

# Simulation of Cathode Plasma Expansion in Magnetically Insulated Transmission Lines (MITLs)\*

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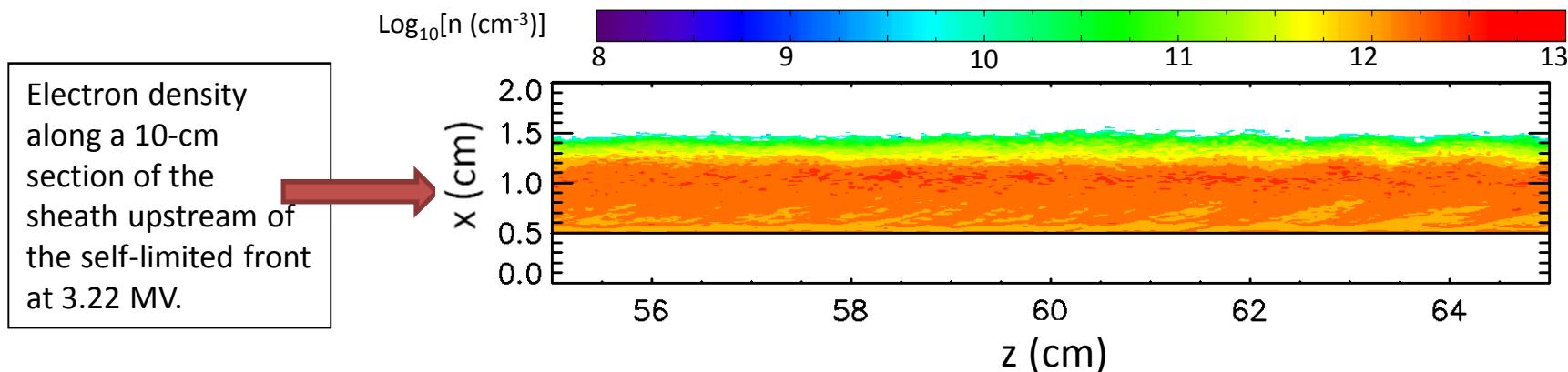
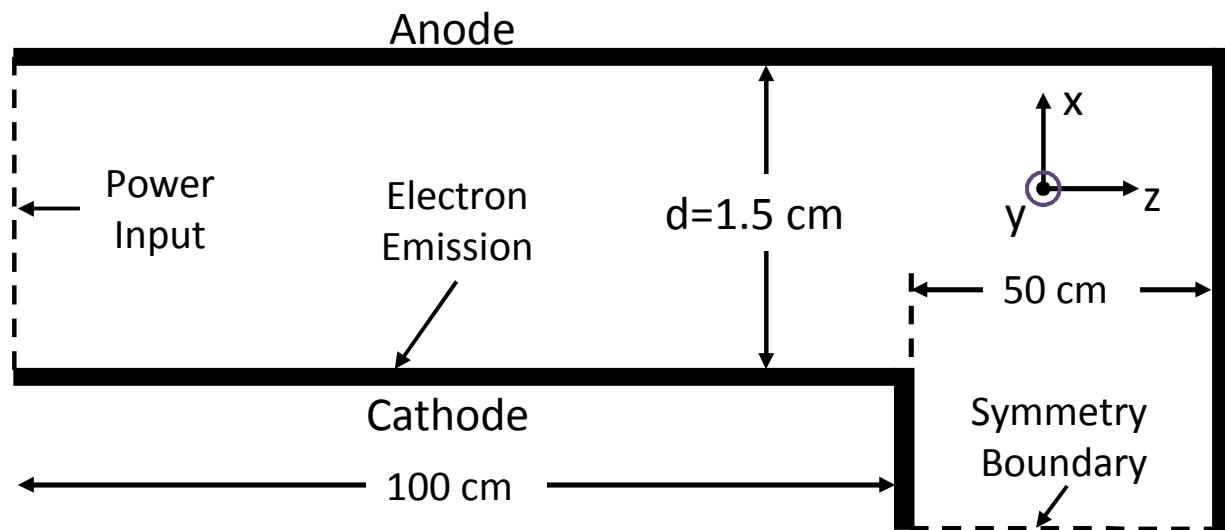
# Outline/Overview

