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Title: Multidimensional Discrete Ordinates Transport on Massively Parallel Architectures

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Multidimensional Discrete Ordinates Transport on Massively Parallel Architectures

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NC State University

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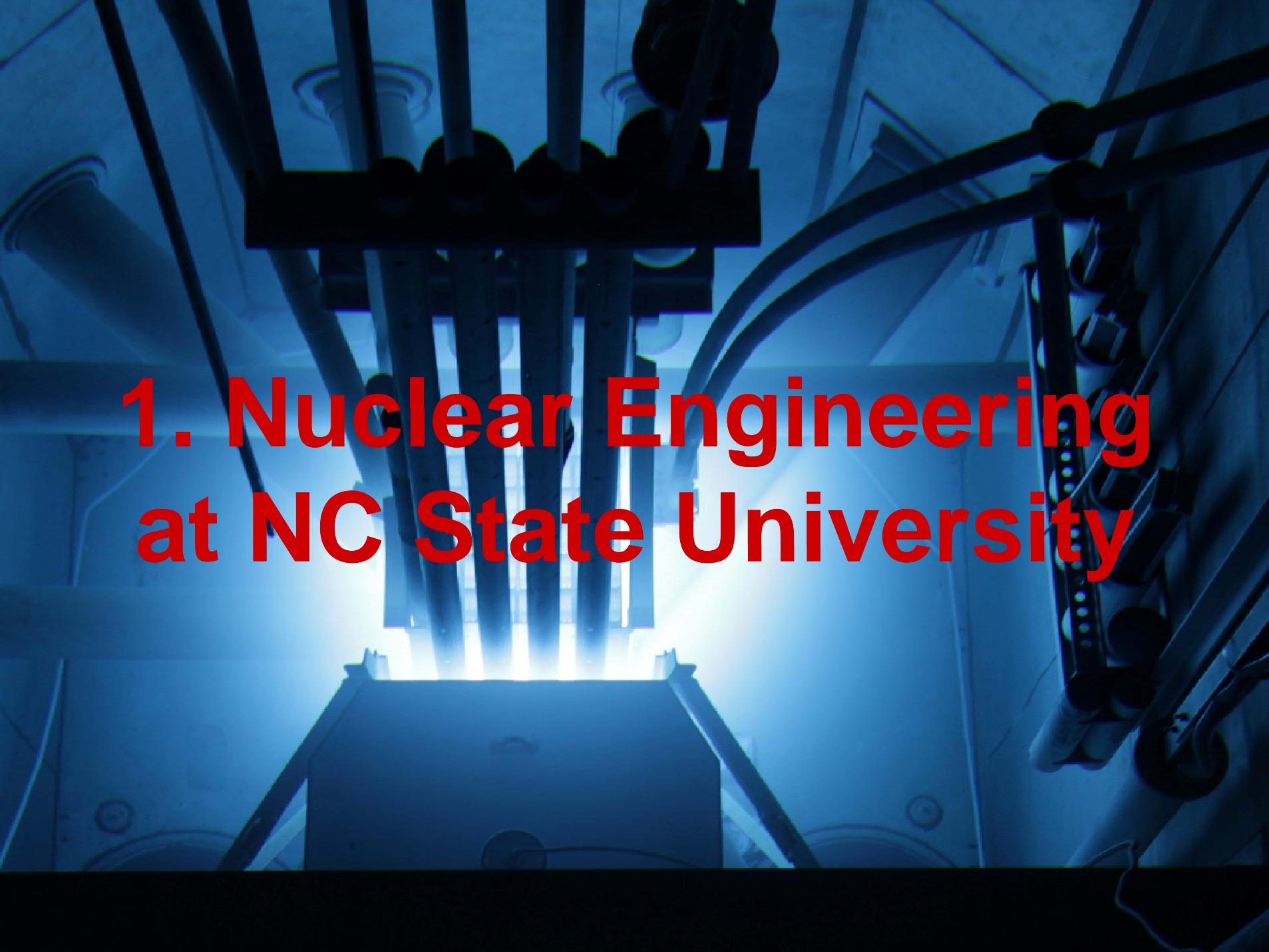
Research Staff
Los Alamos National Laboratory



0. Outline

- 1. Nuclear Engineering at NC State University**
- 2. Nuclear Computational Science Group**
- 3. Transport Theory in a Nutshell**
- 4. Solution of the Transport Equation**
- 5. Alternative SDD Parallelizations**
- 6. Measured Parallel Performance**
- 7. Conclusions**





1. Nuclear Engineering at NC State University

1. Department's Brief History

- 1950** Established as graduate program in Physics Dept
- ~1950** First non-governmental university-based research reactor
- 1955** Two PhDs awarded
- 1962** Department of Nuclear Engineering established
- 1965** Rapid growth from 4 to 9 faculty; thrust areas: (1) Fission power reactors; (2) Radiation applications
- 1973** 1MW PULSTAR operational (4th on-campus reactor)
- 1983** Added Plasma/fusion graduate track
- 1994** Combined five-year BS/MNE degree established
- 2008** *Master of Nuclear Engineering* degree via Distance Ed



1. NCSU's Nuclear Engineering Today

□ Our Faculty:

- ❖ 8 active faculty in 2007 ⇒ 15 today
- ❖ 2 open positions currently in search
- ❖ 2 endowed chairs (*Progress Energy & Duke Energy*): PE Chair in search
- ❖ Multiple Joint Faculty Appointments with ORNL and INL
- ❖ Pivotal role in CASL: Turinsky Chief Scientist, Doster Ed Programs
- ❖ Gilligan: Director of NEUP

□ Our Students:

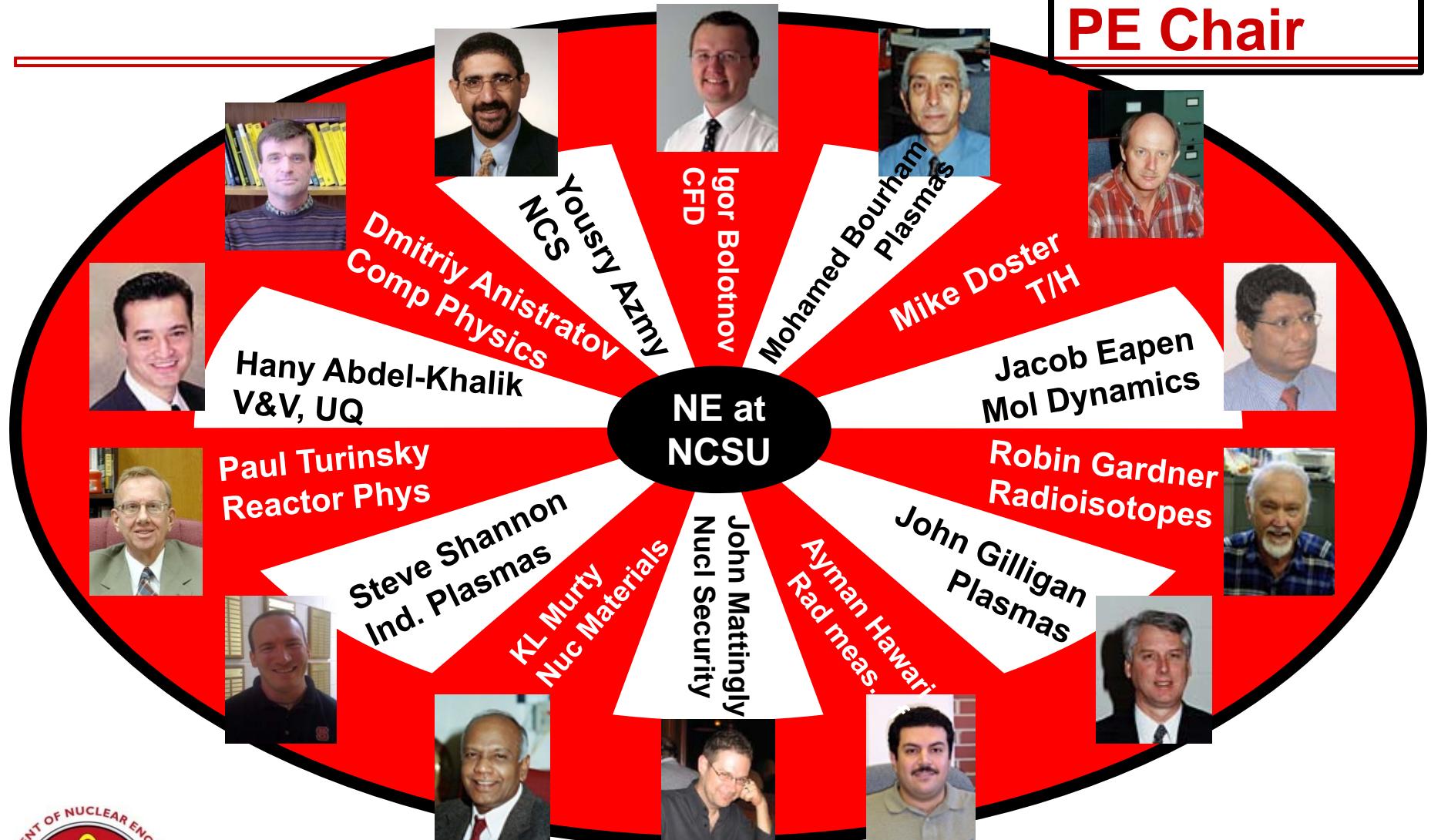
- ❖ Enrolments surpassed 200 UGs & 100 Grads
- ❖ Won Mark Mills Award (best PhD) 9 times in Award's 53 years
- ❖ ~10% win one or more award, scholarship, or fellowship

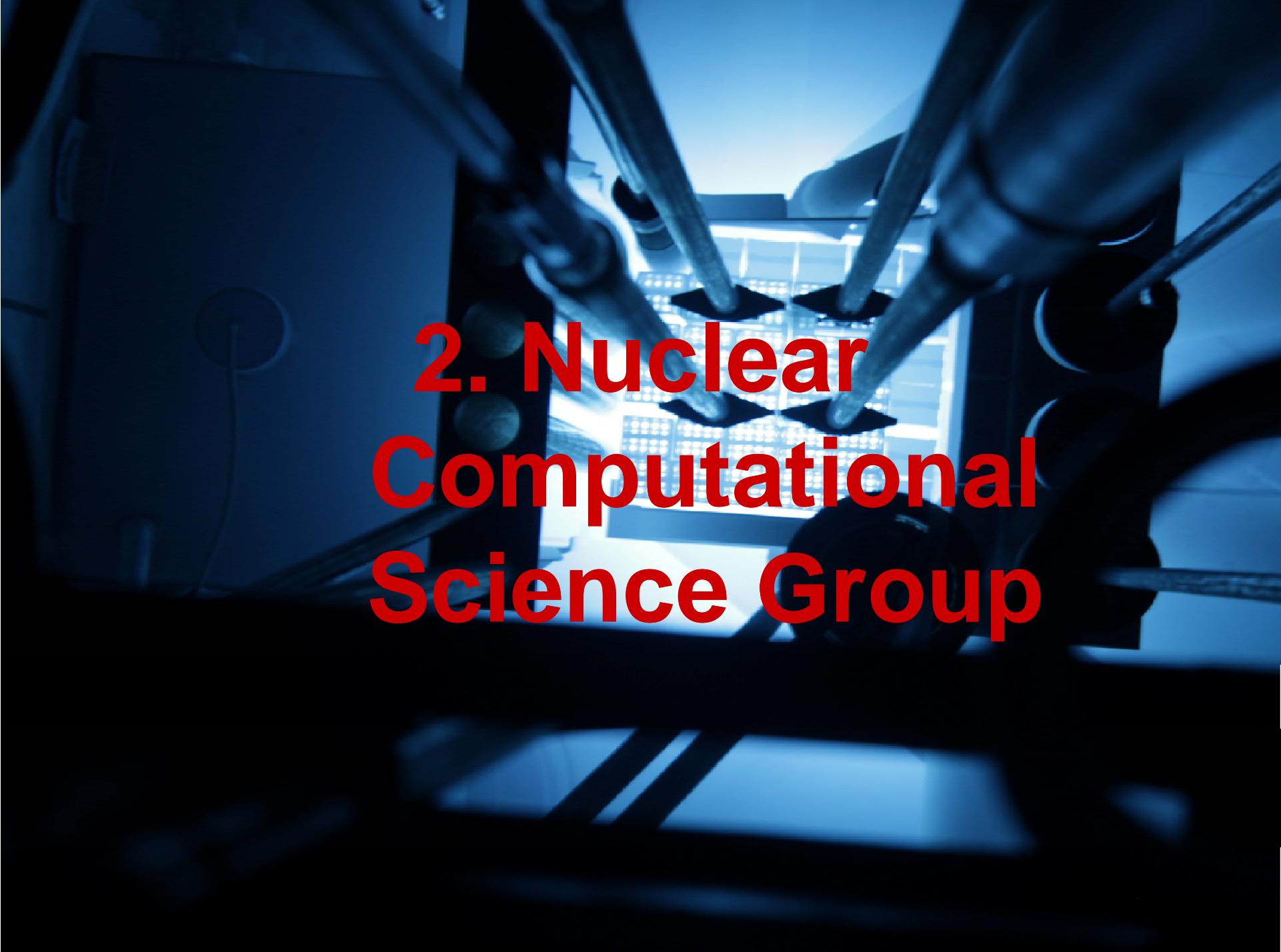
□ Space:

- ❖ Increased by more than 50% since 2008
- ❖ Future move to new building on Centennial Campus



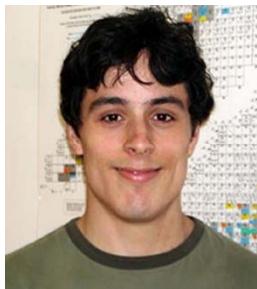
Add Nam &
PE Chair





2. Nuclear Computational Science Group

2. NCSG at NCSU



Sebastian Schunert: PhD Student

Thesis: Comparing Various Spatial Discretization Schemes Based on a Method of Manufactured Solutions Benchmark Suite

Sean O'Brien: PhD Student

MS Thesis: A Posteriori Error Estimators for the Discrete Ordinates Approximation of the One-Speed Neutron Transport Equation



Sameer Vhora: MNE Student

Thesis: Spectral Analysis of Parallel Block Jacobi Iterations for Solving the Discrete Ordinates Equations with the ITMM Approach



2. NCSG at NCSU



Brian Powell: MS Student

Thesis: *Efficient Computation of Subdomain Operators Employed in the Integral Transport Matrix Method (ITMM)*

Noel Nelson: MS Student

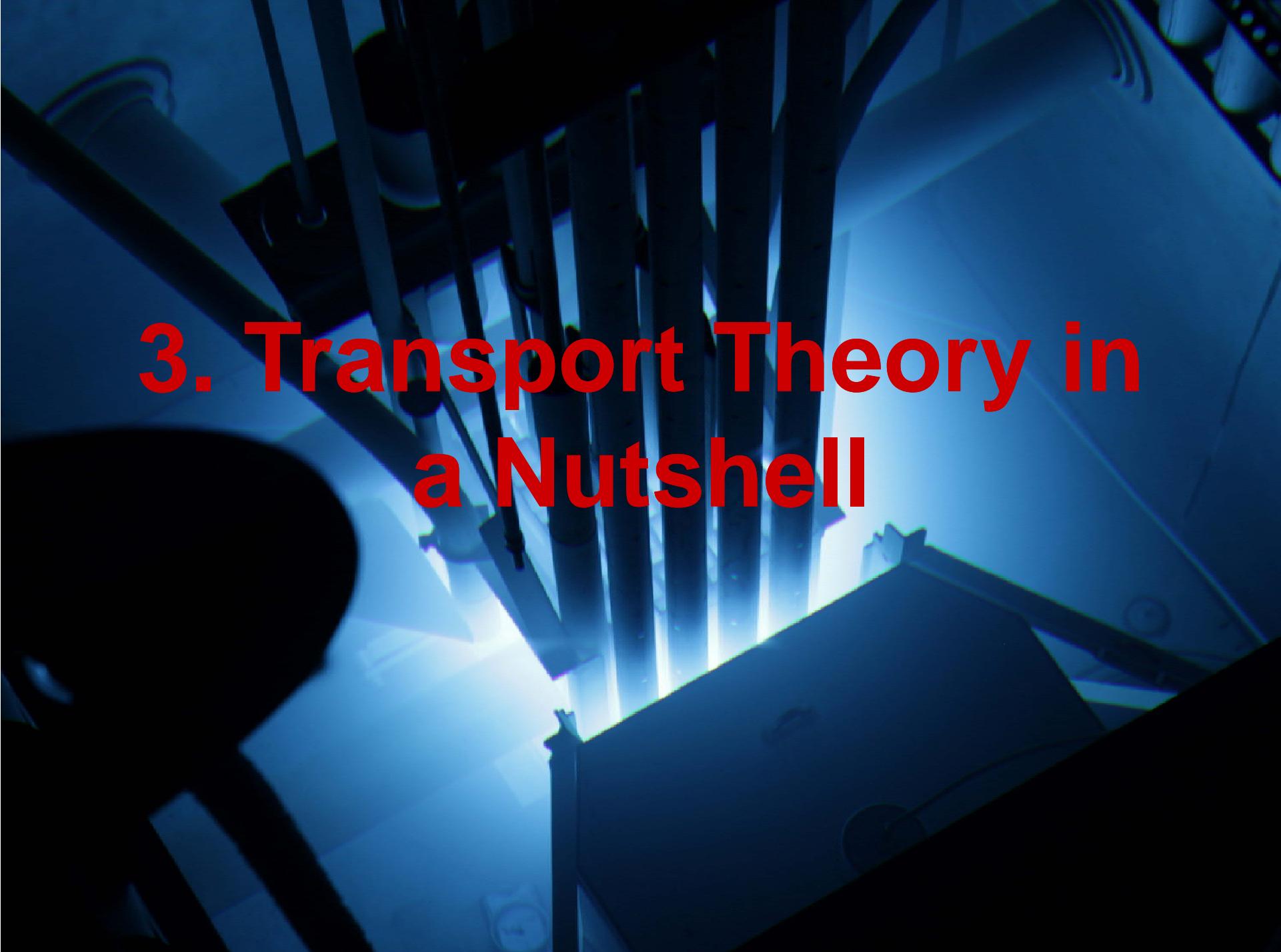
MS Thesis: *Accurate Holdup Calculations with Predictive Modeling & Data Integration*



