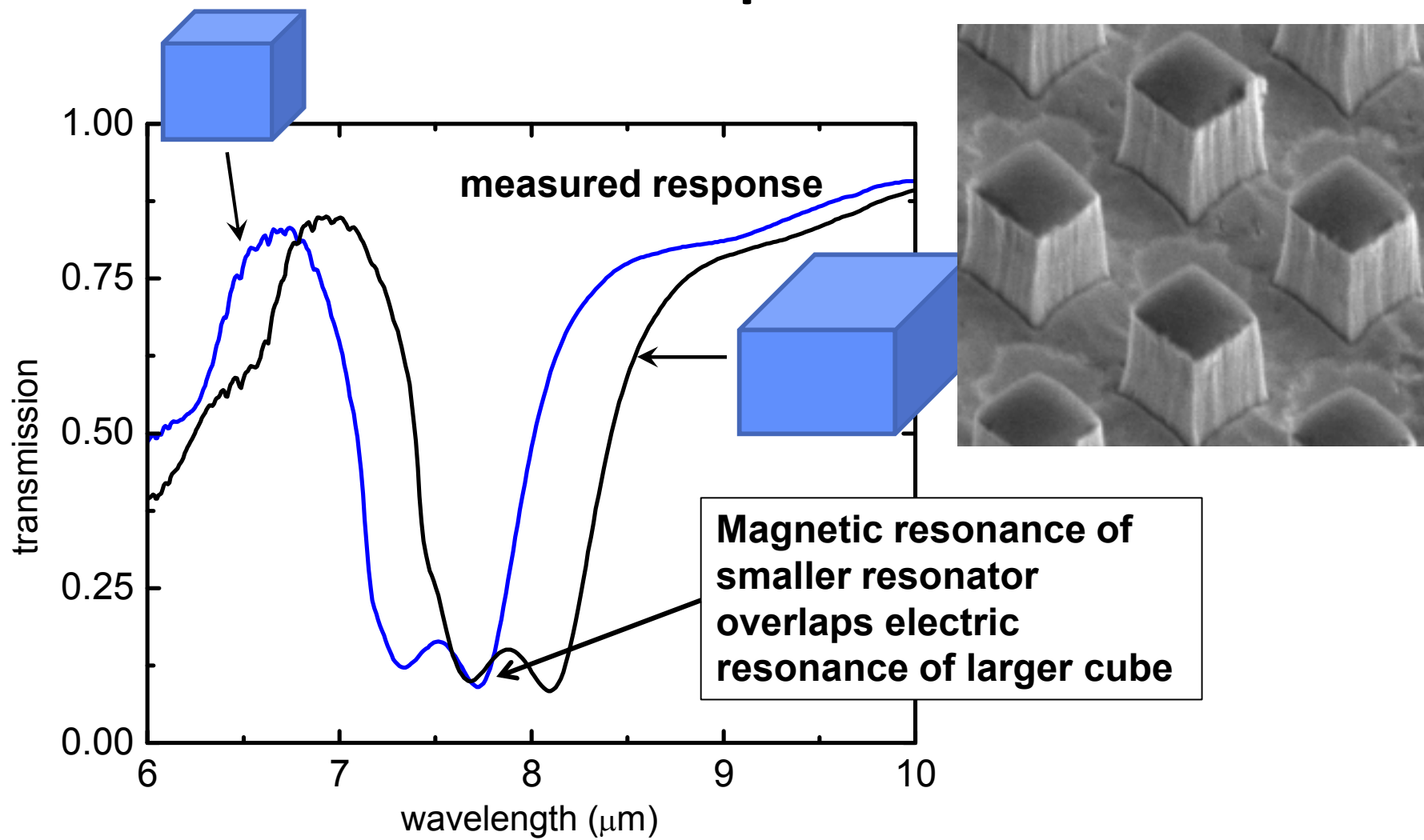


$$\mu < 0$$



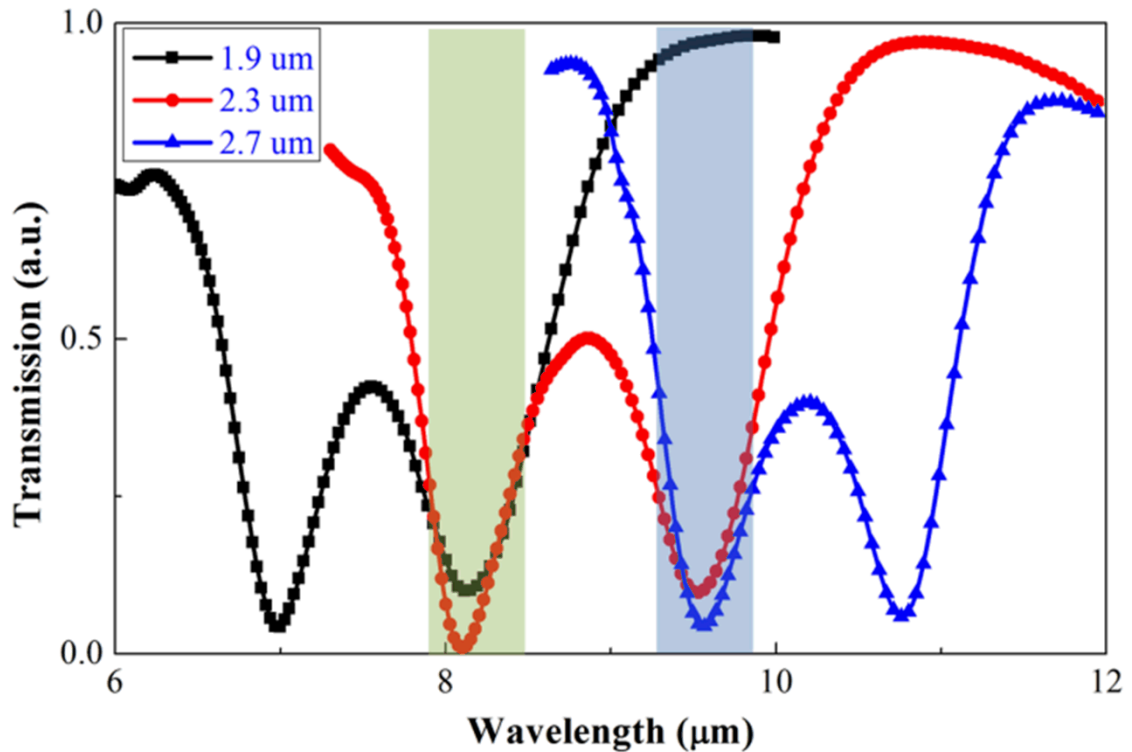
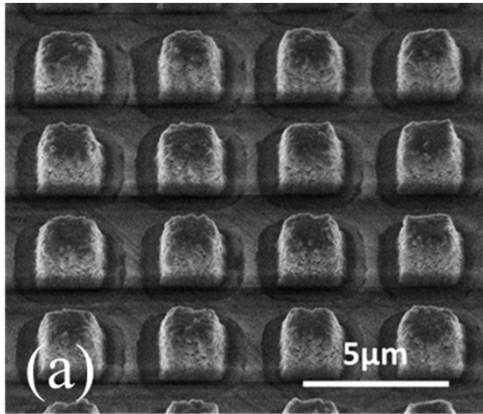


Te Resonators: Electric and Magnetic Mode Overlap



Magnetic and Electric resonances can overlap: different size cubes, or cubes with “cuts”

Te Resonators: Overlapping Electric and Magnetic Resonances



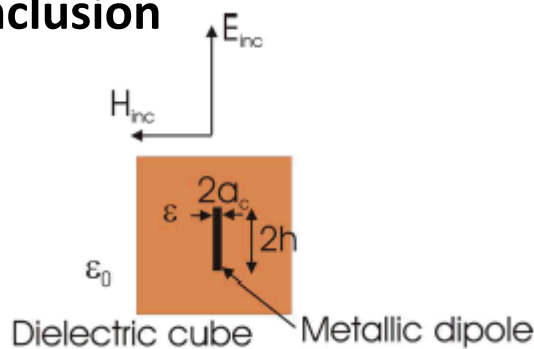
Te resonators using a single multiple deposition & liftoff process.

- Side: 1.7 μm , 2.3 μm and 2.7 μm with 50% duty cycle (height constant at $\sim 1.8 \mu\text{m}$).
- Shaded areas are the spectral regions where magnetic and electric resonances overlap.

Adjusting the Resonance Frequencies: Perturbation Approach

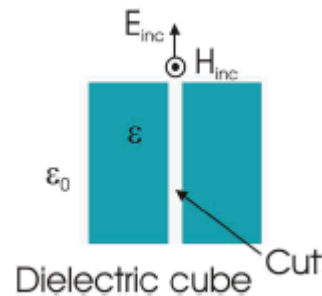
- One approach to alignment of resonances is to place perturbations within the resonator volume

