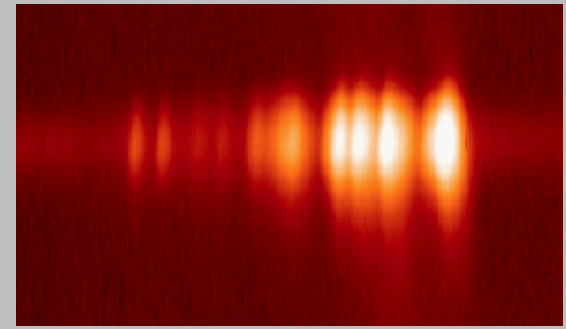
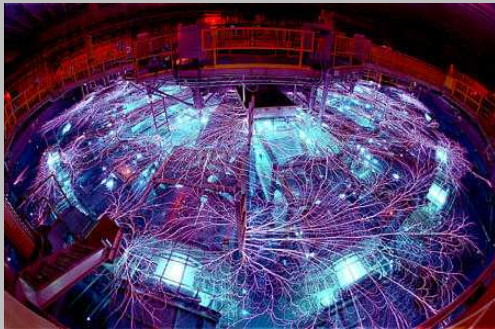


Exceptional service in the national interest



X-ray Thomson Scattering (XRTS) Development for Z

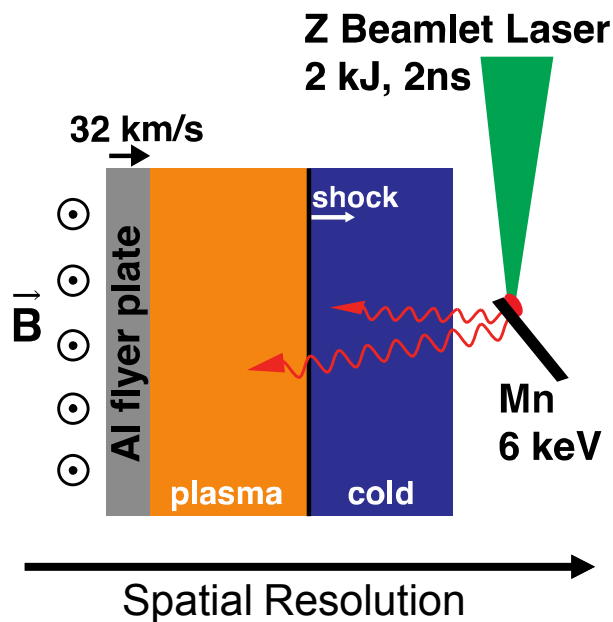
E.C. Harding¹, T. Ao¹, J.E. Bailey¹, D.B. Sinars¹, G.A. Rochau¹, S.B. Hansen¹, M.P. Desjarlais¹, L.P. Mix¹, M. Geissel¹, I.C. Smith¹, R.W. Lemke¹, G. Gregori²

¹Sandia National Laboratories, Albuquerque, NM

²University of Oxford, Oxford, UK

Why do XRTS on Z?

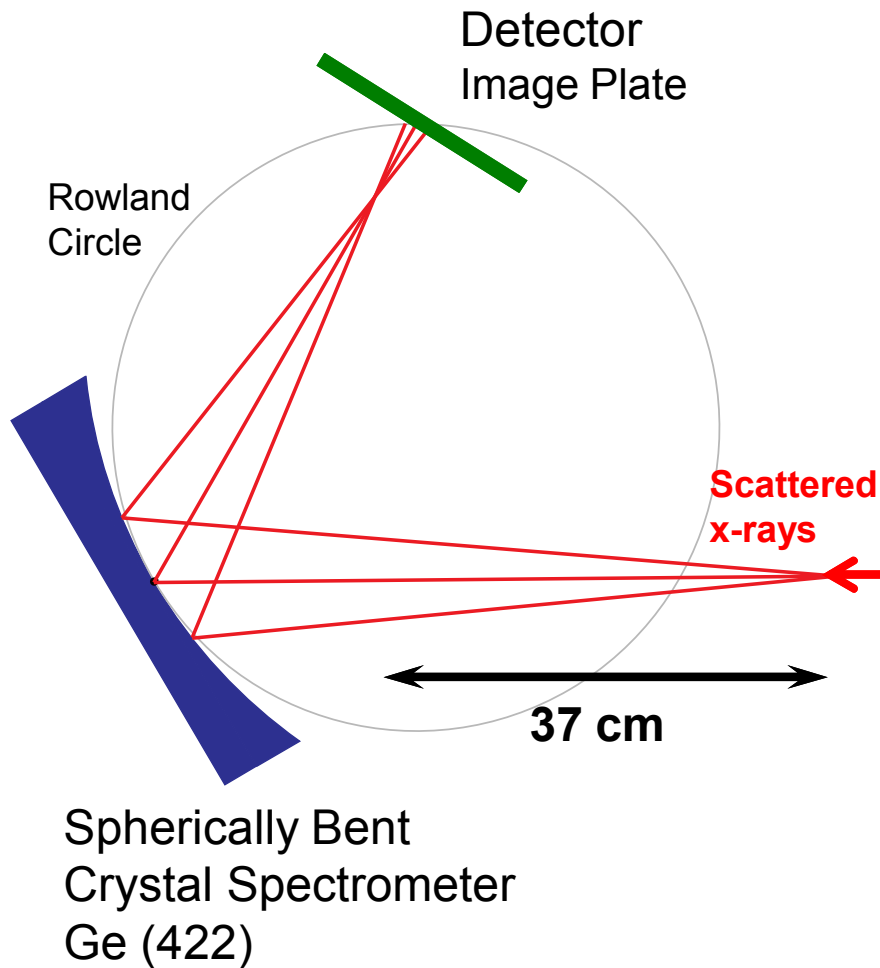
- Need to constrain EOS models with T_e measurements.
- Flyer plate driven targets offer some advantages:
 - Big targets (many mm sized)
 - Long lived steady shocks (90 ns in foam experiment)
 - No x-ray or hot-electron preheat



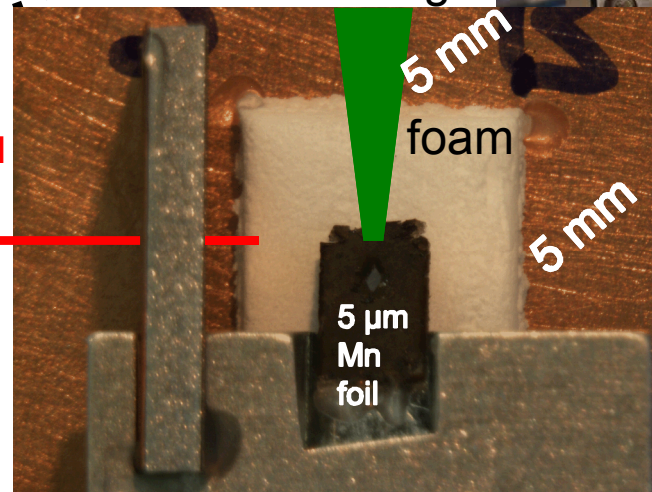
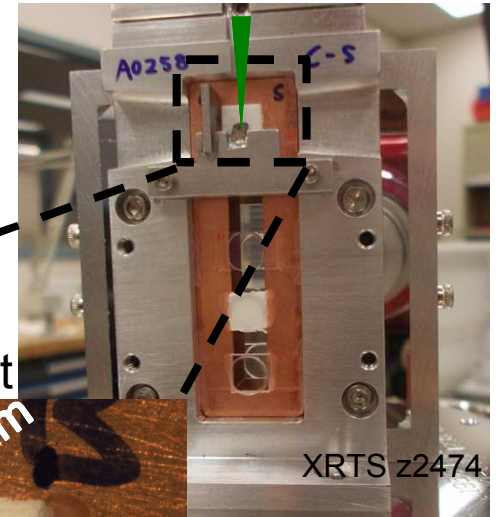
Warm Dense Matter on Z

	CH ₂ Foam	⁶ LiD	Be
T_e	7 eV	4 eV	3 eV
ρ	0.5 g/cc	2.55 g/cc	4.9 g/cc
P	1.1 Mbar	5 Mbar	9 Mbar
Zbar	0.6	0.6	2.0

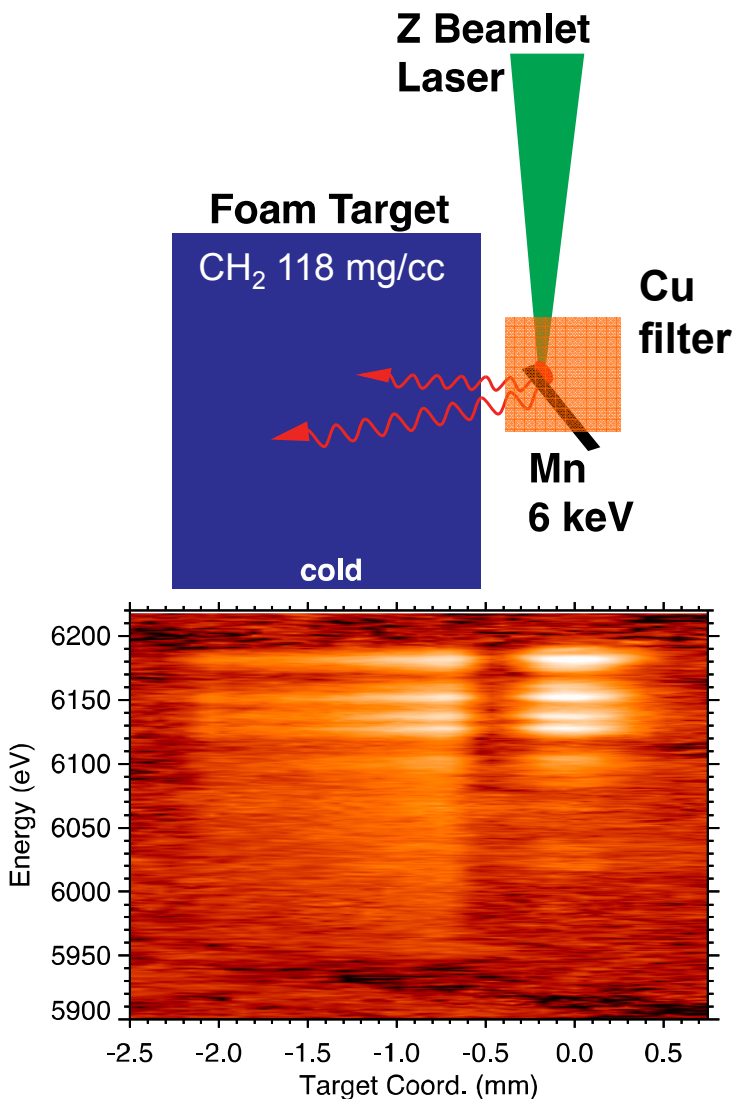
A new spectrometer and target were developed to enable space-resolved XRTS measurements.



