



Comparison of thermal conductivity and thermal boundary conductance sensitivities in continuous-wave and ultrashort pulsed thermoreflectance analyses

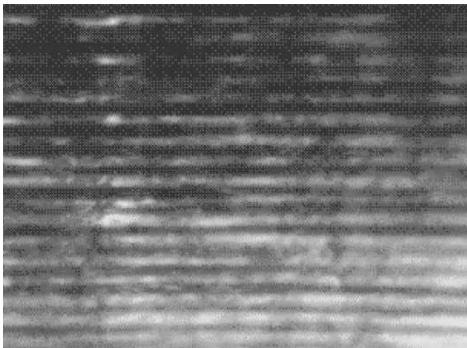
Patrick E. Hopkins, Justin R. Serrano, Leslie M. Phinney

Sandia National Laboratories, Albuquerque, NM

**17th Symposium on Thermophysical Properties
Monday, June 22, 2008**

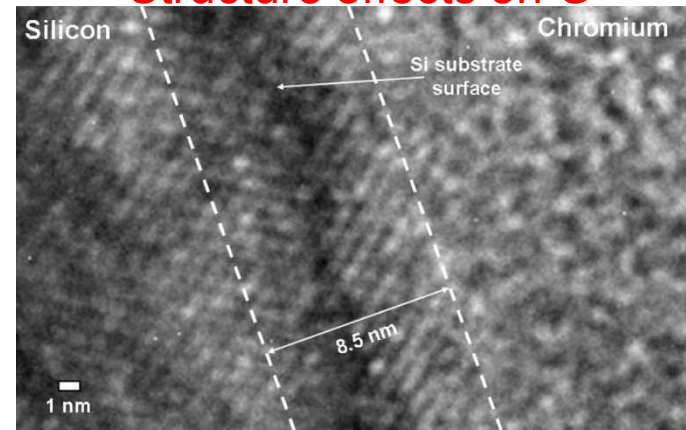
Sensitive measurements of thermal properties of nanostructures

Superlattice structures for TE



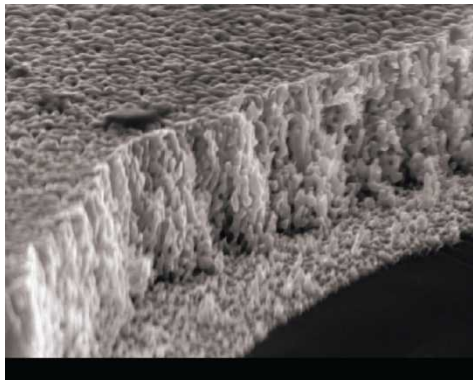
Touzelbaev *et al.*, JAP, 90, 763 (2001)

Structure effects on G



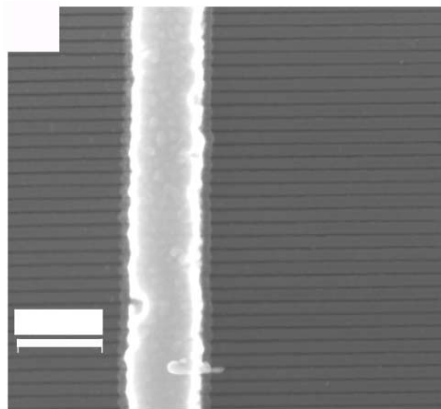
Hopkins *et al.*, JHT, 130, 062402 (2008)

Nanoporous films



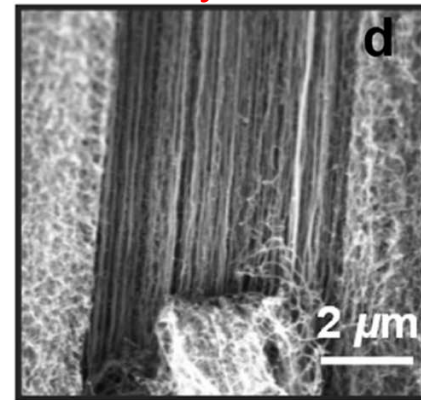
Hopkins *et al.*, Journal of Nanomaterials, **2008**, 418050 (2008)

Nanowires



Boukai *et al.*, Nature, **451**, 168 (2008)

