

Stimulated Brillouin scattering in silicon photonics

SAND2013-5631C

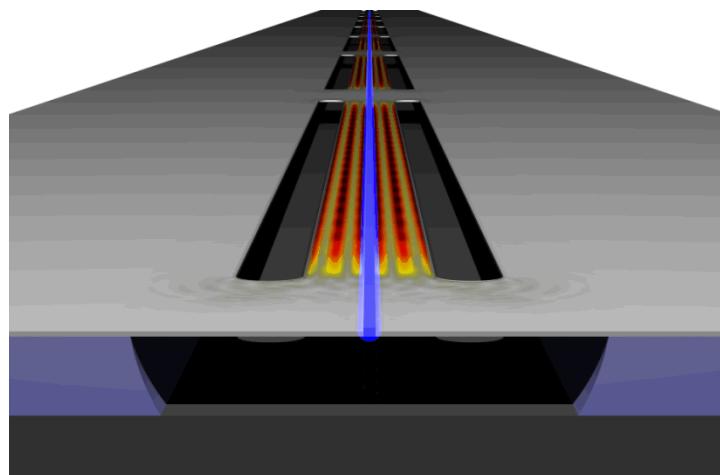
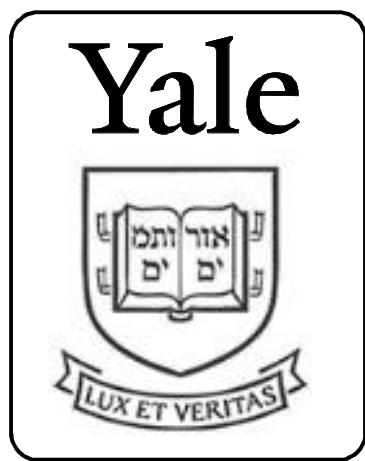
Integrated Photonics Research: IT4A

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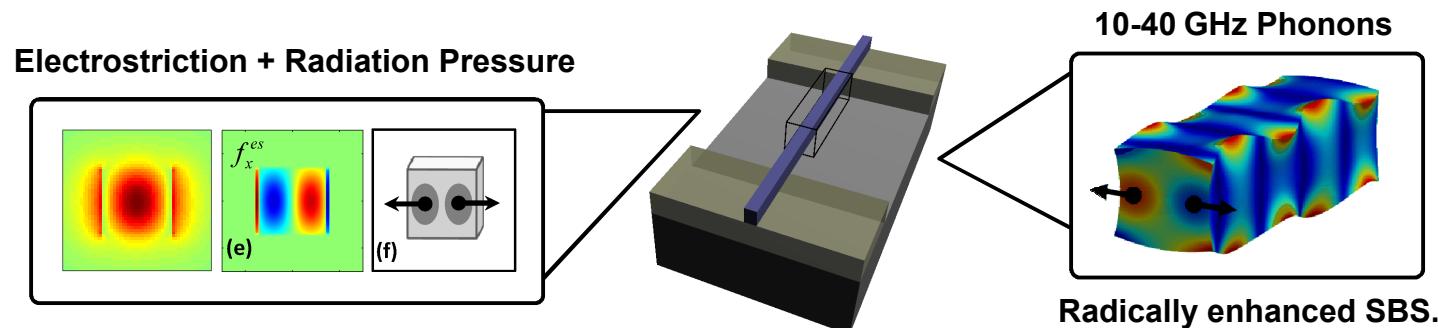
² *Applied Photonic Microsystems Group, Sandia National Labs, NM, USA*

³ *University of Texas at Austin, Electrical Engineering Department, TX, USA*

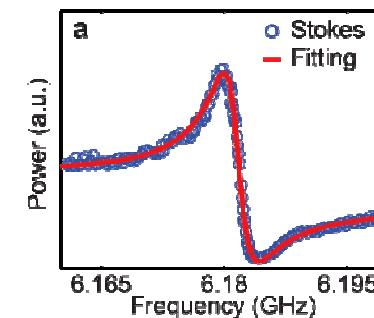
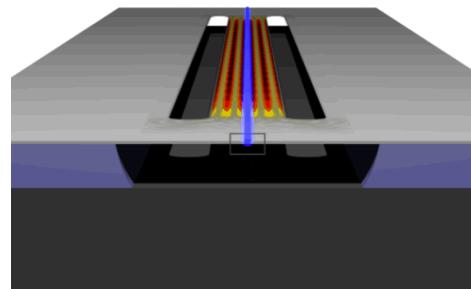
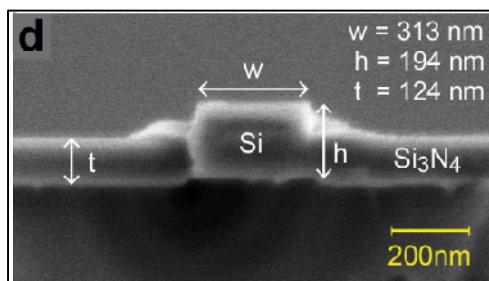


Outline of Presentation:

Part I: Physics of Stimulated Brillouin Scattering (SBS) at Nanoscales.

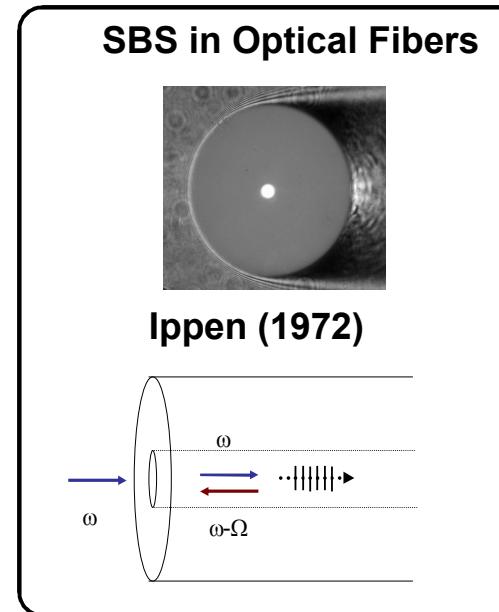
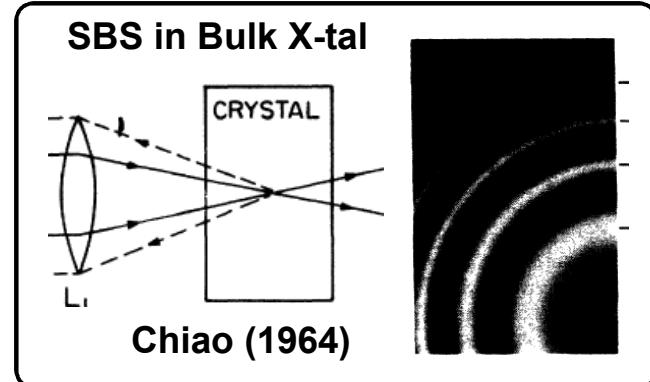


Part II: Experimental Demonstration of Forward-SBS in Silicon.

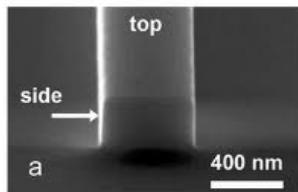


1. Bulk SBS with MASER (Chaio 1964)
2. Guided-wave SBS (Ippen, 1972)
3. Engineered SBS in fibers. (Russel 2007)

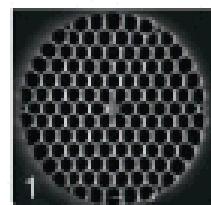
↓ Micro ← Macro



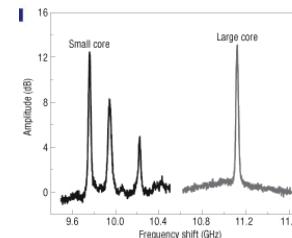
SBS in Nanophotonics?



Micro-structured Fiber

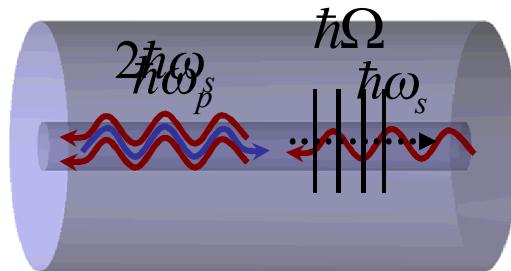


Russel (2007)

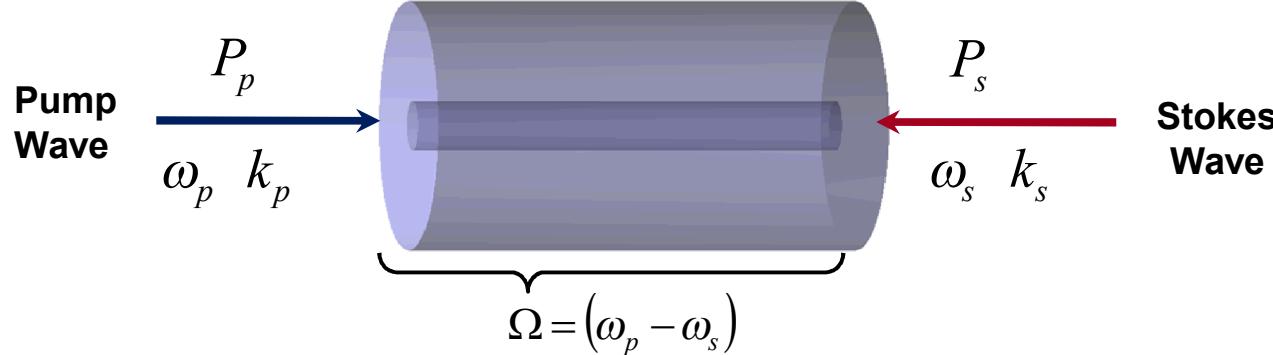


Trend: Brillouin interactions become stronger at smaller length-scales

Backward-SBS

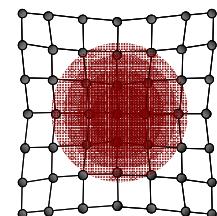


How does backward-SBS work?

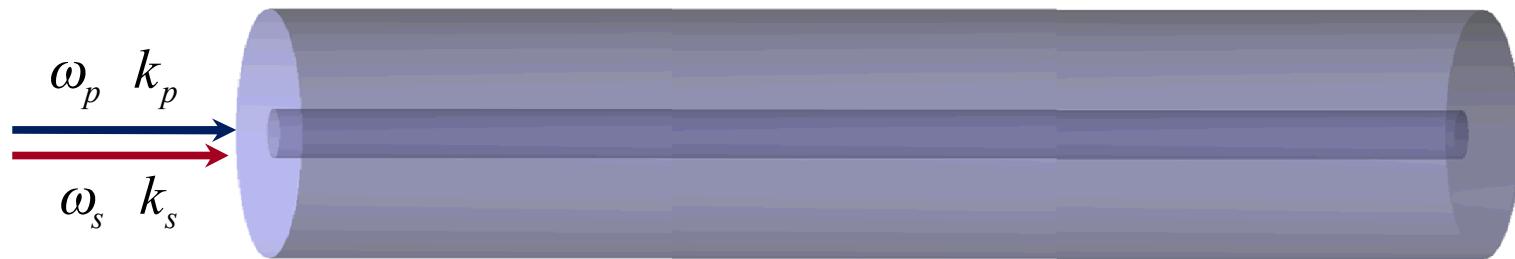


Electrostrictive forces compress medium

Electrostriction:



From dynamic material response.

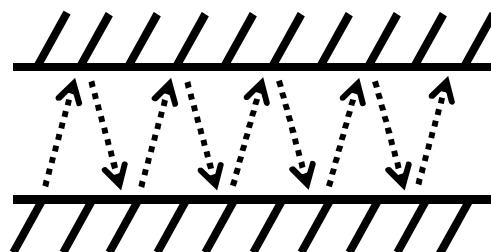


- Fringes advance at group velocity of light!

Q: Can the phase velocity of sound ever match the group velocity of light?

A: Yes, if the modes are highly dispersive:

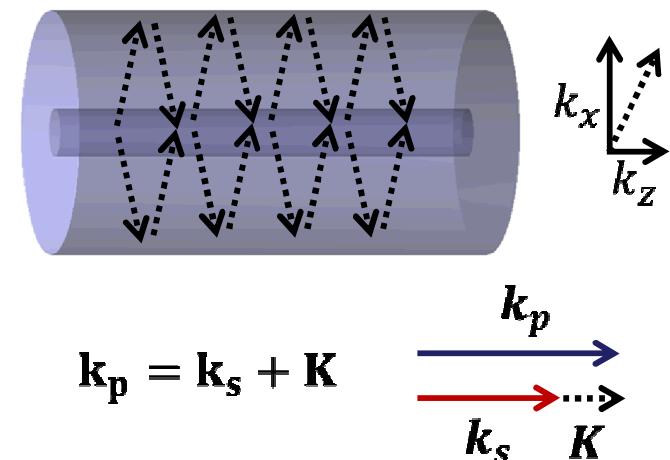
Reflecting boundaries:



Result:

$$\left\{ \begin{array}{l} v_p \approx 10^8 \text{ m/s} \\ v_g \approx 1 \text{ m/s} \end{array} \right\}$$

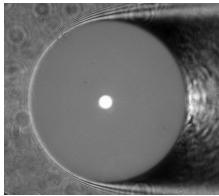
$$v_g \cdot v_p = v_a^2$$



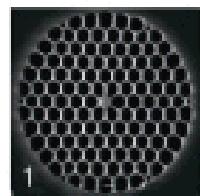
Problem: since phonon modes delocalized, forward-SBS exceedingly weak.

Various Brillouin Systems:

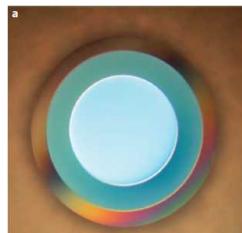
Backward-SBS Systems:



Ippen (1978)



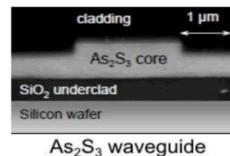
Russel Group



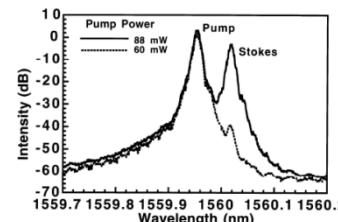
Vahala Group



Carmon Group

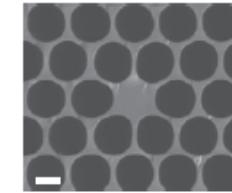


Eggleton Group

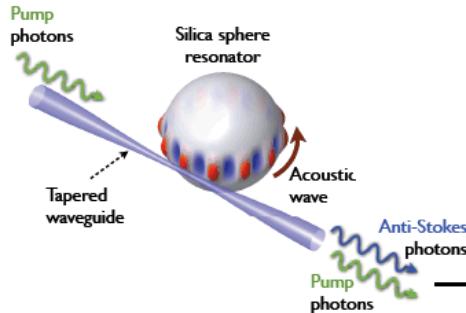


K. Abedin (2005)

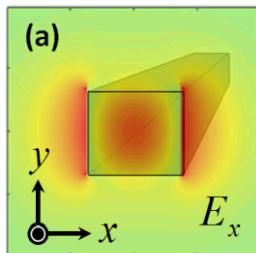
Forward-SBS Systems:



(Kang et. al Nat. Phys. 2009)



(Bahl et. al Nat. Phys. 2012)

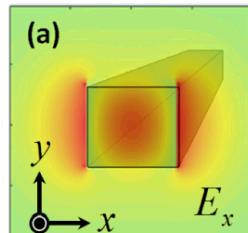


Can we engineer
SBS in silicon
nanophotonics?