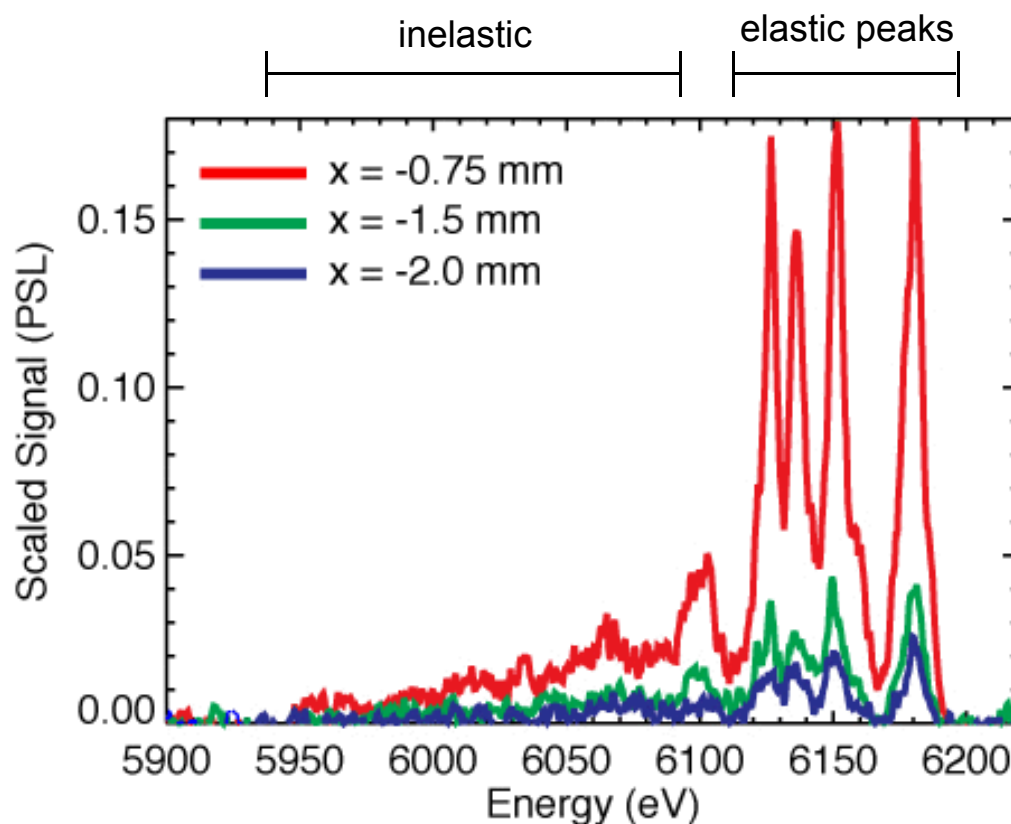
Z target assembly



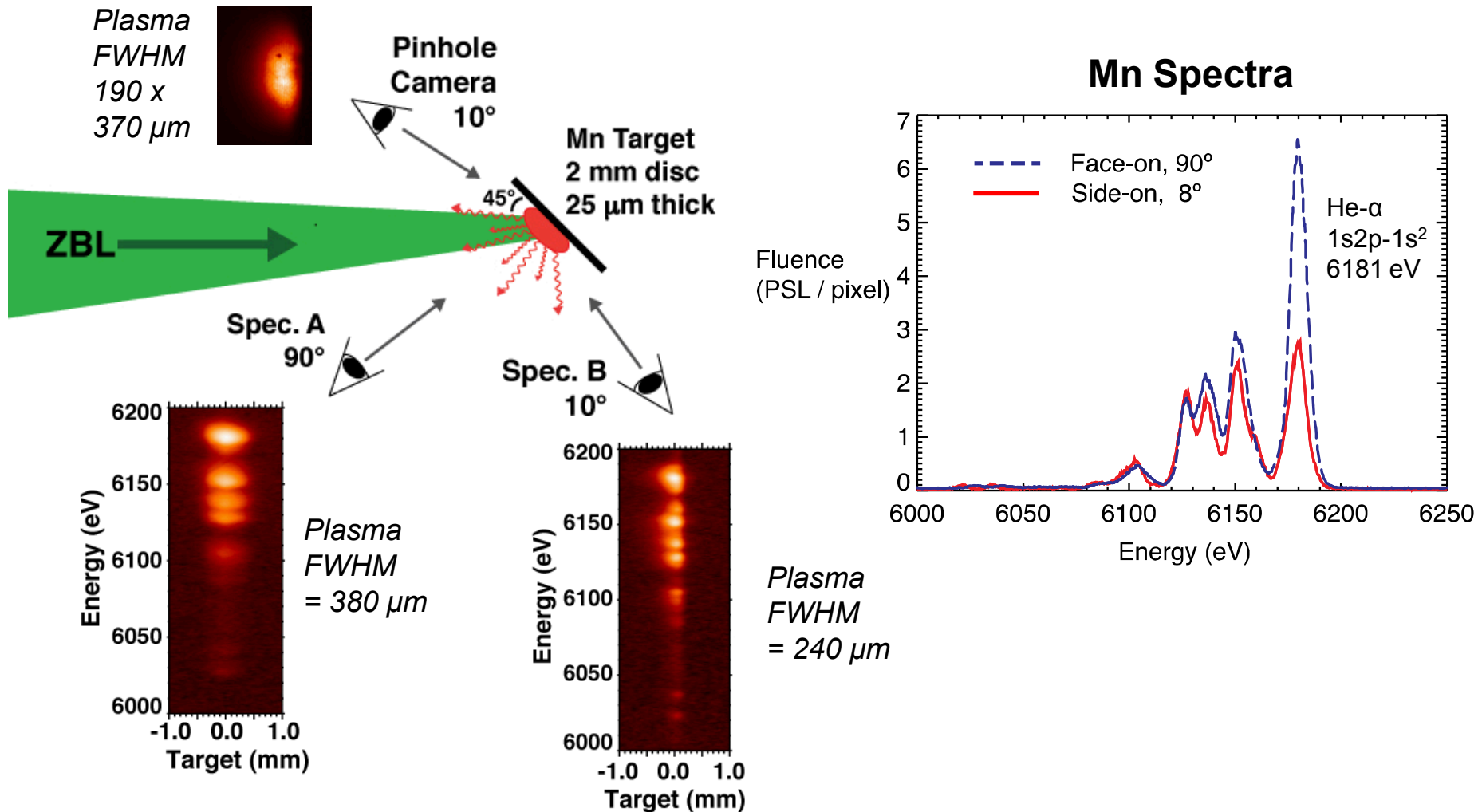
The first full integrated Z shot successfully observed space-resolved, x-ray scattering from unshocked foam.



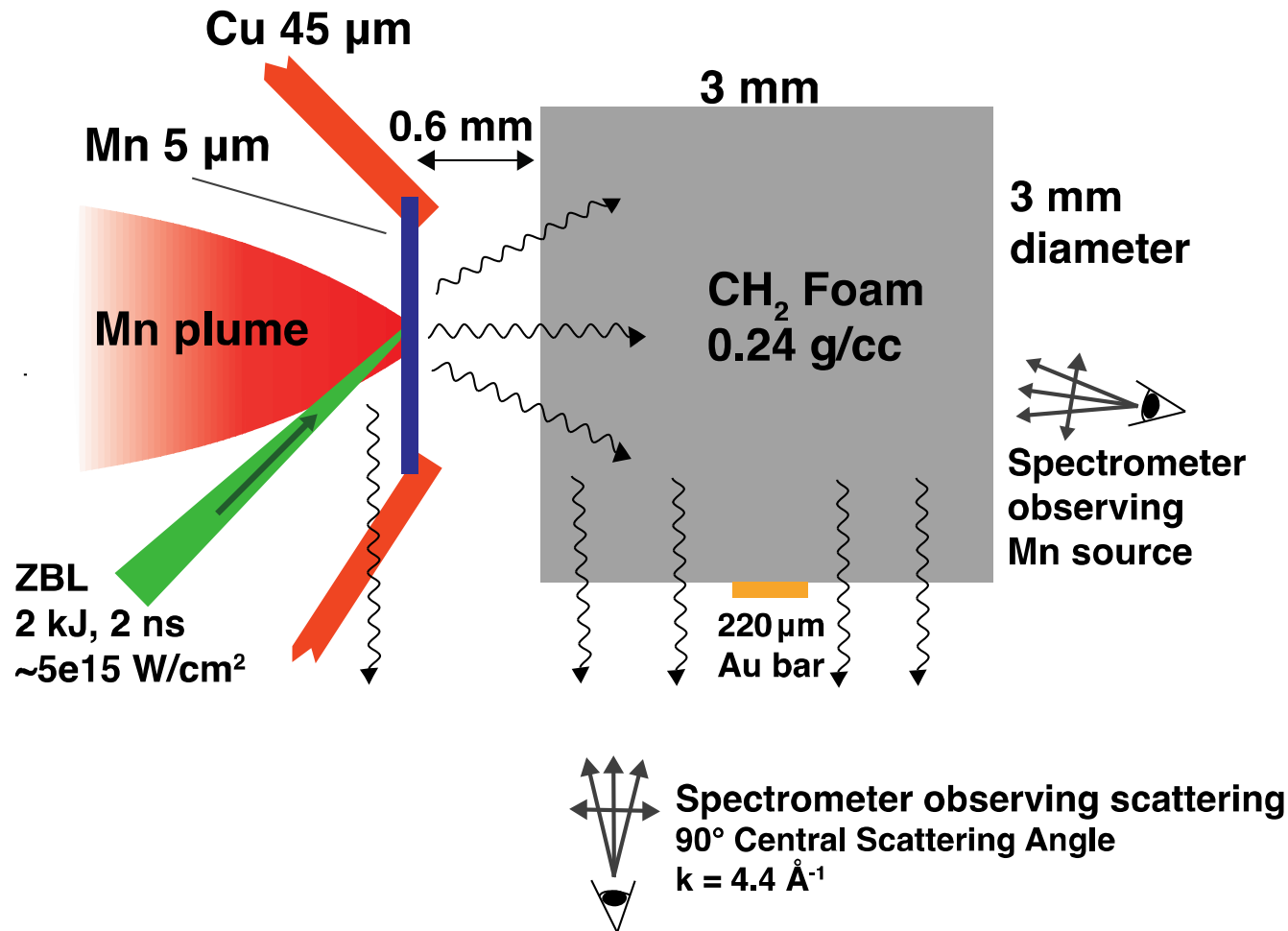
Lineouts show high-resolution elastic and inelastic scattering spectra.



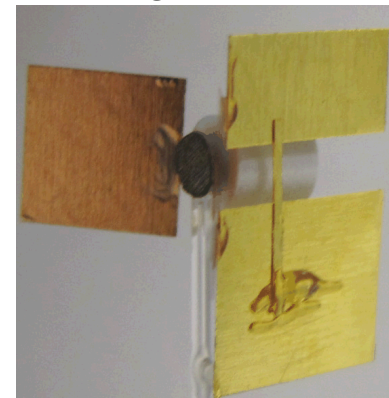
We have observed a variation in the shape and intensity of Dedicated Mn source experiments.



