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Kinetic Simulations of Particle Acceleration at Shocks

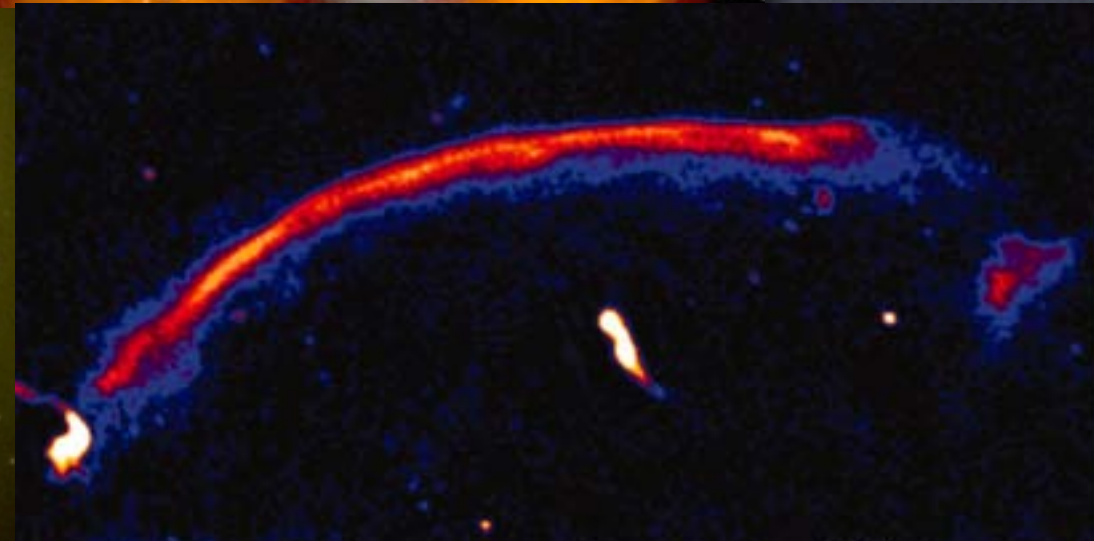
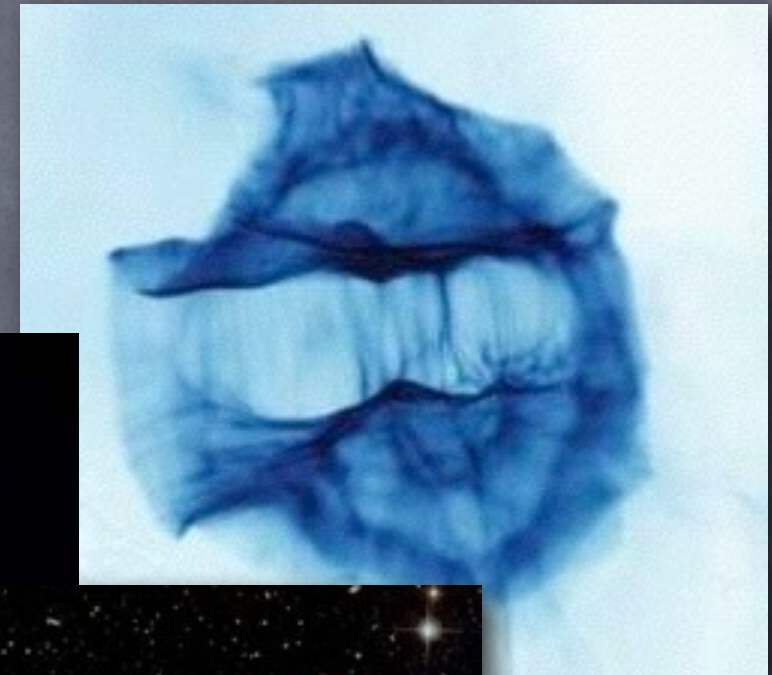
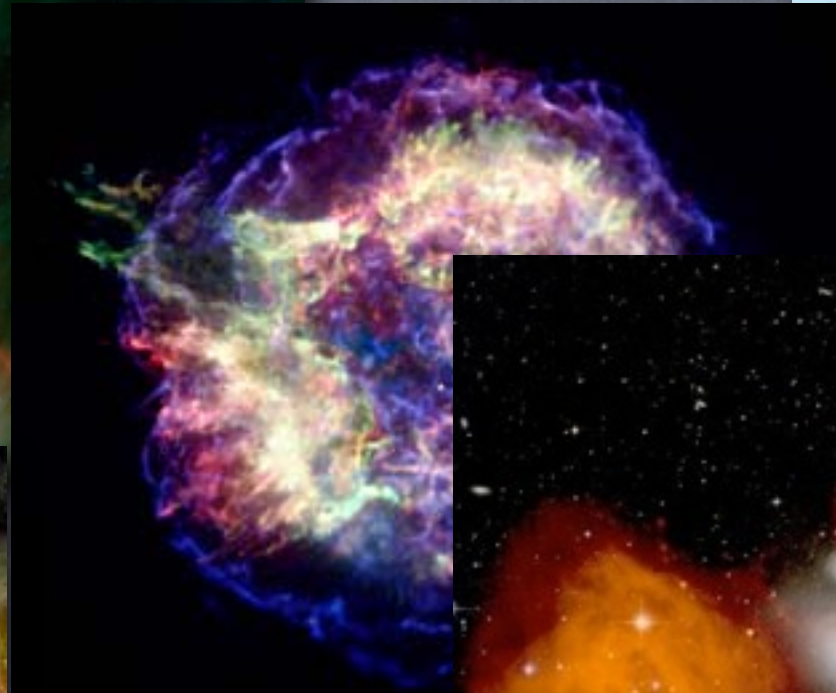
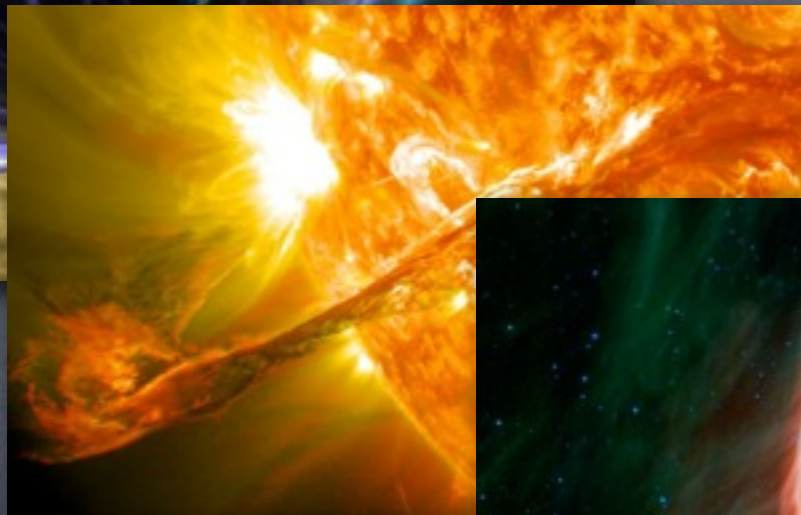
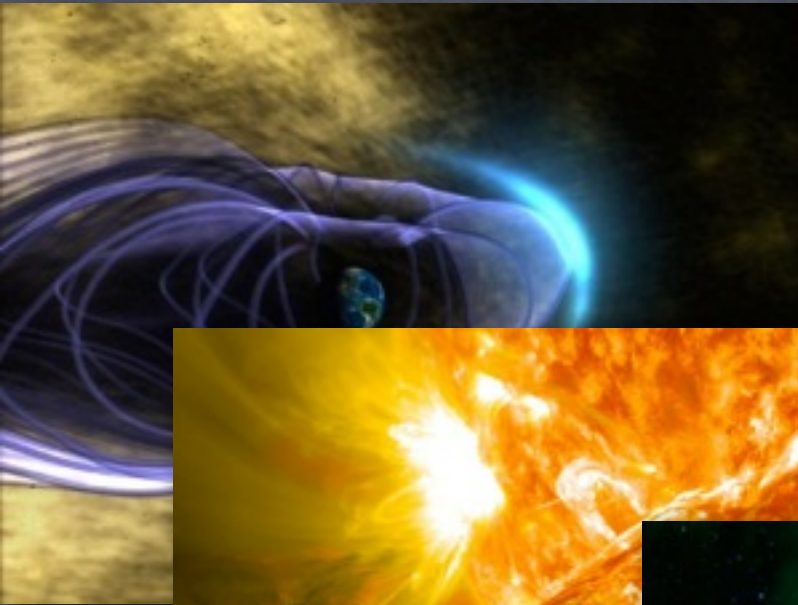
Damiano Caprioli
Princeton University



Collisionless shocks



- Mediated by **collective** electromagnetic interactions
- Sources of **non-thermal** particles and emission
- Reproducible in laboratory



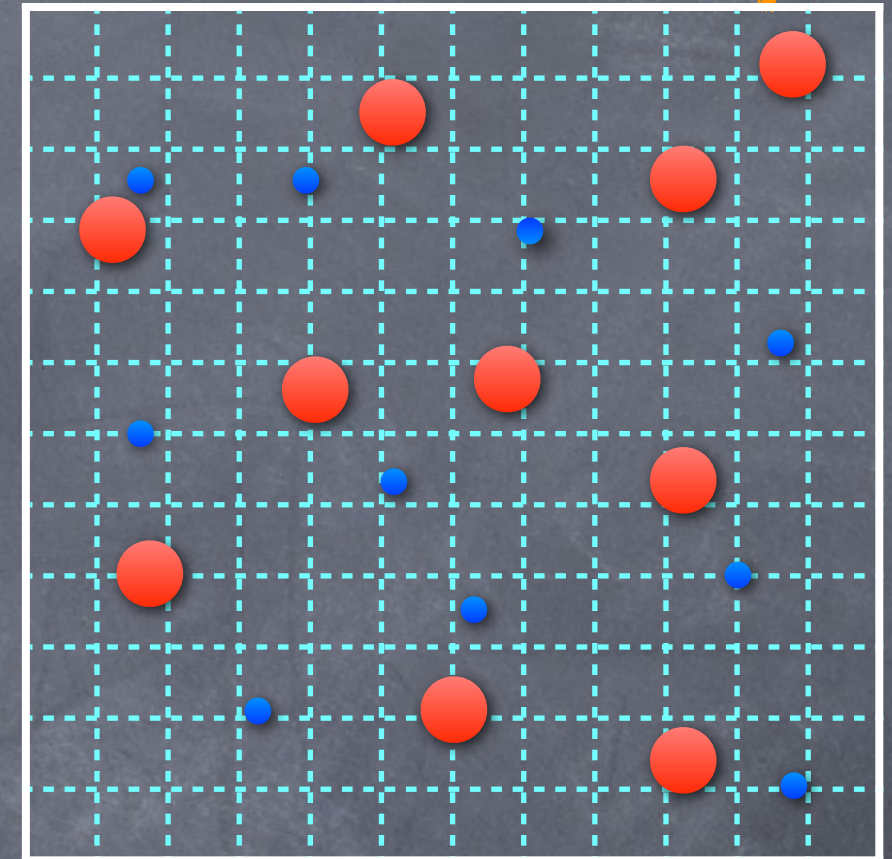
Acceleration from first principles



- **Full particle in cell** approach

(..., Spitkovsky 2008, Niemiec+2008, Stroman+2009, Riquelme & Spitkovsky 2010, Sironi & Spitkovsky 2011, Park+2012,2015, Niemiec+2012, Guo+2014, DC+15...)

- Define electromagnetic field on a **grid**
- Move particles via **Lorentz force**
- Evolve fields via **Maxwell equations**
- Computationally very challenging!

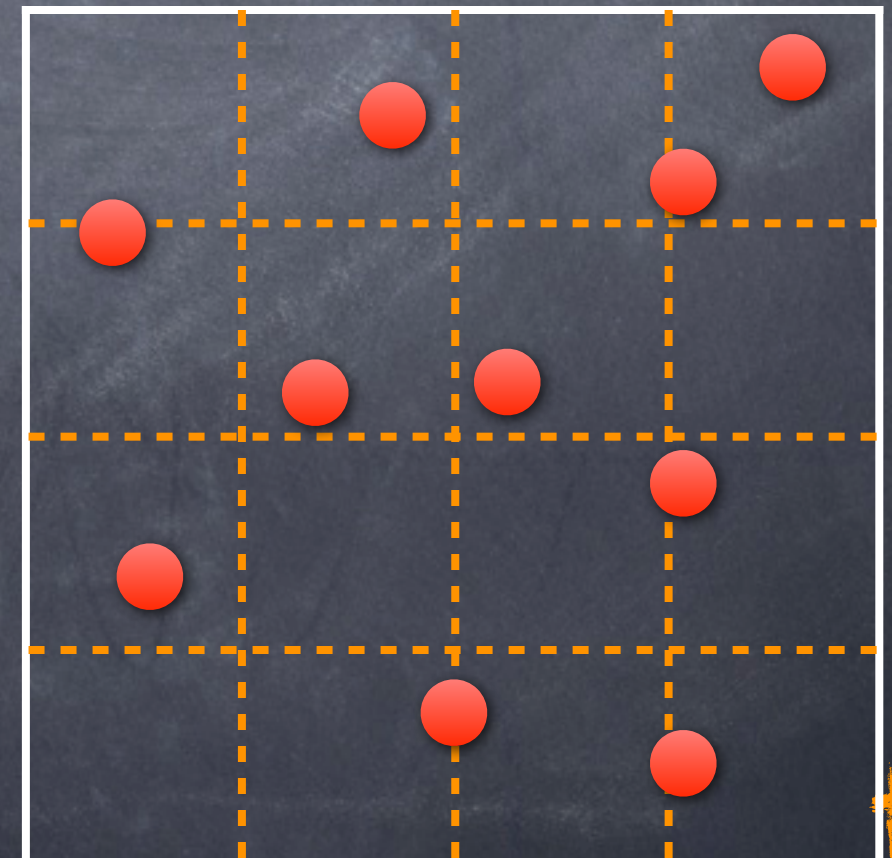


- **Hybrid approach:**

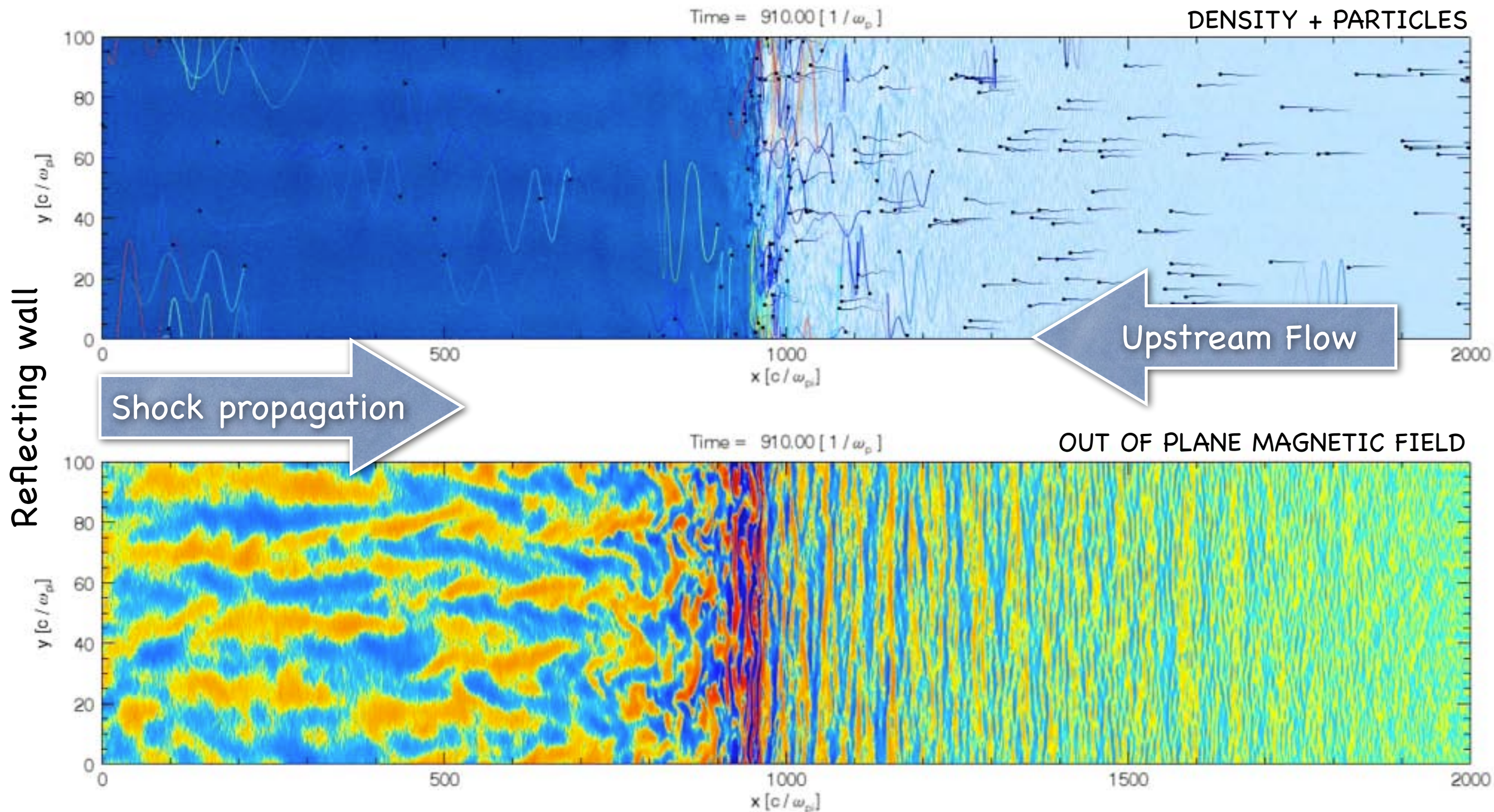
Fluid electrons – Kinetic protons

(Winske & Omid; Lipatov 2002; Giacalone et al.; Gargat  & Spitkovsky 2012, DC & Spitkovsky 2013–2015,...)

