

Damping Models for Shear Beams with Applications to Spacecraft Wiring Harnesses

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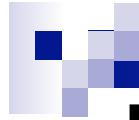
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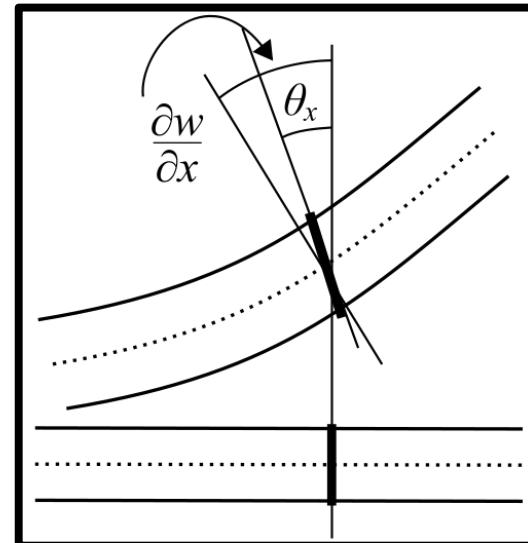
Presenting a new viscous damping model for shear beams

Spacecraft Cable Applications

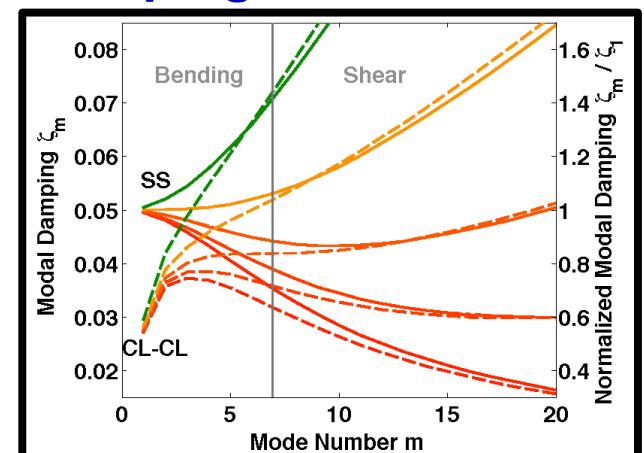


that yields approximately
constant modal damping

Shear Beam & Model Development

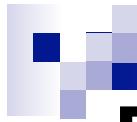


Damping Model Results



Ardelean et al. (2010)

http://upload.wikimedia.org/wikipedia/commons/thumb/e/e4/Plate_theory.svg/500px-Plate_theory.svg.png



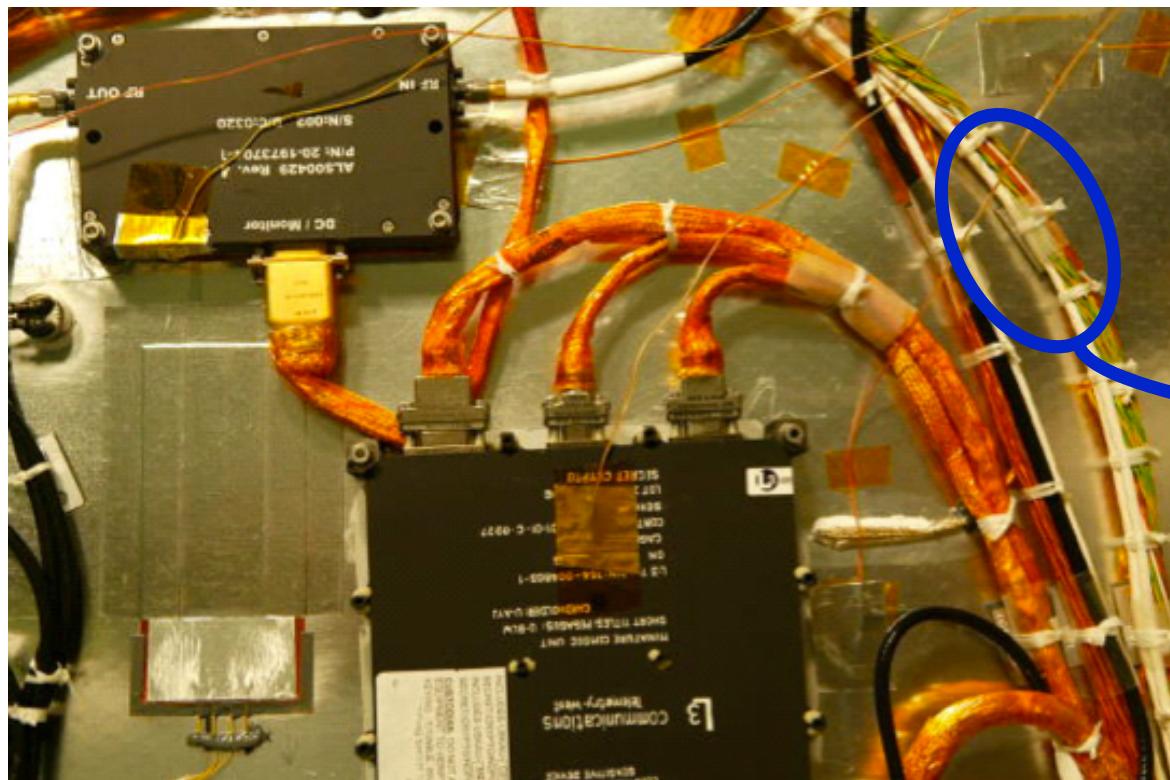
Power & data cables modify spacecraft dynamics, especially at high frequency

- Cabling can account for 30% of spacecraft dry mass!
 - Increasing power / data reqts
 - Decreasing density of structure
- Accurate dynamics model is essential for spacecraft design
 - Launch loads
 - Precision control
- Current models (structure only) over-predict response levels
 - Cables add damping
- Ground testing can augment models, but is incomplete

