

LA-UR- 11-04195

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Title: Representing Shear Localization under Dynamic Loading Conditions

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Intended for: THERMEC 2011
Quebec City, Canada



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Representing Shear Localization under Dynamic Loading Conditions

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The localization and adiabatic shear banding process is studied for both annealed tantalum and 316L stainless steel in the annealed, rolled, and pre-shocked conditions. Tantalum does not form adiabatic shear bands and the dynamic response of this material is examined using both continuum and crystal plasticity computational tools to examine the large deformation response of this material and the corresponding microstructural evolution with deformation. It is found that models which are based upon thermally activated physics of dislocation motion represent well the response of this material under the conditions examined. Comparisons between numerical simulations and experimental data (i.e. flow stress, crystallographic texture, metallographic strain measurements) made on the deformed samples compare reasonably well to the experimental results. Results also suggest that complex and statistically varying evolution of stress state is predicted. Adiabatic shear banding does occur in the three different initial state 316L stainless steel materials. These experiments are considered in the context of both traditional high resolution numerical representation and also a sub-grid technique which is implemented within an explicit Lagrangian framework. This sub-grid method is based on the assumed-strain approach of Fish and Belytschko (1988), which allows for localization bands to be embedded within an element, thereby alleviating mesh sensitivity and reducing the required computational effort for representation of this very small scale process zone. Computational results are presented and compared to both the tantalum and 316L stainless steel experimental results (LA-UR-10-07418).

Presentation: Oral

Representing Shear Localization under Dynamic Loading Conditions

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THERMEC 2011
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2 August, 2011



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Outline

- Motivation and challenge
- Macro-scale experiments
- Continuum isotropic constitutive model
 - Anisotropic MTS
- Thermo-viscoplastic single crystal model
 - Thermally activated slip
- Embedded polycrystal forced shear simulations
- Comparison to experimental results
- Summary & comments



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Explosively Loaded Sample Demonstrates Ductile Damage and Failure Physics

T. Mason



Explosively Loaded Tantalum Experiment
6 mm thick PETN Beneath Sample – Center Detonated
Soft Sample Recovery



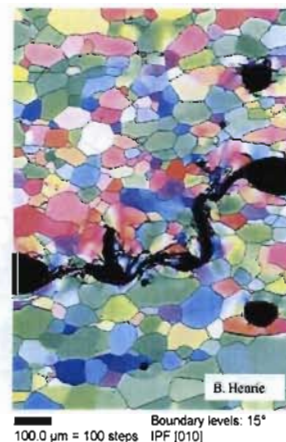
'Void sheets' are forming along grain boundaries in sample

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Motivation

- Large-deformation high-rate ductile failure involves
 - Strain localization
 - Void nucleation
 - Void growth
 - Void coalescence
- Predictive modeling of failure requires a good representation of each of these mechanisms
- Modeling strategy
 - Sub-grid computational technique
 - Constitutive / micro-mechanical models



Ta plate impact

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Strain Localization in Fragmentation Problems



Cross-sectional metallography points to considerable amount of localized plastic work before failure



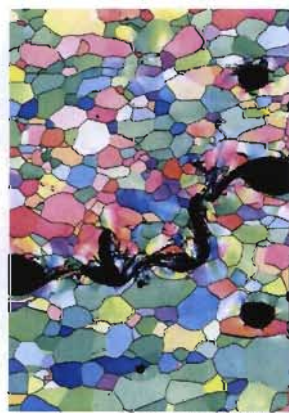
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Challenge of linking microstructure to performance

Ta Plate Impact



100.0 μm = 100 steps

Boundary levels: 15°
IPF [010]



B. Henrie

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Cu Plane Strain Compression

Undeformed

$\epsilon = -1.5$

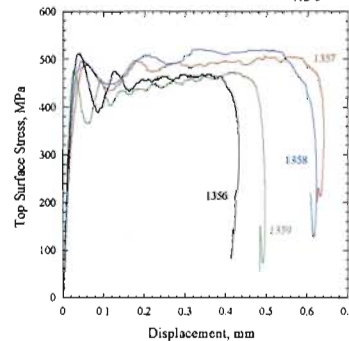
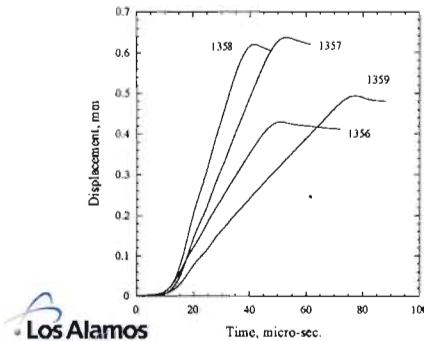
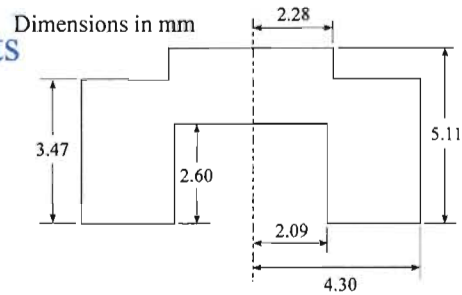


- The ductile failure process generally involves localization, porosity initiation, porosity growth, and coalescence dominated by localized deformation.
- These events occur at the length scale of the single crystal.



Forced shear experiments

- Axisymmetric, Split-Hopkinson
- Top-to-bottom loading, 10-25 m/s
- Tantalum plate, residual texture (Chen & Gray, 1996; Maudlin et al., 1999), grain size $42\text{ }\mu\text{m}$
- $-100\text{ }^{\circ}\text{C}$ and $25\text{ }^{\circ}\text{C}$

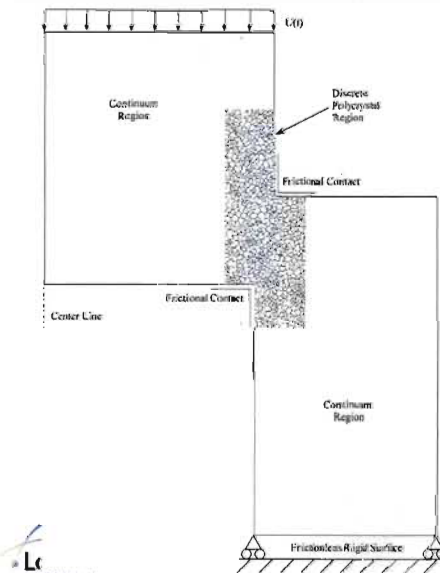


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Numerical model with boundary conditions

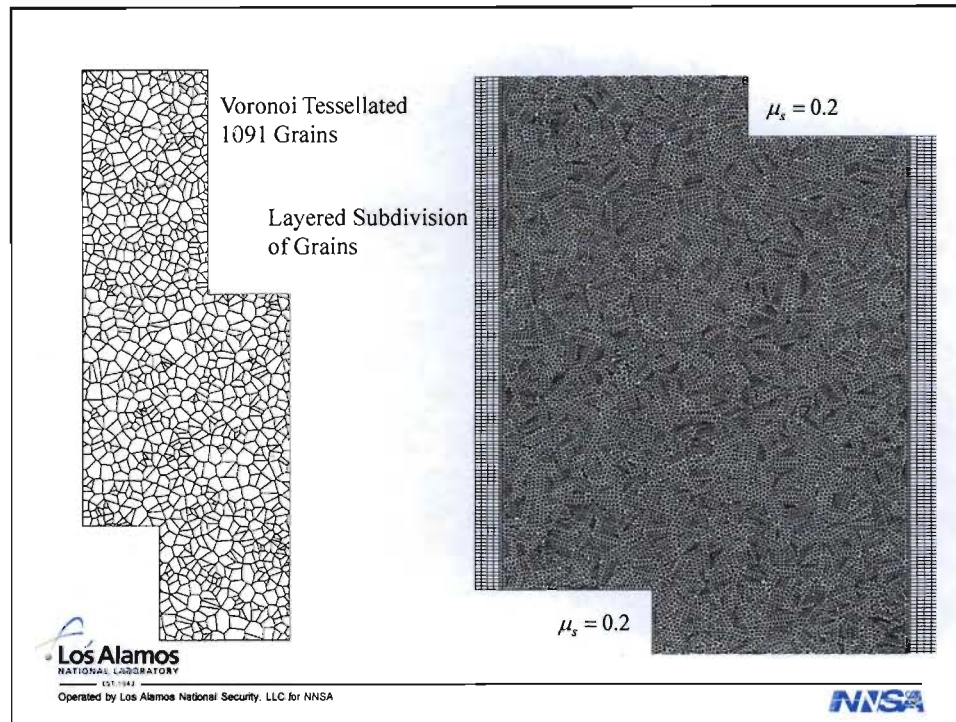


- Axisymmetric, 3 node linear elements for shear zone
- Adiabatic
- Embedded 1091 grain polycrystal region
- $40\text{ }\mu\text{m}$ grain size
- $\sim 7\text{ }\mu\text{m}$ element size
- Rate-dependent isotropic continuum regions
- Frictionless contact surfaces at corners
- ABAQUS - implicit used but dynamic displacement rate applied
- 3 crystallographic realizations

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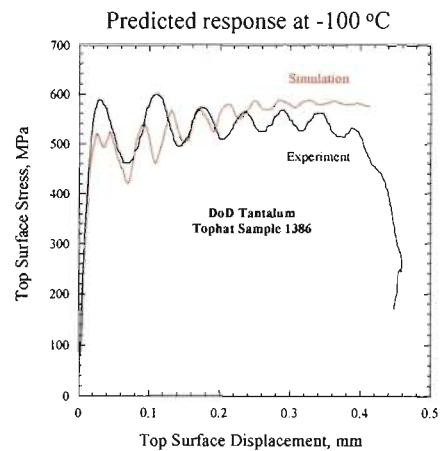
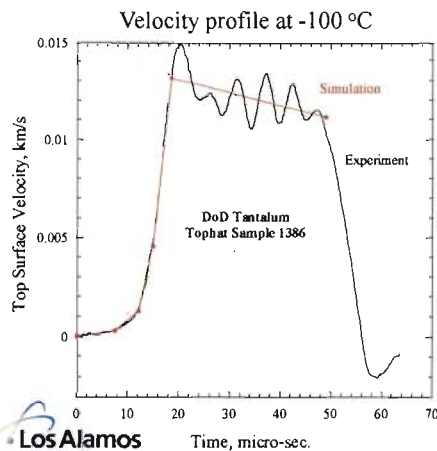
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Localization - Ta

- Dynamic response of the SHPB system is not represented
- Piecewise linear velocity profile applied uniformly to the top surface
- Top surface stress response is weighted average of top row of elements



Small strain continuum model for the inner and outer regions

Stress

$$\dot{\sigma} = \underline{C} \dot{\epsilon}^e$$

Strain

$$\dot{\epsilon} = \dot{\epsilon}^e + \dot{\epsilon}^p$$

Normality Flow

$$\dot{\epsilon}^p = \frac{3}{2} \frac{\dot{\epsilon}^p}{\bar{\sigma}} \frac{\sigma'}{\sigma}$$

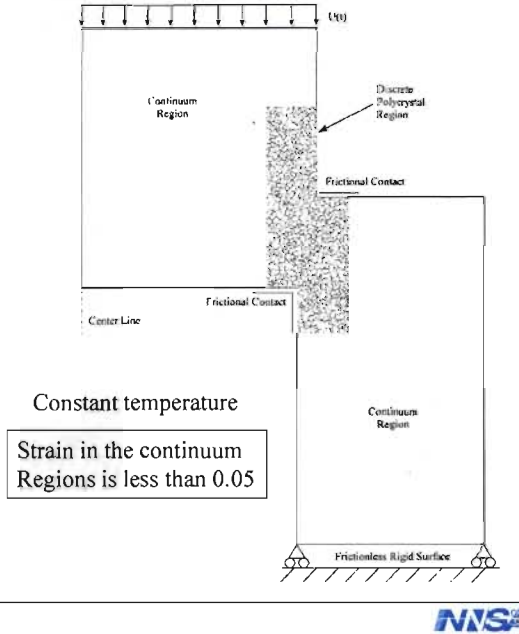
Flow Surface

$$\bar{\sigma} - \sigma_f(\dot{\epsilon}^p, \theta) = 0$$

$$\bar{\sigma} = \sqrt{\frac{3}{2} \sigma' \cdot \sigma'}$$



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MTS model - continuum

$$\sigma_f = \sigma_a + \frac{\mu}{\mu_0} (S_i(\dot{\epsilon}, \theta) \dot{\sigma}_i + S_e(\dot{\epsilon}, \theta) \dot{\sigma}_e)$$

$$\sigma_a = 40 \text{ MPa} \quad \dot{\sigma}_i = 1203, 167 \text{ MPa} \quad \dot{\sigma}_e = 0$$

$$\mu = \mu_0 - \frac{D_0}{\exp\left(\frac{\theta_0}{\theta}\right) - 1}$$

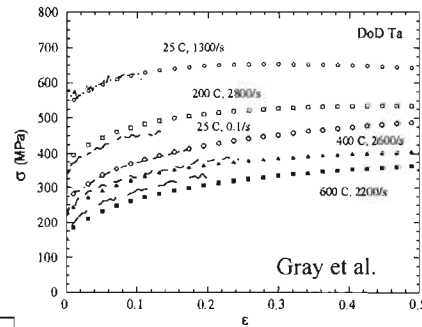
$$\mu_0 = 65.25 \text{ GPa} \\ D_0 = 0.38 \text{ GPa} \\ \theta_0 = 40^\circ \text{K}$$

$$S_i(\dot{\epsilon}, \theta) = \left(1 - \left[\frac{k\theta}{\mu b^3 g_{0i}} \ln\left(\frac{\dot{\epsilon}_{0i}}{\dot{\epsilon}}\right) \right]^{q_i} \right)^{1/p_i}$$

$$S_e(\dot{\epsilon}, \theta) = \left(1 - \left[\frac{k\theta}{\mu b^3 g_{0e}} \ln\left(\frac{\dot{\epsilon}_{0e}}{\dot{\epsilon}}\right) \right]^{q_e} \right)^{1/p_e}$$

$$b = 2.863 \text{ \AA} \\ g_{0i} = 0.1236, 5.1463 \\ \dot{\epsilon}_{0i} = \dot{\epsilon}_{0e} = 10^7 \text{ s}^{-1} \\ q_i = 3/2 \\ p_i = 1/2 \\ q_e = 1 \\ p_e = 2/3$$

