



Pulsed-power, Z-pinches, and Applications

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Special Thanks to
Mark Herrmann (SNL)

Thanks to John Lindl, and Ed Moses, M. Keith Matzen and Dan Sinars, Steve Slutz, Bill Stygar, Sasha Velikovich, Marcus Knudson, Ray Lemke for viewgraphs



Outline

- Pulsed-power technology produces large currents (15-26 MA) in a short pulse (100-600 ns) on the Z machine
- Large currents generate large magnetic fields = tremendous pressure
- Large pressures enable access to High Energy Density regimes
- There are several interesting applications

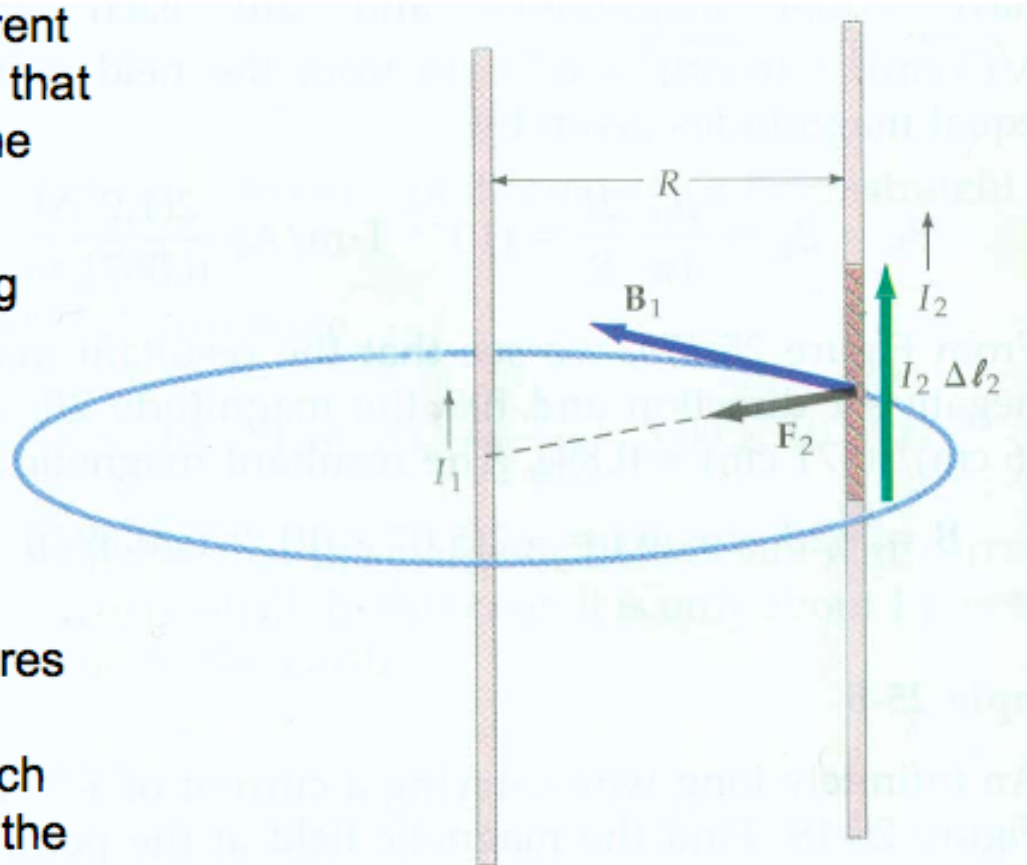


1st year physics refresher

A single wire carrying current produces a magnetic field that encircles it according to the right-hand rule

Two parallel wires carrying current along the same direction will attract each other (Biot-Savart Law, “ $\mathbf{J} \times \mathbf{B}$ force”)

Definition of an Ampere:
If two very long parallel wires 1 m apart carry equal currents, the current in each is defined to be 1 A when the force/length is 2×10^{-7} N/m





We can incorporate the effect of magnetic fields into our fluid equations

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0$$

mass conservation

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla P$$

momentum conservation
(F=ma) cgs

For slowly varying fields we can approximate:

$$\nabla \times \mathbf{B} = \frac{4\pi \mathbf{J}}{c}$$

$$\mathbf{J} \times \mathbf{B} = \frac{c}{4\pi} (\nabla \times \mathbf{B}) \times \mathbf{B} = -\frac{c}{4\pi} \mathbf{B} \times (\nabla \times \mathbf{B})$$



The magnetic field can be included as an effective pressure

From vector identities

$$\mathbf{B} \times (\nabla \times \mathbf{B}) = \frac{1}{2} \nabla (\mathbf{B} \cdot \mathbf{B}) - \mathbf{B} \cdot \nabla \mathbf{B} = \nabla \left(\frac{B^2}{2} \right) - \mathbf{B} \cdot \nabla \mathbf{B}$$

$$\mathbf{J} \times \mathbf{B} = \frac{c}{4\pi} \left(\mathbf{B} \cdot \nabla \mathbf{B} - \nabla \left(\frac{B^2}{2} \right) \right)$$

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = \frac{\mathbf{J} \times \mathbf{B}}{c} - \nabla P = \frac{1}{4\pi} \mathbf{B} \cdot \nabla \mathbf{B} - \nabla \left(P + \frac{B^2}{8\pi} \right)$$

magnetic tension

magnetic pressure

In the case of an axisymmetric Z current (theta field) the magnetic tension is zero



How strong is this pressure?

$$P_m(\text{dyne} / \text{cm}^2) = \frac{B(\text{G})^2}{8\pi}$$

A typical refrigerator magnetic is 100 gauss ~ 400 dyne/cm²

A 5000 G (0.5 T) magnetic field ~ 10⁶ dyne/cm² = 1 atmosphere = 1 Bar

A 5x10⁶ G (500 T) magnetic field ~ 1 Million atmospheres = 1 Megabar (MB)=
High energy density physics

A 5x10⁹ G magnetic field ~ 1 Trillion atmospheres = 1 Terabar (TB) > pressure in
the center of the sun

Note that high explosives have pressure ~ few 100,000 atmospheres



Large currents can create large fields!

$$\nabla \times \mathbf{B} = \frac{4\pi\mathbf{J}}{c}$$

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \frac{4\pi}{c} \iint_S \mathbf{J} \cdot d\mathbf{S} \quad \text{Ampere's law}$$

For an axial current I :

$$2\pi r B_\theta = \frac{4\pi}{c} I$$

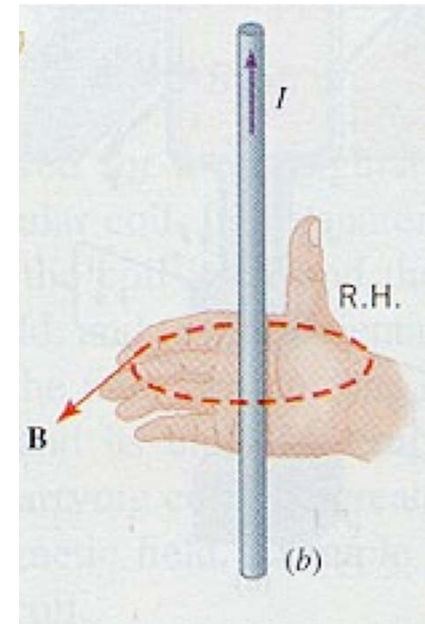
$$B_\theta = \frac{2}{c} \frac{I}{r} \quad (cgs)$$

$$B_\theta (G) = \frac{I(A)}{5 r(cm)}$$

100 A at 2 mm radius is 100 G

10^7 A at 4 mm radius is 5×10^6 G = 1 MBar of pressure

2.5×10^7 A at 1 mm radius is 5×10^7 G = 100 MBar of pressure!!





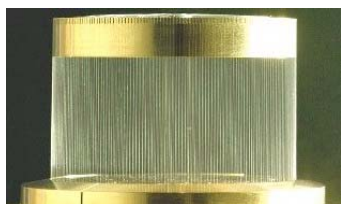
We can use high currents to push plasmas in different ways

High Current

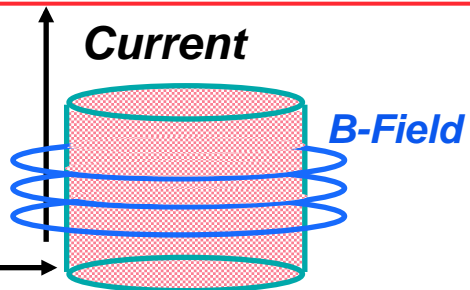
Cylindrical compression:
Z-pinch x-ray source

Planar compression:
Materials experiments

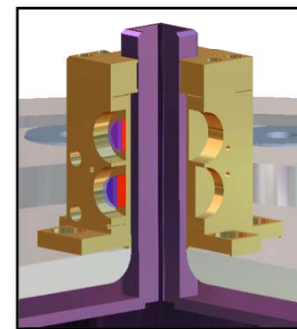
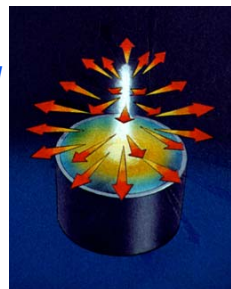
wire array



Current

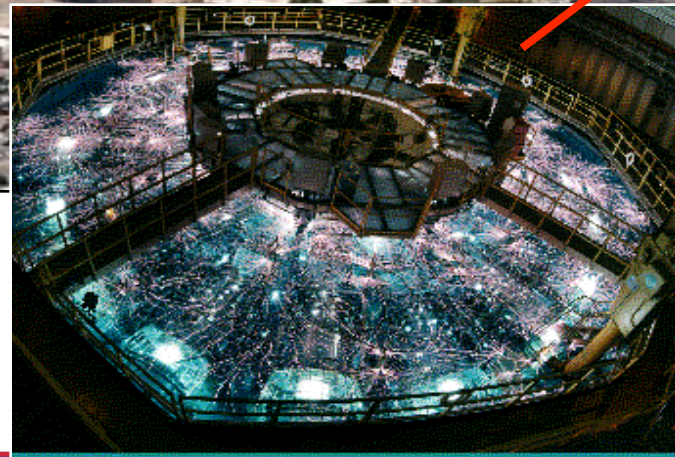


$J \times B$ Force





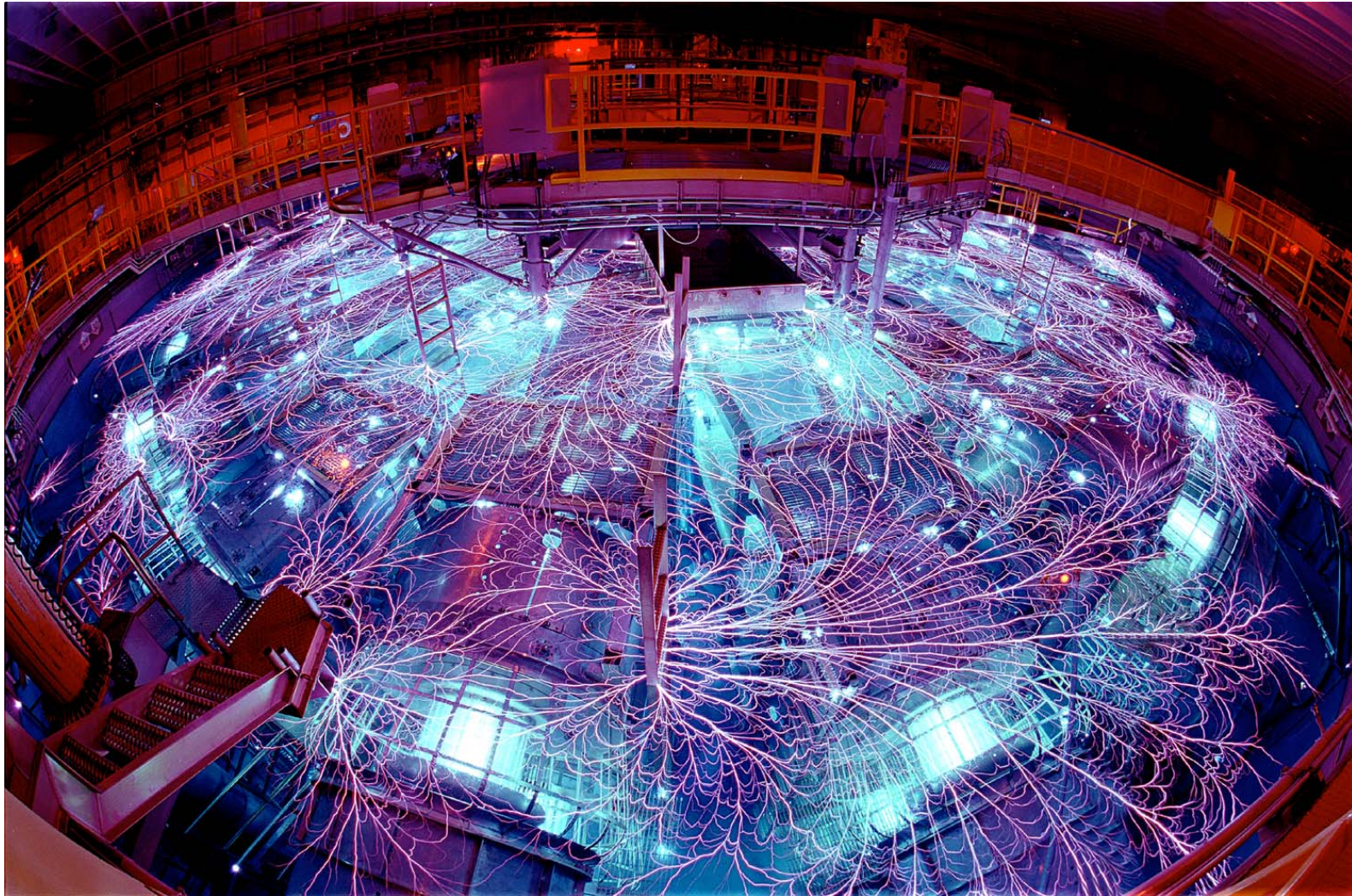
The “Z” pulsed power facility is located at Sandia National Laboratories in Albuquerque, New Mexico



Sandia National Laboratories

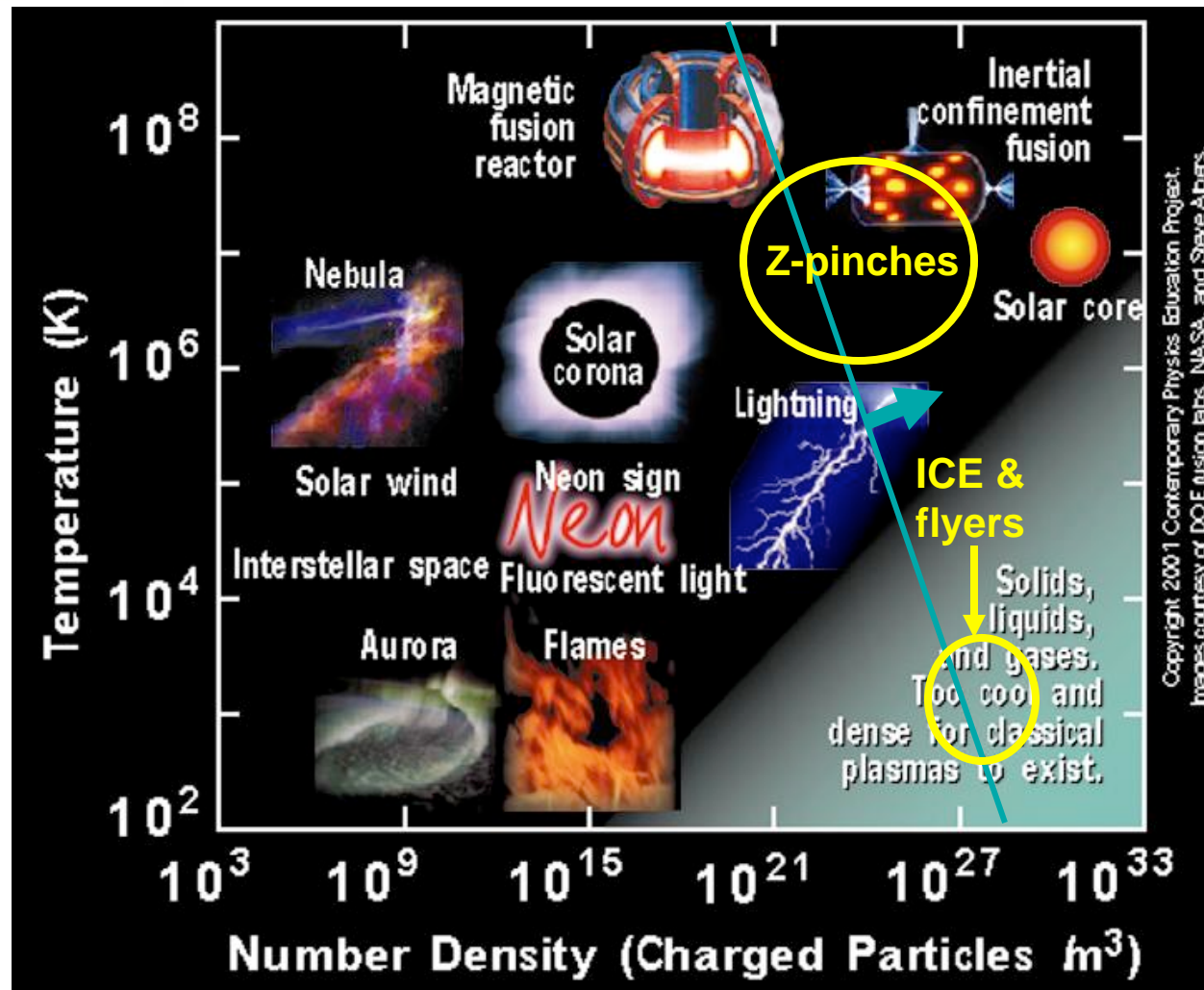


“Z” is the world’s largest pulsed power facility





Regimes of high energy density are typically associated with energy density $10^5 \text{ J/cm}^3 = 1 \text{ Mbar}$



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Images courtesy of DOE fusion labs, NASA, and Steve Ables.

