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Title: Crystal plasticity based constitutive laws for Be:
effect of strain rate, twinning and texture on the mechanical
response

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Crystal plasticity based constitutive laws for Be: effect of strain rate, twinning and texture on the mechanical response

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Abstract

Because of its intrinsic anisotropy and variety of slip and twin modes, Beryllium displays a mechanical behavior that strongly depends upon texture, temperature and strain rate. As a consequence, accounting for crystal plasticity mechanisms and directionality is imperative for developing reliable constitutive laws for Be. In this presentation we describe a self-consistent polycrystal model and apply it to the prediction of internal stress of thermal origin, mechanical response of textured and non-textured Be, deforming in compression at strain rates ranging from 10^{-3} s^{-1} to 10^4 s^{-1} . We also discuss the possibility of measuring dislocation densities using the technique of Peak Shape Analysis.

To be presented at Khariton Readings, March 14-18, 2011, Sarov, Russia

**Crystal plasticity based constitutive laws for Be:
effect of strain rate, twinning and texture on the
mechanical response**

Carlos Tome¹, Irene Beyerlein², Donald Brown¹

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² Theoretical Division

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Visco-plastic grain model

Grain deformation is based on crystallographic shear on slip and twinning systems.

The **shear-rate** of each system s is given by a viscous (rate sensitive) law

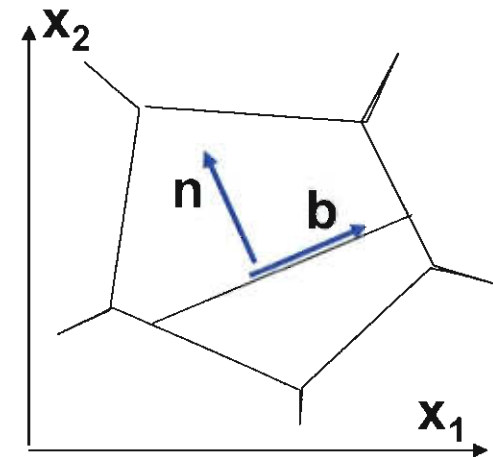
$$\dot{\gamma}^s = \dot{\gamma}_0 \left(\frac{\tau_{\text{resolved}}}{\tau^s} \right)^n = \dot{\gamma}_0 \left(\frac{\mathbf{m}^s : \boldsymbol{\sigma}}{\tau^s} \right)^n$$

The threshold stress describes the hardening of slip systems

The crystallographic shears determine crystal reorientation and texture evolution

The **strain-rate** of the grain is given by superposition of shear rates on all active systems

$$\dot{\epsilon}'_{ij} = \sum_s \frac{n_i^s b_j^s + n_j^s b_i^s}{2} \dot{\gamma}^s = \sum_s m_{ij}^s \dot{\gamma}^s$$



Visco Plastic Self Consistent (VPSC) Polycrystal Model

- Each grain is a visco-plastic **anisotropic ellipsoidal inclusion** embedded in a visco-plastic anisotropic Homogeneous Effective Medium (HEM).
- Linearized constitutive response: M^g and \bar{M} \rightarrow grain and effective medium compliance tensors

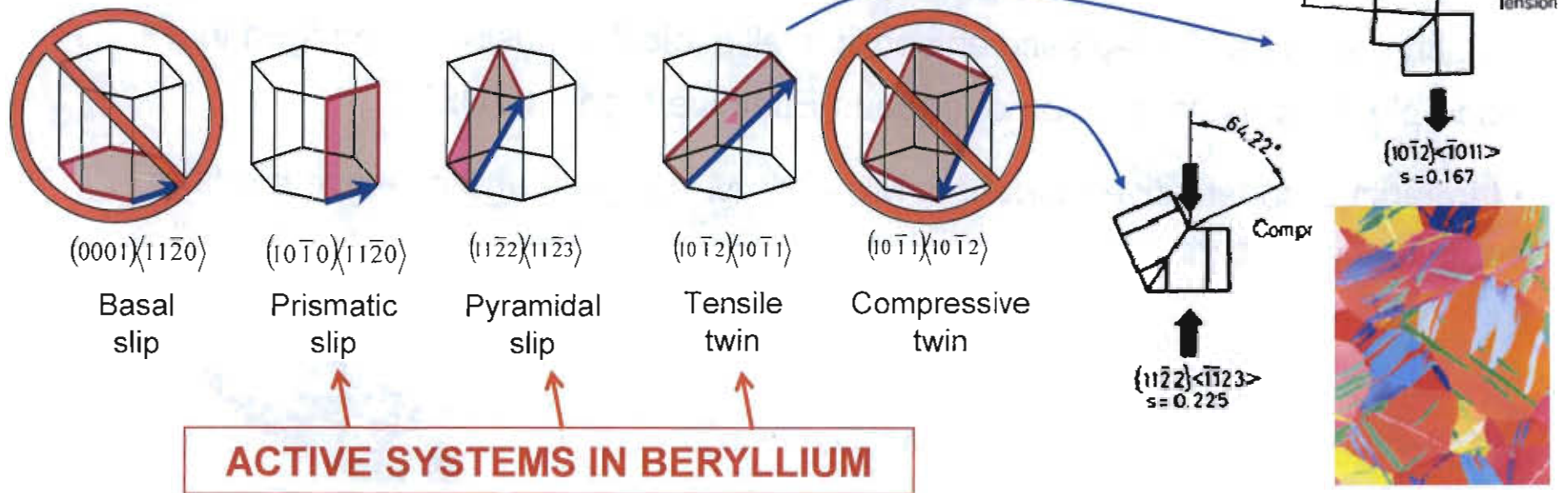
grain: $\dot{\epsilon}_{ij}^g = \dot{\gamma}_0 \sum_s m_{ij}^s \left(\frac{m^s : \sigma^g}{\tau^s} \right)^n = M_{ijkl}^g \sigma_{kl}^g + \dot{\epsilon}_{ij}^{og}$

medium: $\bar{\dot{\epsilon}}_{ij} = \bar{M}_{ijkl} \bar{\sigma}_{kl} + \bar{\dot{\epsilon}}_{ij}^o$

Solve equilibrium equation for inclusion in homogeneous medium: $\sigma_{ij,j} = 0$

Eshelby result \rightarrow stress and strain-rate are uniform inside the inclusion
 ...but different from the macroscopic stress and strain rate !

