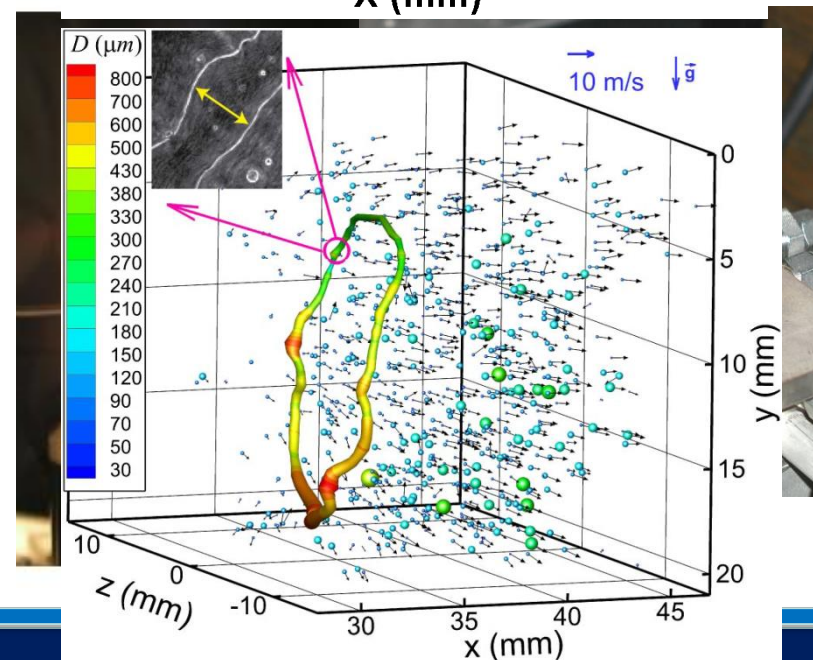
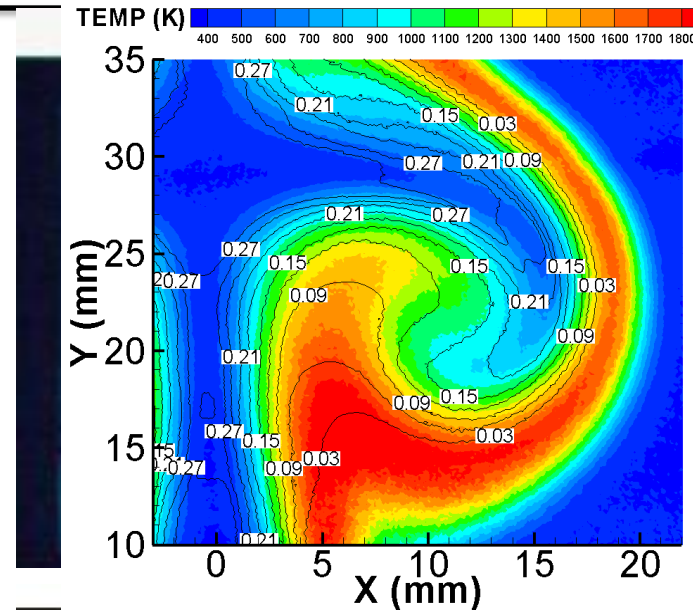


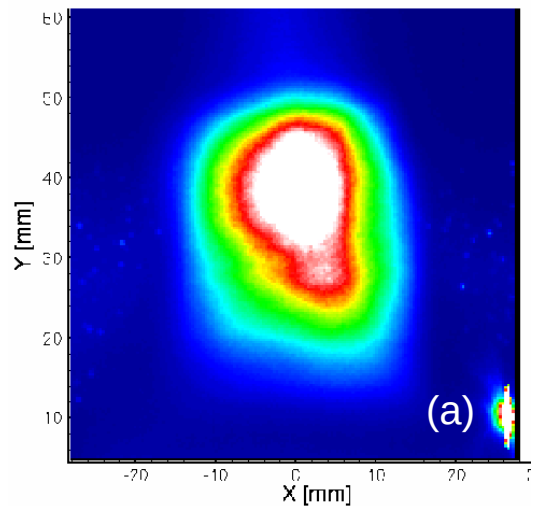
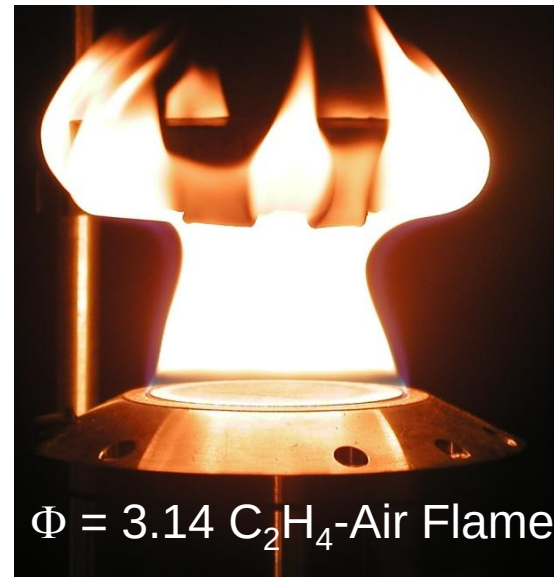
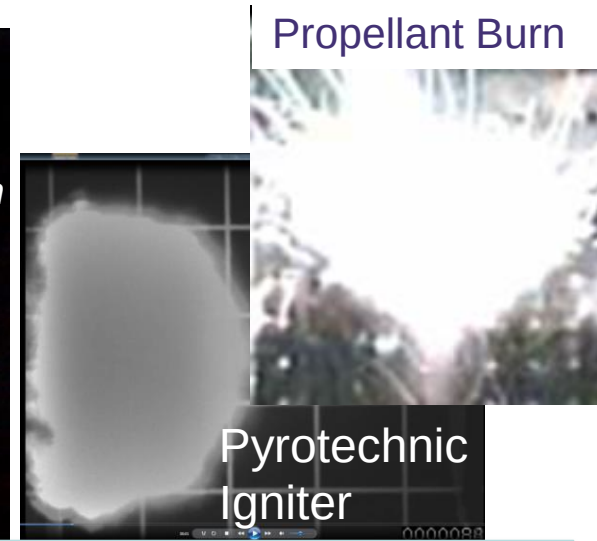
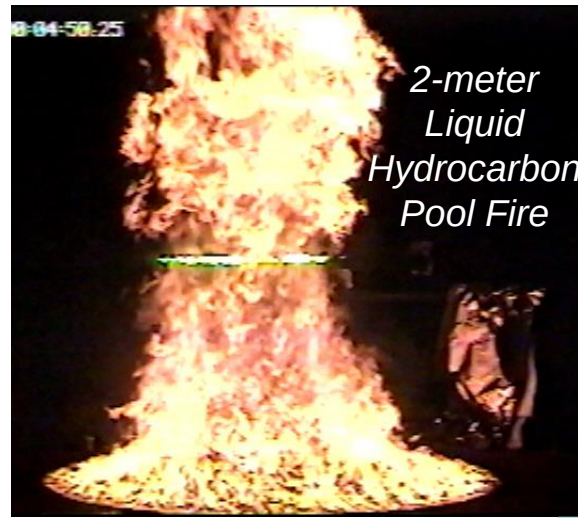
- Background – 5 mins
 - Laser combustion diagnostics
 - Sandia applications in “hostile” environments
- Coherent anti-Stokes Raman scattering (CARS) – 10 mins
 - Fundamentals with nanosecond laser pulses
 - Application to meter-scale fire measurements
- Ultrafast (fs/ps) CARS development – 30 mins
 - Why ultrafast?
 - Time domain Raman processes and proof-of-concept
 - Game-changing advance: SHBC
 - Flame measurements: accuracy and precision
 - Application to composite fire problem
- New ultrafast applications to shock physics (5 mins)

Laser-based diagnostics empower combustion research

- Non-perturbing
- Free of radiation and insertion errors
- 2-D or even 3-D **quantitative** imaging
- Multiple parameters (T/species/soot/velocity...)
- High temporal resolution – 10 ns or better
- High spatial resolution – 10^{-4} – 10^{-5} cm³
- Most effective in clean, laboratory flames

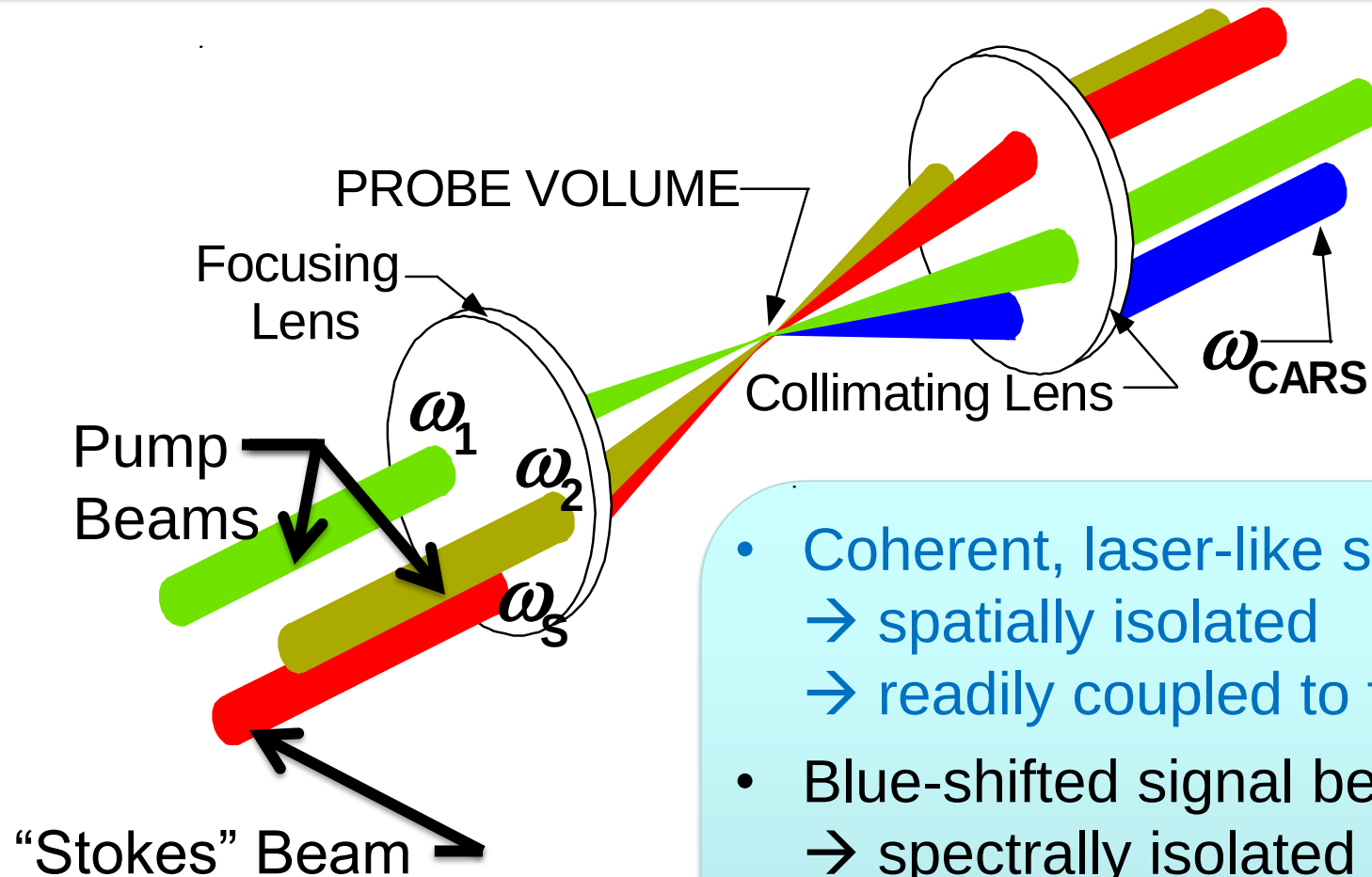


- “Dirty” environments
 - Fire research
 - Energetic materials
- Soot, aluminum particulate
- Luminosity
- Scattering
- Absorption/optical thickness
- **Large-scale** of combustion systems



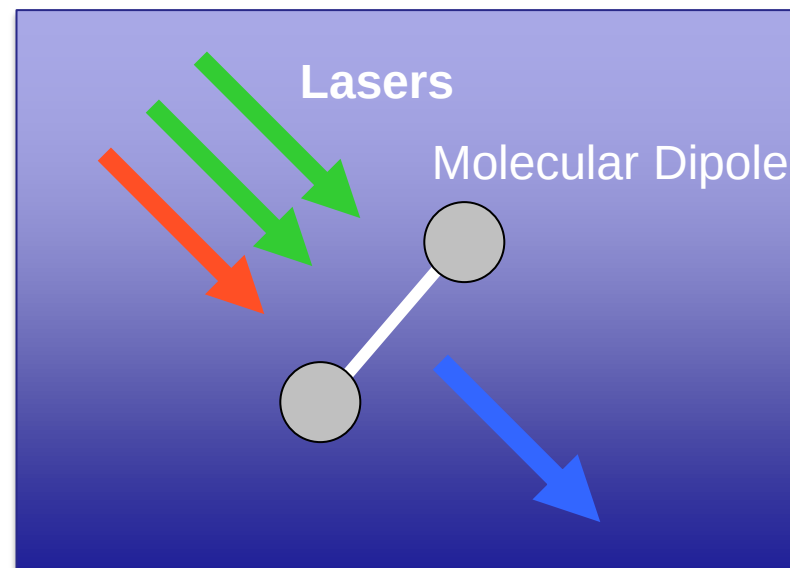
Sooting

Coherent anti-Stokes Raman Scattering (CARS)

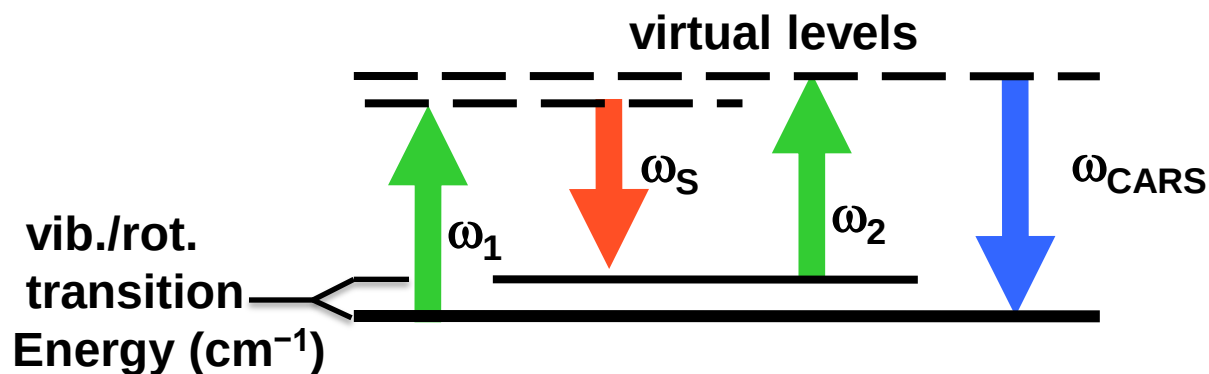


- Coherent, laser-like signal beam
→ spatially isolated
→ readily coupled to fibers
- Blue-shifted signal beam
→ spectrally isolated
- Orders of magnitude stronger than incoherent scattering

- A 'polarization' or induced dipole is prepared by pump and Stokes beams
- This polarization scatters the second pump wave
- Constructive interference in one phase-matched direction only

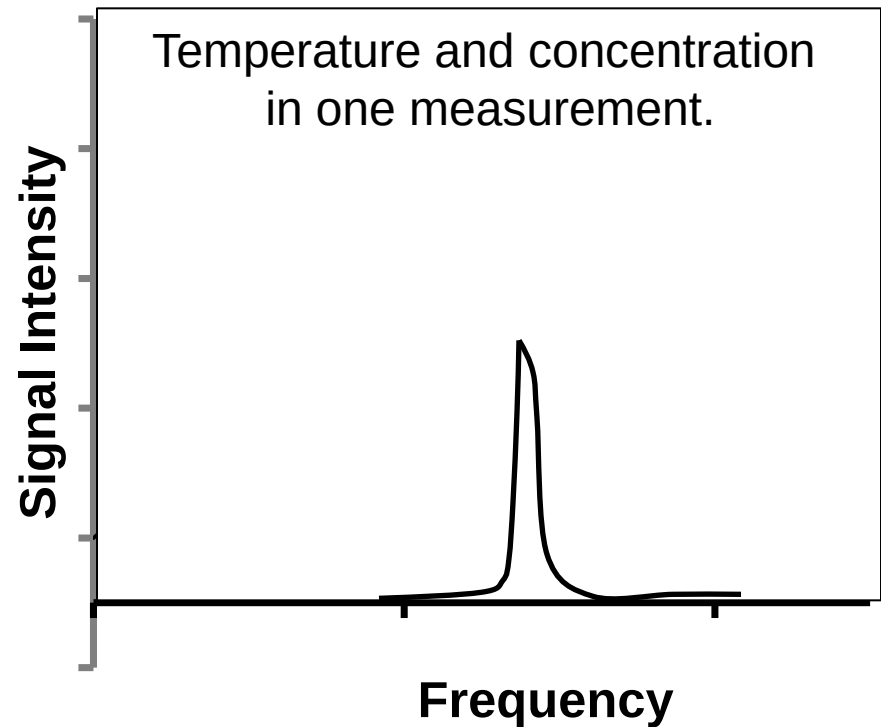
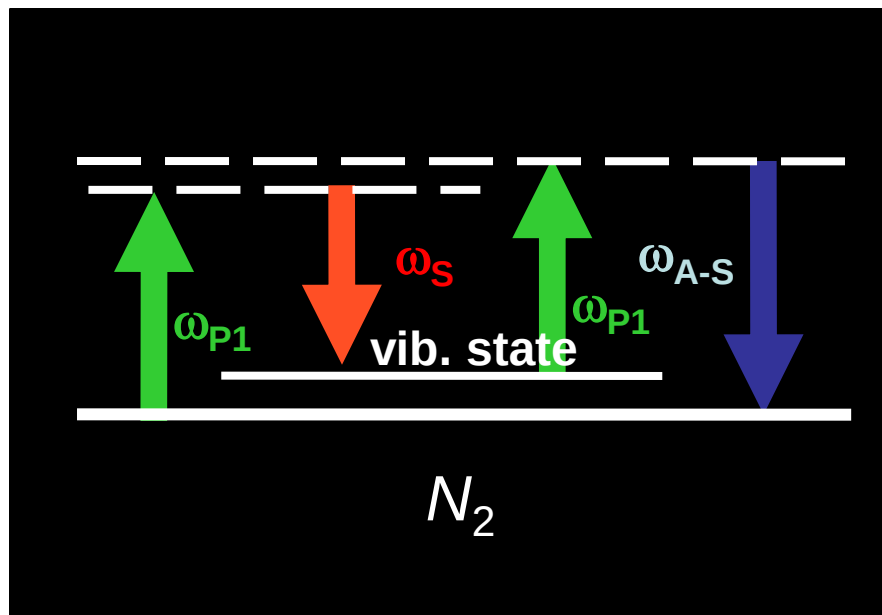