High-permittivity inclusion



Frequency downshift
of electric resonance

Low-permittivity inclusion



Frequency upshift
of magnetic resonance

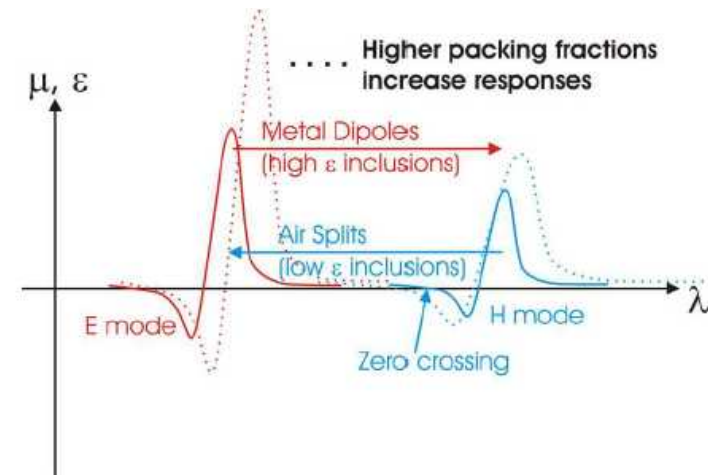
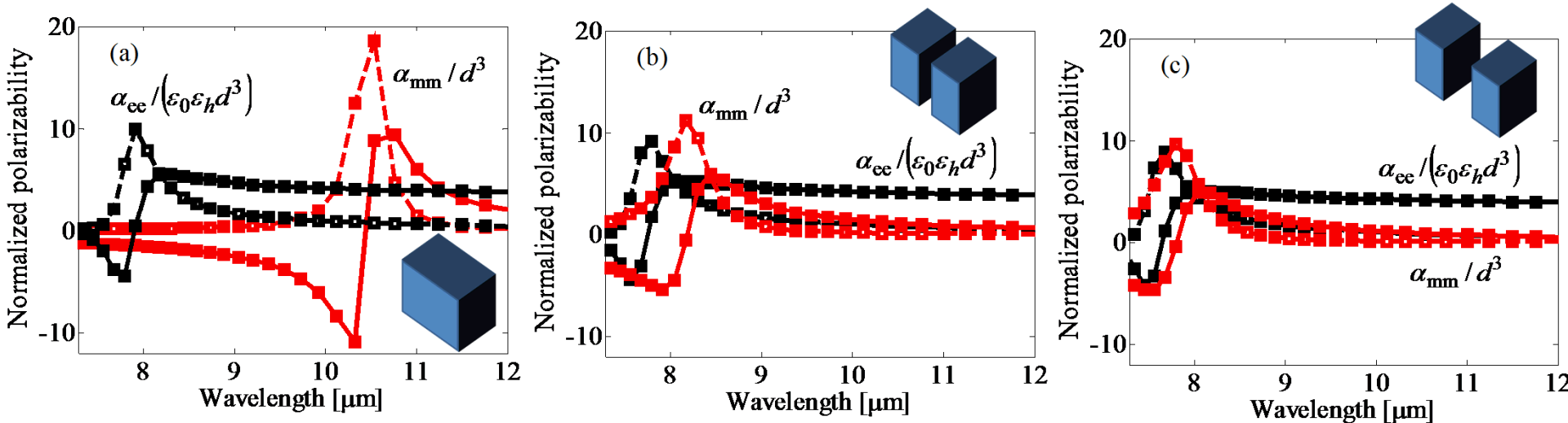


Illustration of resonance
shifting using these
methods

Perturbation theory to control relative location of electric and magnetic resonances

- Apply perturbation theory to a dielectric cube e.g. by introducing a split
Warne et al. *PIER B* 44, 1-29 (2012); *IEEE Trans. Antennas Propagat.* 61, 2130-2141 (2013)
- Retrieve electric and magnetic polarizabilities of single scatterers using full-wave simulations

Basilio et al. *IEEE Antennas Wireless Propag. Lett.* 10, 1567 (2011); Rockstuhl et al. *Phys. Rev. B* 83, 245119 (2011)

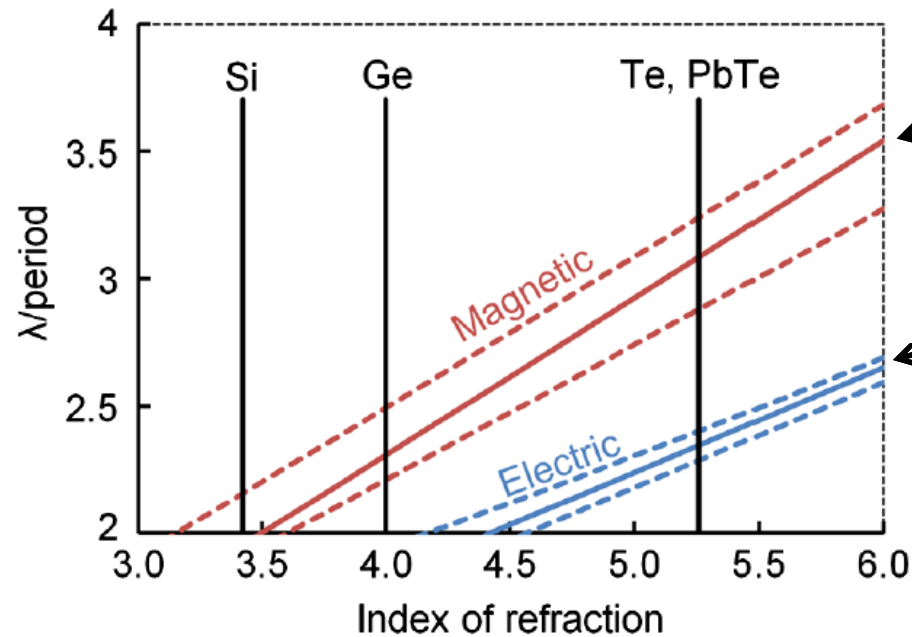


- Electric and magnetic polarizabilities can be easily engineered to overlap at a given frequency

