



Investigations of stagnated plasma conditions and opacity for K-shell x-ray sources at the Z accelerator

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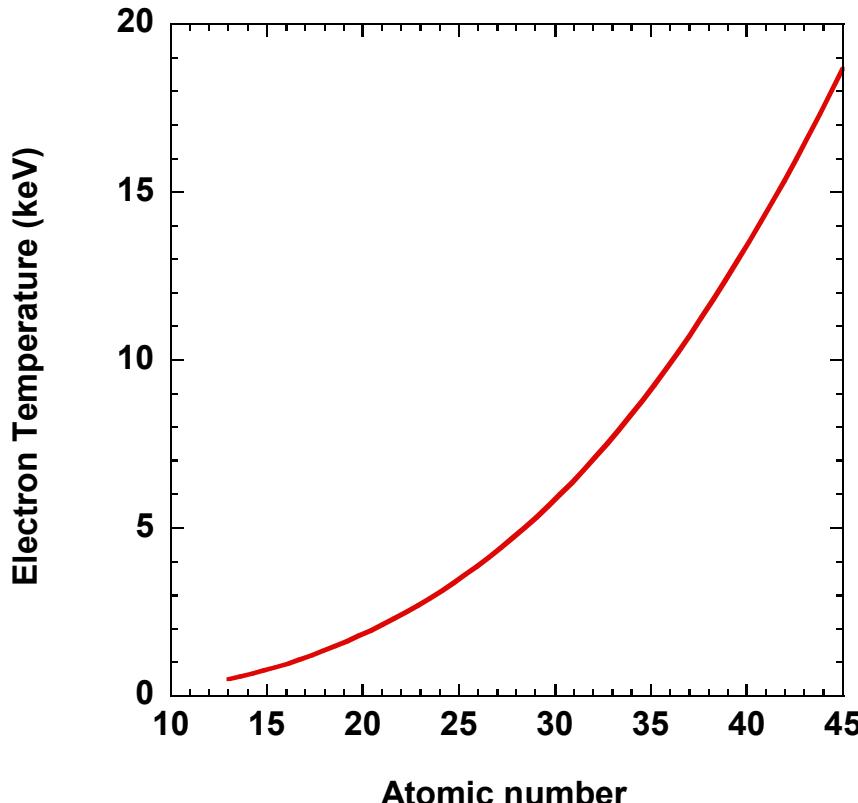
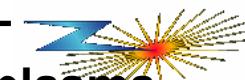


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Temperature requirements for K-shell emission from z pinches can be difficult to attain

- K-shell scaling theories have shown that the temperature of the plasma necessary to produce the K-shell varies with atomic number
 - This suggests that for Cu, T_e must be > 5 keV

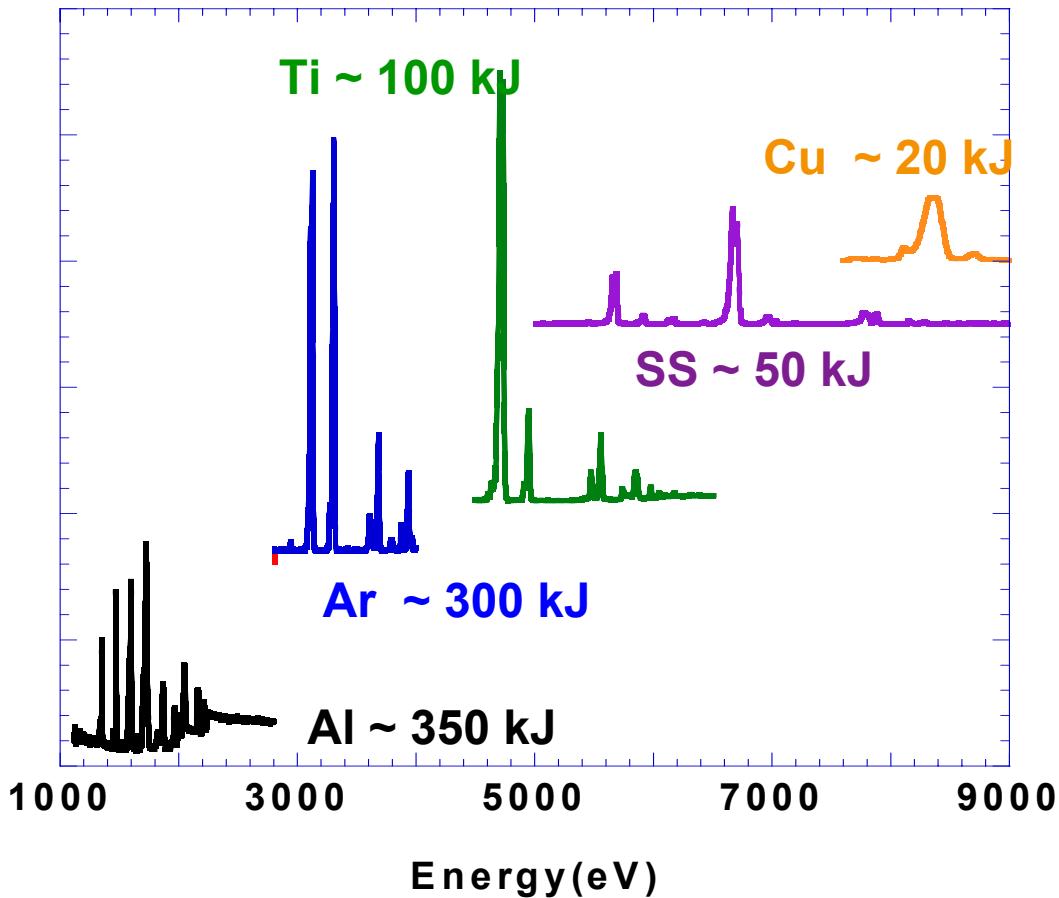
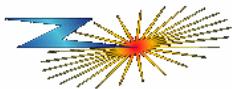


- As the load diameter increases, the $j \times B$ coupled energy is higher at comparable implosion times (lower mass)
- Designing loads that can take advantage of this energy to meet the T_e necessary for K-shell is challenging



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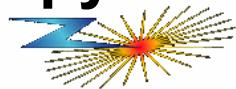
Summary of the Z K-shell sources



- These are peak outputs
- Initial Load diameters:
 - Al: 40mm on 20mm nested
 - Ar: 1234 nozzle
 - Ti: 50mm on 25mm nested
 - SS: 55mm on 27.5mm nested
 - Cu: 60mm on 30mm nested



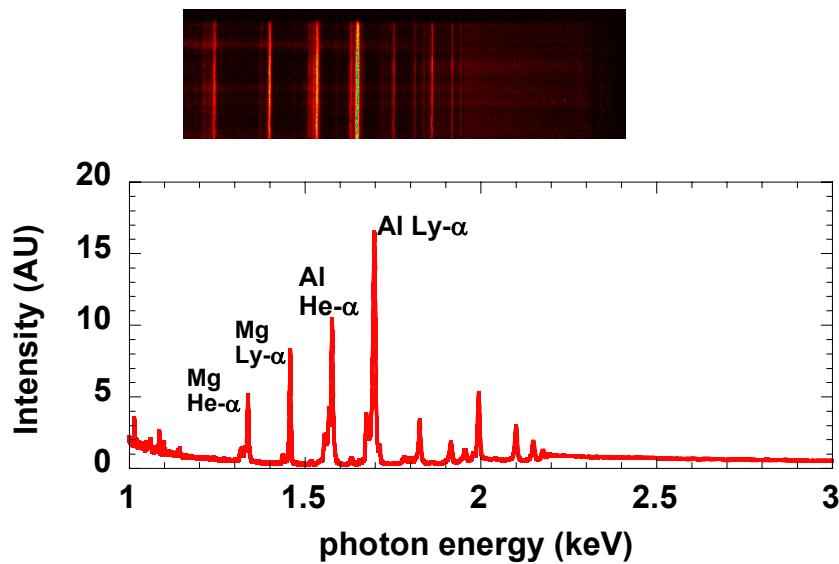
The K-shell x-ray sources can be used to study stagnated z-pinch plasmas through spectroscopy



