

# Deuterium gas puff z-pinch simulation with LSP

**D. R. Welch and D. V. Rose**  
**Voss Scientific, Albuquerque, NM 87108 USA**  
**W. A. Stygar and R. J. Leeper**  
**Sandia National Laboratories, Albuquerque, NM 87185**  
**USA**

**June 1, 2009**  
**International Conference on Plasma Science 2009**  
**San Diego, CA USA**

\* Sandia is a multiprogram laboratory operated by Sandia Corporation, a Lockheed Martin Company, for the United States Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000.

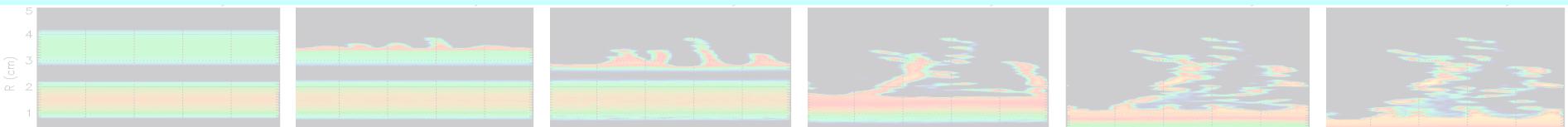
# Goal: Address 50 year old question of neutron production mechanism in z-pinch

- Thermonuclear neutron production scaling as  $\text{I}^4$  is quite promising for fusion energy\*
- Deuterium pinches have been examined experimentally for many years
  - Significant neutron production
  - Early conclusions had been that the majority of neutron yield is from non-thermal deuterium population that does not scale favorably
- We have built a fully kinetic electromagnetic model that includes both nonthermal and thermal processes to address this question

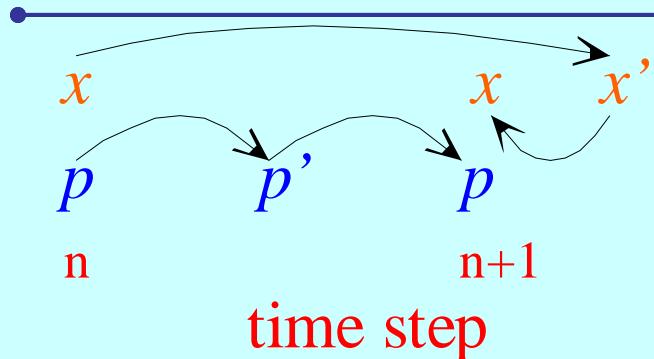
\*J. Ise, Jr. and R. V. Pyle, in Conference on Controlled Thermonuclear Reactions, Princeton Univ., 17-20 October 1955 (TID-7503, USAEC, 1955), p.218; Velikovich, et al. Phys. of Plasmas 14 022701 (2007).

# Elements of LSP computational model

- Fully electromagnetic solution to allow non neutral plasmas (Rayleigh Taylor bubble and spike potentials)
  - EM energy-conserving **direct-implicit** algorithm
- Electron and ion inertial effects
  - Fully kinetic particle-in-cell treatment
- Plasma interpenetration
  - Coulomb scattering between all particles
  - binary ion-ion Coulomb collisions
  - approx. e-ion, e-e based on drifting Maxwellian
- Fusion reaction calculation for arbitrary reactant energy distribution function
  - Binary fusion model



# Direct Implicit Algorithm\*



particle  $p$  advance using half of the old and new electric field

- First, particle  $p$  and  $x$  advance with  $E_{n+1}(x_{n+1}) = 0$
- The EM fields are then pushed using linear correction terms to predict the effect of  $E_{n+1}(x_{n+1})$  on the perturbed current  $\delta J = \langle S \rangle \bullet E_{n+1}(x_{n+1})$  ( $\langle S \rangle$  is the susceptibility)

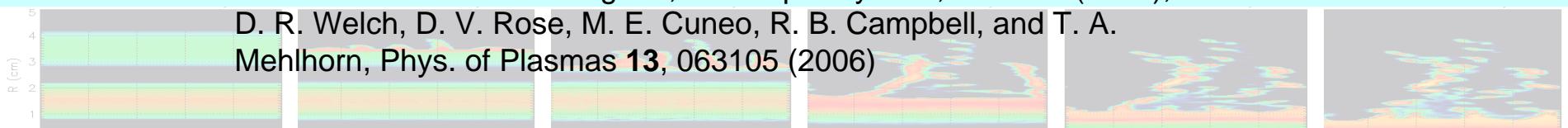
$$\frac{\partial E}{\partial t} = \frac{1}{\epsilon} \left( \nabla \times \frac{B}{\mu} \right) - J - \langle S \rangle \cdot E$$

