

# Optimization-Based Modeling

## A new strategy for predictive simulations of multiscale, multiphysics problems

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# Outline

- **What is an optimization-based modeling (OBM)?**
- **A generic multiphysics OBM framework**
- **Application to Atomistic-to-Continuum (AtC) coupling**
  - Optimization-based formulation of AtC (OB-AtC)
  - Analysis of OB-AtC
  - Preliminary results

## Key Collaborators



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# A Metaphor

Suppose you want to solve these problems:

**P1 = Airplane**



**P2 = Car**



**There are many excellent solutions to each problem!**



# A Metaphor

Now, suppose you want to solve the **composite problem**  $P=P_1+P_2$ .

Even though there are far better solutions for  $P_1$  and  $P_2$

**How is this at all relevant to solving multiphysics problems?**

**As model complexity increases we encounter this scenario more and more often!**

most likely we will end up with this

**Lesson 1:** Good solutions are **easier to find** for **the constituent components**

**Lesson 2:** Merging them into a **good solution** for the composite problem is hard!



# Multiphysics & multiscale challenges

Typically  $L_1$  and  $L_2$  have different **mathematical structures**

$$(L_1 \& L_2)(u) = f$$

Goal: stable, accurate, discrete model,  
which preserves key physical properties

$$(L_1 \& L_2)^h(u^h) = f^h$$

$$L_1^h \quad L_2^h$$

**Stable compatible methods** may exist for  $L_1$   
and  $L_2$  but not for the composite problem:

$$(L_1 \& L_2)^h \neq L_1^h \& L_2^h$$

$$(L_1^h)^{-1} \quad (L_2^h)^{-1}$$

**Efficient solvers** may exist for  $L_1$  and  $L_2$   
but not for the composite problem

$$\left( (L_1 \& L_2)^h \right)^{-1} = ?$$

The dominant strategy for the past 30 years, based on first-order-accurate operator splitting and decoupled nonlinear methods, is approaching a **point of diminishing returns** because it

- 1) Lacks the stability properties for simulations over **dynamic scales of interest**,
- 2) Often **relies on heuristics** to control the splitting errors
- 3) Is prone to **non-intuitive instabilities** and fragile solutions



# Preservation of Physical Properties

## Challenges:

$$\partial_t u = Lu$$

Many **physical properties are not preserved** automatically under discretization, even with stabilization/regularization

$$\partial_t u^h = L^h u^h$$

$$\underline{C} \leq Cu \leq \bar{C}$$

In multiphysics codes the solution is **input** for another physics component

~~$$\underline{C} \leq Cu^h \leq \bar{C}$$~~

$$Bu = b$$

.....

**Automatic preservation** of maximum principle, local and global bounds, is required for robust, predictive simulations

~~$$Bu^h = b$$~~

.....

- **Preservation of structure:** “**easy**” - requires “**discrete vector calculus**”
- **Preservation of features:** “**hard**” - **burdens** discretization with tasks it is not well-equipped to handle and **imposes restrictions** on grid and variables



# Optimization-based modeling (OBM)

**Our strategy: use optimization and control ideas to**

- ☛ **Synthesize multiphysics models** from models for their constituent components
- ☛ **Build multiphysics solvers** from solvers for the constituent components
- ☛ **Preserve physical properties** that are difficult to impose directly in the discretization

## Potential payoffs

- ✓ **Mathematically rigorous** operator splitting and reconnection process through **reformulation into an equivalent optimization problem:**
  - Elimination or reduction of splitting errors
  - Precise stability and accuracy conditions obviate need for heuristics.
- ✓ **Relieve discretization** from difficult tasks such as preservation of physical properties
  - natural incorporation of physical properties as optimization constraints
- ✓ **Generality:** applicable to FE, FV and FD, particle methods, arbitrary grids...
- ✓ **Enable code reuse:** solvers, optimization tools,...



# OBM in a nutshell

## Generic optimization “harness”

$$\begin{array}{llll} & & \text{objective} & \rightarrow u \text{ (state)} \\ \text{minimize} & J(u) + \frac{\varepsilon}{2} \|\theta\|^2 & \leftarrow \text{regularization} & \rightarrow \theta \text{ (control)} \\ \text{subject to} & \begin{cases} L(u, f; \theta) = 0 \\ \underline{C} \leq Cu \leq \overline{C} \end{cases} & \begin{array}{l} \leftarrow \text{equality} \\ \leftarrow \text{inequality} \end{array} & \rightarrow \text{constraints} \end{array}$$

**Key question:** how to reformulate (map) a given problem into an optimization problem

**3 basic steps:**

- (1) identify **states & controls**;
- (2) identify **objective**;
- (3) identify **constraints**:

# Preservation of properties as an optimization problem: basic idea

Objective	Target	Constraints
minimize the distance between the solution and a suitable <i>target</i>	stable and accurate solution, not required to possess all desired physical properties	provided by the desired <b>physical properties</b>

$$\|u^h - u^\top\|$$

$$\partial_t u^\top = L^h u^\top$$

$$\begin{aligned} \underline{C} &\leq Cu^h \leq \bar{C} \\ Bu^h &= b \\ &\dots\dots\dots \end{aligned}$$