CNT arrays for TIMS



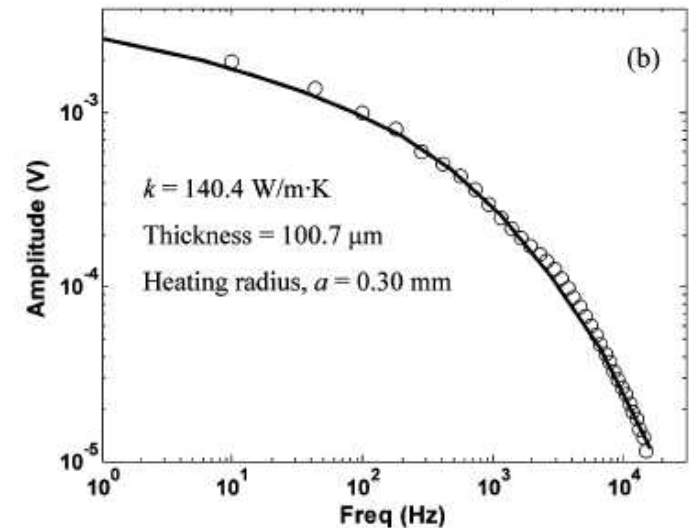
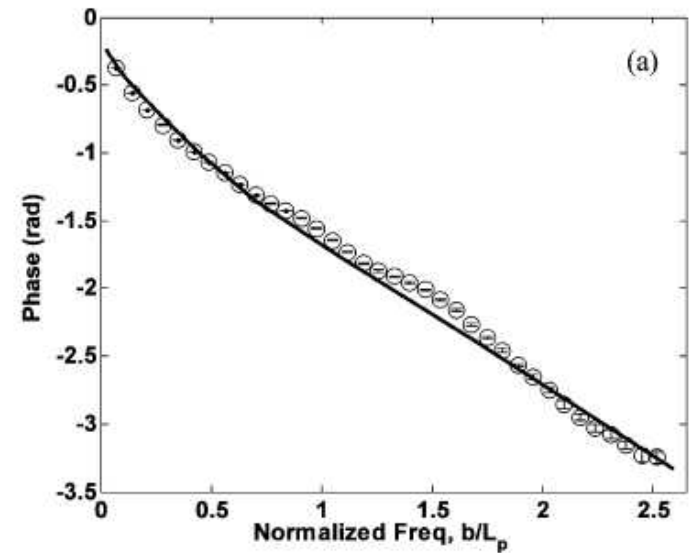
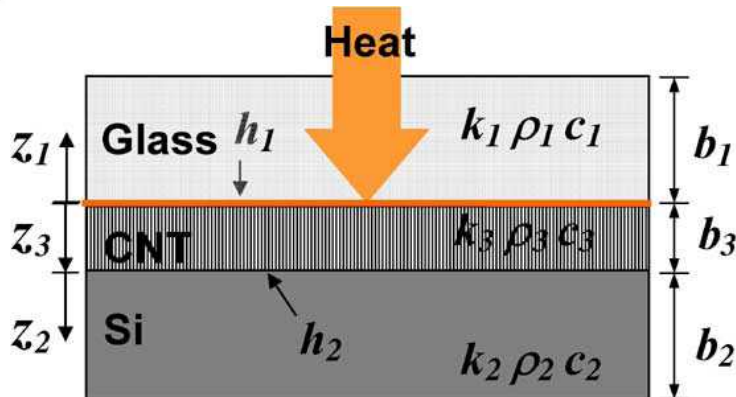
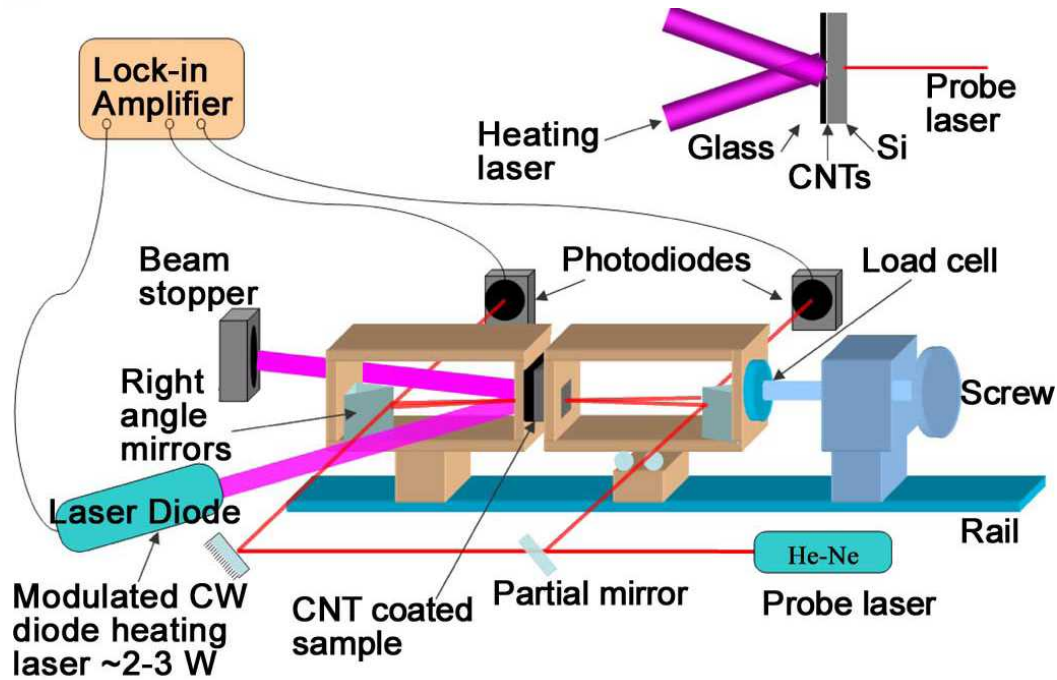
Tong *et al.*, IEEE Trans. Comp. & Pack. Tech, **30**, 92 (2007)



