

# Propagating Uncertainty from Simulation Parameters and Sampling Noise through Coupled Atomistic-to-Continuum Systems

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Supported by the US Department of Energy, Office of Advanced  
Scientific Computing Research

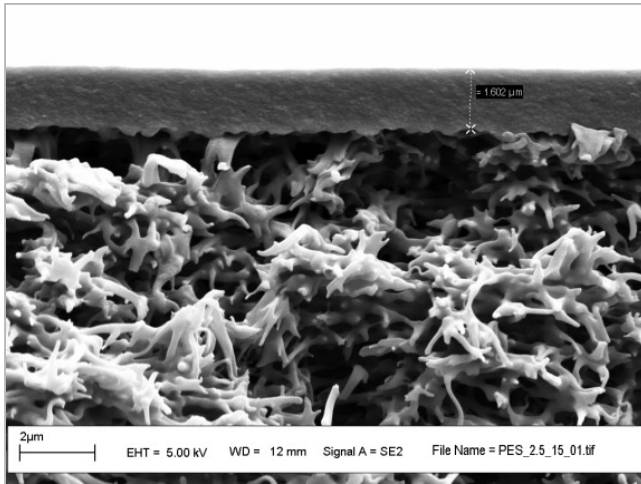
# Quantifying prediction fidelity in multiscale multiphysics simulations: the overall team

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  - Helgi Adalsteinsson
  - Bert Debusschere
  - Maher Salloum
  - Reese Jones
  - Khachik Sargsyan
- Johns Hopkins University
  - Omar Knio
  - Francesco Rizzi
- Texas Tech University
  - Kevin Long
  - Kaleb McKale
  - Jed Gohlke
  - Simon Rush
- Massachusetts Institute of Technology
  - Youssef Marzouk
  - Jinglai Li

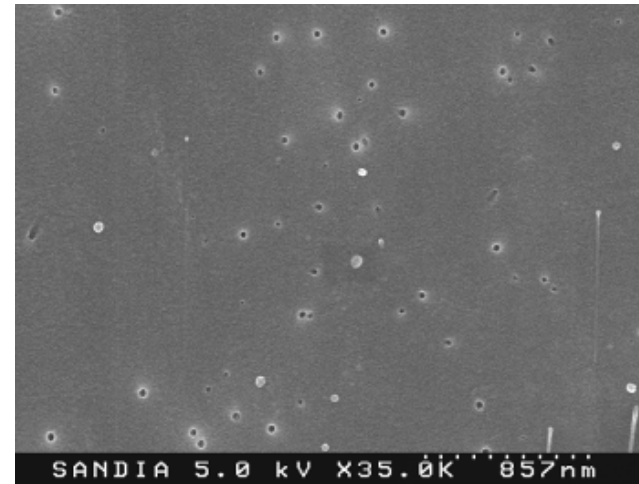
# Overview

- Introduction and motivation
- Couette flow test case
- Multiscale coupling approach
  - Atomistic sampling noise
  - Parametric uncertainty
- Ongoing work
- Conclusions

# Multiscale methods needed to account for phenomena coupled over wide ranges of time and length scales



**Nanoporous polyaniline membrane**

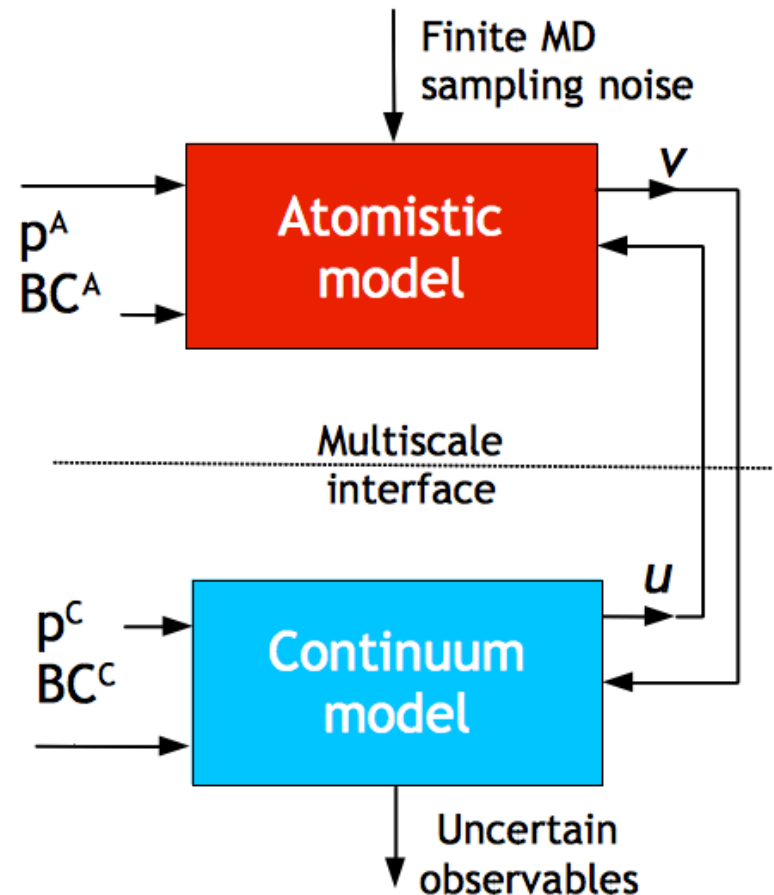


**Polycarbonate track etched membrane**

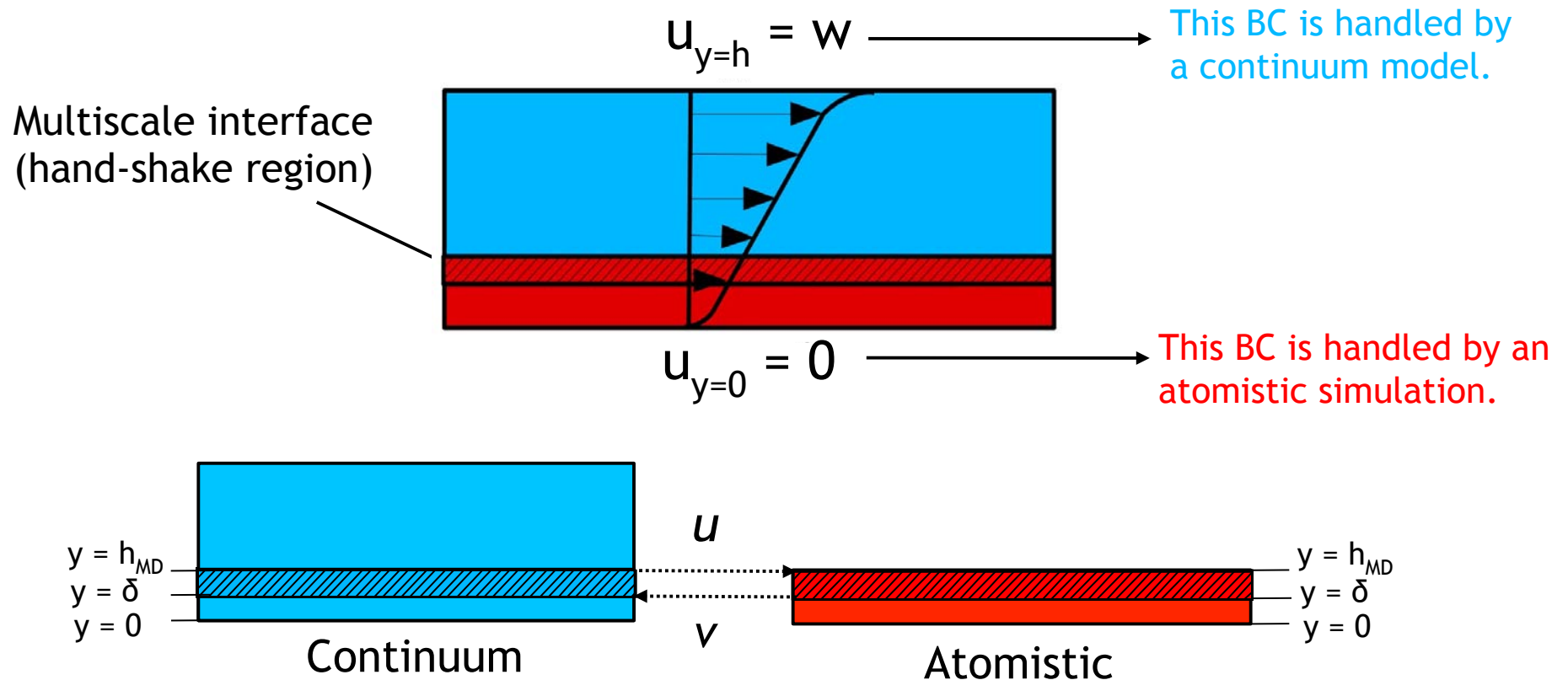
- Many key applications have macroscale behavior driven by microscale phenomena
  - E.g. ionic flux through nanopores in water desalination
- Multiscale simulations resolve key physics on different scales
  - Uncertainties on all scales
  - Uncertainties in coupling

# Predictive multiscale simulation requires quantification of the many sources of uncertainty

- This talk focuses on
  - Assessment of sampling noise on the atomistic level
  - Propagation of parametric uncertainty and sampling noise across scales
- Related work presented at this meeting
  - Forward propagation and inference of parametric uncertainty on atomistic level: Knio *et al.*, poster 1
  - Propagation of uncertainty through the continuum level: Long *et al.*, poster 48

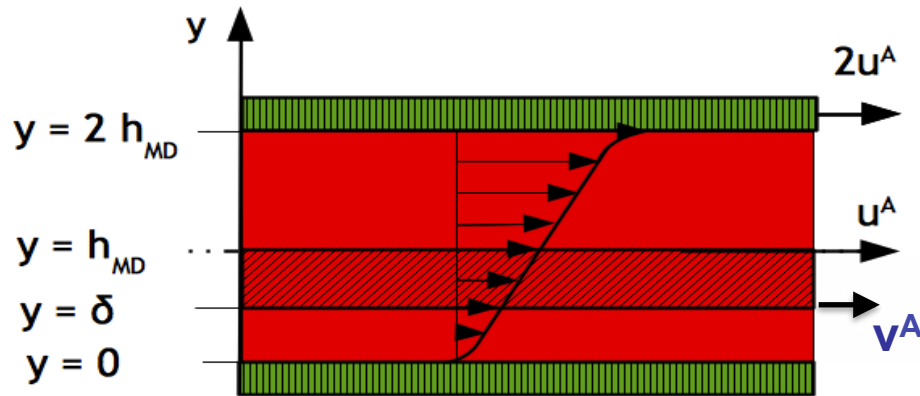


# Canonical plane Couette flow is used as model problem for algorithm development



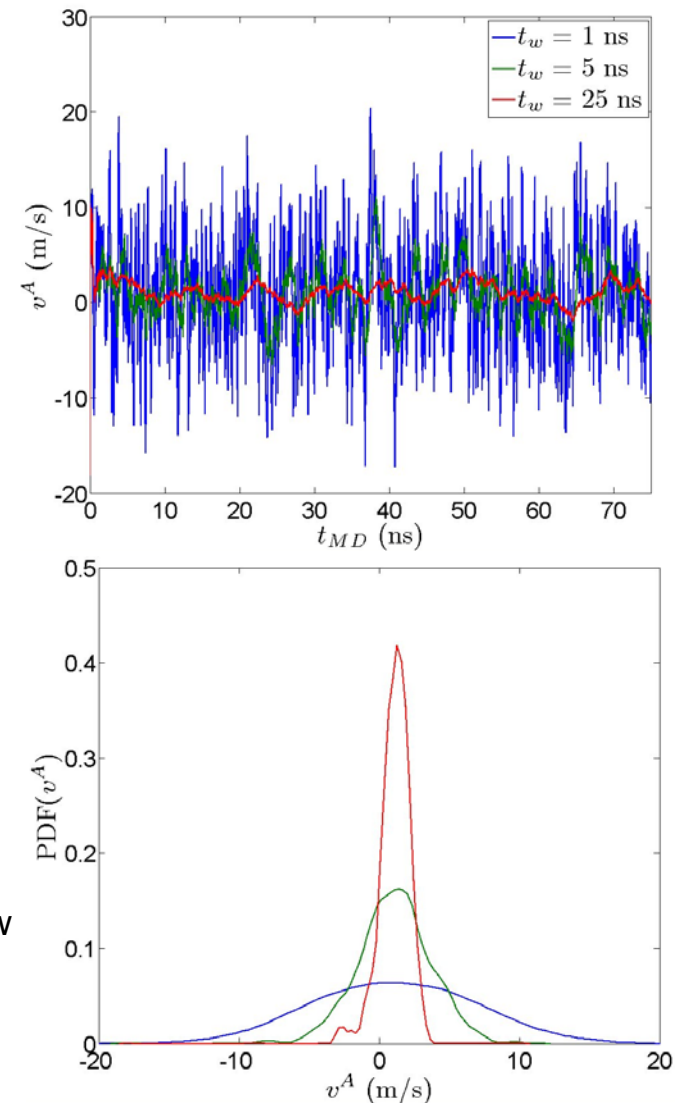
- Complex enough to illustrate key challenges in coupling atomistic to continuum with uncertainty
- Simple enough to make computations tractable

