

Phonon considerations in the reduction of thermal conductivity in phononic crystals

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Controlling thermal conductivity through phonon transport

Phonon engineering for thermoelectric applications

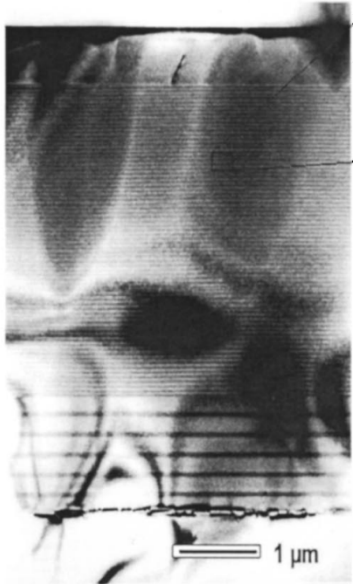
Thermoelectric

Figure of Merit

$$Z = \frac{S^2 \sigma}{\kappa}$$

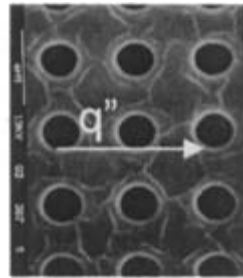
κ

Superlattices



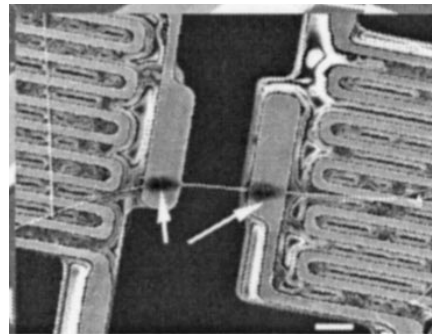
Cahill *et al.* Journal of Heat Trans. **124**, 233 (2002)

Microporous materials



Song and Chen, APL **84**, 687 (2004)

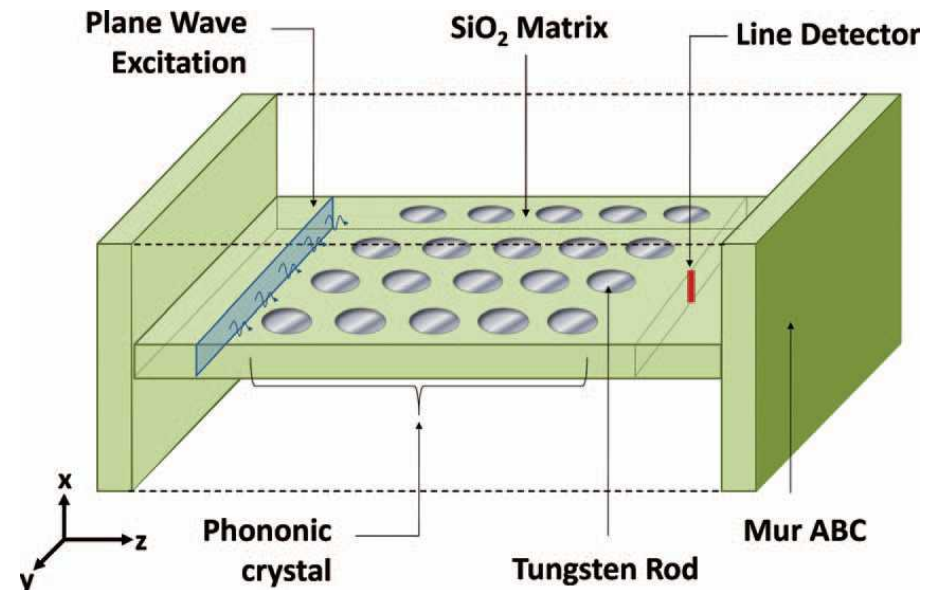
Nanowires



Li *et al.*, APL **83**, 2934 (2003)

Recent realization of GHz PnC

Su, Olsson, Leseman, and El-Kady, APL **96**, 05311 (2010)



More details in this session at
15:10 (Roy H. Olsson III)