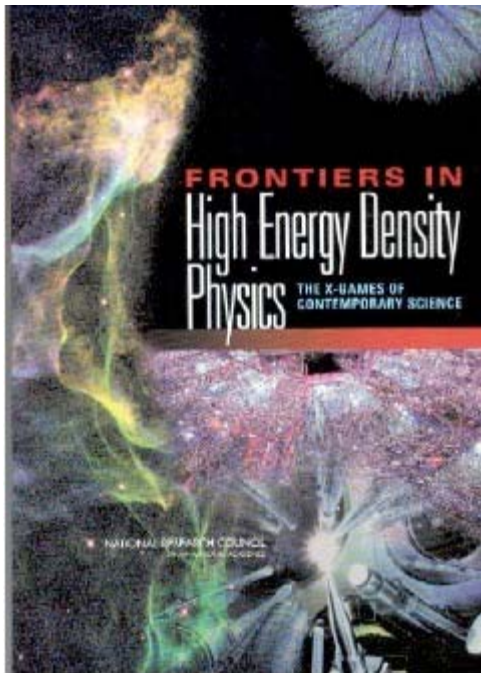
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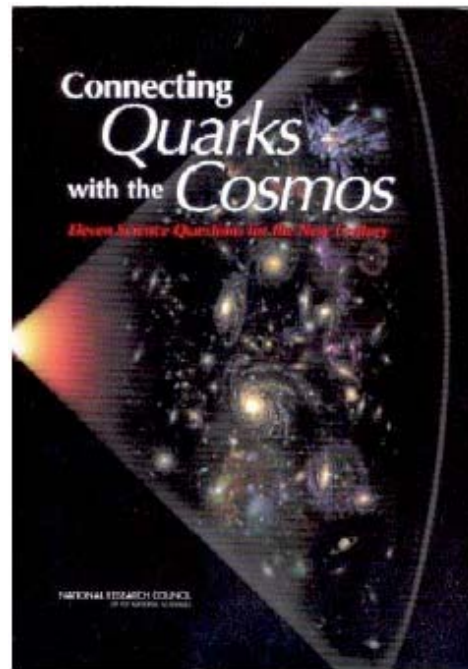
Sandia National Laboratories



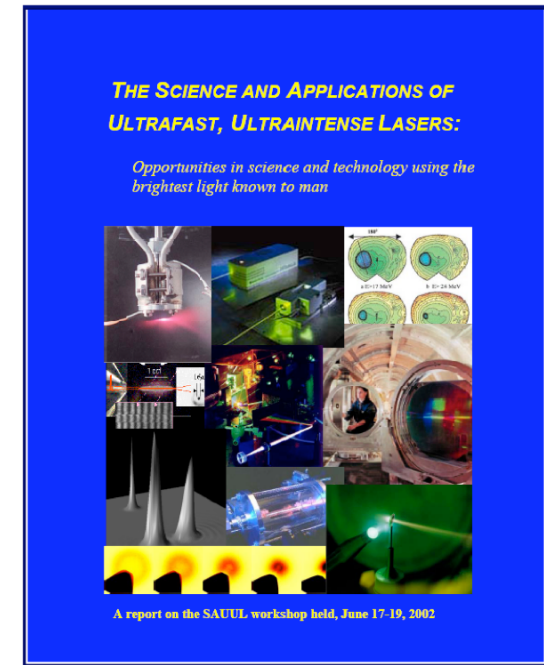
Several recent studies have highlighted High Energy Density Science



**“Frontiers in High Energy Density Physics”,
R. Davidson et al. 2004**



**“Connecting Quarks with the Cosmos: Eleven Science Questions for the New Century”,
M Turner et al. 2002**



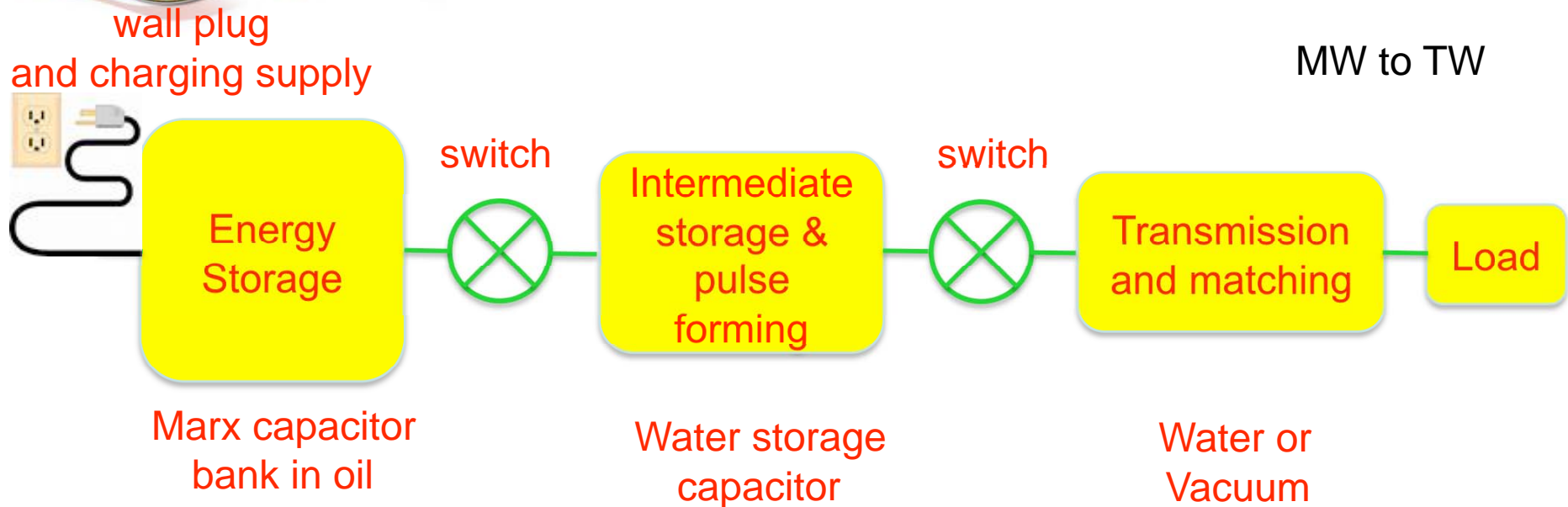
Science and Applications of Ultrafast, Ultraintense Lasers (SAUUL)



Scales of energy density

- 0.001 MJ/kg water at 100 m dam height
- 0.5 MJ/kg Li ion battery
- 1.968 MJ/kg water
- 7.5 MJ/kg stick of dynamite
- 33 MJ/kg Low Earth Orbit
- 45 MJ/kg gasoline
- 170,000 MJ/kg typical z-pinch implosion at 20 MA
- 3.5 million MJ/kg fission of 3.5% enriched U-235
- 337 million MJ/kg DT fusion
- 645 million MJ/kg hydrogen fusion (Sun)
- 89.9 billion MJ/kg ($E=mc^2$, antimatter-matter annihilation)
- (see http://en.wikipedia.org/wiki/Energy_density)

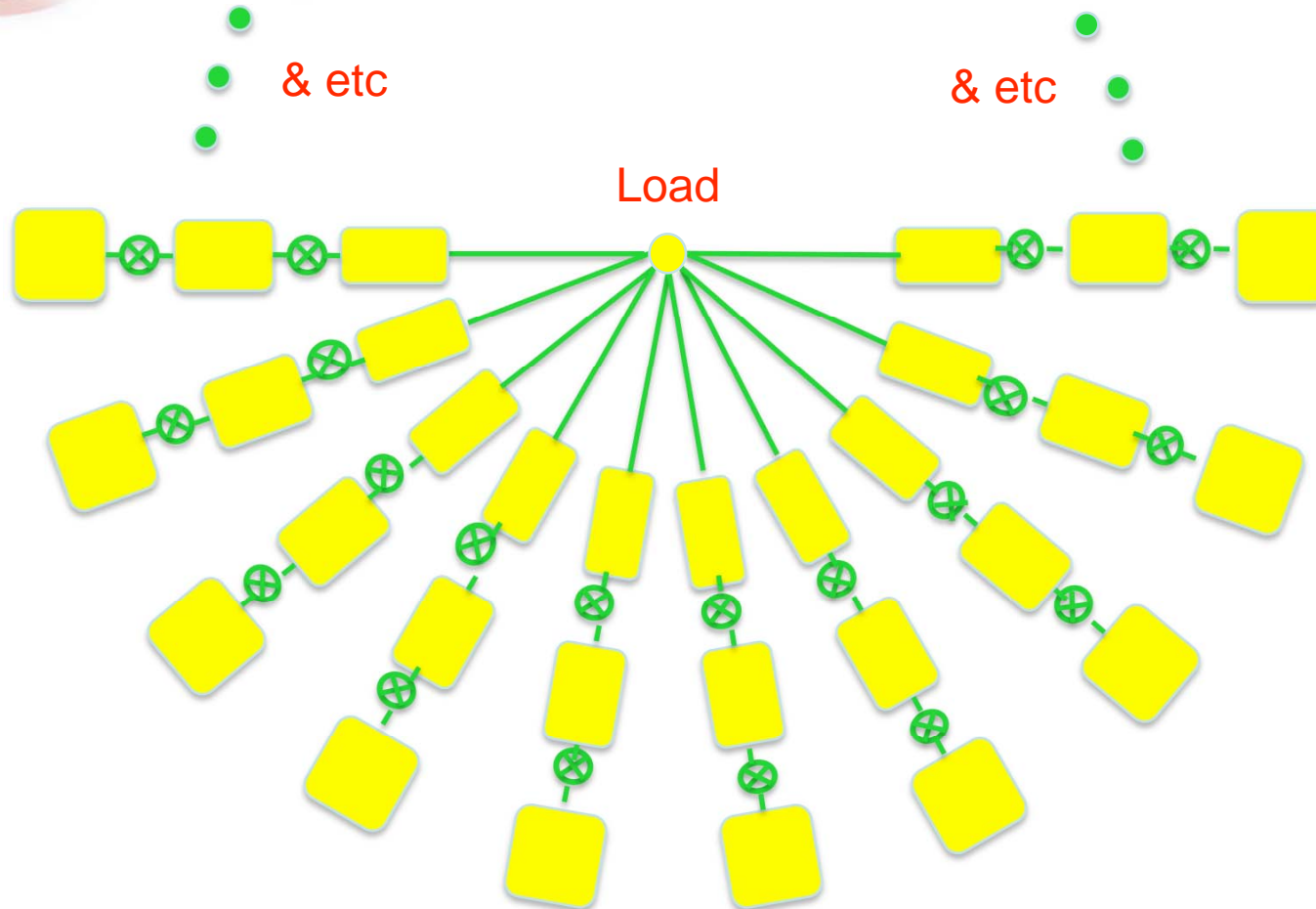
Pulsed-power technology produces high electrical power using rapid switching and pulse compression



- Typically the pulse is compressed in both space and time
- The load produces the last step in pulse compression and power gain



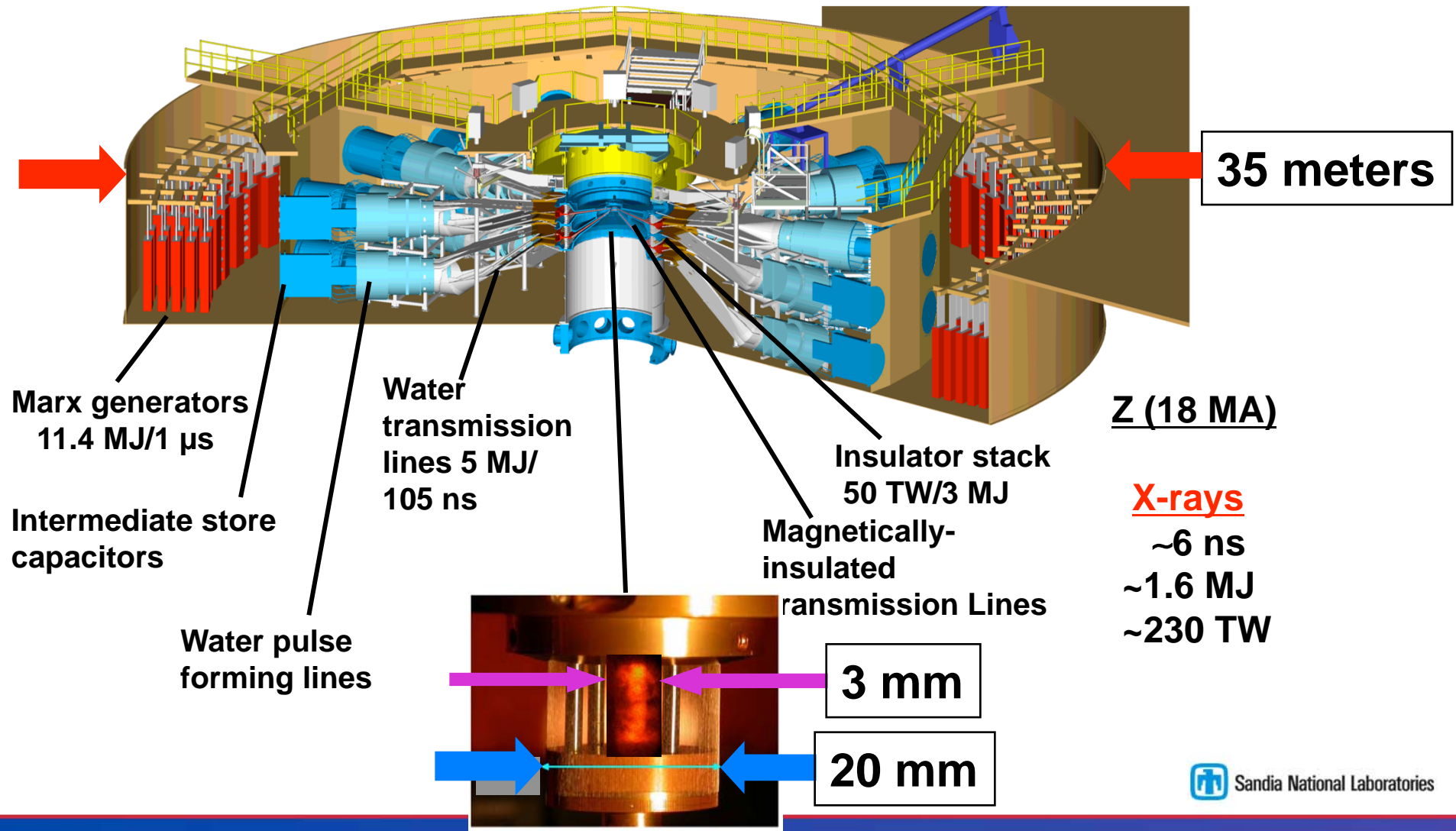
Multiple modules are used to achieve the highest powers



- Laser triggered gas switches are used to synchronize the pulses from the individual modules to within a few ns

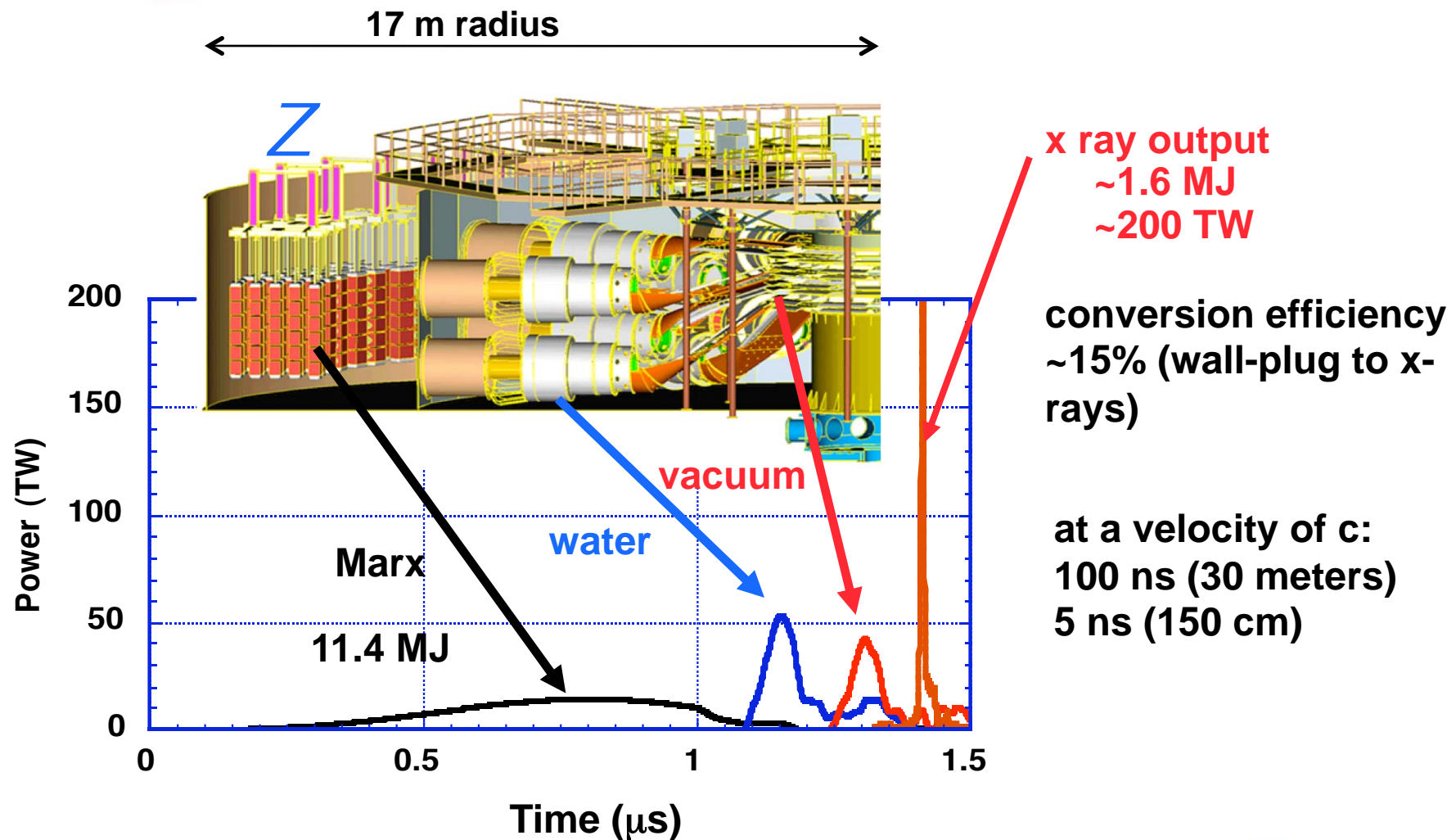


Z creates a ~100 ns, ~25 MA pulse to drive the world's most energetic x-ray source





Pulsed-power provides compact, efficient, power amplification

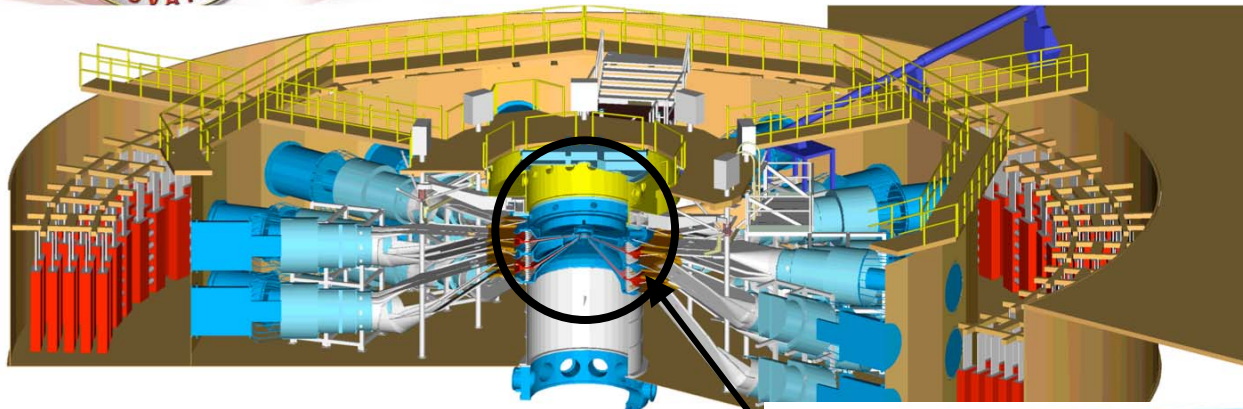




We compress energy in space and time on Z

- Compression in height X 625
 - 625 cm tank Marx height to 1.0 cm hohlraum height
- Compression in radius X 1375
 - 1650 cm in Marx tank radius to 1.2 cm hohlraum radius
- Compression in time X 1.7×10^{10}
 - 2 minute Marx charge to 7 ns pinch output
- Total power density compression factor ~ 1.5×10^{16}
 - ($\eta \times \text{area} \times \text{time}$, $\eta \sim 0.4$)

Self magnetically insulated transmission lines (MITLs)
allow power flow at electric fields far above breakdown



OIL

WATER

VACUUM

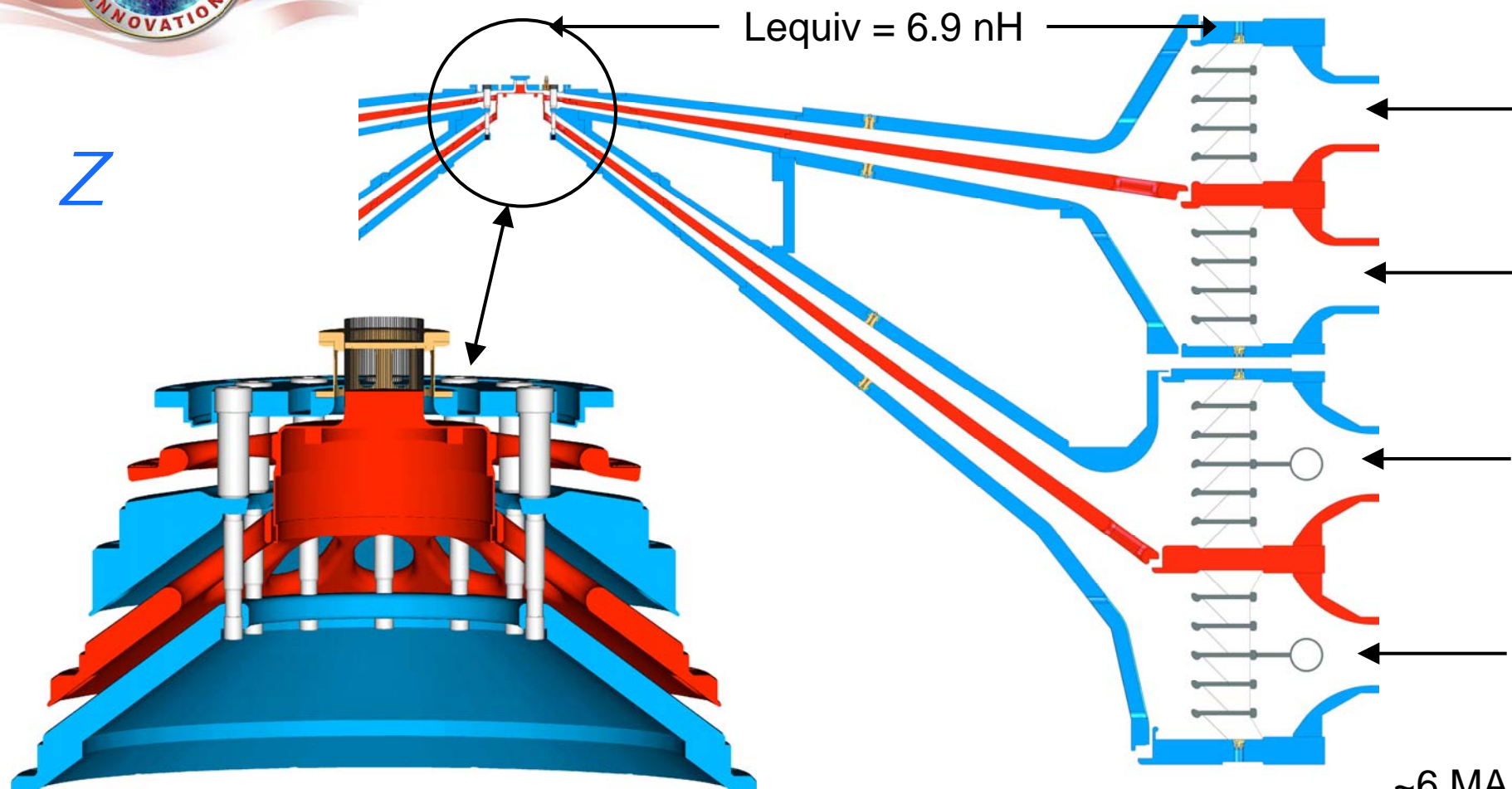
50 TW
electrical
power

$E \sim 10 \text{ MV/cm}$
 $B > 40 \text{ MG}$
 $E \sim 100 \text{ kV/cm}$
 $B \sim 6 \text{ kG}$
Vacuum Insulator Stack