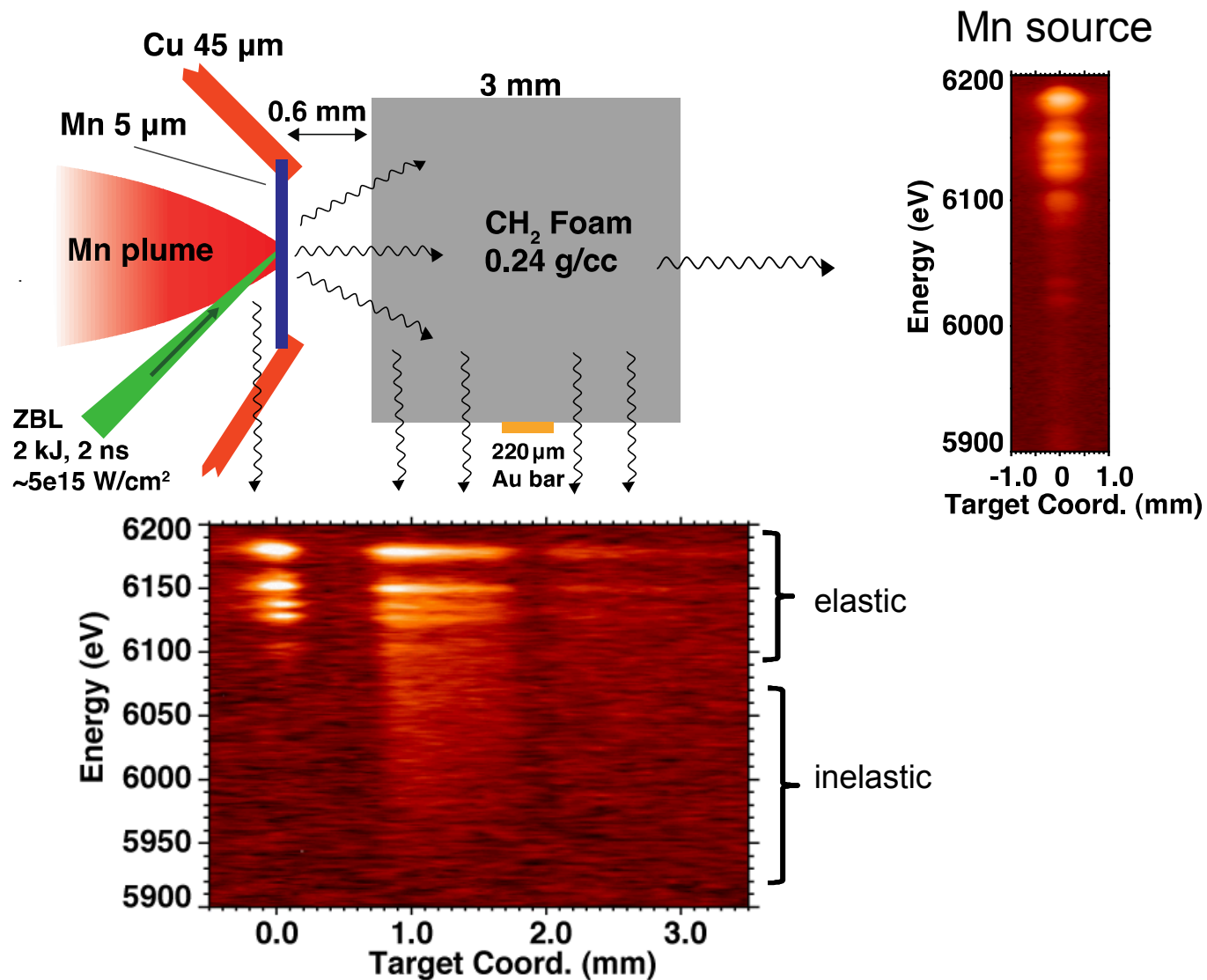
A laser only target was tested with two identical spectrometers that observed the scattering and the source as witnessed by the foam.



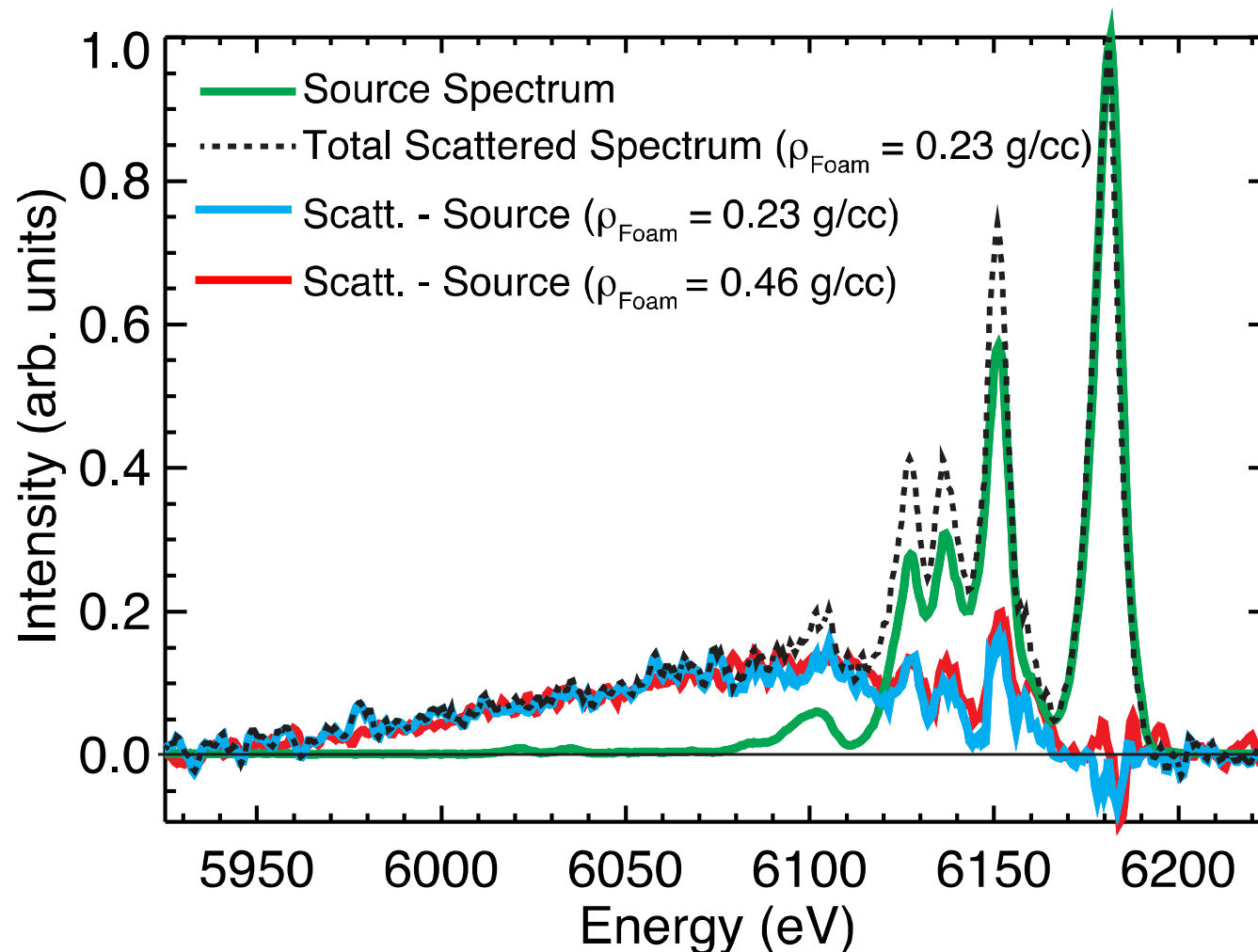
Target Photo



Foam experiments observed detailed scattering and source spectra in the same shot.

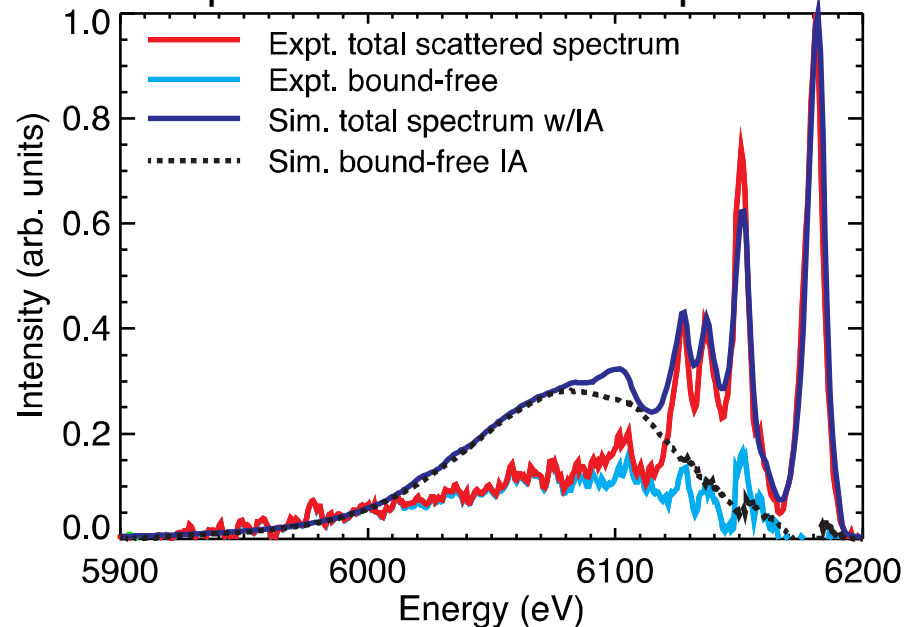


The inelastic scattering is isolated by subtracting the source spectrum from the total scattered spectrum.

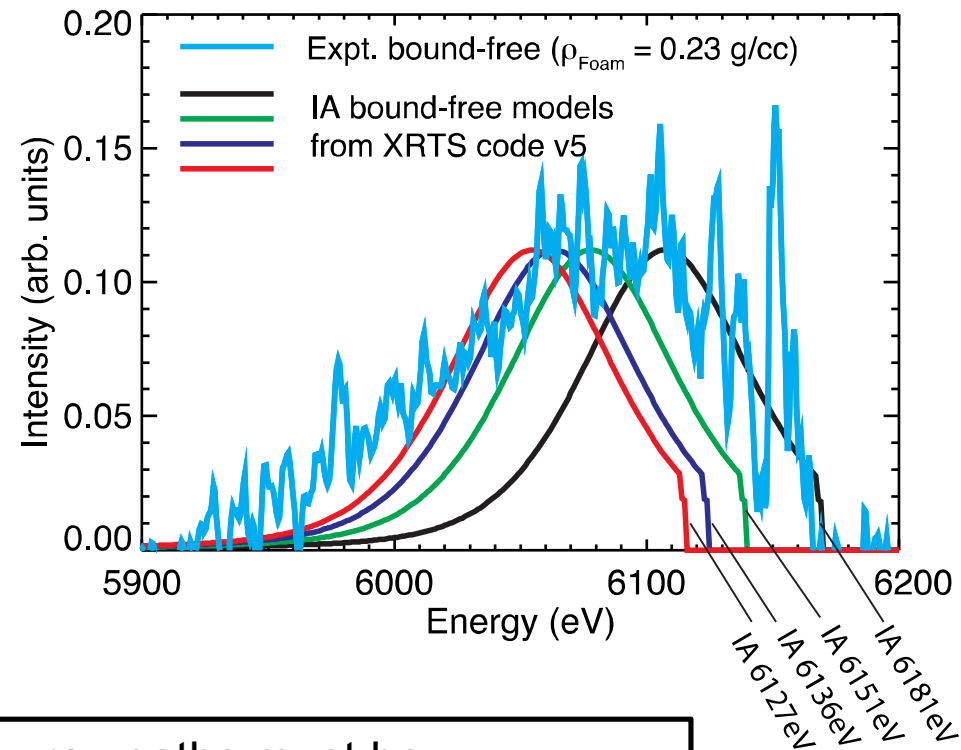


The relative line intensities of the scattered spectrum depend on the binding energies of the Carbon 2p electrons.

Comparison of simulated and experimental scattered spectra.



