

# Coarse Graining in Peridynamics

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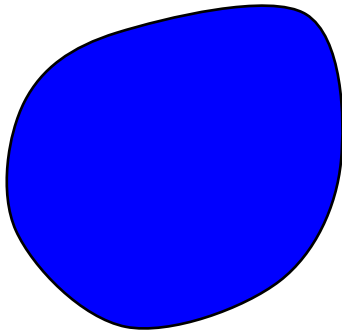
SES2011  
Evanston, IL

October 12, 2011

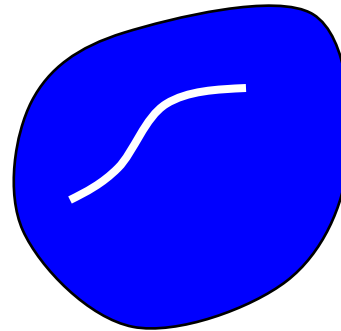


# Purpose of peridynamics

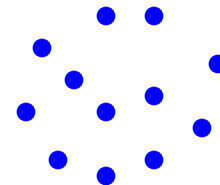
- To unify the mechanics of continuous and discontinuous media within a single, consistent set of equations.



Continuous body

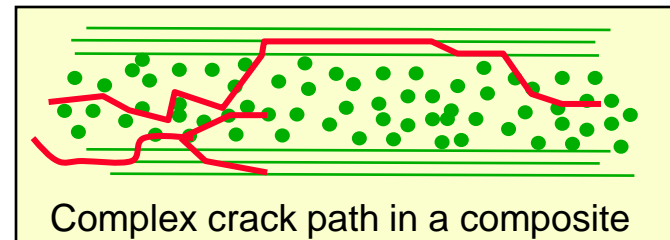


Continuous body  
with a defect



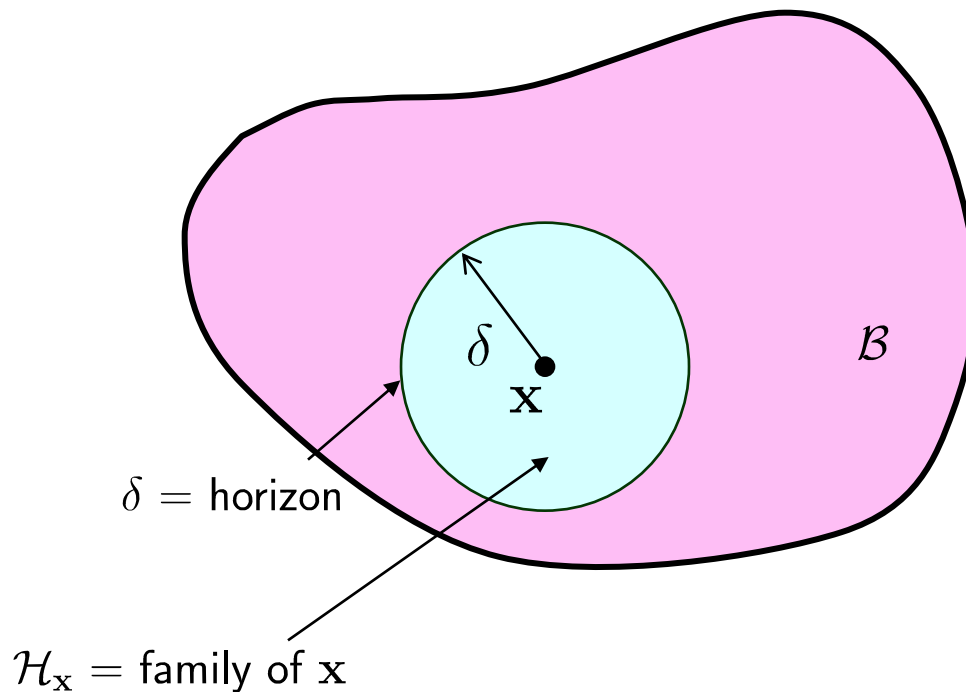
Discrete particles

- Why do this?
  - Avoid coupling dissimilar mathematical systems (A to C).
  - Model complex fracture patterns.
  - Communicate across length scales.



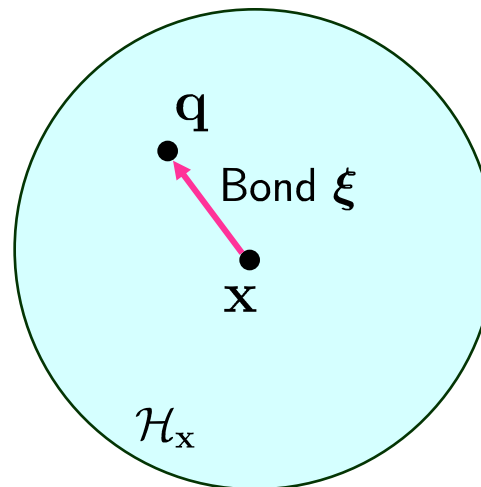
# Peridynamics basics: Horizon and family

- Any point  $\mathbf{x}$  interacts directly with other points within a distance  $\delta$  called the “horizon.”
- The material within a distance  $\delta$  of  $\mathbf{x}$  is called the “family” of  $\mathbf{x}$ ,  $\mathcal{H}_{\mathbf{x}}$ .



# States

- A *state* is a function defined on all the bonds in  $\mathcal{H}_x$ .
- Notation for states:  $\underline{\mathbf{A}}\langle \xi \rangle$ .
- *Vector state*:  $\underline{\mathbf{A}}\langle \cdot \rangle : \mathcal{H} \rightarrow \mathbb{R}^3$  where  $\mathcal{H}$  is the set of all the bonds in  $\mathcal{H}_x$ .