Follansbee & Kocks, 1988; Chen & Gray, 1996; Maudlin et al., 1999



$$\frac{d\hat{\sigma}_e}{d\epsilon} = h_0 \left(1 - \frac{\hat{\sigma}_e}{\hat{\sigma}_{e2}} \right)^\kappa$$

$$\frac{\hat{\sigma}_{e2}}{\hat{\sigma}_{e0}} = \left(\frac{\dot{\epsilon}}{\dot{\epsilon}_{0e}} \right)^{\frac{kT}{\mu b^3 g_{0e}}}$$

$$h_0 = 2.0 \text{ GPa} \\ \kappa = 3 \\ \hat{\sigma}_{e2} = 350 \text{ MPa} \\ \dot{\epsilon}_{0e} = 10^7 \\ g_{0e} = 1.6$$



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Shear zone single crystal model

Asaro & Rice (1977), Acharya & Beaudoin (2000),
Kothari & Anand (1998), Busso et al. (2000), Kocks (1976)
Kalidindi et al. (1992), Bronkhorst et al. (1992), Anand (1998)

Stress

$$\mathbf{T}^* = \underline{\mathbf{L}} \left[\mathbf{E}^* - \mathbf{A} (\theta - \theta_0) \right] \quad \mathbf{T}^* \equiv \mathbf{F}^{*-1} (\det \mathbf{F}^*) \mathbf{J} \mathbf{F}^{*-T}$$

$$\mathbf{E}^* \equiv \frac{1}{2} (\mathbf{F}^{*T} \mathbf{F}^* - \mathbf{I}) \quad \mathbf{F}^* = \mathbf{F} \mathbf{F}^{\rho-1}$$

Texture

$$\mathbf{m}^\alpha = \mathbf{F}^* \mathbf{m}_0^\alpha$$

$$\mathbf{n}^\alpha = \mathbf{F}^{*-T} \mathbf{n}_0^\alpha$$

Hardening

$$\dot{s}^\alpha = \sum_\beta h^{\alpha\beta} |\dot{\gamma}^\beta| \quad h^{\alpha\beta} = [r + (1-r) \delta^{\alpha\beta}] h^\beta$$

$$h^\beta = h_0 \left(\frac{s_s^\beta - s_0^\beta}{s_s^\beta - s_0^\beta} \right)$$

$$s_s^\beta = \hat{s}_s^\beta (\dot{\gamma}, \theta) = s_{s_0} \left(\frac{\dot{\gamma}^\beta}{\dot{\gamma}_0} \right)^{\frac{k\theta}{A}}$$



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Voronoi
tessellated

1091 grains



Shear zone single crystal model

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Flow Rule

$$\mathbf{L}^\rho = \dot{\mathbf{F}}^\rho \mathbf{F}^{\rho-1} = \sum_\alpha \dot{\gamma}^\rho \mathbf{S}_0^\alpha \quad \mathbf{S}_0^\alpha \equiv \mathbf{m}_0^\alpha \otimes \mathbf{n}_0^\alpha$$

$$\dot{\gamma}^\alpha = \dot{\gamma}_0 \exp \left[-\frac{F_0}{k\theta} \left\langle 1 - \left\langle \frac{|\tau^\alpha| - s^\alpha \frac{\mu}{\mu_0}}{s_s^\alpha \frac{\mu}{\mu_0}} \right\rangle^p \right\rangle^q \right] \text{sgn}(\tau^\alpha)$$

$$\tau^\alpha \equiv (\mathbf{C}^* \mathbf{T}^*) \cdot \mathbf{S}_0^\alpha$$

$$\mathbf{C}^* = \mathbf{F}^{*T} \mathbf{F}^*$$

$$\frac{\mu}{\mu_0} \equiv \frac{C_{12}}{C_{12e}} = 1 + \frac{m_{12}}{C_{12e}} \theta$$

Adiabatic Heating

$$\dot{\theta} = \frac{\eta}{\rho c_p} \sum_\alpha \tau^\alpha \dot{\gamma}^\alpha$$



Simmons & Wang, 1971

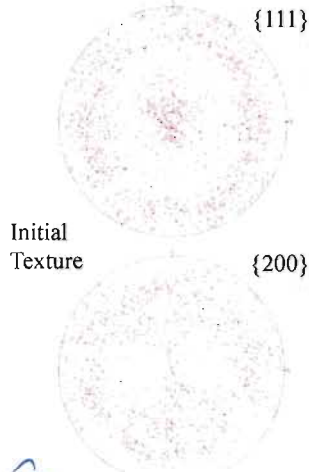
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900 μm between corners



Single crystal model - Ta

Single Crystal Model



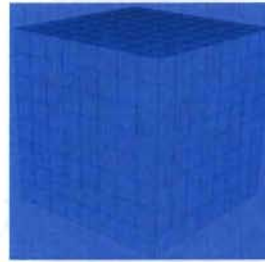
$\rho = 16640 \text{ kg/m}^3$
 $c_p = 150 \text{ J/kg-K}$
 $\alpha = 6.5 \text{ } \mu\text{m/m-K}$
 $\eta = 0.0, 0.95$
 $m_{11} = -24.5 \text{ MPa/K}$
 $C_{110} = 268.5 \text{ GPa}$
 $m_{12} = -11.8 \text{ MPa/K}$
 $C_{120} = 159.9 \text{ GPa}$
 $m_{44} = -14.9 \text{ MPa/K}$
 $C_{440} = 87.1 \text{ GPa}$
 $r = 1.4$
 $\dot{\gamma}_0 = 10^7 \text{ sec}^{-1}$
 $s_0 = 50 \text{ MPa}$
 $s_i = 550 \text{ MPa}$
 $F_0 = 2.1 \times 10^{-19} \text{ J}$
 $p = 0.34$
 $q = 1.66$
 $s_s = 125 \text{ MPa}$
 $h_s = 300 \text{ MPa}$
 $A = 10^{-18} \text{ J}$

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Material Parameter Evaluation 512 Elements/Grains

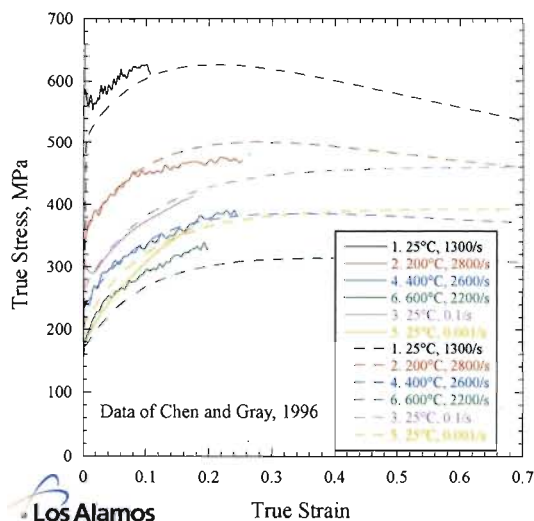
Ta - 24 BCC systems
{110}<111>, {112}<111>



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Single crystal model - Ta

Single Crystal Model



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Material Parameter Evaluation

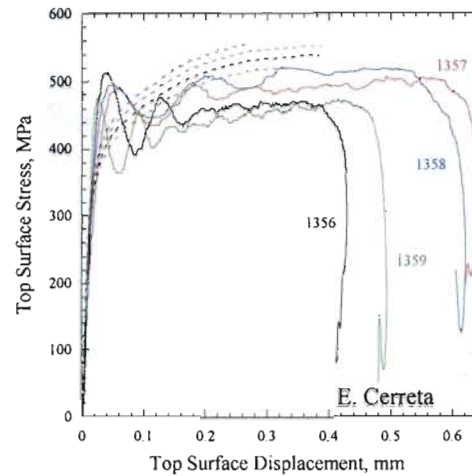
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Experimental load-displacement response is over-predicted

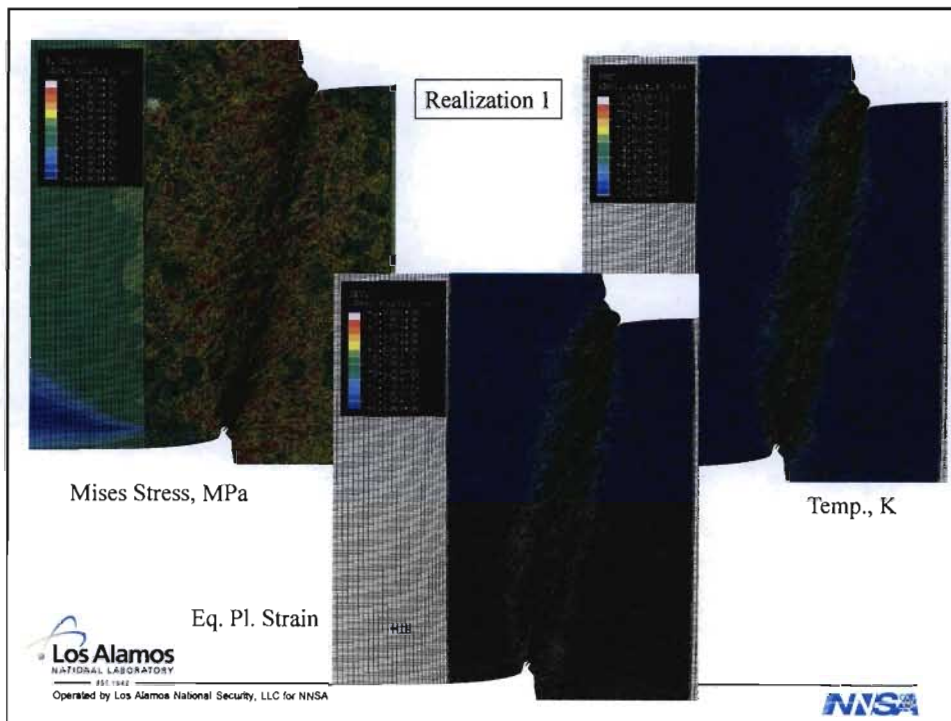
- Although the loading rate difference is described, overall the model over-predicts the magnitude of load required to deform the sample.
- This is believed to be caused primarily by 2D representation.
- The single crystal model could also be inadequate to well represent this level of detail.