Coherent Anti-Stokes Raman



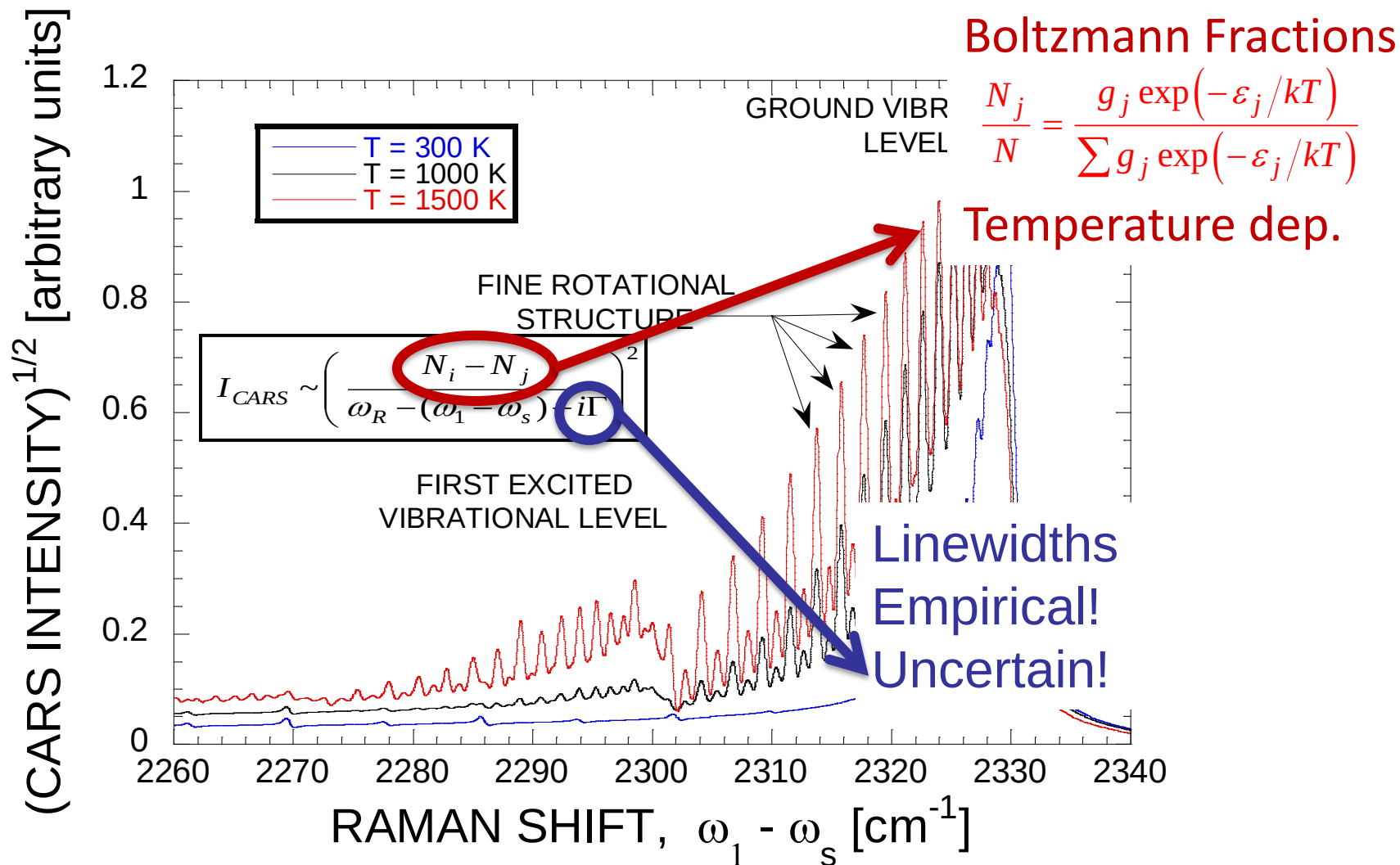
A broadband source permits single-shot detection

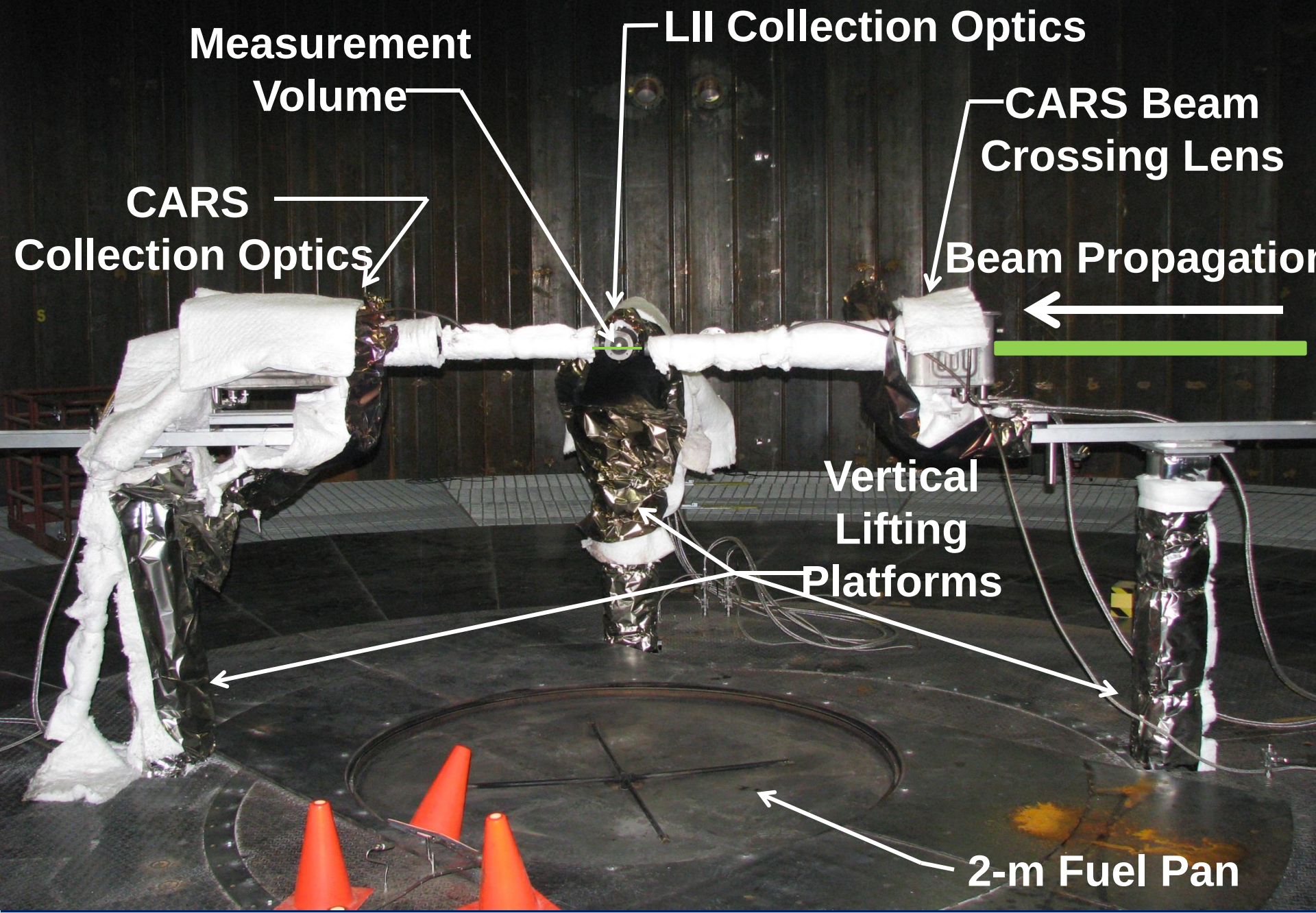
- If all lasers are narrowband one energy level is probed



- If one (or more) laser is broadband then a range of energy levels differences are probed

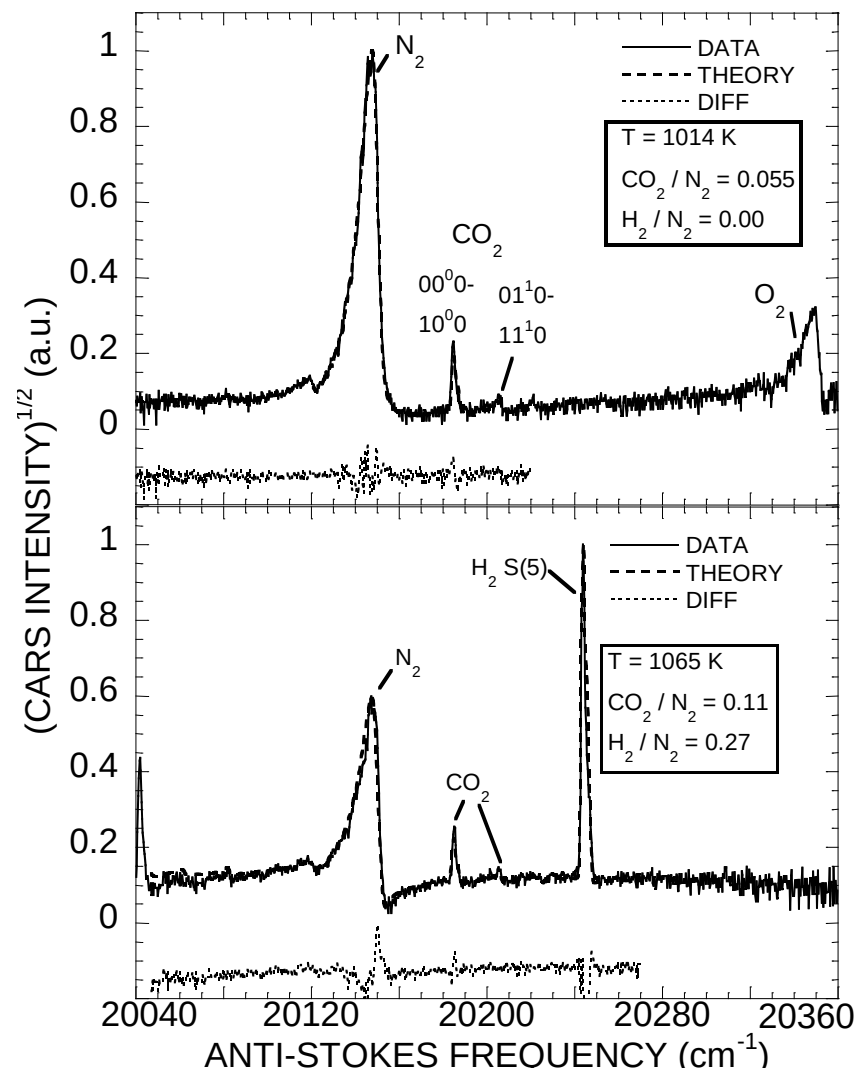
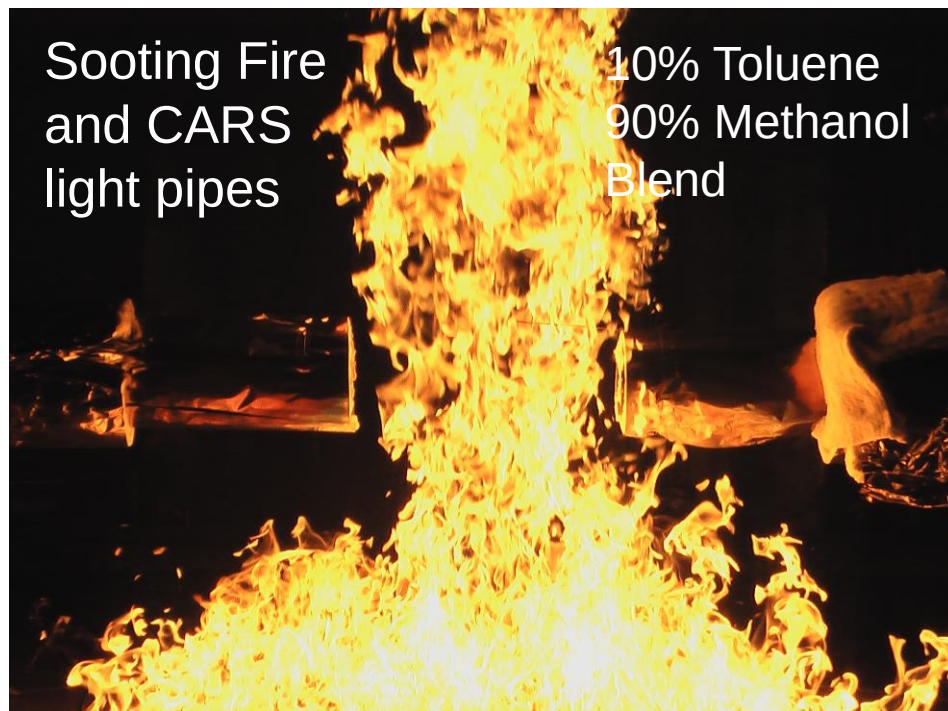
Temperature sensitivity comes from the spectral shape





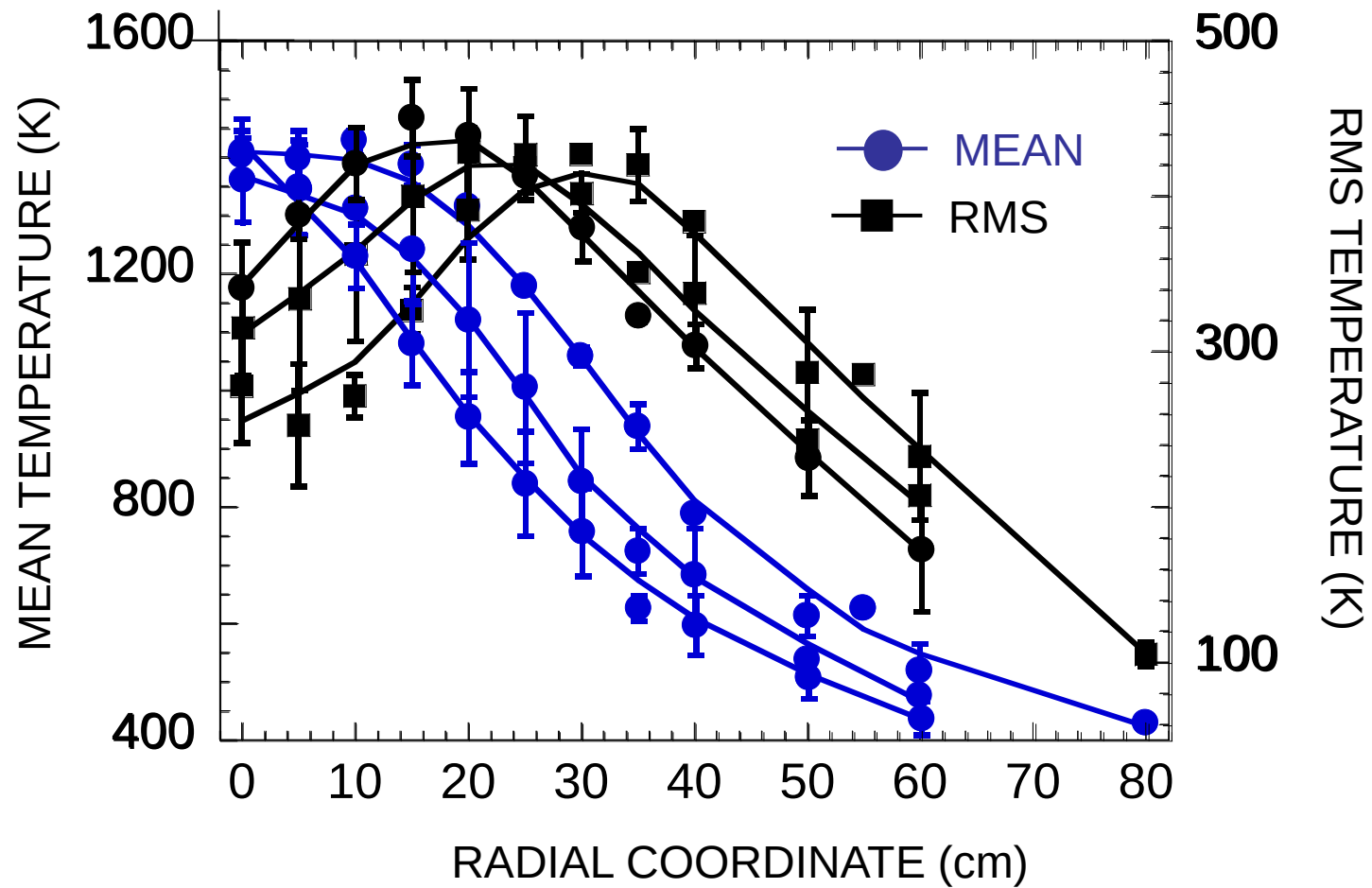
Single-Shot Spectra Provide Simultaneous Temperature/Species Information in Sooting Fire

- CARS spectra from sooting fire show N_2 , CO_2 , H_2 , and O_2
- Full ensemble of species data not available as of yet
- Two representative spectra with theoretical fits (Sandia CARSFT code) shown here



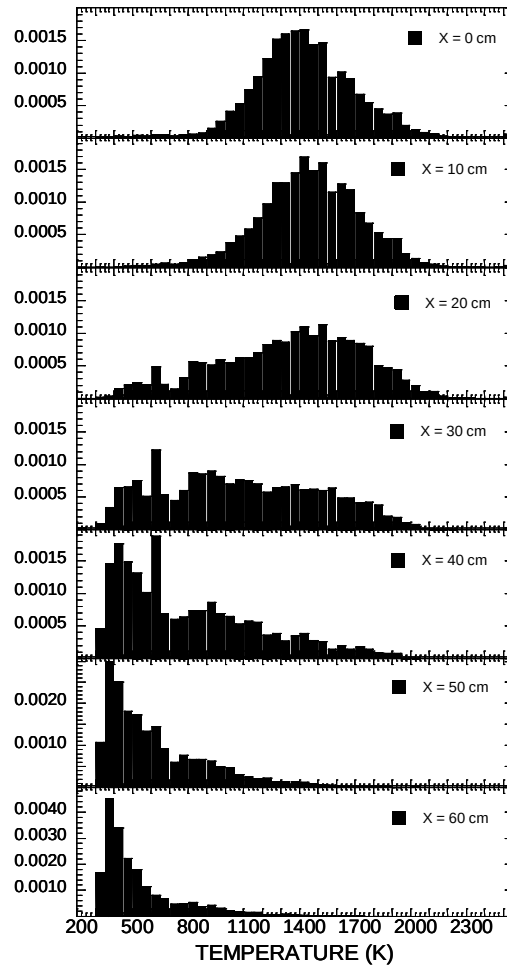
Kearney *et al.*, Proc. Combust.
Inst. **32**, 871-878, 2009.