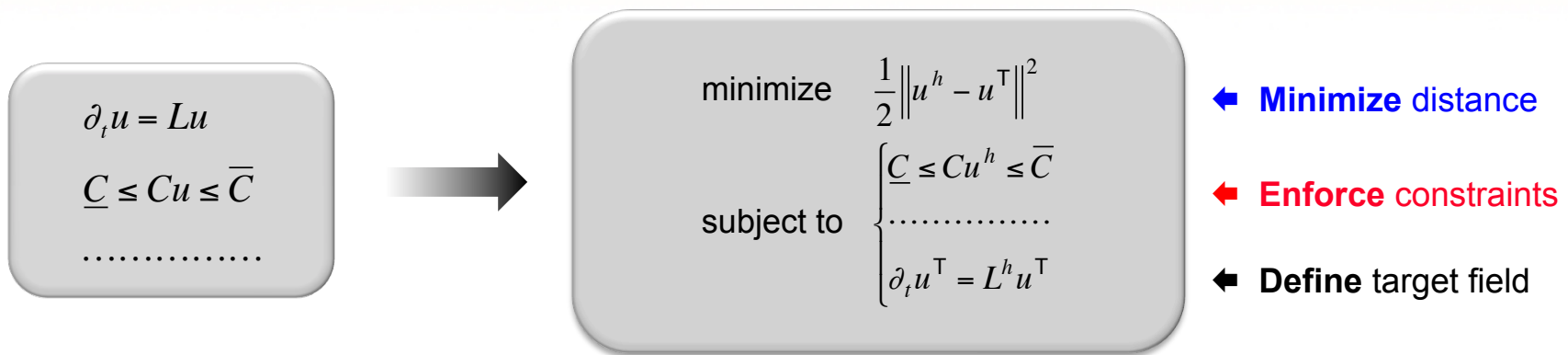
## Related work

- “*Optimization-based synchronous remap (FCR)*”, Liska et al JCP, 2010
- “*Optimization-based remap*”, Bochev, Ridzal, Scovazzi, Shashkov JCP, 2011, FCT Book 2012
- “*Optimization-based transport*”, Parts 1-3, Bochev, Peterson, Ridzal, Young, LNCS 2012
- “*Enforcing DMP for 2<sup>nd</sup> order elliptic problems*”, Liska, Shashkov, Comm. Comp. Phys, 2008
- “*Enforcement of constraints and DMP in VMS*”, Evans, Hughes, Sangali, CMAME 2009



# Preservation of properties as an optimization problem: advantages

## Generic formulation

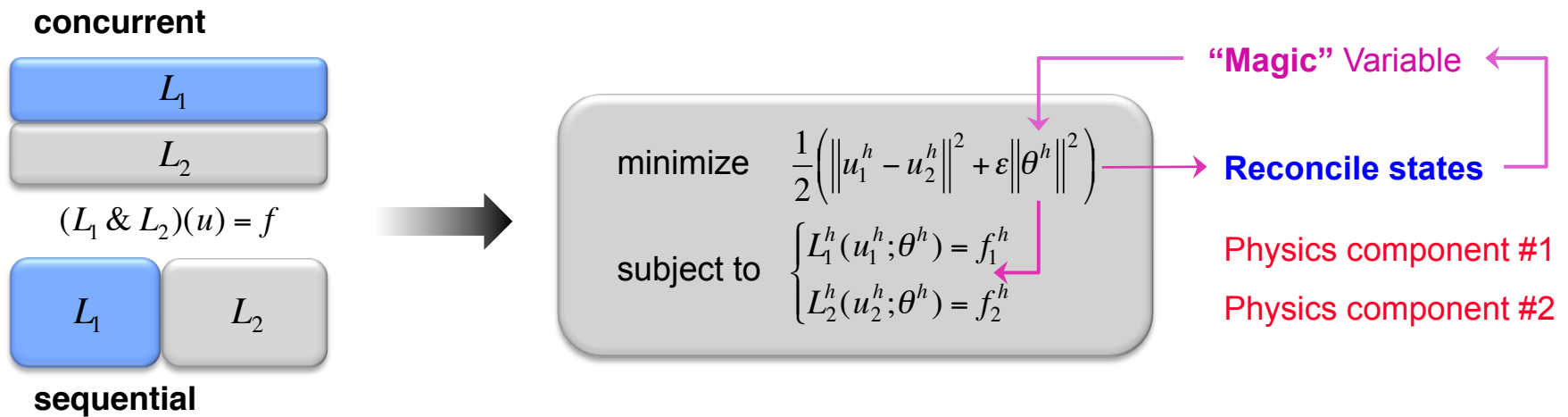


- ⇒ **Solution is a globally optimal state**: the best possible, relative to the target, state that has the desired physical properties
- ⇒ **Decoupling of the target** from the enforcement of the properties allows to adapt the formulation by choosing the most appropriate target and objective for a given problem
- ⇒ **Enforcement of properties** as constraints is impervious to grid structures and discrete field representations: enables feature preserving methods on arbitrary grids.



# Multi-physics coupling as an optimization problem

Objective	Constraints	Controls
<b>Minimize</b> error between the states of the constituent physics components	<b>Enforce</b> constituent physics components <b>independently</b>	<b>Enable</b> variability of the states, required to minimize the <b>objective</b>



- Sequential: “*Optimization-based domain decomposition*”, M. Gunzburger, 1997
- Concurrent: “*Decomposition of everything*”, J.L. Lyons, 2001, “*Additive operator-splitting*” Bochev&Ridzal, 2008-11



# An overview of AtC coupling

## AtC objectives

Combine the efficiency of continuum methods with the accuracy of (more expensive) atomistic models necessary to resolve local features such as cracks and dislocations.

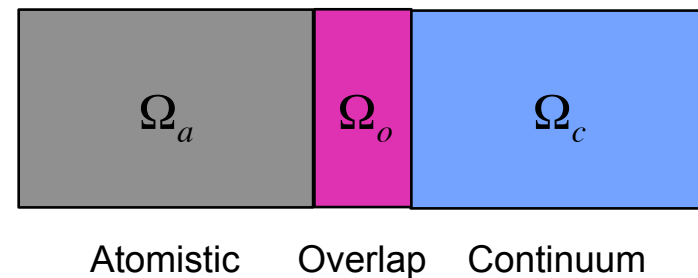
## Typical AtC configuration

$A$  – atomistic operator: valid in  $\Omega$

$C$  – continuum operator: valid in  $\Omega_o$  and  $\Omega_c$

$f_c$  – external force at continuum points

$f_a$  – external force at atom locations



## Key challenge:

- **How to merge A and C into a stable and accurate AtC formulation?**

## “Blending” is a dominant AtC strategy

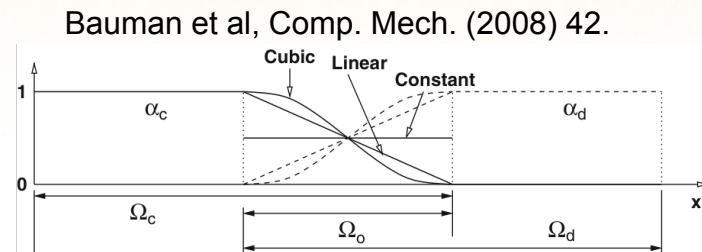
- **Energy:** blend atomistic and continuum energy functionals over  $\Omega_o$  (Arlequin, QC)
- **Force:** blend atomistic and continuum Euler-Lagrange equations over  $\Omega_o$



# Drawbacks of blending

## Typical energy-blended AtC method

$$\begin{cases} \text{minimize} & \alpha E_a(u_a) + (1 - \alpha) E_c(u_c) \\ \text{subject to} & \|\Pi u_a - u_c\|_{\Omega_o} = 0 \end{cases}$$