# Finite sampling in extracting macroscale observables from atomistic simulations results in uncertainty



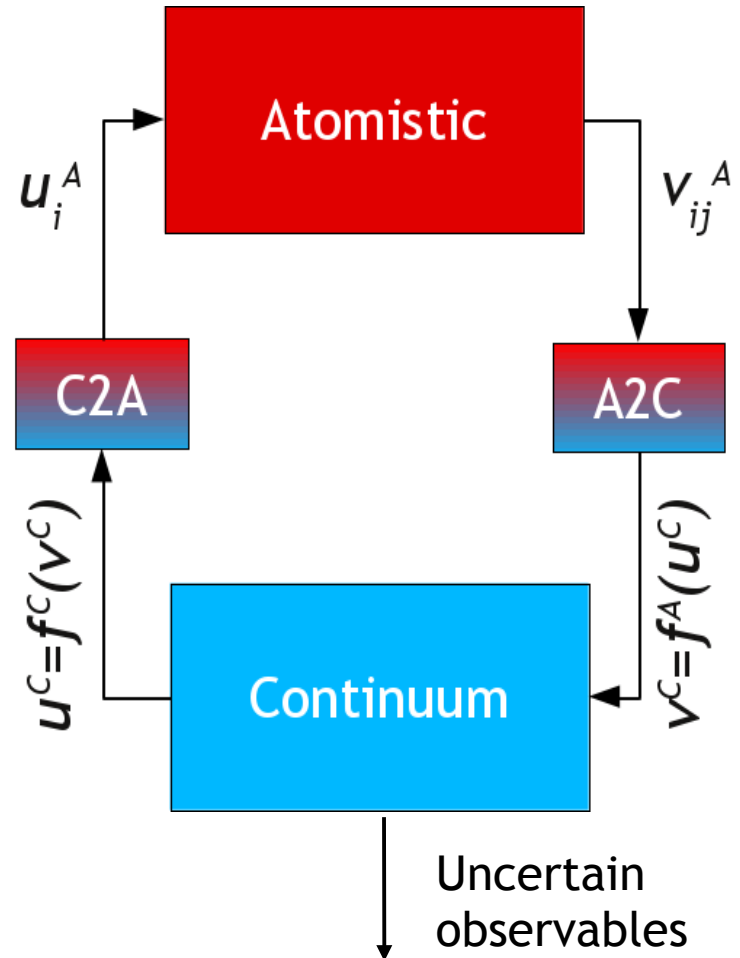
- Molecular dynamics deterministic
  - Lennard Jones
  - LAMMPS
- Sampling noise due to
  - Finite size of domain
  - Finite size of averaging time window  $t_w$
- For non-trivial MD simulations, we can not sample our way out of this

$$\phi_{ij} = 4\phi_0 \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right]$$



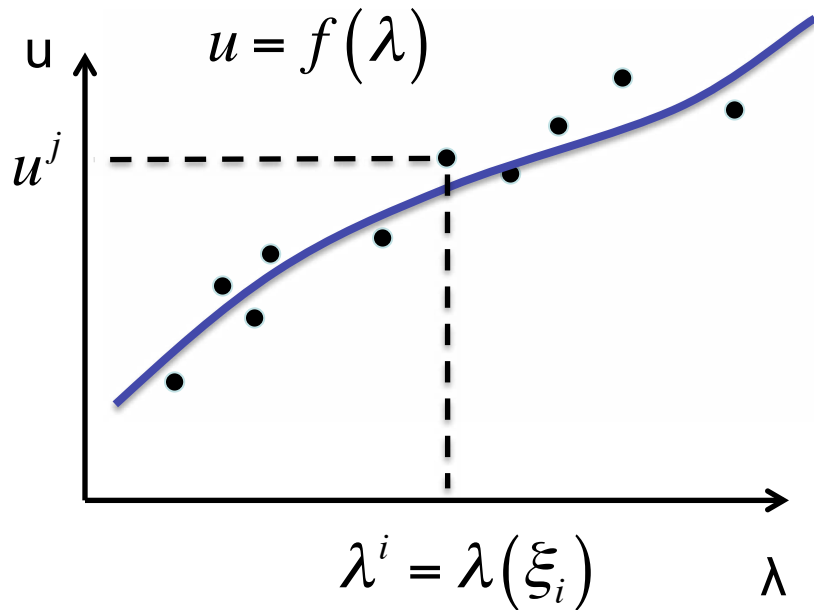
# Building blocks in atomistic to continuum coupling

- Characterizing sampling noise in atomistic output
- Uncertain inputs and sampling noise in atomistic simulation
- Continuum simulation with uncertain inputs
- Coupled atomistic and continuum simulation
  - Sampling noise only
  - Sampling noise and uncertain inputs





# Bayesian inference of response surfaces



Analytical expression for  $\mathbf{u}$  as student-t process if:

- Gaussian noise model
- Infinite (improper) uniform prior on  $u_k$
- Jeffreys prior for  $s^2$
- Marginalize over  $s^2$

$$\lambda = \sum_{k=0}^P \lambda_k \psi_k(\xi) \quad u = \sum_{k=0}^P u_k \psi_k(\xi)$$

$$\mathbf{d} = \{u^j\}_{j=1}^N \quad \mathbf{u} = \{u_k\}_{k=0}^P$$

$$u^j = f(\lambda^i) + s \eta_{ij} \quad \eta_{ij} \sim \mathcal{N}(0,1)$$

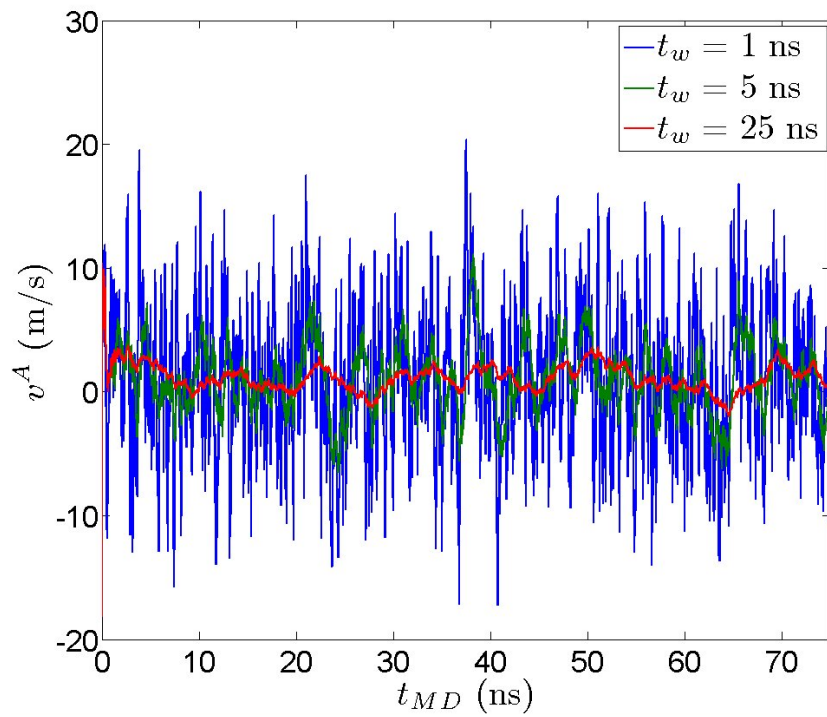
$$P(\mathbf{u}, s^2 | \mathbf{d}) \propto P(\mathbf{d} | \mathbf{u}, s^2) P(\mathbf{u}, s^2)$$

$$\mathbf{u} \sim St(\bar{\mathbf{u}}, S, \gamma)$$

$$u = \psi(\xi)^T \bar{\mathbf{u}} + \zeta \sqrt{\psi(\xi)^T S \psi(\xi)}$$

$$\zeta \sim St(0,1,\gamma)$$

# Quantification of sampling noise in atomistic simulations



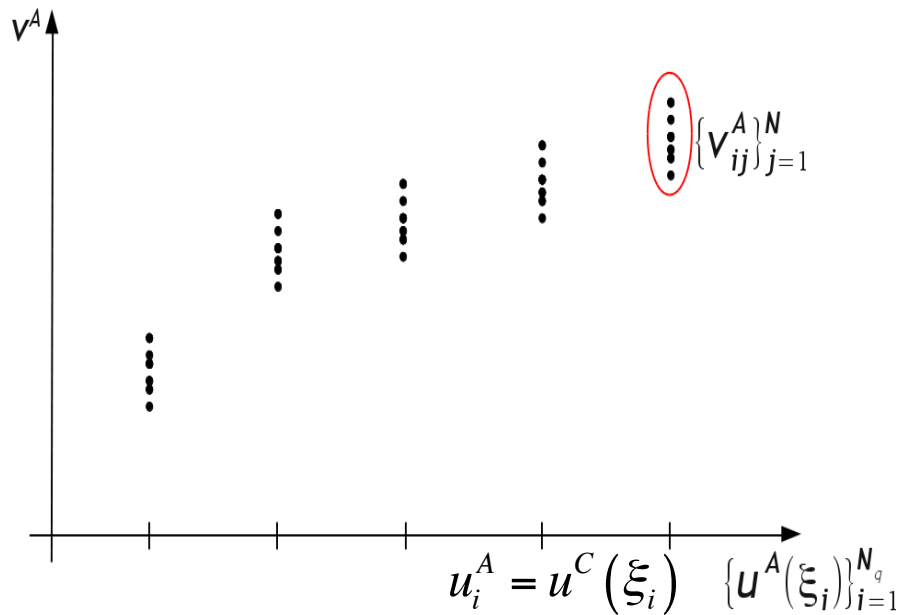
$$\mathbf{d} = \left\{ v_j^A \right\}_{j=1}^N$$

$$v_j^A = v^C + s \eta_j$$

$$v^C = \bar{v}^C + \sqrt{S} \xi \quad \xi \sim St(0, 1, \gamma)$$

- Infer  $v^C$  from  $N$  short-term averaged MD velocity samples  $v_j^A$ 
  - Gaussian model for data noise due to Central Limit Theorem (CLT)
  - Analytical solution gives  $v^C$  as student-t random variable
- Averaging over longer time window or more data reduces sampling noise

# Propagating parametric uncertainty and sampling noise through atomistic simulations



$$v_{ij}^A = v^C(\xi) + s \eta_{ij}$$

$$u^C = \sum_{k=0}^P u_k^C \psi_k(\xi)$$

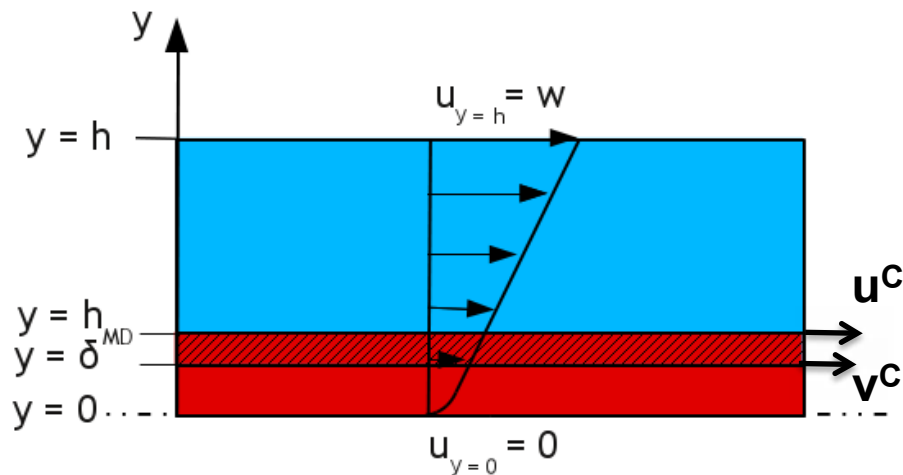
$$v^C = \sum_{k=0}^P v_k^C \psi_k(\xi)$$

$$v^C = \psi(\xi)^T \bar{\mathbf{v}}^C + \xi \sqrt{\psi(\xi)^T S \psi(\xi)}$$

$$\xi \sim St(0, 1, \gamma)$$

- Infer  $v^C$  from  $N$  short-term averaged MD velocity samples  $v_{ij}^A$ 
  - Sampled over a range of input velocities  $u_i^A$
  - Gaussian model for data noise due to Central Limit Theorem (CLT)
  - Analytical solution gives  $v^C$  as student-t process over input velocity uncertainty

# Propagating uncertainty through continuum



$$u^C = w - \frac{h - h_{MD}}{h - \delta} (w - v^C)$$

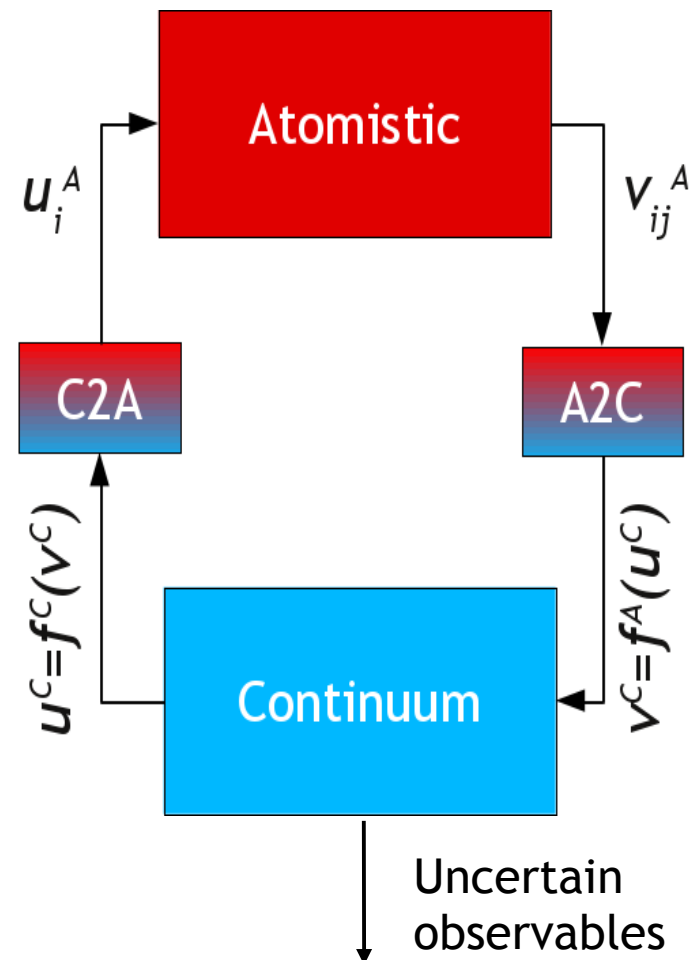
$$u^C = \sum_{k=0}^P u_k^C \psi_k(\xi) \quad v^C = \sum_{k=0}^P v_k^C \psi_k(\xi)$$

$$u_k^C = \begin{cases} w + \frac{(w - v_k^C)(h_{MD} - h)}{h - \delta}, & \text{for } k = 0 \\ \frac{-v_k^C(h_{MD} - h)}{h - \delta}, & \text{for } 0 < k \leq P \end{cases}$$