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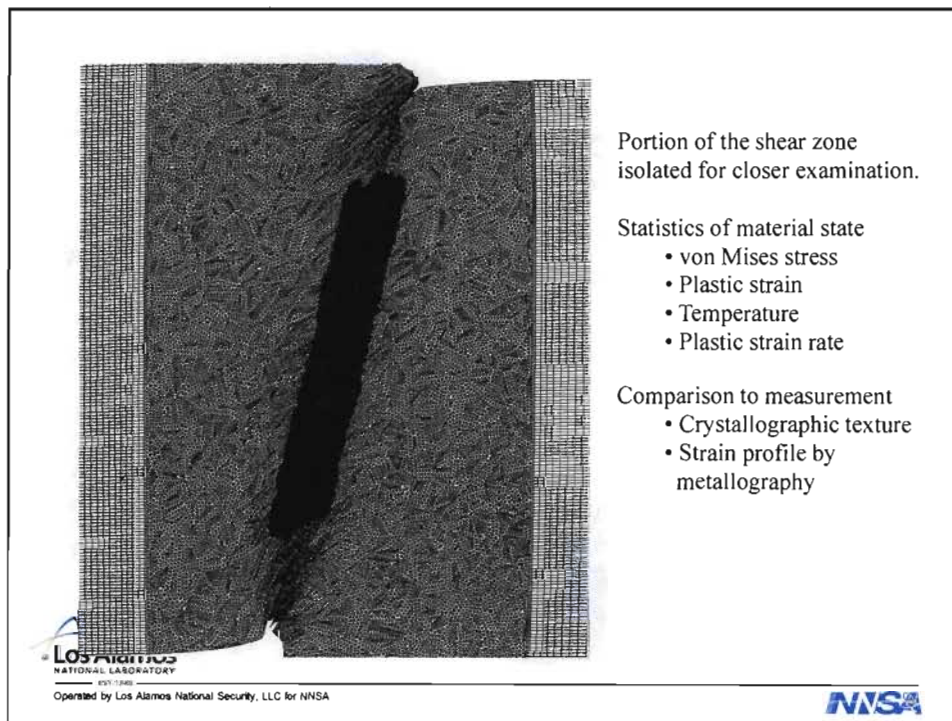
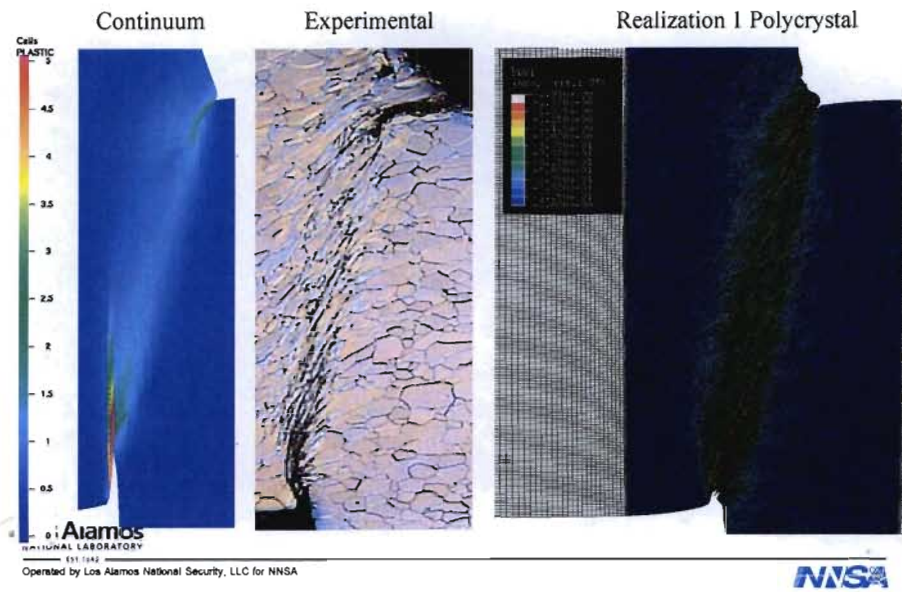


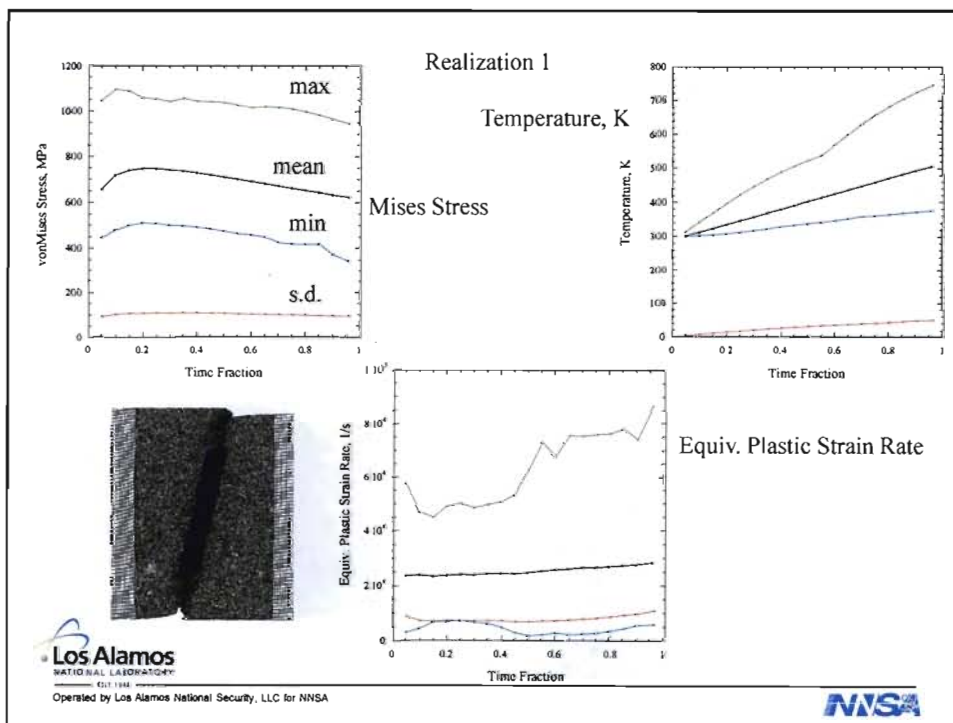
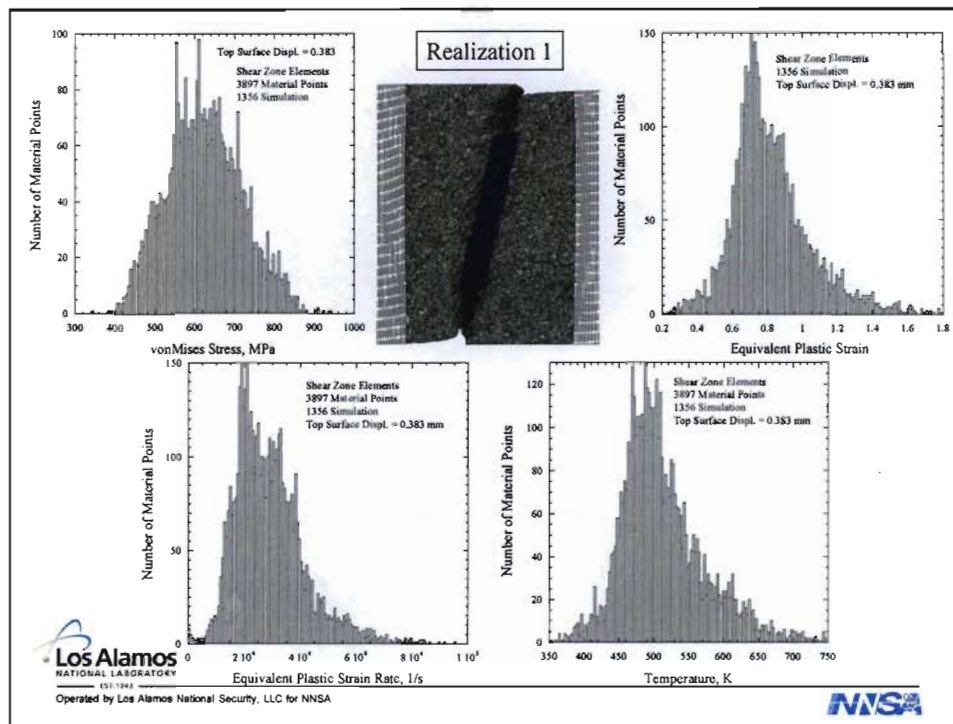
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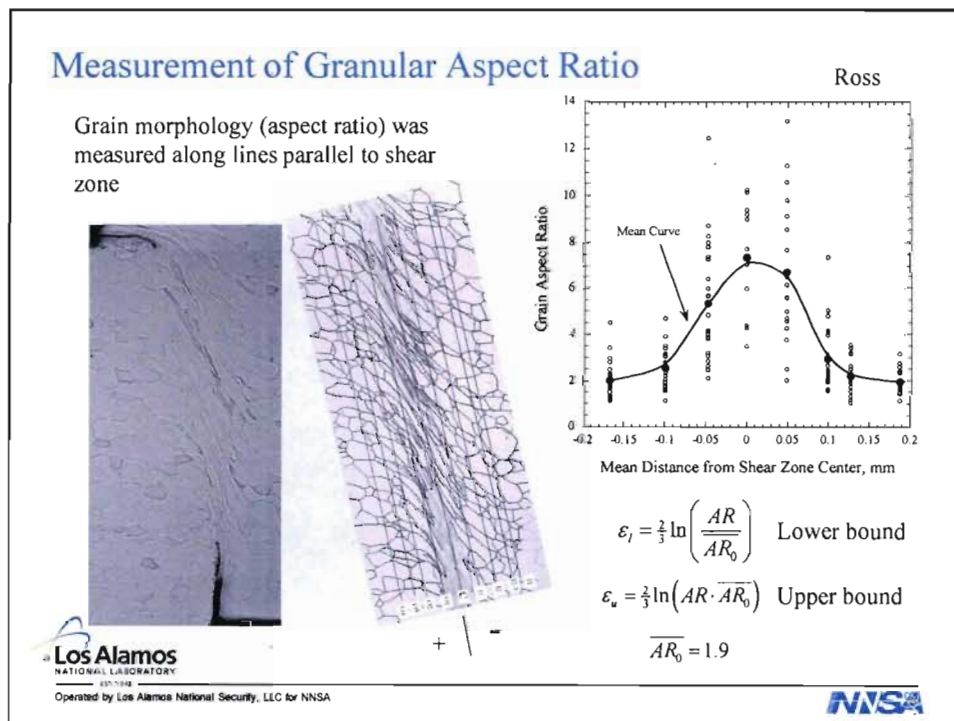
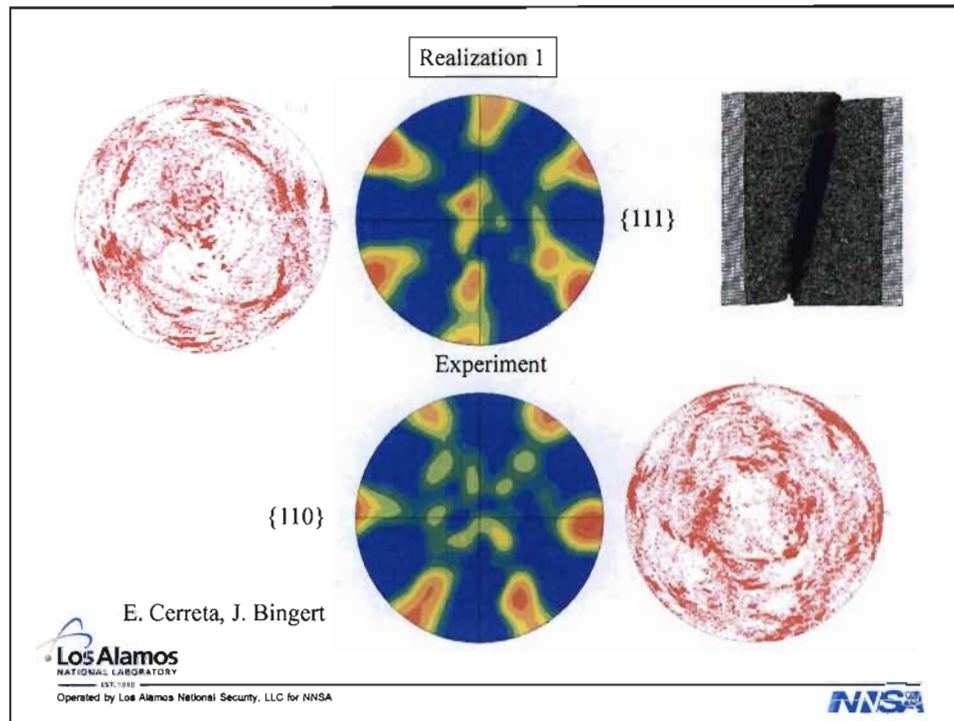
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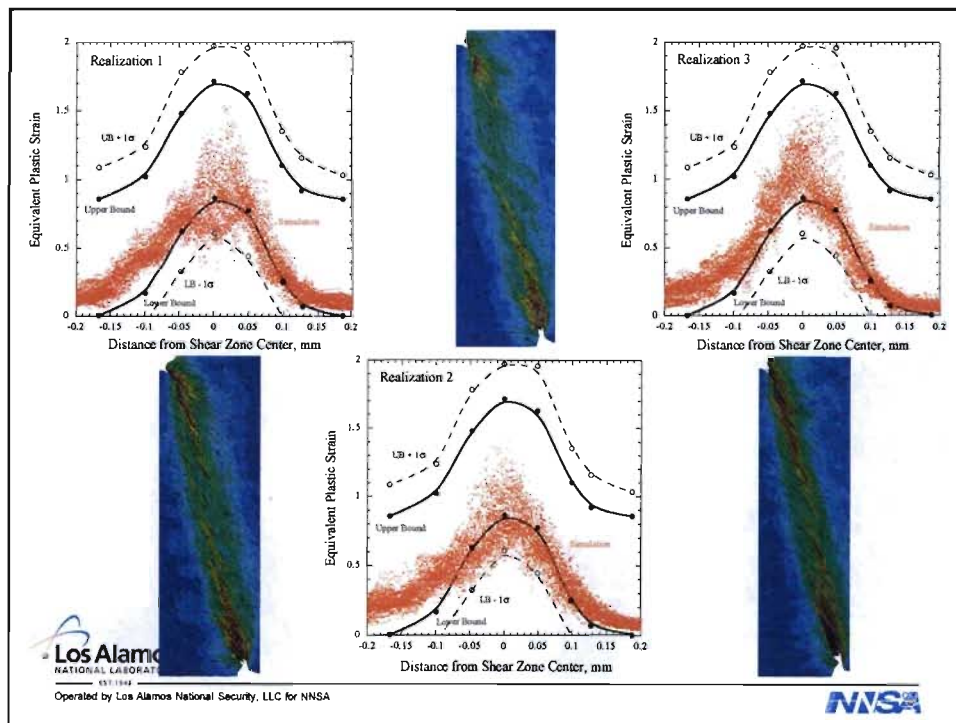
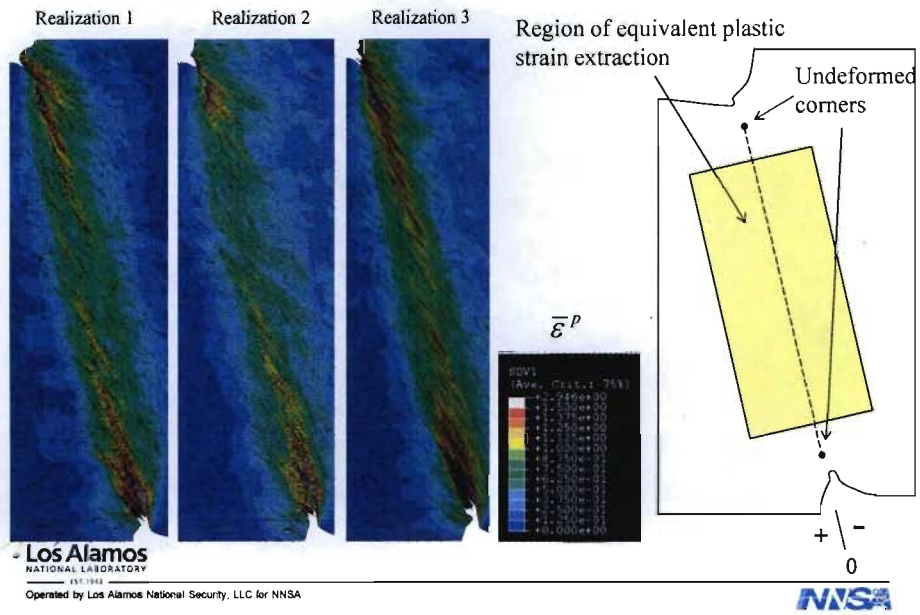
Localization - Ta