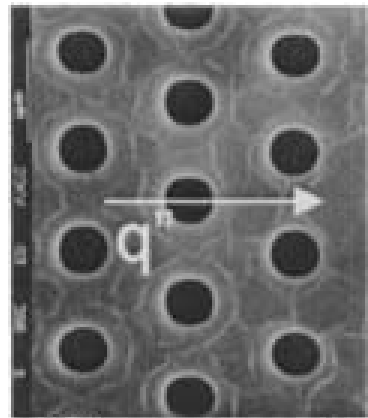
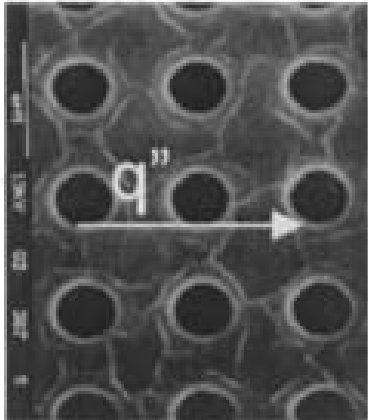
Outline

Goal: Study phonon scattering and thermal conductivity reduction in PnC

- Predicting the thermal conductivity in periodically arranged porous solids
- Thermal conductivity reduction in the PnC
- Thermal conductivity measurements of PnC with time-domain thermoreflectance

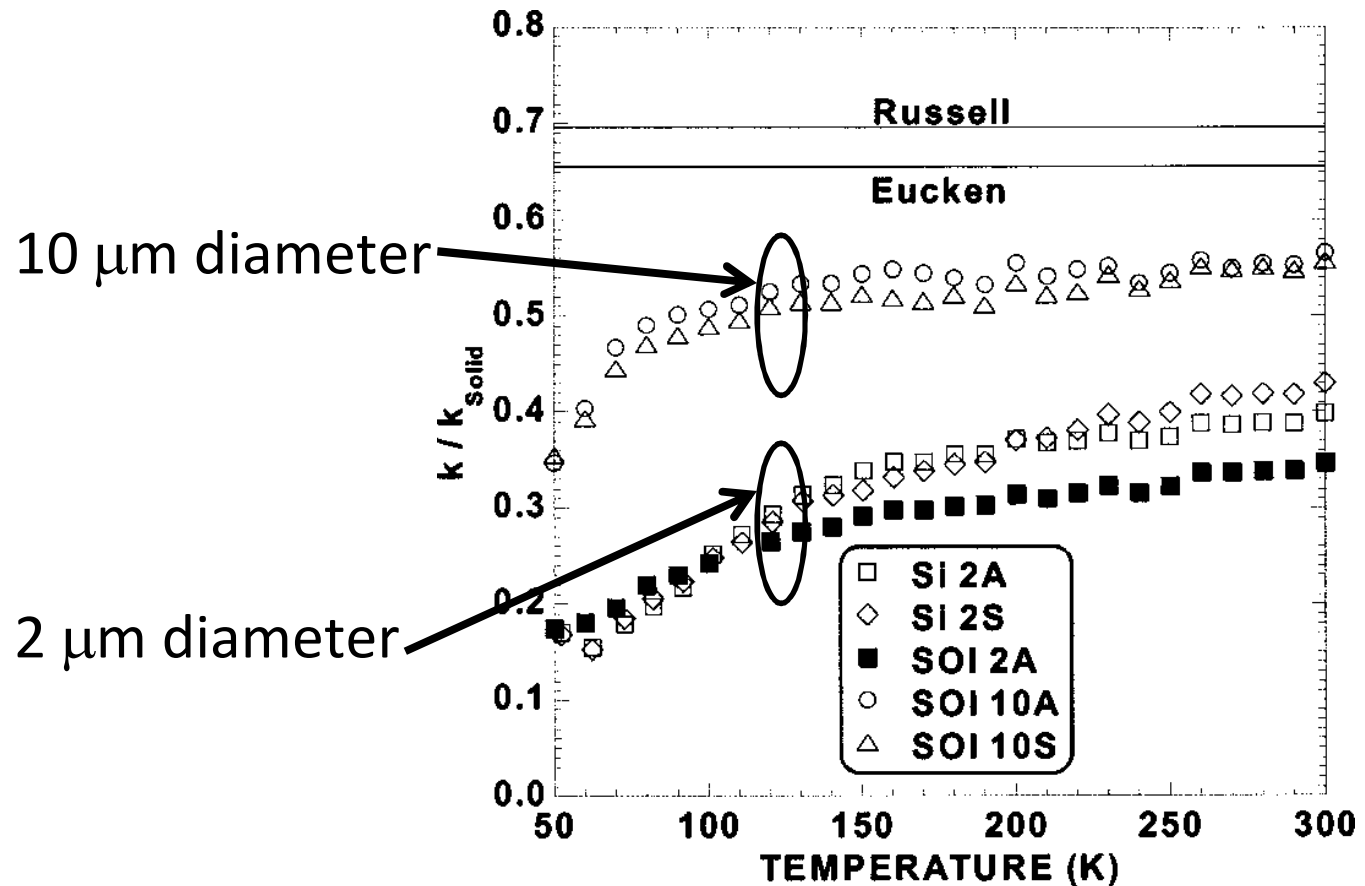
Thermal conductivity in periodic, porous Si