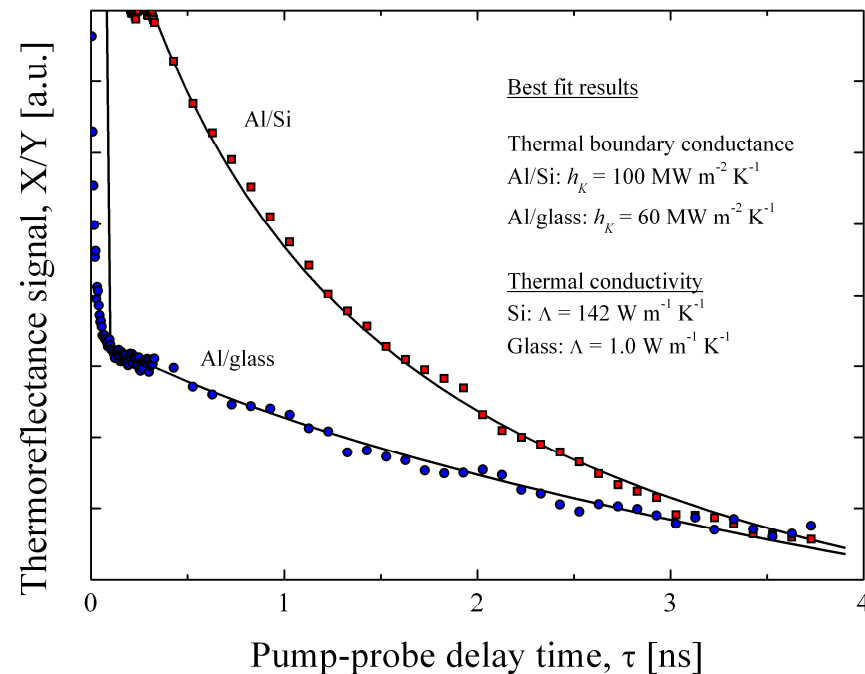
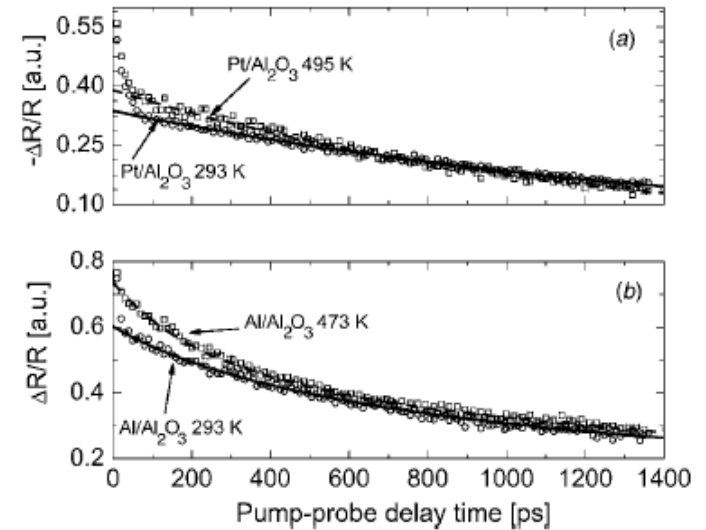
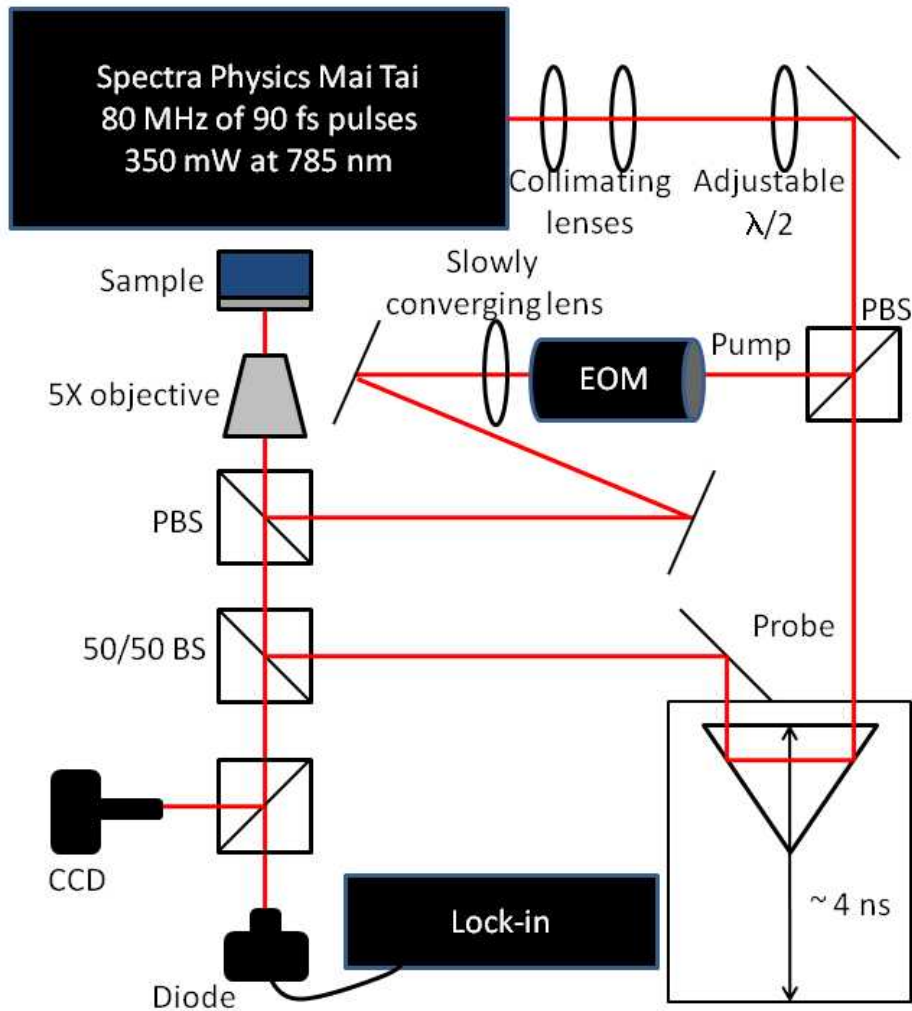
Outline