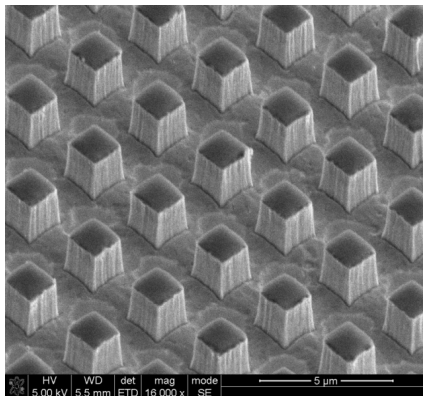
Scaling Dielectric Resonators to the Near IR

Design metric for
1:1 CDR
metamaterials

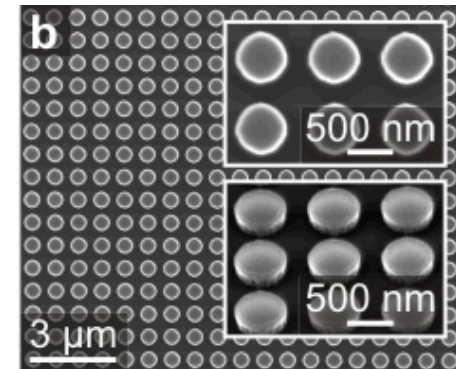
PRL 108, 097402 (2012)



Peak of
M-
resonance
Peak of E-
resonance



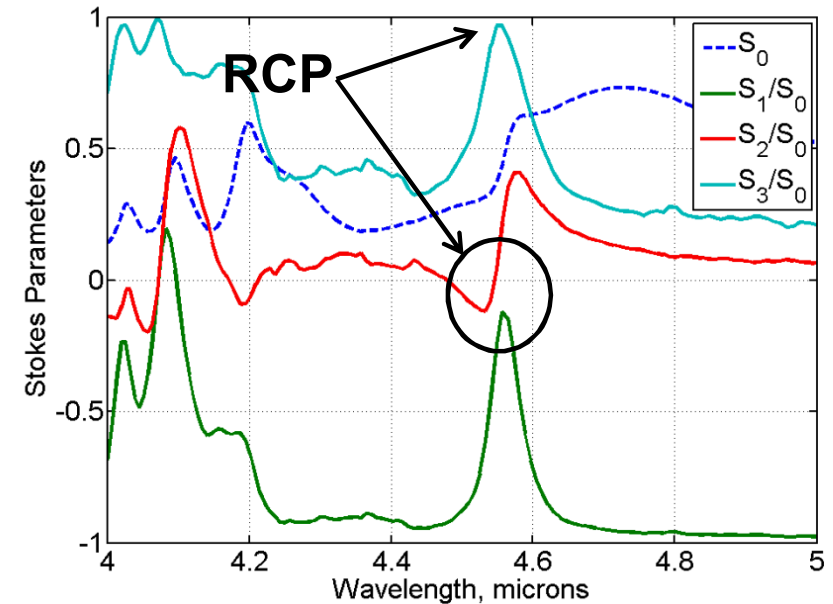
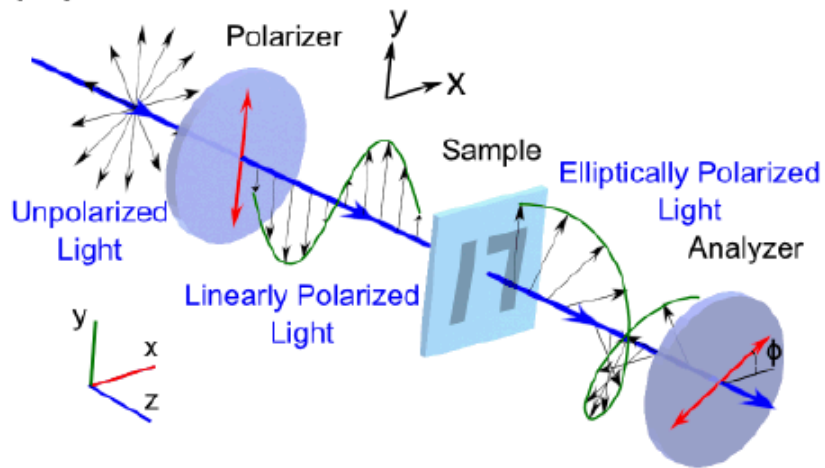
Tellurium: $n \sim 5$, Size $\sim 1.5 \mu\text{m}$
 $\lambda > 5 \mu\text{m}$



Silicon: $n \sim 3.5$, Size $< 200\text{-}500\text{nm}$
 $\lambda > 1 \mu\text{m}$