Necessary Conditions:
1. Large optical forces.
2. Phonon confinement.

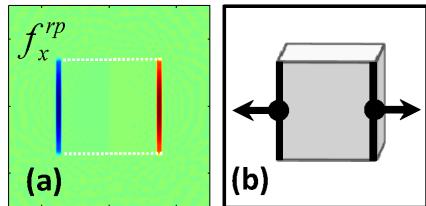
Optical Forces at Nano-scales

Silicon Waveguide.

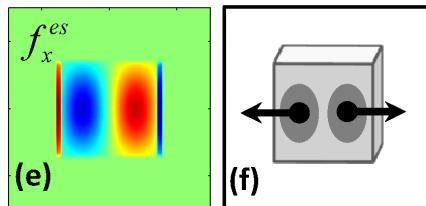


300nm x 300nm

Radiation Pressure



Electrostrictive Forces

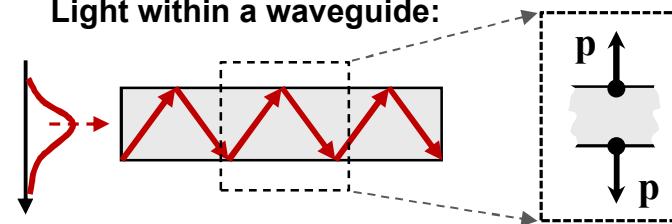


Both: Radically enhanced at nanoscales.

Radiation Pressure.

Radiation Pressure: Scattering from boundaries.

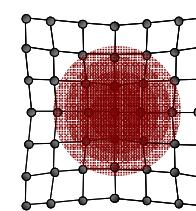
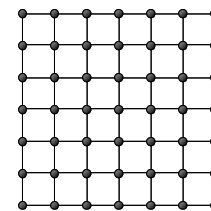
Light within a waveguide:



Entirely depends on geometry.

Electrostrictive Forces

Electrostriction: dynamic response of media to light.



Photoelastic Response

Depends primarily on material properties.

- What factors dictate the magnitude of the force?
- How can we enhance forces and SBS?

What Governs Strength of Electrostrictive Force?

Dependence on material properties:

$$\left\{ \begin{array}{ll} \epsilon_{kl} & \text{(Dielectric Tensor)} \\ p_{jkmn} & \text{(Elasto-optic Tensor)} \end{array} \right\}$$

Electrostrictive Stress in Cubic Crystal (Si):

$$\sigma_{kl}^{es} = -\frac{1}{2} \epsilon_0 \cdot n^4 \cdot p_{ijkl} \cdot E_i E_j$$

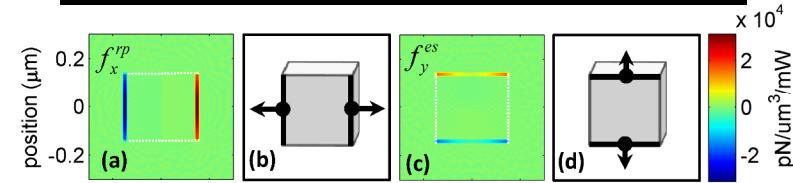
Force Density: $\mathcal{F}_j = -\partial_i \sigma_{ij}$

Important Properties of Stress/Force

1. Increases as n^4 .
2. Proportional to p_{ijkl} .
3. Sign & magnitude depends on p_{ijkl}

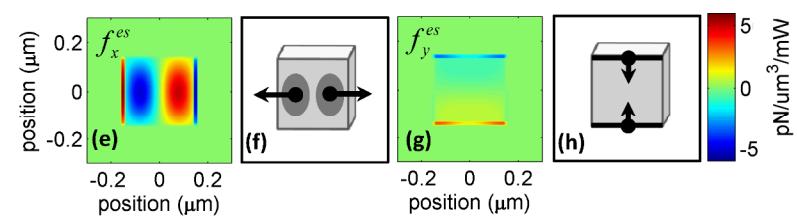
Example: Silicon waveguide.

Radiation Pressure: Si waveguide



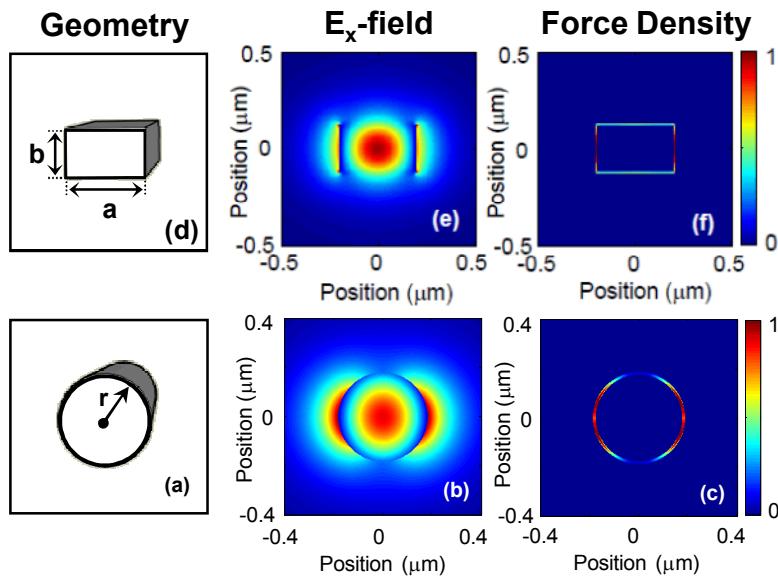
1. Force localized to boundary.
2. Directed outward.

Electrostriction: Si waveguide



1. Force distributed within volume.
2. Directed outward or inward.

[1] P. T. Rakich, P. Davids, and Z. Wang, "Tailoring Optical Forces in Waveguides Through Radiation Pressure and Electrostriction," Opt. Express 18, 14439-14453 (2010)

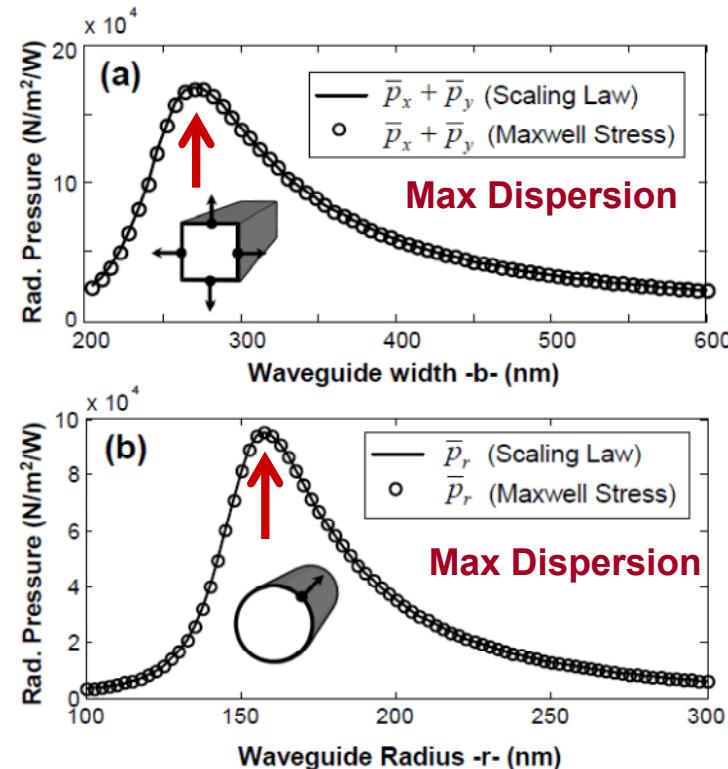


Radiation Pressure:

$$\bar{p}^{rp} = \frac{P_{opt}}{c \cdot A} \cdot (n_g - n_p)$$

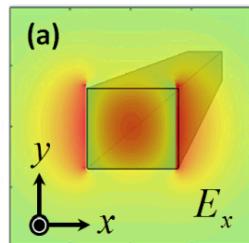


Comparison of Maxwell Stress & Scaling Law



Relation Holds for Any Dielectric Waveguide and Any Guided Mode!

[1] P. T. Rakich, Z. Wang, and P. Davids "Scaling of Optical Forces in Dielectric Waveguides: Rigorous Connection Between Dispersion and Radiation Pressure," *Optics Letters*.

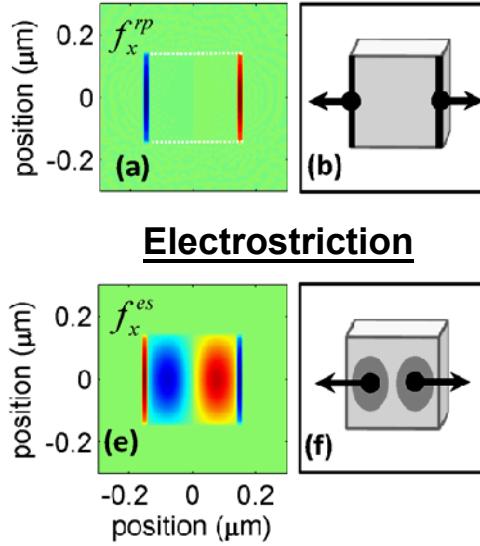


How Large are Forces?

Material	Pressure (Pwr = 100mW)
Si	$\sim 5 \times 10^4 \text{ N/m}^2$

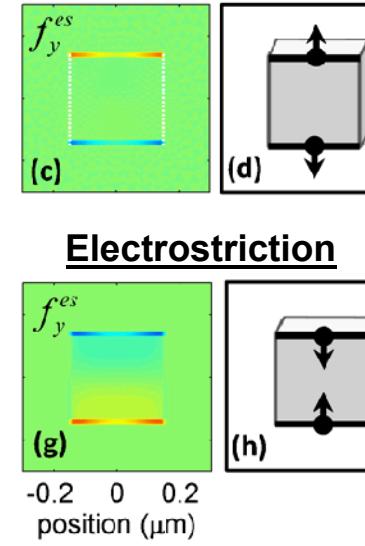
5 People standing on manhole cover

Radiation Pressure



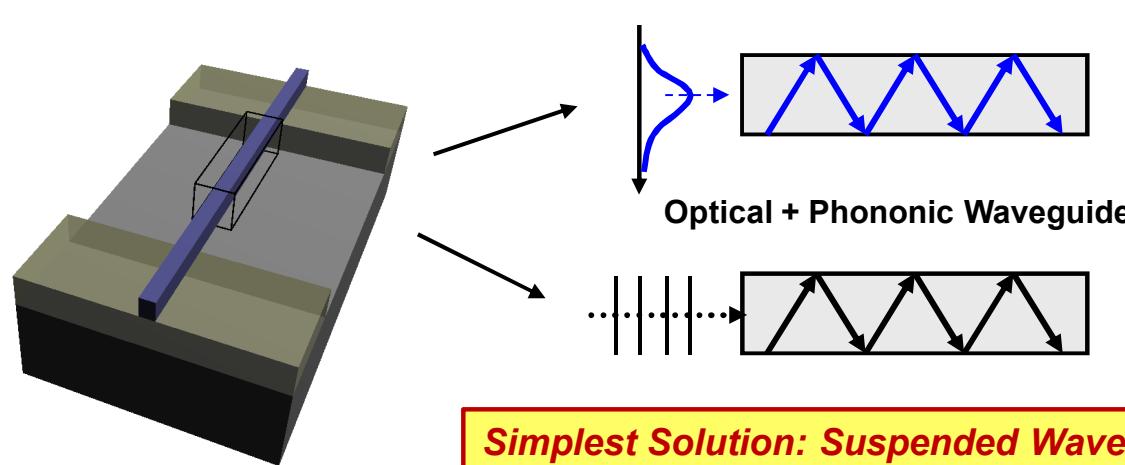
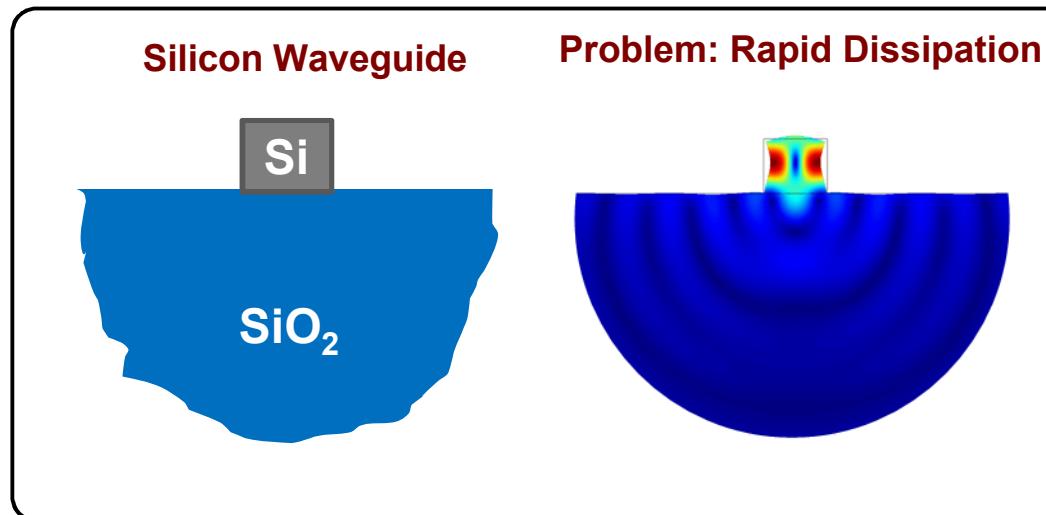
Forces Enhanced

Radiation Pressure

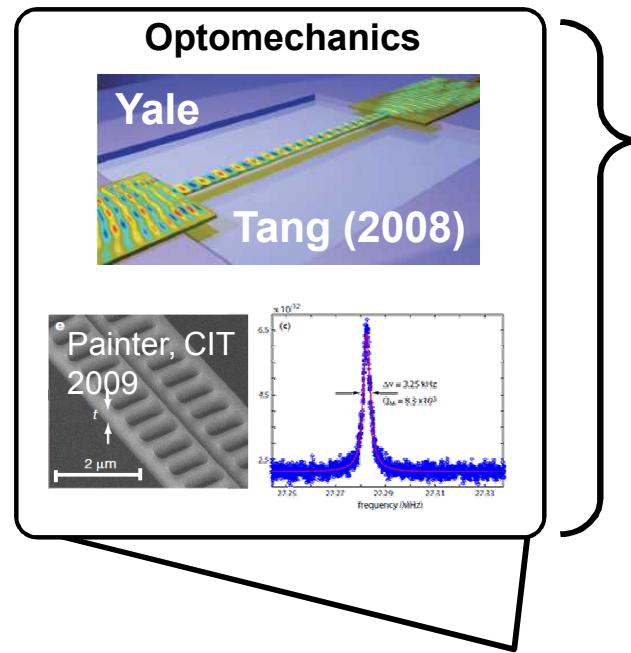


Forces Cancel

Q: Why hasn't SBS been observed in Silicon?

