

# **Modeling Differential Diffusion of Soot with Conditional Moment Closure**

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**Sandia National Laboratories**

**Workshop on Conditional Moment Closure  
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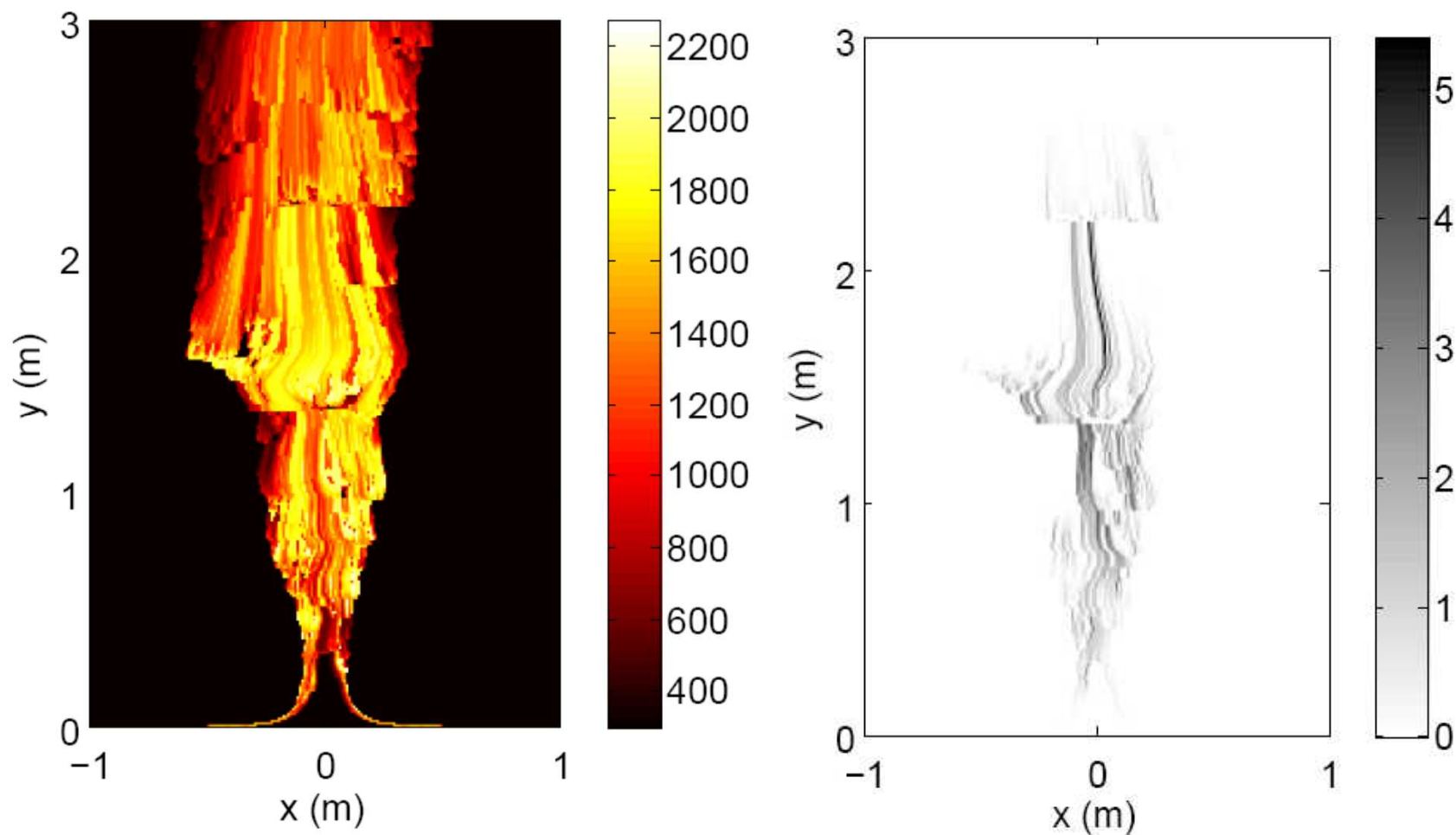


# Introduction

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- **Conserved scalar approaches: Flamelets, CMC**
  - Typically derived for unity Lewis numbers
- **Differential diffusion alters relationship between reacting and conserved scalars.**
  - Previous work by Kronenberg (1997), Pitsch (1998)
  - Focus here on soot.
- **New look at CMC formulation**
  - Choices made in derivation can give slightly different formulations, require different modeling.
  - Different terms can be easier or more difficult to model depending on the scalar of interest.
- **A priori analysis looks at significant terms and models for different terms.**

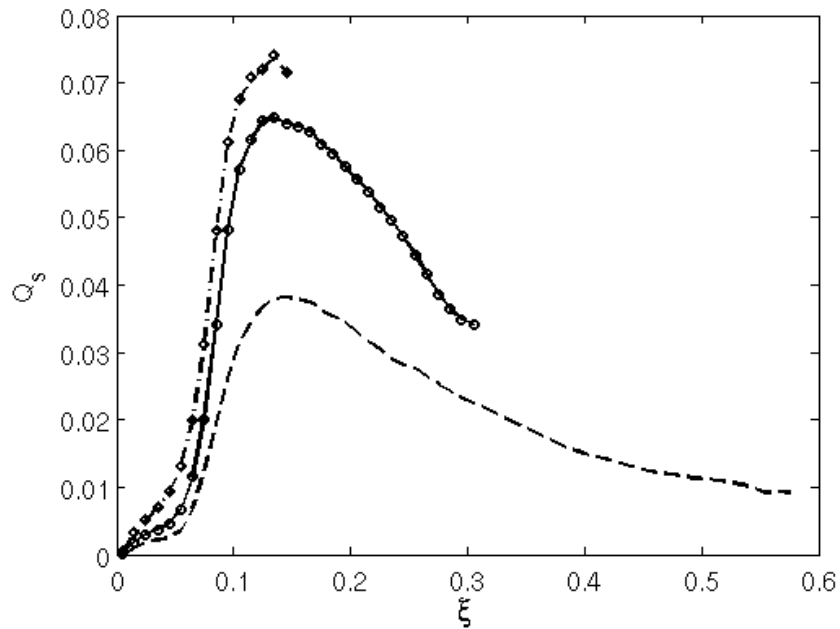
## Higher fidelity data to check closures



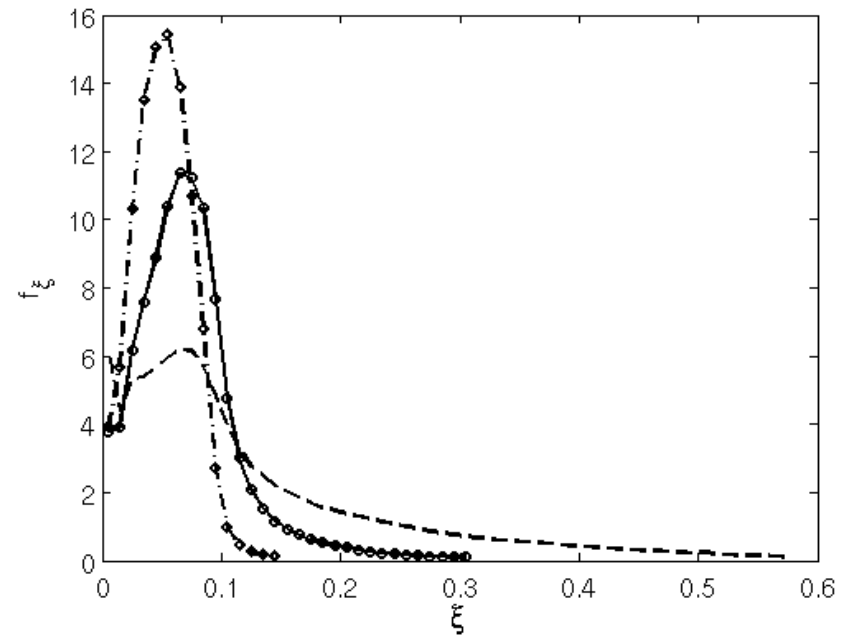


# ODT results

Soot mass fractions



Mixture fraction PDF



heights are 0.9, 1.4 and 1.9 source widths



# General species conservation

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$$\frac{\partial \rho Y_i}{\partial t} + \nabla \cdot (\rho \vec{u} Y_i) - \nabla \cdot (\rho D_i \nabla Y_i) + \nabla \cdot (\rho \vec{V}_i' Y_i) = \rho \omega_i$$