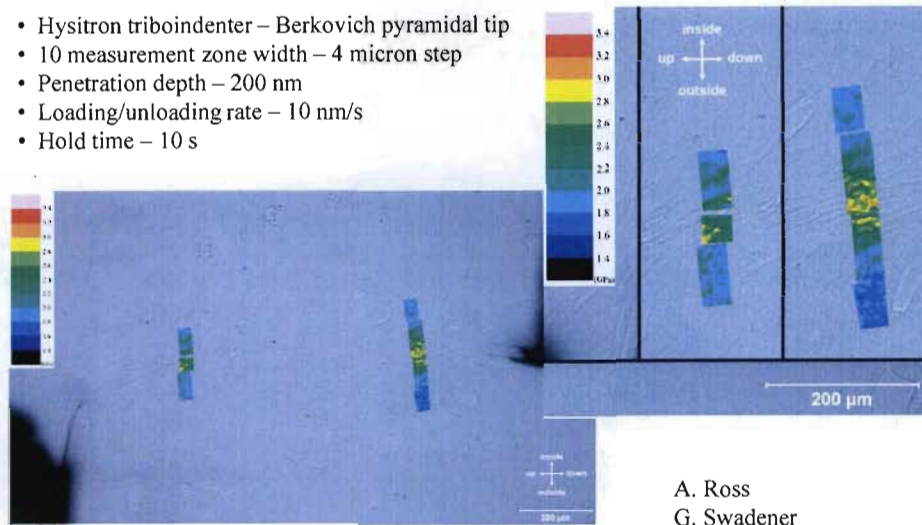


3 Differing Crystallographic Realizations



Shear zone region interrogation by nano-indentation

- Hysitron triboindenter – Berkovich pyramidal tip
- 10 measurement zone width – 4 micron step
- Penetration depth – 200 nm
- Loading/unloading rate – 10 nm/s
- Hold time – 10 s



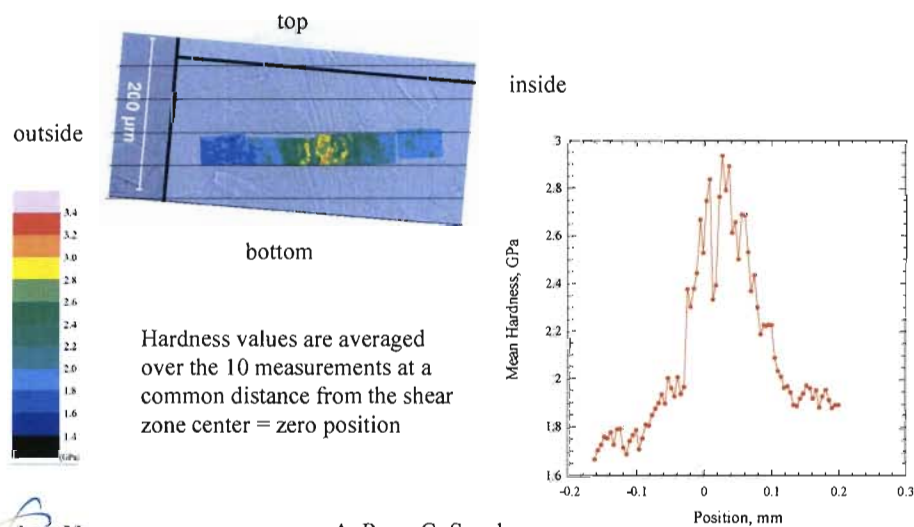
A. Ross
G. Swadener

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Shear zone interrogation by nano-indentation



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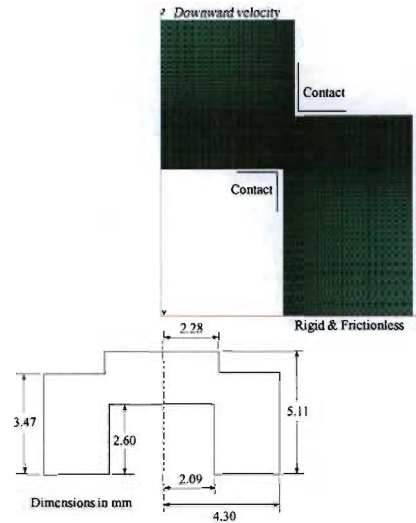
A. Ross, G. Swadener

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Top Hat Problem: Experiments vs. Computations

- Experiments
 - Split-Hopkinson pressure bar
 - Top-to-bottom loading (10-25 m/s)
- Computations
 - 46,000 triangular elements in EPIC
 - 2,000 elements in shear section
 - Axisymmetric, adiabatic conditions
 - Frictional self-contact at both corners
 - Piece-wise linear (in time) downward velocity profile applied to top surface
 - Rigid frictionless base

Compare response (average stress vs. deflection at top surface)

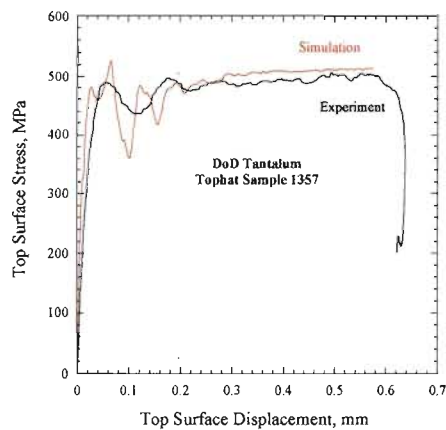


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Top Hat Problem: Tantalum

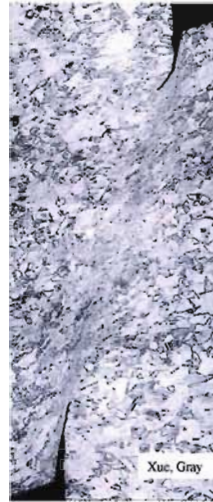
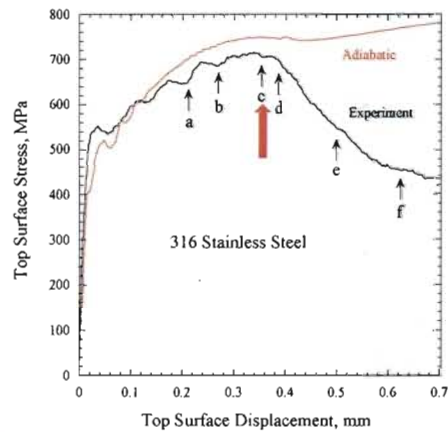


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Top Hat Problem: 316L Stainless Steel

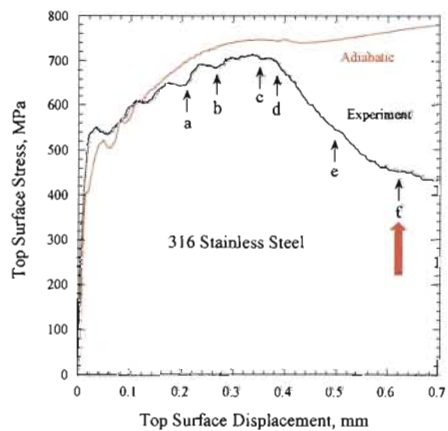


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Top Hat Problem: 316L Stainless Steel



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Sub-grid Computational Approach



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Modeling Localization: Conventional vs. Embedded-zone Approach

- Localization phenomena involve a **discontinuous strain field**
- Conventional computational techniques
 - Can represent jumps in strain only **at inter-element boundaries**, not inside elements
 - Represent a localization zone as a band of elements, spanned by a single element in width, and thus yield **mesh-dependent** results
 - Necessitate resolving localization zones explicitly, which is **prohibitively expensive, and requires a priori knowledge of the band location and orientation**
- Sub-grid computational technique
 - Allows part of the localization band to be embedded **inside an element**, obviating the need for excessive mesh refinement
 - Allows localization band width to be **specified as a material parameter**, instead of being dictated by the mesh size
 - Allows band orientation to be determined based on material **stability analysis**
 - Allows **smooth transition** from uniform to localized deformation, and facilitates use of **different material models inside and outside the band**



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Localization: Bifurcation Analysis

- Based on Hadamard (1903), Hill (1962)
- Localization zone R_L characterized by thickness b , unit normal n and strain jump direction m .

- Compatibility

$$[[\epsilon_{ij}]] = \frac{1}{2} g(m_i n_j + m_j n_i)$$

- Equilibrium

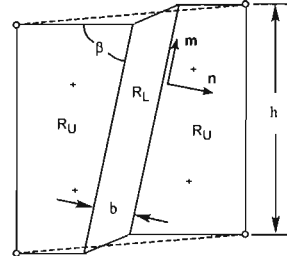
$$[[t_j]] = [[n_i \sigma_{ij}]] = n_i [[\sigma_{ij}]] = 0$$

- Constitutive relations

$$\dot{\sigma}_{ij} = D_{ijkl} \dot{\epsilon}_{kl}$$

- Combined, they lead to: Find n and m such that

$$(n_i D_{ijkl} n_l) m_k = 0 \implies \det(n_i D_{ijkl} n_l) = 0$$



Enhanced Assumed Strain (B-bar) Method

- Based on the work of Belytschko et al (1988), Fish and Belytschko (1988)

- Conventional FEM approximation

$$u_i = N_i^A d^A, \quad \epsilon_{ij} = B_{ij}^A d^A$$

- Strain field is **continuous** within elements

$$\epsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right), \quad B_{ij}^A = \frac{1}{2} \left(\frac{\partial N_i^A}{\partial x_j} + \frac{\partial N_j^A}{\partial x_i} \right)$$

- Introduce assumed strain field $\bar{\epsilon}$ such that

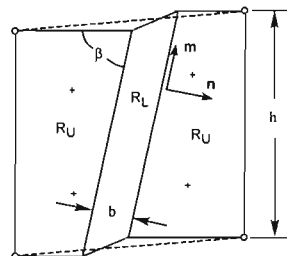
$$\Delta \bar{\epsilon}_{ij} = \bar{B}_{ij}^A \Delta d^A$$

- Constitutive relations

$$\sigma = \Sigma(\bar{\epsilon})$$

- Internal force vector in the explicit Central Difference Method

$$M^{AB} \ddot{d}^B = (f_{\text{ext}}^A - f_{\text{int}}^A), \quad f_{\text{int}}^A = \int_{\Omega} \bar{B}_{ij}^A \Sigma_{ij}(\bar{\epsilon}) d\Omega$$



Enhanced Assumed Strain (B-bar) Method

- Weak discontinuity is characterized by

$$T = \frac{1}{2}(m \otimes n + n \otimes m)$$
- Projection of $\Delta \varepsilon$ onto discontinuous mode

$$(\Delta \varepsilon : T)T$$
- The assumed strain field takes the form

$$\Delta \bar{\varepsilon}_{ij} = \Delta \varepsilon_{ij} + \alpha (\Delta \varepsilon_{kl} T_{kl}) T_{ij}$$
- Recall that, by definition, $\Delta \bar{\varepsilon}_{ij} = \bar{B}_{ij}^A \Delta d^A$
- The B-bar operator has the form

$$\bar{B}_{ij}^A = (\delta_{ik} \delta_{jl} + \alpha T_{kl} T_{ij}) B_{kl}^A$$

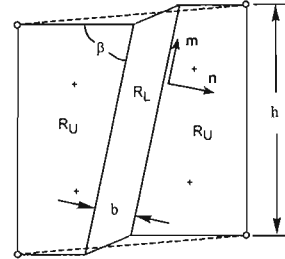
where $\alpha = \alpha^L$ in R_L , and $\alpha = -\alpha^N$ in R_U .

