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*Title:* Crystal plasticity based constitutive laws for Be:  
effect of strain rate, twinning and texture on the mechanical  
response

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# **Crystal plasticity based constitutive laws for Be: effect of strain rate, twinning and texture on the mechanical response**

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## **Abstract**

Because of its intrinsic anisotropy and variety of slip and twin modes, Beryllium displays a mechanical behavior that strongly depends upon texture, temperature and strain rate. As a consequence, accounting for crystal plasticity mechanisms and directionality is imperative for developing reliable constitutive laws for Be. In this presentation we describe a self-consistent polycrystal model and apply it to the prediction of internal stress of thermal origin, mechanical response of textured and non-textured Be, deforming in compression at strain rates ranging from  $10^{-3}$  s $^{-1}$  to  $10^4$  s $^{-1}$ . We also discuss the possibility of measuring dislocation densities using the technique of Peak Shape Analysis.

To be presented at Khariton Readings, March 14-18, 2011, Sarov, Russia

## **Crystal plasticity based constitutive laws for Be: effect of strain rate, twinning and texture on the mechanical response**

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<sup>2</sup> Theoretical Division

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## Visco-plastic grain model

Grain deformation is based on crystallographic shear on slip and twinning systems.

The **shear-rate** of each system  $s$  is given by a viscous (rate sensitive) law

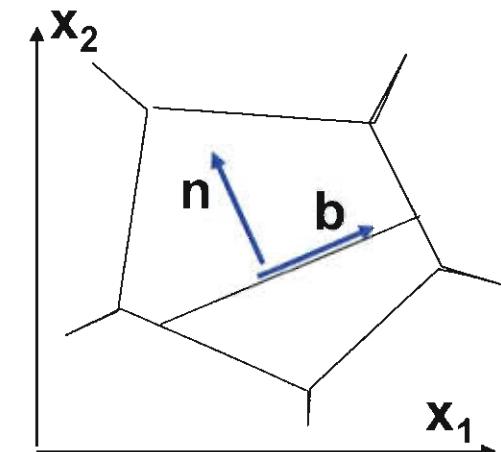
$$\dot{\gamma}^s = \dot{\gamma}_0 \left( \frac{\tau_{\text{resolved}}}{\tau^s} \right)^n = \dot{\gamma}_0 \left( \frac{m^s : \sigma}{\tau^s} \right)^n$$

The threshold stress describes the hardening of slip systems

The crystallographic shears determine crystal reorientation and texture evolution

The **strain-rate** of the grain is given by superposition of shear rates on all active systems

$$\dot{\varepsilon}'_{ij} = \sum_s \frac{n_i^s b_j^s + n_j^s b_i^s}{2} \dot{\gamma}^s = \sum_s m_{ij}^s \dot{\gamma}^s$$