SS, Askari, Epton, Wecker, & Xu, Peridynamic States and Constitutive Modeling, J. Elast (2007)

# Vector states take the place of tensors

- The *deformation state* maps each bond to its deformed image:

$$\underline{\mathbf{Y}}\langle \mathbf{q} - \mathbf{x} \rangle = \mathbf{y}(\mathbf{q}, t) - \mathbf{y}(\mathbf{x}, t).$$

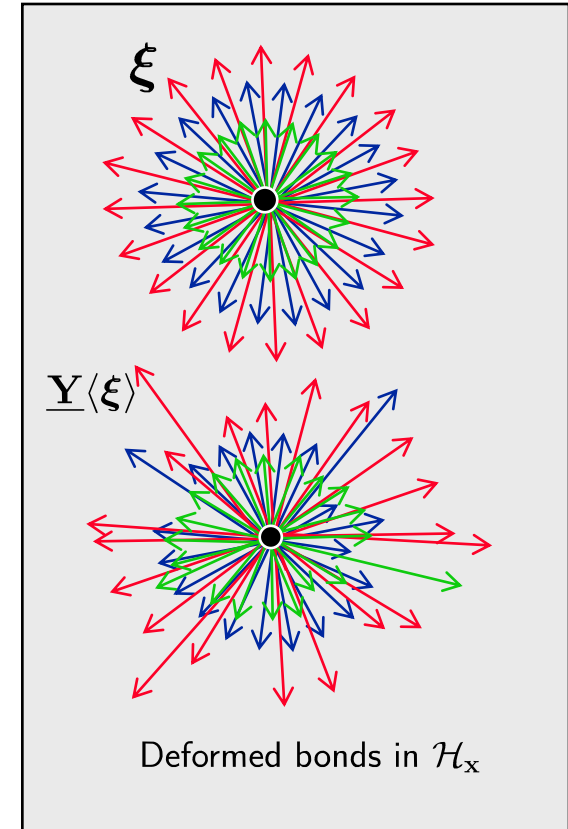
- The *force state* maps a bond to its *bond force*:

$$\underline{\mathbf{T}}\langle \mathbf{q} - \mathbf{x} \rangle = \mathbf{t}(\mathbf{q}, \mathbf{x}, t).$$

- Constitutive model is a state-valued function of a state:

$$\underline{\mathbf{T}} = \hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}})$$

- Compare standard theory:  $\boldsymbol{\sigma} = \hat{\boldsymbol{\sigma}}(\mathbf{F})$ .





## Summary of nonlinear, state based PD

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- Equation of motion:

$$\rho(\mathbf{x})\ddot{\mathbf{y}}(\mathbf{x}, t) = \mathbf{L}(\mathbf{x}, t) + \mathbf{b}(\mathbf{x}, t)$$

where the internal force density at  $\mathbf{x}$  is given by

$$\mathbf{L}(\mathbf{x}, t) = \int_{\mathcal{H}_{\mathbf{x}}} \mathbf{f}(\mathbf{q}, \mathbf{x}, t) dV_{\mathbf{q}}$$

and  $\mathbf{b}$  is the body force density.

- Pairwise force density is given by

$$\mathbf{f}(\mathbf{q}, \mathbf{x}, t) = \mathbf{t}(\mathbf{q}, \mathbf{x}, t) - \mathbf{t}(\mathbf{x}, \mathbf{q}, t)$$

where  $\mathbf{t}$  is determined by the constitutive model:

$$\mathbf{t}(\mathbf{q}, \mathbf{x}, t) = \underline{\mathbf{T}}[\mathbf{x}, t] \langle \mathbf{q} - \mathbf{x} \rangle, \quad \underline{\mathbf{T}} = \hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}}).$$

Force state at  $\mathbf{x}, t$

Constitutive model



# Elastic materials

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- Strain energy density:

$$W(\mathbf{x}, t) = \hat{W}(\underline{\mathbf{Y}}[\mathbf{x}, t])$$

- Compare standard theory:

$$W(\mathbf{x}, t) = \hat{W}(\mathbf{F}(\mathbf{x}, t)), \quad \mathbf{F} := \frac{\partial \mathbf{y}}{\partial \mathbf{x}}.$$

- Bond force:

$$\mathbf{t}(\mathbf{q}, \mathbf{x}, t) = \underline{\mathbf{T}} \langle \mathbf{q} - \mathbf{x} \rangle, \quad \underline{\mathbf{T}} = \hat{W}_{\underline{\mathbf{Y}}}(\underline{\mathbf{Y}})$$

where  $\hat{W}_{\underline{\mathbf{Y}}}$  is the Fréchet derivative of  $\hat{W}$  with respect to  $\underline{\mathbf{Y}}$ .



# Linearized peridynamics

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- Suppose  $|\mathbf{u}| \ll \delta$ . Equilibrium equation becomes

$$\mathbf{L}(\mathbf{x}) + \mathbf{b}(\mathbf{x}) = \mathbf{0}$$

$$\mathbf{L}(\mathbf{x}) = \int_{\mathcal{N}} \mathbf{C}(\mathbf{x}, \mathbf{q}) (\mathbf{u}(\mathbf{q}) - \mathbf{u}(\mathbf{x})) dV_{\mathbf{q}}$$

where  $\mathbf{C}$  is the tensor-valued *micromodulus* function and  $\mathcal{N}$  has twice the radius of  $\mathcal{H}$ .

- $\mathbf{C}$  involves the second Fréchet derivatives of strain energy density  $\hat{W}_{\underline{\mathbf{Y}\mathbf{Y}}}$ :

$$\mathbf{C}(\mathbf{x}, \mathbf{q}) = \int_{\mathcal{B}} \left( \hat{W}_{\underline{\mathbf{Y}\mathbf{Y}}}[\mathbf{x}] \langle \mathbf{p} - \mathbf{x}, \mathbf{q} - \mathbf{x} \rangle - \hat{W}_{\underline{\mathbf{Y}\mathbf{Y}}}[\mathbf{p}] \langle \mathbf{x} - \mathbf{p}, \mathbf{q} - \mathbf{p} \rangle + \hat{W}_{\underline{\mathbf{Y}\mathbf{Y}}}[\mathbf{q}] \langle \mathbf{x} - \mathbf{q}, \mathbf{p} - \mathbf{q} \rangle \right) dV_{\mathbf{p}}$$

- $\mathbf{C}$  has the following symmetry:

$$\mathbf{C}^T(\mathbf{x}, \mathbf{q}) = \mathbf{C}(\mathbf{q}, \mathbf{x}) \quad \forall \mathbf{x}, \mathbf{q}.$$

- Linearized equilibrium equation is same as in Kunin's nonlocal theory (1983).