**Minimize blended energy subject to constraints forcing the equality of the atomistic and continuum states**

- Blending of two distinct physical models is inherently **ambiguous**
- The notion of “equality” depends on the projection  $\Pi$ , adding **more ambiguity**
- Resulting AtC formulations typically suffer from spurious effects, e.g., **ghost forces**

## An optimization-based AtC approach

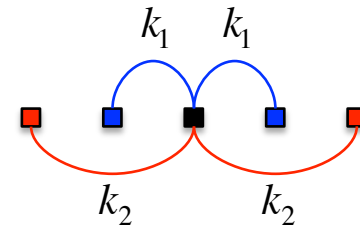
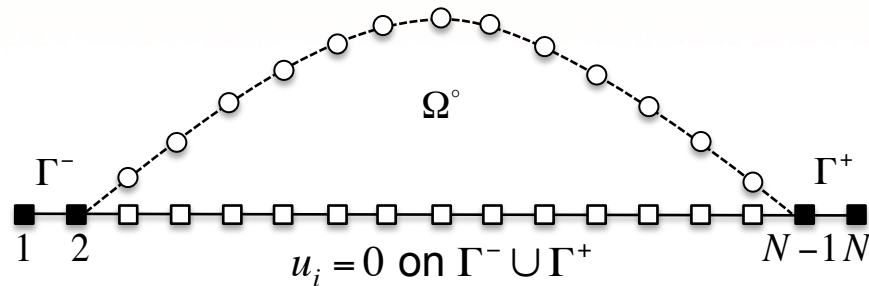
A **divide-and-conquer** strategy, which **completely separates** the atomistic and continuum models, and **reverses the roles** of constraints and objectives:

**Minimize the discrepancy between the atomistic and continuum states**  
subject to the **two models acting** independently in  $\Omega_a$  and  $\Omega_c$



# The global atomistic model

Linear mass-spring system with next nearest neighbor interactions



The total atomistic energy

$$E^a(u) = \sum_{i=0}^{N-1} \frac{k_1}{2} (u_{i+1} - u_i)^2 + \sum_{i=1}^{N-1} \frac{k_2}{2} (u_{i+1} - u_{i-1})^2 - (f, u)_{\ell^2}$$

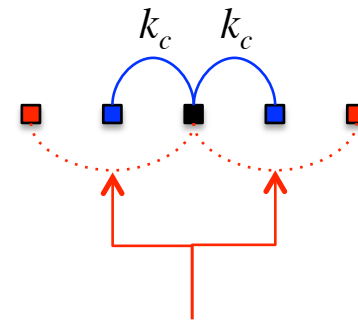
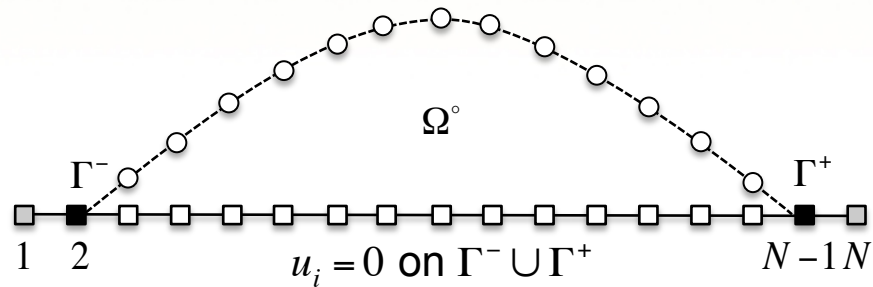
The atomistic problem

$$\tilde{u}^a = \arg \min E^a(u) \Rightarrow \partial E^a(\tilde{u}^a) = 0 \Rightarrow \begin{cases} A\tilde{u}^a = f & \text{in } \Omega^\circ \\ \tilde{u}^a = 0 & \text{on } \Gamma^- \cup \Gamma^+ \end{cases} \longrightarrow A = -(k_1\Delta_1 + k_2\Delta_2)$$



# The global continuum approximation

## Linear mass-spring system with nearest neighbor interactions



## The total continuum energy

Cauchy-Born rule:  $u_i \approx \frac{1}{2}(u_{i+1} + u_{i-1})$   $\Rightarrow (u_{i+1} - u_{i-1})^2 \approx 2(u_{i+1} - u_i)^2 + 2(u_i - u_{i-1})^2$

$E^c(u) = \sum_{i=0}^{N-1} \frac{k_c}{2} (u_{i+1} - u_i)^2 - (f, u)_{\ell^2(\Omega)}$   $\Rightarrow k_c = k_1 + 4k_2$

## The continuum problem

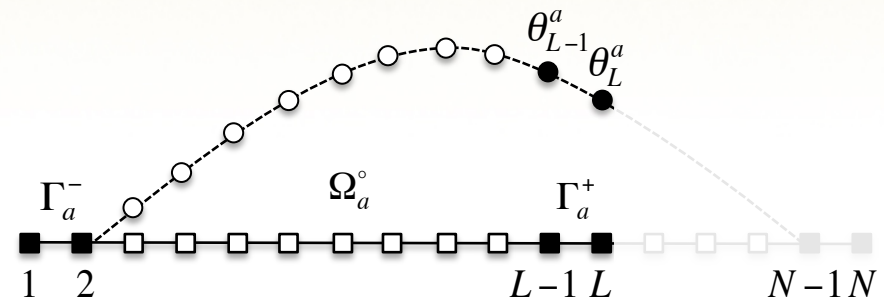
$\tilde{u}^c = \arg \min E^c(u) \Rightarrow \partial E^c(\tilde{u}^c) = 0 \Rightarrow \begin{cases} C\tilde{u}^c = f & \text{in } \Omega^\circ \\ \tilde{u}^c = 0 & \text{on } \Gamma^- \cup \Gamma^+ \end{cases} \Rightarrow C = -(k_c \Delta_1)$



# Optimization-based AtC. Step 1: splitting

## The atomistic subdomain problem

$$\begin{cases} Au^a = f^a & \text{in } \Omega_a^\circ \\ u^a = \theta^a & \text{on } \Gamma_a^+ \\ u^a = 0 & \text{on } \Gamma_a^- \end{cases}$$