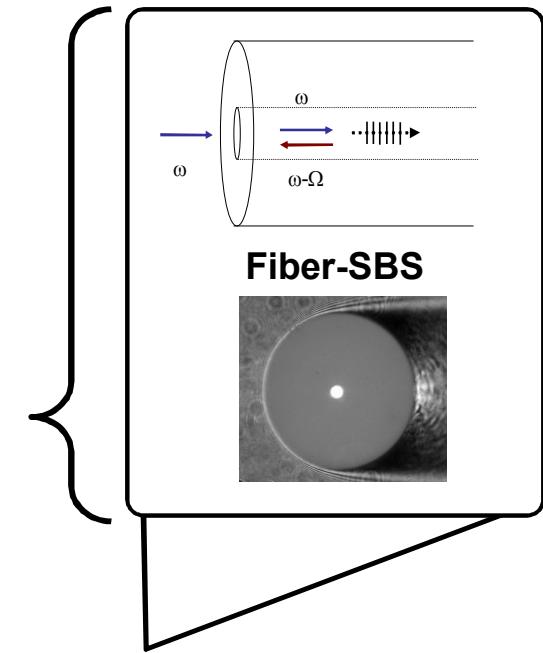


Revised Paradigm for SBS at Nanoscales:



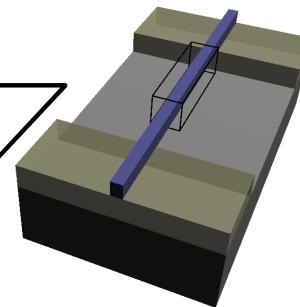
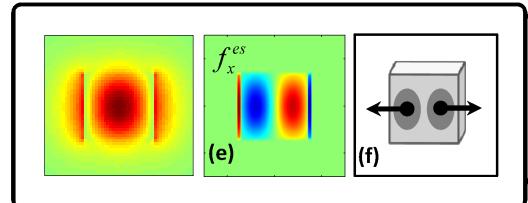
Nanoscales:
Radiation pressure
mediated photon-
phonon coupling.

Microscales:
Electrostrictively
mediated photon-
phonon coupling.

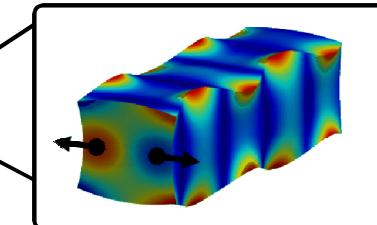


Result: Giant Enhancement of Stimulated Brillouin Scattering at Nanoscales

Electrostriction + Radiation Pressure

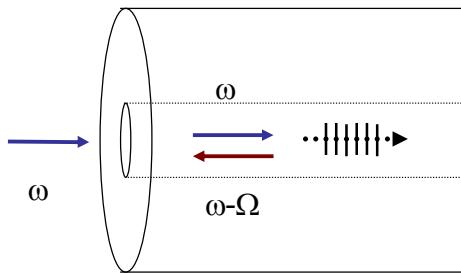


10-40 GHz Phonons



Radically enhanced SBS.

- No longer bulk material nonlinearity.
- Challenges 40-yr SBS paradigm.
- Geometry introduces new nonlinearities.

Microscale SBS Theory:

$$g_B = \frac{2\pi n^7 p_{12}^2}{c \lambda_p^2 \rho v_a \Delta V_B}$$

- Neglects radiation pressure.
- Neglects boundary induced polarization currents

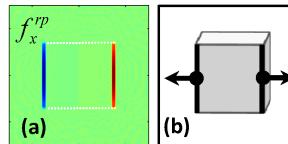
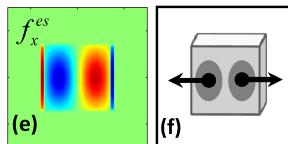
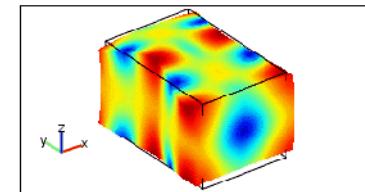
Not valid at nanoscales.

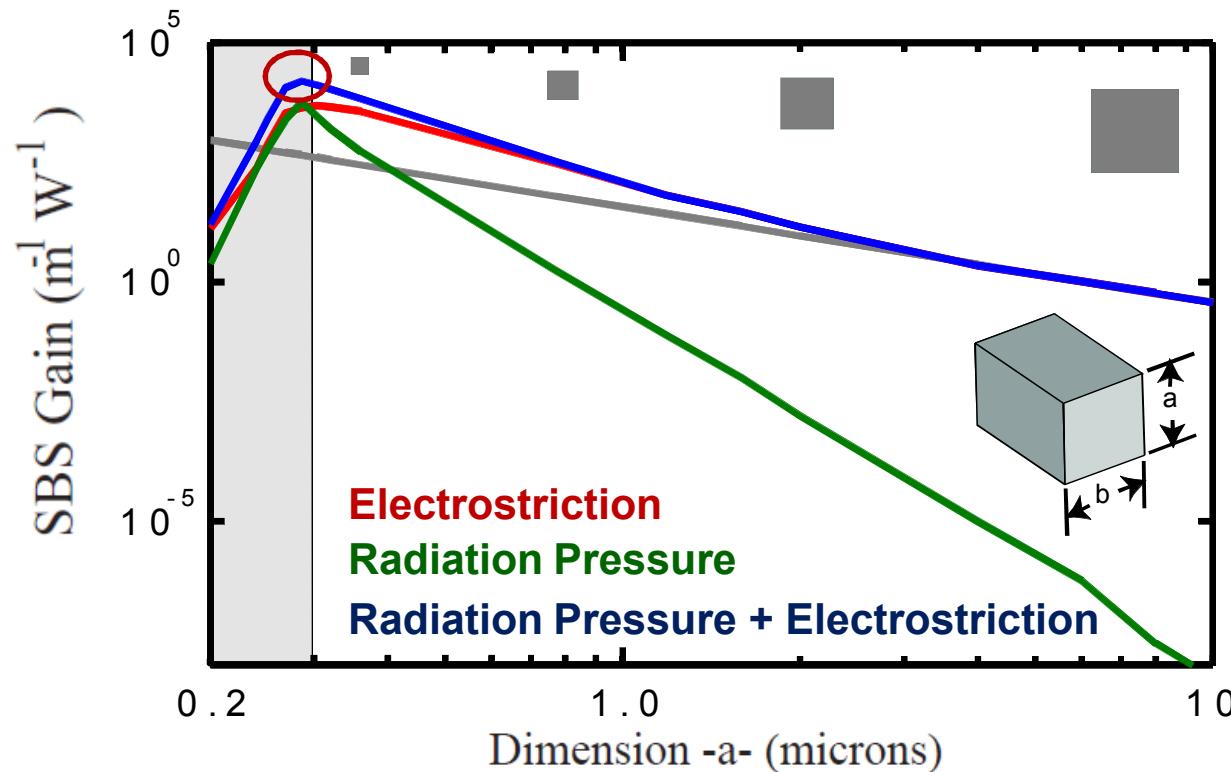
Full Vectorial Treatment: Valid at any Length-Scale**Canonical relation for SBS gain:**

$$dP_s/dz = G_B \cdot P_p P_s$$

SBS gain

$$G_B(\Omega) = \frac{1}{\delta z} \frac{\omega_s}{\Omega} \cdot \frac{1}{P_p \cdot P_s} \int_{\delta V} \langle \mathbf{f}_\Omega(\mathbf{r}, t) \cdot \dot{\mathbf{u}}(\mathbf{r}, t) \rangle \cdot dV,$$

Radiation Pressure:Time harmonic forceElectrostriction:Velocity distribution

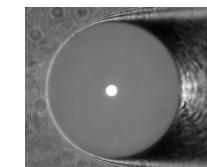


Conventional Treatment of SBS.

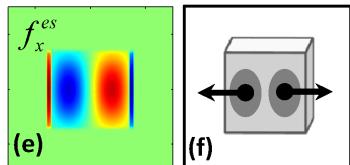
$$g_B = \frac{2\pi n^7 p_{12}^2}{c \lambda_p^2 \rho v_a \Delta \nu_B}$$

Conventional theory:

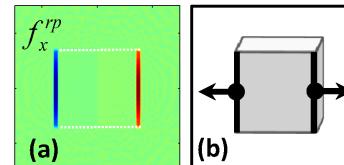
1. Silicon material properties used.
2. Perfect agreement from 2-10 microns.



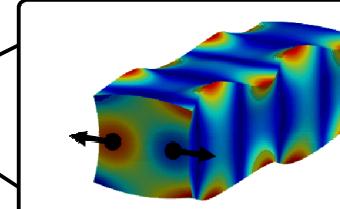
Electrostrictive forces



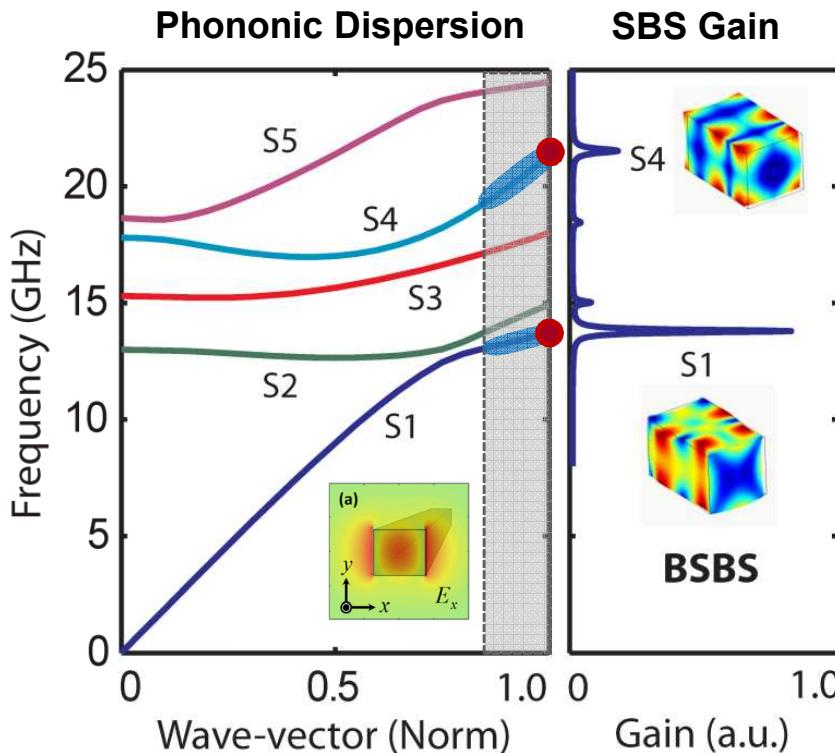
Radiation Pressure



Giant Stimulated Brillouin Scattering (SBS)



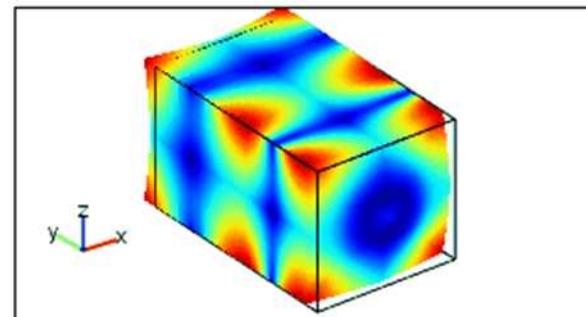
100x larger than predicted by scalar theories



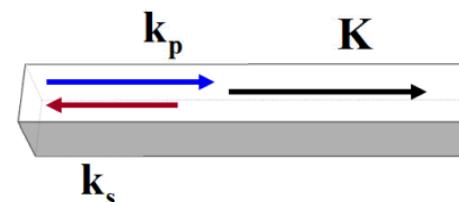
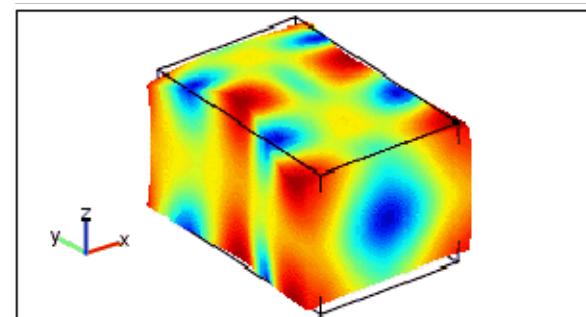
Nano-optomechanical backward-SBS:

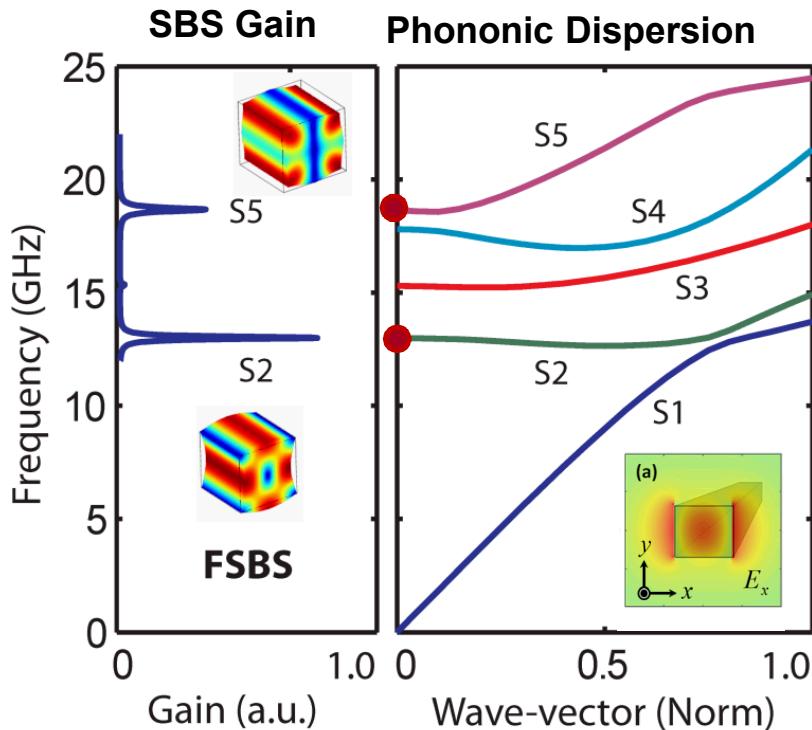
1. Narrow-band ultra-high frequency phonons.
2. Gain is 10^4 x Larger than in Fibers.
3. SBS occurs for any optical wavelength.
4. 20% frequency tunable phonon emission.

