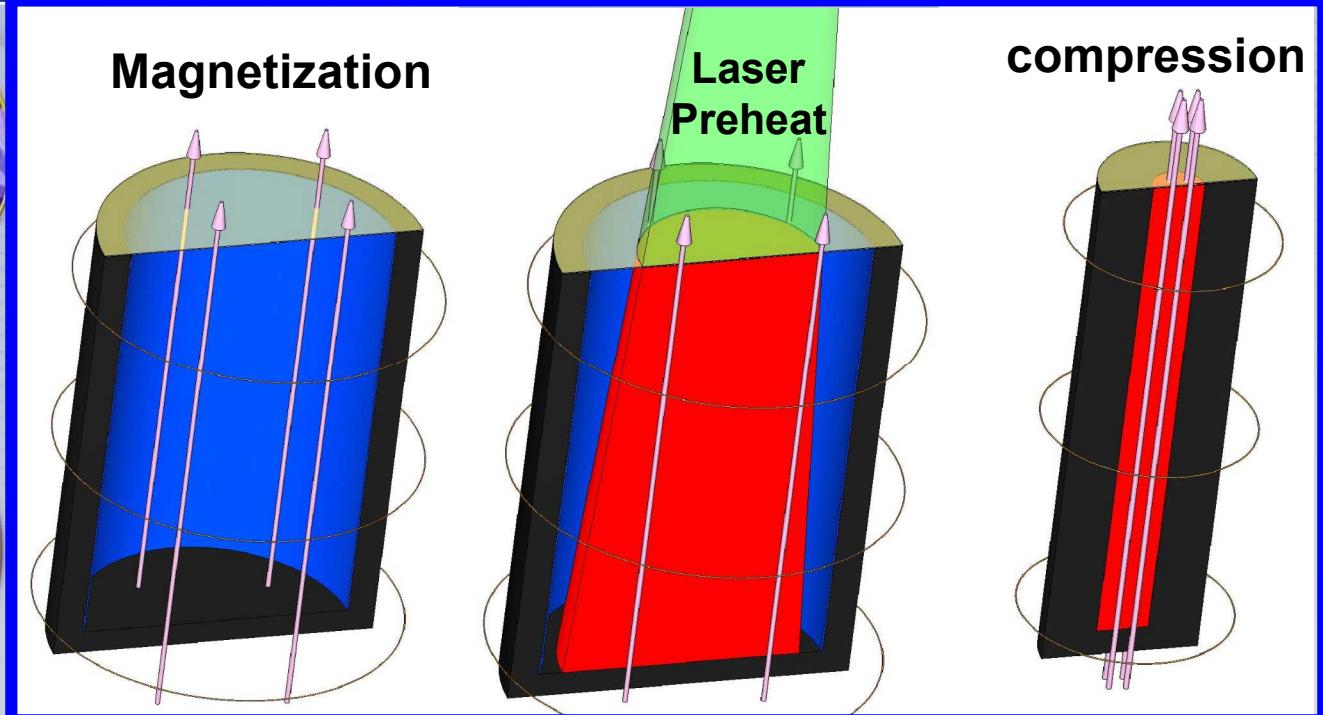




Blast Deflection of High-Yield Magnetized Liner Fusion Explosions

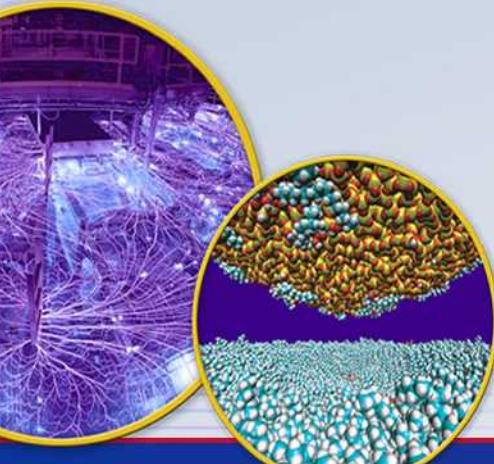
SAND2011-6420C



IFSA, Bordeaux September 11-16, 2011

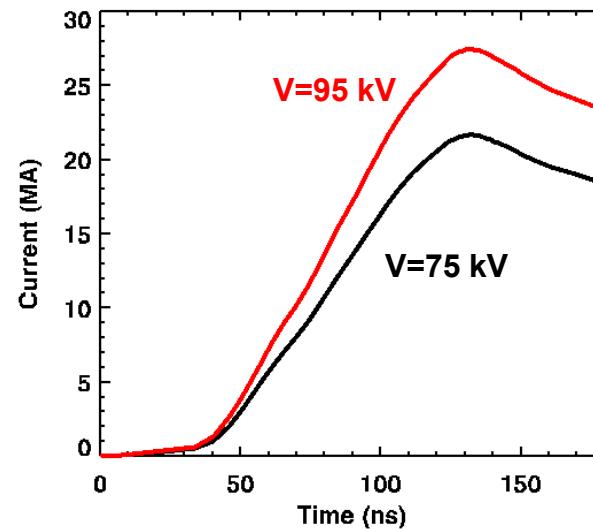
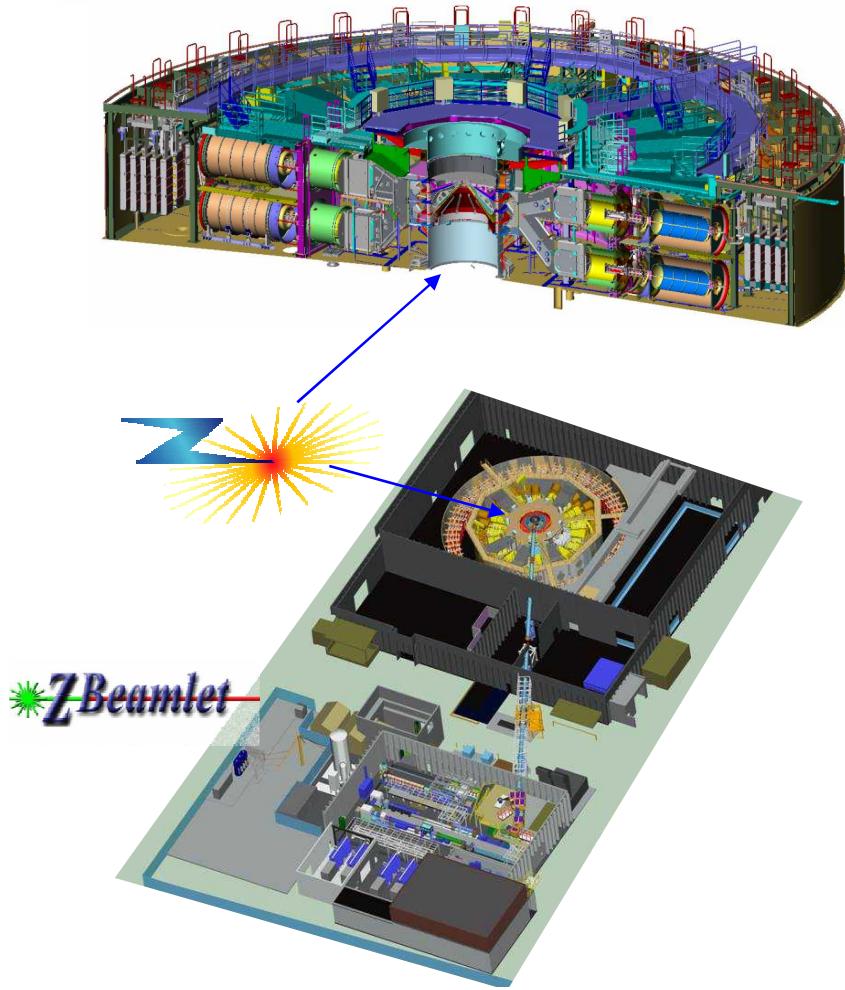
Stephen A. Slutz, and Roger A. Vesey