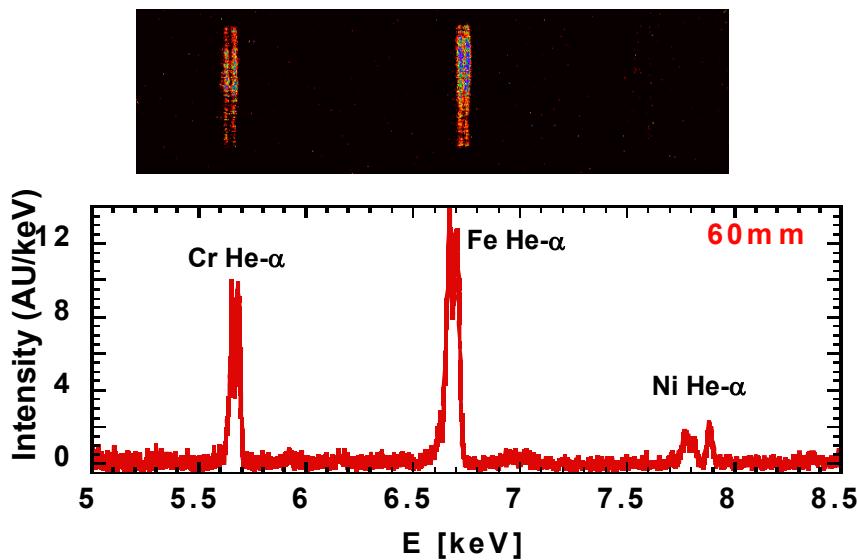
- Variety of K-shell z-pinch experiments have been performed
- Variations are observed in T_e and n_i from the different load materials
 - Plasma conditions are inferred from time-integrated, spatially-resolved spectra, and spatially-integrated, time-resolved spectra, in conjunction with radiated K-shell power and pinch diameter (*J.P. Apruzese et. al., J. Quant. Spectrosc. Radiat. Transfer 57, 41 (1997)*)
 - Measured T_e confirm the scaling theory predictions, although in some cases the conditions are achieved only in isolated regions of the pinch
- The impact of opacity on the K-shell emissions has been directly observed by comparing the line intensities from low level dopant materials with those of the main load constituents
 - Al loads show significant opacity
 - Cu loads appear to be optically thin

K-shell sources offer opportunities to study plasma conditions through spectroscopy

Al (Al 5056), 40 mm nested



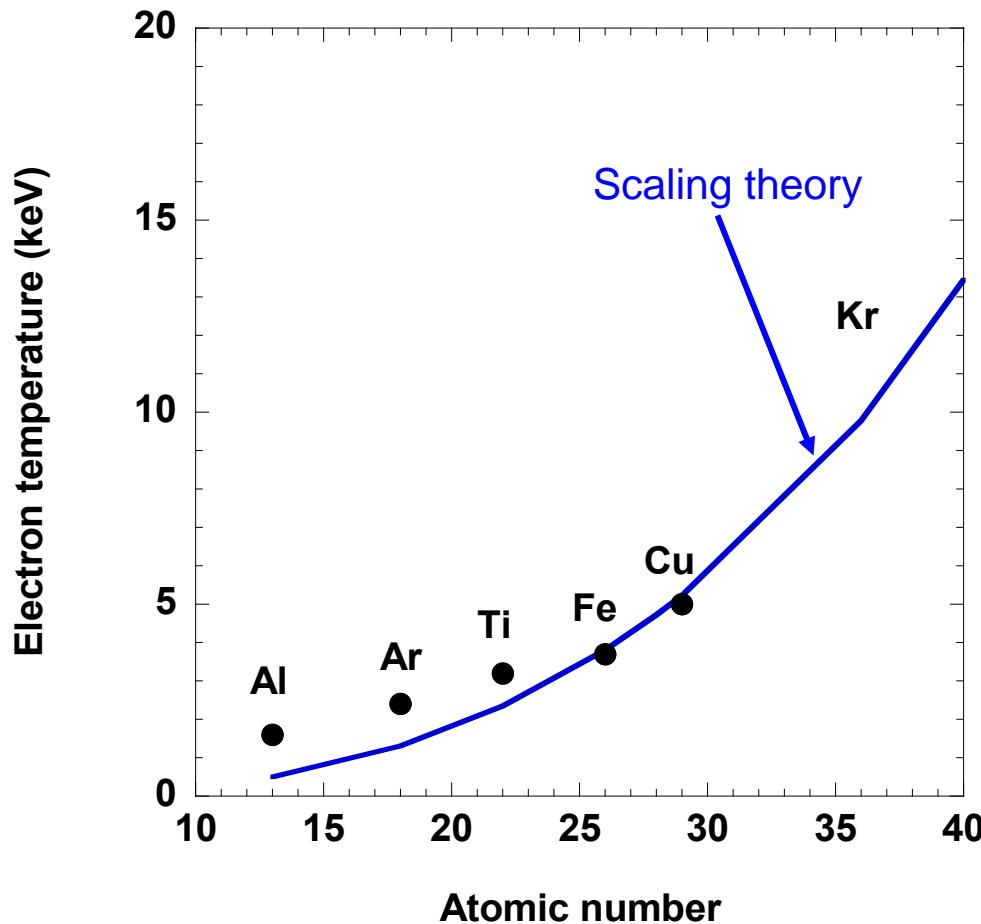
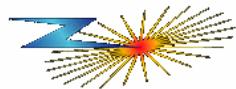
SS, 60 mm nested



Lower Z materials show higher levels of uniformity along full length of pinch, as anticipated based on the less demanding plasma conditions necessary for production of the K-shell x-rays

Axially averaged spectra used to estimate plasma conditions

Inferred electron temperatures agree well with scaling theory trend



Scaling theory:

$$T_e = 0.3 \cdot Z^{2.9} \text{ eV}$$

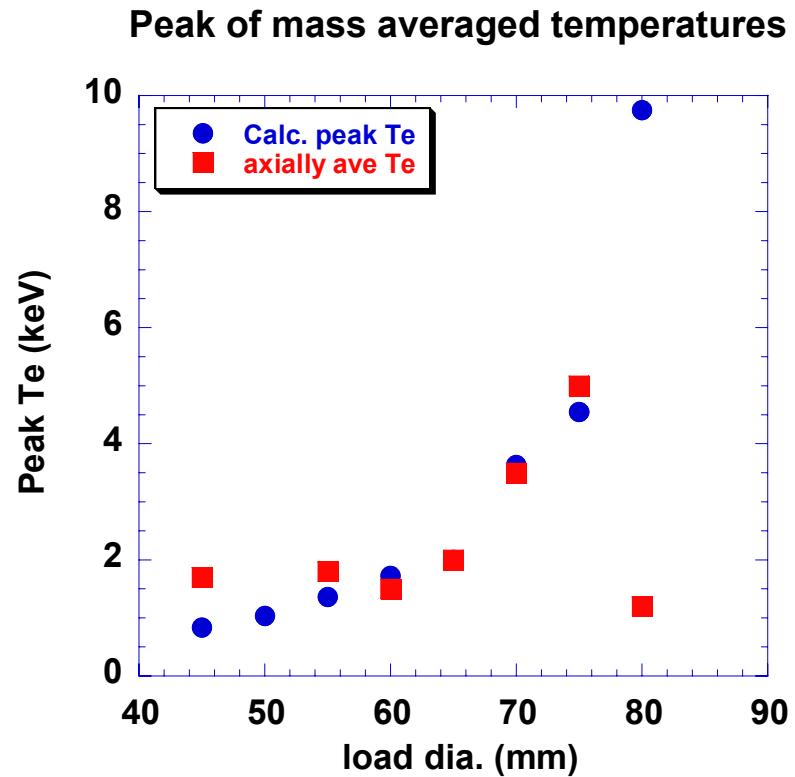
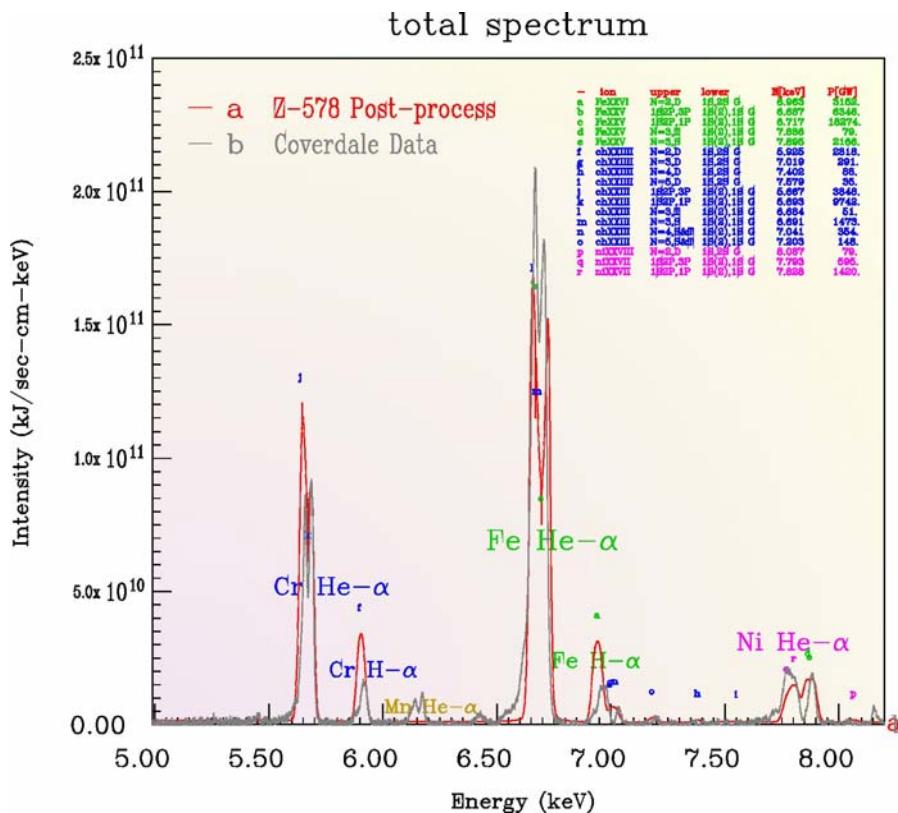
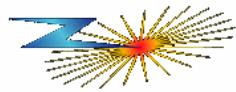
Electron temperatures are obtained from axially averaged, time-integrated spectral line ratios from “optimized” K-shell x-ray sources