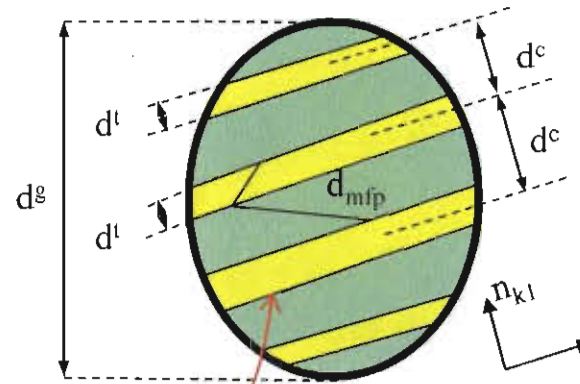
Plasticity of hexagonal materials



MECHANICAL RESPONSE OF HEXAGONALS

- Several slip and twin modes with different temperature and rate dependencies
- Anisotropy of single crystal \rightarrow difficult to deform along the c-axis
- Need to account for slip-slip and for slip-twin interactions
- Twinning activity affects texture, hardening and anisotropy

Composite-Grain Twin Model



Treat grain as an evolving stack of twin-matrix layers on the plane of the predominant twin system.

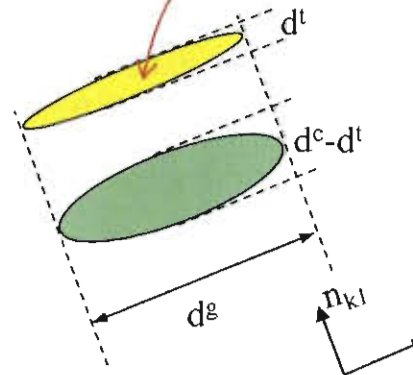
COUPLED

Continuity of stress and strain across the twin-matrix interface

$$\sigma_{13}^{parent} = \sigma_{13}^{twin}$$

$$\sigma_{23}^{parent} = \sigma_{23}^{twin}$$

$$\dot{\epsilon}_{12}^{parent} = \dot{\epsilon}_{12}^{twin}$$



UNCOUPLD

Separate treatment of parent and twin as ellipsoidal inclusions.

Account for ellipsoid orientation and thickness evolution with deformation

Proust & Tome, Acta Materialia **55** (2007) 2137
Proust et al, Int. J. of Plasticity **25** (2009) 861

Rate Dependent Experiments Prompt Dislocation Based Hardening Model

(α = slip mode)

$$\frac{\partial \rho^\alpha}{\partial \gamma^\alpha} = \frac{\partial \rho_{storage}^\alpha}{\partial \gamma^\alpha} - \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} = k_1^\alpha \sqrt{\rho^\alpha} - k_2^\alpha(\dot{\epsilon}, T) \rho^\alpha$$

$$\frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} \begin{cases} \nearrow (1 - f^\alpha) \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} & \text{ANNIHILATION} \\ \searrow f^\alpha \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} & \text{SUBSTRUCTURE } \rho_{sub} \end{cases}$$

- Temperature and strain rate dependence enter in rate of annihilation.

Dislocation-based constitutive law for HCP (temperature and rate dependent)

Dislocation densities in the slip systems (α = basal, prism or pyramidal) evolve with crystal shears via generation, annihilation (removal) and substructure build-up

$$d\rho_{\text{gener, forest}}^{\alpha} = k_1^{\alpha} \sqrt{\rho_{\text{forest}}^{\alpha}} d\gamma^{\alpha} \quad \leftarrow \text{Dislocation generation}$$

$$d\rho_{\text{remov, forest}}^{\alpha} = k_2^{\alpha} (\dot{\epsilon}, T) \rho_{\text{forest}}^{\alpha} d\gamma^{\alpha} \quad \leftarrow \text{Dislocation annihilation}$$

$$d\rho_{\text{substructure}} = \sum_{\alpha} f^{\alpha}(T) d\rho_{\text{remov, forest}}^{\alpha} \quad \leftarrow \text{Substructure formation}$$

The threshold to propagate dislocations is related to the density via Taylor's law

$$\tau_{\text{slip}}^{\alpha} = b^{\alpha} \chi \mu \sqrt{\sum_{\alpha'} D^{\alpha\alpha'} \rho_{\text{forest}}^{\alpha'}} \quad \leftarrow \text{Forest hardening}$$

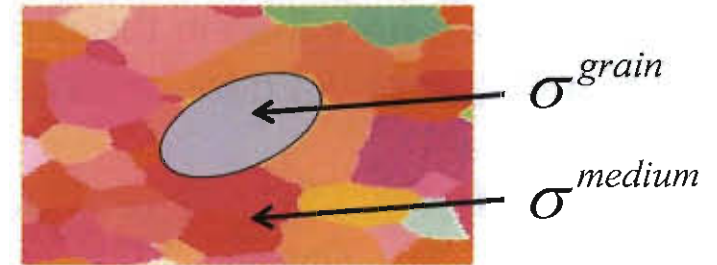
$$\tau_{\text{twin}}^{\beta} = \mu \sum_{\alpha} C^{\beta\alpha} b^{\alpha} b^{\beta} \rho^{\alpha} \quad \leftarrow \text{Twin hardening by dislocations}$$

Multiscale modeling

Visco Plastic (VPSC) and Elasto Plastic (EPSC) Polycrystal Aggregate Models

VPSC → Lebensohn & Tome (1993); *Proust et al (2009)*

EPSC → Turner & Tome (1994); *Clausen et al (2008)*

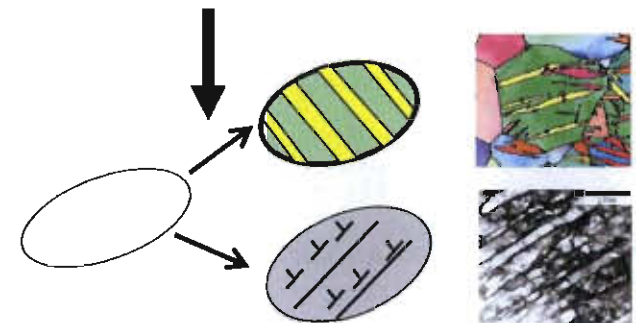


Composite Grain Model

Proust et al. (2007)

Dislocation-based Hardening Model

Beyerlein & Tome (2008)



Shear rates & hardening for slip & twinning
controlled by evolving Threshold Stress