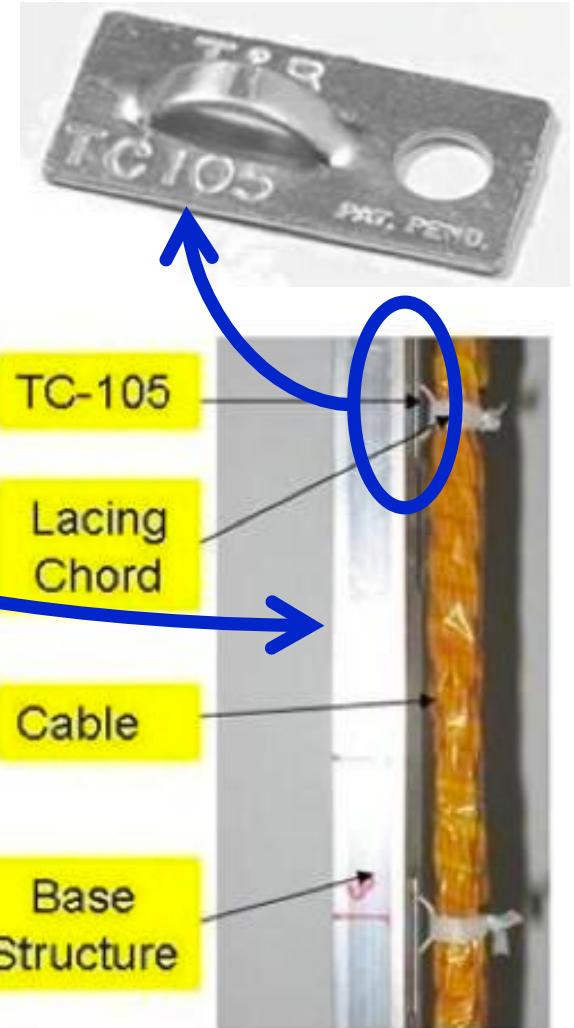


Ardelean et al. (2010)

Spacecraft & cable dynamics are coupled through cable tiedowns

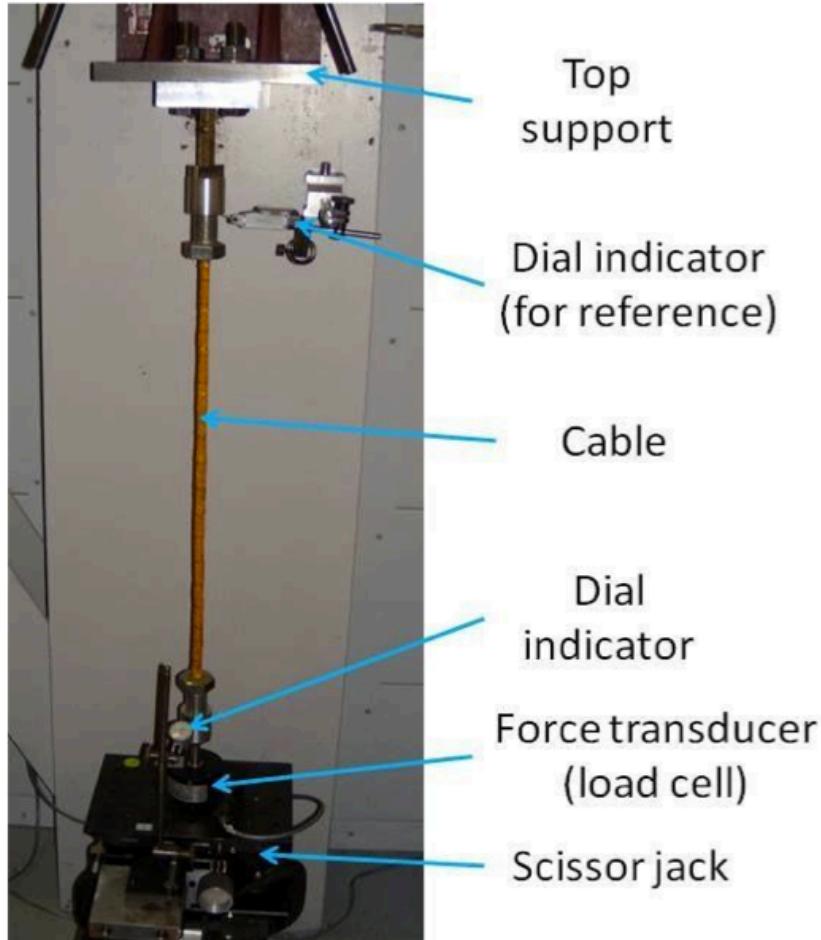


Goodding (2008)
Ardelean et al. (2010)



Cables are modeled using effective stiffnesses determined experimentally

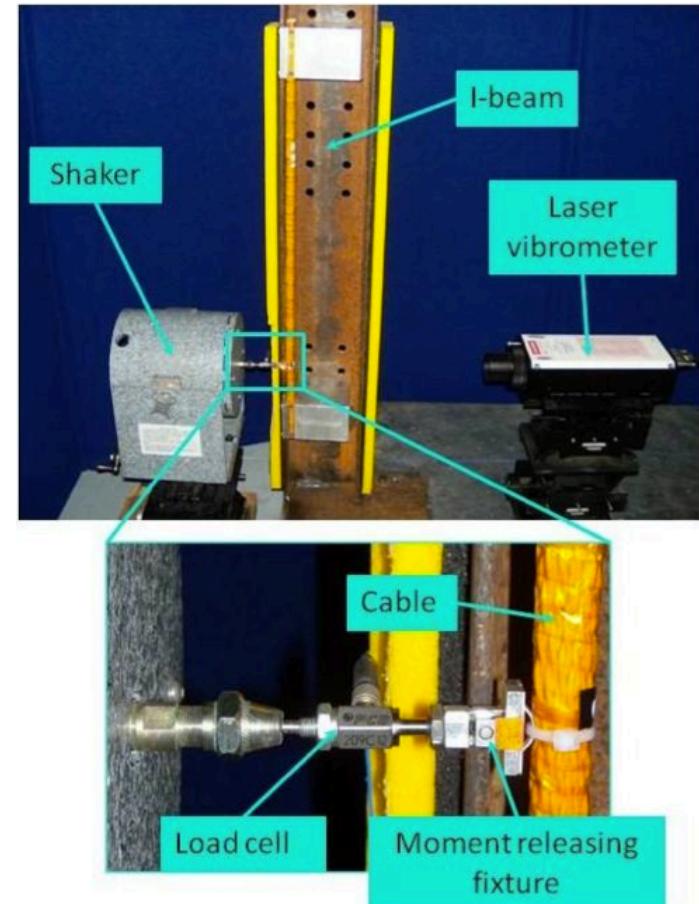
- Extension testing: EA



Sandia NL — Ardelean et al. (2010)