Z



High current is achieved by adding four low inductance conical transmission lines in parallel



- Z produced currents up to 20 MA in 100 ns
- Z-R has produced currents up to 26 MA in 100 to 250 ns



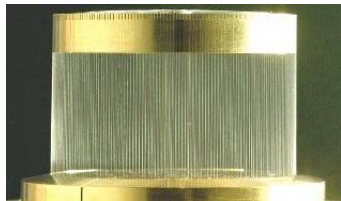
We can use high currents to push plasmas in different ways for different applications

High Current

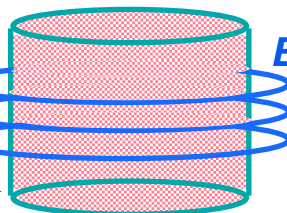
Cylindrical magnetic pressure

Planar magnetic pressure

wire array

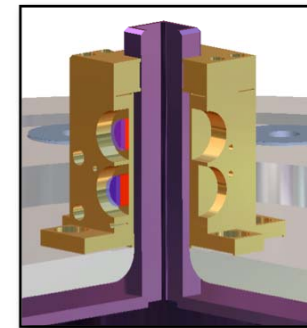
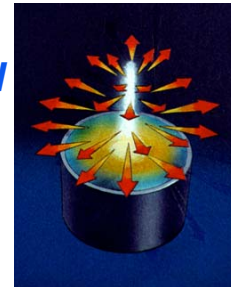


Current

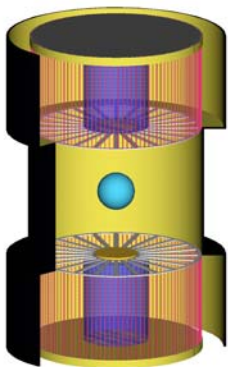


B-Field

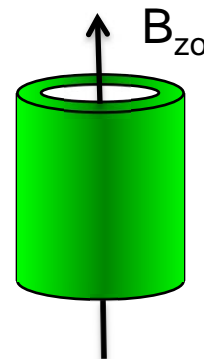
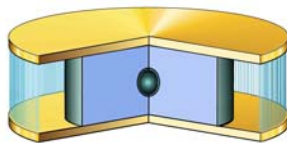
$J \times B$ Force



Inertial Confinement Fusion (ICF)

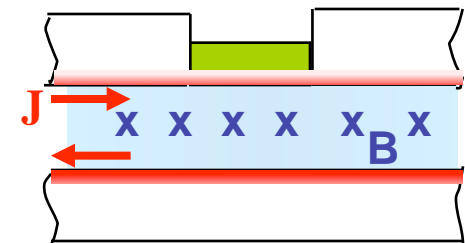


Indirect



Direct (MTF)

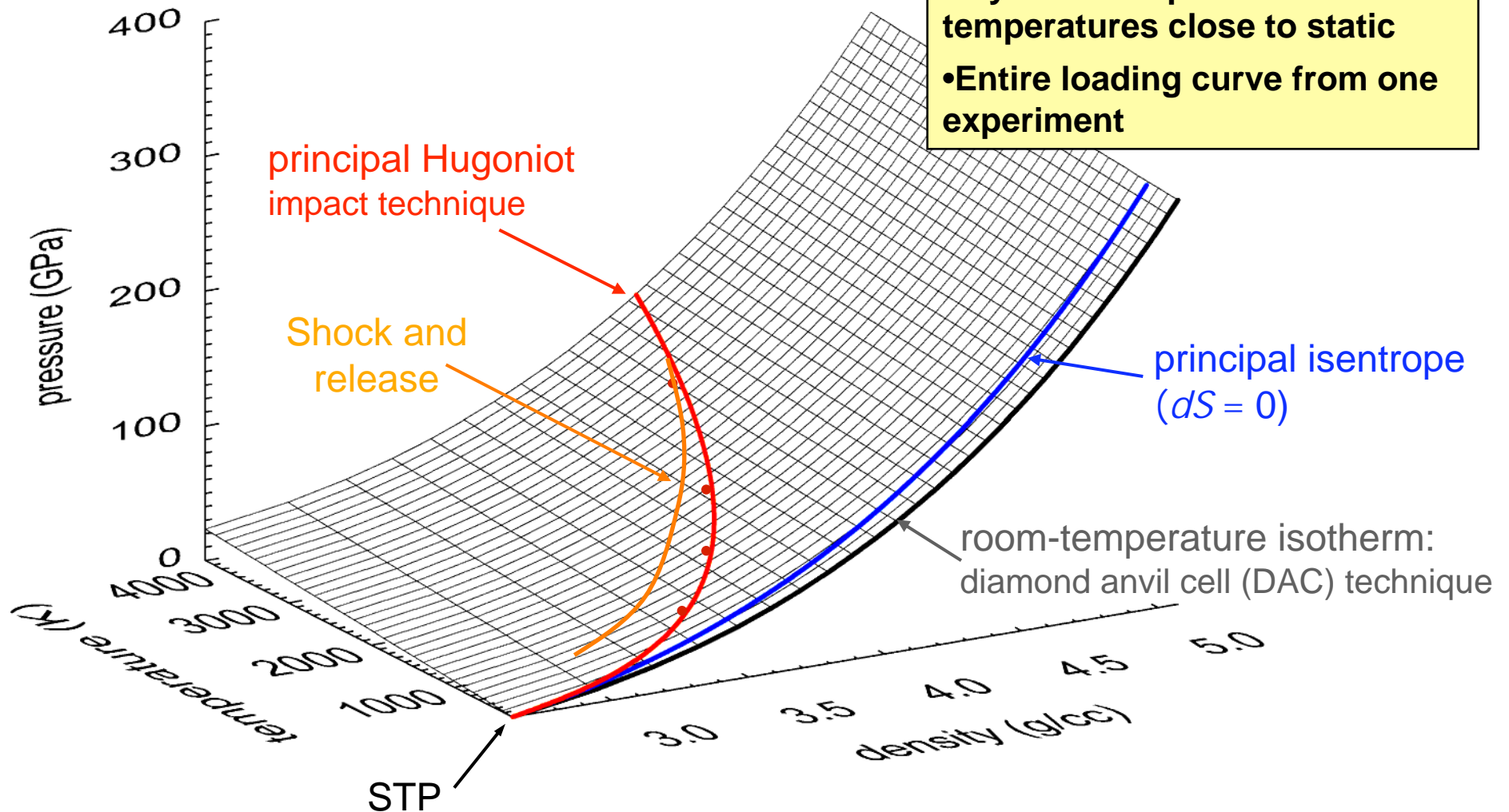
Material Properties





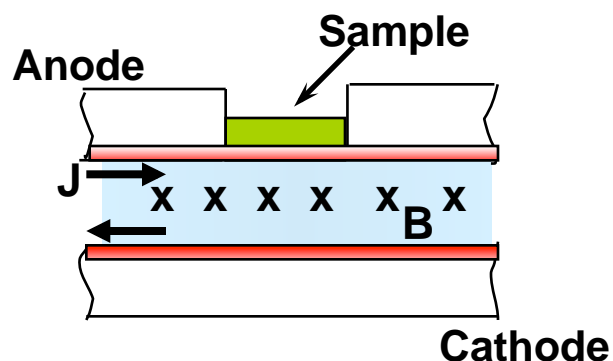
Magnetic pressure can be used to examine fundamental material properties at HED conditions

- Dynamic response at temperatures close to static
- Entire loading curve from one experiment





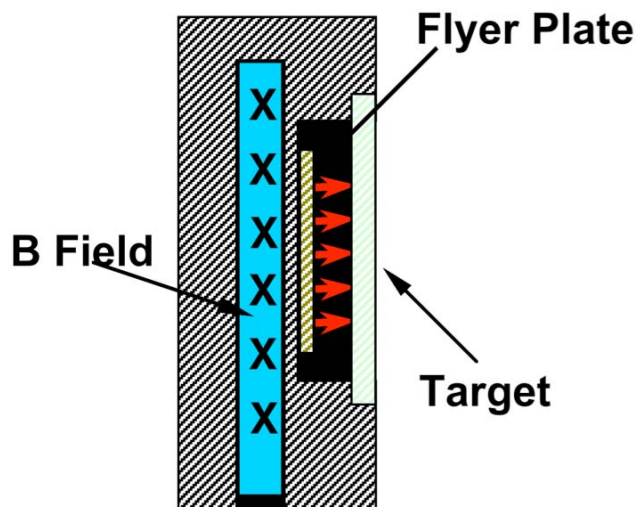
Techniques developed on Z have enabled major advances for accurate EOS studies



Isentropic Compression Experiments (ICE)*

Magnetically produced Isentropic Compression Experiments (ICE) to provide measurement of continuous compression curves to ~3 Mbar
- previously unavailable at Mbar pressures

* Developed with LLNL



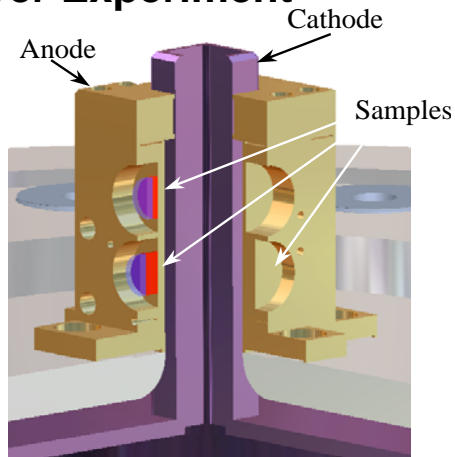
Magnetically launched flyer plates

Magnetically driven flyer plates for shock Hugoniot experiments at velocities to ~ 34 km/s
- exceeds gas gun velocities by ~ 4X and pressures by ~ 7-8X with comparable accuracy

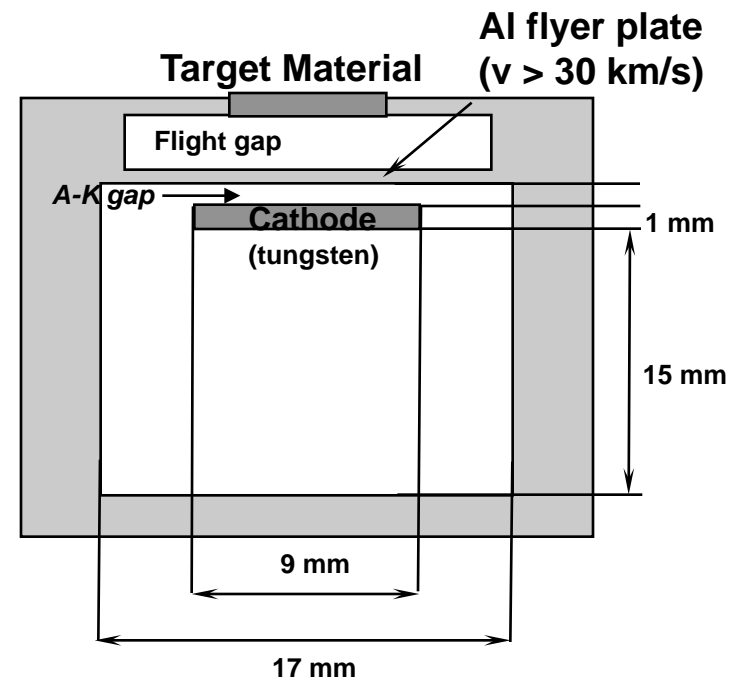


Numerical simulations enable the design of high velocity flyer plate experiments

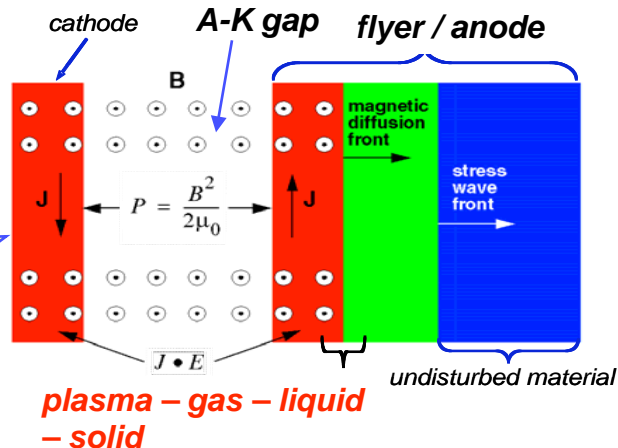
Setup for High Velocity Flyer Experiment



2D Model for Numerical Simulation



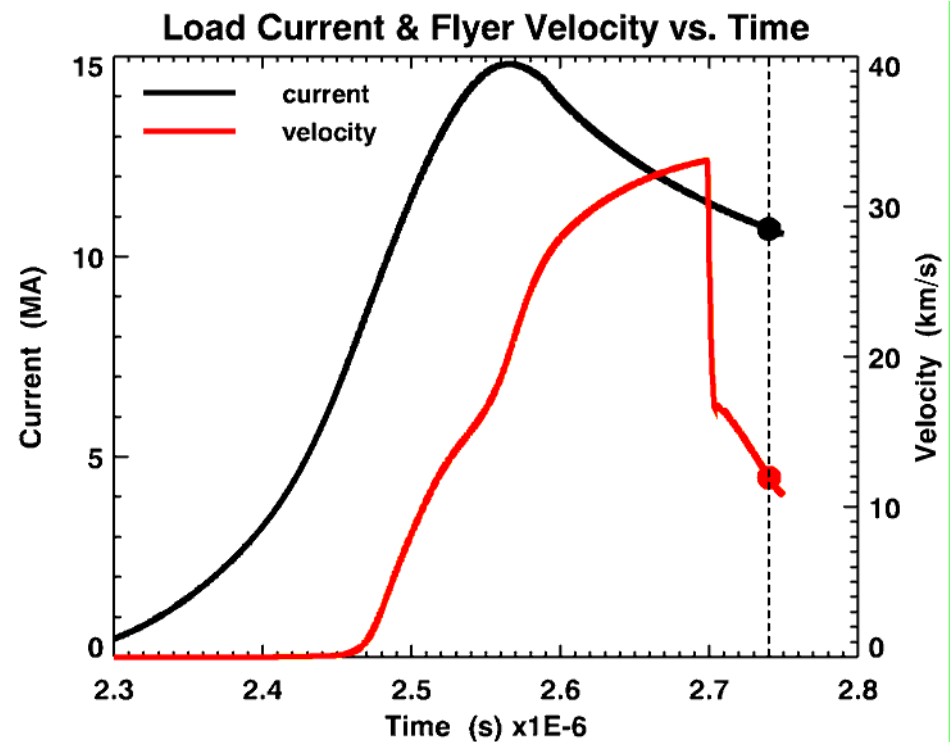
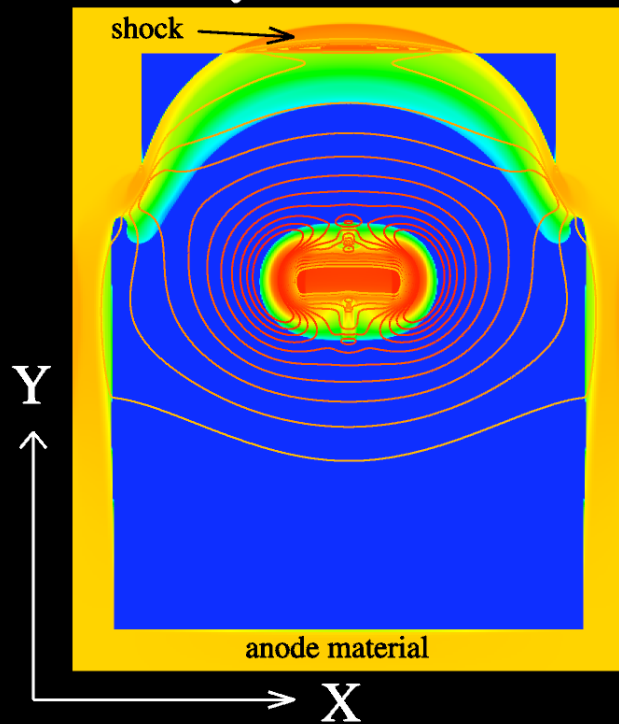
Physics of magnetically accelerated flyer in 1D.





Snapshots from 2D numerical simulation show dynamics of flyer in high velocity

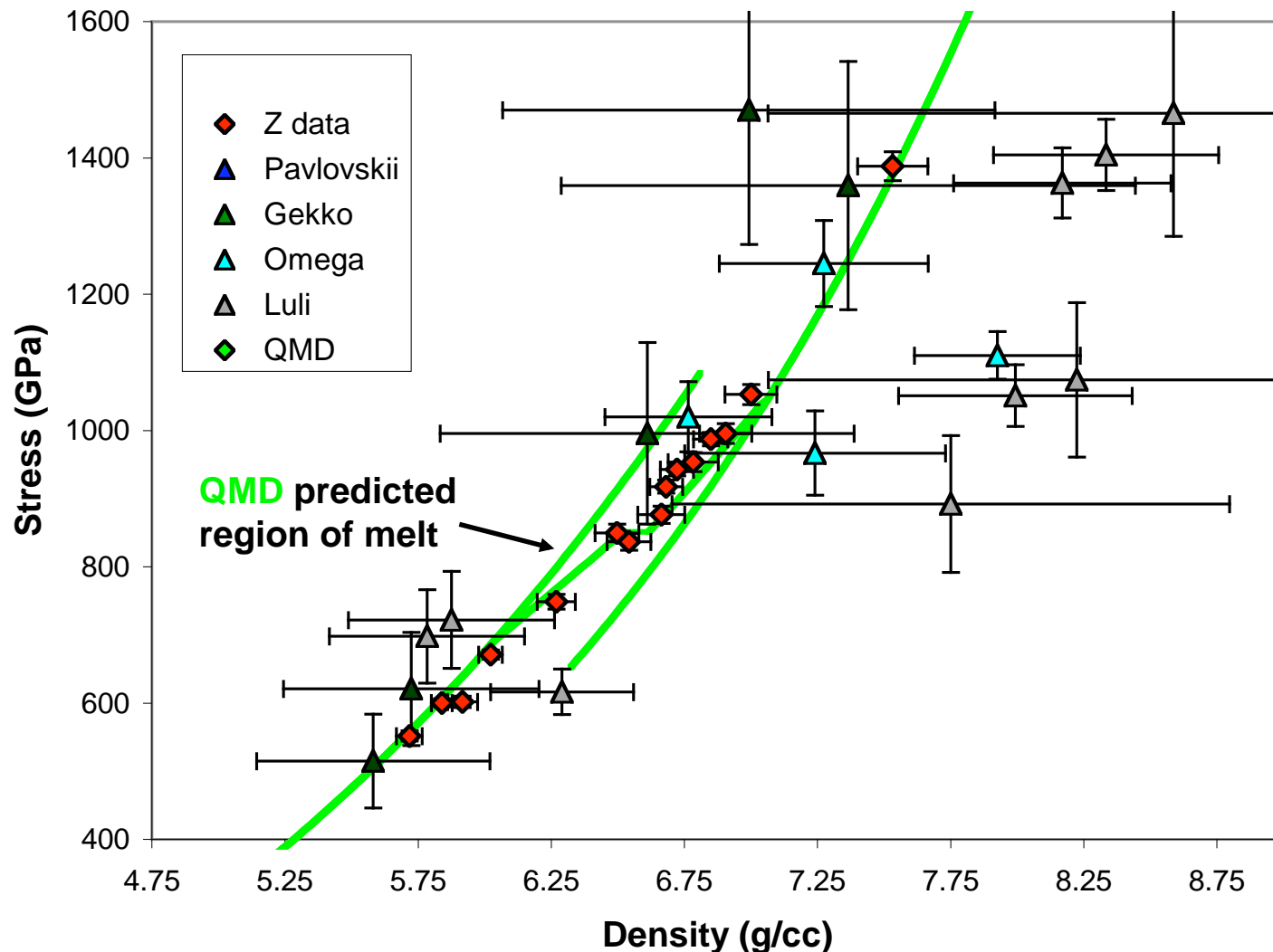
34 km/s Flyer $t = 2.7400\text{e-}06$ s





Z answered important questions about the properties of diamond at high pressure

stress versus density for diamond



- The Z data was obtained in 1 week
- Measurements on Z have a accuracy of $\leq 1\%$



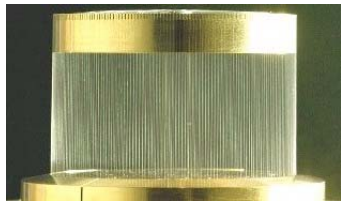
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High Current