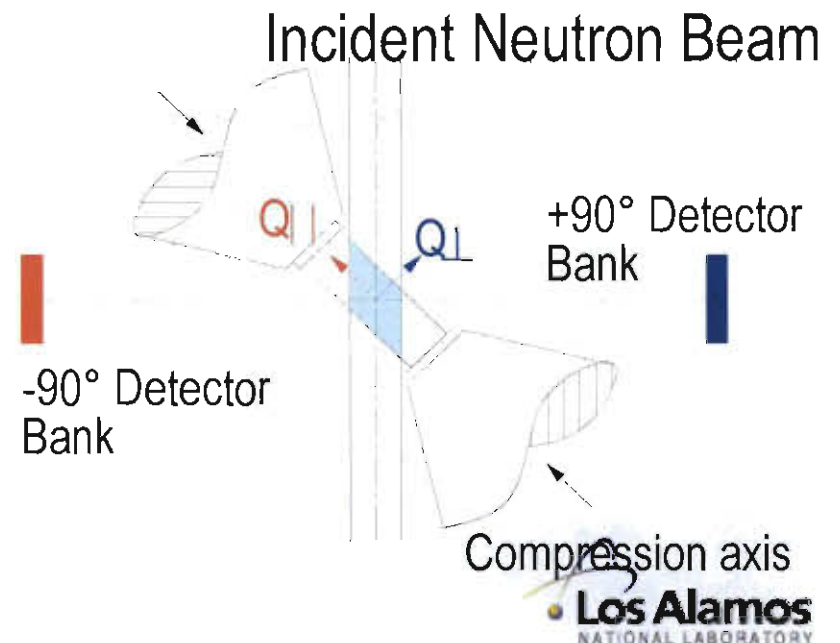
$$\dot{\gamma}^s = \dot{\gamma}_0 \left(\frac{\tau_{resolved}}{\tau_{threshold}^s} \right)^n$$

Neutron diffraction facilities: a tool for 'in-situ' & 'in-bulk' stress measurements

SMARTS Diffractometer
LANSCE - LANL



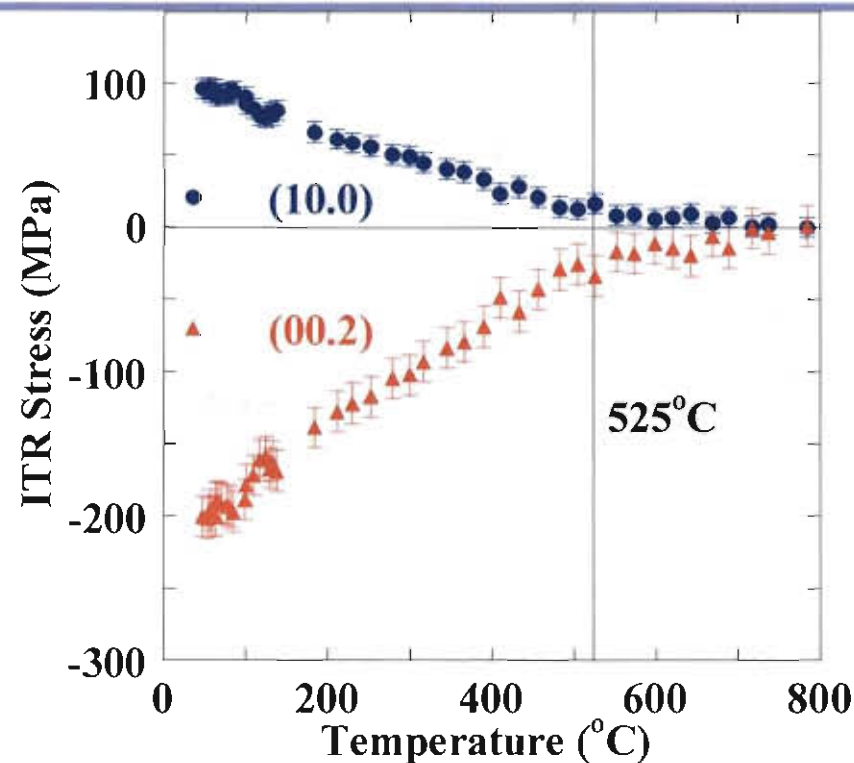
- Crystal diffraction measures lattice spacing in the bulk
- Lattice spacing provides information about elastic strains
- Elastic strains are related to internal stress via Hooke's law
- Diffracting region: $\sim 5\text{mm}$



Modeling of Thermal and Mechanical Response of Be

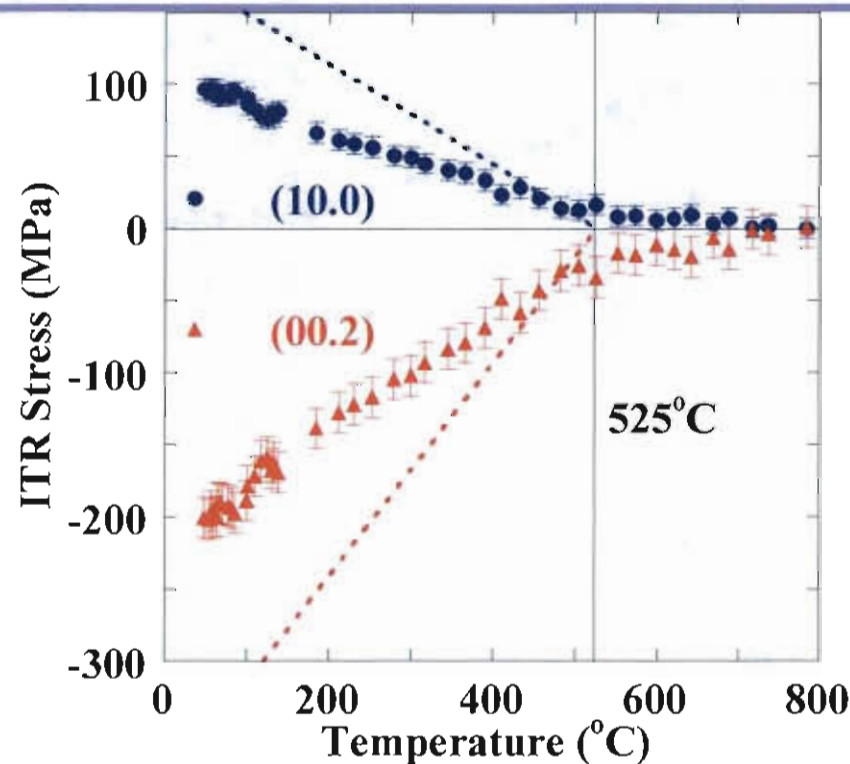
- Use Elasto-Plastic Polycrystal model to calculate internal stress inside grains of non-textured Be.
→ Compare with neutron diffraction results
- Use Visco-Plastic Polycrystal model to calculate Stress-Strain response of rolled and non-textured Be.
→ Use an unique Single Crystal constitutive law, function of temperature and strain-rate
- Calculate dislocation densities using Peak Profile Analysis on neutron diffraction peaks

Intergranular Thermal Residual Stresses in Beryllium are Significant.



