

LA-UR- 11-02358

Approved for public release;
distribution is unlimited.

Title: Boettger-Wallace Phase Transition Model Revisited

Author(s): Jonathan C. Boettger, XCP-5, LANL

Intended for: JOWOG 32 Mat
May 9 - 13, 2011
Los Alamos National Laboratory



Los Alamos National Laboratory, an affirmative action/equal opportunity employer, is operated by the Los Alamos National Security, LLC for the National Nuclear Security Administration of the U.S. Department of Energy under contract DE-AC52-06NA25396. By acceptance of this article, the publisher recognizes that the U.S. Government retains a nonexclusive, royalty-free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government purposes. Los Alamos National Laboratory requests that the publisher identify this article as work performed under the auspices of the U.S. Department of Energy. Los Alamos National Laboratory strongly supports academic freedom and a researcher's right to publish; as an institution, however, the Laboratory does not endorse the viewpoint of a publication or guarantee its technical correctness.

Boettger-Wallace Phase Transition Model Revisited

Jonathan C. Boettger*
X Computational Physics Division
Los Alamos National Laboratory
Los Alamos, New Mexico 87545

The Boettger-Wallace phase transition model is reviewed. The nuts and bolts of implementing the metastability and kinetics models in a simulation code are described in detail. The limitations of the BW model are discussed.

*Work supported by the U.S. Dept. of Energy under contract DE-AC52-06NA25396.

Boettger-Wallace Phase Transition Model Revisited

Jonathan C. Boettger

XCP-5, LANL

JOWOG 32 MAT

Los Alamos, New Mexico

May 9-13, 2011

Motivation

Why talk about work that was completed more than 10 years ago?

The nuts and bolts of implementing the Boettger-Wallace phase transition model (**metastability** and **kinetics**) in a numerical simulation code never has been presented in a comprehensive form.

Metastability Model

Assumptions

- There are two phases, a low P phase (**1**) and a high P phase (**2**).
- The equations of state (**EOSs**) for both phases are known.
- LTE is maintained during transition; **$P = P_1 = P_2$** and **$T = T_1 = T_2$** .
- The instantaneous mass fraction of the 2 phase is denoted by λ .
- The metastable mass fraction of the 2 phase is denoted by λ_m .
- The metastable state is **long-lived** compared to the duration of the experiment.

Metastability Under Loading

Incremental load increase \rightarrow incremental ($\Delta G = G_2 - G_1$) increase $d\Delta G$
 \rightarrow incremental λ_m increase $d\lambda_m$. Material transformed also should be proportional to $(1 - \lambda_m)$. Try

$$B d\lambda_m = (1 - \lambda_m) d\Delta G.$$

Integrating yields the **metastability equation**.

$$\lambda_m = 1 - \exp[1 - \exp(A - \Delta G / B)]$$

$A = (+)$ energy barrier or $(-)$ shear stress enhancement

$B =$ energy width of the transition.

An equivalent expression can be deduced for the reverse transition.

Solving for P , T , and λ_m Given V and E in a Cell

1. **Estimate P , T , and λ_m** – Taken from previous time step.
2. **Determine P and T with λ_m fixed** – Solve four equations in four unknowns (V_1 , V_2 , E_1 , E_2) **iteratively** using Newton's method
$$0 = (1 - \lambda_m) V_1 + \lambda_m V_2 - V$$
$$0 = (1 - \lambda_m) E_1 + \lambda_m E_2 - E$$
$$0 = P_1(V_1, E_1) - P_2(V_2, E_2)$$
$$0 = T_1(V_1, E_1) - T_2(V_2, E_2)$$
3. **Determine λ_m with P and T fixed** – Solve metastability equation.

Problem – **Nested do loops** for each cell and time step. Outer loop over λ_m is unstable unless damped and then has **slow convergence**.

Step 2: Solving the set of 4 equations

Write equations in matrix form as

$$\mathbf{F}(\mathbf{x}) = \mathbf{0}$$

$$\mathbf{F}(\mathbf{x}) = \begin{bmatrix} (1-\lambda)V_1 + \lambda V_2 - V \\ (1-\lambda)E_1 + \lambda E_2 - E \\ P_1 - P_2 \\ T_1 - T_2 \end{bmatrix} \quad \mathbf{x} = \begin{bmatrix} V_1 \\ E_1 \\ V_2 \\ E_2 \end{bmatrix}$$

Make an initial guess for unknowns, \mathbf{x}_0 , and expand $\mathbf{F}(\mathbf{x})$ about \mathbf{x}_0 .

$$\mathbf{F}(\mathbf{x}) = \mathbf{F}(\mathbf{x}_0) + \mathbf{A} \cdot (\mathbf{x} - \mathbf{x}_0)$$

Step 2: Solving the set of 4 equations

$$\mathbf{A}(\mathbf{x}) = \begin{bmatrix} (1-\lambda) & 0 & \lambda & 0 \\ 0 & (1-\lambda) & 0 & \lambda \\ \frac{\partial P_1}{\partial V_1} & \frac{\partial P_1}{\partial E_1} & -\frac{\partial P_2}{\partial V_2} & -\frac{\partial P_2}{\partial E_2} \\ \frac{\partial T_1}{\partial V_1} & \frac{\partial T_1}{\partial E_1} & -\frac{\partial T_2}{\partial V_2} & -\frac{\partial T_2}{\partial E_2} \end{bmatrix}$$

Slow for tabular EOS

Solve for \mathbf{X} .

$$\mathbf{x} = \mathbf{x}_0 - \mathbf{A}^{-1} \cdot \mathbf{F}(\mathbf{x}_0).$$

Now replace \mathbf{x}_0 with \mathbf{x} and iterate to convergence.