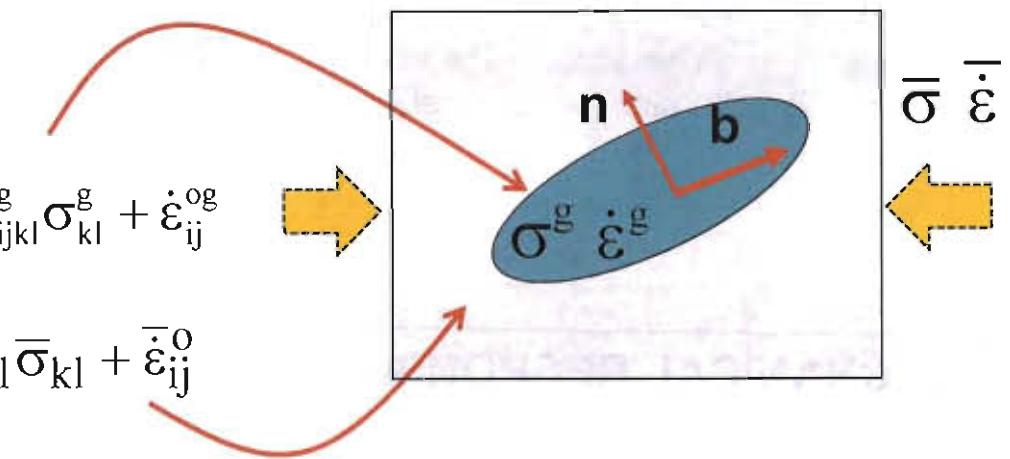


## Visco Plastic Self Consistent (VPSC) Polycrystal Model

- Each grain is a visco-plastic **anisotropic ellipsoidal inclusion** embedded in a visco-plastic anisotropic Homogeneous Effective Medium (HEM).
- Linearized constitutive response:  $M^g$  and  $\bar{M}$   $\rightarrow$  grain and effective medium compliance tensors

grain:  $\dot{\varepsilon}_{ij}^g = \dot{\gamma}_o \sum_s m_{ij}^s \left( \frac{m^s : \sigma^g}{\tau^s} \right)^n = M_{ijkl}^g \sigma_{kl}^g + \dot{\varepsilon}_{ij}^{og}$

medium:  $\bar{\varepsilon}_{ij} = \bar{M}_{ijkl} \bar{\sigma}_{kl} + \bar{\varepsilon}_{ij}^0$

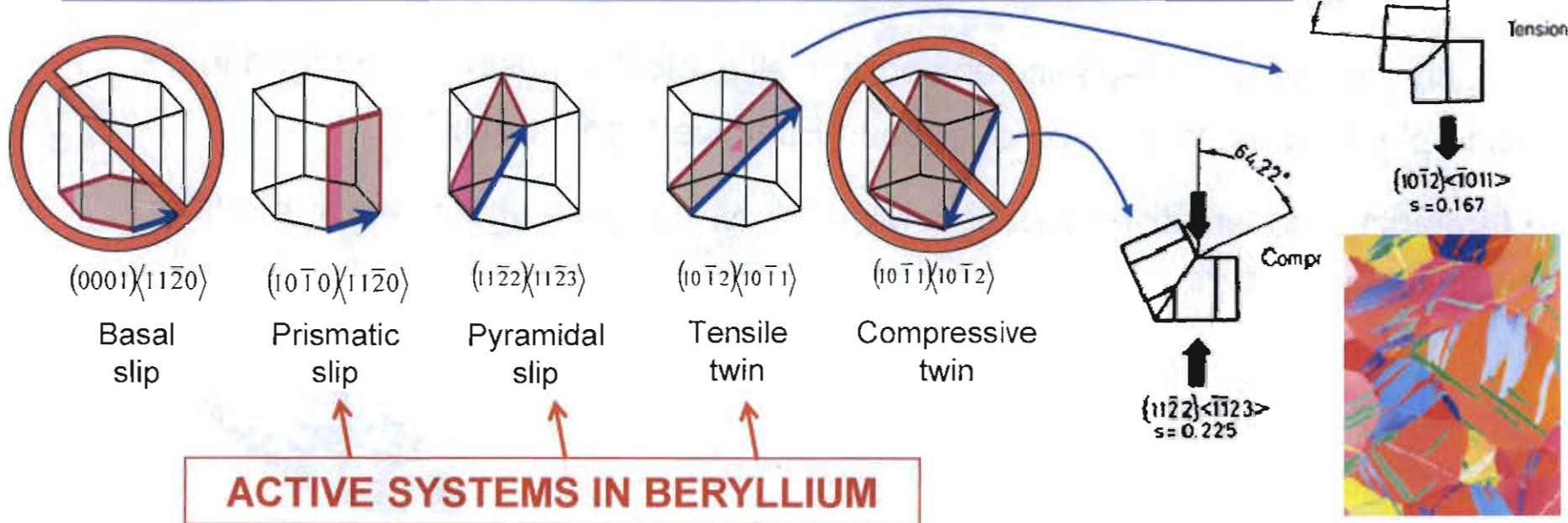


Solve equilibrium equation for inclusion in homogeneous medium:  $\sigma_{ij,j} = 0$

Eshelby result  $\rightarrow$  stress and strain-rate are uniform inside the inclusion

...but different from the macroscopic stress and strain rate !