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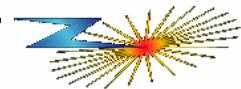
1-D calculations of SS variations fielded reproduce the measured spectra

R. Clark & J. Davis, NRL

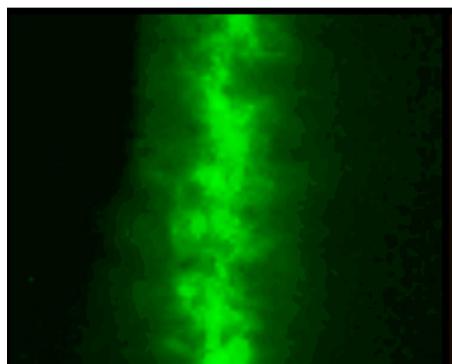


- Overlap of strong He- α with He intercombination line leads to higher emission on the low energy side of the lines

Imaging of K-shell emission near peak radiated output show significant non-uniformity

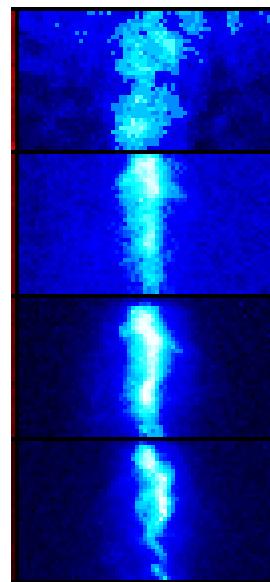


Al 40mm nested



-2.3 ns

SS 55mm nested



-1 ns

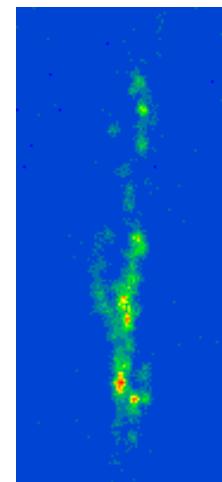
0 ns

+1 ns

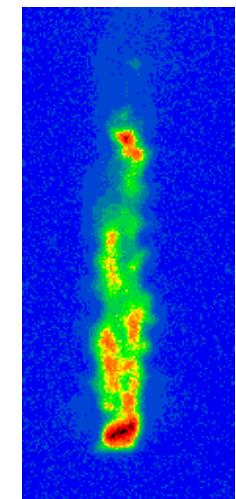
+2 ns

Cu

55mm single



55mm nested



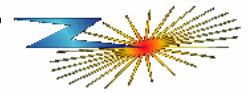
- Lower Z shows instability and structure, but more uniformly distributed intensity along the pinch length
 - Higher Z materials show more substantial variations
- Use of a nested configuration clearly helps with K-shell emission for higher Z materials



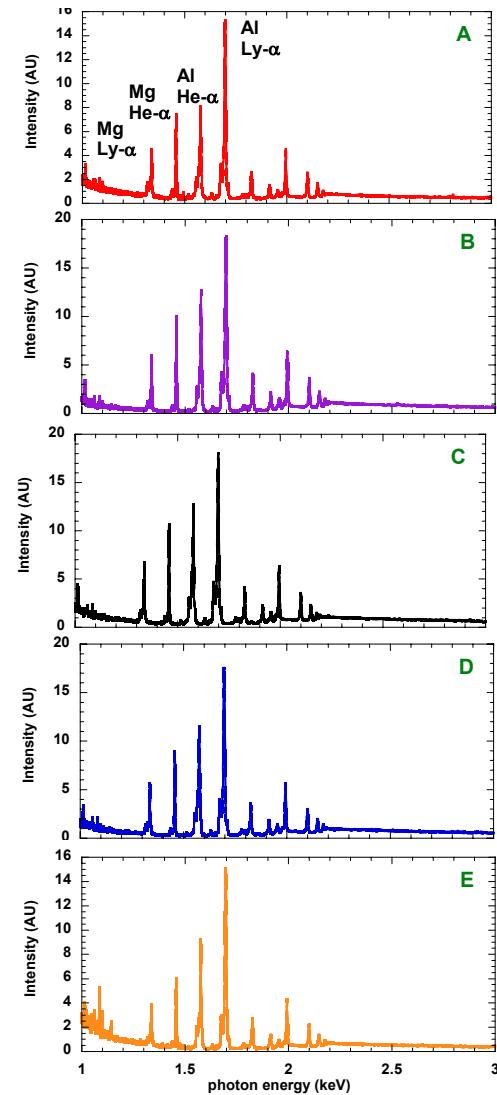
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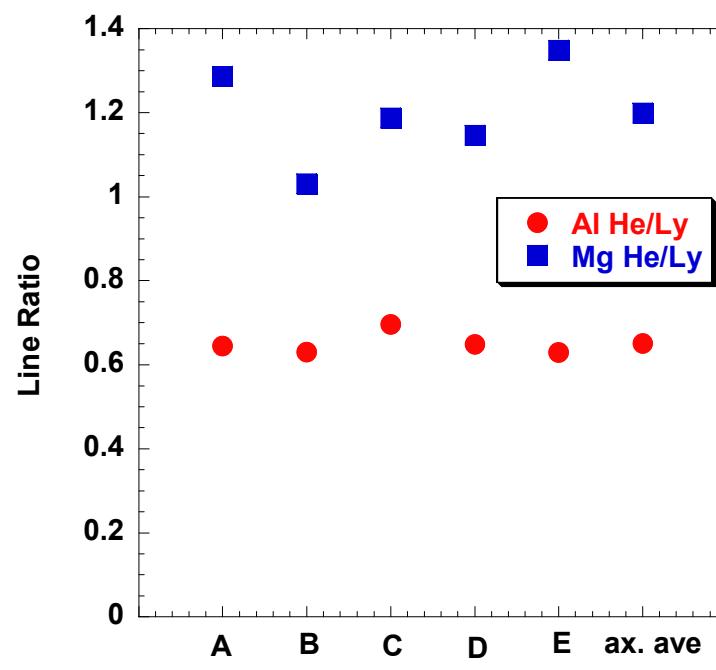
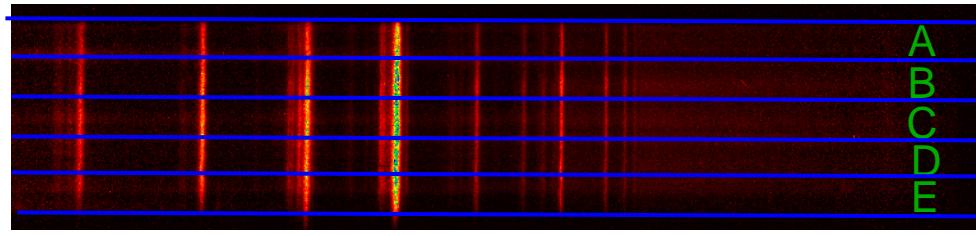
Al z pinch is overdriven, and temperature and density vary little over the pinch length



Z1520, Al



40mm nested; Al 5056 (5% Mg)