- Particle  $p$  and  $x$  are then corrected in final push

\*See

D. W. Hewitt and A. B. Langdon, J. Comp. Phys. **72**, 121-155 (1987);

D. R. Welch, D. V. Rose, M. E. Cuneo, R. B. Campbell, and T. A. Mehlhorn, Phys. of Plasmas **13**, 063105 (2006)



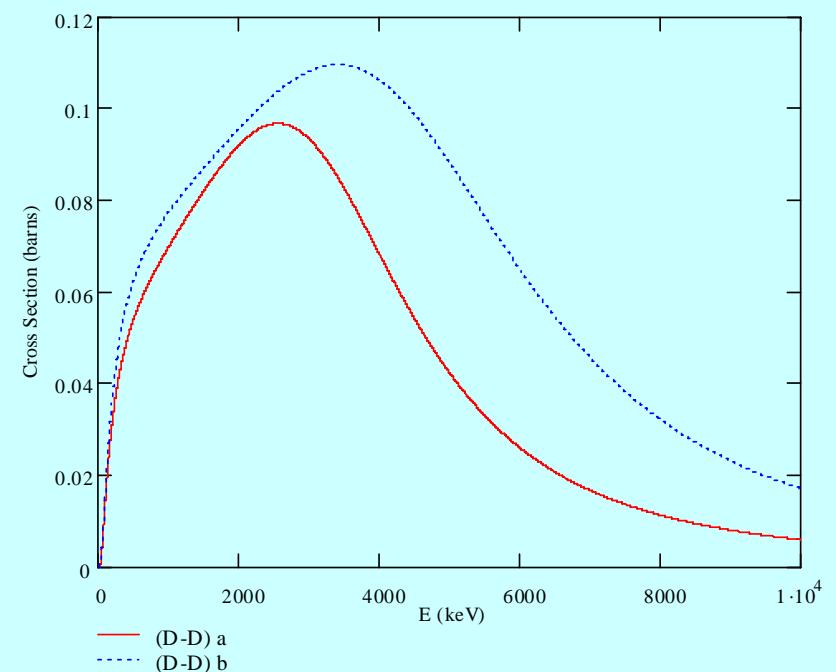
# Binary collision algorithm\* adapted for Coulomb, fusion reactions

- Particles are sorted to respective cells and each pair is scattered with only assumption that of Coulomb logarithm

Here we consider fusion reactions:



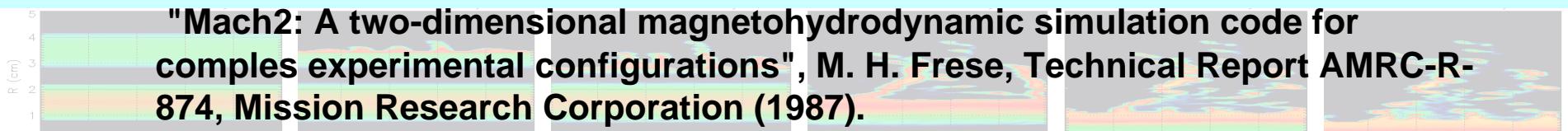
*Non-relativistic kinematics good for  $E < 10 \text{ MeV}$*



# Detailed PIC comparison with MHD

---

- MHD simulations with Mach2 code\*
- Same numerical spatial grid (100  $\mu\text{m}$  resolution)
- Internal Mach2 timestep longer than LSP
- D-D neutron production assuming Maxwellian ion temperature
- Electrode effects ignored (periodic BC in z)

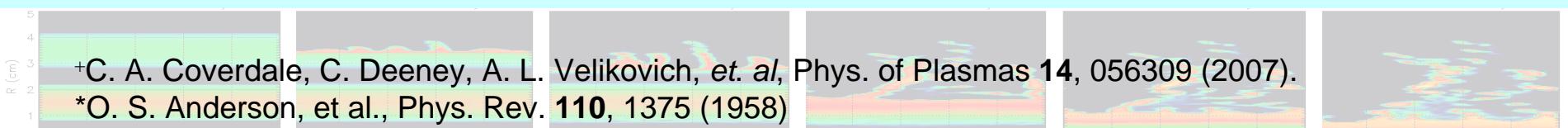


# 1D, 2D simulations with varying current test neutron yield scaling

- Pinch experiments on Z performed in the 15 MA range offer a good starting point<sup>+</sup>
- Consider same 2 ring configuration
  - Outer puff 222.5 mg, 3-4 cm radius
  - Inner puff 222.5 mg, 1-3 cm radius
  - Assuming  $I_o = 15$  MA with a linear  $\tau = 100$  ns rise,  $R = 4$  cm
  - Keep time to first bounce\* fixed at 100 ns to pinch
  - Lengths scale as  $I_o^{1/2}$  with fixed gas density