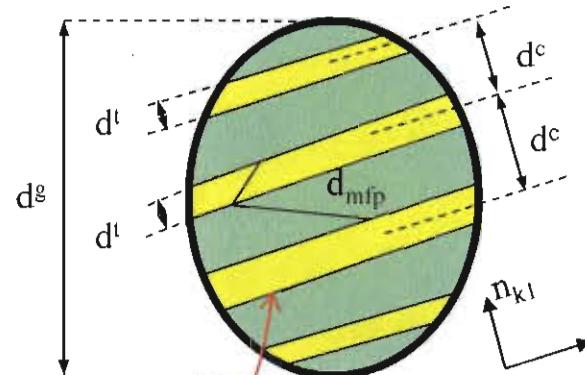
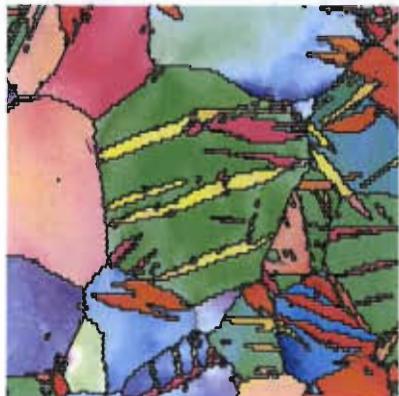
# Plasticity of hexagonal materials



## MECHANICAL RESPONSE OF HEXAGONALS

- Several slip and twin modes with different temperature and rate dependencies
- Anisotropy of single crystal  $\rightarrow$  difficult to deform along the c-axis
- Need to account for slip-slip and for slip-twin interactions
- Twinning activity affects texture, hardening and anisotropy

# Composite-Grain Twin Model



Treat grain as an evolving stack of twin-matrix layers on the plane of the predominant twin system.

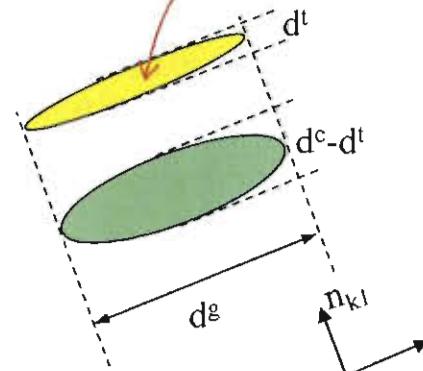
## COUPLED

Continuity of stress and strain across the twin-matrix interface

$$\sigma_{13}^{parent} = \sigma_{13}^{twin}$$

$$\sigma_{23}^{parent} = \sigma_{23}^{twin}$$

$$\dot{\epsilon}_{12}^{parent} = \dot{\epsilon}_{12}^{twin}$$



## UNCOPLED

Separate treatment of parent and twin as ellipsoidal inclusions.

Account for ellipsoid orientation and thickness evolution with deformation

Proust & Tome, Acta Materialia **55** (2007) 2137  
Proust et al, Int. J. of Plasticity **25** (2009) 861

## Rate Dependent Experiments Prompt Dislocation Based Hardening Model

$$\frac{\partial \rho^\alpha}{\partial \gamma^\alpha} = \frac{\partial \rho_{storage}^\alpha}{\partial \gamma^\alpha} - \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} = k_1^\alpha \sqrt{\rho^\alpha} - k_2^\alpha(\dot{\epsilon}, T) \rho^\alpha \quad (\alpha = \text{slip mode})$$

$$\frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} \quad \begin{array}{l} \xrightarrow{\hspace{10em}} (1 - f^\alpha) \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} \quad \text{ANNIHILATION} \\ \xrightarrow{\hspace{10em}} f^\alpha \frac{\partial \rho_{removed}^\alpha}{\partial \gamma^\alpha} \quad \text{SUBSTRUCTURE} \quad \rho_{sub} \end{array}$$

- Temperature and strain rate dependence enter in rate of annihilation.