# EMU (and LAMMPS) numerical method

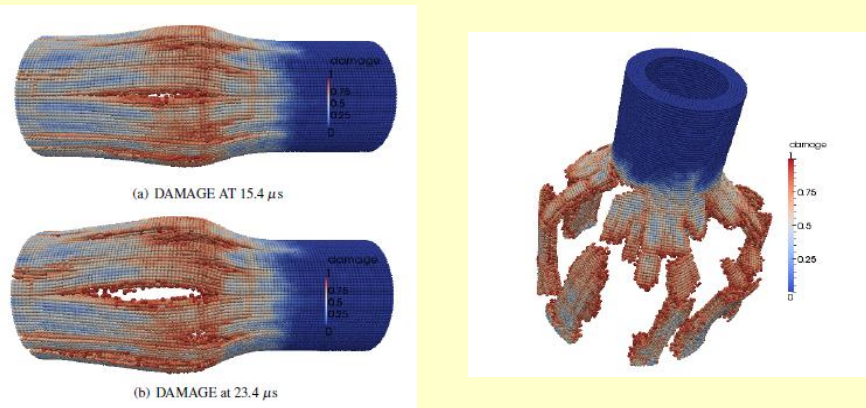
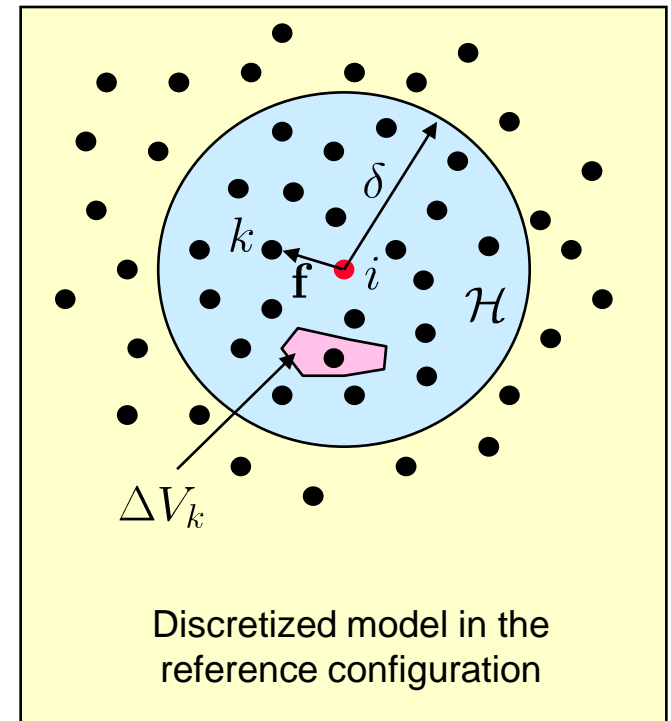
- Integral is replaced by a finite sum: resulting method is [meshless](#) and [Lagrangian](#).

$$\rho \ddot{\mathbf{y}}(\mathbf{x}, t) = \int_{\mathcal{H}} \mathbf{f}(\mathbf{q}, \mathbf{x}, t) dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x}, t)$$

↓

$$\rho \ddot{\mathbf{y}}_i^n = \sum_{k \in \mathcal{H}} \mathbf{f}(\mathbf{x}_k, \mathbf{x}_i, t) \Delta V_k + \mathbf{b}_i^n$$

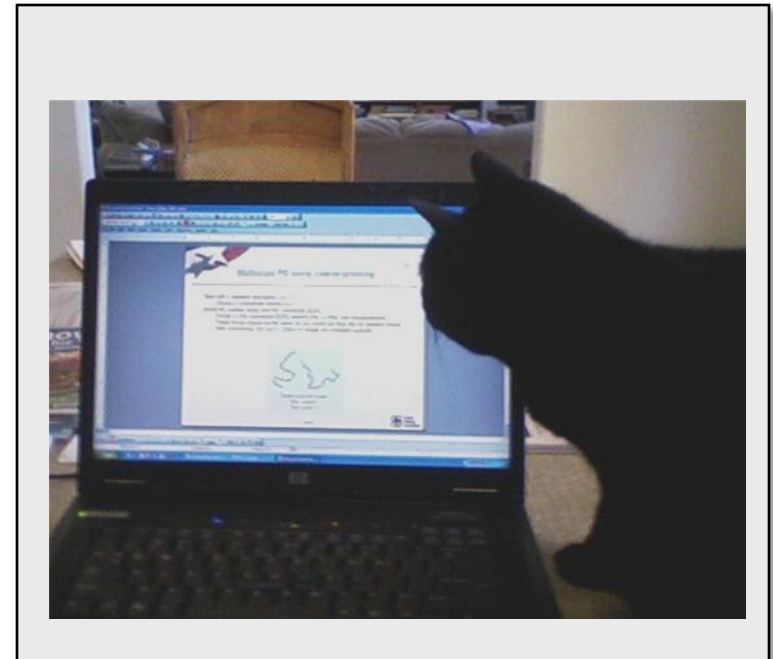
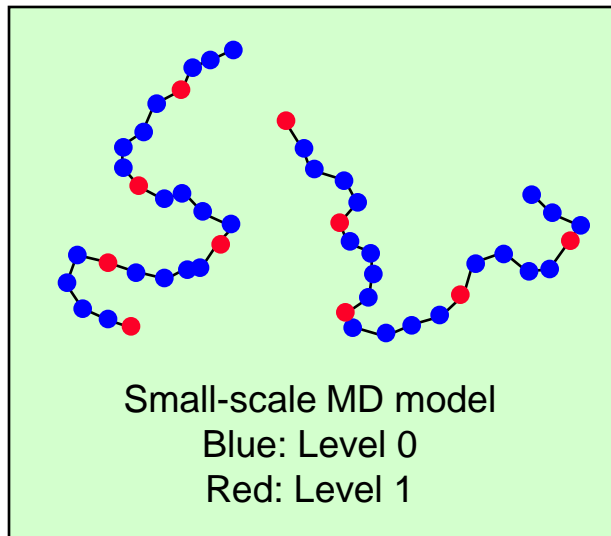
- Looks a lot like MD!
- LAMMPS implementation by M. Parks & . Seleson



Method is also available in Sierra (D. Littlewood)

# Multiscale PD using coarse-graining

- Start with a detailed description (**level 0**).
  - Choose a *coarsened* subset (**level 1**).
- Model the system using only the coarsened DOFs...
  - Forces on **level 1** DOFs depend only on their own displacements.
  - These forces should be the same as you would get from the full detailed model.
  - The **level 0** DOFs no longer appear explicitly.