Sandia is a multiprogram laboratory operated by Sandia Corporation, a Lockheed Martin Company, for the United States Department of Energy's National Nuclear Security Administration under Contract DE-AC04-94AL85000.





The Z facility contains the worlds largest pulsed power machine and the Z-Beamlet and Z-Petawatt lasers



Magnetically-Driven Cylindrical Implosion

$$P = \frac{B^2}{2\mu_0} = 140 \left(\frac{I_{MA}/30}{R_{mm}} \right)^2 \text{ MBar}$$

140 MBar is generated by
300 eV radiation drive



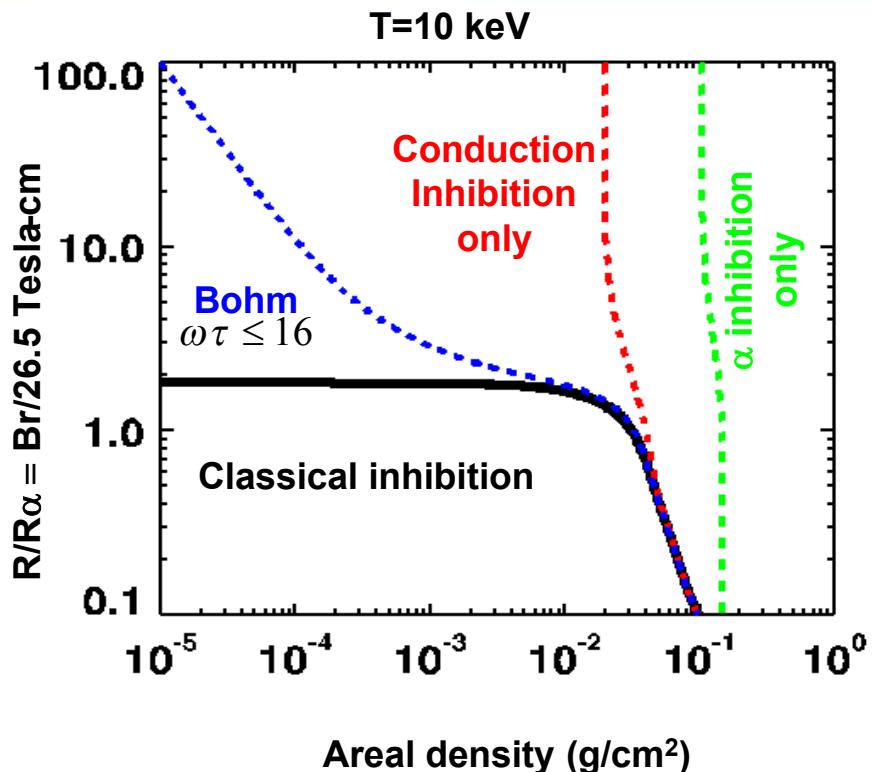
Direct magnetically driven implosion ~ 50 times more efficient than radiation drive

Natural geometry is cylindrical

- reduced volume compression (ρr and T_{ig} difficult)
- implosion velocity is slow $V_{imp} \sim 12 \text{ cm}/\mu\text{s}$ for instability-robust liners

Fuel magnetizing and preheating is a potential solution

- the attainment of ignition conditions with slow implosions and modest radial convergence



Axial α -trapping requires closed field lines or moderately high fuel density so that