LP-CP Metasurface Conversion

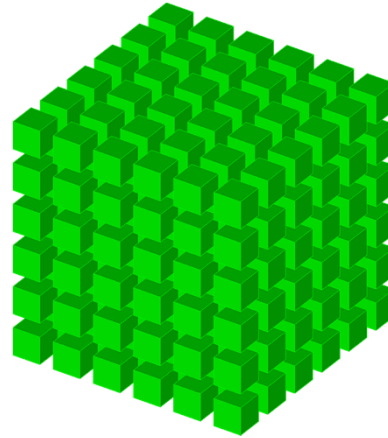
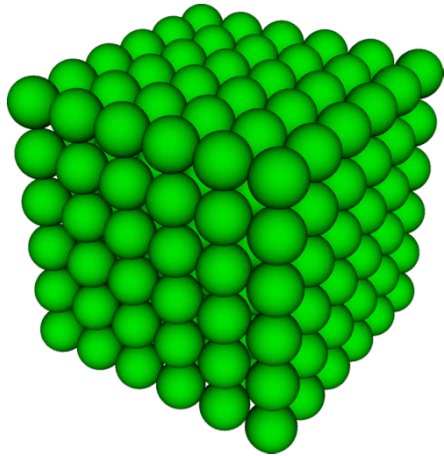
Rotating analyzer Stokes polarimetry



- LP \rightarrow CP conversion by 1 micron thick MS with high efficiency (50%) at normal incidence
- Experimentally measured quality factor
- Close to 100% Degree of Circular Polarization

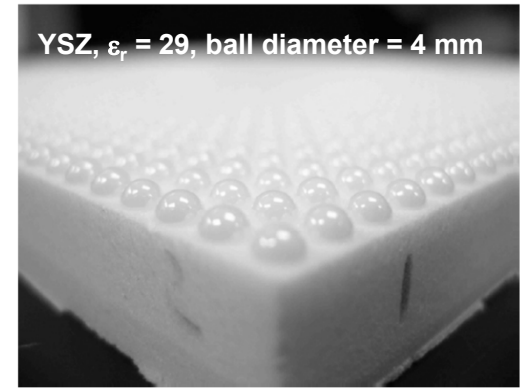
Could use a chiral metasurface for developing a CP thermal emitter

Arrays of High ϵ Resonators: Low Loss Metamaterials



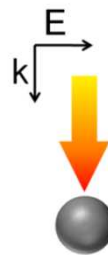
Works nice in the RF

YSZ, $\epsilon_r = 29$, ball diameter = 4 mm



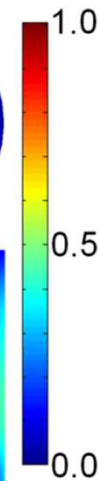
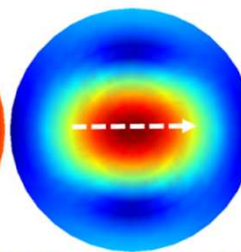
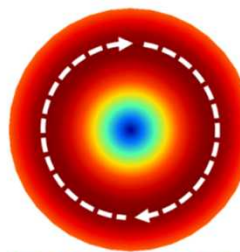
(negative ϵ , negative μ)

Spherical

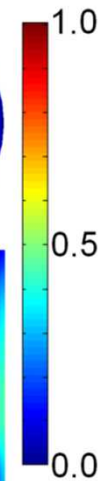
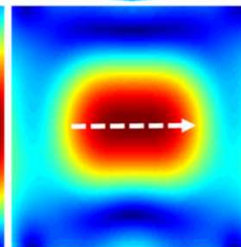
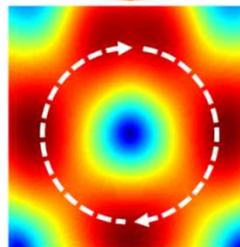
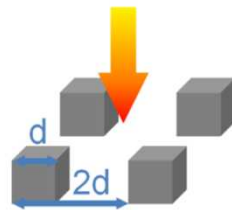