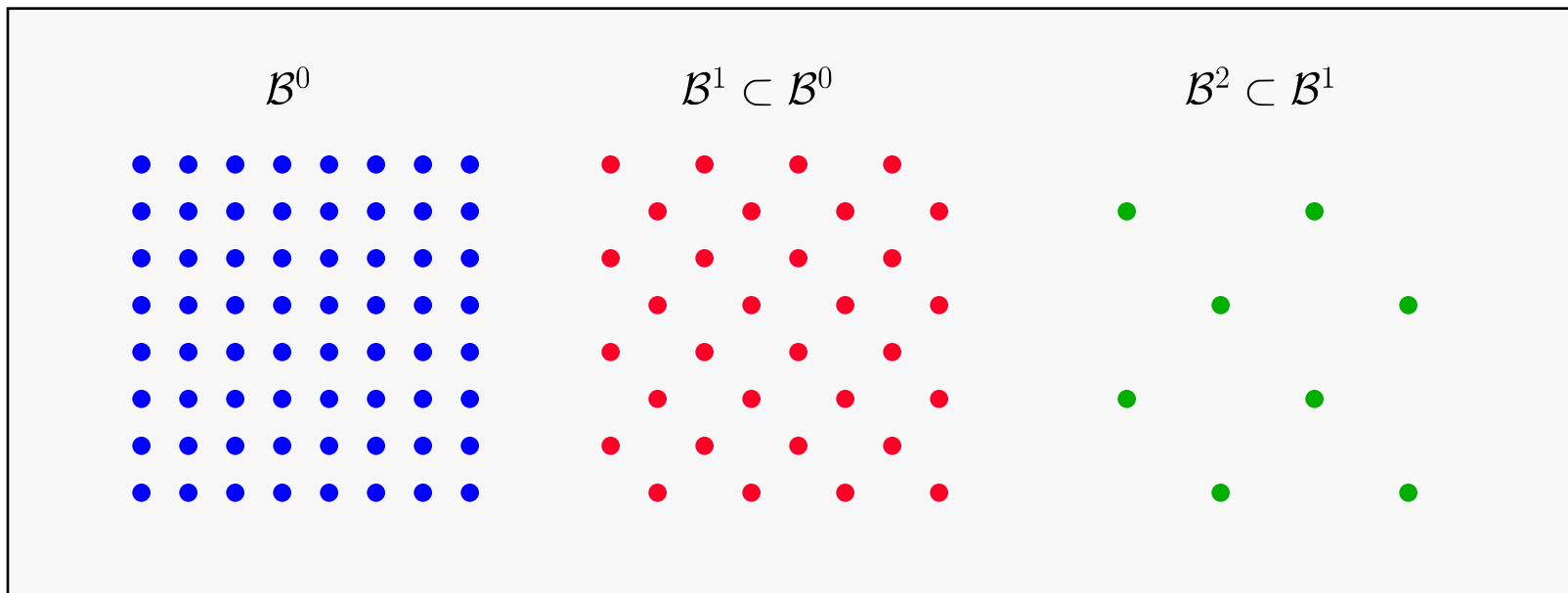




# Coarse-graining: Reduce the number of degrees of freedom

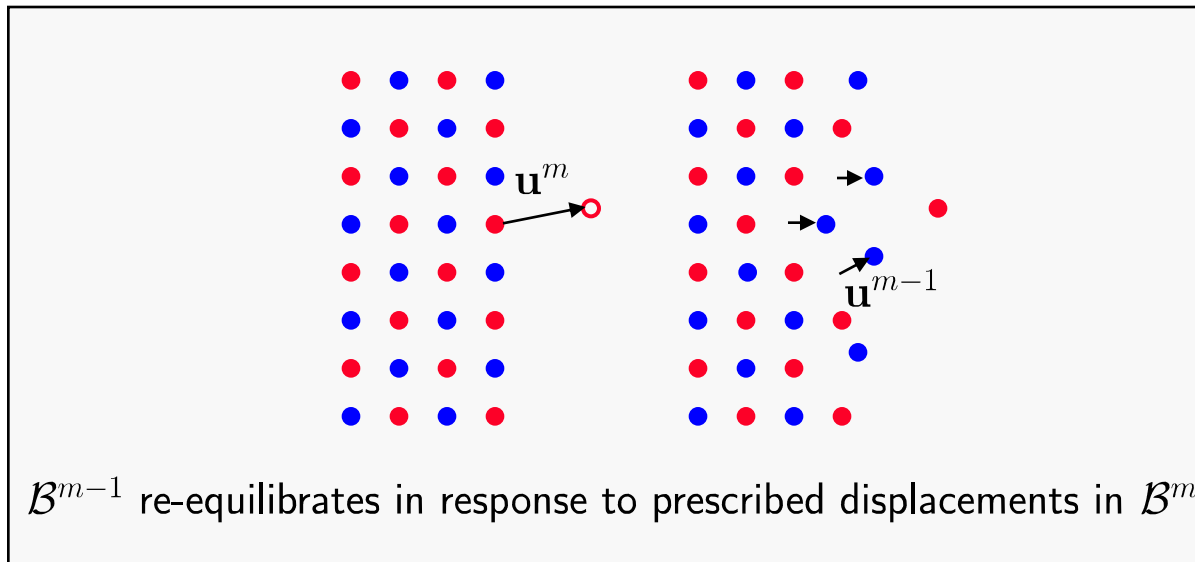
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- Start with a detailed “level 0” body  $\mathcal{B}^0$ .
- Level 0 can be either continuous or discrete.
- Choose a sequence of  $M$  coarsened levels:  $\mathcal{B}^M \subset \mathcal{B}^{M-1} \subset \dots \subset \mathcal{B}^1 \subset \mathcal{B}^0$ .