## Dislocation-based constitutive law for HCP (temperature and rate dependent)

Dislocation densities in the slip systems ( $\alpha = \text{basal, prism or pyramidal}$ ) evolve with crystal shears via generation, annihilation (removal) and substructure build-up

$$d\rho_{\text{gener, forest}}^{\alpha} = k_1^{\alpha} \sqrt{\rho_{\text{forest}}^{\alpha}} dy^{\alpha} \quad \leftarrow \text{Dislocation generation}$$

$$d\rho_{\text{remov, forest}}^{\alpha} = k_2^{\alpha}(\dot{\epsilon}, T) \rho_{\text{forest}}^{\alpha} dy^{\alpha} \quad \leftarrow \text{Dislocation annihilation}$$

$$d\rho_{\text{substructure}} = \sum_{\alpha} f^{\alpha}(T) d\rho_{\text{remov, forest}}^{\alpha} \quad \leftarrow \text{Substructure formation}$$

The threshold to propagate dislocations is related to the density via Taylor's law

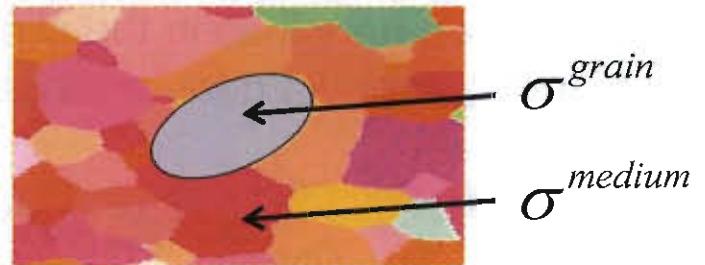
$$\tau_{\text{slip}}^{\alpha} = b^{\alpha} \chi \mu \sqrt{\sum_{\alpha} D^{\alpha\alpha} \rho_{\text{forest}}^{\alpha}} \quad \leftarrow \text{Forest hardening}$$

$$\tau_{\text{twin}}^{\beta} = \mu \sum_{\alpha} C^{\beta\alpha} b^{\alpha} b^{\beta} \rho^{\alpha} \quad \leftarrow \text{Twin hardening by dislocations}$$

# Multiscale modeling

Visco Plastic (VPSC) and Elasto Plastic (EPSC) Polycrystal Aggregate Models

VPSC → Lebensohn & Tome (1993); Proust et al (2009)  
EPSC → Turner & Tome (1994); Clausen et al (2008)



Composite Grain Model

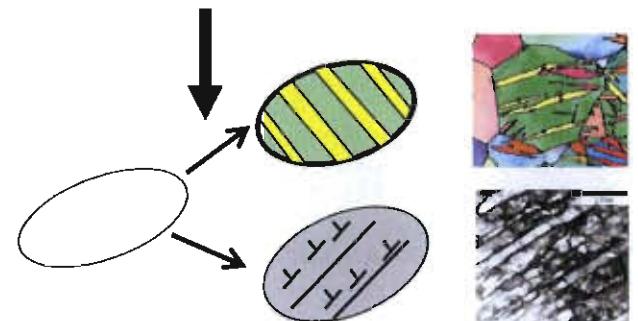
Proust et al. (2007)

Dislocation-based Hardening Model

Beyerlein & Tome (2008)



Shear rates & hardening for slip & twinning controlled by evolving Threshold Stress



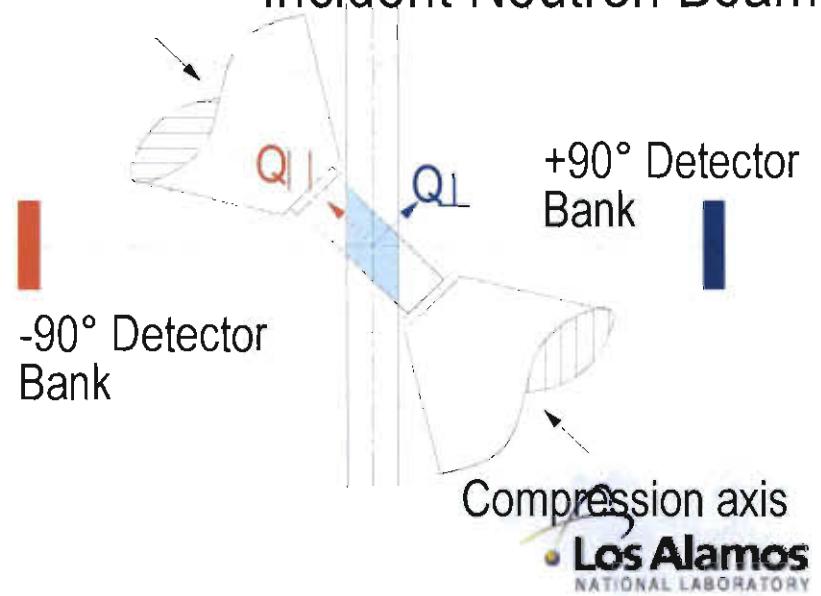
$$\dot{\gamma}^s = \dot{\gamma}_0 \left( \frac{\tau_{resolved}}{\frac{\tau^s}{\tau_{threshold}}} \right)^n$$