Magnetic

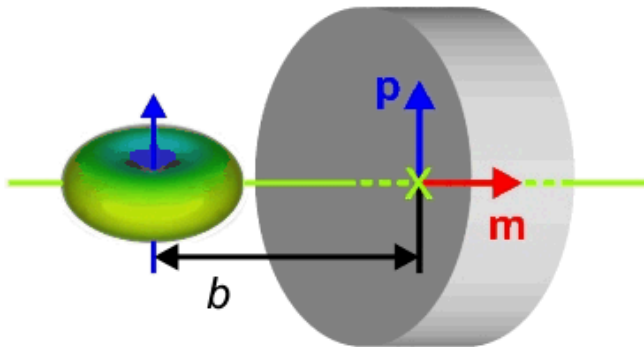
Electric



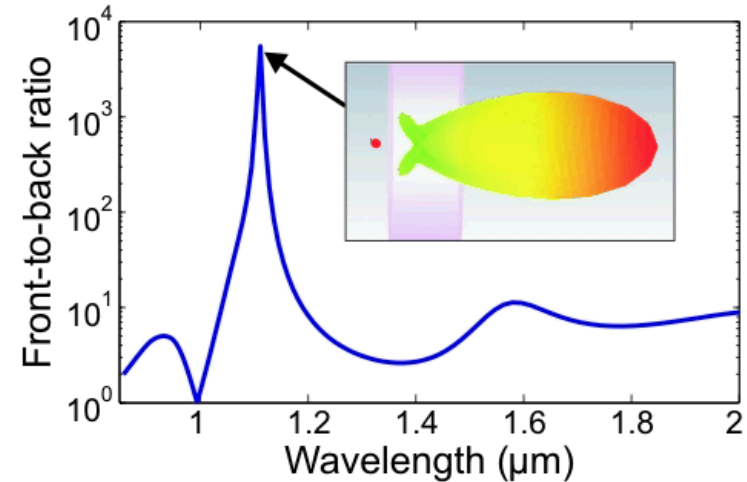
Cubic



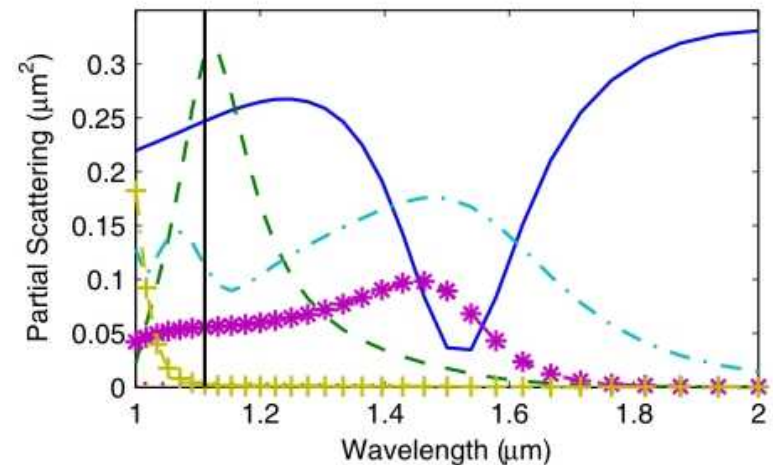
Importance of the Quadrupole Mode



- Silicon nanodisks: highly directional nanoantennas with giant front-to-back ratio
- Electric quadrupole mode essential to achieve high directivity



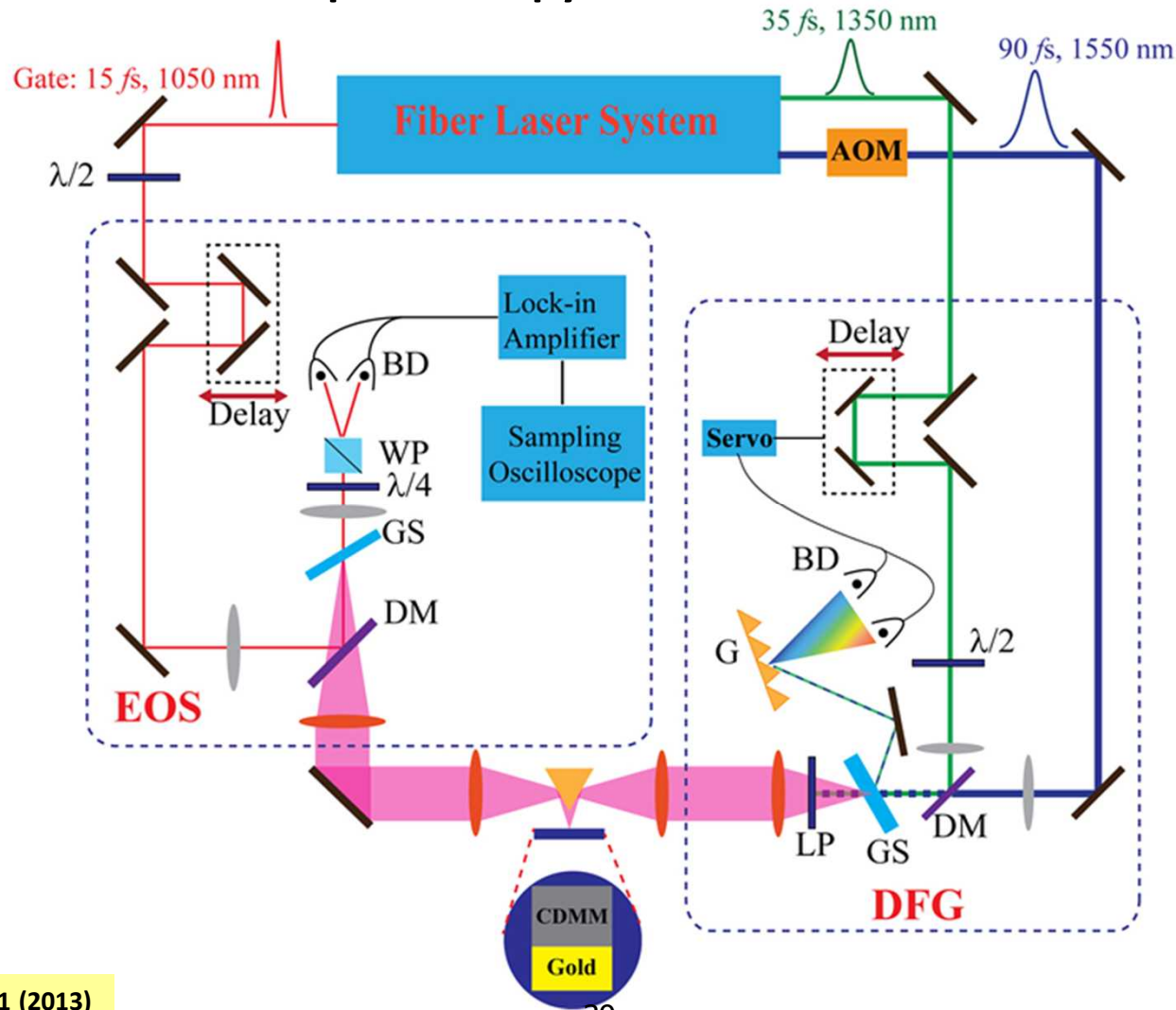
I. Staude *et al.*, *ACS Nano* 7, 7824, 2013.