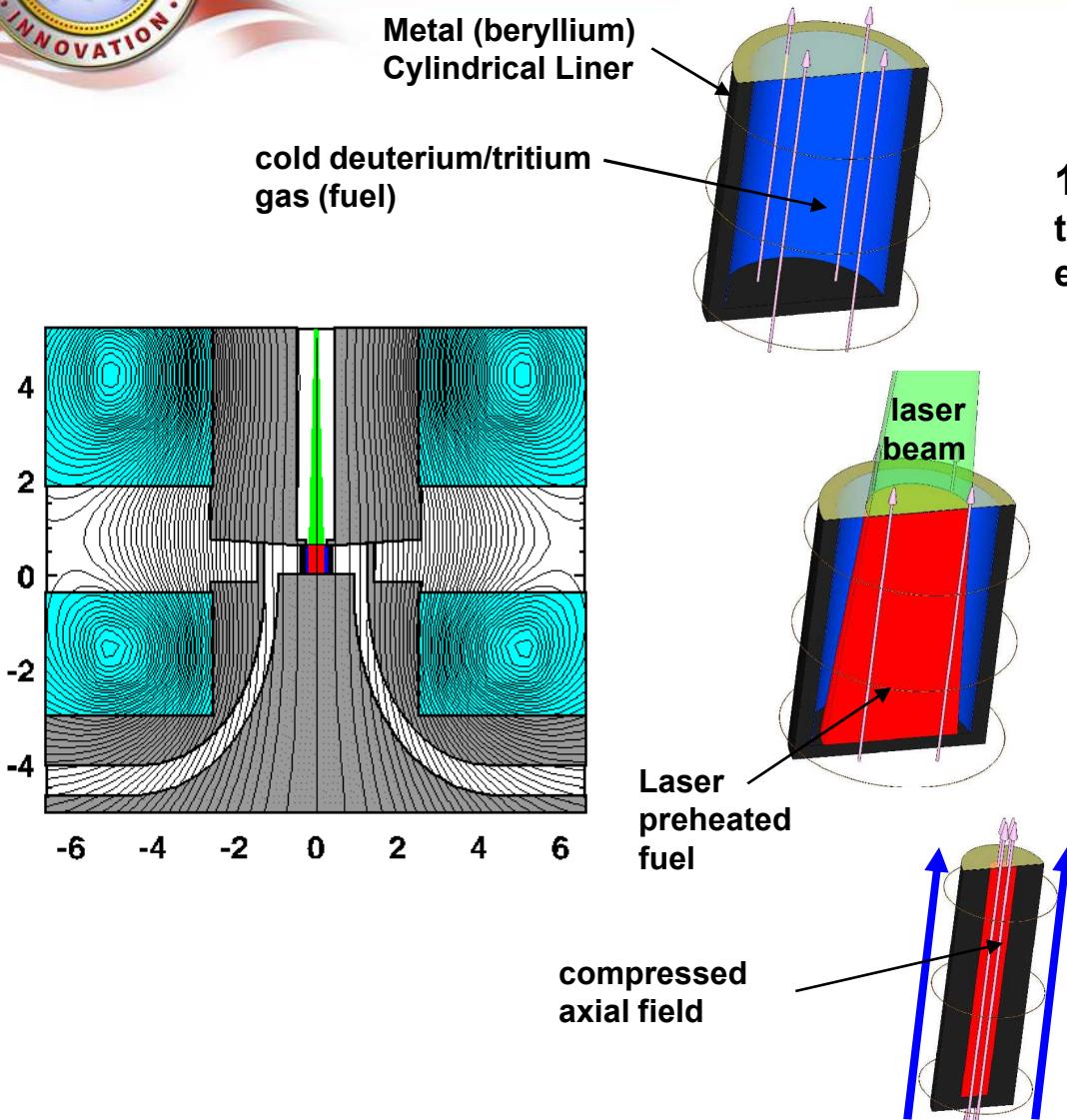
$$\rho\Delta z > 0.5 \text{ g}/\text{cm}^2$$



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The Z facility provides a unique opportunity to test the benefits of fuel magnetization and preheat*



* S.A. Slutz et al. *Physics of Plasmas* 17, 056303 (2010)

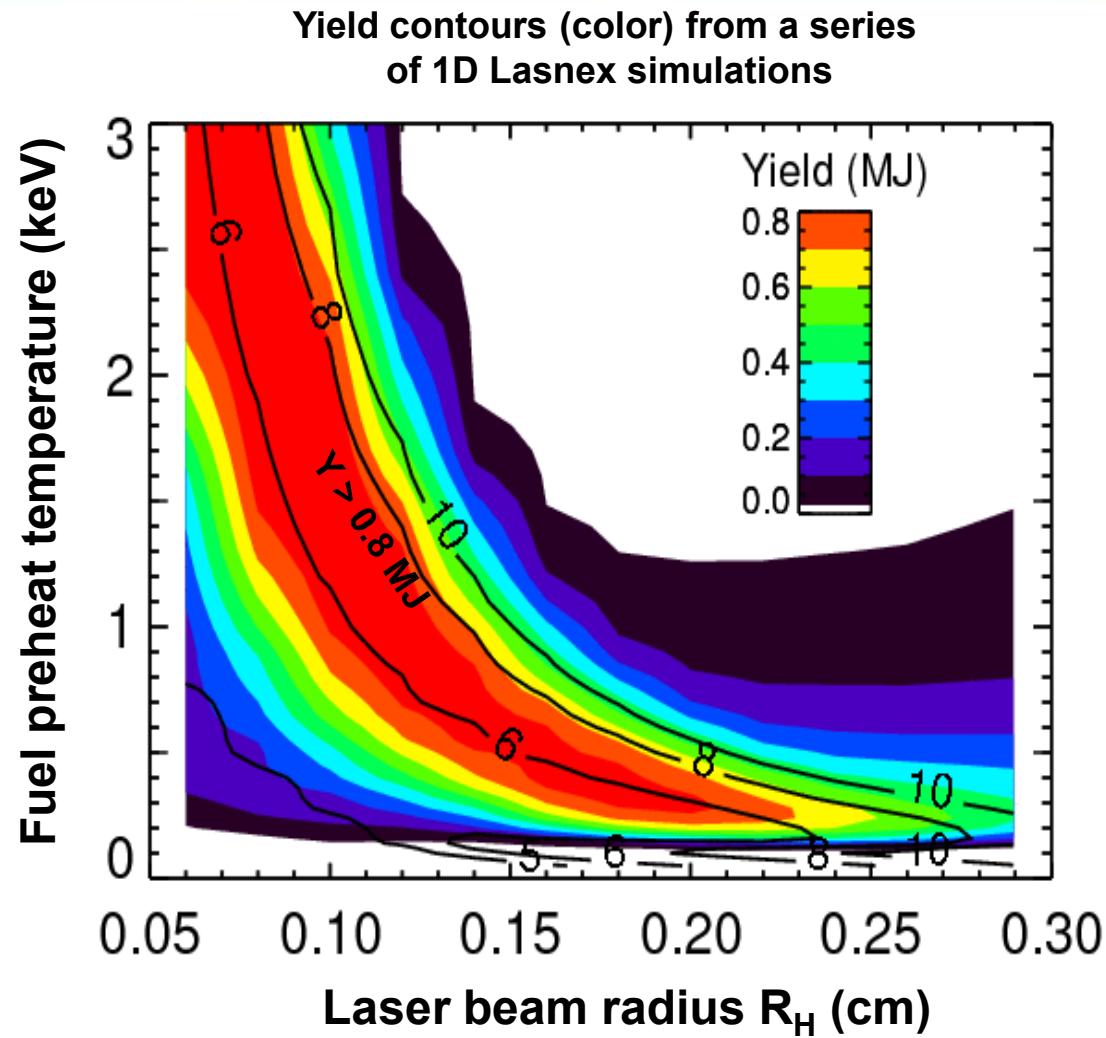
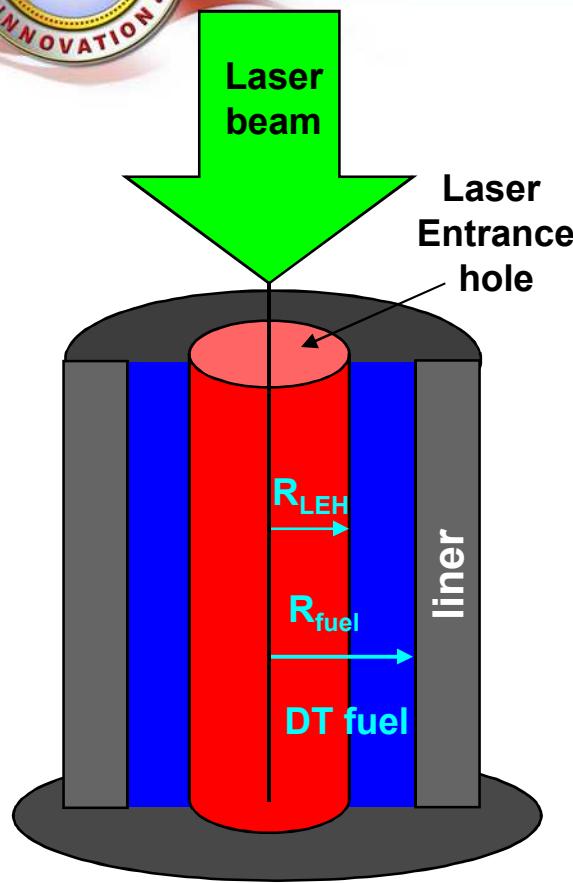
1. An axial magnetic field is applied to inhibit thermal conduction and enhance alpha particle deposition