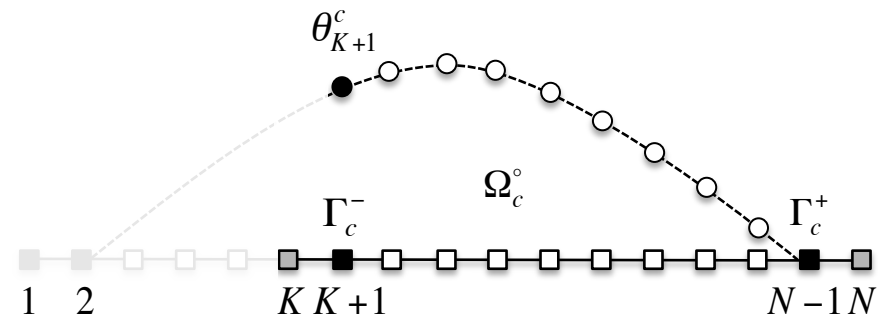


$\Gamma_a^+$  – an artificial atomistic domain boundary due to splitting

$\Theta^a$  – unknown boundary value on  $\Gamma_a^+$

## The continuum subdomain problem

$$\begin{cases} Cu^c = f^c & \text{in } \Omega_c^\circ \\ u^c = \theta^c & \text{on } \Gamma_c^- \\ u^c = 0 & \text{on } \Gamma_c^+ \end{cases}$$



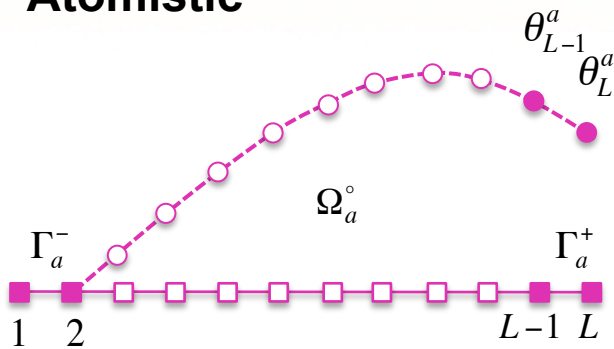
$\Gamma_c^-$  – an artificial atomistic domain boundary due to splitting

$\Theta^c$  – unknown boundary value on  $\Gamma_c^-$



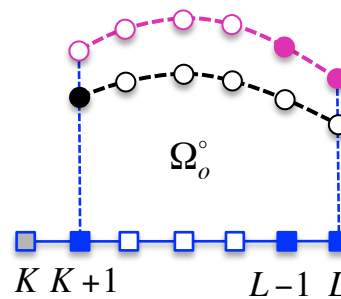
# Optimization-based AtC. Step 2: reconnection

## Atomistic



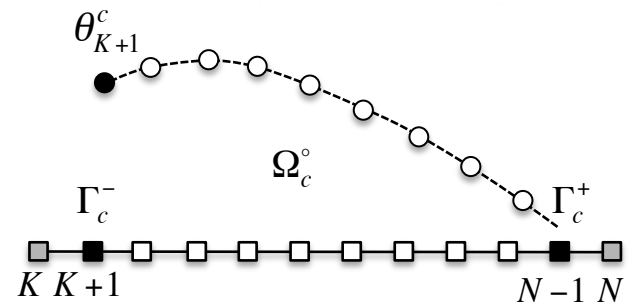
$$\begin{cases} Au^a = f^a & \text{in } \Omega_a^\circ \\ u^a = \theta^a & \text{on } \Gamma_a^+ \\ u^a = 0 & \text{on } \Gamma_a^- \end{cases}$$

## Overlap



$$\underbrace{\|\theta^c - u^a\|_{\Gamma_c^-}^2 + \|u^a - u^c\|_{\Omega_o^\circ}^2 + \|\theta^a - u^c\|_{\Gamma_a^+}^2}_{\|u^a - u^c\|_{\Omega_o}^2}$$

## Continuum



$$\begin{cases} Cu^c = f^c & \text{in } \Omega_c^\circ \\ u^c = \theta^c & \text{on } \Gamma_c^- \\ u^c = 0 & \text{on } \Gamma_c^+ \end{cases}$$

- Atomistic and continuum problems are **separated** rather than blended
- Their solutions communicate through the **artificial boundary conditions**
- The overlap terms measure an **artificial “mismatch” energy**



# Optimization-based AtC formulation (OB-AtC)

$$\text{minimize } \frac{1}{2} \left( \|u^a - u^c\|_{\Omega_o}^2 + \|\theta^a - u^c\|_{\Gamma_a^+}^2 + \|\theta^c - u^a\|_{\Gamma_c^-}^2 \right)$$

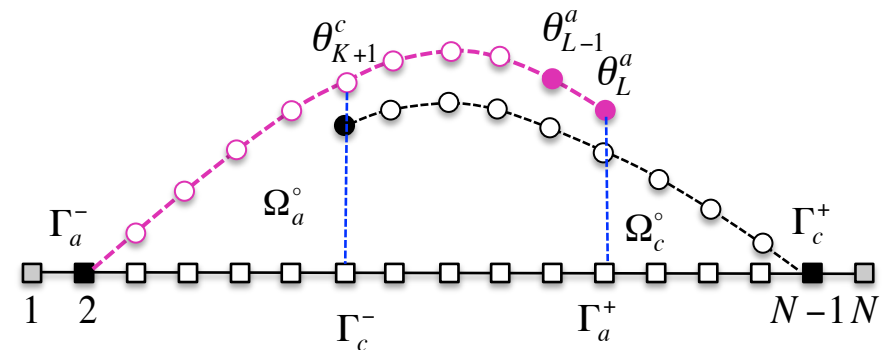
$$\text{subject to } \begin{cases} Au^a = f^a & \text{in } \Omega_a^\circ \\ u^a = \theta^a & \text{on } \Gamma_a^+ \\ u^a = 0 & \text{on } \Gamma_a^- \end{cases} \text{ and } \begin{cases} Cu^c = f^c & \text{in } \Omega_c^\circ \\ u^c = \theta^c & \text{on } \Gamma_c^- \\ u^c = 0 & \text{on } \Gamma_c^+ \end{cases}$$

Minimize the artificial mismatch energy subject to the atomistic and continuum force balance equation