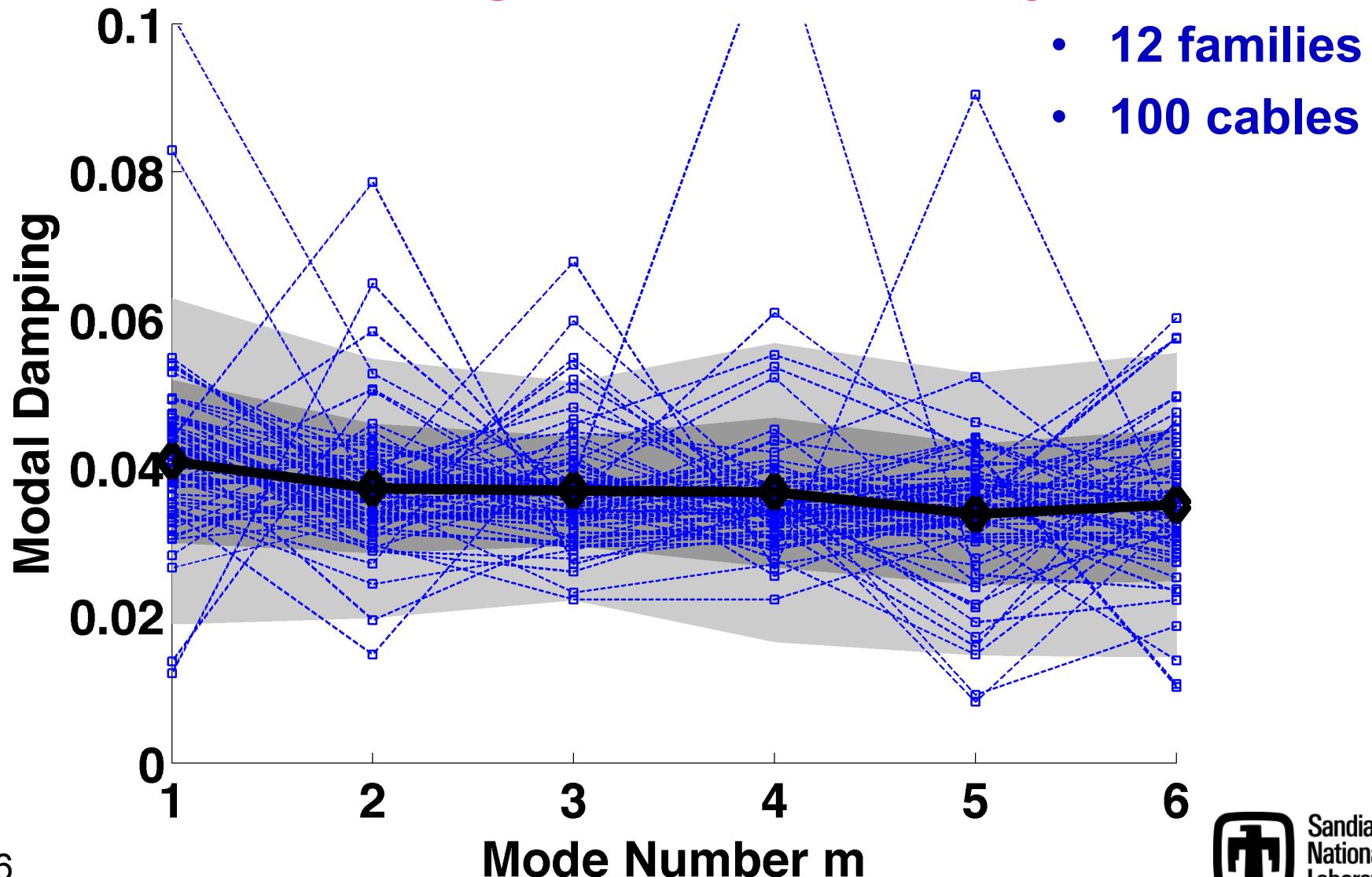
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- Lateral testing: EI & κG



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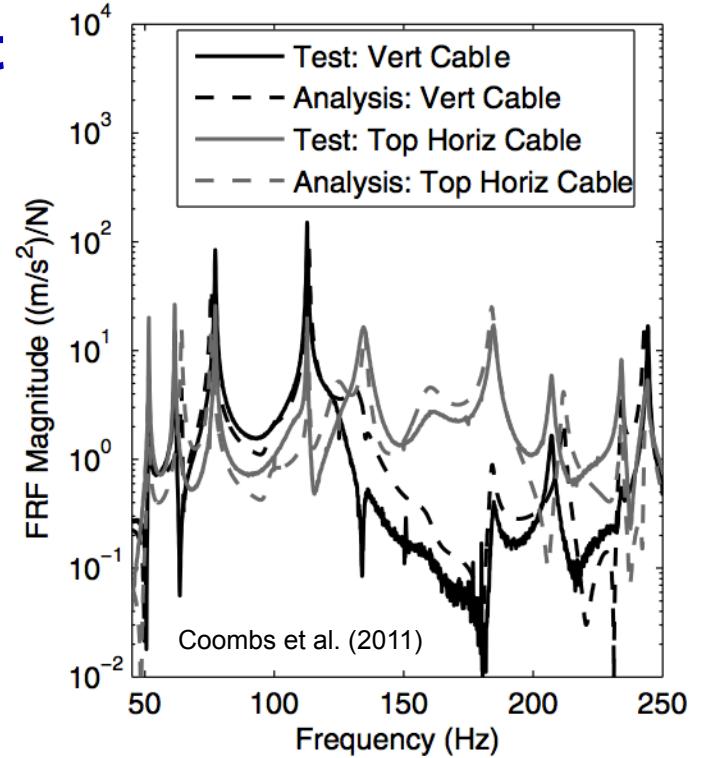
Experimental results show that
modal damping is approximately constant

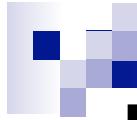




Cables modeled as shear beams initially with “structural” damping

- By researchers at Sandia / AFRL / CSA Engineering / Schafer Corp.
 - Goodding, Ardelean, Babuška, Coombs, et al. (2008-2011)
- Predicts natural frequencies, but damping model is inadequate
- Time-domain model essential
 - Transients & impact response
 - Nonlinearities
- Ideal: ~constant damping
 - Higher damping in higher modes to reduce response
- Need better understanding of physical mechanisms

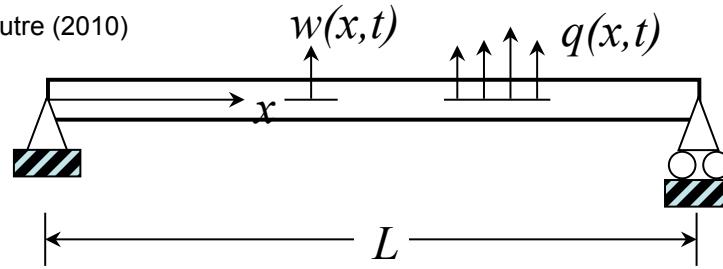




Frequency-independent modal damping is possible for Euler-Bernoulli beams

- Especially for SS BCs

Lesieutre (2010)



$$\rho A \ddot{w} - \alpha_G \dot{w}'' + EI w'''' = q$$

BCs: $w(0,t) = w(L,t) = w''(0,t) = w''(L,t) = 0$

$$w_m = a_m \sin \frac{m\pi x}{L}$$

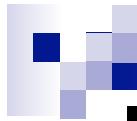
- Modal EOM

$$\ddot{a}_m \rho A + \dot{a}_m \alpha_G \left(\frac{\pi}{L} \right)^2 m^2 + a_m \left(\frac{\pi}{L} \right)^4 m^4 = 0$$