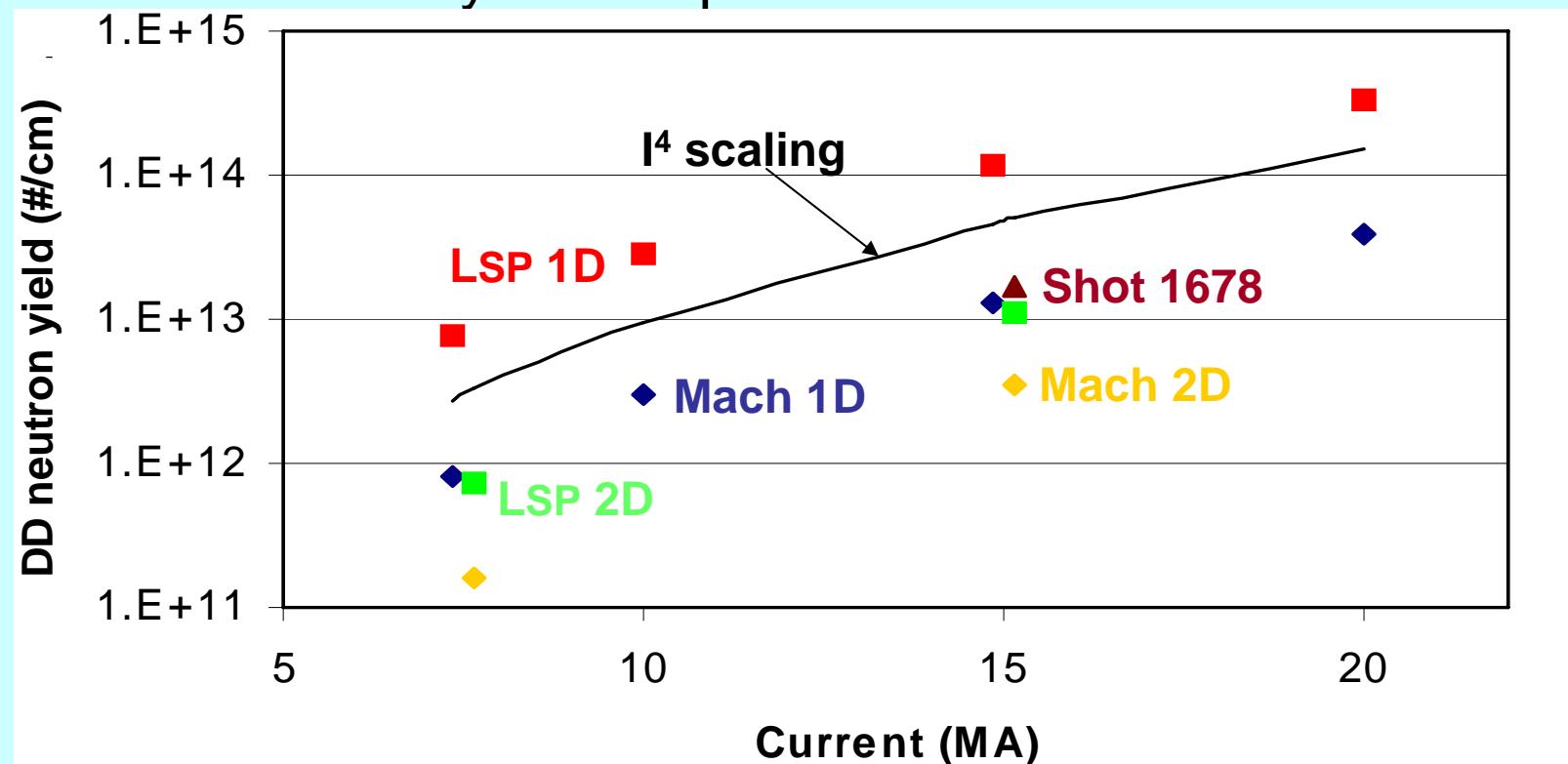
$$t = 1.43R \left[ \frac{100\pi\rho\tau^2}{I_o^2} \right]^{1/4}$$

Assumes uniform distribution of gas density  $\rho$  with  $R$



# Neutron yield scaling roughly $I^4$

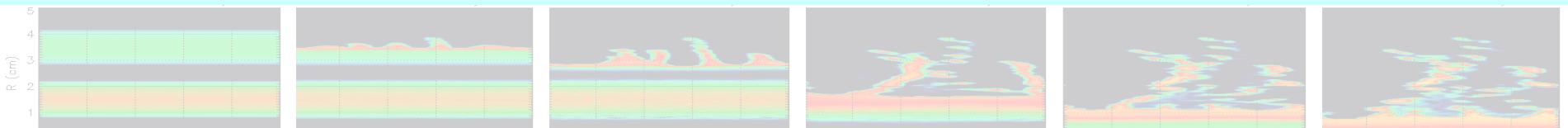
- Roughly thermonuclear scaling from 7-20 MA regardless of dimensionality or computation model



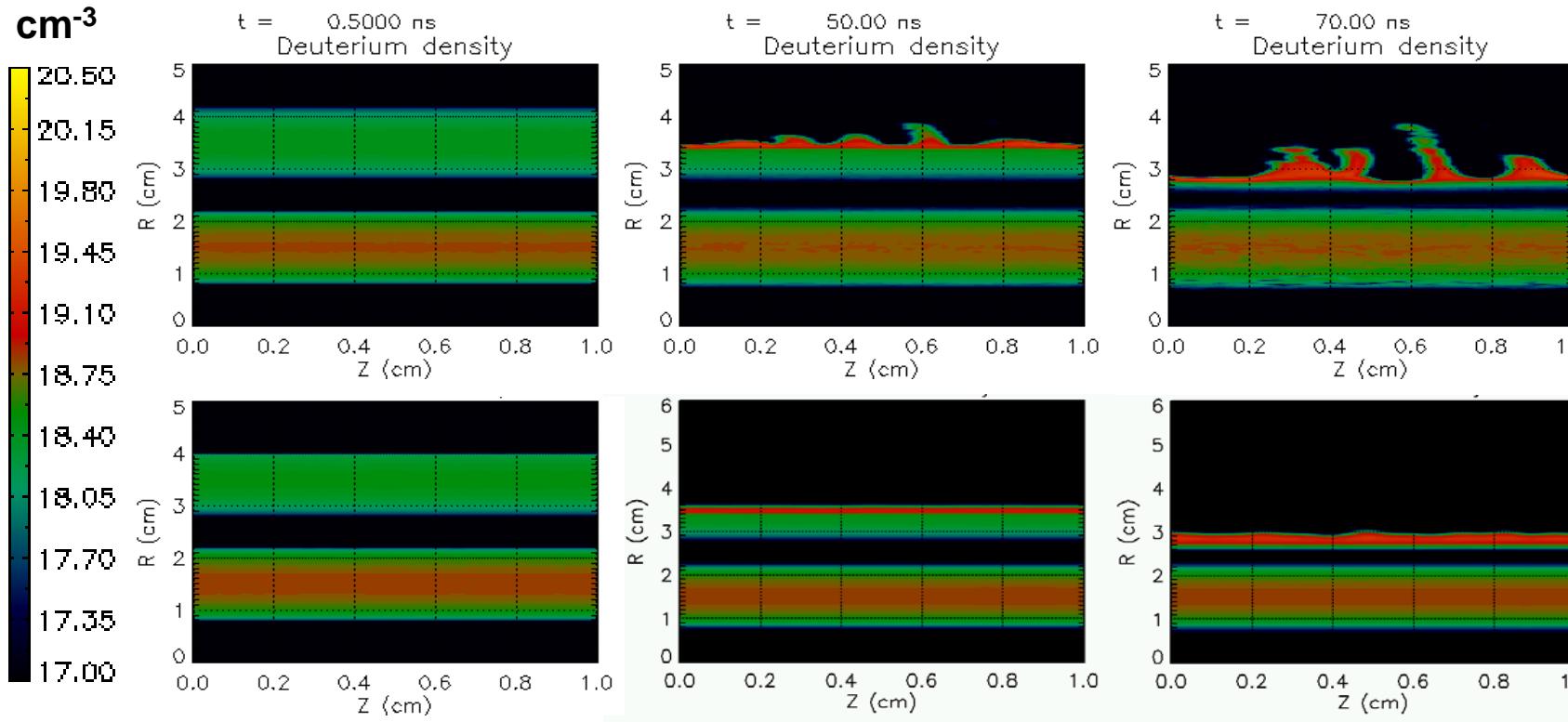
# Careful 2D comparison at 15 MA between PIC and MHD