Cylindrical magnetic pressure

Planar magnetic pressure

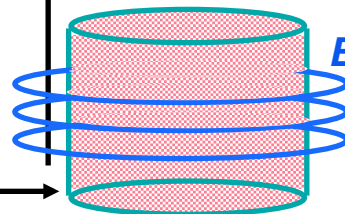
wire array



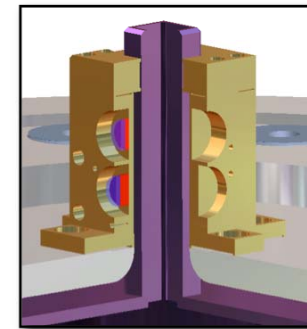
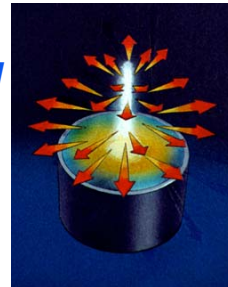
Current



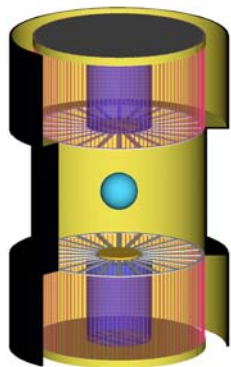
$J \times B$ Force



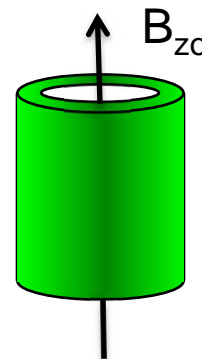
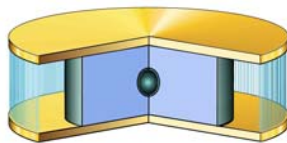
B-Field



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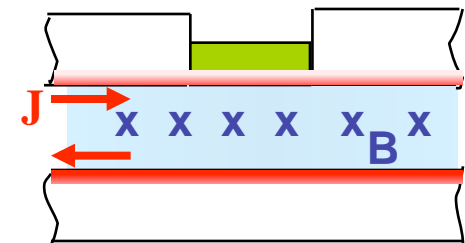


Indirect



Direct (MTF)

Material Properties



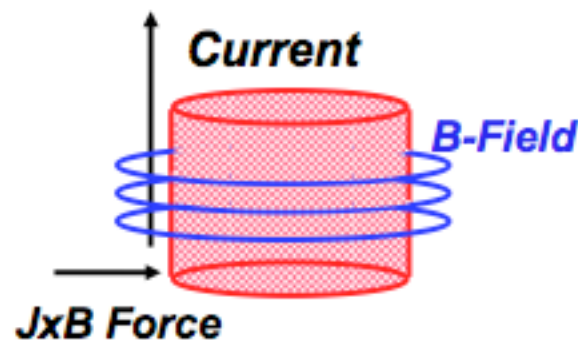
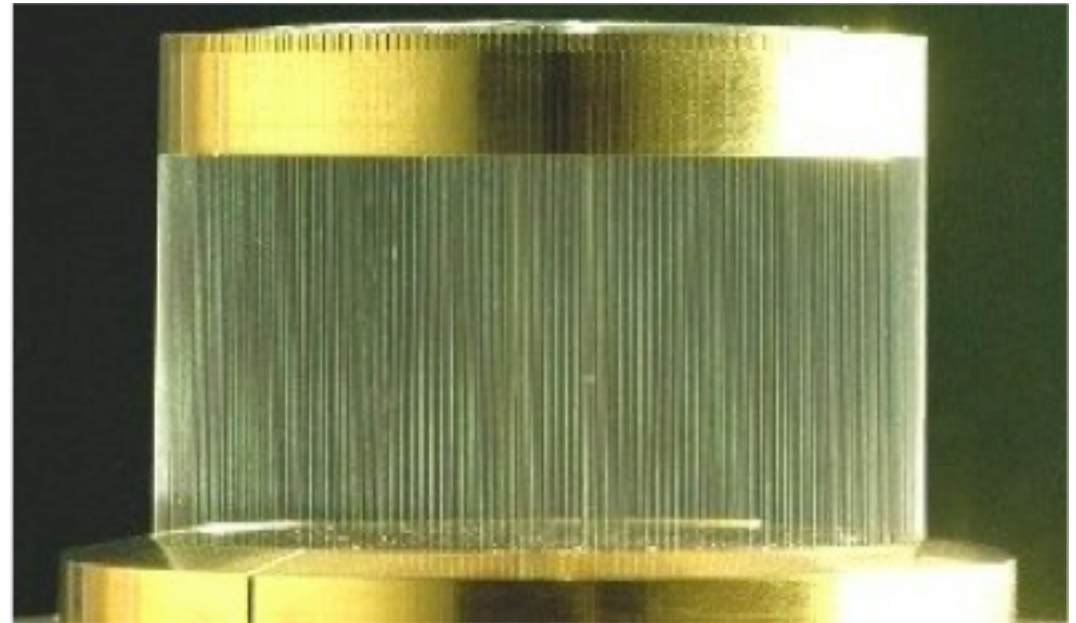


Wire arrays are a “simple” extension of the two wire problem

Instead of 2 wires, use ~300 wires in a cylindrical array. The $J \times B$ force accelerates the wires toward the array axis.

Instead of 1 mA or 1 A, use 20,000,000 Amperes of current in the array, delivered in a ~100 ns current pulse.

The result is the creation of soft x rays (~0.1-10 keV) with 10-15% efficiency from the stored electrical energy

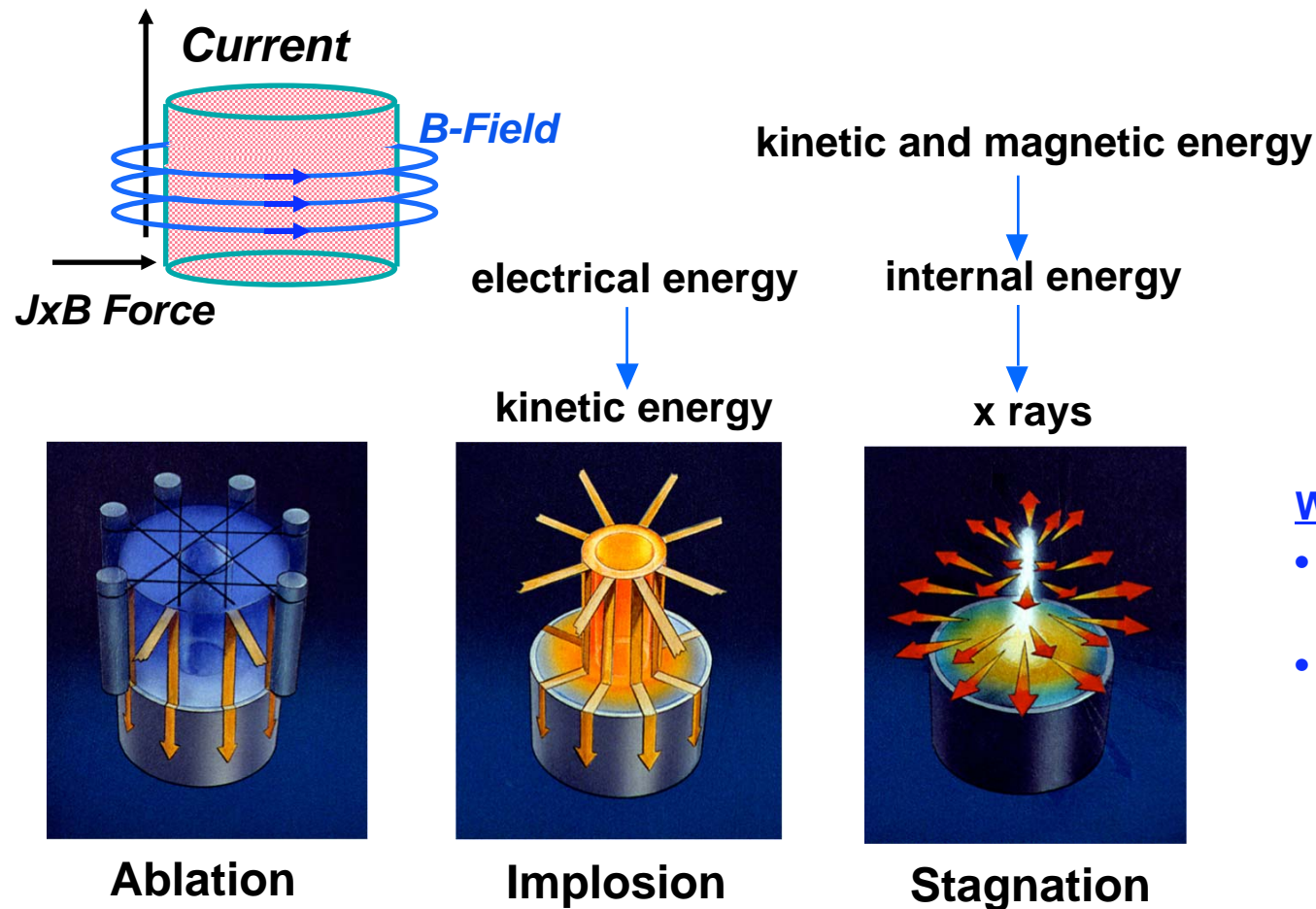


Z-pinch loads:

Wire Array
Gas Puffs
Foil/Liner
Foam



Magnetically-driven z-pinch implosions efficiently convert electrical energy into radiation

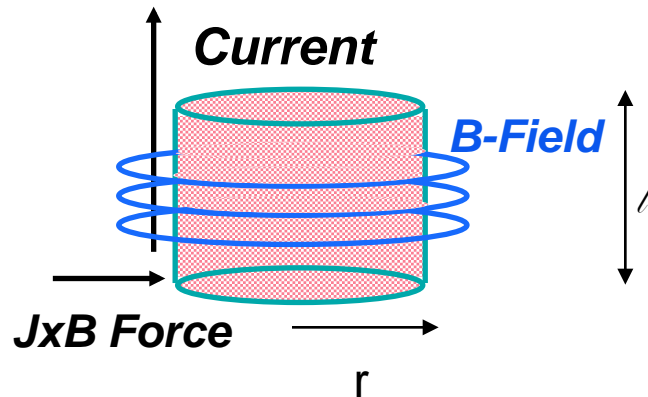


Wire z-pinch loads:

- Energy: x-ray \approx 15% of stored electrical
- Power: x-ray \approx 2-5 x electrical



How much magnetic energy can we put into kinetic energy ?



According to the MHD equations we can treat the magnetic field working on the wires as a magnetic pressure $\sim B^2$

For a thin shell with all the current on the outside:

$$ma = F$$

$$m\ddot{r} = PA \quad A = 2\pi rl \quad P = \frac{B^2}{8\pi} \quad B \approx \frac{I}{5r}$$

cgs units

$$\frac{m}{l} \ddot{r} \approx \frac{I^2}{100r}$$

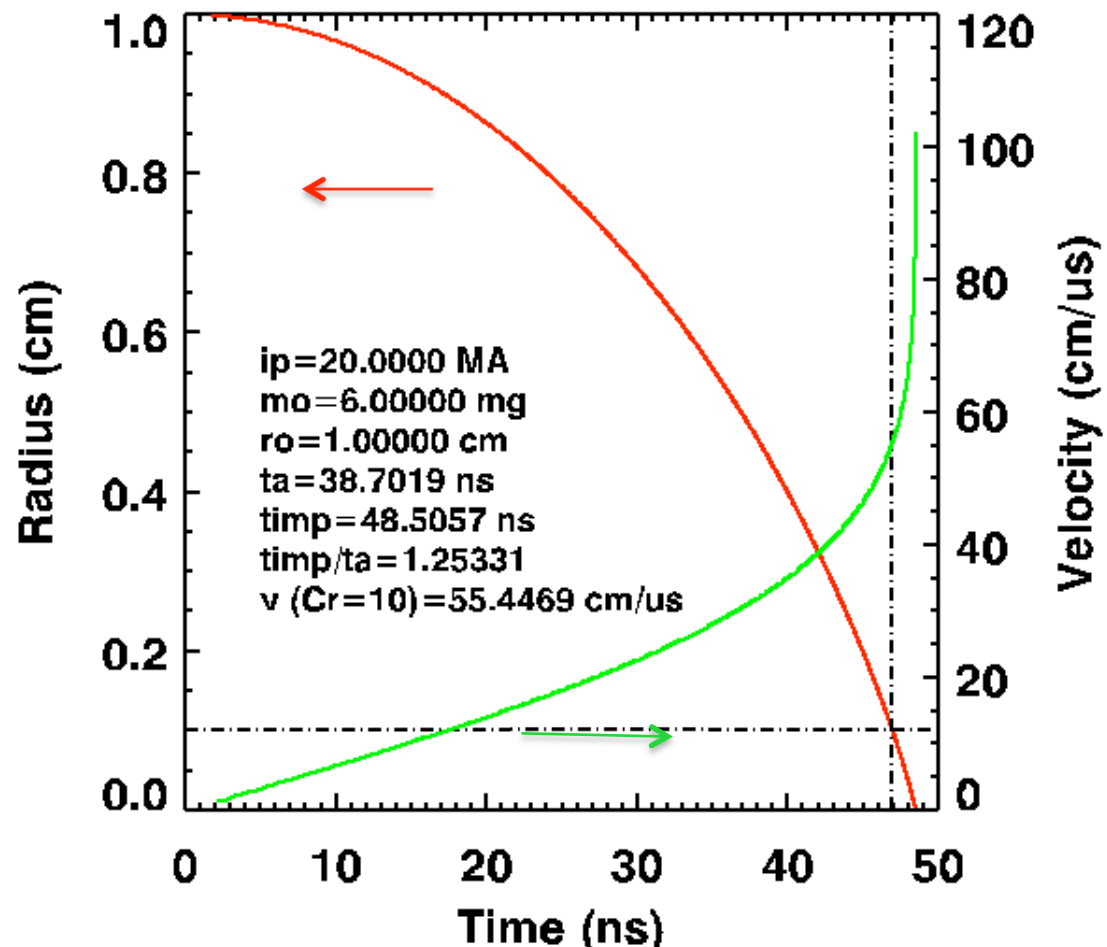
Note magnetic pressure increases for decreasing r

We can integrate this equation to get the kinetic energy given to the imploding shell



The analytic solution with $I \sim \text{constant}$ shows increasing magnetic pressure accelerating the pinch

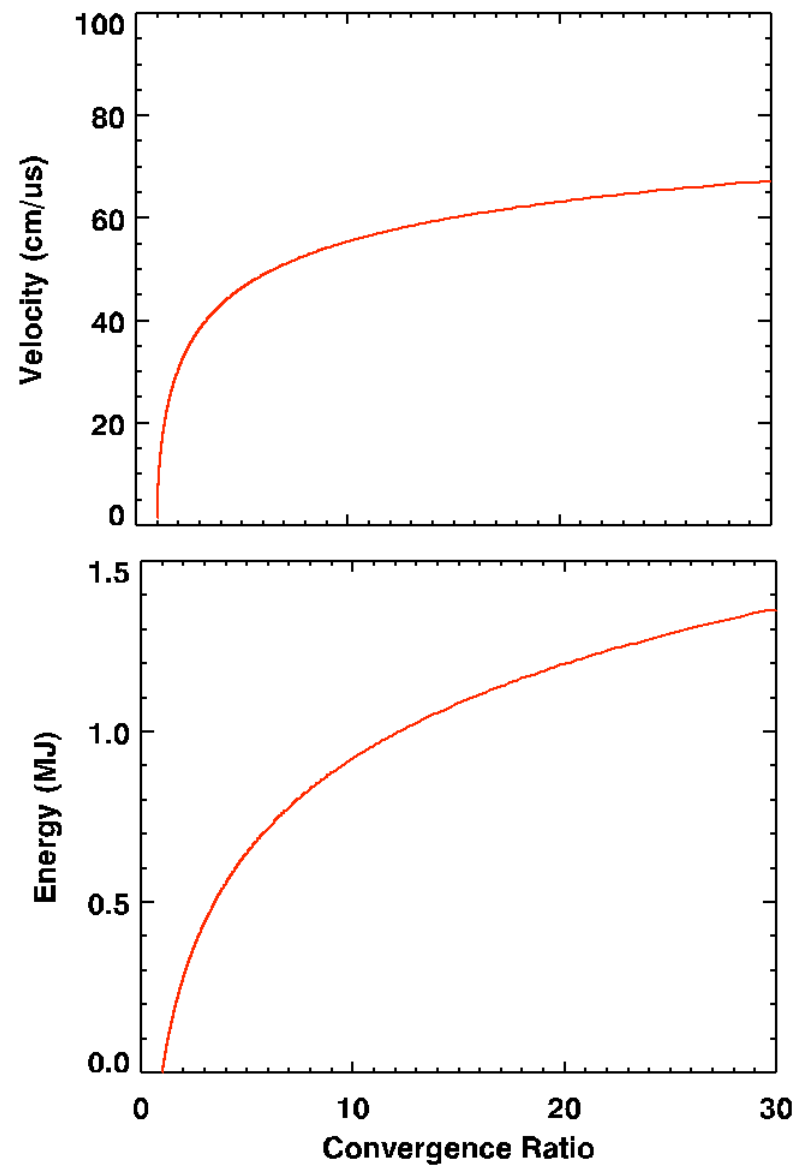
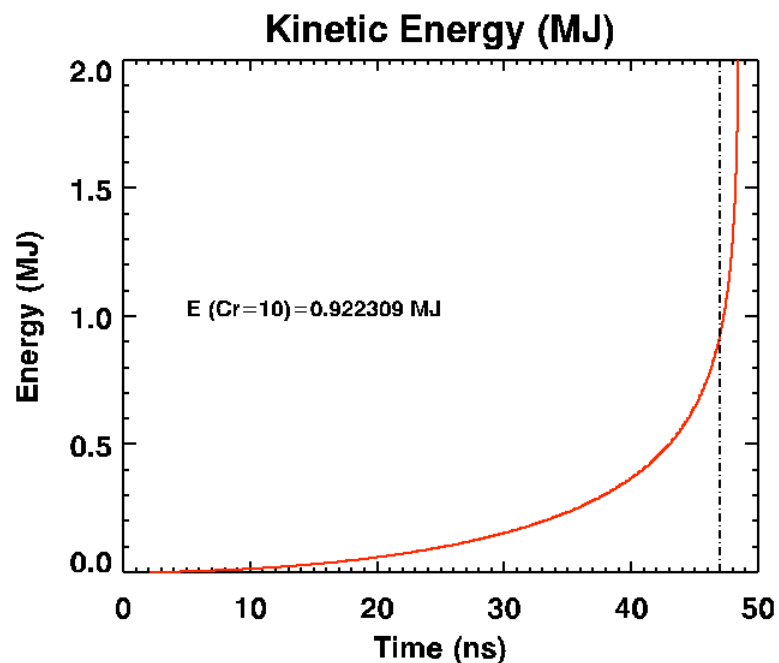
$$P = \frac{B^2}{8\pi} \sim \left(\frac{I}{r(t)} \right)^2$$



- Velocity at convergence ratio of 10 ~ 50 cm/μs ~ 1,100,000 mph
- Princeton to LA in 8 seconds
- 1/600th of the speed of light



Significant kinetic energies are coupled to the axis





We employ kinetic energies of ~ 1 MJ in every day objects



- $m_{F150} = 2950$ kg
- $v_{F150} = 94$ km/hour (58 mph)
- $E = 1$ MJ
- In a typical z-pinch, this 1 MJ is released in 5 ns

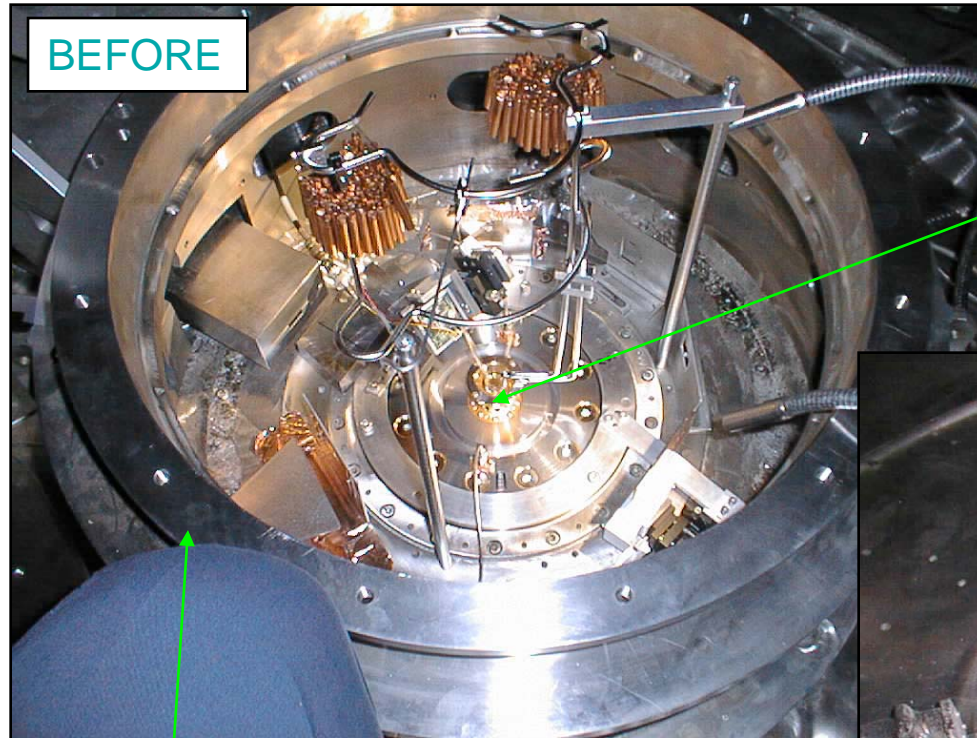


An energy of 1 MJ

- Kinetic energy of F150 at 60 mph
- 0.48 x energy in a stick of dynamite
- 100 W light bulb uses 1 MJ in 2.8 hours