# Neutron diffraction facilities: a tool for 'in-situ' & 'in-bulk' stress measurements

SMARTS Diffractometer  
LANSCE - LANL



Incident Neutron Beam



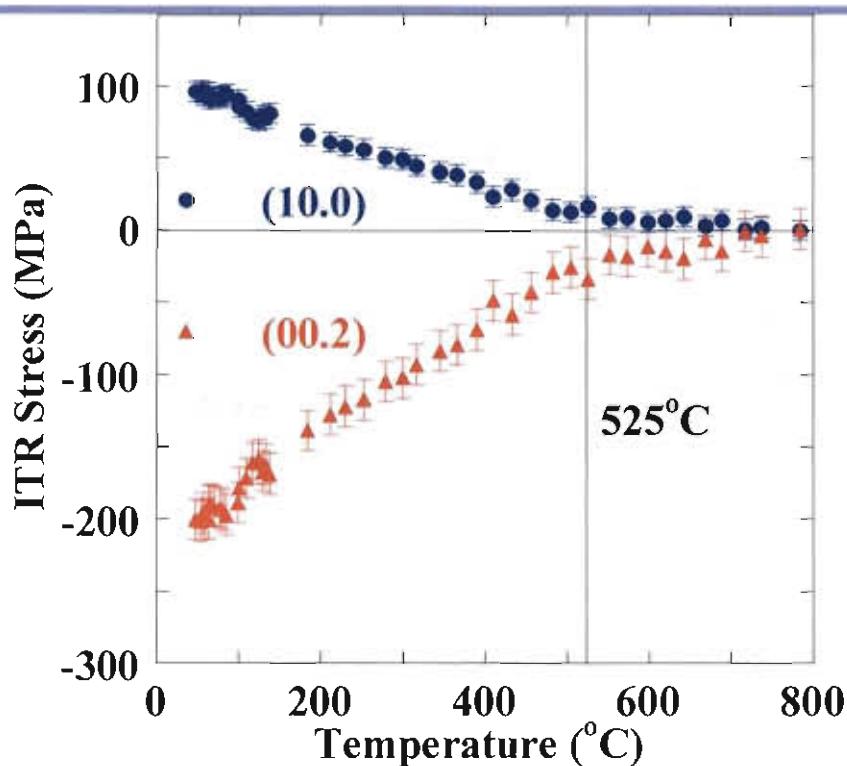
- Crystal diffraction measures lattice spacing in the bulk
- Lattice spacing provides information about elastic strains
- Elastic strains are related to internal stress via Hooke's law
- Diffracting region: ~ 5mm