---

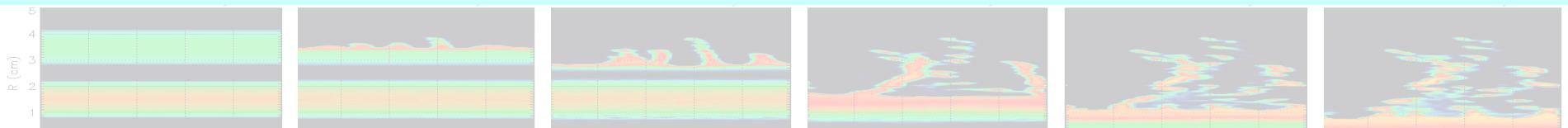
- 1% random axial mass variation drives more rapid Rayleigh Taylor growth in PIC
- PIC shows annular character of plasma maintained thru coalesce on axis, no shocks as seen in MHD
- PIC has larger current near axis, drives stronger electric fields, fast ions



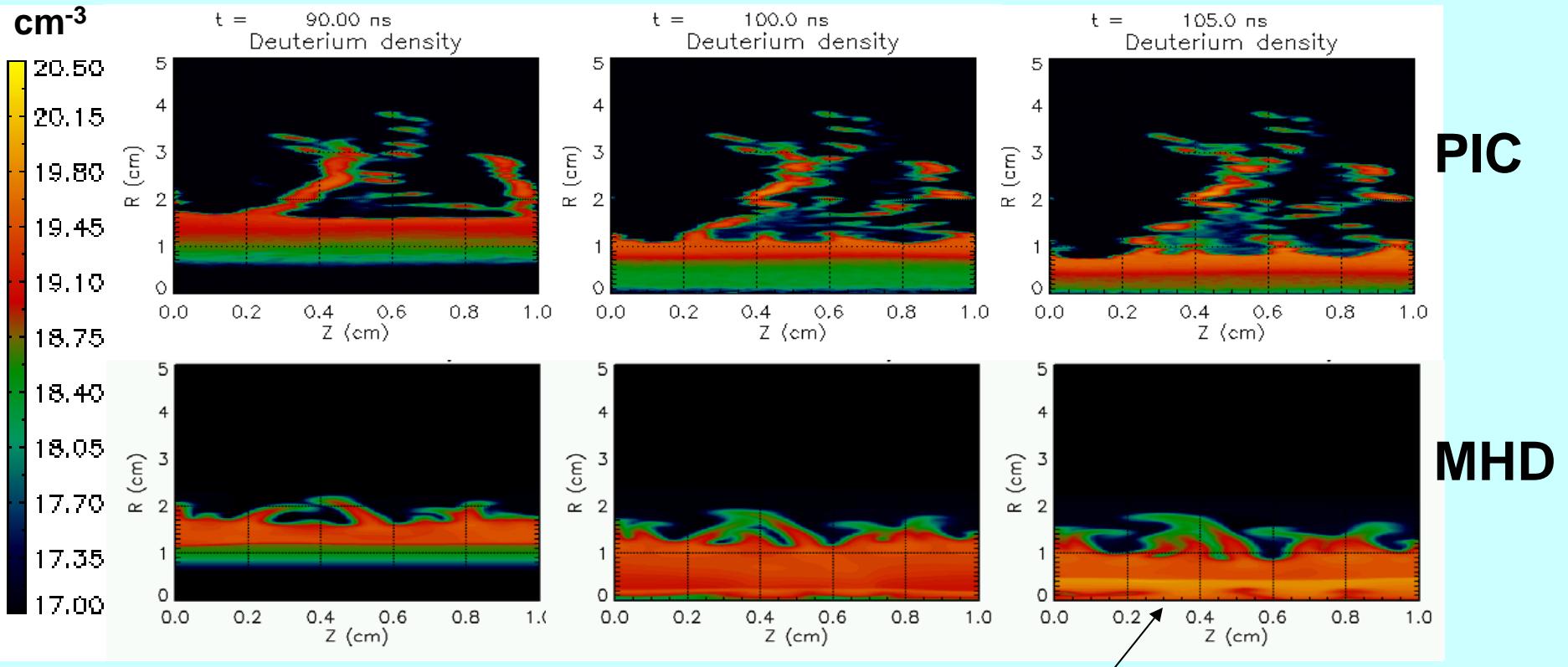
# Early phase – PIC exhibits rapid Rayleigh Taylor growth