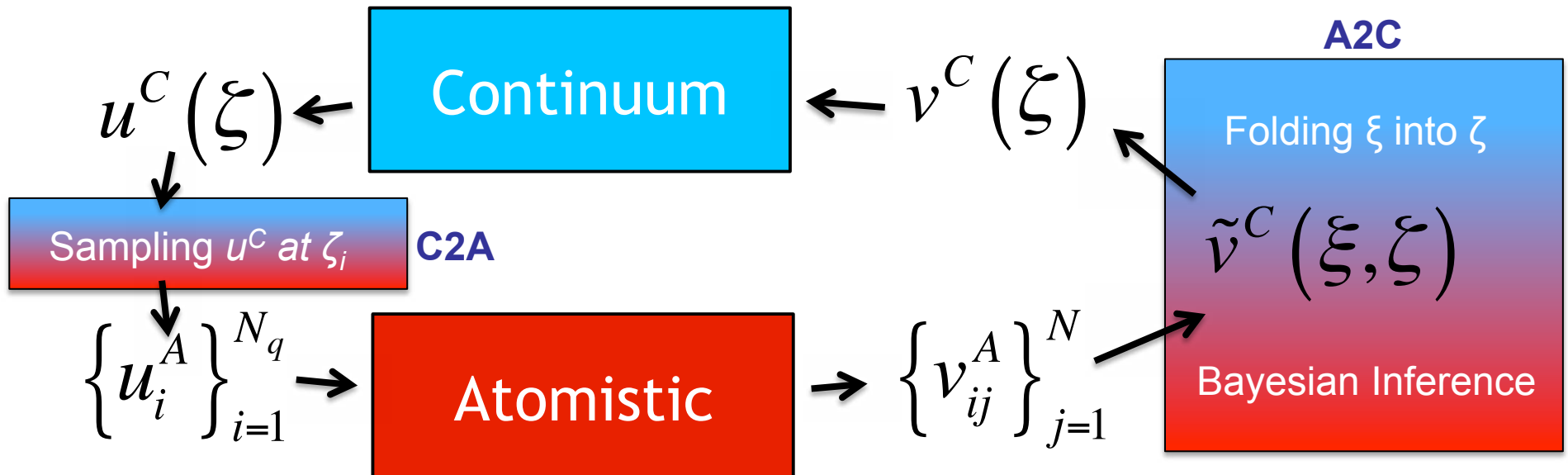
- Steady state, linear velocity profile
  - Allows analytical propagation of uncertainties

# Coupled atomistic to continuum simulation accounting for sampling noise on atomistic level

- Sampling noise is only source of uncertainty
  - All external inputs deterministic
- Sampling noise in atomistic outputs propagated through coupling
  - Uncertain continuum simulation
  - Uncertain atomistic inputs
- Different approaches
  - Fixed point iteration on atomistic level
  - Intersecting sampled continuum response surfaces
  - Fixed point iteration on intersecting uncertain continuum response surfaces



## Fixed point iteration on atomistic level



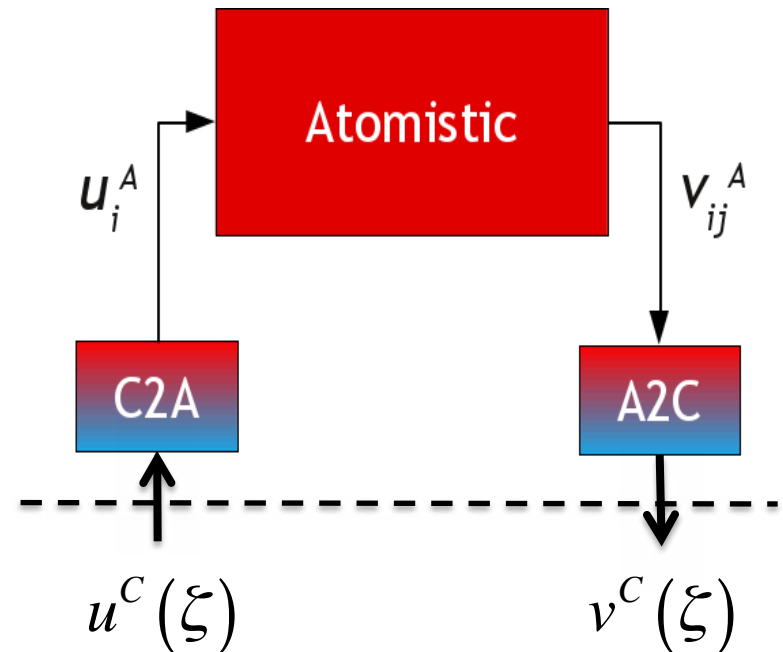
- Additional sampling noise introduced with every atomistic simulation
  - Merged with uncertainty present in atomistic input velocity
- Requires many atomistic simulations at nearby inputs
  - Expensive unless surrogate model used
- Salloum *et al.*, SIAM MMS, submitted 2011

# Response surfaces make atomistic simulation available to the macroscale

$$v^C = \psi(\xi)^T \bar{v}^C + \xi \sqrt{\psi(\xi)^T S \psi(\xi)}$$

$$\xi \sim St(0,1,\gamma)$$

$$v^C = f^A(u^C, \xi)$$

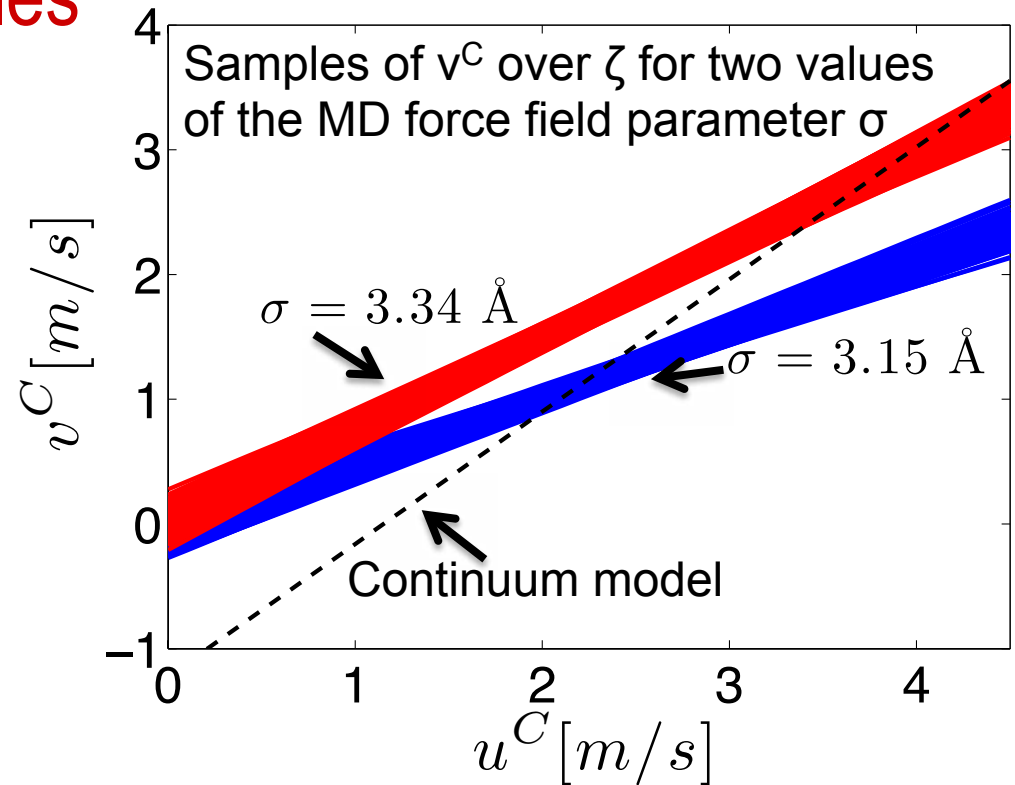


- Same formulation as propagation of parametric uncertainty
- Can be evaluated instead of running MD simulations
- Uncertainty due to sampling noise captured by  $\zeta$ 
  - Can be reduced by adding more MD data

# Intersecting sampled response surfaces readily provides coupling variables

$$v^C = f^A(u^C, \xi)$$

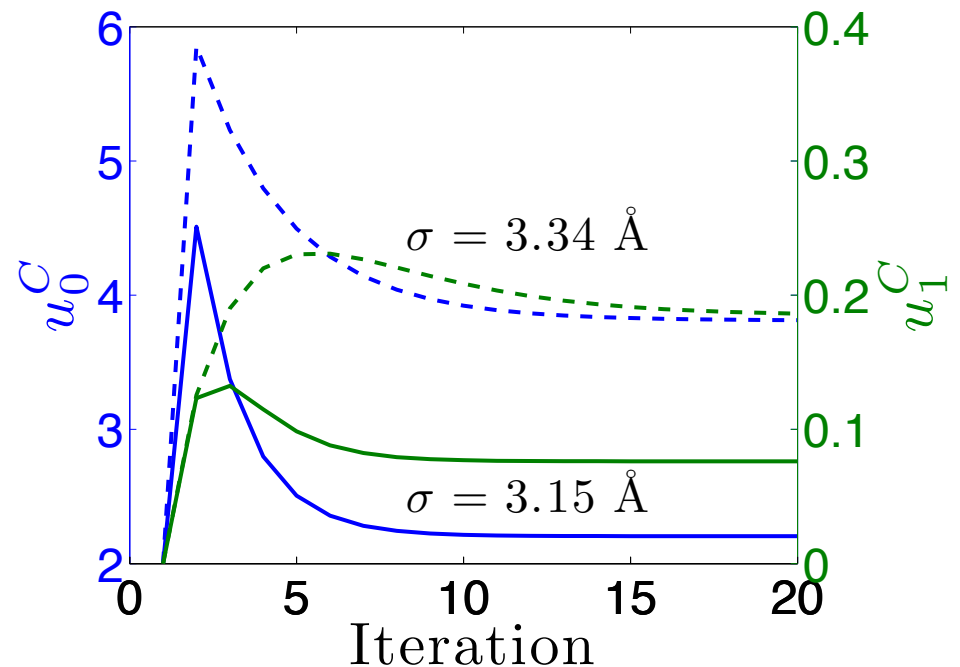
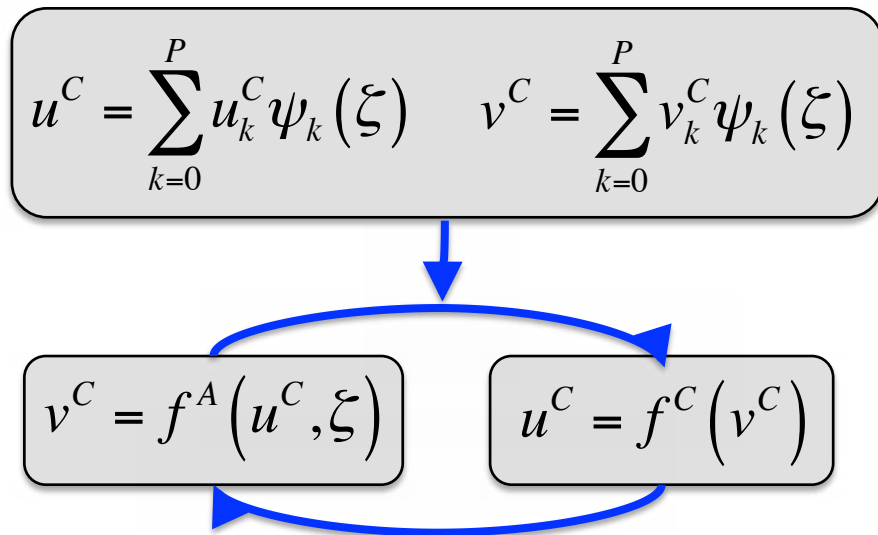
$$u^C = f^C(v^C)$$



- Sample atomistic response surface over intrinsic variability  $\zeta$ 
  - Deterministic intersection with continuum response surface
- Project resulting samples of  $v^C$  onto PC basis
  - Mapping to PC random variables using inverse CDF of sampled  $v^C$