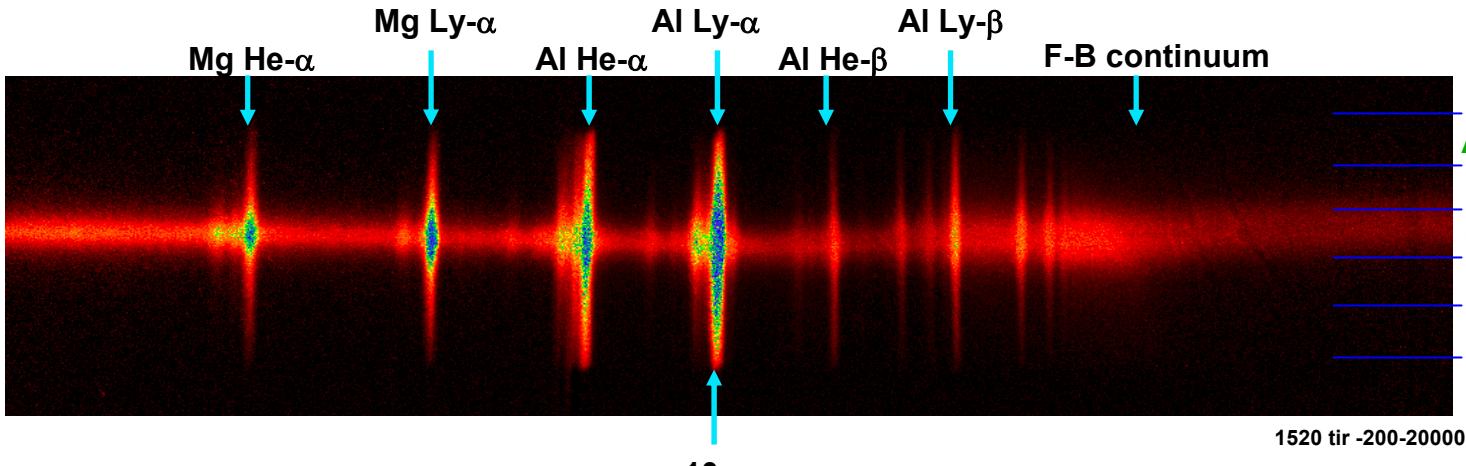


- Brightest line is Al Ly- α
 - indicates overdriven plasma
- Analysis shows:
 - $T_e = 0.5 \text{ keV}$
 - $N_i = 2.4 \times 10^{20}/\text{cm}^3$
 - K-shell Mass Fraction = 20%



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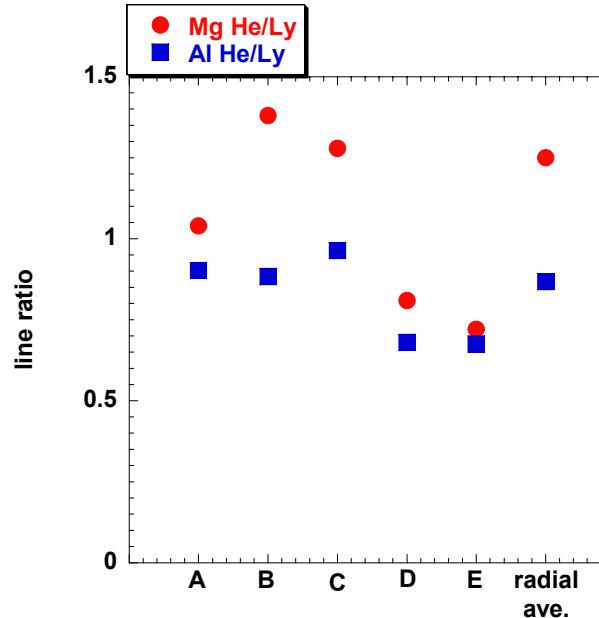
Radially resolved Al5056 spectrum shows radial gradients exist



FWHM (mm)

Al He- α :	6.1
Al Ly- α :	8.0
Mg He- α :	2.7
Mg Ly- α :	3.4
Free-Bound cont.:	2.1

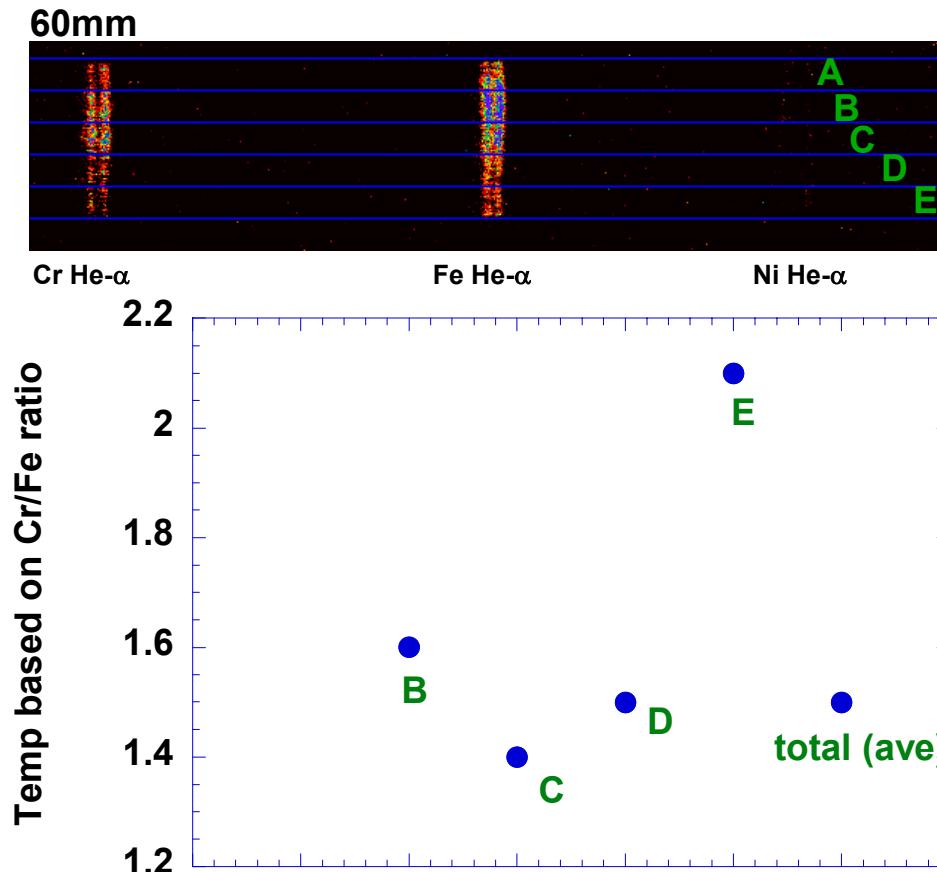
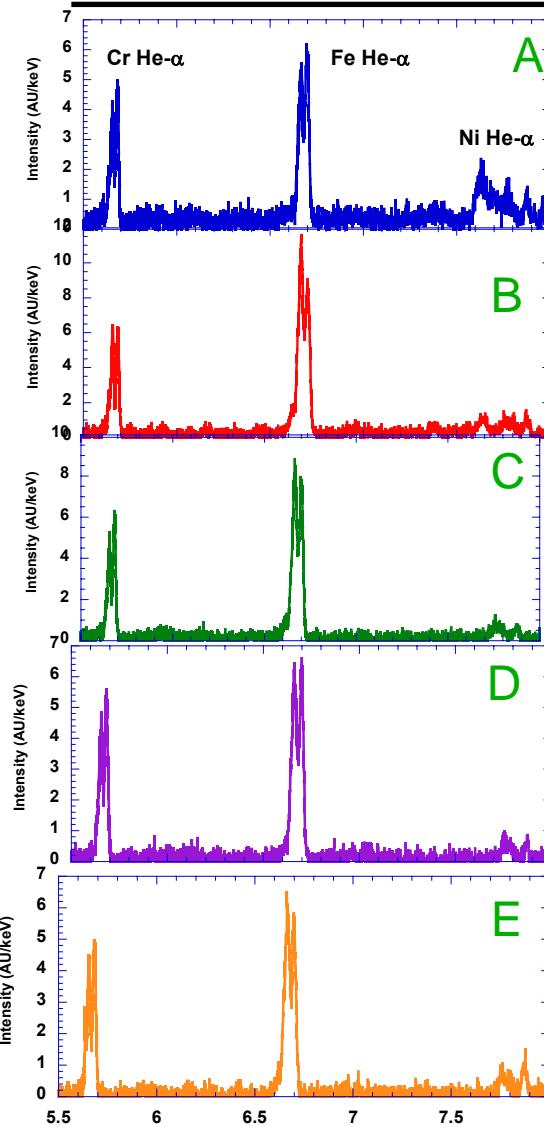
- Full viewing of the plasma is limited by the width of the slots in the current-return can.
- Lower density plasma may be resonantly scattering line radiation at large radii, but difficult to fully assess with time-integrated data



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More significant variation along the pinch length is observed for SS z pinches

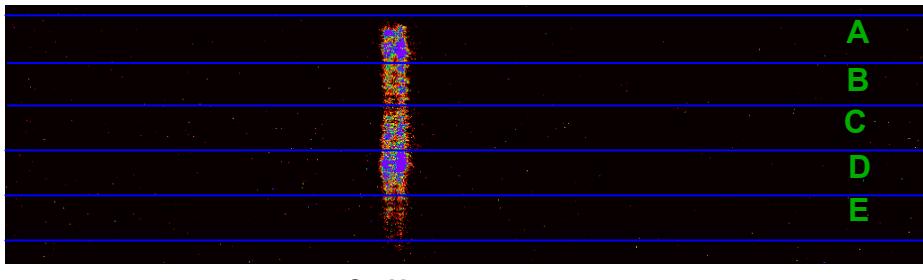
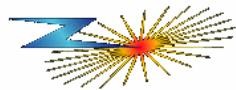