Radial Temperature Profiles – Sooting Pool Fire

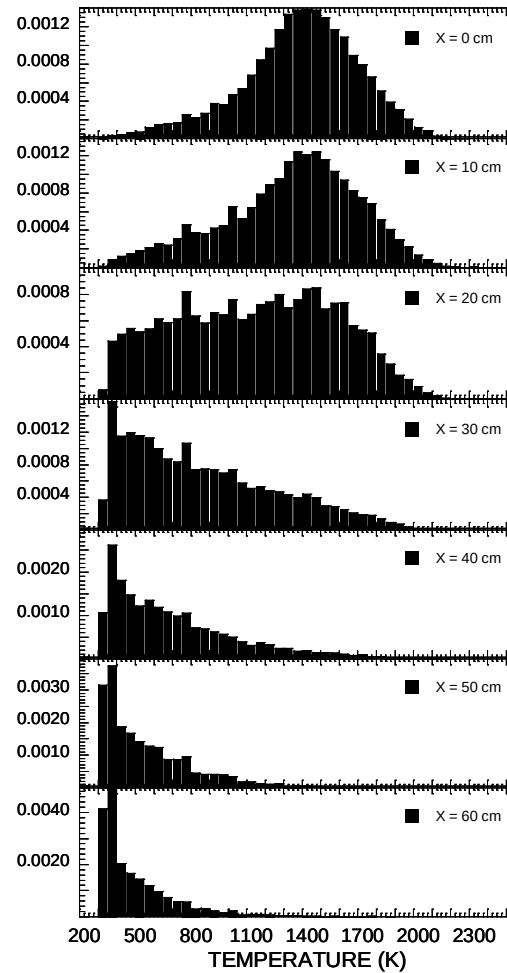


Temperature PDF – Sooting Pool Fire

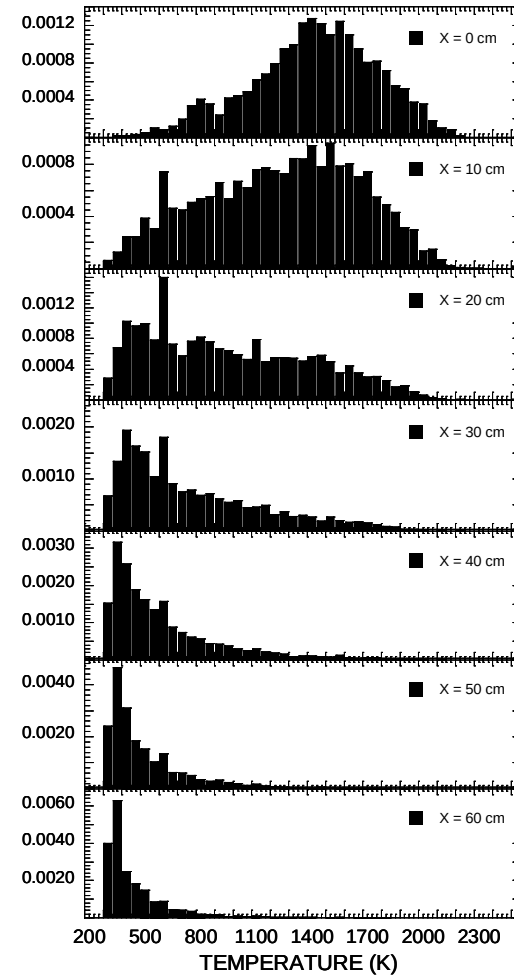
Y = 0.5 m



Y = 1.0 m



Y = 1.5 m



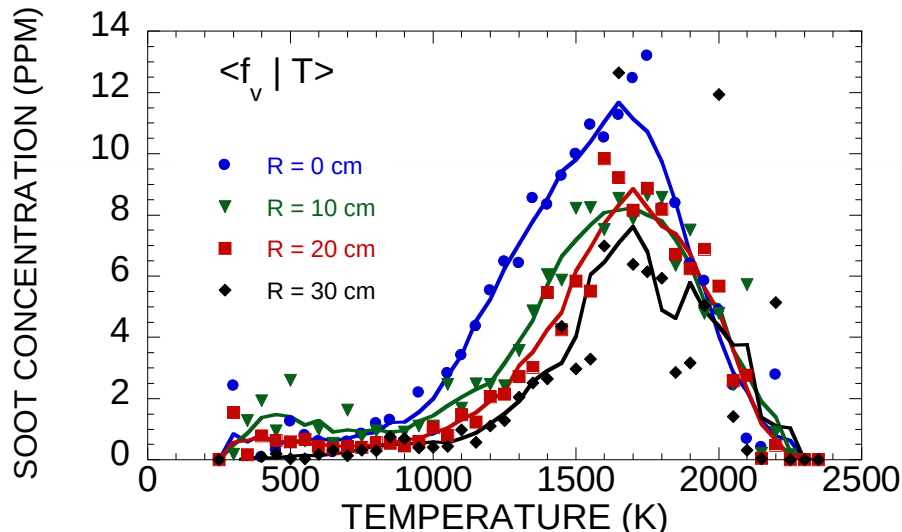
Joint/Temperature Soot Statistics for Emission

Radiative Transfer Equation

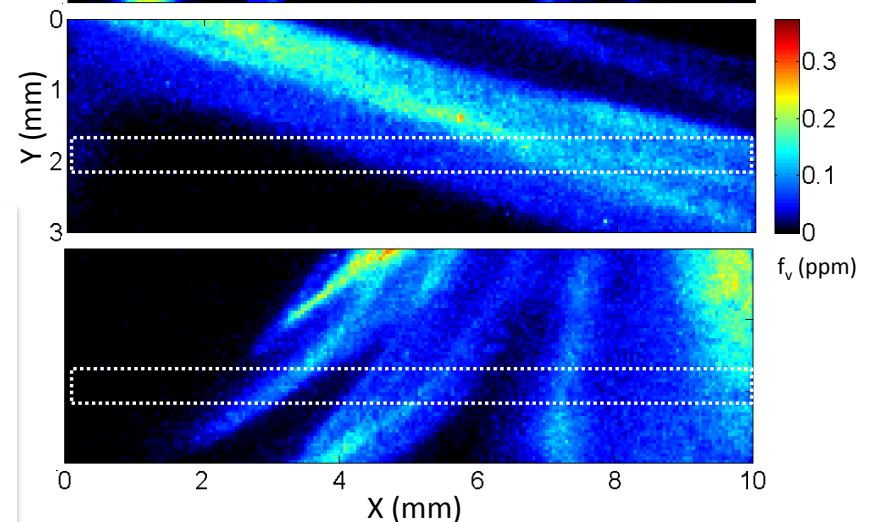
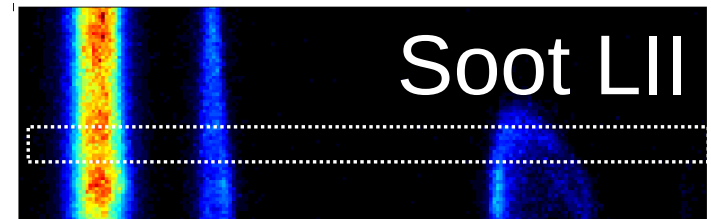
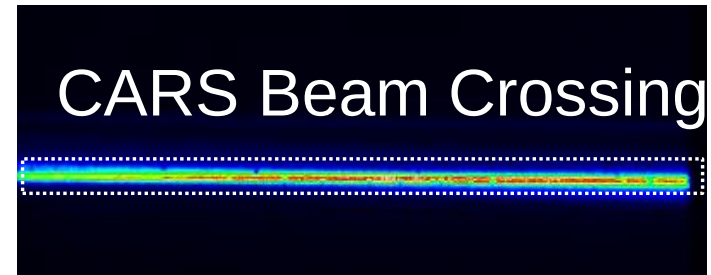
$$\frac{d\bar{I}_\lambda}{ds} = \overline{\mu_\lambda I_{\lambda,b}(T)} - \overline{\mu_\lambda I_\lambda}$$

Joint Temp. Soot
Statistics Desired

- CARS system combined with LII soot detection
- Average soot in 10^{-5} cc CARS volume correlated with enthalpy-pooled temperature



Mean Temperature Conditioned on Soot

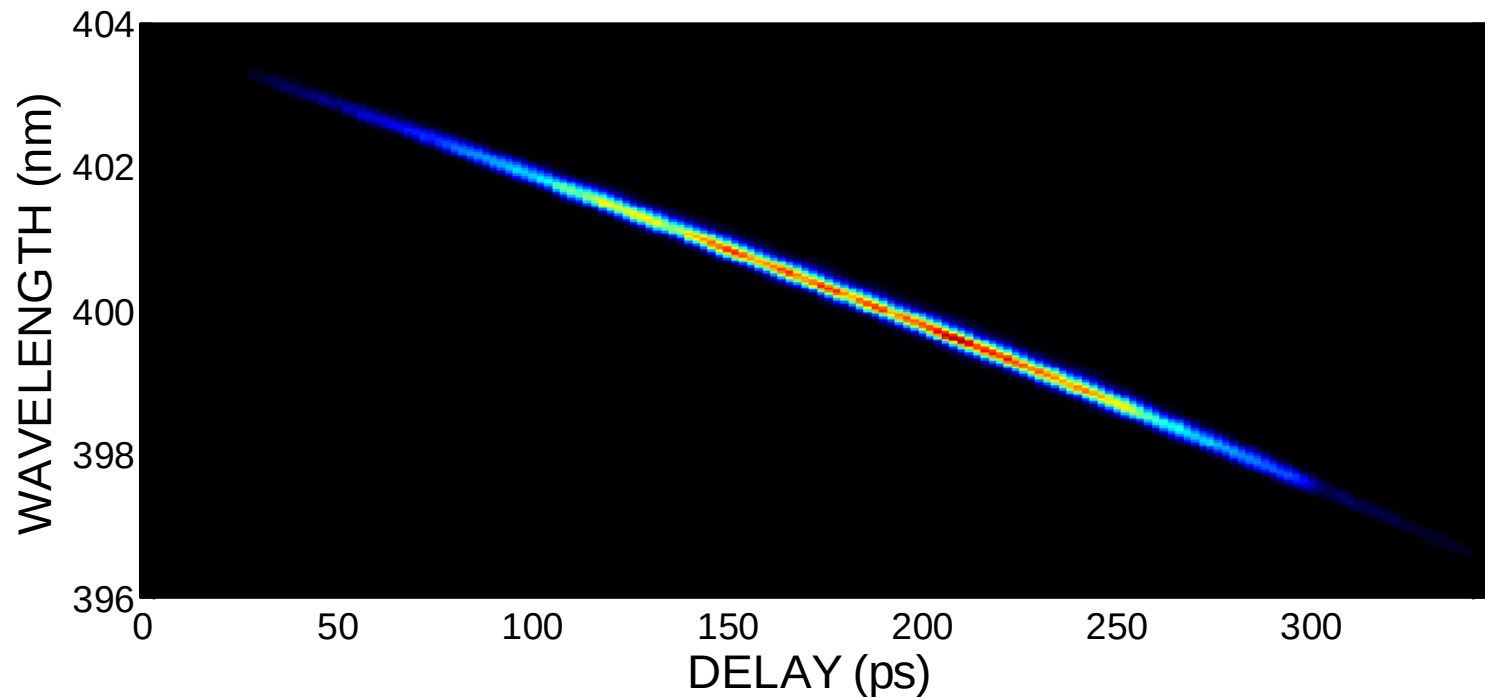


Frederickson and Kearney, Appl. Opt. **50** (2011)
Kearney and Pierce, Combust. Flame **159** (2012)

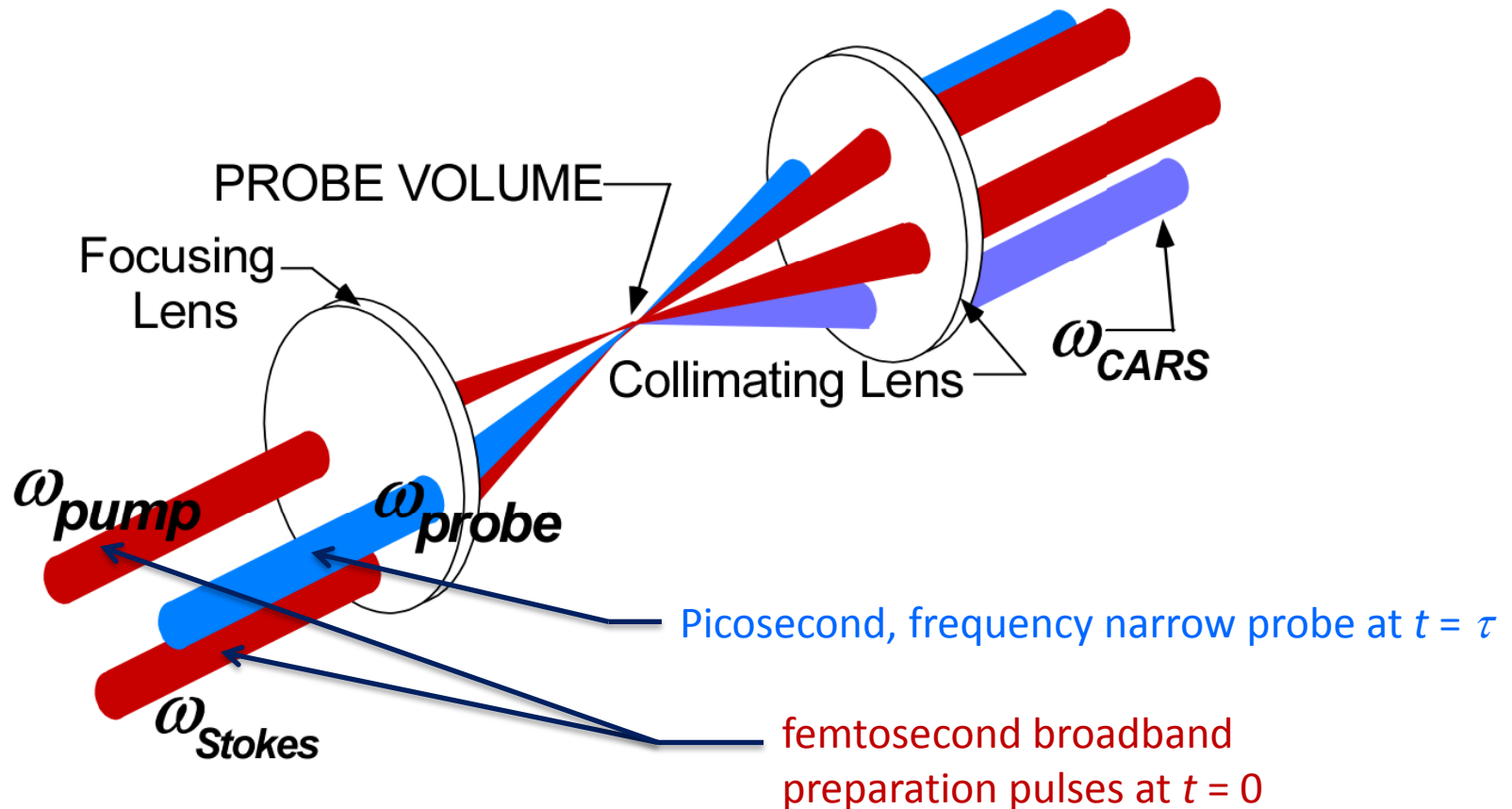
Femtosecond CARS

Why Ultrafast?

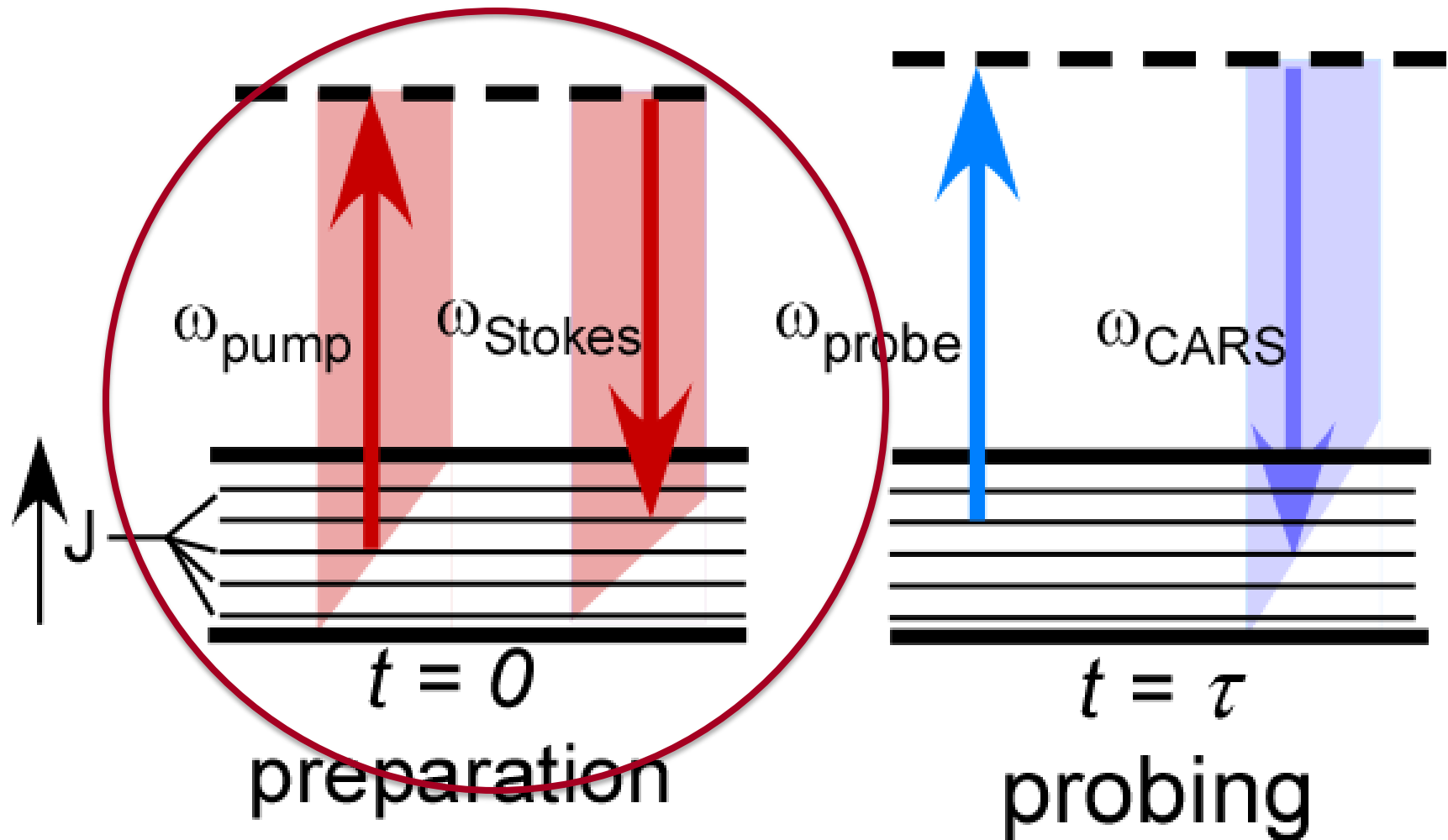
- High-quality (transform-limited) **broadband** sources
$$\Delta\tau \Delta\nu \geq \text{const.}$$
- High repetition rates (kHz vs 10 Hz)
- Transient vs. steady state measurements → no linewidths!