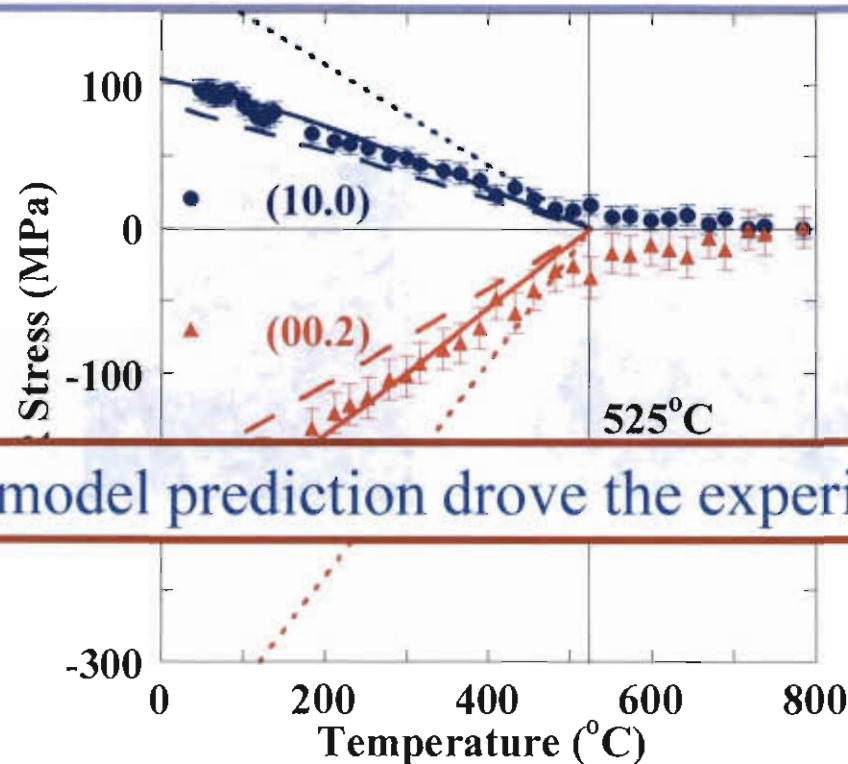
- Intergranular stresses relax on time scale of measurement above 525C.
 - Zero strength temperature.
 - Important for modeling studies.

Intergranular Thermal Residual Stresses in Beryllium are Significant.



- Upper bound model, assuming grain in rigid matrix over-predicts stresses by $\sim x2$

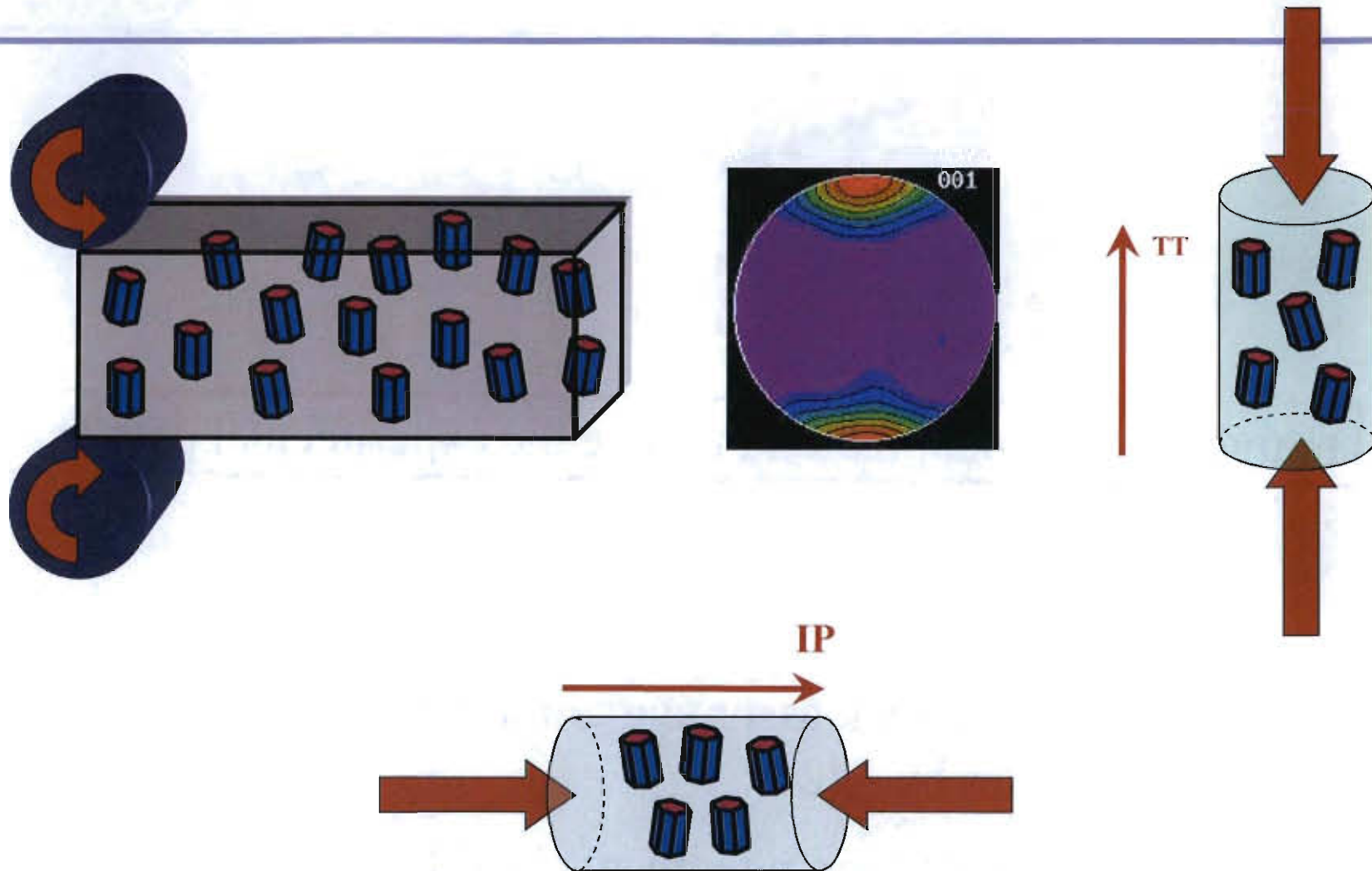
Intergranular Thermal Residual Stresses in Beryllium are Significant.



In this case, model prediction drove the experimental path !