- **Noncontact experimental techniques for measuring Λ and G in nanostructures – PSTR and TTR**
- **Thermal models during laser heating**
- **Sensitivity of PSTR and TTR signals to Λ and G**

PSTR technique



TTR technique



Thermal diffusion due to modulated heat source

Heat equation in cylindrical coordinates due to laser heating

$$C \frac{\partial \theta(z, t)}{\partial t} = \Lambda_z \frac{\partial^2 \theta(z, t)}{\partial z^2} + \frac{\Lambda_r}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \theta(z, t)}{\partial r} \right)$$

Transform

$$\frac{\partial^2 \theta(\omega)}{\partial z^2} = q^2 \theta(\omega) \quad q^2 = \frac{\Lambda_r k^2 + iC\omega}{\Lambda_z} \quad q^2 = k^2 + \frac{iC\omega}{\Lambda_z}$$

Feldman algorithm (Feldman, High Temperatures-High Pressures, **31**, 293 (1999))

$$F(k) = \frac{1}{\gamma_1} \left(\frac{F_{T1}^+ + F_{T1}^-}{F_{T1}^- - F_{T1}^+} \right) \quad \begin{bmatrix} F_{T1}^+ \\ F_{T1}^- \end{bmatrix} = \begin{bmatrix} \exp[-q_1 d_1] & 0 \\ 0 & \exp[q_1 d_1] \end{bmatrix} \begin{bmatrix} F_{B1}^+ \\ F_{B1}^- \end{bmatrix}$$

$$\theta(\omega) = \frac{A}{2\pi} \int_0^\infty F(k) \exp \left[\frac{-k^2 (w_m^2 + w_r^2)}{8} \right] k dk$$

Need to specify a back side boundary condition

TTR technique – thermal assumptions

Thermal penetration depth

$$D = \sqrt{\Lambda / (C \pi f)}$$

Thermal penetration time

$$\tau = L^2 C / \Lambda$$

Condition for TTR

$$\tau = 1 / (\pi f)$$

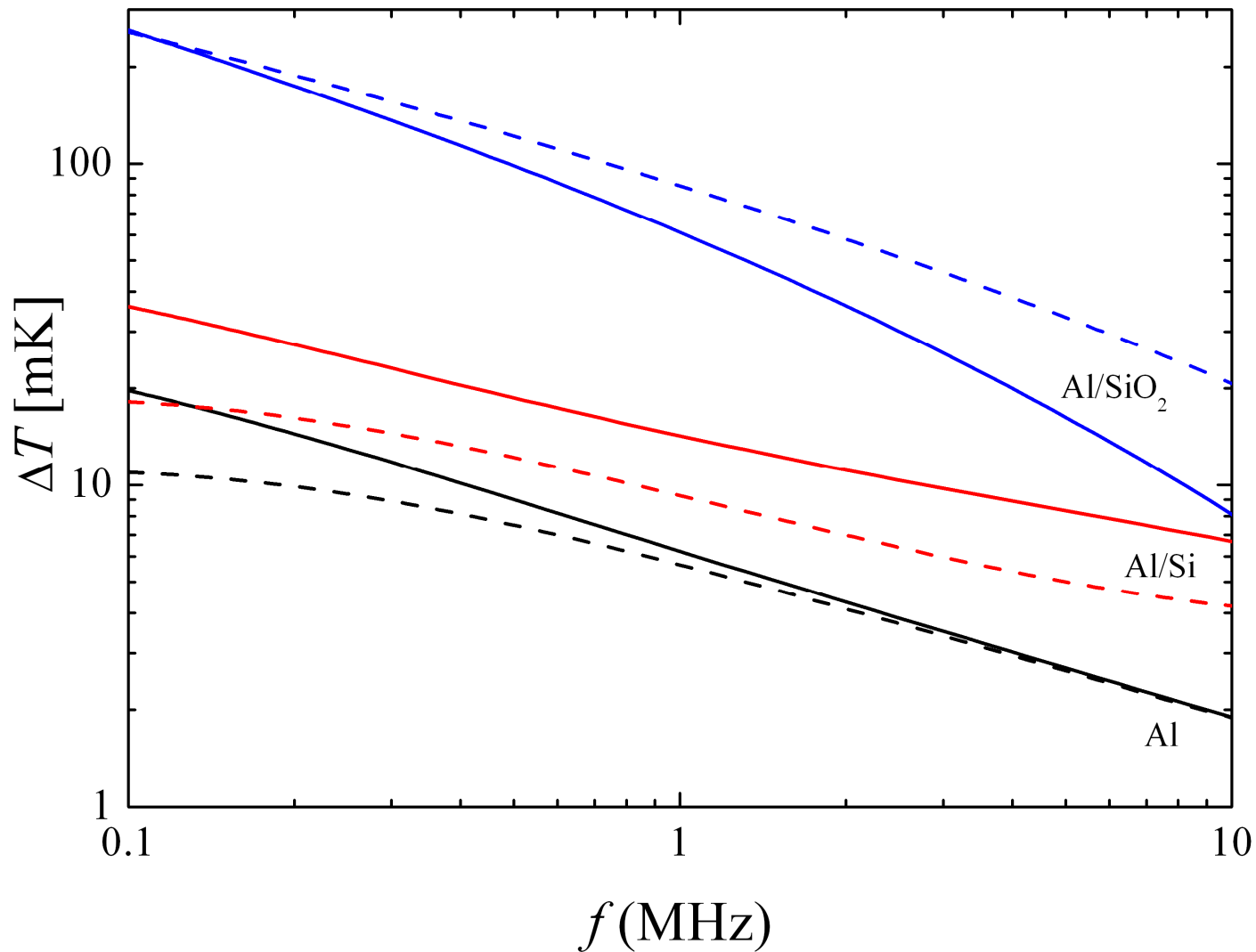
Over 8 ns pump-probe delay time (τ), at typical TTR modulation frequencies for film on substrate systems, assume back side of substrate is semi-infinite