- massless electrons for more **macroscopical** time/length scales



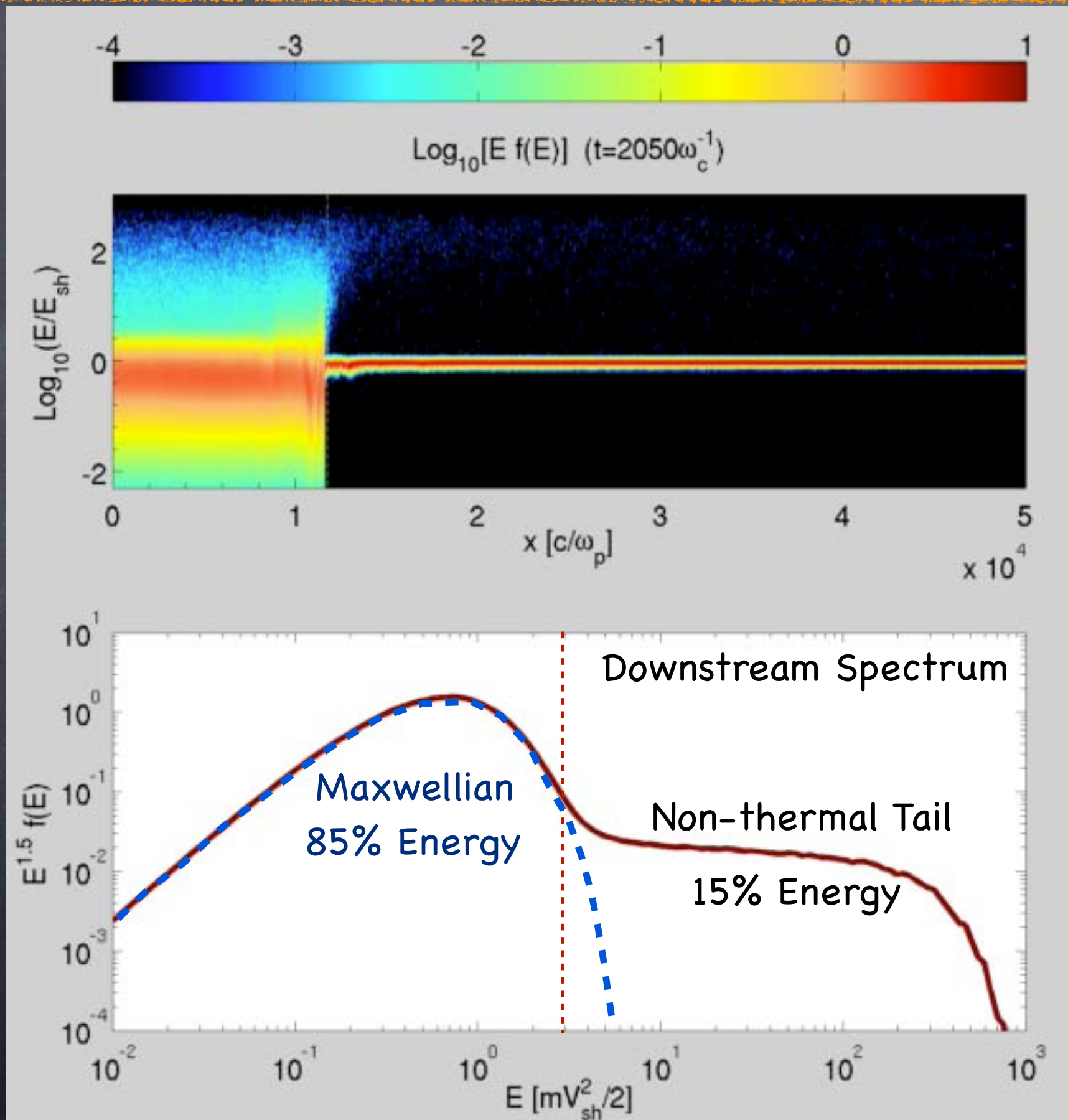
Hybrid simulations of collisionless shocks



 **dHybrid** code (Gargaté et al, 2007)

Initial B field

Spectrum evolution



First-order Fermi acceleration:

$$f(p) \propto p^{-4}$$

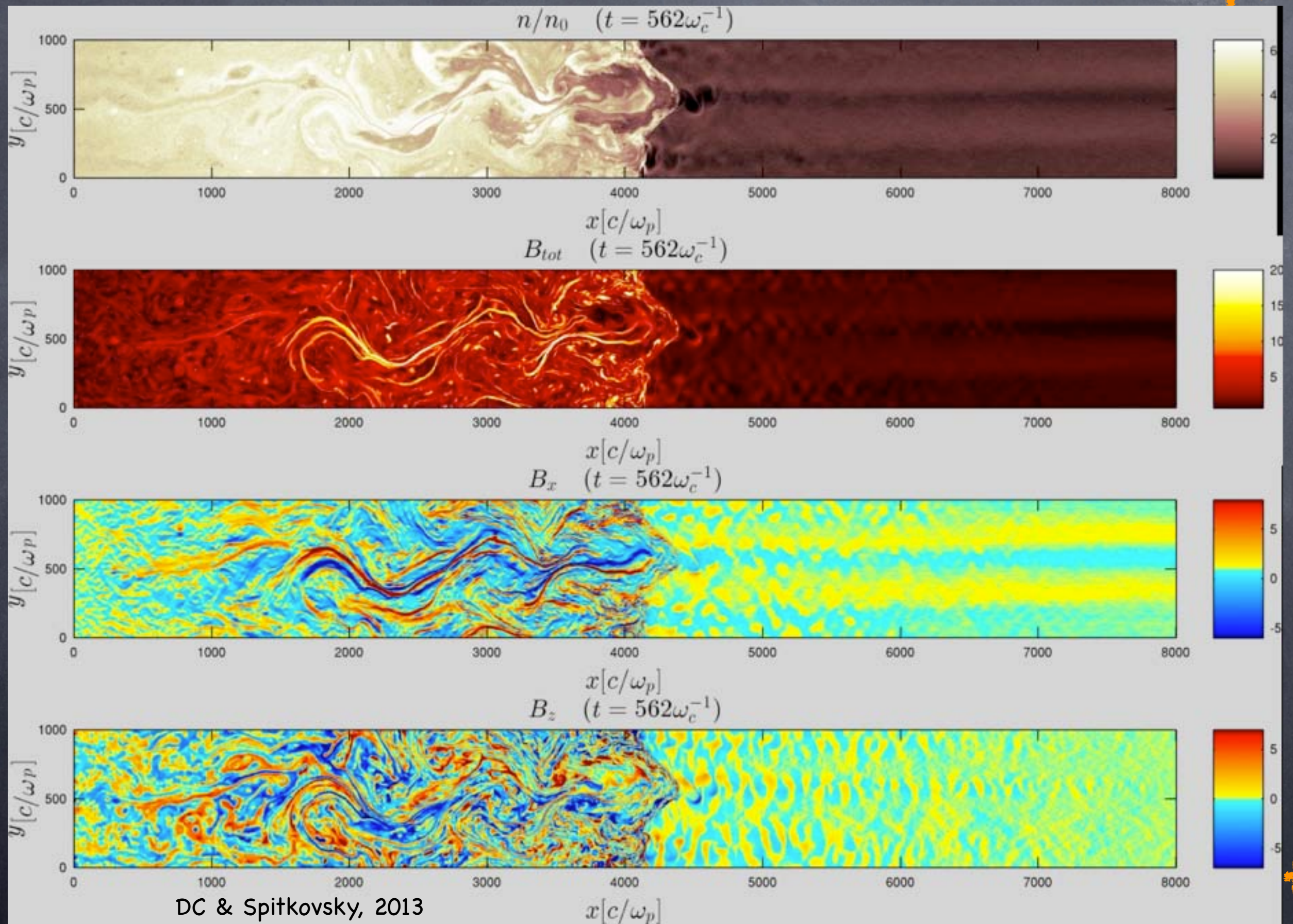
$$4\pi p^2 f(p) dp = f(E) dE$$



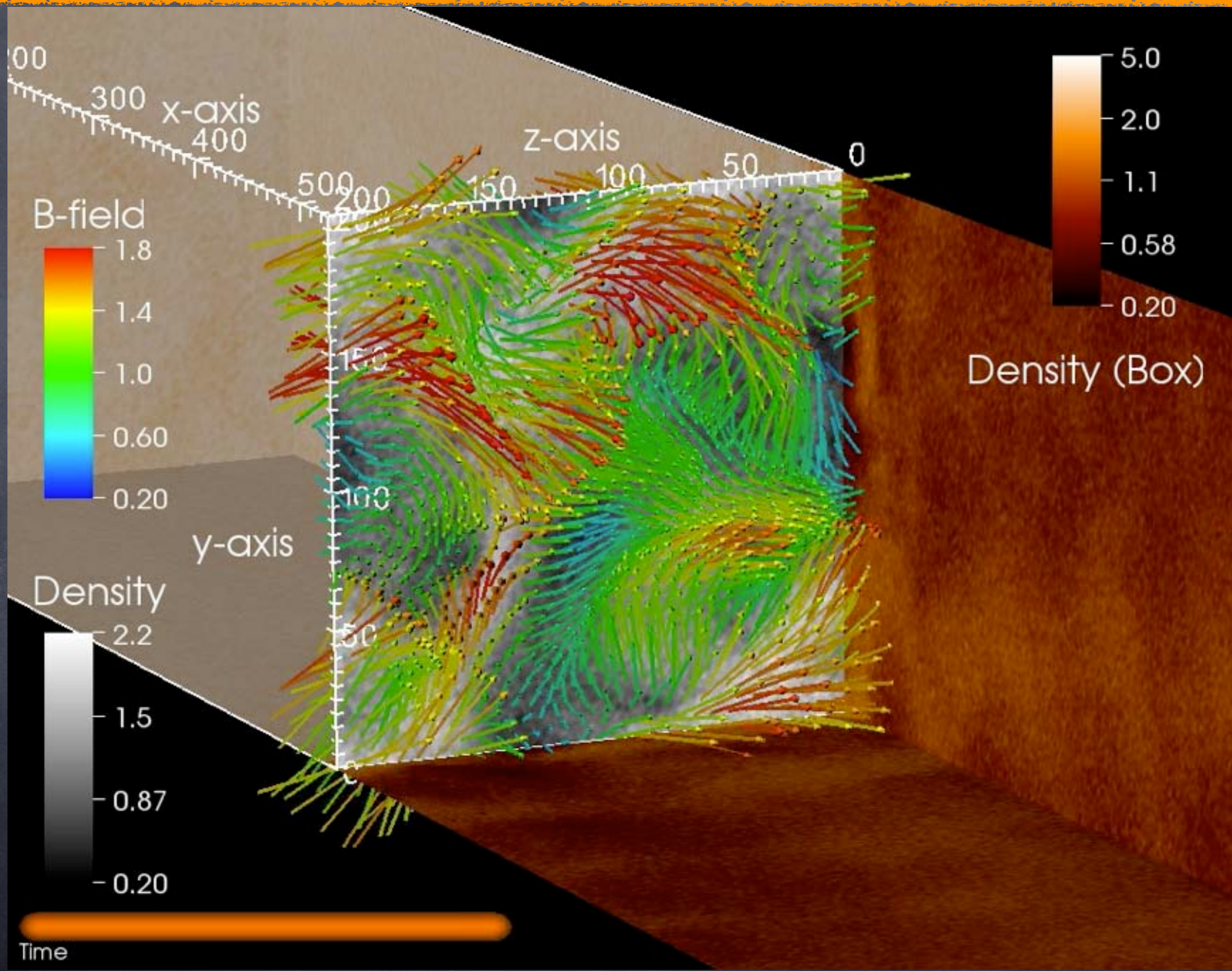
$$f(E) \propto E^{-2} \text{ (relativ.)}$$

$$f(E) \propto E^{-1.5} \text{ (non rel.)}$$

CR-driven instability

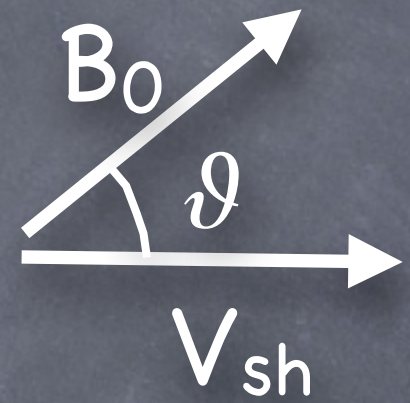
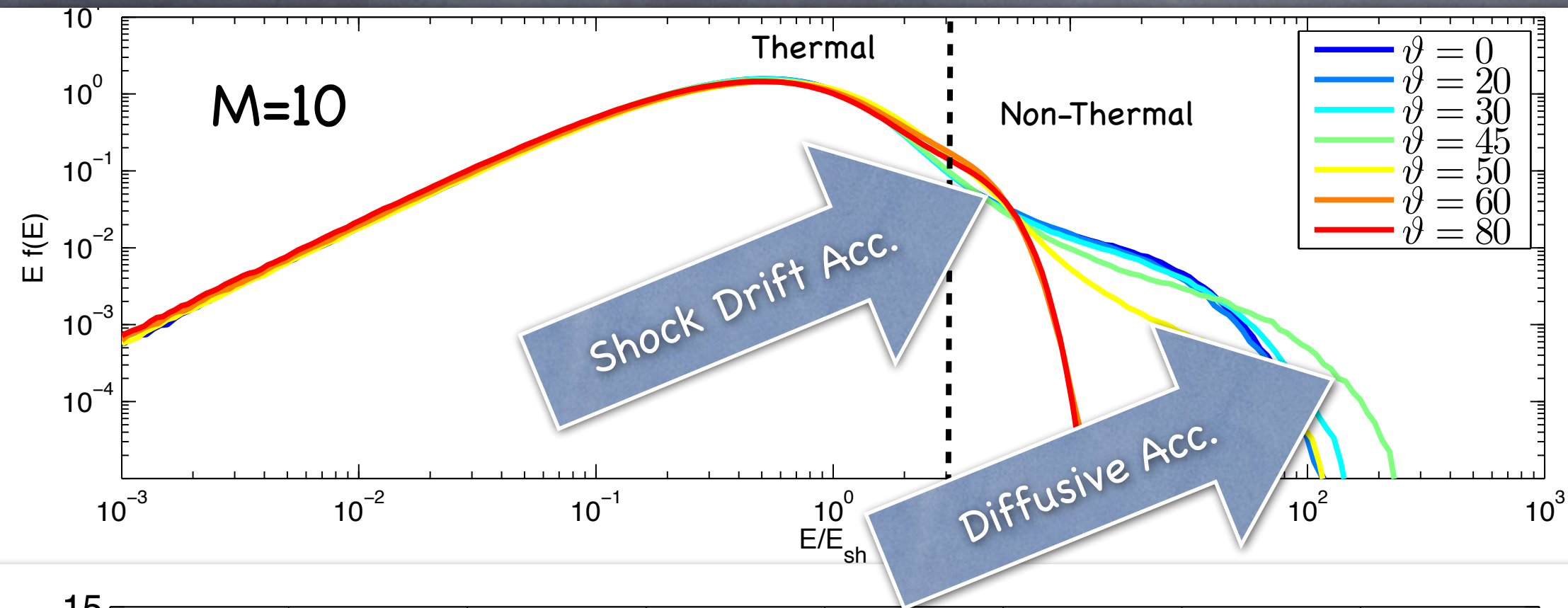