2. Work Presented Today

- Dr. Joe Zerr is a former member of my research group
- Presently Research staff at Los Alamos National Laboratory
- Received PhD 2010,
Penn State University
- Recipient of 2010
Mark Mills Award



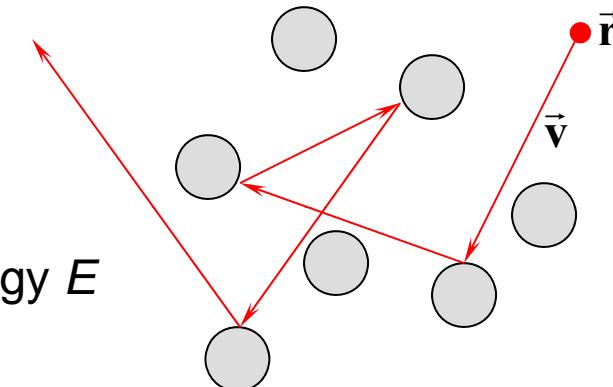


3. Transport Theory in a Nutshell

3. Fundamentals

- Classical point particle is fully described by independent variables:

- ❖ Time: t
- ❖ Space: \vec{r}
- ❖ Direction of motion: $\hat{\Omega}$
- ❖ Energy: E
- ❖ **Note:** the unit vector $\hat{\Omega}$ & energy E are equivalent to velocity \vec{v}



- Observable quantities (e.g. heating) depend on reaction rate:

- ❖ Large number of interacting particles (neutrons/photons)
- ❖ Much larger number of host targets (nuclei/electrons)
- ❖ Impractical to solve dynamic system for individual particles/targets
⇒ statistical model: **mean** collision density rate
- ❖ Proportional to density of interacting particles & host targets



3. Statistical Model

□ Particle Angular Density:

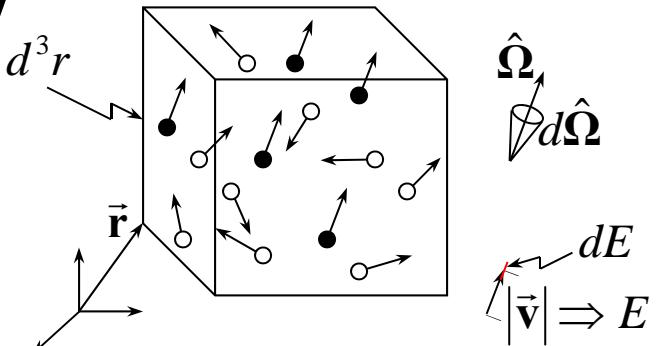
- ❖ $n(\vec{r}, \hat{\Omega}, E, t) d^3r dE d\hat{\Omega} \equiv$ mean number of neutrons at time t in d^3r at \vec{r} , with energy in $[E, E+dE]$ traveling in the directional cone $d\hat{\Omega}$ at $\hat{\Omega}$

□ Particle Density:

- ❖ $N(\vec{r}, E, t) d^3r dE \equiv$ mean number of neutrons at time t in d^3r at \vec{r} , with energy in $[E, E+dE] \Rightarrow N(\vec{r}, E, t) \equiv \int_{4\pi} d\hat{\Omega} n(\vec{r}, \hat{\Omega}, E, t)$

□ Particle Flux: Speed \times particle density

- ❖ Angular flux: $\psi(\vec{r}, \hat{\Omega}, E, t) \equiv v n(\vec{r}, \hat{\Omega}, E, t)$
- ❖ Scalar flux: $\varphi(\vec{r}, E, t) \equiv v N(\vec{r}, E, t)$
- ❖ Leakage rate: $\psi(\vec{r}, \hat{\Omega}, E, t) \hat{\Omega} \cdot \hat{n} dA$
- ❖ Reaction rate density: $\sum_j \varphi(\vec{r}, E, t)$
 - $\sum_j \equiv$ Probability reaction j / path length



Need to compute $\varphi(\vec{r}, E, t) = \int_{4\pi} \psi(\vec{r}, \hat{\Omega}, E, t) d\hat{\Omega}$



3. Neutron Transport Equation

- Special case of Boltzmann equation: First-order integro-differential
 - ❖ Neutral particles \Rightarrow no electro-magnetic forces
 - ❖ Low particle densities \Rightarrow ignore neutron-neutron collisions \Rightarrow linear
- Balance over infinitesimal element in phase space: $(\vec{r}, \hat{\Omega}, E)$
- Dependent variable: Angular flux $\psi(\vec{r}, \hat{\Omega}, E, t)$

Transient	Streaming	Total collision
$\frac{1}{v} \frac{\partial \psi(\vec{r}, E, \hat{\Omega}, t)}{\partial t} + \hat{\Omega} \cdot \nabla \psi(\vec{r}, E, \hat{\Omega}, t) + \Sigma_t(\vec{r}, E) \psi(\vec{r}, E, \hat{\Omega}, t)$		
Scattering	$= \int_{4\pi} d\hat{\Omega}' \int_0^\infty dE' \Sigma_s(E', \hat{\Omega}' \rightarrow E, \hat{\Omega}) \psi(\vec{r}, E', \hat{\Omega}', t)$	
Fission	$+ \frac{\chi(E)}{4\pi} \left[\int_0^\infty dE' \nu(E') \Sigma_f(\vec{r}, E') \int_{4\pi} d\hat{\Omega}' \psi(\vec{r}, E', \hat{\Omega}', t) \right] + \underbrace{s(\vec{r}, E, \hat{\Omega}, t)}_{\text{Fixed source}}$	



3. Interface & Boundary Conditions

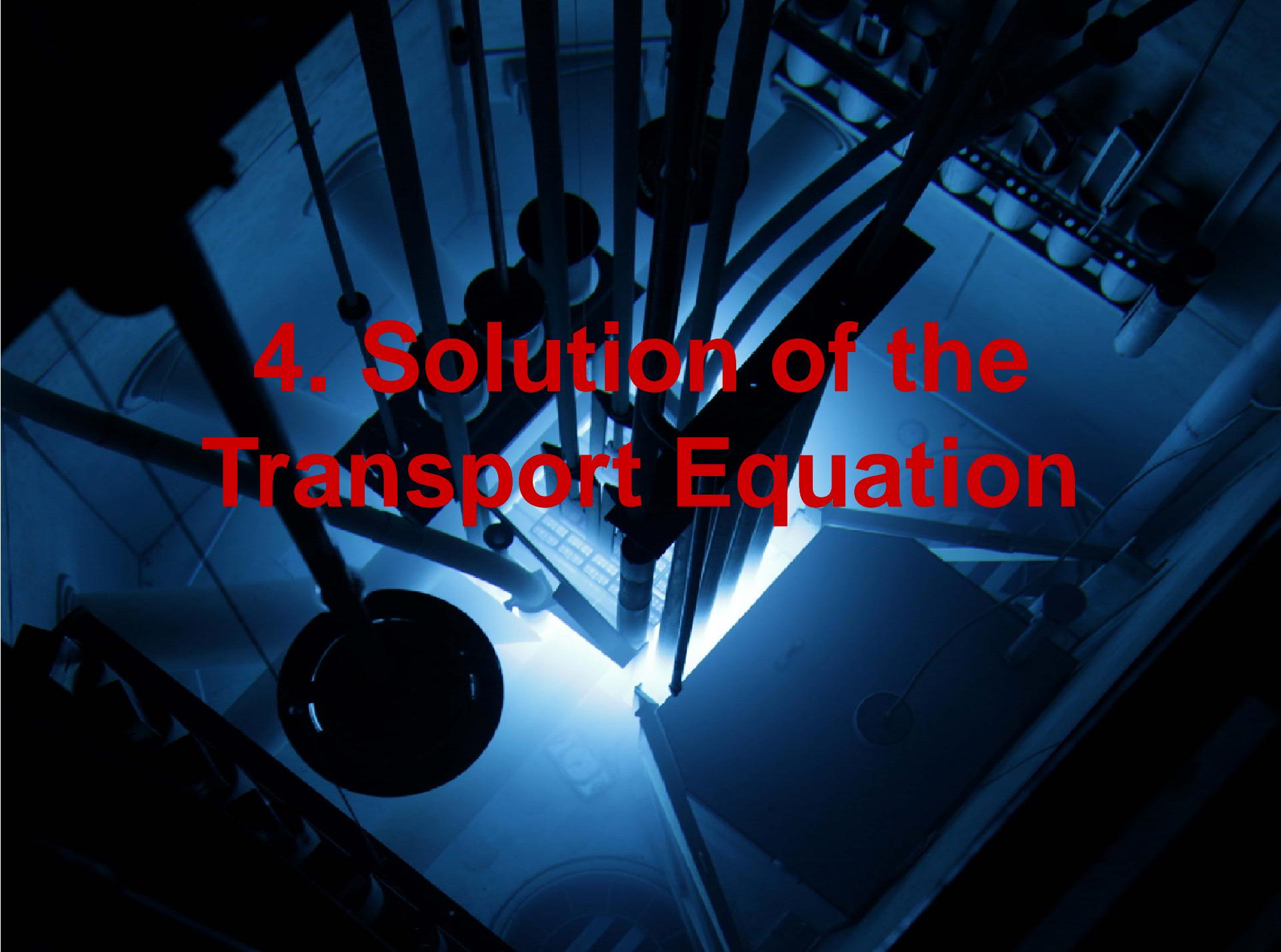
- **Steady state:** Time derivative vanishes
- **Interface condition:** Angular flux continuous along direction of motion, $\hat{\Omega}$, across material boundaries
- **Physical intuition:** Can specify what goes into a system
 - ❖ What comes out is a consequence of the transport process inside
 - ❖ Example: shining light into crystal
 - Can choose color/intensity of incoming light
 - Can't choose color/intensity of outgoing light: depends on what happens inside
- **Typical Boundary Condition (BC):**
 - ❖ Set **incoming** flux $\psi(\vec{r}_s, E, \hat{\Omega}, t) = \psi_{in}(\vec{r}_s, E, \hat{\Omega}, t)$ for:
 - All times t
 - All energies: $E \in [0, \infty]$
 - Each \vec{r}_s on the boundary S
 - Each incoming angle: $\hat{\Omega} \cdot \hat{e}_s < 0$, \hat{e}_s is the normal unit vector pointing out
 - ❖ The function $\psi_{in}(\vec{r}_s, E, \hat{\Omega}, t)$ can be specified explicitly or implicitly
 - Vacuum BC: $\psi_{in}(\vec{r}_s, E, \hat{\Omega}, t) = 0$



3. Discretization of Transport Equation

- Implementation on digital computer \Rightarrow discretize independent variables & consequently dependent variables
- Energy: Multigroup \Rightarrow discretization into *bins* (E_g, E_{g-1})
 - ❖ Victory, 1985: Total & scattering cross section fluctuations diminish with refinement of energy group structure
 - \Rightarrow **Multigroup solution \rightarrow exact solution**
- Angle: Discrete-ordinates \Rightarrow discretization along discrete $\hat{\Omega}_n$
 - ❖ Madsen, 1971: Quadrature formula converges with increasing order
 - \Rightarrow **Discrete Ordinates solution \rightarrow exact one-speed solution**
- Space: Multitude of methods discretize $\nabla\psi$ on spatial mesh
 - ❖ Madsen, 1972: Exact solution has bounded 3rd derivatives
 - \Rightarrow **Diamond Difference solution \rightarrow exact Discrete Ordinates solution**
 - ❖ Smoothness hypothesis unrealistic for most applications





4. Solution of the Transport Equation

4. Traditional Solution Algorithms

- **Difficulty of solving the transport equation (partial list):**
 - ❖ Steady-state 3-D problems: phase space is 6-D \Rightarrow 100 discrete variables per phase space dimension yields 10^{12} unknowns
 - ❖ Neutron cross sections sensitive to energy & nuclide composition
 - ❖ Source (fission & scattering) depends on solution $\psi \Rightarrow$ iterate
- **Nested loops:**
 - ❖ Outer Iteration: converge fission/scattering source
 - ❖ Loop over energy groups: from highest to lowest E
 - ❖ In each group sum fission + in-scattering source guess into q
 - ❖ Inner (or Source) Iteration: reconcile within group source & flux
 - Starting with guess for group φ within-group scattering source
 - Invert 1st order PDE on source distribution \Rightarrow group ψ
 - Integrate over angle \Rightarrow new φ : test if too different from starting guess
 - Yes: Repeat unless already used too many iterations
 - No: solution converged successfully to group flux for in-group sources



4. Cell Equations

- **Kernel operation: Solving 1st order PDE for given source**

- ❖ Conducted via *mesh sweep* algorithm: 1 cell, 1 angle at a time

- **Discretized equations per energy group/angle/cell:**

- ❖ Balance: Include fission & in-scattering from other groups in q_n

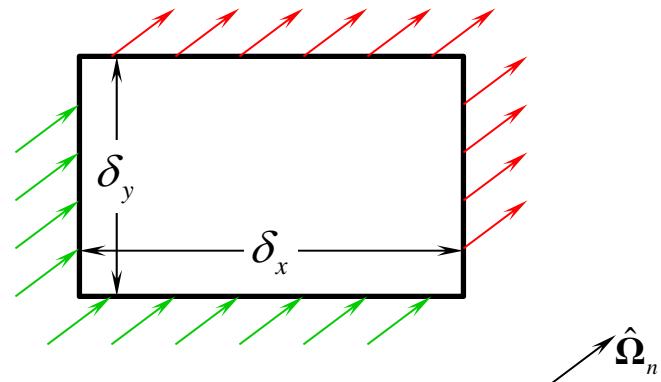
$$\frac{\mu_n}{\delta_x} (\psi_{n,out}^x - \psi_{n,in}^x) + \frac{\eta_n}{\delta_y} (\psi_{n,out}^y - \psi_{n,in}^y) + \frac{\xi_n}{\delta_z} (\psi_{n,out}^z - \psi_{n,in}^z) + \sigma_t \bar{\psi}_n = \sigma_s \bar{\phi} + q_n$$

- ❖ Auxiliary: Method-dependent, simplest
is Diamond Difference

$$\bar{\psi}_n = \frac{1}{2} (\psi_{n,out}^u + \psi_{n,in}^u), \quad u = x, y, z$$

- ❖ Quadrature formula:

$$\bar{\phi} = \sum_{n=1}^N \omega_n \bar{\psi}_n$$



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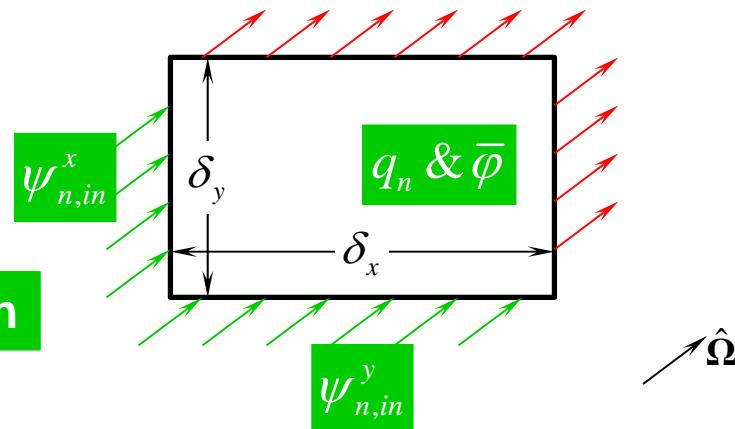
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known