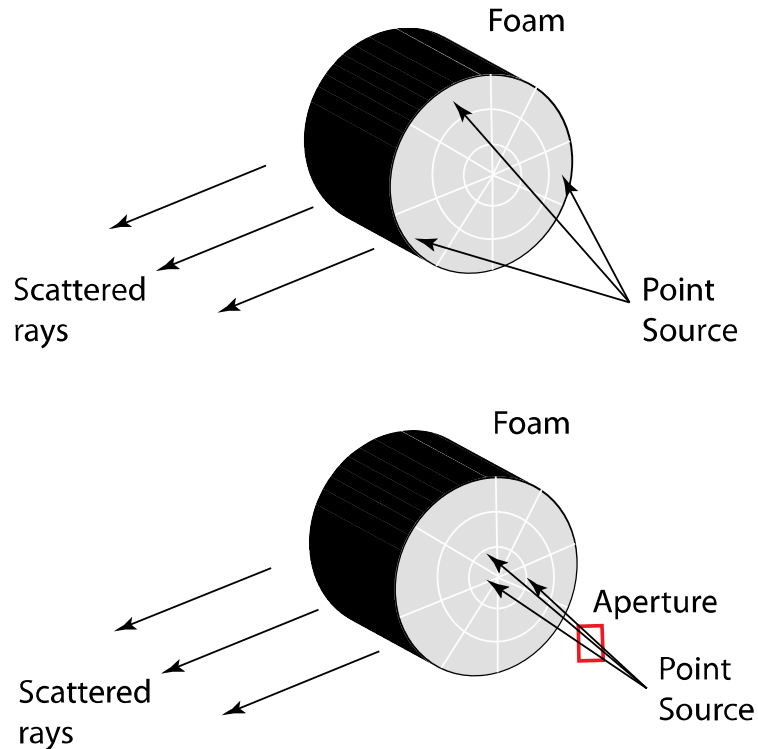
Each spectral line begins to contribute to the bound free spectrum the Carbon 2p² binding energy is exceed.



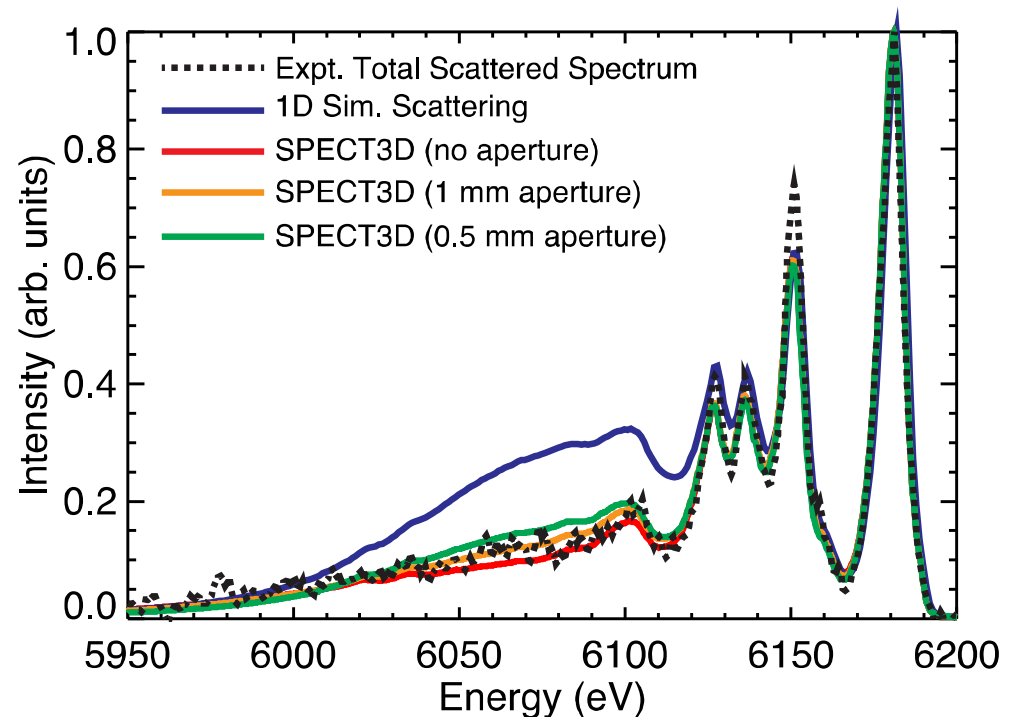
With no input aperture multiple x-ray paths must be considered

The XRTS code was added to SPECT3D in order to model scattering from 3D, non-homogeneous plasmas.

SPECT3D Geometry



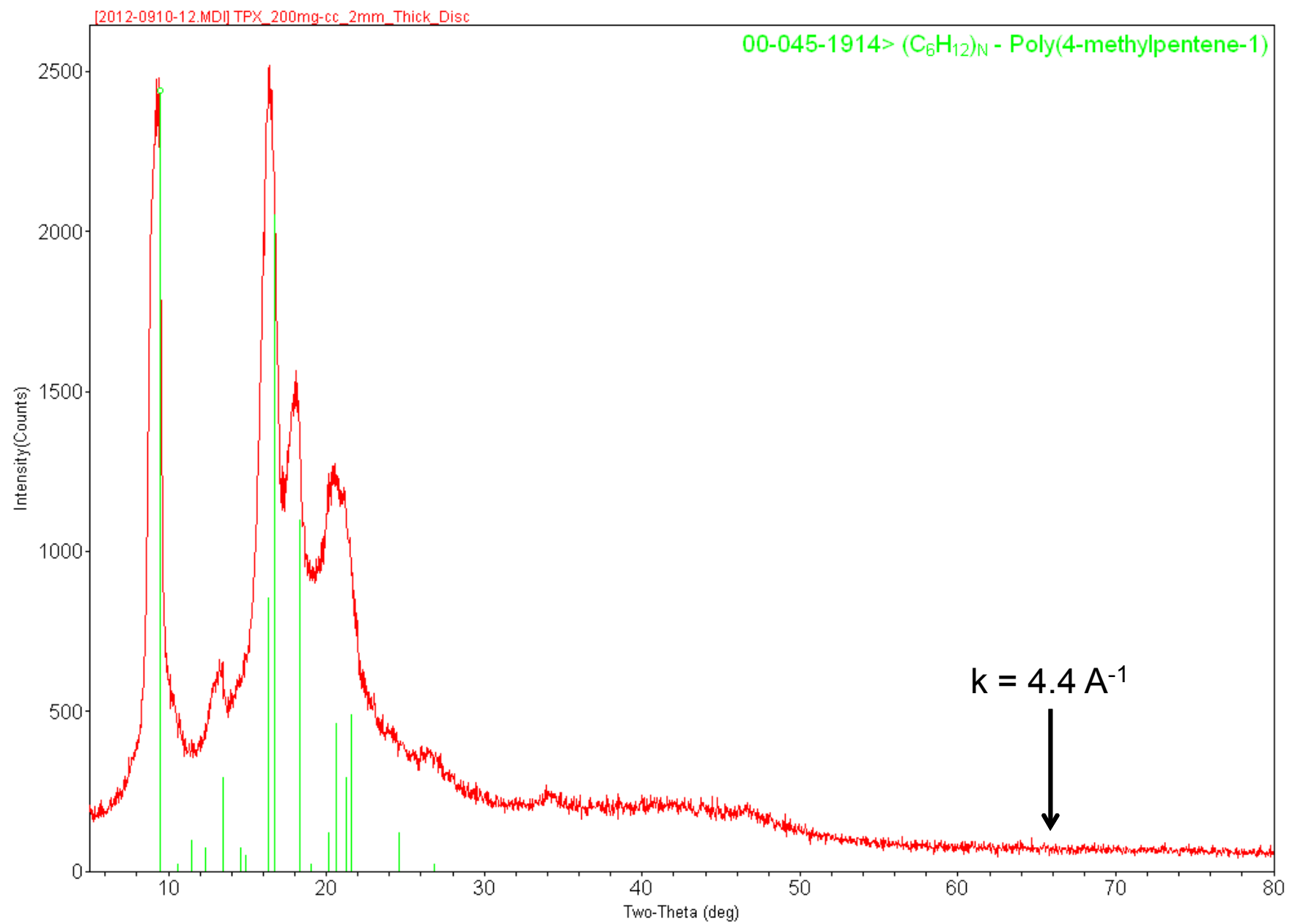
Comparison of SPECT3D simulated scattering and experiment.



SPECT3D simulated scattering fits data well with $S_{ij} = 1.5$. However, the relative line intensities are not reproduced.

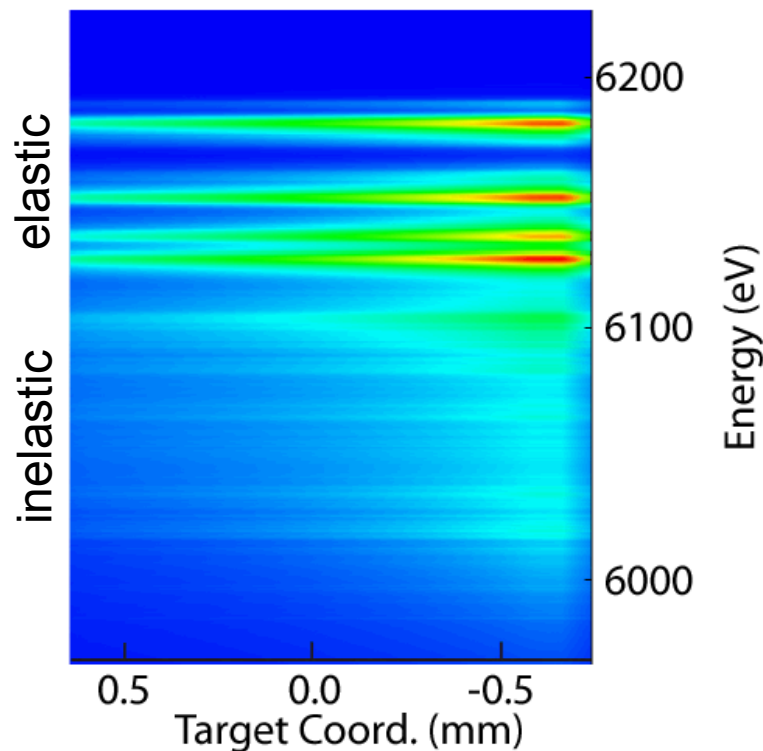
Summary

Sample Name: **TPX 200mg/cc 2mm Thick Disc** identified as Poly(4-methylpentene-1)

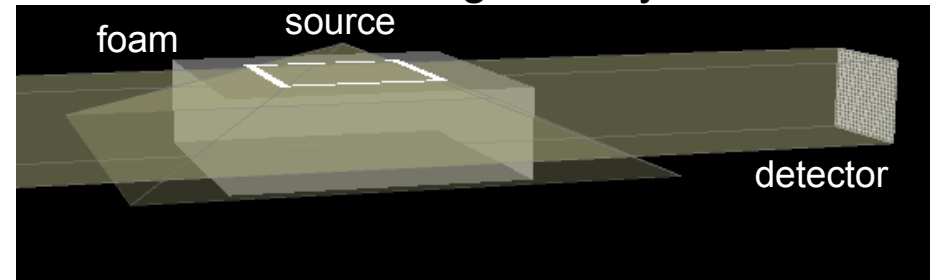


Gregori's XRTS code was added to SPECT3D in order to model scattering from 3D, non-homogenous plasmas¹.