Step 2 is slow but the outer do loop over λ can be accelerated.

Accelerating the Outer Loop Over λ_m

Use an iterative procedure similar to Newton's method

1. Input $\lambda_{i1} = \lambda_m$ from last time step (or 0) \rightarrow output = λ_{o1} .
2. Input $\lambda_{i2} = 0.5 (\lambda_{i1} + \lambda_{o1}) \rightarrow$ output = λ_{o2} .
3. Assume $\lambda_o = \alpha + \beta \lambda_i$, with α and β from last two iterations.
Input λ_{i3} from intercept with $\lambda_o = \lambda_i \rightarrow$ output = λ_{o3} .
4. Now iterate to convergence, $\Delta\lambda < 0.01$.

In most cases this procedure converges after 3 iterations.

The overall metastability calculation is still very slow.

Metastability Model

Some observations

- Must track whether each cell is **loading or unloading**. I tracked V using a tolerance of $\Delta V = 0.00005$ to avoid spurious switching due to ringing or numerical noise; i.e., reversal by less than ΔV is ignored.
- Calculations are much faster for analytic EOS than for tabular EOS, but are still **slow** compared to single phase EOS look-up.

Metastability Model

Could make metastability tractable with tabular EOS

- Pre-calculate $\lambda_m(V,E)$, $P_m(E,V)$, and $T_m(E,V)$ for loading and unloading on a dense mesh, then interpolate using standard tabular EOS software.

Phase Transition Kinetics

Characteristic assumptions of BW model

- System drives toward **metastable** state not equilibrium state.
- The transition rate is **linear** in distance from metastable state.
- Driving force is $\Delta G(P,T) \rightarrow$ distance measured at **fixed (P,T)**.

This gives the **kinetics equation**.

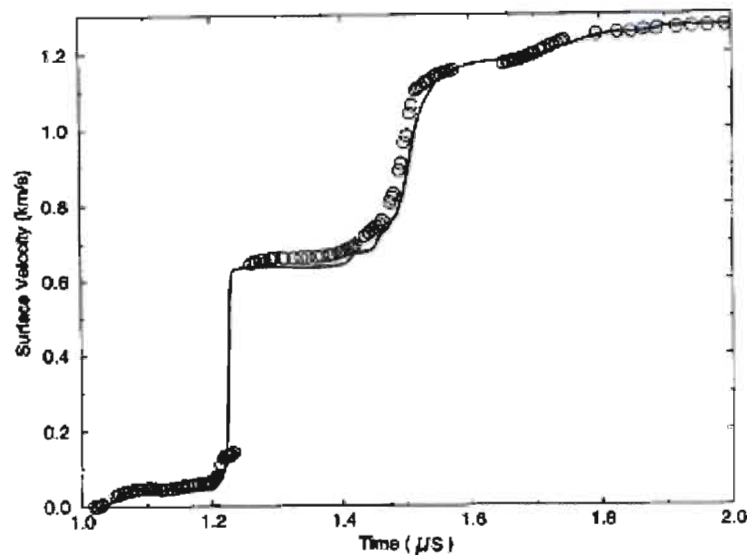
$$\frac{d\lambda}{dt} = \frac{\lambda_m(P,T) - \lambda(P,T)}{\tau}$$

Phase transition kinetics

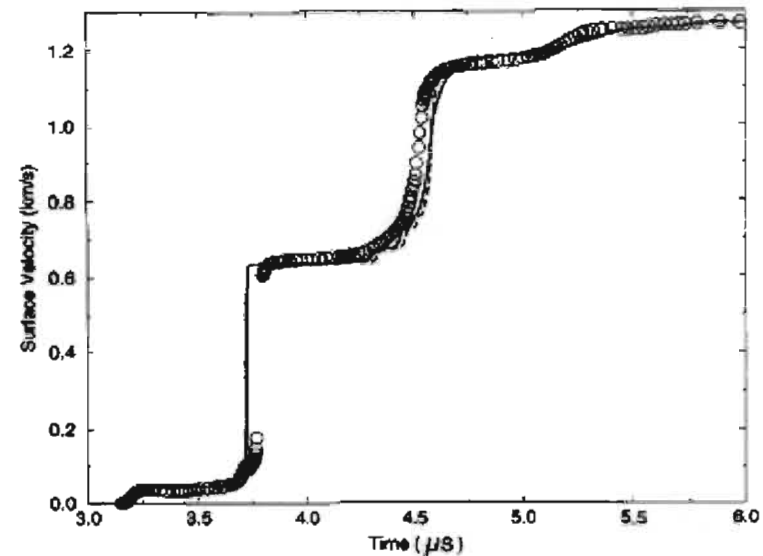
Some observations

- The maximum amount of material transformed in a given time step Δt is calculated as $\Delta\lambda = (d\lambda/dt)_{\text{avg}} \Delta t$, where $(d\lambda/dt)_{\text{avg}}$ is an average of rates calculated before and after the time step.
- Adding kinetics to the metastability model **does not slow down** the simulations.
- The kinetics model is however **fundamentally flawed**.

Application to Fe free surface experiments



Peak stress 23.6 GPa, sample thickness 6.370 mm – Best fit with $\tau = 30$ ns.



Peak stress 23.7 GPa, sample thickness 19.14 mm – Best fit with $\tau = 50$ ns.

Phase transition kinetics

A major problem with the kinetics equation

- For Fe, the relaxation time τ appears to depend on **peak stress** and **sample thickness**, and has to be fitted for each experiment.
- τ is not a constant as in BW model or a thermodynamic function $\tau(P,T)$ as in some other models.
- τ is a **constitutive property** that depends on microstructure and strain rate.