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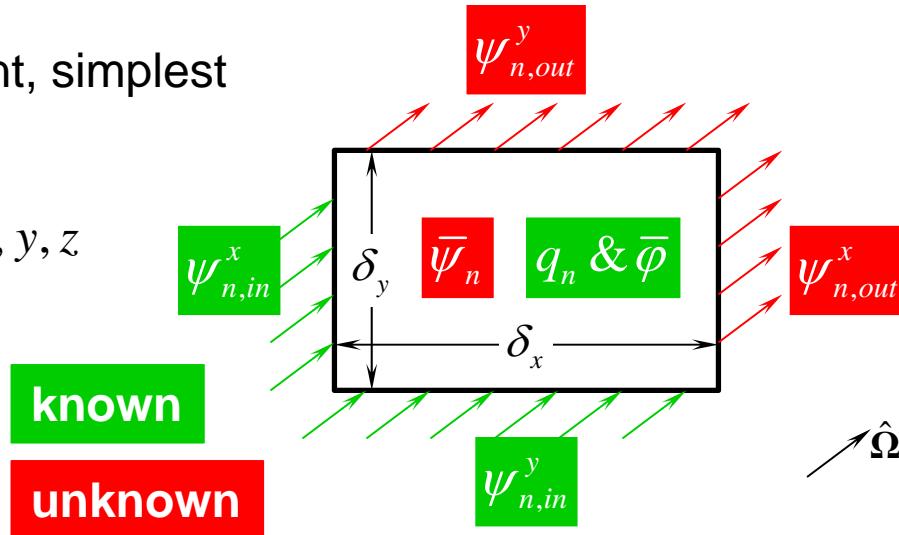
$$\frac{\mu_n}{\delta_x} (\psi_{n,out}^x - \psi_{n,in}^x) + \frac{\eta_n}{\delta_y} (\psi_{n,out}^y - \psi_{n,in}^y) + \frac{\xi_n}{\delta_z} (\psi_{n,out}^z - \psi_{n,in}^z) + \sigma_t \bar{\psi}_n = \sigma_s \bar{\phi} + q_n$$

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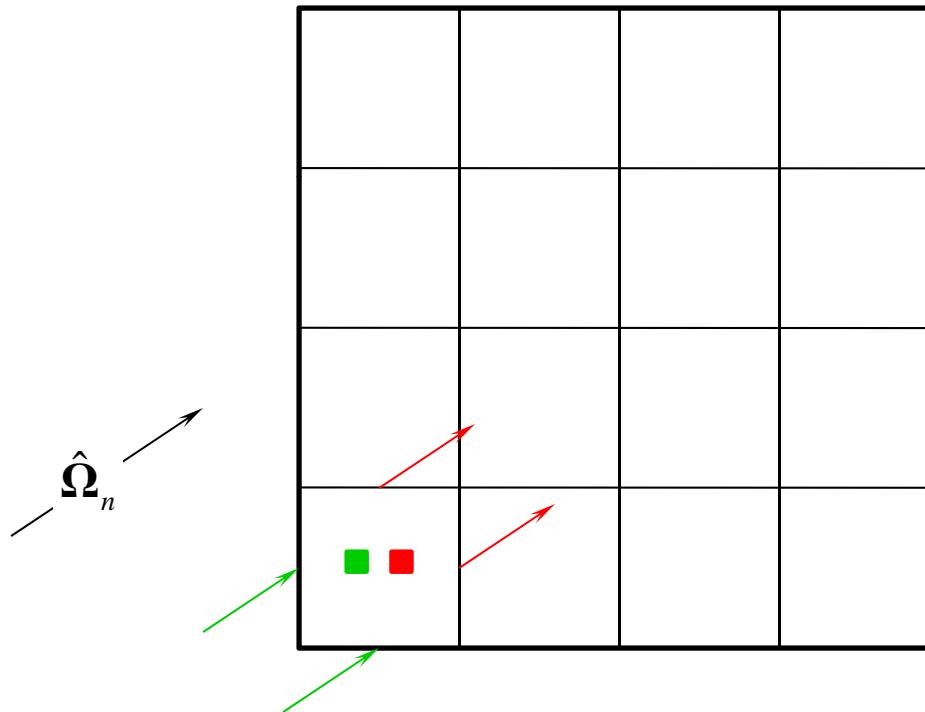
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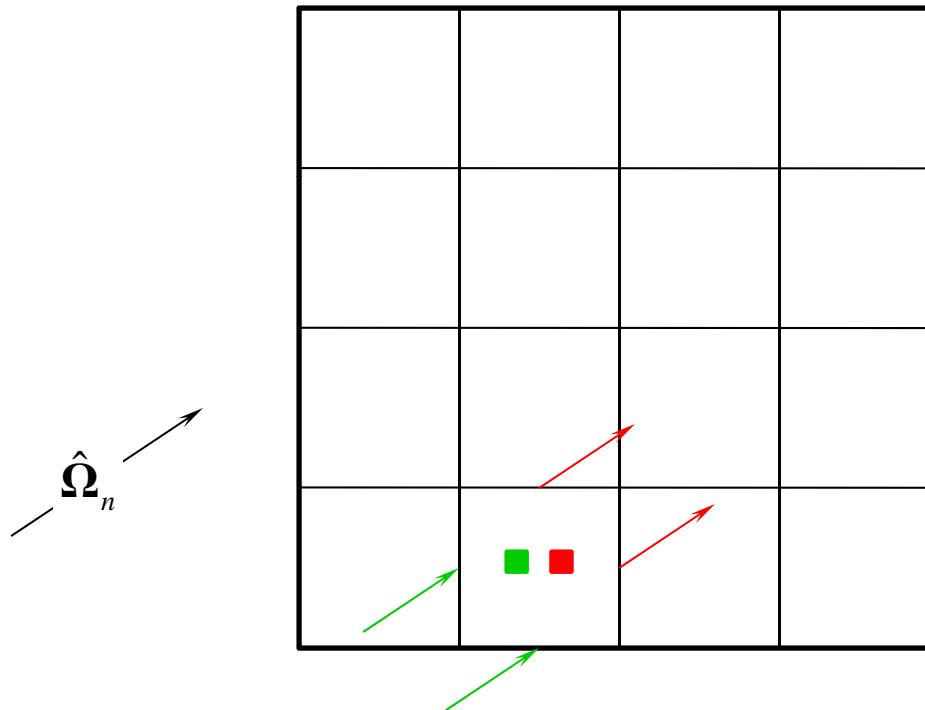
4. Mesh Sweep

- For each discrete ordinate sweep cells along $\hat{\Omega}_n$

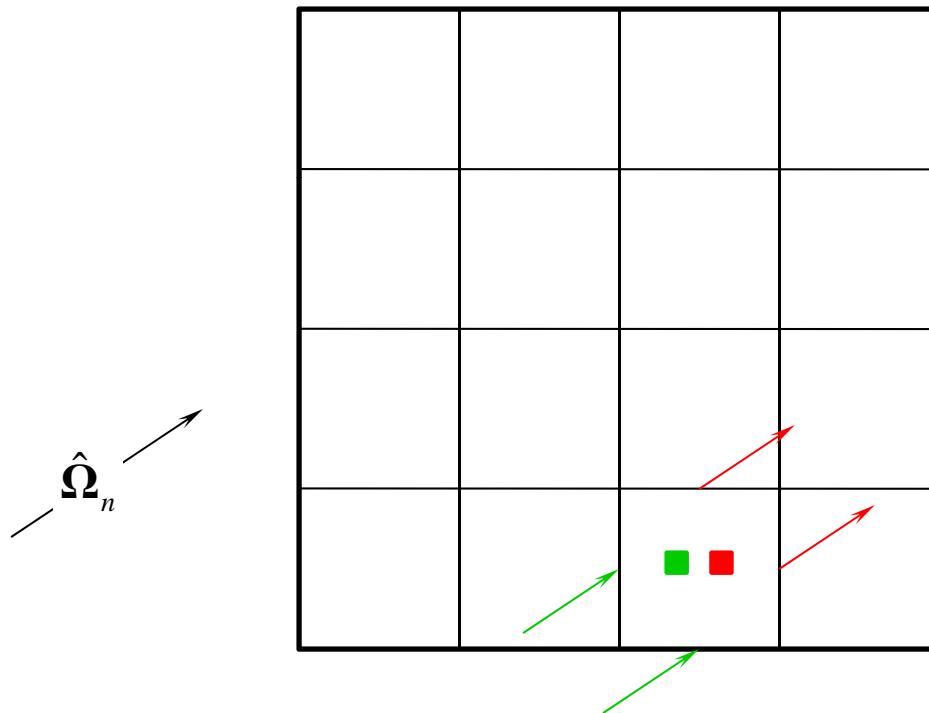


4. Mesh Sweep

- Interface angular fluxes couple neighboring cells

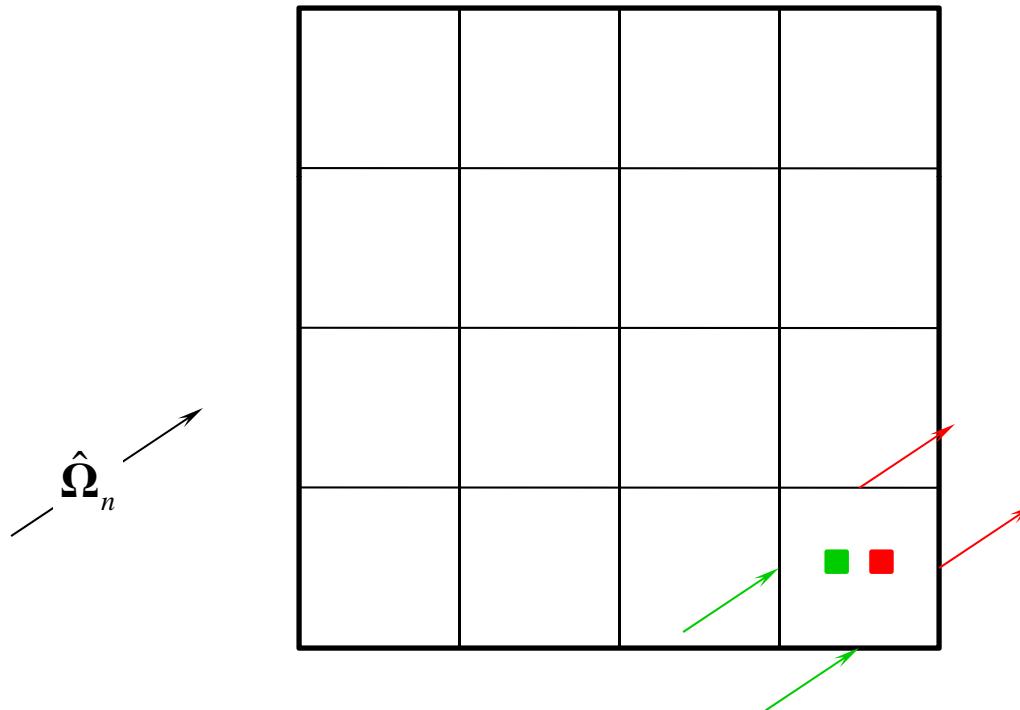


4. Mesh Sweep



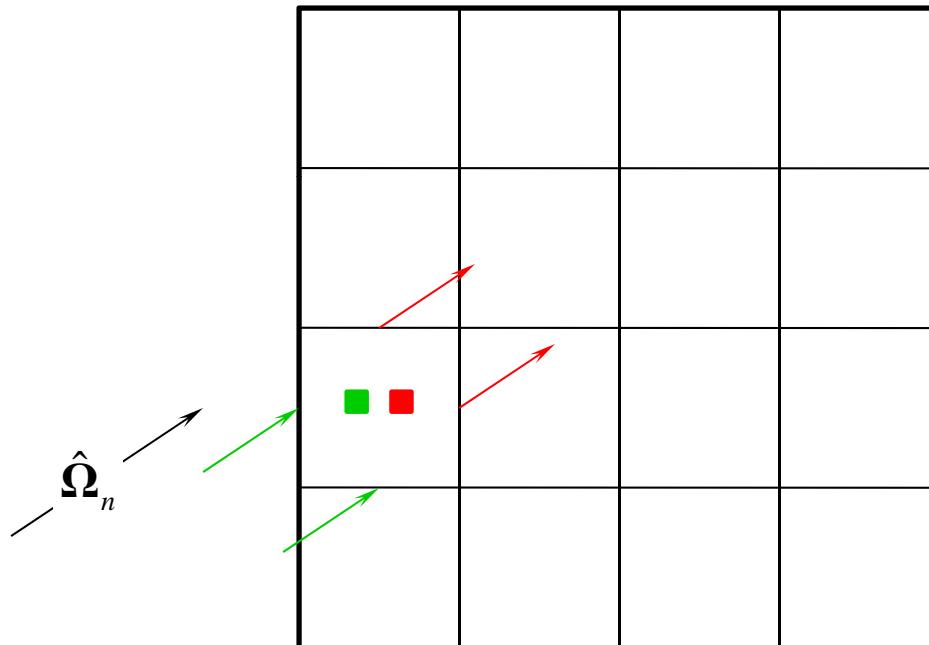
4. Mesh Sweep

- Upon reaching end of row go to next row along $\hat{\Omega}_n$

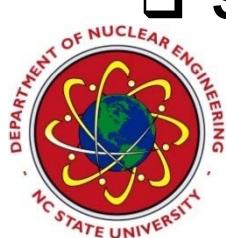


4. Mesh Sweep

- Note sequential nature: must compute upstream cell first



- Slow convergence if SI often demands acceleration (DSA)



4. Parallelization Strategies

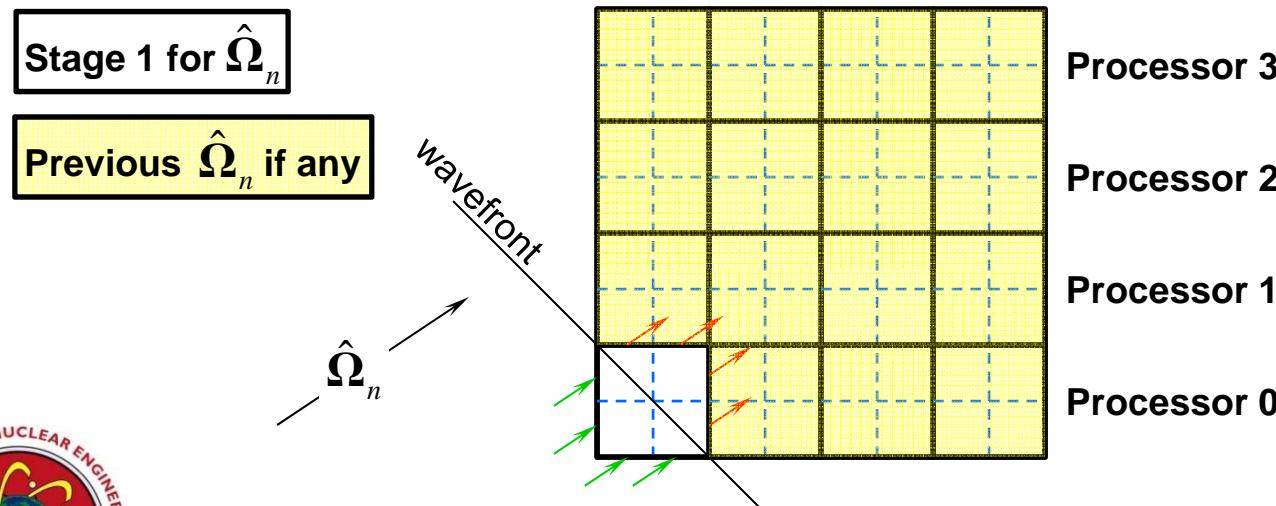
- **Domain decomposition:** split range of phase space variable in P subdomain & assign each to different process
 - ❖ Perfect parallelization \Rightarrow reduce execution time by factor P
 - ❖ **Synchronous** DD: processes fully independent \Rightarrow computation (e.g. number of iterations) independent of P
 - ❖ **Asynchronous** DD: coupled processes \Rightarrow work P dependent
- **Possible DD for S_N algorithms:**
 - ❖ Energy: speedup limited by # groups (hundreds at most)
 - Coarse-grain \Rightarrow easy to implement & high parallel efficiency
 - Asynchronous & pointless unless there is fission &/or upscattering
 - ❖ Angle: speedup limited by # angles (few thousands at most)
 - Medium-grain \Rightarrow easy to implement & good parallel efficiency
 - Synchronous in Cartesian coordinate system
 - ❖ Spatial: speedup limited by # cells (potentially many millions)
 - Fine-grain \Rightarrow difficult to implement & OK parallel efficiency
 - Synchronous implies retaining sequential order among subdomains



4. Spatial Domain Decomposition

☐ Koche- Baker-Alcouffe SDD: synchronous

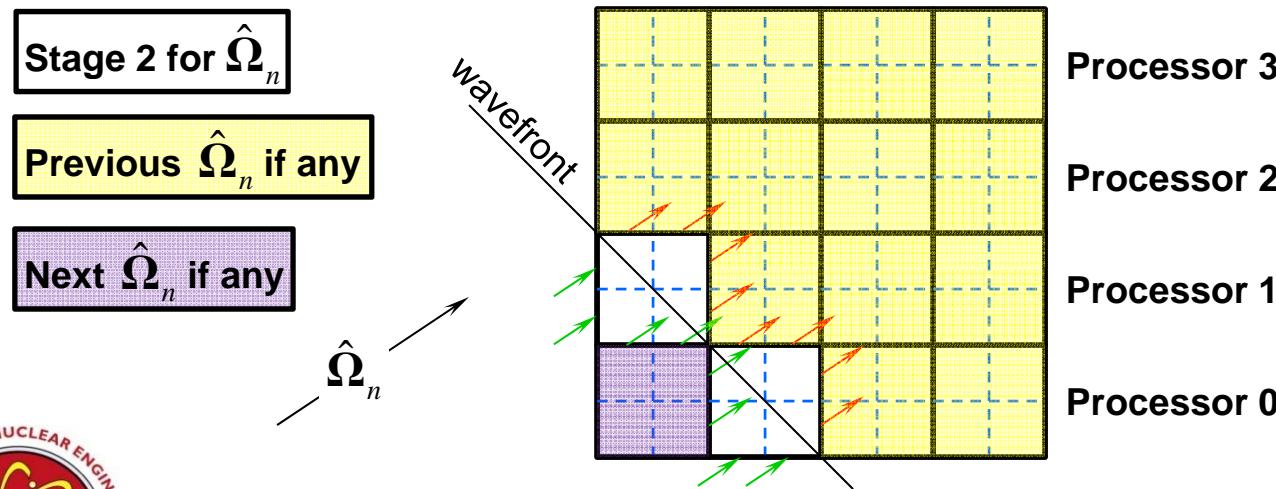
- ❖ Maps 3D mesh onto 2D processor topology
- ❖ Sweep mesh by subdomain in natural sequence per $\hat{\Omega}_n$
- ❖ Concurrently sweep ready subdomains (on waverfront) \Rightarrow SDD
- ❖ Pipeline angles to reduce processor idleness
- ❖ Communicate outgoing interface ψ to neighbors across waverfront