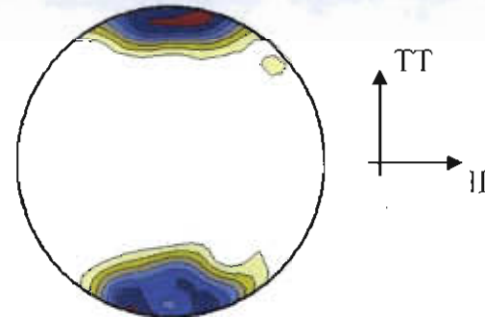
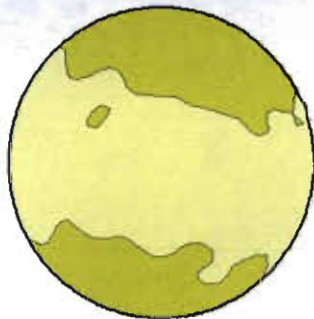
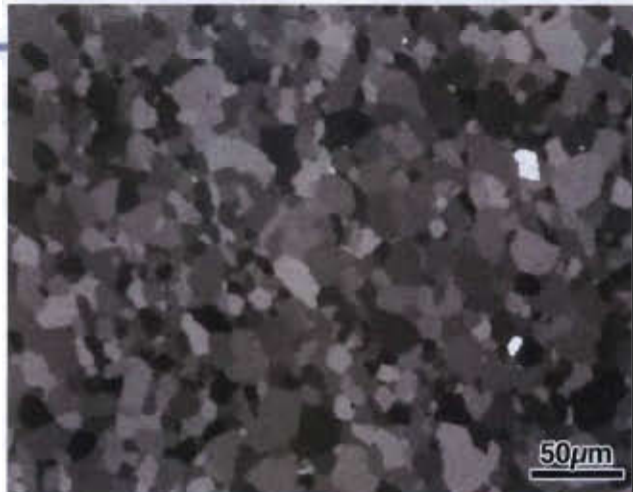
- Upper bound model, assuming grain in rigid matrix over-predicts stresses by $\sim \times 2$.
- EPSC model with room temperature CTE's under-predicts stresses slightly.
- EPSC with temperature dependent CTE matches data well.

Rolling of Beryllium Results in Strong Crystallographic Texture



- Orientation of crystals with respect to load axis is critical.

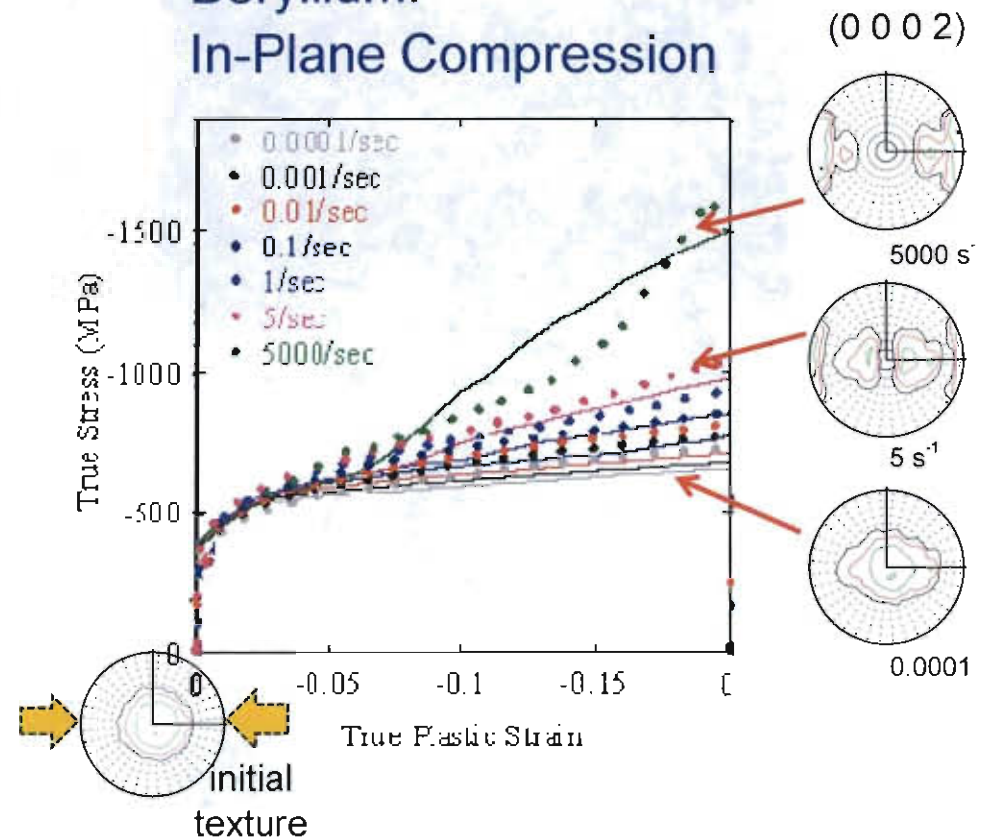
Rate and Temperature Dependent Deformation of Beryllium



- Hot-pressed (random texture) and rolled plate (strong texture) beryllium.
- In-plane (IP) and through-thickness (TT) compression of rolled plate.
 - Normal and parallel to dominant basal texture, respectively.
- Strain rates from 0.0001/sec to 5000/sec.