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- We have developed a cathode plasma model for the PIC code Lsp and applied it 1D and 2D simulations magnetically insulated transmission lines (MITLs).
- The numerical implementation of the algorithm is detailed.
- The algorithm is tested by using pre-calculated MITL equilibria into which we introduce a thin cathode plasma which is allowed to evolve in time and study
- The Lsp results are qualitatively similar to a simple model in which a cathode plasma diffusion equation is coupled to laminar flow theory.
- 1D simulations exhibit simple scaling of cathode plasma thickness,  $x_s \sim t^{1/2}$ .
- In 2D, with cyclotron orbits in the simulation plane, a flute instability can develop from transverse charge separation. In this case the plasma thickness can grow much more rapidly than for a purely diffusive process.

# High-Power 2D Parallel-Plate MITL Simulations\*

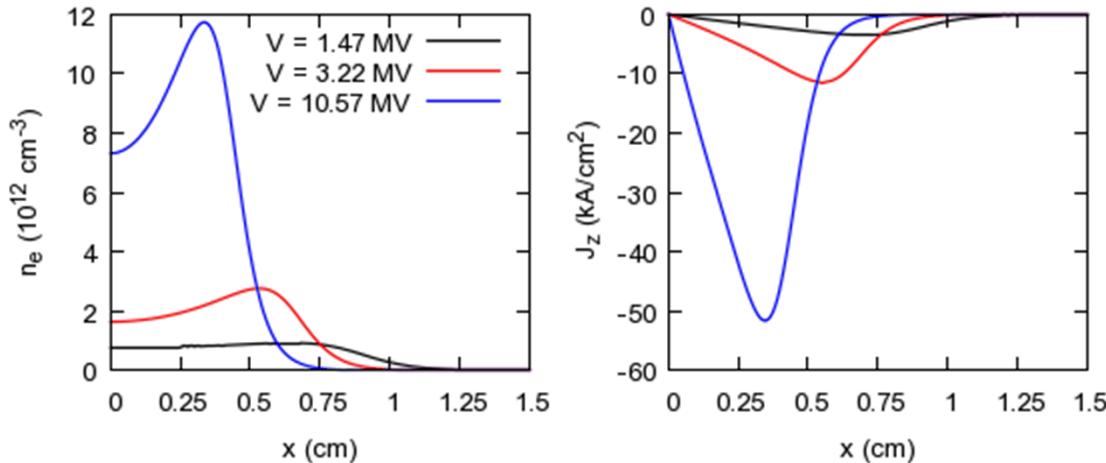
- Explicit EM, kinetic particles, 1 – 10 MV
- High resolution across the AK gap (150 cells)
- Electron flow initiated by SCL emission from cathode, electron flow develops a quasi-equilibrium sheath profile along the MITL.



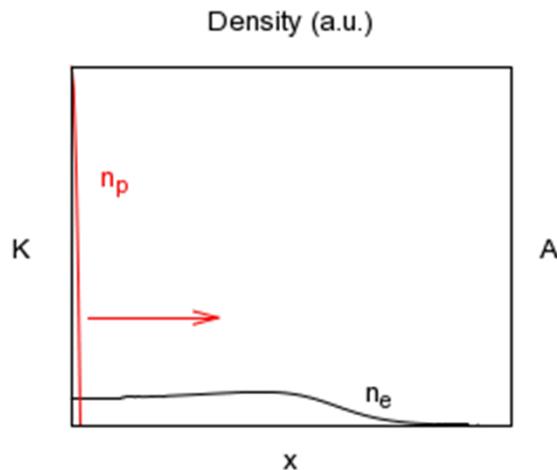
\*D.V. Rose, *et al.* "Electron Flow Stability in Magnetically Insulated vacuum transmission lines", Phys. of Plasmas **18**, 033108 (2011).

