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Terahertz plasmonic cavities and crystals: coherent phenomena in 2D electronic systems



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PECS-X 2012, Wednesday June 6

This work is supported by the U.S. AFOSR under Contract No. FA9550-09-C-0168 and Physical Sciences Inc. (PSI) under Agreement No. FI011090528 and does not reflect the views of the U.S. AFOSR or PSI.

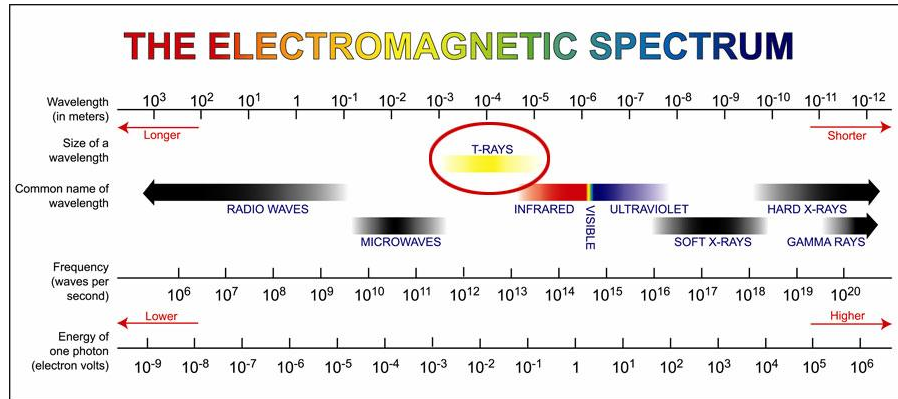


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Outline

- THz Plasmonics
- Integrated 2D plasmonic THz and sub-THz detectors
 - Grating gated & plasmonic cavity detectors
 - Figures of merit: $\text{NEP} < 50 \text{ pW/Hz}^{1/2}$, Responsivity $> 100 \text{ kV/W}$
- Generalized 2D Plasmonic Transmission Line Model
 - Hydrodynamic and Finite Element approaches
 - Connecting screened and unscreened limits via TL model
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 - Finite plasmonic crystal
- Conclusions

Plasmonics in the Terahertz Band



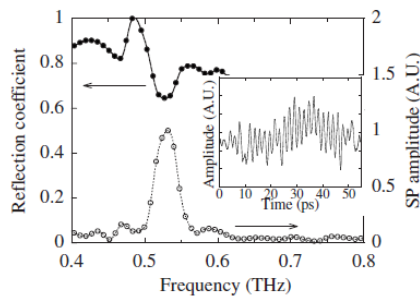
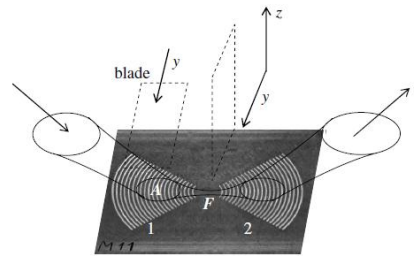
Electronics: RC limited bandwidth

Photonics: kT limited below 1 THz

Plasmonics: $\omega\tau$ cutoff

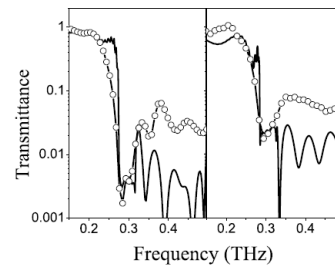
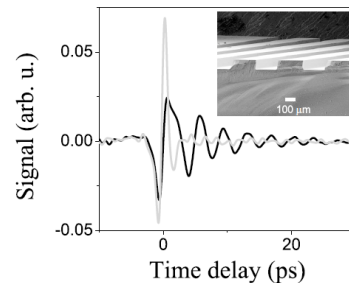
from http://web.njit.edu/~barat/RBB_research.html

Metal Gratings



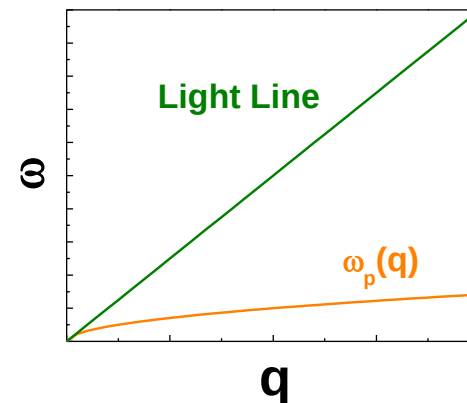
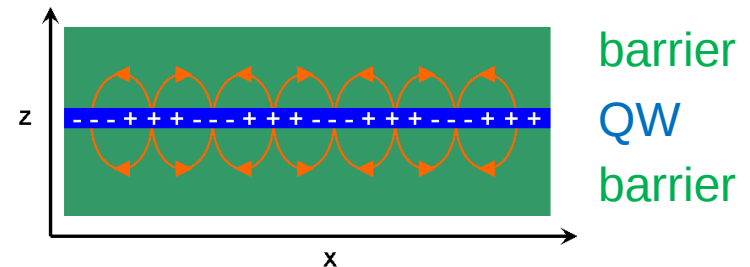
G. Gaborit et. al., *APL*
94, 231108 (2009)

Semiconductor Gratings



J. G. Rivas et. al., *PRL*
93, 256804 (2004)

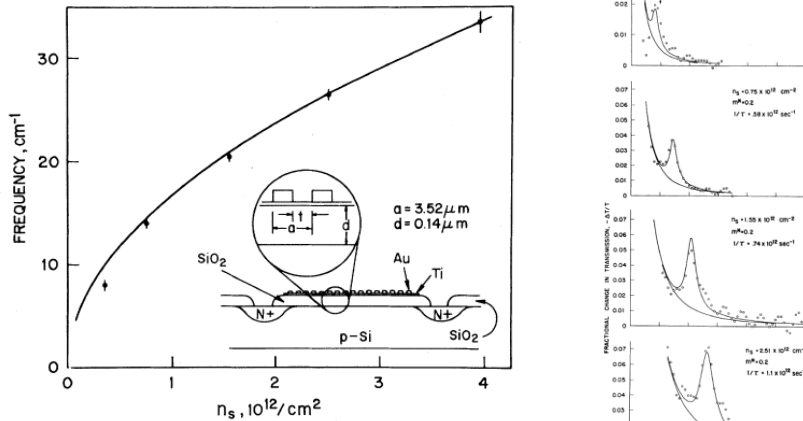
Semiconductor 2DEGs



$$\omega_p = \sqrt{\frac{n_{2D} e^2 q}{2\epsilon m^*}}$$

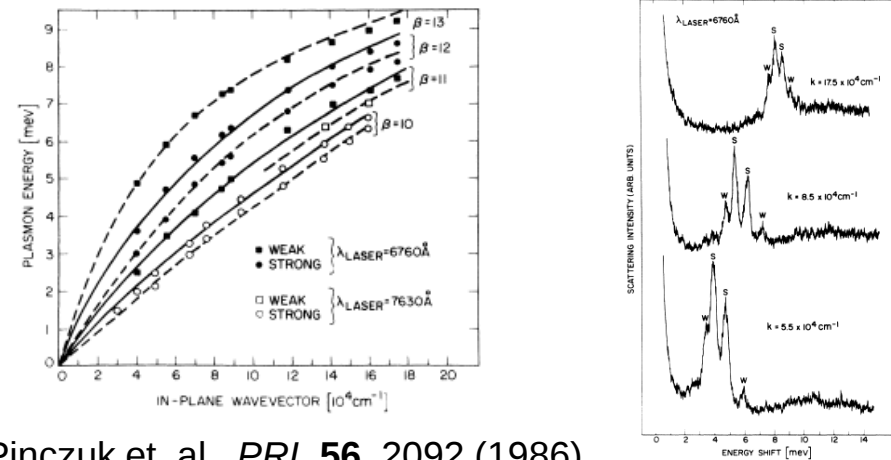
THz 2D Plasmonic Systems

Grating Gated Si MOSFETs



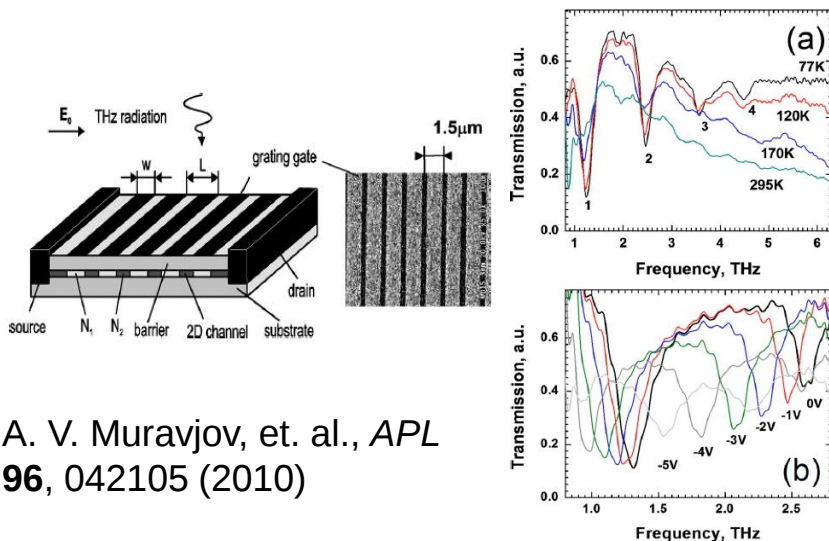
S. J. Allen et. al., *PRL* **38**, 980 (1977)

GaAs/AlGaAs Heterostructures



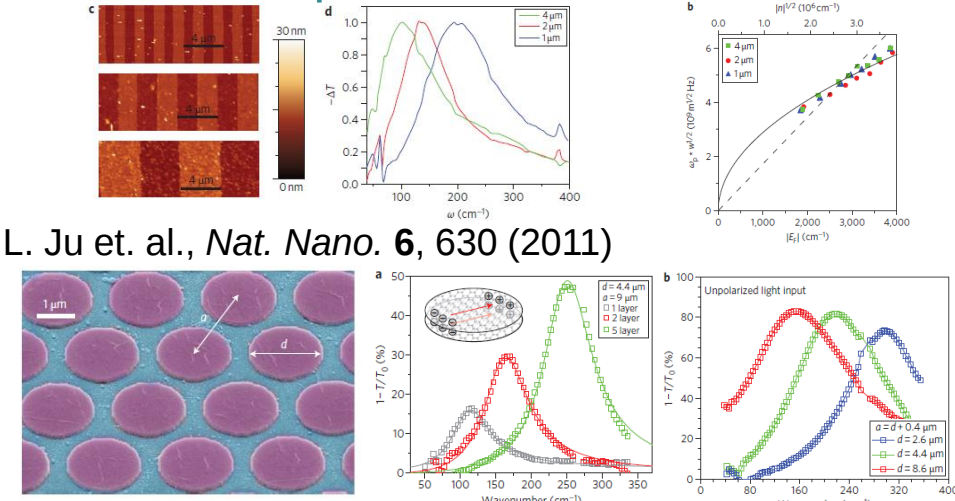
A. Pinczuk et. al., *PRL* **56**, 2092 (1986)