# Each level's displacements are determined by the next higher level

- Assumption: If  $\mathbf{u}^m$  is prescribed, the excluded DOFs in  $\mathbf{u}^{m-1}$  “float” (respond by changing their equilibrium displacements).



- Let  $\mathbf{S}^m$  be the solution operator that gives  $\mathbf{u}^{m-1}$  in terms of  $\mathbf{u}^m$ :

$$\mathbf{u}^{m-1}(\mathbf{x}) = \int_{\mathcal{B}^m} \mathbf{S}^m(\mathbf{x}, \mathbf{q}) \mathbf{u}^m(\mathbf{q}) dV_{\mathbf{q}}.$$

# Each level has the same mathematical structure

- Equilibrium equation for level  $m$  involves only the level  $m$  displacements:

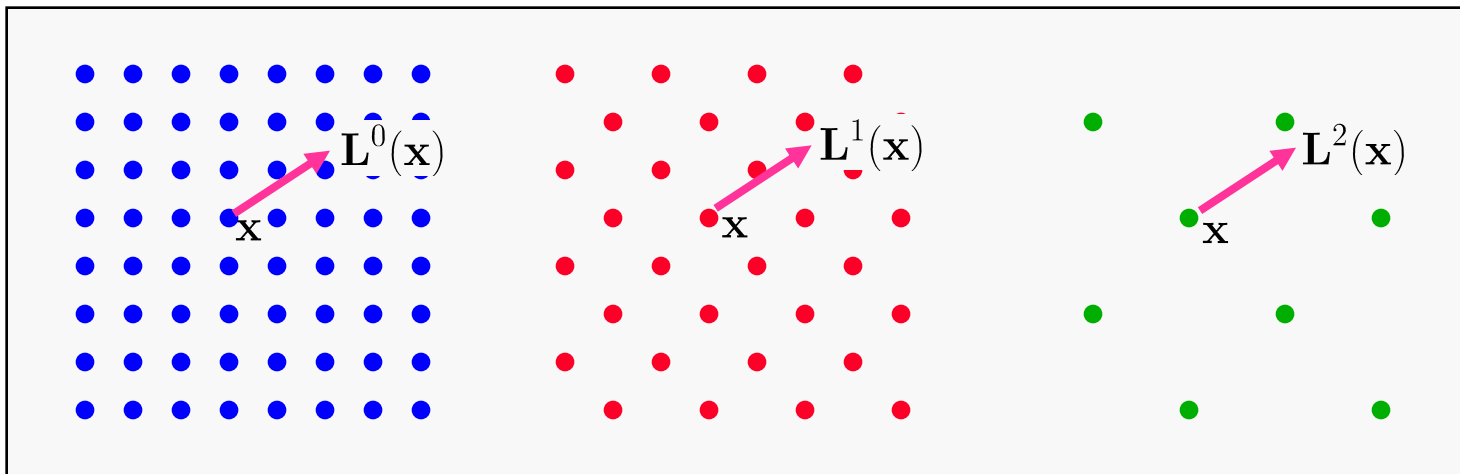
$$\mathbf{L}^m(\mathbf{x}) = \int_{\mathcal{B}^m} \mathbf{C}^m(\mathbf{x}, \mathbf{q})(\mathbf{u}^m(\mathbf{q}) - \mathbf{u}^m(\mathbf{x})) dV_{\mathbf{q}} = 0$$

where the level  $m$  micromodulus is found from

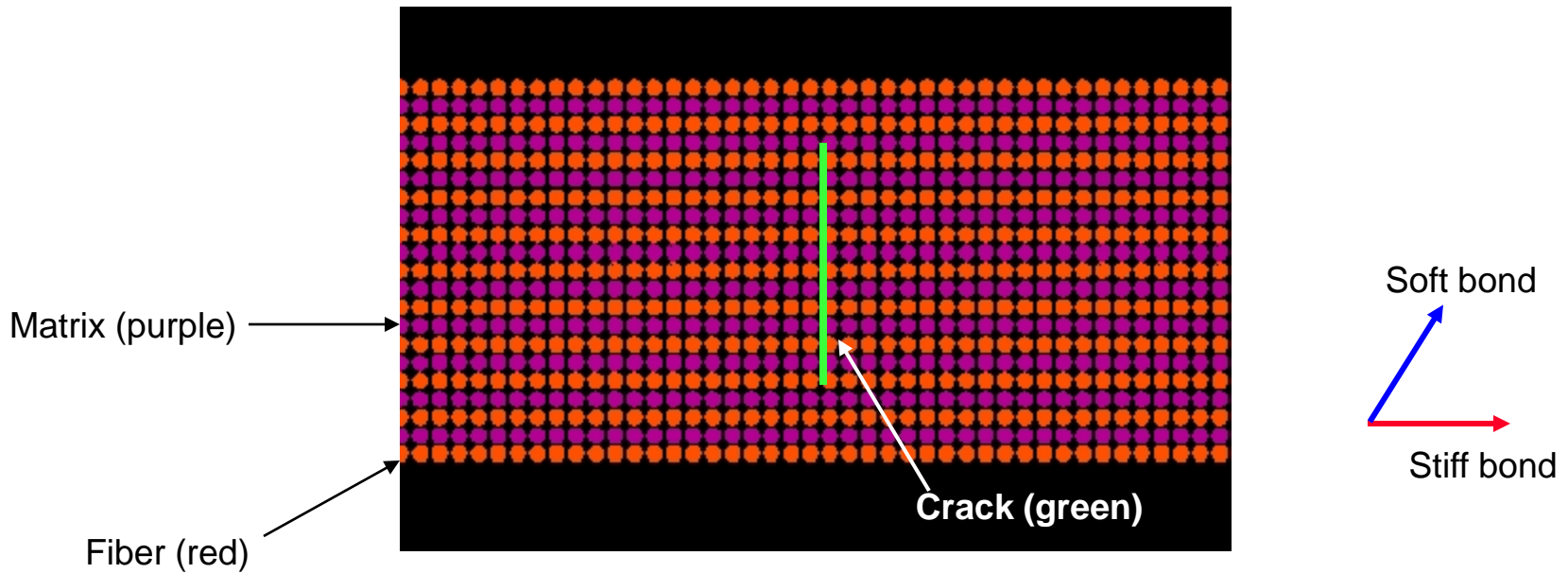
$$\mathbf{C}^m(\mathbf{x}, \mathbf{q}) = \int_{\mathcal{B}^{m-1}} \mathbf{C}^{m-1}(\mathbf{x}, \mathbf{p}) \mathbf{S}^m(\mathbf{p}, \mathbf{q}) dV_{\mathbf{p}}.$$