2. Z Beamlet preheats the fuel

3. The Z accelerator efficiently drives a z-pinch implosion

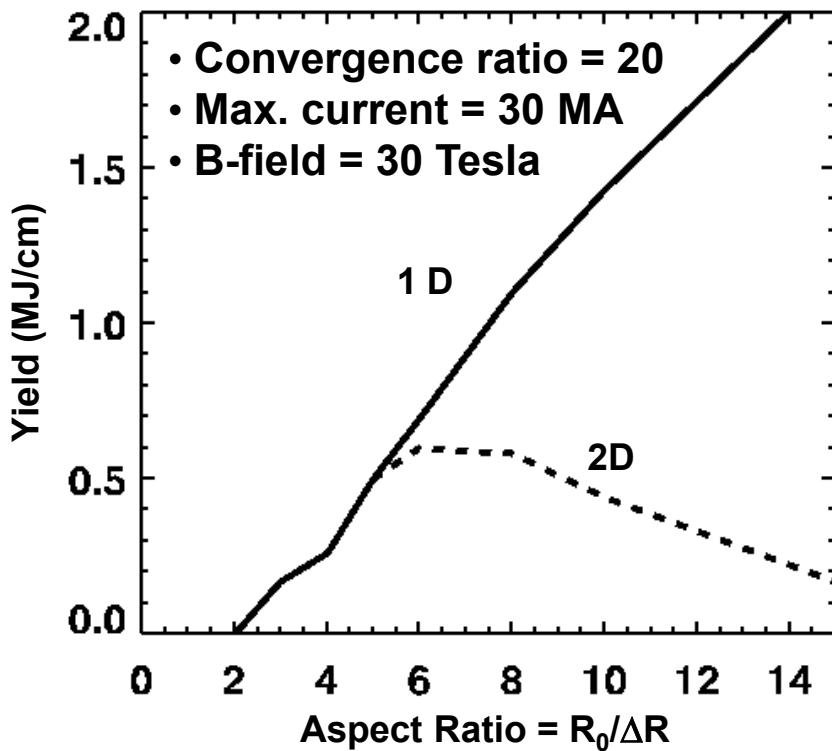


A small laser entrance hole (LEH) increases the yield

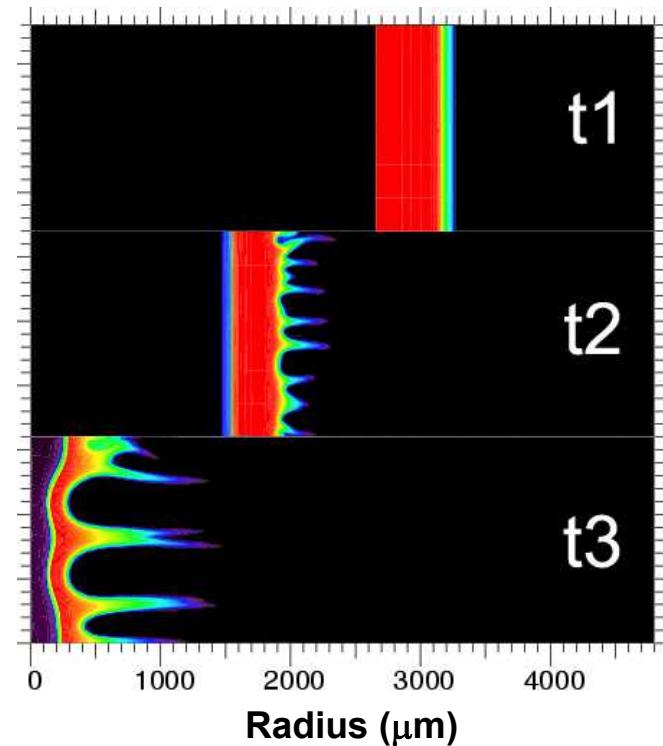




The MRT instability determines performance of MagLIF and other implosion systems

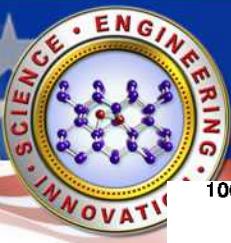


The Magneto-Rayleigh-Taylor instability degrades the yield as the aspect ratio is increased due to decreased liner ρr

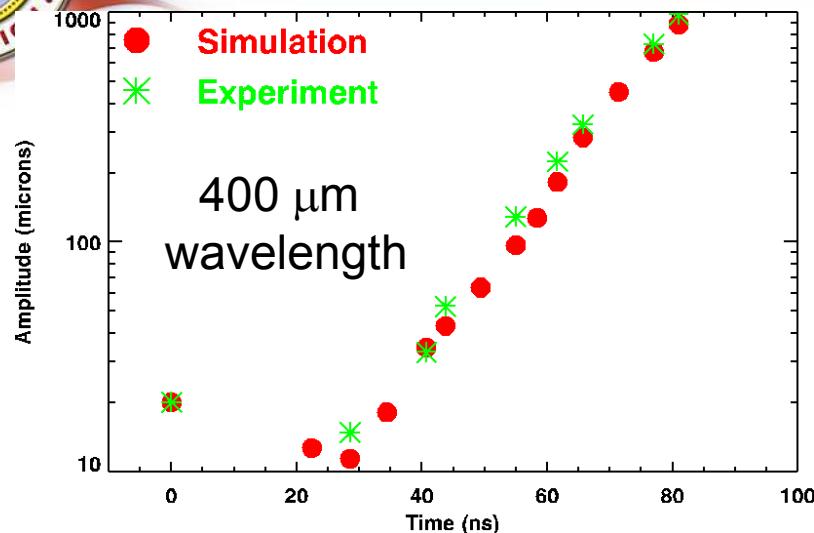


2D Lasnex Be liner simulations AR=6

- 60 nm surface roughness
- resolve waves down to $\sim 80 \mu\text{m}$
- wavelengths of 200-400 μm near stagnation