# Simulated MITL Equilibria for Voltages from 1-10 MV

- Equilibrium electron number and current density profiles for Voltages of 1.47, 3.22, 10.57 MV.
- Averaged over 10-cm section of MITL in 2D simulation

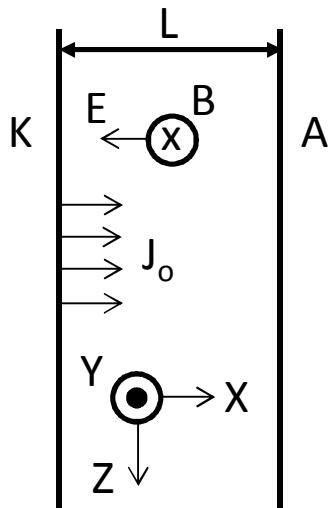


## Schematic illustration of 1D Cartesian Cathode Plasma Simulation



- Load in an initial MITL equilibrium.
- Introduce a thin cathode plasma and allow it to diffuse into the AK gap.

# Quasi-Laminar MITL Equilibria



Compare LSP with Quasi-Laminar theory of Ron, et al.\*

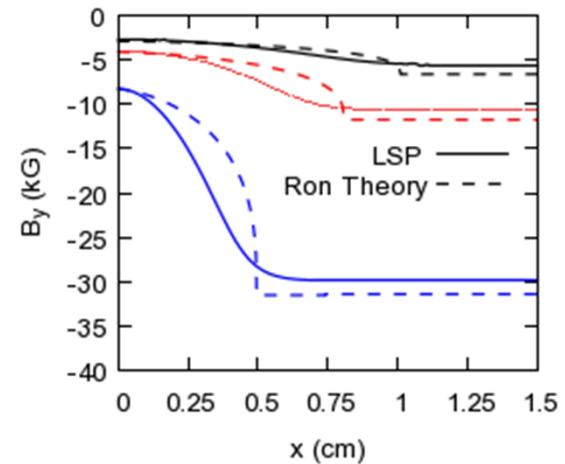
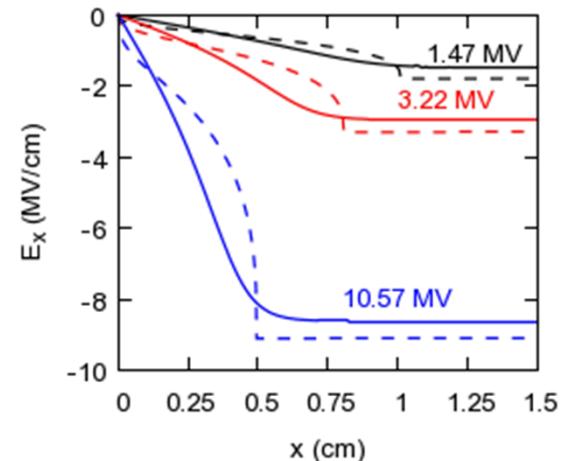
Theory provides equilibrium solutions for all fields in normalized coordinates, e.g.,

$$\eta = x \sqrt{\frac{2e\mu_o J_o}{m_e c}} \quad \beta = B \sqrt{\frac{\epsilon_o}{2m_e c \mu_o J_o}}$$

For a specified voltage and AK gap ( $L = 1.5$  cm), compare Ron equilibria and LSP for same cathode current,  $I_K$ .

Voltage (MV)	$\beta_K$	$\eta_A$	$I_K$ (kA/cm)	PIC $I_A$ (kA/cm)	RON $I_A$ (kV/cm)
1.47	2.33	1.04	2.19	4.48	4.99
3.22	2.233	1.635	3.30	8.47	9.35
10.57	2.1512	3.399	6.61	23.76	25.04

Anode currents differ by  $\sim 10\text{-}20\%$



Note that LSP fields lack discontinuities of Ron model at sheath edge

\*A. Ron, et al. "Equilibria for Magnetic Insulation", IEEE Trans. Plasma Sci **PS-1**, 85 (1973).