fs/ps Rotational CARS Experimental Arrangement



Time-Domain Rotational Raman: $\tau_{\text{laser}} \ll \tau_{\text{molecule}}$



"The story here is really in the time domain"

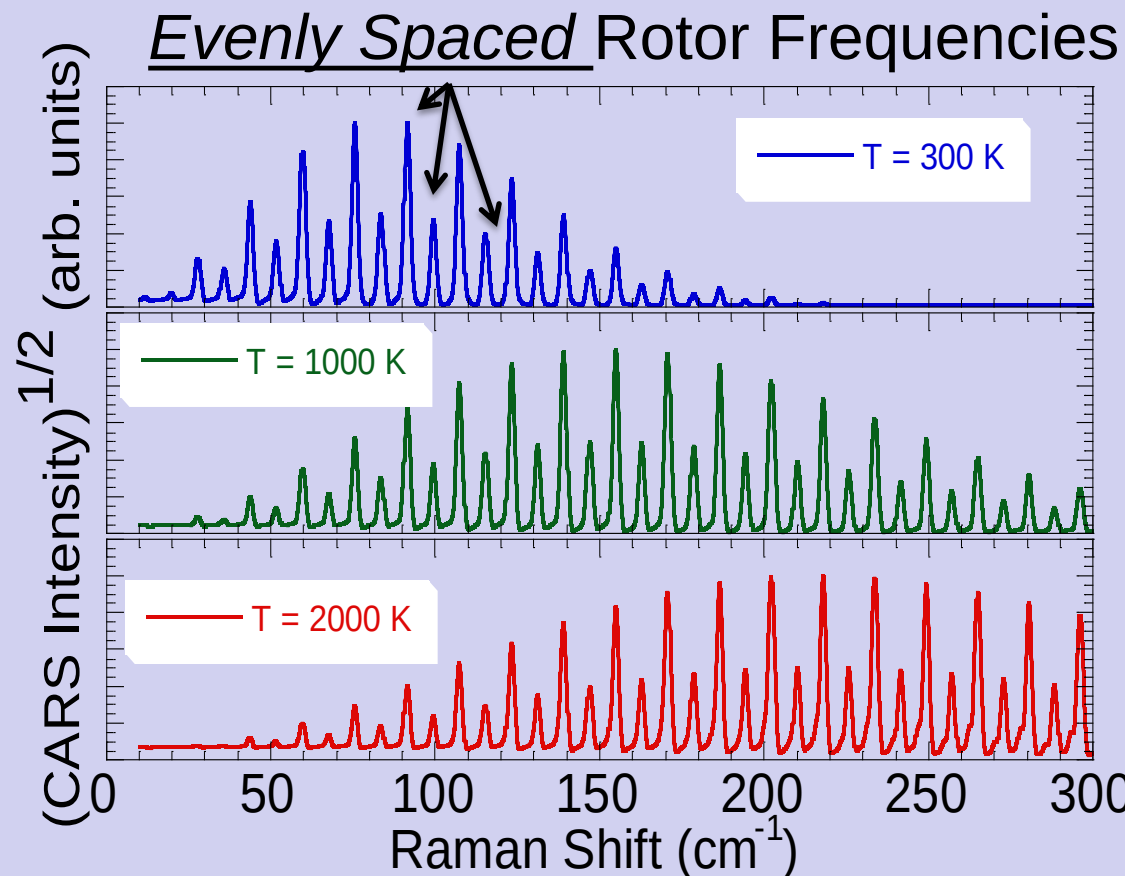
Pump/Stokes Preparation

- Impulsive molecular alignment at time $t = 0$

- Assem

- Rotor

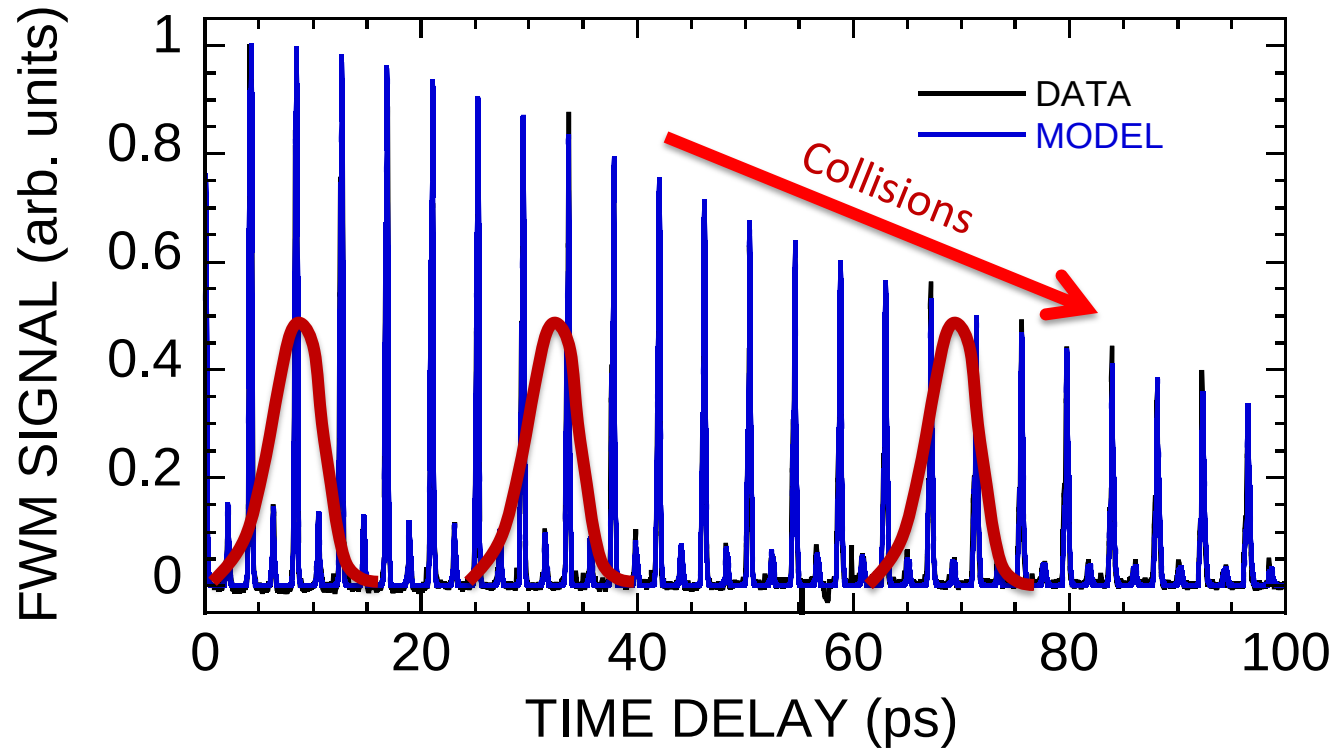
- Rotors



- Assem

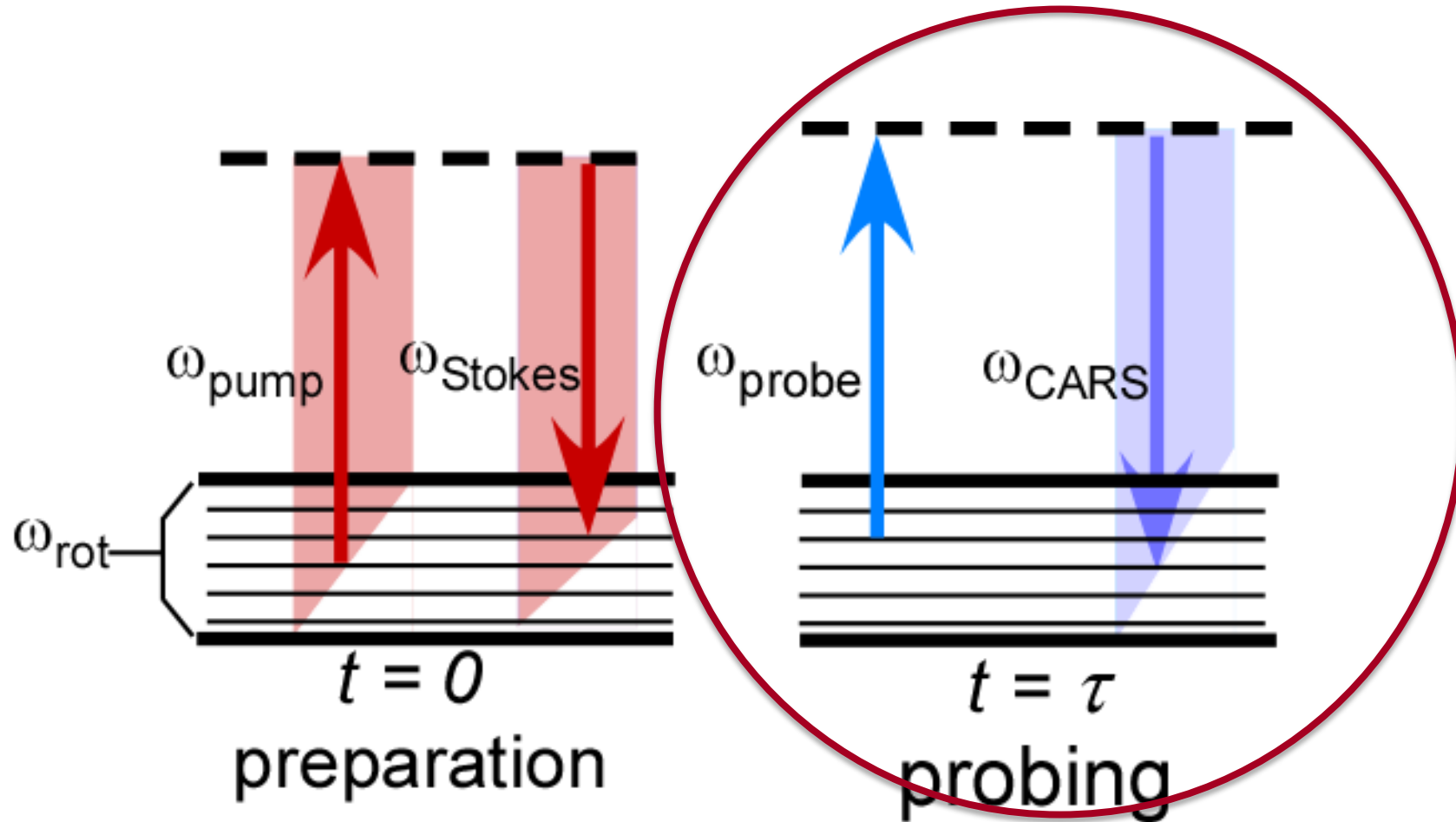
- The result is a periodically recurring and long-lived Raman polarization

Measured Response in N₂ at T = 300 K



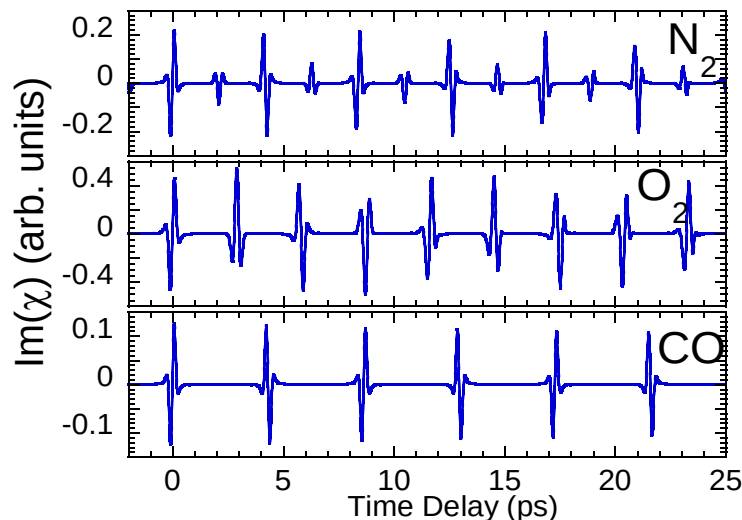
$$\chi = \sum_J W_J \exp(i \omega_J t) \exp(-\Gamma_J t)$$

Probe Step

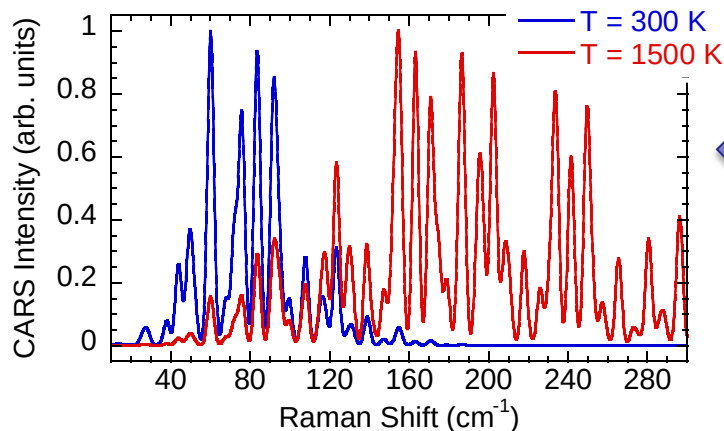


Probe Step and Spectral Synthesis

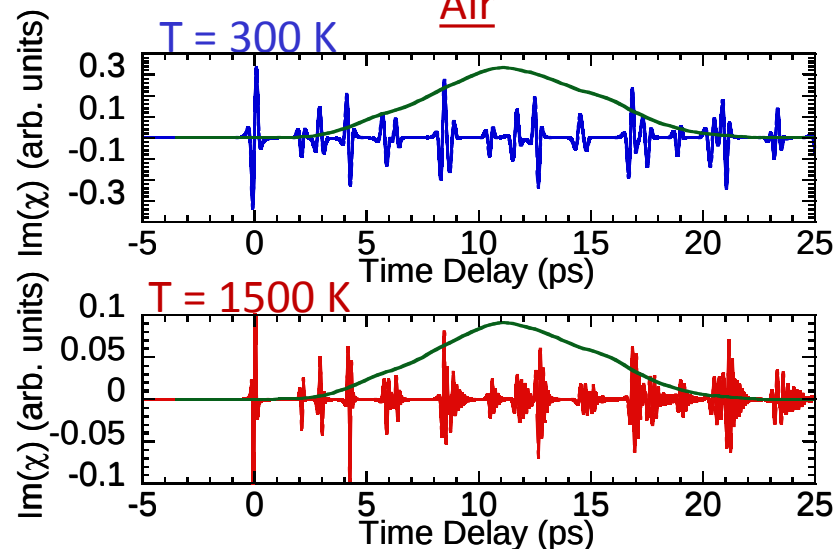
Pure Gases, $T = 300\text{ K}$



$$\chi = \sum_J W_J \exp(i \omega_J t) \exp(-\Gamma_J t)$$



Air

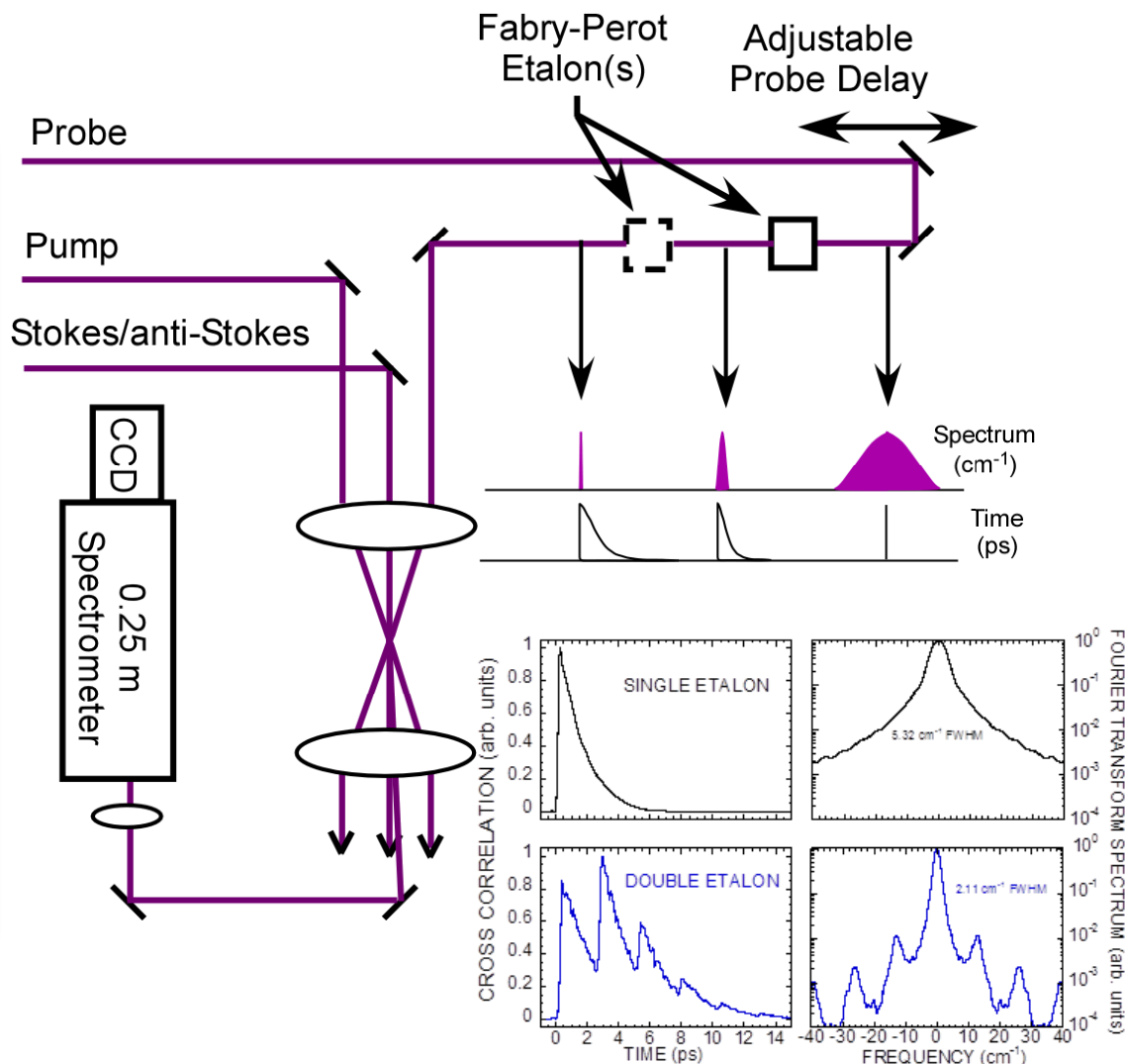


$$\chi = \sum_k X_k \sum_J W_J^{(k)} \exp(i \omega_J^{(k)} t) \exp(-\Gamma_J^{(k)} t)$$