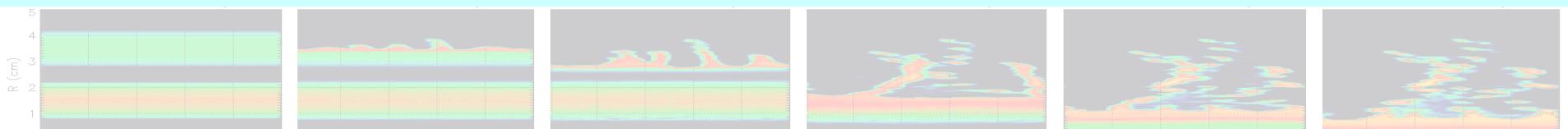
PIC growth could be enhanced by particle noise



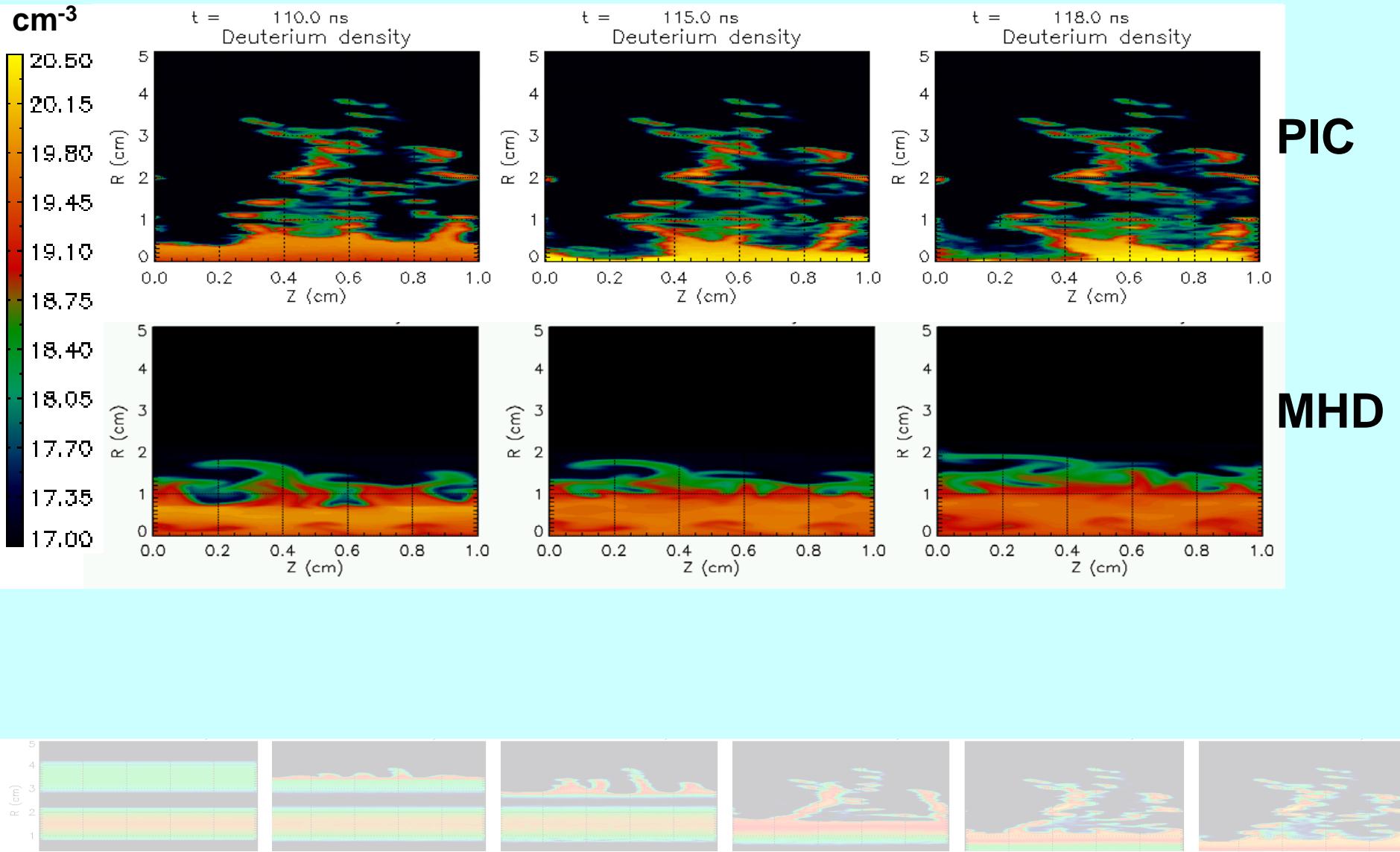
# Pinch phase – coalescence in PIC is delayed