- **Allows for different diffusion models, thermophoresis, etc. for species, enthalpy, aerosols (soot).**



## Joint-pdf equation

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$$\frac{\partial \rho Y_i}{\partial t} + \nabla \cdot (\rho \vec{u} Y_i) - \nabla \cdot (\rho D_i \nabla Y_i) + \nabla \cdot (\rho \vec{V}_i' Y_i) = \rho \omega_i$$

- From species equation, derive equation for joint-pdf evolution

$$\begin{aligned} \frac{\partial \left[ \left\langle \rho Y_k \middle| \vec{Z} \right\rangle P(\vec{Z}) \right]}{\partial t} + \nabla \cdot \left[ \left\langle \rho \vec{u} Y_k \middle| \vec{Z} \right\rangle P(\vec{Z}) \right] = & \sum_i \frac{\partial}{\partial Z_i} \left[ \left\langle \rho \omega_i \middle| \vec{Z} \right\rangle P(\vec{Z}) \right] \\ & - \sum_i \frac{\partial}{\partial Z_i} \left[ \left\langle \nabla \cdot (\rho D_i \nabla Y_i) \middle| \vec{Z} \right\rangle P(\vec{Z}) \right] \\ & + \sum_i \frac{\partial}{\partial Z_i} \left[ \left\langle \nabla \cdot (\rho \vec{V}_i' Y_i) \middle| \vec{Z} \right\rangle P(\vec{Z}) \right] \end{aligned}$$



# Replacing diffusion terms with dissipation terms

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- **Mathematical identity used to replace diffusion terms in joint-pdf equation with scalar dissipation terms.**

$$\nabla \cdot \left[ \left\langle \rho D_k \middle| \vec{Z} \right\rangle \nabla P(\vec{Z}) \right] = - \sum_i \frac{\partial}{\partial Z_i} \left[ \left\langle \nabla \cdot (\rho D_k \nabla Y_i) \middle| \vec{Z} \right\rangle P(\vec{Z}) \right] \\ + \sum_i \sum_j \frac{\partial^2}{\partial Z_i \partial Z_j} \left[ \left\langle D_k \nabla Y_i \nabla Y_j \middle| \vec{Z} \right\rangle P(\vec{Z}) \right]$$

- **Can use any diffusion coefficient in this expression to replace any of the terms in the joint-pdf equation.**
- **Different choices result in different forms of conditional moment equations when differential diffusion is allowed.**



# Analyzing closure requirements for traditional CMC

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- **Conditional moment equation derivation**
  - **Multiply joint PDF by scalar of interest,  $Y_k$  ; integrate over PDF except mixture fraction,  $\eta$**

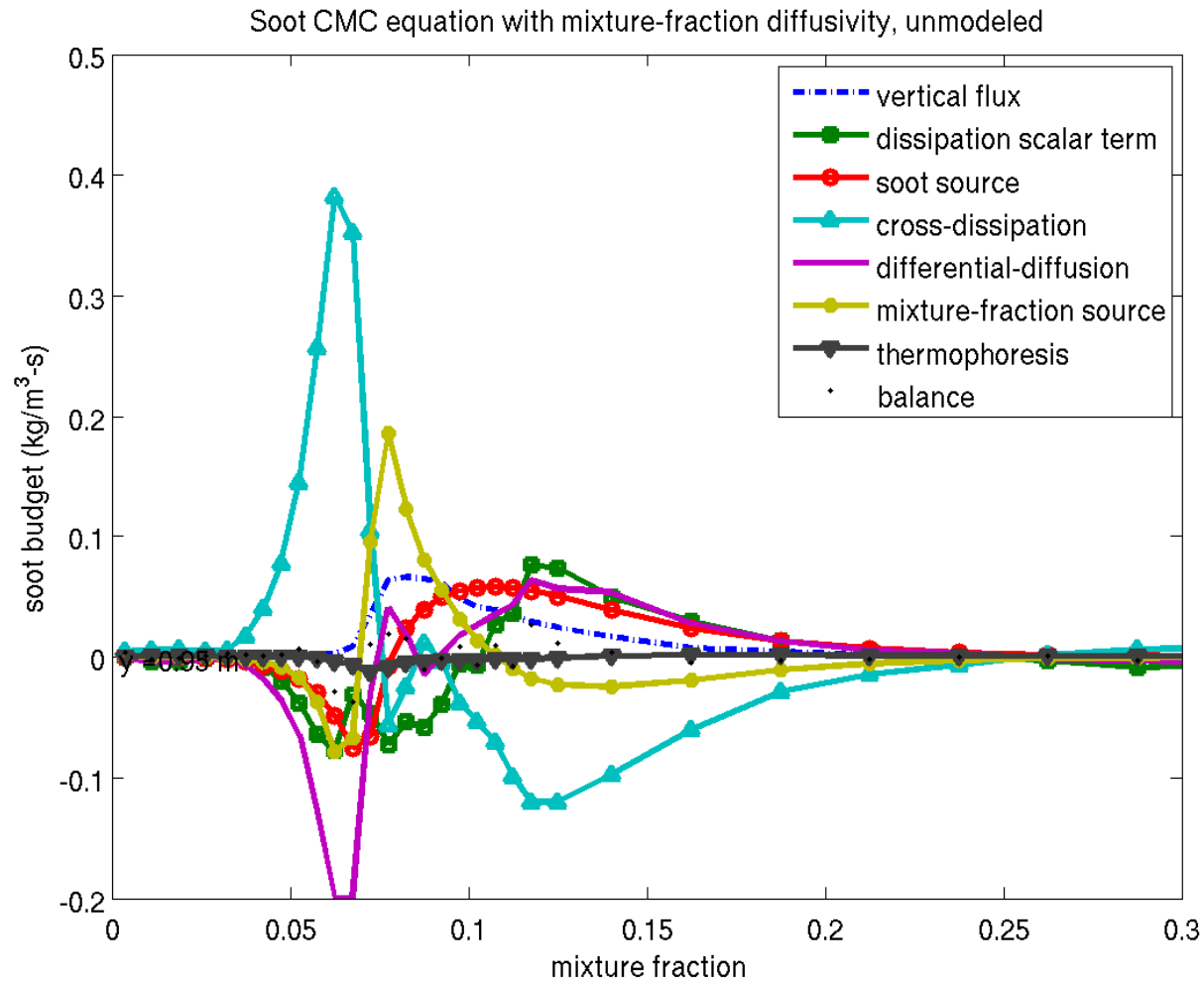
$$\begin{aligned}\nabla \cdot [\langle \rho \bar{u} Y_k | \vec{\eta} \rangle P(\vec{\eta})] &= \langle \rho \omega | \vec{\eta} \rangle P(\vec{\eta}) \\ &+ \frac{\partial}{\partial \eta} [\langle \rho D_\xi \nabla \xi \nabla Y_k | \vec{\eta} \rangle P(\vec{\eta})] \\ &- \frac{\partial^2}{\partial \eta^2} [\langle \rho D_\xi (\nabla \xi)^2 Y_k | \vec{\eta} \rangle P(\vec{\eta})] \\ &+ \langle \nabla \cdot (\rho D_T Y_k \nabla \ln T) | \vec{\eta} \rangle P(\vec{\eta}) \\ &- \langle \nabla \cdot (\rho (D_k - D_\xi) \nabla Y_k) | \vec{\eta} \rangle P(\vec{\eta})\end{aligned}$$

- **Equations are exact here: no modeling**





# Significant terms in traditional CMC



# Cross-dissipation term in traditional CMC

- Standard approach for closure

$$\langle XY|\eta\rangle \approx \langle X|\eta\rangle\langle Y|\eta\rangle$$

$$\frac{\partial Y_s}{\partial x_i} = \frac{\partial Y_s}{\partial \xi} \frac{\partial \xi}{\partial x_i}$$