# Cathode Plasma Diffusion Equation

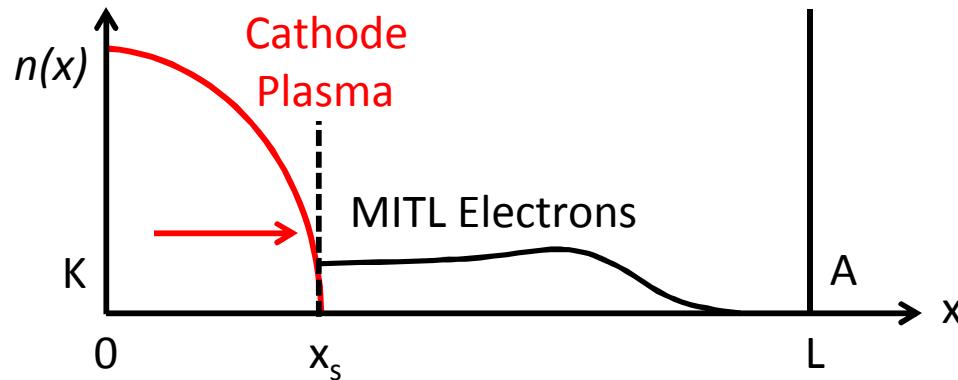
Cathode plasma: Fully ionized, Isothermal, Quasi-neutral plasma  $n = n_i = n_e/Z$

$$\frac{\partial n}{\partial t} = -\frac{\partial \Gamma}{\partial x}$$

$$\Gamma_e \square \Gamma_i = \Gamma = -\tilde{D} \frac{\partial n}{\partial x}$$

$$\tilde{D} = \frac{v_{ei}}{Zm_e\omega_{ce}^2} (T_i + ZT_e)$$

$$\frac{\partial B_y}{\partial x} = \mu_0 Ze \left( \frac{\omega_{ce}}{v_{ei}} \right) \tilde{D} \frac{\partial n}{\partial x}$$



Plasma is coupled to Ron equilibrium for MITL electrons at  $x = x_s$ .

## Numerical Implementation

Initial flat-top density profile, width  $x_s$ .

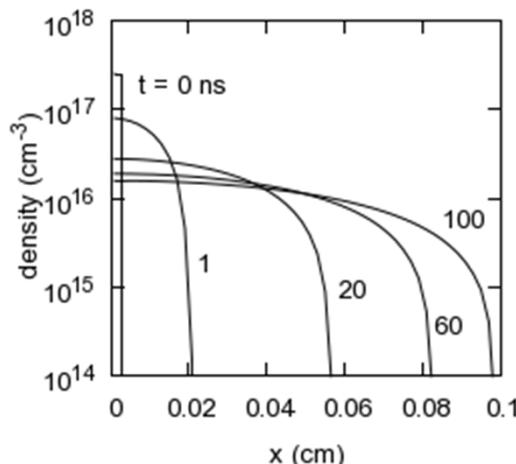
Diffusion equation is advanced.

$$\text{If: } n(x_s) \sqrt{T_e/m_e} \geq \left( \frac{m_e c}{2e \mu_0} \right) \left( \frac{\eta_A}{L - x_s} \right)^2$$

then  $x_s$  is incremented by  $\Delta x$ .

Integrate from  $x_s$  to get magnetic field in plasma, where

$$B_y(x_s) = -\left( \frac{m_e c}{e} \right) \frac{\beta_K \eta_A}{L - x_s}$$



$L = 1.5 \text{ cm}$   
 $m_i = 1$   
 $Z = 1$   
 $T_e = T_i = 5 \text{ eV}$   
Initial plasma width:  $40 \mu\text{m}$   
 $\Delta x = 20 \mu\text{m}$ ,  $\Delta t = 0.001 \text{ ns}$   
 $V = 10.57 \text{ MV}$  ( $\beta_K = 2.15$ ,  $\eta_A = 3.399$ )

$$x_s(t) \sim 0.01 \text{ cm} (t/\text{ns})^{1/3}$$

Gap closure time ( $x_s \sim L$ )  $\sim 0.1 \text{ ms}$

Roughly independent of voltage.