$$F(k) = \frac{1}{\gamma_1} \left(\frac{F_{T1}^+ + F_{T1}^-}{F_{T1}^- - F_{T1}^+} \right) \quad \begin{bmatrix} F_{T1}^+ \\ F_{T1}^- \end{bmatrix} = \begin{bmatrix} \exp[-q_1 d_1] & 0 \\ 0 & \exp[q_1 d_1] \end{bmatrix} \begin{bmatrix} F_{B1}^+ \\ F_{B1}^- \end{bmatrix}$$

$$\begin{bmatrix} F_{B1}^+ \\ F_{B1}^- \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 + \frac{\gamma_2}{\gamma_1} - \frac{\gamma_2}{G_{12}} & 1 - \frac{\gamma_2}{\gamma_1} + \frac{\gamma_2}{G_{12}} \\ 1 - \frac{\gamma_2}{\gamma_1} - \frac{\gamma_2}{G_{12}} & 1 + \frac{\gamma_2}{\gamma_1} + \frac{\gamma_2}{G_{12}} \end{bmatrix} \begin{bmatrix} F_{T2}^+ \\ F_{T2}^- \end{bmatrix} \quad \begin{bmatrix} F_{T2}^+ \\ F_{T2}^- \end{bmatrix} = \begin{bmatrix} 0 \\ \exp[-u_2 d_2] \end{bmatrix}$$

Semi-infinite substrate

TTR technique – temperature change

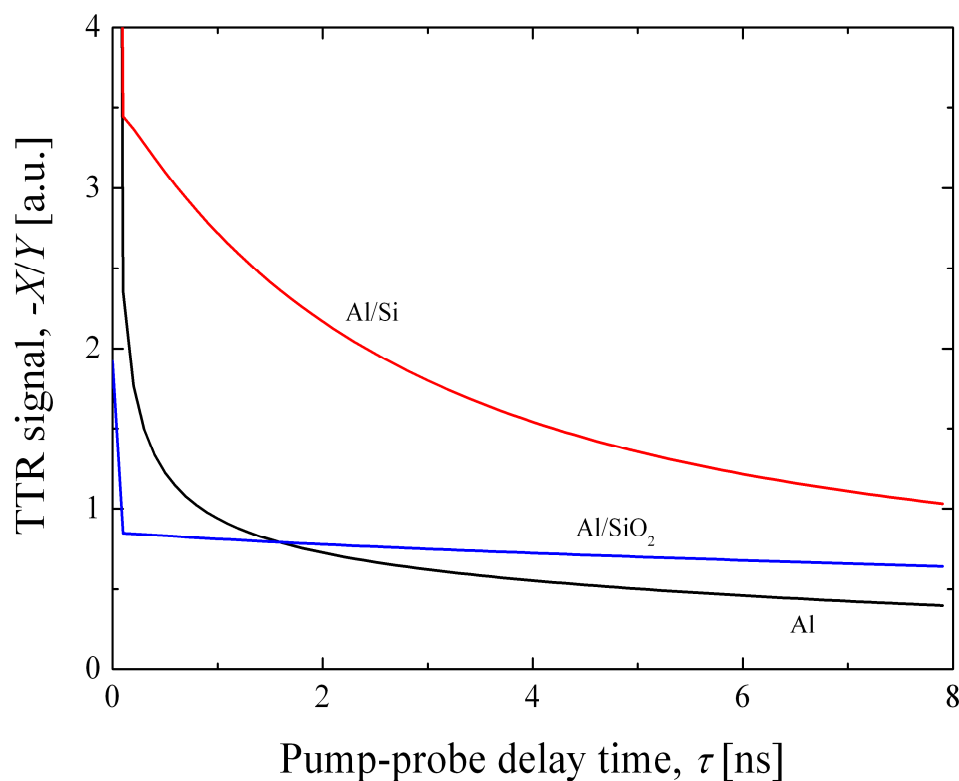


TTR technique – lock in signal

$$R \exp[i(\omega_m t + \phi)] = Z(\omega_m) \exp[i\omega_m t] \quad Z(\omega_0) = \frac{(2\pi)^2 \chi}{\omega_s^2} \sum_{M=-\infty}^{\infty} \theta(\omega_0 + M\omega_s) \exp[iM\omega_s \tau]$$

$$X = \text{Re}[Z(\omega_0)], Y = \text{Im}[Z(\omega_0)]$$

Cahill, RSI, **75**, 5119 (2004) & Schmidt *et al.*, RSI, **79**, 114902 (2008)