**Objective:** measures the mismatch energy

**States:** atomistic and continuum solutions

**Controls:** artificial boundary conditions



Atomistic and continuum problems are patch test consistent individually

⇒ **OB-AtC passes patch test by construction**



# OB-AtC is well-posed

## Reduced space problem

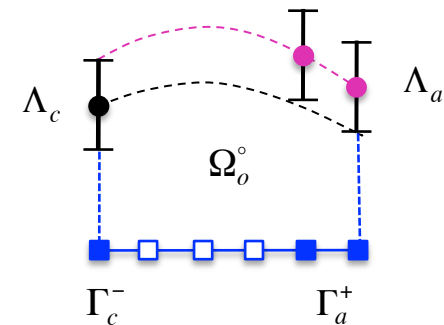
We switch to an equivalent **reduced space problem** in terms of the virtual controls only

$$\begin{cases} Au^a = f^a & \text{in } \Omega_a^\circ \\ u^a = \theta^a & \text{on } \Gamma_a^+ \\ u^a = 0 & \text{on } \Gamma_a^- \end{cases} \rightarrow u^a = u^a(\theta^a)$$

$$\begin{cases} Cu^c = f^c & \text{in } \Omega_c^\circ \\ u^c = \theta^c & \text{on } \Gamma_c^- \\ u^c = 0 & \text{on } \Gamma_c^+ \end{cases} \rightarrow u^c = u^c(\theta^c)$$

$$\underset{\Lambda_a \times \Lambda_c}{\text{minimize}} \frac{1}{2} \|u^a(\theta^a) - u^c(\theta^c)\|_{\Omega_o}^2$$

Unconstrained minimization problem over the atomistic and continuum trace spaces  $\Lambda_a$  and  $\Lambda_c$

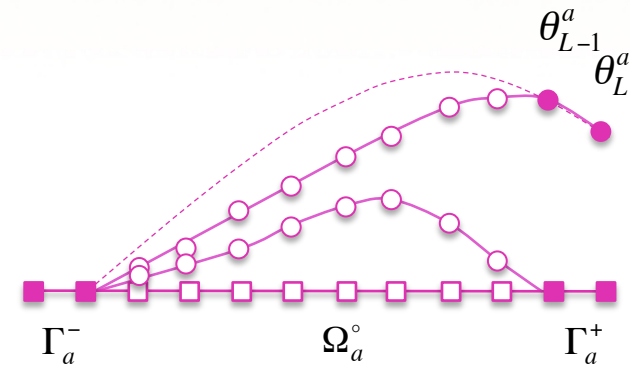


# Existence and uniqueness of optimal solutions

## Decomposition of the states

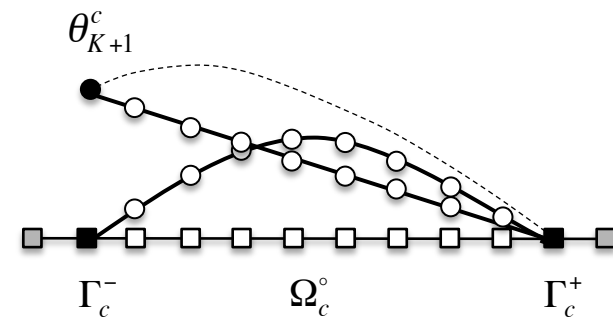
$$u^a(\theta^a) = v^a(\theta^a) + u^{a,0} \longrightarrow \begin{cases} Av^a = 0 & \text{in } \Omega_a^\circ \\ v^a = \theta^a & \text{on } \Gamma_a^+ \end{cases}$$

$$\begin{cases} Au^{a,0} = f^a & \text{in } \Omega_a^\circ \\ u^{a,0} = 0 & \text{on } \Gamma_a \end{cases}$$



$$u^c(\theta^c) = v^c(\theta^c) + u^{c,0} \longrightarrow \begin{cases} Cv^c = 0 & \text{in } \Omega_c^\circ \\ v^c = \theta^c & \text{on } \Gamma_c^- \end{cases}$$

$$\begin{cases} Cu^{c,0} = f^c & \text{in } \Omega_c^\circ \\ u^{c,0} = 0 & \text{on } \Gamma_c \end{cases}$$



This decomposition follows the idea of Gervasio et al., Num. Math. 2001



# The dual inner product and norm

**Theorem.** The following form defines an inner product on the trace space  $\Lambda_a \times \Lambda_c$ .

$$\langle \{\theta^a, \theta^c\}, \{\mu^a, \mu^c\} \rangle := \left( v^a(\theta^a) - v^c(\theta^c), v^a(\mu^a) - v^c(\mu^c) \right)_{\ell^2(\Omega_o)}$$

**Reduced space problem in terms of the dual norm**

$$\underset{\Lambda_a \times \Lambda_c}{\text{minimize}} \frac{1}{2} \left\| \{\theta^a, \theta^c\} \right\|_{\ell^*}^2 + \left( v^a(\theta^a) - v^c(\theta^c), u^{a,0} - u^{c,0} \right)_{\ell^2(\Omega_o)} + \frac{1}{2} \left\| u^{a,0} - u^{c,0} \right\|_{\ell^2(\Omega_o)}$$

**The Euler-Lagrange equation**

$$\langle \{\theta^a, \theta^c\}, \{\mu^a, \mu^c\} \rangle := - \left( u^{a,0} - u^{c,0}, v^a(\mu^a) - v^c(\mu^c) \right)_{\ell^2(\Omega_o)} \quad \forall \{\mu^a, \mu^c\} \in \Lambda_a \times \Lambda_c$$

**Corollary.** The reduced space problem has a unique solution  $\{\theta^a, \theta^c\} \in \Lambda_a \times \Lambda_c$ .



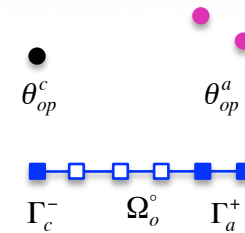
# The OB-AtC approximation

## Solution procedure

### 1. Solve the reduced space problem:

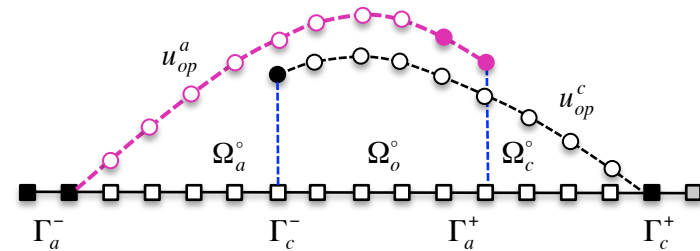
$$\{\theta_{op}^a, \theta_{op}^c\} \in \Lambda_a \times \Lambda_c$$

approximate solution traces  $\rightarrow$



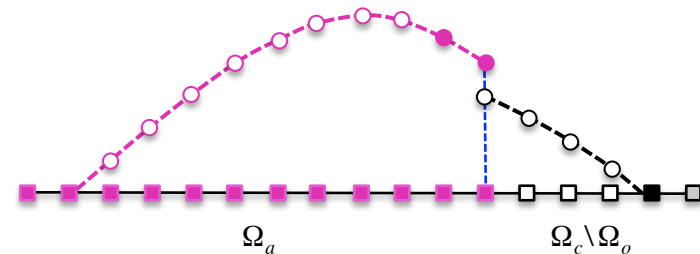
### 2. Recover the optimal states

$$u_{op}^a = u^{a,0} + v^a(\theta_{op}^a) \quad \text{and} \quad u_{op}^c = u^{c,0} + v^c(\theta_{op}^c)$$