Excitation: 21.6 GHz Phonons

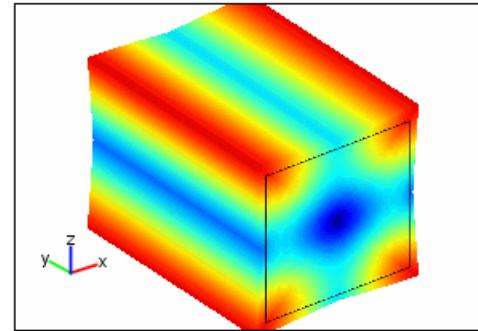


Excitation: 13.8 GHz Phonons

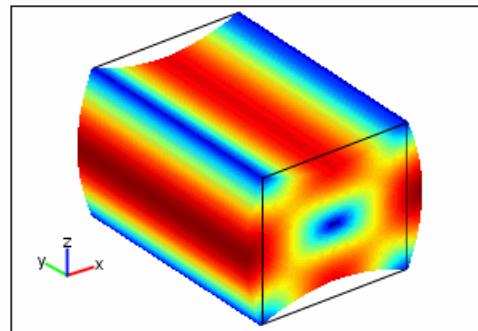




Excitation: 18.6 GHz Phonons

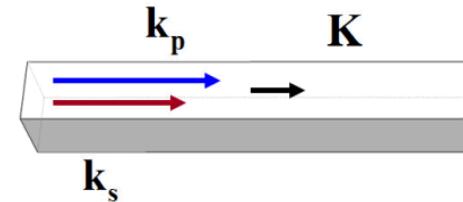


Excitation: 13.0 GHz Phonons

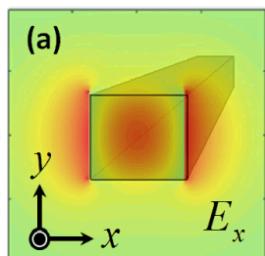


Nano-optomechanical Forward-SBS:

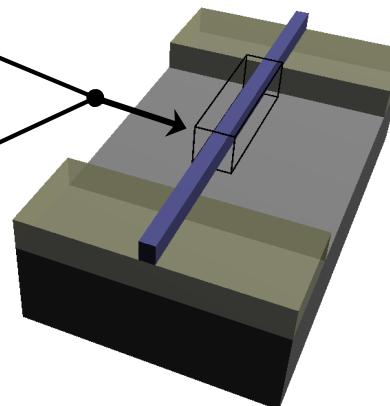
1. Generally forbidden in guided wave-systems.
2. Negligible anchoring losses \rightarrow intrinsic Mech. Q.
3. Ultra-low threshold parametric oscillation possible.



Yale Narrow-Band High-frequency RF Oscillators:

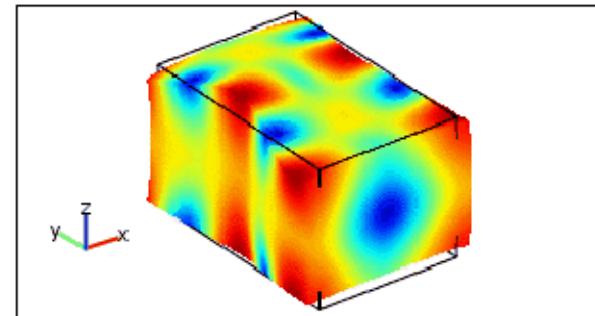


Guided mode within suspended dielectric waveguide. (300x300nm)



Suspended waveguide: $L = 100$ microns

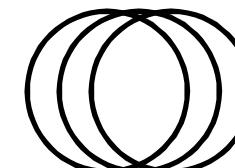
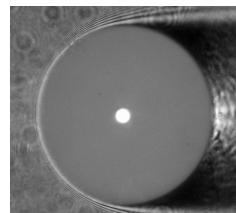
Excitation: 13.8 GHz Phonons



= Equivalent nonlinearity of 10-1000 meters of fiber

New Technologies:

- Slow light & information storage
- Narrow-band signal amplification and lasers.
- Tailorable nonlinear susceptibilities.



Fiber optic: $L = 10-1000$ meters

P. Rakich, C. Reinke, R. Camacho, P. Davids, and Z. Wang, "Giant Enhancement of Stimulated Brillouin Scattering in the Subwavelength Limit," *Physical Review X*, vol. 2, no. 1, pp. 1-15, Jan. 2012.

Slow and Fast Light via Backward Stimulated Brillouin Scattering

[1] K. Y. Song, K. S. Abedin, K. Hotate, M. González Herráez, and L. Thévenaz, "Highly efficient Brillouin slow and fast light using As₂Se₃ chalcogenide fiber," *Optics Express*, vol. 14, no. 13, pp. 5860-5865, 2006.

[2] Y. Okawachi et al., "Tunable all-optical delays via Brillouin slow light in an optical fiber," *Physical review letters*, vol. 94, no. 15, p. 153902, 2005.

[3] Ravi Pant, Adam Byrnes, Christopher G. Poulton, Enbang Li, Duk-Yong Choi, Steve Madden, Barry Luther-Davies, and Benjamin J. Eggleton, "Photonic-chip-based tunable slow and fast light via stimulated Brillouin scattering," *Opt. Lett.* 37, 969-971 (2012)

[4] R. W. Boyd, "Slow and fast light: fundamentals and applications," *Journal of Modern Optics*, vol. 56, no. 18-19, pp. 1908-1915, Oct. 2009.

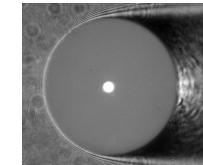
Narrow Band SBS Lasers

[1] L. F. Stokes, M. Chodorow, and H. J. Shaw, "All-fiber stimulated Brillouin ring laser with submilliwatt pump threshold," *Optics letters*, vol. 7, no. 10, pp. 509-11, Oct. 1982.

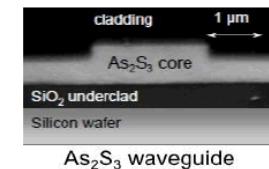
RF Signal Processing and Acousto-Optics via Forward-SBS

[1] M. S. Kang, A. Nazarkin, A. Brenn, and P. S. J. Russell, "Tightly trapped acoustic phonons in photonic crystal fibres as highly nonlinear artificial Raman oscillators," *Nature Physics*, vol. 5, no. 4, pp. 276-280, Mar. 2009.

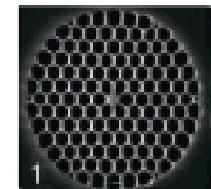
[2] M. S. Kang, a. Butsch, and P. S. J. Russell, "Reconfigurable light-driven opto-acoustic isolators in photonic crystal fibre," *Nature Photonics*, no. August, pp. 1-5, Aug. 2011.



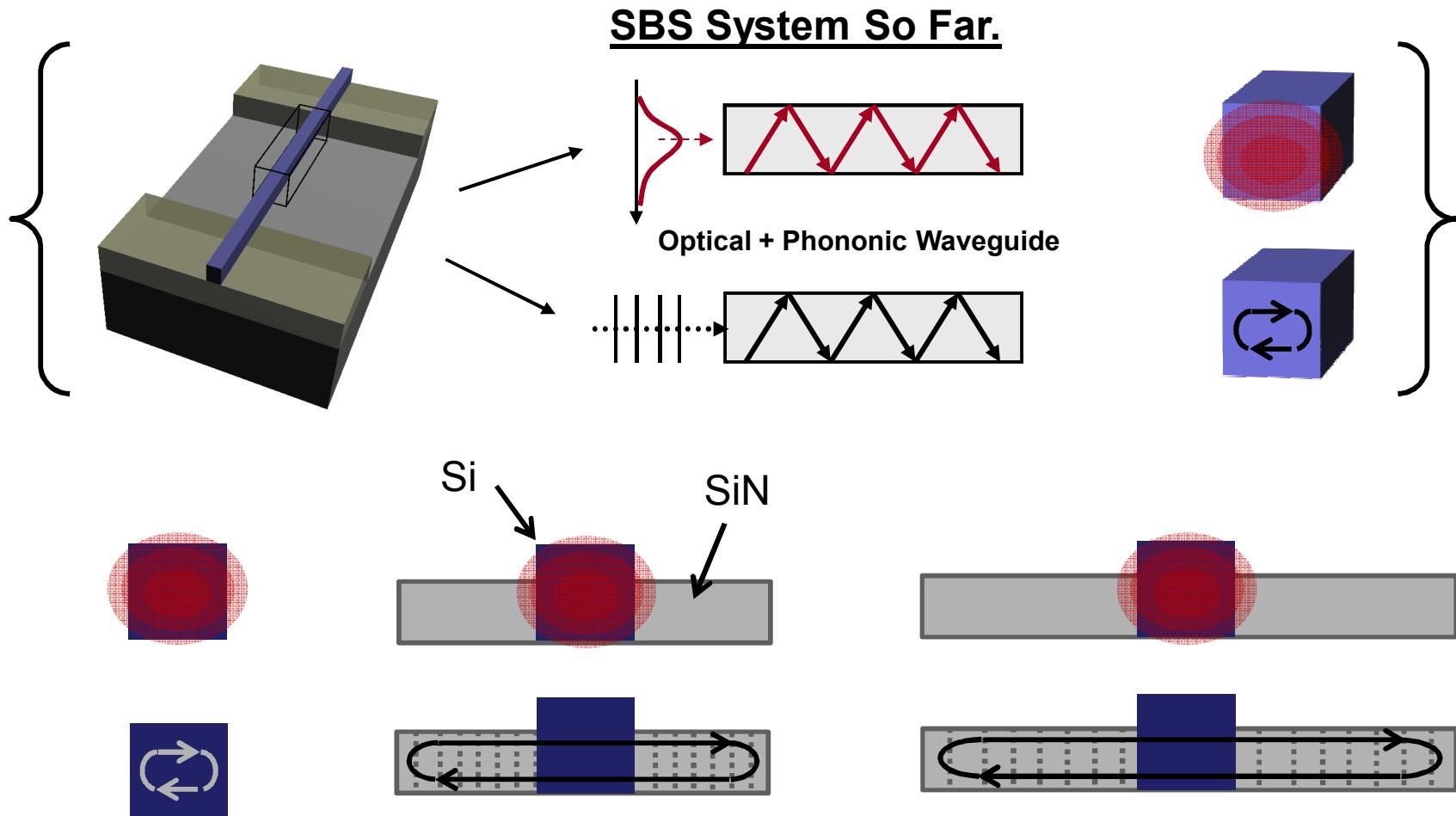
Eggleton



Carmon

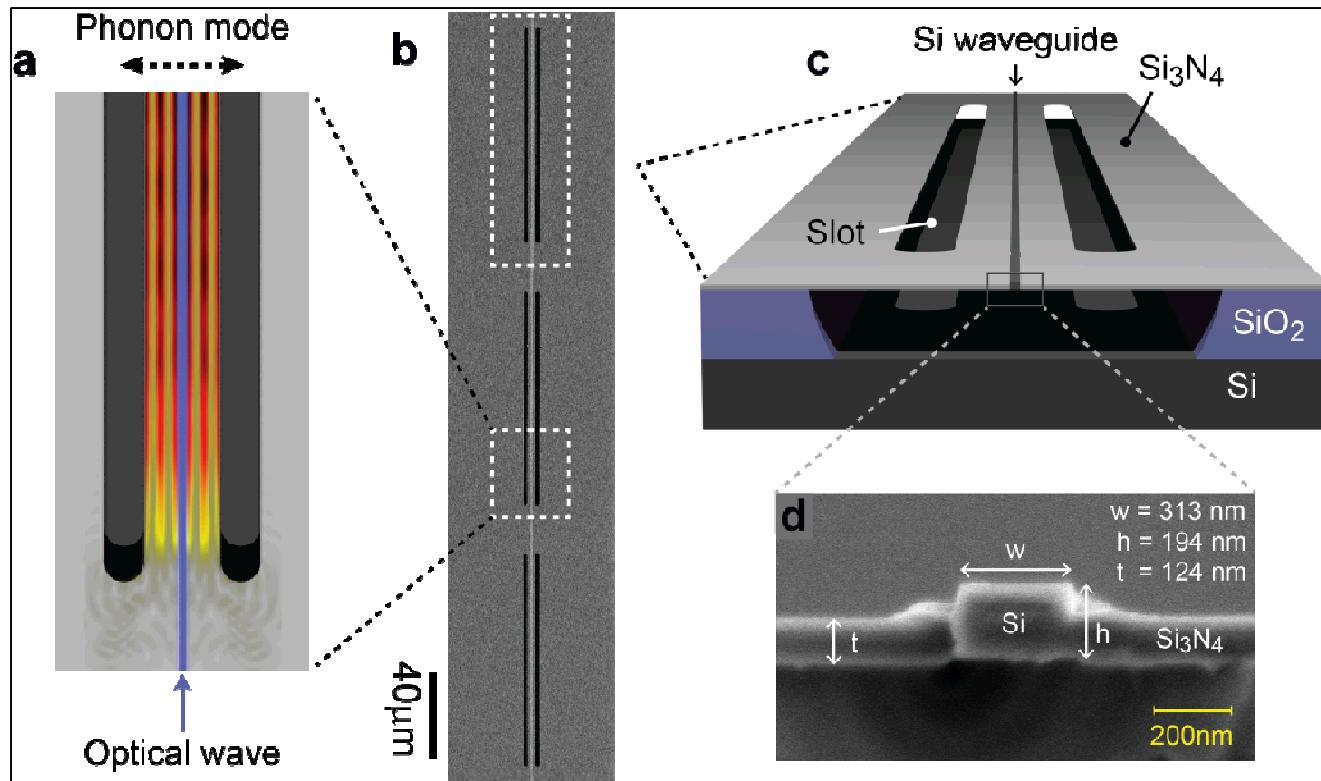


Russel



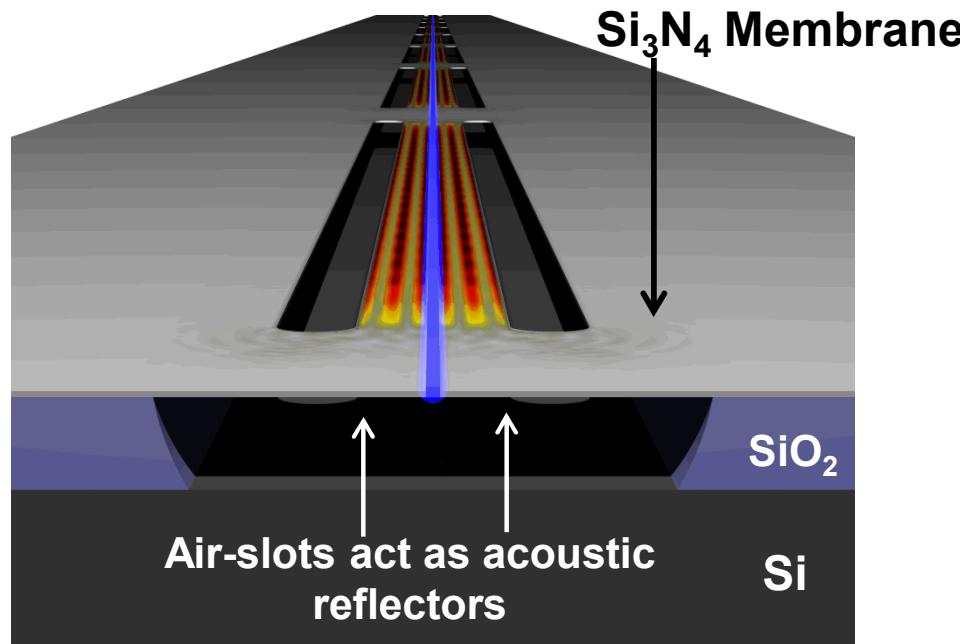
Experimental system: Can tailor phononic guidance independent of optical forces.