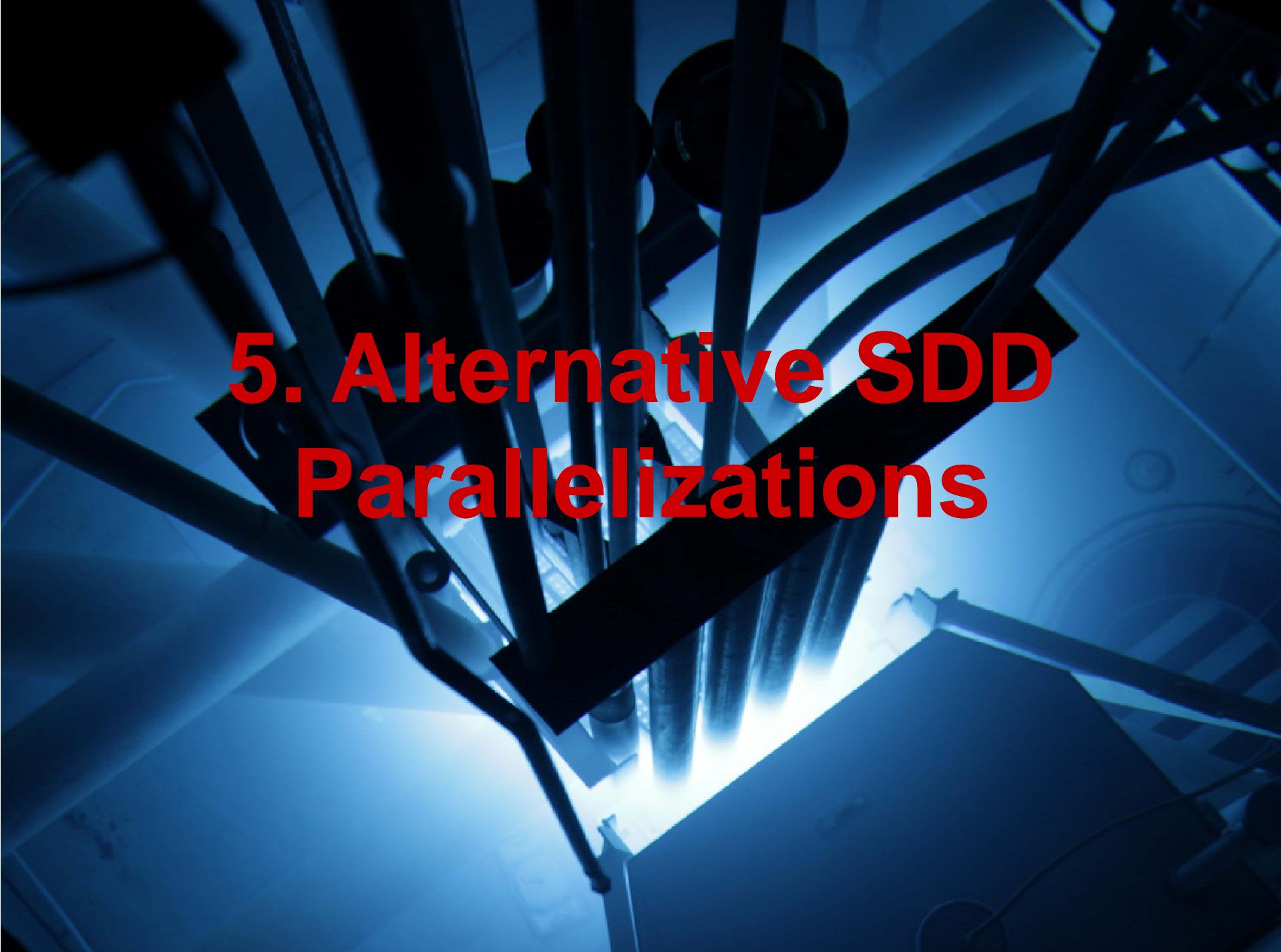


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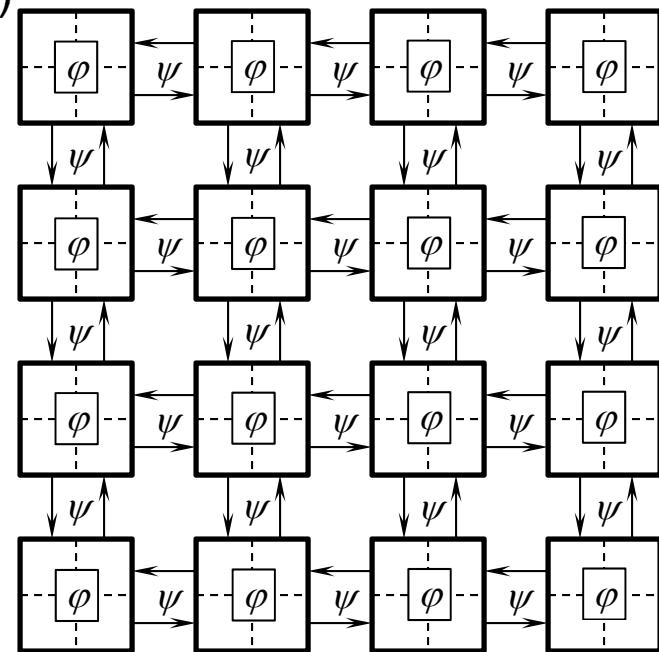
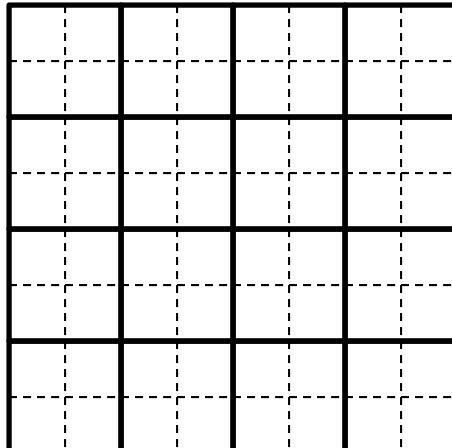


5. Alternative SDD Parallelizations

5. Parallel Block Jacobi (PBJ)

□ Asynchronous SDD alternative to KBA:

- ❖ Eliminate processor idleness & increase concurrent processes
- ❖ Combine all angles computations in subdomain via ITMM operators
- ❖ Replace SI with iterations on subdomain interface fluxes:
Communicate ψ between iterations
- ❖ Need $\varphi = \Phi(\psi_{in}, \mathbf{q})$ & $\psi_{out} = \Psi(\varphi, \psi_{in}, \mathbf{q})$



5. Outline of ITMM

- Express SI as a mapping of flux ℓ iterate into $\ell+1$ iterate:

$$\varphi^{(\ell+1)} = \mathbf{J}_\varphi \left(\varphi^{(\ell)} + \Sigma_s^{-1} \mathbf{q} \right) + \mathbf{K}_\varphi \psi_{in} \quad \Rightarrow \quad \mathbf{J}_\varphi = \partial \varphi^{(\ell+1)} / \partial \varphi^{(\ell)}$$

- ❖ Upon iterative convergence:

$$(\mathbf{I} - \mathbf{J}_\varphi) \varphi^\infty = \mathbf{J}_\varphi \Sigma_s^{-1} \mathbf{q} + \mathbf{K}_\varphi \psi_{in} \quad \rightarrow \quad \Phi(\psi_{in}, \mathbf{q})$$

- ❖ For full domain where ψ_{in} is known from BCs:

$$\psi_{out} = \mathbf{J}_\psi \left(\varphi^\infty + \Sigma_s^{-1} \mathbf{q} \right) + \mathbf{K}_\psi \psi_{in} \quad \rightarrow \quad \Psi(\varphi, \psi_{in}, \mathbf{q})$$

- Apply to subdomain: ψ_{in} is not known requires iteration

- ITMM operators are response matrices:

- ❖ \mathbf{J}_φ : cell-averaged scalar flux due to unit distributed source
 - ❖ \mathbf{K}_φ : cell-averaged scalar flux due to unit incident angular flux
 - ❖ \mathbf{J}_ψ : outgoing angular flux due to unit distributed source
 - ❖ \mathbf{K}_ψ : outgoing angular flux due to unit incident angular flux



5. Construction of ITMM Operators

- **Differential Mesh Sweep (DMS): concurrently in all subdomains**
 - ❖ Perform single sweep to compute elements of \mathbf{J}_φ via $\mathbf{J}_\varphi = \partial\varphi^{(\ell+1)} / \partial\varphi^{(\ell)}$
 - ❖ Compute elements of other operators along the way
- **Dense operators: memory limitations as size grows**
 - ❖ Operators sizes grow superlinear with # cells, linear with # angles
 - ❖ Full coupling of cell- and face-fluxes
 - ❖ Expensive to invert for large subdomains
- **Applicable to high-order spatial discretizations & anisotropic scattering ($\varphi \Rightarrow$ angular moments)**



5. PBJ Algorithm

□ For each energy group, fully concurrent:

- ❖ Perform DMS per subdomain/processor \Rightarrow 4 ITMM operators
- ❖ Need to effect $(\mathbf{I} - \mathbf{J}_\phi)^{-1} \Rightarrow$ LU factorization only once then store
- ❖ Execute PBJ iterations on subdomain interface angular fluxes:
 - Given $\psi_{in}^{(\lambda)}$ compute for each subdomain $\phi^{(\lambda+1)} = \Phi(\psi_{in}^{(\lambda)}, \mathbf{q})$
 - Test convergence of scalar flux: $|1 - \phi^{(\lambda+1)} / \phi^{(\lambda)}|$ small?
 - If converged go to next group (if any)
 - Otherwise start new iteration with: $\psi_{out}^{(\lambda+1)} = \Psi(\phi^{(\lambda+1)}, \psi_{in}^{(\lambda)}, \mathbf{q})$
 - Communicate $\psi_{out}^{(\lambda+1)}$ as $\psi_{in}^{(\lambda+1)}$ to neighboring subdomains

□ Observations:

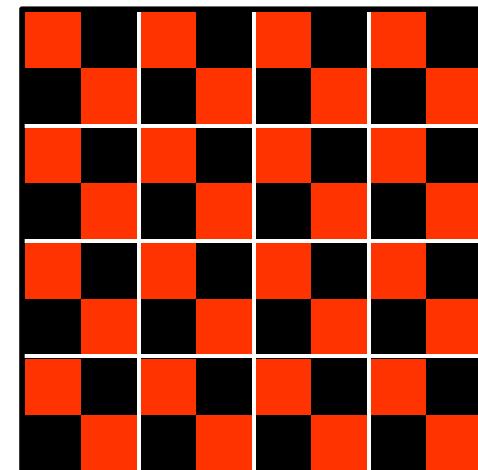
- ❖ Reduced local computations to matrix-vector multiplies & adds
- ❖ Sources of parallel inefficiency:
 - Increasing # iterations with P & tighter subdomain coupling: $c \uparrow$ & $\Sigma_t h \downarrow$
 - Network contention: Communicate subdomain interface angular fluxes

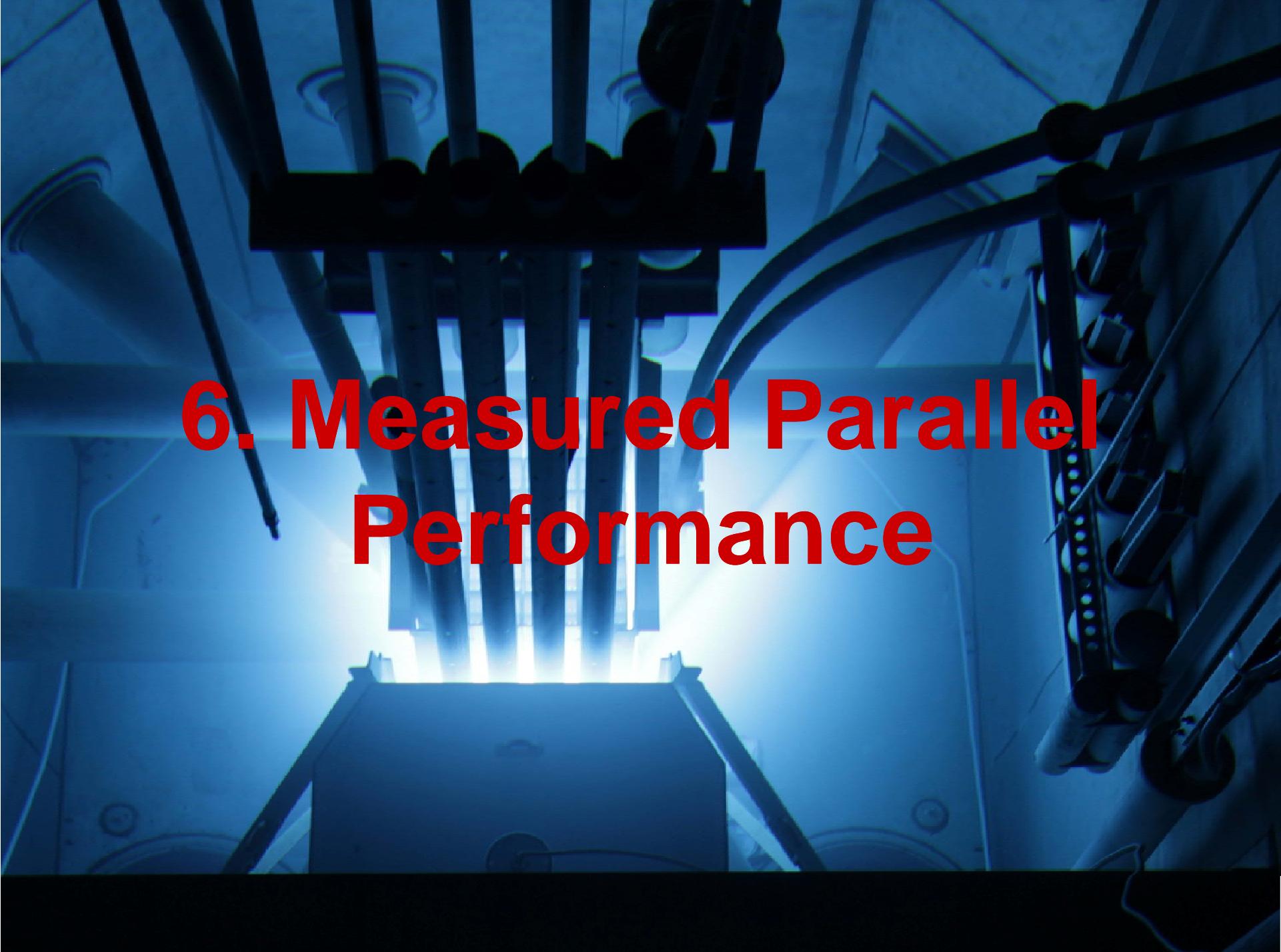


5. Parallel Gauss-Seidel (PGS)

□ Red/Black splitting of each subdomain:

- ❖ Each sub-subdomain is either **red** or **black**
- ❖ Operations per global iteration over interface angular flux:
 - Solve local ITMM system for φ & ψ_{out}
 - Copy/send $\psi_{out} \rightarrow \psi_{in}$ to intra-/iter-subdomain neighbor
 - Solve local ITMM system for φ & ψ_{out}
 - Copy/send $\psi_{out} \rightarrow \psi_{in}$ to intra-/iter-subdomain neighbor
- ❖ Pros:
 - Memory requirement \downarrow super-linearly with \downarrow number of cells
 - Typically Gauss-Seidel convergence rate better than Block Jacobi
- ❖ Cons:
 - Smaller ITMM subdomains \Rightarrow slower convergence of global iterations