### 3. Define the AtC approximation

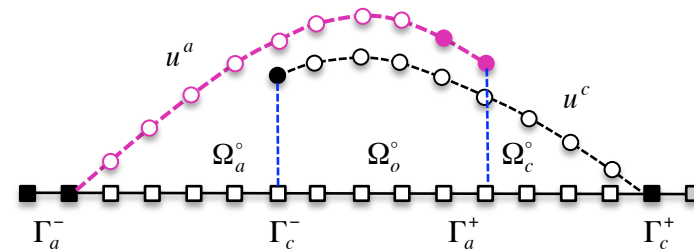
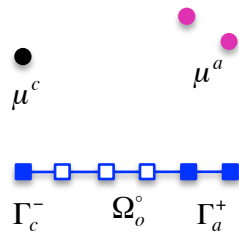
$$u^{atc} = \begin{cases} u_{op}^a & \text{in } \Omega_a \\ u_{op}^c & \text{in } \Omega_c \setminus \Omega_o \end{cases}$$



# The AtC approximation error

Define the trace lifting operator

$$P(\{\mu^a, \mu^c\}) = U_0 + Q(\{\mu^a, \mu^c\}) \quad \text{where} \quad U_0 = \begin{pmatrix} u^{a,0} \\ u^{c,0} \end{pmatrix} \quad \text{and} \quad Q(\{\mu^a, \mu^c\}) = \begin{cases} v^a(\mu^a) & \text{in } \Omega_a \\ v^c(\mu^c) & \text{in } \Omega_c \setminus \Omega_o \end{cases}$$



The AtC approximation error

$$u^{atc} = P(\{\theta_{op}^a, \theta_{op}^c\}) \quad \longrightarrow \quad \left\| \tilde{u}^a - u^{atc} \right\|_{\ell^2(\Omega)} = \left\| \tilde{u}^a - P(\{\theta_{op}^a, \theta_{op}^c\}) \right\|_{\ell^2(\Omega)}$$

Where  $\tilde{u}^a$  is the global atomistic solution.



# The AtC error estimate

Define the trace operator  $r(u) = \left\{ \begin{pmatrix} u_{L-1} \\ u_L \end{pmatrix}; (u_K) \right\} = \{r(u^a), r(u^c)\}$

## Error splitting

$$\begin{aligned} \|\tilde{u}^a - u^{atc}\|_{\ell^2(\Omega)} &= \left\| \tilde{u}^a - P\left(\{\theta_{op}^a, \theta_{op}^c\}\right) \right\|_{\ell^2(\Omega)} = \left\| \tilde{u}^a - P\left(r(\tilde{u}^a)\right) + P\left(r(\tilde{u}^a)\right) - P\left(\{\theta_{op}^a, \theta_{op}^c\}\right) \right\|_{\ell^2(\Omega)} \\ &= \left\| \tilde{u}^a - P\left(r(\tilde{u}^a)\right) \right\|_{\ell^2(\Omega)} + \left\| Q\left(r(\tilde{u}^a)\right) - Q\left(\{\theta_{op}^a, \theta_{op}^c\}\right) \right\|_{\ell^2(\Omega)} \\ &\leq \left\| \tilde{u}^a - P\left(r(\tilde{u}^a)\right) \right\|_{\ell^2(\Omega)} + \|Q\| \cdot \left\| r(\tilde{u}^a) - \{\theta_{op}^a, \theta_{op}^c\} \right\|_{\ell^*} \end{aligned}$$

## The error components

$$\left\| \tilde{u}^a - P\left(r(\tilde{u}^a)\right) \right\|_{\ell^2(\Omega)}$$

Consistency error

$$\left\| r(\tilde{u}^a) - \{\theta_{op}^a, \theta_{op}^c\} \right\|_{\ell^*}$$

Approximation error

$$\|Q\| = \sup_{\{\mu^a, \mu^c\}} \frac{\|Q\{\mu^a, \mu^c\}\|_{\ell^2(\Omega)}}{\|\{\mu^a, \mu^c\}\|_{\ell^*}}$$

Lifting operator norm



# Estimate of the consistency error

## Assumptions

$|\Omega_a| \leq |\Omega|^{1/p}$  for some  $0 < p$  ← The atomistic domain  $\Omega_a$  is **small** compared to  $\Omega$

$|\Omega_o| = (1 - \gamma)|\Omega_a|$  for some  $0 < \gamma < 1$  ← The overlap domain  $\Omega_o$  remains **proportional** to  $\Omega_a$

**Lemma.** Let  $\tilde{u}^a$  be the exact atomistic solution and  $u^c$  - the continuum lifting of its trace

$$\left\| \tilde{u}^a - P(r(\tilde{u}^a)) \right\|_{\ell^2(\Omega)} \leq \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c \setminus \Omega_o)}$$

*Proof.*

$$P(r(\tilde{u}^a)) = P\left(\left\{r^a(\tilde{u}^a), r^c(\tilde{u}^a)\right\}\right) = \begin{cases} \tilde{u}^a & \text{in } \Omega_a \\ u^c & \text{in } \Omega_c \setminus \Omega_o \end{cases}$$

← Recovers the exact atomistic solution in  $\Omega_a$   
← Continuum lifting of the exact atomistic trace

$$\left\| \tilde{u}^a - P(r(\tilde{u}^a)) \right\|_{\ell^2(\Omega)} \leq \left\| \tilde{u}^a - \tilde{u}^a \right\|_{\ell^2(\Omega_a)} + \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c \setminus \Omega_o)} = \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c \setminus \Omega_o)}$$

**Consistency error is confined to the pure continuum region**



# Estimate of the approximation error

**Lemma.** Let  $\tilde{u}^a$  be the exact atomistic solution and  $u^c$ -the continuum lifting of its trace

$$\left\| r(\tilde{u}^a) - \{\theta_{op}^a, \theta_{op}^c\} \right\|_{\ell^*} \leq \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_o)}$$