Fabricated Brillouin Active Waveguides.

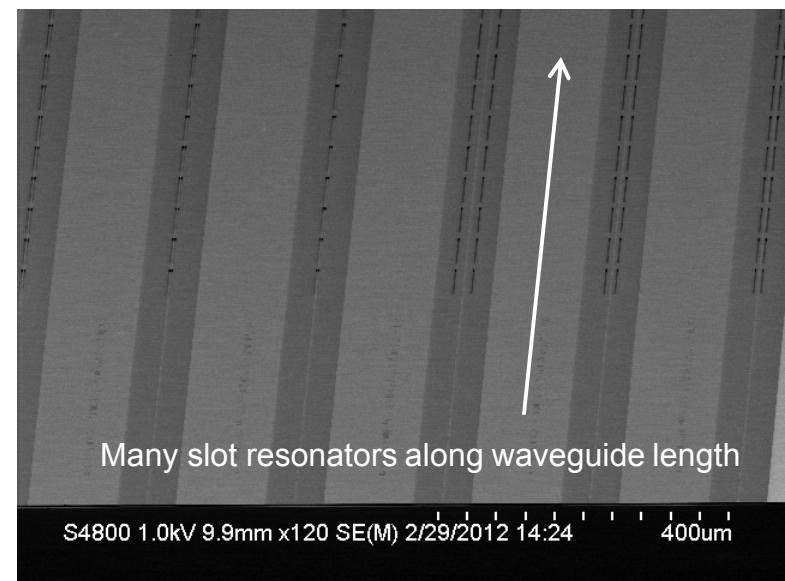


- 5mm-long devices realized with a concatenation of 26 Brillouin active waveguide segments.
- Nonlinear response coherently adds to yield tremendous aggregate Brillouin gain.

Brillouin Active Silicon Waveguides:



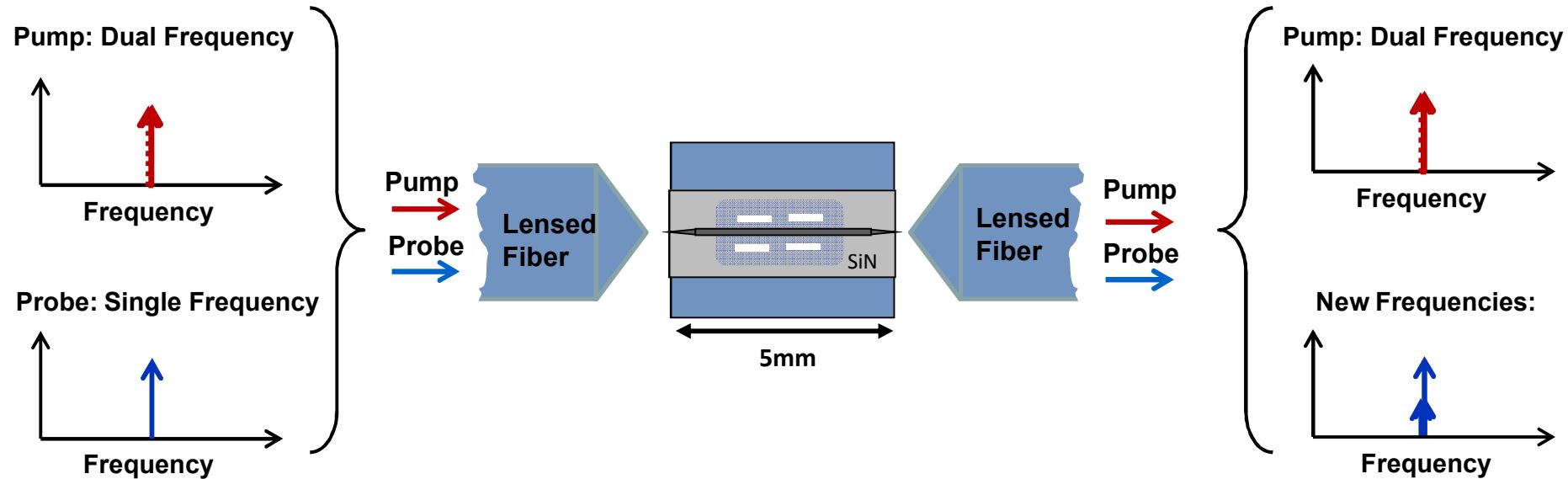
SEM of fabricated waveguides



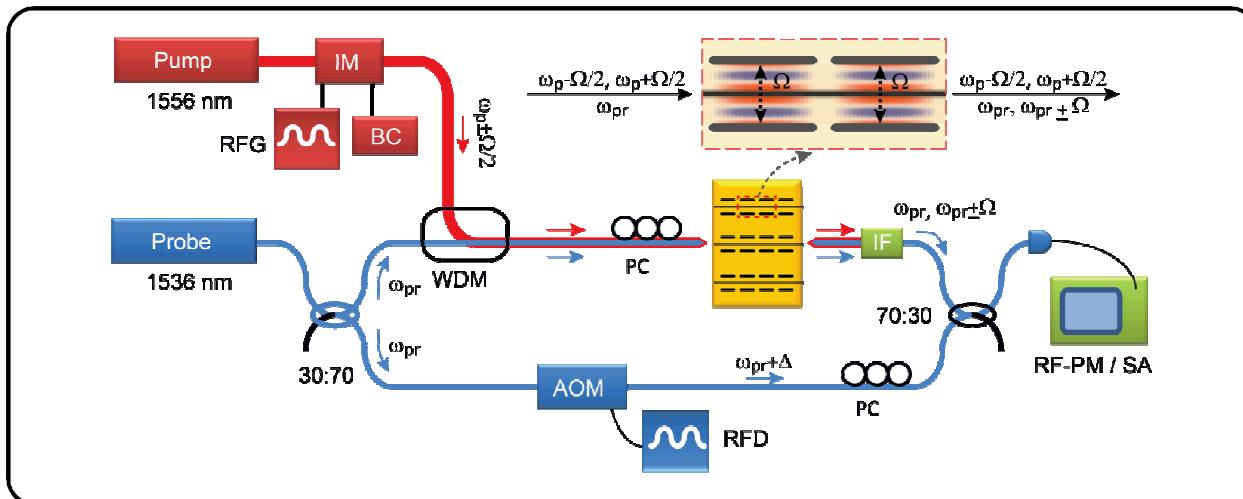
- Light guided via total internal reflection.
- 5mm-long devices realized with a concatenation of 26 Brillouin active waveguide segments.

- Before we present any data, we should say a few things about:
 1. The measurement apparatus.
 2. Coherent addition of nonlinearities.

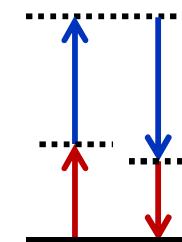
SBS Measurements:



Heterodyne cross-phase modulation apparatus.



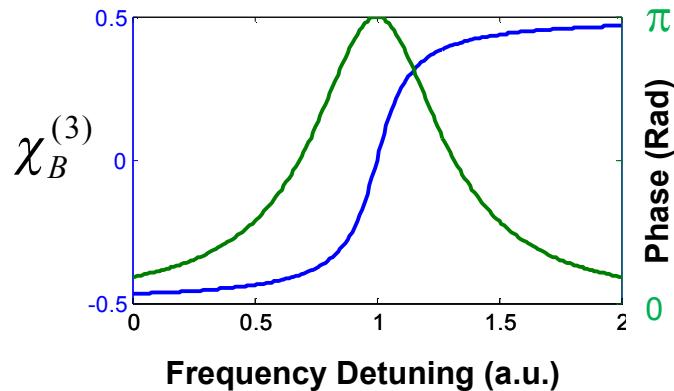
Measurement of
nonlinear susceptibility



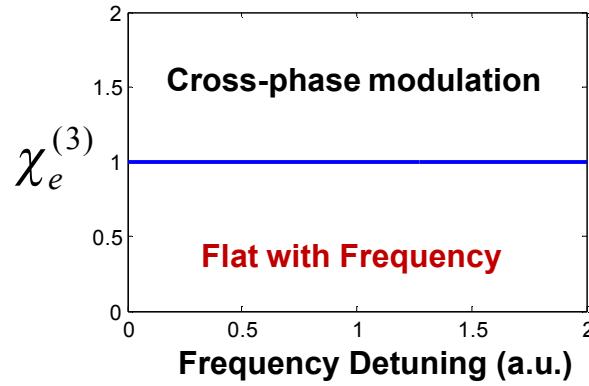
Four-Wave Mixing Transition

Total Nonlinear Susceptibility:

Resonant Brillouin Susceptibility

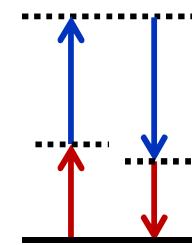


Non-Resonant Electronic Response

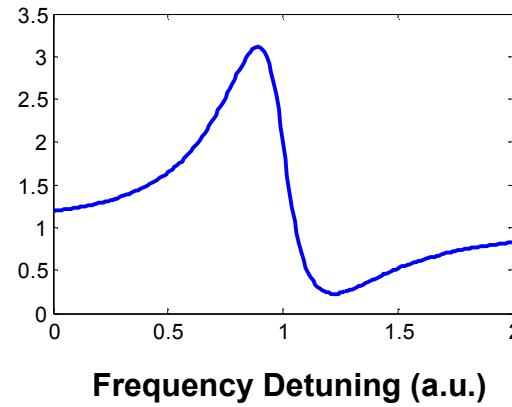


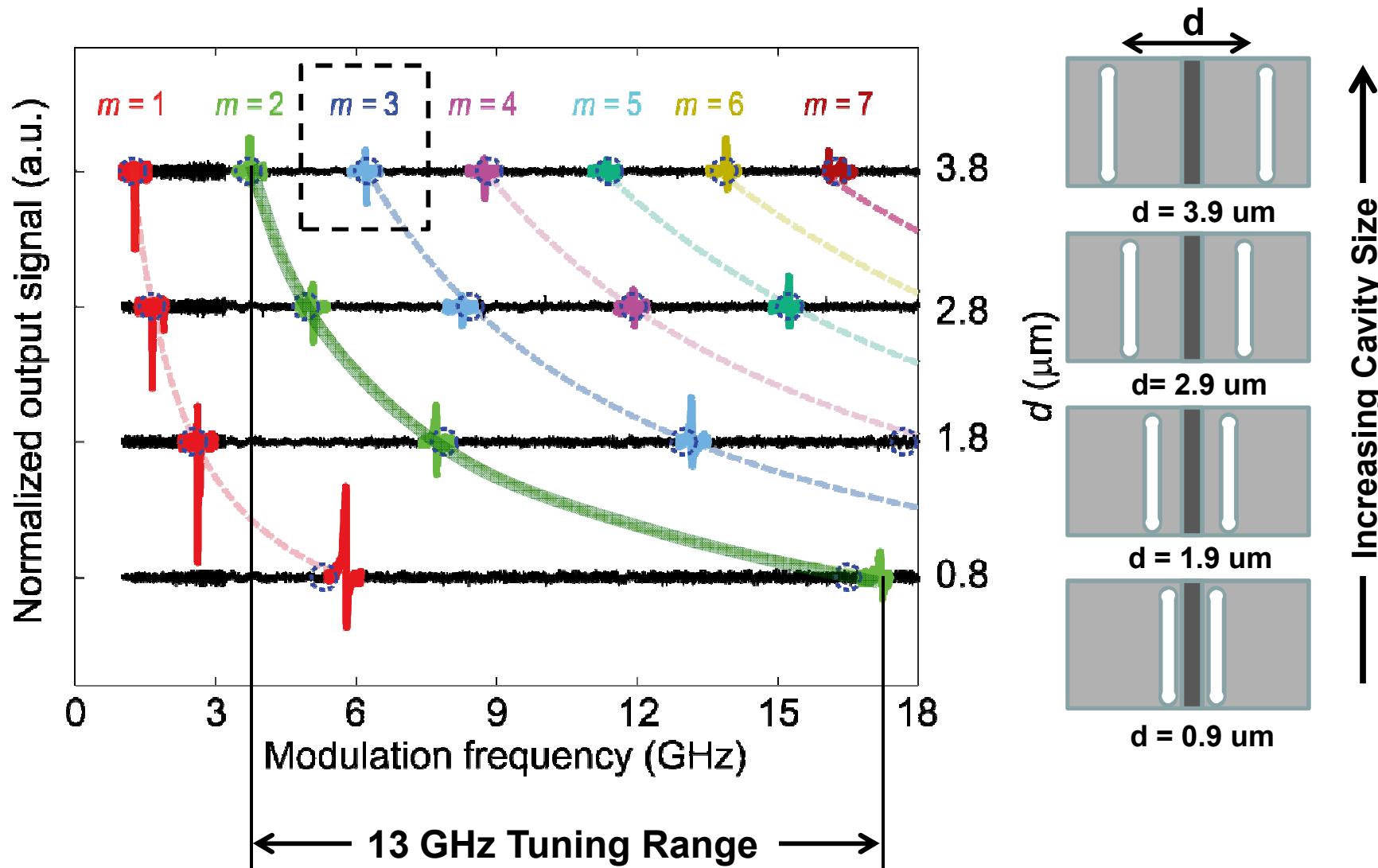
+

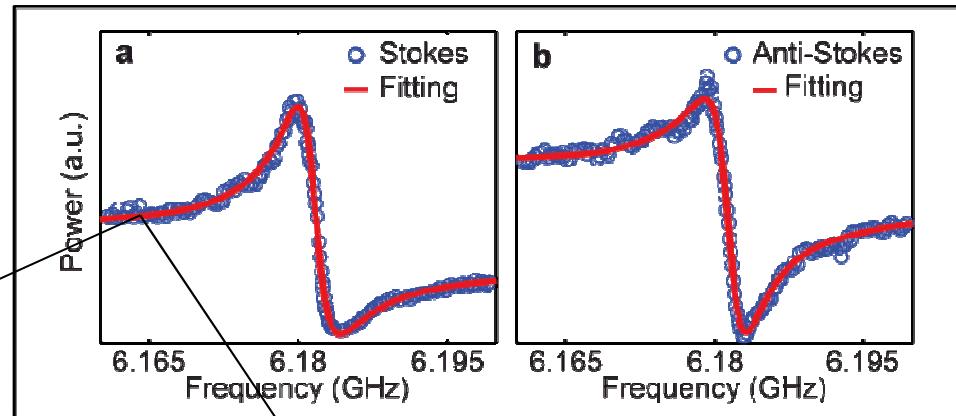
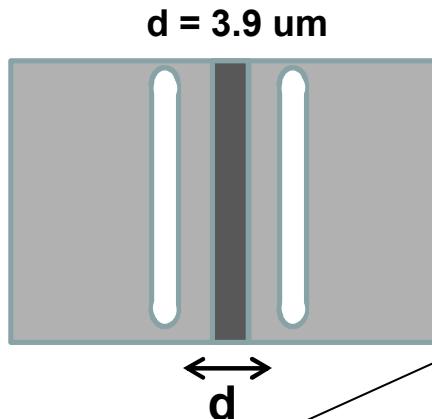
FWM Transition