Grating Gated GaN/AlGaAs HEMTs



A. V. Muravjov, et. al., *APL* **96**, 042105 (2010)

Periodic Graphene Structures



L. Ju et. al., *Nat. Nano.* **6**, 630 (2011)

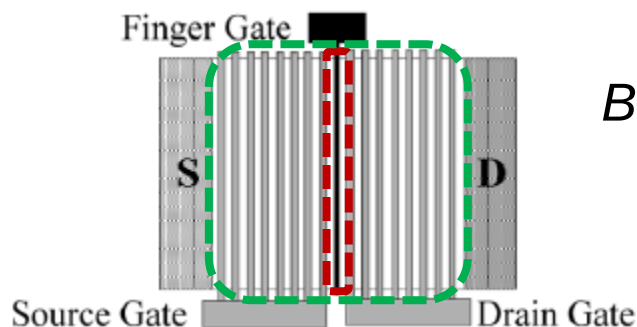
H. Yan et. al., *Nat. Nano. adv. online pub.* (2012)

Outline

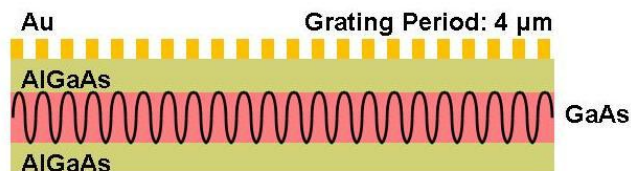
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Integrated 2D Plasmonic Detectors

Grating Gated HEMT Detector



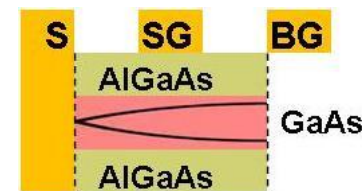
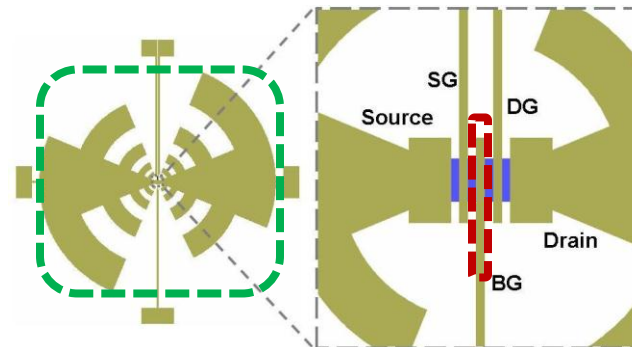
Bolometric Sensor
THz Coupling
2D Plasmons



$$\omega_p \propto \sqrt{n_{2D} q}$$

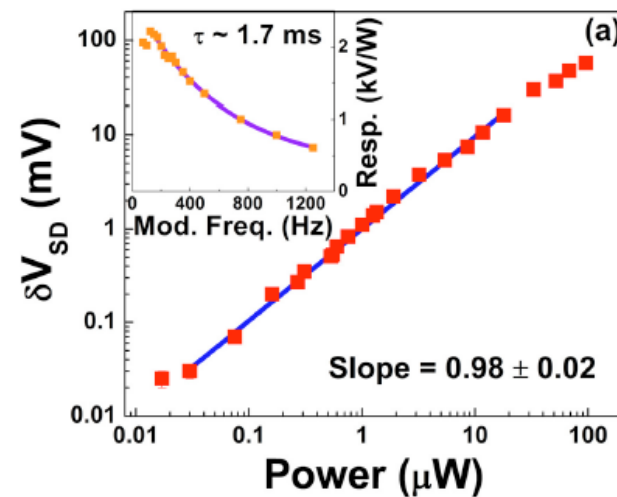
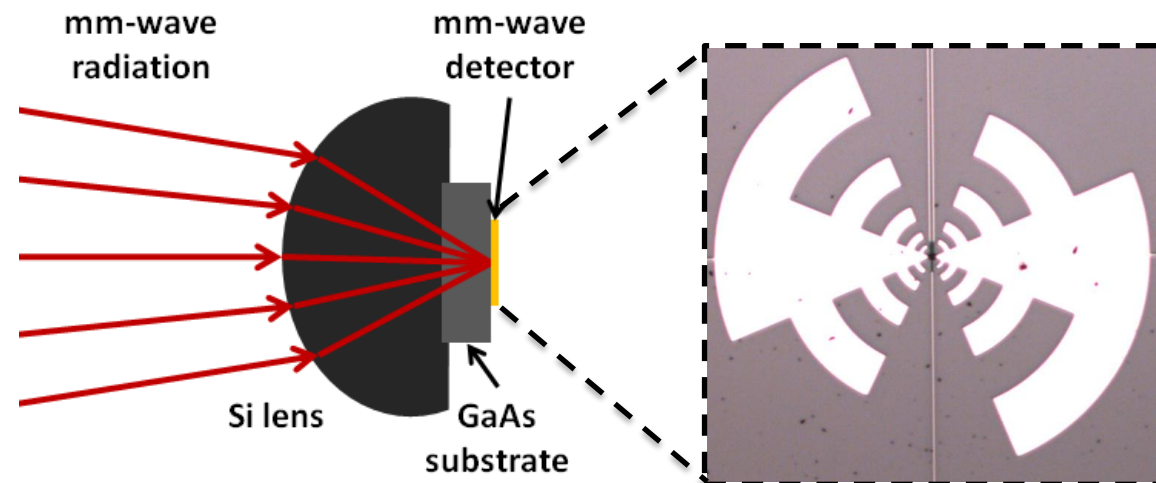
- 1 mm square channel
- 4 μm grating period, 50% duty cycle
- Detection: bolometric sensor
- THz coupling: grating gate
- Resonant absorption: 2D plasmon defined by grating period

Antenna Coupled HEMT Detector



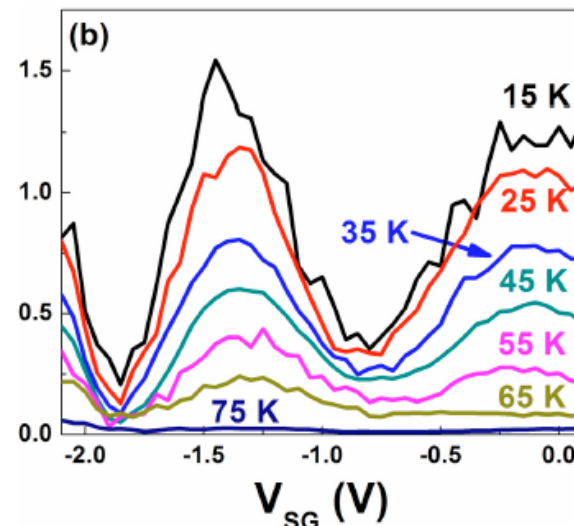
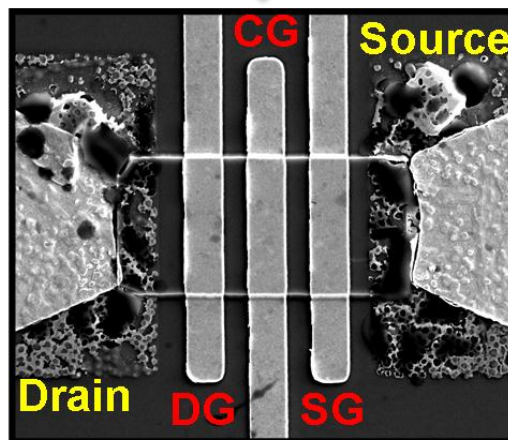
- 1.5 mm diameter antenna
- 10 x 14 μm channel, 2 μm gates
- Detection: bolometric sensor
- THz coupling: log-periodic antenna
- Resonant absorption: 2D plasmon cavity at antenna vertex