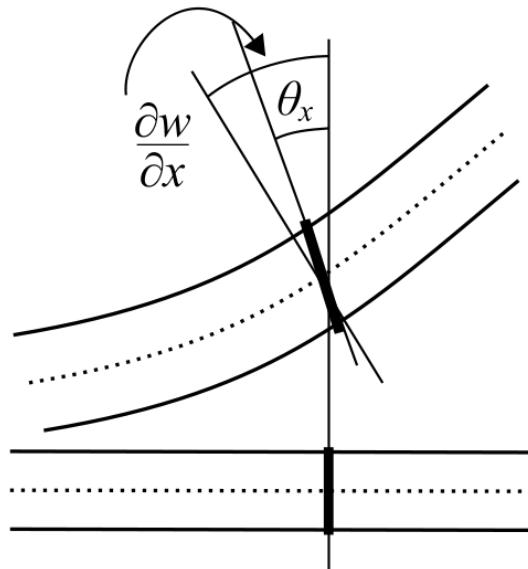
$$\omega_m = \sqrt{\frac{EI}{\rho A}} \left(\frac{m\pi}{L} \right)^2$$

$$\zeta_m = \frac{\alpha_G}{2\sqrt{\rho AEI}} = \text{constant!}$$

Can this be extended to a shear beam?



Including first-order transverse shear requires two variables from three choices



w = transverse displacement

φ = rotation due to bending

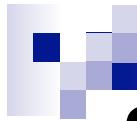
β = shear angle

$$\beta = w' - \varphi$$

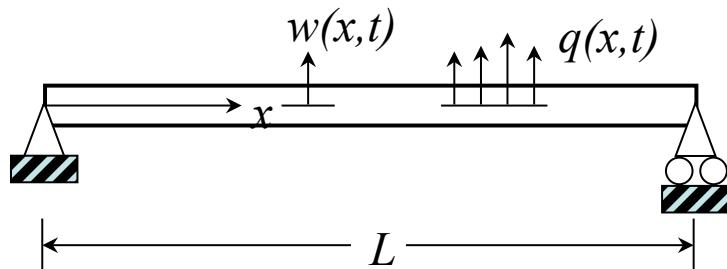
- **Shear strain contains correction factor κ**
 - Actual shear strain is not constant through the thickness
 - Shear force related to nominal shear strain

$$\varepsilon_{xz} = \kappa \beta = \kappa (w' - \varphi)$$

$$\int_A \tau \, dA = \kappa A G (w' - \varphi) = -EI\varphi''$$



Simply-supported BCs provide valuable insight into behavior



$$\rho A \ddot{w} - \frac{\rho A E I}{\kappa A G} \ddot{w}'' + E I w'''' = q - \frac{E I}{\kappa A G} q''$$

$$w(0,t) = w(L,t) = w''(0,t) = w''(L,t) = 0$$

- Mode shape is integer number of half-sine waves

$$w_m = a_m \sin \frac{m\pi x}{L}$$

- Modal EOM and natural frequencies

$$\ddot{a}_m \rho A \left(1 + \varepsilon m^2\right) + a_m \left(\frac{\pi}{L}\right)^4 m^4 = 0 \quad \rightarrow \quad \omega_m = \underbrace{\sqrt{\frac{E I}{\rho A}} \left(\frac{m\pi}{L}\right)^2}_{\omega_{m,BE}} \frac{1}{\sqrt{1 + \varepsilon m^2}}$$

- Shear parameter relates bending & shear stiffness

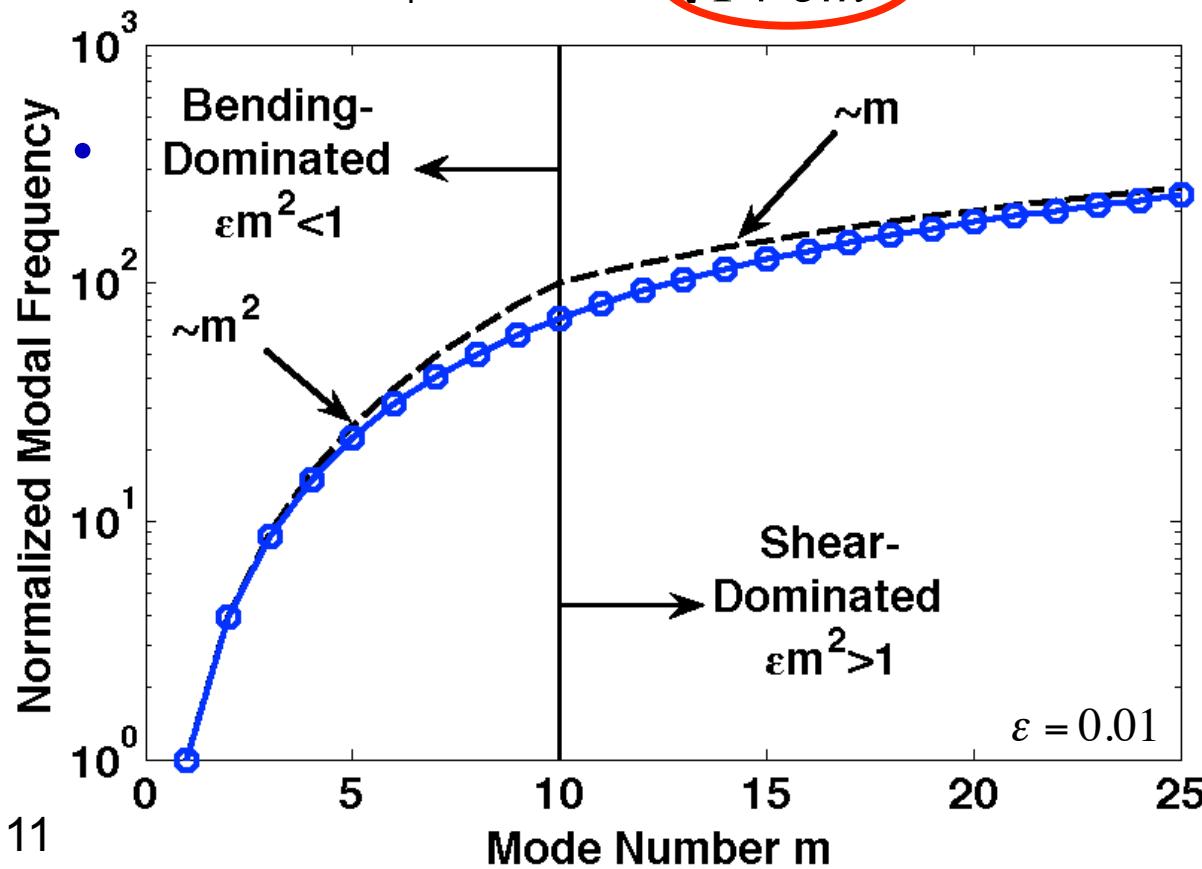
$$\varepsilon = \frac{E I \pi^2}{\kappa A G L^2}$$