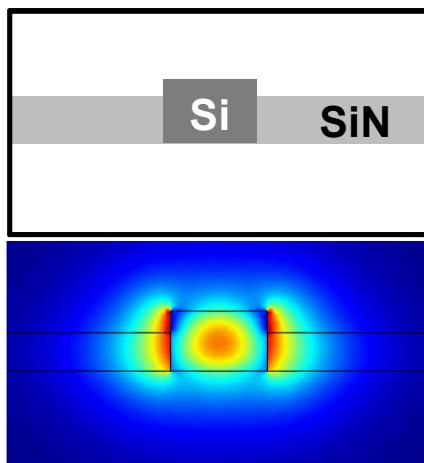
Combined Response







Si Waveguide



Kerr Nonlinearity

Intrinsic Si Nonlinearity:

$$n_2 = 4.5 \cdot 10^{-18} \text{ } m^2/W$$

NL Waveguide Coefficient:

$$\gamma_k = 188 \text{ } W^{-1}m^{-1}$$

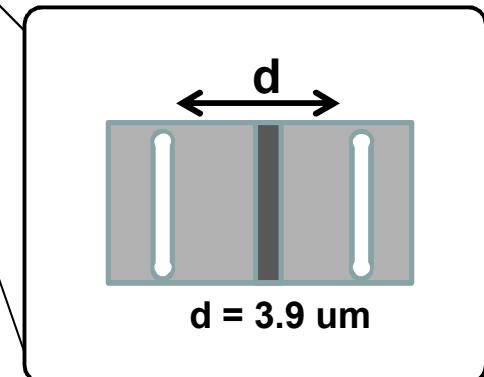
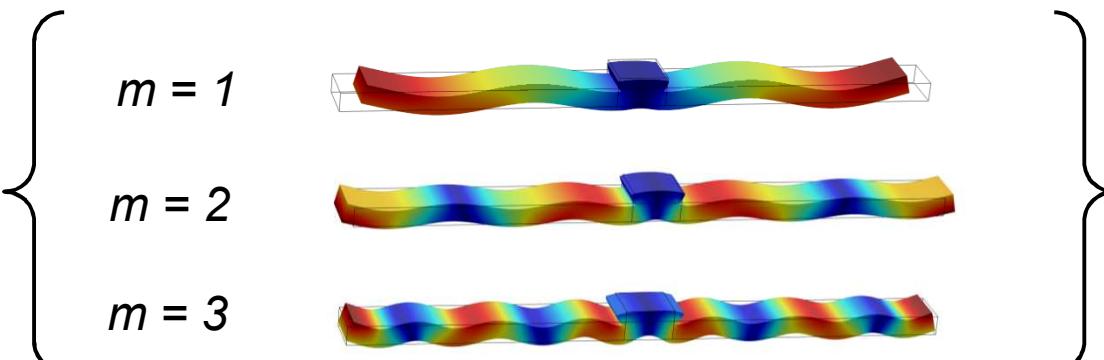
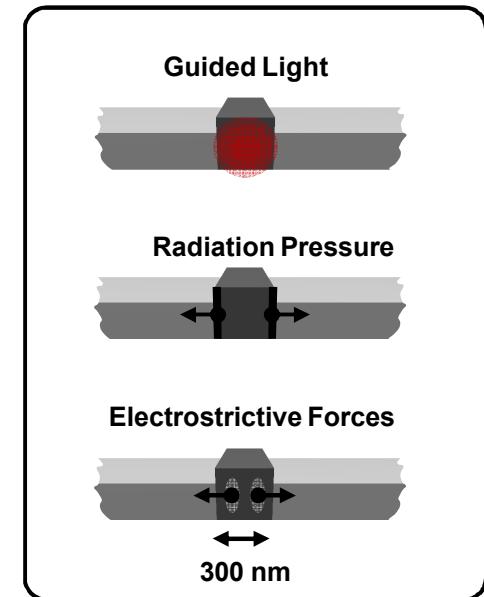
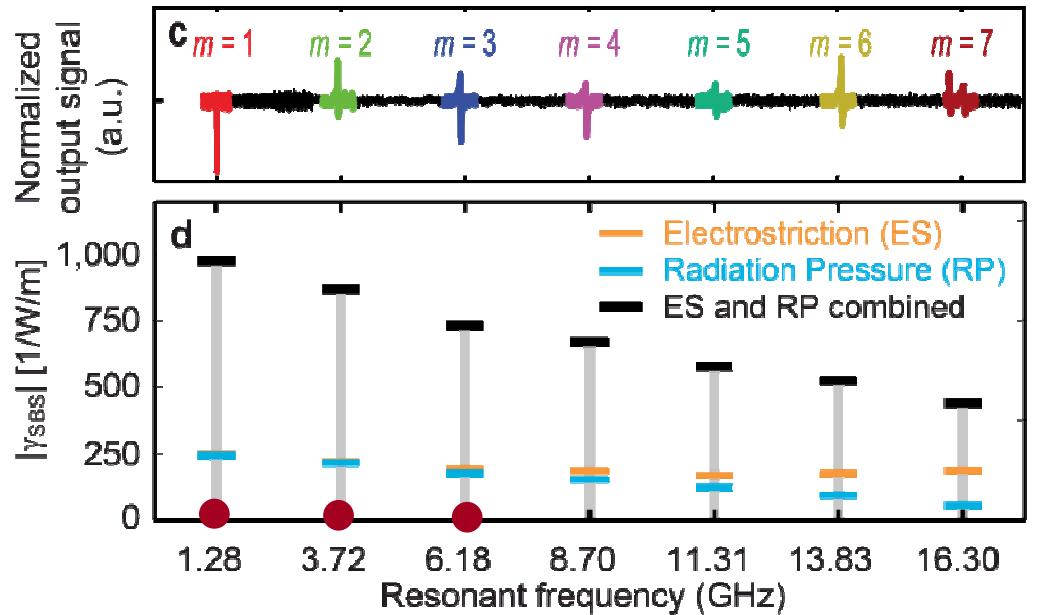
Brillouin nonlinearities are quantified relative Kerr nonlinearities.

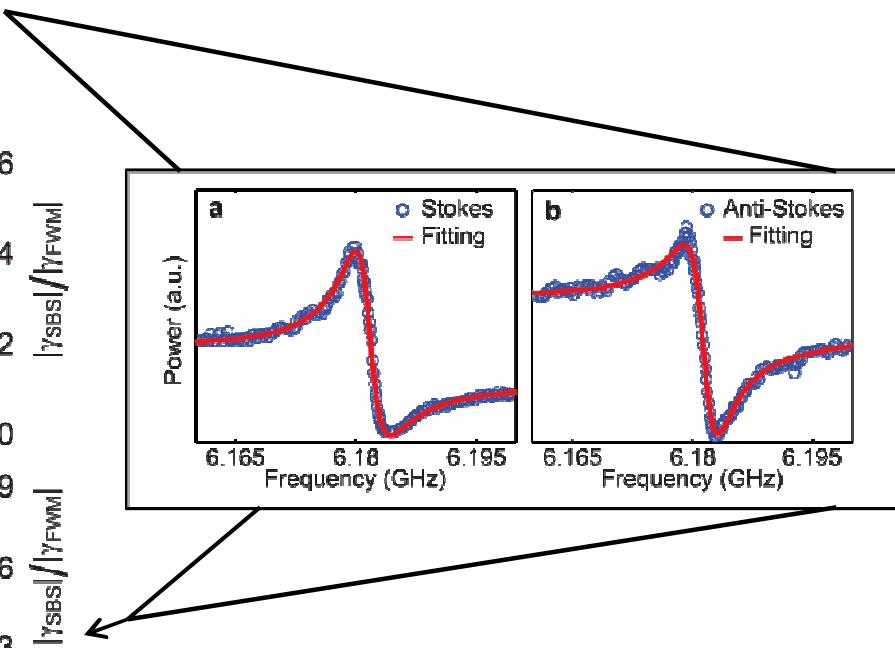
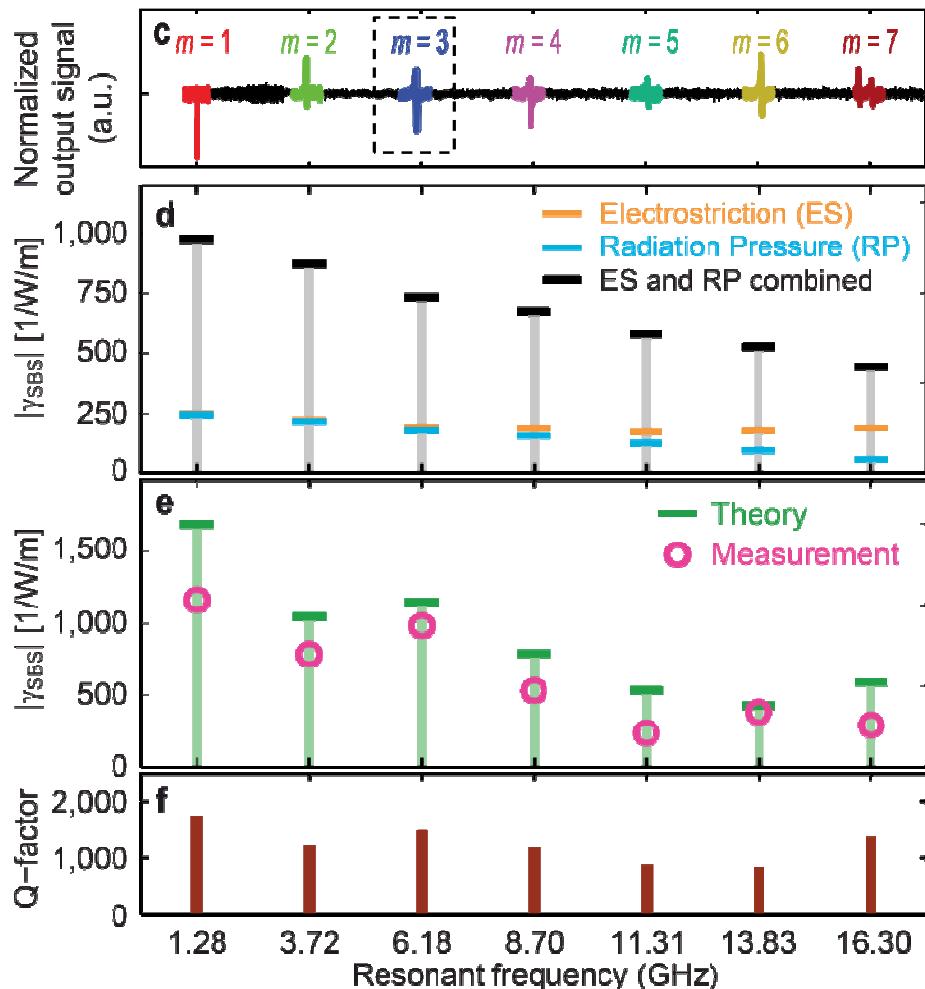
Line-Shape Analysis:

$$G_B/\gamma_k \approx 10.43$$

$$\begin{aligned} G_B &\approx 1960 \text{ } W^{-1}m^{-1} \\ Q &\approx 1561 \end{aligned}$$

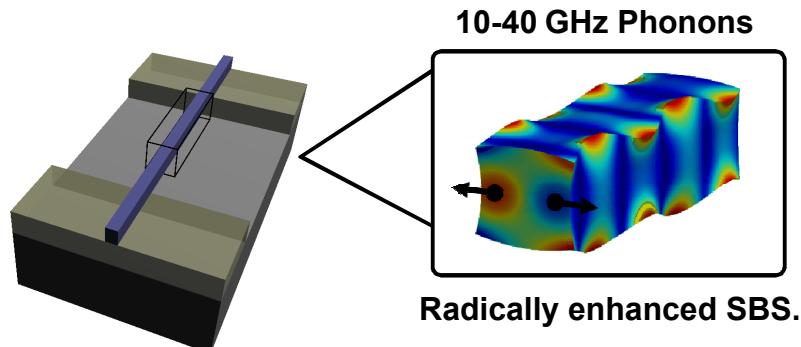
- 1000 X larger than any prior forward-SBS system.
- First SBS in Silicon!





- Efficient transduction 1-18 GHz frequencies.
- Brillouin resonances observed at higher (24 GHz) frequencies
- Electrostriction and Radiation pressure are clearly playing important role.

Theory: photon-phonon coupling.



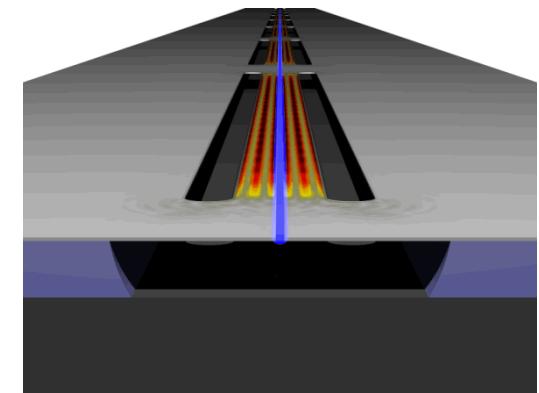
Part I Summary:

- Radically enhanced SBS processes found.
- SBS is no longer a bulk nonlinearity.
- Boundaries alone responsible for SBS.