- Basic term

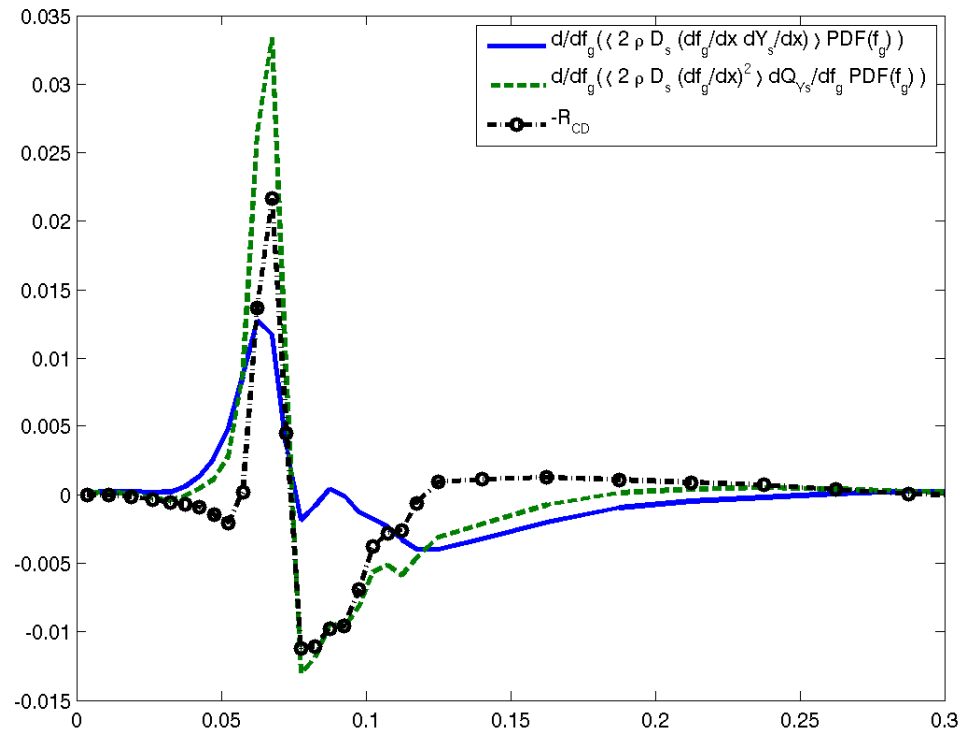
$$\frac{\partial}{\partial \eta} \left[ \left\langle \rho D_\xi \nabla \xi \nabla Y_k \right| \vec{\eta} \right] P(\vec{\eta})$$

- Standard closure

$$\frac{\partial}{\partial \eta} \left[ \left\langle \rho D_\xi (\nabla \xi)^2 \right| \vec{\eta} \right] \frac{\partial Q_k}{\partial \eta} P(\vec{\eta})$$

- Residual (difference between these) is as big as standard closure

Compare standard closure  
and residual term  
(after differentiation)



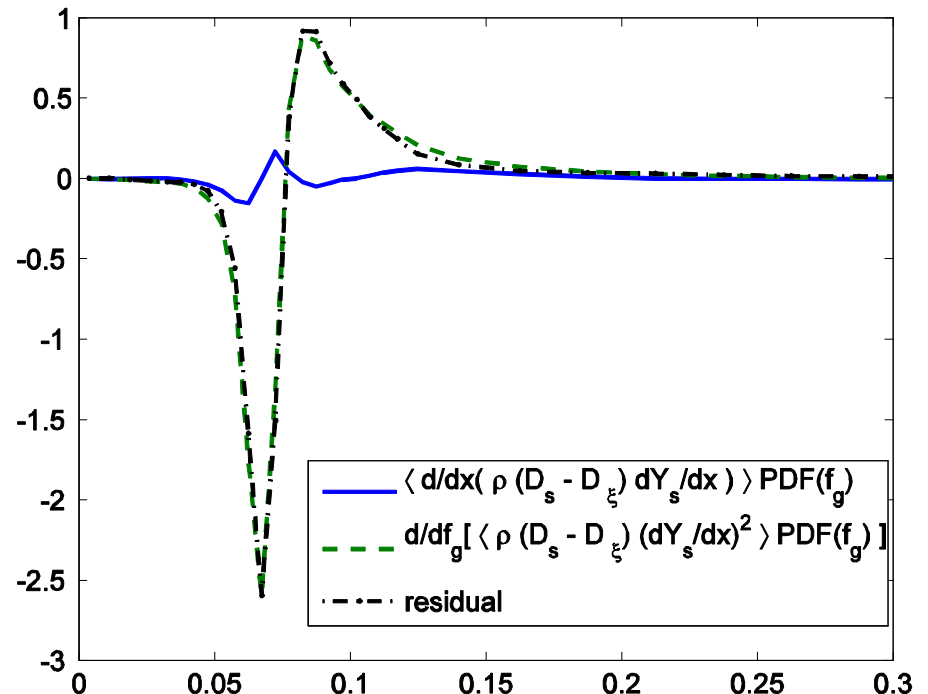
# Differential-diffusion term in traditional CMC

$$\left\langle \nabla \left[ \rho (D_k - D_\xi) \nabla Y_k \right] \middle| \vec{\eta} \right\rangle P(\vec{\eta})$$

- **Standard closure would use soot-dissipation rate.**

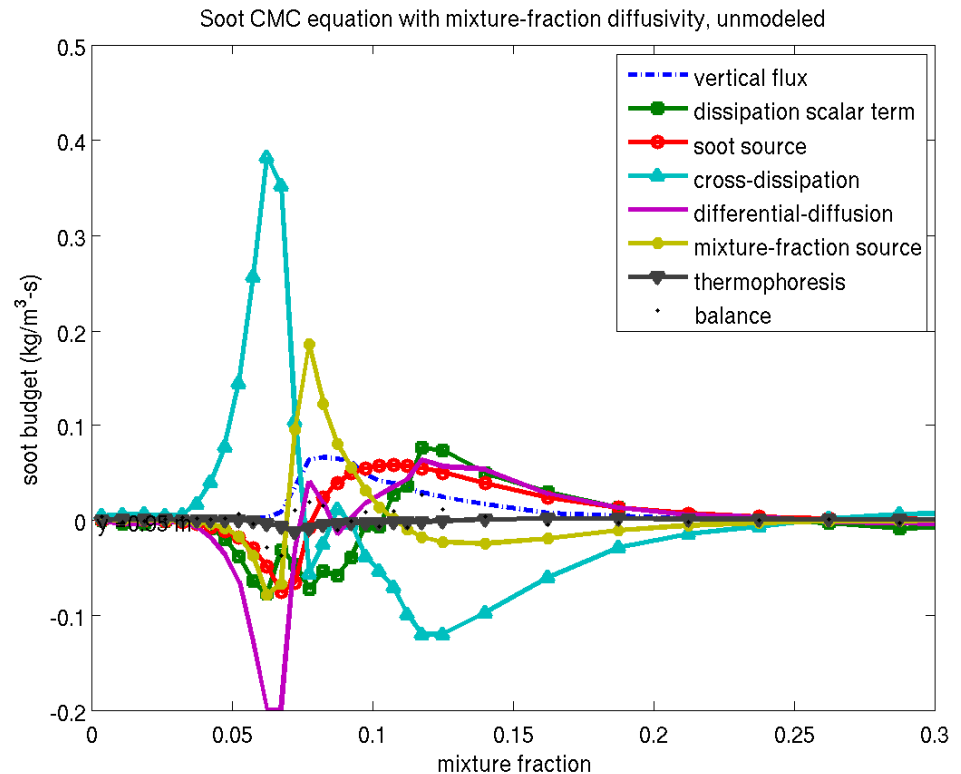
$$\begin{aligned} & \left\langle \nabla \left[ \rho (D_k - D_\xi) \nabla Y_k \right] \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \\ & \approx \frac{\partial}{\partial \eta} \left[ \left\langle \rho (D_k - D_\xi) (\nabla Y_k)^2 \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right] \end{aligned}$$

- **Statistics for high-Schmidt-number reacting-scalar dissipation rates are less available.**



# So what about traditional CMC

- **Standard closures generally result in many large residual terms.**
  - At least some of these residual terms will be difficult to model.
  - Fortunately, there is some cancellation of the residuals.