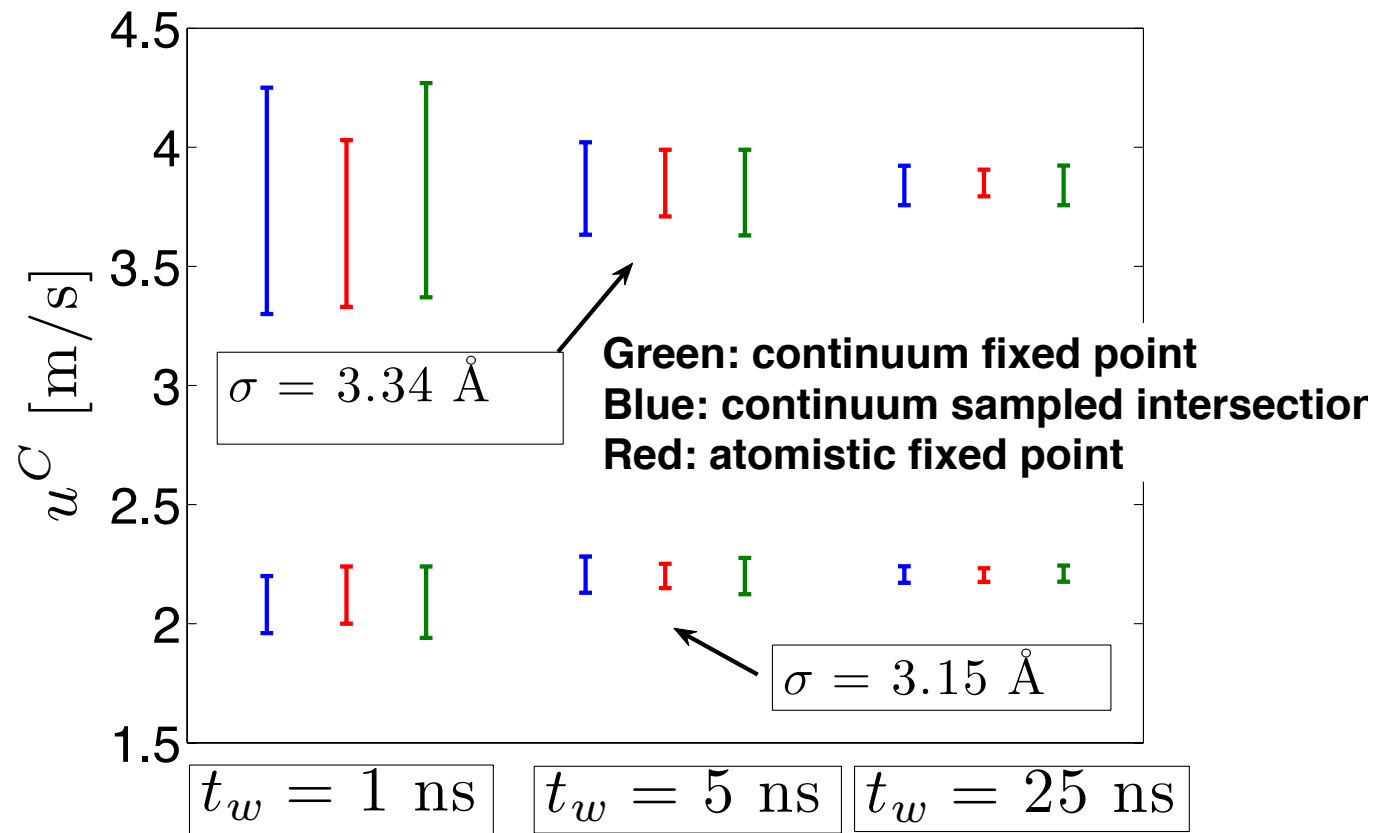


# Fixed point iteration on uncertain response surfaces



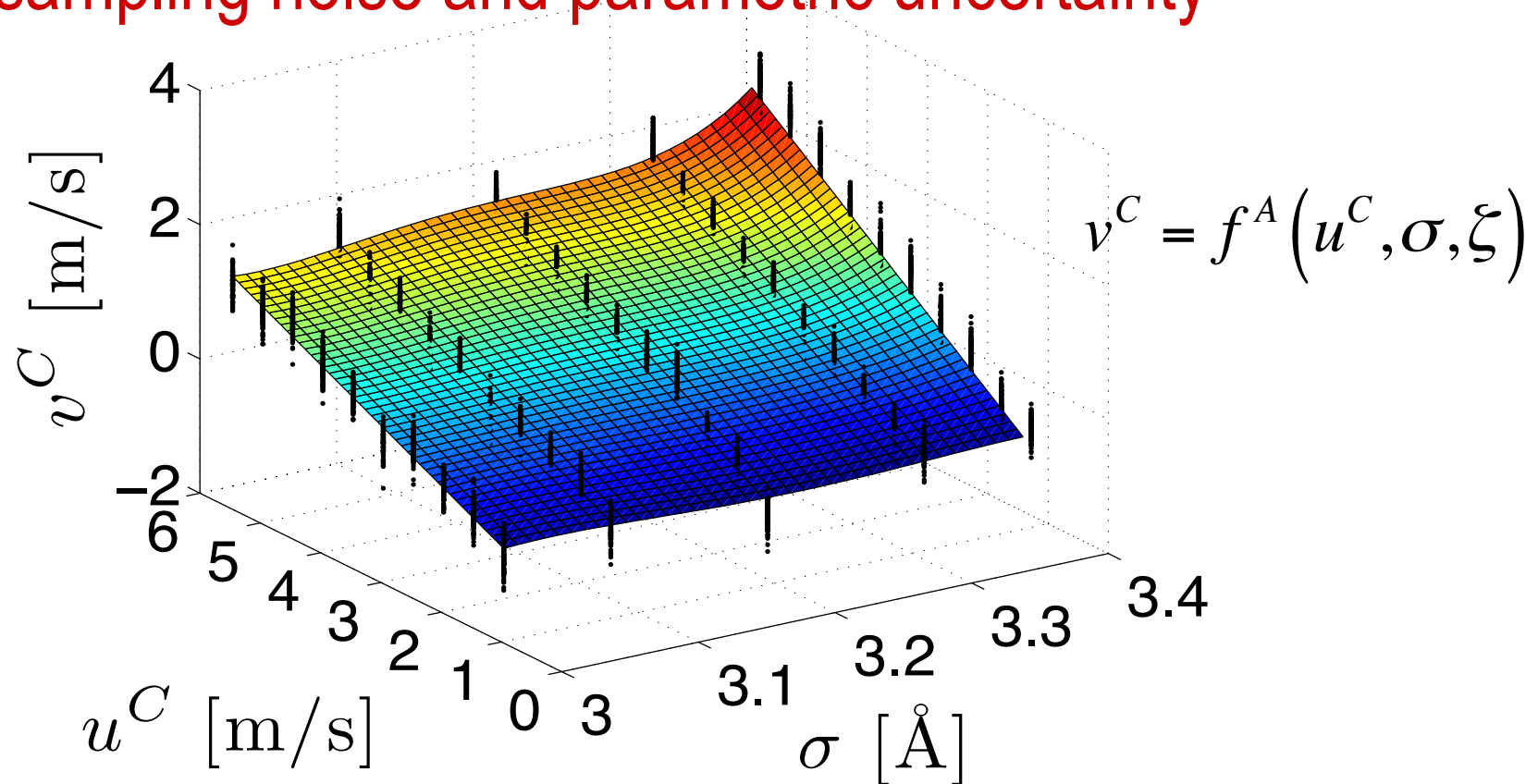
- Assume PC expansion for coupling variables to represent sampling noise
- Substitute PCEs into atomistic and continuum response surfaces
- Starting from an initial guess, iterate till convergence

# The three coupling approaches are in agreement



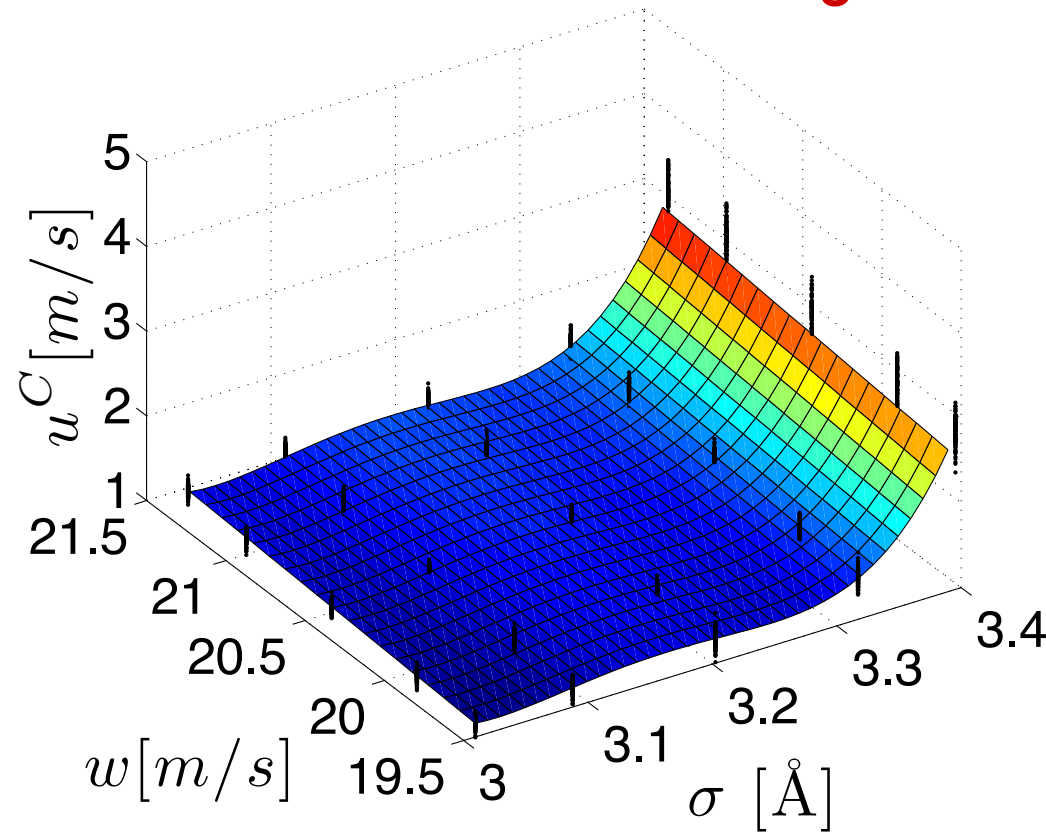
- Larger  $\sigma$  gives more uncertainty
- Noise can be reduced through longer time averaging

# Coupled atomistic to continuum simulation with sampling noise and parametric uncertainty



- Response surface  $f^A$  as function of input parameters
  - Generalization of case with sampling noise only
  - Inferred from MD data at sampled parameter values
  - Sampling noise represented as student-t process

# Response surface intersection through sampling



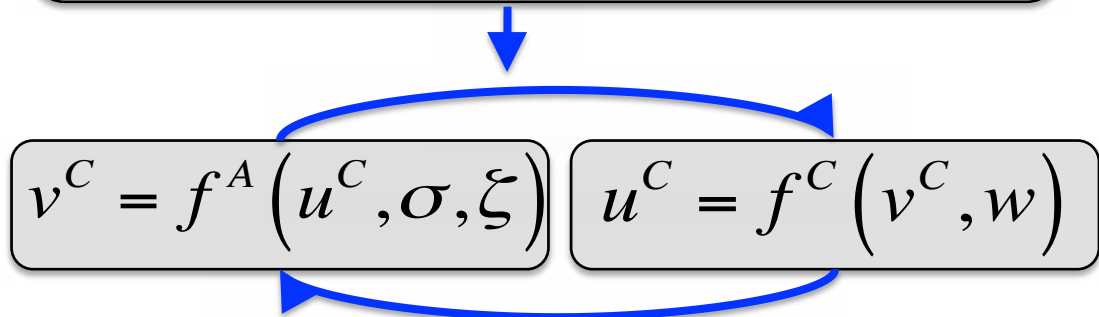
- Intersect response surfaces at specific parameter values
  - Using previously discussed approaches
- Infer polynomial surface through those points

# Intersection through fixed point iteration on uncertain continuum response surfaces

$$w = \sum_{k=0}^P w_k \psi_k(\xi)$$

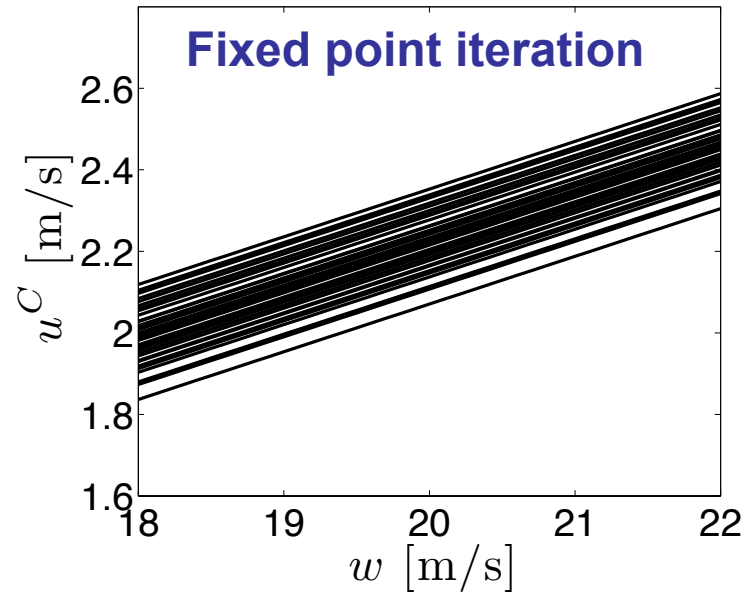
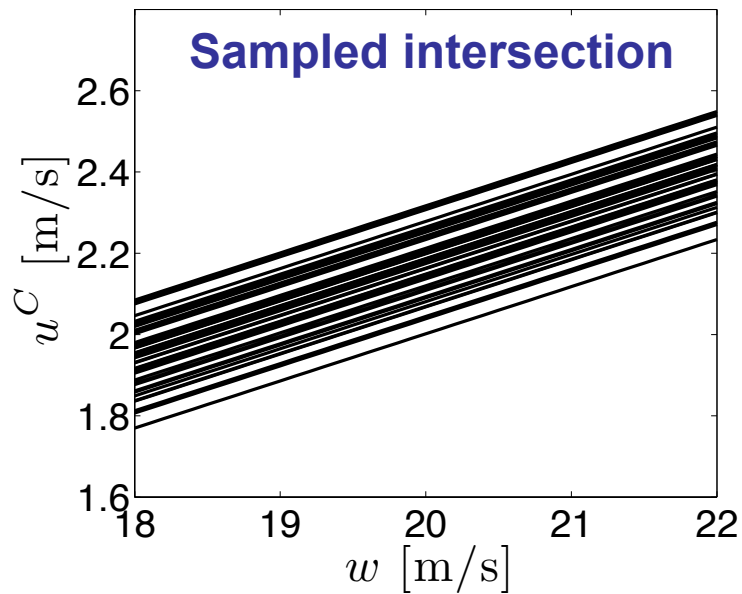
$$\sigma = \sum_{k=0}^P \sigma_k \psi_k(\xi)$$

$$u^C = \sum_{k=0}^P u_k^C \psi_k(\xi, \zeta) \quad v^C = \sum_{k=0}^P v_k^C \psi_k(\xi, \zeta)$$



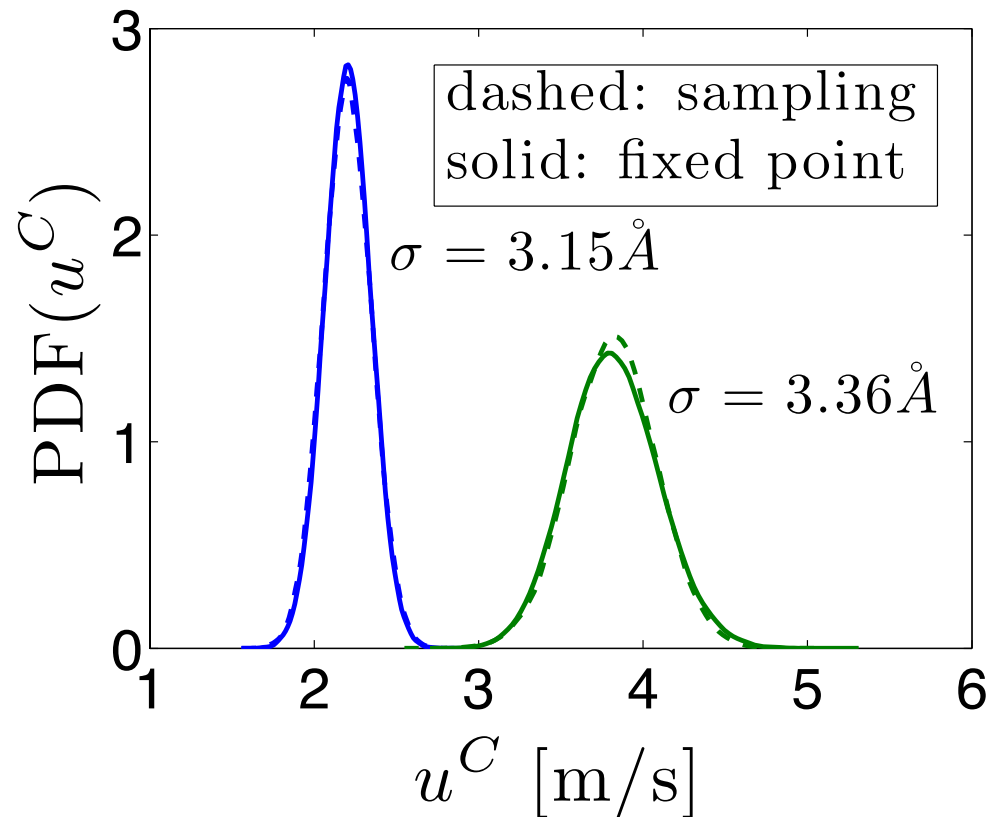
- Assume known uncertainties in  $w$  and  $\sigma$
- Substitute PC expansions into response surfaces
- Iterate on expansions for  $u^C$  and  $v^C$