- If loads or displacements in  $\mathcal{B}^M$  are prescribed, the forces are invariant through the levels:

$$\mathbf{L}^0 = \mathbf{L}^1 = \dots = \mathbf{L}^{M-1} = \mathbf{L}^M$$



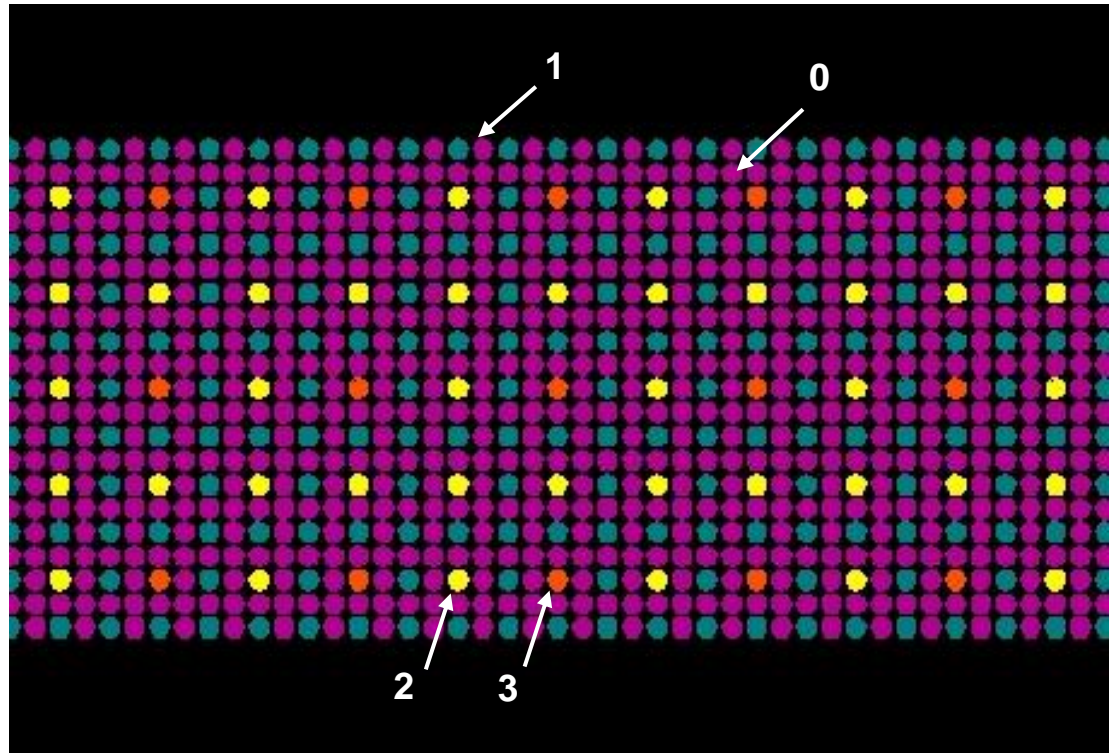
## Example: Composite with a crack



- Crack is inserted into the level 0 model: bonds crossing the crack are ignored.
- $C(\mathbf{x}, \mathbf{q}) = 0$  whenever  $\mathbf{x}$  and  $\mathbf{q}$  are on opposite sides of the crack.

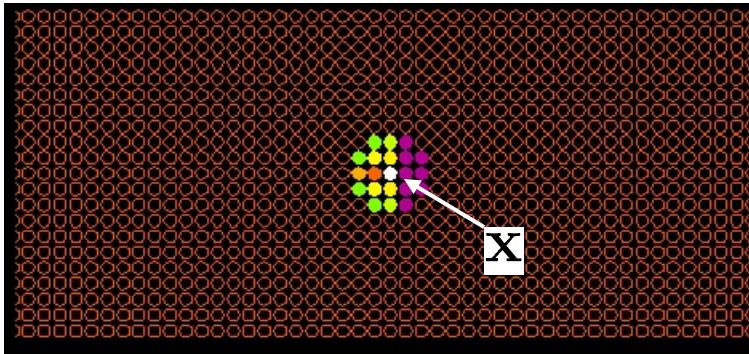
## 3 levels of coarsening

- Purple: Level 0
- Green: Level 1
- Yellow: Level 2
- Red: Level 3
- Each level  $m + 1$  has  $1/4$  as many nodes as level  $m$ .