Aligned



Staggered

Song and Chen, APL **84**, 687 (2004)



How can we model this thermal conductivity reduction?

Thermal conductivity in bulk Si

Thermal conductivity

$$\kappa = \frac{1}{6\pi} \sum_j \int_q \frac{\hbar^2 \omega_j^2(q)}{k_B T^2} \frac{\exp\left[\frac{\hbar \omega_j(q)}{k_B T}\right]}{\left(\exp\left[\frac{\hbar \omega_j(q)}{k_B T}\right] - 1\right)^2} v_j^2(q) \tau_j(q) q^2 dq$$

Fit thermal conductivity model to experimentally measured κ on bulk Si to determine τ (A, B, C , and E)

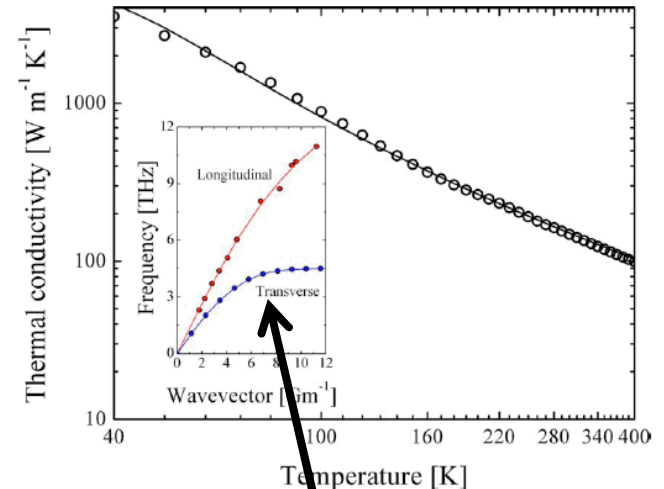
Scattering time

$$\frac{1}{\tau_j(q)} = \frac{1}{\tau_{\text{Umklapp},j}} + \frac{1}{\tau_{\text{impurity},j}} + \frac{1}{\tau_{\text{boundary},j}}$$

$$\frac{1}{\tau_{\text{Umklapp},j}} = BT \omega_j^2(q) e^{\frac{C}{T}}$$

$$\frac{1}{\tau_{\text{Impurity},j}} = A \omega_j^4(q) \quad \frac{1}{\tau_{\text{Boundary},j}} = \frac{\partial \omega_j(q)}{\partial q} \frac{1}{E}$$

Similar approach to that described by Mingo, PRB **68** 113308 (2003)