6. Measured Parallel Performance

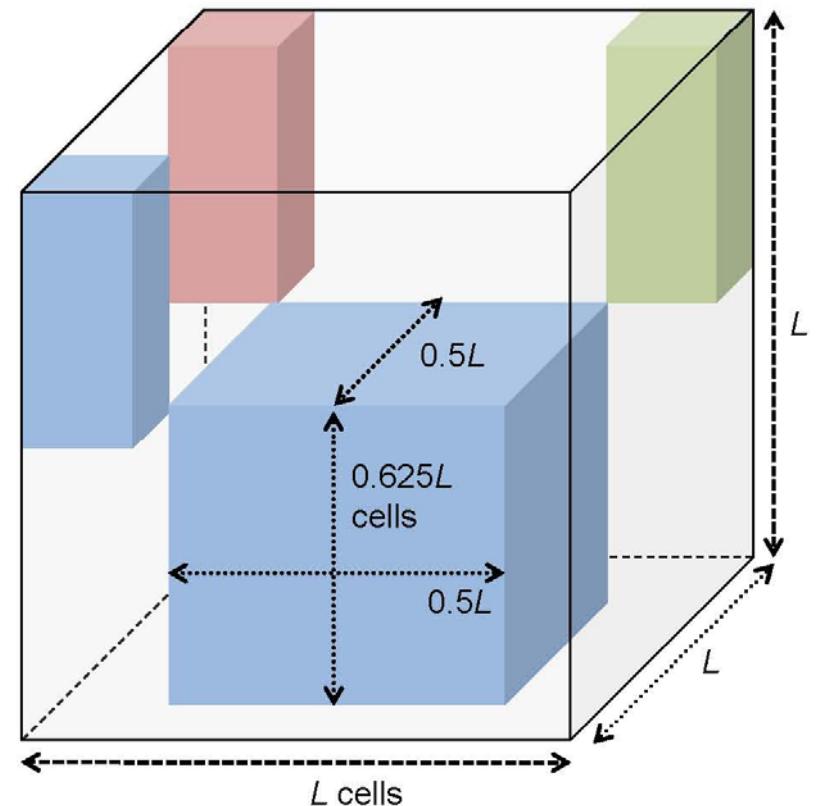
6. Implementation in PIDOTS

- Implemented PBJ & PGS in *Parallel Integral Discrete Ordinates Transport Solver* (PIDOTS)
- All tests performed with one-group, DD, with isotropic scattering
- Preliminary testing on:
 - ❖ LANL's Yellowrail:
 - 139 compute nodes each with 8 Processing Elements (PE) & 16 GB memory
 - Runs up to $P = 256$ processes
 - ❖ LANL's Redtail:
 - 1,834 compute nodes each with 8 PEs & 32 GB memory
 - Runs up to $P = 1,024$
 - ❖ ORNL's JaguarPF: All results here are for this platform
 - 18,688 compute nodes each with 12 PEs & 16 GB memory
 - Runs up to $P = 32,768$



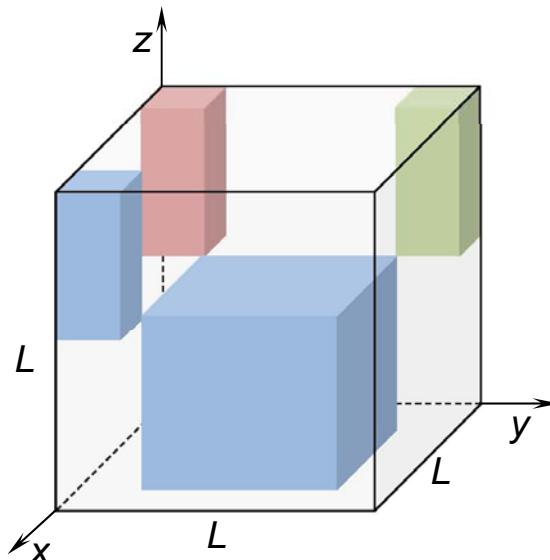
6. Weak Scaling Tests

- Evaluate parallel performance as problem size grows with P
- Weak scaling \Rightarrow Number of cells per processor fixed:
 - ❖ Start with $L \times L \times L$ cubic-cells domain
 - ❖ Comprised of 4 materials: no symmetries
 - ❖ Examine effect on number of iterations \Rightarrow execution time:
 - Cell size h set to $\{0.1, 1.0, 10.0\}$ cm
 - Scattering ratio c set to $\{0.9, 0.99\}$
 - ❖ Choice of L & S_N order is memory-limited

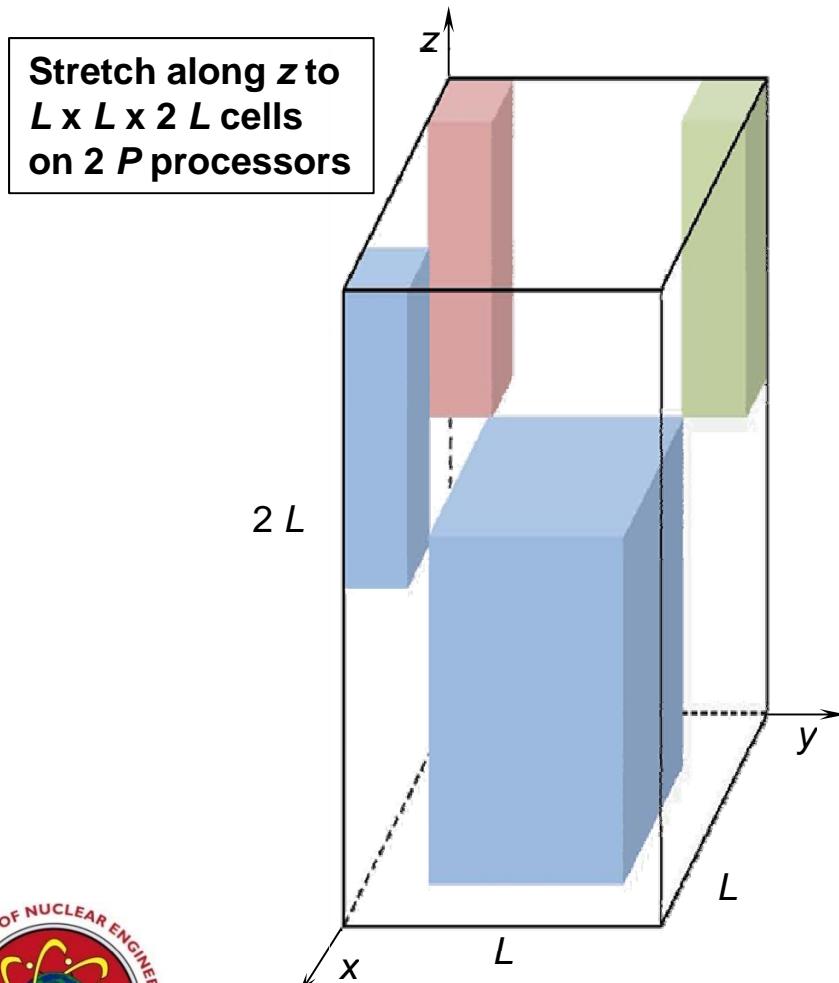


6. Growing Problem Size

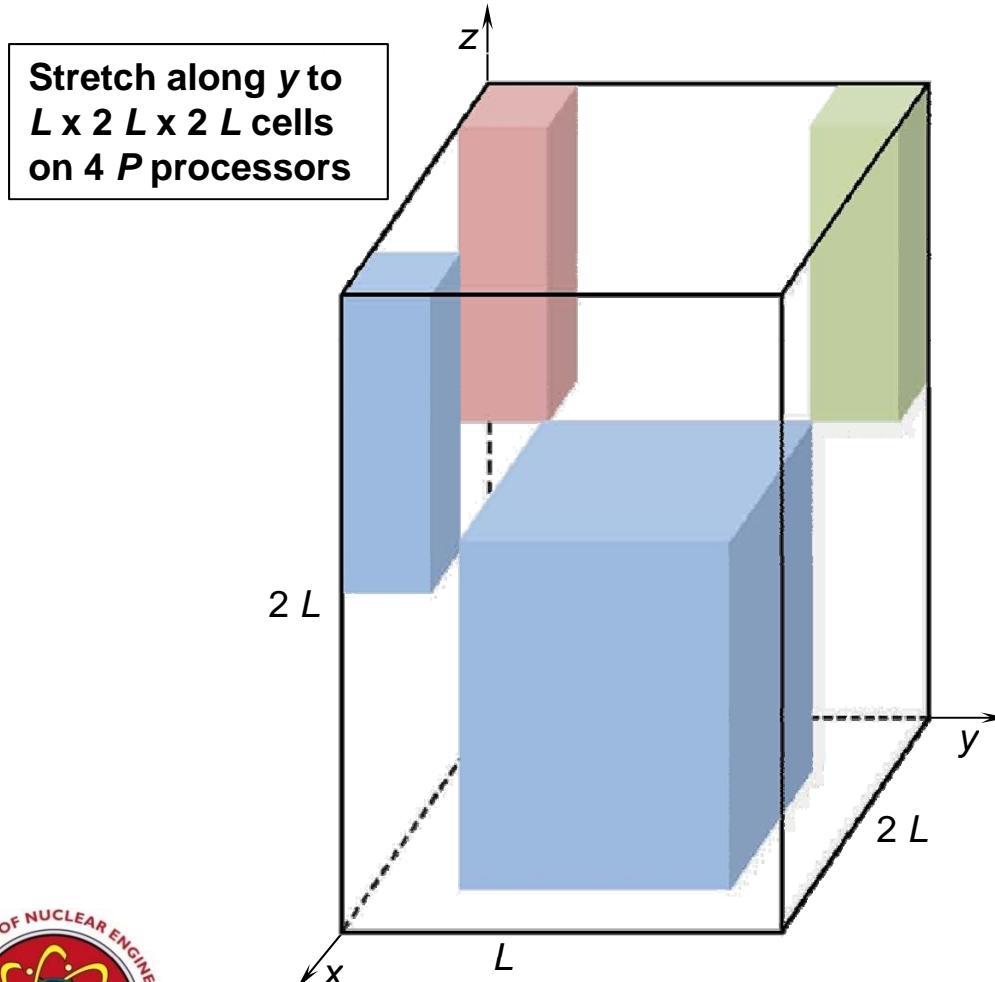
Starting with the
base case:
 $L \times L \times L$ cells
on P processors



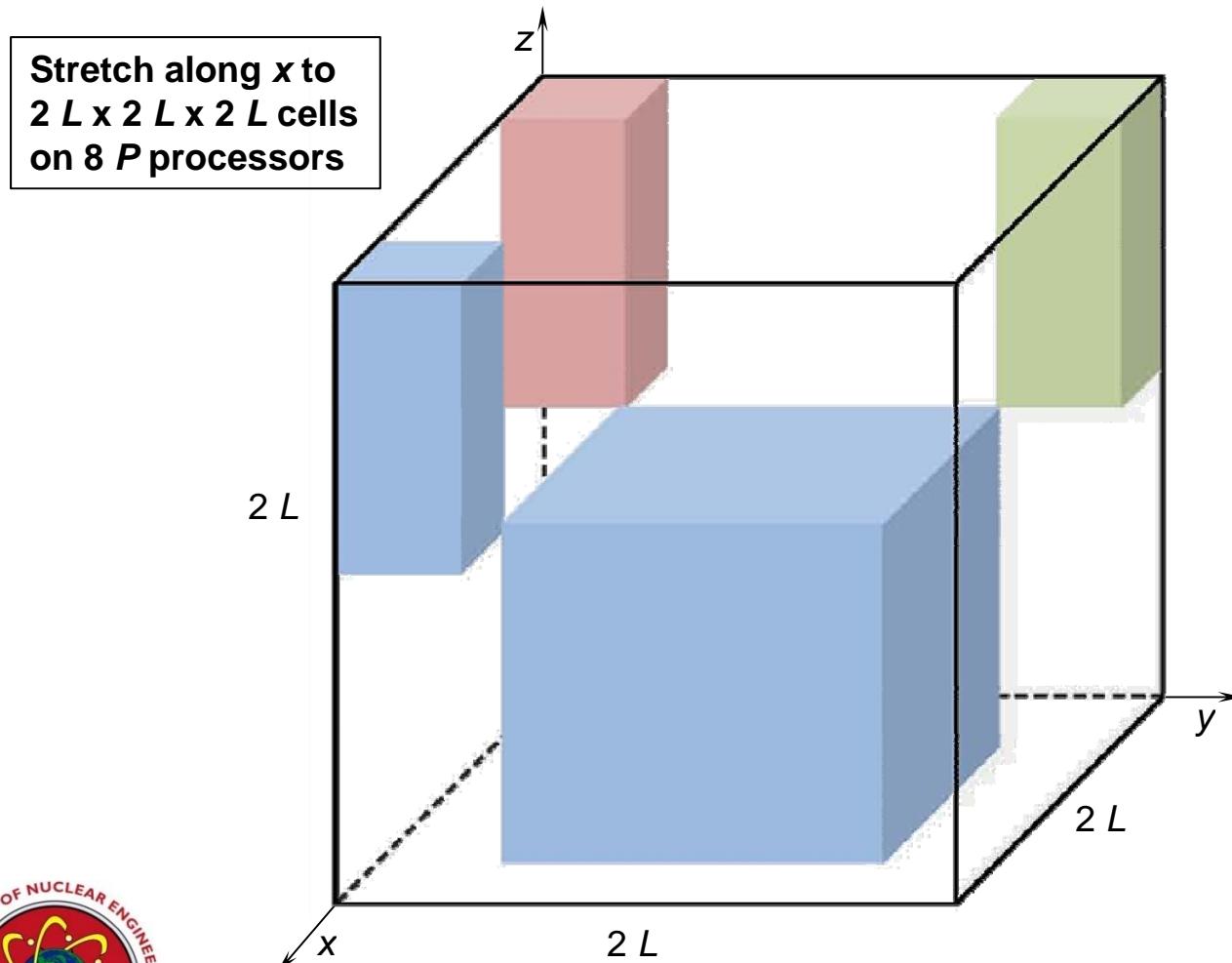
6. Growing Problem Size



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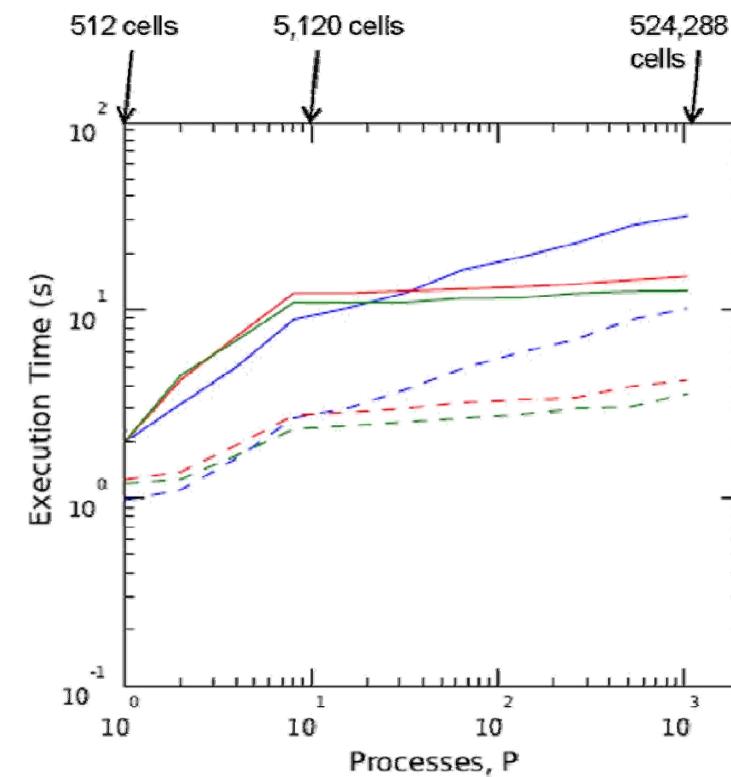
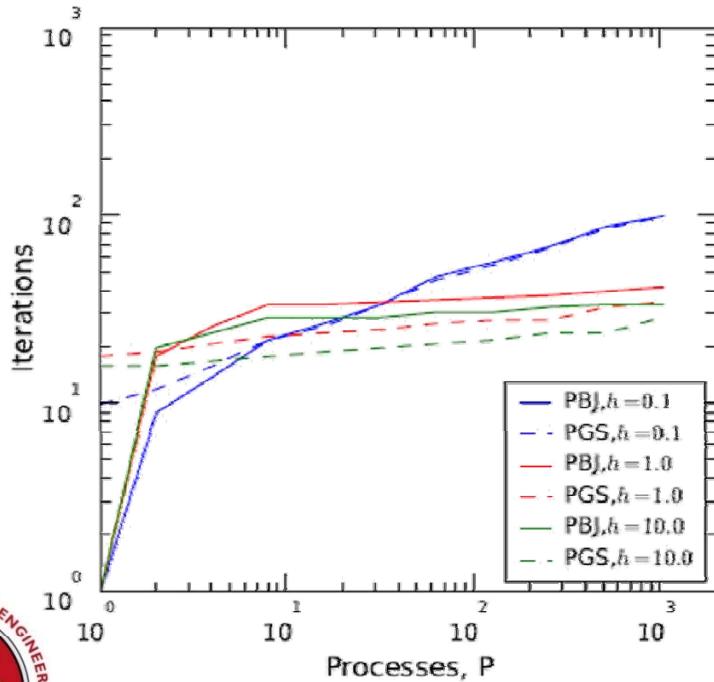
6. Growing Problem Size