- Modified Gauss quadrature scheme

$$\int_{\Omega^*} f(x) d\Omega \approx \sum_q^{N_q} f(x_q) J(\xi_q) W_q \quad \rightarrow \quad \int_{\Omega^*} f(x) d\Omega \approx \sum_q^{N_q} \left[\frac{b}{h} f^L(x_q) + \left(1 - \frac{b}{h}\right) f^N(x_q) \right] J(\xi_q) W_q$$



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Enhanced Assumed Strain (B-bar) Method

- The unit normal vector n is known. Must determine the unit strain jump vector m , and the scalars α^L and α^N
- Equivalence between strain fields

$$\int_{\Omega} (\bar{B}^A - B^A) d\Omega = 0$$

- Definition of m as a unit vector

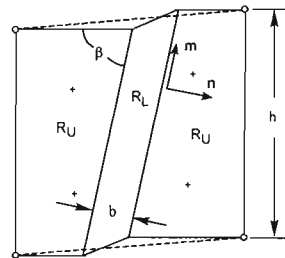
$$\|m\| = 1$$
- Equilibrium at the band-matrix interface

$$[[t]] = n \cdot [[\sigma]] = 0$$

- Traction continuity condition is enforced at each Gauss point, at the end of each time step, using an iterative solution scheme.



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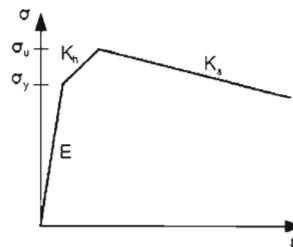


Numerical Examples

Strain Softening: A Prerequisite for Localization

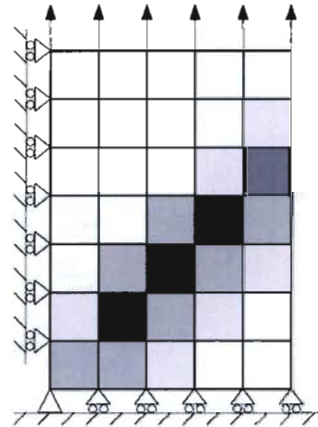
- Softening gives rise to material instability

$$\det(n_i D_{ijkl} n_l) = 0$$
- Mechanisms of material softening include
 - Thermal softening: heating due to plastic work
 - Damage (voids and/or cracks)
 - Dynamic recrystallization
- For simplicity, we choose an isotropic elastic-plastic material model, with linear isotropic hardening/softening:
 - Elastic parameters: $E = 200 \text{ GPa}$, $\nu = 0.29$
 - Strength parameters:
 $\sigma_y = 300 \text{ MPa}$, $\sigma_u = 310 \text{ MPa}$
 $K_h = 220 \text{ MPa}$, $K_s = -66 \text{ MPa}$
 - Mass density, $\rho = 7850 \text{ kg/m}^3$

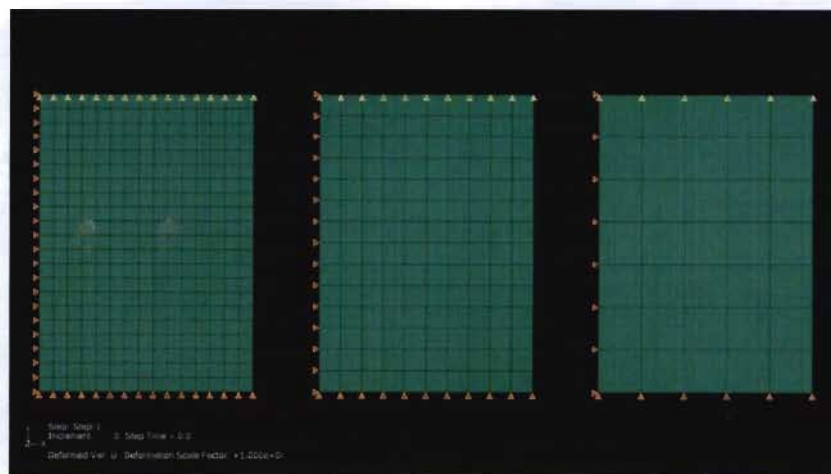


Extension of Rectangular Block

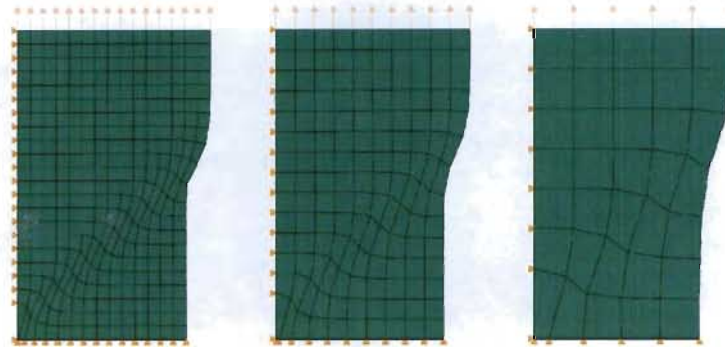
- Dimensions: $L = 0.14$ m, $w = 0.10$ m
- Upper-right quadrant modeled using three meshes to test results for mesh dependence
 - Mesh 1:
5 × 7 embedded-zone elements, $b/h = 1/3$
 - Mesh 2:
10 × 14 embedded-zone elements, $b/h = 2/3$
 - Mesh 3:
15 × 21 **conventional** elements, $b/h = 1$
- Band width $b = 3.33$ mm is assumed
- Bottom-left element is weakened
- Shaded area represents localization zone detected by bifurcation analysis



Extension of Rectangular Block



Extension of Rectangular Block



- Location and orientation of band, and its effect on overall deformation, are captured reasonably well, even by coarsest mesh.

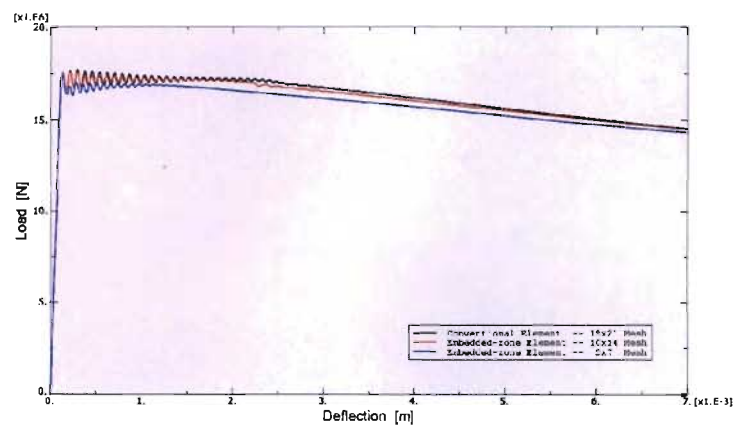
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Extension of Rectangular Block



- Correct structural response is captured throughout loading

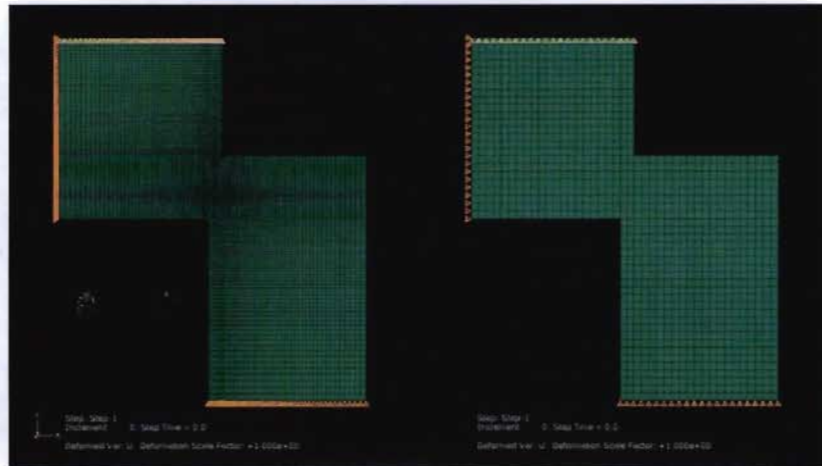
process
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Top Hat Problem: Stainless Steel



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Concluding Remarks

- A sub-grid finite element formulation for modeling strain localization was presented, and its implementation within an explicit Lagrangian framework was described
- The method can treat
 - Plane-strain and axisymmetric problems
 - Small-deformation and finite-deformation problems
 - Compressible and incompressible material behavior
 - Problems involving highly-dynamic loading
- Focus of current and future work
 - Accounting properly for softening mechanisms like damage (e.g. Johnson and Addessio, 1988), thermal softening (e.g. Anand, 1985), and dynamic recrystallization

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Summary & comments

- Damage and failure in polycrystalline metallic materials is strongly dependent upon microstructural details – modern lower length scale tools are necessary to assist in our learning about these complex physical processes. This can then properly motivate the development of physically based higher length scale models for use in applications.
- In general, we do not have the ability to adequately represent the topology of polycrystalline microstructures – this is especially true in 3D. Sub-granular initial state is also a concern.
- The modeling results compare OK with experimental measurements. Comments are:
 - Stress response is consistently over-predicted.
 - Texture prediction clearly shows 2D deformation gradient.
 - Shear zone strain profile is well predicted.
 - Due to tessellation limitations, triangular elements were required.

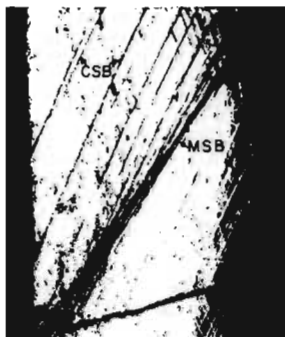


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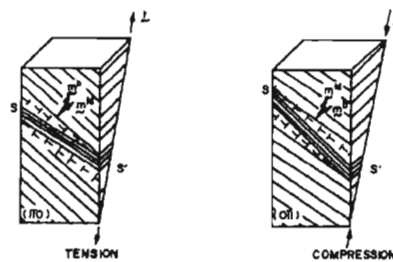


Single Crystal Work

Tension



Asaro and Rice, 1977
Asaro, 1979
Chang and Asaro, 1981



- Single crystals of Al-2.8Cu with θ' and θ precipitate particles
- Course slip bands were aligned with the active slip systems
- Macroscopic shear bands were preceded by the formation of coarse slip bands but were not aligned with active slip systems.



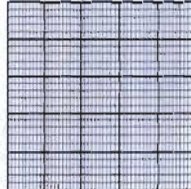
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Polycrystal Work

Plane Strain Compression

1	2	3	4	5
6	7	8	9	10
11	12	13	14	15
16	17	18	19	20
21	22	23	24	25



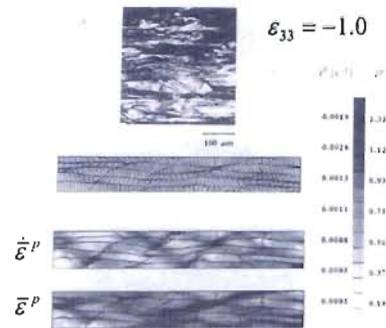
- Each of the 25 crystals was divided into 50 equal-sized elements.
- Localization into zones oriented 30-40° from horizontal.
- Shear banding is a natural part of inhomogeneous large deformation process and is closely linked with evolution of crystallographic texture.



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Bronkhorst et al. (1992)
Anand & Kalidindi (1994)

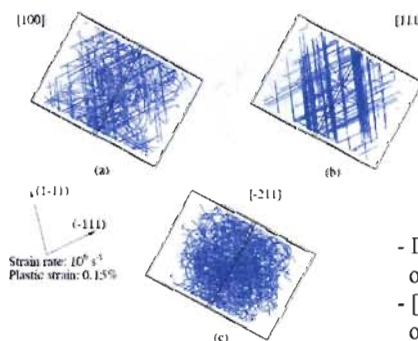
OFHC Cu



Single Crystal Work

Wang, Beyerlein, LeSar (2006, 2007)

Talk on **Tuesday 8:00**



- Dislocation dynamics on Cu at rates of deformation $10^4 - 10^6 \text{ s}^{-1}$.
- [100], [111], [-211] tensile loading orientations.
- Inhomogeneous dislocation accumulation leads to "directional" microstructure.
- Slip bands developing on the more active slip systems, particularly for [111] orientation.



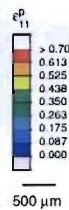
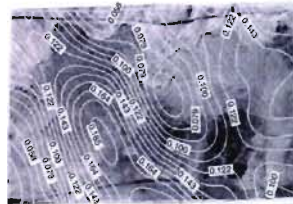
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Single/Polycrystal Work

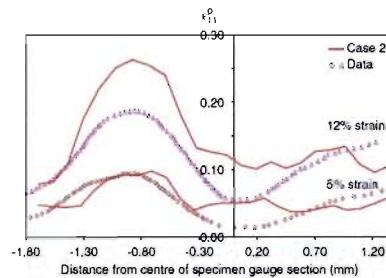
Cheong & Busso (2006)

Tension



(h) Case 2 (max $\epsilon^p_{11} = 0.36$)

- Modeled the tension response of large grain "sheet" sample.
- Able to represent the qualitative profile of localized deformation.
- Required inclusion of misorientation distribution within the large crystals to represent the magnitude correctly.