PSTR technique – thermal assumptions

Thermal penetration depth

$$D = \sqrt{\Lambda / (C\pi f)}$$

At modulation frequencies typical in PSTR, thermal penetration depth is same order as material system thickness

$$F(k) = \frac{1}{\gamma_1} \left(\frac{F_{T1}^+ + F_{T1}^-}{F_{T1}^- - F_{T1}^+} \right) \quad \begin{bmatrix} F_{T1}^+ \\ F_{T1}^- \end{bmatrix} = \begin{bmatrix} \exp[-q_1 d_1] & 0 \\ 0 & \exp[q_1 d_1] \end{bmatrix} \begin{bmatrix} F_{B1}^+ \\ F_{B1}^- \end{bmatrix}$$

$$\begin{bmatrix} F_{B1}^+ \\ F_{B1}^- \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 + \frac{\gamma_2}{\gamma_1} - \frac{\gamma_2}{G_{12}} & 1 - \frac{\gamma_2}{\gamma_1} + \frac{\gamma_2}{G_{12}} \\ 1 - \frac{\gamma_2}{\gamma_1} - \frac{\gamma_2}{G_{12}} & 1 + \frac{\gamma_2}{\gamma_1} + \frac{\gamma_2}{G_{12}} \end{bmatrix} \begin{bmatrix} F_{T2}^+ \\ F_{T2}^- \end{bmatrix} \quad \begin{bmatrix} F_{T2}^+ \\ F_{T2}^- \end{bmatrix} = \begin{bmatrix} \exp[-q_1 d_1] & 0 \\ 0 & \exp[q_1 d_1] \end{bmatrix} \begin{bmatrix} F_{B2}^+ \\ F_{B2}^- \end{bmatrix}$$

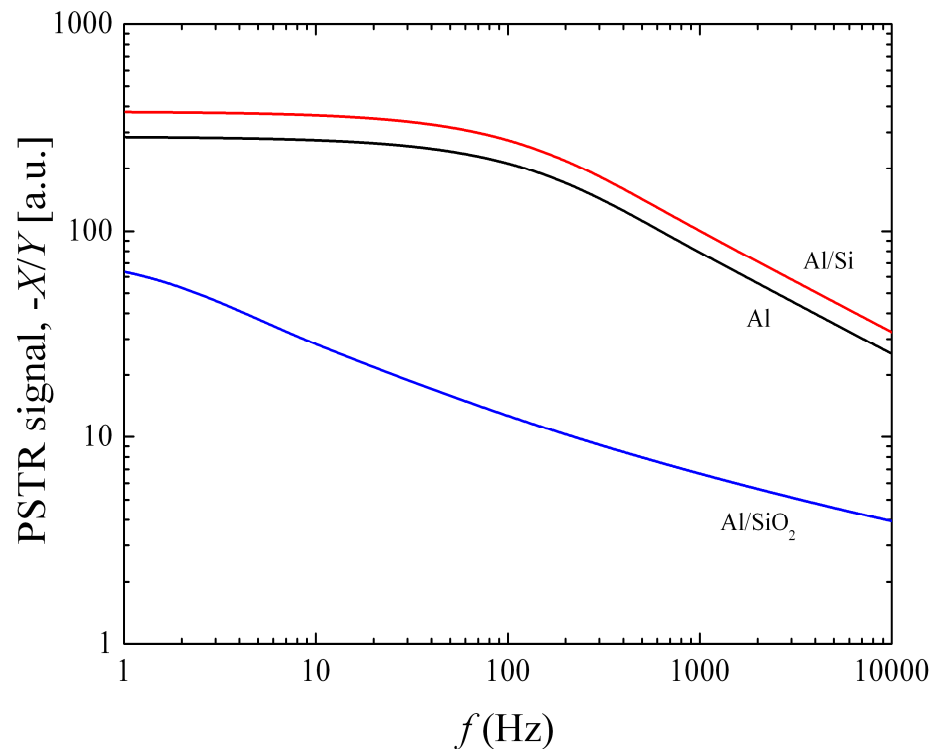
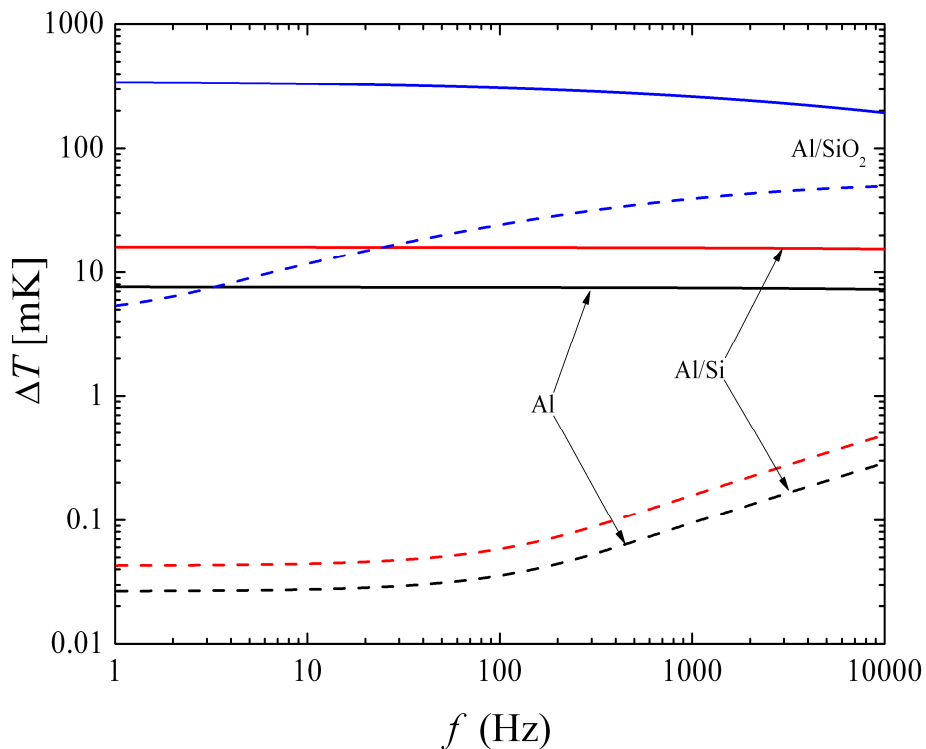
Assume insulative boundary