Vibration modes can be separated into bending- and shear-dominated regimes

- Transition described by shear parameter

$$\frac{\omega_m}{\omega_1} = \sqrt{1 + \varepsilon} \frac{m^2}{\sqrt{1 + \varepsilon m^2}}$$

$\varepsilon m^2 = 1 \Rightarrow m = \frac{1}{\sqrt{\varepsilon}}$



Limiting behavior

$$\frac{\omega_m}{\omega_1} \sim \begin{cases} m^2 & \varepsilon m^2 \ll 1 \\ m & \varepsilon m^2 \gg 1 \end{cases}$$

Two common viscous damping models yield unrealistic damping behavior

- **Coupled EOM:**

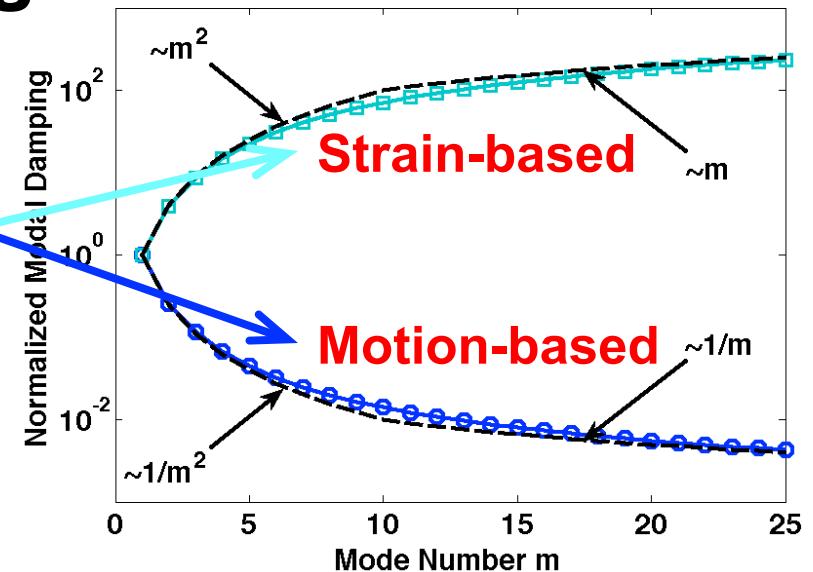
$$\begin{aligned} -\rho A \ddot{w} + \kappa A G (-\varphi' + w'') &= -q + \alpha_M \dot{w} \\ EI \varphi'' + \kappa A G (-\varphi + w') &= -\alpha_{EI} \ddot{w}''' \end{aligned}$$

- **Combined in single EOM:**

$$\begin{aligned} \rho A \ddot{w} - \frac{\rho A E I}{\kappa A G} \ddot{w}'' + E I w'''' + \\ + \alpha_M \dot{w} - \alpha_M \frac{E I}{\kappa A G} \dot{w}'' + \alpha_{EI} \dot{w}''' &= q - \frac{E I}{\kappa A G} q'' \end{aligned}$$

- **Resulting damping:**

$$\zeta_m = \frac{1}{2\sqrt{\rho A E I}} \left[\alpha_M \frac{\sqrt{1 + \varepsilon m^2}}{\left(\frac{\pi}{L}\right)^2 m^2} + \alpha_{EI} \frac{\left(\frac{\pi}{L}\right)^2 m^2}{\sqrt{1 + \varepsilon m^2}} \right]$$





Shear- and bending-related damping terms yield good results

- Introduce two internal shear forces for damping
 - Associated with time rate of change of shear & bending angles

$$V = -\alpha_\beta \dot{\beta} - \alpha_\varphi \dot{\varphi}$$

- EOM with damping

$$-\rho A \ddot{w} + \kappa A G (-\varphi' + w'') = -q - \alpha_\beta \dot{\beta}' - \alpha_\varphi \dot{\varphi}'$$

$$EI\varphi'' + \kappa A G (-\varphi + w') = 0$$

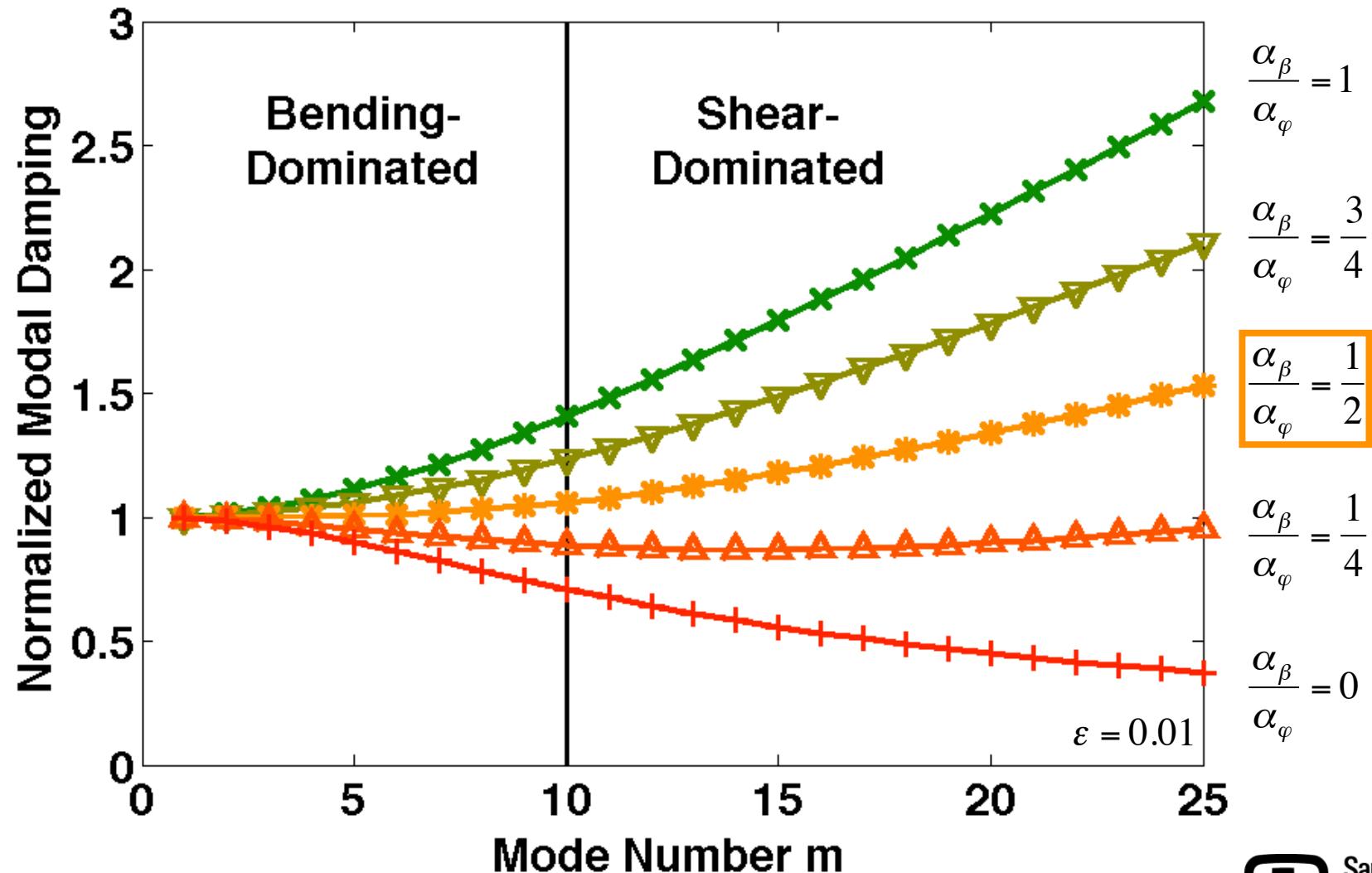
$$\rho A \ddot{\varphi} - \frac{\rho A E I}{\kappa A G} \ddot{\varphi}'' - \alpha_\varphi \dot{\varphi}'' + \alpha_\beta \frac{E I}{\kappa A G} \dot{\varphi}'''' + E I \varphi'''' = q'$$