3D simulations of a parallel shock



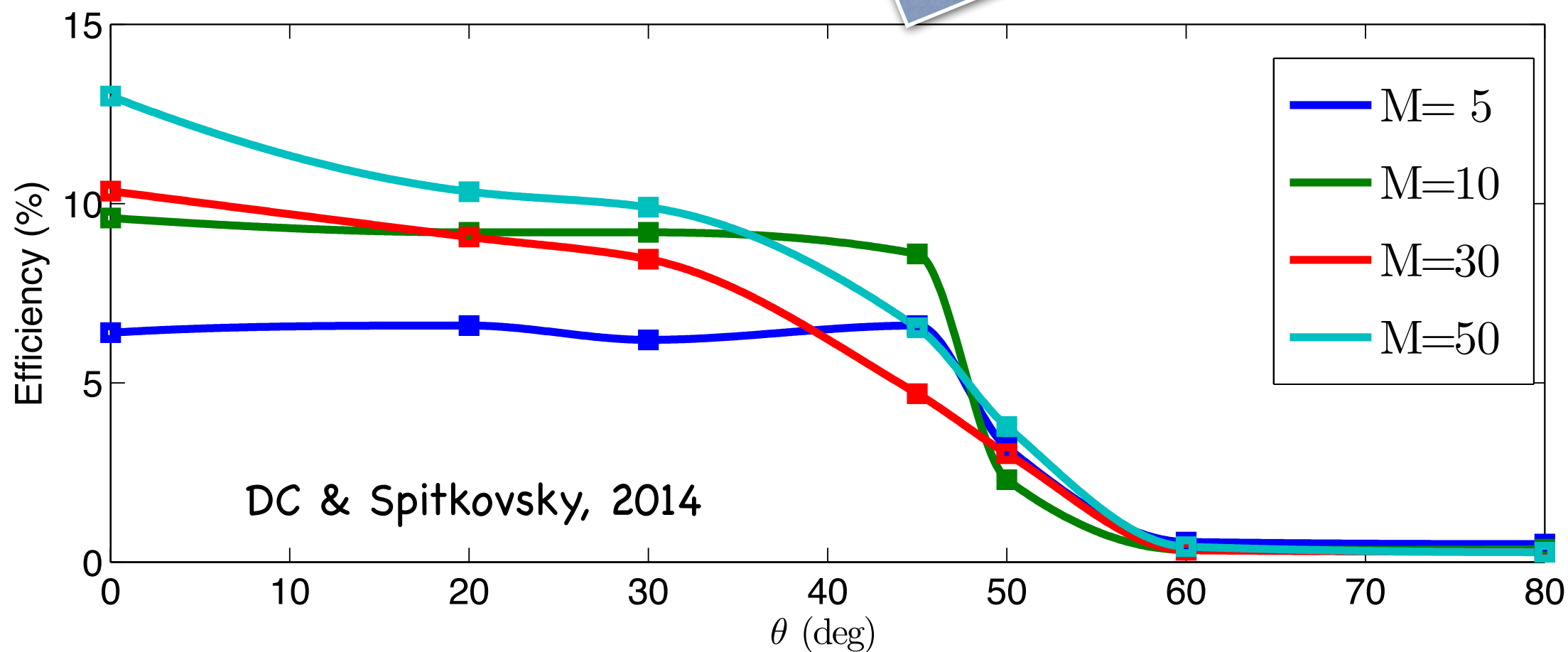
DC &
Spitkovsky,
2014a

Parallel vs Oblique shocks

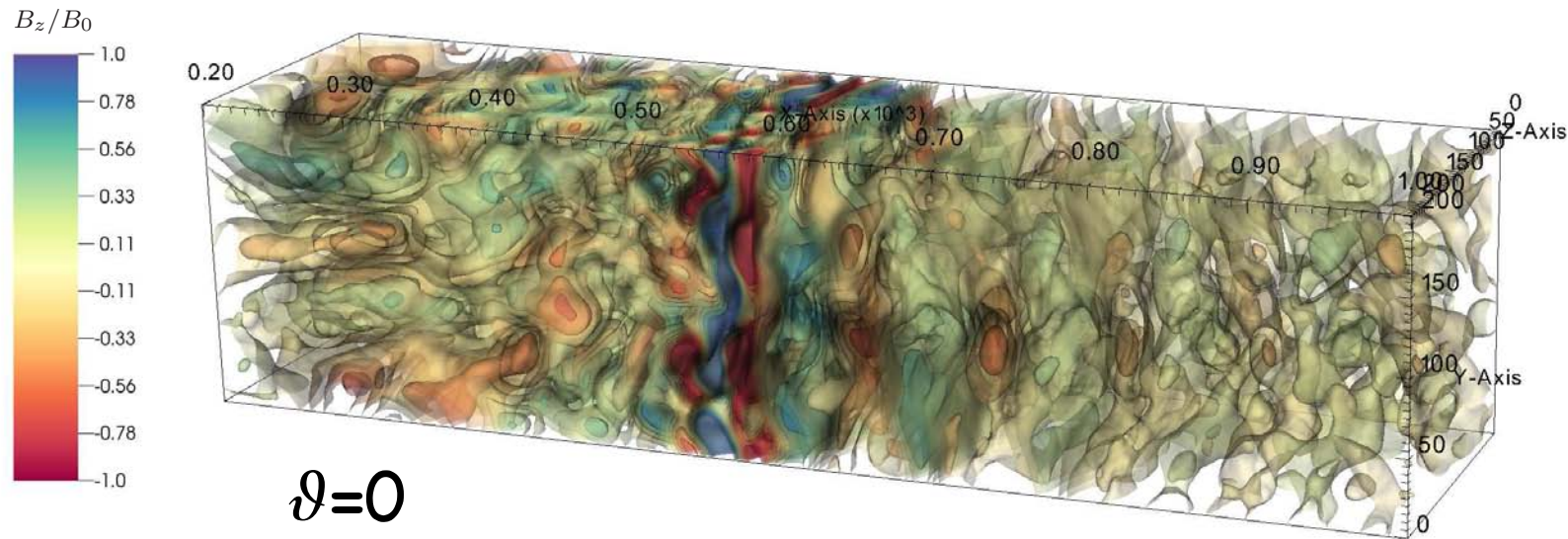


Each point corresponds to a simulation with about 10^9 particles

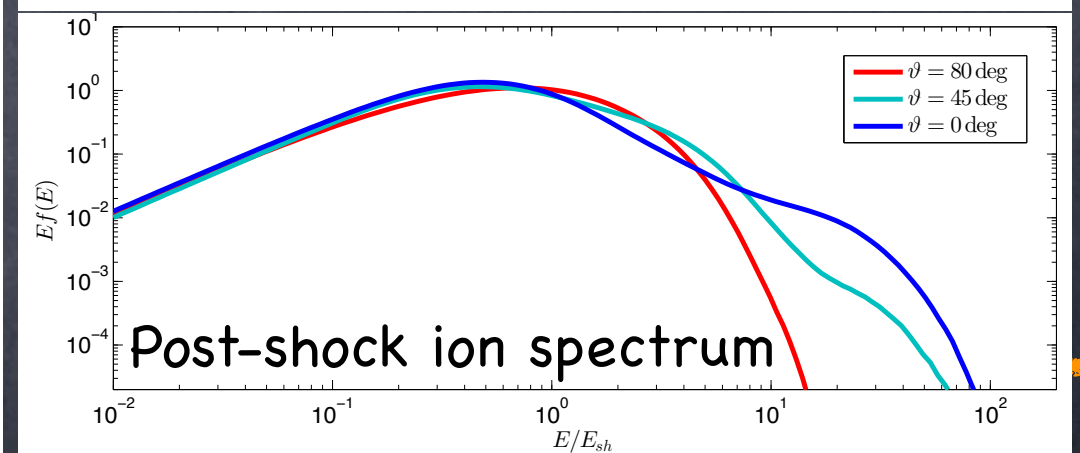
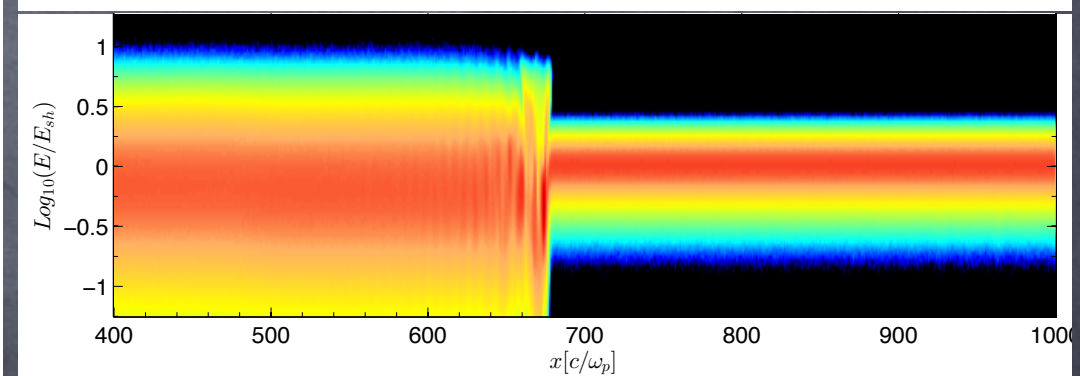
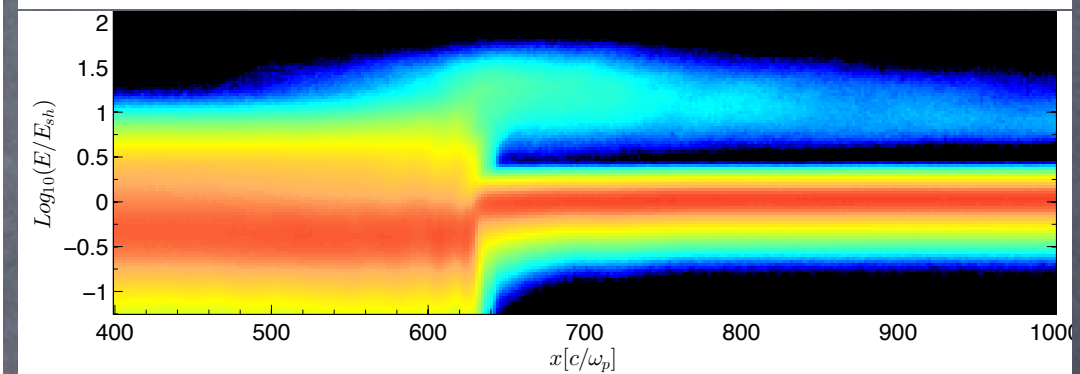
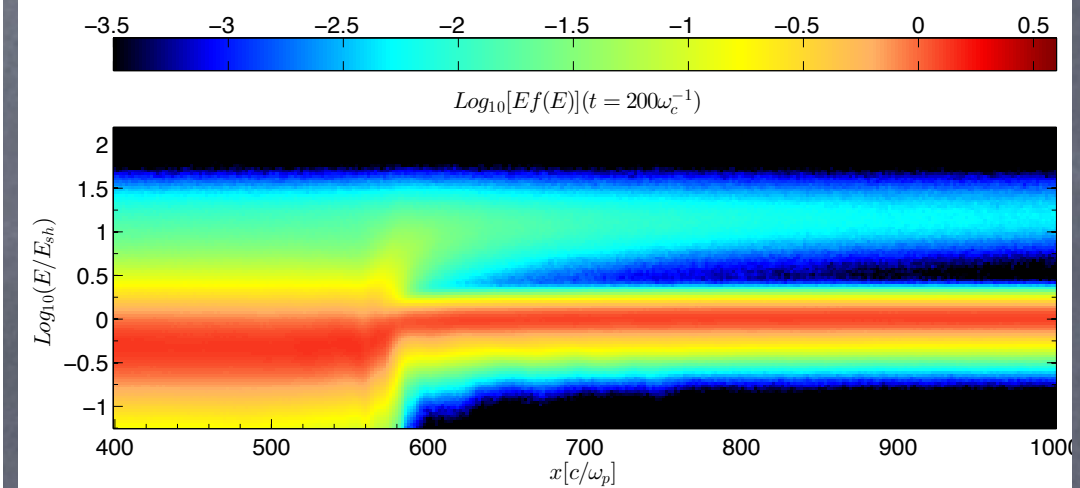
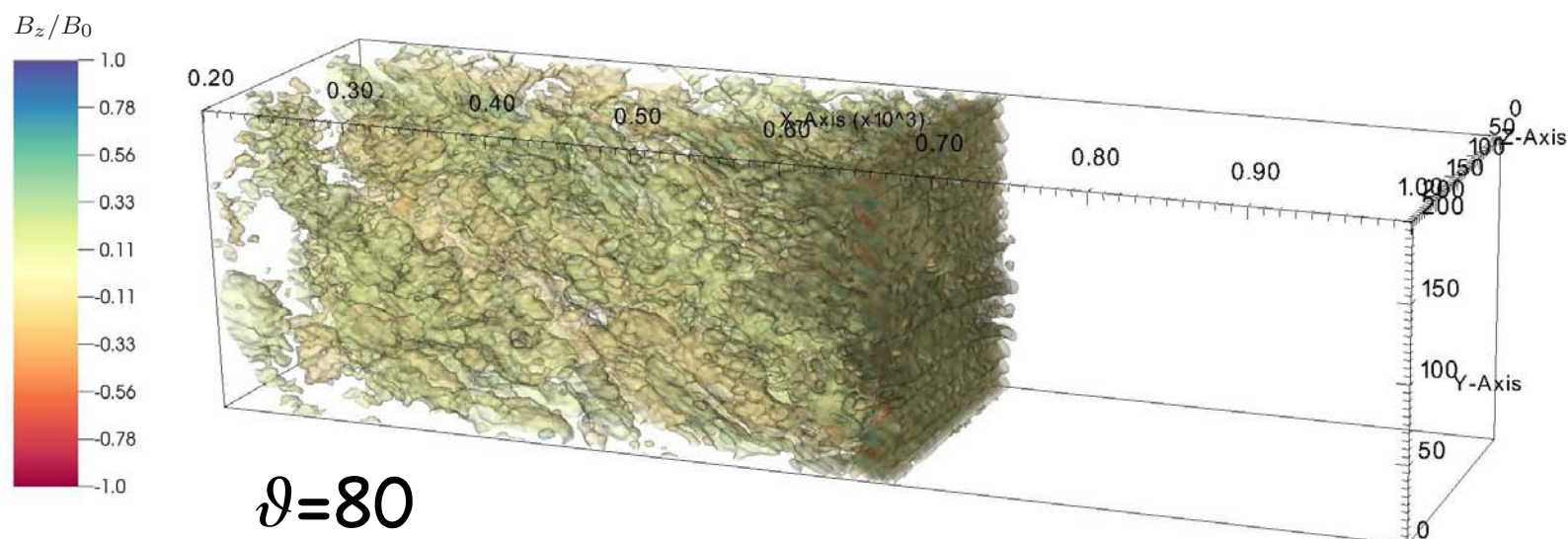
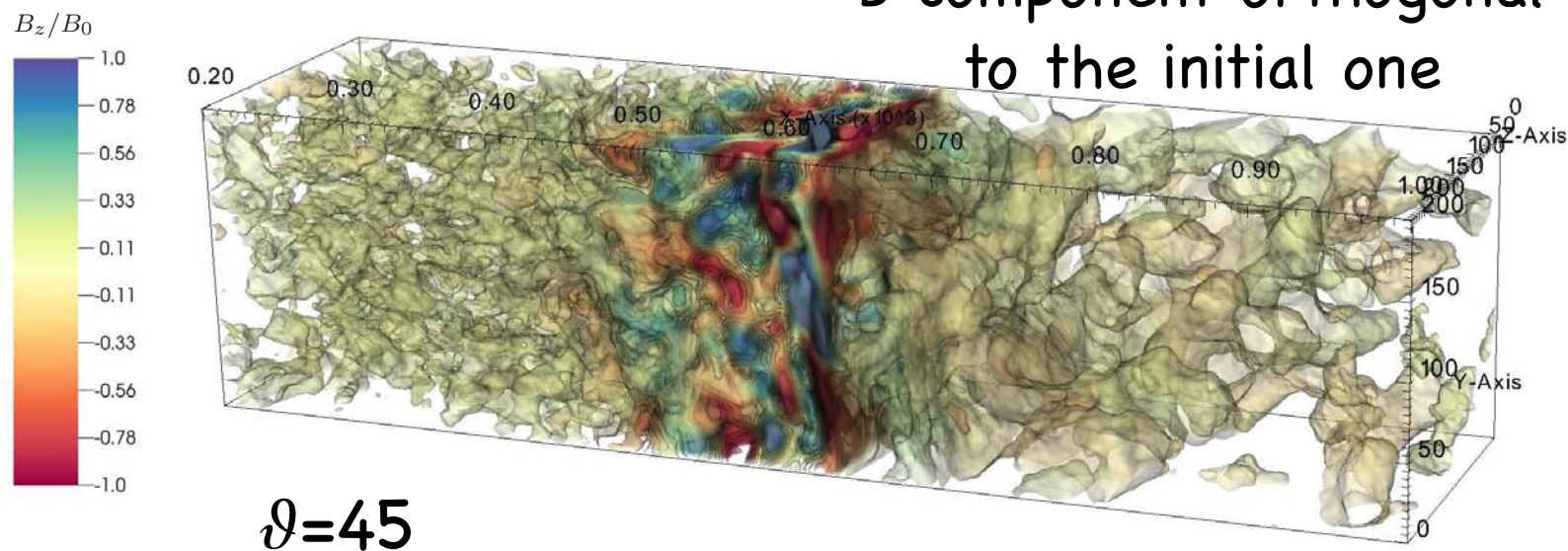
Computation time: almost 2×10^6 cpu h



3D simulations



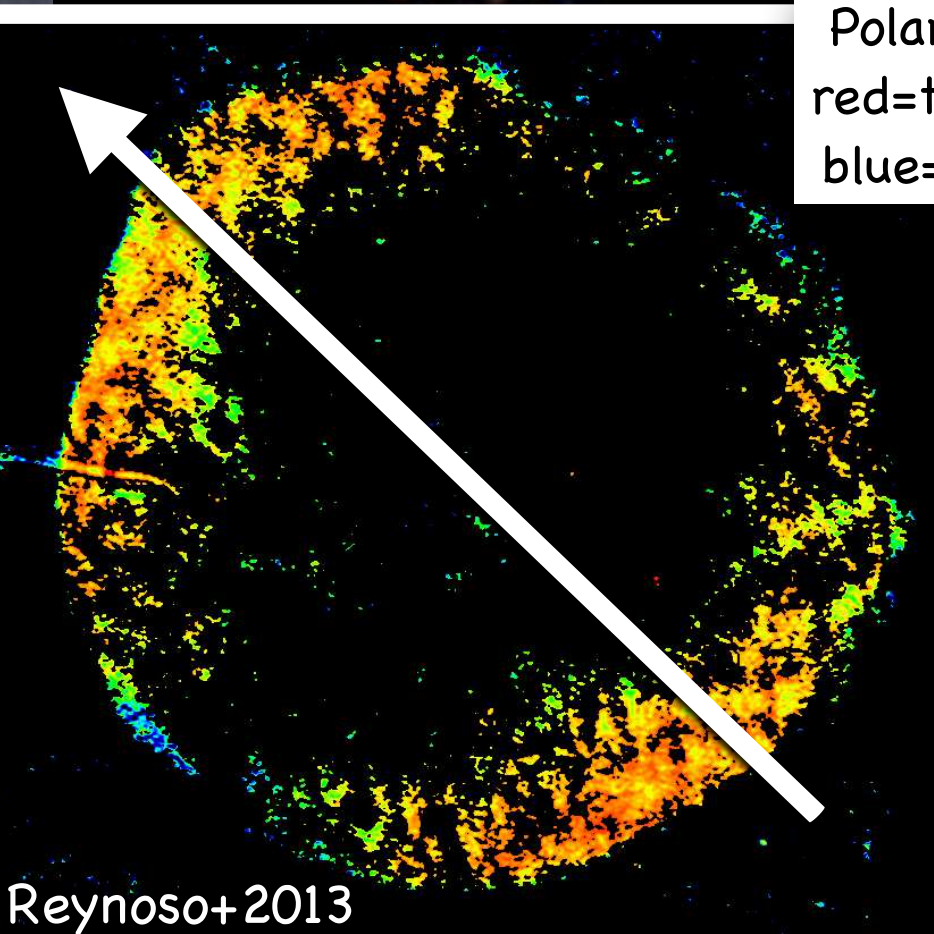
B component orthogonal
to the initial one