*Proof.* Using that  $\{\theta_{op}^a, \theta_{op}^c\}$  solves the Euler-Lagrange equation of the reduced problem

$$\left\| r(\tilde{u}^a) - \{\theta_{op}^a, \theta_{op}^c\} \right\|_{\ell^*} = \sup_{\{\mu^a, \mu^c\}} \frac{\left\langle r(\tilde{u}^a), \{\mu^a, \mu^c\} \right\rangle + \left( u^{a,0} - u^{c,0}, v^a(\mu^a) - v^a(\mu^a) \right)_{\ell^2(\Omega_o)}}{\left\| \{\mu^a, \mu^c\} \right\|_{\ell^*}}$$

$$\left\langle r(\tilde{u}^a), \{\mu^a, \mu^c\} \right\rangle + \left( u^{a,0} - u^{c,0}, v^a(\mu^a) - v^a(\mu^a) \right)_{\ell^2(\Omega_o)}$$

**Using definition of the dual inner product**

$$= \left( \tilde{u}^a - u^c, v^a(\mu^a) - v^a(\mu^a) \right)_{\ell^2(\Omega_o)}$$

$$\leq \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_o)} \left\| v^a(\mu^a) - v^a(\mu^a) \right\|_{\ell^2(\Omega_o)} \leq \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_o)} \left\| \{\mu^a, \mu^c\} \right\|_{\ell^*}$$

**Approximation error is confined to the overlap region**



# Estimate of the operator norm

**Lemma.** The norm of  $Q$  is inversely proportional to the size of the overlap domain:

$$\|Q\| \leq \frac{C}{1-\gamma} \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2}$$

*Proof.* Definition of the operator norm implies that the statement is equivalent to

$$\|Q\{\mu^a, \mu^c\}\|_{\ell^2(\Omega)} \leq \frac{C}{1-\gamma} \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2} \|\{\mu^a, \mu^c\}\|_{\ell^*} \quad \forall \{\mu^a, \mu^c\} \in \Lambda_a \times \Lambda_c$$

Which in turn is equivalent to

$$\|v^a(\mu^a)\|_{\ell^2(\Omega_a)}^2 + \|v^c(\mu^c)\|_{\ell^2(\Omega_c \setminus \Omega_o)}^2 \leq \frac{C}{1-\gamma} \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2} \|v^a(\mu^a) - v^c(\mu^c)\|_{\ell^2(\Omega_o)}^2 \quad \forall \{\mu^a, \mu^c\} \in \Lambda_a \times \Lambda_c$$

We use the structure of the atomistic and continuum solutions to prove this bound.

Equivalent to the existence of a **strong Cauchy-Schwartz** inequality for  $(v^a(\mu^a), v^c(\mu^c))$ .



# The AtC error estimate

**Lemma.** Let  $\tilde{u}^a$  be the exact atomistic solution and  $u^c$ -the continuum lifting of its trace

$$\left\| \tilde{u}^a - u^{atc} \right\|_{\ell^2(\Omega)} \leq C \left( 1 + \frac{1}{1-\gamma} \right) \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2} \left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c)}$$

**AtC accuracy is determined by 2 independent factors:**

**Continuum modeling error**

$$\left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c)}$$

- Small if solution is smooth in  $\Omega_c$  i.e., there are no defects (basic premise of AtC)
- Independent of the choice of coupling mechanism

**Optimization based coupling**

$$\left( 1 + \frac{1}{1-\gamma} \right) \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2}$$

- Inversely proportional to  $\Omega_c$  i.e., better for larger overlap regions
- Independent of the continuum modeling error



# The AtC error estimate

**Theorem.** Let  $\tilde{u}^a$  be the exact atomistic solution and  $u^c$  -the continuum lifting of its trace

$$\left\| \tilde{u}^a - u^{atc} \right\|_{\ell^2(\Omega)} \leq C \left( 1 + \frac{1}{1-\gamma} \right) \cdot \left( \frac{|\Omega|}{|\Omega_o|} \right)^{1/2} |\Omega_c| \cdot \left\| \Delta_1^2 \tilde{u}^a \right\|_{\ell^2(\Omega_c)}$$

*Proof.* Follows from the upper bound on the modeling error

$$\left\| \tilde{u}^a - u^c \right\|_{\ell^2(\Omega_c)} \leq |\Omega_c| \cdot \left\| \Delta_1^2 \tilde{u}^a \right\|_{\ell^2(\Omega_c)}$$

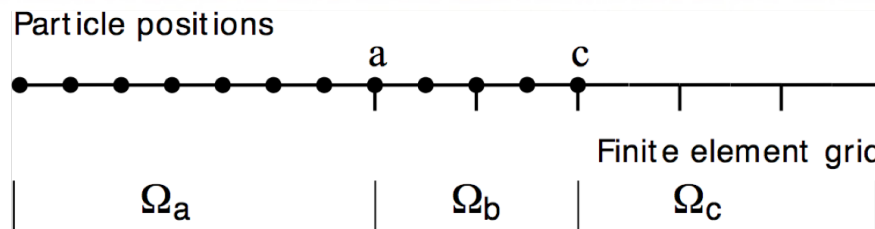
**Corollary.** Set  $\varepsilon=1/N$  and rescale the lattice according to  $i \rightarrow \varepsilon i$ . The AtC error then is:

$$\left\| \tilde{u}^a - u^{atc} \right\|_{\ell^2(\Omega)_\varepsilon} \leq C \frac{\varepsilon^2}{|\Omega_o|^{1/2}} \cdot \left\| \Delta_1^2 \tilde{u}^a \right\|_{\ell^2(\Omega_c)_\varepsilon}$$



# Numerical results

## Model one-dimensional AtC problem (Curtin, Miller, 2003)



$$\Omega_a = [0, 74]$$

$$\Omega_c = [30, 105]$$

$$\Omega_b = [30, 74]$$

A – mass spring network: next nearest neighbor interaction

$$Q_a = -k_1(u_{i-1} - 2u_i + u_{i+1}) - k_2(u_{i-2} - 2u_i + u_{i+2})$$

C – linear elasticity with averaged spring constant

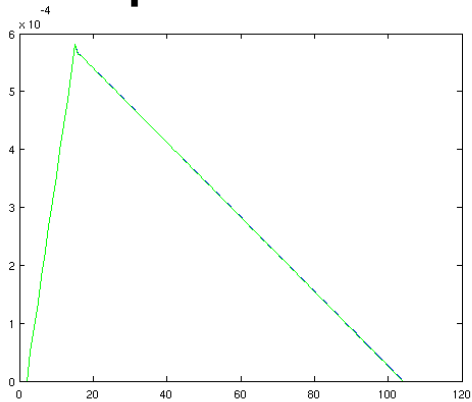
$$Q_c = -k_c(u_{i-1} - 2u_i + u_{i+1}) \quad k_c = k_1 + 4k_2$$