# The sampled intersection and fixed point iteration approaches agree well



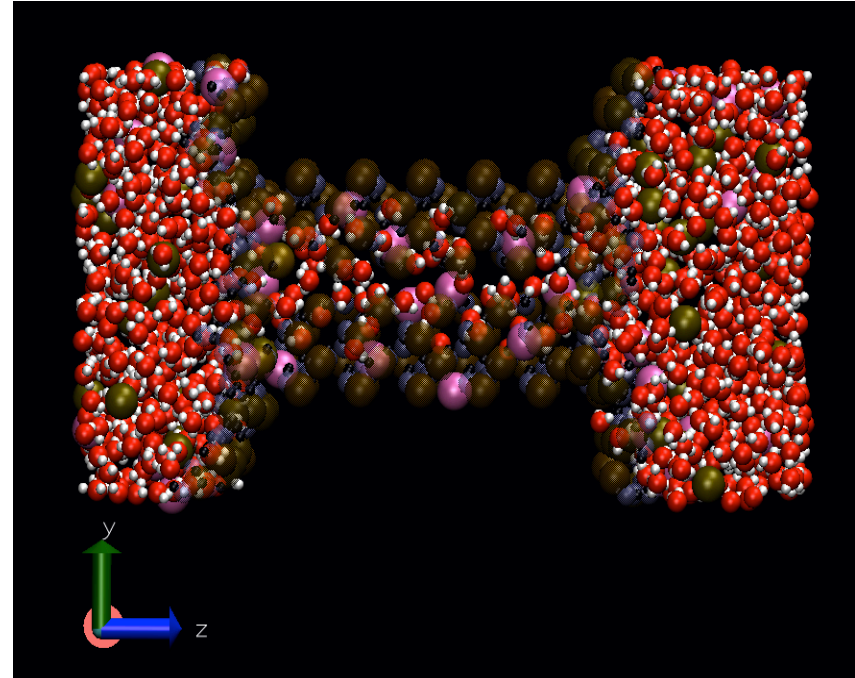
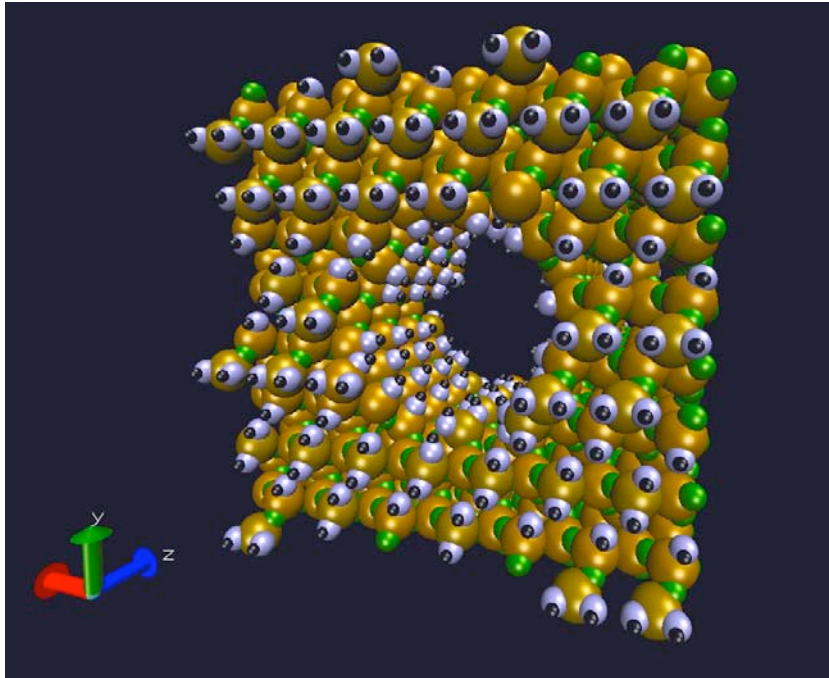
- Implemented for linear dependence on  $w$  at fixed  $\sigma$ 
  - Good agreement with sampled intersection approach
- General non-linear case being implemented

## Forward propagation of uncertainty in driving velocity



- Assume uncertain driving velocity  $w$ 
  - Gaussian with mean 20 m/s and standard deviation 1 m/s
- Uncertainty from  $w$  and MD sampling noise propagated into coupling variables

# Application to more challenging multiscale problems



- Ionic fluxes (NaCl) through Silica nanopores
- MD concentration boundary conditions set by continuum
- Continuum flux boundary conditions set by MD
- Work in progress



# Conclusions

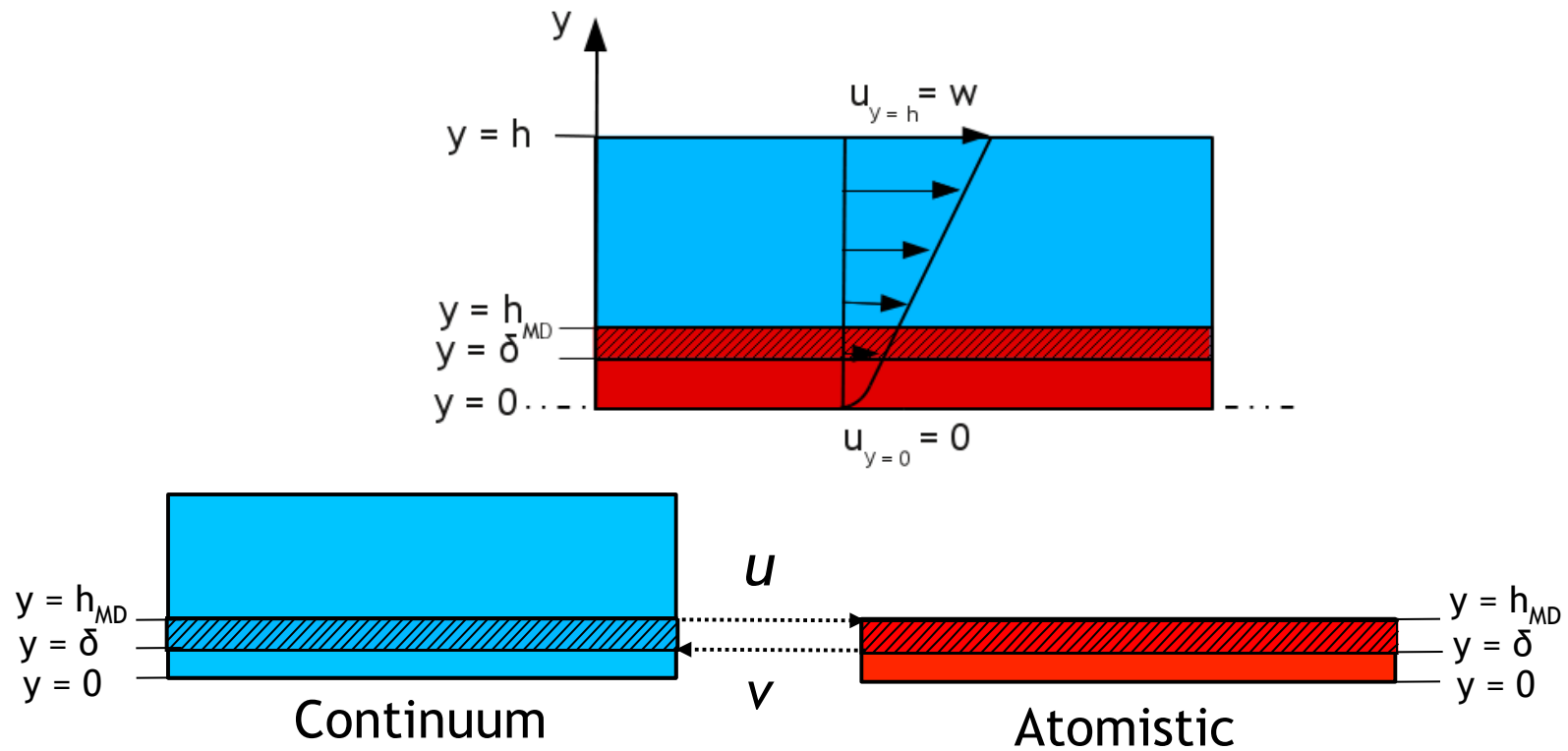
- Bayesian methods are used to quantify sampling noise in macroscale observables extracted from atomistic simulations
- Stochastic multiscale coupling approach accounts for sampling noise and parametric uncertainty
- Response surfaces for atomistic simulations allow coupling on the macroscale level
- Simple model problem here often allows for analytical solutions, but formulation is generally applicable
  - Application to nanopore ionic fluxes in progress
- More details
  - Salloum et al., SIAM MMS, submitted
  - Rizzi et al., J. Comp. Phys, submitted (Part I and II)

# Papers

- Rizzi, F., Najm, H.N., Debusschere, B.J., Sargsyan, K., Salloum, M., Adalsteinsson, H., Knio, O.M., “Uncertainty Quantification in MD Simulations. Part I: Forward propagation”, *J. Comp. Phys.*, submitted, 2011
- Salloum, M., Sargsyan, K., Najm, H.N., Debusschere, B., Jones, R., Adalsteinsson, H. “A Stochastic Multiscale Coupling Scheme to account for Sampling Noise in Atomistic-to-Continuum Simulations” *SIAM Multiscale Modeling and Simulation*, submitted, 2011
- Rizzi, F., Najm, H.N., Debusschere, B.J., Sargsyan, K., Salloum, M., Adalsteinsson, H., Knio, O.M., “Uncertainty Quantification in MD Simulations. Part II: Inference of Force Field Parameters”, *J. Comp. Phys.*, submitted, 2011

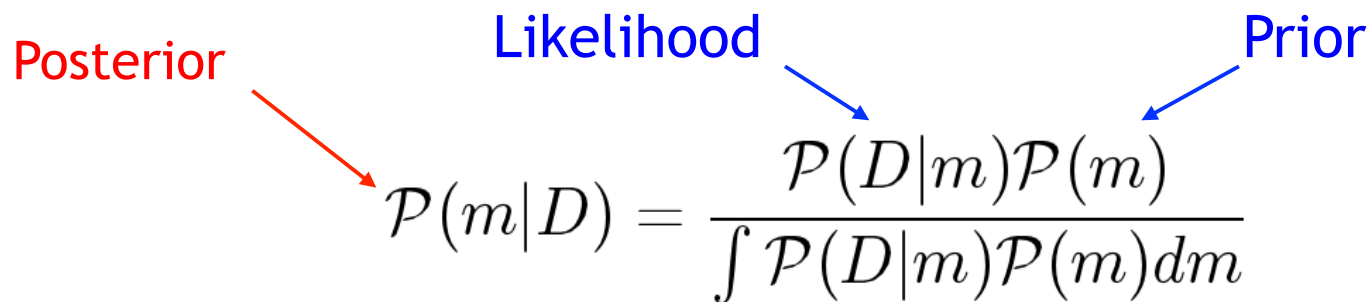
Extra material

# Variables are exchanged across scale interfaces



# Bayesian Inference

Let  $m$  be a hypothesis and  $D$  observed data.



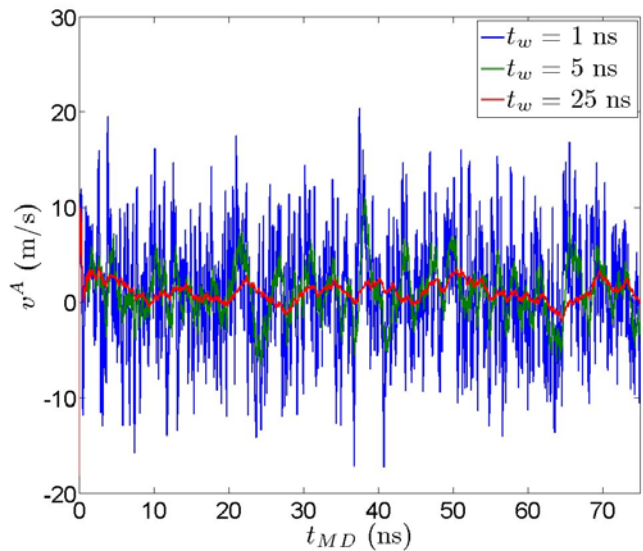
The diagram shows the Bayesian Inference formula with three labels and arrows pointing to specific parts of the equation:

- Posterior** (red text) with a red arrow pointing to  $\mathcal{P}(m|D)$ .
- Likelihood** (blue text) with a blue arrow pointing to  $\mathcal{P}(D|m)$  in the numerator.
- Prior** (blue text) with a blue arrow pointing to  $\mathcal{P}(m)$  in the numerator.

$$\mathcal{P}(m|D) = \frac{\mathcal{P}(D|m)\mathcal{P}(m)}{\int \mathcal{P}(D|m)\mathcal{P}(m)dm}$$

- The prior expresses the initial knowledge about the hypothesis  $m$  (e.g. uniform distribution, expert's knowledge...)
- The likelihood is the probability of observing the data  $D$  given the hypothesis  $m$ . It encompasses the forward model of  $m$ .
- The denominator is a normalization constant.
- The posterior is the probability of the hypothesis  $m$  given the data  $D$  : **offers an enhanced knowledge of  $m$ .**

# Quantification of sampling noise in atomistic simulations



$$v_j^A = v^C + s\eta_j \quad d = \{v_j^A\} \quad j = 1, \dots, N$$

$$P(v^C, s^2 | d) \propto P(d | v^C, s^2) P(v^C, s^2)$$

$$\mathcal{P}(d | v^C, s^2) = (2\pi)^{-N/2} (s^2)^{-N/2} \exp\left(-\frac{\boldsymbol{\epsilon}^T \boldsymbol{\epsilon}}{2s^2}\right)$$