# Lsp Cathode Plasma Model

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- Lay in initial MITL electron kinetic particles from fitted equilibria.
- Initial static and magneto-static solve to get  $E_x$  and  $B_y$ .
- Introduce thin cathode plasma ( $\sim 1$  ML over 1-2 cells  $\Delta x \sim 10 \mu\text{m}$ ).
- The plasma is modeled by fluid electron and ion species, which are periodically remapped onto the Eulerian grid to control the particle number and avoid noise.<sup>1</sup>
- Any ion flux leaving exiting the cathode is replaced by “repopulation” particles. Reflect normal component of velocity, lose transverse. Repop particles get local fluid temp.<sup>1</sup>
- Energetic fluid electron particles are converted to kinetic electrons to resupply the electron flux at at the plasma edge (the “effective” cathode).<sup>2</sup>
- The simulations shown in the following were performed using direct implicit field solution, due to the high initial cathode plasma density ( $\omega_{pe} \Delta t \sim 1$ ).<sup>3</sup>

<sup>1</sup>C. Thoma, *et al.* Phys. Plasmas **18**, 103507 (2011) (1973).

<sup>2</sup> D. R. Welch, *et al.* Phys. Plasmas **13**, 063105 (2006).

<sup>3</sup> D. R. Welch

# 1D MITL/Cathode Plasma Simulation (10.57 MV)

$V = 10.57$  MV

$L = 1.5$  cm

$m_i = 1$

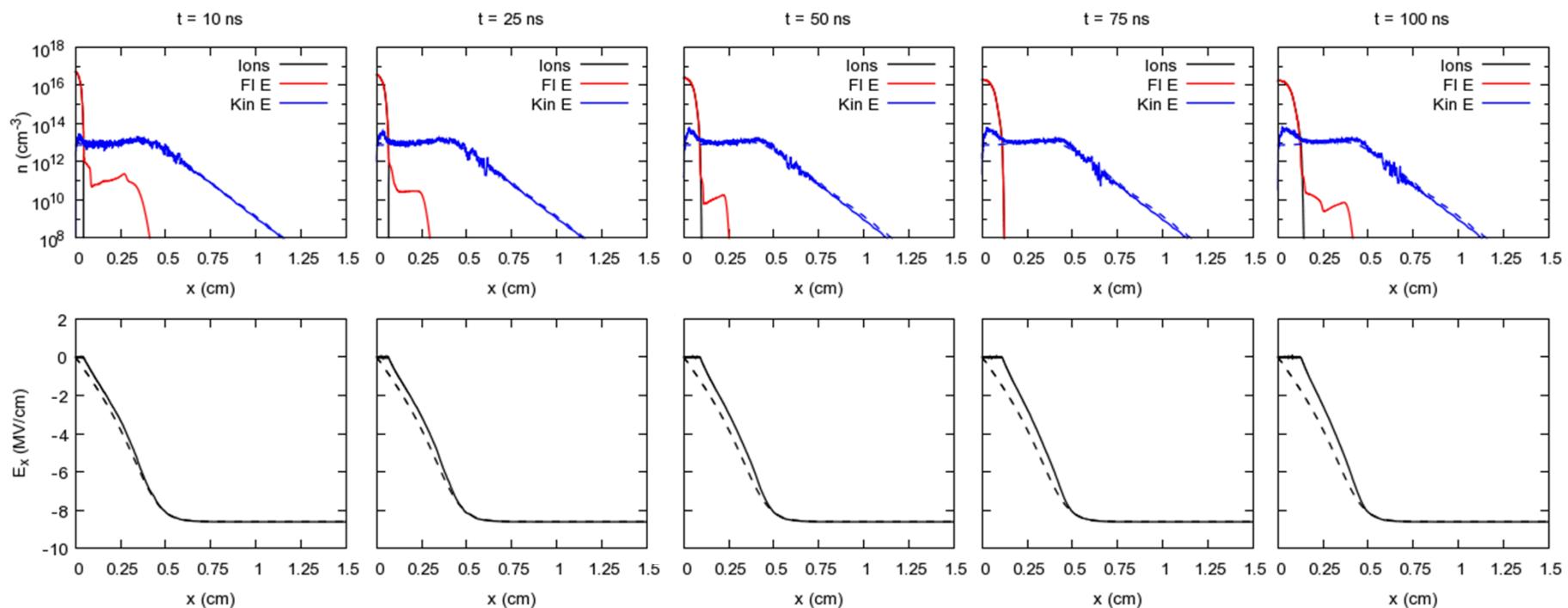
$Z = 1$

$T_e = T_i = 5$  eV

Initial plasma width: 10  $\mu$ m (1 ML)