## Modeling of Thermal and Mechanical Response of Be

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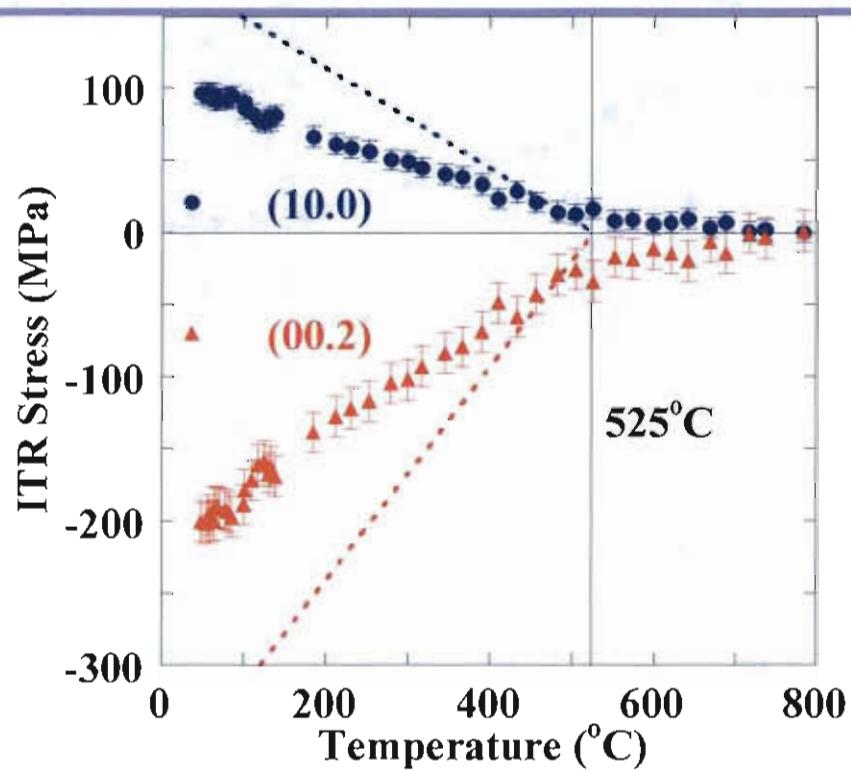
- Use Elasto-Plastic Polycrystal model to calculate internal stress inside grains of non-textured Be.  
→ Compare with neutron diffraction results
- Use Visco-Plastic Polycrystal model to calculate Stress-Strain response of rolled and non-textured Be.  
→ Use an unique Single Crystal constitutive law, function of temperature and strain-rate
- Calculate dislocation densities using Peak Profile Analysis on neutron diffraction peaks

## Intergranular Thermal Residual Stresses in Beryllium are Significant.



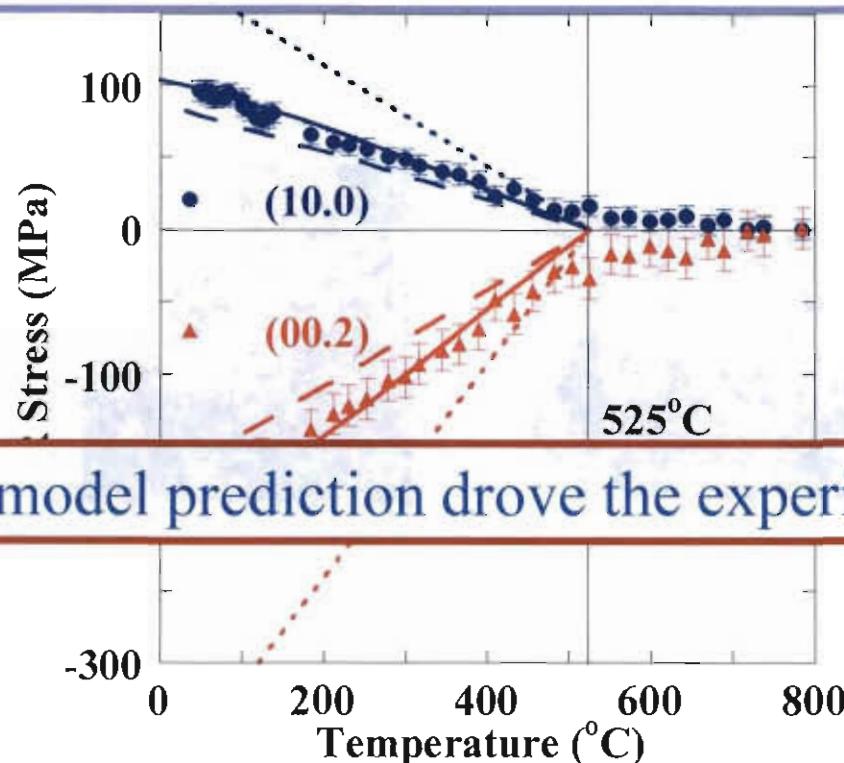
- Intergranular stresses relax on time scale of measurement above 525C.
  - Zero strength temperature.
  - Important for modeling studies.