$$\zeta_m = \frac{\alpha_\varphi}{2\sqrt{\rho A E I}} \frac{1 + \frac{\alpha_\beta}{\alpha_\varphi} \varepsilon m^2}{\sqrt{1 + \varepsilon m^2}}$$

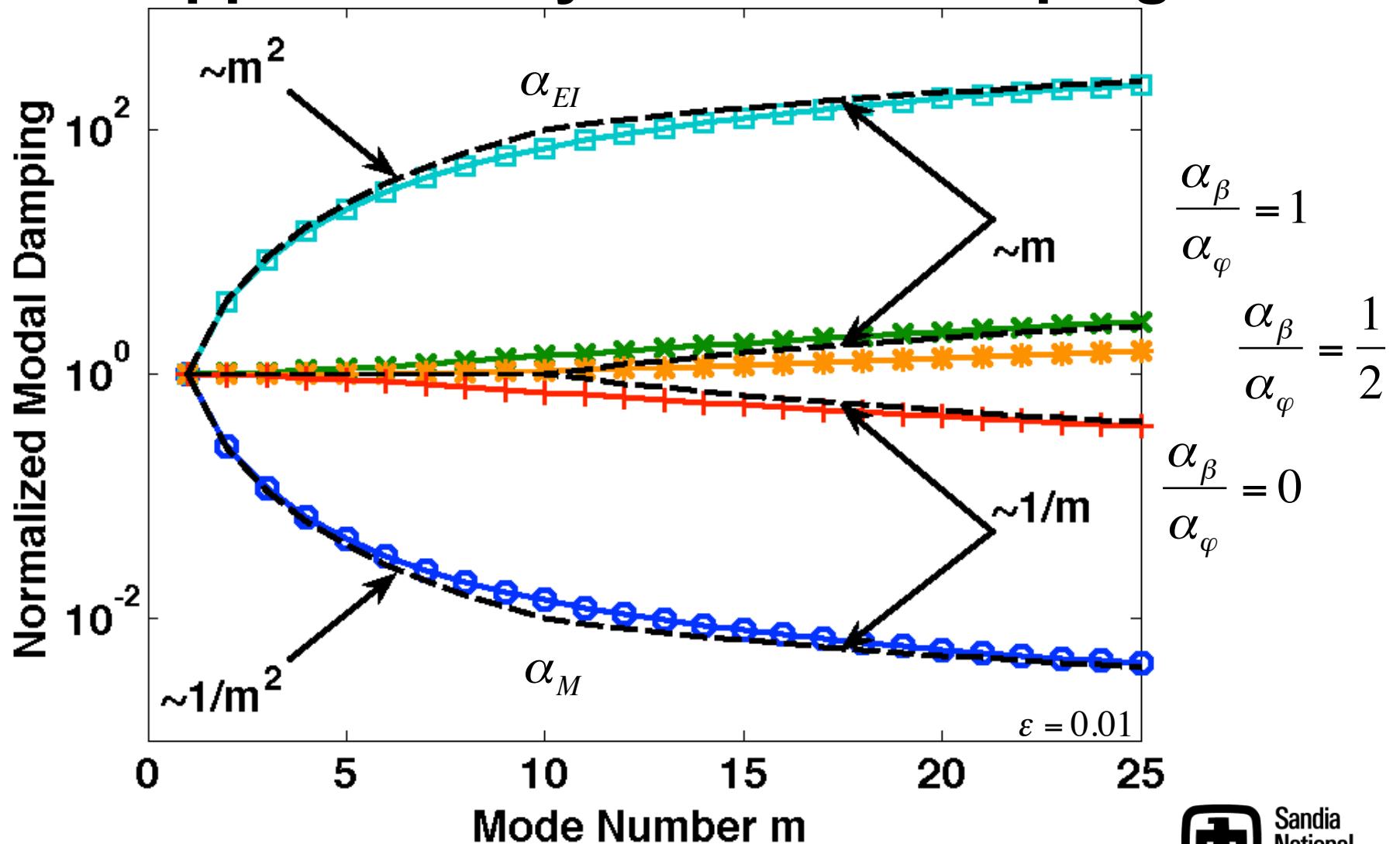
- Shear- and bending-related damping contributions are explicitly separated

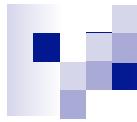
$$\zeta_m \sim \begin{cases} \alpha_\varphi \varepsilon m^2 \ll 1 \\ \alpha_\beta m \varepsilon m^2 \gg 1 \end{cases}$$