MHD exhibits a shock behavior not observed in PIC

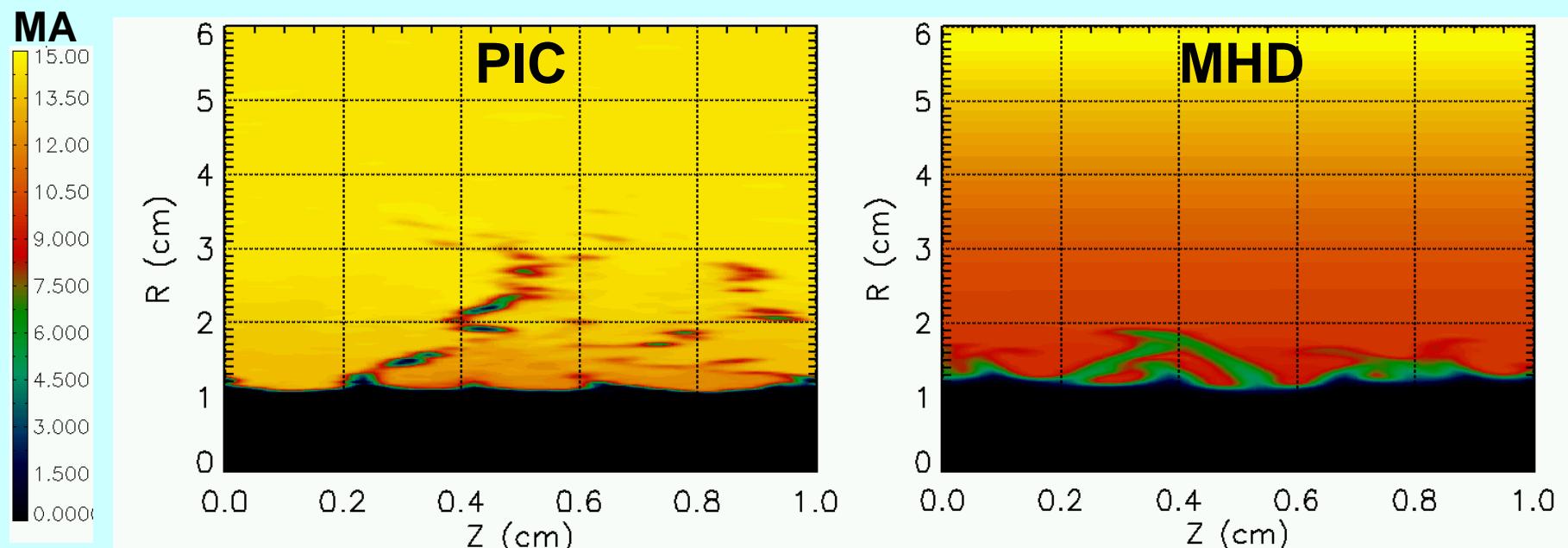


# Late time – some mass still coalescing, filament forms in PIC

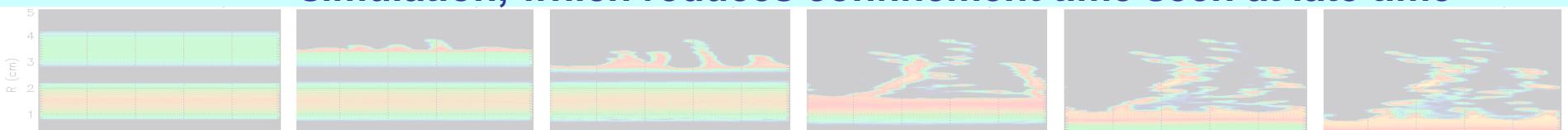


# Hotter corona region carries more current in MHD

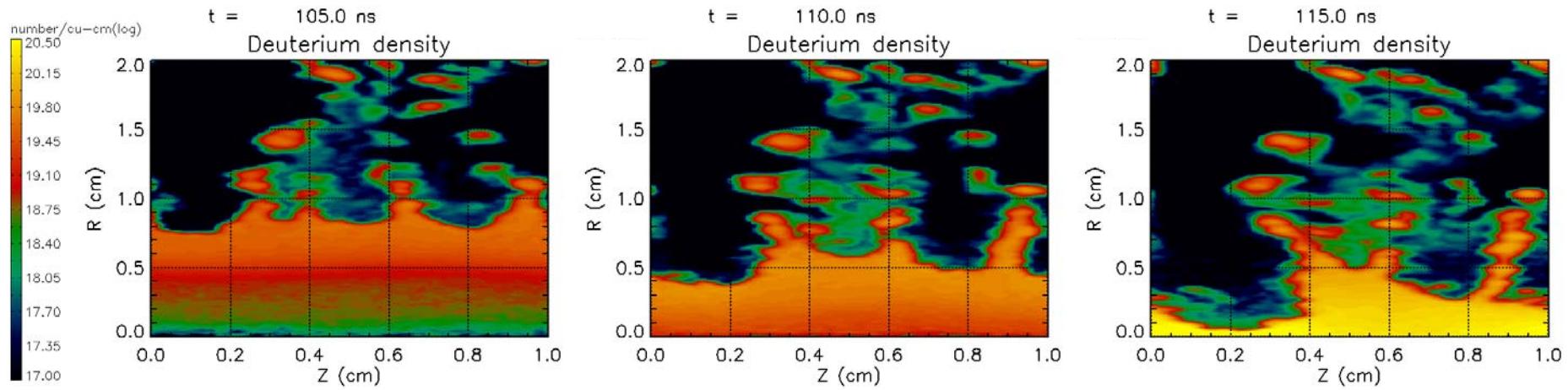
- Comparison is at 100 ns into implosion