SN 1006: a parallel accelerator

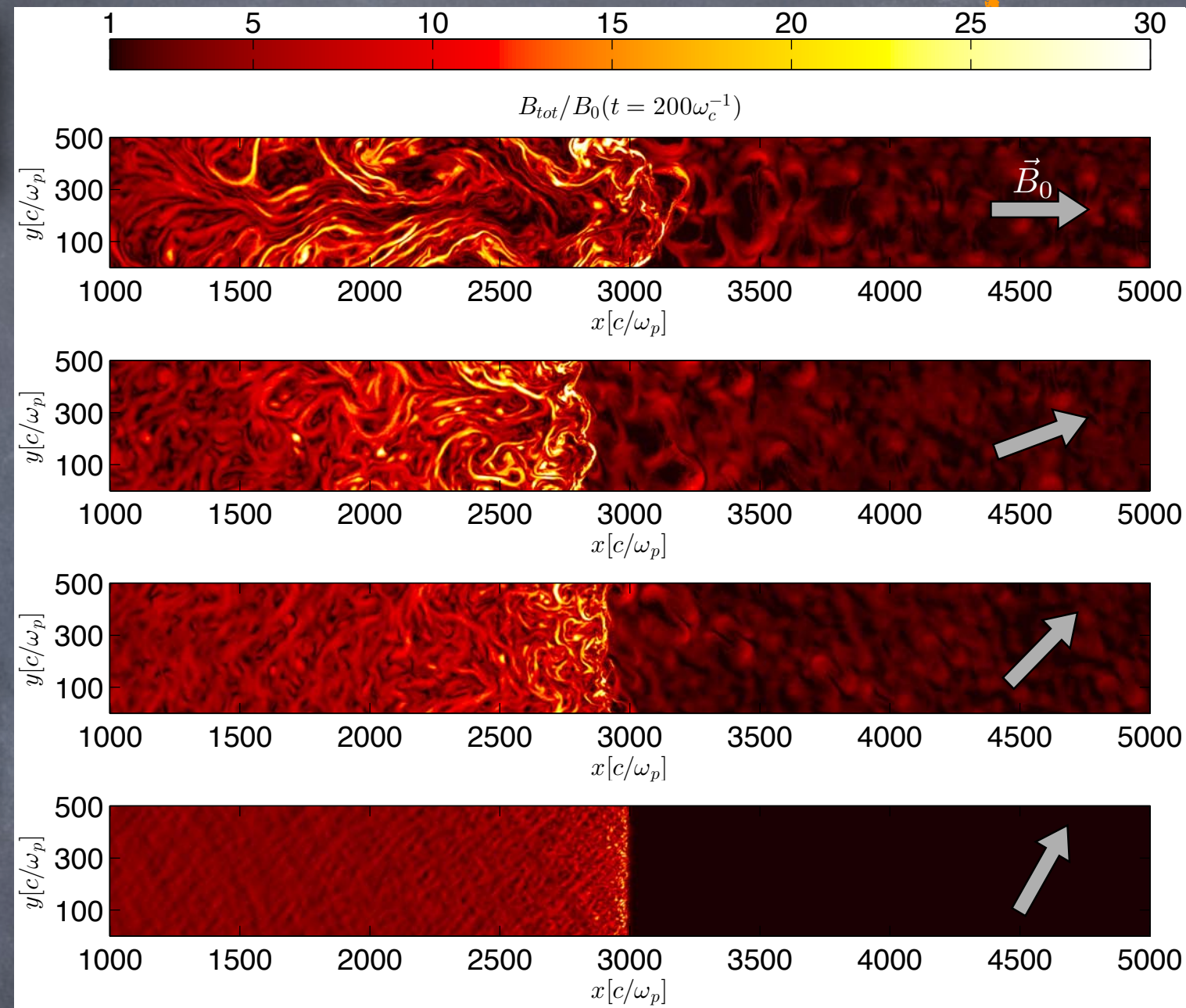


X-ray emission:
red=thermal
white=synchrotron



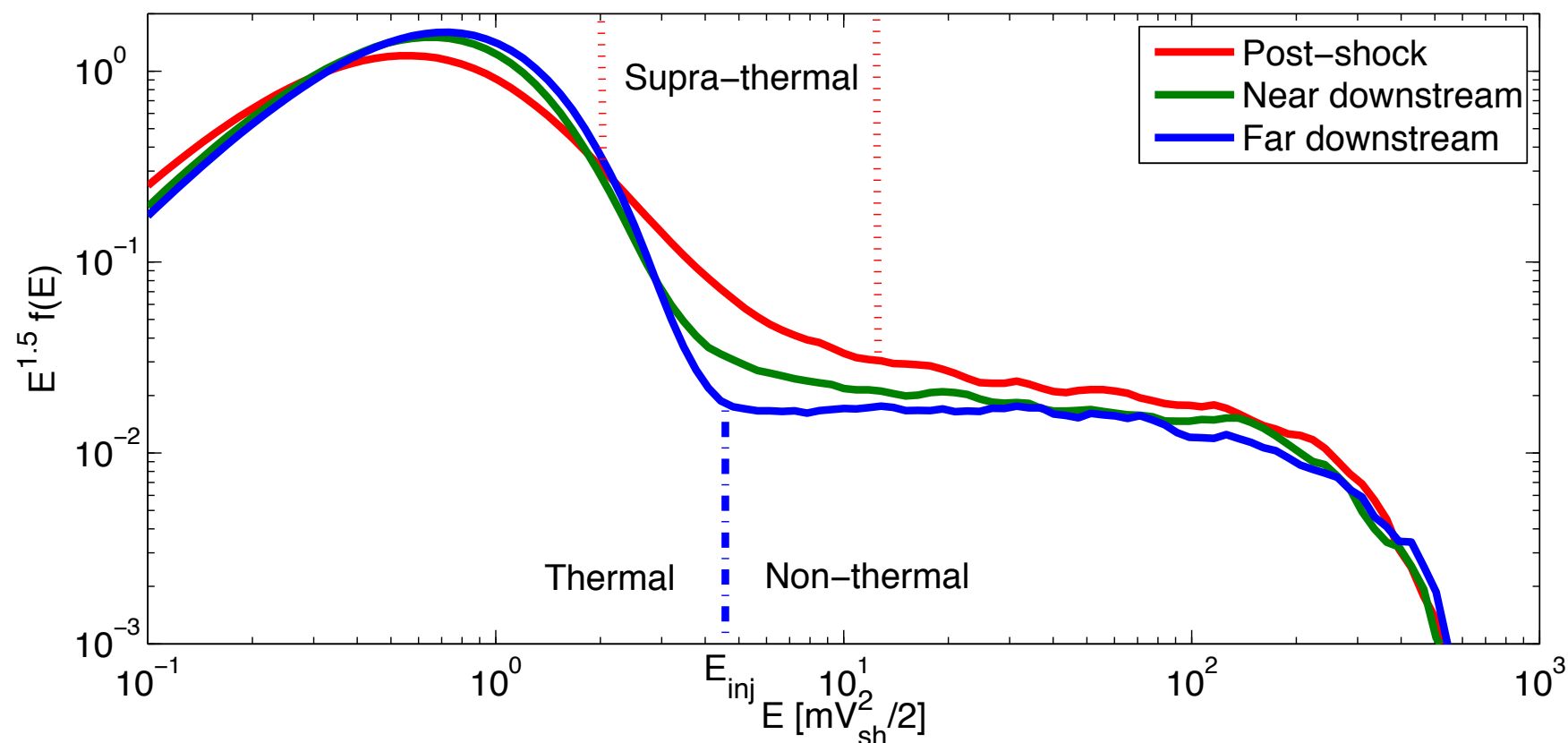
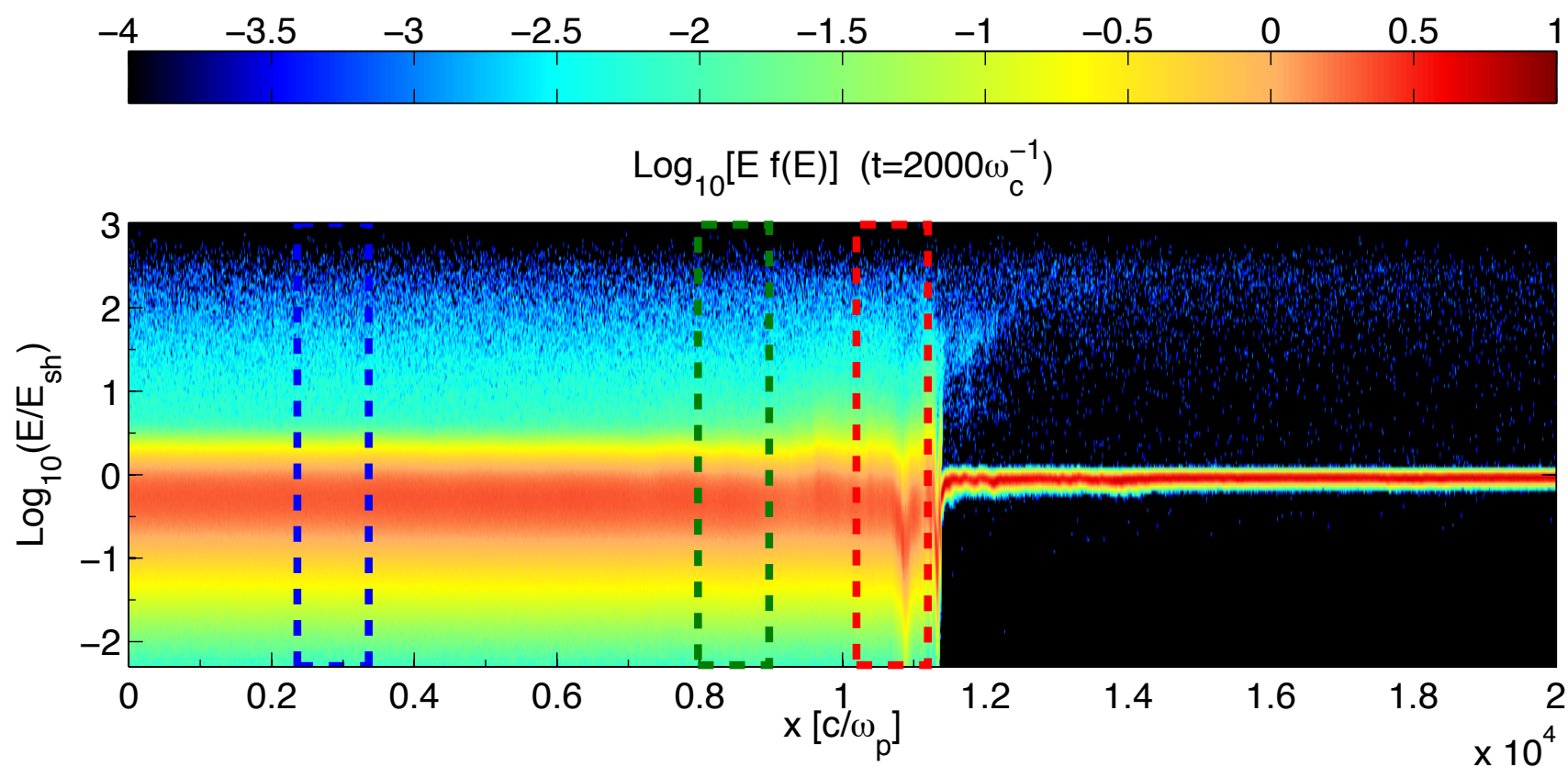
Polarization:
red=turbulent
blue=ordered

Reynoso+2013



Ion acceleration and B
field amplification where
the shock is **parallel**

Supra-thermal ions

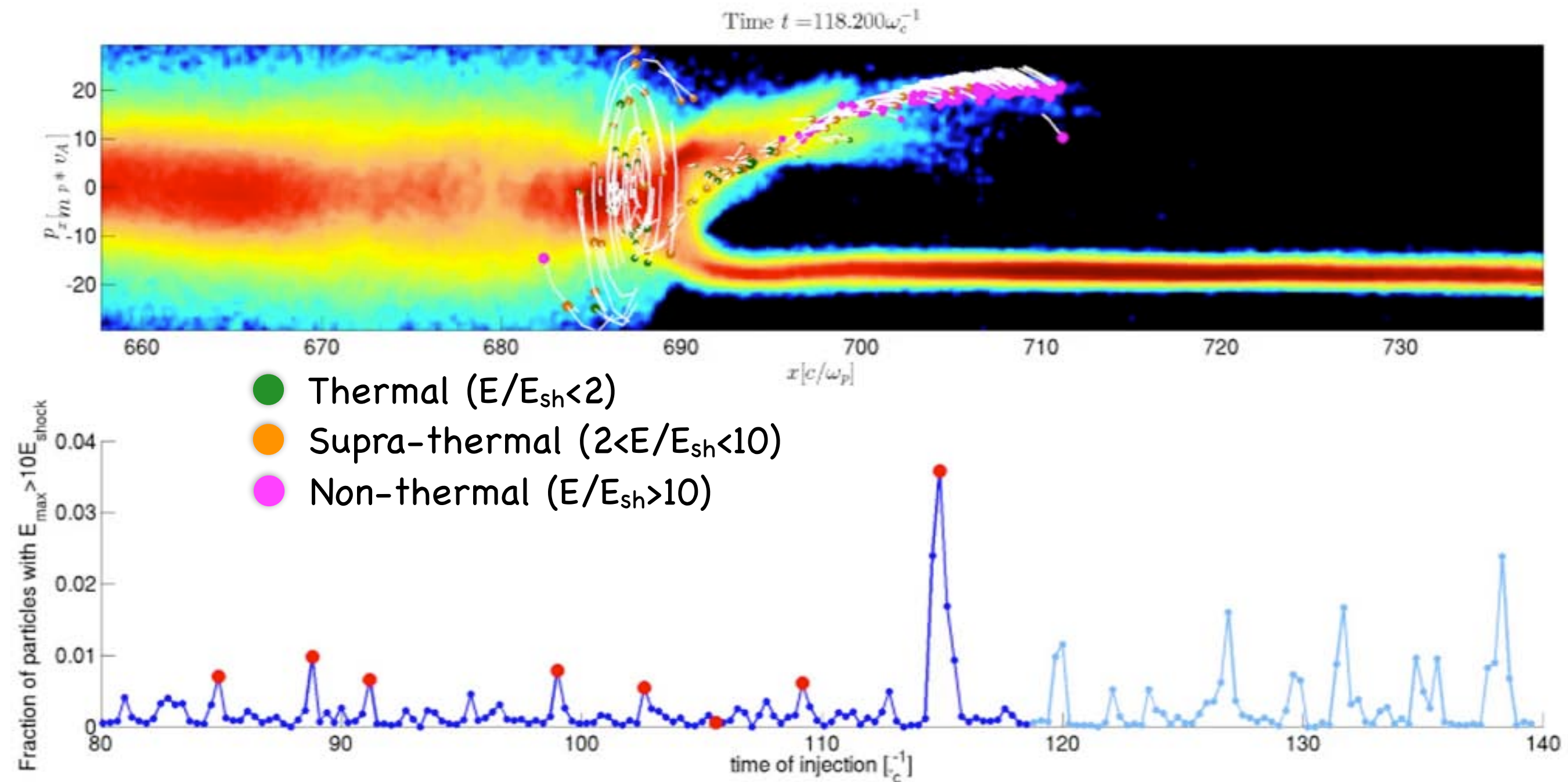


- Steep "bridge" immediately behind the shock
- Contains information on injection and thermalization
- The DSA power-law starts at $p_{inj} \sim 3-4 p_{th,d}$

Particle Injection – Simulations



DC, Pop & Spitkovsky, 2015





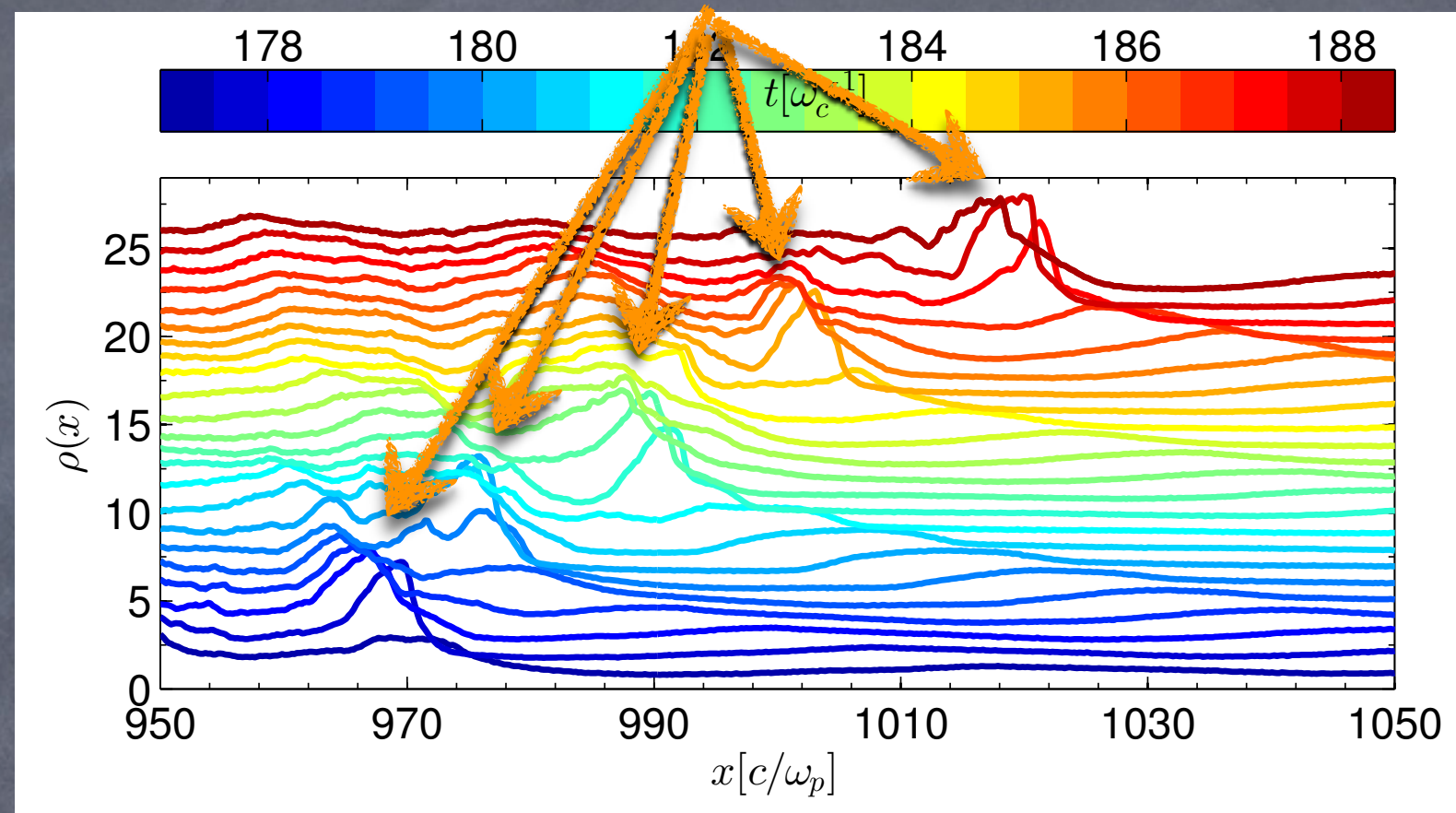
The shock potential barrier

- Quasi-parallel shocks exhibit a **time-varying potential barrier**, with an overshoot

- The jump in the electron pressure produces an **electric field E_x directed upstream** (microscopically due to the charge separation induced by different gyroradii of e and p)

$$E_x \simeq -\frac{\Delta\Phi}{\Delta x} \simeq -\frac{1}{ne} \frac{\Delta P_e}{\Delta x}$$

Overshoot



- Ions impinging on the shock barrier can be **reflected** back because of such a E_x (also, because of the compressed B_{trans})
- Barrier stationary in the **downstream frame!**

Encounter with the shock barrier

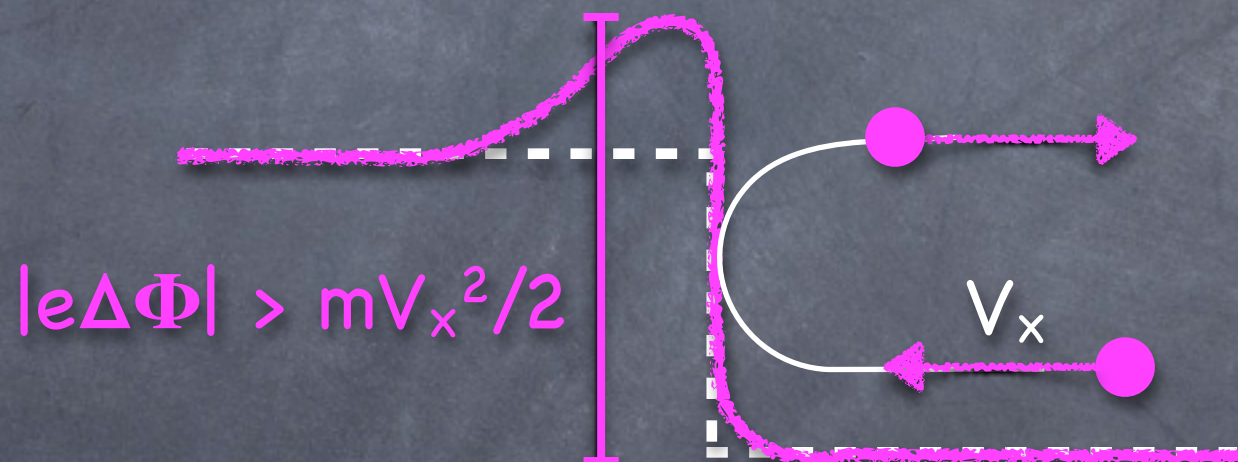


- Low barrier (shock reforming)



Particles are advected downstream, and **thermalized**

- High barrier (overshoot)



Particles are **reflected** upstream, and **energized** via Shock Drift Acc.