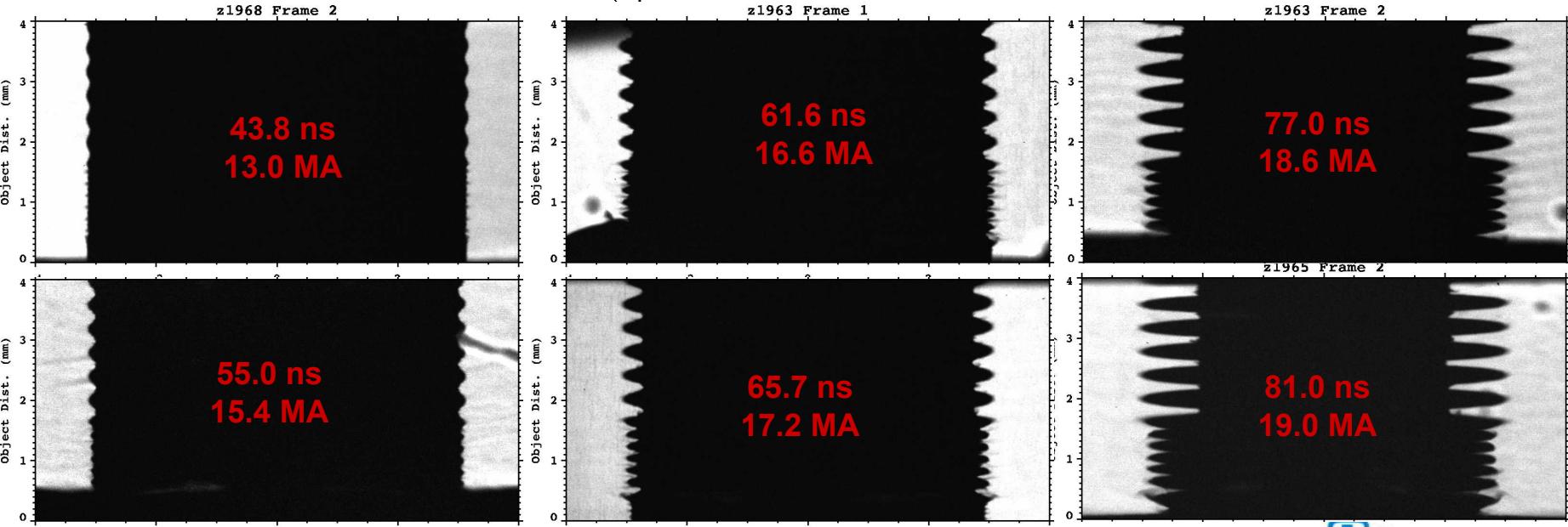


We have initiated studies of liner implosion stability, e.g. D. Sinars et al. PRL 105, 185001 (2010)



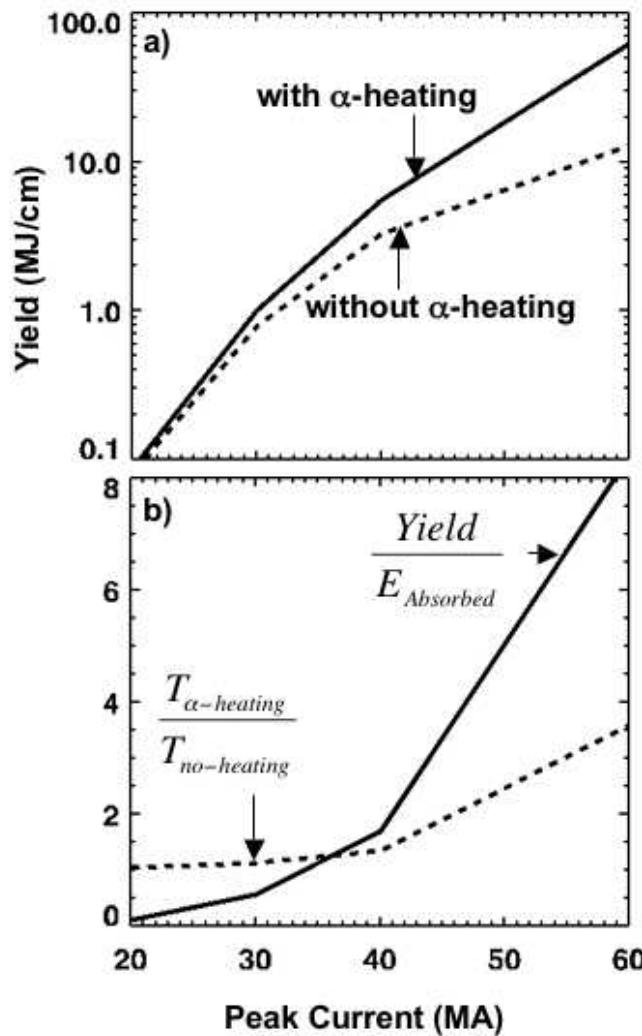
Experiments are progressing with beryllium liners

- Some information about the inner surface can be obtained due to the lower opacity
- Higher energy back lighters are under development





The yield is a moderate function of drive current ($Y \propto I^7$)



Liner parameters:

$R_0/\Delta R = 6$, CR = 20, B=30, Preheat temp~250 eV, Initial fuel density 2 - 5 mg/cc

A large amount of energy is needed to raise DT to the ignition temperature

Heating all of the fuel to ignition temperature “*Batch burn*” doesn’t lead to high gain!

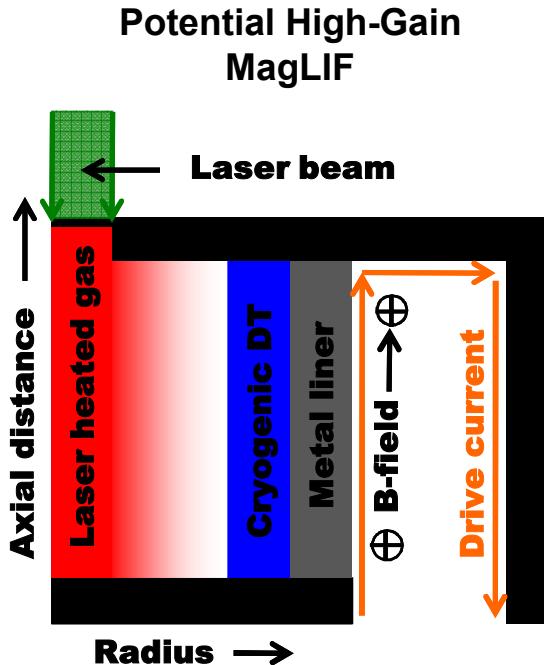
We need to determine if the burn can be propagated radially into a cold dense fuel layer.