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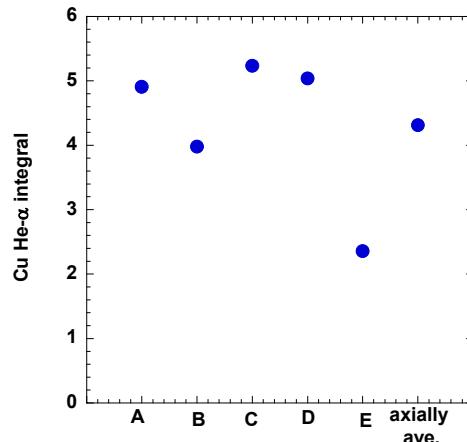
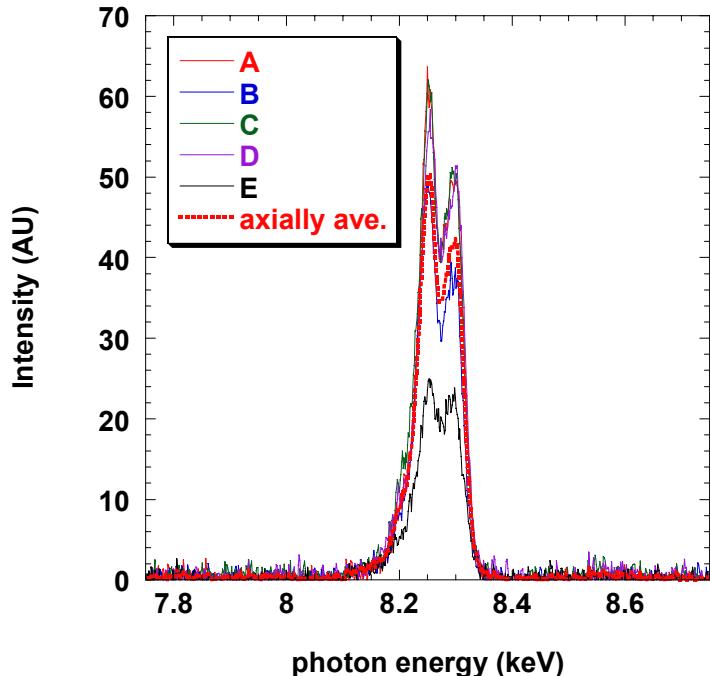
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Cu arrays also show substantial variation along the length of the pinch

Z1270, 60mm nested



Cu He- α



Difficult to obtain Te, ni for this shot due to lack of other spectral lines

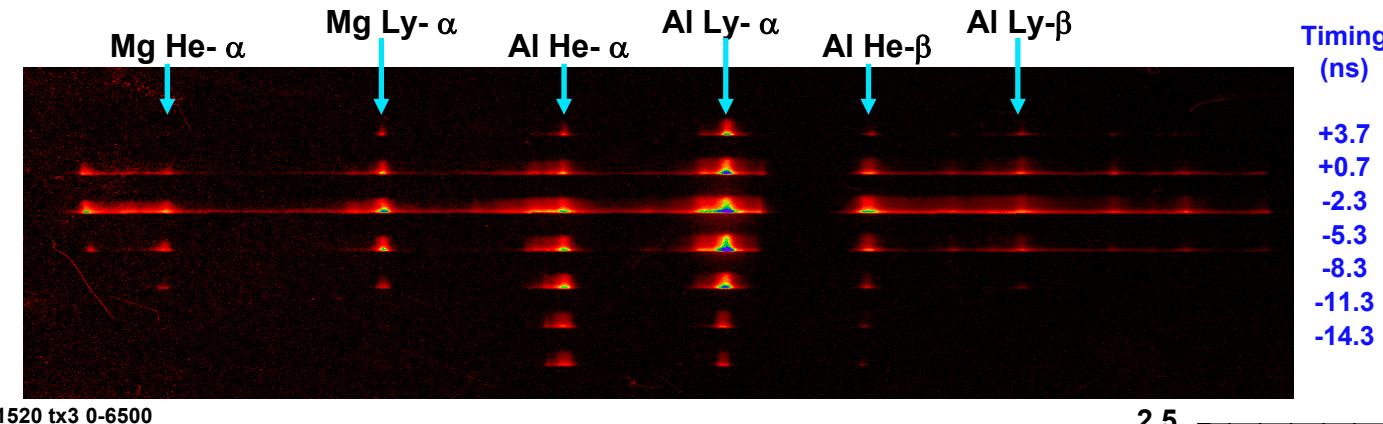
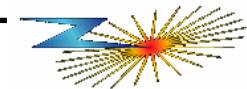
- Z1270: 2% Ni on inner array, pure Cu on outer array
- Note that other work at Saturn (C. Deeney et al., Phys. Rev. Lett. 8 155001 (2004)) and Z (C.A. Coverdale, in preparation) has shown that most of the radiation in the wire arrays originates from material on the outer array, which had no dopant for this array.



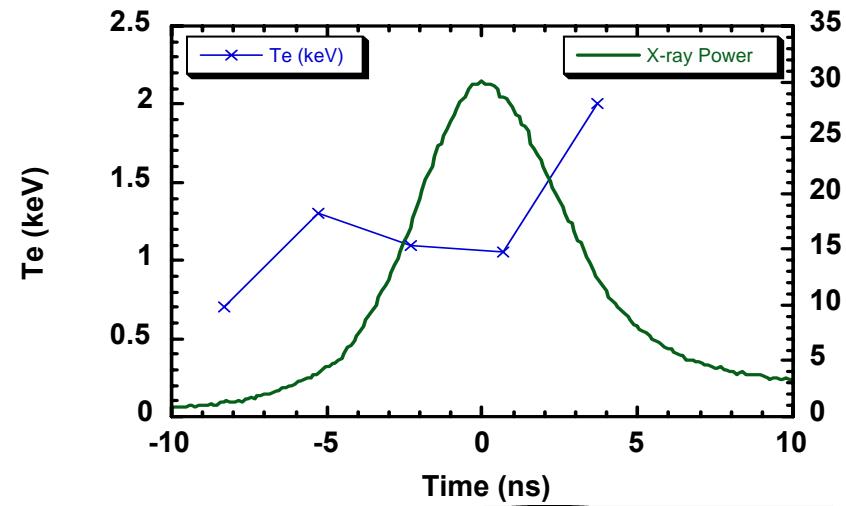
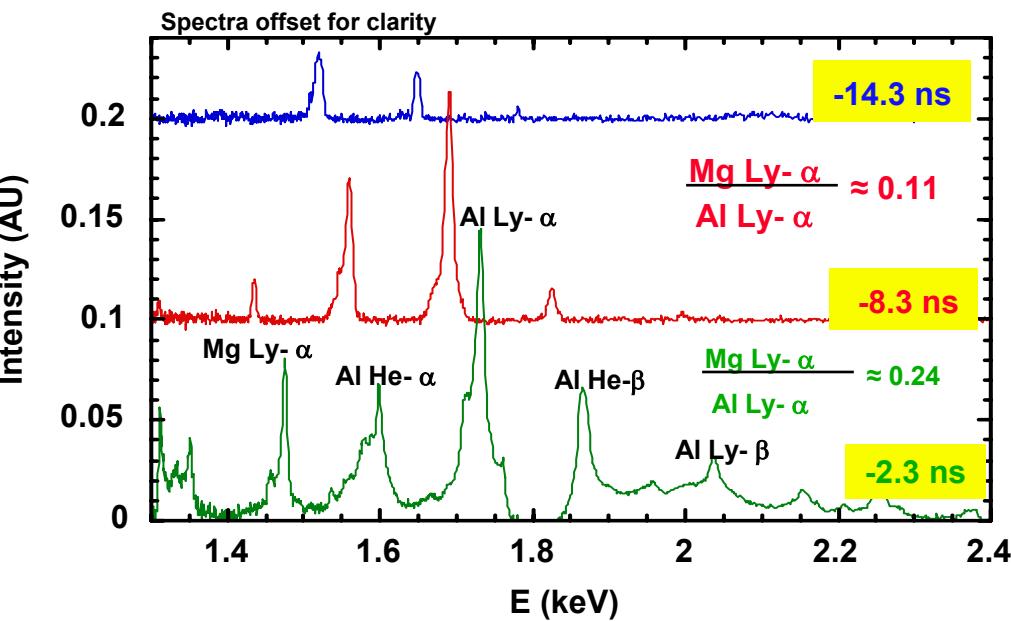
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Time-resolved Al5056 spectrum shows varying plasma conditions during the implosion