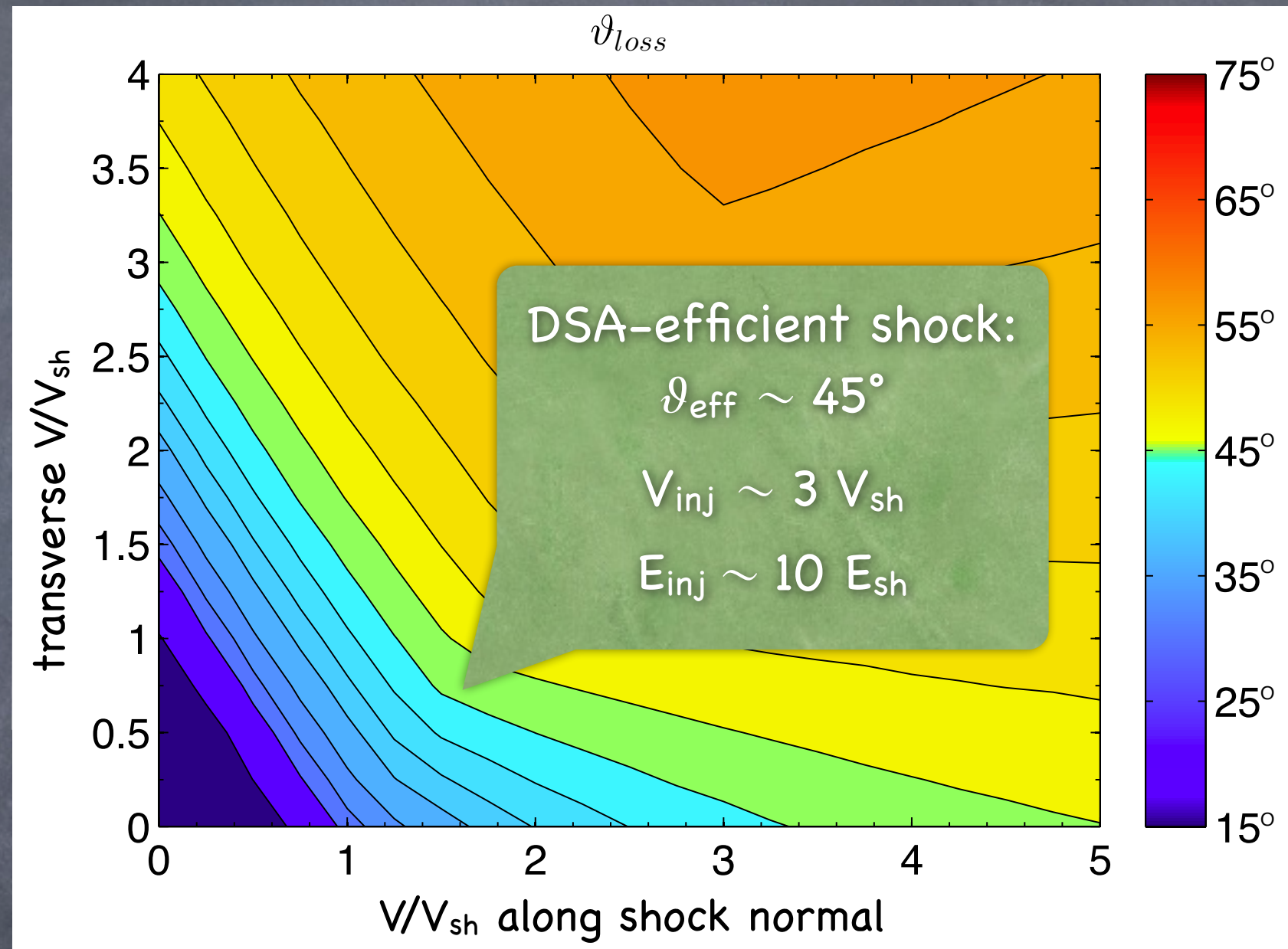
- Coherent reflection induces **shock reformation**
- Acceleration because of a transverse **(motional) E field**
- What is the particle **fate** after reflection?

Ion Injection – Theory



Max ϑ allowing reflection upstream

- Ion fate determined by
 - barrier duty-cycle ($\sim 25\%$)
 - pre-reflection V
 - shock inclination
- If $\vartheta < \vartheta_{\text{loss}}$, ions escape upstream, and are injected into DSA
- Otherwise, they experience SDA, return to the shock (with larger V), and may be either reflected or advected
 - After N SDA cycles, only a fraction $\eta \sim 0.25^N$ survives
 - For $\vartheta_{\text{eff}} \sim 45^\circ$, $N \sim 3 \rightarrow \eta \sim 1\%$



DC, Pop & Spitkovsky, 2015

E_{inj} is larger at oblique shocks:
injection requires more SDA cycles,
and fewer particles can achieve E_{inj}

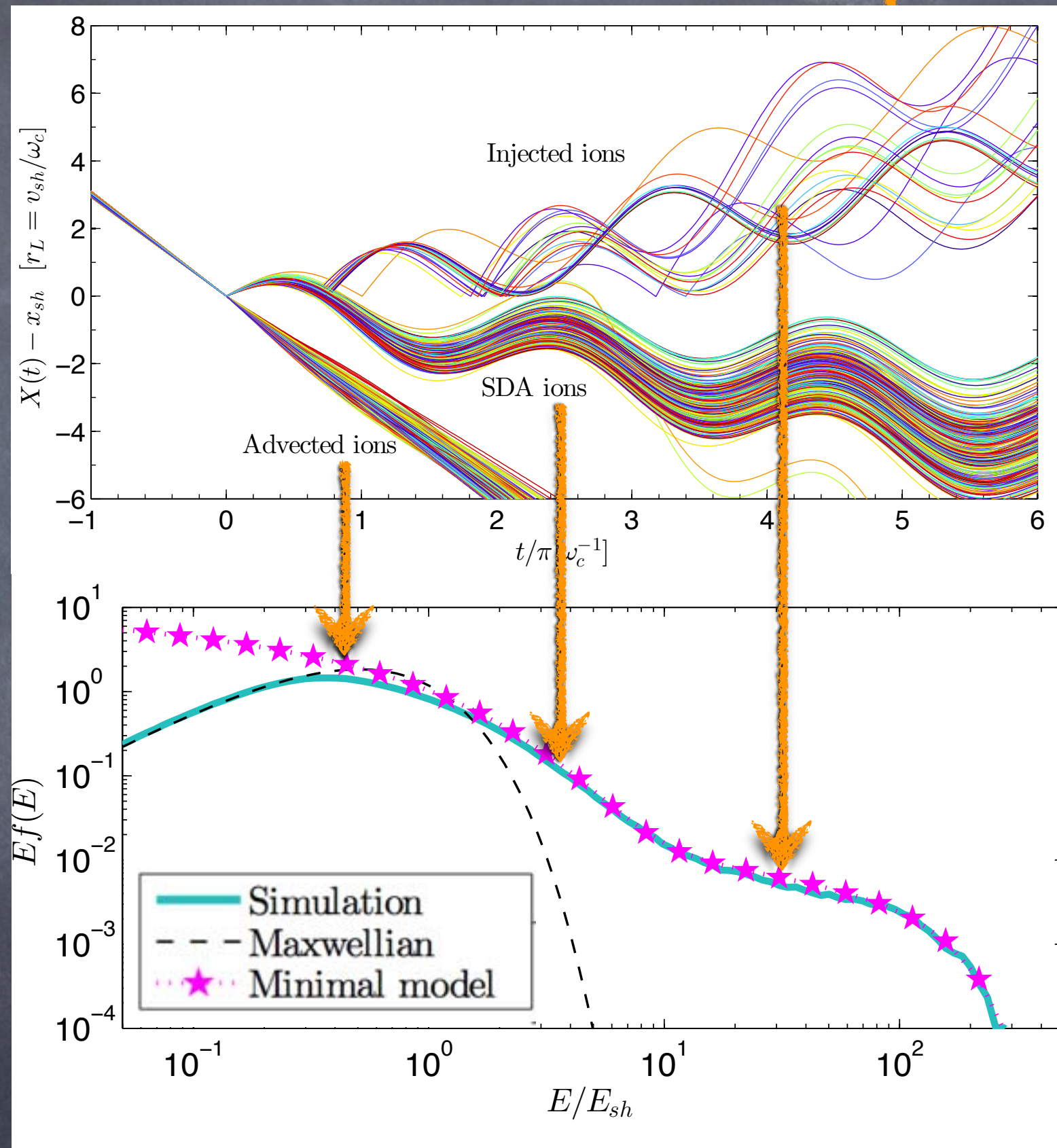
Minimal Model for Ion Injection



- Time-varying potential barrier
 - High state (duty cycle 25%)
 - Reflection
 - Shock Drift Acceleration
 - Low-state → Thermalization
- Spectrum à la Bell (1978)

$$f(E) \propto E^{-1-\gamma}; \quad \gamma \equiv -\frac{\ln(1 - \mathcal{P})}{\ln(1 + \mathcal{E})}$$

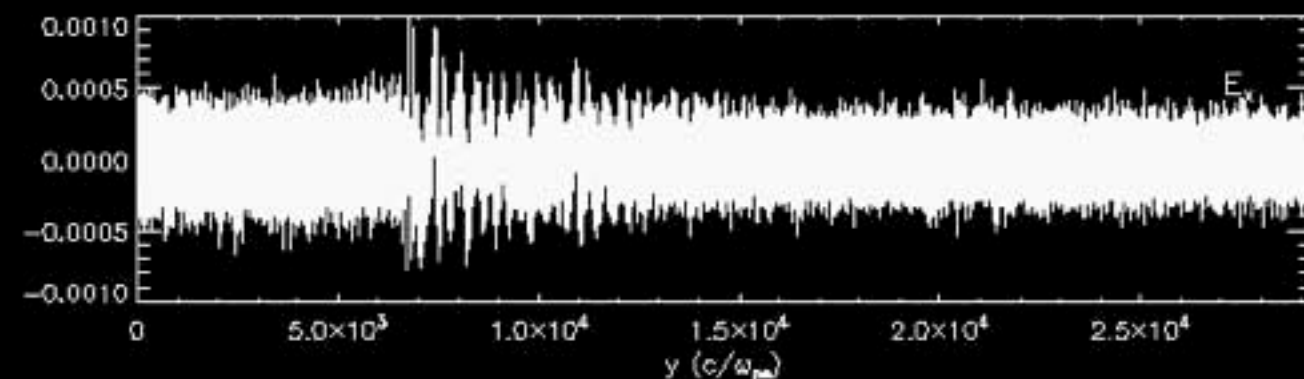
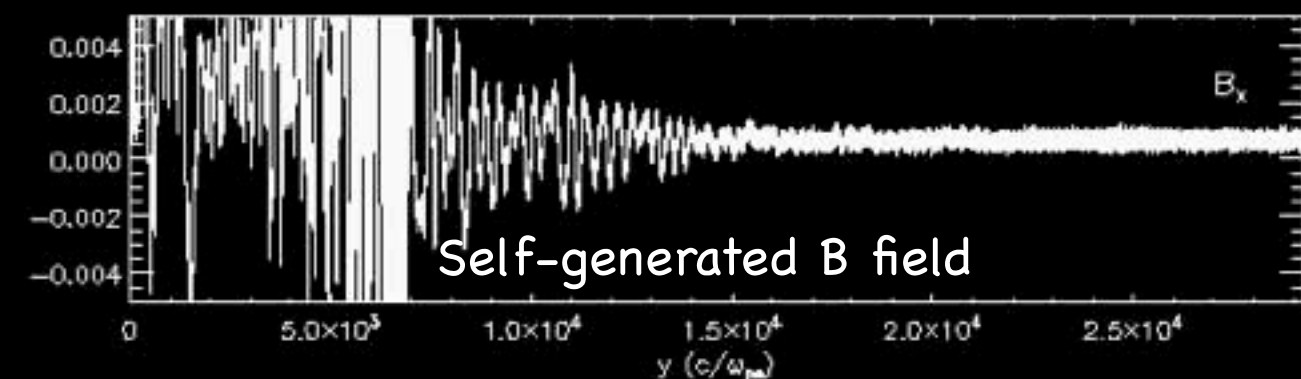
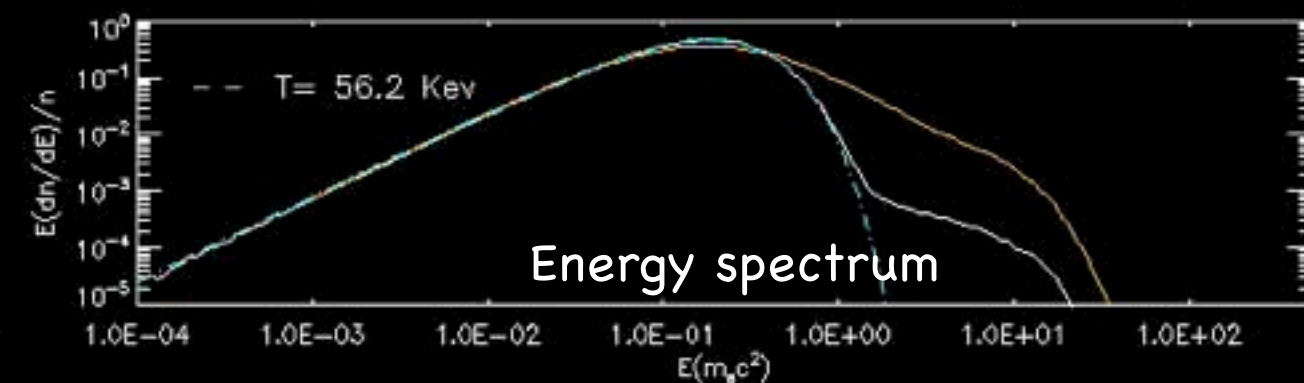
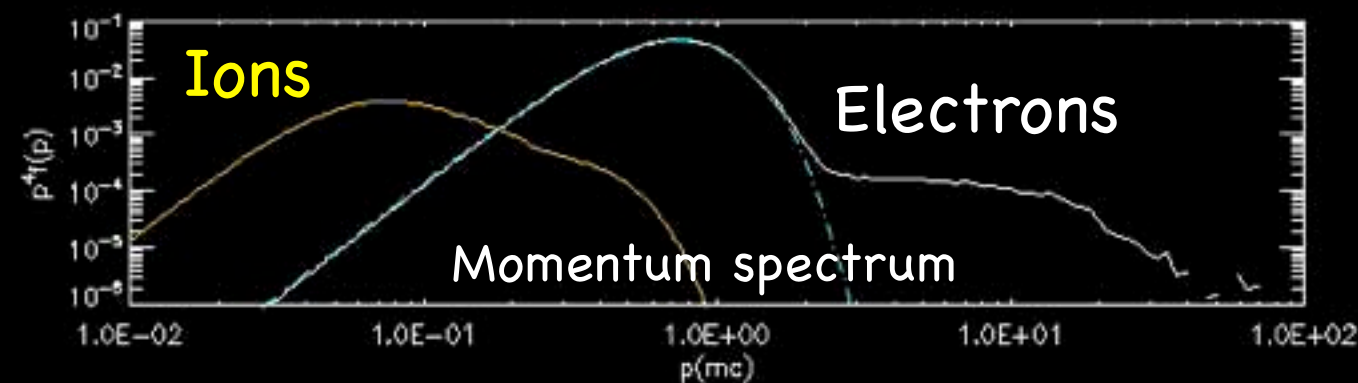
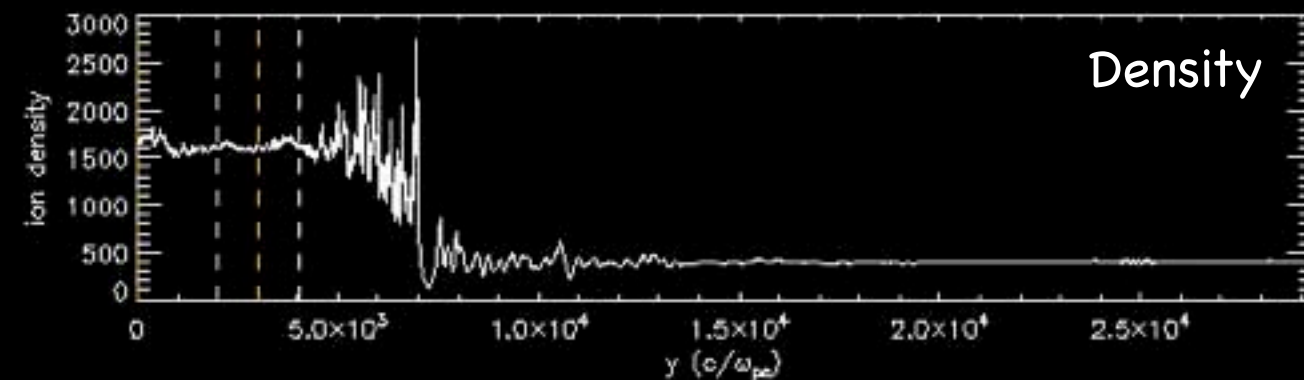
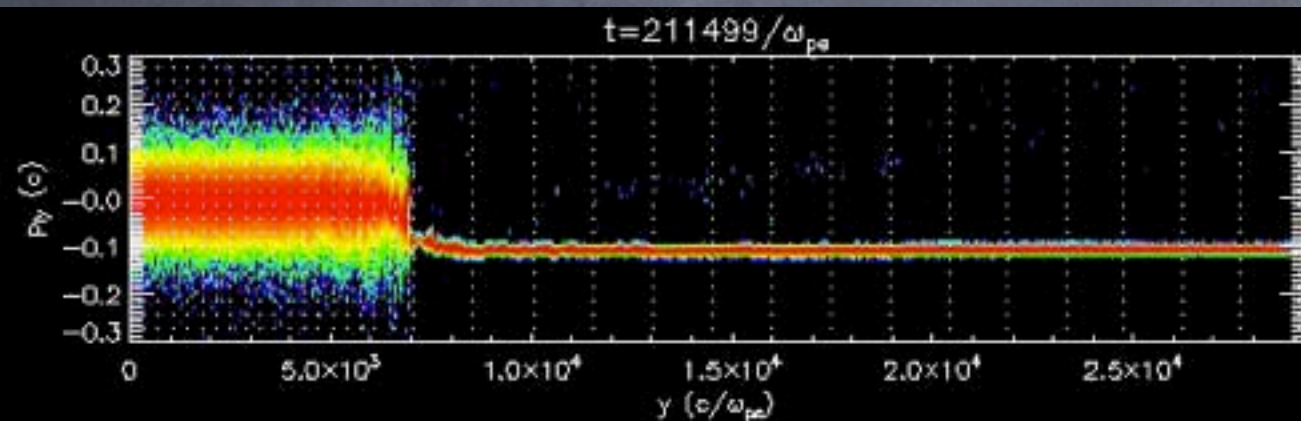
 - $\mathcal{P}$ = probability of being advected
 - $\mathcal{E}$ = fractional energy gain/cycle