Fit polynomial to measured dispersion in bulk Si

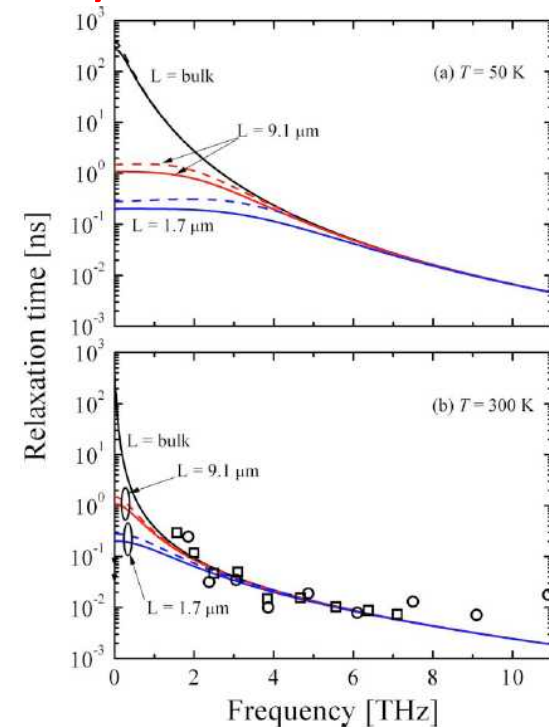
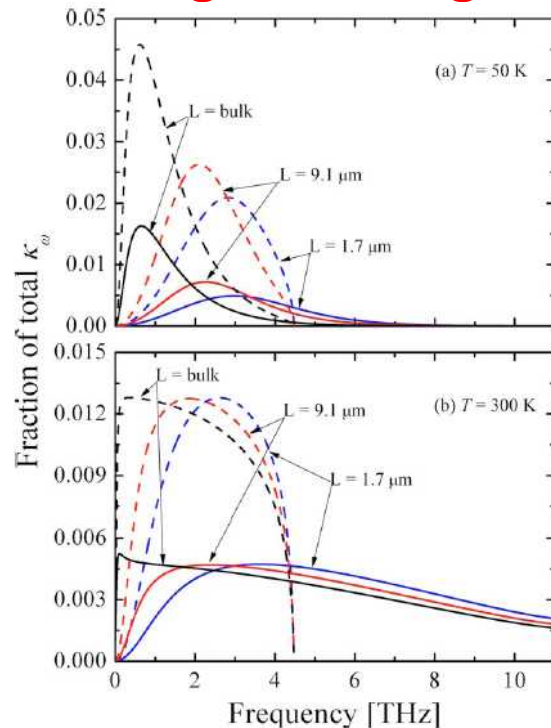
Origin of reduction of κ in microporous Si

Scattering at pore boundaries

$$\frac{1}{\tau_{Boundary,j}} = \frac{\partial \omega_j(q)}{\partial q} \frac{1}{L}$$