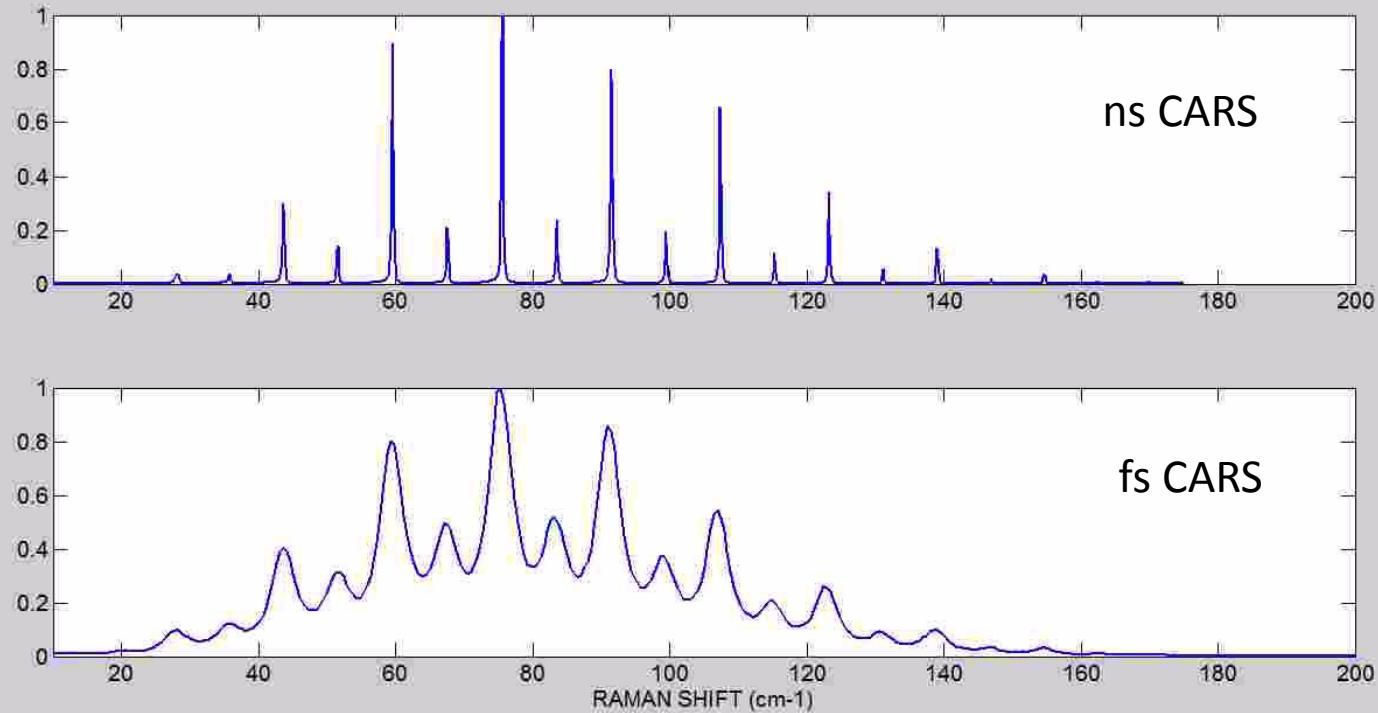
$$E_{CARS}(t) = E_{pr}(t - \tau) \exp[i \omega_o(t - \tau)] \times \chi(t)$$

Proof-of-Concept Experiments in Air

- “Bandwidth-Carving” to generate ps probe pulse
- Raman lines “resolved”
 $\Delta t \Delta \omega \leq \text{const.}$
- Two different probe resolutions investigated with single- and double-etalon configurations
- Very inefficient (0.8 to 2.4% or less transmission)
- Atmospheric air spectra in tube-furnace up to 800 K for probe delays up to 20 ps

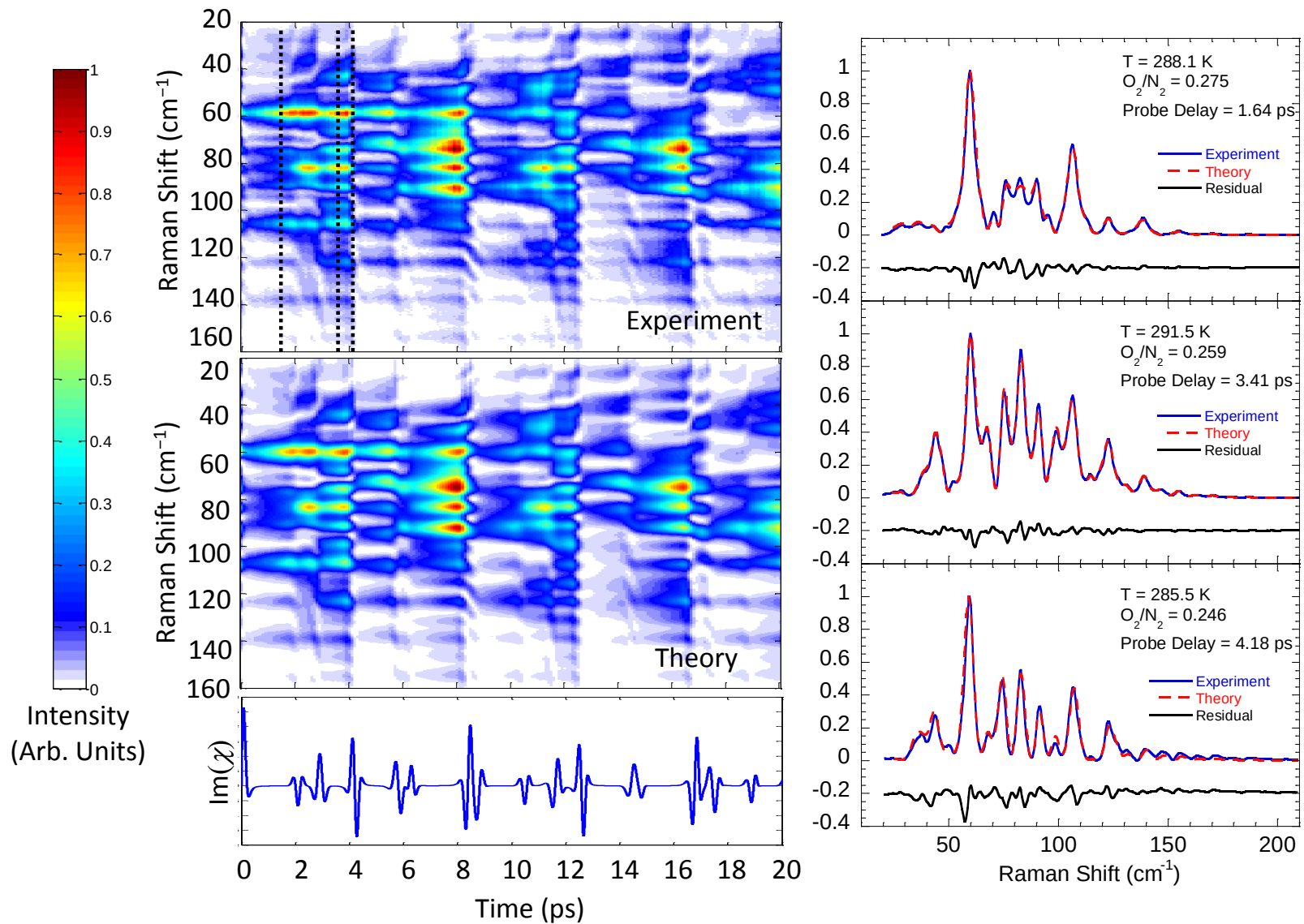


Low noise fs preparation pulses result can result in higher single-laser-shot precision

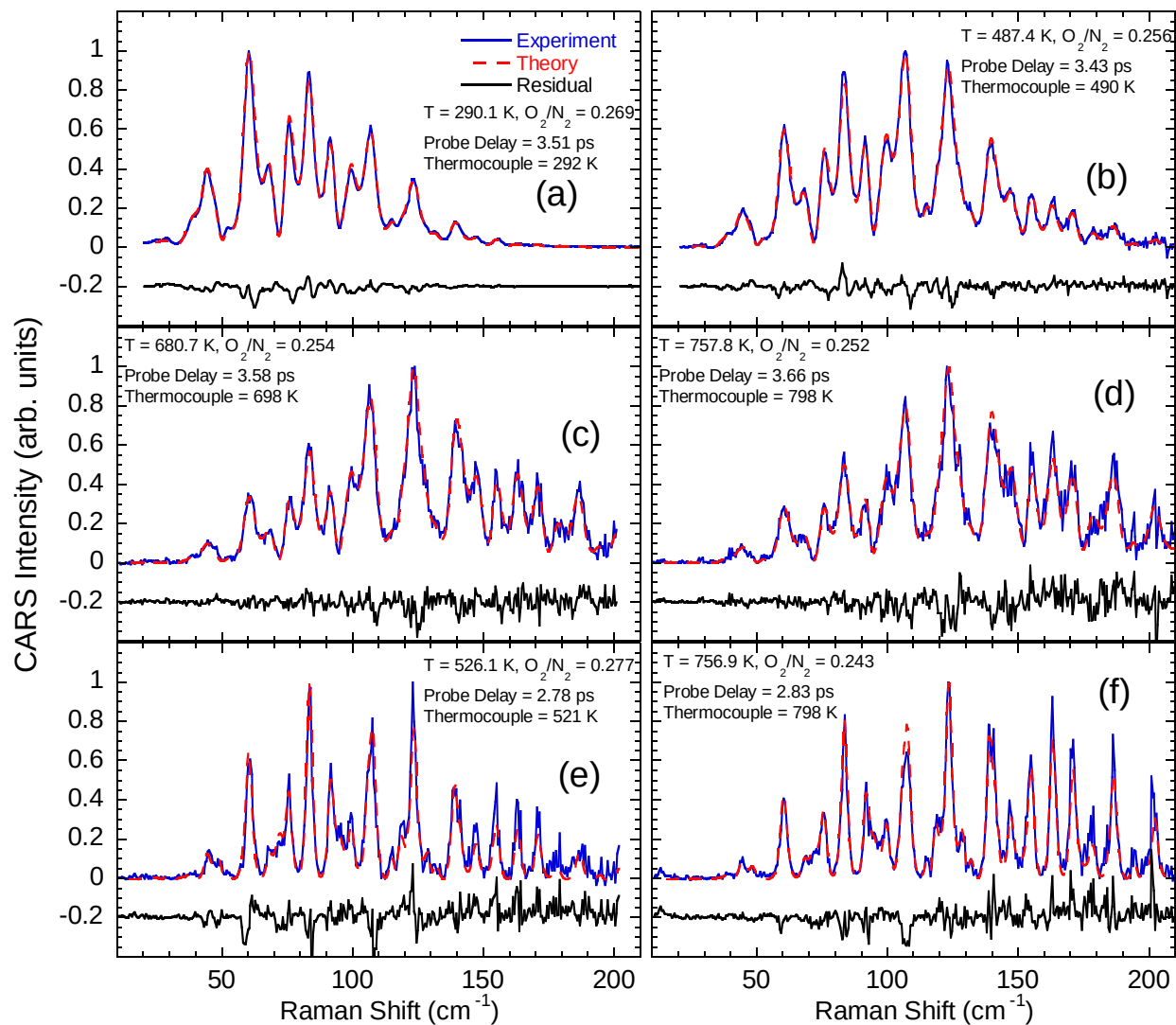


Room-temperature N_2 spectra

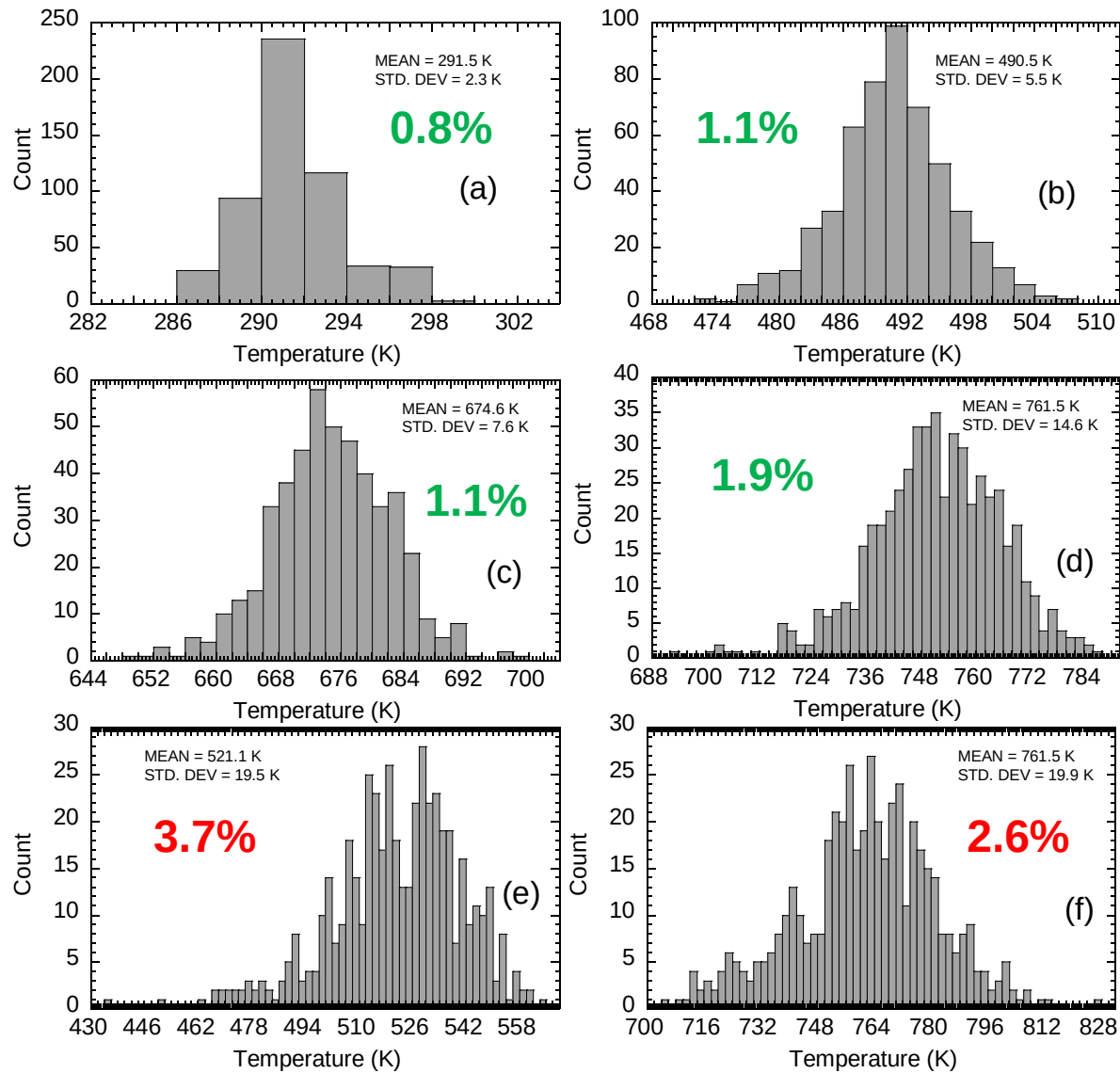
Model validation in room-temperature air: 1.5-ps probe



Collision-free spectra acquired at 1 kHz rate!

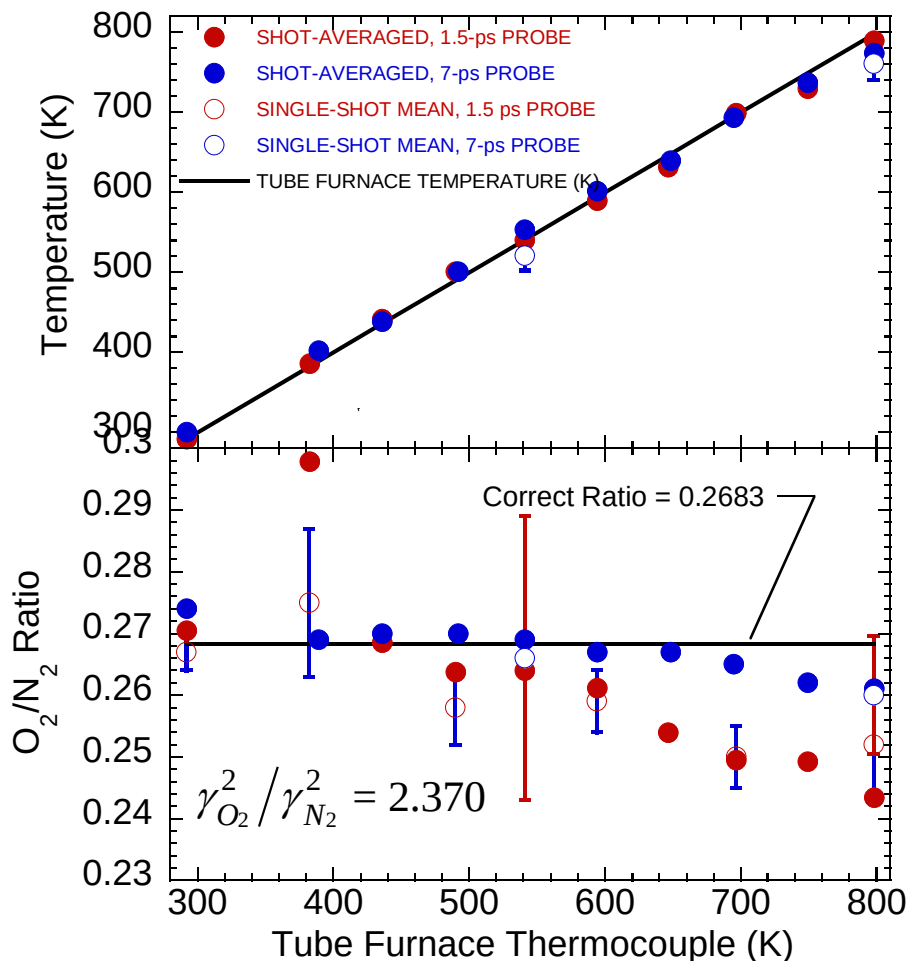


Single-shot temperature histograms



Proof-of-concept: Results and observations

- We have demonstrated single-shot thermometry and concentration measurements in air at temperatures up to 800 K
- Air temperatures are generally within 2.5% of the tube-furnace control thermocouple, with some exceptions at $T = 800$ K.
- Single-shot temperature precision is 1-2%. This superior to all but the best results obtained with ns-CARS
- Excellent shot-to-shot spectral repeatability from low-noise fs preparation pulses
- O₂ measurements need additional refinement
- Probe pulse energy must be increased to reach flame temperatures



S.P. Kearney *et al.*, *Optics Express* **21** (2013)