Electron/ion acceleration



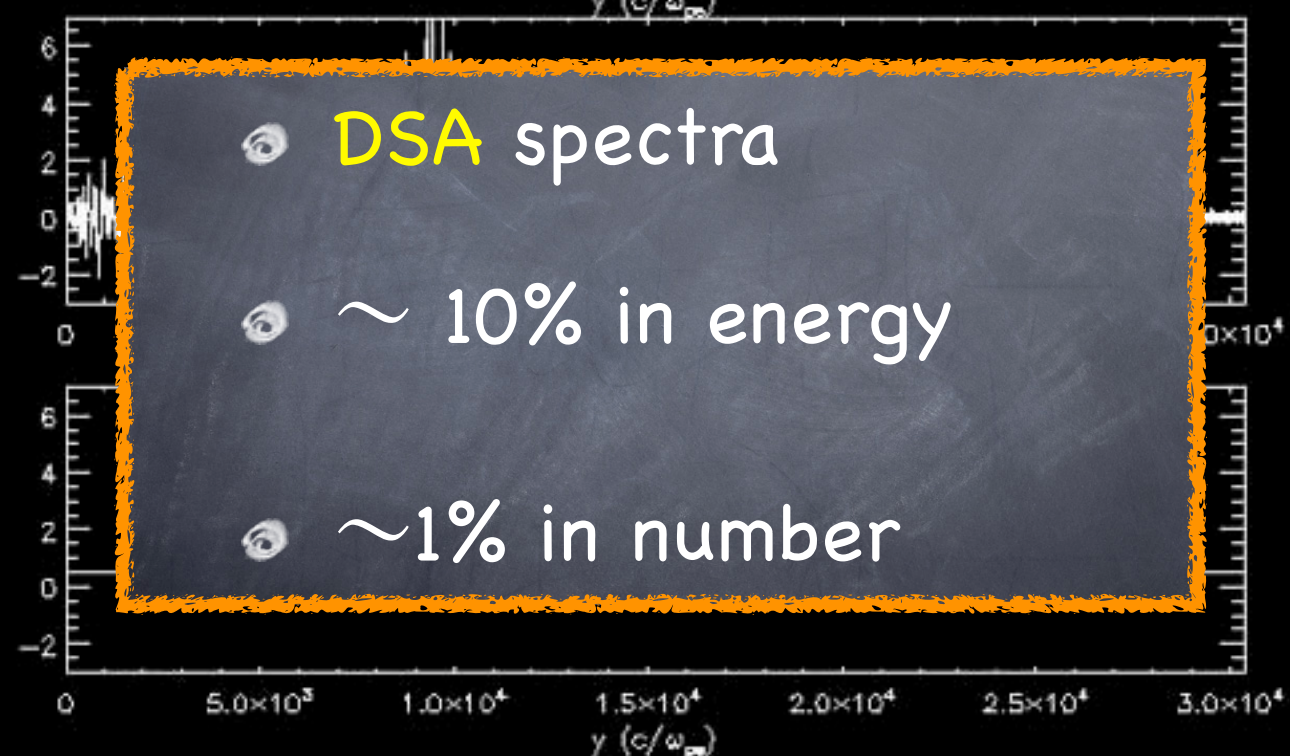
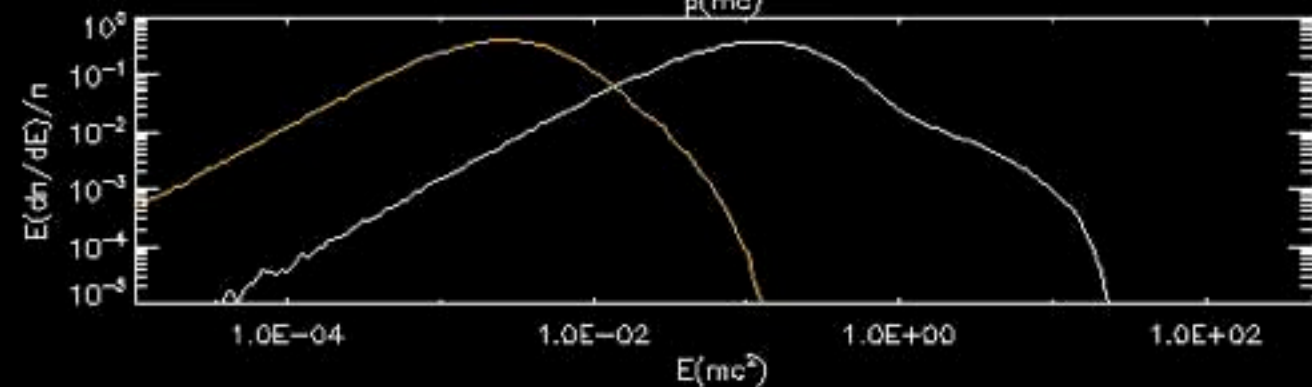
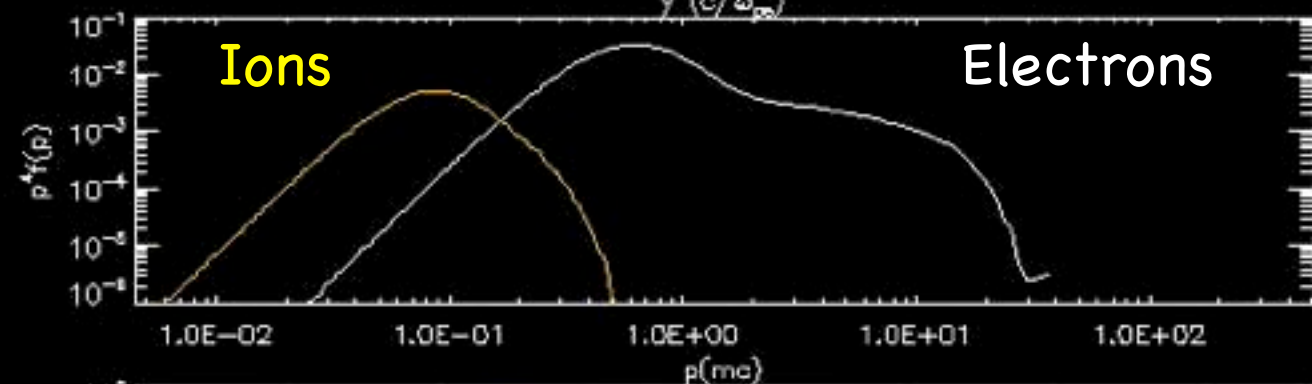
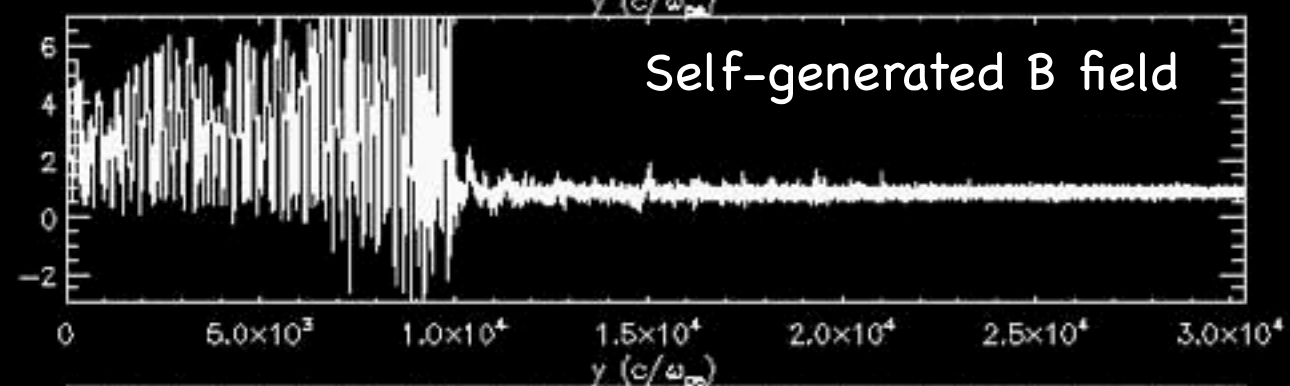
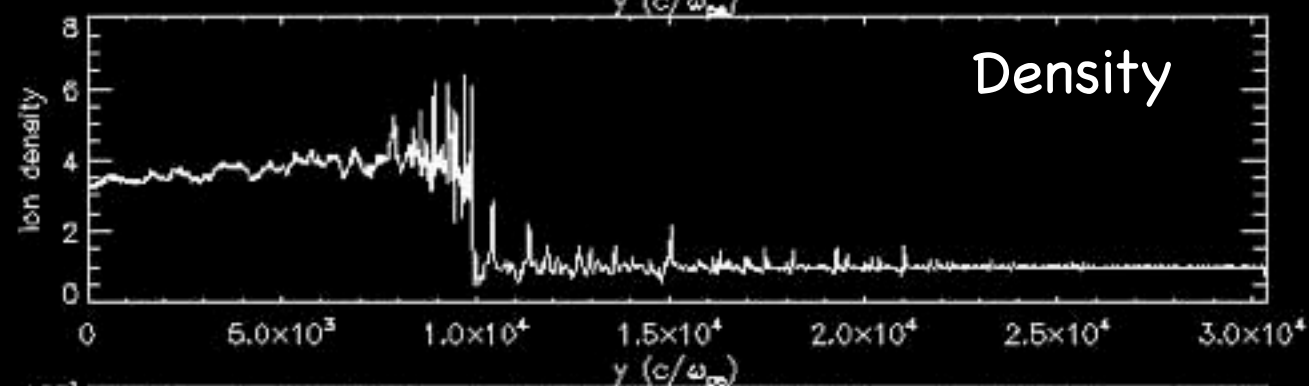
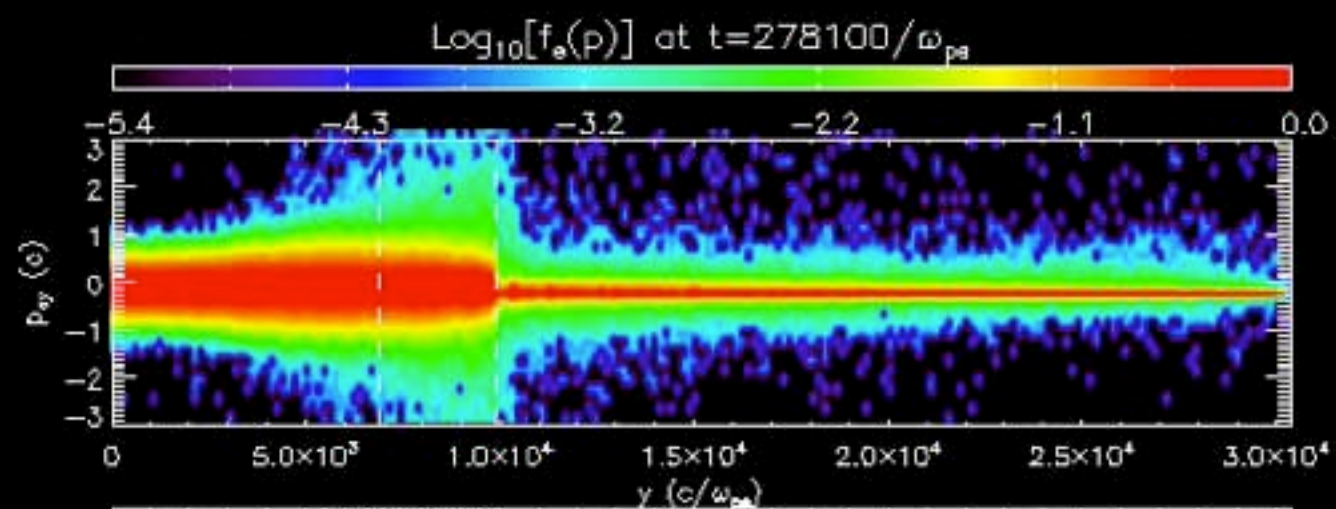
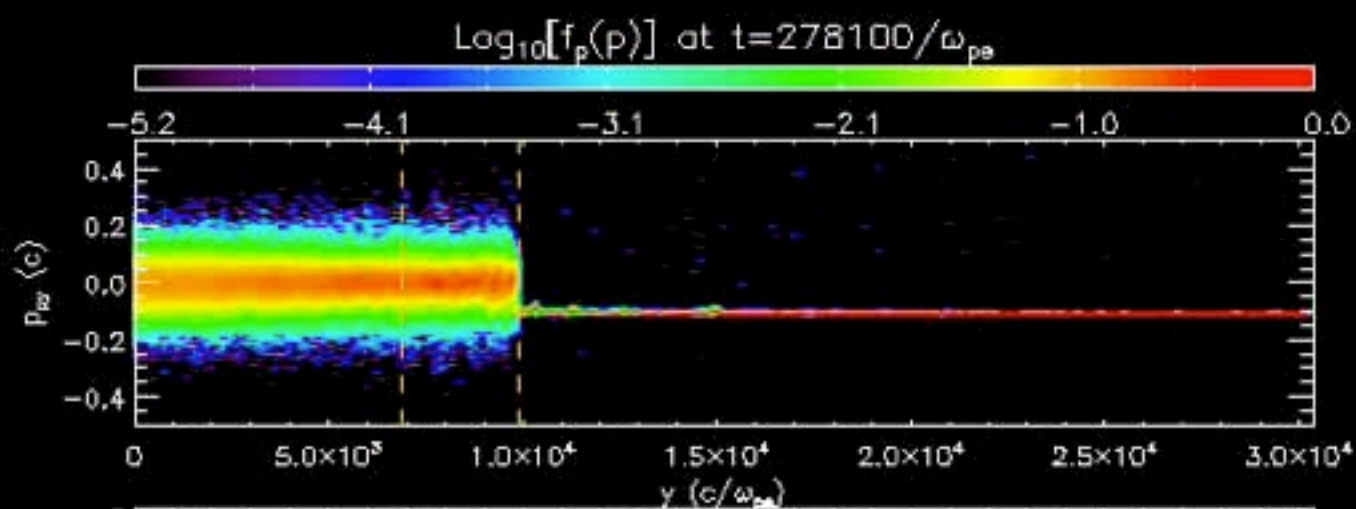
- Full **PIC simulations**: Tristan-MP (Park, DC, Spitkovsky 2015, PRL)
 - $M=20$, $v_{sh}=0.1c$, **quasi-parallel** shock
 - Electrons are accelerated, but **ele/proton ratio is a few %**