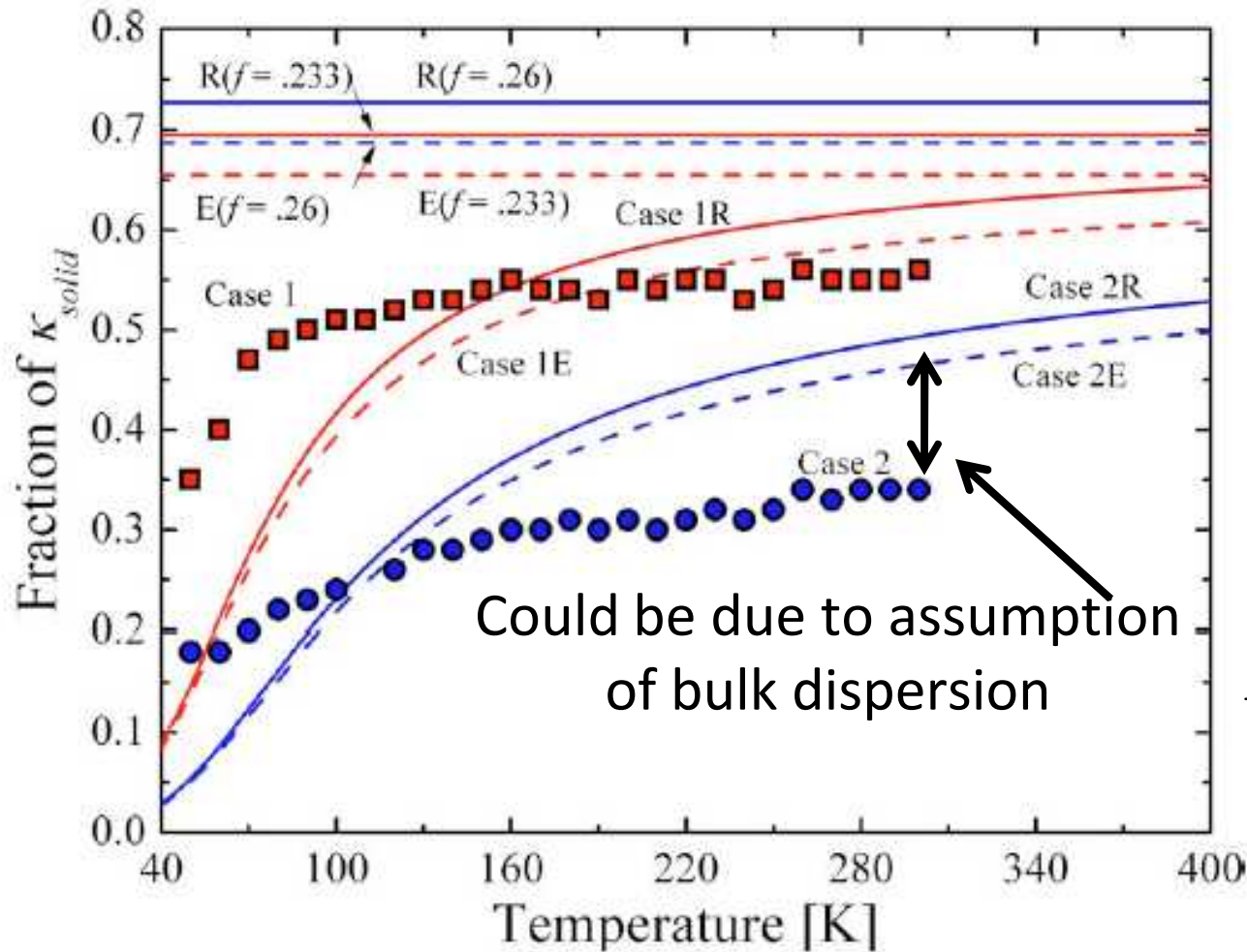
← Pore-edge to pore-edge spacing

Long wavelength modes scattered by boundaries



Predicting κ in microporous Si

Porosity accounted for by “Eucken” and “Russel” models



“E” model