$\Delta x = 5$   $\mu$ m,  $c \Delta t = 0.25 \Delta x$

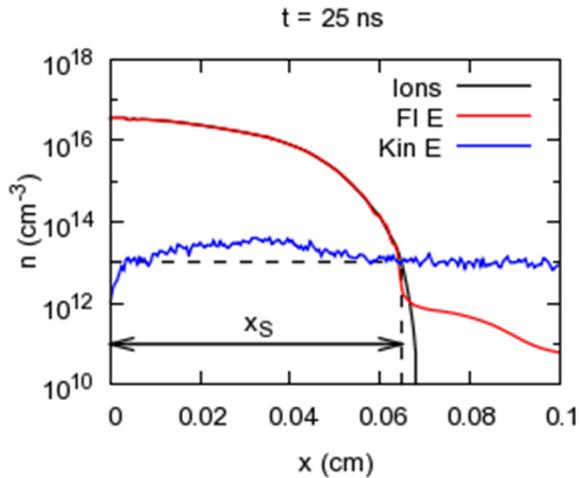
- Cathode plasma diffuses into AK gap, qualitatively similar to diffusion model.
- Electric field is screened out in plasma. Plasma edge becomes “effective” cathode.
- Electron flow in MITL sheath is only slightly perturbed for  $x_s/L \ll 1$ .



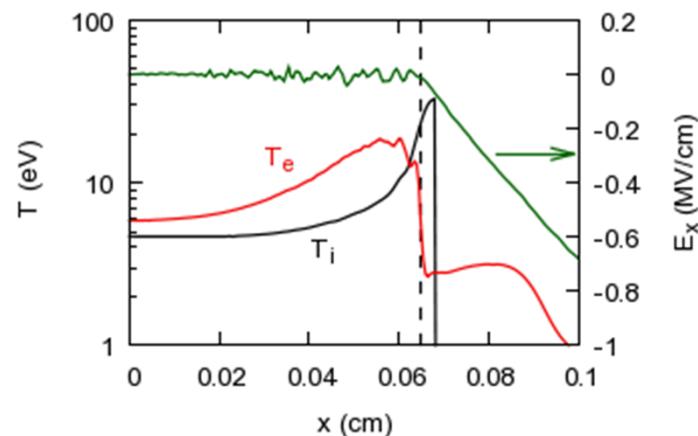
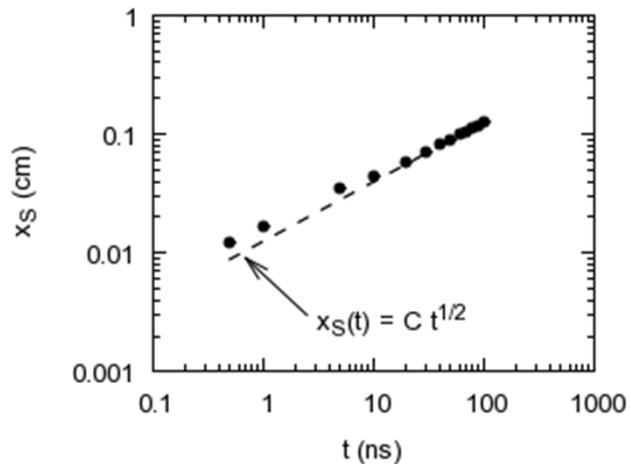
Fluid electrons with total KE > 300 eV are converted to kinetic electrons.