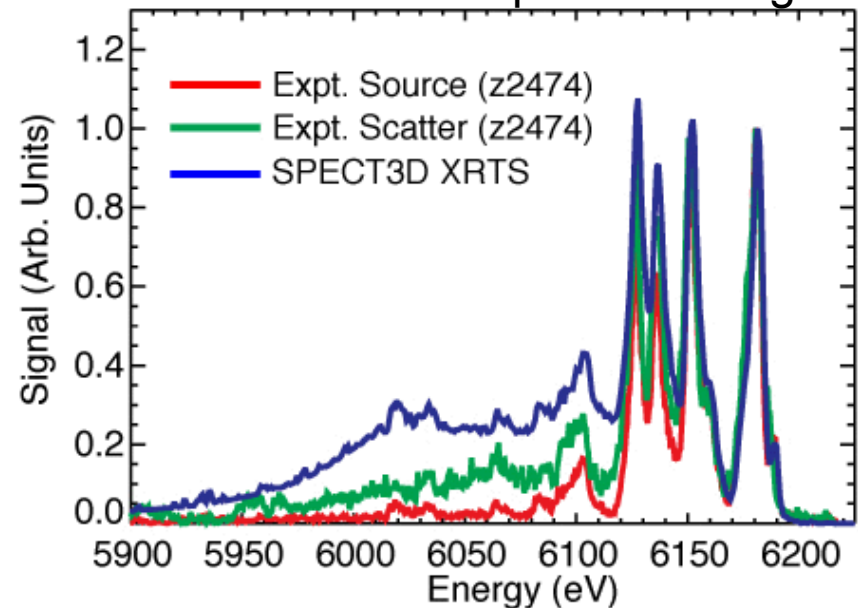
Simulated Scattering Spectra



SPECT3D geometry



Simulated vs. Expt. Scattering



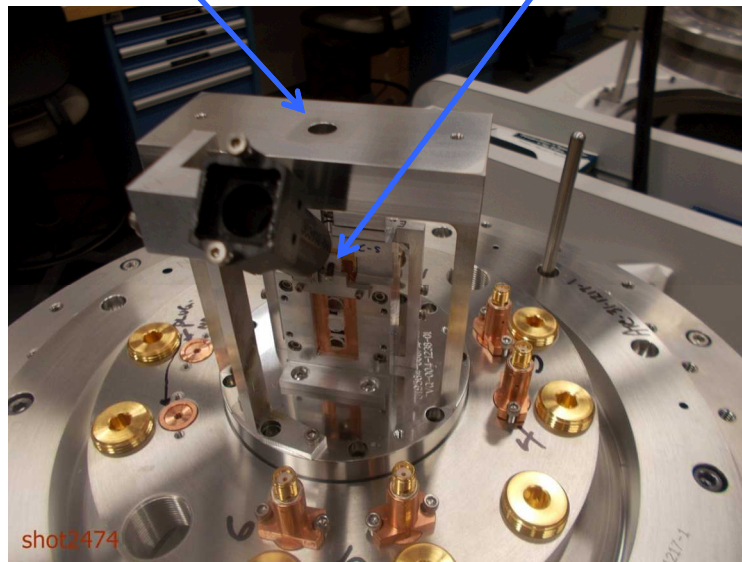
¹Igor Golovkin *et. al.*, JQSRT (*submitted*)

Limiting apertures were designed to protect both the crystal and laser focusing optics.

- Engineering design of apertures informed by NASA micro-meteorite data from impacts on glass

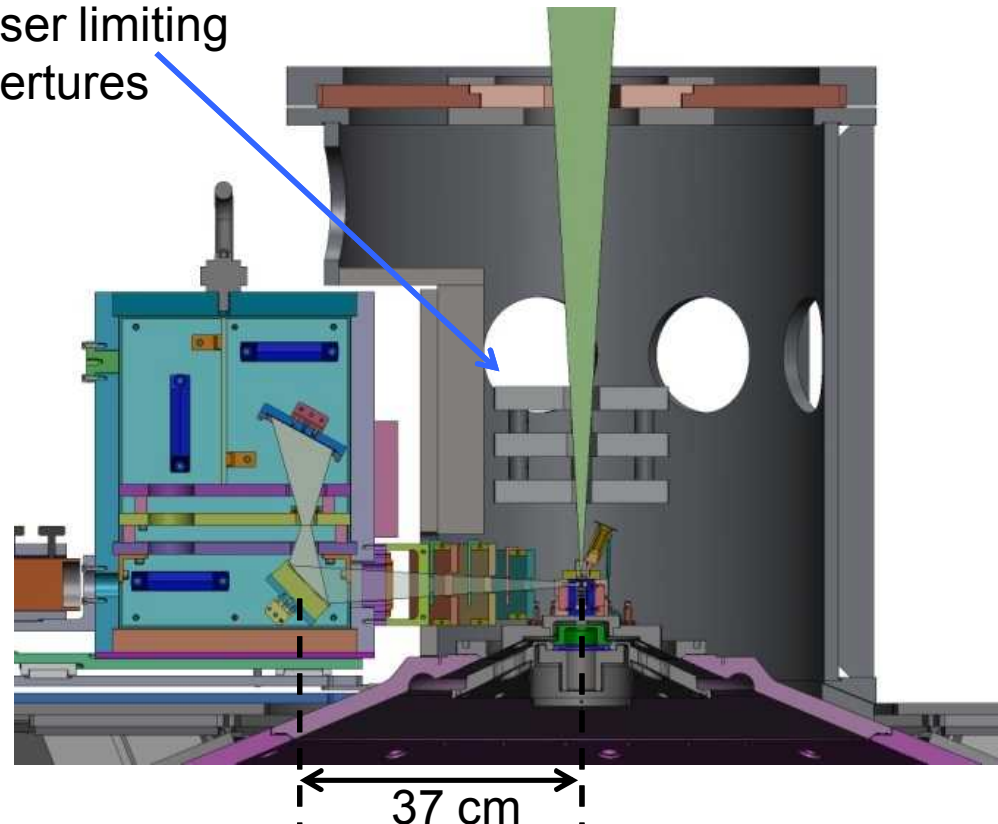
Final laser
aperture
(5.8 mm)

XRTS
target

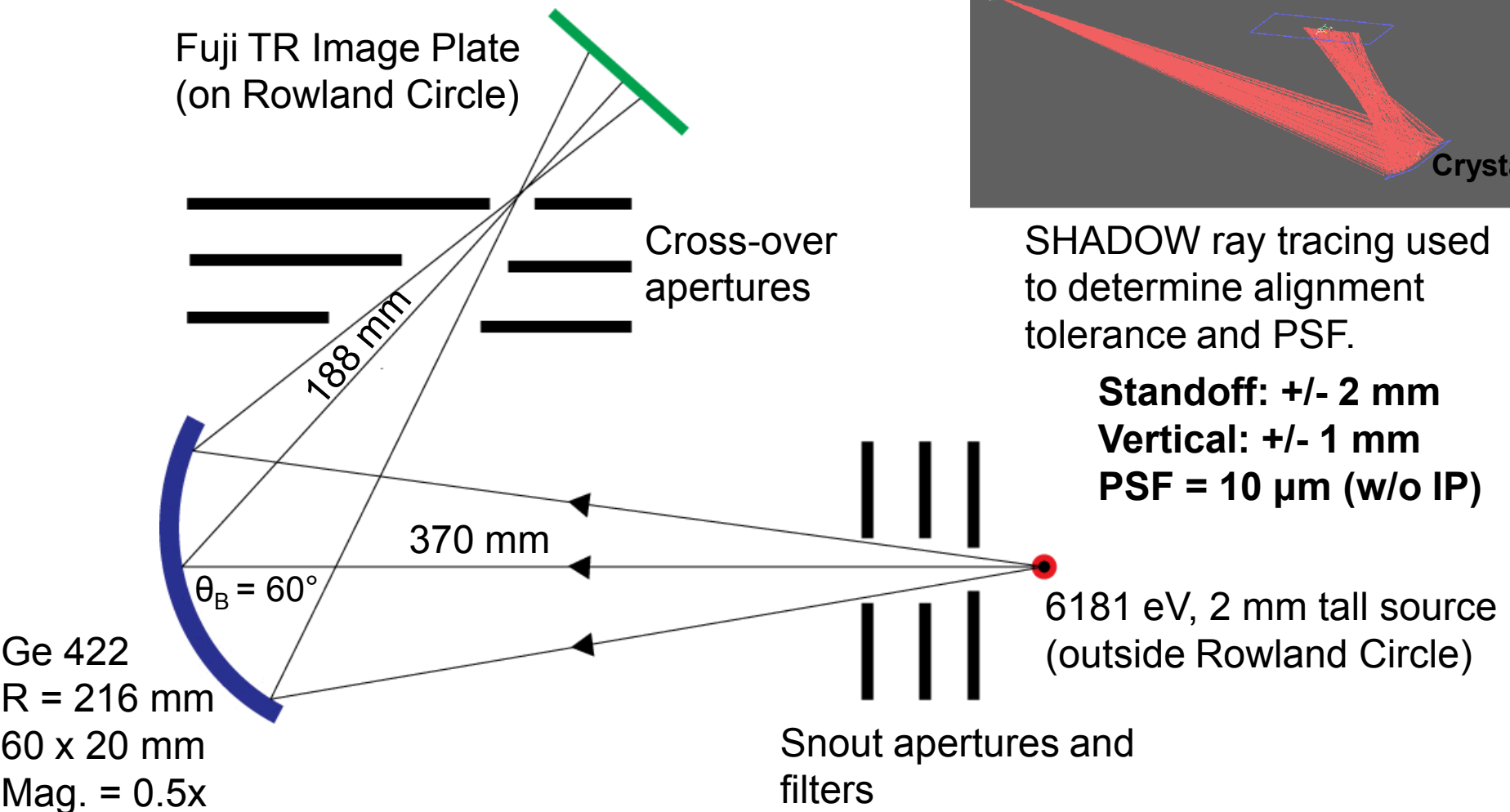


Laser limiting
apertures

ZBL laser



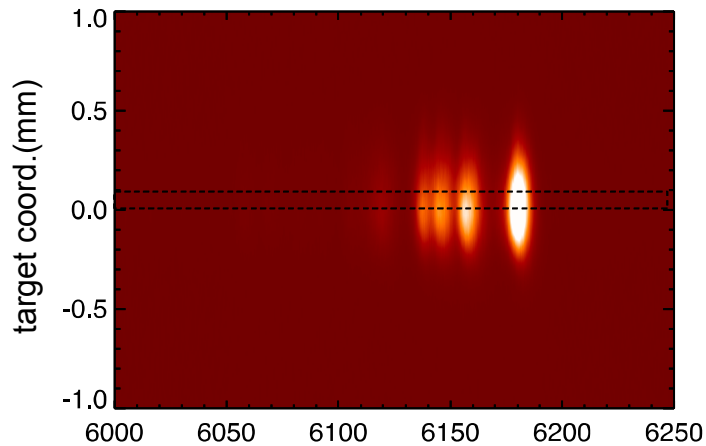
The spectrometer point design uses a spherical Ge 422 crystal ($R = 216$ mm) and has a spatial magnification = 0.5x.