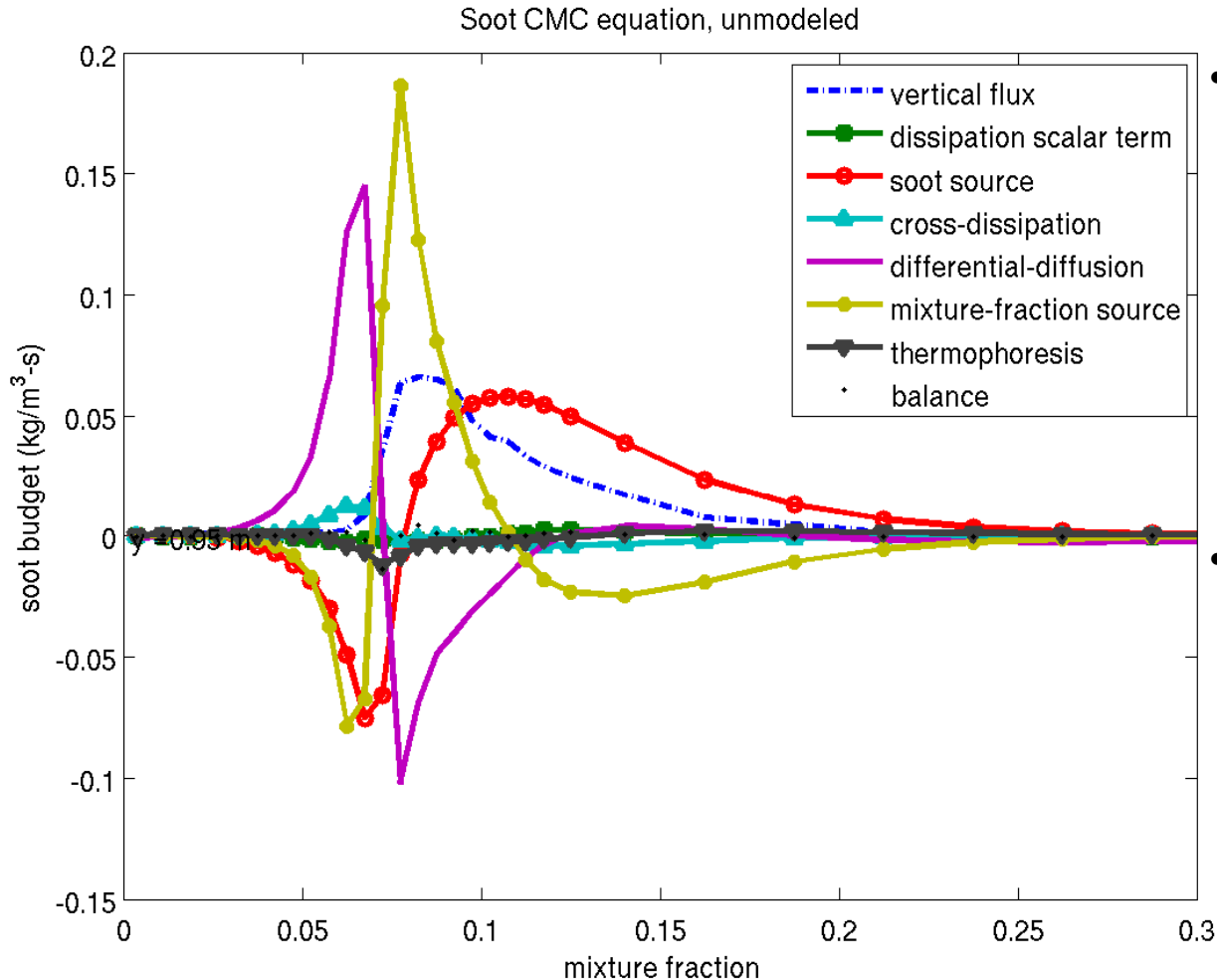


## A alternate form of the CMC equations

- It is possible to eliminate different terms in the process of deriving the conditional moment equations.
- Resulting CM equation has  $D_k$ , (for soot) in most places.
  - Many diffusive terms will be small.

$$\begin{aligned}\nabla \cdot \left[ \left\langle \rho \vec{u} Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right] = & \left\langle \rho \omega \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \\ & + \frac{\partial}{\partial \eta} \left[ \left\langle \rho D_k \nabla \xi \nabla Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right] \\ & - \frac{\partial^2}{\partial \eta^2} \left[ \left\langle \rho D_k (\nabla \xi)^2 Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right] \\ & + \left\langle \nabla \cdot (\rho D_T Y_k \nabla \ln T) \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \\ & - \frac{\partial}{\partial \eta} \left[ \left\langle \nabla \cdot (\rho (D_\xi - D_k) \nabla \xi) Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right]\end{aligned}$$

# Significant terms in alternate CMC



- **Soot diffusivity causes dissipation-scalar, cross-dissipation terms to be small.**
  - Residuals remain but are trivially small.
- **Differential diffusion is still substantial.**
  - Standard model involves mixture fraction dissipation.

# Differential-diffusion term in alternate CMC

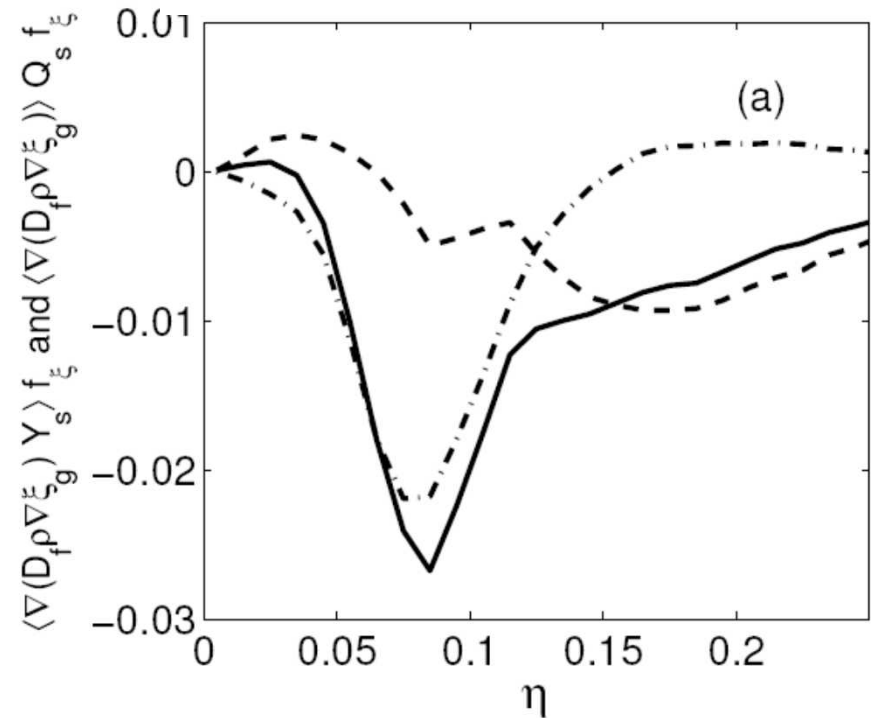
$$\frac{\partial}{\partial \eta} \left[ \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right]$$

- **Standard closure:**

$$\begin{aligned} & \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \\ & \approx \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] \middle| \vec{\eta} \right\rangle P(\vec{\eta}) Q_k \\ & = \frac{\partial}{\partial \eta} \left( \left\langle \rho (D_\xi - D_k) (\nabla \xi)^2 \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right) Q_k \end{aligned}$$

- **Residual is significant.**

**Compare standard closure  
(before differentiation)  
and residual**





## So what about alternate CMC

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- **Standard closures generally result in fewer large residual terms.**
  - **Need to look at differential diffusion terms. What is the physics behind these terms?**