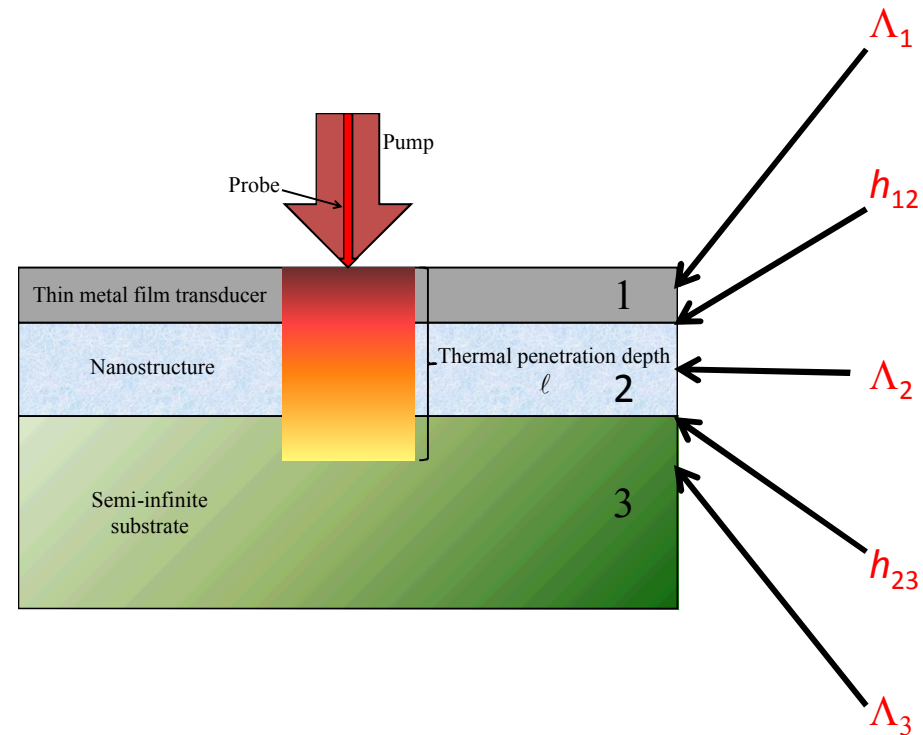
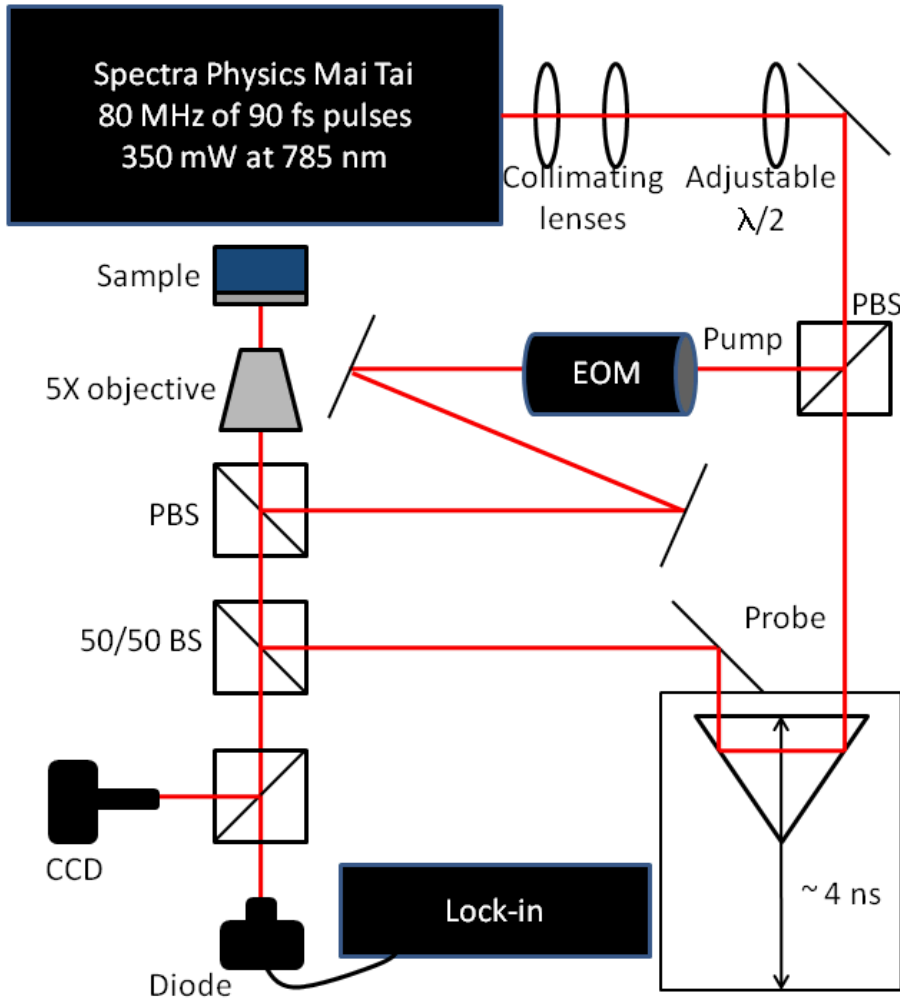
$$\frac{\kappa_{porous}}{\kappa_{solid}} = \frac{1-p}{1+\frac{p}{2}}$$

“R” model

$$\frac{\kappa_{porous}}{\kappa_{solid}} = \frac{1-p^{\frac{2}{3}}}{1-p^{\frac{2}{3}}+p}$$