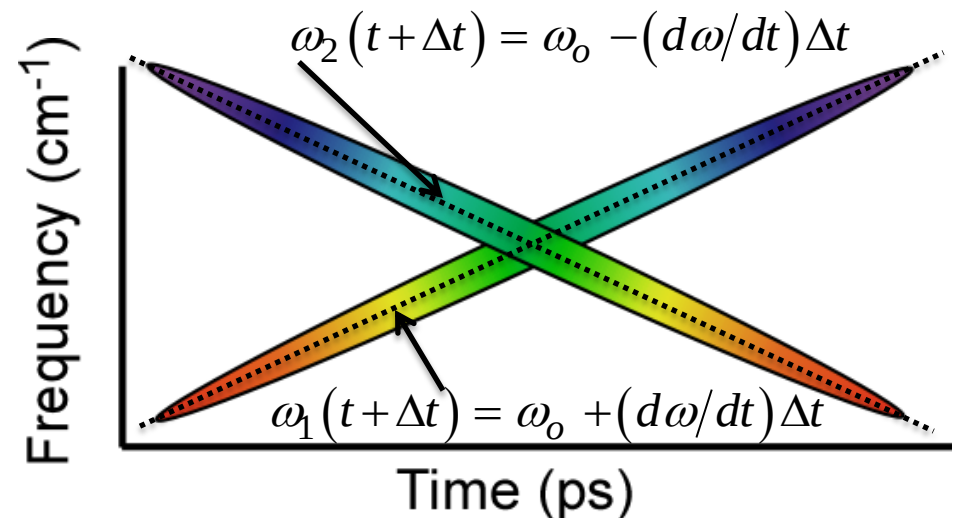
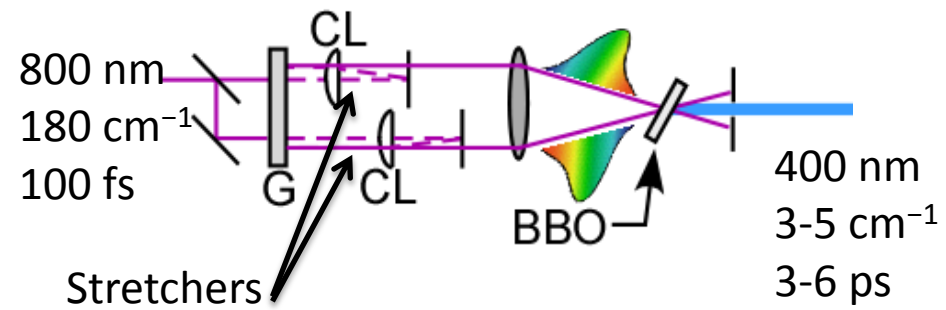
Second-Harmonic Bandwidth Compression (SHBC)

- Commercial device (Light Conversion)
- Converts fs radiation at 800 nm to ps radiation at 400 nm
- Grating pulse stretchers
- **Phase-conjugate** temporal chirps imparted upon broadband fs pumps
- Sum-frequency generation in BBO
- Output linewidth 3.5-4.0 cm^{-1}

$$\Delta\omega_{sfg} \sim d\omega/dt$$

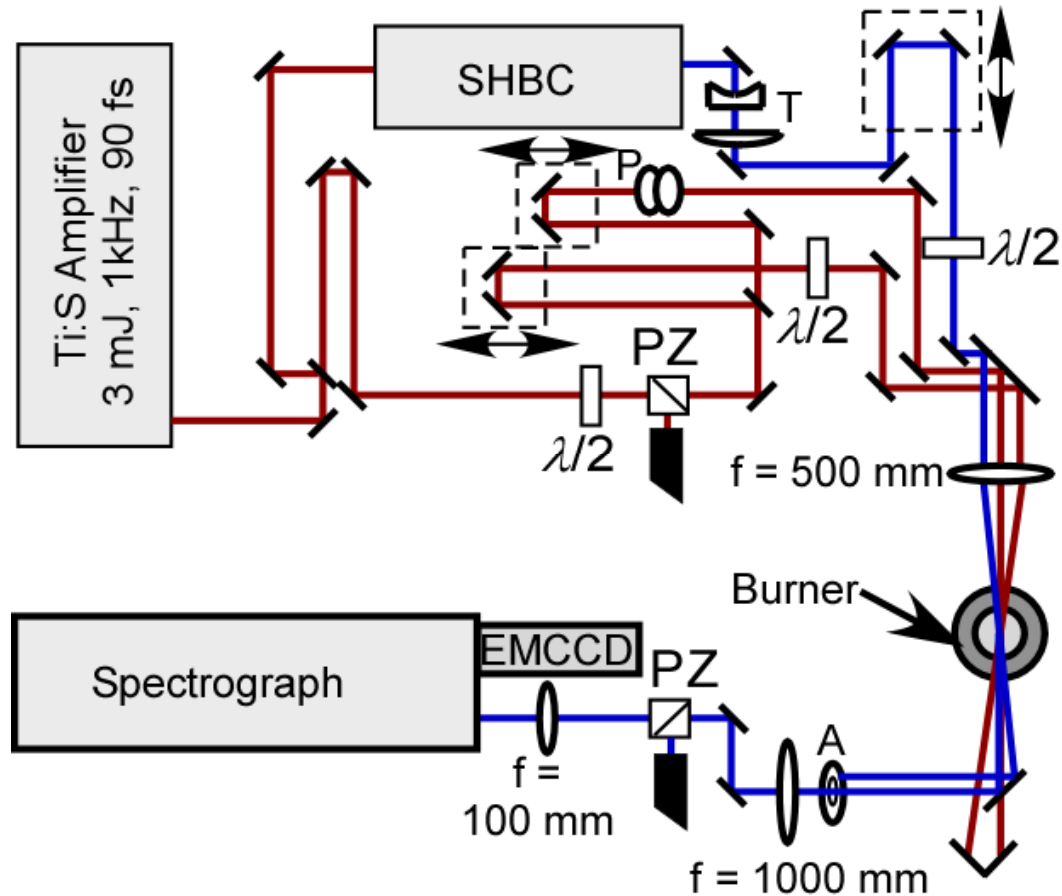
$$\Delta\omega_{sfg} \sim (\Delta t)^{-1}$$

- Conversion efficiency: 35-50%!
- Output pulse energy: 1-1.4 mJ!



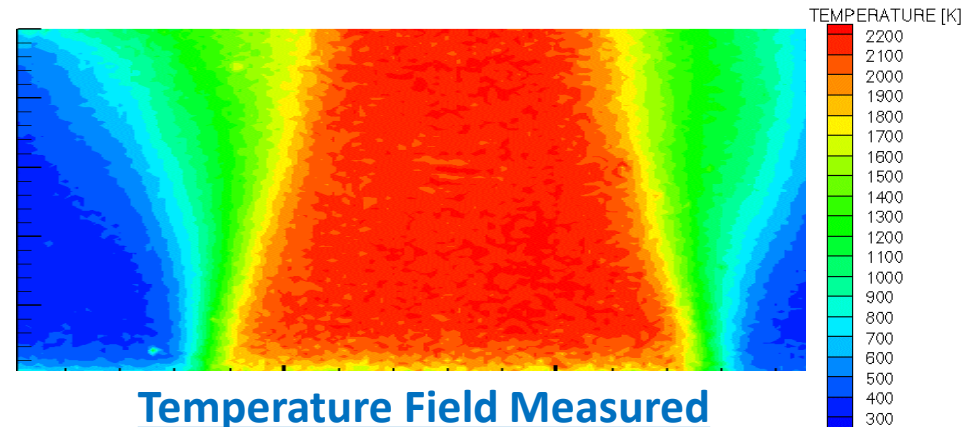
$$\omega_{sfg} = \omega_1 + \omega_2 = 2\omega_0$$

fs/ps CARS Instrument

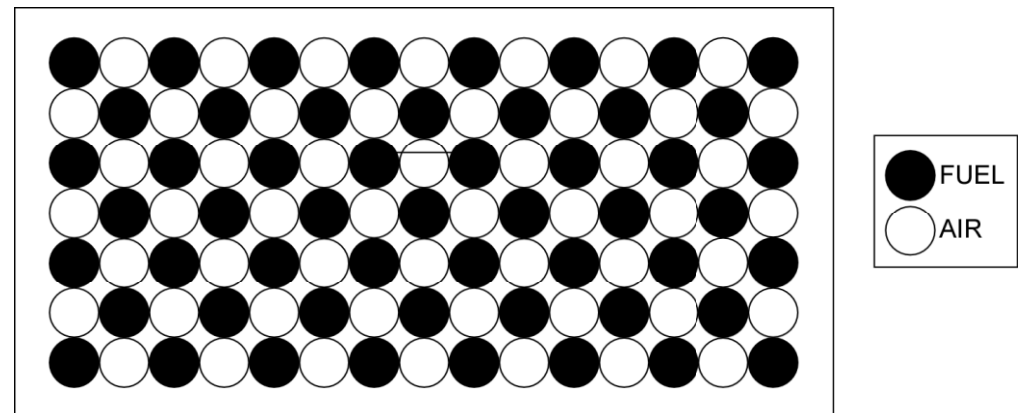


Hencken Burner

- Slightly lifted flat flame
- Flow rates 98 to 116 SLPM
- Non-premixed
- Provides *nearly adiabatic flames*
- Temperature and major species mole fractions *calculated from equilibrium*

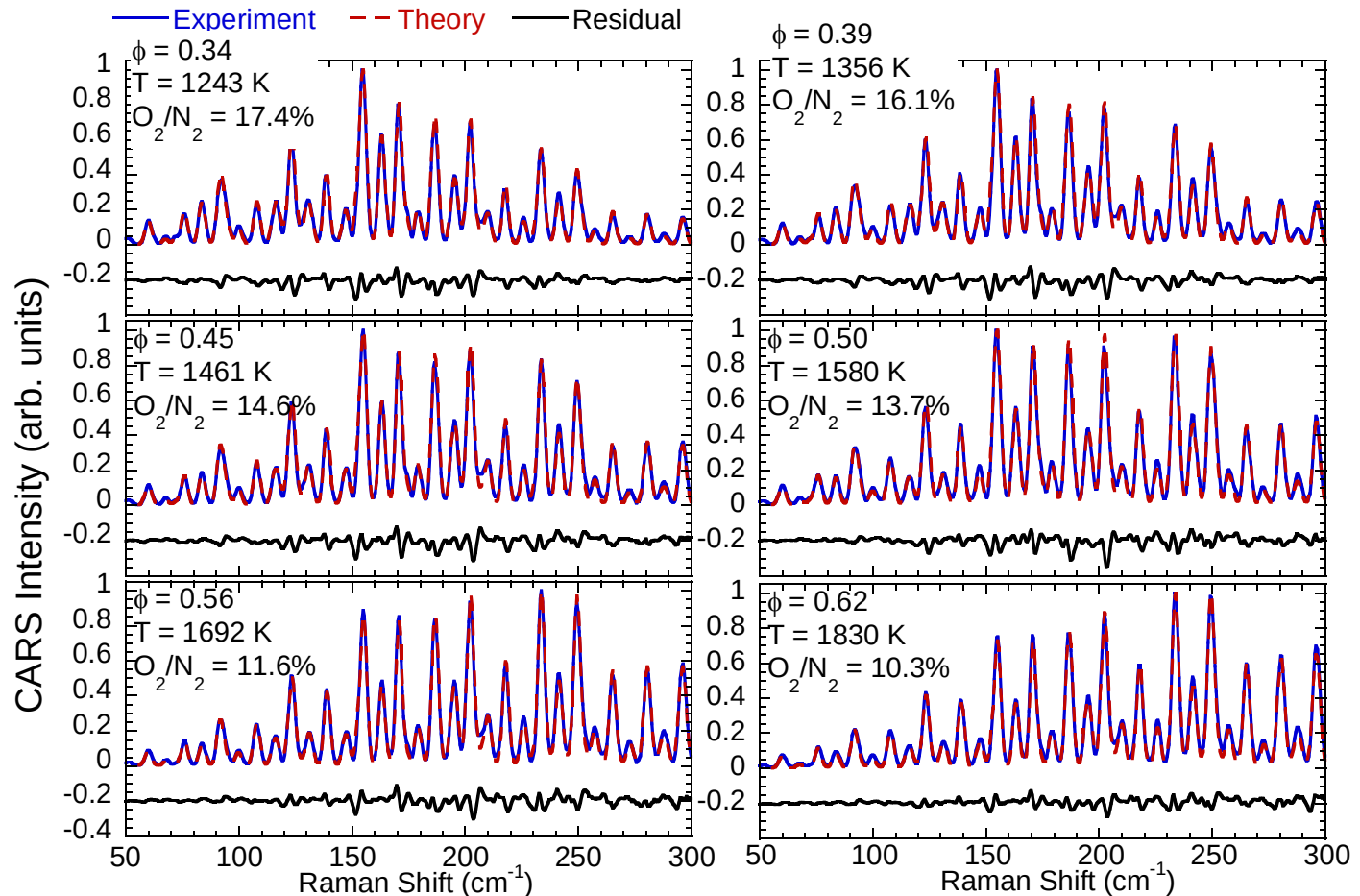


Temperature Field Measured
by Rayleigh Scattering
(CH₄/air)



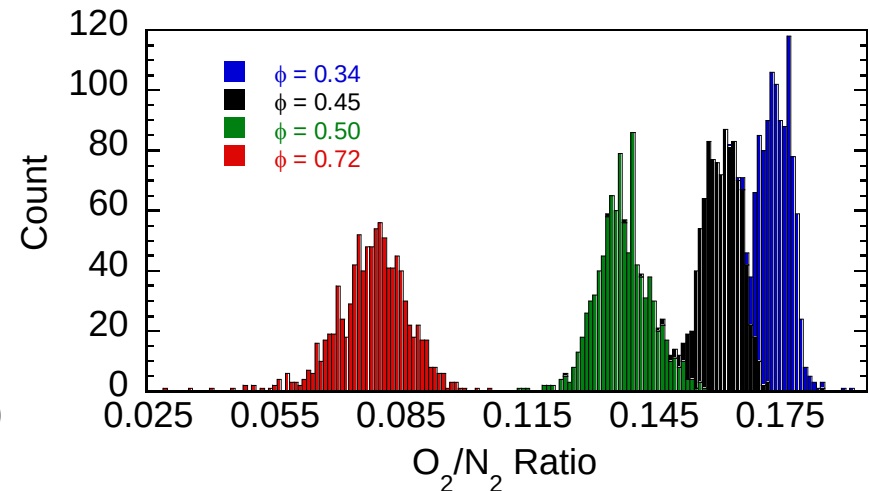
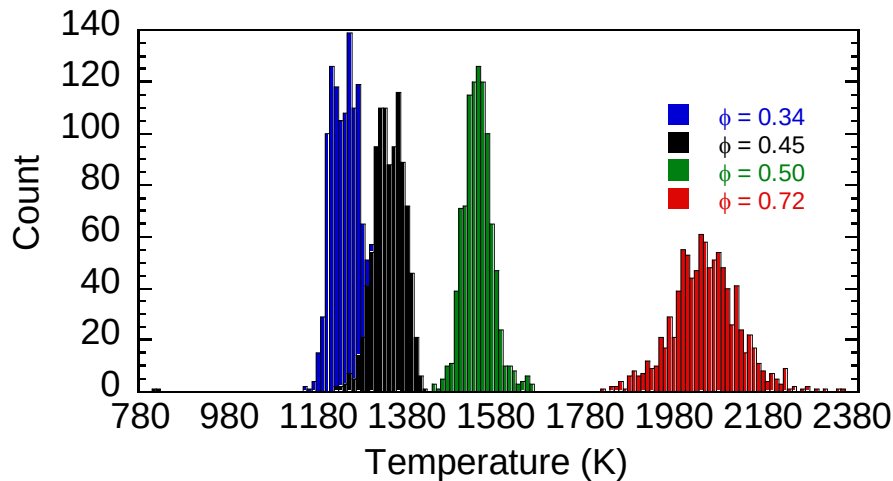
Burner Hole Configuration

Shot-Averaged Spectra from Near-Adiabatic H₂/air flame



- Spectra averaged for several thousand laser shots to optimize SNR
- N₂ contributions dominate all spectra
- O₂ sensitivity arises from alteration of spectral “envelope” and subtle line shifts

Temperature and O₂/N₂ Measurement Precision

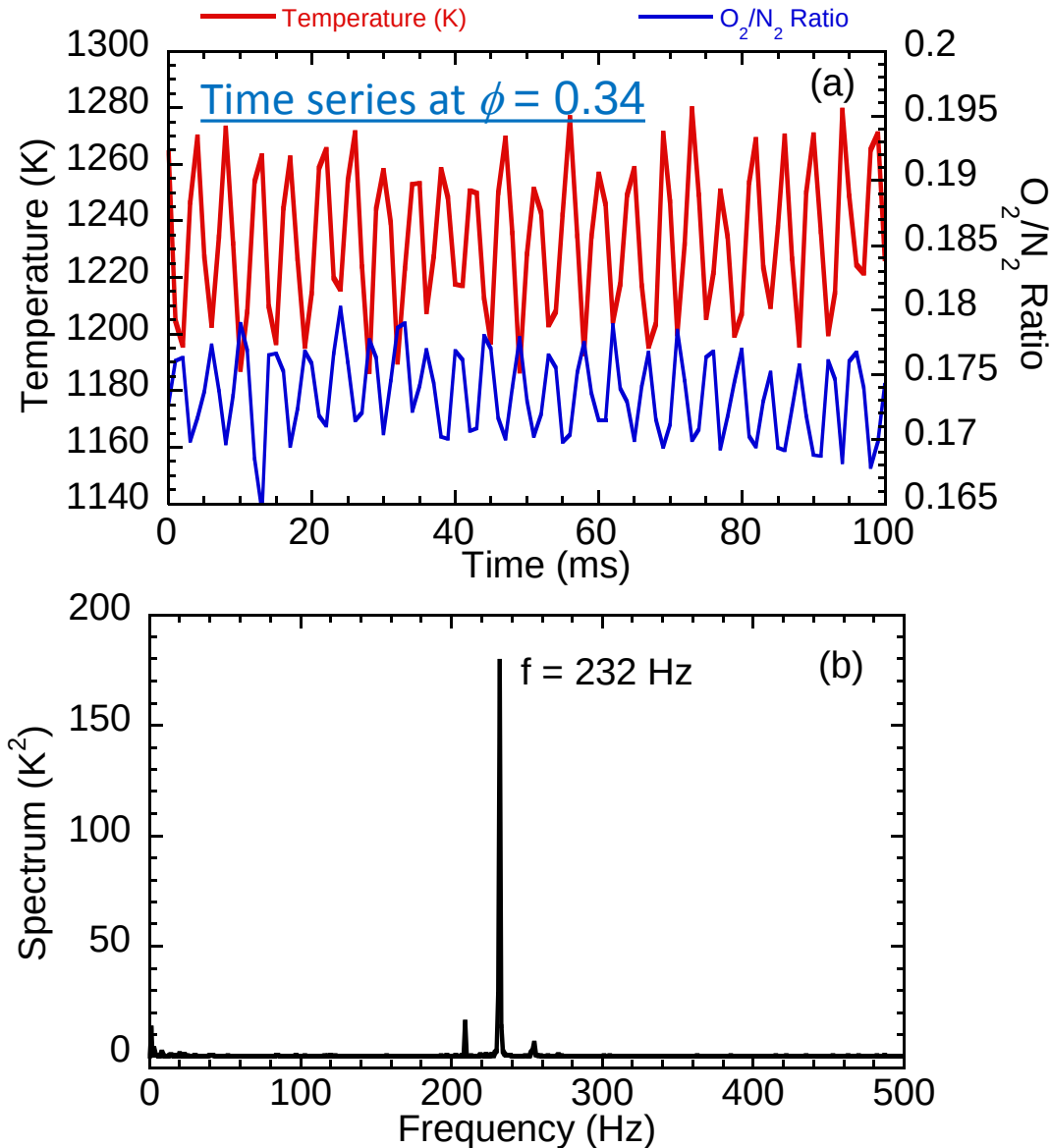


ϕ	T(K)	σ_T (K/%)	O ₂ /N ₂	σ_{O_2} (-/%)
0.34	1235	26.8/2.2%	17.4%	0.34/1.95%
0.39	1338	40.7/3.0%	16.1%	0.46/2.85%
0.50	1539	32.8/2.1%	13.7%	0.64/4.67%
0.72	2051	75.2/3.6%	7.8%	0.85/10.9%