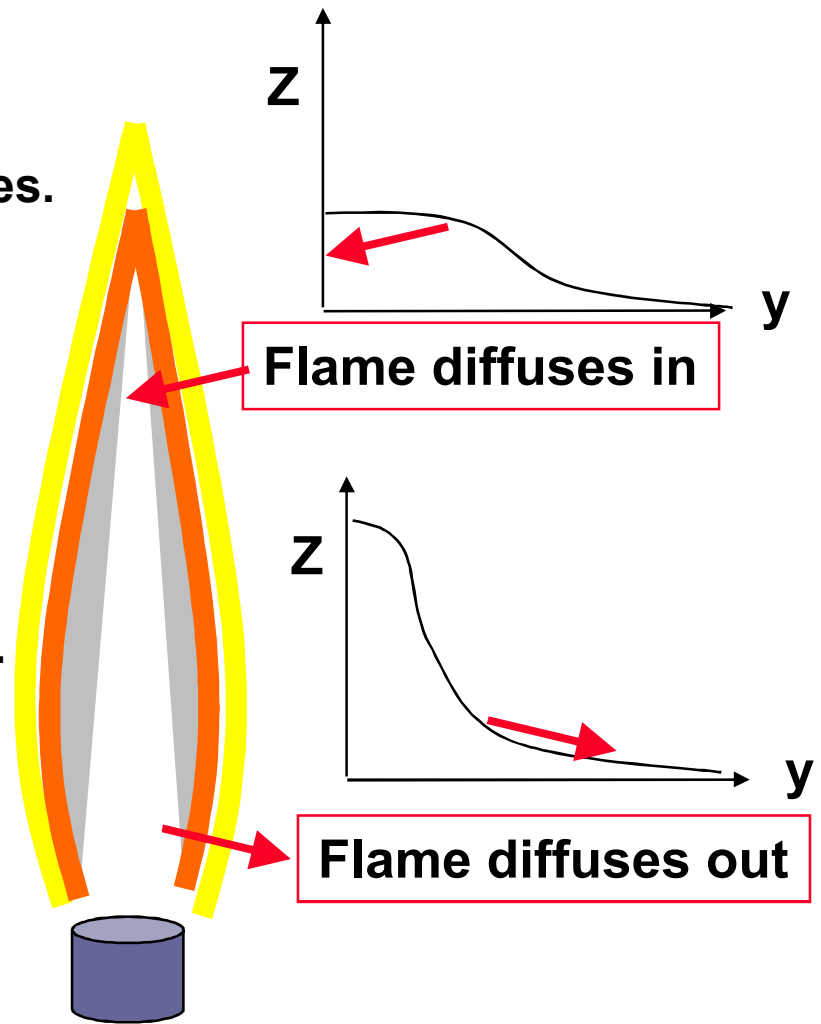
$$\frac{\partial}{\partial \eta} \left[ \left\langle \nabla \square (\rho(D_\xi - D_k) \nabla \xi) Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right] \approx \frac{\partial}{\partial \eta} \left[ \left\langle \nabla \square (\rho(D_\xi - D_k) \nabla \xi) \middle| \vec{\eta} \right\rangle P(\vec{\eta}) Q_k \right]$$

$$R_{DD} = \frac{\partial}{\partial \eta} \left[ \left\langle \nabla \square (\rho(D_\xi - D_k) \nabla \xi) Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) - \left\langle \nabla \square (\rho(D_\xi - D_k) \nabla \xi) \middle| \vec{\eta} \right\rangle P(\vec{\eta}) Q_k \right]$$

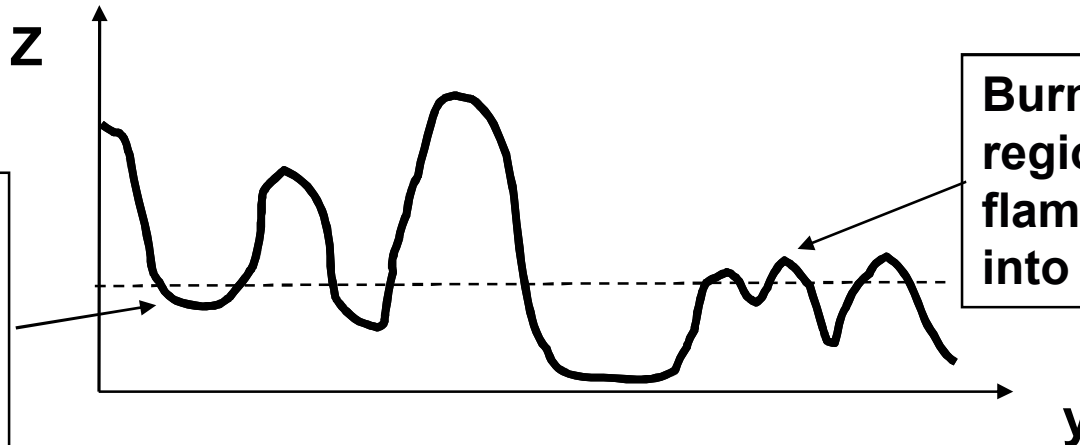
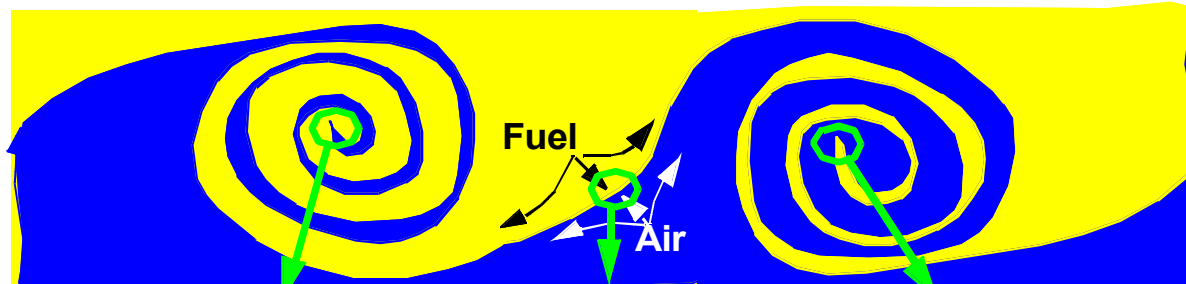


# Differential diffusion of soot in flames due to mean mixture fraction evolution

- Soot diffusion is
  - Slow relative to gas-phase species.
  - Affected by thermophoresis, etc.
- Soot is convected along with everything else.
- Flames diffuse!
  - Diffuse towards location of nearest mixture fraction extrema.
  - Candle analogy suggests mean behavior: Related to evolution of mixture fraction PDF.



# Differential diffusion of soot in flames due to mixture fraction fluctuations



**Flame motion across fluid/soot elements is random walk about a mean**

# Advection due to mixture fraction fluctuations

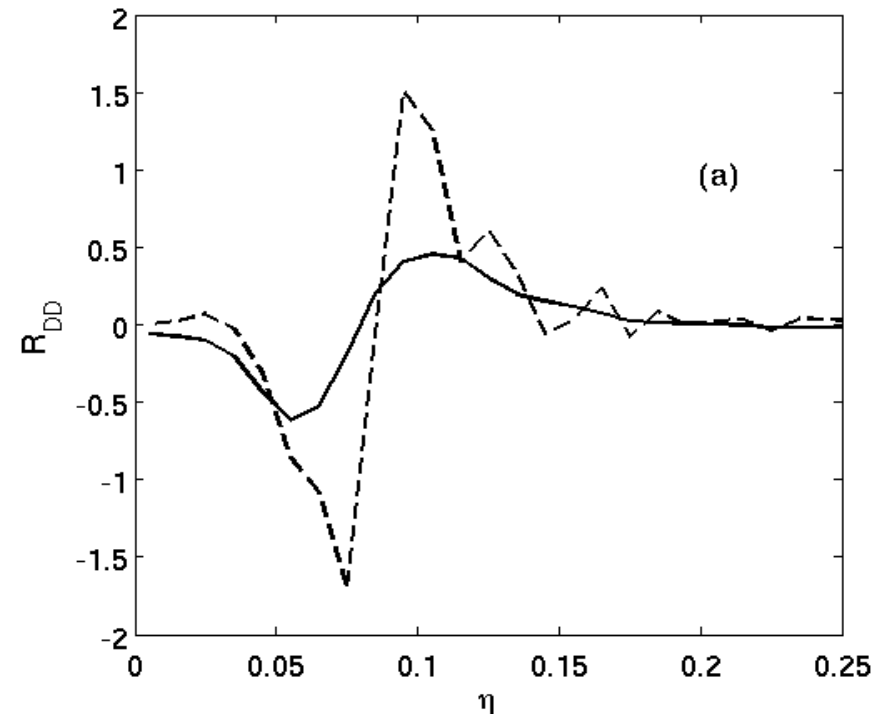
- Additional terms for diff-diff closure:

$$R_{DD} = \frac{\partial}{\partial \eta} \left[ \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) - \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] \middle| \vec{\eta} \right\rangle P(\vec{\eta}) Q_k \right]$$

- The residual term describes the effect of fluctuating “advective” flows in mixture fraction space.