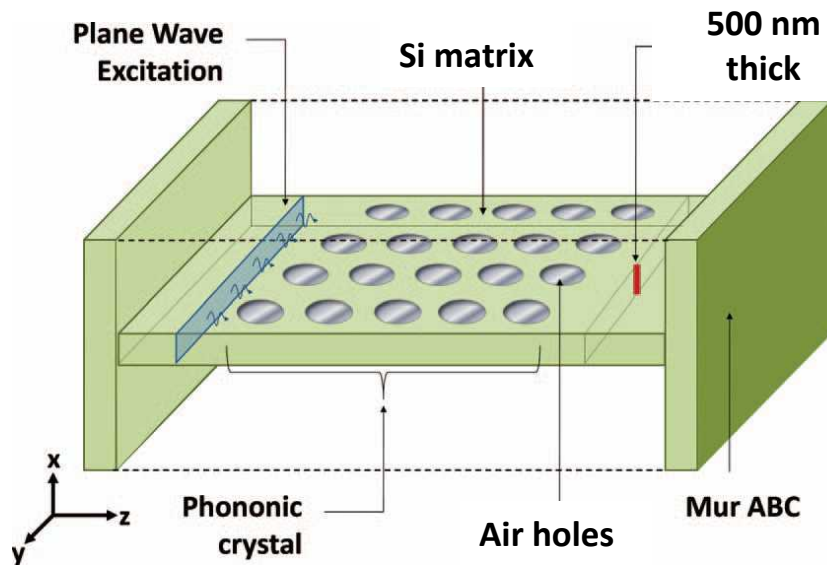
Consider bulk approach as “maximum” κ , or smallest reduction in κ