$f_a$  – point force at lattice location #25

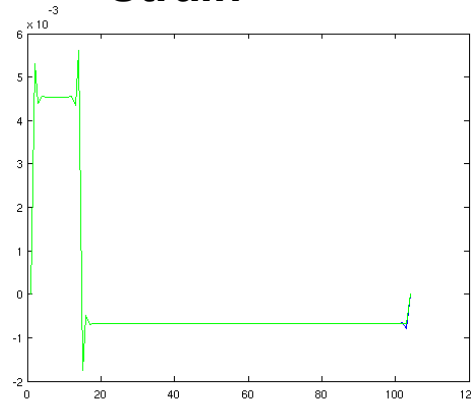


# Numerical results

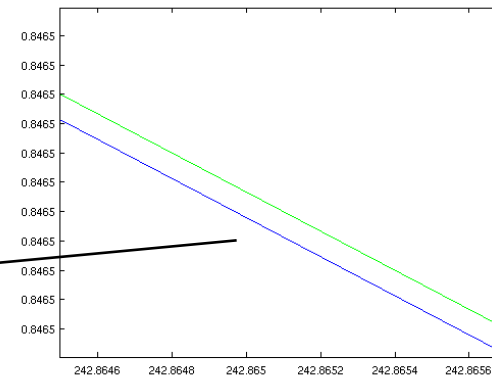
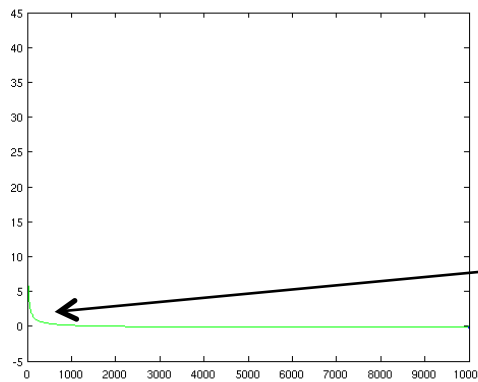
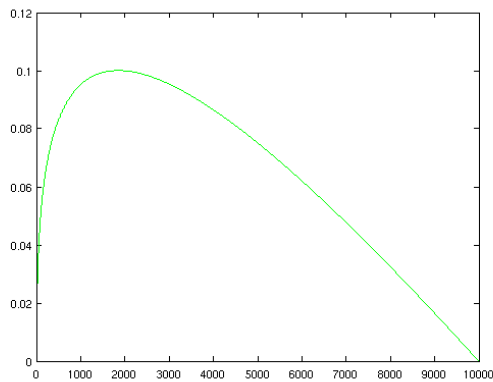
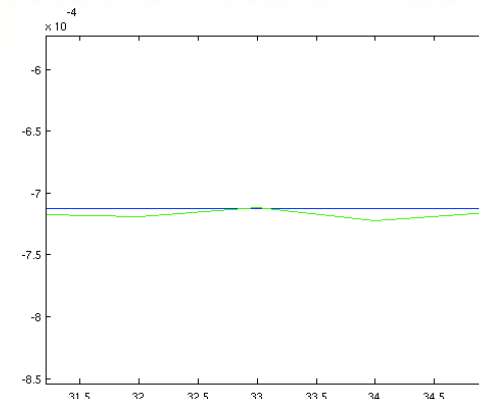
## Displacement



## Strain



## Strain in the overlap region

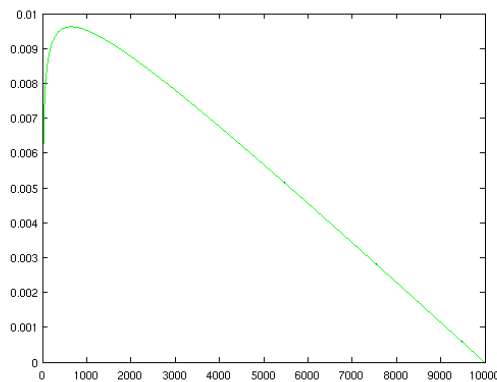


**Blue** = exact solution  
**Green** = OBM solution

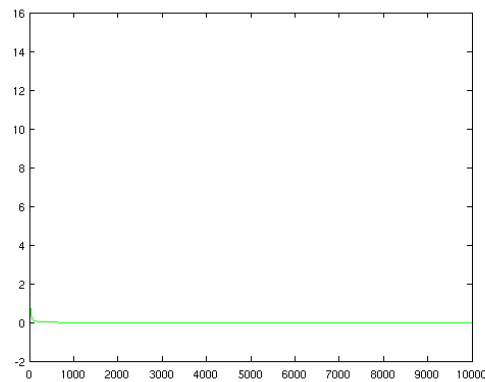


# Numerical results

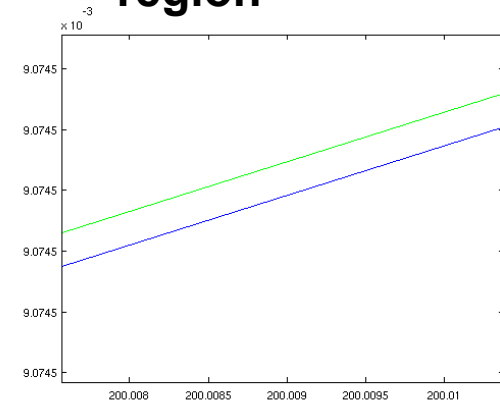
## Displacement



## Strain



## Displacement in the overlap region



**Blue** = exact solution  
**Green** = OBM solution



# Summary

We use optimization and control ideas to **automatically manage discretization tasks** that are **difficult to accomplish directly**:

- Reconcile solutions from constituent components of multi-physics problems
- Reconcile different representations of the same field in physically consistent manner
- Automatically figure out the needed “ghost forces” to balance the interfaces

## Advantages

- **Increase concurrency** by exposing constituent physics components
- **Remove order limitations** (reformulation yields **nearly equivalent** problems)
- **Rigorous mathematical foundations** inherited from rich optimization theory
- **Restore** physical constraints that were lost during discretization
- **Reuse of software** components through “synthesis” of solvers and discretizations