# Sheath edge calculation and Electrical Properties

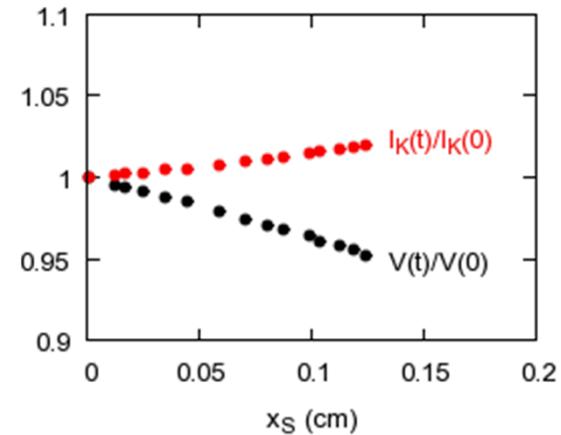
## Procedure used to calculate $x_s(t)$ .



- After a few ns  $x_s(t) \sim t^{1/2}$ , as in a classical diffusion process.
- Approximate gap closure time ( $x_s = 1.5 \text{ cm}$ )  $\sim 10 \mu\text{s}$



- Both  $V$  and  $I_K$  vary linearly with  $x_s$  for  $x_s \ll L$ .



Recall that diffusion model used isothermal approximation.

# 1D MITL/Cathode Plasma Simulations

$V = 1.47, 3.22$ , and  $10.57$  MV.