6. PBJ vs PGS Performance

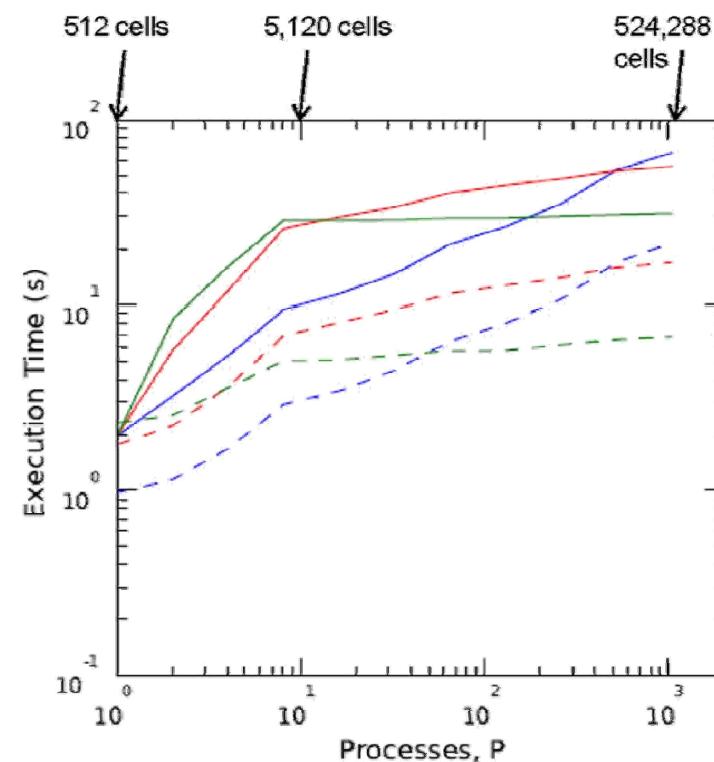
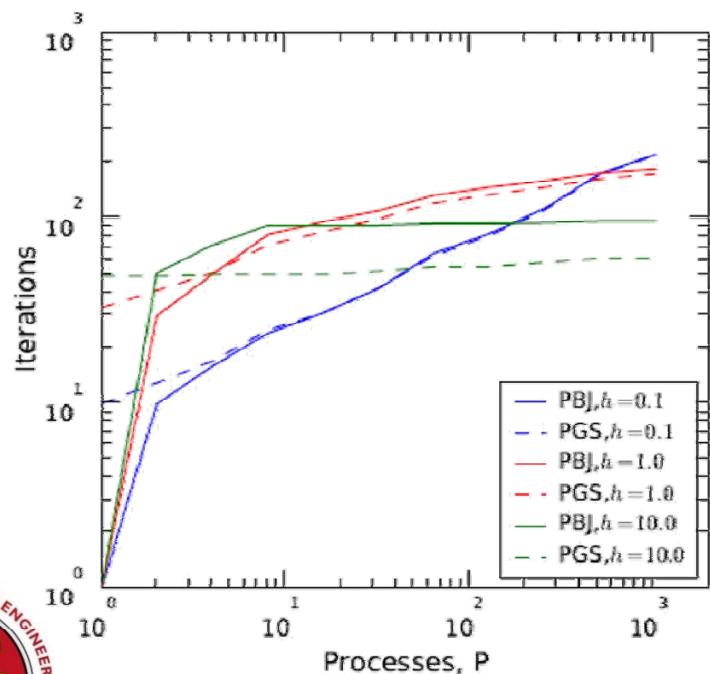
- 8x8x8-cell model per P , S_{16} , P up to 1,024 on JaguarPF
- PGS: eight 4x4x4-cell sub-subdomains per $P \Rightarrow$ shorter construction & per-iteration times

$c = 0.9$ results:



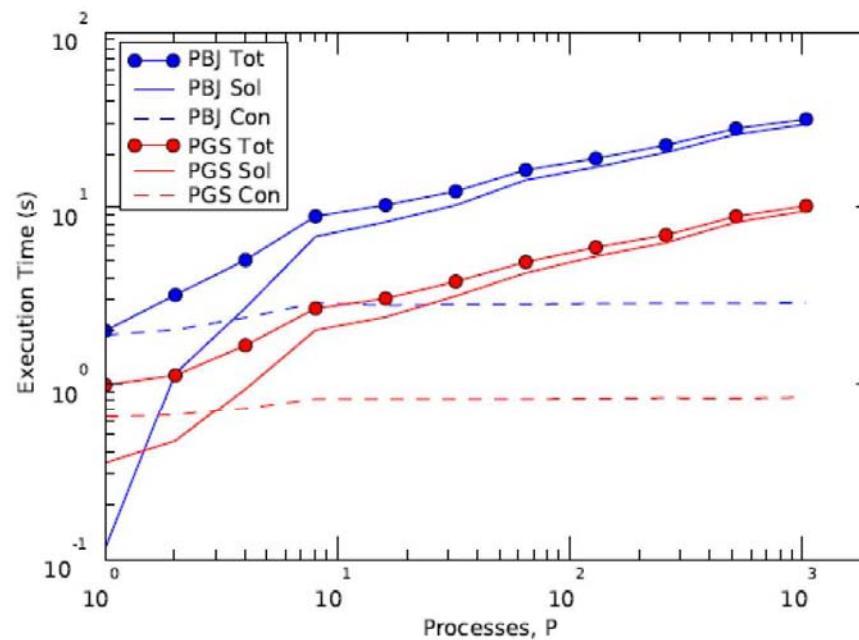
6. PBJ vs PGS Performance

$c = 0.99$ results:



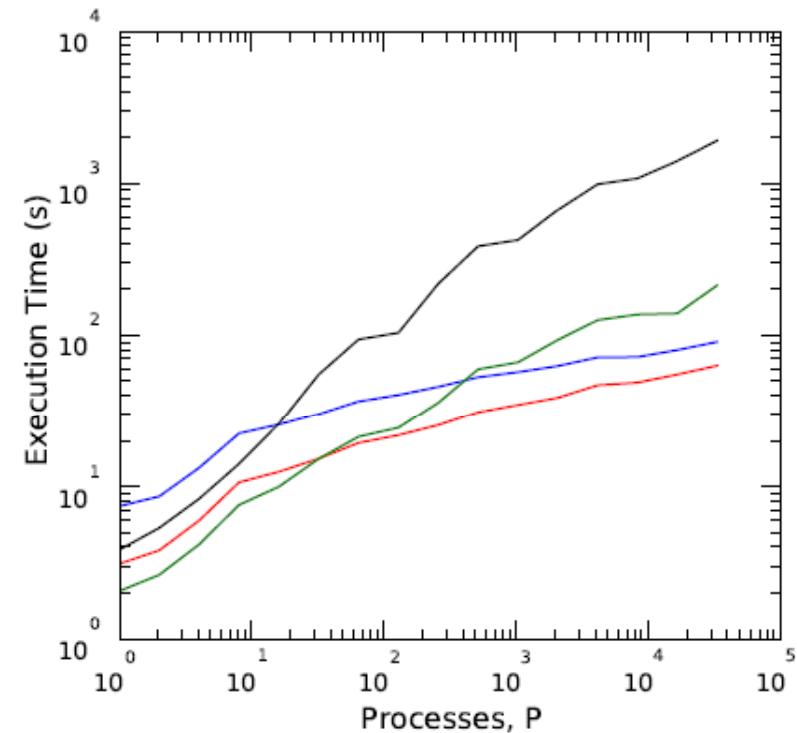
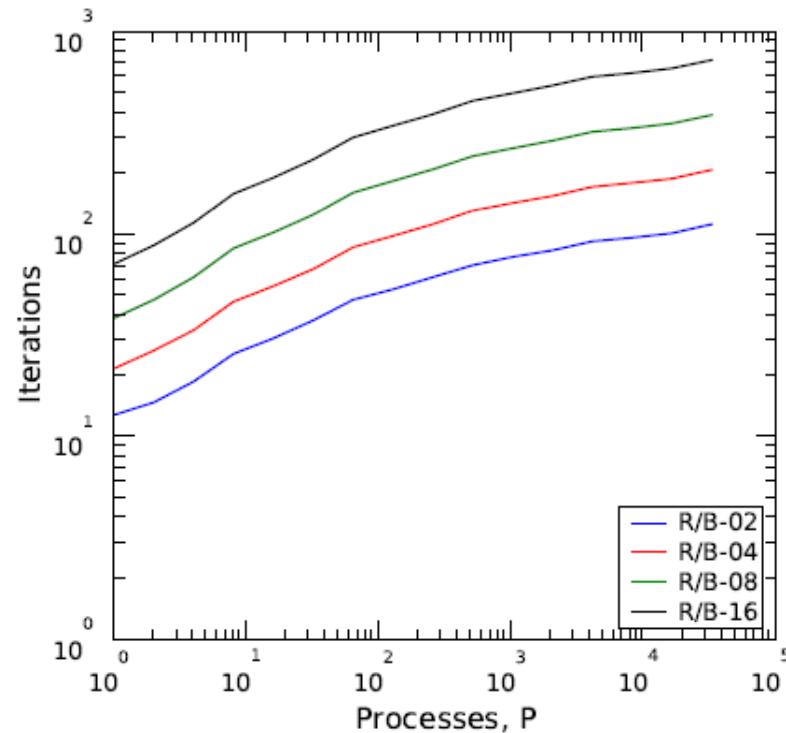
6. Construction vs Solution Time

- Total execution time = construction time + iterative solution time:
 - ❖ Construction time: independent of c & $h \Rightarrow$ average over all cases
 - ❖ Iterative solution depicted for $c = 0.9$ & $h = 0.1$ cm as example



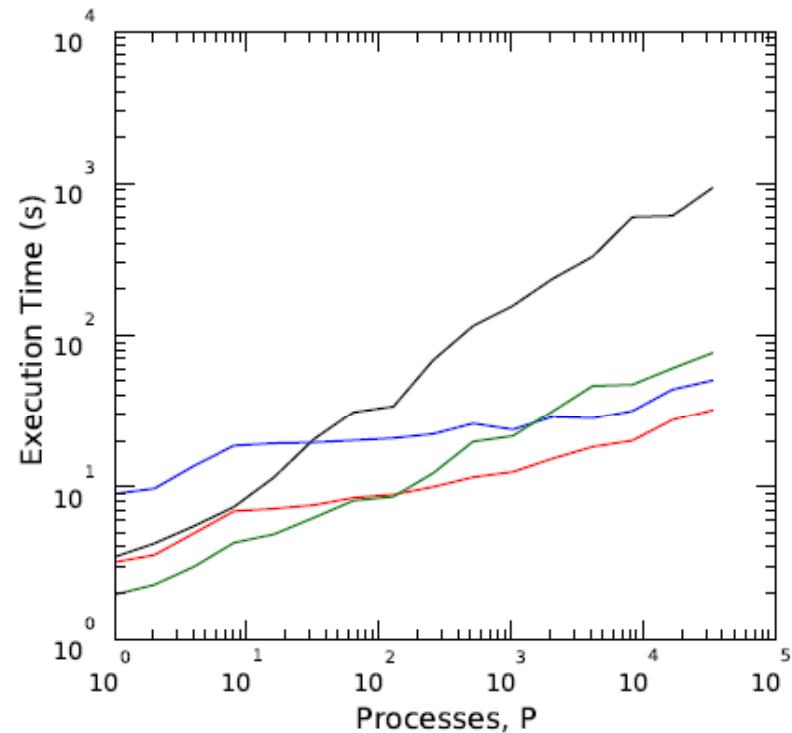
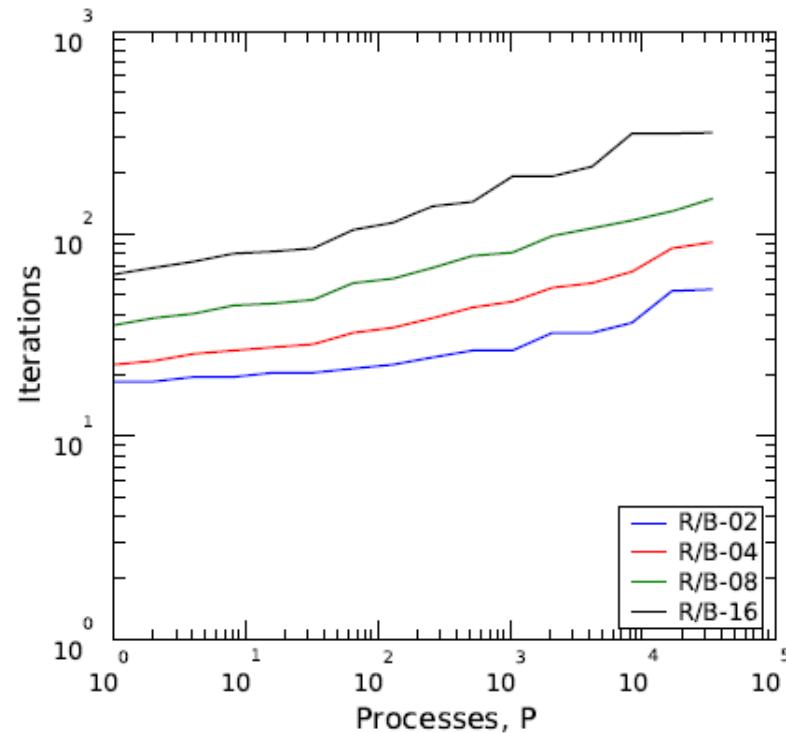
Weak Scaling Results

□ $c=0.9, h=0.1$



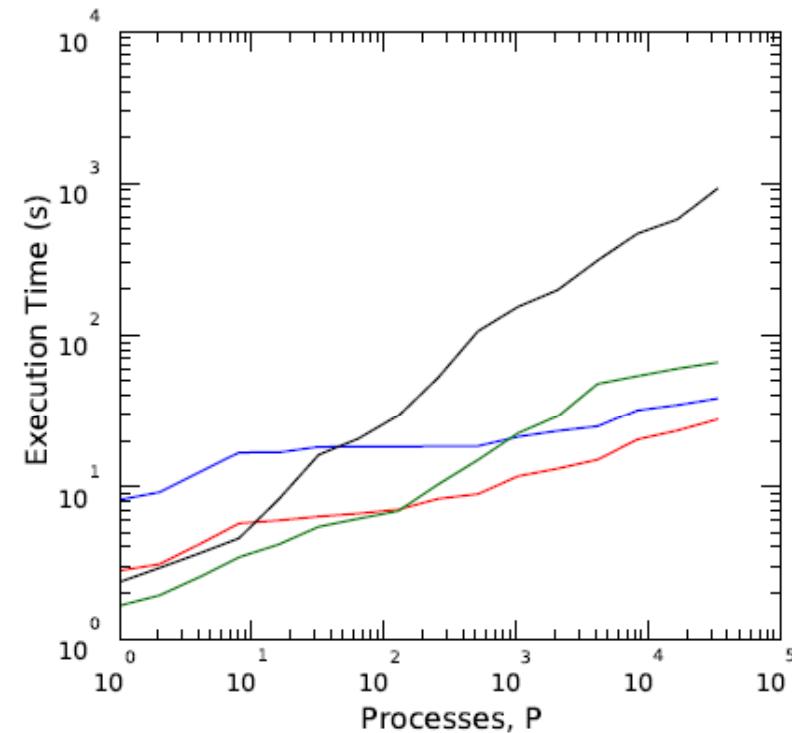
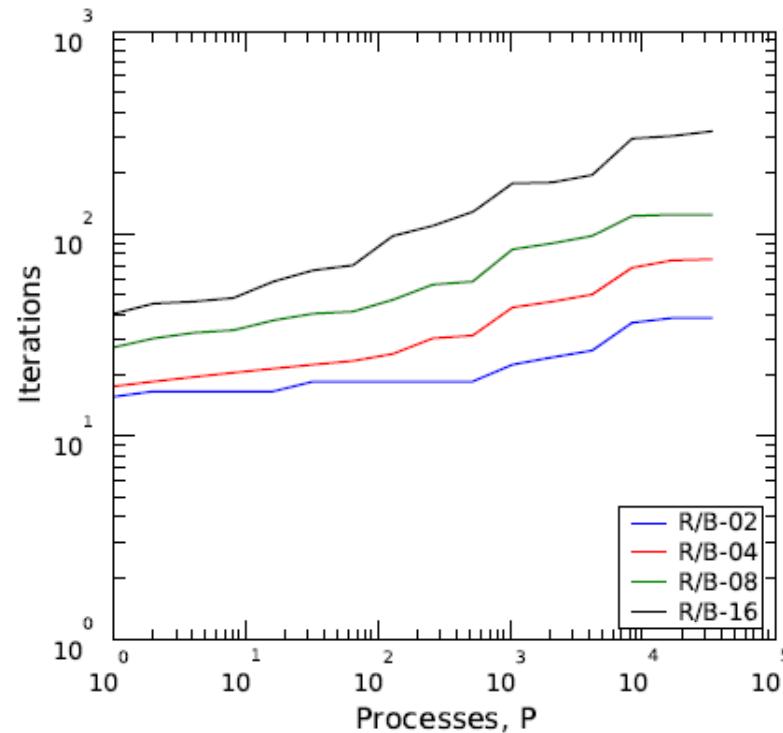
Weak Scaling Results