A range of damping trends available from choice of shear & bending terms



Proposed model provides realistic and approximately constant damping





A numerical (FE) approach can use conventional K, K_G, & M matrices

- **Coupled EOM**

$$-\rho A \ddot{w} + \kappa A G (-\varphi' + w'') = -q - \alpha_\beta \dot{\beta}' - \alpha_\varphi \dot{\varphi}'$$

$$EI\varphi'' + \kappa A G (-\varphi + w') = 0$$

- **Combined in single EOM in φ**

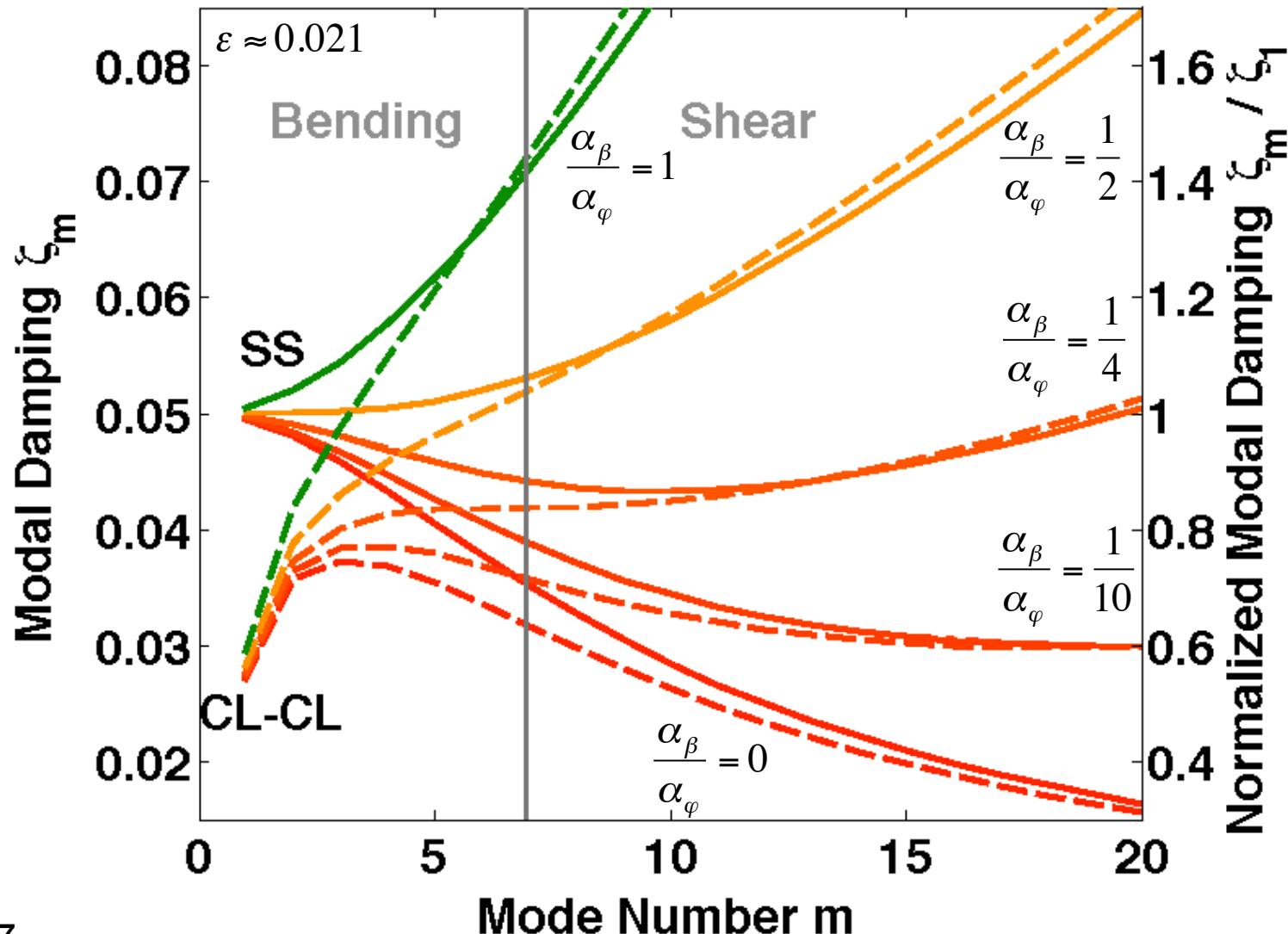
$$\rho A \ddot{\varphi} + \frac{\rho A E I}{\kappa A G} \ddot{\varphi}'' - \alpha_\varphi \dot{\varphi}'' + \alpha_\beta \frac{E I}{\kappa A G} \dot{\varphi}'''' + E I \varphi'''' = q'$$

$$\left[[M] + \frac{\rho A E I}{\kappa A G} [K_G] \right] \{ \ddot{\varphi} \} + \left[\alpha_\varphi [K_G] + \alpha_\beta \frac{1}{\kappa A G} [K] \right] \{ \dot{\varphi} \} + [K] \{ \varphi \} = \{ q' \}$$

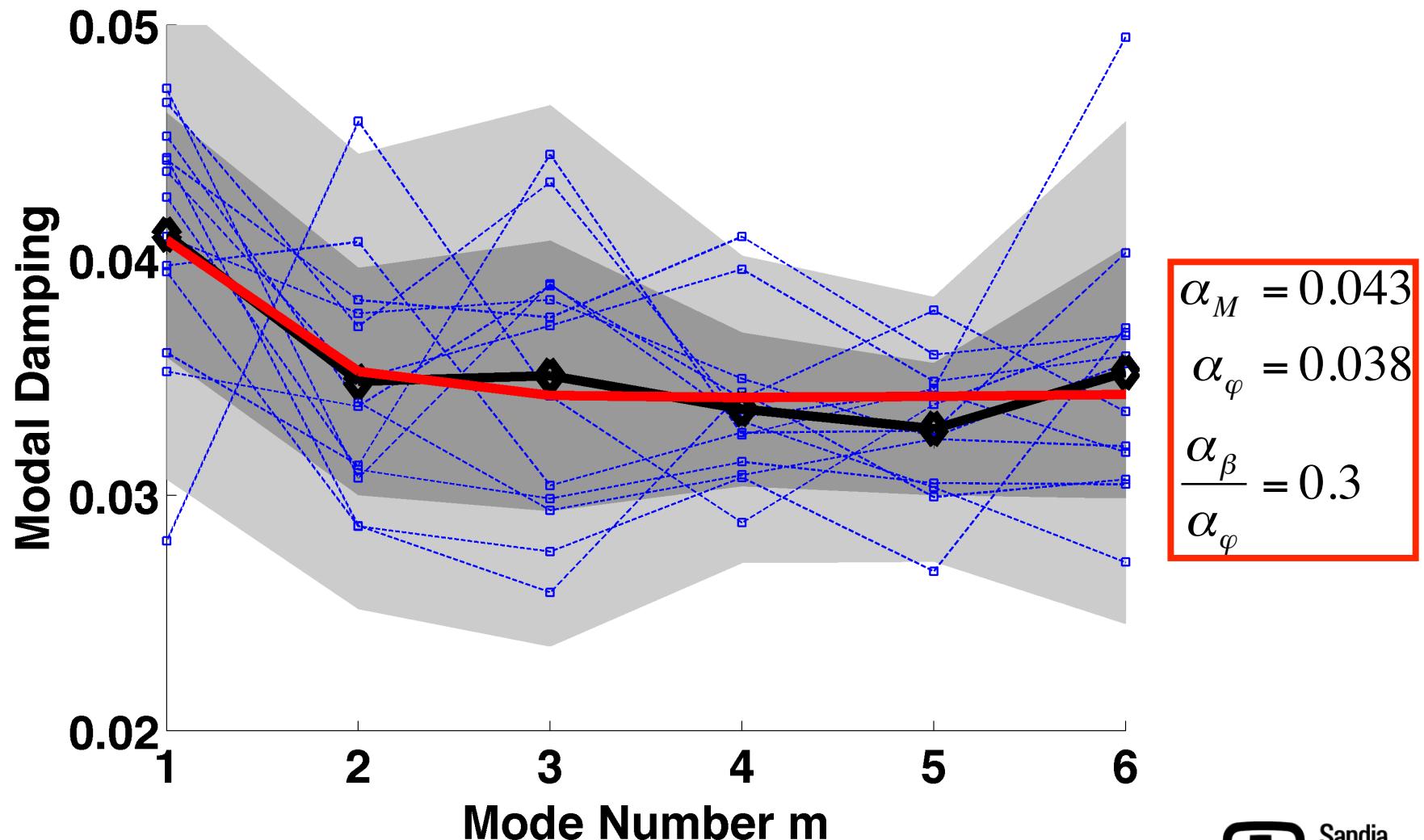
- **Forcing might be known; it could be finite-differenced**