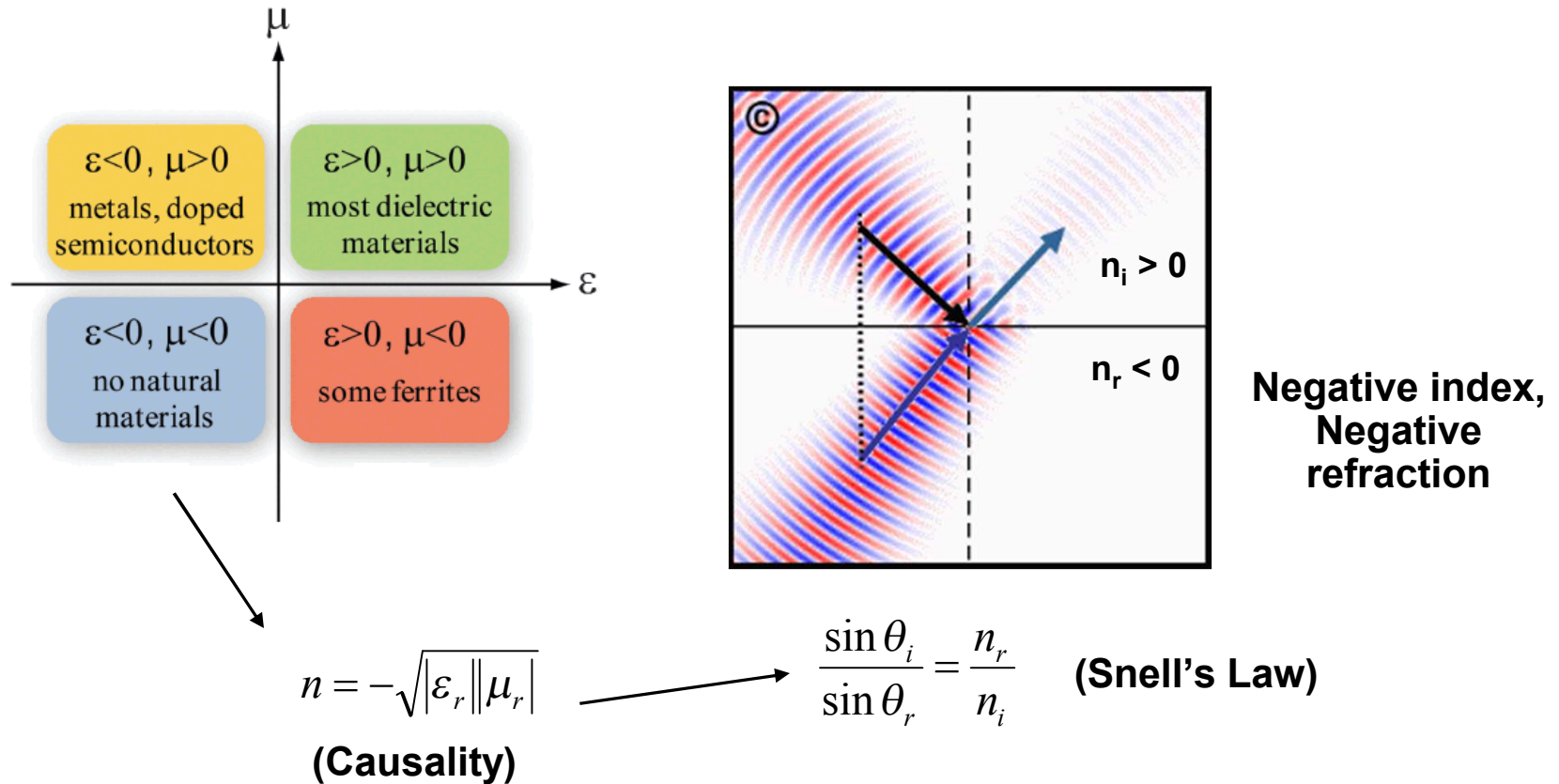
E. Rusak *et al.*, *submitted* (2014).