Log Periodic Antenna & Si Lens

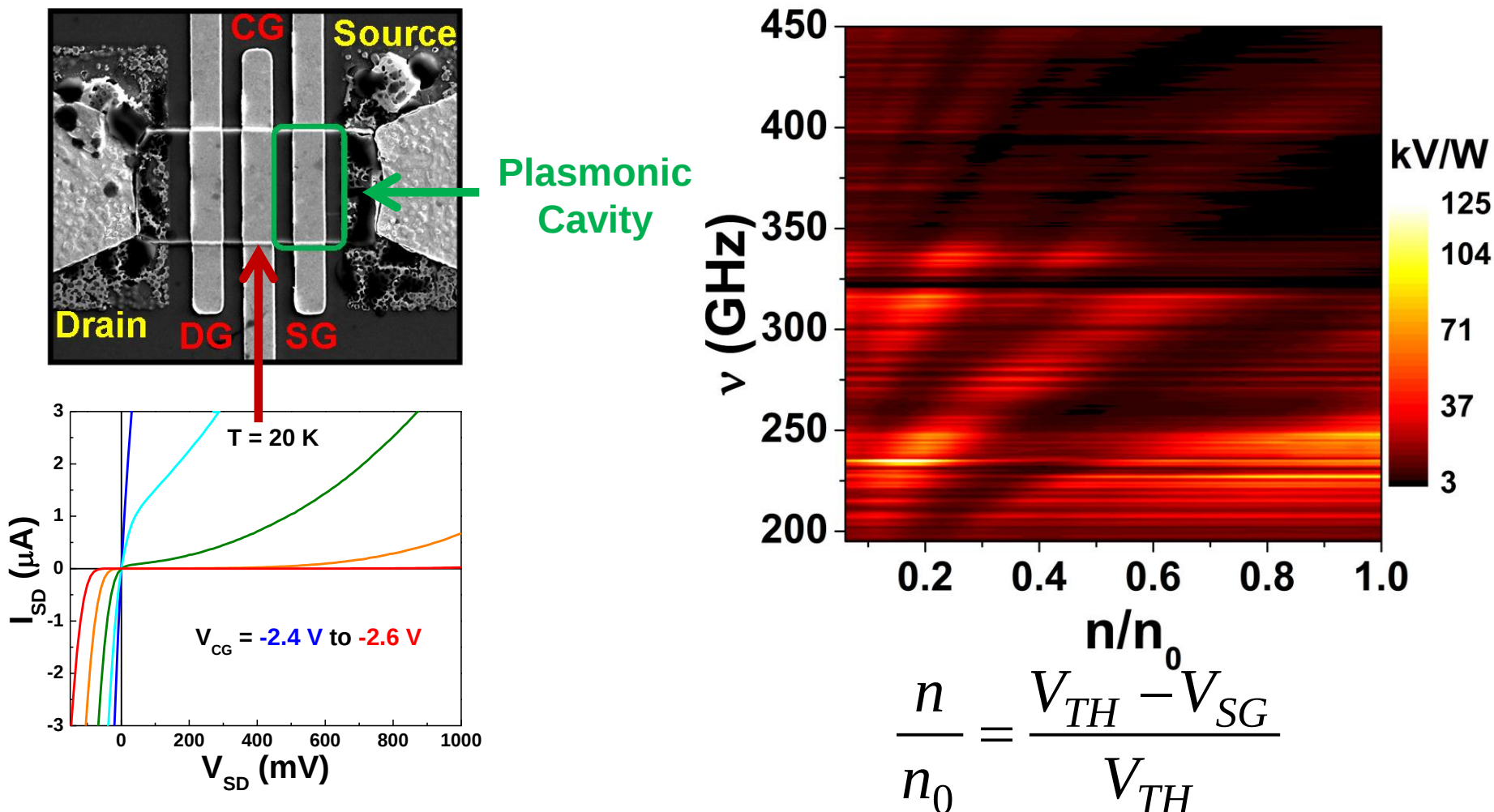


Five terminal HEMT at antenna vertex: S, D and gates (SG, CG and DG)

Self complementary antenna: impedance of $\sim 72 \Omega$ on substrate

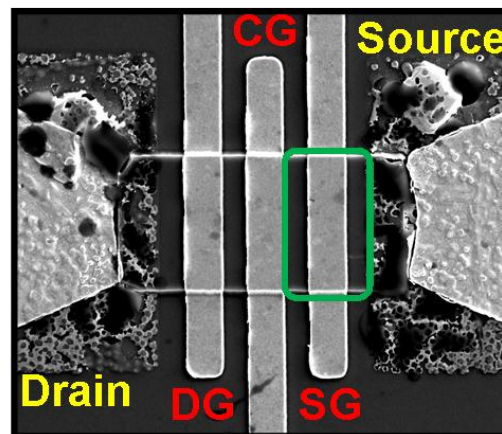
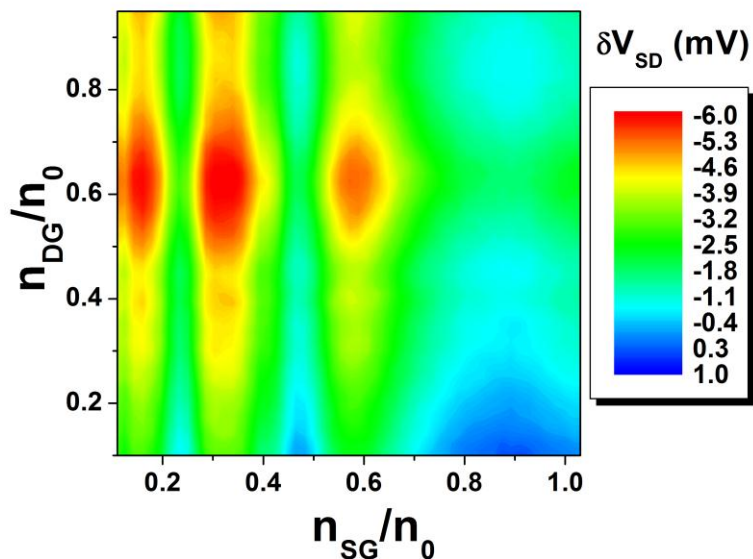


Gate Controlled Plasma Cavity



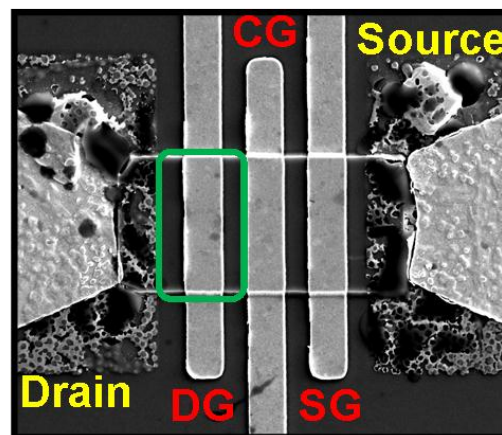
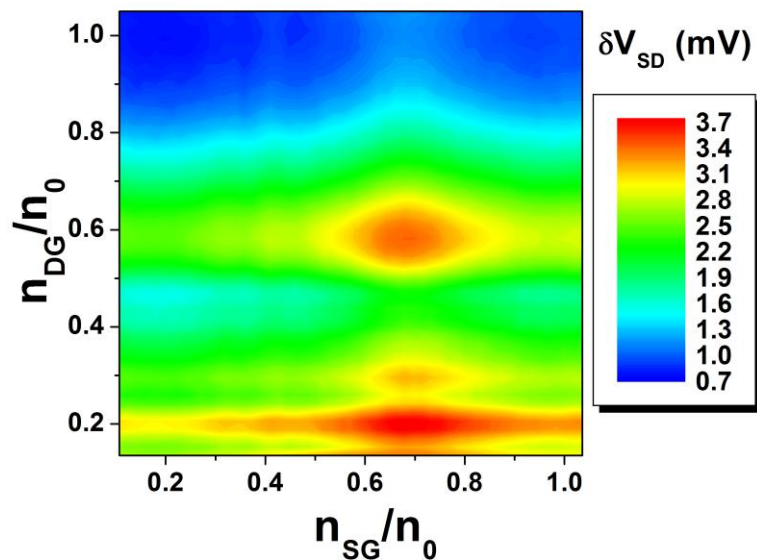
- Source gate (SG) tunes plasma resonance with center gate (CG) beyond threshold

S-D Bias Cavity Tuning



370 GHz

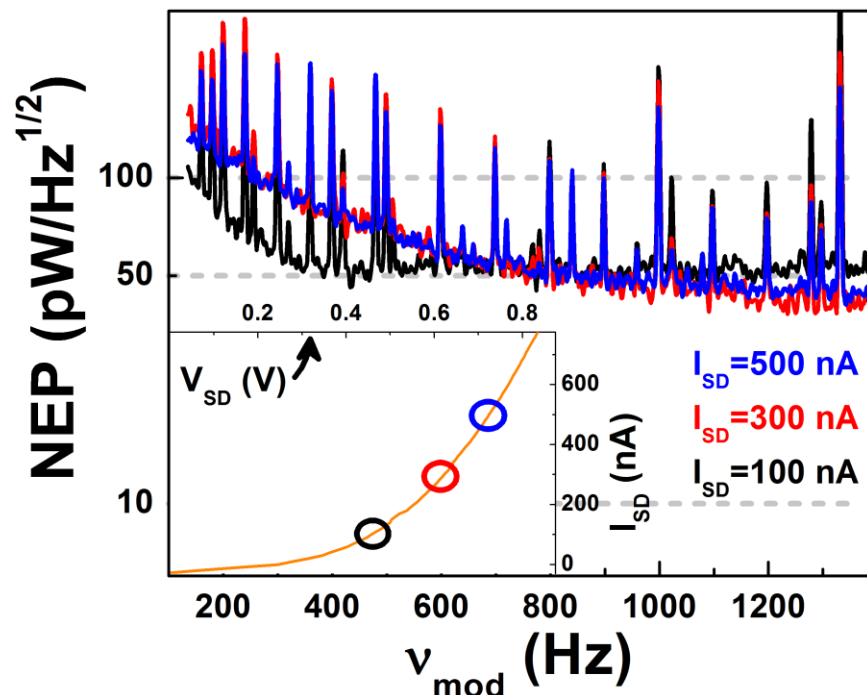
$I = +500$ nA



$I = -500$ nA

G. C. Dyer et. al., *APL*
100, 083506 (2012)

Figures of Merit

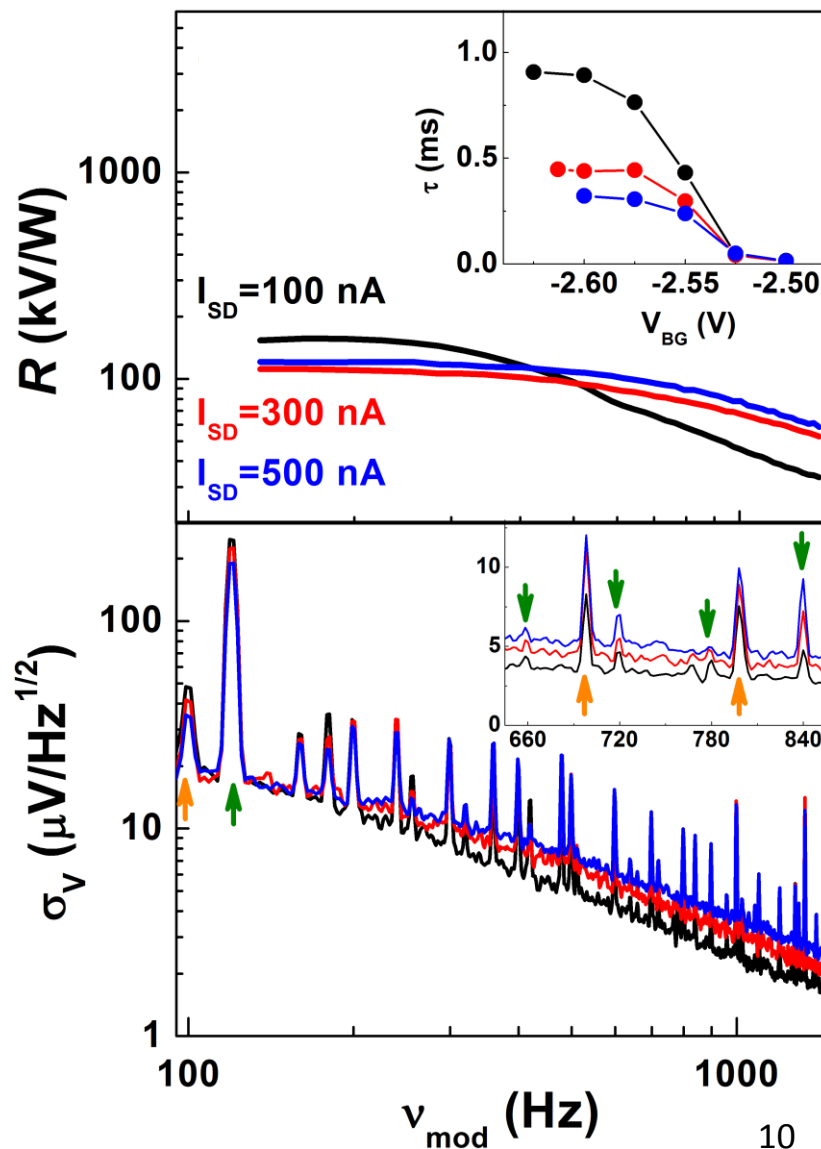


11K Operation, 243 GHz

$R > 100 \text{ kV/W}$

$NEP < 50 \text{ pW/Hz}^{1/2}$

G. C. Dyer et. al., *APL* **100**, 083506 (2012)



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Hydrodynamic Model

$$\frac{\partial \delta v}{\partial t} + \delta v \frac{\partial \delta v}{\partial x} = \frac{e}{m} \frac{\partial \phi}{\partial x} \quad \& \quad \frac{\partial n}{\partial t} + \frac{\partial (n \delta v)}{\partial x} = 0$$

Self-consistent soln. of Euler's hydrodynamic eqn. of motion with continuity eqn.