$$\begin{bmatrix} F_{B2}^+ \\ F_{B2}^- \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} F_{T3}^+ \\ F_{T3}^- \end{bmatrix} \quad \begin{bmatrix} F_{T3}^+ \\ F_{T3}^- \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

PSTR technique – lock in signal

$$Z(\omega_0) = \chi\theta(\omega_0)$$

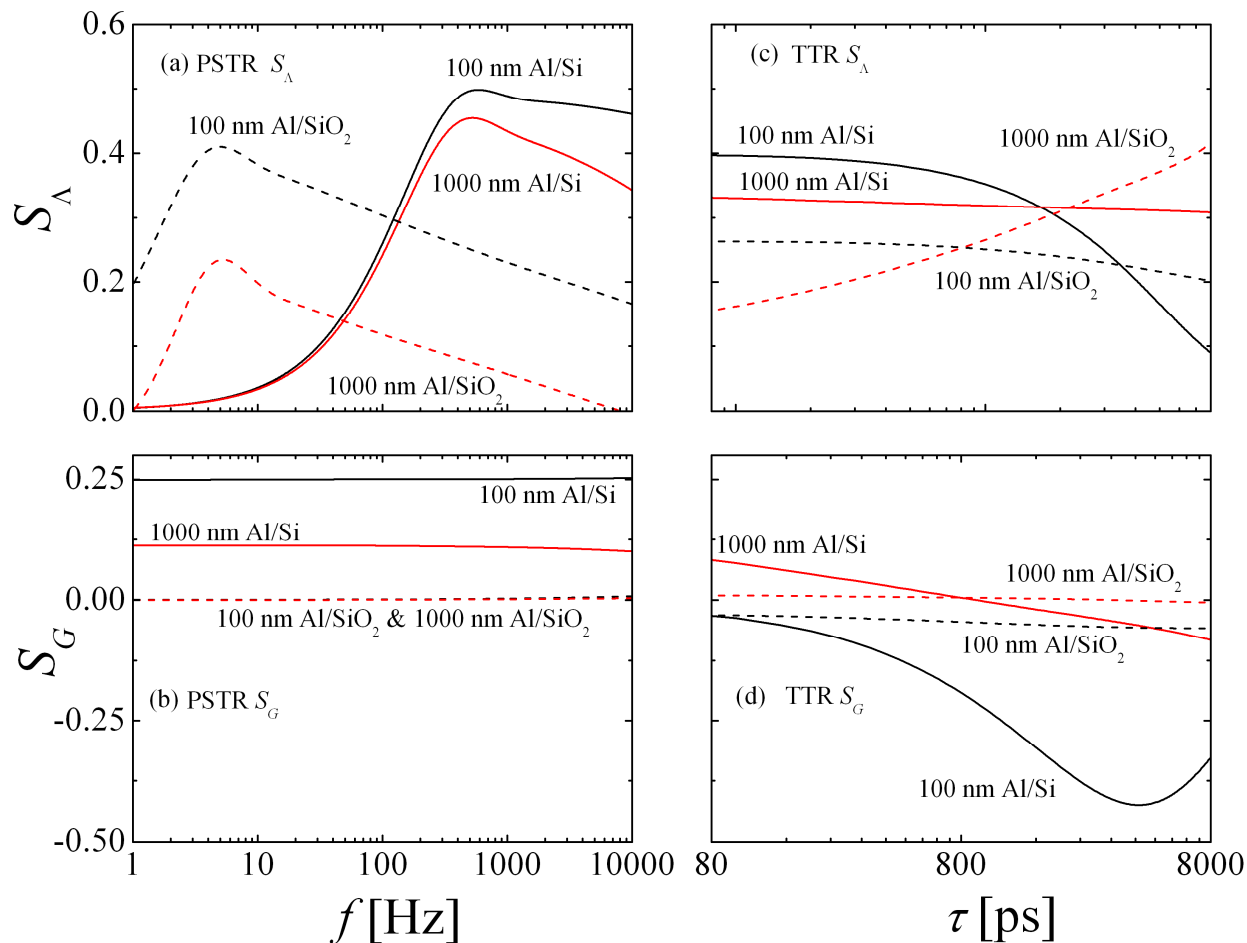
$$X = \text{Re}[Z(\omega_0)], Y = \text{Im}[Z(\omega_0)]$$