□ $c=0.9, h=1.0$



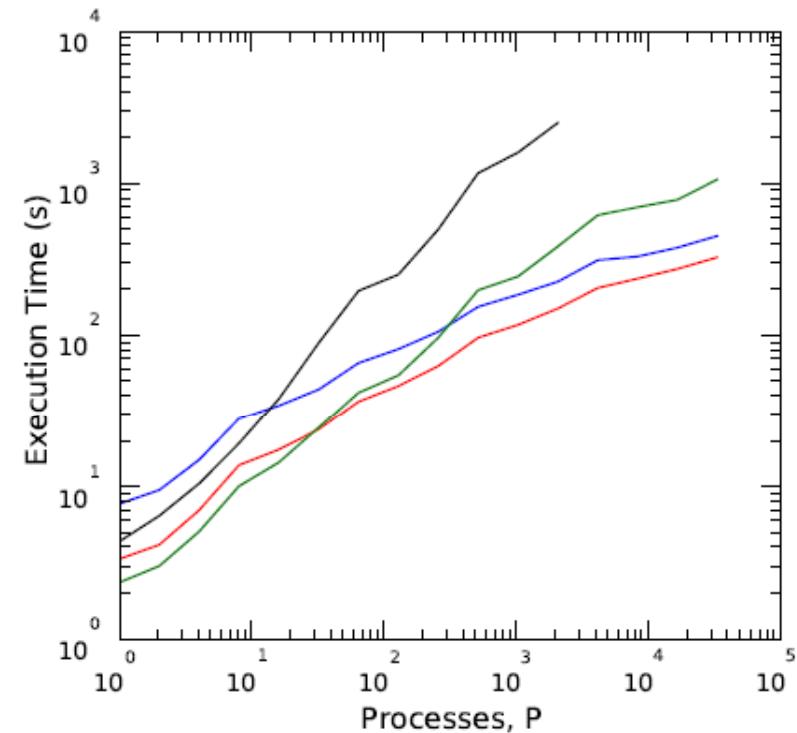
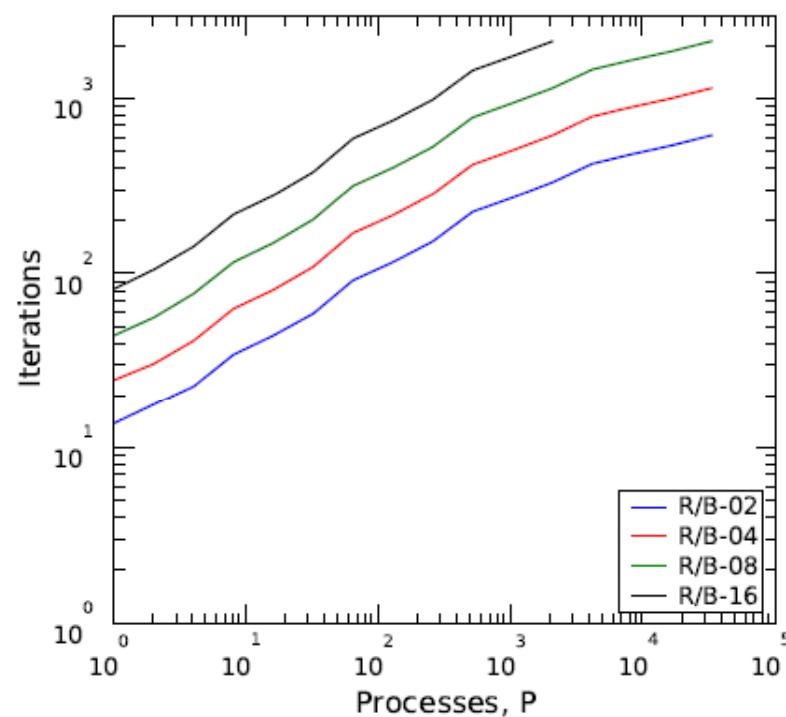
Weak Scaling Results

□ $c=0.9, h=10.0$



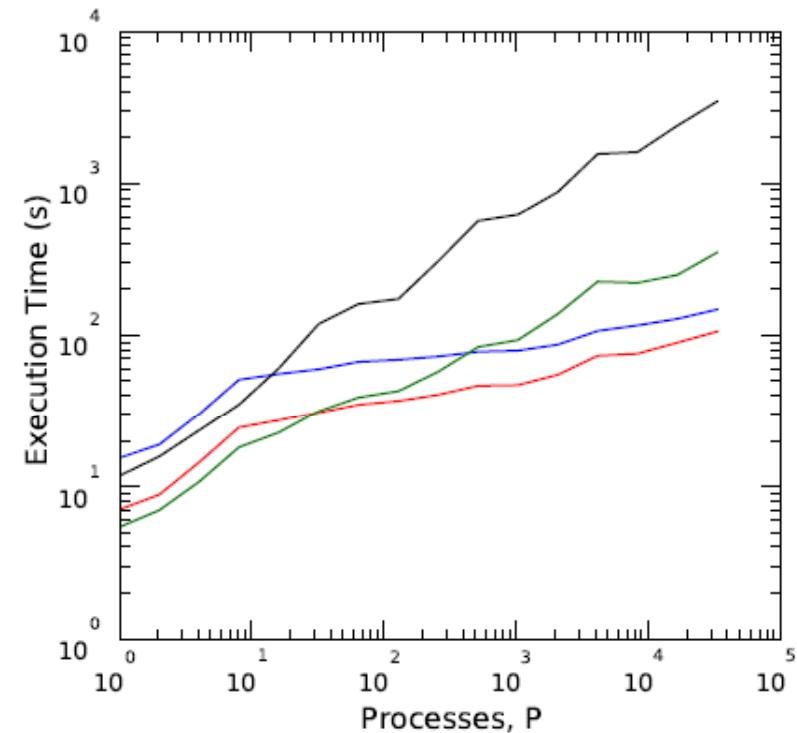
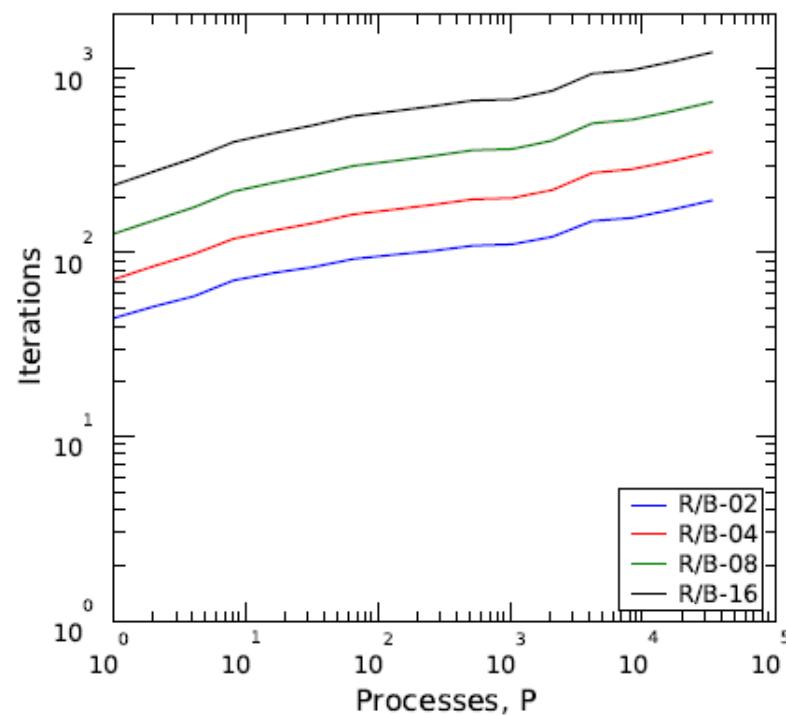
Weak Scaling Results

□ $c=0.99, h=0.1$



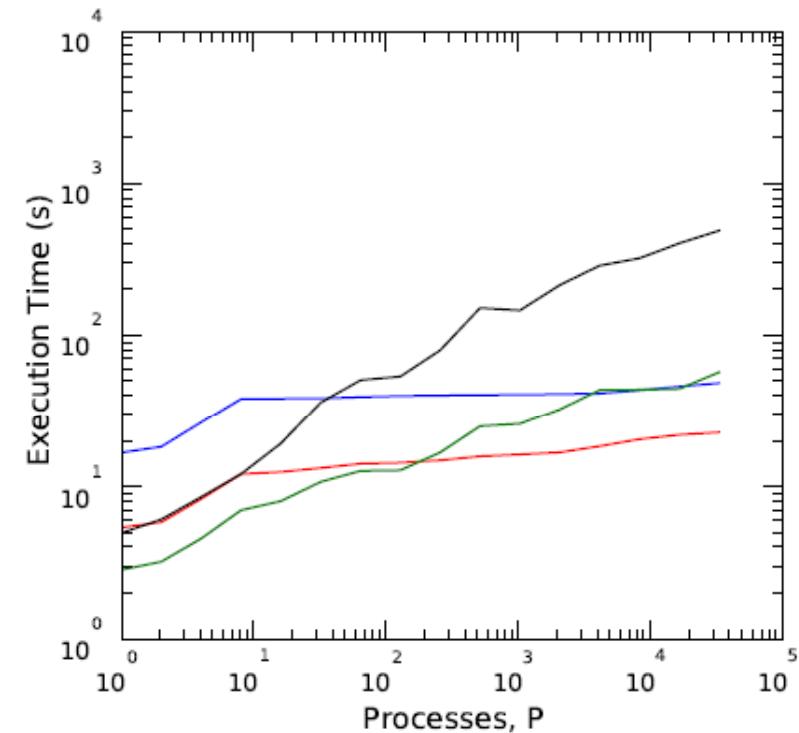
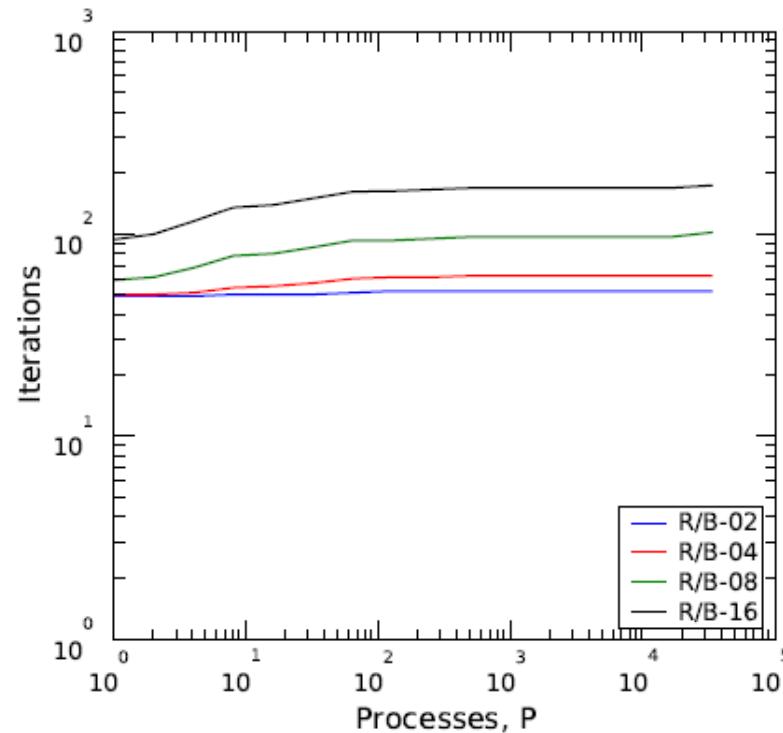
Weak Scaling Results

□ $c=0.99, h=1.0$



Weak Scaling Results

□ $c=0.99, h=10.0$



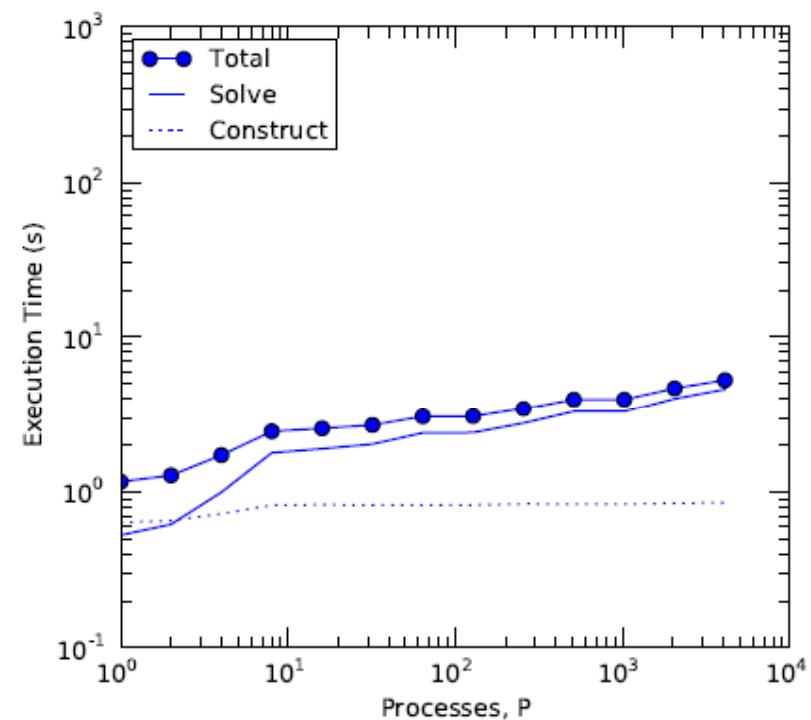
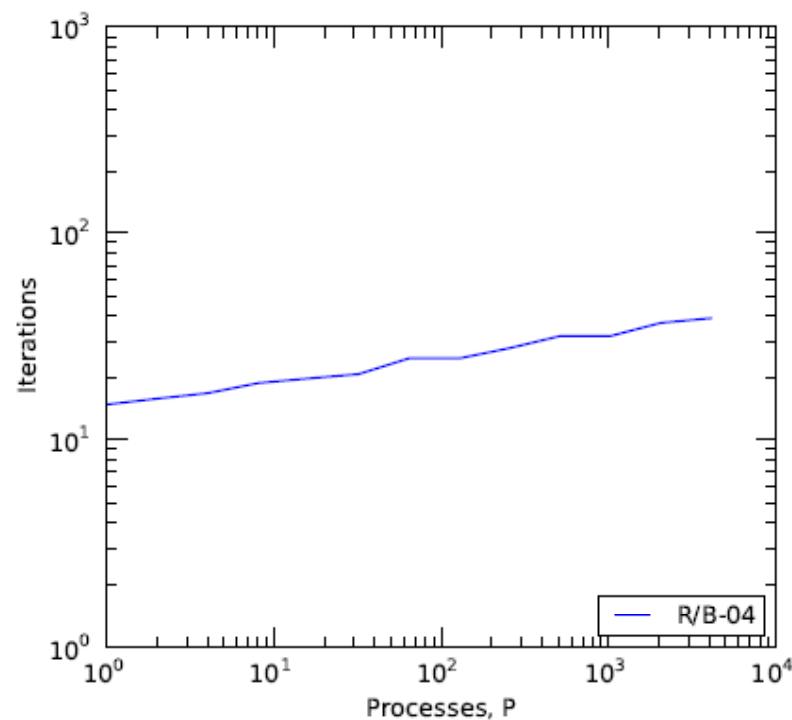
Periodic Heterogeneous Layers

- Alternating layers of optically thin and thick materials
- Known challenge for SI-DSA convergence
- ITMM explicitly couples thick and thin materials
- Starting with $h=1$, increase every other layer by a factor of 'a' and decrease the other layers by a factor of 'a'



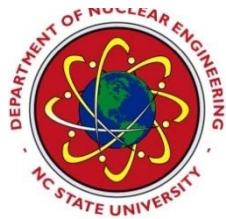
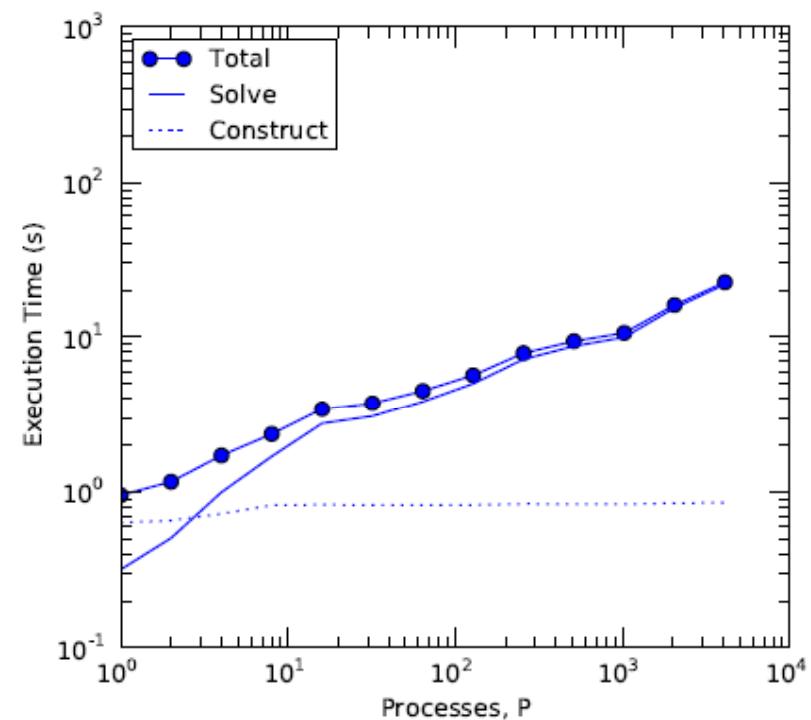
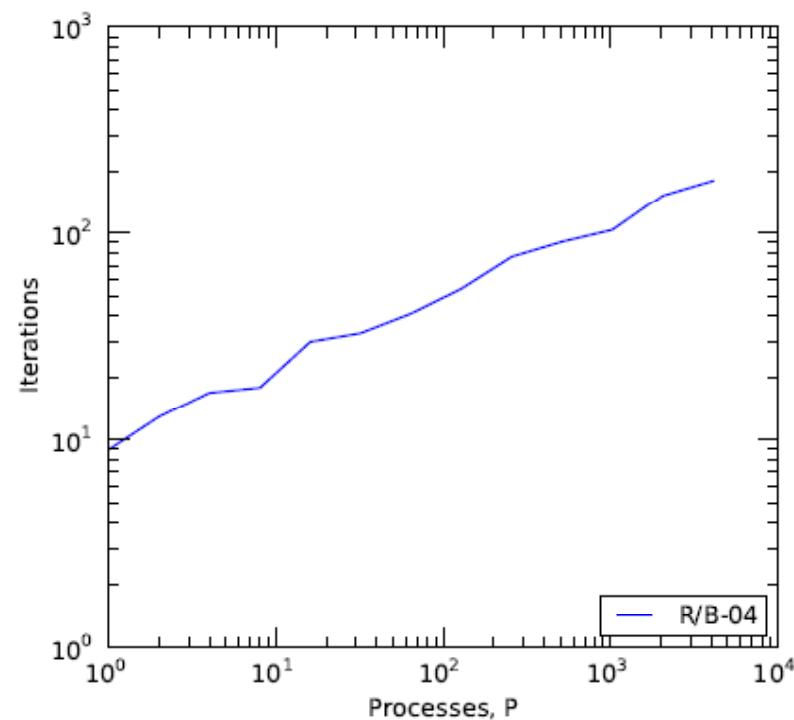
PHL Weak Scaling Results