Part II Summary:

- First demonstration of SBS in silicon.
- Exhibits 3000x stronger forward-SBS nonlinearity than any prior system.
- Many applications to come.

Experiment: Forward-SBS in Silicon



Our Team:

Yale

Peter Rakich
Heedeuk Shin



Sandia
National
Laboratories



Jonathan Cox
Ryan Camacho
Troy Olsson
Rob Jarecki

Zheng Wang
Wenjun Qiu (MIT)

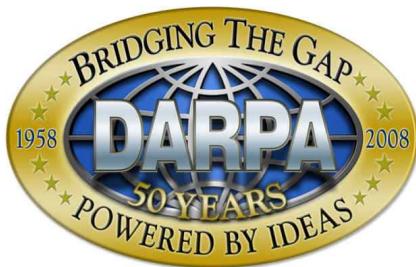


Acknowledgments:

Funding Agencies:

DARPA—MTO (PM: Jeff Rogers).

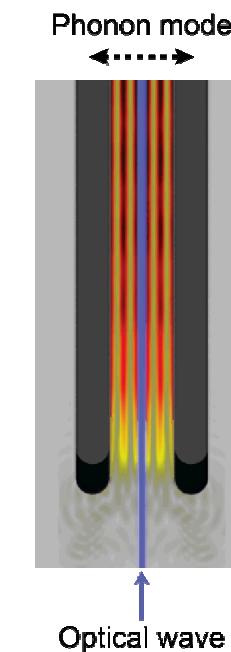
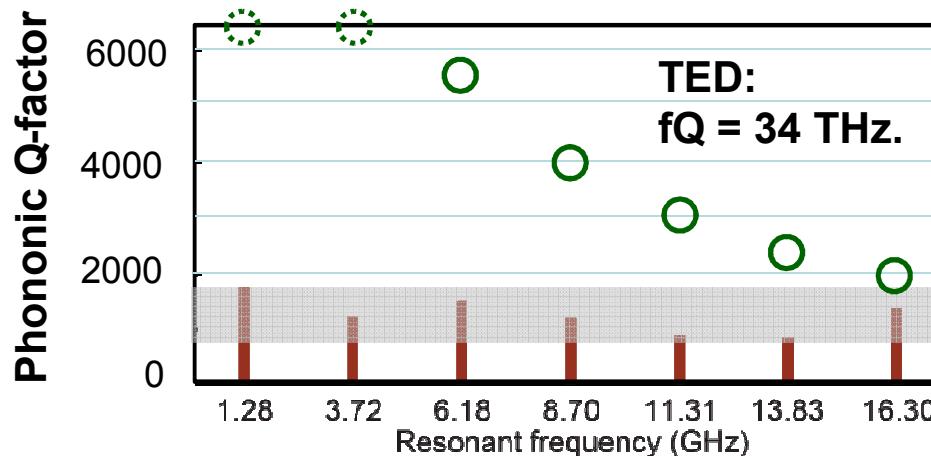
DOE—Laboratory Directed Research and Development funding.



Sandia is a multi-program laboratory operated by Sandia Corporation, a Lockheed Martin Company, for the United States Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000. This work was supported in part by the office of the Director of Defense Research and Engineering under Air Force contract FA8721-05-C-0002.

Q: How to obtain ultimate efficiency and noise performance?

A: Enhance parametric gain further.



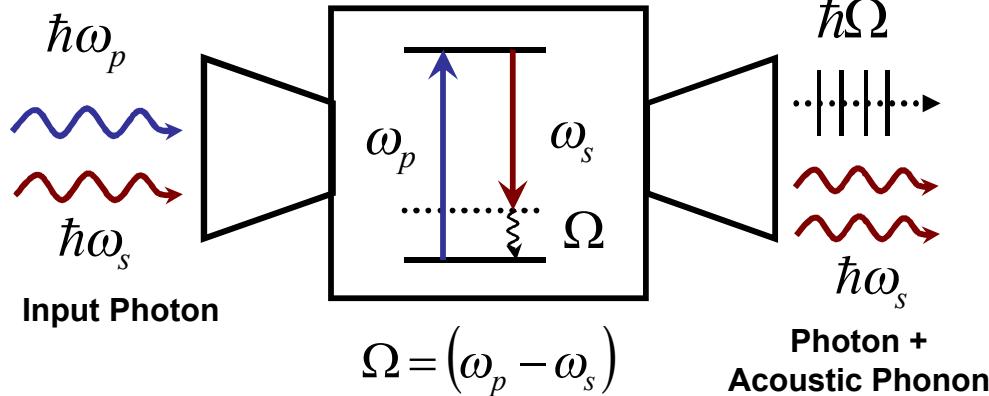
- Mean Q-factor ~ 1500
- Some improvement before thermoelastic damping limit
- Possible culprits: Anchoring losses, roughness.

D. Wilson, C. Regal, S. Papp, and H. Kimble, "Cavity optomechanics with stoichiometric SiN films," *Physical review letters*, vol. 103, no. 20, p. 207204, 2009.

1. **First demonstration of wideband (1-18 GHz) stimulated Mach-wave phonon emission.**
2. **First-ever demonstration *chip-scale* Forward Stimulated Brillouin Scattering (SBS).**
 - > First demonstration of SBS in silicon.
 - > 3,000 x stronger SBS than any known system.
 - > Demonstrated tailorable phonon emission from 1GHz-18GHz.
 - > Demonstrated tailorable nonlinear susceptibility from the coherent interference of Kerr and Brillouin nonlinearities.
3. **Demonstrated device physics for tailorable bandwidth phononic crystal Mach-wave emitter.**

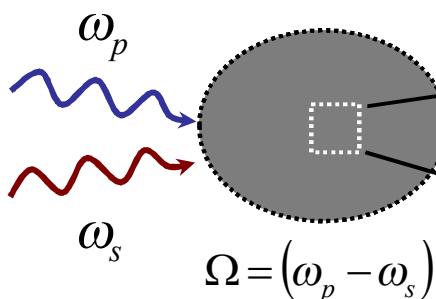
Physics of Brillouin Scattering:

Nonlinear System

Brillouin Process:

- Coupling between *acoustic phonons* and photons.
- At the heart of all optomechanical interactions.

Micro-scale origins of parametric process:



Interference yields intensity “Beat Note”.

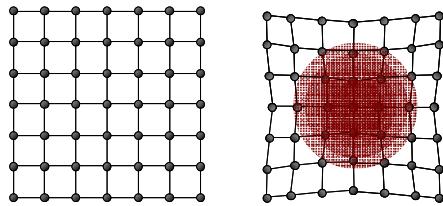
$$|E(t)|^2 = 2E_p E_s \cdot \cos(\Omega \cdot t) + C$$

Optical Force: Proportional to Intensity.

$$F(t) \approx \alpha \cdot \sqrt{P_p \cdot P_s} \cdot \cos(\Omega \cdot t)$$

Time Varying Forces Excite Phonons.

Origin of electrostrictive forces: dynamic material response.



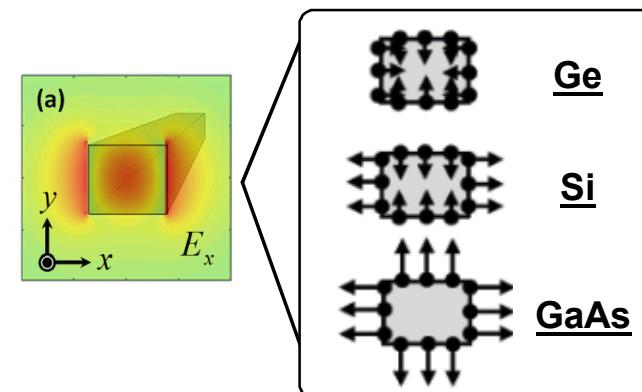
$$S = \alpha \cdot E_i + \beta \cdot E_i E_j$$

Strain
Piezo Coeff.
Electrostrictive Coeff.

Electrostriction = Material induced optical forces.

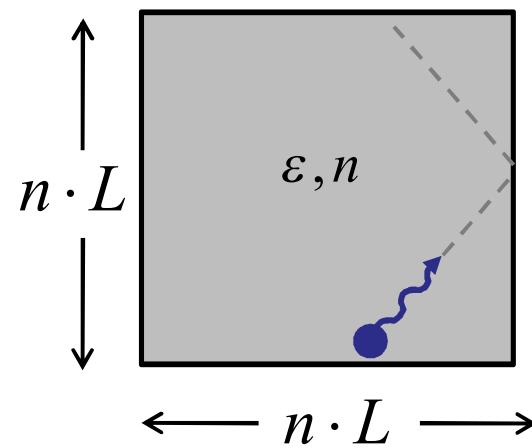
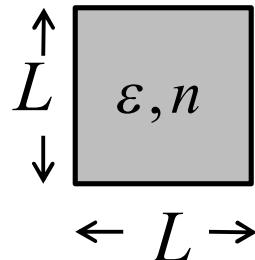
- All dielectrics exhibit electrostriction (not piezo electricity).
- Sign and magnitude are tailorable by choice of material

Medium properties dictate force dist.



Box: Photon's Perspective

Box: Real Space



From Photon's Perspective:

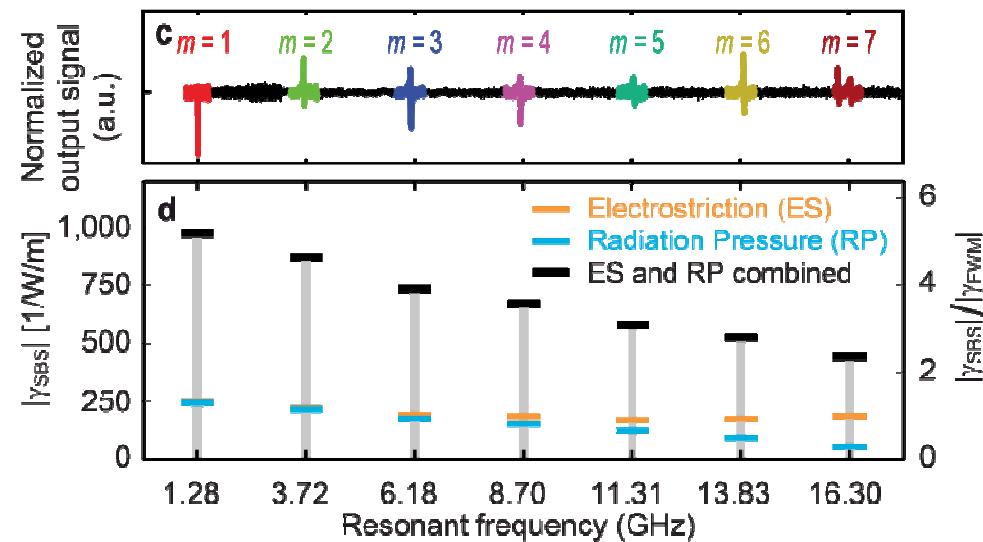
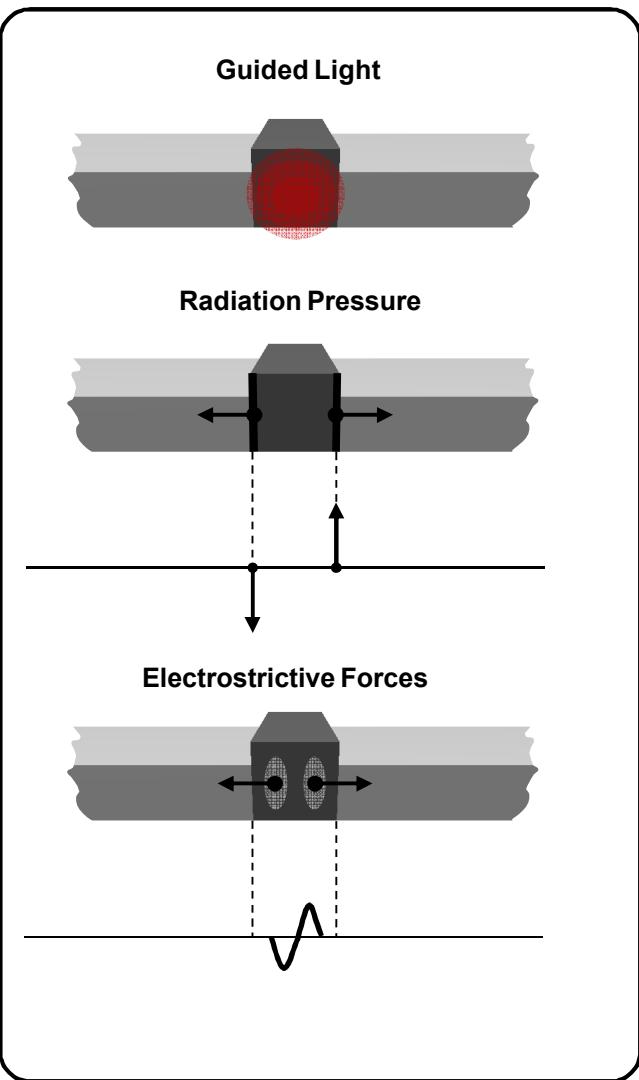
- Space is now quite different.

$$V \Rightarrow V' = V \cdot n^3$$

Box seems much bigger.

Oddities don't end here...