- Correlation between random mixture fraction fluctuations and soot moments.
- Model with diffusive process?

$$R_{DD} \approx \frac{\rho_\eta \chi_\eta f_\xi}{2Le_{DD,t}} \frac{\partial^2 Q_s}{\partial \eta^2}$$



# Advection due to mixture fraction fluctuations

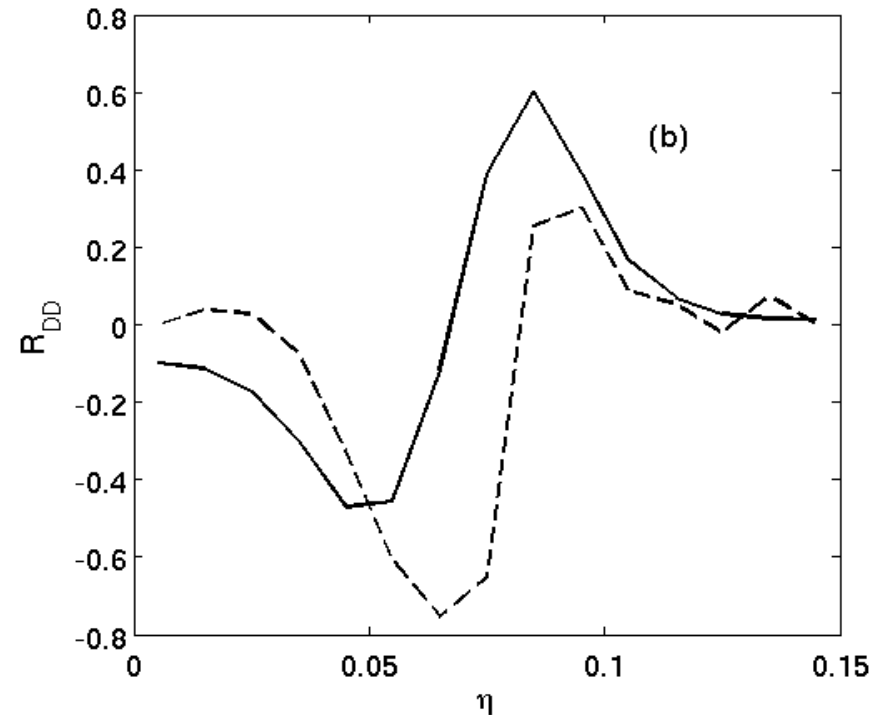
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# Advection due to mixture fraction fluctuations

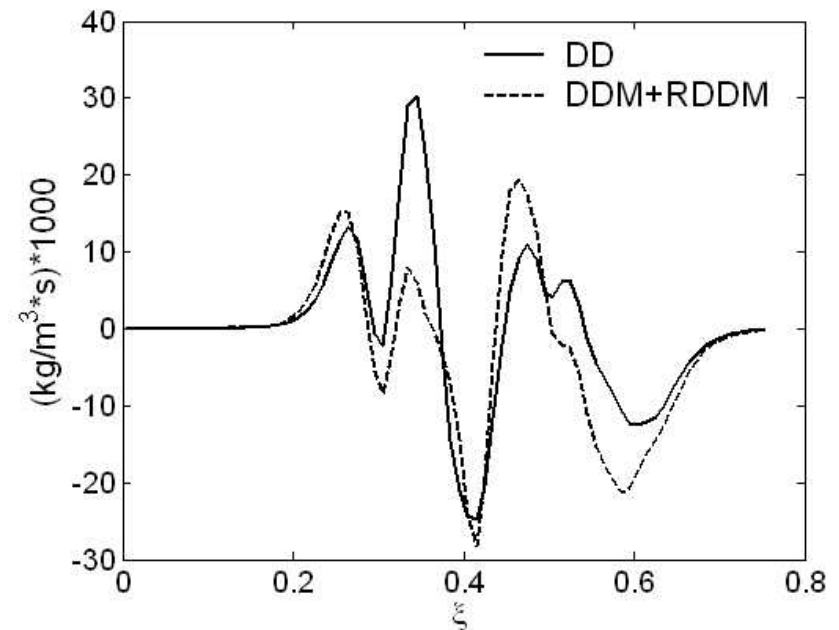
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- The residual term describes the effect of fluctuating “advective” flows in mixture fraction space.

- Correlation between random mixture fraction fluctuations and soot moments.
- Model with diffusive process?

$$R_{DD} \approx \frac{\rho_\eta \chi_\eta f_\xi}{2Le_{DD,t}} \frac{\partial^2 Q_s}{\partial \eta^2}$$





# Summary

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- **Conditional moment equations extended to allow general diffusion models.**
- **This opens options as to what terms need to be modeled.**
- **For soot at least, the alternate approach presented here appears preferable:**
  - **Differential diffusion of soot relative to mixture fraction dominates.**
  - **Mean component of transport related to PDF(Z) evolution.**
  - **Fluctuating component related to dissipative processes using new model.**
- **Results with ODT study suggest model captures relevant physics.**





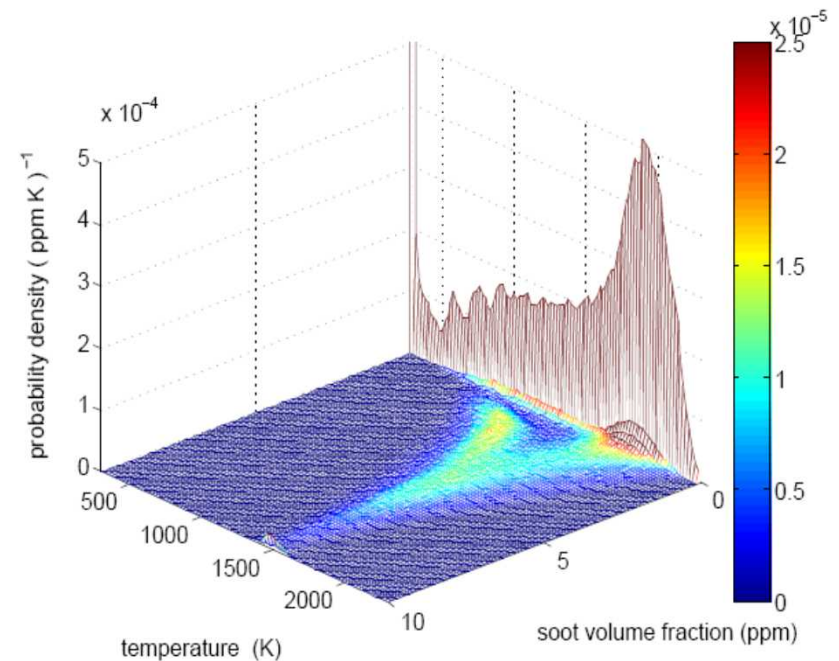
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# Backup slides



# ODT simulations provide high-fidelity data to evaluate closures

- Buoyant 1 m wide ethene plume (line fire) spatially evolving ODT simulation.
- Simple soot model (Fairweather *et al.* 1992) with steady laminar flamelet source terms tabulated by enthalpy and mixture fraction.
- Generate statistical quantities like soot-temperature joint PDF.







## Additional key closure requirements for CMC

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- **For particles, differential diffusion is key physics.**
- “New” CMC formulation retains a residual fluctuation term.
- Retains conservative form with PDF

Conditionally  
Averaged

$$\begin{aligned} \frac{\partial(\rho_\eta Q_i f_\xi)}{\partial t} + \nabla \cdot (\langle \rho \vec{u} | \eta \rangle Q_i f_\xi) = & \langle \rho \omega_i | \eta \rangle f_\xi + \frac{\rho_\eta \chi_\eta f_\xi}{2Le_i} \frac{\partial^2 Q_i}{\partial \eta^2} - \frac{\partial^2}{\partial \eta^2} \left( \frac{\rho_\eta \chi_\eta f_\xi}{2} \right) Q_i \\ & - \left( \frac{D_\xi - D_i}{D_\xi} \right) \frac{\partial}{\partial \eta} \left( \frac{\rho_\eta \chi_\eta f_\xi}{2} \right) \frac{\partial Q_i}{\partial \eta} - R_{DD} \end{aligned}$$

- These terms describe the effect of the mean and fluctuating “convective” flows in mixture fraction space.
  - Diffusion of mixture fraction relative to soot



## Additional key closure requirements for CMC

---

- For particles, differential diffusion is key physics.
- “New” CMC formulation retains a residual fluctuation term.