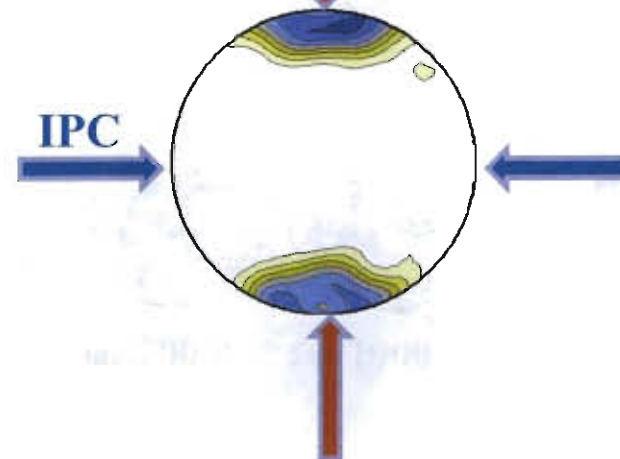
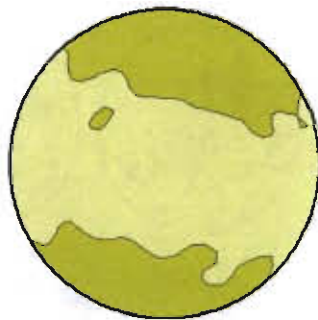
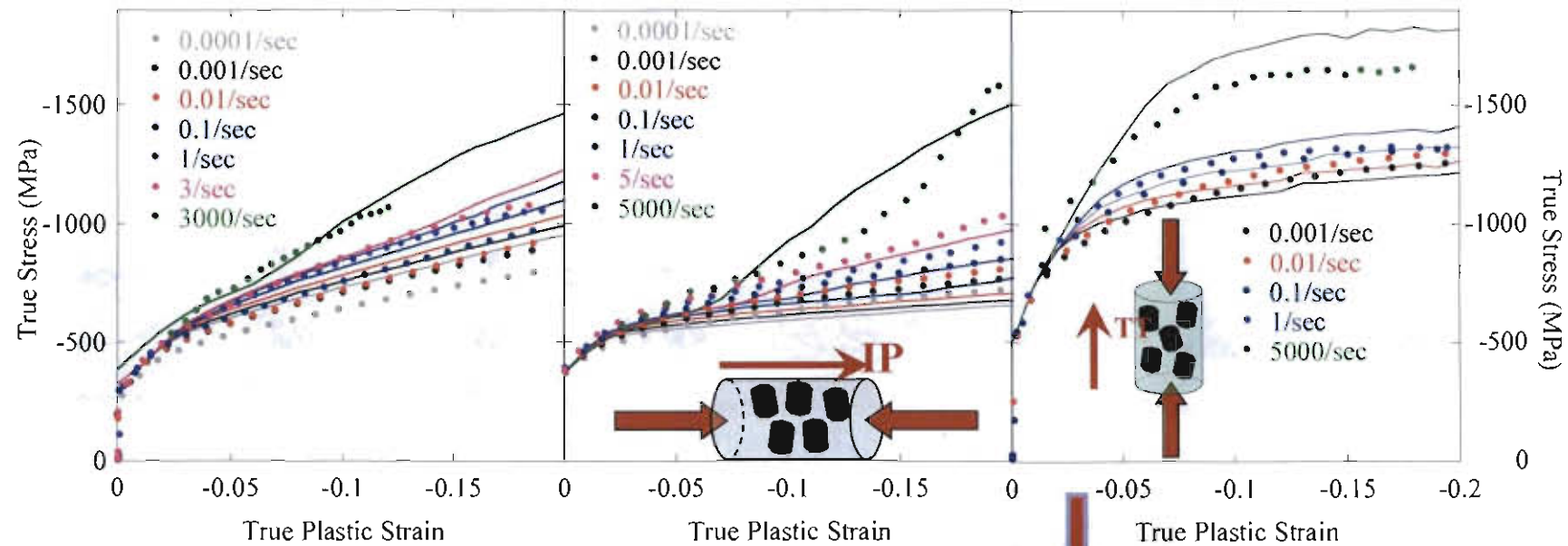
Be: slip vs twinning competition

Beryllium: In-Plane Compression



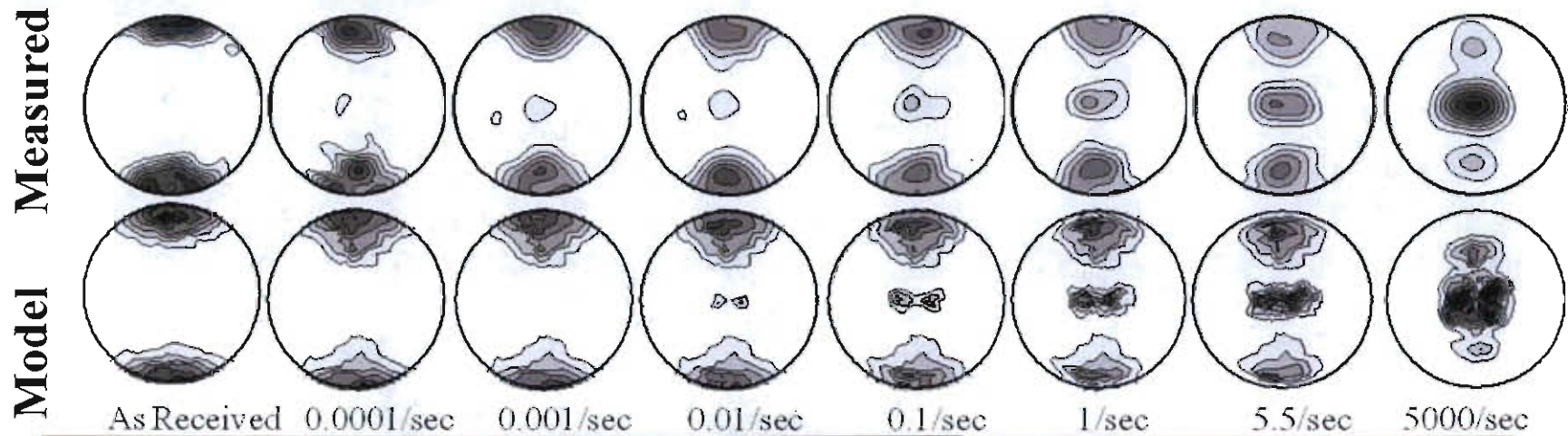
Rapid transition from dislocation-controlled to twin-controlled hardening as the temperature decreases (Zr) or as the strain rate increases (Be)