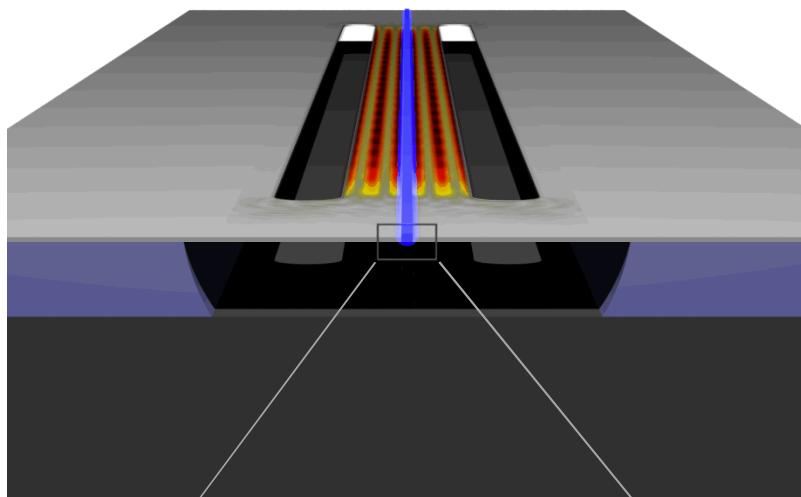
Origins of Wide-Bandwidth Coupling:



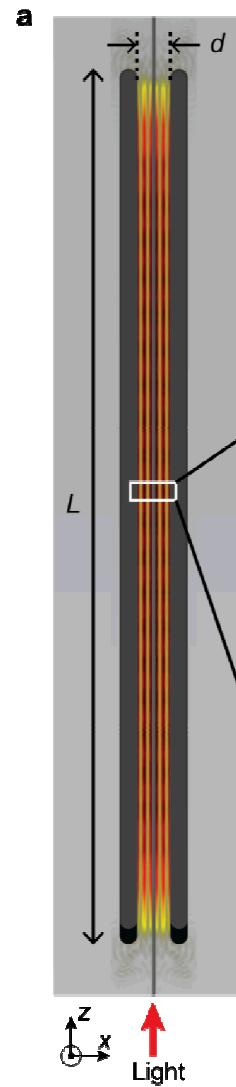
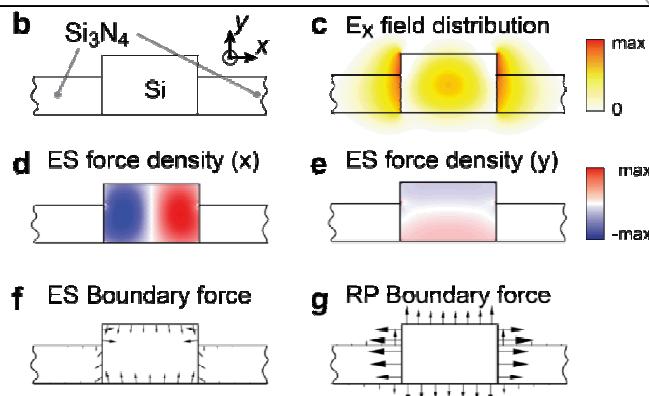
- <3dB variation in total gain from 1-18 GHz.
- Transductive bandwidth of RP is ~10 GHz.
- Transductive bandwidth of ES > 18 GHz.

Q: What is responsible for large bandwidth?

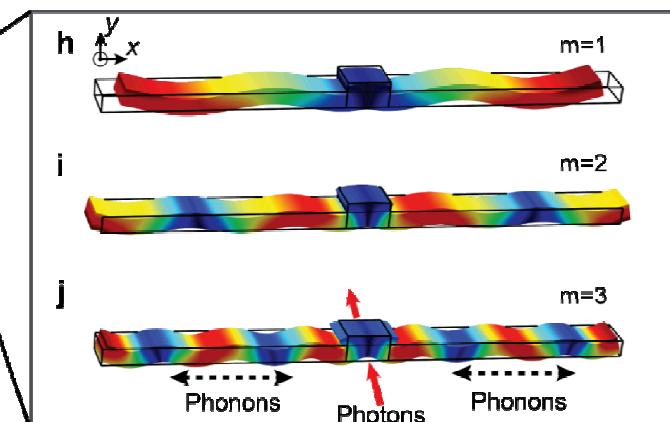
A: The higher spatial fundamental frequency associated with electrostriction more efficiently excite waves to higher bandwidths.



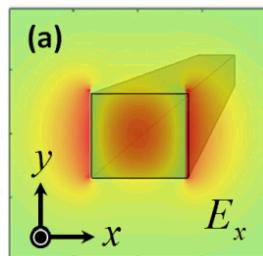
Mode and Optical Forces



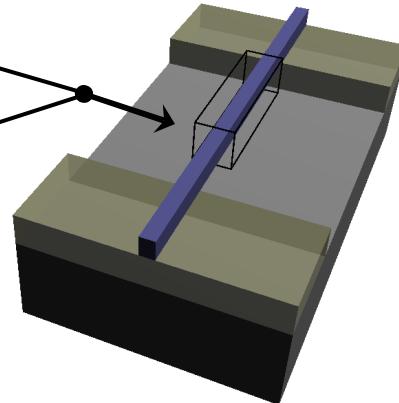
Brillouin Active Phonon Modes



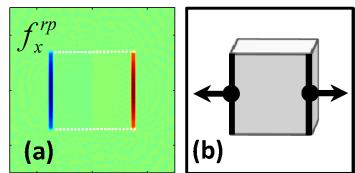
- Tailorable Brillouin mode spectrum
- Supports 1-100 GHz phonon modes.
- Phononic modes and optical forces independently controlled.



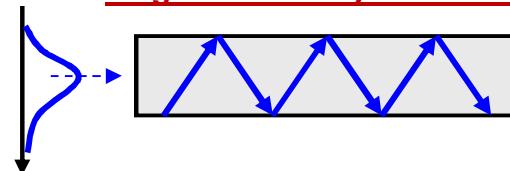
Guided mode within suspended dielectric waveguide. (300x300nm)



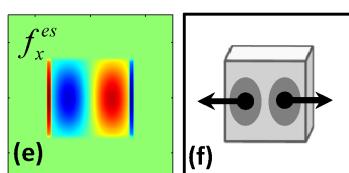
Radiation Pressure



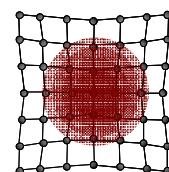
Origin: Boundary Scattering



Electrostrictive forces



Origin: Dynamic Material Response

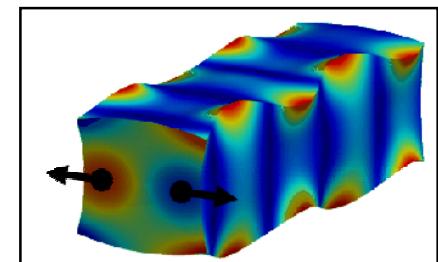


Akin to piezoelectricity

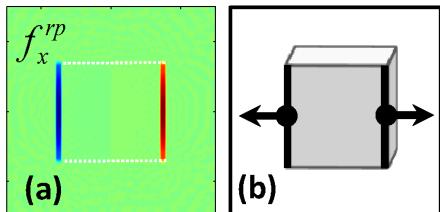
$$S = \alpha \cdot E_i + \beta \cdot E_i E_j$$

Strain Piezo Coeff. Electrostrictive Coeff.

photon-phonon coupling:



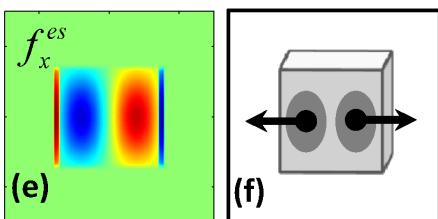
Efficient ultra-high frequency phonon emission.

Radiation Pressure

Rad. Pressure: $\bar{p}^{rp} = \frac{P_{opt}}{c \cdot A} \cdot (n_g - n_p) = \frac{P_{opt}}{c \cdot A} \cdot \alpha^{rp}$

Simp

Electrostriction: $\bar{p}^{es} \cong \frac{P_{opt}}{c \cdot A} \cdot n_g n^2 (p_{11} + 2p_{12}) / 2 = \frac{P_{opt}}{c \cdot A} \cdot \alpha^{es}$

Electrostrictive forces

Material	Symmetry	p_{11}	p_{12}	$p_{11} + 2p_{12}$	n	α^{rp}	α^{es}
Si	cubic	-0.09	+0.017	-0.056	3.5	-5	-1.7
Ge	cubic	0.27	0.235	0.74	4.2	-6.4	+40
GaAs	cubic	-0.165	-0.14	-0.445	3.4	-4.8	-12
Silica	amorphous	0.121	0.27	0.661	1.45	-0.89	+1.0
As ₂ S ₃	amorphous	0.25	0.24	0.73	2.4	-2.8	+6.5
As ₂ Se ₃	amorphous	-	-	-	2.8	-3.6	-

How Large are Forces?

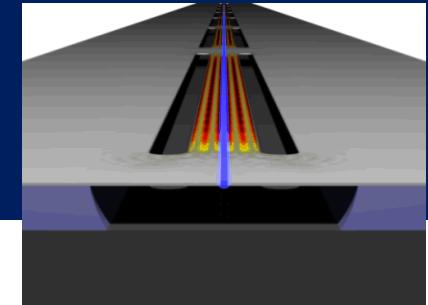
Material	Pressure (Pwr = 100mW)
Si	$\sim 5 \times 10^4 \text{ N/m}^2$
Ge	$\sim 10^6 \text{ N/m}^2$

5-50 People
standing on
manhole
cover

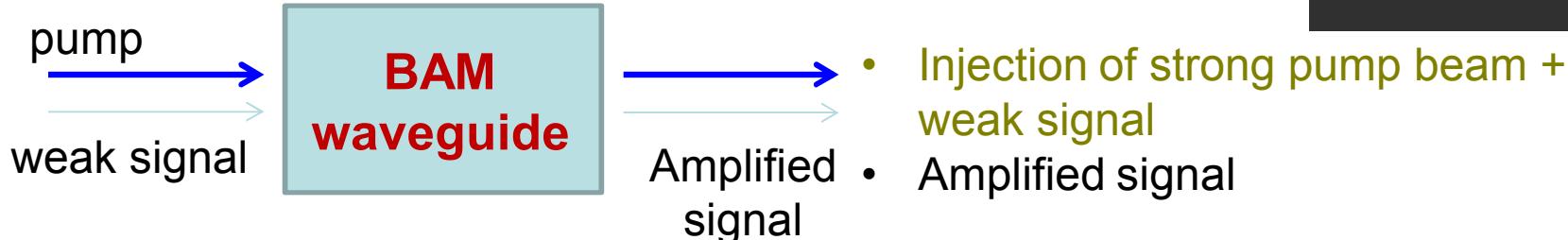
Material	Pressure (Pwr = 1kW)
Si	$\sim 5 \times 10^8 \text{ N/m}^2$
Ge	$\sim 10^{10} \text{ N/m}^2$

Stresses
Approach
Material Yield
Strength

How large is the SBS gain?



How do we quantify the SBS gain?

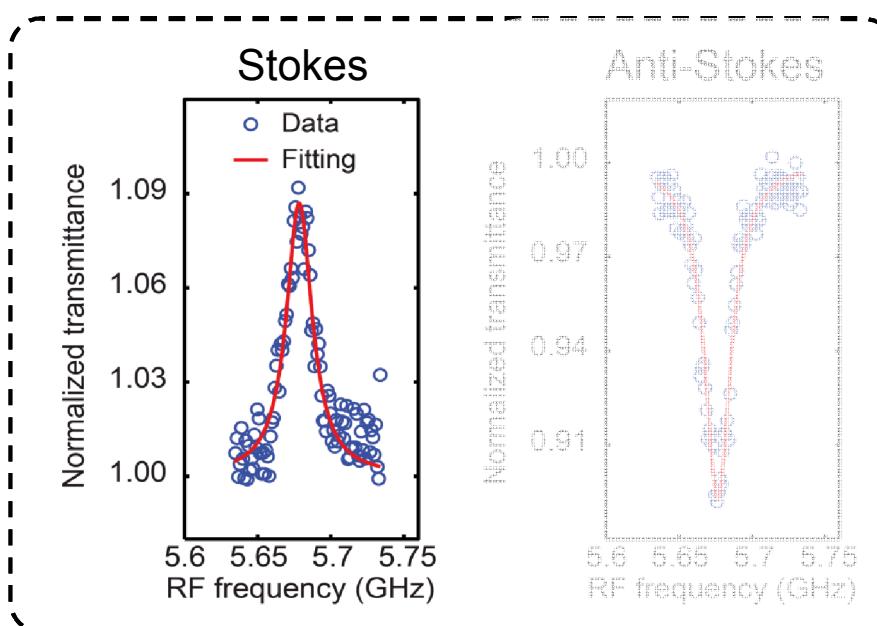


What do these gain measurements tell us?

- 10 % amplification/depletion for Stokes/anti-Stokes fields
- Consistent with our previous heterodyne measurements
- Large gain coefficient, but overall gain is modest.

What limits gain?

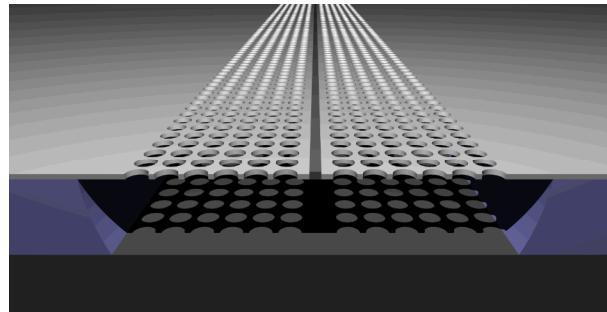
- Gain limited by power handling & linear loss in waveguide



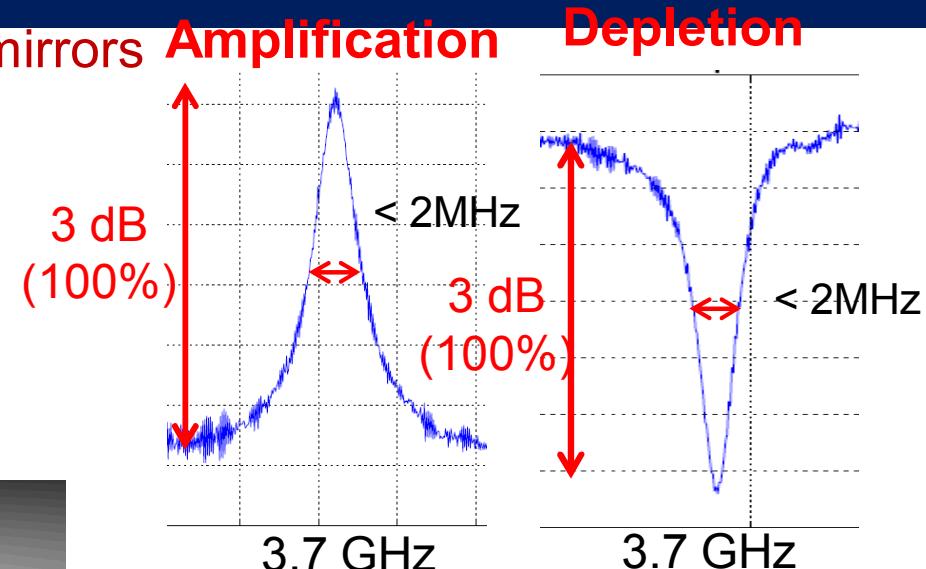
Gain coefficient: $\sim 2800 \text{ W}^{-1}\text{m}^{-1}$

BAM waveguide with phononic mirrors

- Same concept (hybrid photon-phonon waveguide)
- Phonon guided through Bragg reflection
- Wider silicon waveguide (1 um)

**Results**

- Higher power handling
- Lower propagation loss
- Longer effective interaction length
- Better photon-phonon coupling at low frequency



- 10 times enhancement than previous result
- We believe there is opportunity to enhance performance further.
- Expecting many applications