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High Gain is required for Inertial Fusion Energy (IFE)



Fusion energy requires $G\eta_E f_{RP}\eta_D = 1$
 $G > 50$

G = Gain=Yield/liner absorbed energy
 η_E = electricity production efficiency~0.4
 f_{RP} = fraction of power recirculated~0.25
 η_D = driver efficiency ~ 0.2 for magnetic implosions

High Gain is obtained by only heating a small central portion of the fuel during compression

- minimizes compressive heating
- burn wave propagates into cold fuel as a deflagration

B_z decreases radial heat transport

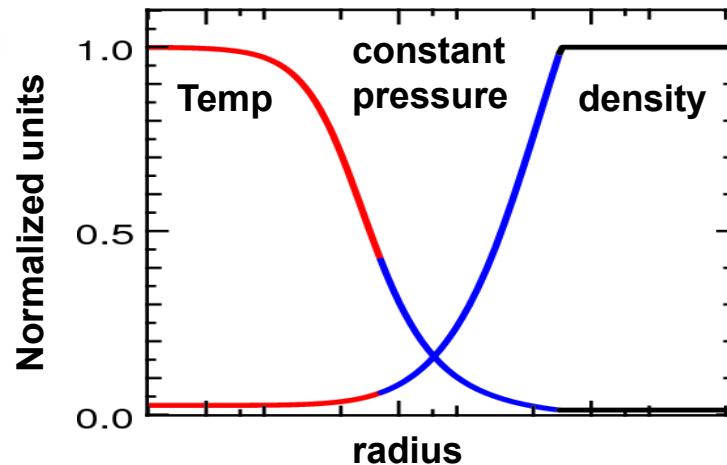
- increases self-heating in the hot spot
- decreases burn wave propagation into cold fuel

Jones and Mead, Nuclear Fusion 1986 suggested magnetized fuel is not compatible with propagating burn!

Is there a regime that allows both enhancement of hot spot ignition and burn wave propagation?



Radial propagation in magnetized fuel was studied with a simple stagnation model



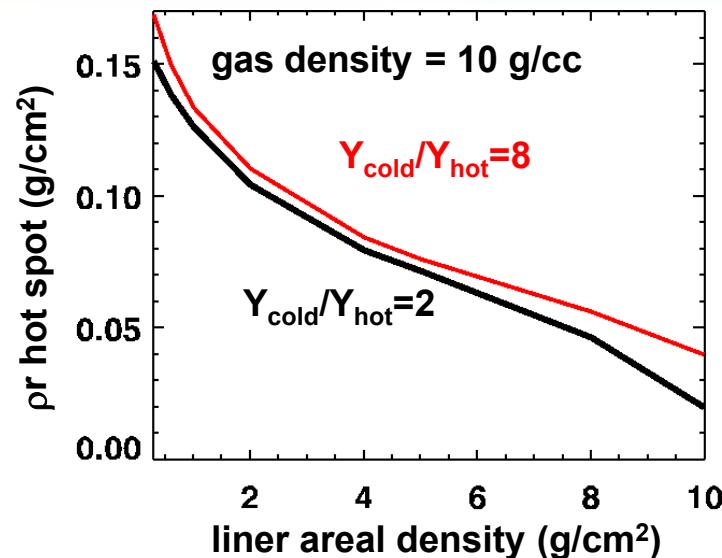
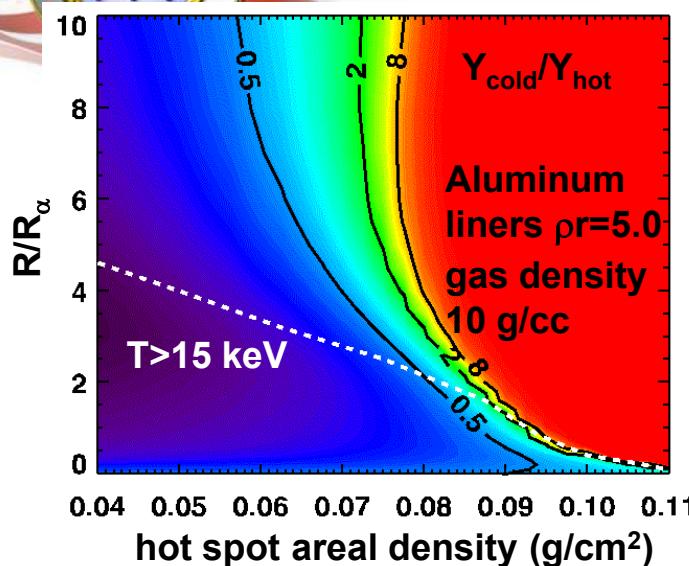
1D Lasnex simulations were performed starting at an idealized stagnated state with the following attributes:

- pressure equilibrium
- central hot spot $T_{ion}=10$ keV @ center
- specific areal density of the hot spot and the surrounding metal liner were varied

The ratio of the yield from the cold ice layer to the yield from the hot spot is a measure of the propagation, i.e. Y_{cold}/Y_{hot}