Electron acceleration



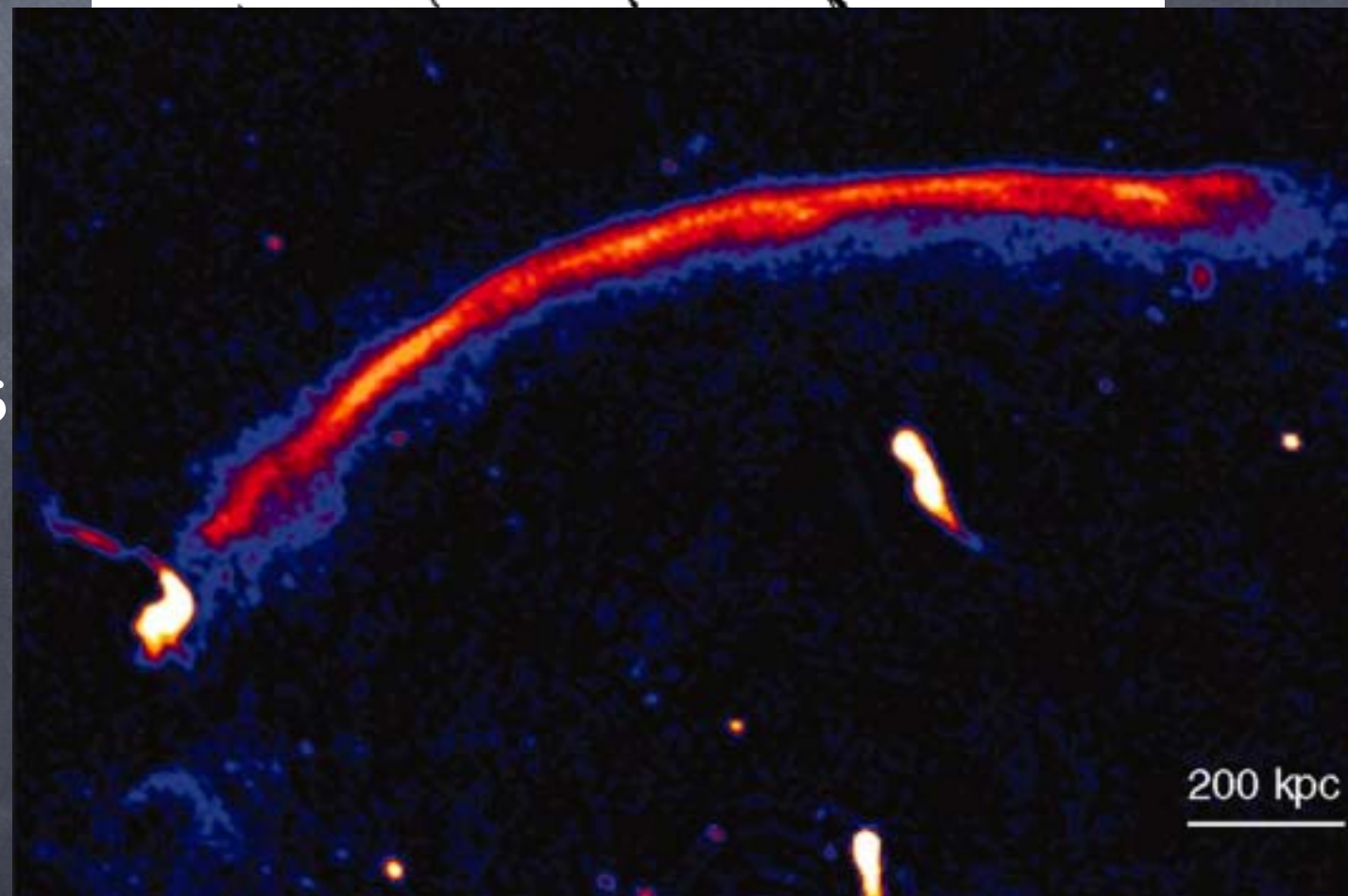
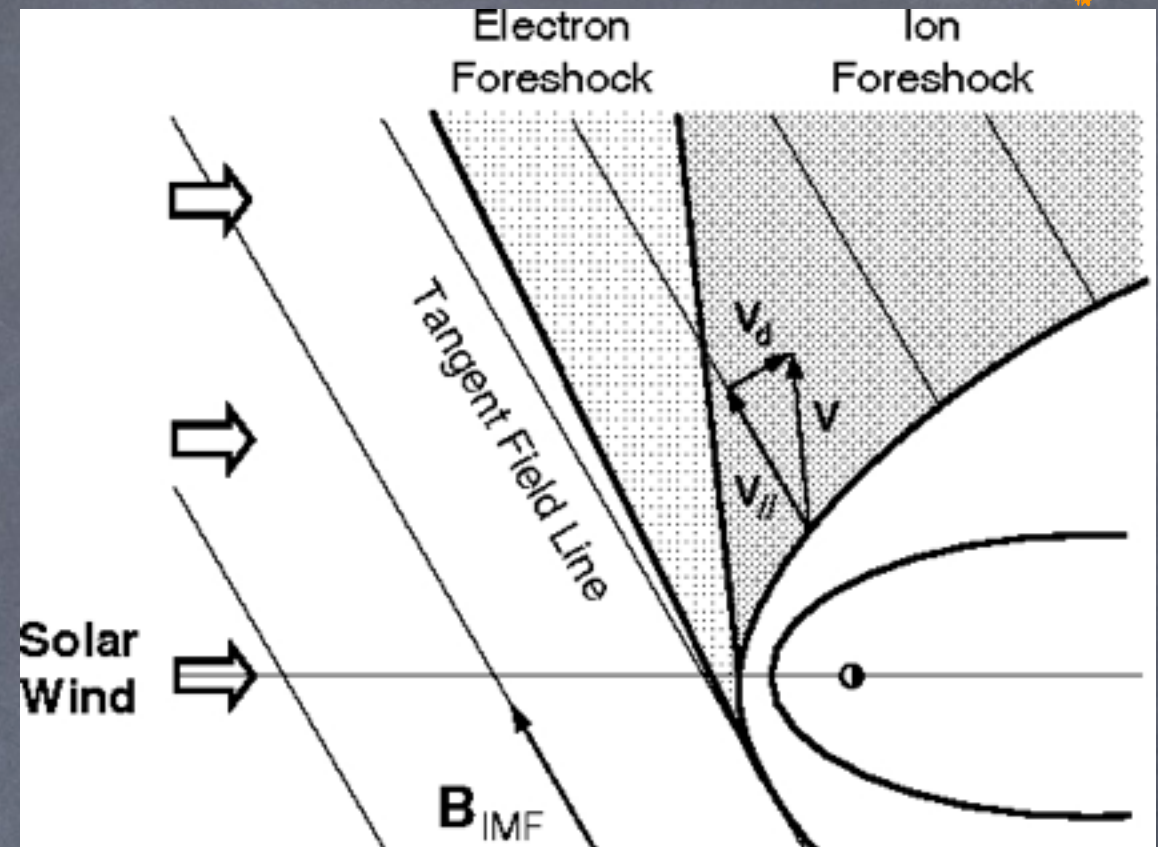
- PIC simulations of **oblique shocks** ($\vartheta=60^\circ$): DC, Park, Spitkovsky, in prep.



Electron acceleration



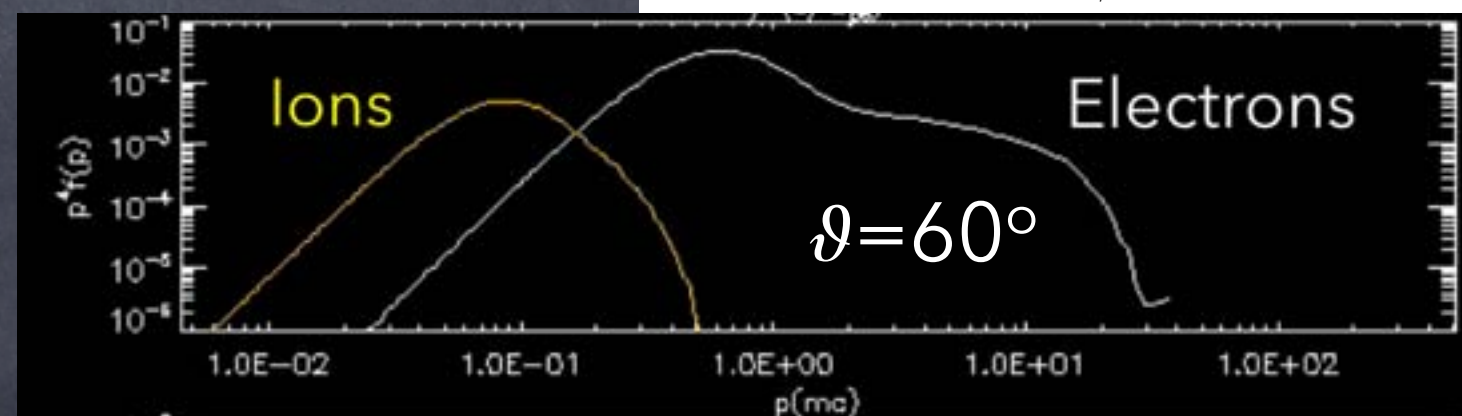
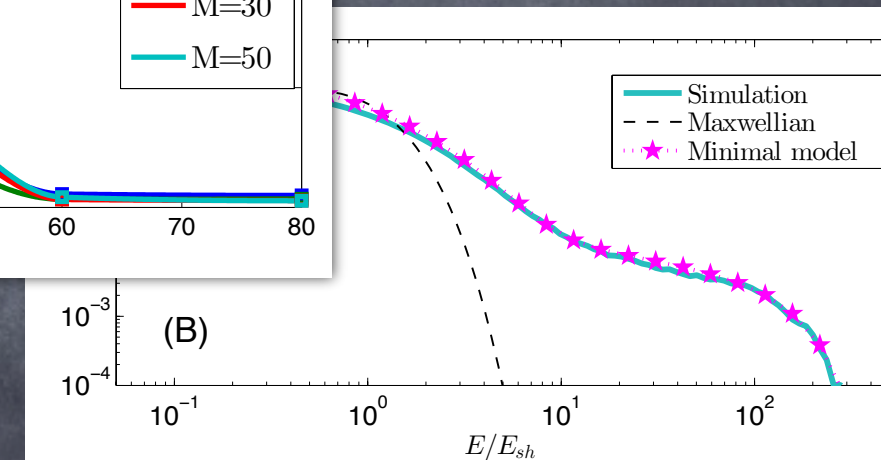
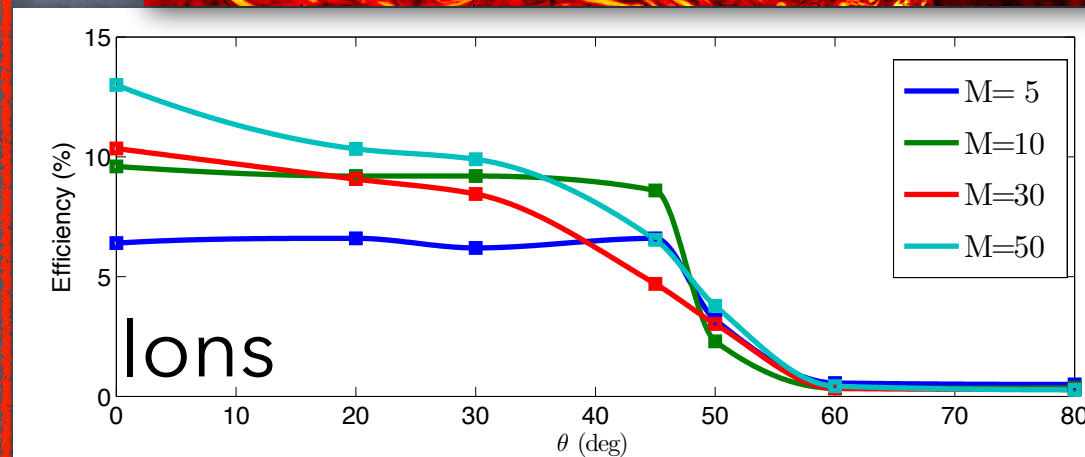
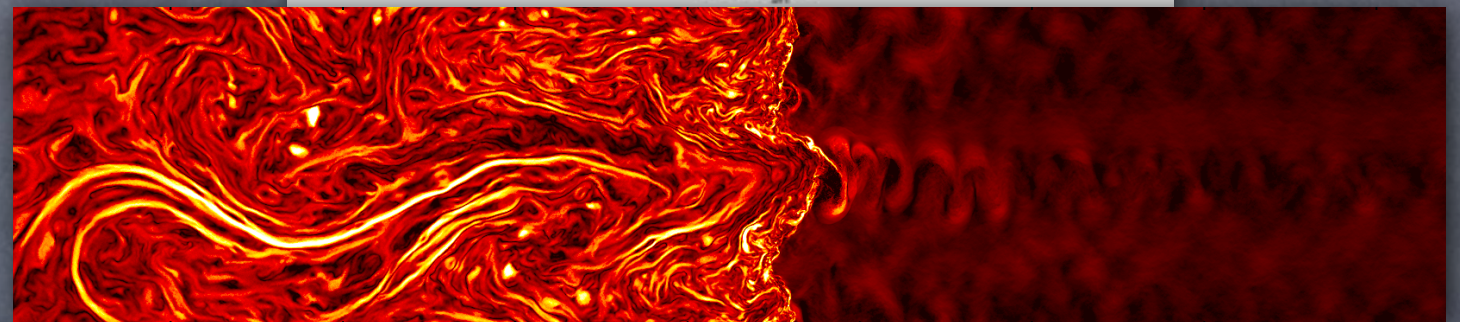
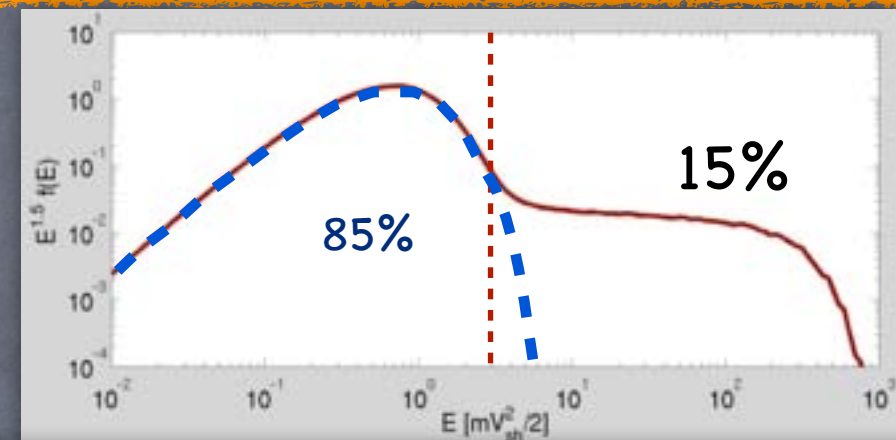
- **Planetary bow shocks**
(Earth, Venus, Saturn,...)
- **In situ** measurements:
Geotail, Venus Orbiter, Polar,
SoHO, WIND, Cluster, ...
- **Radio relics** in **galaxy clusters**
(sausage, toothbrush,...)
- Extended polarized structures
- **Fermi-LAT** limits on γ -ray
emission: constrain **e/p** ratio!



Conclusions



- Acceleration at **shocks** can be **efficient**: >15%
- CRs amplify B field via **streaming instability**
- Ion DSA efficient** at **parallel**, strong shocks
- Ions are **injected** via reflection and shock drift acceleration
- Electron DSA efficient** at **oblique** shocks



Exchanges between Fluid and Kinetic Scales

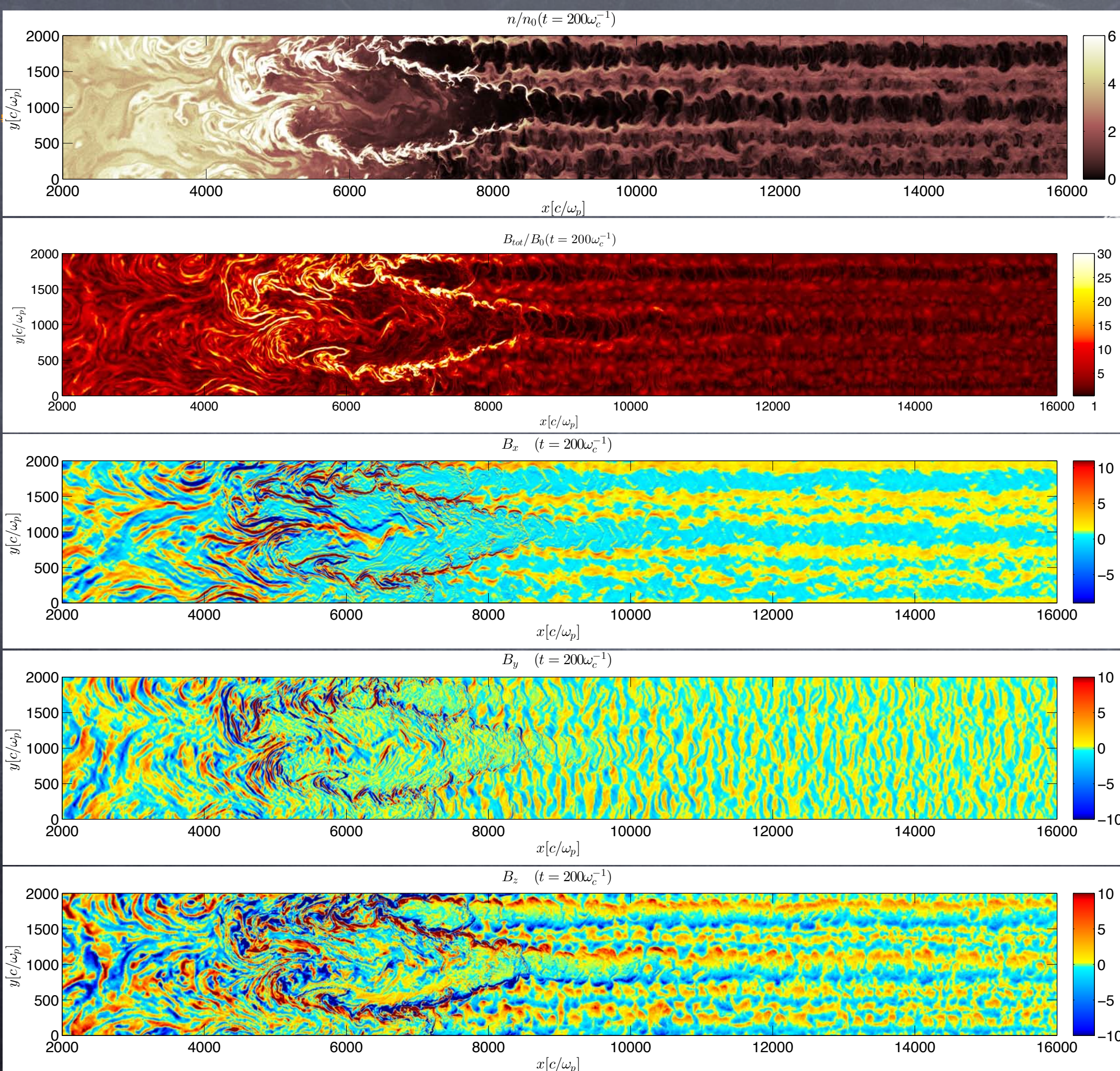
- Towards real shocks: going **bigger and faster**
 - **Super-Hybrid** (Bai, DC, Sironi, Spitkovsky 2014)
- Embedding **microphysics in hydro/MHD** simulations
 - **CRAFT**: CR Analytic Fast Tool (DC et al., in prep)



Hybrid
simulations of a
very strong
shock: $M=100$

Total $\delta B/B$
larger than 10
in the
precursor!