$${}^n \{ q' \} \approx \frac{{}^{n+1} \{ q \} - {}^{n-1} \{ q \}}{2 L_{el}}$$

FE analysis confirms damping variation with mode number for various BCs

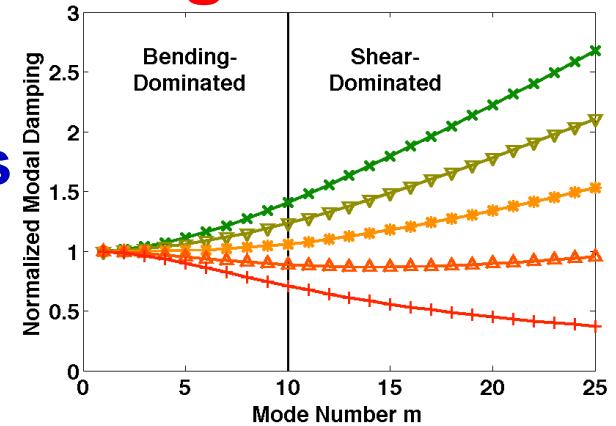


Proposed model with motion-, shear-, and bending-based terms fits data well

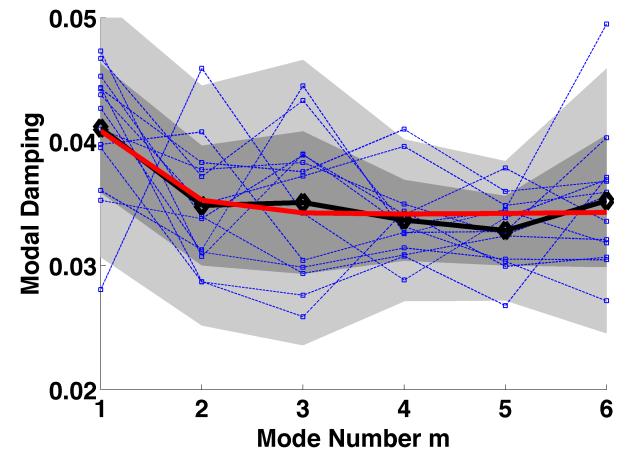


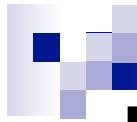
Time-domain damping model for shear beams captures dynamics of spacecraft cabling

- Behavior can be separated into bending- and shear-dominated regimes
 - Corresponding physical understanding
$$-\rho A \ddot{w} + \kappa A G (-\varphi' + w'') = -q - \alpha_\beta \dot{\beta}' - \alpha_\varphi \dot{\varphi}'$$
- Freq-independent modal damping achievable in bending region
 - Can control damping in shear regime
 - Can achieve best possible freq-indep
- Damping model can be readily implemented using FEM
 - Uses conventional K , K_G , & M matrices
- Model predictions agree well with experimental data



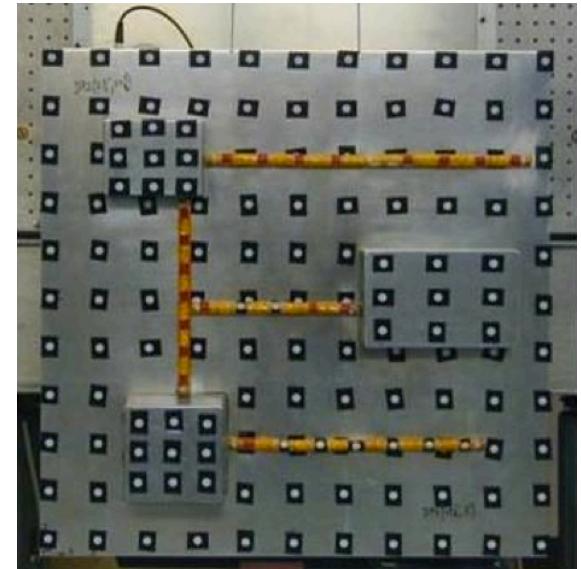
$$\zeta = \frac{\alpha_\varphi}{2\sqrt{\rho AEI}} \frac{1 + \frac{\alpha_\beta}{\alpha_\varphi} \varepsilon m^2}{\sqrt{1 + \varepsilon m^2}}$$





Future work will examine realistic BCs and consider Timoshenko beam model

- **Typical spacecraft cabling configuration**
 - Not really SS or CC BCs
 - Model as adjustable rotational stiffness

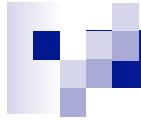


Ardelean et al. (2010)

- **Timoshenko model adds rotatory inertia**

$$-\rho A \ddot{w} + \kappa A G (-\varphi' + w'') = -q - \alpha_\beta \dot{\beta}' - \alpha_\varphi \dot{\varphi}'$$

$$-\rho I \ddot{\varphi} + EI \varphi'' + \kappa A G (-\varphi + w') = 0$$



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