□ $c=0.9, a=10$



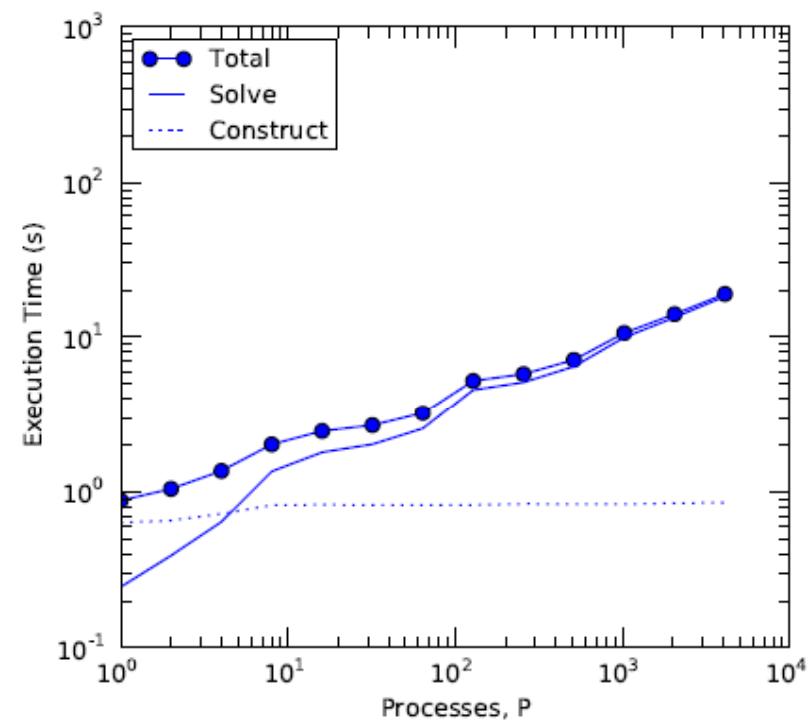
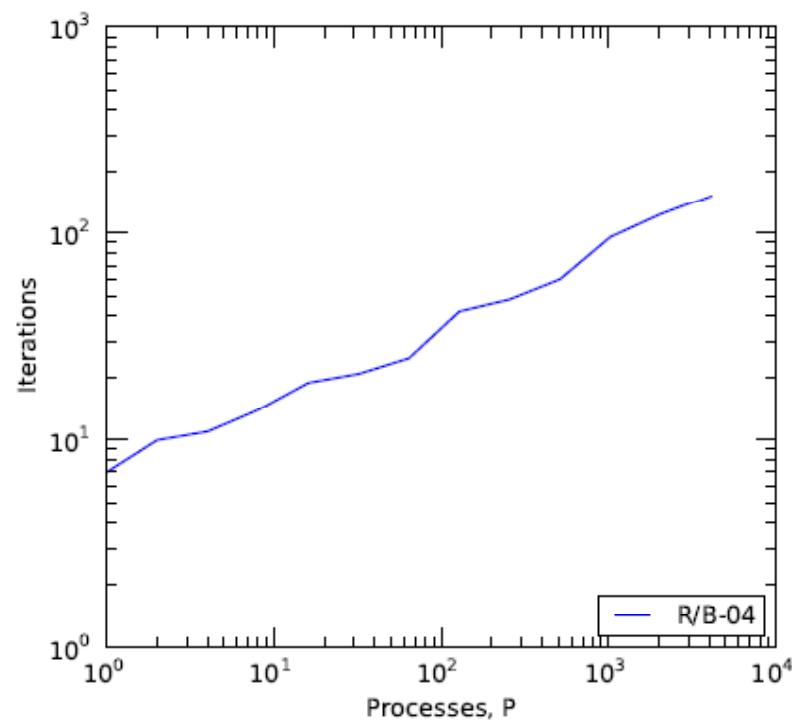
PHL Weak Scaling Results

□ $c=0.9, a=100$



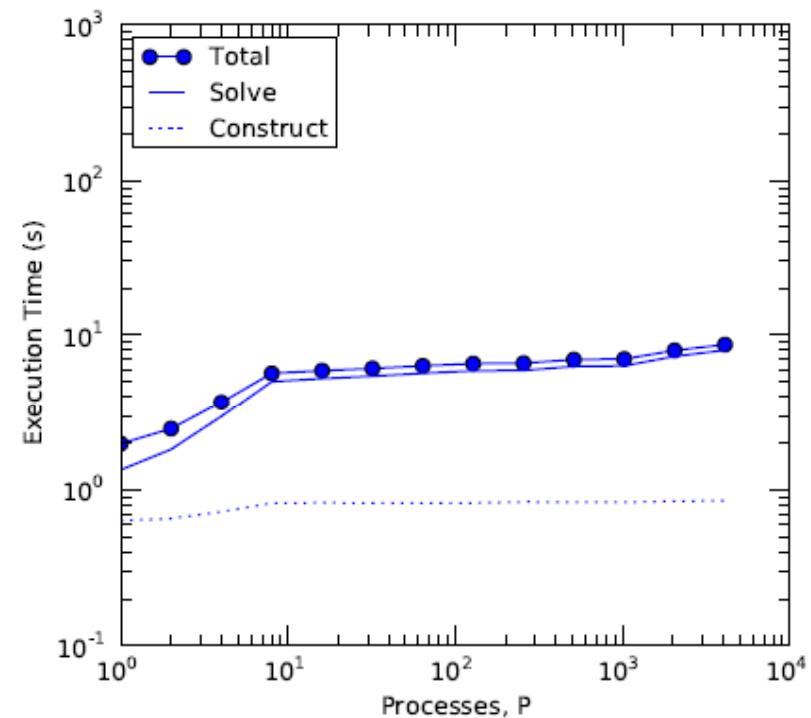
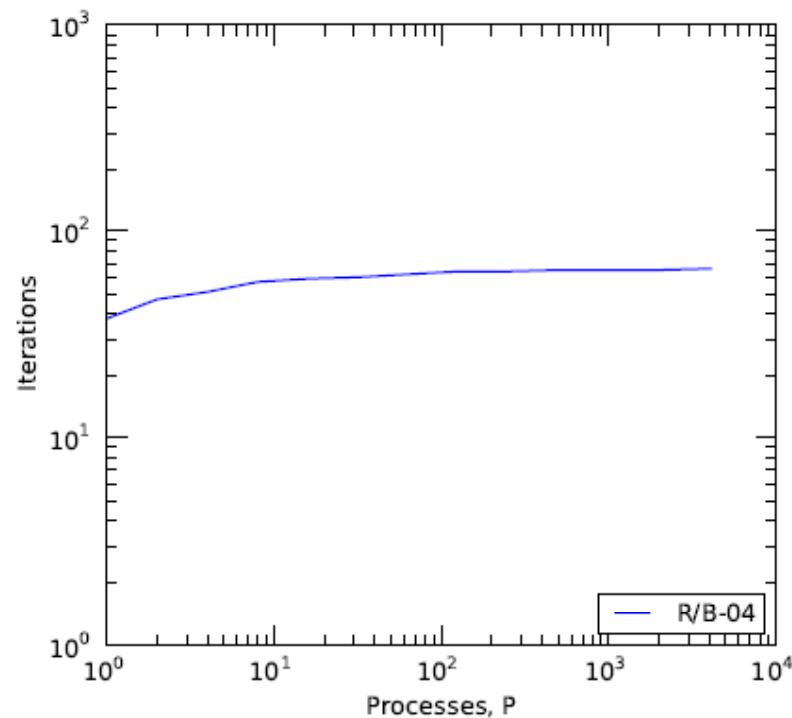
PHL Weak Scaling Results

□ $c=0.9, a=1000$



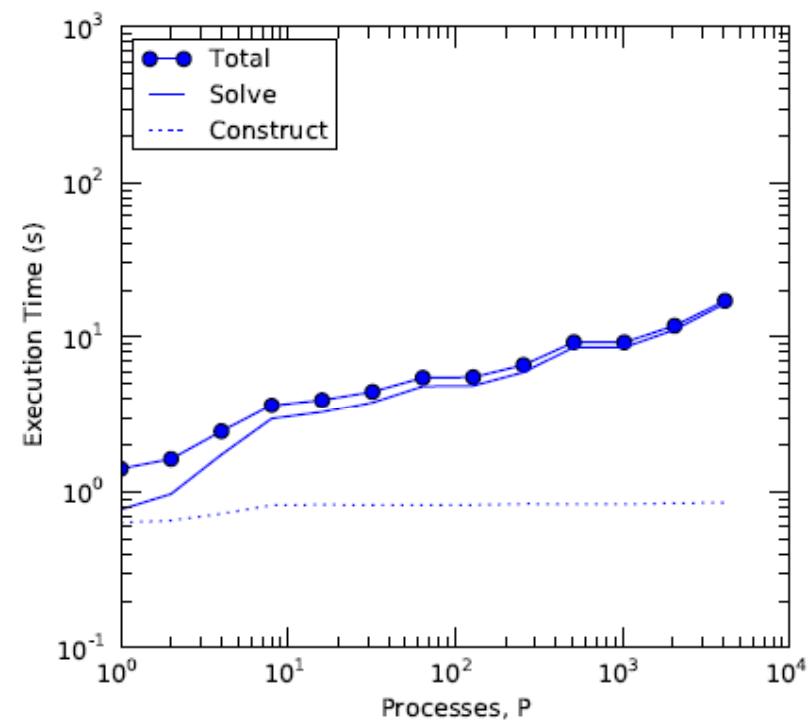
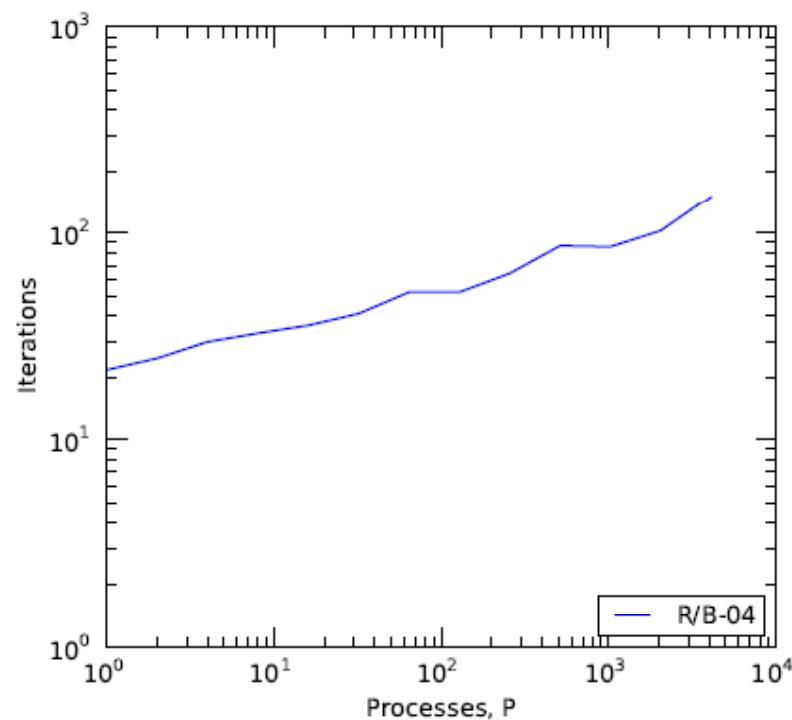
PHL Weak Scaling Results

□ $c=0.99, a=10$



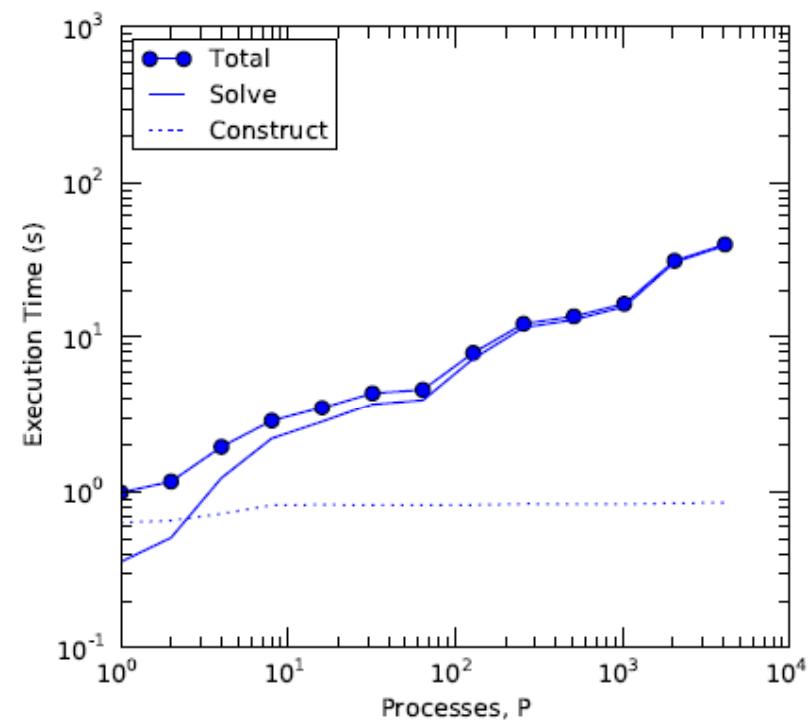
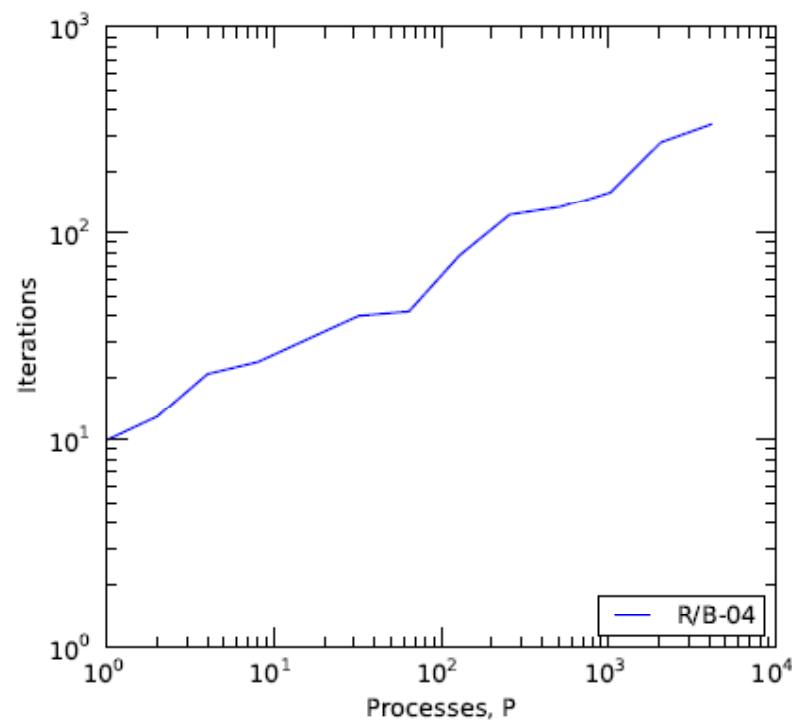
PHL Weak Scaling Results

□ $c=0.99, a=100$



PHL Weak Scaling Results

□ $c=0.99, a=1000$



7. Conclusions



7. Conclusions

- The Nuclear Computational Science Group at NC State is engaged in broad span of topics
- Topic illustrated today: Multiprocessing strategies particularly suited for massively parallel architectures
 - ❖ ITMM SDD avoids sequential mesh-sweeps
 - ❖ Considered BJ & GS parallelizations: PGS bests PBJ
 - ❖ Compared PGS to traditional KBS in PARTISN:
 - Very large differences when SI is accelerated with DSA
 - Gap closes as optical thickness and scattering ratio are increased \Rightarrow most difficult SI problems
 - SI & SI-DSA demonstrate larger growth in execution time as $P \uparrow$
 - Conclusions must be validated for $P > 1,024$
- PGS performance should improve with suitable preconditioner: difficult