## Intergranular Thermal Residual Stresses in Beryllium are Significant.



- Upper bound model, assuming grain in rigid matrix over-predicts stresses by ~x2

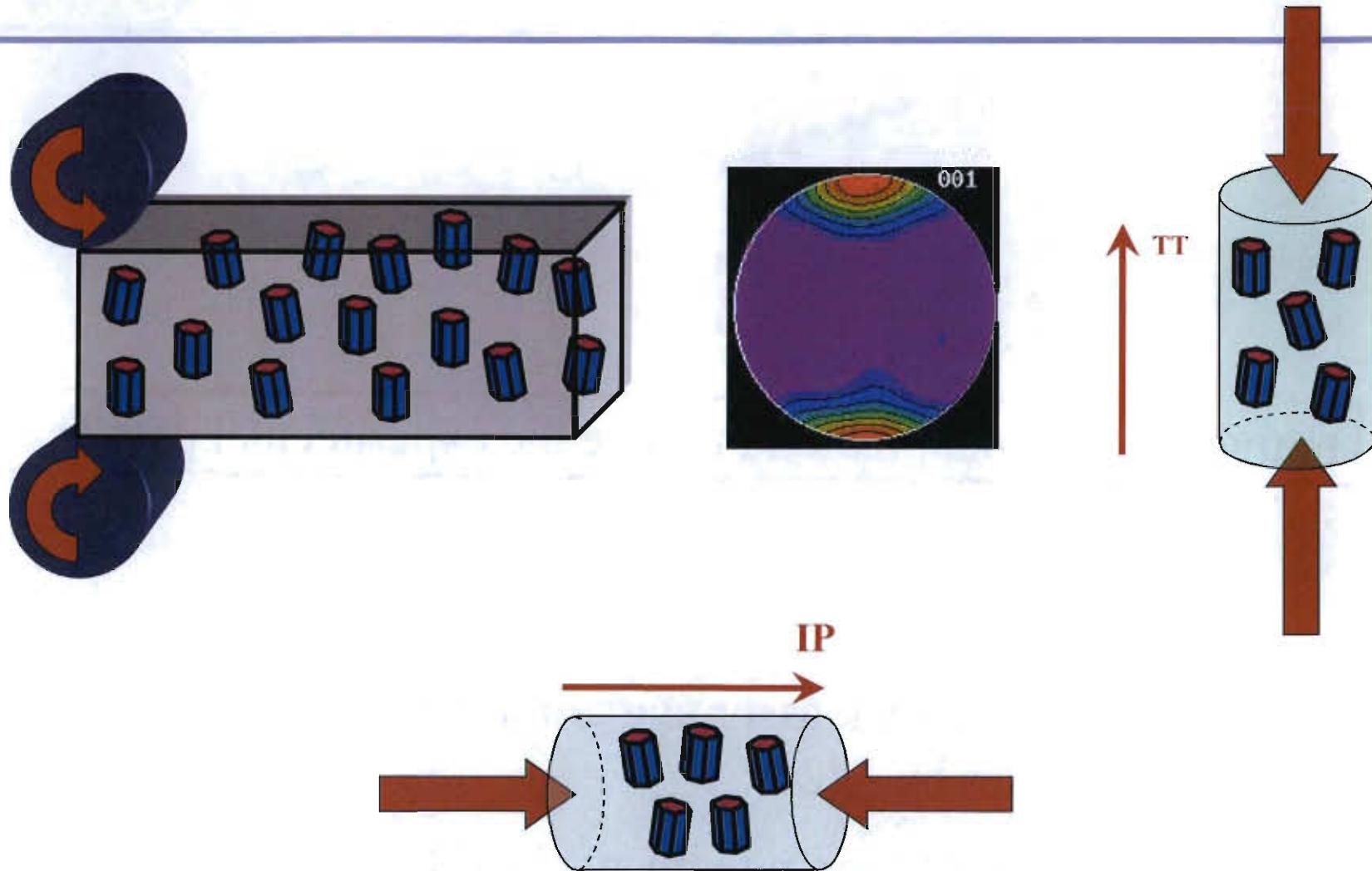
## Intergranular Thermal Residual Stresses in Beryllium are Significant.