- **Pros:**
 - little computational overhead
 - accurately predicts both linear and non-linear phenomena (absorption, mixing, emission)
- **Cons:**
 - non-trivial to phrase in Maxwell or circuit formalism
 - must choose boundary conditions
 - assumptions about screening
 - simple geometry

Detection, Mixing, and Frequency Multiplication of Terahertz Radiation by Two-Dimensional Electronic Fluid

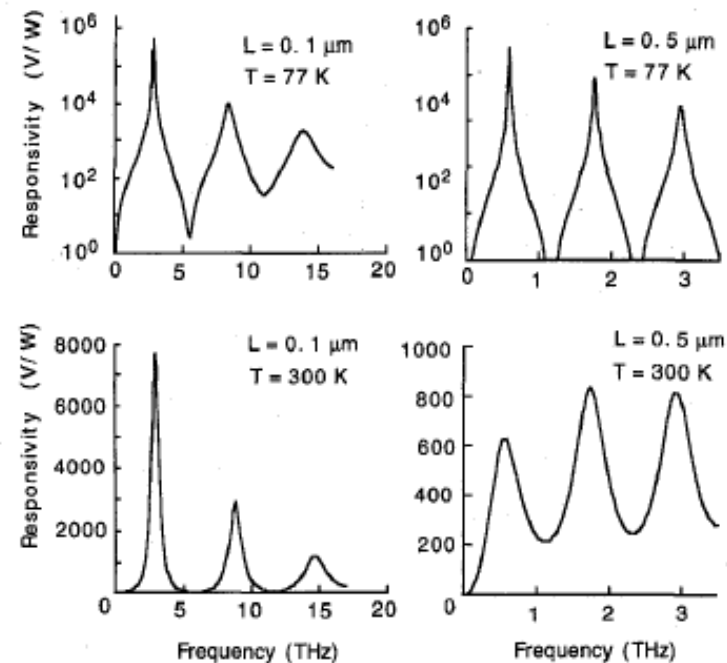
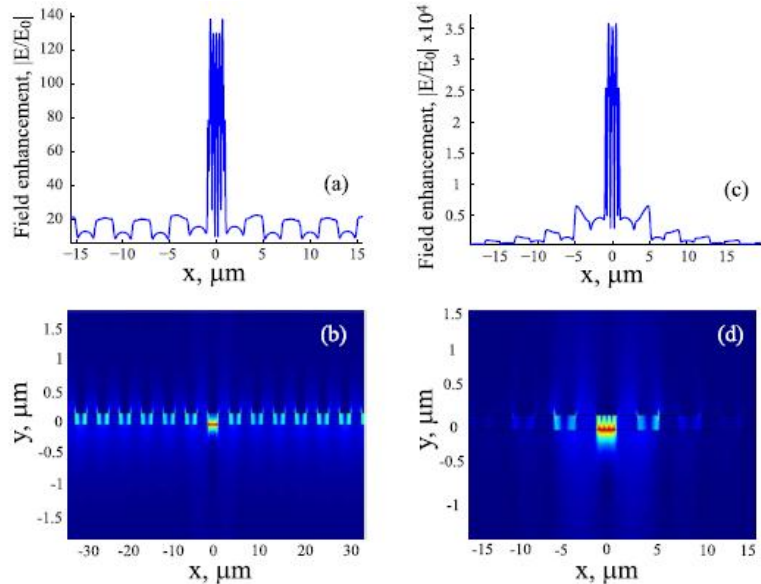


Fig. 5. Calculated responsivity as a function of frequency for 0.1 μm and 0.5 μm gate HEMT's at 77 K and 300 K. Parameters used in the calculation: $U_0 = 0.5$ V, $G = 1.5$, $\nu = 15$ cm²/s, $m = 0.063m_0$, $\mu = 300,000$ cm²/Vs at 77 K, $\mu = 8,000$ cm²/Vs at 300 K.

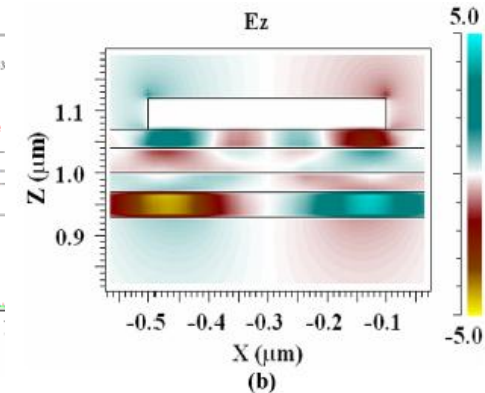
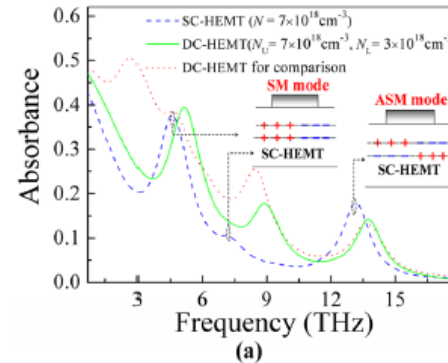
Finite Element Analysis (Full EM)

Tailoring Terahertz Near-Field Enhancement via Two-Dimensional Plasmons



A. R. Davoyan et. al., *PRL* **108**, 127401 (2012)

The plasmonic resonant absorption in GaN double-channel high electron mobility transistors



L. Wang et. al., *APL* **99**, 063502 (2011)

Pros:

- direct solution of EM parameters (fields, absorption, transmission, etc.)
- can embed 2DEG in complex geometry including material parameters
- Maxwell imposes boundary conditions

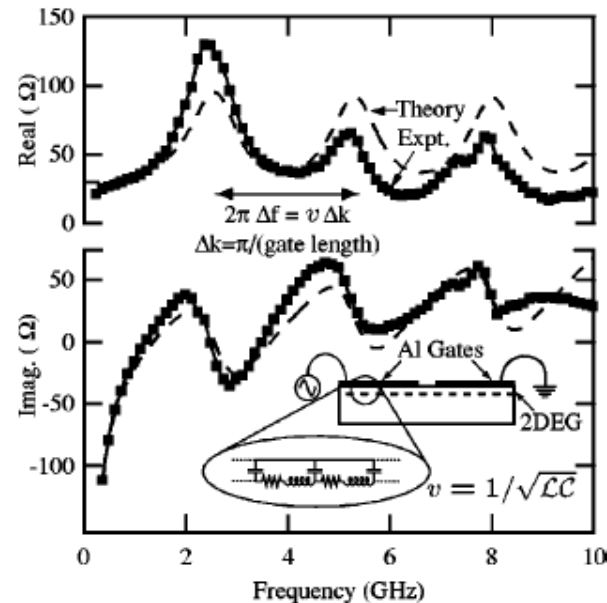
Cons:

- large computational overhead
- expensive
- “intuition before computation”

Plasmon Transmission Line Model

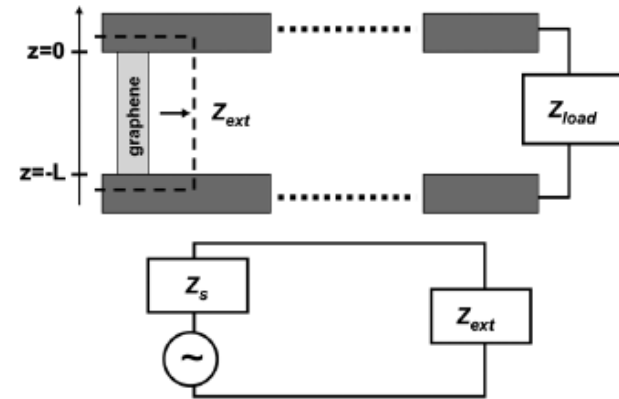
High frequency conductivity of the high-mobility two-dimensional electron gas

Graphene Terahertz Plasmon Oscillators



$$Z_0 = \sqrt{\frac{Z_{2DEG}}{i\omega C(\omega)}}$$

$$iq = \sqrt{i\omega C(\omega)Z_{2DEG}}$$



$$\lim_{qd \rightarrow 0} C(\omega) = \frac{W\epsilon}{d}$$

(fully screened limit)

$$Z_{2DEG} = R_{2DEG} + i\omega L_{2DEG}$$

$\rightarrow$ from Drude model
 L is kinetic inductance

$$\lim_{qd \rightarrow \infty} C(\omega) = 2W\epsilon q$$

(unscreened limit)