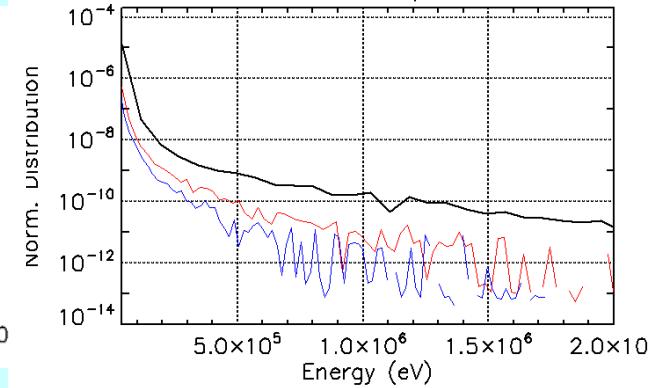
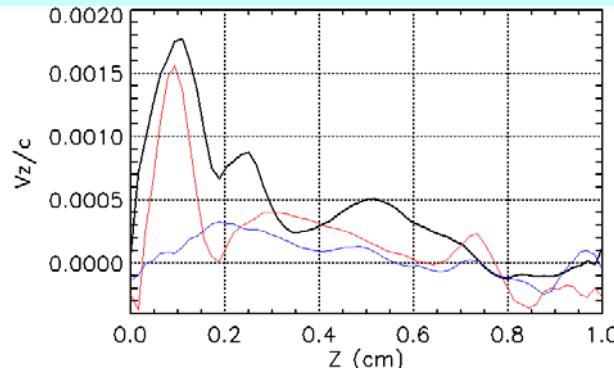
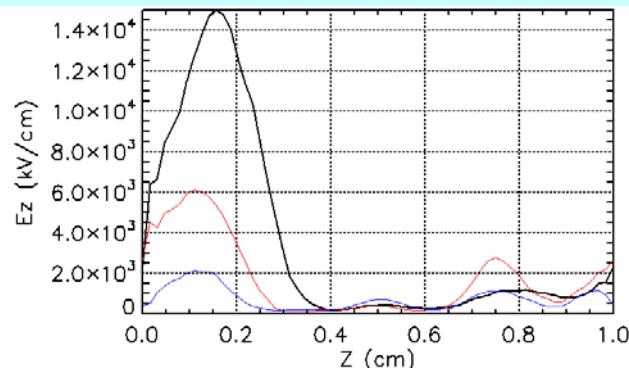
The effective pinch current at 2 cm is 2 MA less in MHD simulation, which reduces confinement time seen at late time



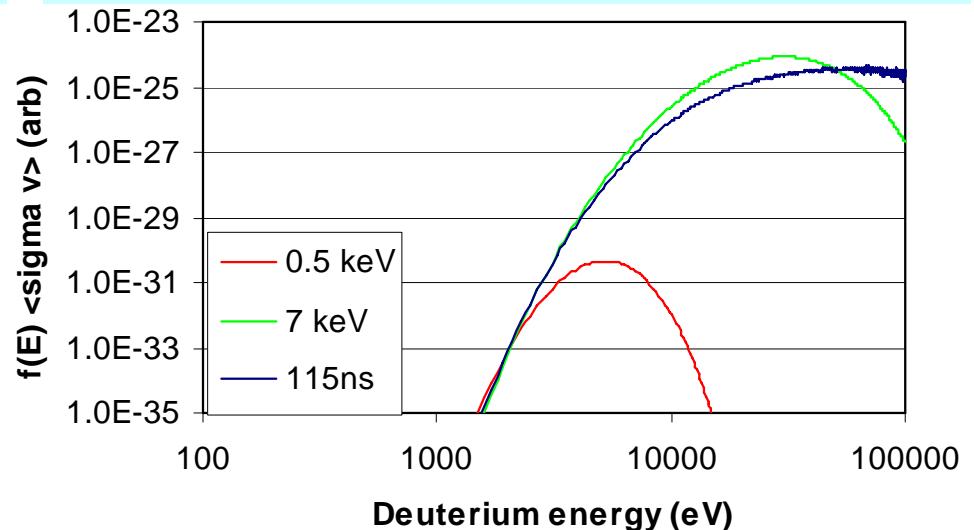
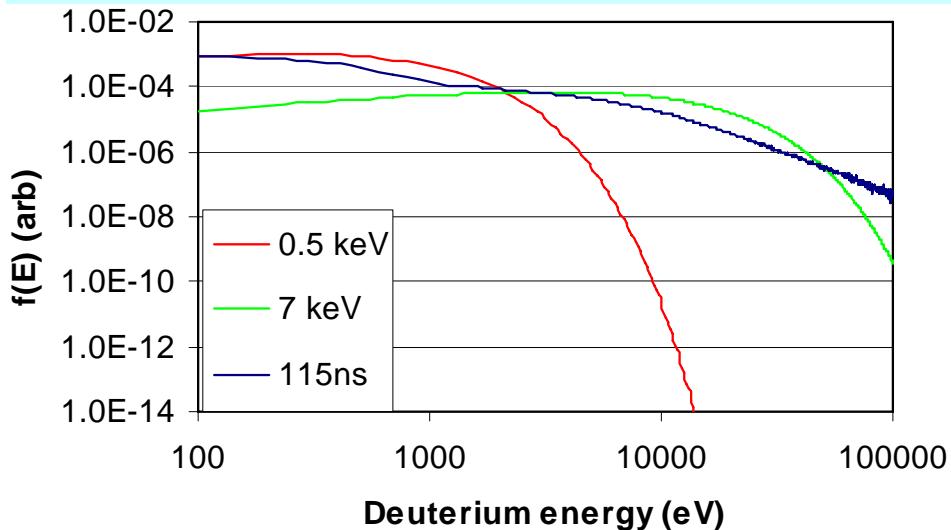
# Electric fields large between spikes, drive ions to higher energy