Time-domain thermoreflectance



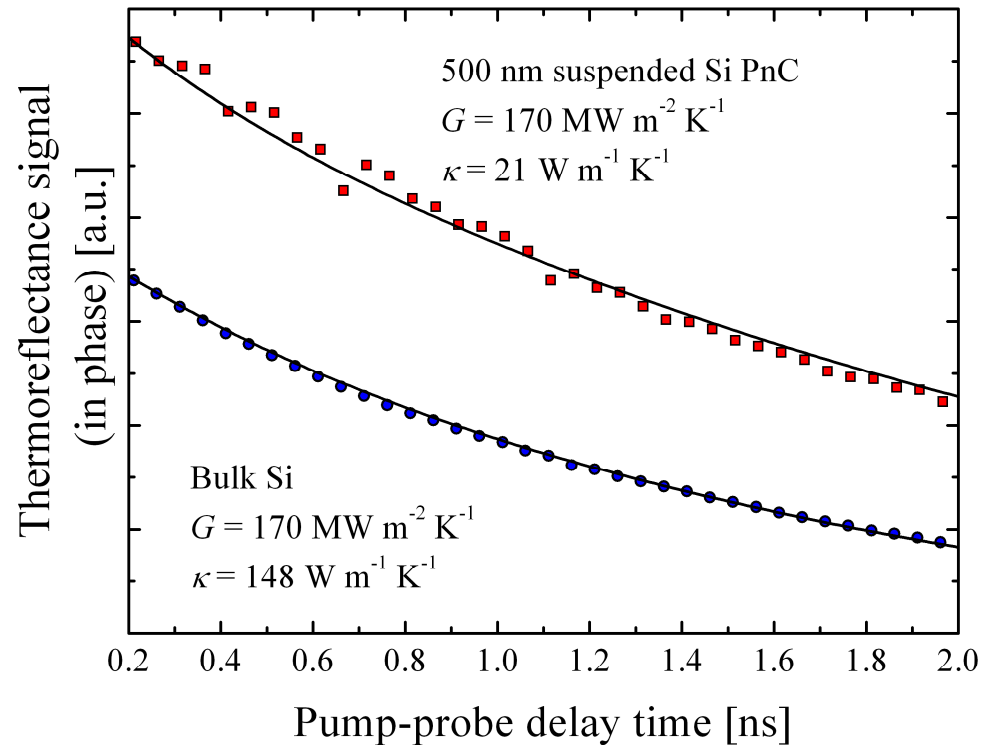
Measurement of κ of the PnC

Thermal conductivity measurements on a 500 nm thick
suspended PnC structure

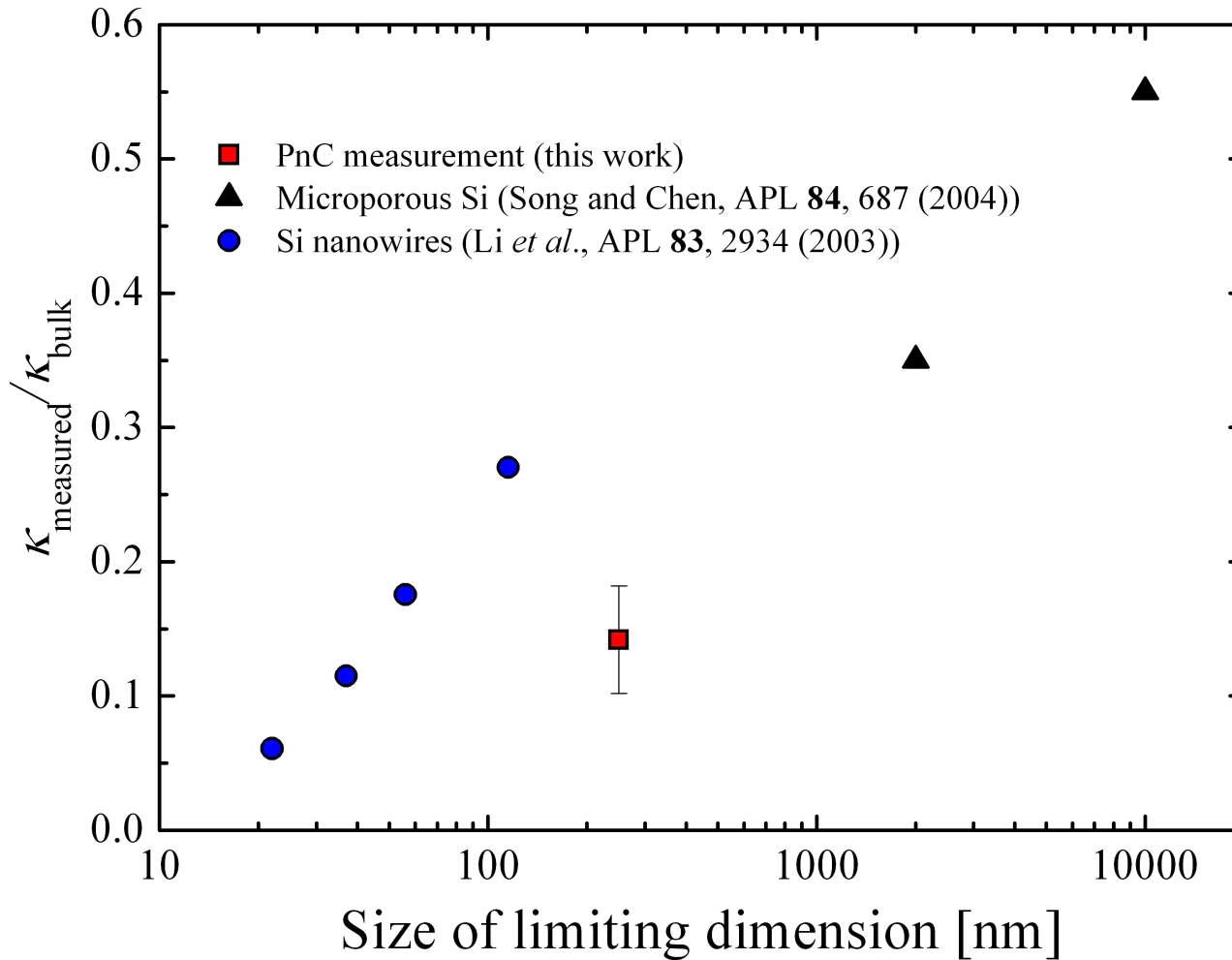


$$\frac{r}{a} = 0.32 \quad a = 700 \text{ nm}$$

$$L = 250 \text{ nm}$$



Phonon engineering of κ in nanostructures





Conclusions

Goal: Study phonon scattering and thermal conductivity reduction in PnC

- Reduction in thermal conductivity of microporous, periodic solids due to long wavelength phonon scattering
- PnC further reduces thermal conductivity through wave interference of phonon modes in periodic lattice
- TDTR experiments show reduced thermal conductivity in PnC beyond that from boundary scattering considerations

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