$L = 1.5$  cm

$m_i = 1$

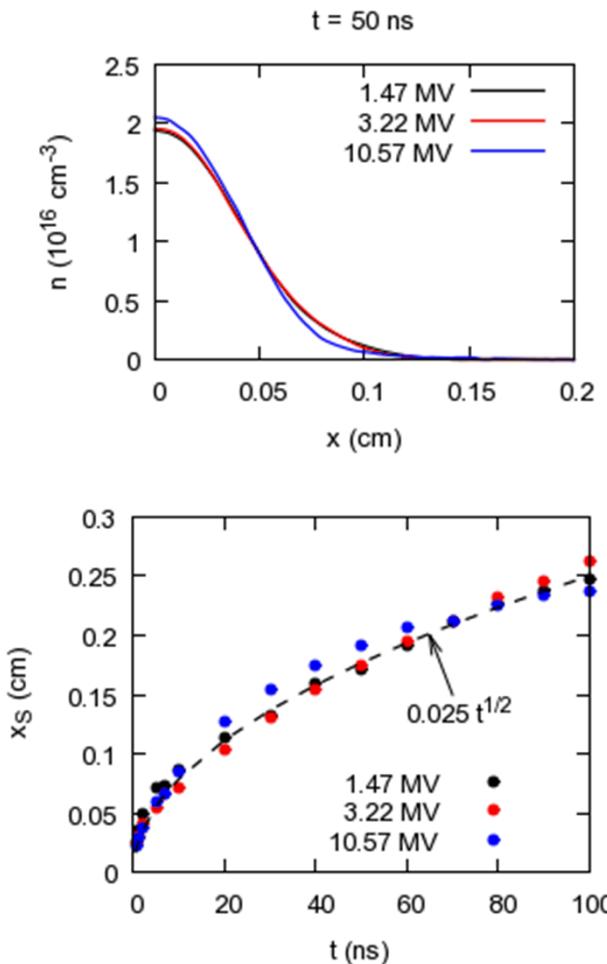
$Z = 1$

$T_e = T_i = 8$  eV

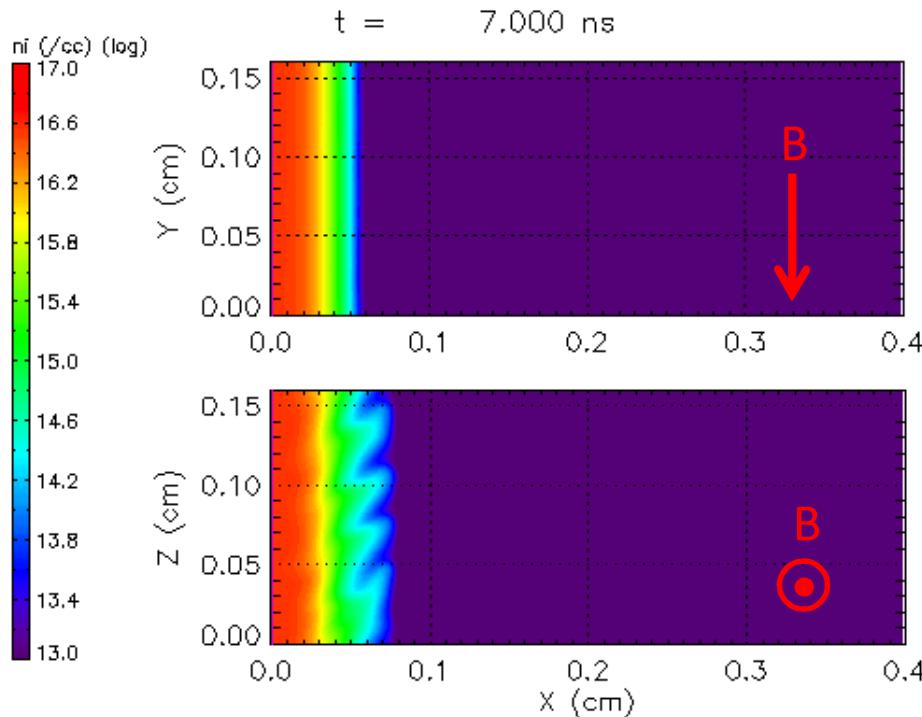
Initial plasma width:  $40$   $\mu$ m (1 ML)

$\Delta x = 20$   $\mu$ m,  $c \Delta t = 0.25 \Delta x$

- Very little variation in plasma density profile with voltage (in the bulk of the cathode plasma and at edge,  $E = 0$ ).
- For all three voltages  $x_s(t) \sim C t^{1/2}$ .
- Constant  $C$  can vary slightly with migration parameters, and begins to increase with cell size for  $x > 10-20$   $\mu$ m. But  $t^{1/2}$  scaling is very robust.

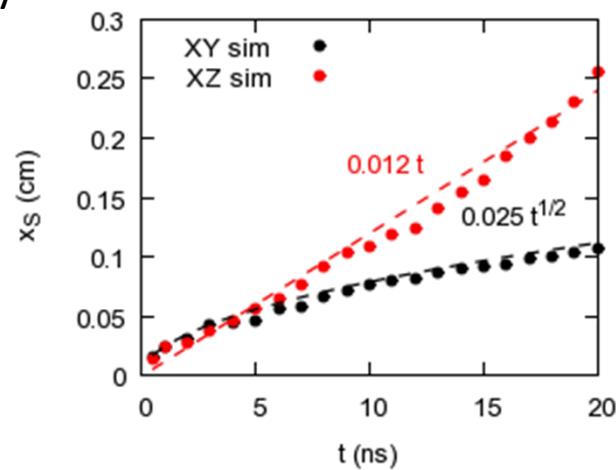


# 2D Cartesian Simulations: XY and XZ



- This instability in multiple dimensions allows much faster expansion of the cathode plasma than the 1D diffusive process predicts.

- Repeat 10.57 simulation in 2D Cartesian coordinates: Both XY and XZ.
- Periodic boundaries in Y(Z).
- YZ coordinates look essentially like 1D analog.
- In XZ coordinates, charge separation due to opposite cyclotron motions for electrons and ions is captured in the simulation plane and can seed a flute instability.



# Conclusions

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- We have described the developed a cathode plasma model for the PIC code Lsp and applied it 1D and 2D simulations magnetically insulated transmission lines (MITLs).
- The detailed PIC algorithm is in qualitative agreement with a simplified diffusion model, and predicts plasma expansion of  $t^{1/2}$ .
- The algorithm is tested by using pre-calculated MITL equilibria into which we introduce a thin cathode plasma which is allowed to evolve in time and study
- The Lsp results are qualitatively similar to a simple model in which a cathode plasma diffusion equation is coupled to laminar flow theory.
- 1D simulations exhibit simple scaling of cathode plasma thickness,  $x_s \sim t^{1/2}$ .
- In 2D simulations, with cyclotron orbits in the simulation plane, a flute instability can develop from transverse charge separation. In this case the plasma thickness can grow much more rapidly than for a purely diffusive process.
- In future we hope to consider larger-scale 2D and 3D simulations with both anode and cathode plasmas modeled self-consistently.