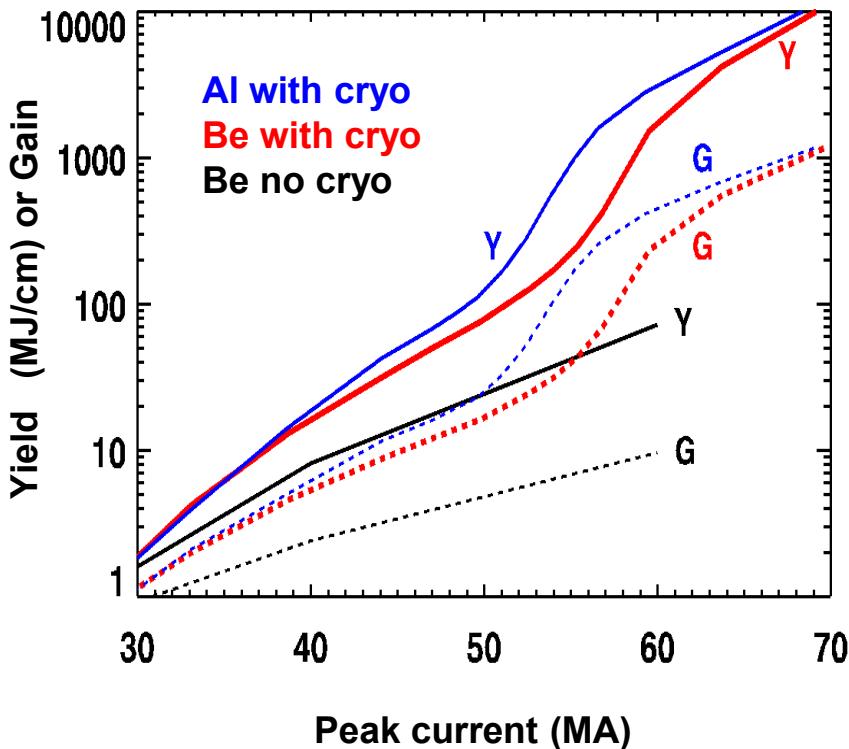
There is an optimal magnetic field strength for radial propagation



- The inhibition of thermal conduction and alpha transport lowers the hot spot areal density required for ignition
- Too large a B-field inhibits propagation into the cold fuel
- Minimum hot spot areal density is a weak function of gas density, but is significantly affected by the areal density of the liner
- More confinement time allows slower burn propagation!



Gains >1000 are possible with MagLIF!



Further optimization is possible

- fuel is not on a low isentrope
- no current pulse shaping

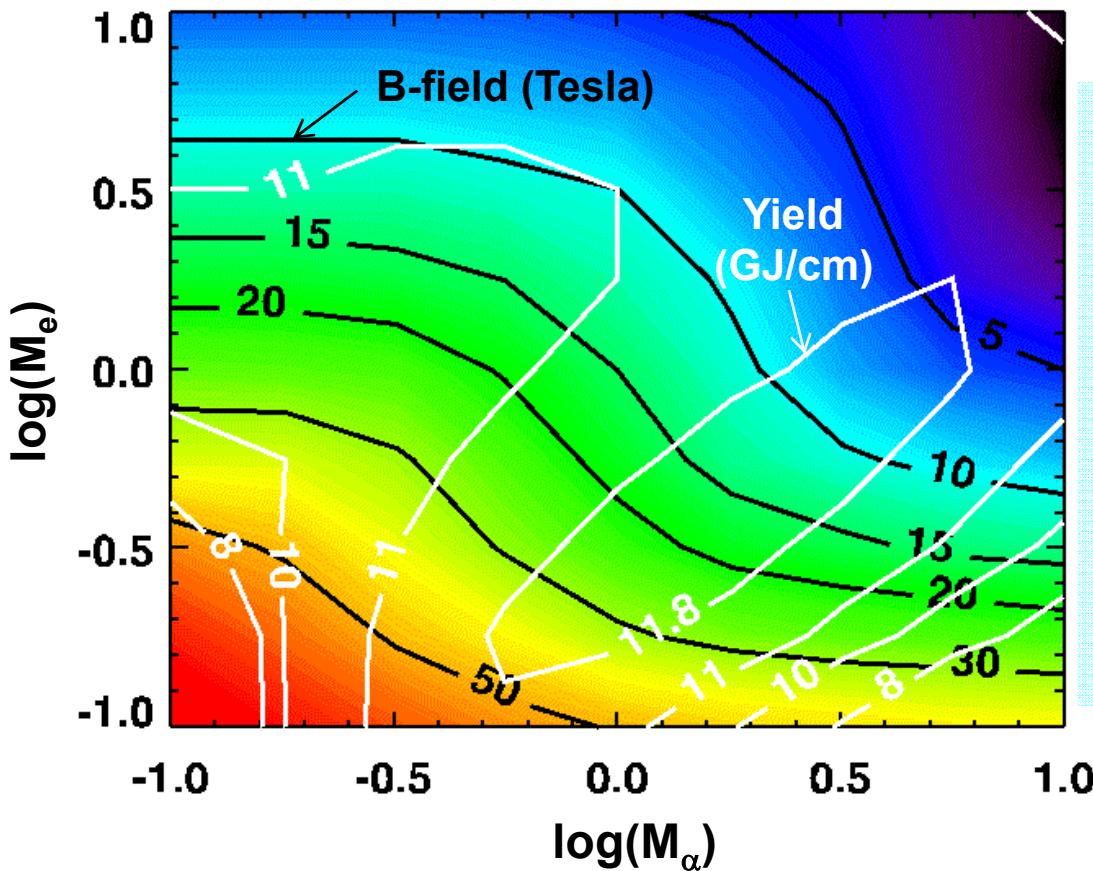
Margin increases with current
• >50% as $I \Rightarrow 70$ MA

Initial gas density increases with current
• not at vapor pressure
• laser energy => 22 KJ
• pulse length => 30 ns





High Gain MagLIF is robust to errors in magnetic inhibition modeling



Magnetic inhibition of thermal conduction depends on $\omega_e \tau_e$

$$K_{cond} = K_0 F(M_e \omega_e \tau_e)$$

Magnetic inhibition of α -particle transport depends on $\omega_\alpha \tau_{\alpha e}$

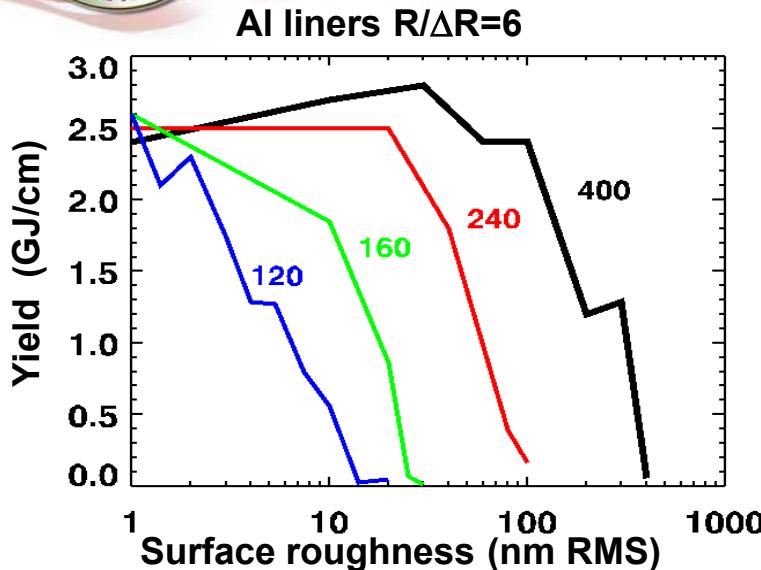
$$D_\alpha = D_{\alpha 0} G(M_\alpha \omega_\alpha \tau_{\alpha e})$$