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Polycrystal Work

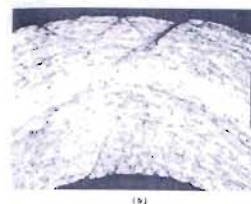
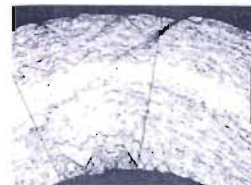
Becker (1992)

Bending

Aluminum



- Taylor-type model applied to each integration point using 8 grains chosen at random from a distribution of 800 grains.
- Resistance to shear band growth by grains which are not favorably oriented.
- Overall localization planes are formed at reasonable strain levels and at appropriate orientations.



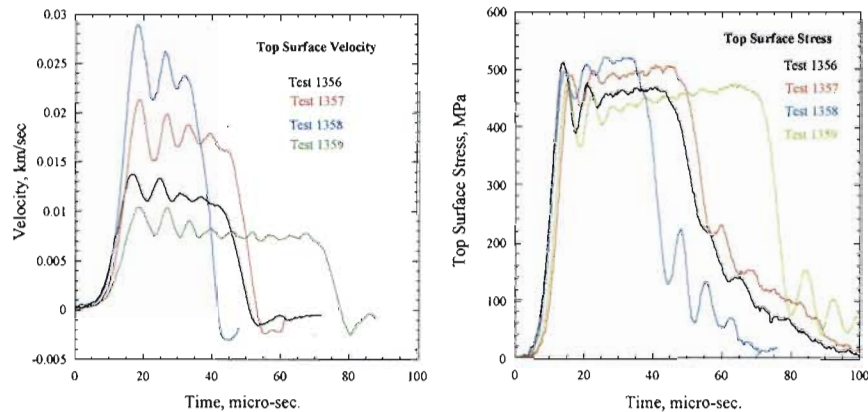
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Tophat Experiments

Results from tophat experiments performed at 25 C



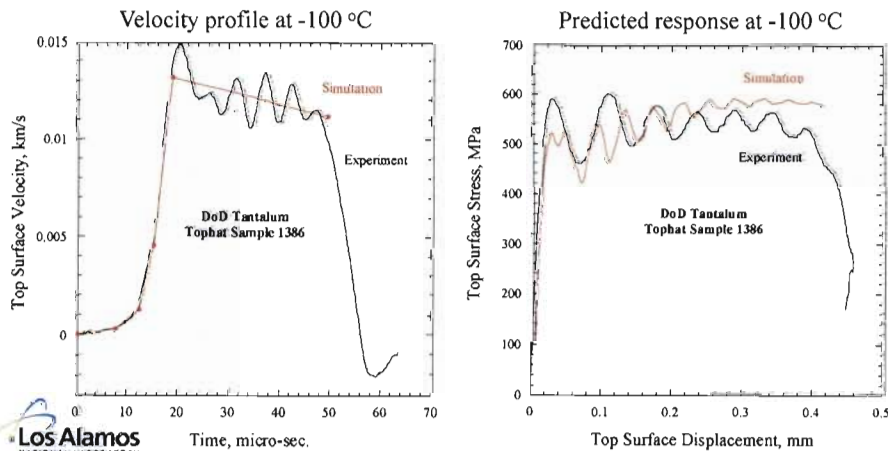
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Localization - Ta

- Dynamic response of the SHPB system is not represented
- Piecewise linear velocity profile applied uniformly to the top surface
- Top surface stress response is weighted average of top row of elements



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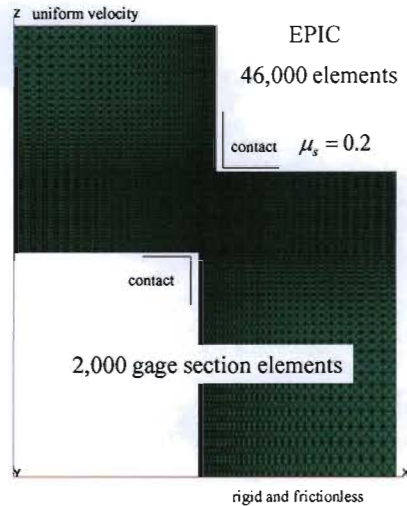
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Localization

- FE code EPIC - explicit
- Axisymmetric and adiabatic
- 46,000 triangular elements
- 2000 shear section elements
- Contact friction surfaces in two corners
- Velocity profile applied to top surface
- Rigid and frictionless base
- Stress response at top surface compared to experiments

Low density mesh shown



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Continuum Model

Addessio and Johnson, 1993; Maudlin et al., 1999
Mason, 1999, 2004

Stress

$$\mathbf{T} = \mathbf{M}\tilde{\mathbf{T}} \quad \mathbf{M} = (1 - \phi)\mathbf{I}$$

$$\dot{\mathbf{T}} = \mathbf{M}\mathbf{L}^{-1}(\mathbf{D}' - \mathbf{D}^{p'}) - \mathbf{M}\left(-\frac{\partial \tilde{P}}{\partial \beta} \text{tr} \mathbf{D}^e + \frac{\tilde{P}}{\rho} \Gamma \mathbf{T} \cdot \mathbf{D}^p\right)\mathbf{I} + \mathbf{M}\mathbf{M}^{-1}\mathbf{T}$$

Strain Rate

$$\mathbf{D} = \mathbf{D}^e + \mathbf{D}^p = \mathbf{D}^e + \mathbf{D}^d + \mathbf{D}^{p'}$$

Plastic Flow Rate

$$\mathbf{D}^p = \frac{1}{\tau_r}(\mathbf{T} - \mathbf{T}^{p'oj})$$

$$\begin{aligned} \tau_r &= 70 \text{ Pa} \cdot \text{s} \\ q_1 &= 1.5 \\ q_2 &= 0.625 \\ q_3 &= 2.25 \end{aligned}$$

$$\tau - \sigma_f(\dot{\epsilon}, \theta)^2 [1 + q_3 \phi^2 - 2q_1 \phi \cosh \delta] = 0$$

$$\tau = \frac{1}{2} \mathbf{T}' \cdot \underline{\underline{\alpha}} \cdot \mathbf{T}'$$

$$\delta = -\frac{3q_2 \tilde{P}}{2\sigma_s} \quad \tilde{P} = -\frac{1}{3} \text{tr} \tilde{\mathbf{T}}$$

Damage Rate

$$\text{tr} \mathbf{D}^p > 0$$

$$\mathbf{D}^d = \frac{\dot{\phi}}{3(1 - \phi)} \mathbf{I} \quad \dot{\phi} = (1 - \phi) \text{tr} \mathbf{D}^p$$

$$\phi_0 = 0.0003$$

Equation of State

$$\tilde{P} = (K_1 \tilde{\beta} + K_2 \tilde{\beta}^2 + K_3 \tilde{\beta}^3)(1 - \Gamma \tilde{\beta}/2) + \Gamma \tilde{E}_s(1 + \tilde{\beta})$$

$$\tilde{\beta} = \frac{\tilde{P}}{\tilde{P}_0} - 1$$

$$\rho = (1 - \phi) \tilde{\rho}$$

$$\begin{aligned} K_1 &= 196.8 \text{ GPa} & \Gamma &= 1.60 \\ K_2 &= 259.8 \text{ GPa} & \rho_0 &= 16,640 \text{ kg/m}^3 \\ K_3 &= 256.6 \text{ GPa} \end{aligned}$$

Plastic Work

$$\dot{\theta} = \frac{\eta}{\rho C_p} E_s$$

$$\begin{aligned} \eta &= 0.95 \\ C_p &= 54.4 \text{ J/kg} \cdot \text{K} \end{aligned}$$

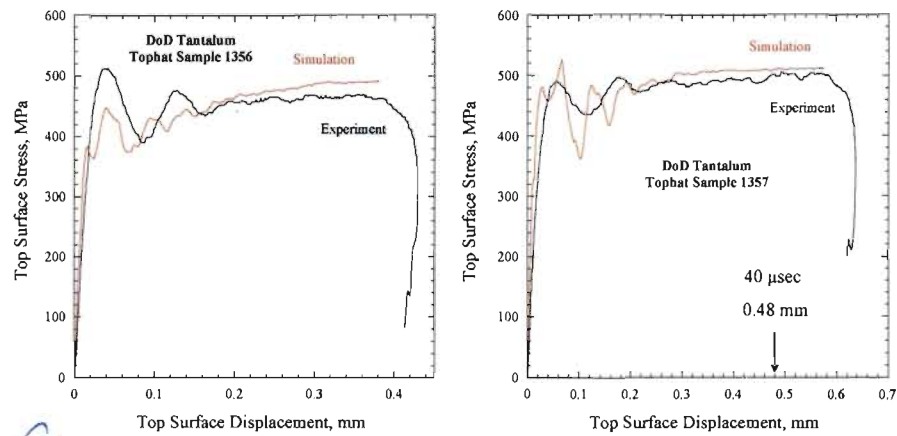
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Localization - Ta

- Predicted responses for sample 1356 and 1357 experiments performed at 25 °C



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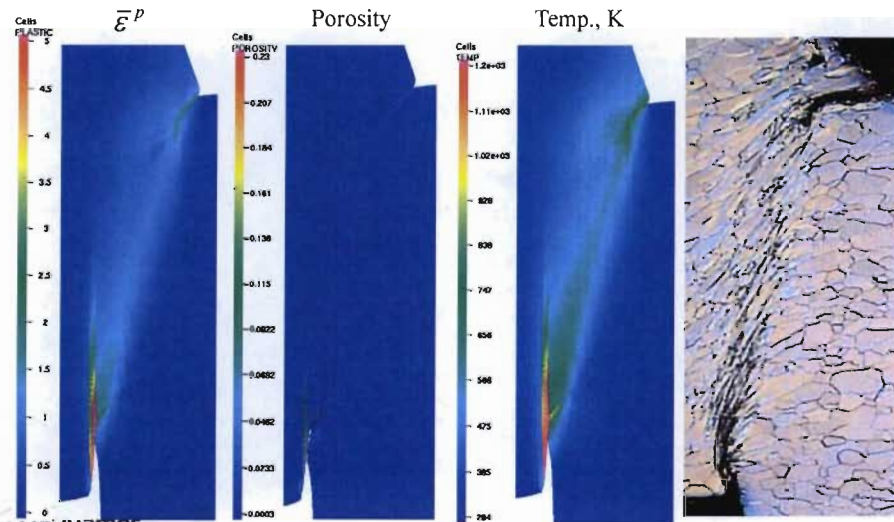
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Continuum Model – Sample 1357

Time = 40 micro-sec.

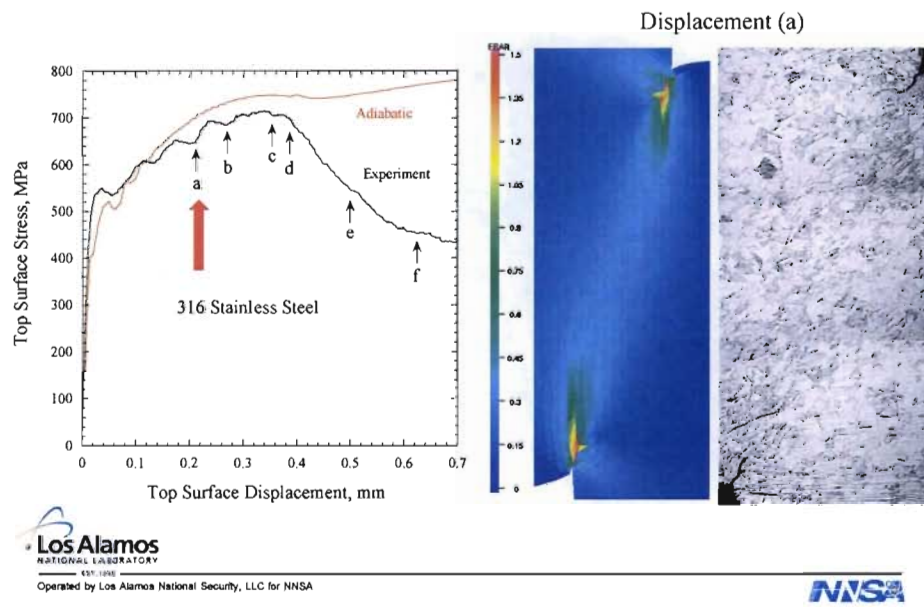
- Contact friction significant in top corner
- Damage plays no role
- Shear band does not form