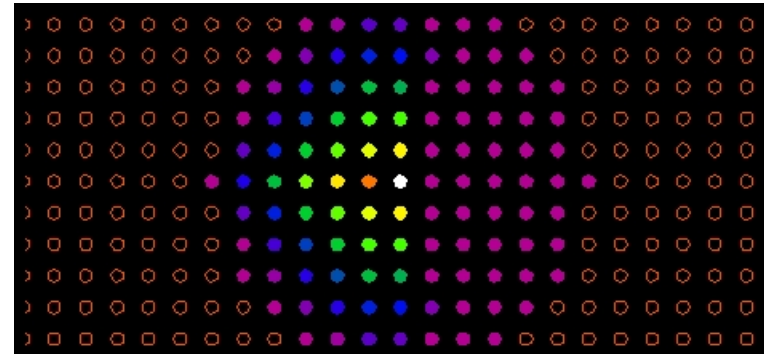




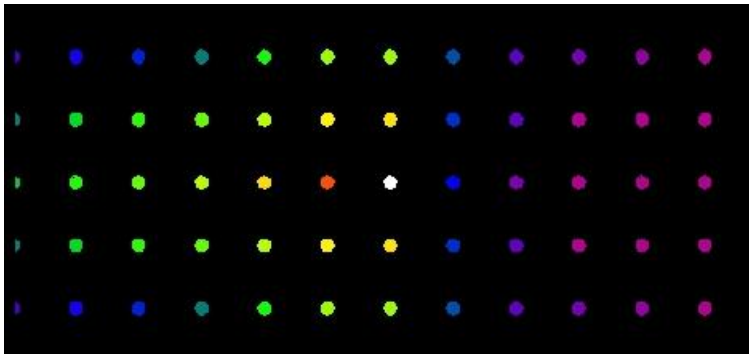
# Composite bar with defect: Coarsened micromodulus



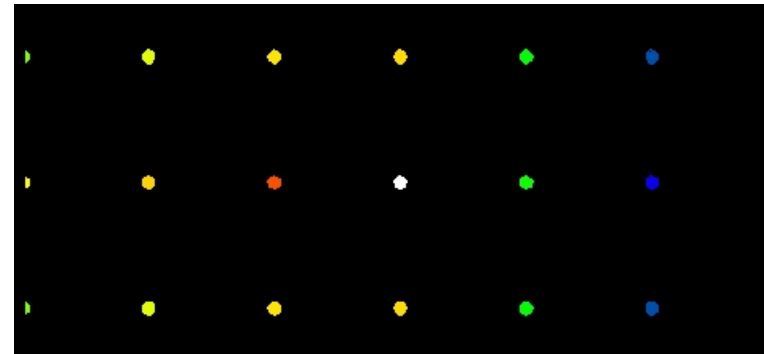
Level 0



Level 1



Level 2

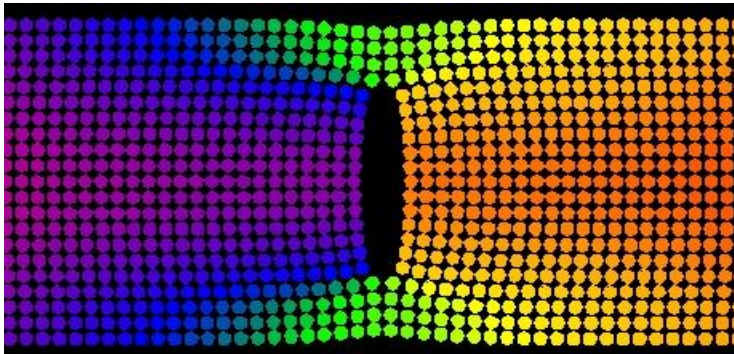


Level 3

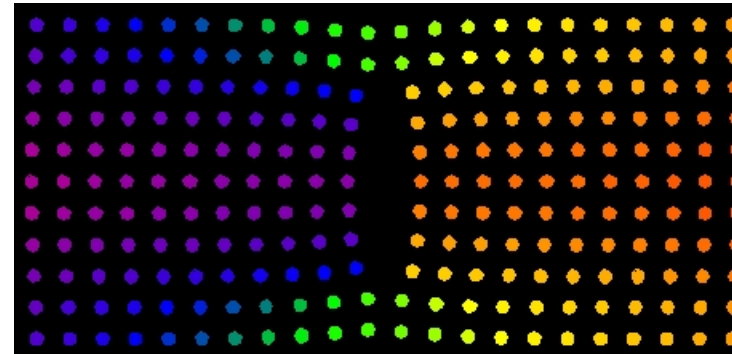
- Figures show contours of  $\log |C^m(\mathbf{x}, \cdot)|$  where  $\mathbf{x}$  is the white dot,  $m = 0, 1, 2, 3$ .
- This  $\mathbf{x}$  is near the crack surface.
- The effect of the crack on the micromodulus is visible in all levels.



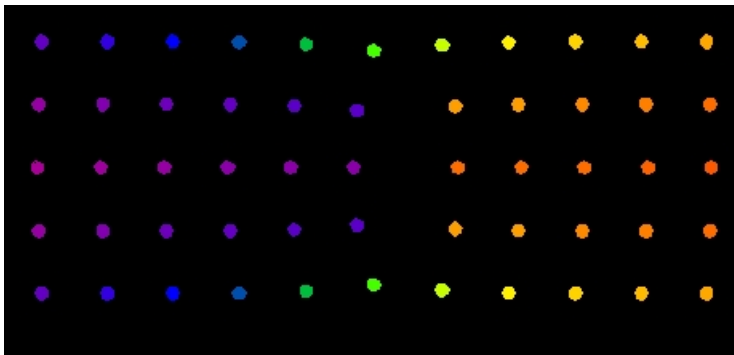
# Composite bar with defect: Displacements