Initial B-field is varied to obtain maximum yield for (M_e, M_α)





The Magneto-Rayleigh-Taylor instability is always a concern



The simulated yields are not robust to short wavelengths

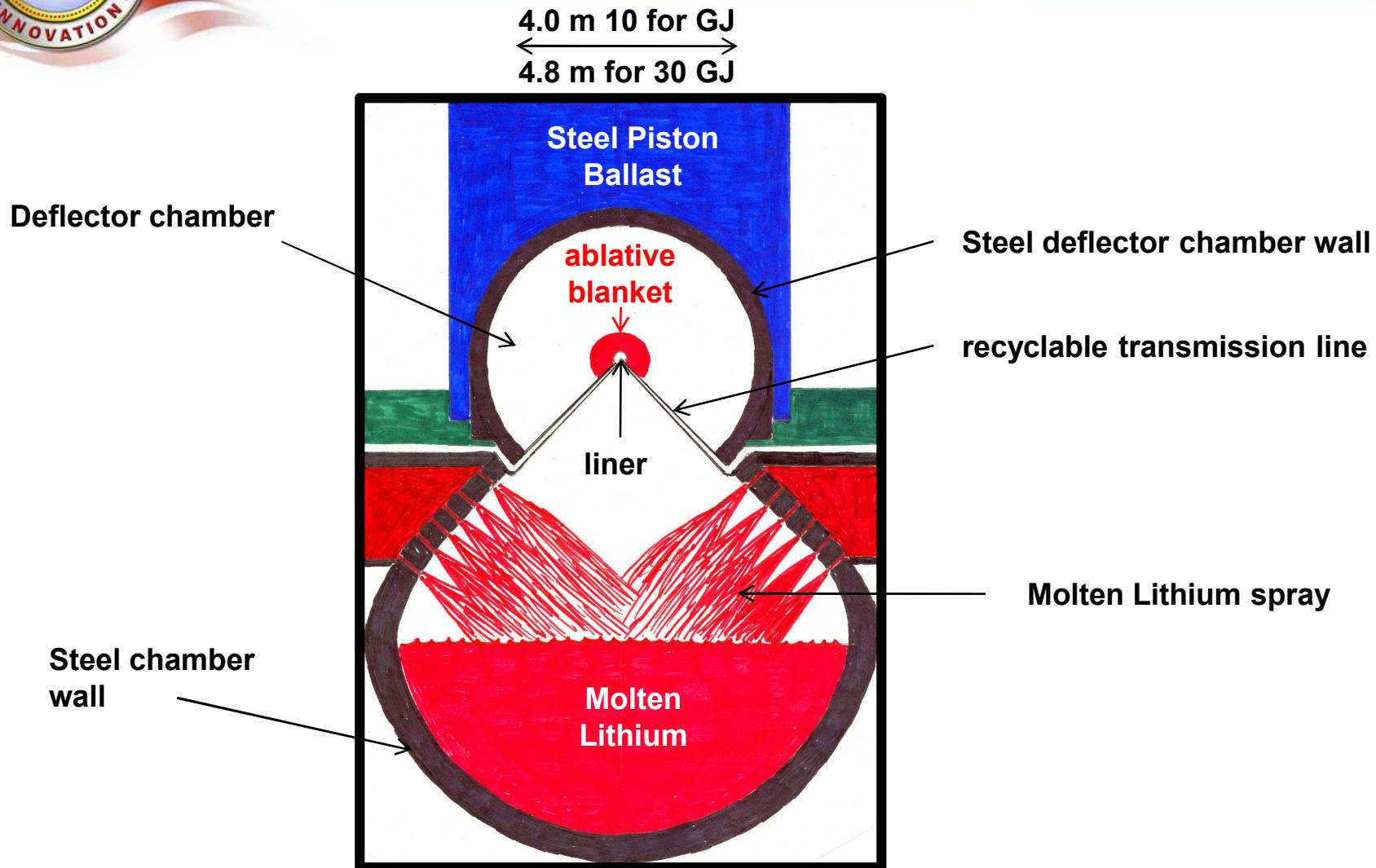
- 2D simulations had stronger growth than experiments at short wavelengths (Sinars et al. PRL 105, 185001)
- 2D may be pessimistic at short wave lengths
- 3D simulations and experiments needed for clarification
- RMS amplitudes less than 20 nm may be required and achieved with polishing

There are several possible ways to mitigate the effect of the MRT

- The axial B-field reduces the growth rate on the fuel/liner interface (Ryutov)
- An initially low density liner would have a large B-field when compressed which could suppress the instability
- Nested wire arrays mitigate the MRT. Multi-shell liners may also be effective
- Pulse shaping can avoid shocks and retain material strength



Large yields (10-30 GJ) can be contained in compact chambers using the concept of a deflector



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Blast containment in a spherical vessel

Tritium breeding and neutron protection can be provided by an ablative blanket (AB) surrounding the capsule, 40 cm of LiH is sufficient

Lasnex simulations indicate that nearly 100% of the capsule yield is converted to kinetic energy of the AB, which delivers a momentum pulse to the containment vessel

$$\Pi \approx 2(2YM_{AB})^{1/2}$$

The vibration of the steel shell with aspect ratio, A_R , Young's modulus, E , density, ρ radius, r_0 , shell is given by the equation

$$\ddot{S} + \frac{2E}{\rho r_0^2} S = \frac{PA_R}{\rho r_0^2} \quad S = \frac{\Delta r}{r}$$

The maximum amplitude for a tensile strength, T is $S_x = \frac{T}{E}$

The deflector geometry allows energy to escape from the AB reducing the momentum by a factor of 3-4. The deflector chamber radius is then

$$R_{def} > 2.2 \left(\frac{Y}{10GJ} \right)^{1/6} \left(\frac{A_R}{6} \right)^{1/3} m$$



Summary

We have presented detailed simulations indicating that High-Gain Magnetized Inertial Fusion is possible

- This may lead to a low cost means to generate electricity with fusion
- High-Gain MagLIF is obtained by adding a cryogenic layer of DT to the inside of the liner
- High-Gain MagLIF with yields of 10 GJ/cm could allow the low rep-rates required by recyclable transmission lines

We have presented the concept of the deflector, which enables the containment of large yields (10-30 GJ) with relatively short recyclable transmission lines (2-3 m).