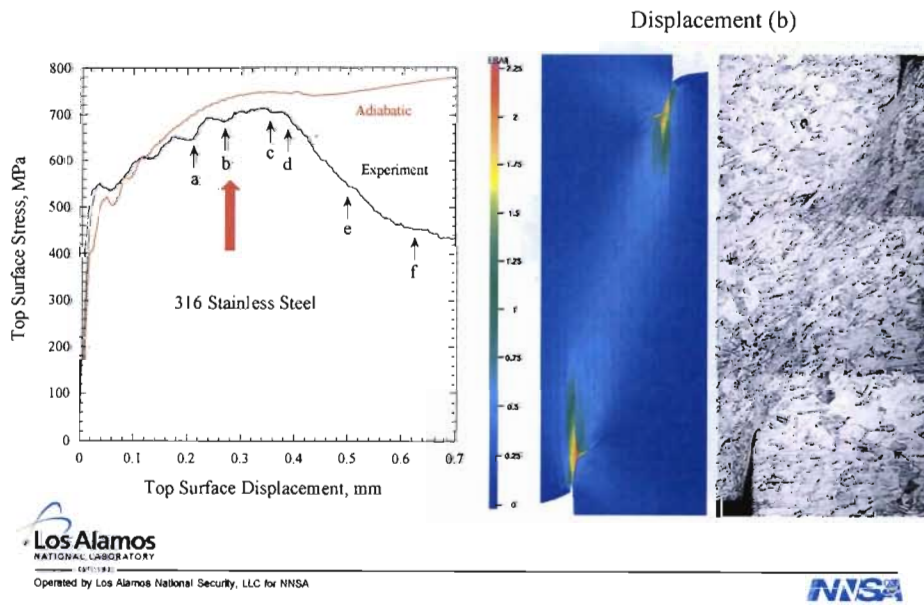
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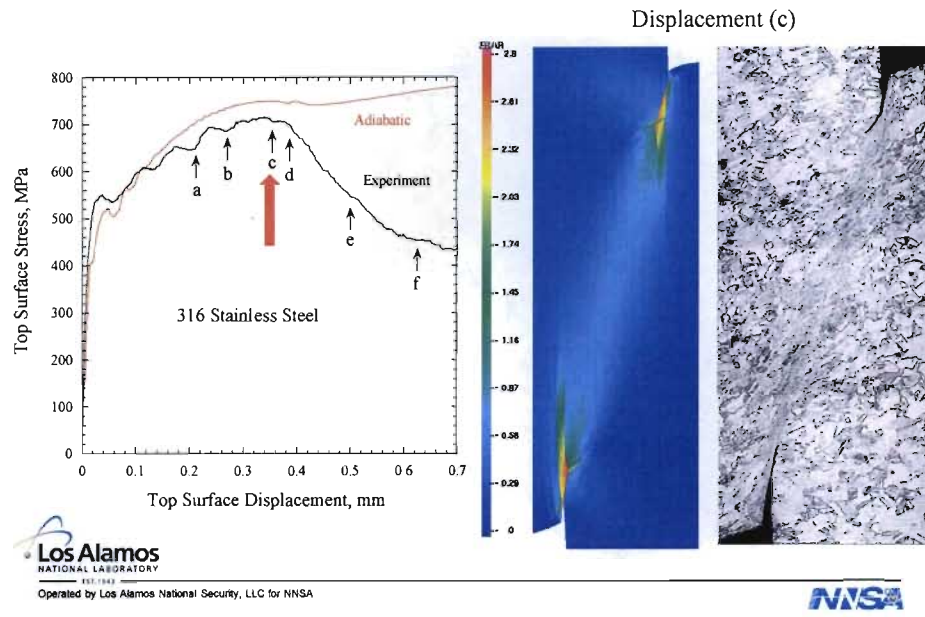
Shear Band – 316 SS



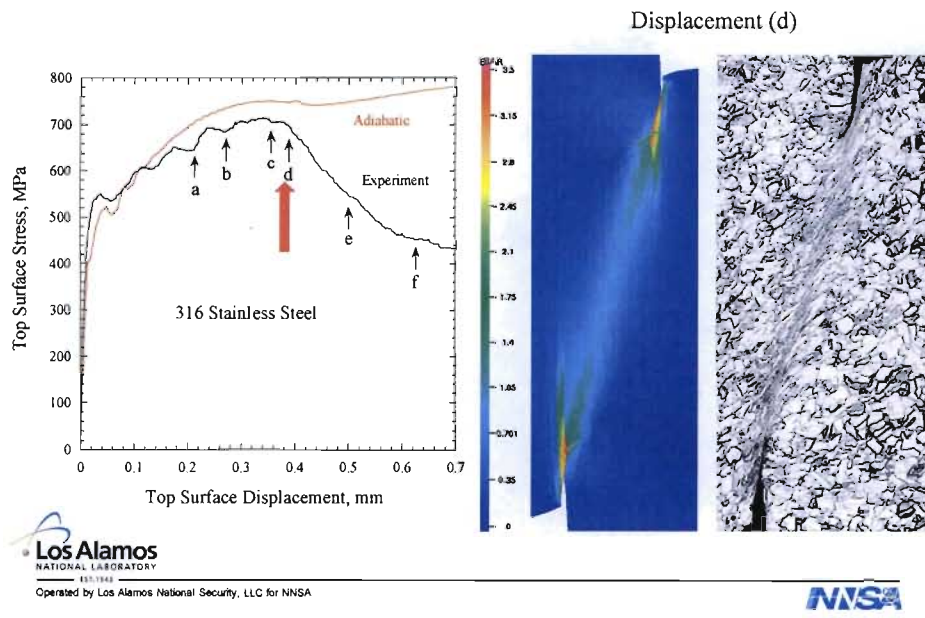
Shear Band – 316 SS



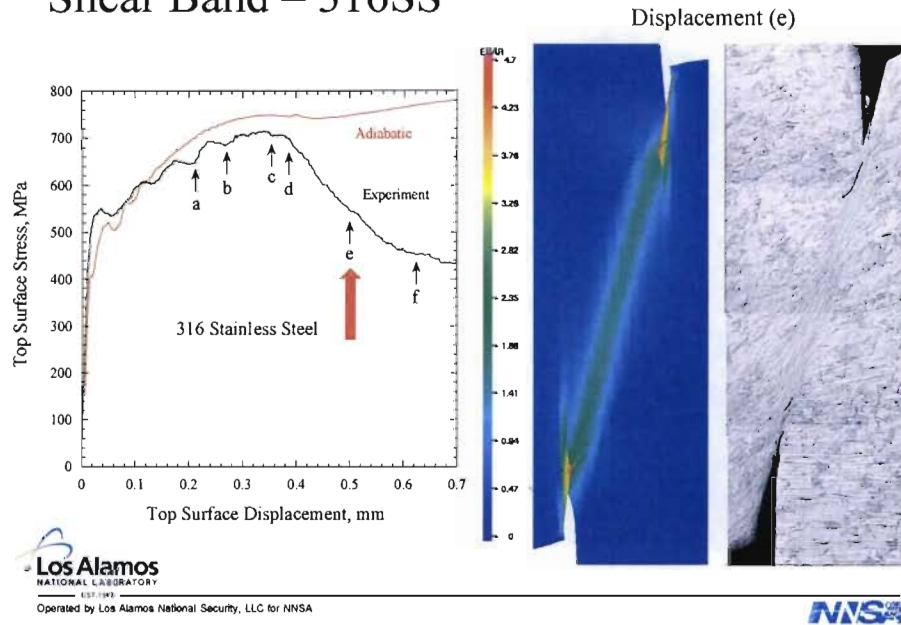
Shear Band – 316 SS



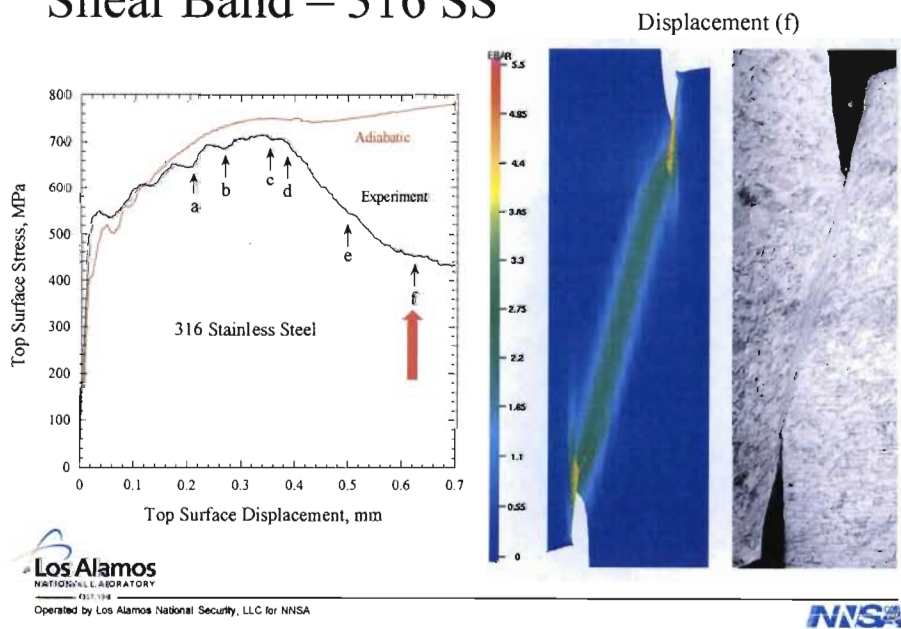
Shear Band – 316SS



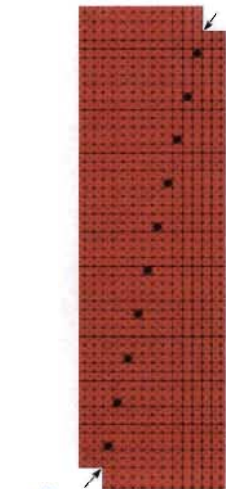
Shear Band – 316SS



Shear Band – 316 SS



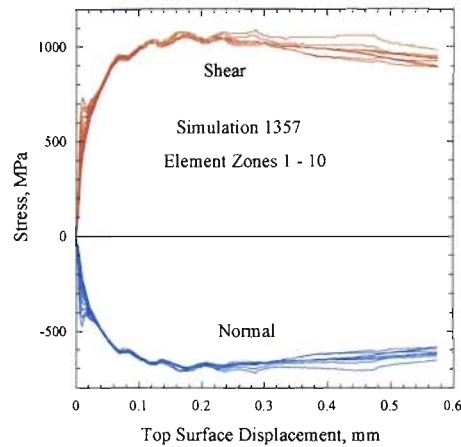
Localization



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- Shear and normal stress defined relative to shear zone line
- Stress in shear zone is quite uniform
- Shear zone line angle changes from 77.2° to 75.8°

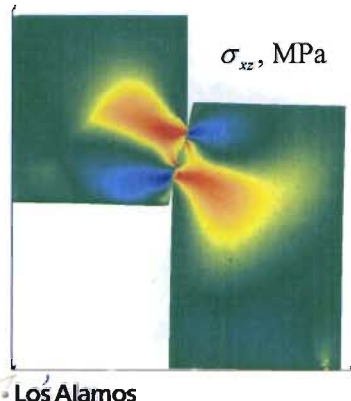


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Localization

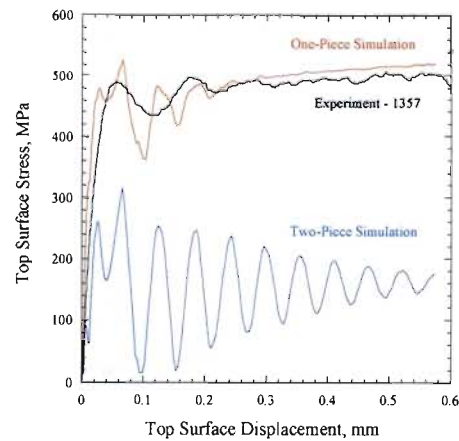
- Shear zone was bisected to form separated “hat” and “brim” pieces – frictionless contact
- There is a resistance to deformation due to the expansion of the outer ring.