1520 tx3 0-6500



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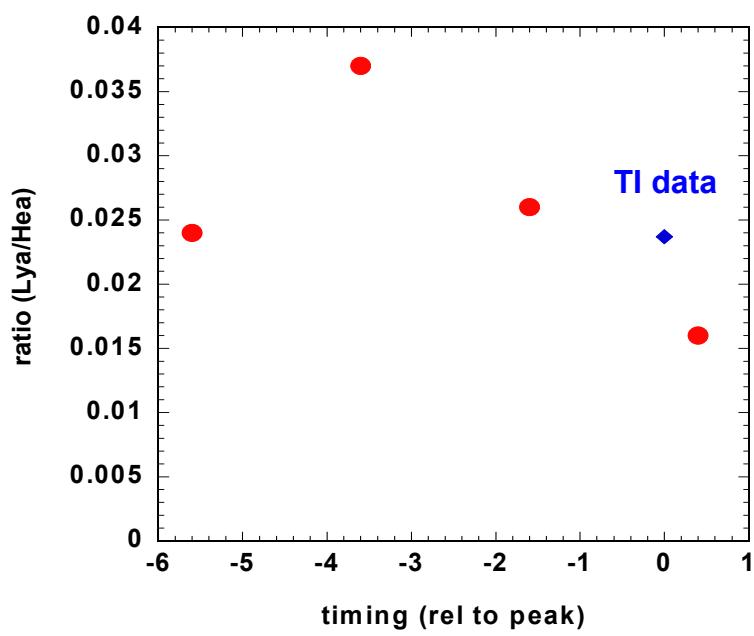
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The plasma conditions are also observed to evolve in Cu z pinches

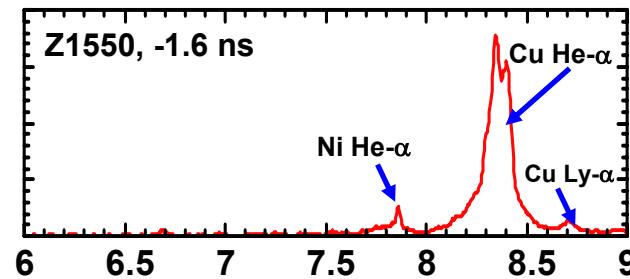
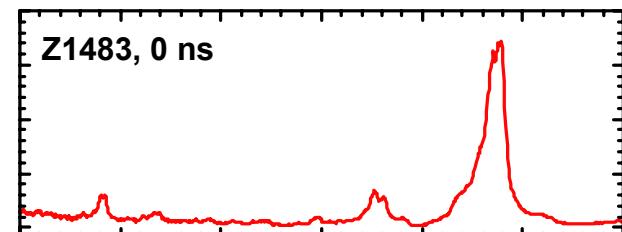
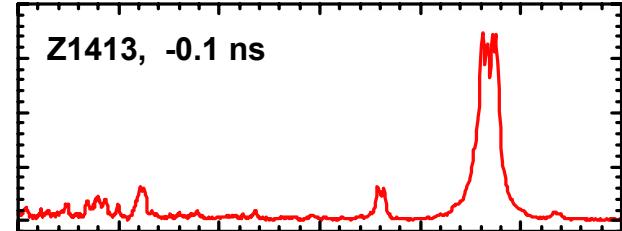
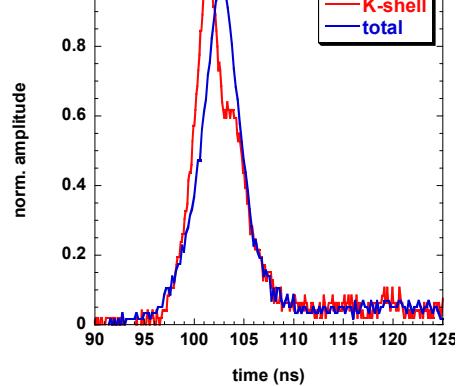
Z1550, 60mm Cu (6% Ni)

Ni He- α Cu He- α Cu Ly- α

-5.6
-3.6
-1.6
+0.4



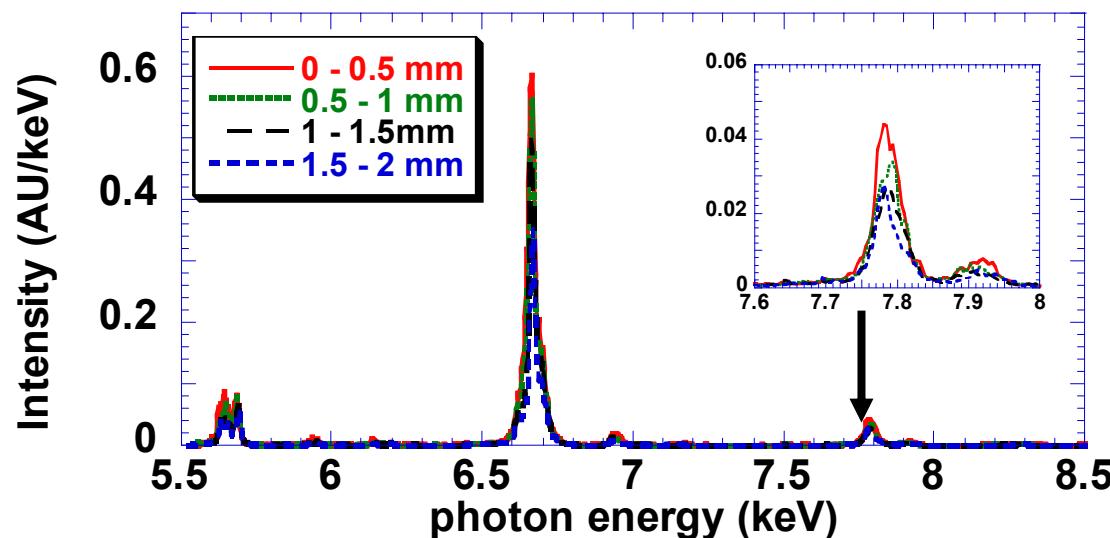
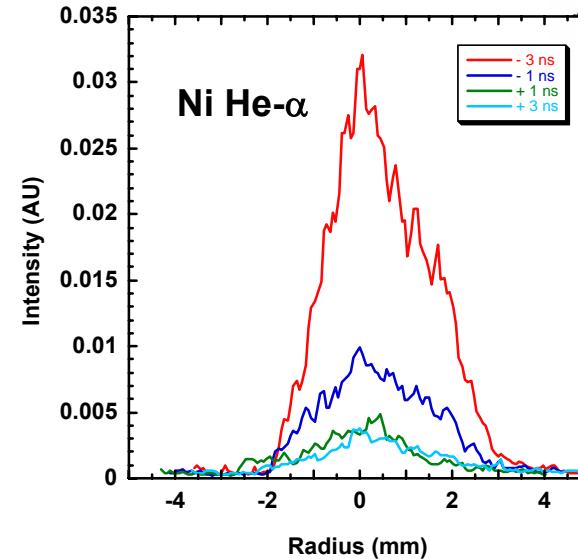
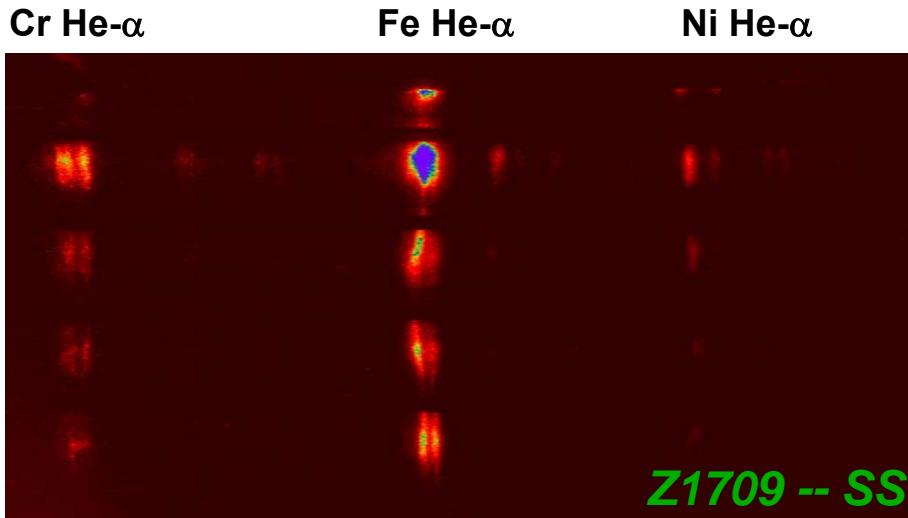
K-shell emission peaked \sim 1.7ns before total emission



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Radial gradients are observed in SS pinches with time and space resolved spectroscopy

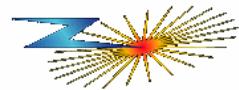


- Radially resolved data shows similar core size for Cr and Ni He- α emissions
- Ni emissions dissipate faster than Cr, however, due to higher T_e needed to produce Ni



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Observance of hot cores surrounded by cooler blankets suggests that opacity can be studied