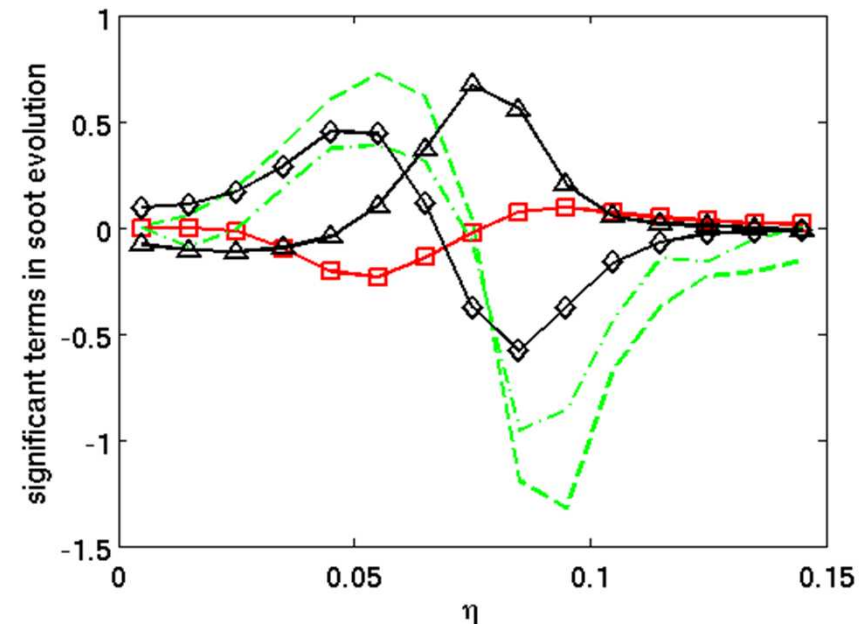
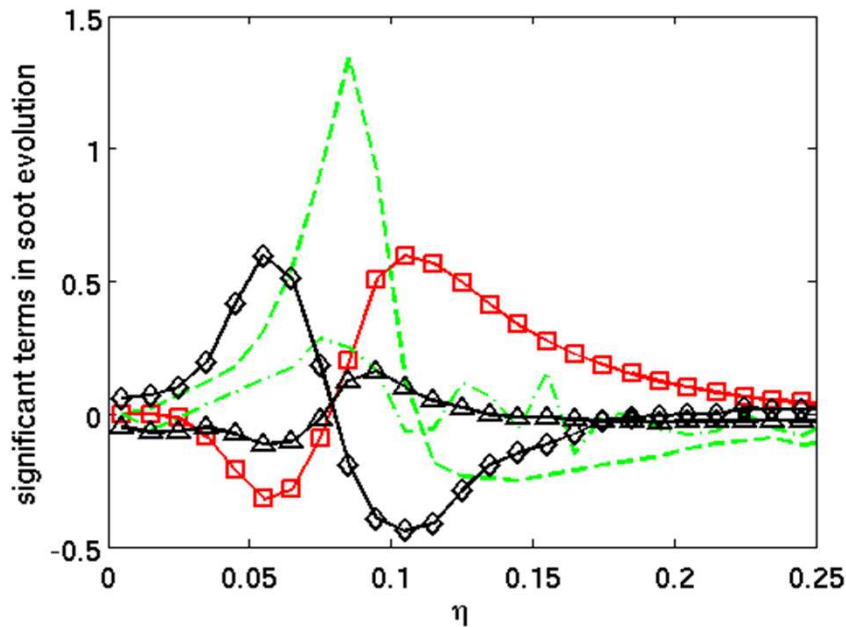
Conditionally  
Averaged

$$\frac{\partial(\rho_\eta Q_i f_\xi)}{\partial t} + \nabla \cdot (\langle \rho \vec{u} | \eta \rangle Q_i f_\xi) = \langle \rho \omega_i | \eta \rangle f_\xi + \frac{\rho_\eta \chi_\eta f_\xi}{2Le_i} \frac{\partial^2 Q_i}{\partial \eta^2} - \frac{\partial^2}{\partial \eta^2} \left( \frac{\rho_\eta \chi_\eta f_\xi}{2} \right) Q_i$$
$$- \left( \frac{D_\xi - D_i}{D_\xi} \right) \frac{\partial}{\partial \eta} \left( \frac{\rho_\eta \chi_\eta f_\xi}{2} \right) \frac{\partial Q_i}{\partial \eta} - R_{DD}$$

- Boxed term represents the mean evolution of the mixture composition

# ODT results

- Terms plotted below for heights in ODT simulations where mixture fraction pdf is centered on production (left) and on oxidation (right).



Advection (dash), pdf flux (dash-dot) -- long-term evolution of soot.

Soot source (squares).

Diff-diff by evolution of pdf (triangles) -- long-time advection in mixture fraction.

Diff-diff fluctuations  $R_{DD}$  (diamonds) -- short-time diffusion in mixture fraction.

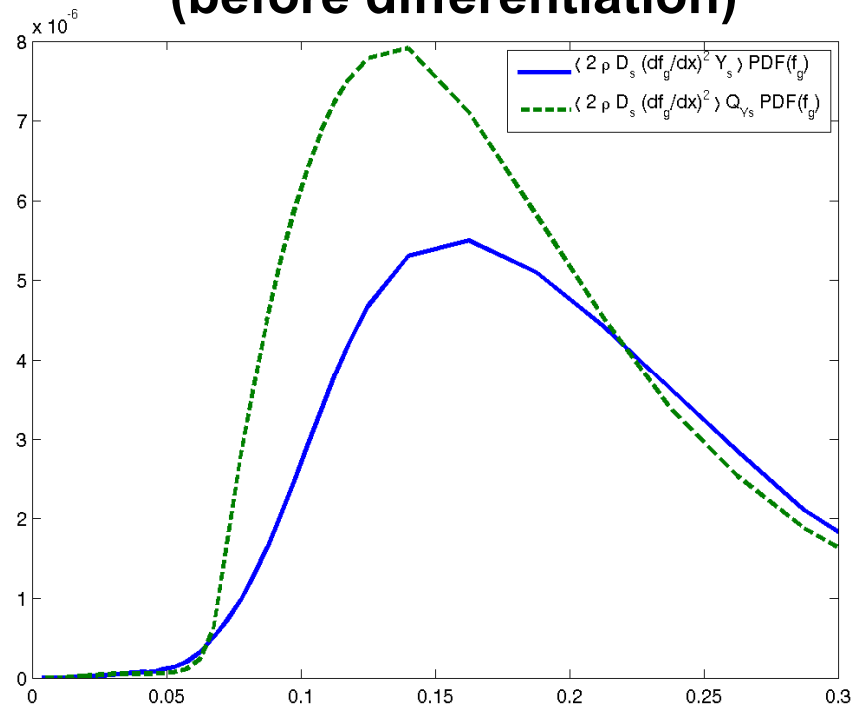
# Dissipation-scalar term in Klimenko-Bilger CMC

$$\frac{\partial^2}{\partial \eta^2} \left[ \left\langle \rho D_\xi (\nabla \xi)^2 Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \right]$$

- **Standard closure**

$$\begin{aligned} & \left\langle \rho D_\xi (\nabla \xi)^2 Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta}) \\ &= \left\langle \rho D_\xi (\nabla \xi)^2 \middle| \vec{\eta} \right\rangle Q_k P(\vec{\eta}) \end{aligned}$$

**Compare standard closure  
(before differentiation)**



# Dissipation-scalar term in Klimenko-Bilger CMC

- Basic term

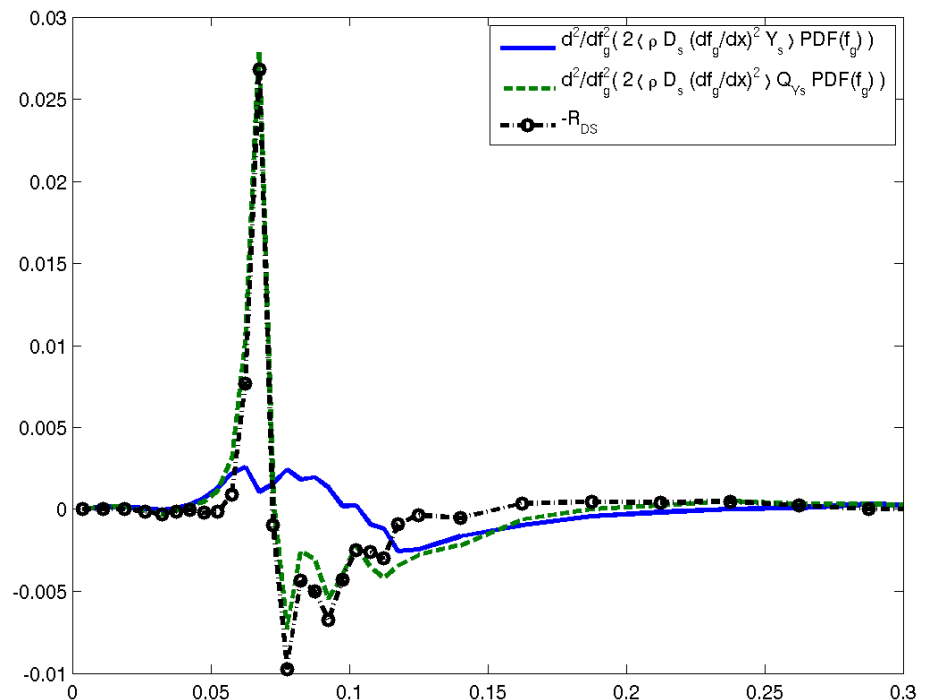
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- Standard closure

$$\frac{\partial^2}{\partial \eta^2} \left[ \left\langle \rho D_\xi (\nabla \xi)^2 \middle| \vec{\eta} \right\rangle Q_k P(\vec{\eta}) \right]$$