The 2-3 MJ of magnetic energy and radiation destroys
by the debris following the radiation burst

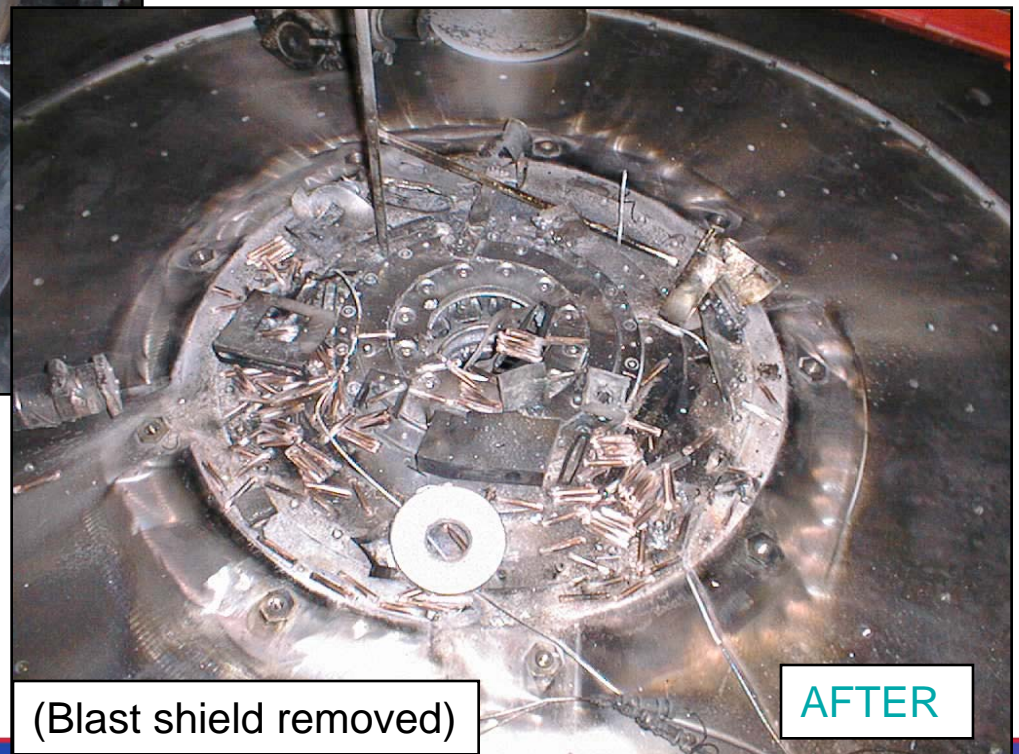


BEFORE

Wire array

Blast shield

Equivalent to 2 lbs high explosive
released in a few ns in <1 cc volume!

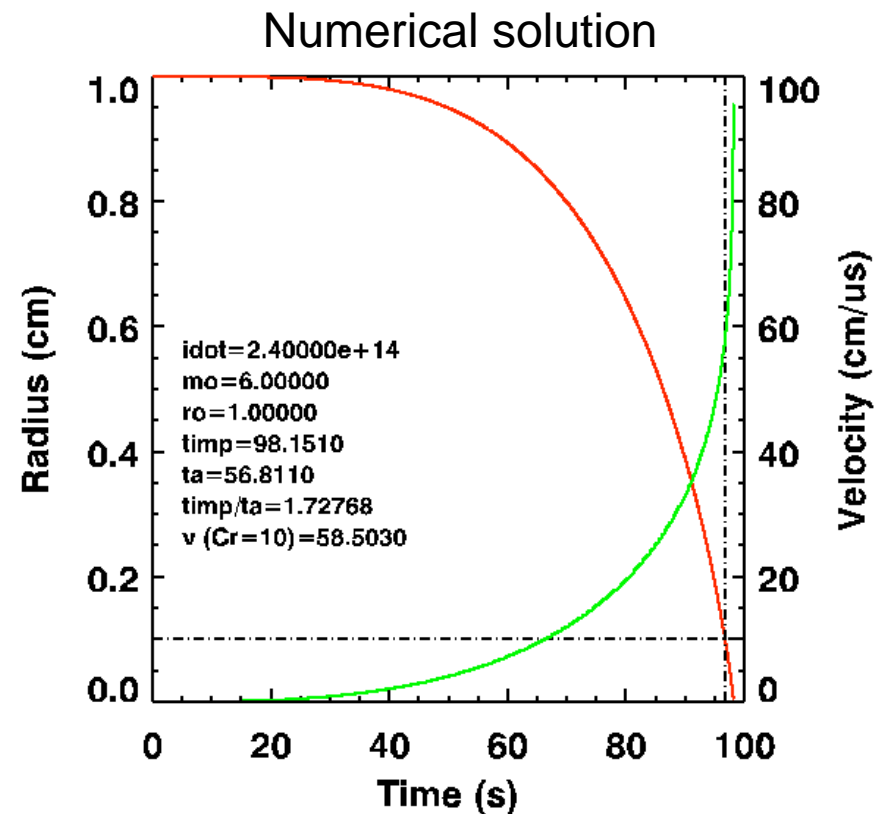
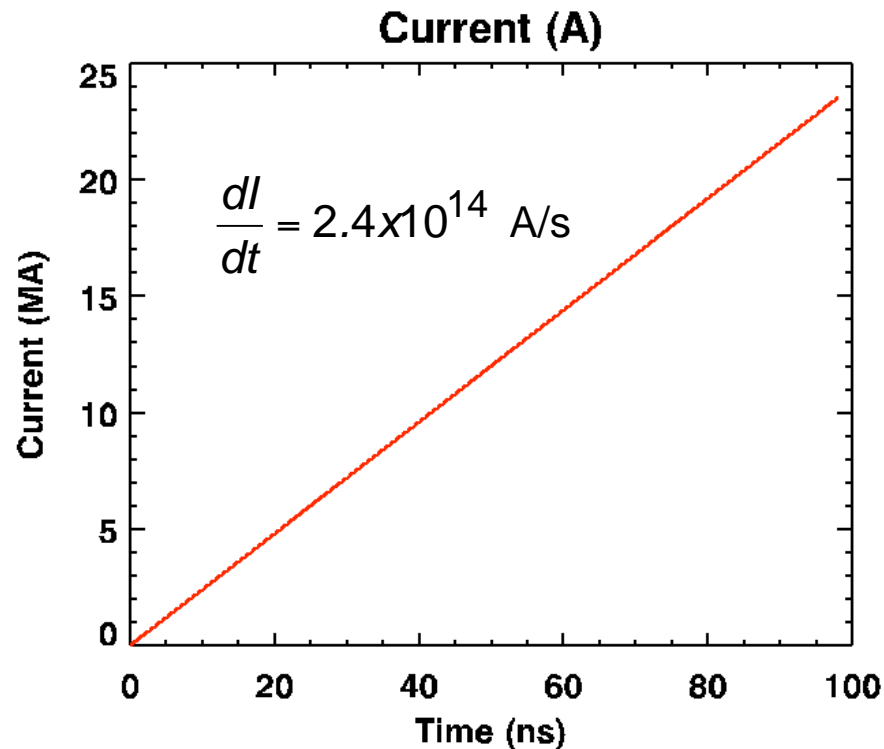


(Blast shield removed)

AFTER



A real pulsed power driver does not produce a constant current

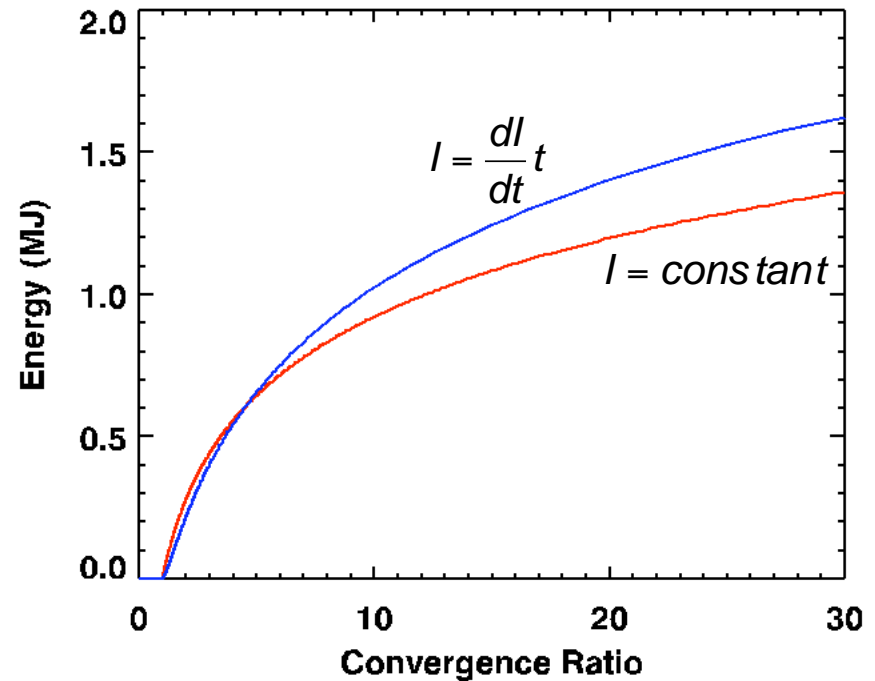
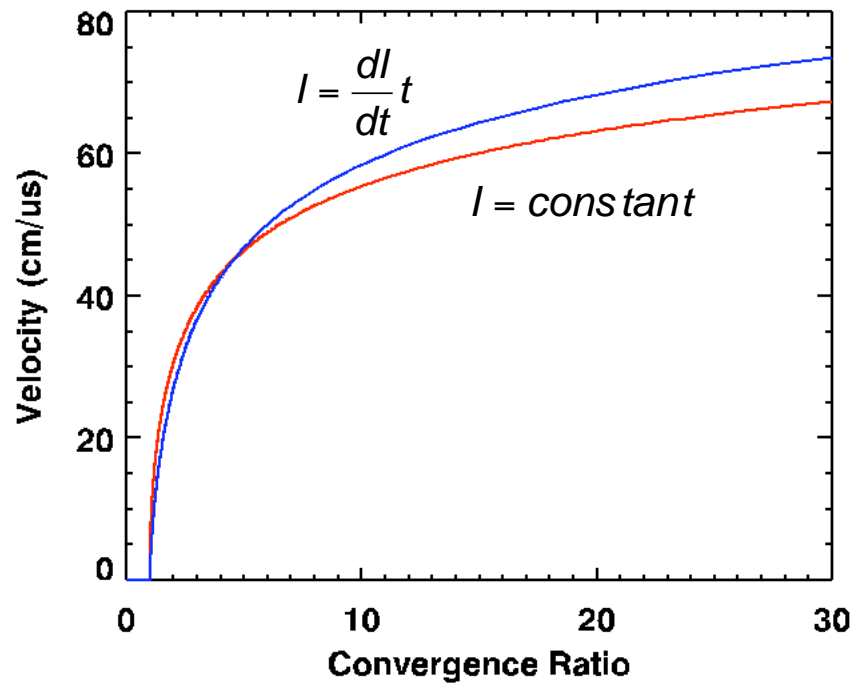


- A pulsed power driver is an inductor charging problem
- Typically current waveforms are linearly rising or $\sim \sin(\omega t)$

$$I(t) = \frac{\int V(t) dt}{L}$$



The kinetic energy just depends on the peak current and the convergence: insensitive to shape of $I(t)$



$$\text{KineticEnergy} = \frac{1}{2}mv^2 \sim \ell l_{\max}^2 \ln\left(\frac{r_o}{r_{fin}}\right)$$

$$\text{ConvergenceRatio} = \frac{r_o}{r_{fin}}$$



The plasma is an efficient radiator at stagnation

- The wire arrays are made out of a high Z material like tungsten ($Z=74$)
- Because of this wire arrays are very good radiators of their kinetic energy (E_{KE}) at stagnation
- The magnetic field does additional work on the plasma at stagnation, some of which is radiated during the main radiation pulse (E^*)
- So we know how much energy will be radiated by a wire array, but the issue for ICF is power density

$$P = \frac{E_{KE} + E^*}{\tau}$$

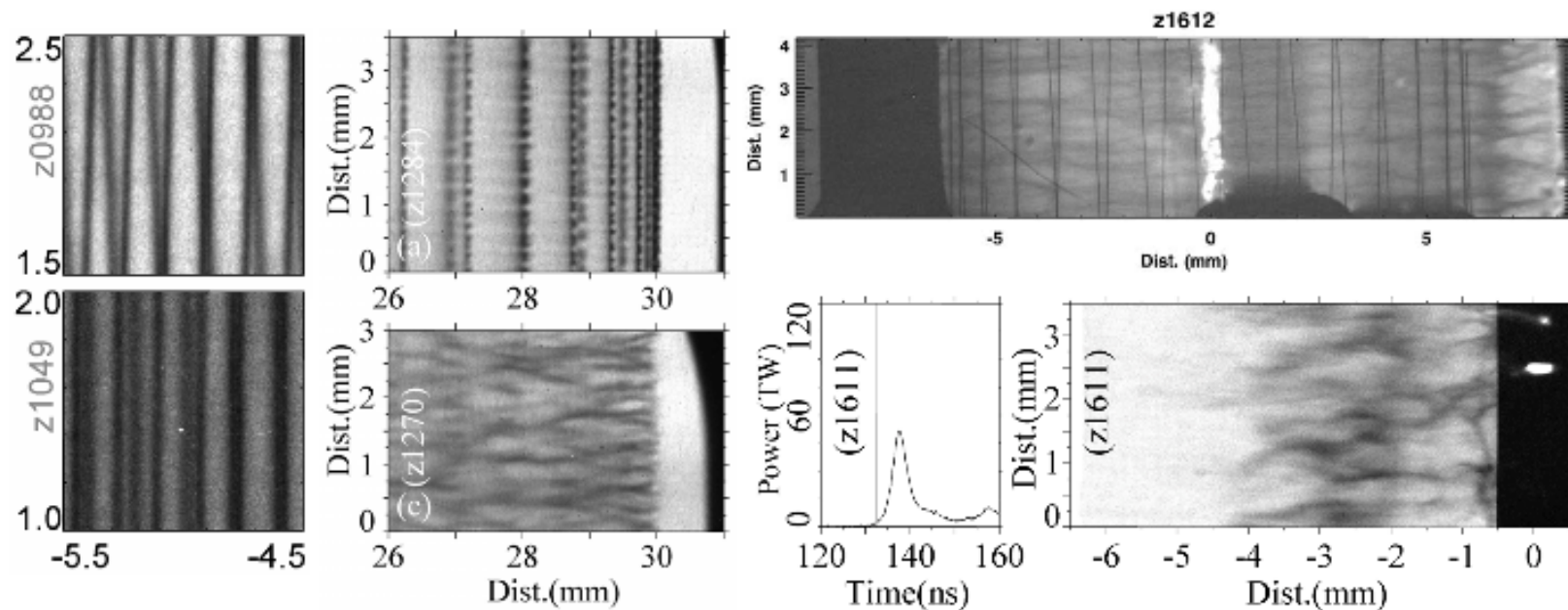
The shorter τ the higher the power will be! What determines τ ?



The magneto-Rayleigh-Taylor instability limits how short τ can be

A good guess for τ :
$$\tau \sim \frac{\delta r}{V}$$

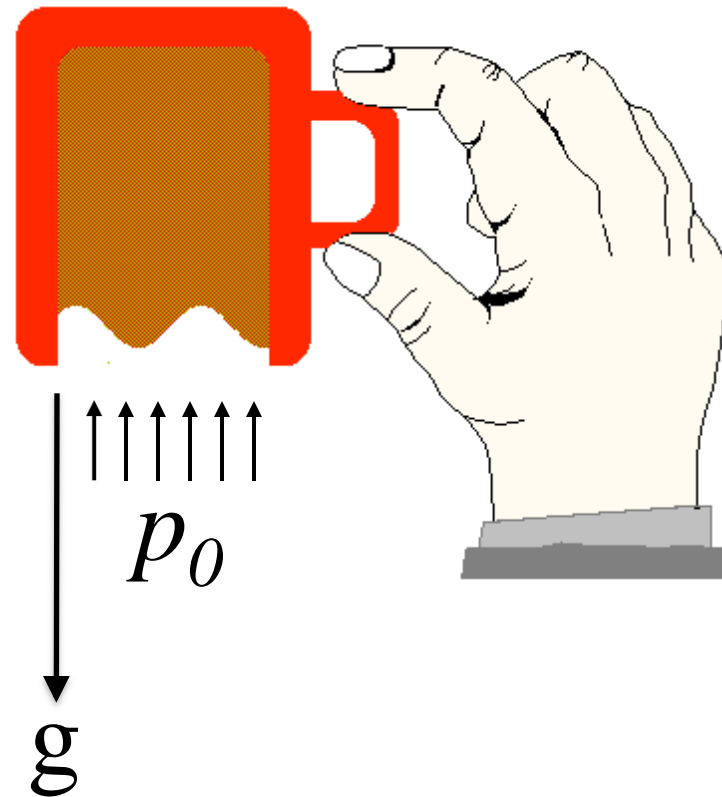
X-ray images of a wire array during its implosion!





A heavy fluid supported by a light fluid is unstable

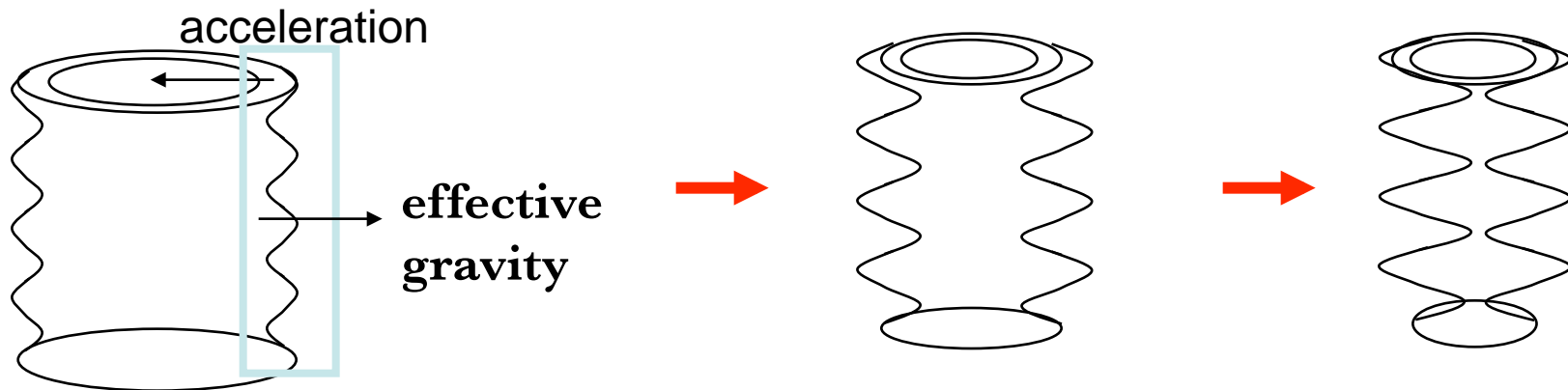
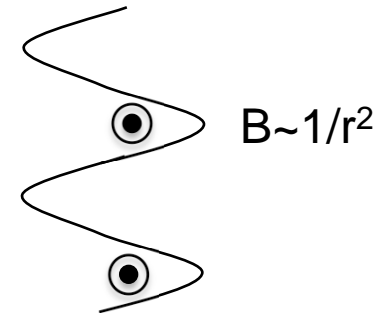
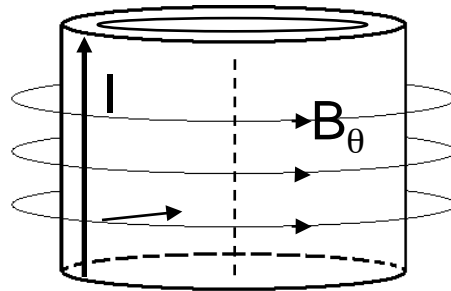
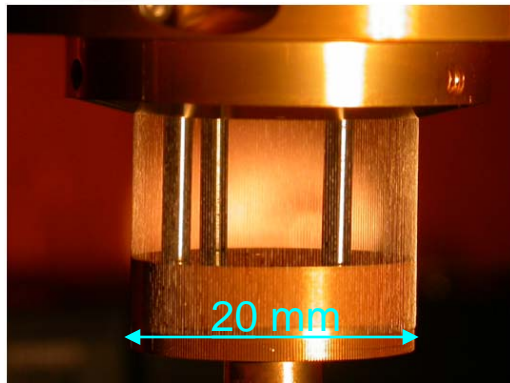
Atmospheric pressure p_0 is sufficient to support a coffee column 33 feet high against gravity. The RT instability helps us daily!