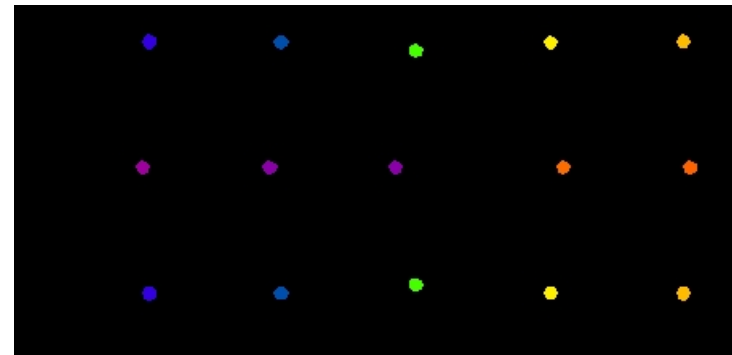
Level 0



Level 1



Level 2



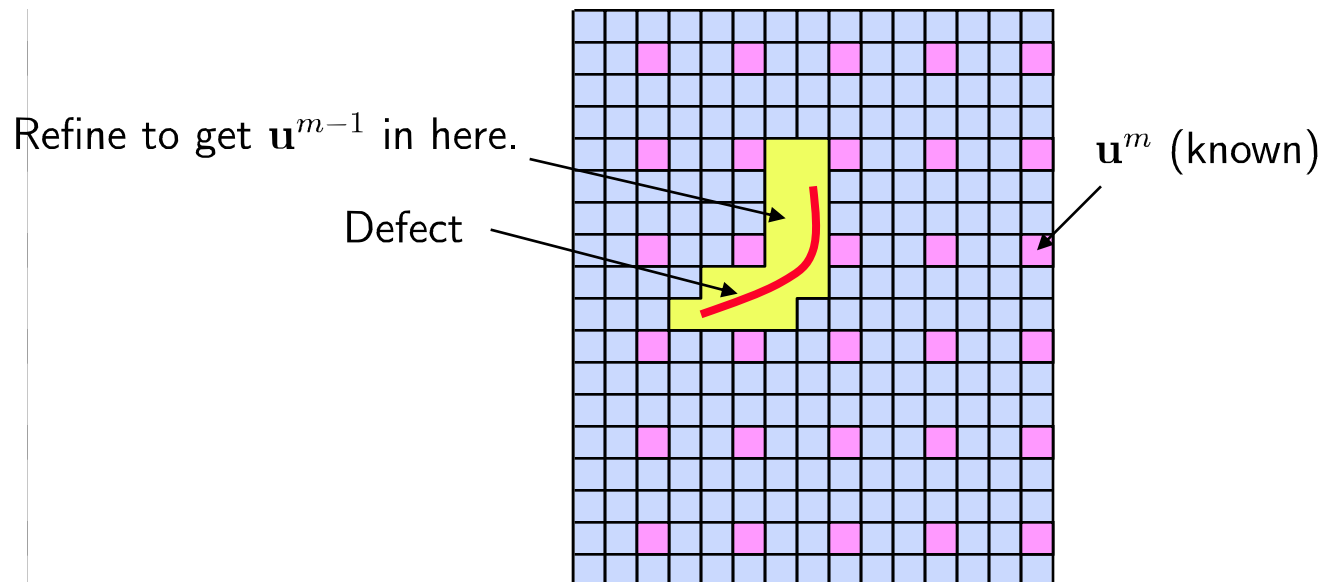
Level 3

- The same boundary value problem is solved at each level  $m$ , where  $m = 0, 1, 2, 3$ .
- Figures show contours of  $u_1^m$ .
- Displacements at all nodes and total force on ends agree between levels.

## Remarks

- In practice,  $S^m$  has to be found numerically by discretizing the level  $m-1$  equilibrium equation.
- We can *refine* in some region of interest by reversing the coarsening procedure.

$$\mathbf{u}^{m-1}(\mathbf{x}) = \int_{\mathcal{B}^m} \mathbf{S}^m(\mathbf{x}, \mathbf{q}) \mathbf{u}^m(\mathbf{q}) dV_{\mathbf{q}}$$



# Coarse graining in dynamics: Strong nonlocality in time

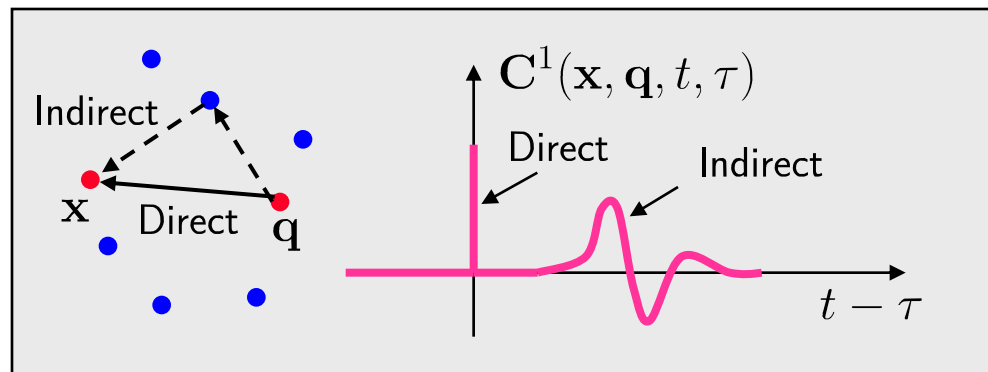
- Prescribe  $\mathbf{u}^1$  over time as well as position.
- Solution operator in  $\mathcal{B}^0$  is now time-dependent:

$$\mathbf{u}^0(\mathbf{x}, t) = \int_{-\infty}^{\infty} \int_{\mathcal{B}^1} \mathbf{S}^{0,1}(\mathbf{x}, \mathbf{q}, t, \tau) \mathbf{u}^1(\mathbf{q}, \tau) dV_{\mathbf{q}} d\tau.$$

- End up with time-dependent coarsened micromodulus:

$$\mathbf{L}^1(\mathbf{x}, t) = \int_{-\infty}^{\infty} \int_{\mathcal{B}^1} \mathbf{C}^1(\mathbf{x}, \mathbf{q}, t, \tau) (\mathbf{u}^1(\mathbf{q}, \tau) - \mathbf{u}^1(\mathbf{x}, t)) dV_{\mathbf{q}} d\tau$$

$$\mathbf{C}^1(\mathbf{x}, \mathbf{q}, t, \tau) = \int_{-\infty}^{\infty} \int_{\mathcal{B}^0} \mathbf{C}^0(\mathbf{x}, \mathbf{p}, t, \sigma) \mathbf{S}^{0,1}(\mathbf{p}, \mathbf{q}, \sigma, \tau) dV_{\mathbf{p}} d\sigma$$





# Discussion

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- Method accomplishes the following:
  - Two-way coupling between length scales (coarsening + refinement).
  - Forces on higher level DOFs exactly reproduce the original level 0 model.
  - Effects at smaller length scales are incorporated into larger length scales.
- Limitations:
  - Many issues remain regarding how to make the method efficient.
  - Need to re-linearize for evolving damage or other nonlinearities.

SS, A coarsening method for linear peridynamics, Int. J. Multiscale Computation Engineering (to appear).