EM Gain ON

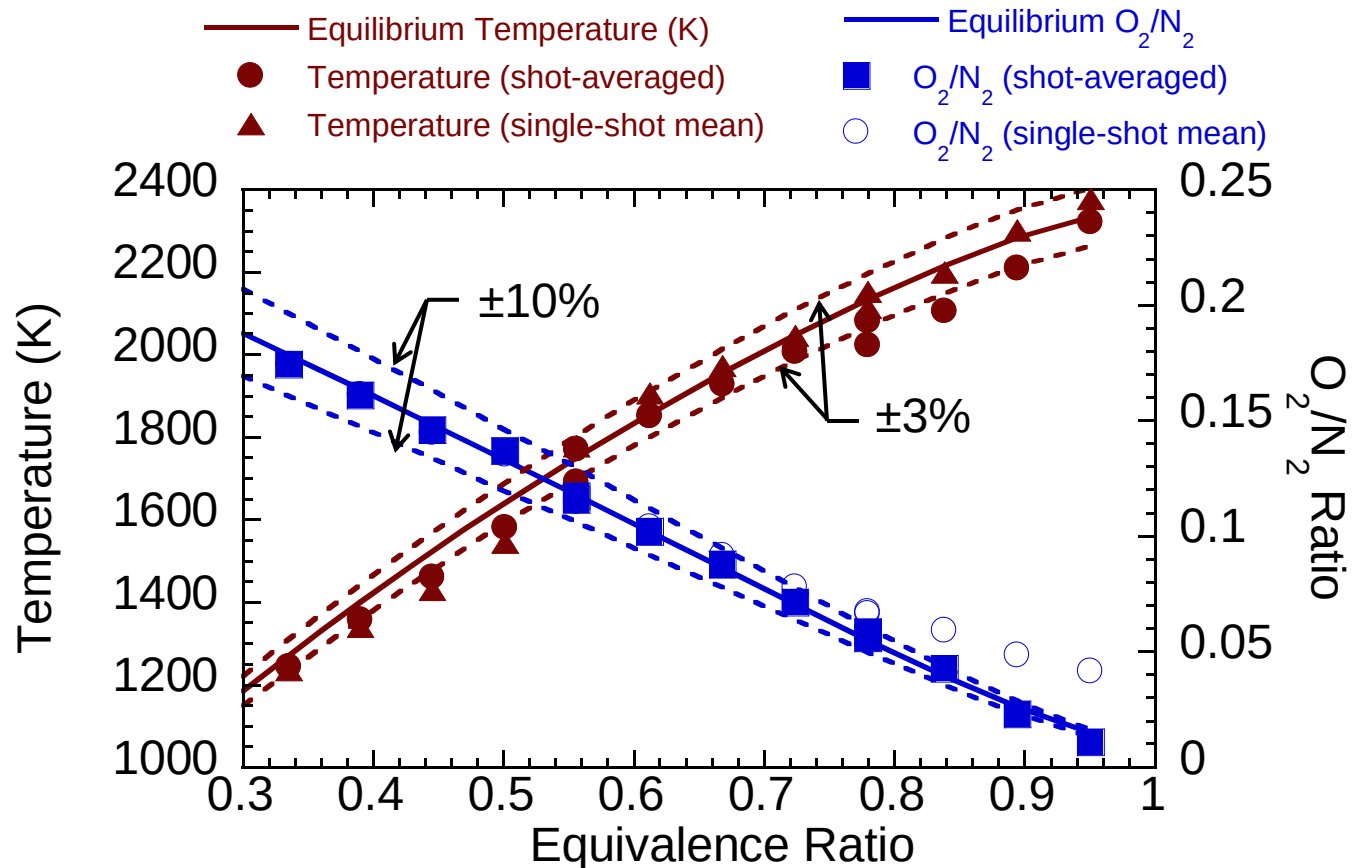
- Temperature precision is 2-3% with EM gain off
- Best results with fs CARS are ~1%.
- ns rotational CARS ~ 3-4%

Acoustic Interaction at ϕ up to 0.55



- Audible “hum” from burner at leanest stoichiometries
- Inspection of T and O_2 data reveals oscillation in both signals
- Sampled at 1 kHz
- Negative correlation is consistent with low-level oscillations in ϕ
- Amplitude of temperature oscillation is ~ 30 K (2.4%)
- Distinct peak in PSD near 232 Hz is consistent with tone heard from the burner
- Precision could be understated

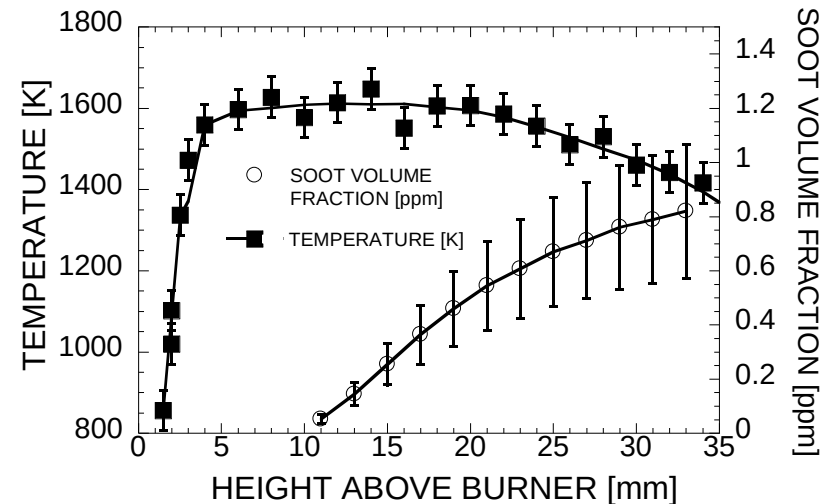
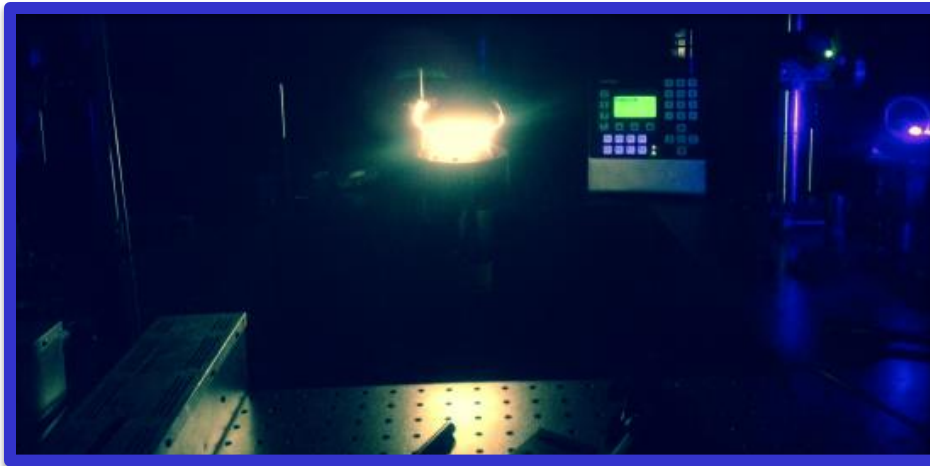
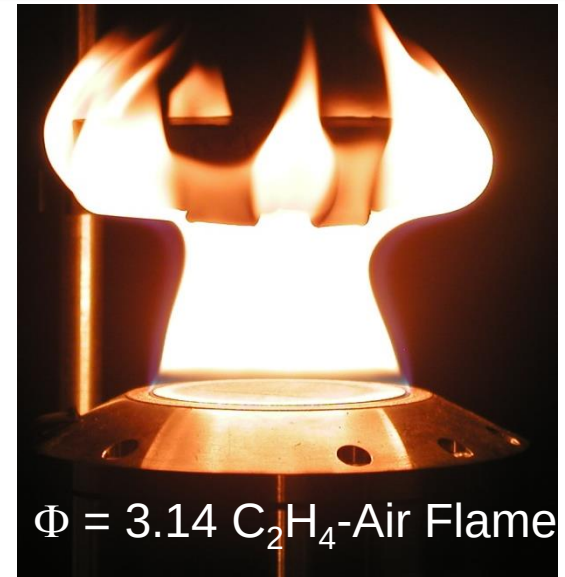
H₂/Air Flame Measurement Accuracy



- Temperature accuracy: -3% when shot-averaged spectra are used
- Temperature accuracy: -3-6% when single-shot means are used
- O₂/N₂ accuracy is ±6%
- Uncertainty due to metered gas flows: Temperature ±3% / O₂/N₂ ±6-12%

C₂H₄/Air Flames on McKenna Burner

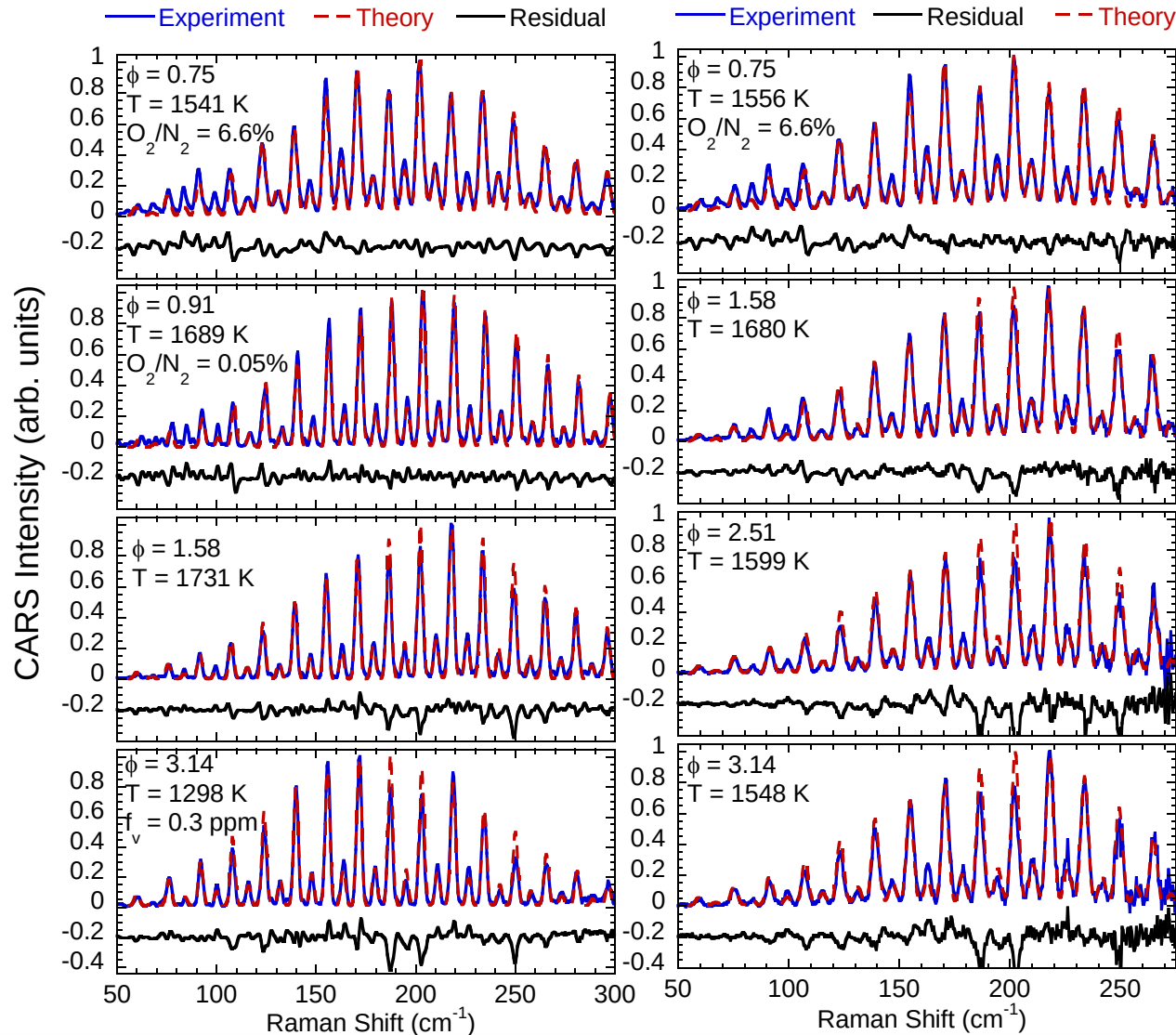
- Premixed hydrocarbon/air flame
- Water-cooled non-adiabatic burner
- Stable region ~5-15 mm above burner
- Previously studied at $\phi = 3.14$ in our lab (and elsewhere!)
- Wide range of stoichiometry, $\phi = 0.75$ to 3.14
- Potential contributions from N₂, O₂, CO, (CO₂ minimized by probe delay)



C₂H₄/Air Flame Spectra

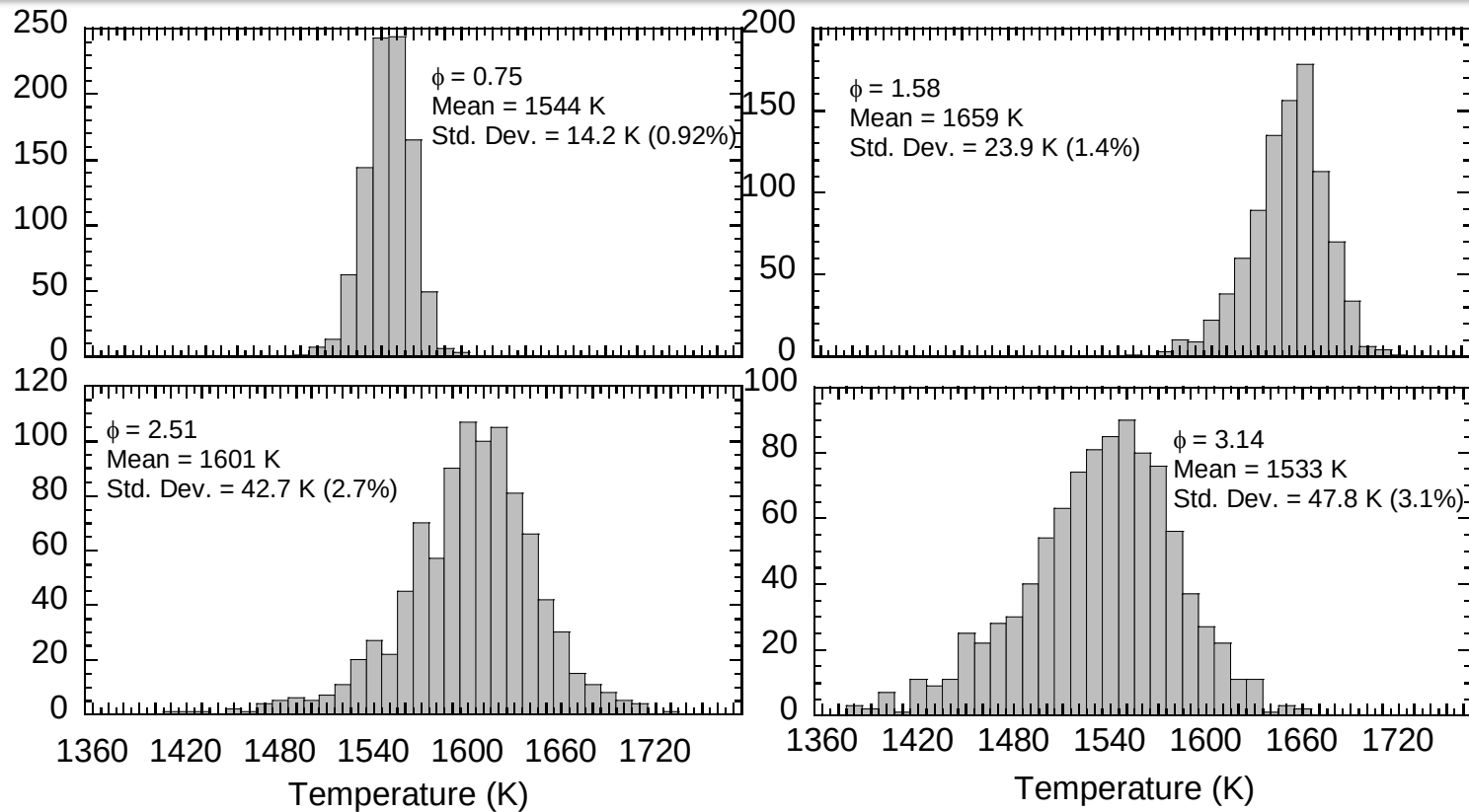
Shot-Averaged

Single-Shot at 1 kHz

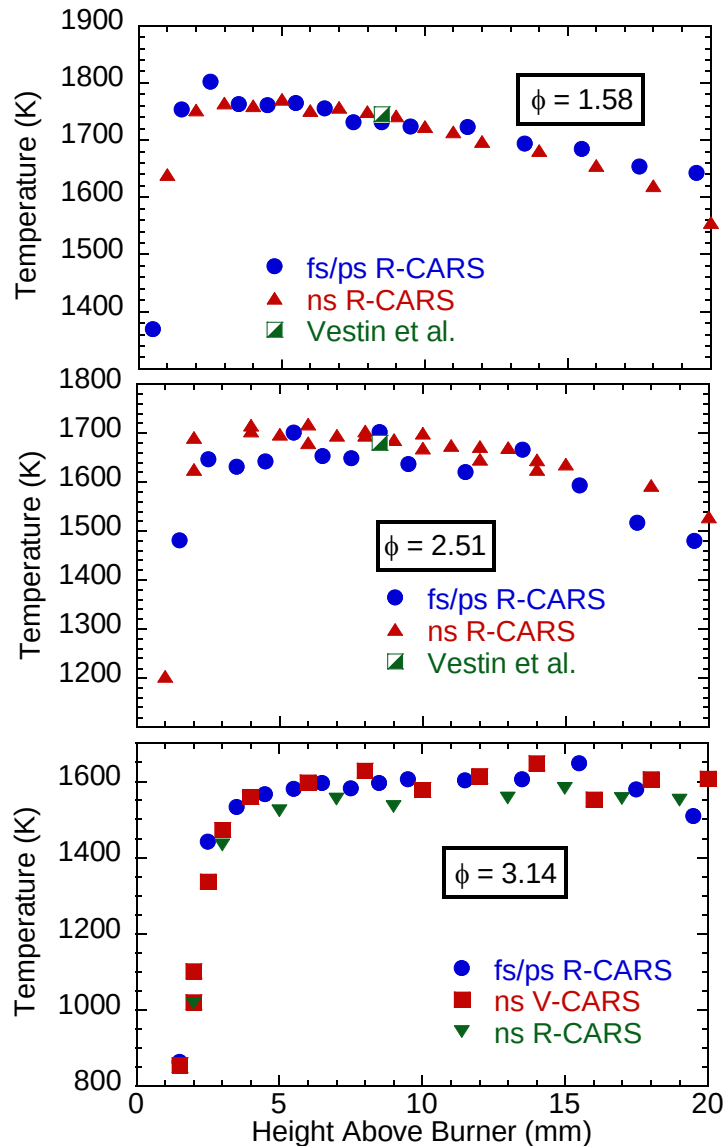


- Spectra acquired for fuel-lean to rich sooting flames
- High-quality fits observed for $\phi < 1$
- Systematic bias toward “underfit” of isolated lines for fuel-rich flames
- Fitted temperature appears to be robust
- Reliable spectra obtained in sooting regions of the flame
- No EM gain used