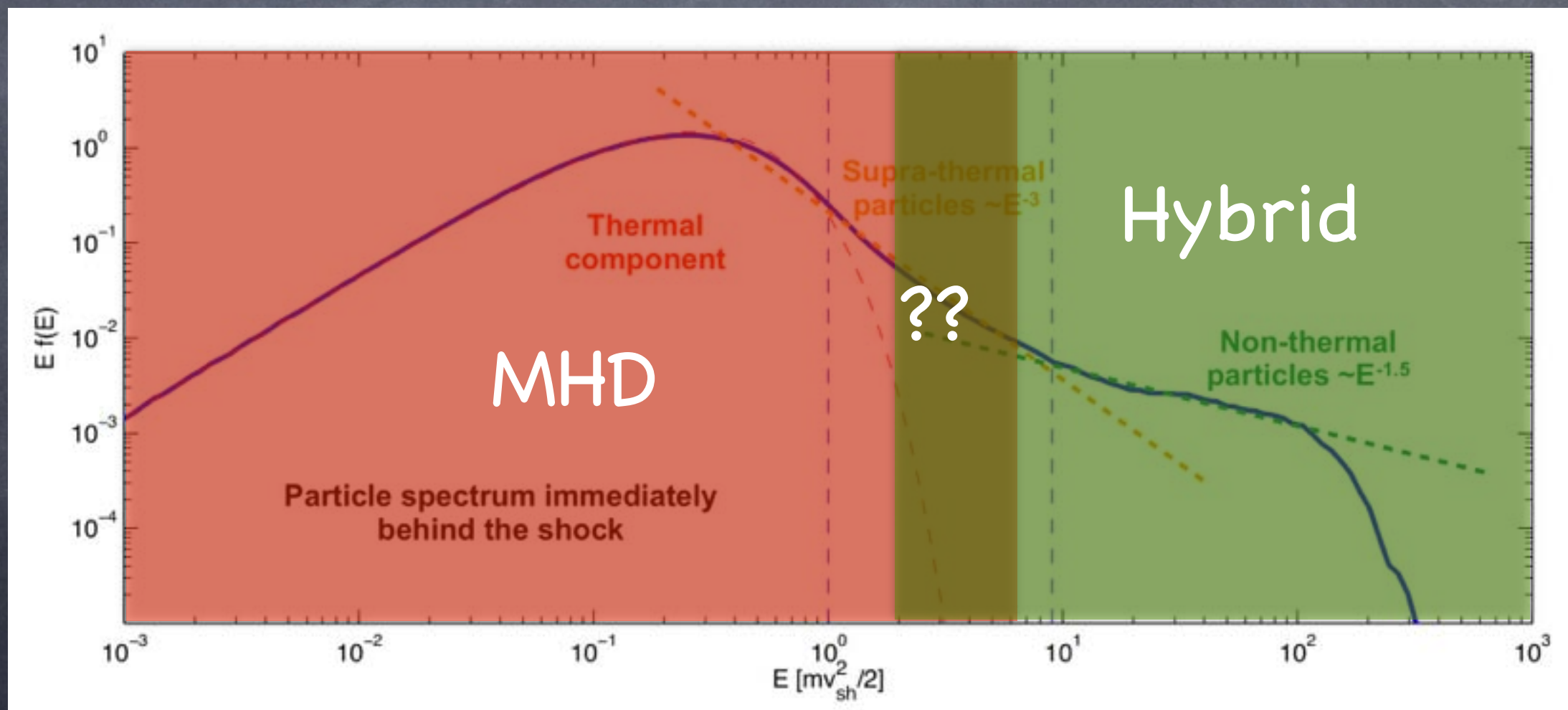
Very expensive
to study in the
hybrid limit



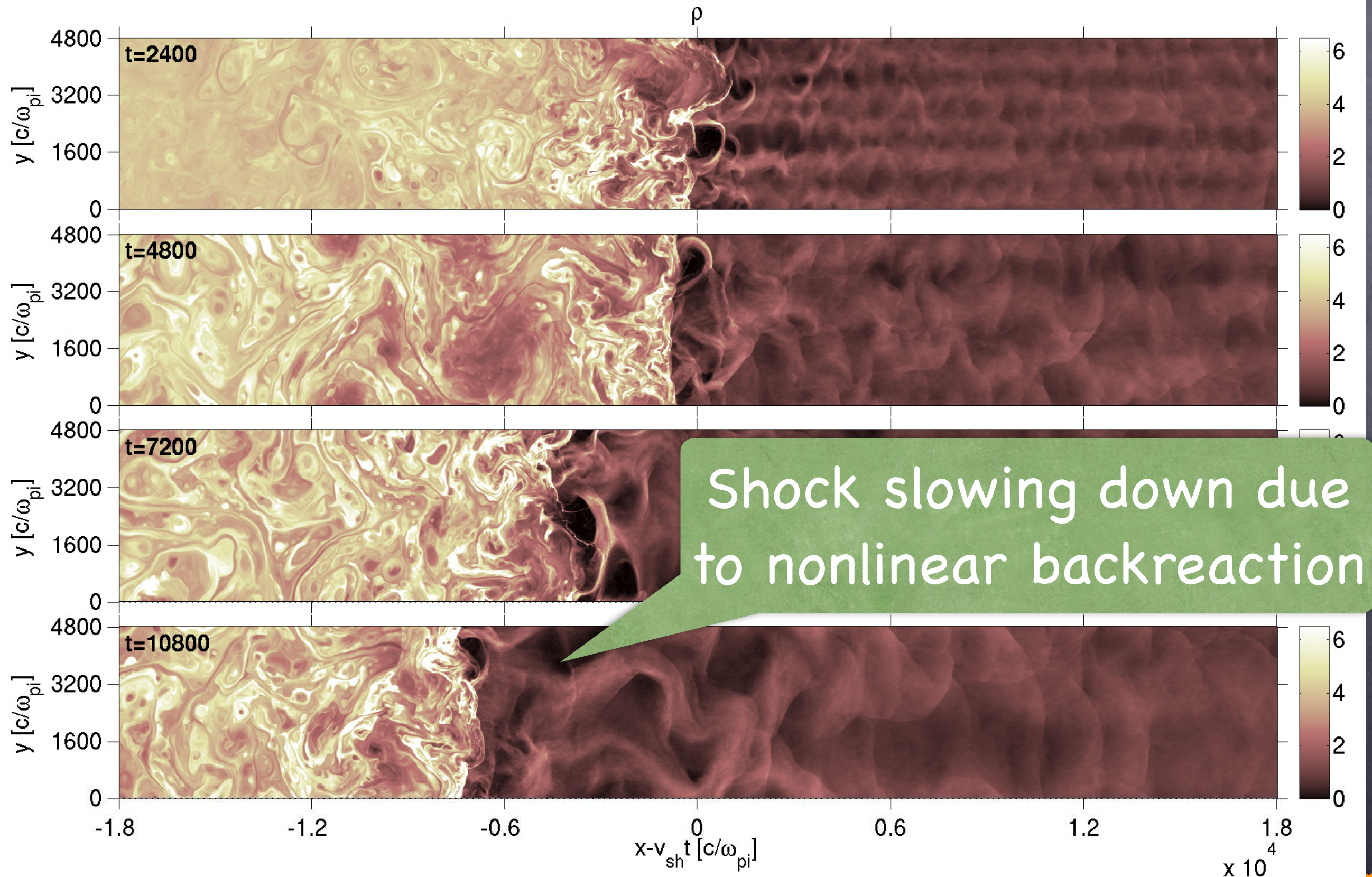


Going bigger: Super-Hybrid

- **MHD** (Athena) + **kinetic ions** (also relativistic)
- Needs **injection** (tuned form hybrid)
- Allows to go to **higher Mach** # and larger scales



Long-term evolution



CRAFT: a Cosmic-Ray Fast Analytic Tool



(Caprioli et al. 2009–2015, to be **publicly released** soon)

- Iterative solution of the **CR transport equation**:

$$\tilde{u}(x) \frac{\partial f(x, p)}{\partial x} = \frac{\partial}{\partial x} \left[D(x, p) \frac{\partial f(x, p)}{\partial x} \right] + \frac{p}{3} \frac{d\tilde{u}(x)}{dx} \frac{\partial f(x, p)}{\partial p} + Q(x, p)$$

$$Q(x, p) = \eta \frac{\rho_1 u_1}{4\pi m_p p_{inj}^2} \delta(p - p_{inj}) \delta(x)$$

Injection

$$f(x, p) = f_2(p) \exp \left[- \int_x^0 dx' \frac{\tilde{u}(x')}{D(x', p)} \right] \left[1 - \frac{W(x, p)}{W_0(p)} \right]$$

$$\Phi_{esc}(p) = -D(x_0, p) \left. \frac{\partial f}{\partial x} \right|_{x_0} = -\frac{u_0 f_2(p)}{W_0(p)};$$

$$W(x, p) = \int_x^0 dx' \frac{u_0}{D(x', p)} \exp \left[\int_{x'}^0 dx'' \frac{\tilde{u}(x'')}{D(x'', p)} \right]$$

$$f_2(p) = \frac{\eta n_0 q_p(p)}{4\pi p_{inj}^3} \exp \left\{ - \int_{p_{inj}}^p \frac{dp'}{p'} q_p(p') \left[U_p(p') + \frac{1}{W_0(p')} \right] \right\}$$

$$U_p(p) = \frac{\tilde{u}_1}{u_0} - \int_{x_0}^0 \frac{dx}{u_0} \left\{ \frac{\partial \tilde{u}(x)}{\partial x} \exp \left[- \int_x^0 dx' \frac{\tilde{u}(x')}{D(x', p)} \right] \left[1 - \frac{W(x, p)}{W_0(p)} \right] \right\}$$

CR distribution function

- Very fast: **a few seconds** on a laptop (vs days on clusters)
- Embeds **microphysics** from kinetic simulations into (M)HD

Mass+momentum
conservation eqs.

$$\frac{p(x)}{\rho(x)^\gamma} = \frac{p_0}{\rho_0^\gamma};$$

$$\rho(x)u(x) = \rho_0 u_0$$

$$\rho(x)u(x)^2 + p(x) + p_{cr}(x) + p_B(x) = \rho_0 u_0^2 + p_{g,0} + p_{B,0}$$

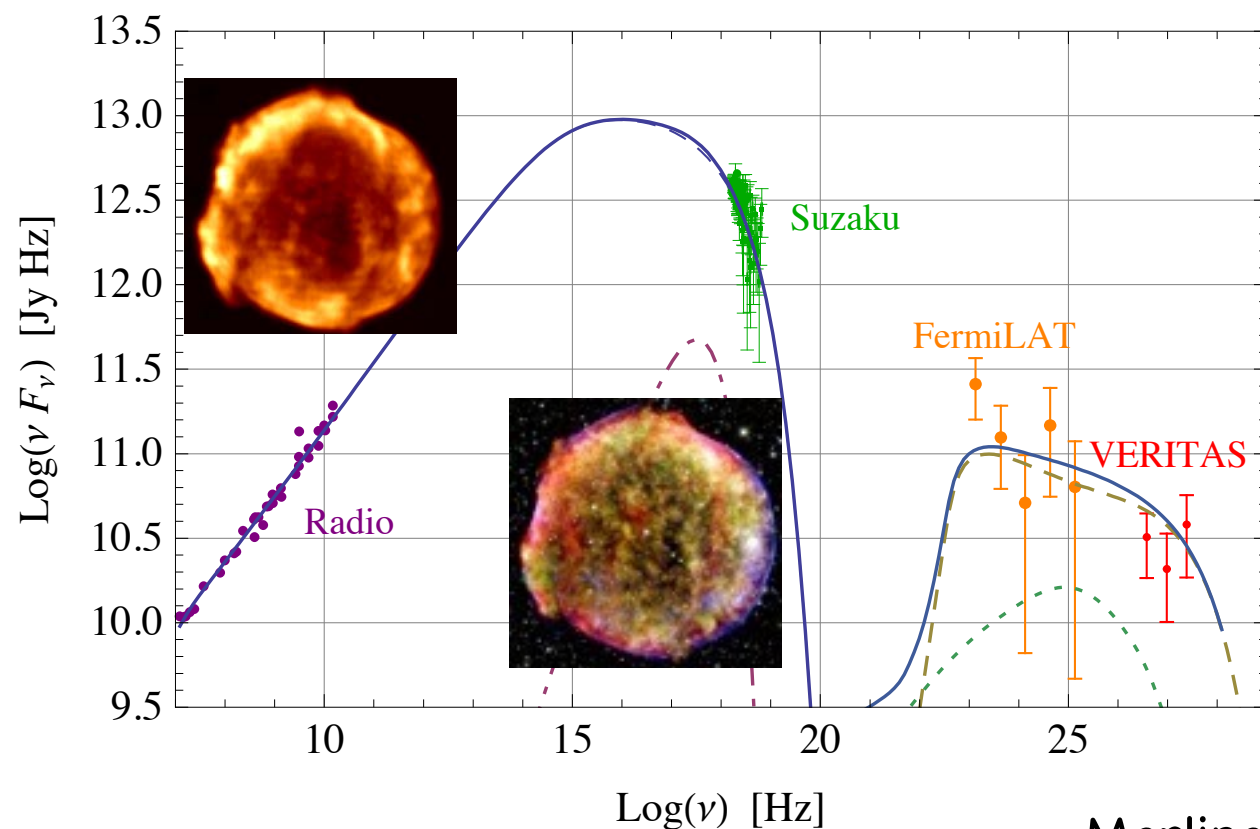
$P_B + P_{cr}$

P_{cr}

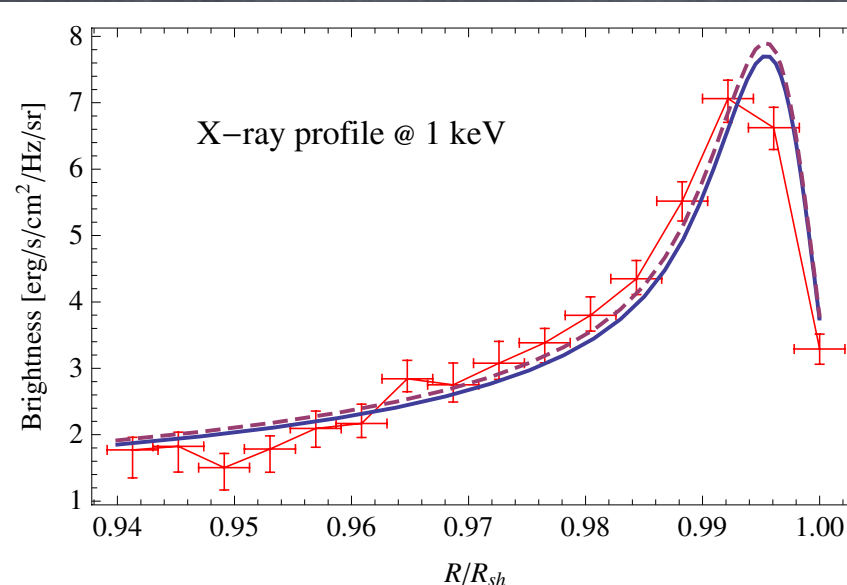
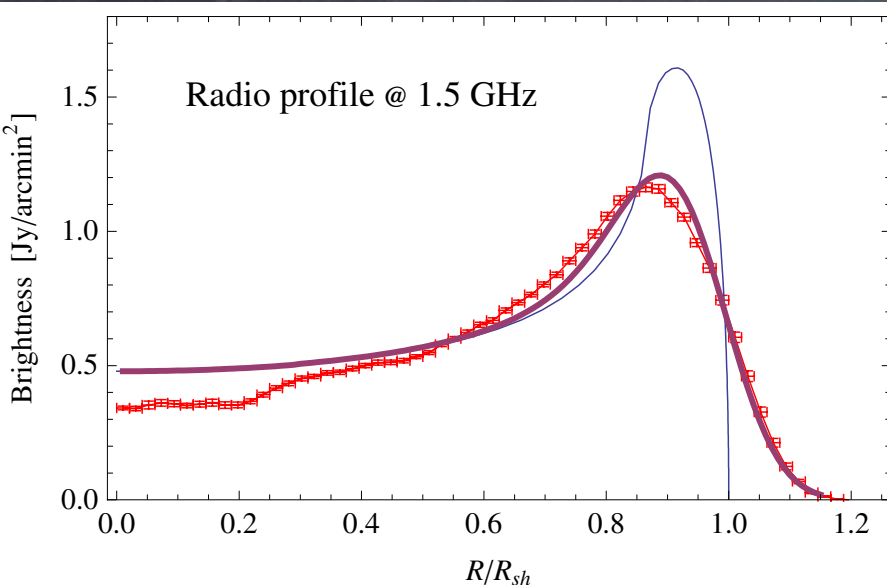
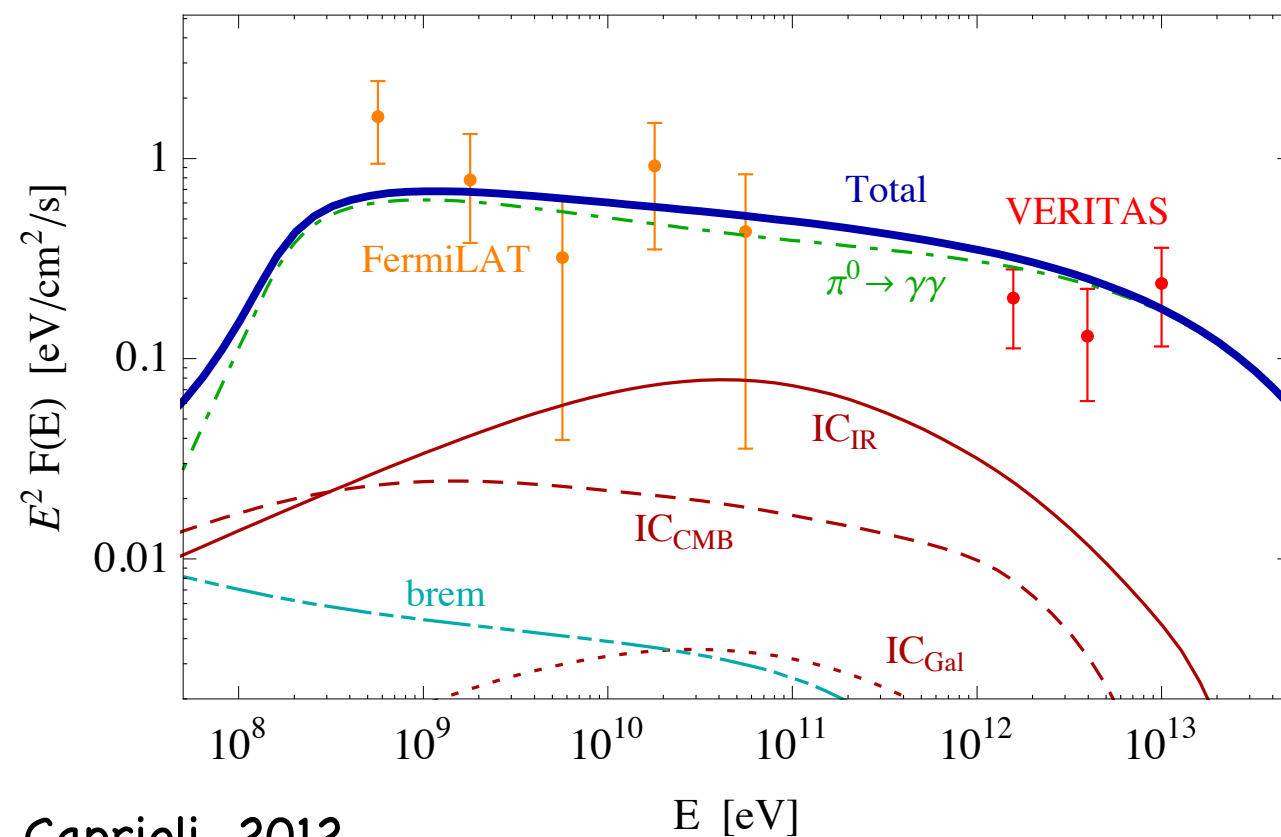
$$2\tilde{u}(x) \frac{dp_B(x)}{dx} = v_A(x) \frac{dp_{cr}(x)}{dx} - 3p_B(x) \frac{d\tilde{u}(x)}{dx}$$

Magnetic turbulence transport eq.

Tycho: a clear-cut hadronic accelerator



Morlino & Caprioli, 2012



- Account for **spectra**, SNR hydrodynamics, and **morphology**
- Hadron acc. eff. **~10%**
- Protons up to **0.5 PeV**

Only two free parameters: **injection efficiency** and **electron/proton ratio**

Thank you!