- This instability is the classical Rayleigh Taylor instability



The plasma is the heavy fluid, the magnetic field is the light fluid

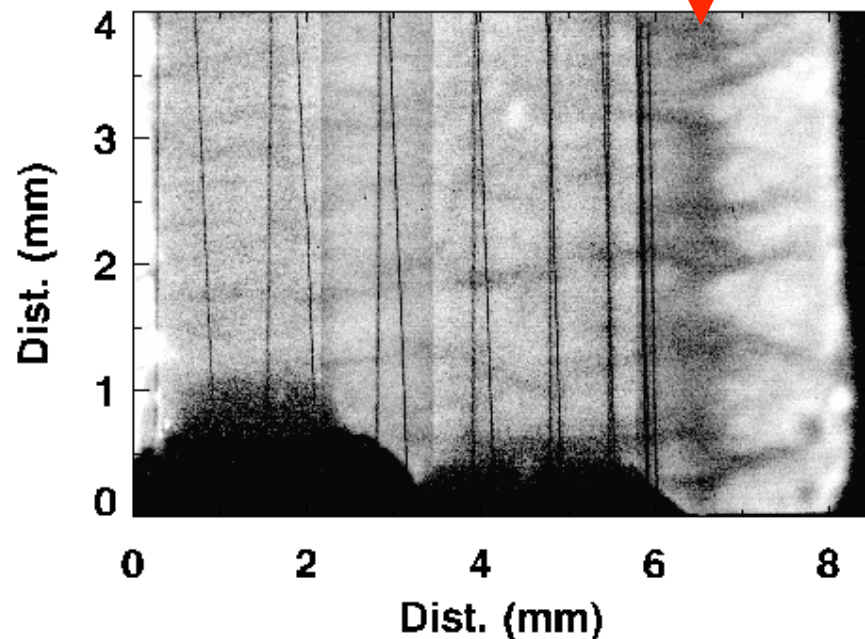


- The increase of $B \sim 1/r^2$ at the bubble position with respect to the spikes might be thought to increase the growth, but ...

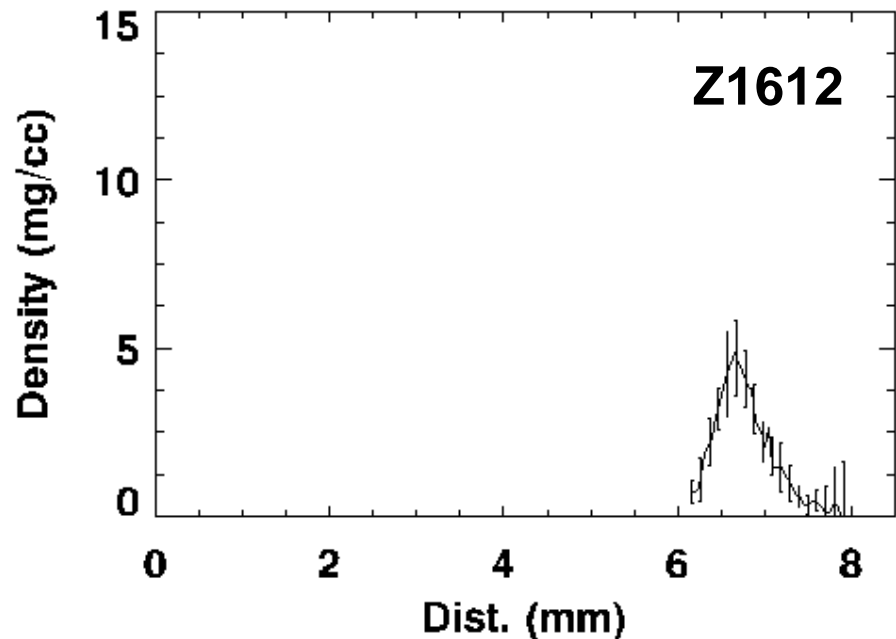


The magneto-Rayleigh-Taylor instability limits how short the pulsewidth τ can be

Mass profile of outer = profile of single

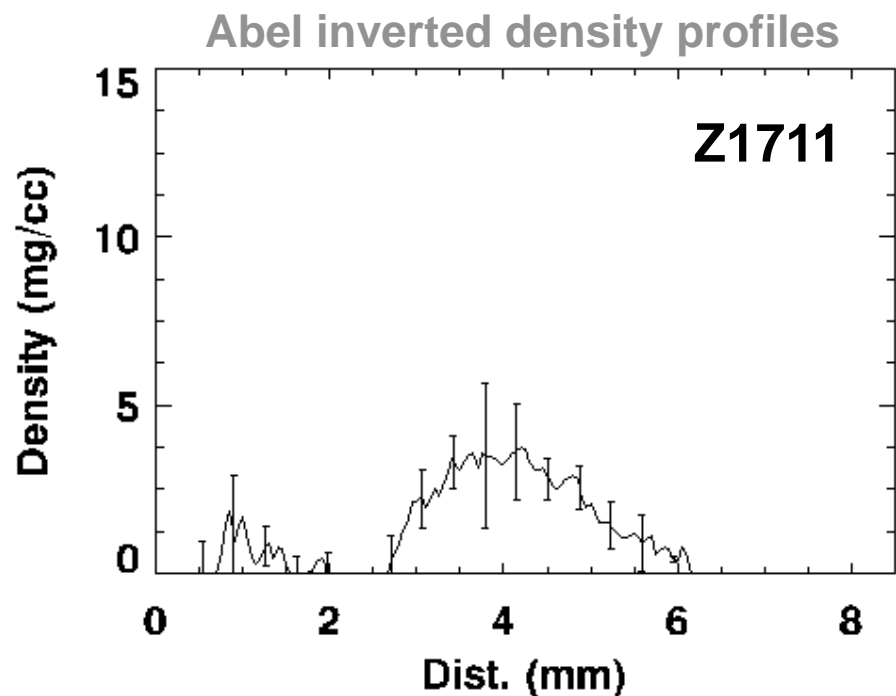
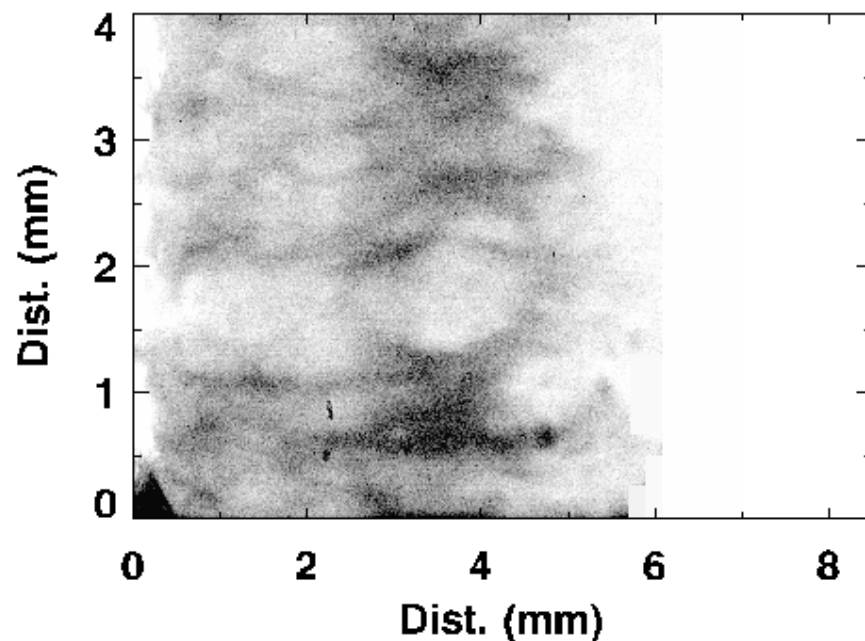


Abel inverted density profiles



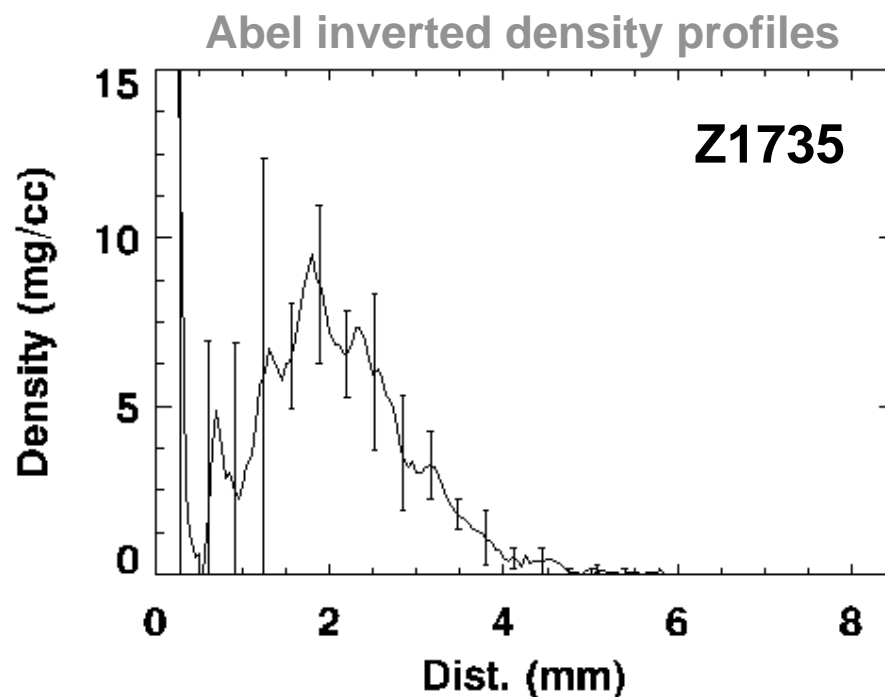
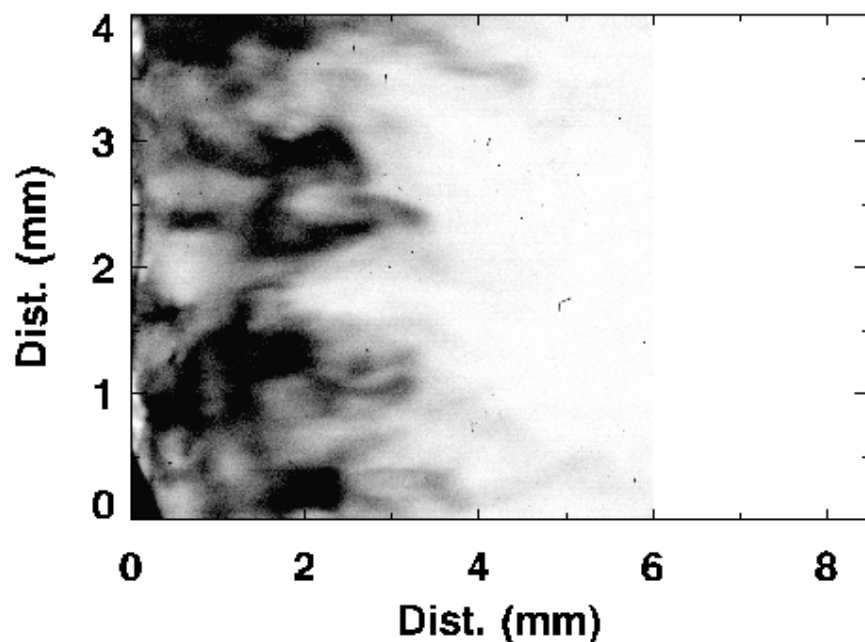


The magneto-Rayleigh-Taylor instability limits how short the pulsewidth τ can be





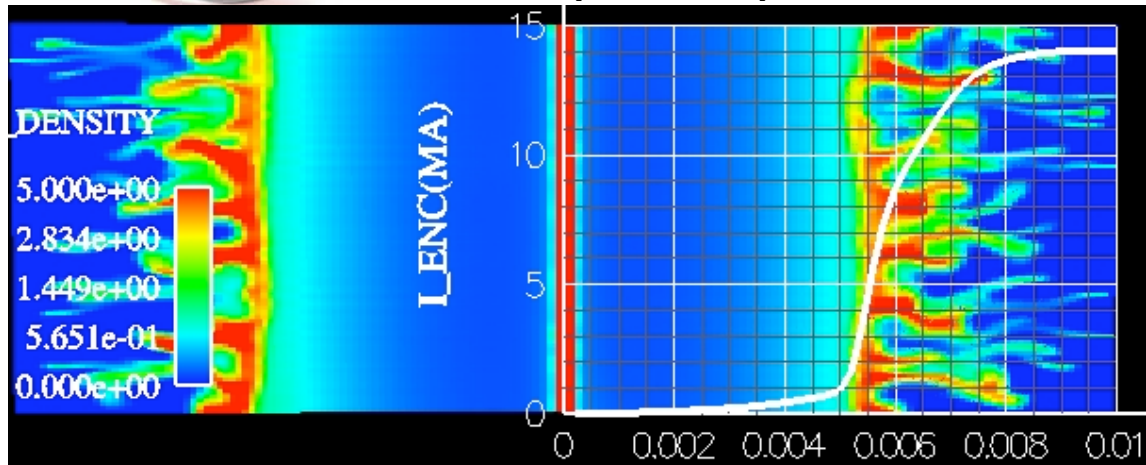
The magneto-Rayleigh-Taylor instability limits how short the pulsewidth τ can be





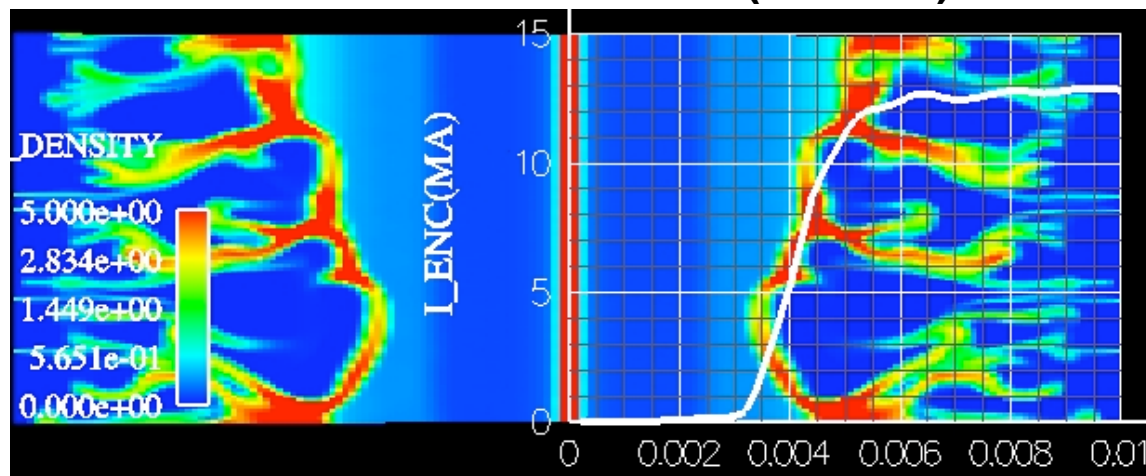
Instability bubble growth is reduced by 3D effects

C=3% (2520 ns)



- In 3% correlated problem, bubble growth is reduced because current can flow azimuthally.

2D Case: C=100% (2520 ns)



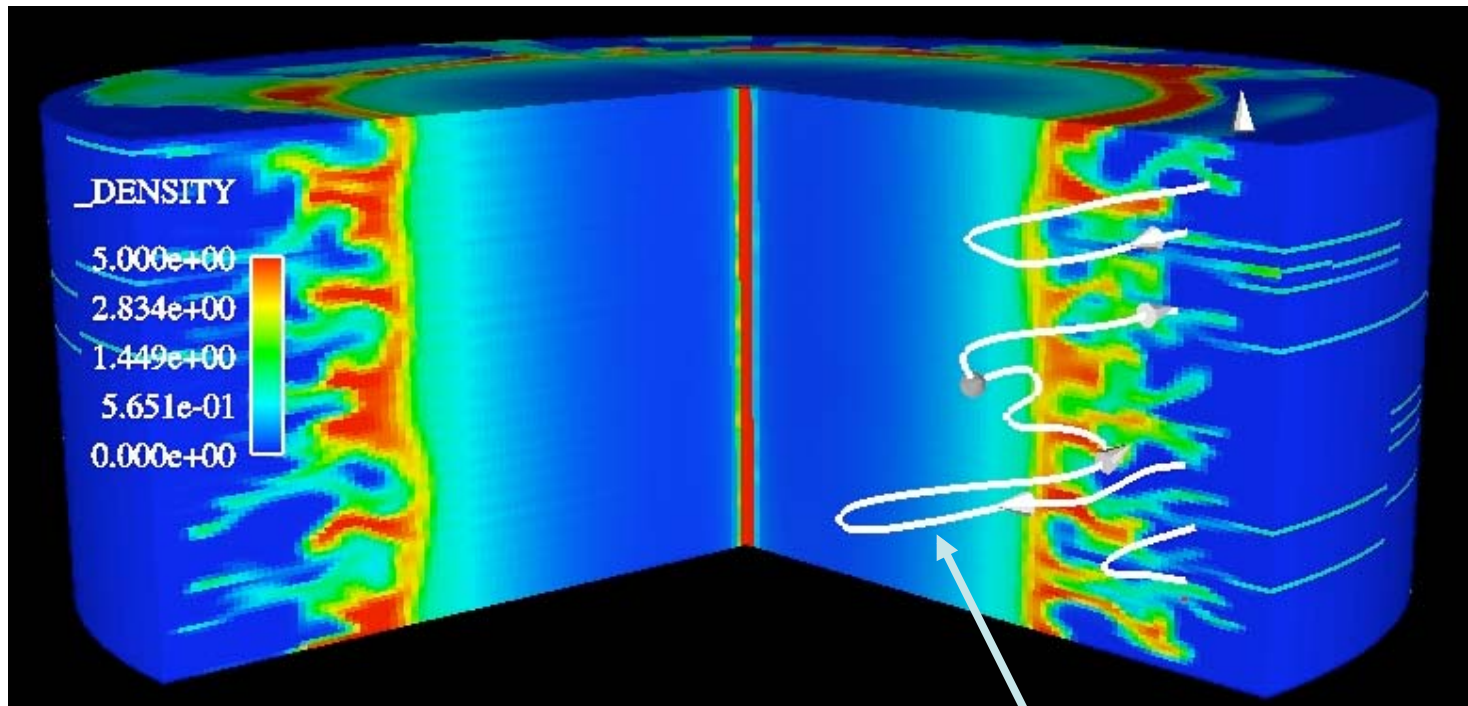
- Mass distribution is left to right asymmetric as observed in radiographs

Yu, et al., Phys. Plasmas (2008)



In 3D, current can travel azimuthally, “self-regulating” bubble growth

$C=3\%$, $t=2518$ ns

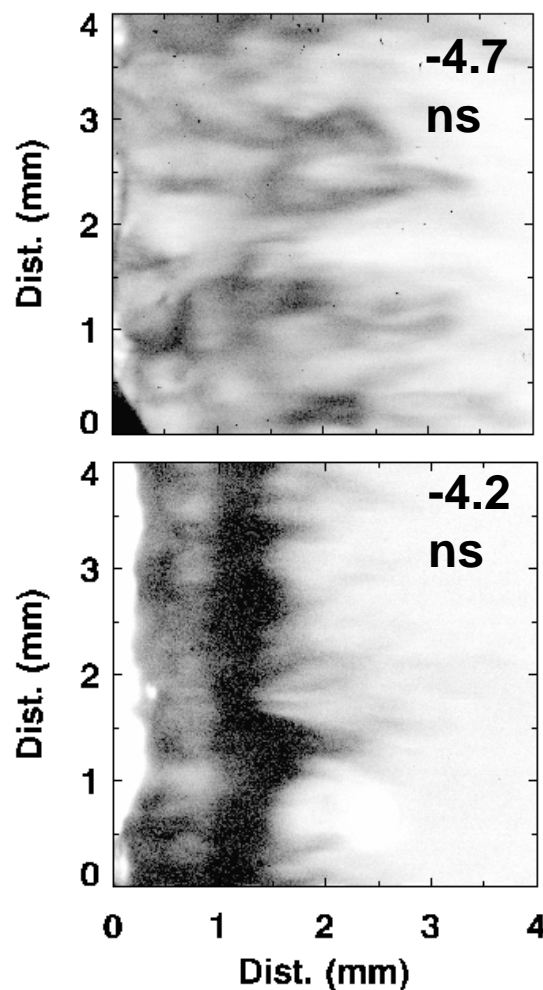


current streamline

Yu, et al., Phys. Plasmas (2008)



We can stabilize the magneto-Rayleigh-Taylor instability and decrease the width of the mass by 3X