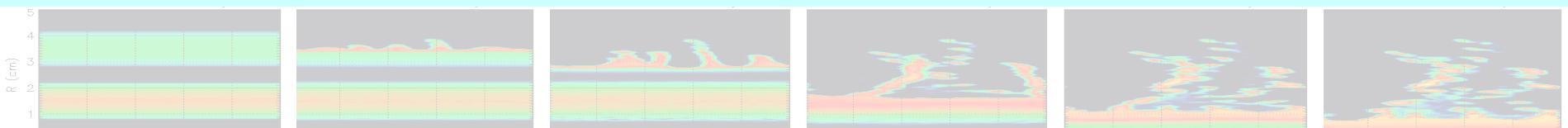
105 ns 110 ns 115 ns



# Deuterium energy distribution exhibits high energy tail

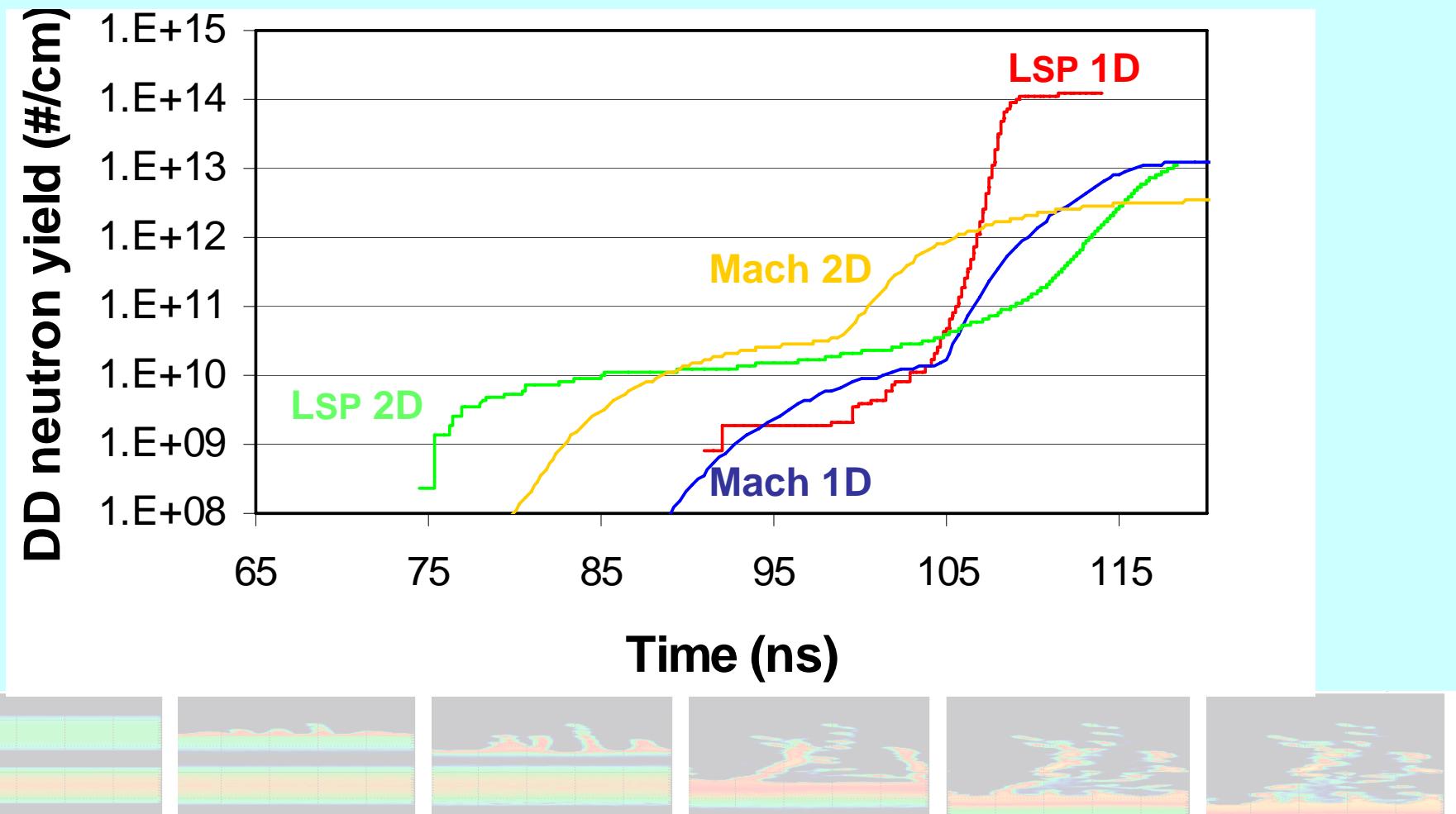


Distribution tail accounts for half the neutron production by 115 ns



# Neutron production at 15 MA

- As in experiment, neutron production in PIC is somewhat longer than MHD with significant nonthermal component at late time



# Kinetic, electromagnetic simulation reveals nonideal pinch behaviors

---

- Non Maxwellian distributions, mean free path effects, non neutral assumption results in strong differences in pinch characteristics
- Strong E fields drive energetic ions which enhance neutron production late in time – may account for half the yield
- Electrode, 3D effects subject of future work