Temperature Histograms C₂H₄/Air Flame



- Measurement volume positioned 11.5 mm above burner where flame is stable
- Minimal variation in temperature in this range of ϕ
- Precision is 0.9 to 1.4% in leanest flames investigated : CARS photon yields are highest
- Precision degrades to ~3% in richest flames
- Precision appears to be correlated to photon yield



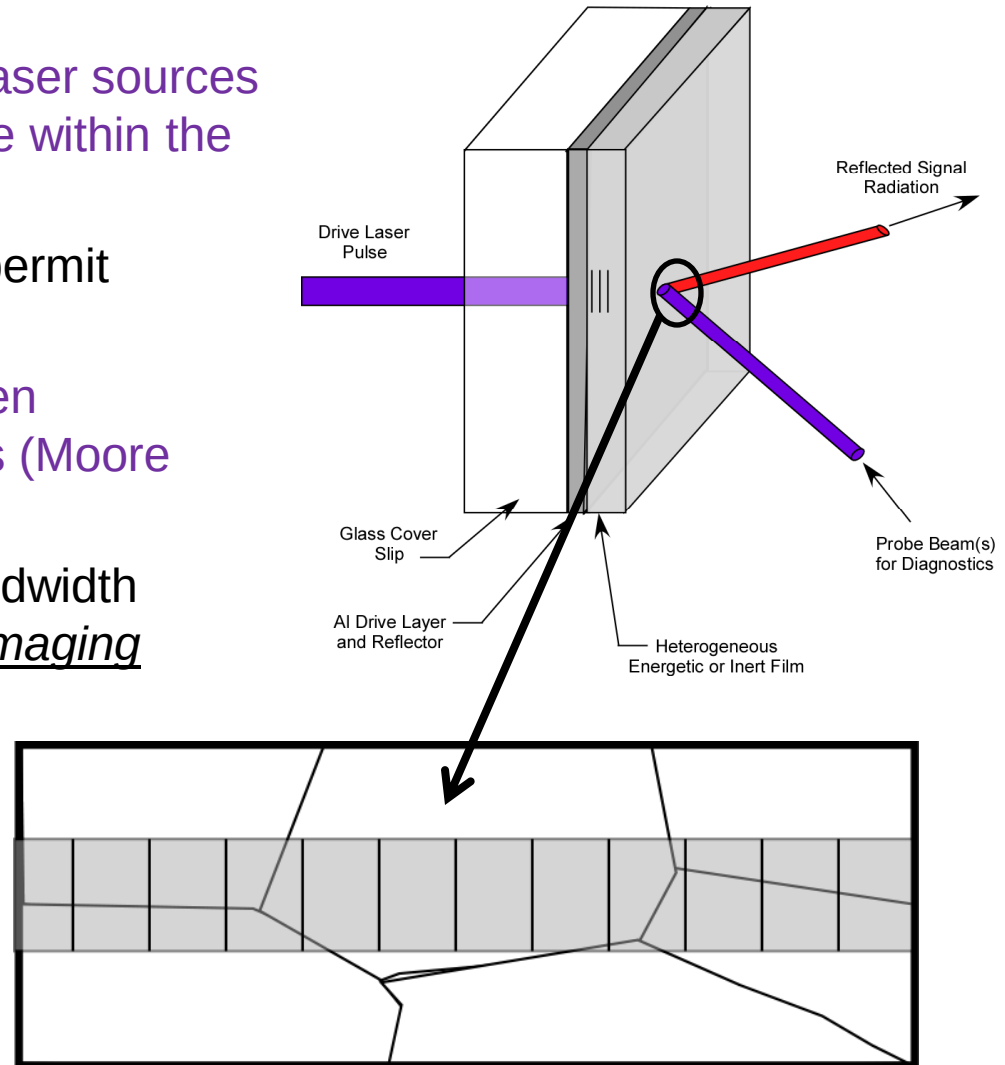
- Comparison is most valid in stable region 5-15 mm height
- ns-CARS temperatures based on fits to shot-averaged spectra
- Agreement of fs/ps and ns CARS is within 1-3% in these *fuel-rich, sooting* flames
- Within reported accuracy of rotational ns-CARS
- ns-CARS measurements
 - Rotational CARS at Sandia 11/2013
 - Vibrational CARS at Sandia (2006)
 - Rotational CARS at Lund (Vestin et al. 2005)

Summary and Conclusions

- High-energy probes generated via SHBC enable flame temperature measurements with rotational fs/ps CARS
- We have performed a **systematic assessment of accuracy and precision**
 - Near-adiabatic H₂/air flames
 - Premixed hydrocarbon air flames for $\phi = 0.75$ to 3.14
- Temperature accuracy is 3% for shot-averaged spectra, 3-6% for single-shot means
- Uncertainty in temperature standards is 3%
- Accuracy of O₂/N₂ ratio is 2-6% in H₂/air flames
- Temperature precision is outstanding!
 - As good as 1%
 - Periodic oscillations of 30 K (2%) amplitude realized in Hencken burner
 - Results suggest correlation of CARS signal photons to precision
- Precision of O₂/N₂ ratio is 1-10% and monotonically increases with O₂/N₂
- O₂/N₂ sensitivity arises from spectral envelope and not resolution of O₂ lines

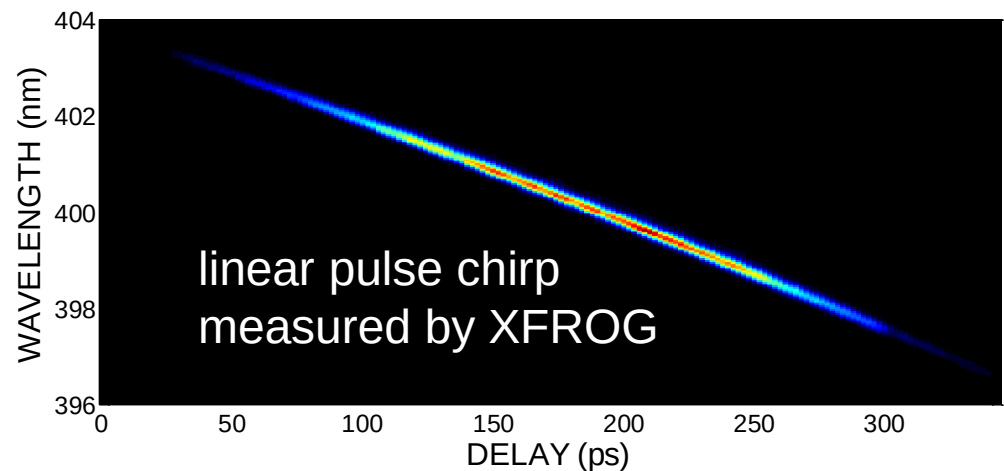
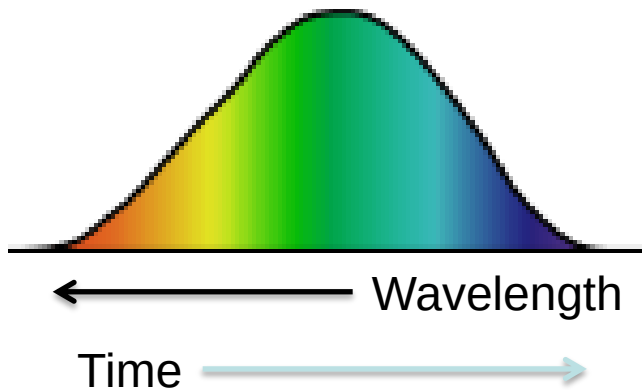
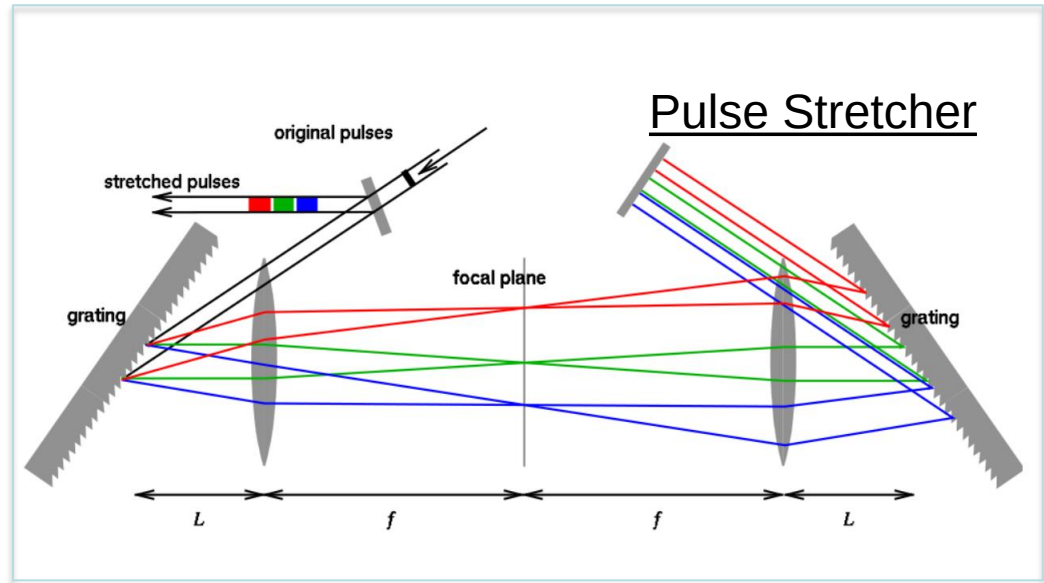
Our solution: Use femtosecond laser pulses for mesoscale-resolved experiments on heterogeneous materials

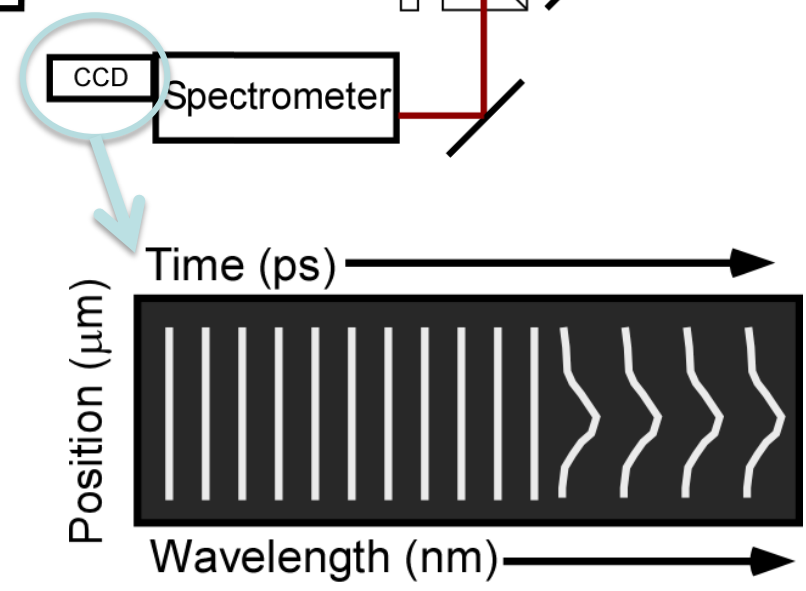
- Reliable femtosecond (fs) pulsed laser sources have become commercially available within the last ~10 years
- Ultrashort 10^{-14} – 10^{-13} sec pulses permit access to ps-scale events
- Well-controlled laser drive has been demonstrated in homogeneous films (Moore group, LANL)
- Broadband yet low-noise laser bandwidth enables multiple spatially resolved imaging spectroscopies
- Spatially resolved measurements on heterogeneous samples at realistic grain sizes



Femtosecond laser pulses have wide frequency spectra that allow us to encode time into frequency (color)

- Femtosecond bandwidths are typically 10s of nanometers
- Can be readily compressed or stretched (chirped) using grating-based stretchers
- Linear chirp is added in our lab to encode wavelength into time with picosecond resolution

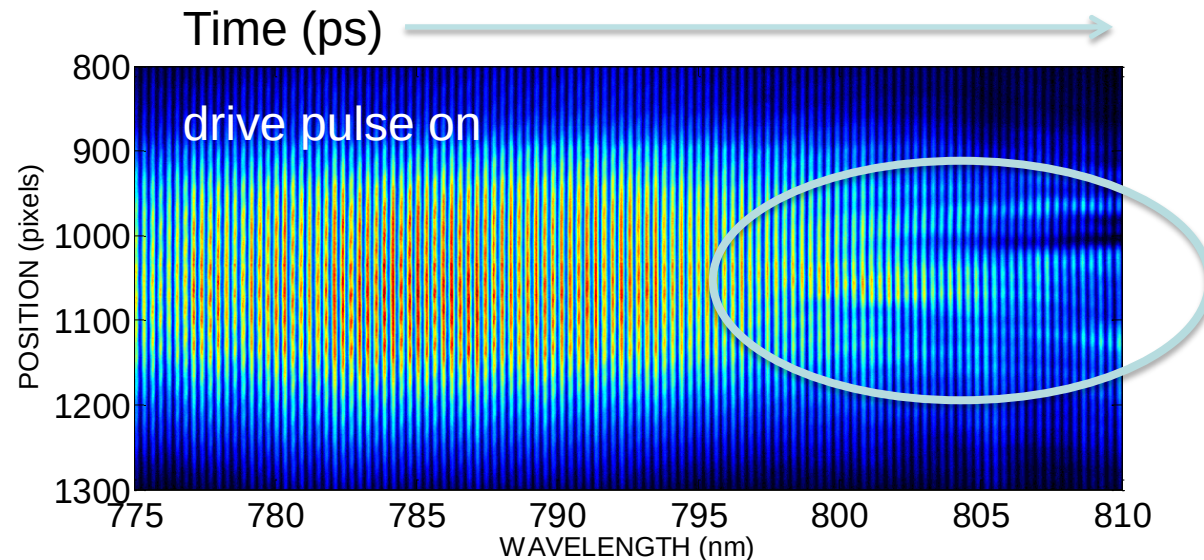
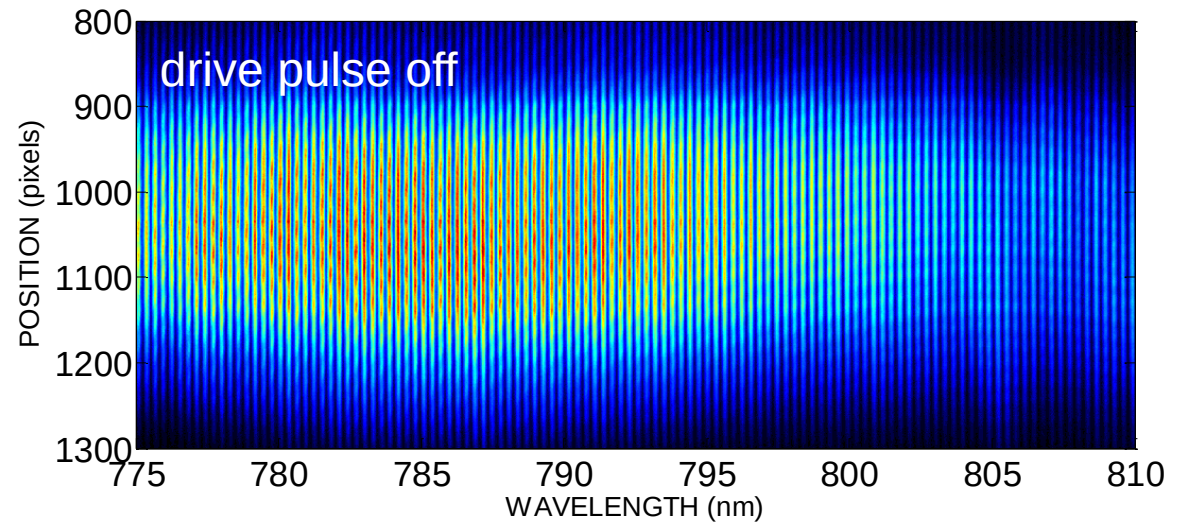
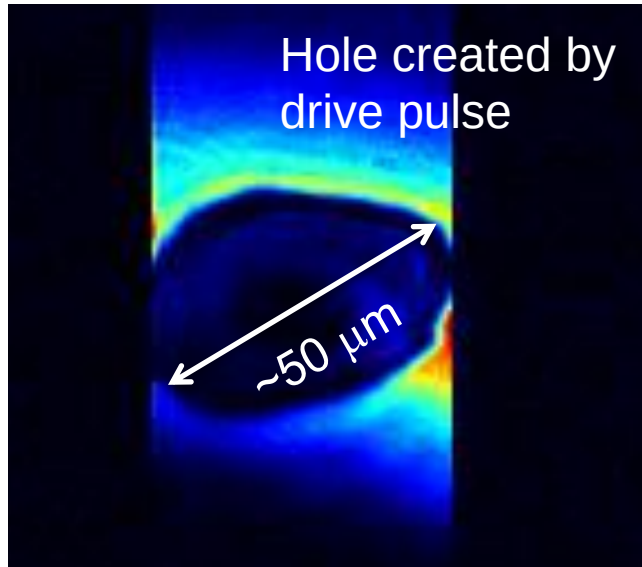




-
- Time (ps)
- Position (μm)
- Wavelength (nm)

USI interferograms: laser-shocked aluminum films

- 1-micron Al film on glass substrate
- Ablator surface motion is tracked
- Surface velocity vs time and shock breakout profiles monitored



USI-measured ablator surface velocity histories