Generalized Plasmonic TL Model

$$\Phi(x, z) = [\Phi_+ e^{-iqx} + \Phi_- e^{+iqx}] \begin{cases} \sinh[q(d - z)] e^{-qd}, & z \geq 0 \\ \sinh[qd] e^{q(z-d)}, & z < 0 \end{cases}$$

$$\frac{I(x)}{W} = j(x, z) = \frac{1}{Z_0} [\Phi_+ e^{-iqx} - \Phi_- e^{+iqx}] \delta(z)$$

$$C(\omega) = \frac{\Phi(x, z=0)}{\rho(x)} = W\epsilon q(1 + \coth[qd])$$



The 2D plasmon is in general a TM wave: $P = \frac{1}{2} VI^*$ is correct only for TEM wave in fully screened limit

$$\begin{aligned} P(x) &= \frac{1}{2} \iint S_x dydz = -\frac{W}{2} \iint E_z H_y dydz \\ &= \frac{1}{2} \Phi(x, 0) I^*(x) \xi^*(\omega, d) \end{aligned}$$

$$P(x) \rightarrow \begin{cases} \lim_{qd \rightarrow 0} P(x) = \frac{1}{2} \Phi(x, 0) I^*(x) \\ \lim_{\tau, qd \rightarrow \infty} P(x) = \frac{1}{4} \Phi(x, 0) I^*(x) \end{cases}$$

$$\xi(\omega, d) = 1 - \frac{q(1 - e^{-2q'd} \cos[2q''d])}{2q'(1 - e^{-2q^*d})} + \frac{q e^{-2q'd} \sin[2q''d]}{2q''(1 - e^{-2q^*d})}$$

Effective V , I & Z

$$\frac{1}{2}\tilde{V}(x)\tilde{I}^*(x) = \frac{1}{2}\Phi(x, 0)I^*(x) \xi^*(\omega, d)$$

$$\tilde{V}(x) \equiv \Phi(x, 0)$$

$$\tilde{I}(x) \equiv I(x)\xi(\omega, d)$$

$$\tilde{Z}_0 \equiv \frac{Z_0}{\xi(\omega, d)}$$

Characteristic Impedance of Microstrip Lines

JOHN R. BREWS, MEMBER, IEEE

FOR non-TEM structures such as microstrip lines it generally is agreed that the characteristic impedance is not unique. It also is agreed that the root of this uncertainty is the longitudinal field components of non-TEM modes. Because of these components, current and voltage are not uniquely defined in terms of the usual TEM path integrals of the fields. Faced with an ambiguity in the meaning of current I and voltage V a variety of results for the characteristic impedance Z_0 is obtained.

J. R. Brews, *IEEE Trans. on Microwave Theory and Techniques* **35**, 30 (1987)

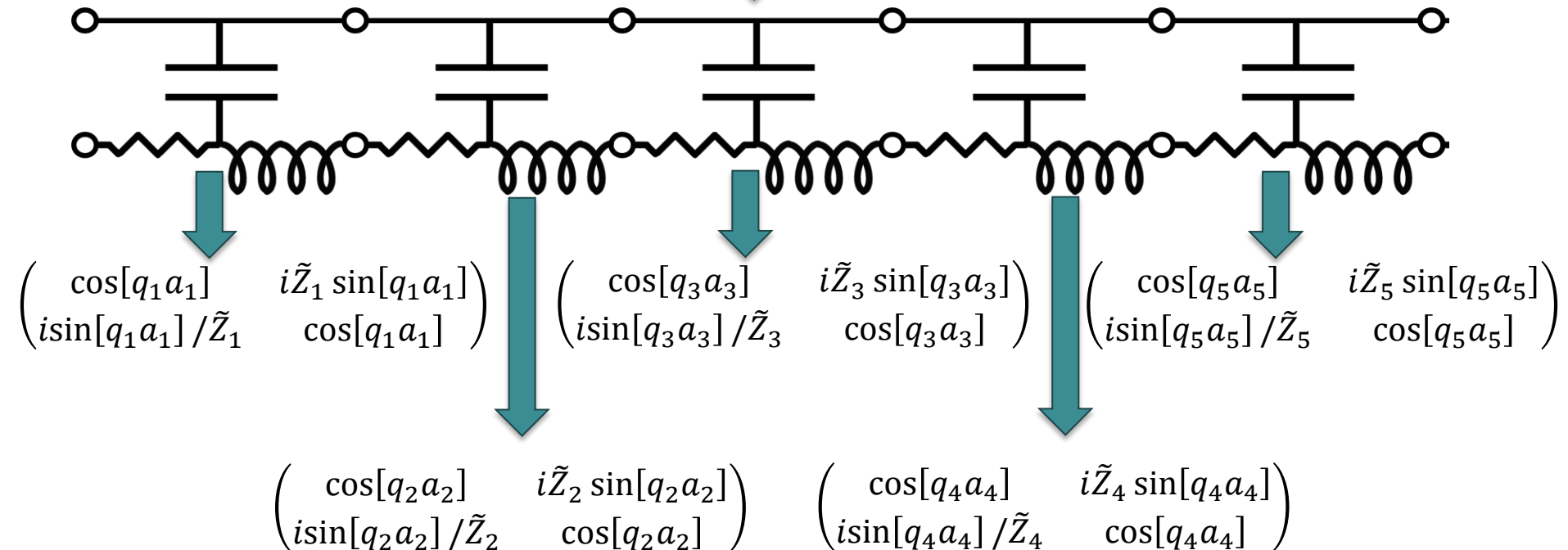
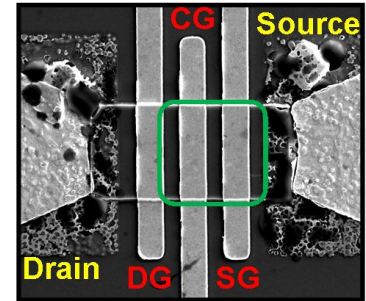
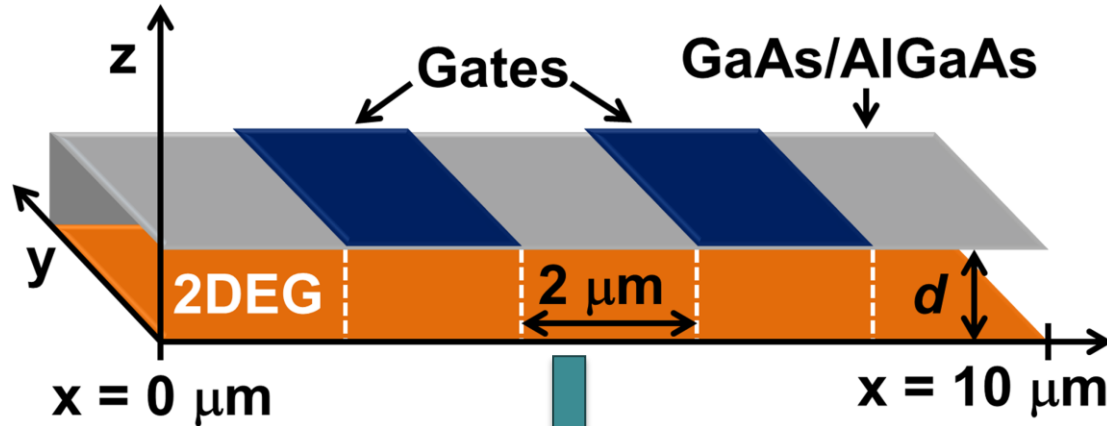
Above definitions are not unique, but accurately address physics:

1. *Effective potential $\tilde{V}(x)$ is continuous*
2. *Power $\tilde{V}(x)\tilde{I}^*(x)$ is continuous*
3. *Discontinuity in real carrier density $\rho(x)$ and real current $I(x)$ is consistent with edge charge accumulation at step-like boundary*

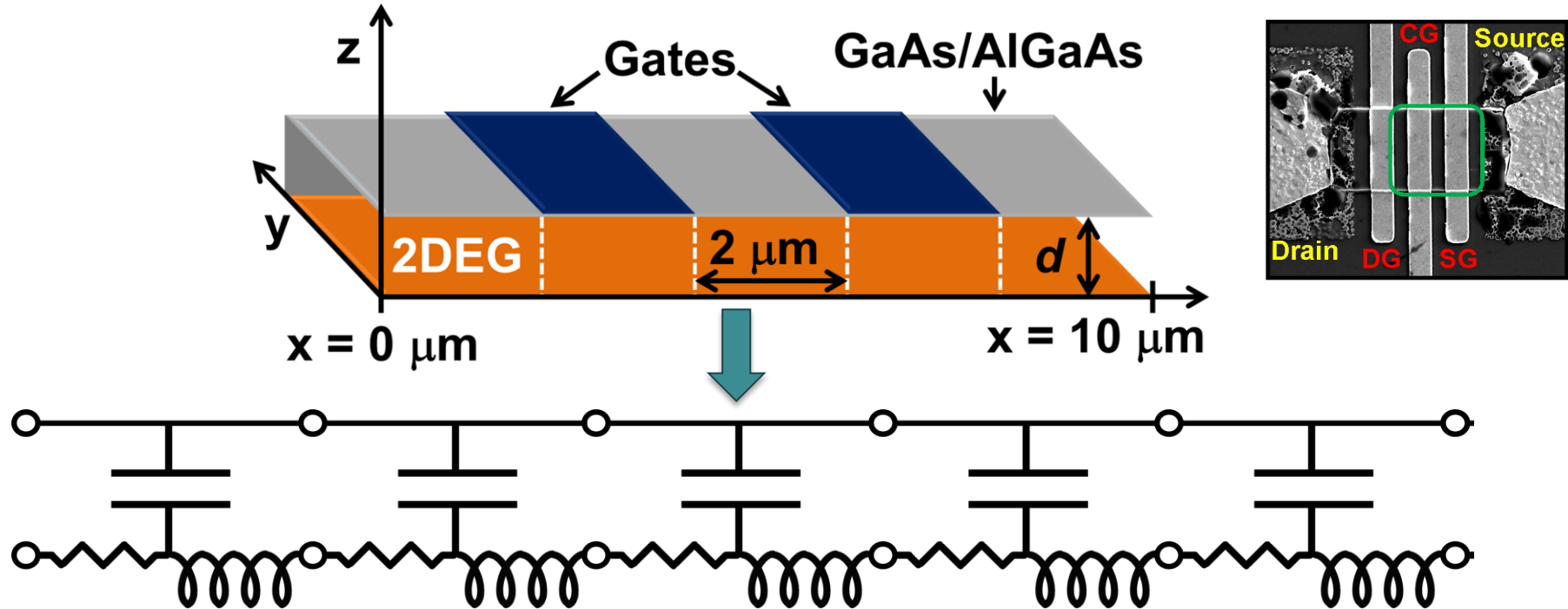
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Stepped Plasmonic TL