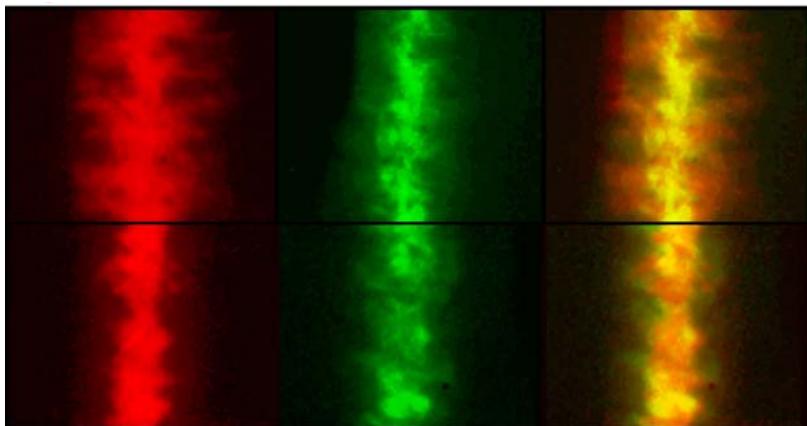


277 eV

1.7 keV
Al K-shell

Overlay

Z1519, Al



-2.3 ns

+3.2 ns

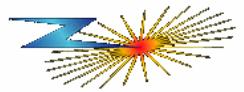
- MLM images show hot, dense column emitting K-shell on axis with colder material still imploding

Radially resolved spectra shown earlier for Al and SS illustrate variations in line emission moving radially outward from central pinch axis

High temperature, dense cores observed with Al arrays at Saturn by T.J. Nash *et al.* (J. Quant. Spectrosc. Radiat. Transfer 60, 97 (1998))



Sources of opacity, with approximate cross sections (cm²), in order of importance



1. Strong lines (oscillator strength ~ 0.5):

$$\sigma(\text{line center}) \approx 6 \times 10^{-18} Z^{0.55}/[\hbar\nu(\text{keV})]$$

-- assumes Doppler broadening with $T_{\text{ion}} = 20 \text{ keV}$
(reasonable for most cases, based on

J. S. Levine *et al.*, Phys. Plasmas 8, 533 (2001)

M.G. Haines *et al.* Phys. Rev. Lett. 96, 075003 (2006))

-- Since $\hbar\nu \sim Z^2$, $\sigma \sim Z^{-1.45}$

2. Bound-free (photoionization per valence electron):

$$\sigma(\text{at threshold}) \approx 1.2 \times 10^{-17} Z^{-2}$$

-- above threshold $\sigma \sim (\hbar\nu)^{-n}$, $n > 1$, typically, $n \sim 3$

3. Free-free (bremsstrahlung, Gaunt factor of 1):

$$\sigma \approx 7.6 \times 10^{-28} N_i(10^{20} \text{ cm}^{-3}) Z^3 T_e(\text{keV})^{-1/2} [\hbar\nu(\text{keV})]^{-3}$$

-- assumes $N_e = N_i Z$



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Absorption and reradiation of photons



- The absorption process of line radiation removes photons from the emission, yielding higher energy electrons, but fewer x-rays
- Step 1: The photon excites an ion of the same species that emitted it, raising an electron from the ground state to an excited level. Example: $\text{Al (1s)} + h\nu \rightarrow \text{Al (2p)}$
- Step 2: The excited ion is then de-excited by superelastic collision(s), restoring the photon's energy to the electron pool. Example: $\text{Al (2p)} + e^- \rightarrow \text{Al (1s)} + e^-$
- Reradiation to the ground state before collisional de-excitation does not remove photons (scattering), no energy loss.
- These two processes compete, but absorption increases with increasing density, eventually dominating.
- Step 1: The probability that a line photon is intercepted by a K shell ion varies as the cross section: $Z^{-1.45}$
- Step 2: The probability of electron quenching is the collisional de-excitation rate divided by the spontaneous decay rate. The collisional rate varies as Z^{-3} ; the radiative decay rate as Z^4 , giving an overall variation as Z^{-7} !

The sharp dependences of the underlying atomic physics factors can produce large variations in K-shell opacity for relatively small differences in Z

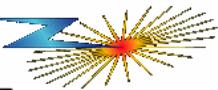
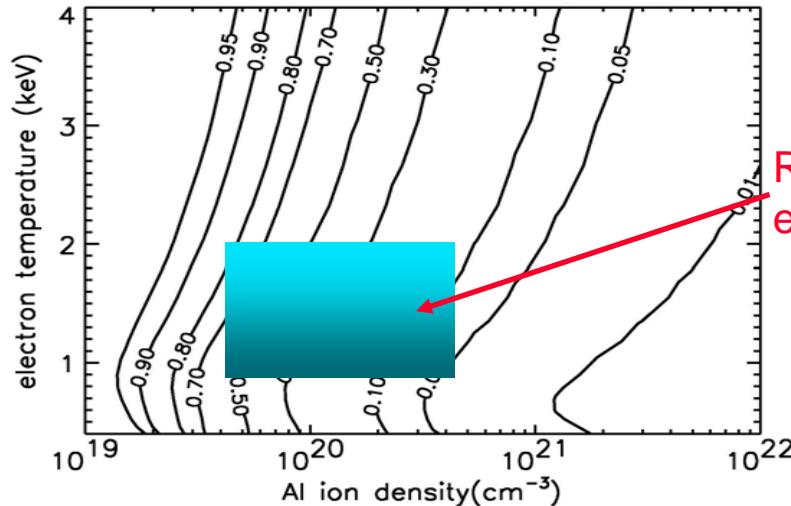


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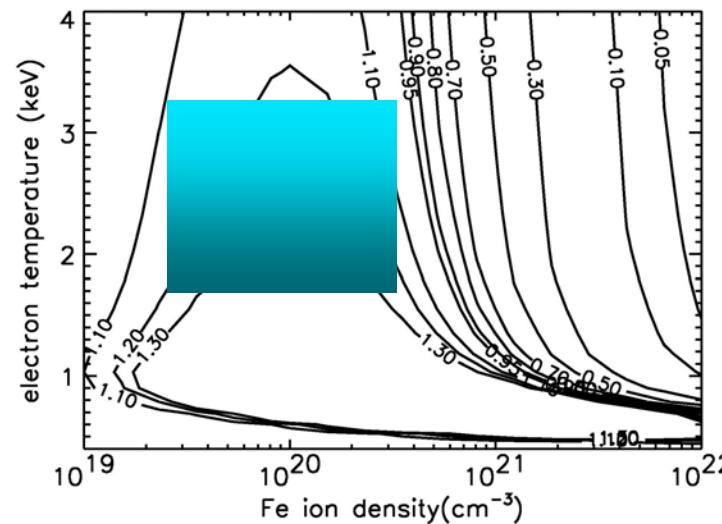
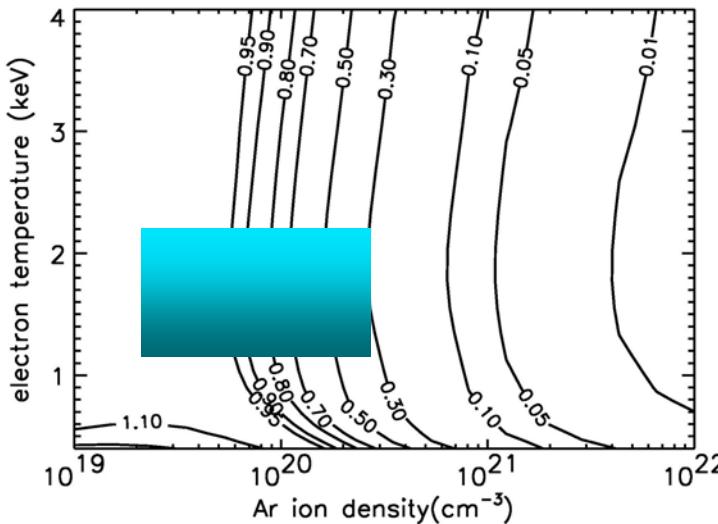
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Thick-to-thin power ratio for K-shell α lines show effects of increasing opacity

- A value of 0.10 indicates an optically thick plasma radiates only 10% that of an optically thin plasma
- Al plasmas on Z show these effects; higher Z materials (Ar, Fe) do not
 - In fact, ladder ionization at low density can produce a ratio > 1