Dislocation Based Hardening Law Allows Us to Predict Rate Dependent Flow curve

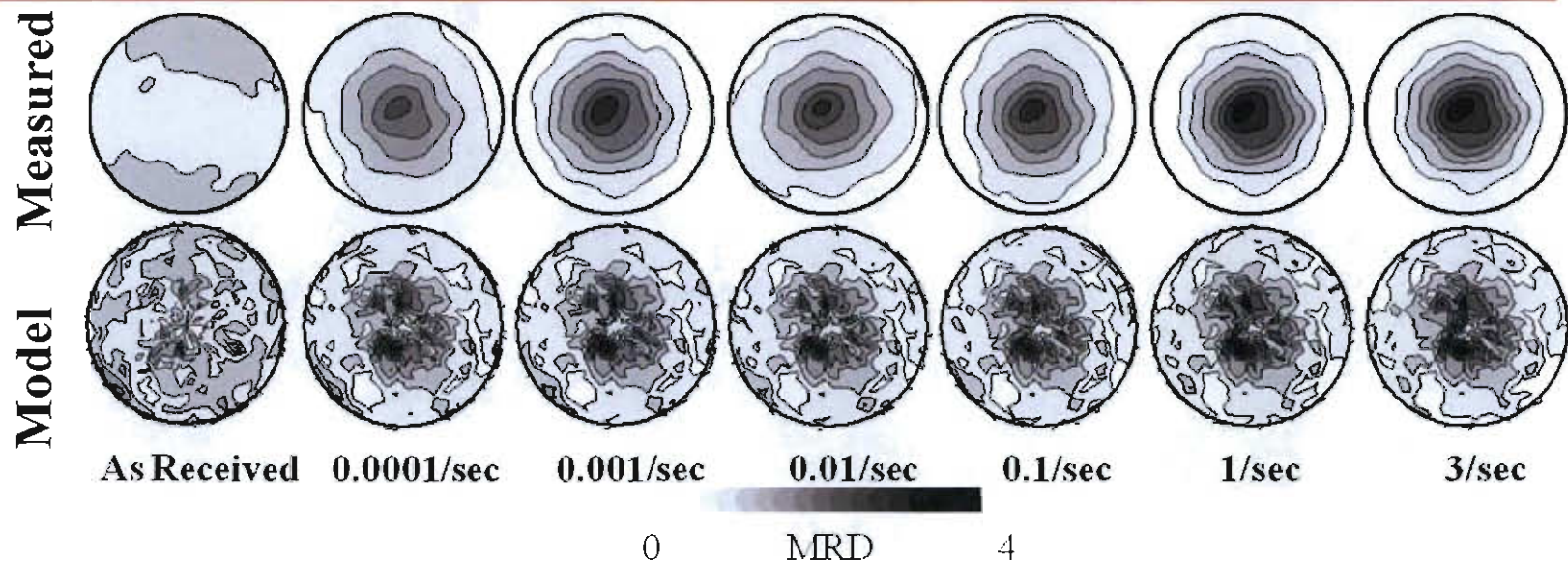


Dislocation Based Hardening Law Allows Us to Predict Rate Dependent Texture

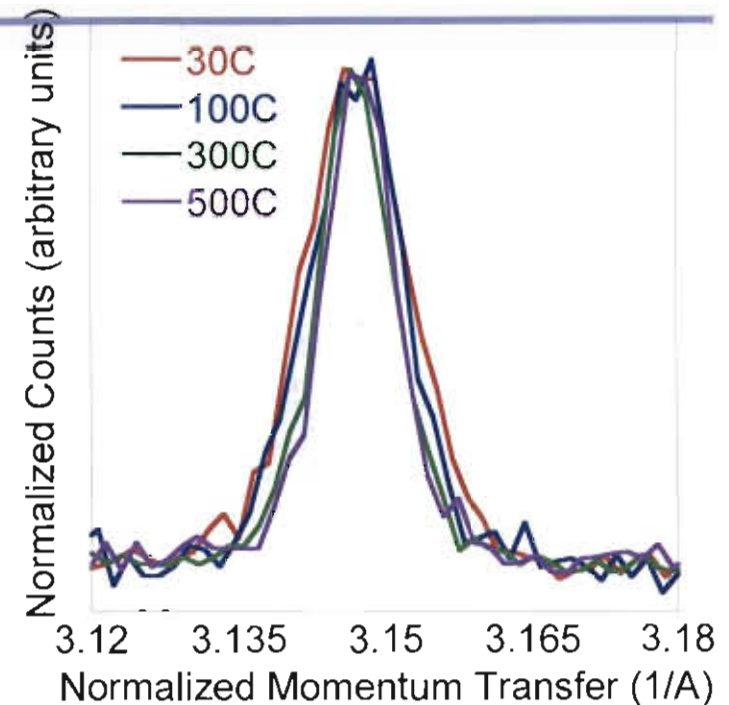
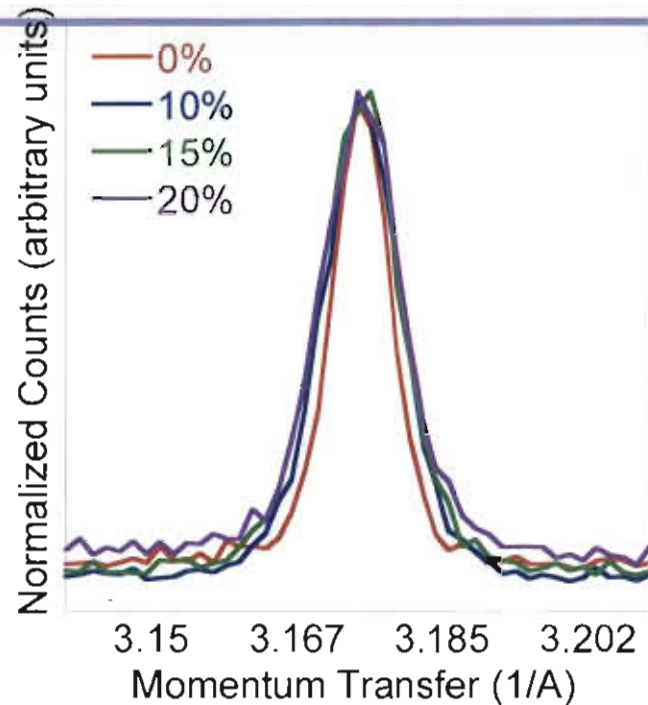
Increased Strain Rate



In this case, the experiment drove the model development.