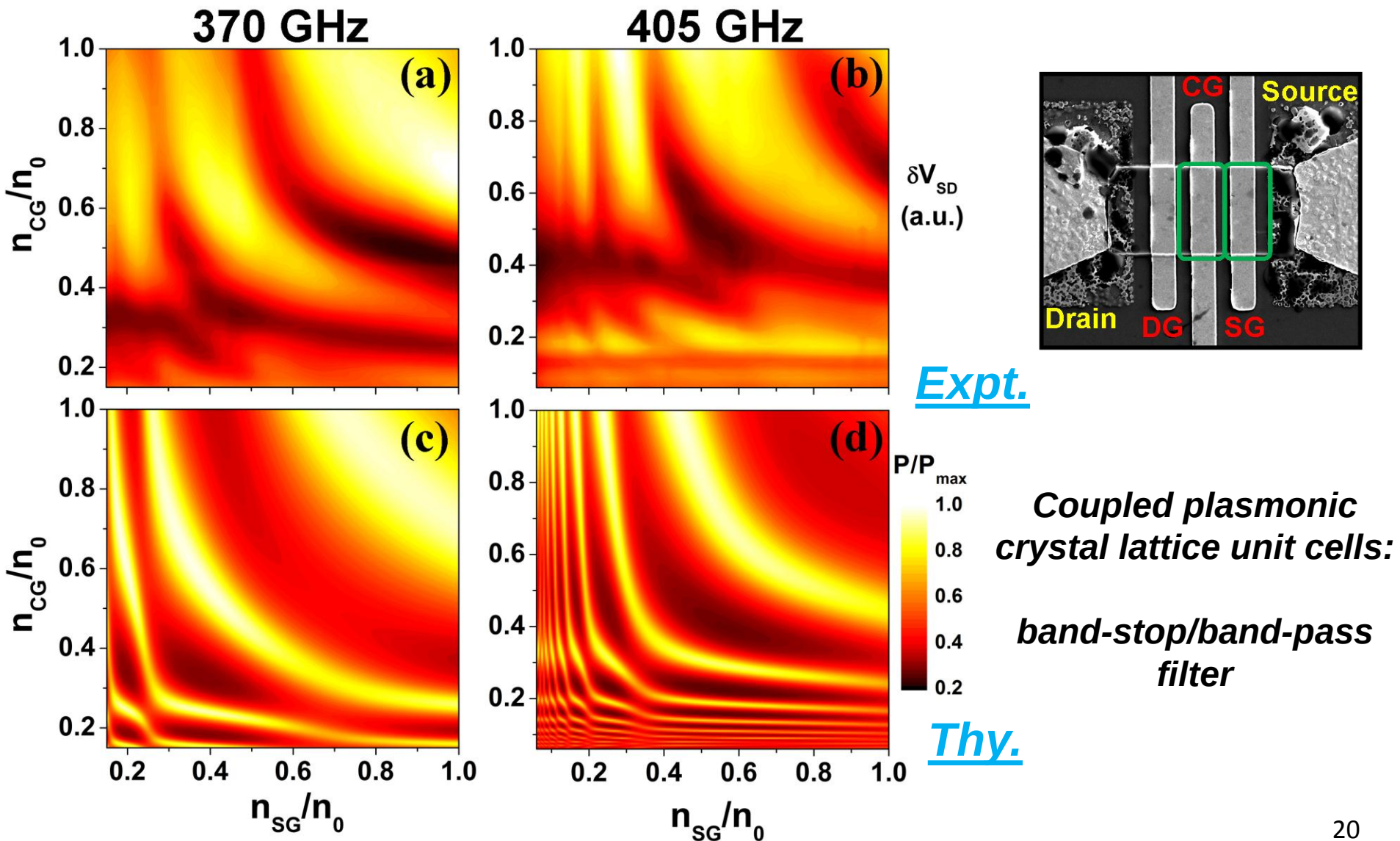
Stepped Plasmonic TL



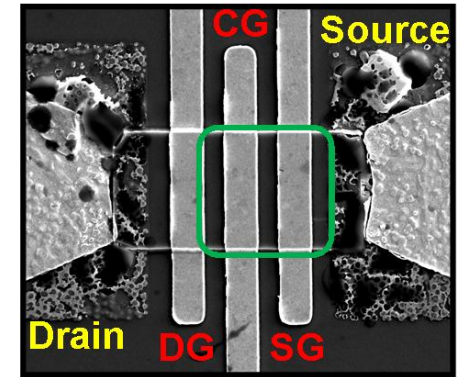
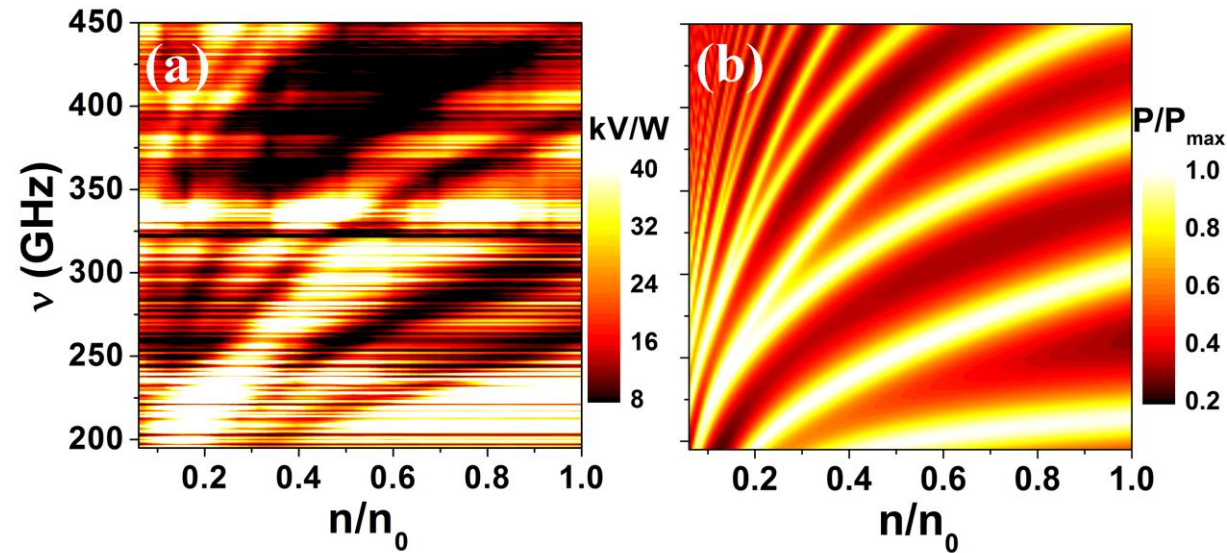
Solution for an inhomogeneous cavity with step-like boundaries between TL elements is reduced to matrix multiplication:

$$\begin{pmatrix} \tilde{V}_m \\ \tilde{I}_m \end{pmatrix} = \left[\prod_m^n \begin{pmatrix} \cos[q_m a_m] & i\tilde{Z}_m \sin[q_m a_m] \\ i\sin[q_m a_m] / \tilde{Z}_m & \cos[q_m a_m] \end{pmatrix} \right] \begin{pmatrix} \tilde{V}_{n+1} \\ \tilde{I}_{n+1} \end{pmatrix}$$

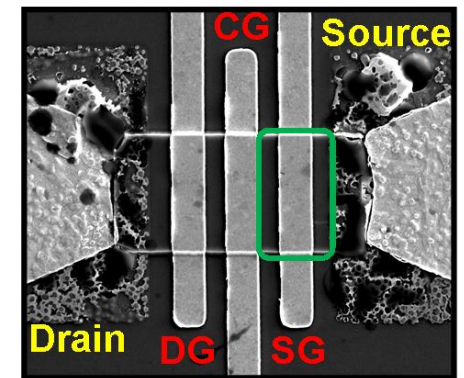
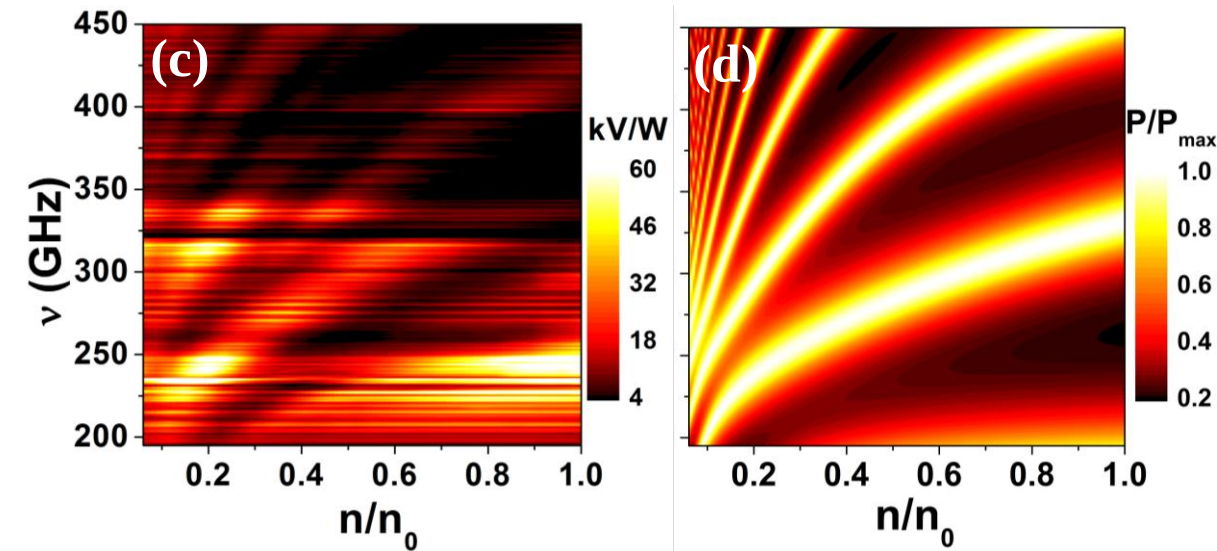
Coherent Sub-Cavity Interaction