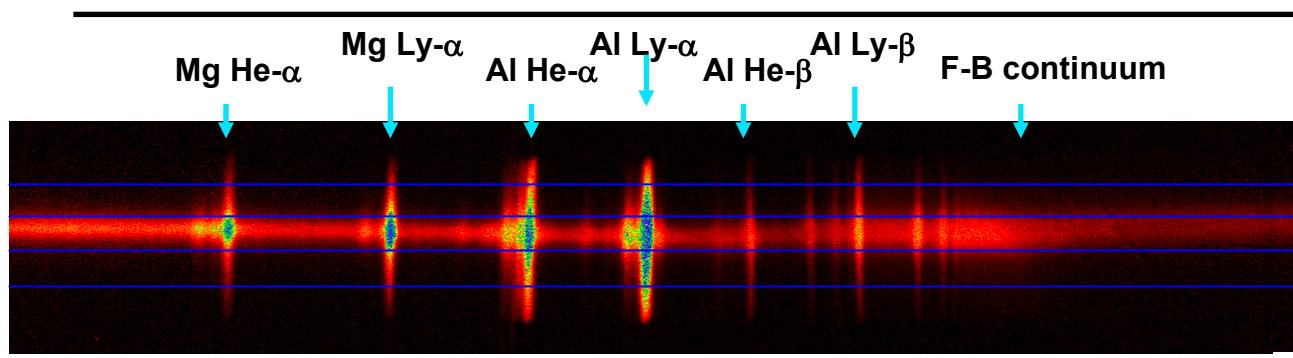
Region of Z
experiments



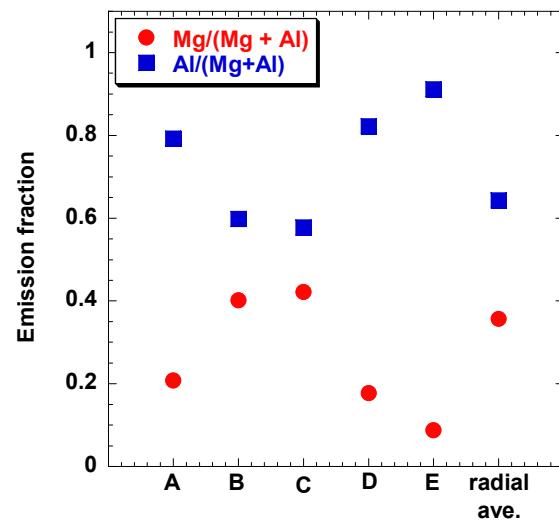
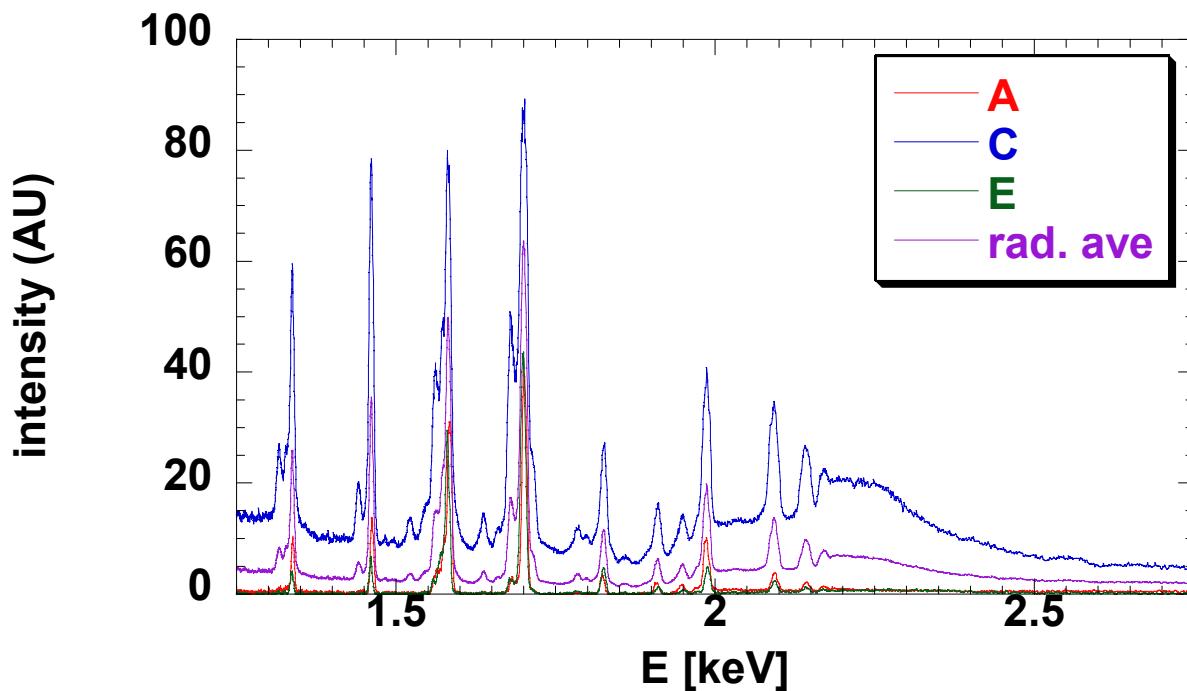
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Al data (time-integrated, radially resolved) shows effects of opacity



Easier for photons produced in cooler blanket to escape

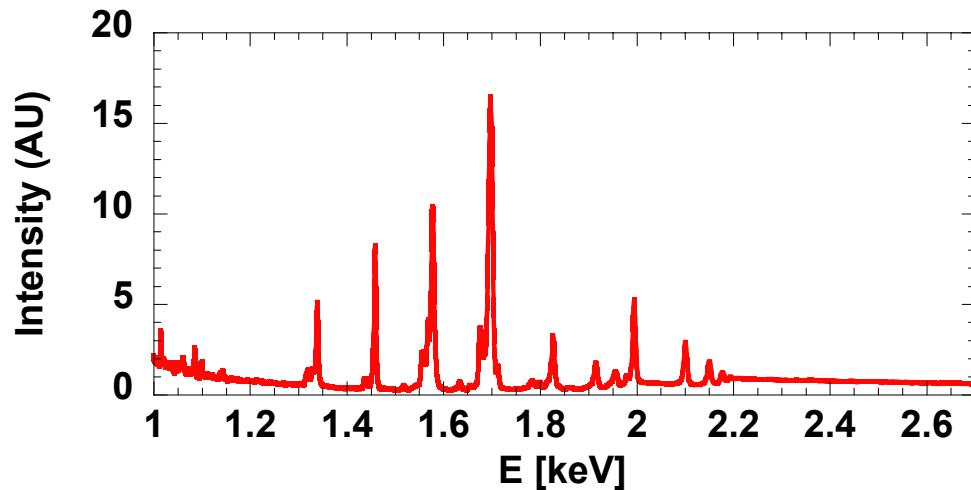


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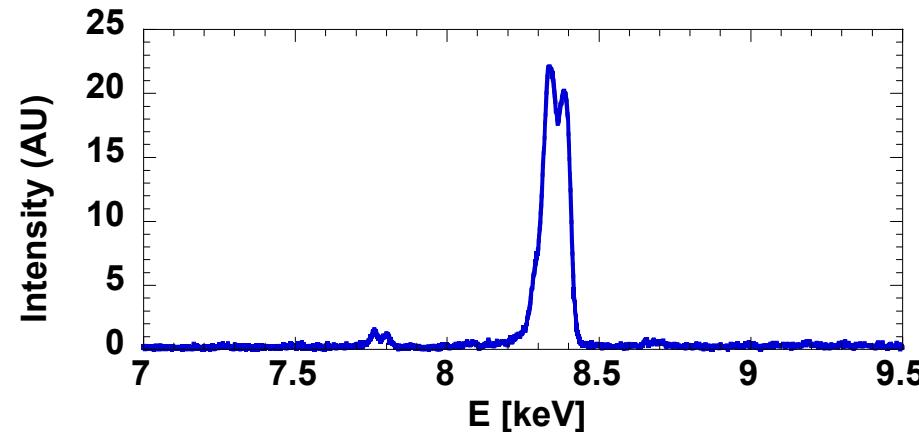
The impact of opacity is evident experimentally in time-integrated spectra; Atomic number trend is consistent with theory

Al (40 mm nested); Al 5056 (5% Mg) on both arrays



Amplitude of Mg relative to Al suggests that mass distribution is **~30% Mg, ~70% Al**
• clear evidence of opacity issues for Al at Z

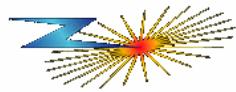
Cu (60 mm nested); 94% Cu / 6% Ni on both arrays



Amplitude of Ni relative to Cu suggests that mass distribution is **~10% Ni, 90% Cu**
• Clearly little if any opacity impact here



Summary



- A variety of K-shell sources have been fielded at Z
 - Wide range of photon energies fielded
 - Assessment of K-shell scaling theories on required temperatures possible
 - Study of temperature gradients possible (axial, radial, and temporal)
- Temperature and density gradients observed
 - Pinhole images suggest significant axial structure present
 - Axially resolved spectra show peak temperature generally higher than axially averaged temperature for higher Z materials
 - Consistent with images showing brighter regions of emission
 - Radially resolved spectra show hot core with cooler surrounding blanket
 - Temporally resolved spectra show varying pinch conditions during radiation pulse
- Effects of opacity clearly evident in measured spectra
 - Opacity comes into play due to atomic physics of absorption and reradiation
 - Expected to be more severe with lower Z
 - Al data clearly shows impact of opacity
 - Cu data appears to be optically thin