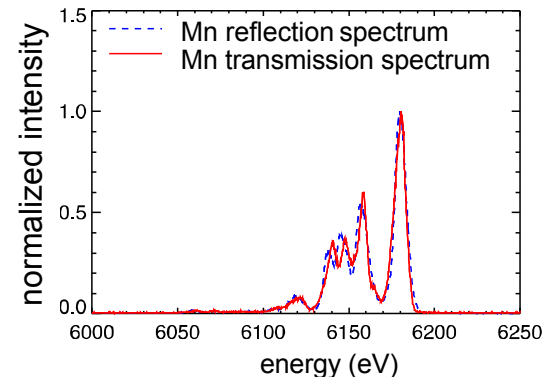
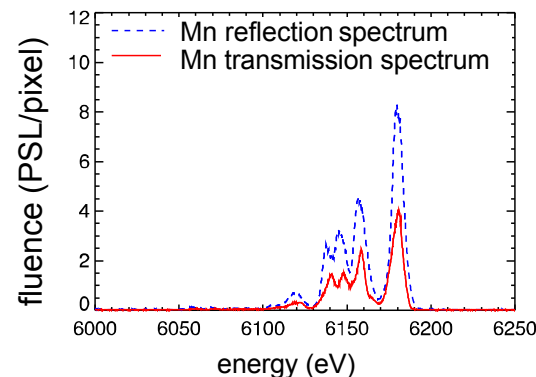
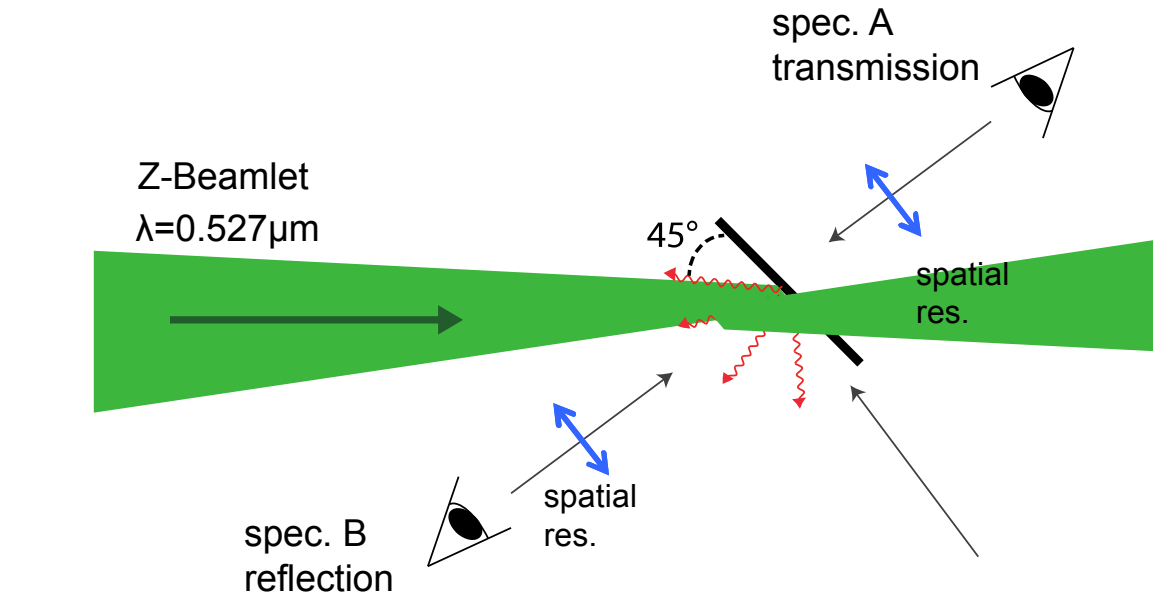
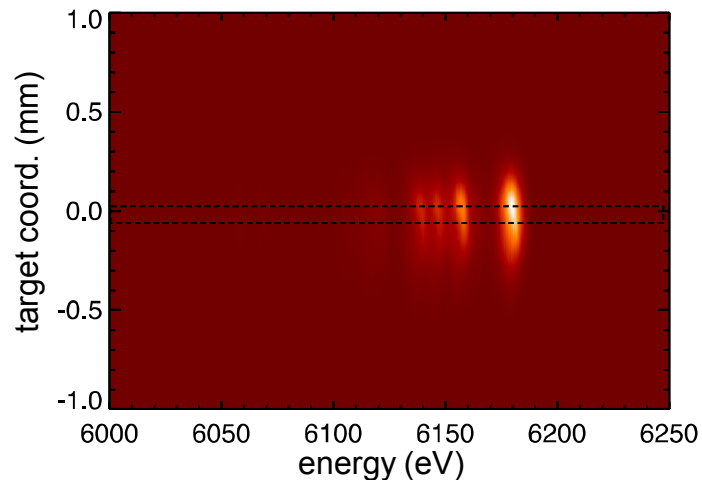
Reflection and transmission spectra

- Spectra similar but transmitted x-rays attenuated by 50% compared to reflected x-rays

Mn reflection image

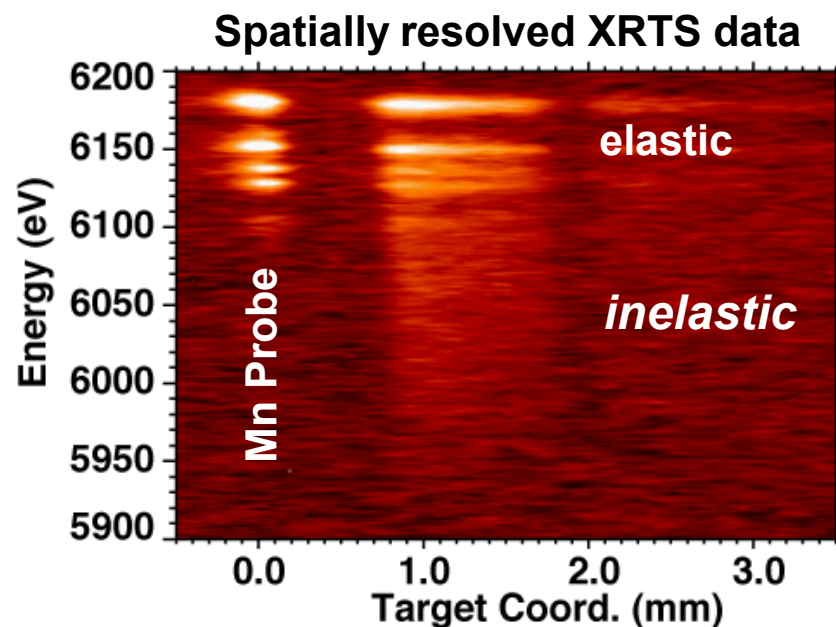
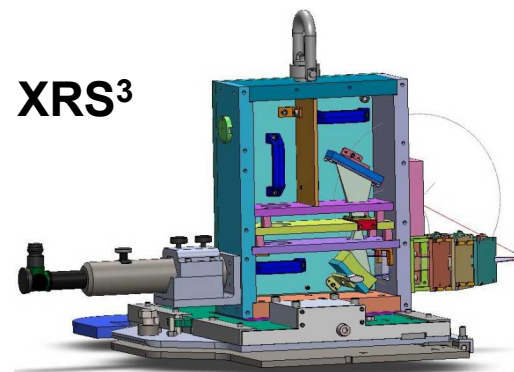


Mn transmission image

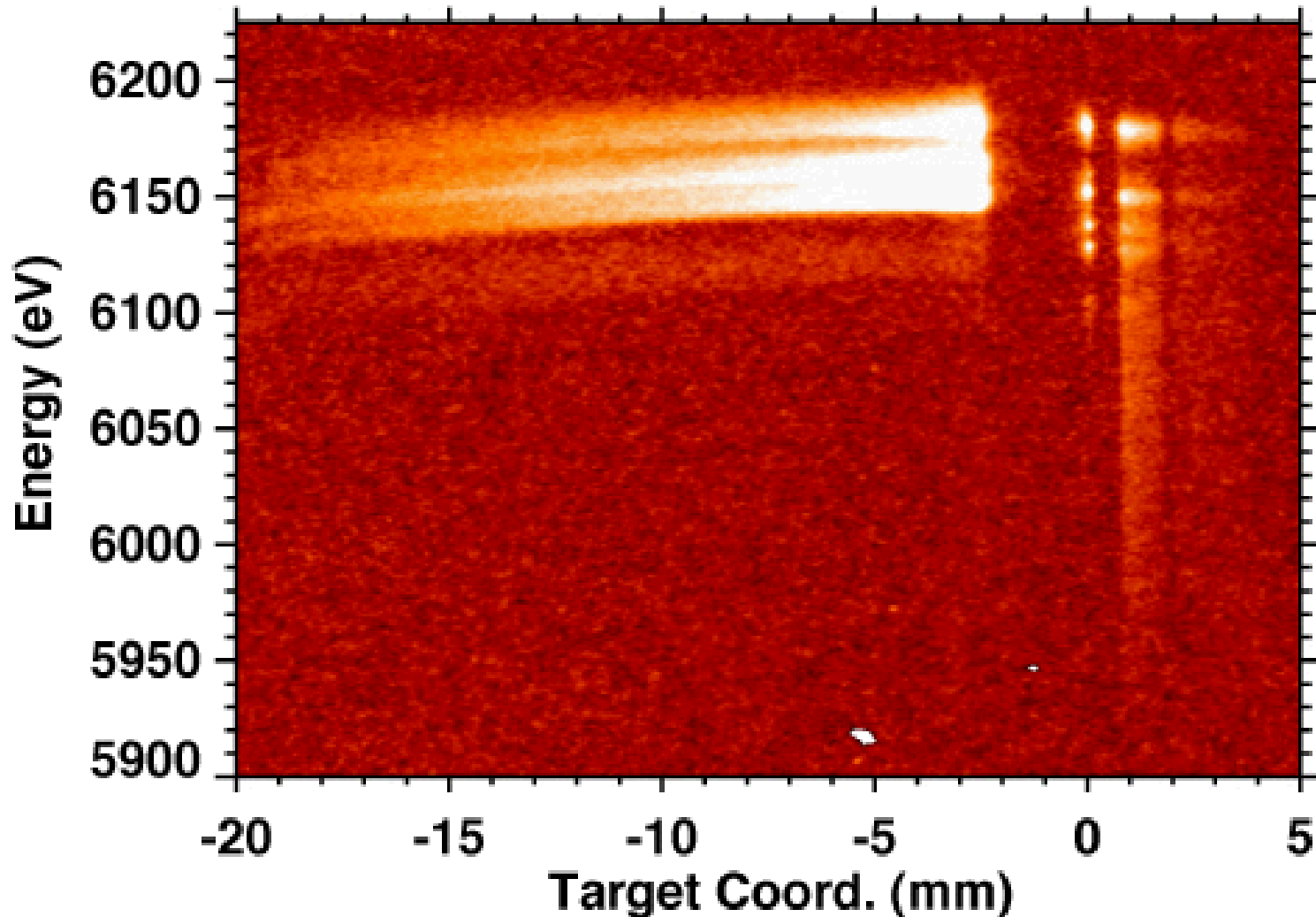


Summary: We have used a spherical spectrometer to record spatially resolved XRTS spectra.

- Goal: XRTS with spatial resolution on the Z machine
- We are using spherically bent Ge 422 crystals
- We are characterizing the probe plasma to improve its effectiveness for XRTS



Full Scattering Data Image



Debris damage continues to surprise us, but mitigation has been successful.