Gate Bias Cavity Tuning



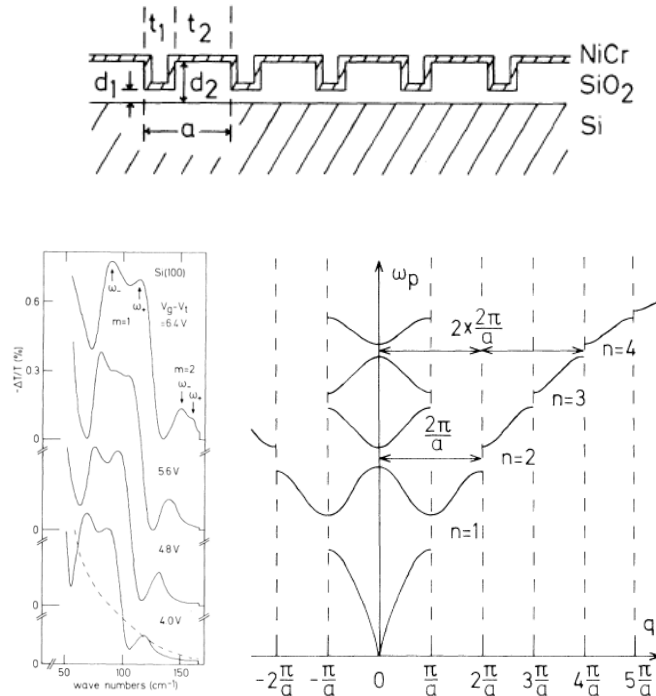
$$L_0 = 10 \mu m$$



$$L_0 = 6 \mu m$$

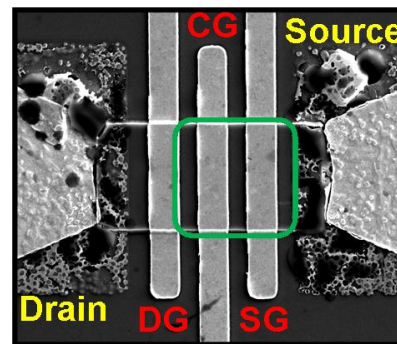
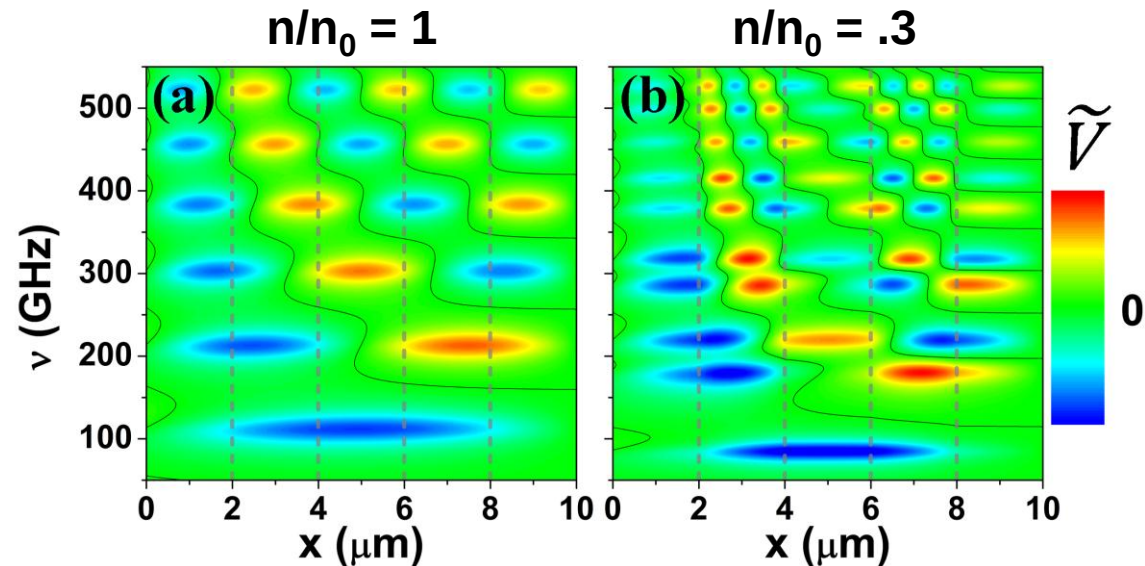
Finite Plasmonic Crystal

Quasi-Infinite Plasmonic Crystal in Si MOSFET



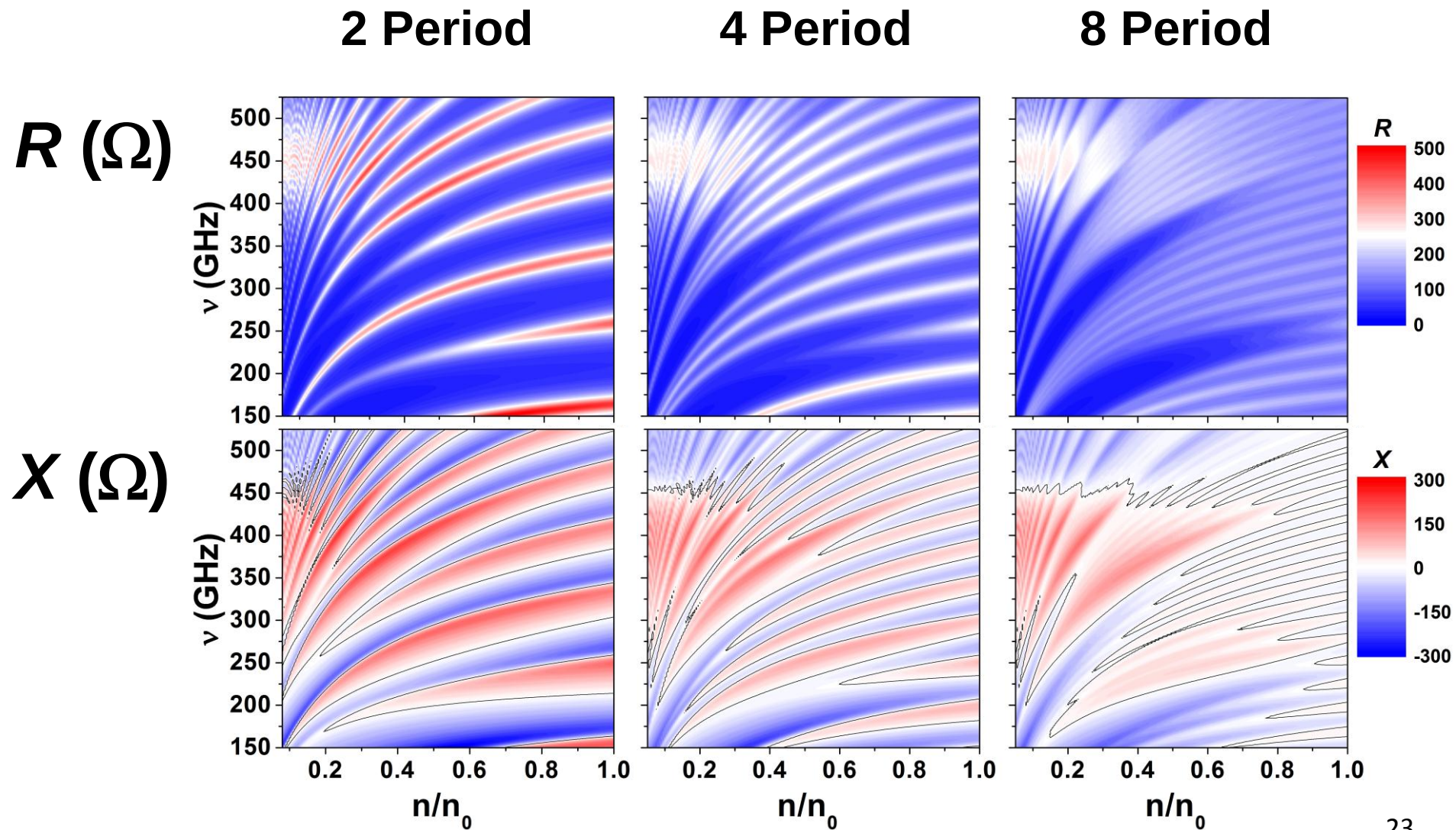
U. Mackens et. al., *PRL* **53**, 1485 (1984)

TL Model Cavity Voltage Distribution

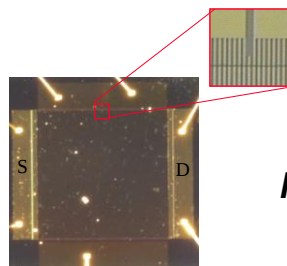


Two-period 1D plasmonic crystal embedded in 10 μm source cavity

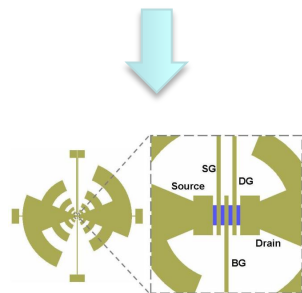
TL Model: 2, 4 & 8 Period Crystals



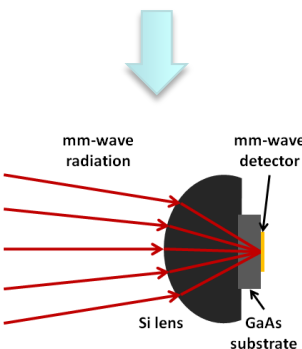
Evolution, Conclusions & Future



$R \sim 50 \text{ V/W}$
 $NEP \sim 100 \text{ nW/Hz}^{1/2}$
 $T = 20 \text{ K}$

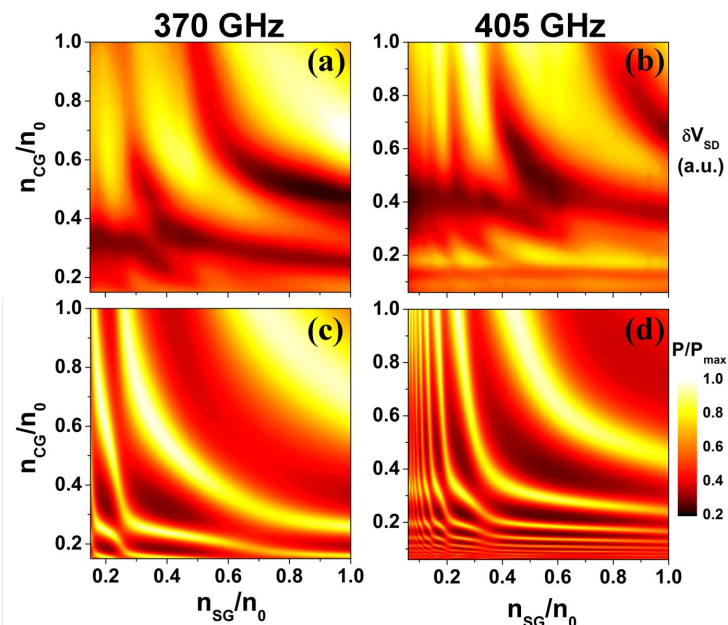
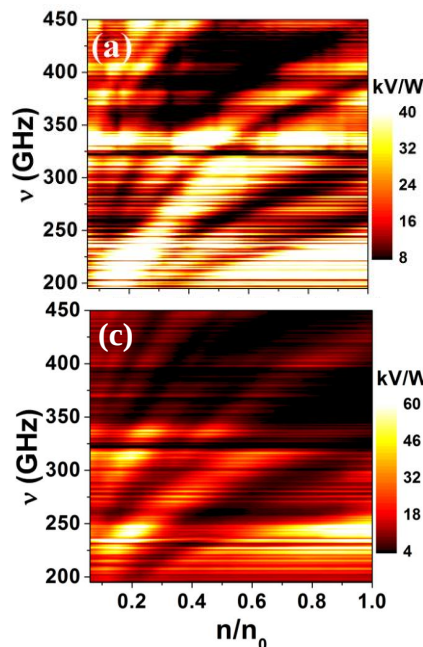