Two-piece, frictionless model

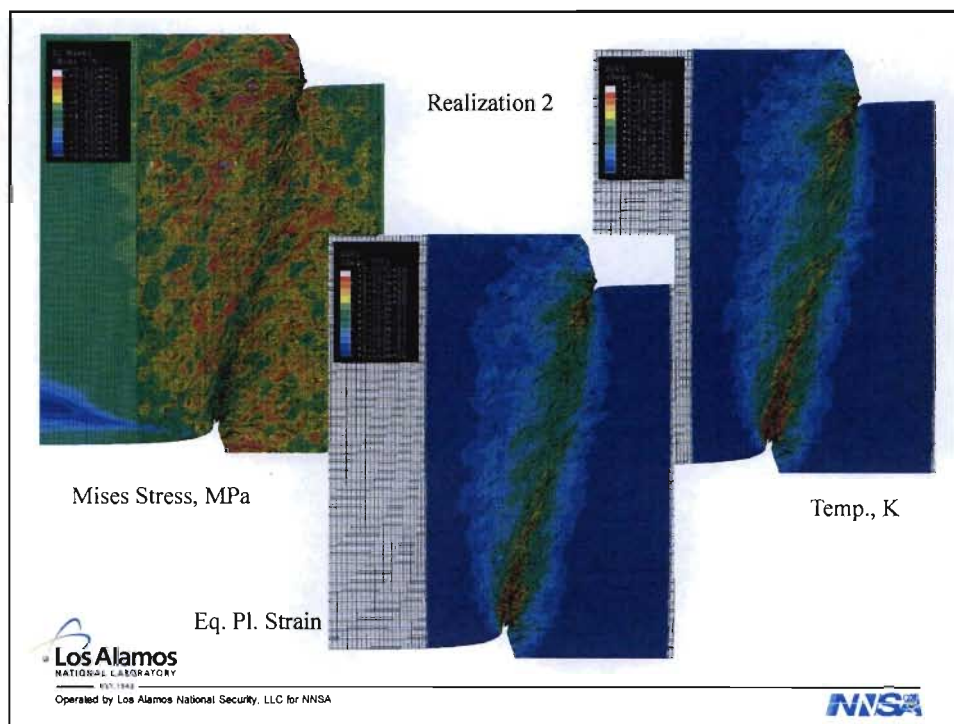
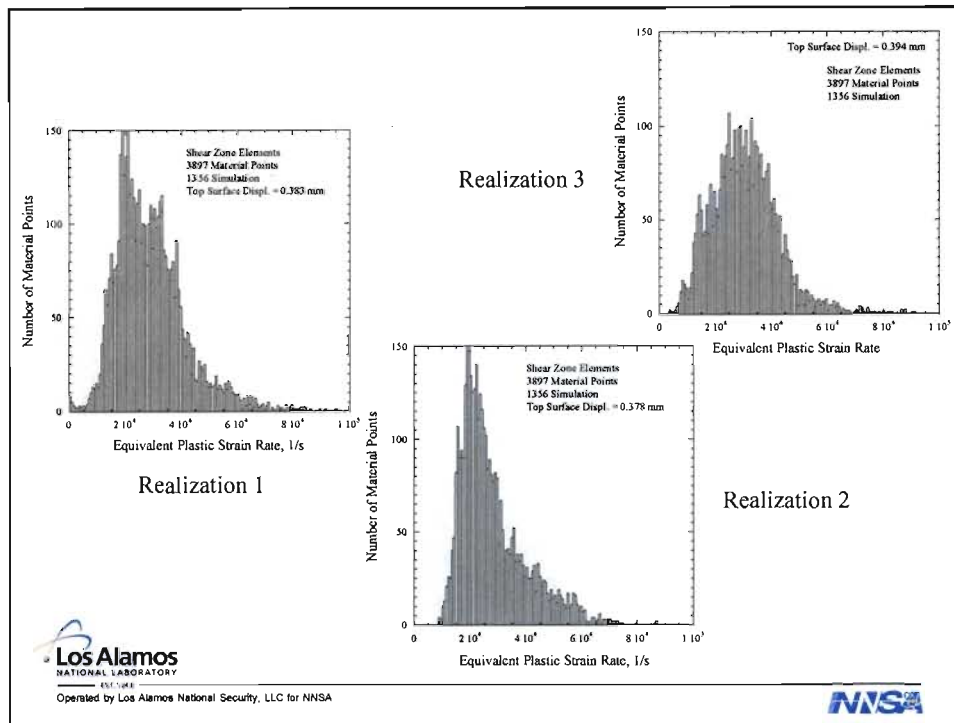


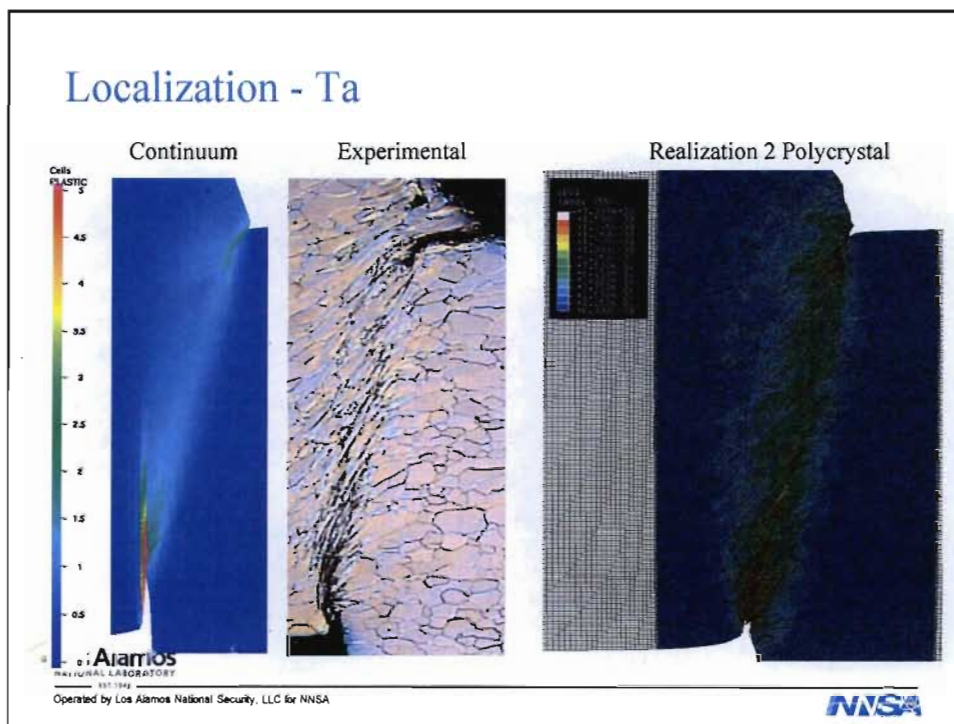
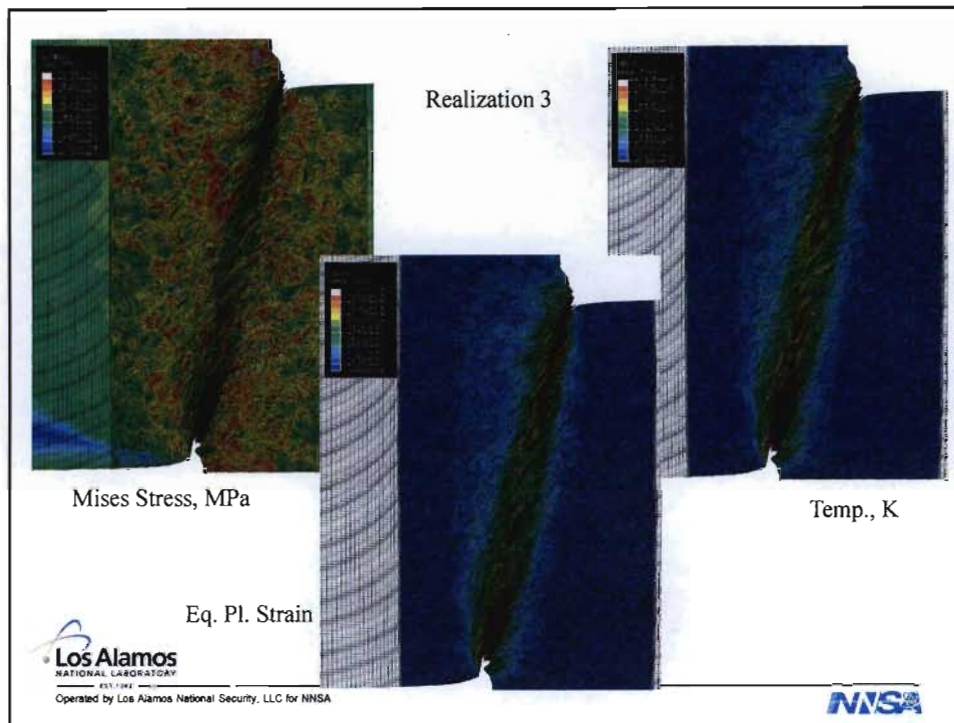
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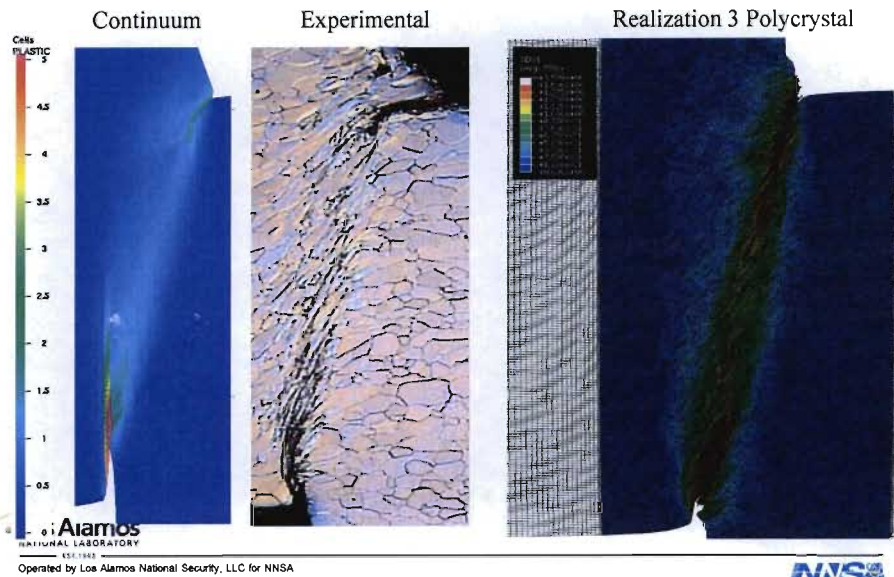


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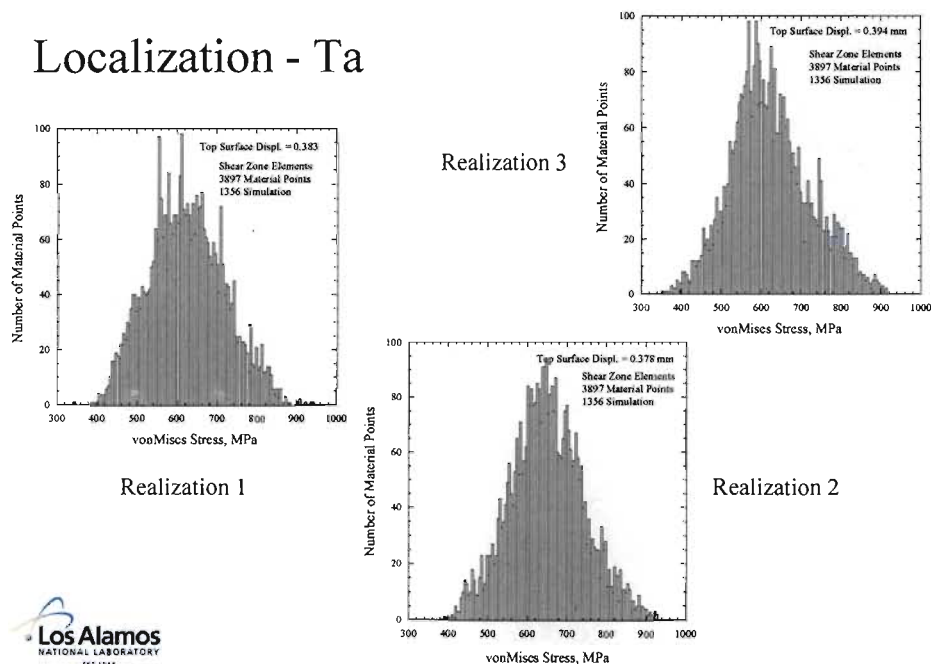




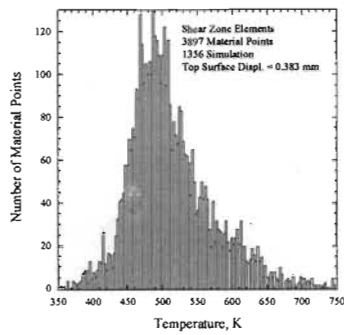
Localization - Ta



Localization - Ta

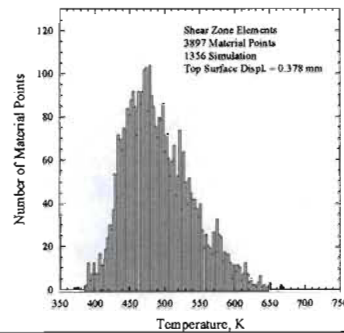
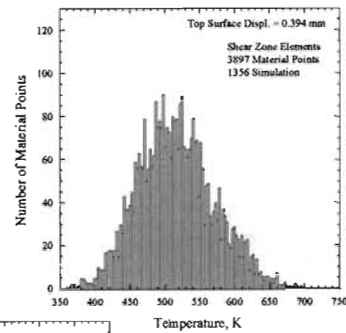


Localization - Ta

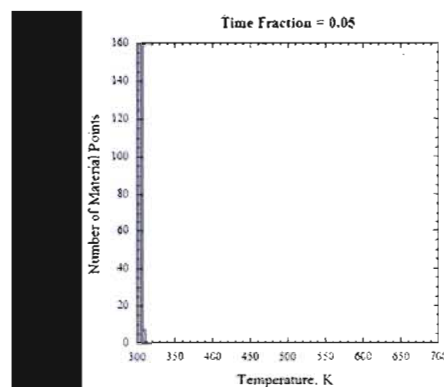


Realization 1

Realization 3



Realization 2



Realization 1