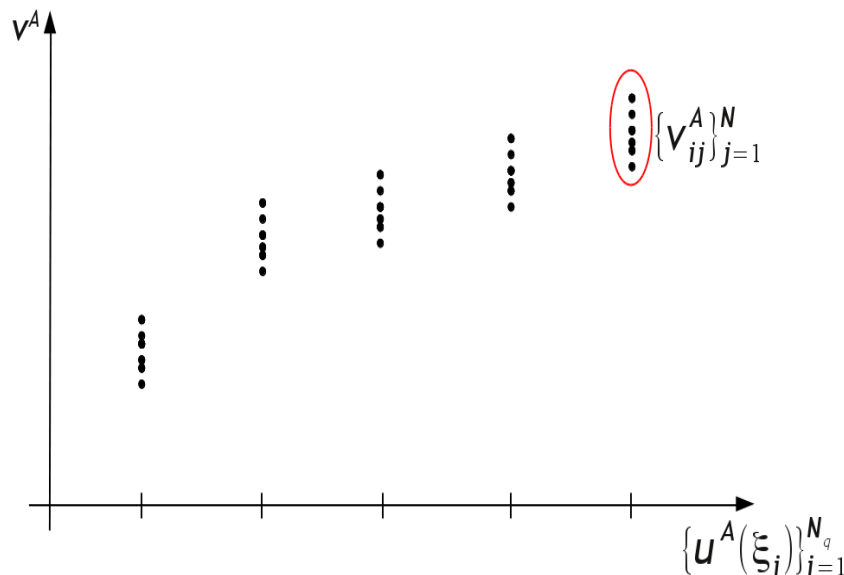
$$\epsilon_j = s\eta_j = v_j^A - v^C$$

$$v^C = \mu + \Sigma \zeta$$

$$\zeta \sim \mathcal{S}(0, 1, \gamma)$$

- Infer  $v^C$  from  $N$  short-term averaged MD velocity samples  $v_j^A$ 
  - Gaussian model for data noise due to Central Limit Theorem (CLT)
  - Analytical solution gives  $v^C$  as student-t random variable
    - $P(v^C)$  marginalized over  $s$  is student-t distributed

# Propagating parametric uncertainty and sampling noise through atomistic simulations



$$v_{ij}^A = v^C(\xi) + s\eta_{ij} \quad v^C = \sum_{k=0}^P v_k^C \psi_k(\xi)$$

$$\begin{aligned} \tilde{v}^C &= \sum_{k=0}^P \tilde{v}_k^C \Psi_k(\xi) \\ &= \Psi(\xi)^T (\lambda + \Lambda \zeta) \end{aligned}$$

$$\zeta \sim \mathcal{S}(0, 1, \gamma)$$

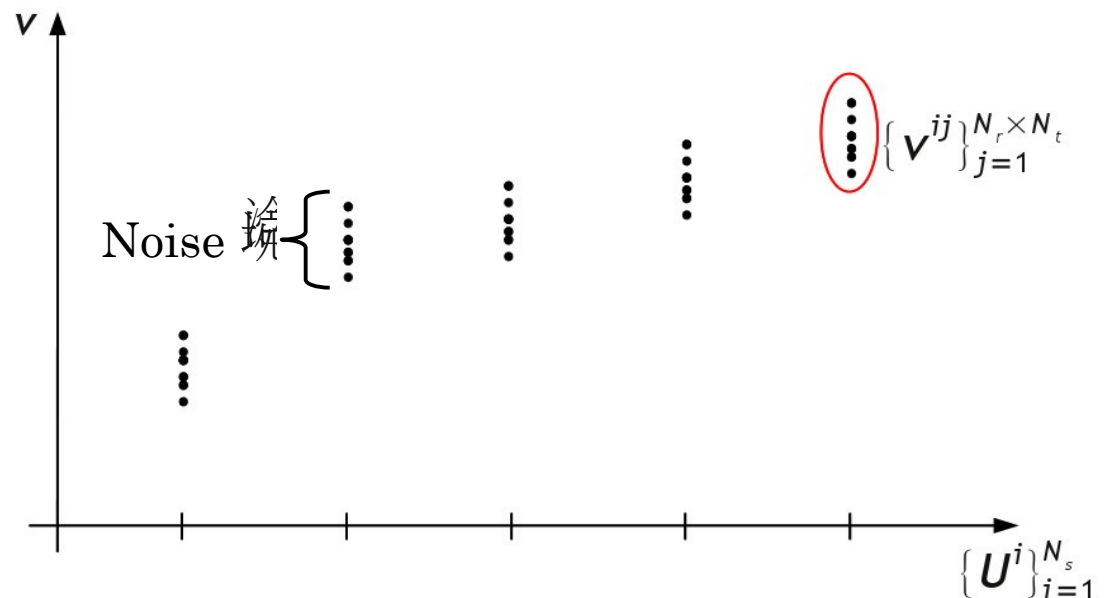
- Infer  $v^C$  from  $N$  short-term averaged MD velocity samples  $v_{ij}^A$ 
  - Sampled over a range of input velocities  $u_i^A$
  - Gaussian model for data noise due to Central Limit Theorem (CLT)
  - Analytical solution gives  $v^C$  as student-t process over input velocity uncertainty

# Inferring the Output Variable

$$\mathcal{P}(\tilde{\mathbf{u}}, \sigma^2 | D_v) \propto \mathcal{P}(D_v | \tilde{\mathbf{u}}, \sigma^2) \mathcal{P}(\tilde{\mathbf{u}}, \sigma^2)$$

$$\mathbf{v} = \sum_{k=0}^P \tilde{u}_k \Psi_k(\xi) + \underbrace{\sigma \hat{\xi}}_{\text{Noise term}}$$

We draw samples from the posterior using Markov Chain Monte Carlo (MCMC) sampling.



$\xi$  relates to the spread in  $U^i$



## Folding the input uncertainty and the sampling noise into one uncertain output

After marginalizing over  $\sigma^2$ , we obtain a joint posterior on the  $\{\tilde{\mathbf{u}}_k\}_{k=0}^P$ :

$$\tilde{u} = \sum_{k=0}^P \tilde{u}_k \Psi_k(\xi)$$

We approximate  $\{\tilde{\mathbf{u}}_k\}_{k=0}^P$  as a Multivariate Normal Distribution (MVN) as follows:

$$\{\tilde{u}_k\}_{k=0}^P \sim \mathcal{MVN}(\boldsymbol{\mu}, \Sigma) = \boldsymbol{\mu} + L\boldsymbol{\zeta} \quad \text{where} \quad L^T L = \Sigma$$

We obtain:

$$\tilde{u} = \boldsymbol{\Psi}(\xi)^T \cdot \boldsymbol{\mu} + \zeta \sqrt{\boldsymbol{\Psi}(\xi)^T \cdot \Sigma \cdot \boldsymbol{\Psi}(\xi)}$$

## Folding the input uncertainty and the sampling noise into one uncertain output

$$\tilde{u} = \mathbf{\Psi}(\xi)^T \cdot \boldsymbol{\mu} + \zeta \sqrt{\mathbf{\Psi}(\xi)^T \cdot \Sigma \cdot \mathbf{\Psi}(\xi)}$$

This expression of  $\tilde{u}$  is “cheap” for sampling in  $\zeta$  and  $\xi$  !

Inverse Cumulative  
Distribution  
Function (CDF)  
transform

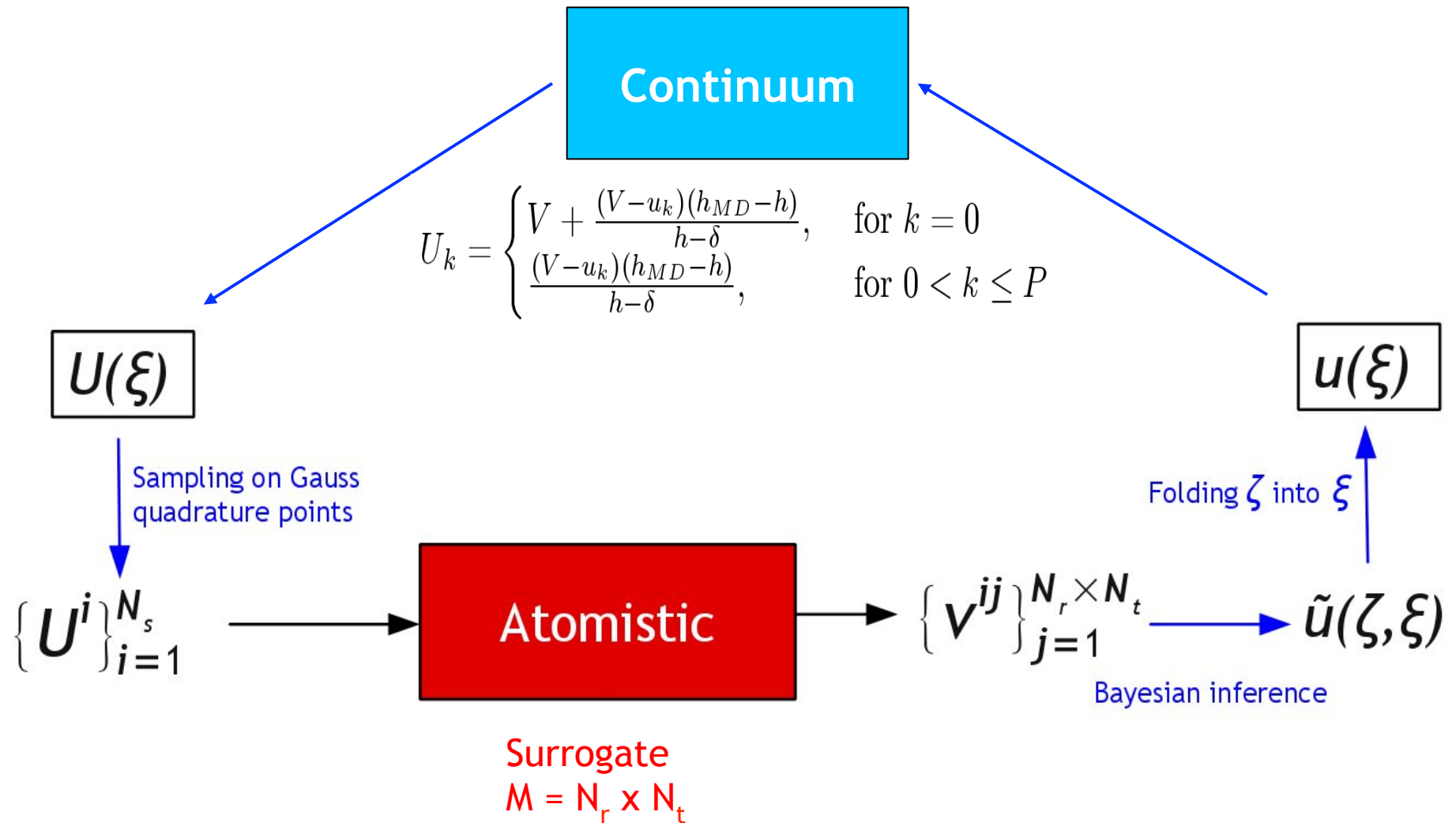
$$\left\{ \begin{array}{l} F(.) \text{ is the CDF of } \tilde{u} \\ u_k = \frac{\langle F^{-1}(\Phi(\xi)) \Psi_k(\xi) \rangle}{\langle \Psi_k(\xi)^2 \rangle} \end{array} \right.$$

$$u = \sum_{k=0}^P u_k \Psi_k(\xi)$$

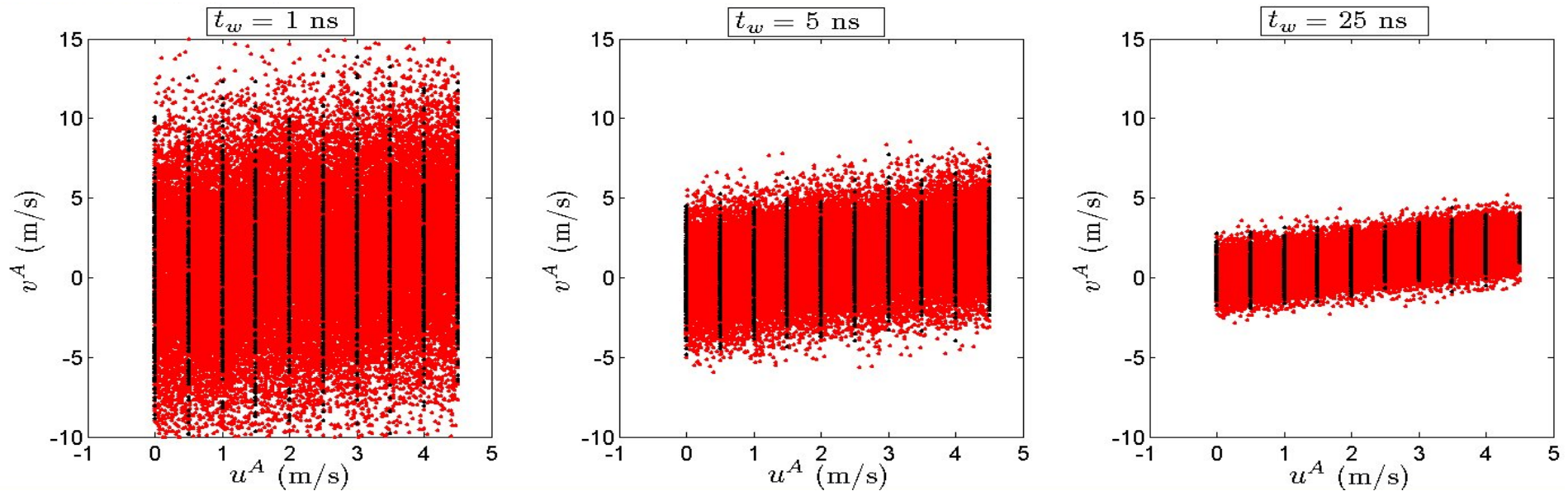
$\xi$  is the degree of  
freedom associated  
with the sampling noise