$R \sim 5 \text{ kV/W}$
 $NEP \sim 500 \text{ pW/Hz}^{1/2}$
 $T = 20 \text{ K}$

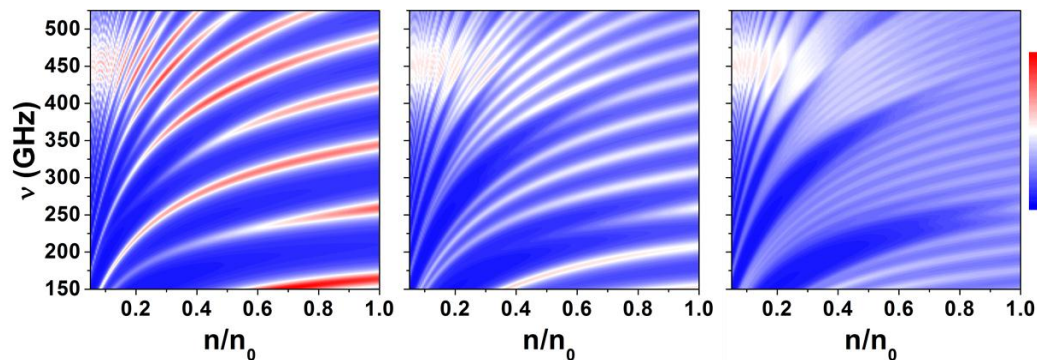


$R > 100 \text{ kV/W}$
 $NEP < 50 \text{ pW/Hz}^{1/2}$
 $T = 11 \text{ K}$

On-chip THz spectroscopy



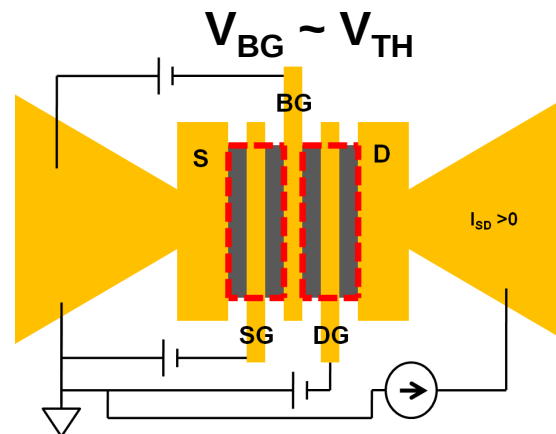
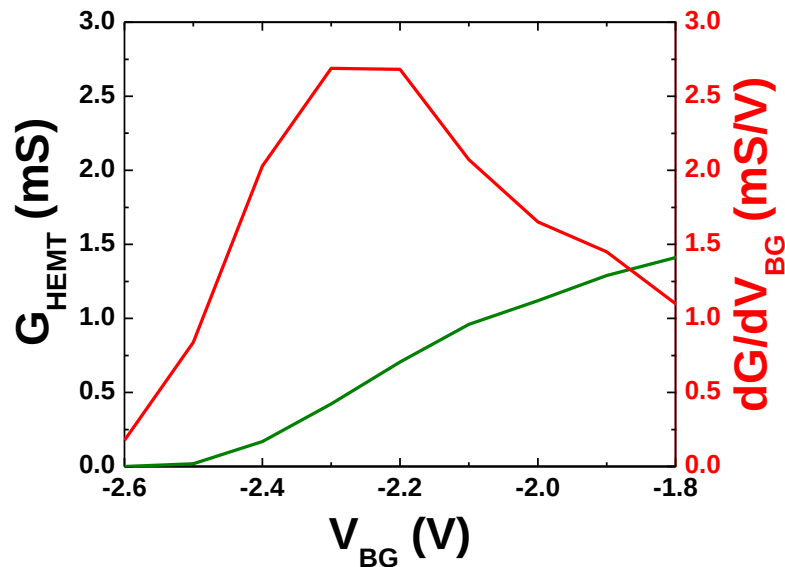
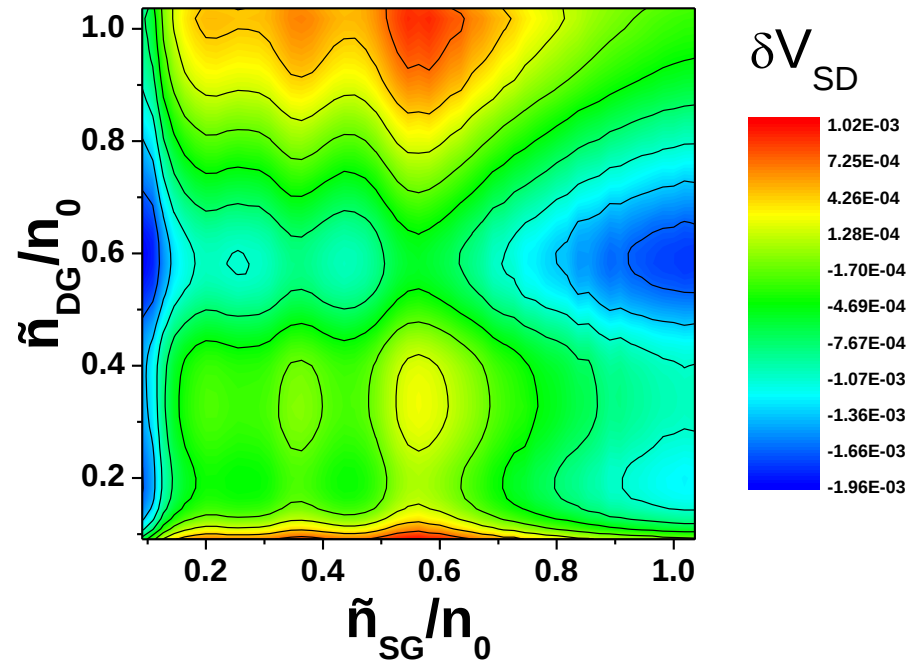
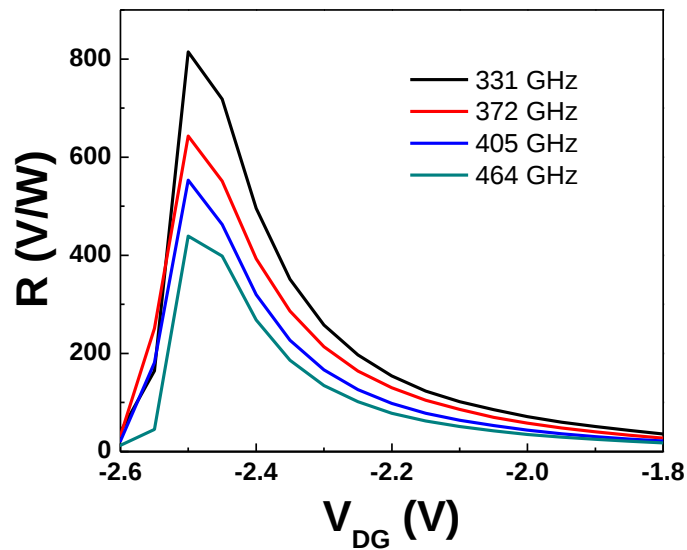
in-situ 2D plasmonic engineering



Finite plasmonic crystal

ADDITIONAL MATERIAL

Resonant Self-mixing Response

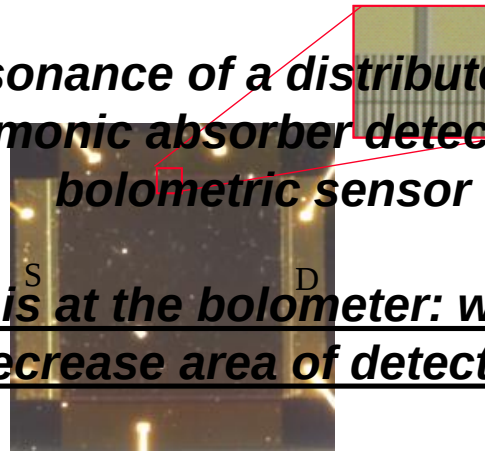


$$\begin{aligned} \langle \delta I_{SD} \rangle &= \frac{\partial G}{\partial V_G} \langle \delta V_{SD} \delta V_G \rangle \\ &= \frac{1}{2} \delta V_{THz} \delta V'_{THz} \frac{\partial G}{\partial V_G} \cos(\phi) \end{aligned}$$

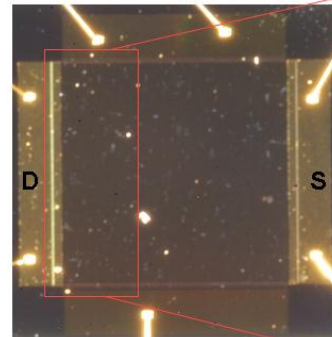
Resonant Terahertz Response

Resonance of a distributed 2D plasmonic absorber detected by bolometric sensor

Action is at the bolometer: what if we decrease area of detector?

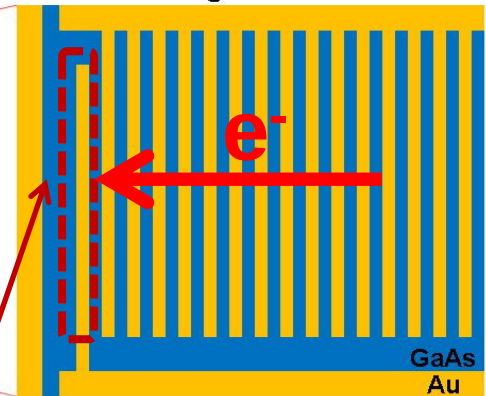


1 mm by 1 mm channel



Grating Period: 4 μm

Tuning Gate Contact



Barrier Gate Contact

Bolometric Element

$$V_{BG} < V_{TH}$$

Split-Gate Detector Photoresponse