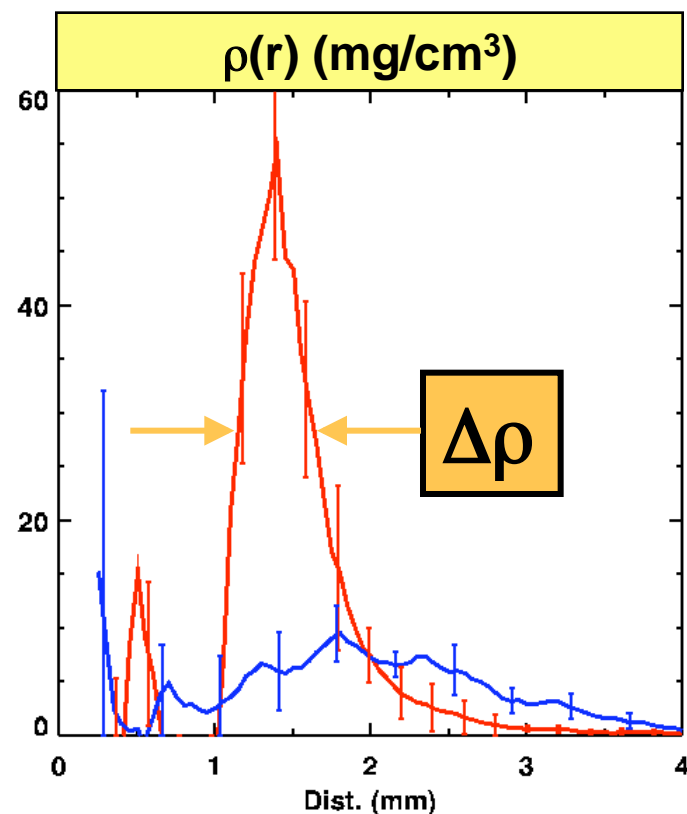


Single

$\Delta\rho=1.58$ mm
 $P=160$ TW
 $\tau_w=3.9$ ns

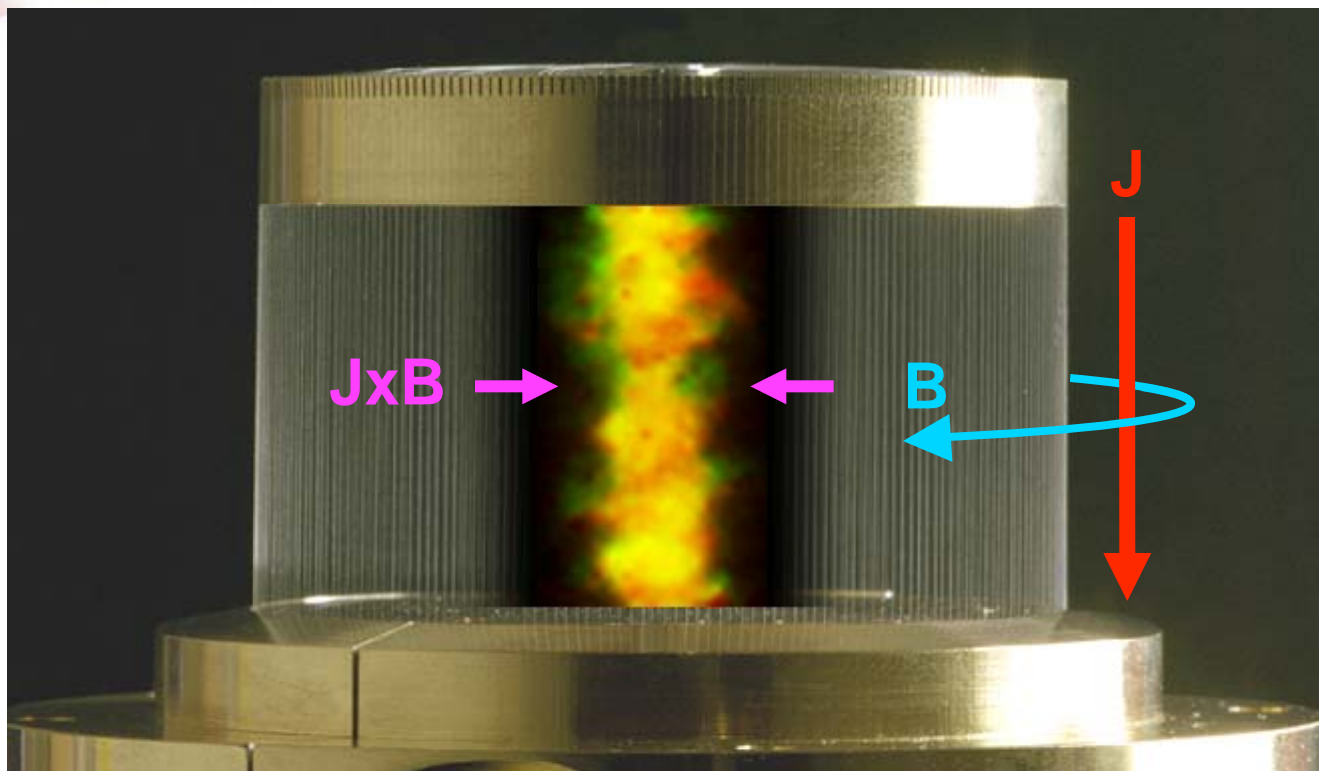
Nested+foam

$\Delta\rho=0.51$ mm
 $P=260$ TW
 $\tau_w=1.7$ ns





J x B force pinches wire array into a dense, radiating plasma column



z-pinch sources

$$n_i \sim 1 - 5 \times 10^{20} \text{ cm}^{-3}, Y_{\text{rad}} \sim 1 - 2 \text{ MJ}$$

$$P_{\text{rad}} \sim 100\text{-}250 \text{ TW} \sim 200 \text{ million million Watts}$$

$$T_{\text{rad}} \sim 200 \text{ eV} \sim 2,300,000 \text{ }^\circ\text{K}$$

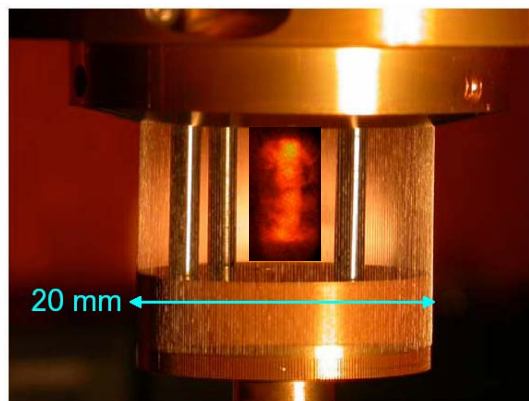


A power of 200 TW

- 15 x global power consumption
- 63 x US annual power consumption
- <http://www.wolframalpha.com/>



How can we use this efficient x-ray source to do ICF?

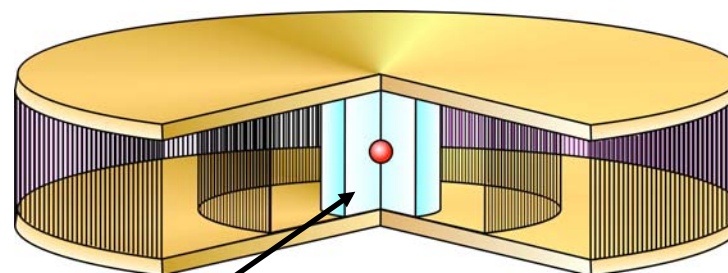


Where do we put the capsule?

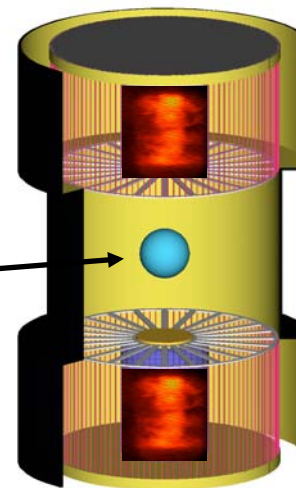
We want high intensity (high T_r)
for high ablation pressure
→ let the capsule see the pinch

We need high uniformity ($\sim 1\%$) in
x-rays the capsule sees for symmetry
→ hide the capsule from the pinch

Dynamic Hohlraum



Double-Ended Hohlraum



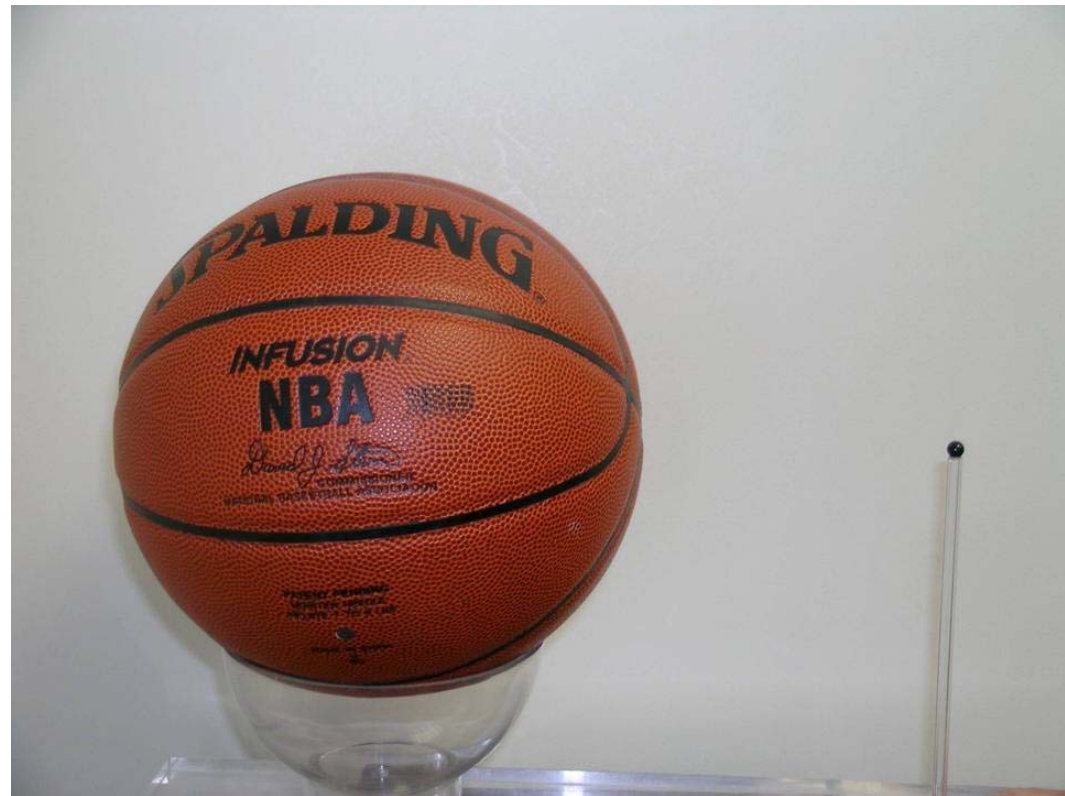
This approach is
the most
conservative



Spherical convergence is used to reach the high densities and pressures

30 to 1 convergence ratio is required

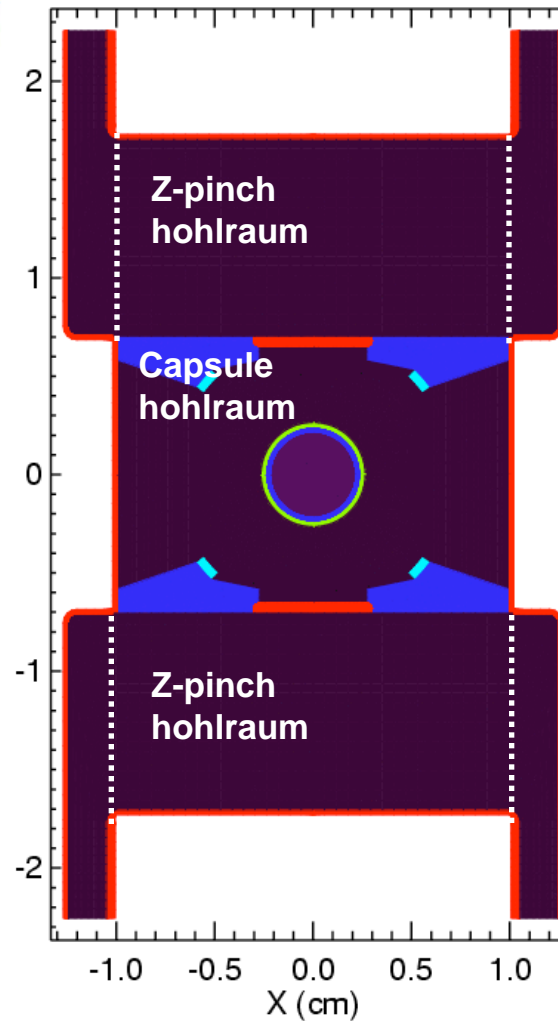
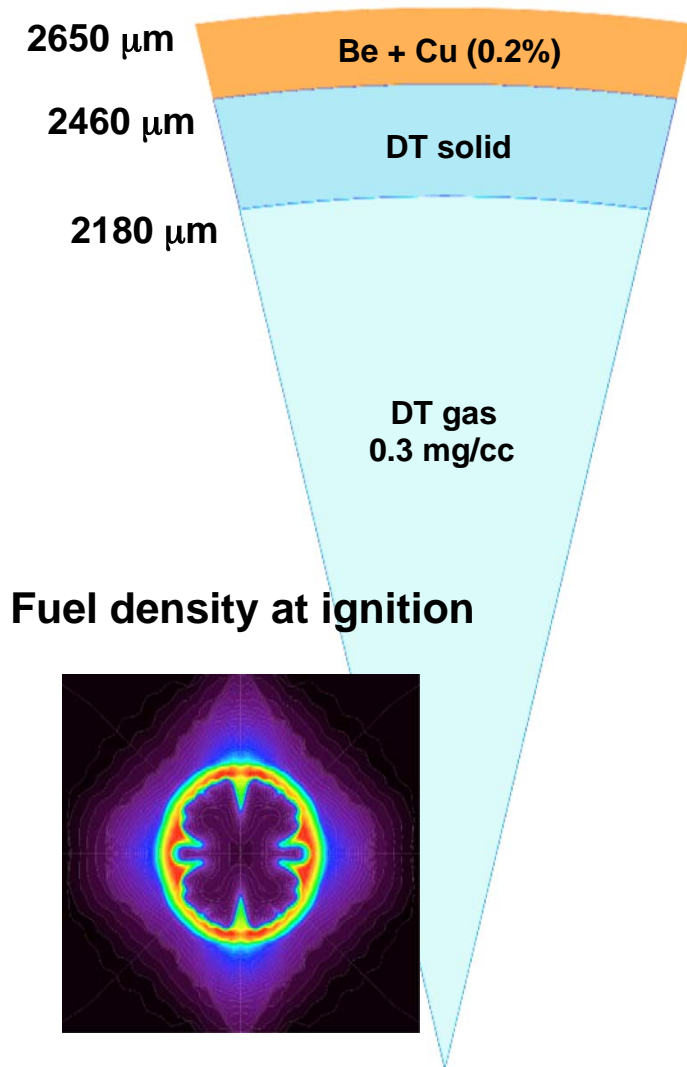
- One requirement is 1-2% deviation from perfect sphere
- This means excellent pressure symmetry
- In turn this means excellent radiation symmetry



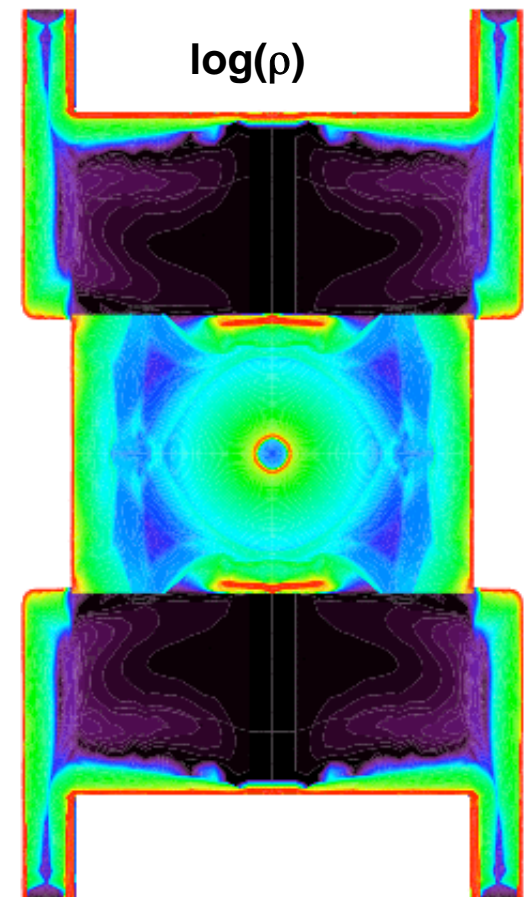


Integrated simulations demonstrate 400+ MJ fusion yield in a pulsed-power z-pinch driven hohlraum

R. A. Vesey et al., Phys. Plasmas (2007)



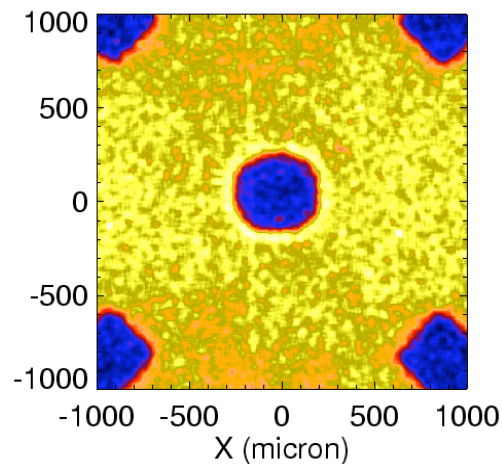
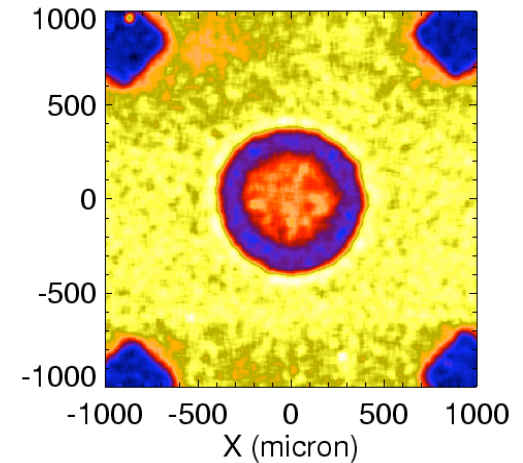
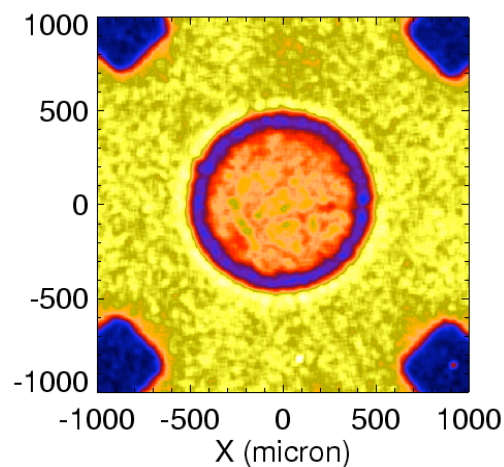
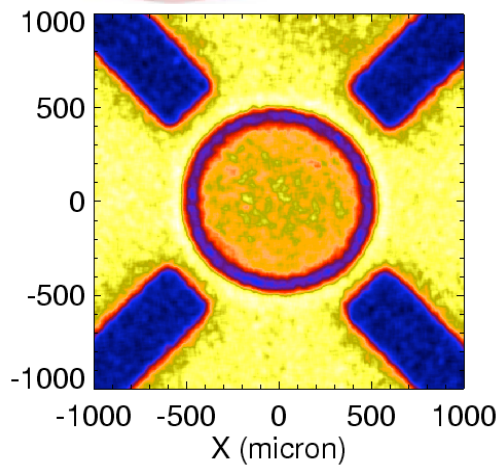
470 MJ 2D yield



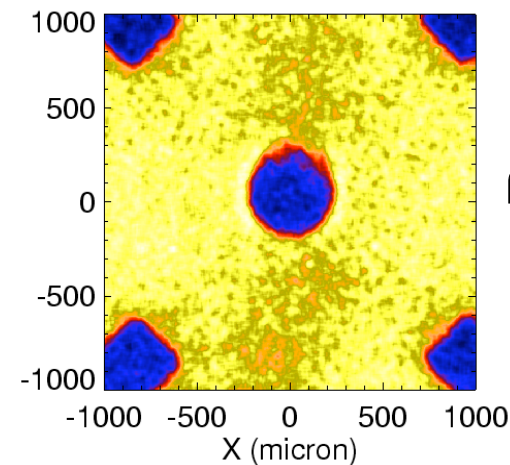
37 papers in 7 years



The double pinch has produced some of the most symmetric capsule implosions, for any driver