# Summary of the Different Steps for Coupling

Laminar Newtonian Couette flow  
*The analytical solution is available*



# The surrogate reflects the original short-time averaged data

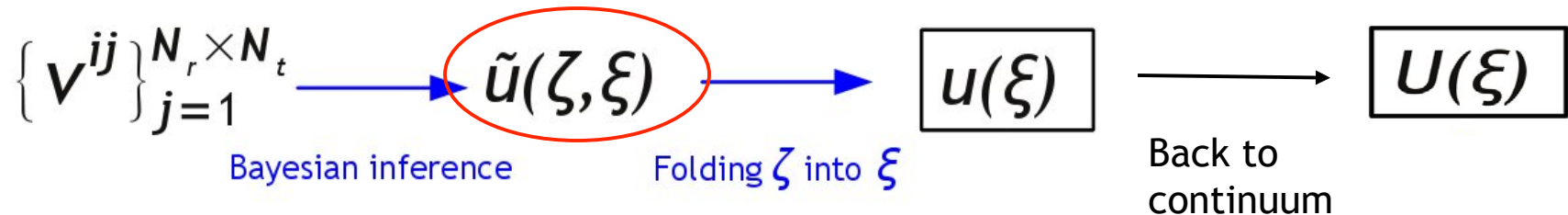
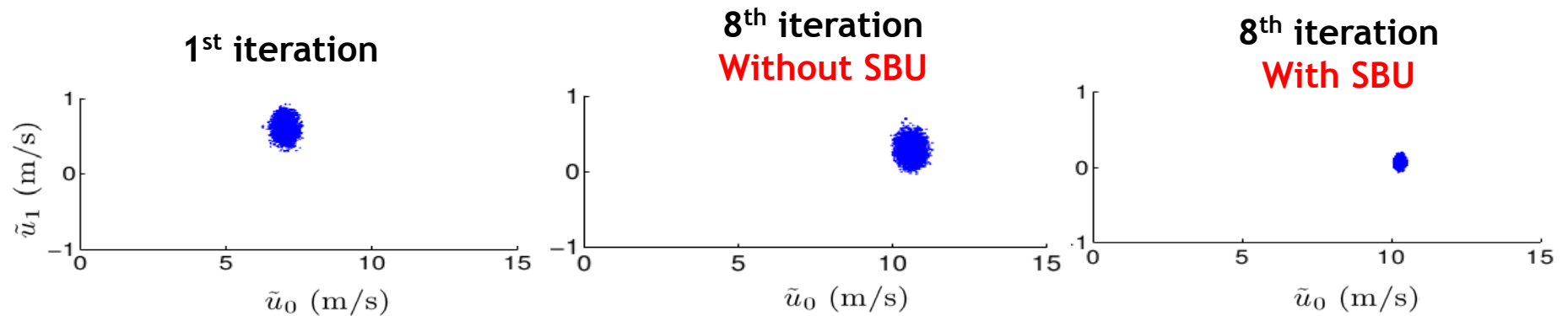


$t_w$  = Time averaging window width

$$v^A = \alpha_0 + \alpha_1 u^A + s\eta \longrightarrow \text{Spread in the short-time averaged MD data}$$

- For a given  $t_w$ , the properties of the joint posterior on  $\{\alpha_0, \alpha_1, s\}$  are tabulated.
- This joint posterior is a surrogate to the original **atomistic scale short-time** averaged MD data.

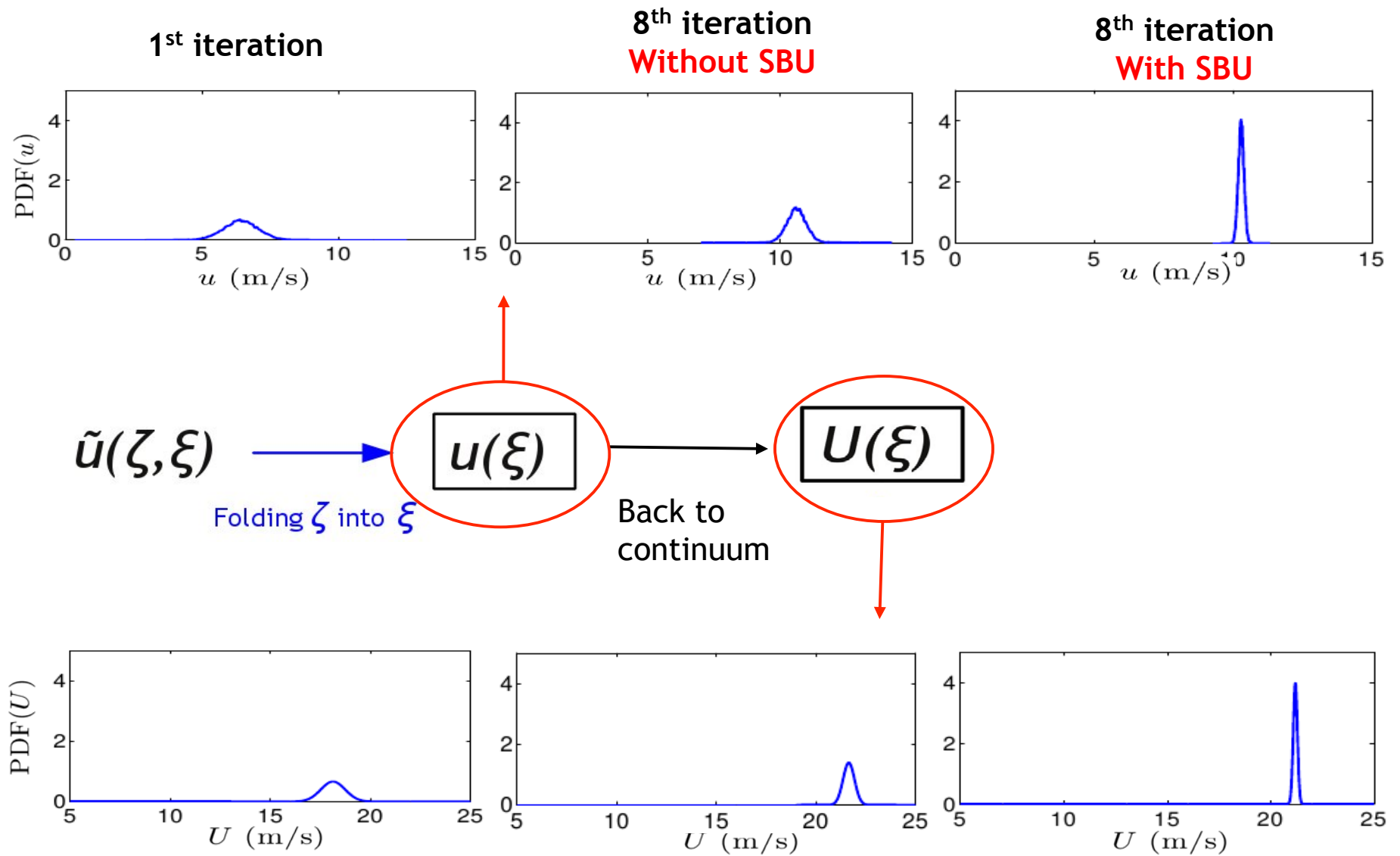
# Joint Posterior of $\{\tilde{u}_0, \tilde{u}_1\}$



## Sequential Bayesian Updating (SBU)

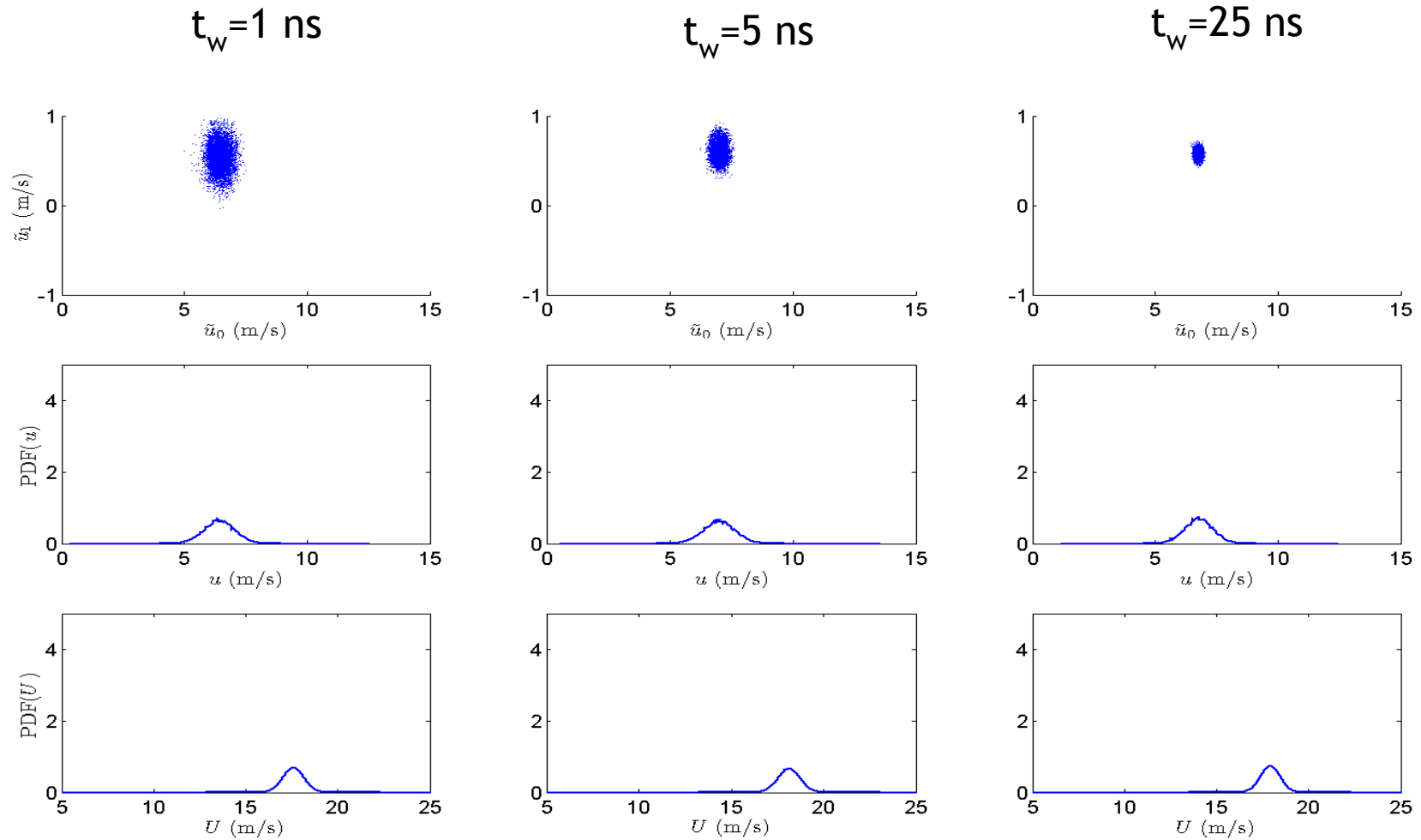
The posterior of the previous iteration is used as the prior in the current iteration.

# PDFs of $u$ and $U$



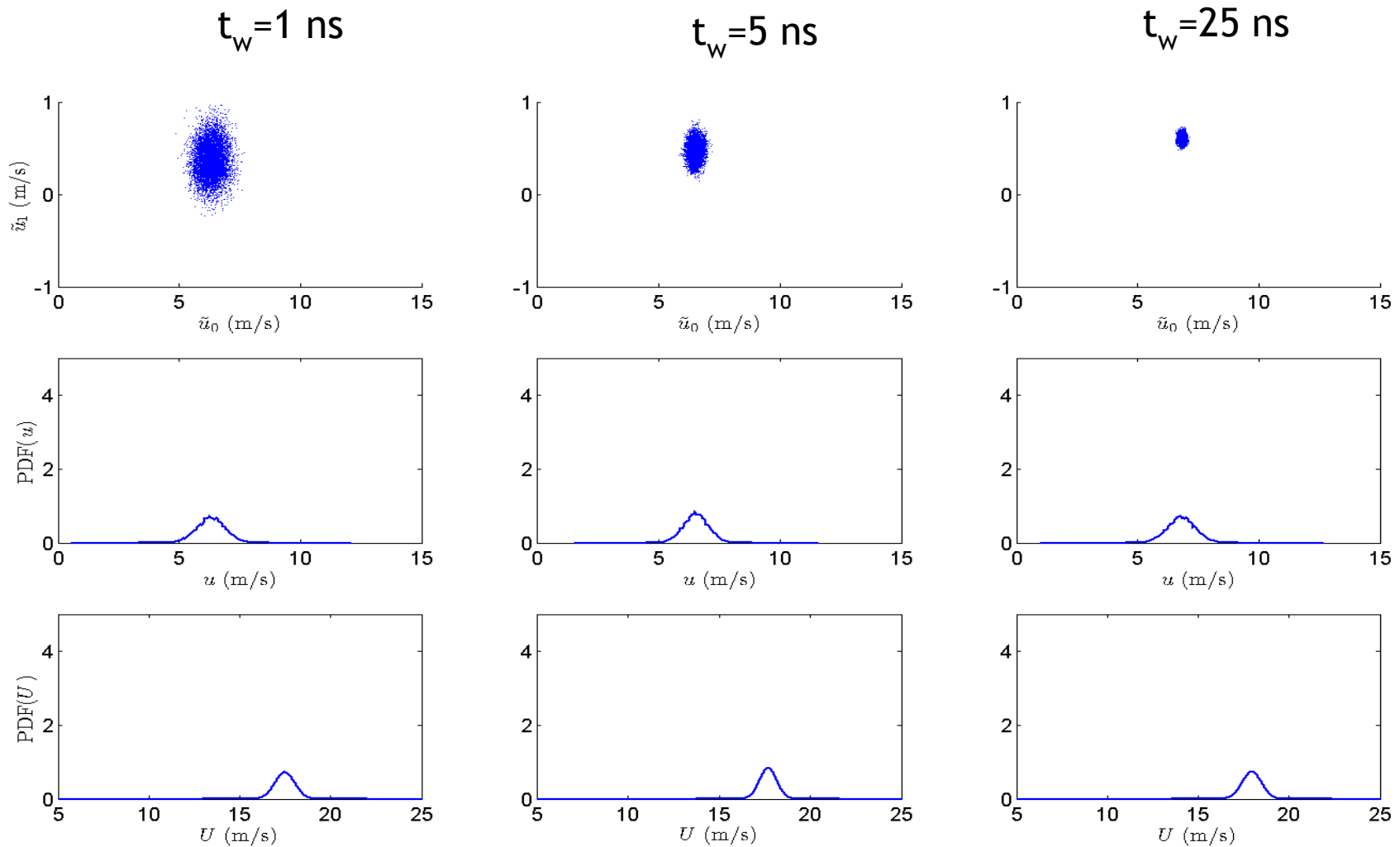
# Stochastic Coupling Algorithm Convergence *Without* SBU