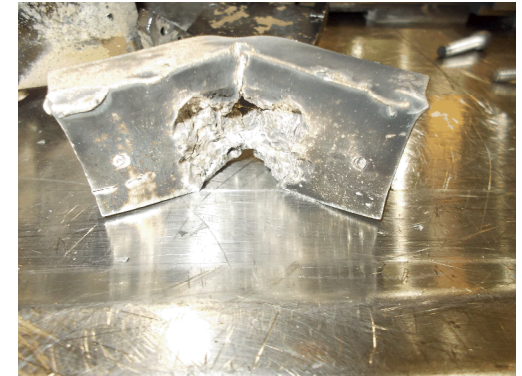
Inside spectrometer after first shot



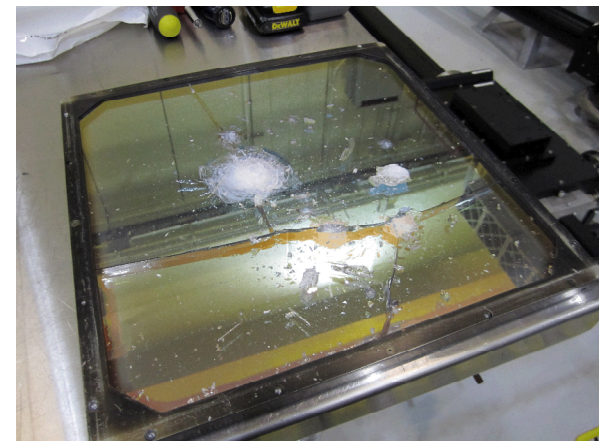
Crystal does not survive

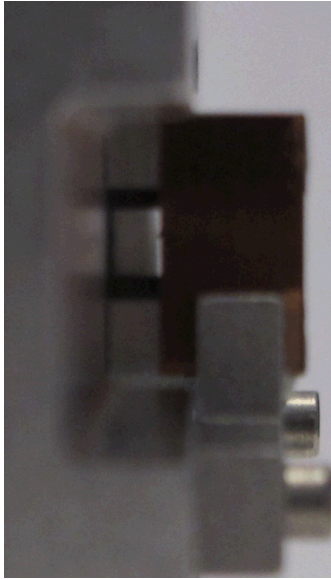


Aperture block successfully mitigated axial debris

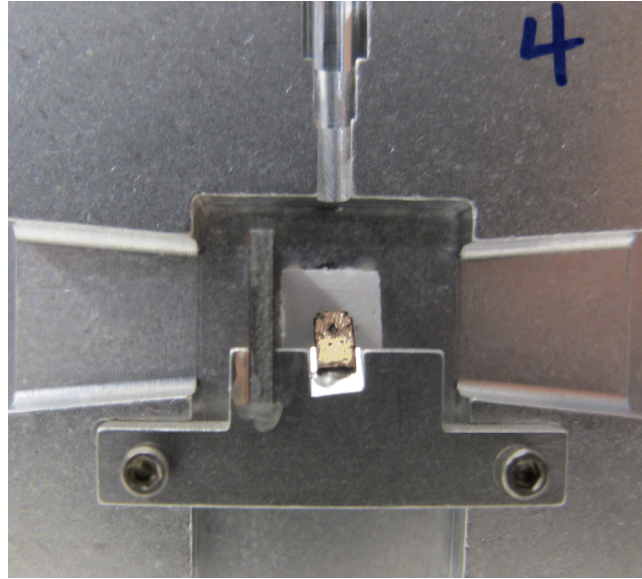


Laser debris shield damaged but not compromised





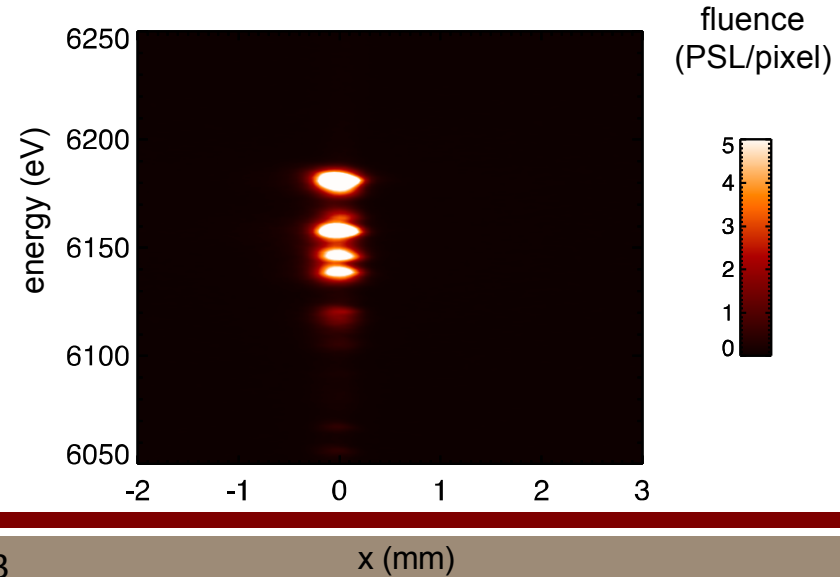
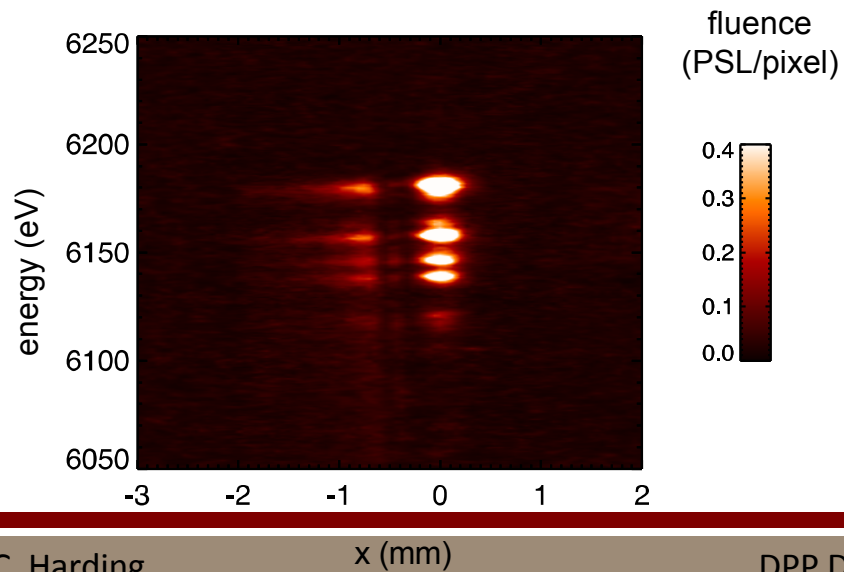
Spec. A (90° view)



ZBL (0° view)

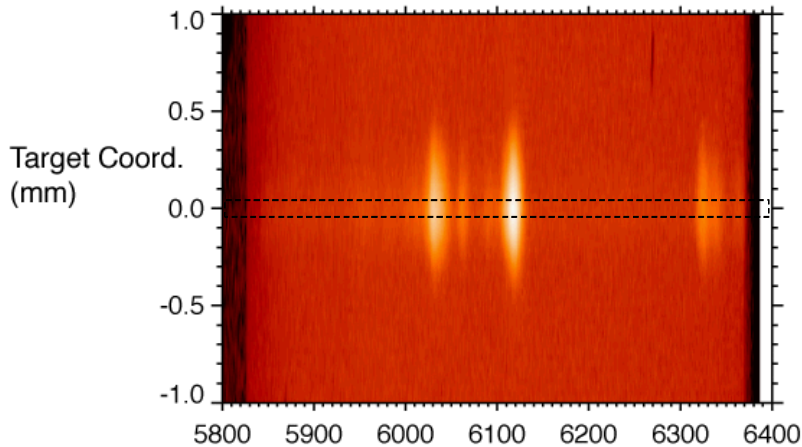


Spec. B (-90° view)

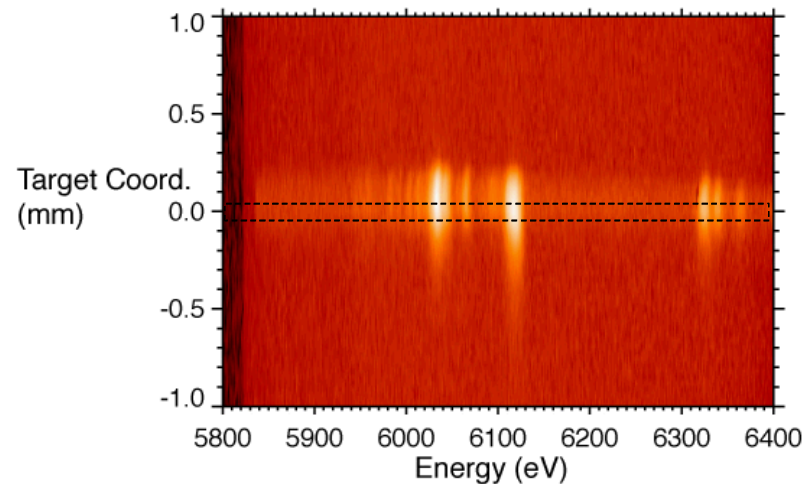


The V He- β spectra also show differences between the two viewing angles.

V Face-on Image

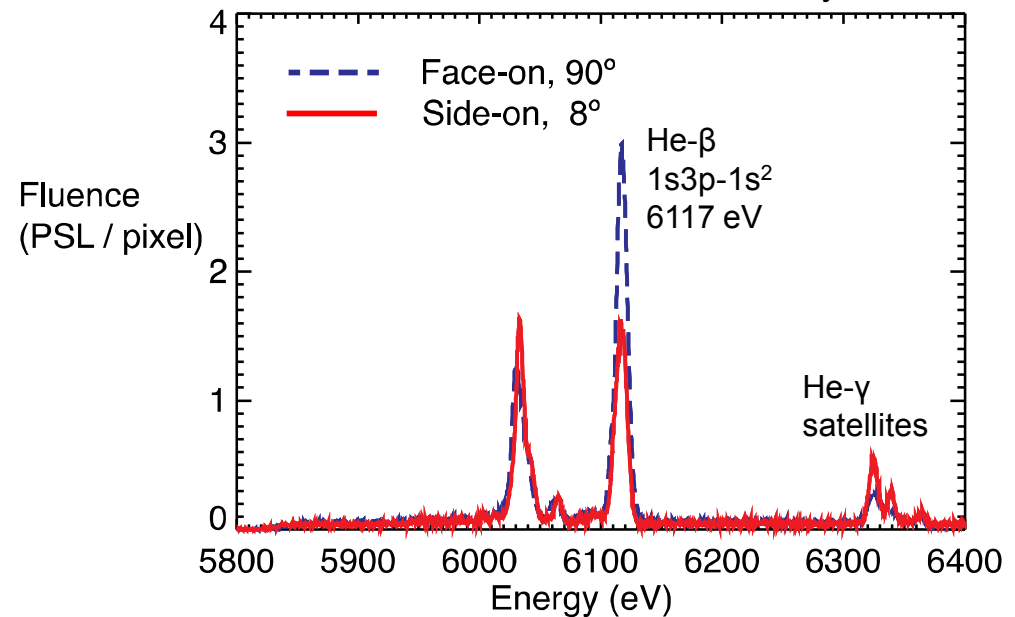


V Side-on Image



V Spectra

Lineout width = 0.1 mm, centered at y = 0 mm



Spatial Information from vertical lineouts

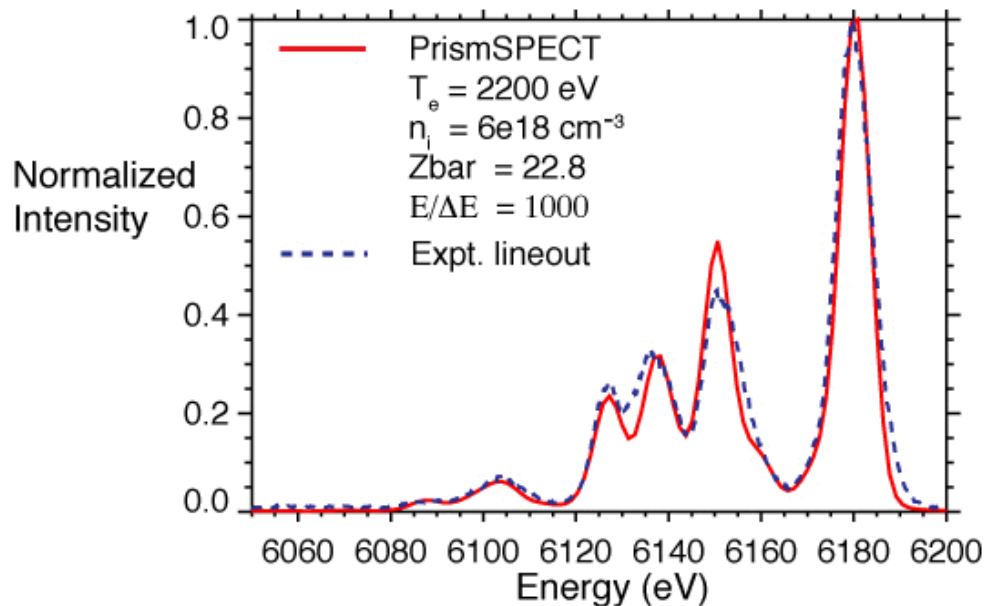
Face-on FWHM: 310 μm @ 6117 eV

Side-on FWHM: 230 μm @ 6117 eV

PrismSPECT NLTE simulations of Mn at $T_e = 2.2$ keV and $n_i = 6 \times 10^{18} \text{ cm}^{-3}$ show good agreement with the Face-on spectrum.