**near peak CR
absolute transmission
~ 0 near center**



$\rho > 40 \text{ g/cc}$

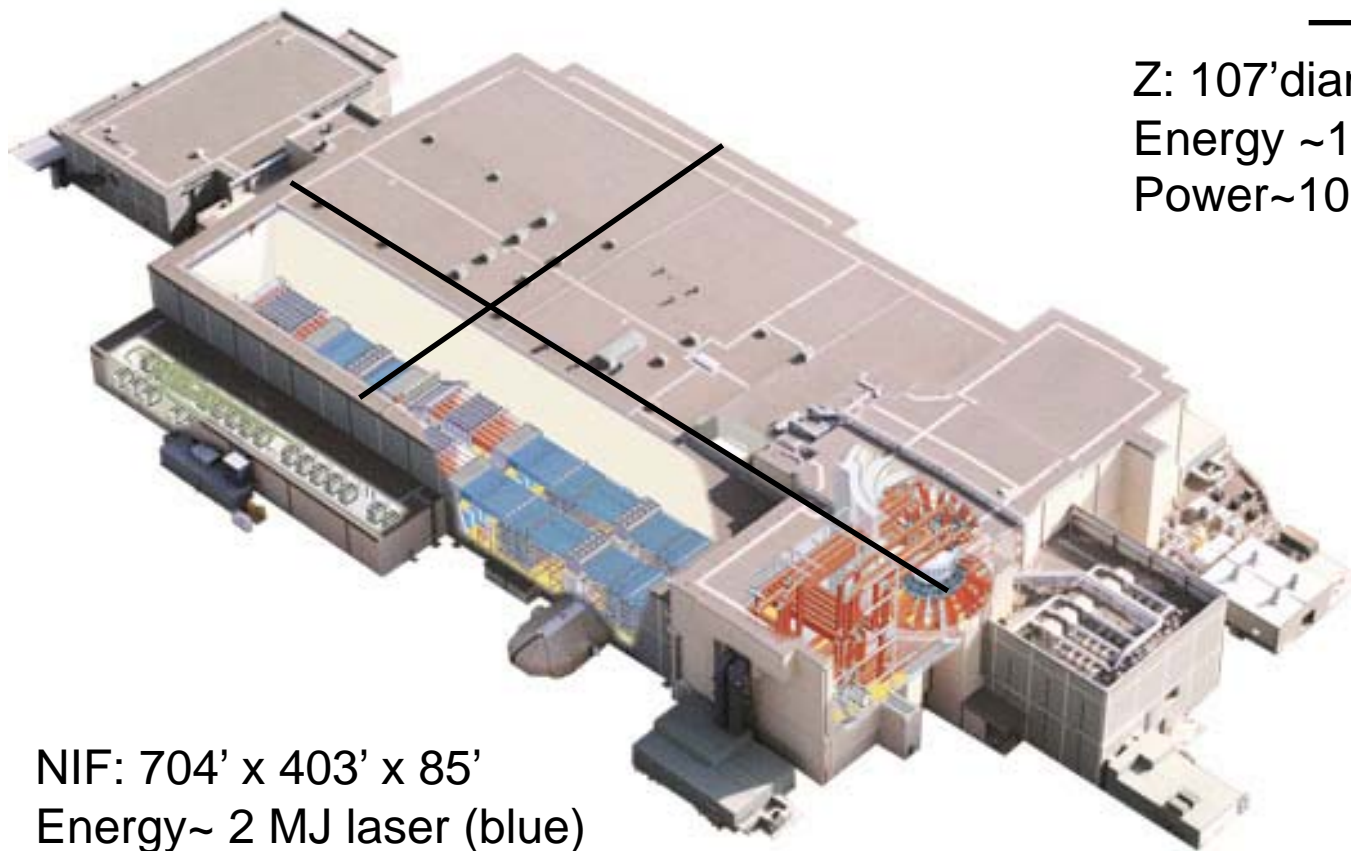
- Record convergence ratios for pulsed power systems



The cost effectiveness and efficiency of pulsed power is evident from a size comparison of Z and NIF



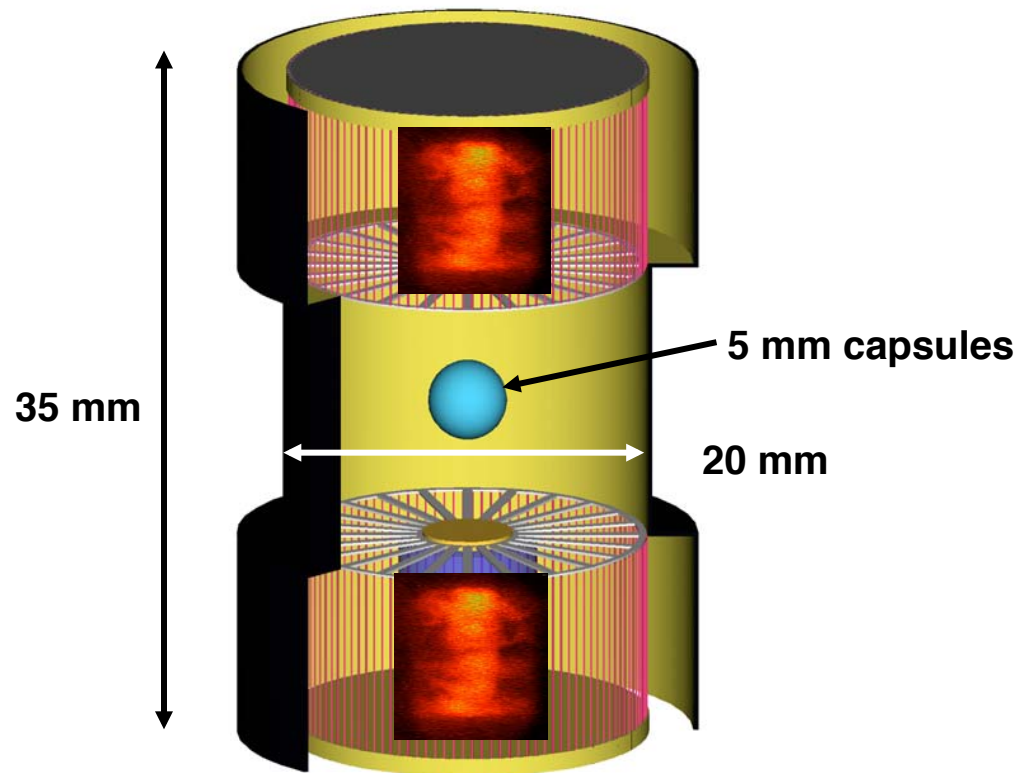
Z: 107'diam x 20' high
Energy ~1 MJ "useable" x-rays
Power ~100-200 TW



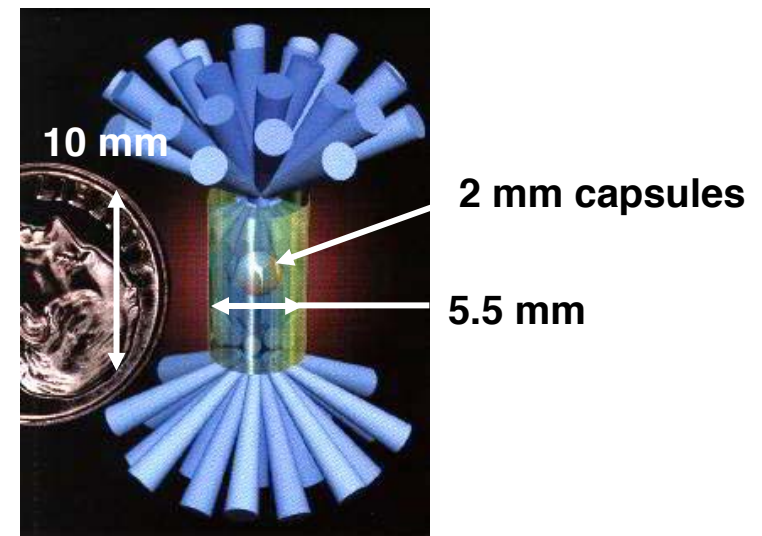
NIF: 704' x 403' x 85'
Energy ~ 2 MJ laser (blue)
Power ~ 500 TW

HOWEVER:

While Z pinches are more efficient radiators, they need more energy to reach ICF conditions because they radiate in bigger volumes

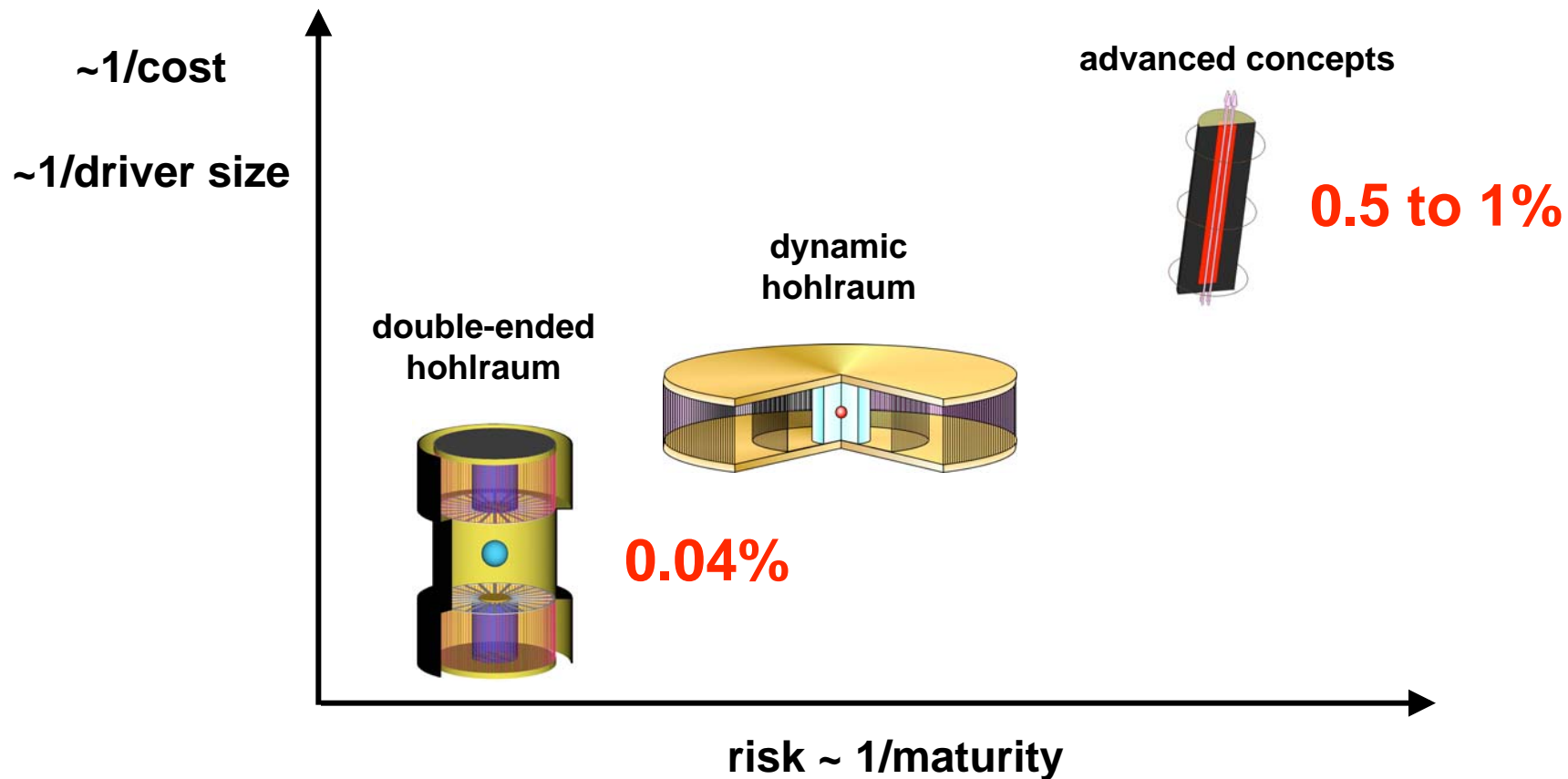


Z-pinch driven hohlraum (2 Z-pinches)
Z: 1-2 MJ X-ray source
High Yield requirement ~ 16 MJ x-ray source



NIF Laser (192 laser beams)
1-2 MJ X-ray source

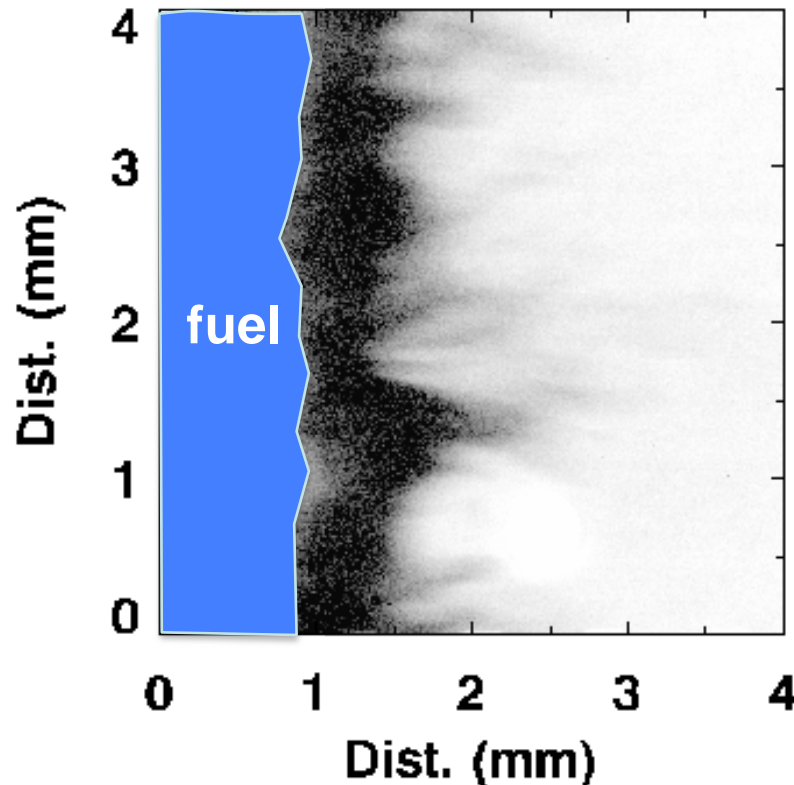
Are there more efficient pulsed power methods for heating and compressing fusion fuel?



- Pulsed power can flexibly drive many targets types
- Direct fuel compression and heating with the magnetic field could be up to 20X more efficient



Direct magnetic compression of fuel could increase coupling efficiency by 20X!



Use stabilized cylindrical implosion to compress fuel

CR ~ 10:1

Magnetic pressure at $I=16$ MA
and $r=1.5$ mm

$P \sim 18$ Mbar

Ultimate convergence ratio of
this pinch ~20:1 with

$P \sim 160$ Mbars

At 60 MA, $P \sim 2.25$ GBars



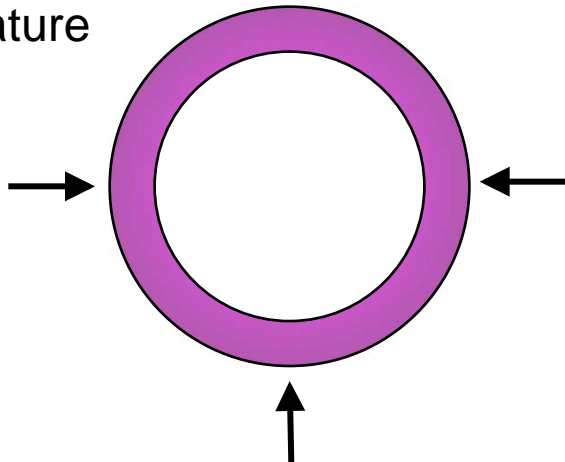
Spherical and cylindrical implosions are used to create high energy density conditions

Radiation-Driven Spherical Implosion
(spherical rocket)

$$P = \frac{(2/5)(1 - \alpha)\sigma T^4}{C_s} = 3T^{3.5} \text{ MBar}$$

140 MBar \approx 2e9 psi
at $T = 300 \text{ eV}$

X-ray drive
temperature
 T



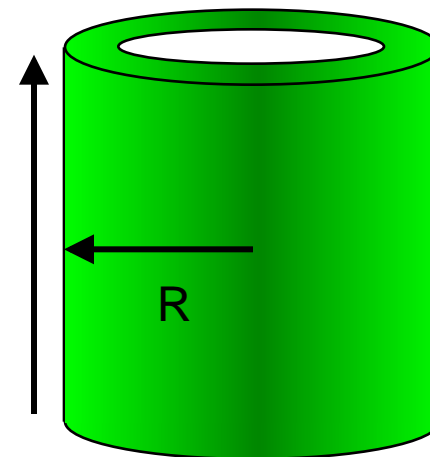
Need fuel pressures of 500 Gbar and fuel ρr of 1 g/cm² for ignition

Magnetically-Driven Cylindrical Implosion

$$P = \frac{B^2}{2\mu_0} = 1.57e - 3 \left(\frac{I_{MA}}{R_{cm}} \right)^2 \text{ MBar}$$

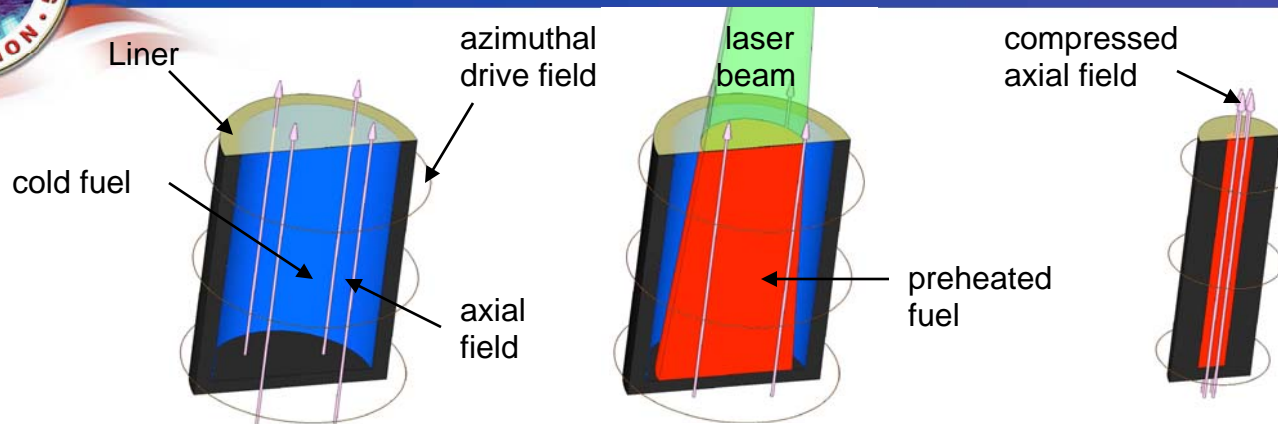
97 MBar \approx 1e9 psi
at $I = 25\text{MA}$, $R = 0.1 \text{ cm}$

Drive
current
 I





Fuel magnetization and preheat provide beneficial effects for liner driven ICF



A magnetic field inhibits thermal conduction and enhances alpha particle deposition within the fuel

Preheating the fuel reduces the compression needed to obtain ignition temperature

...calculations indicate this could be done by the Z Beamlet laser

Simulations indicate significant yields on ZR with modest convergence ratios and implosion velocities

The Magneto-Rayleigh-Taylor instability will be a key issue for this concept

S. Slutz, M. Herrmann



There are many applications of pulsed power technology

- Pulsed electric fields
 - Electroporation
 - Bacterial decontamination
 - Discharges through solids and liquids
- Pulsed magnetic fields
 - Equation of state measurements
 - High energy density physics
- High power beams
 - Electron beams
 - Ion beams
- Intense radiation sources
 - Laser flashlamps
 - Microwave generation
 - Z-pinch soft X-ray sources (< 5 keV)
 - Z-pinch warm x-ray sources (5-10 keV)
 - Hard X-ray sources (>100 keV)

“Pulsed Power Systems” by
H. Bluhm, Springer, (2006)



Summary

- Pulsed power can flexibly drive many different kinds of experiments at high currents and high voltages
- Large magnetic fields create large pressures. Large currents create large fields
- Large pressures are needed to access high energy density regimes
- The Z machine creates large currents and is the world's largest x-ray source
- Lasers have more control than Z-pinches regarding where/how/how long energy is deposited. However Z-pinches are cheaper than lasers.
- The upper limits on Z-pinch performance in achieving high energy densities are not known. There is a lot of room for innovation! (Direct drive)



Some References

- ICF and HEDP
 - “The Physics of Inertial Fusion” by S. Atzeni and J. Meyer-ter-Vehn, Clarendon Press, Oxford (2004).
 - “High-Energy-Density Physics” by R.P. Drake, Springer, (2006).
 - “Inertial Confinement Fusion” by J.D. Lindl, Springer Verlag, NY (1998) ; Phys. Plasmas 2, 3933 (1995)
 - M. D. Rosen, Phys. Plasmas 6, 1690 (1999)
 - “Physics of Shock Waves & High Temperature Hydrodynamic Phenomena” by Ya. B. Zeldovich & Yu. P. Raizer, (edited by W. D. Hayes & R. F. Probstein) Academic Press, NY (1966)
 - HEDP Summer School, August, 2005 and August 2007 (see <http://meetings.ile.rochester.edu/FSCschool>)
- Pulsed power and z-pinch
 - “Pulsed Power Systems” by H. Bluhm, Springer, (2006)
 - “Physics of High-Density Z-pinch Plasmas, Springer, (1999)
 - “The physics of fast z-pinch,” D. Ryutov et al., Rev. Modern Physics, 72, 167 (2000).