$t_w$  = Time averaging window width



# Stochastic Coupling Algorithm Convergence *With* SBU

$t_w$  = Time averaging window width



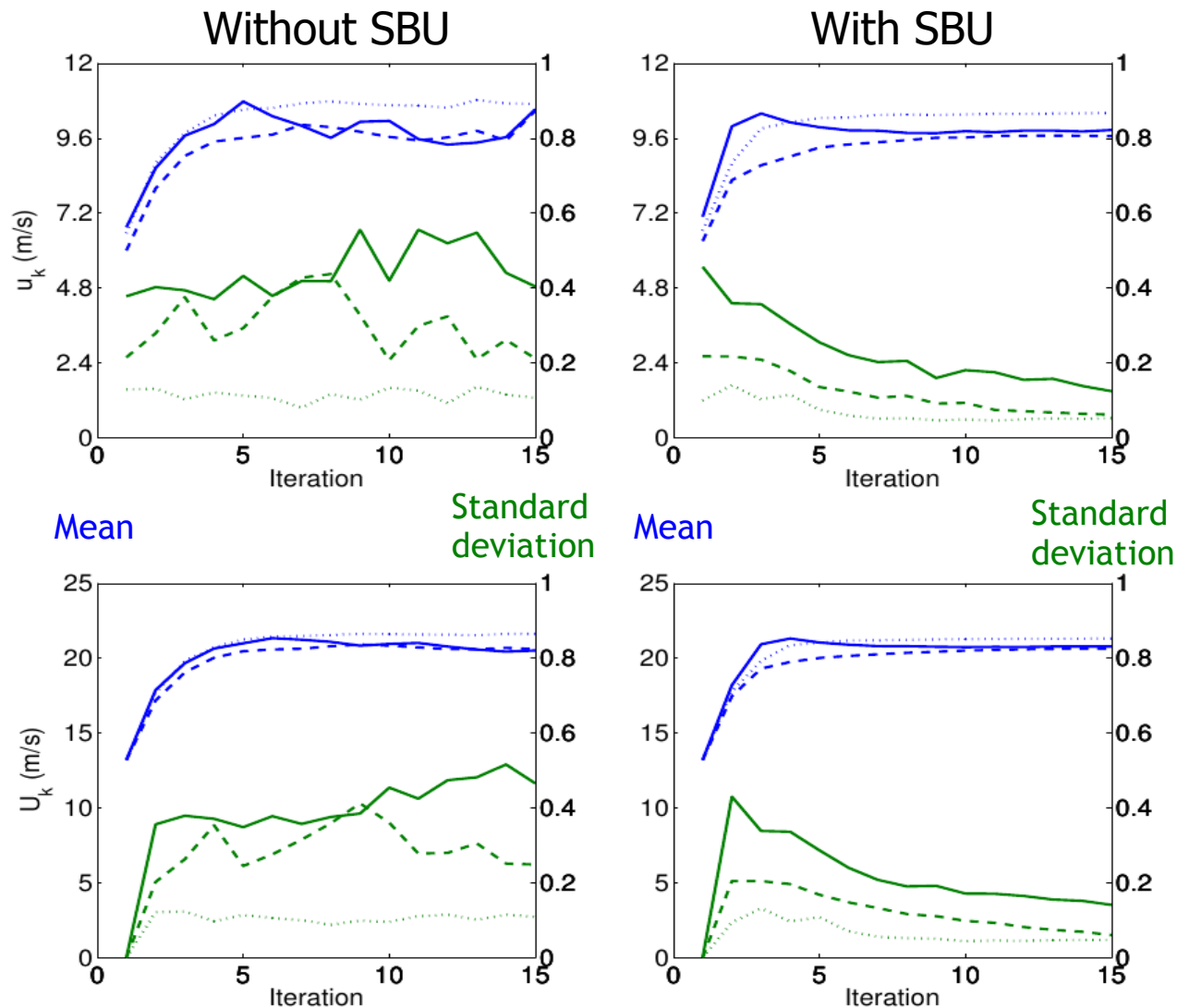


# Effect of the time averaging window $t_w$ on the convergence of the mean and standard deviation

Solid:  $t_w=1$  ns

Dashed:  $t_w=5$  ns

Dotted:  $t_w=25$  ns

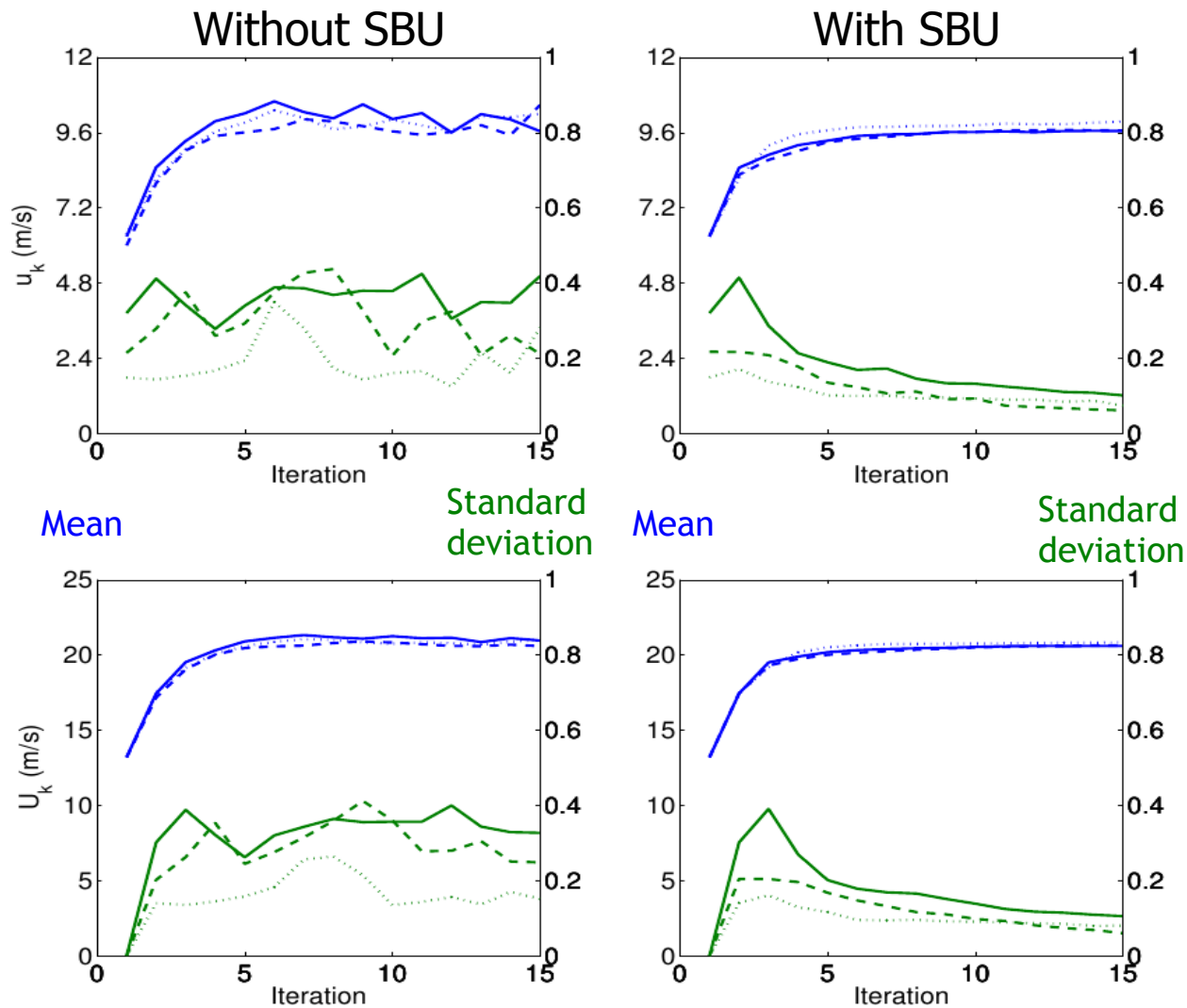


# Effect of the number of samples $M$ on the convergence of the mean and standard deviation

Solid:  $M=10$

Dashed:  $M=20$

Dotted:  $M=40$



# Effect of the time averaging window $t_w$ on the convergence of the mean and standard deviation

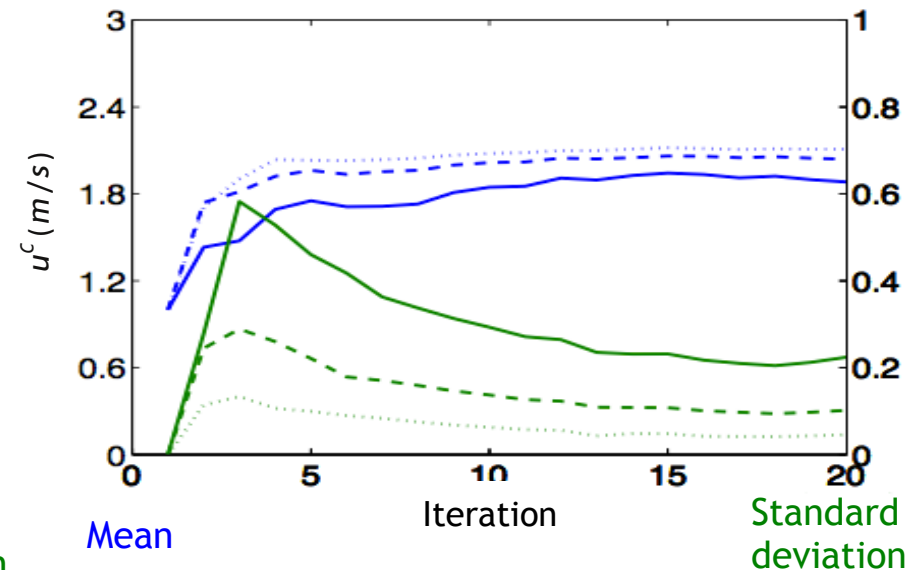
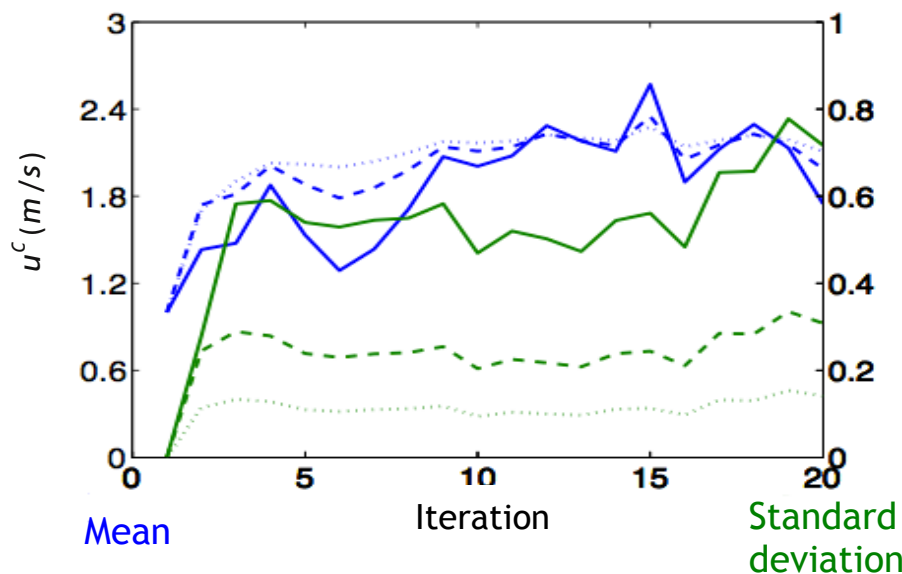
Solid:  $t_w=1$  ns

Dashed:  $t_w=5$  ns

Dotted:  $t_w=25$  ns

Without SBU

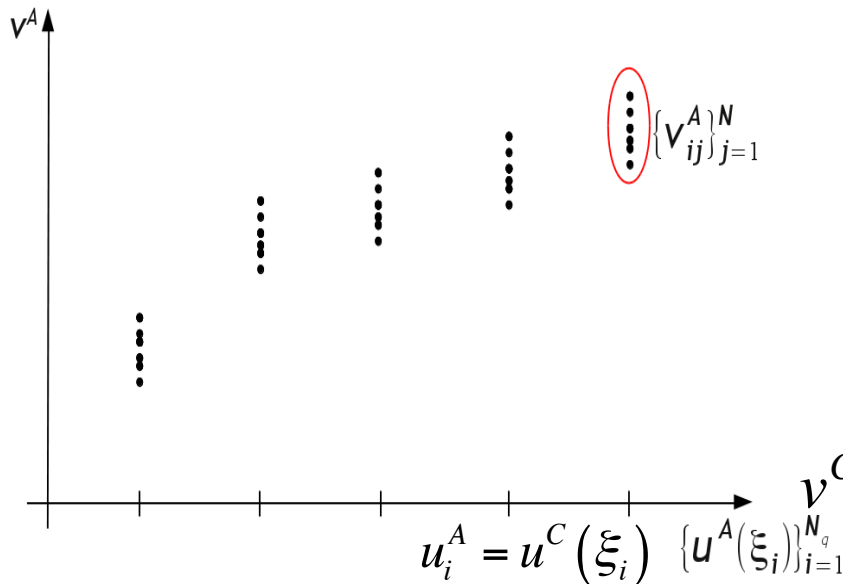
With SBU



## Sequential Bayesian Updating (SBU)

- The posterior of the previous iteration is used as the prior in the current iteration.
- The accuracy in the converged variables by including additional data at each iteration.

# Representation of response surface with student-t process



$$v^C = f^A(u^C, \xi)$$

$$u^C = \sum_{k=0}^P u_k^C \psi_k(\xi) \quad v^C = \sum_{k=0}^P v_k^C \psi_k(\xi)$$

$$v^C = \psi(\xi)^T (\lambda + \Lambda \xi)$$

$$v^C = \psi(\xi)^T \bar{\mathbf{v}}^C + \xi \sqrt{\psi(\xi)^T \Lambda^T \Lambda \psi(\xi)}$$

$$\xi \sim S(0, 1, \gamma)$$

- Sum of  $P+1$  student-t RVs mapped into 1 student-t RV
- Well approximated with Gaussian Process if  $\gamma$  large enough
  - Satisfied if enough MD samples used

# Comparison student-t and Gaussian