Analysis sensitivities

$$S_{\xi} = \frac{\partial \ln \left[\frac{-X}{Y} \right]}{\partial \ln[\xi]} = \frac{\xi}{\left(\frac{-X}{Y} \right)} \frac{\partial \left[\frac{-X}{Y} \right]}{\partial [\xi]}$$

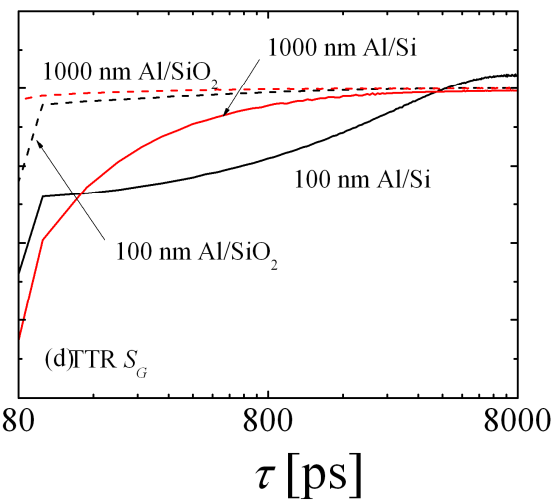
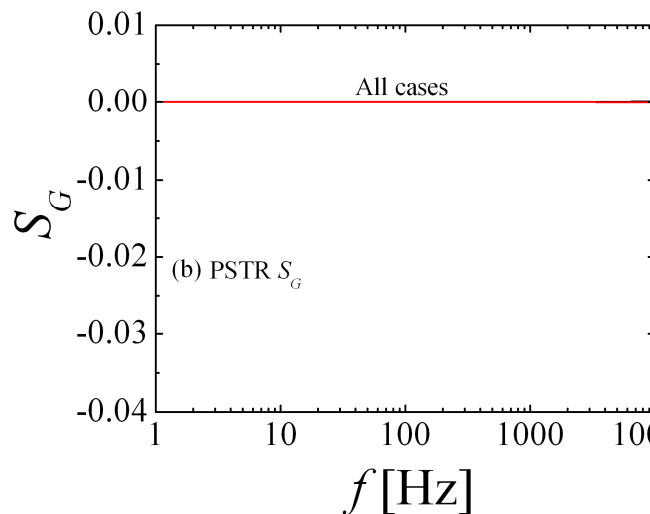
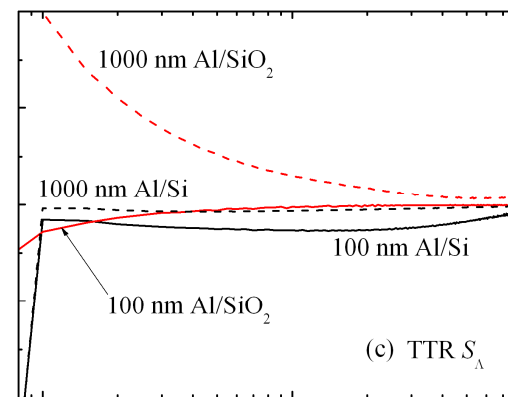
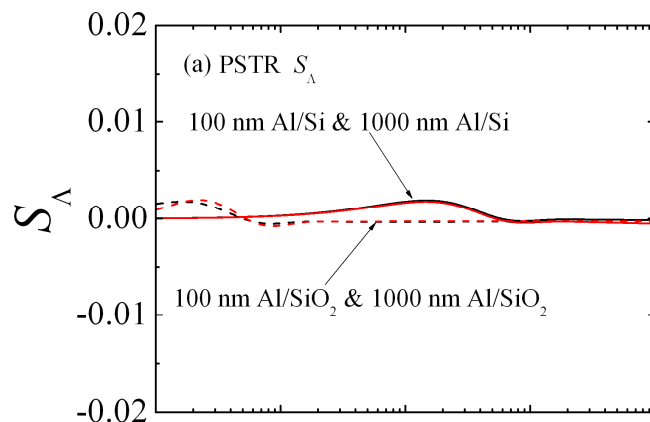
Costescu *et al.* PRB, **67**, 054302 (2002)



Analysis sensitivities – scaling

$$S_{\xi}^{PSTR} = \Delta f \frac{\partial}{\partial f} \left[\frac{\xi}{\left(\frac{-X}{Y} \right)} \frac{\partial \left[\frac{-X}{Y} \right]}{\partial [\xi]} \right]$$

$$S_{\xi}^{TTR} = \Delta \tau \frac{\partial}{\partial \tau} \left[\frac{\xi}{\left(\frac{-X}{Y} \right)} \frac{\partial \left[\frac{-X}{Y} \right]}{\partial [\xi]} \right]$$





Conclusions

- **Feldman algorithm for semi-infinite and insulative substrate system leads to lock-in response - $-X/Y$**
- **TTR and PSTR techniques both sensitive to G and Λ if all experimental parameters are known (spot sizes, absorbed power, electronic gain, etc.)**
- **If parameters are unknown and model must be scaled to data, TTR is far superior except for measuring conductivity of bulk, conductive systems**

Acknowledgements

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