Proving Optical Magnetism: Measure Phase of Reflected Wave

Phase-locked Time Domain Spectroscopy



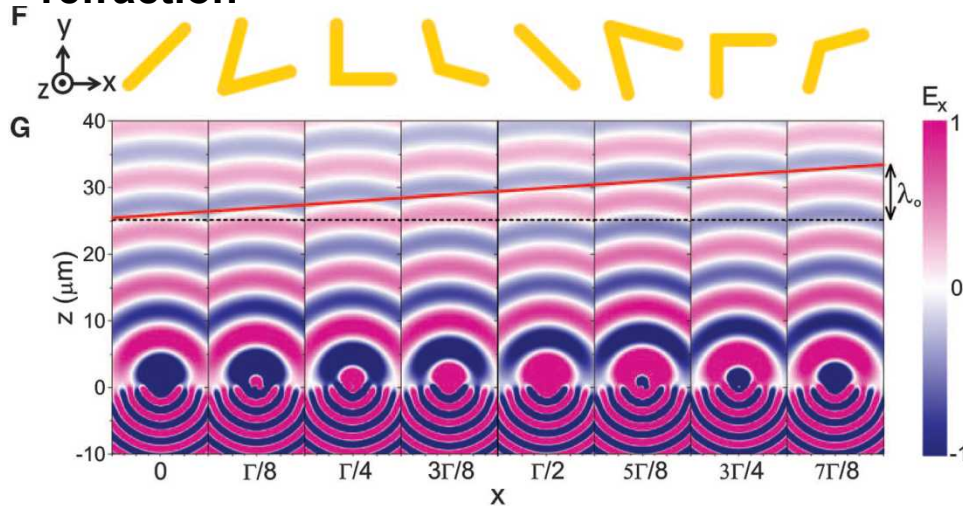
Metamaterials: Exotic optical properties



- Superlensing, Cloaking, Chirality/Optical activity, Perfect absorption
- Enhanced nonlinear interaction, Optical force manipulation, Light emission control

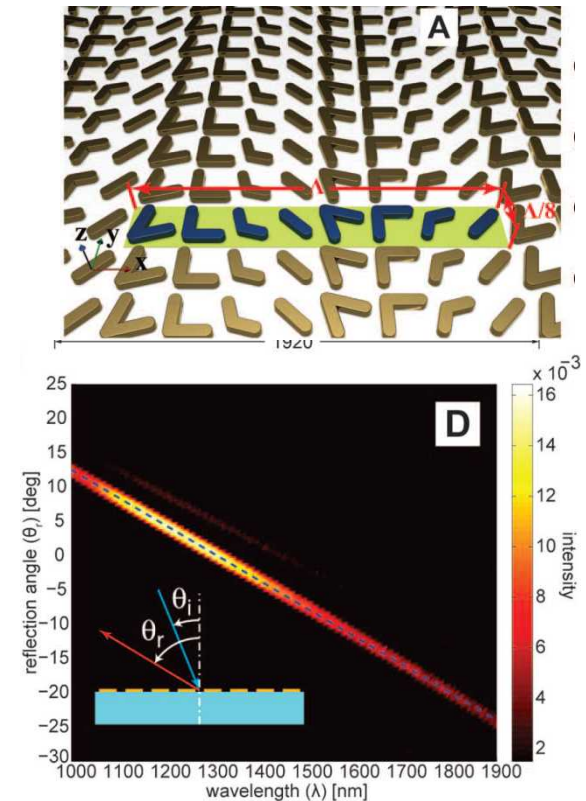
Beam Manipulation With Metasurfaces

Phase gradient to achieve anomalous refraction



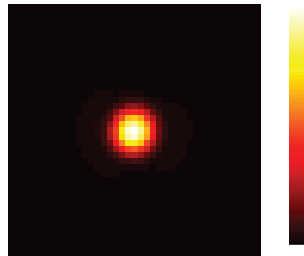
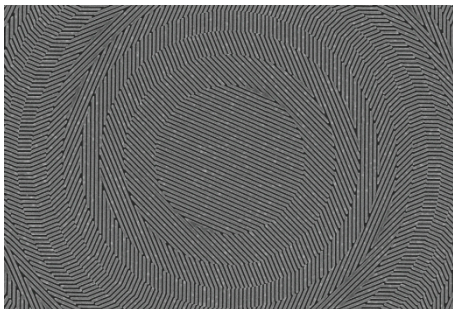
Yu et al. *Science* 334, 333 (2011)

Phase gradient to achieve anomalous reflection



Ni et al. *Science* 335, 427 (2012)

Phase gradient to achieve lensing



Lin et al. *Science* 345, 298 (2014)