Mn PrismSPECT model and Experiment

Intensity normalized to He- α Res. line.

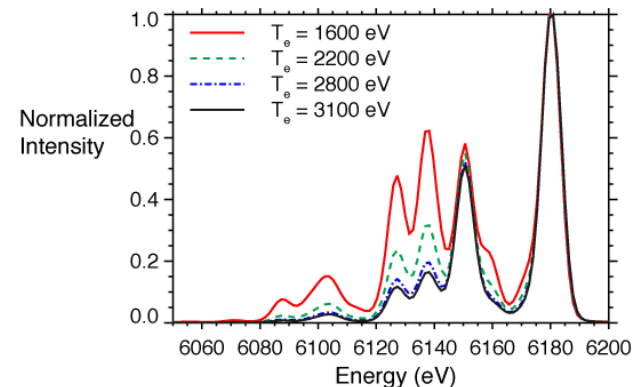


PrismSPECT indicates that our spectrometer has $E/\Delta E = 1000$.

PrismSPECT gives $n_e = 1.4e20 \text{ cm}^{-3}$ which is 30x lower than $n_{\text{cr}} (4e21 \text{ cm}^{-3})$.

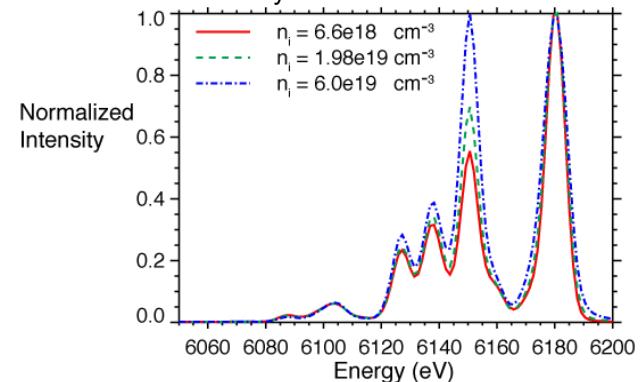
Mn Li-sat. to He- α ratio provides T_e measurement

Plasma slab 200 μm thick, $n_i = 6e18 \text{ cm}^{-3}$
 Intensity normalized to He- α Res. line.



Mn Inter. Comb. line to He- α ratio provides n_i measurement

Plasma slab 200 μm thick, $T_e = 2.2 \text{ keV}$
 Intensity normalized to He- α Res. line.



Assuming a uniform slab geometry PrismSPECT* simulations best fit the experimental spectrum for $T_e = 2.2$ keV and $n_i = 6e18$ cm⁻³.

Mn PrismSPECT Model

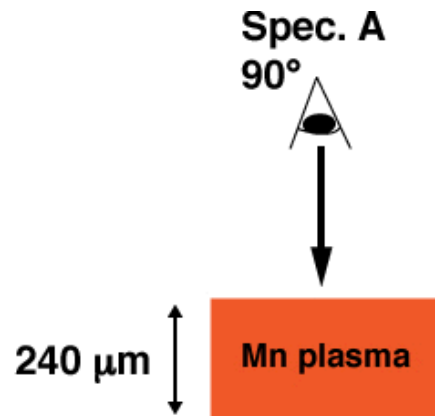
nLTE

1D Slab

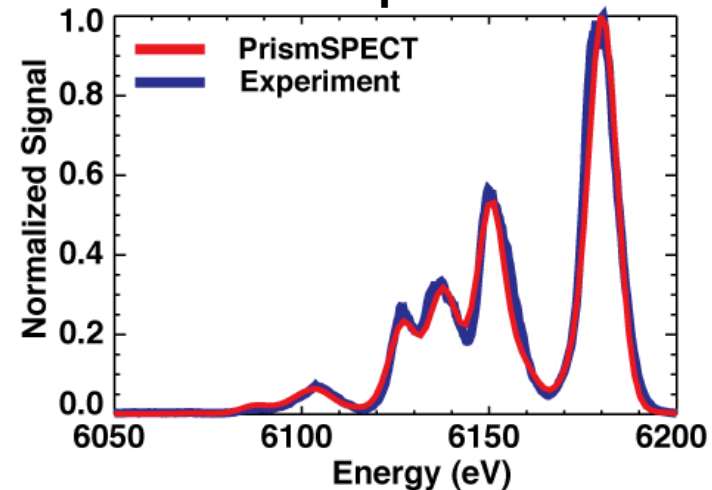
240 μ m thick

$T_e = 2.2$ keV

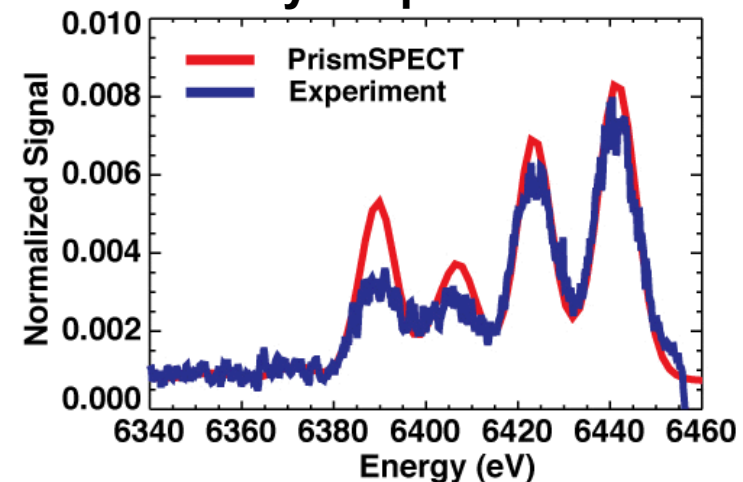
$n_i = 6e18$ cm⁻³



He- α Spectrum

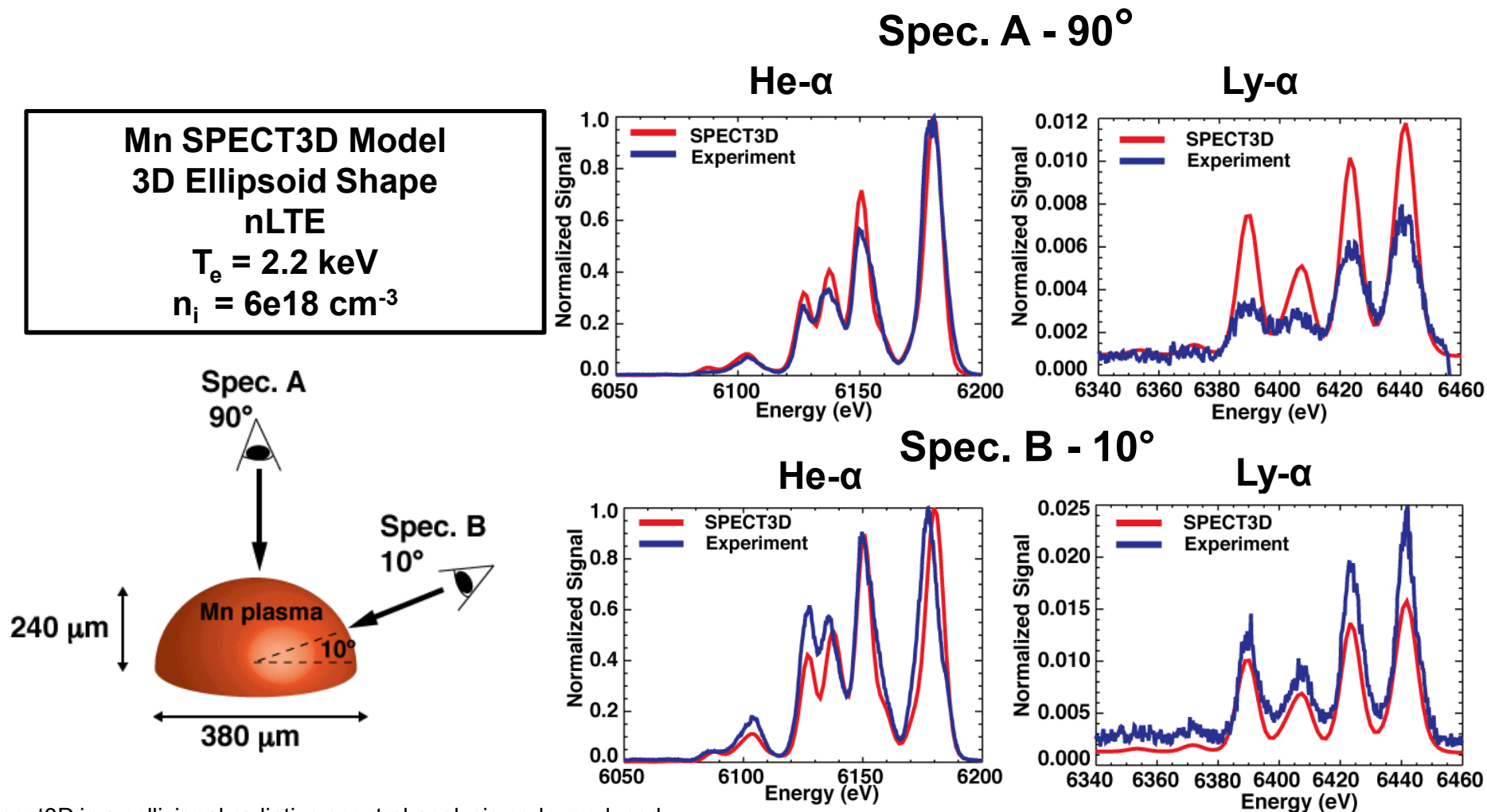


Ly- α Spectrum



*PrismSPECT is a collisional-radiative spectral analysis code produced by Prism Computational Sciences, Inc.

SPECT3D* was used to simulate the spatially resolved spectra from the two spectrometer viewing angles.



*Spect3D is a collisional-radiative spectral analysis code produced by Prism Computational Sciences, Inc.