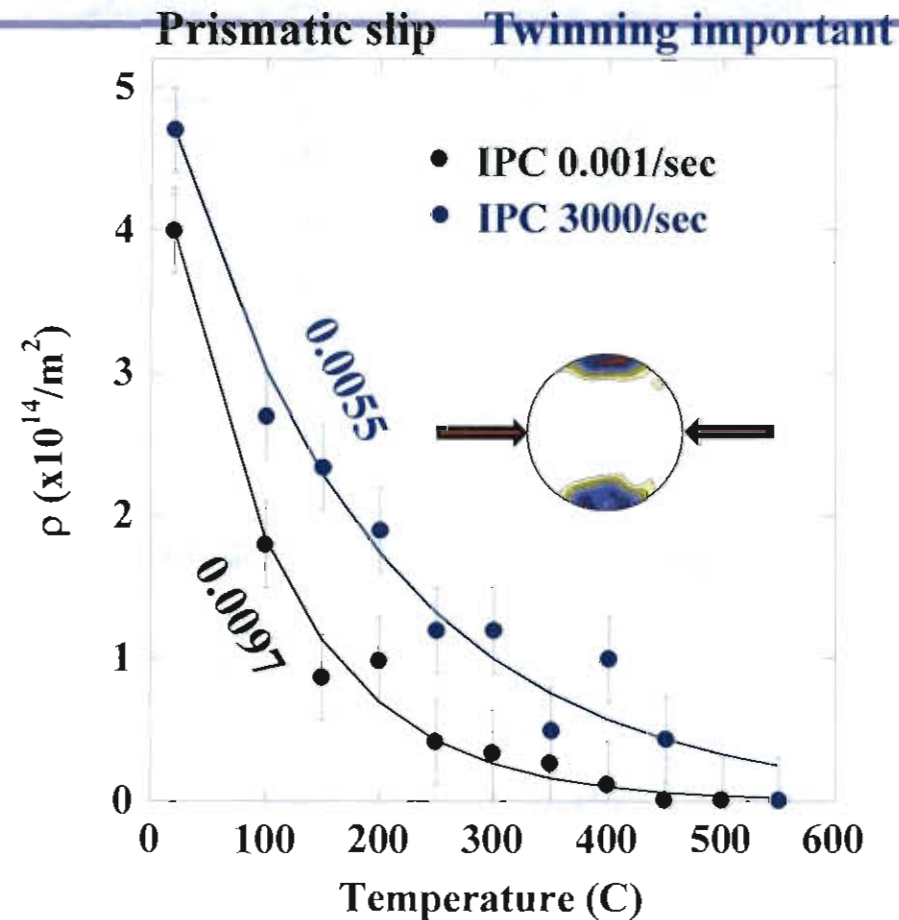


Ability to Calculate Dislocations Drives Us To Measure Dislocations



- **Diffraction peak breadth is related to the dislocation density.**
- **Completed in-situ measurements of peak breadth during deformation and subsequent annealing.**
- **Qualitatively, breadth increases with plastic deformation, decreases with annealing.**

Rate of Recovery Appears Dependent on Deformation Conditions



These experiments are now pushing the model development.

Summary and Conclusions

Polycrystal models are based on crystallinity.

Account for slip and twin systems and their interactions explicitly

Macroscopic anisotropic response is a combined result of: crystal anisotropy, microstructure inside crystal, and texture

Thermal, elastic and plastic regimes can be modeled

Stress-strain response (of grains and PX) can be predicted

Texture evolution (anisotropy evolution) can be predicted

Dislocation density evolution can be predicted

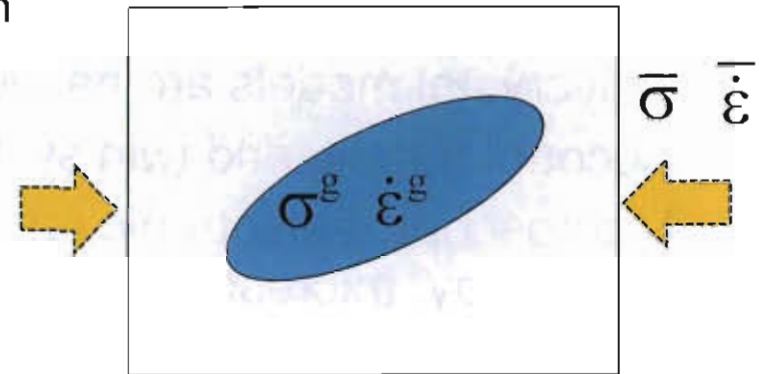
VPSC model → interaction equation, shape effects

Solving stress equilibrium equation for the VP inclusion embedded in the VP medium leads to the interaction equation

$$(\dot{\epsilon}_{ij}^g - \bar{\dot{\epsilon}}_{ij}) = -\tilde{M}_{ijkl} (\sigma_{kl}^g - \bar{\sigma}_{kl})$$

where $\tilde{M}_{ijkl} = (I - S)_{ijmn}^{-1} S_{mnpq} \bar{M}_{pqkl}$

and $S(\bar{M})$ is the Eshelby tensor



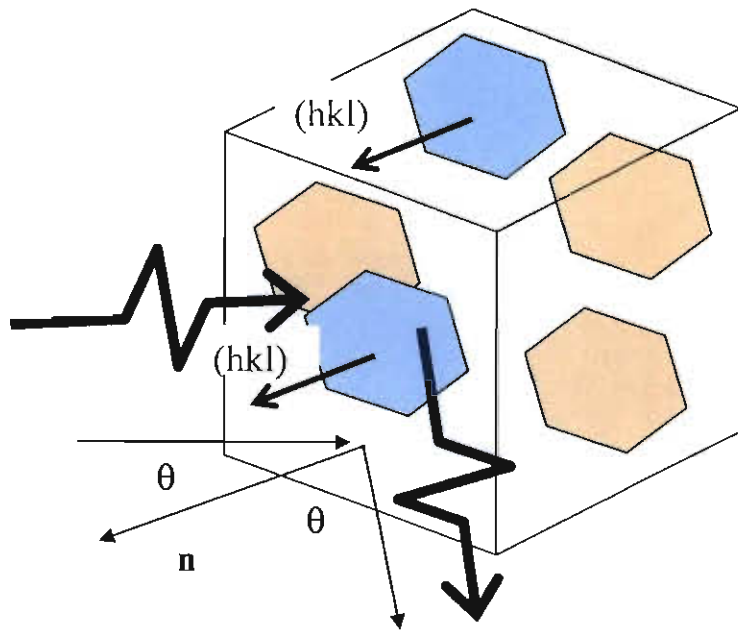
→ \bar{M} large → medium very compliant → $\sigma_{kl}^g = \bar{\sigma}_{kl}$ → lower bound Sachs

→ \bar{M} small → medium very stiff → $\dot{\epsilon}_{ij}^g = \bar{\dot{\epsilon}}_{ij}$ → upper bound Taylor

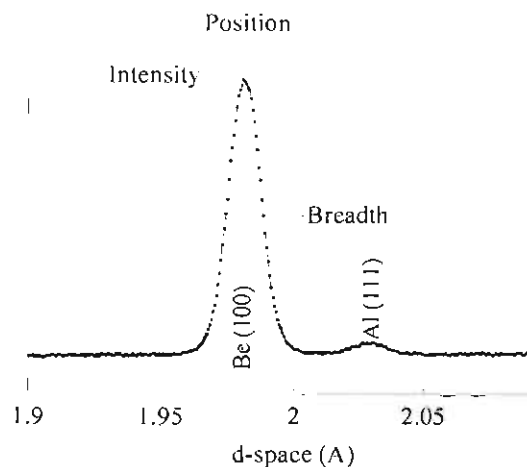
→ \bar{M} components are a mix of small and large → relaxed constraints

→ **self-consistent cases** → intermediate → secant, affine, tangent

Neutron & X-ray diffraction measurements of internal strains



Neutron Diffraction From Be-Al Composite



Subset of diffracting grains with plane (hkl) perpendicular to the diffraction vector \mathbf{n} .

The plane spacing d^{hkl} is given by the **Bragg diffraction condition**: $d^{hkl} \sin \theta = n\lambda$

The difference between stressed and unstressed lattice parameters gives the elastic strain perpendicular to the plane (hkl)

$$\varepsilon^{hkl} = \frac{d^{hkl} - d_o^{hkl}}{d_o^{hkl}}$$

Peak shift → elastic strains → **internal stress**

Peak intensity → volume fraction of grains → **texture**

Peak width → elastic strain fluctuations → **dislocation density**