- Residual (difference between these) is as big as standard closure

Compare standard closure  
and residual term  
(after differentiation)



# Cross-dissipation term in Klimenko-Bilger CMC

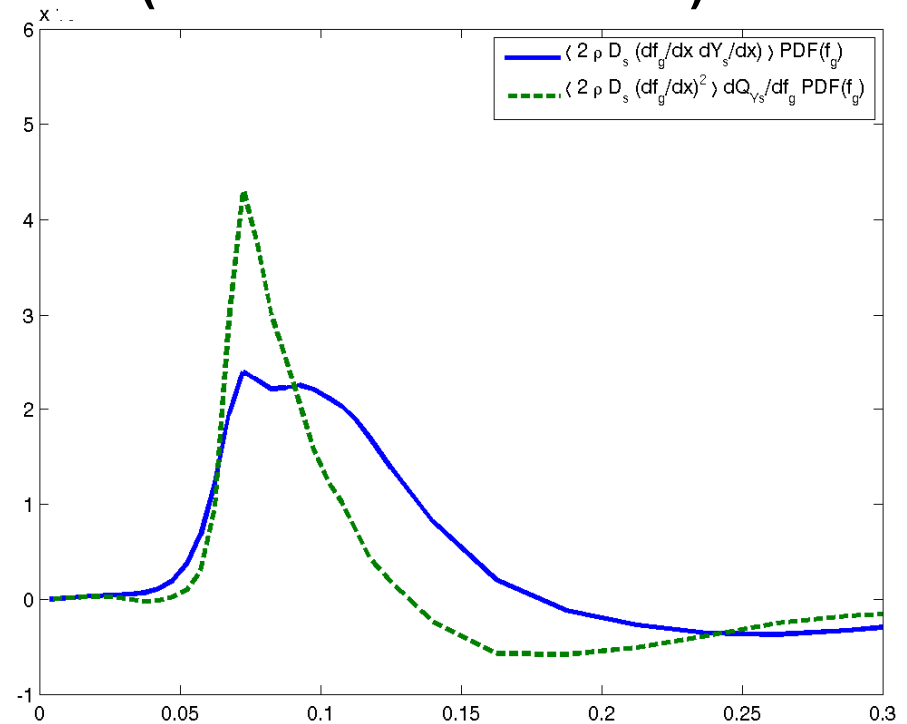
$$\frac{\partial}{\partial \eta} \left[ \left\langle \rho D_{\xi} \nabla \xi \nabla Y_k \mid \vec{\eta} \right\rangle P(\vec{\eta}) \right]$$

- **Standard closure**

$$\left\langle \rho D_{\xi} \nabla \xi \nabla Y_k \mid \vec{\eta} \right\rangle P(\vec{\eta})$$

$$\approx \left\langle \rho D_{\xi} (\nabla \xi)^2 \mid \vec{\eta} \right\rangle \frac{\partial Q_k}{\partial \eta} P(\vec{\eta})$$

**Compare standard closure  
(before differentiation)**



# Cross-dissipation term in Klimenko-Bilger CMC

- Basic term

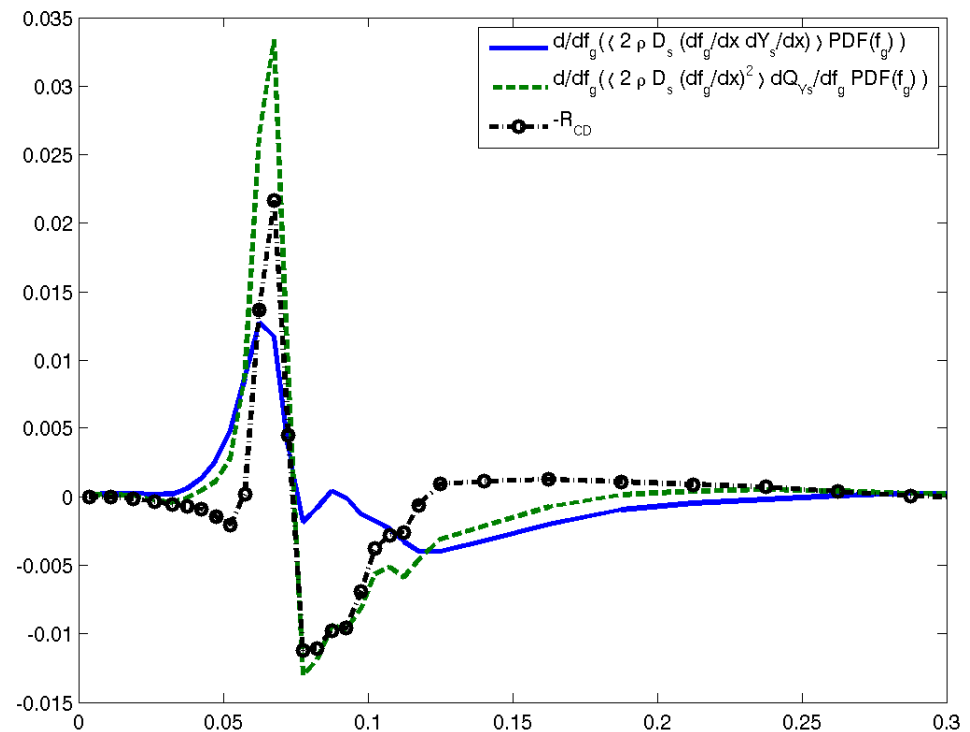
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- Standard closure

$$\frac{\partial}{\partial \eta} \left[ \left\langle \rho D_{\xi} (\nabla \xi)^2 \mid \vec{\eta} \right\rangle \frac{\partial Q_k}{\partial \eta} P(\vec{\eta}) \right]$$

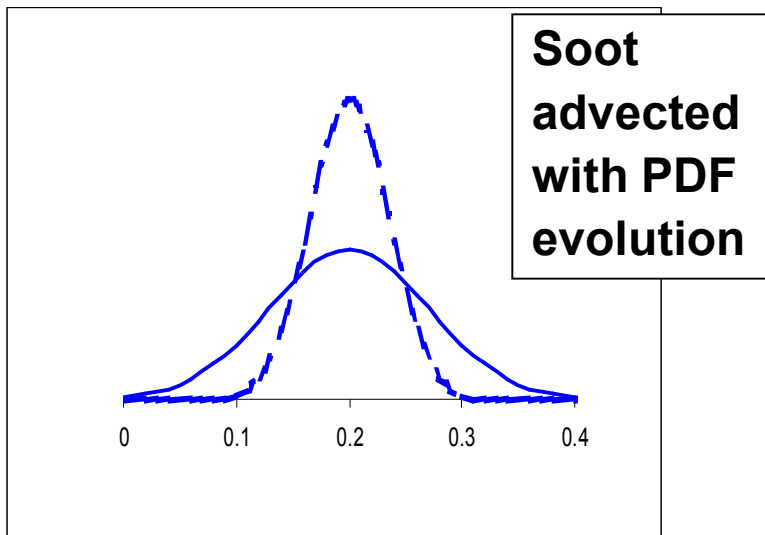
- Residual (difference between these) is as big as standard closure

Compare standard closure  
and residual term  
(after differentiation)



# Soot advection in mixture fraction space

- How does soot advection work in turbulent flows?
  - In CMC-based approach there is advection in the mixture fraction coordinate.
  - Candle analogy suggests mean behavior: Related to evolution of mixture fraction PDF.



$$\frac{\partial}{\partial \eta} \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] Y_k \middle| \vec{\eta} \right\rangle P(\vec{\eta})$$
$$\approx \frac{\partial}{\partial \eta} \left[ \left\langle \nabla \left[ \rho (D_\xi - D_k) \nabla \xi \right] \middle| \vec{\eta} \right\rangle P(\vec{\eta}) Q_k \right]$$