In this case, model prediction drove the experimental path !

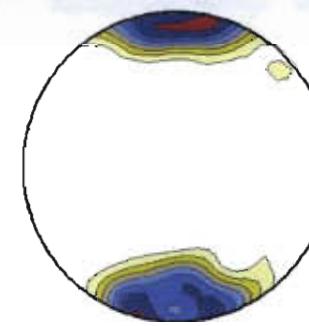
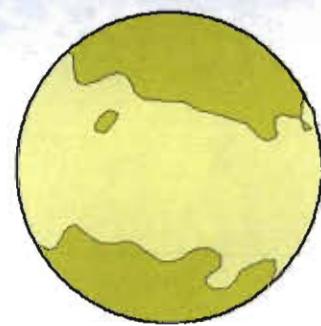
- Upper bound model, assuming grain in rigid matrix over-predicts stresses by  $\sim x2$ .
- EPSC model with room temperature CTE's under-predicts stresses slightly.
- EPSC with temperature dependent CTE matches data 

## Rolling of Beryllium Results in Strong Crystallographic Texture



- Orientation of crystals with respect to load axis is critical.

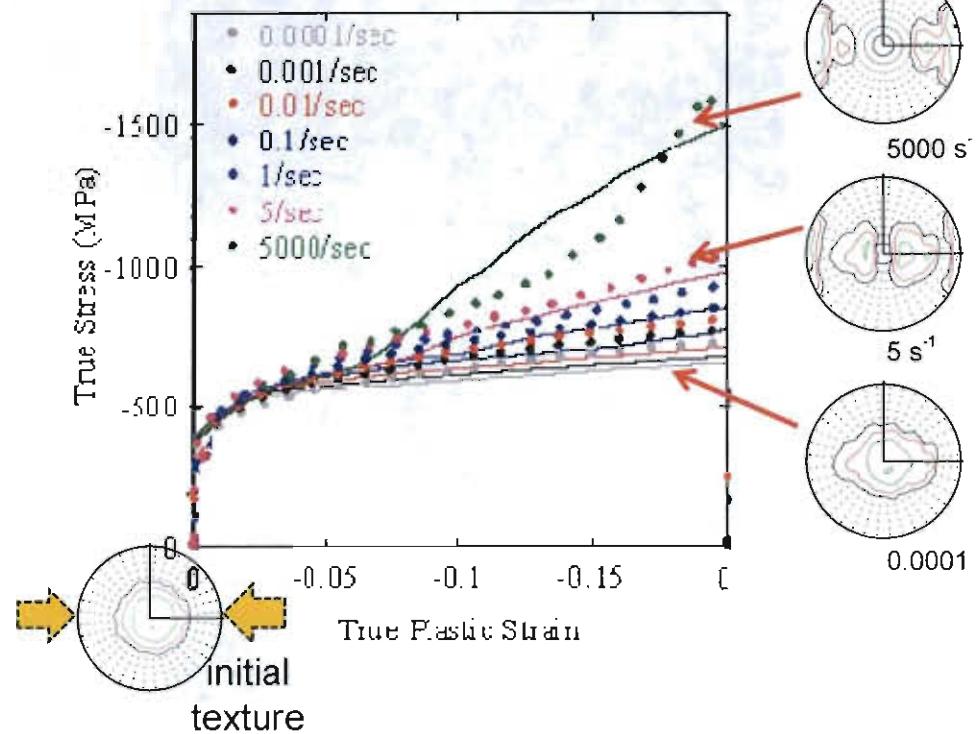
## Rate and Temperature Dependent Deformation of Beryllium



- Hot-pressed (random texture) and rolled plate (strong texture) beryllium.
- In-plane (IP) and through-thickness (TT) compression of rolled plate.
  - Normal and parallel to dominant basal texture, respectively.
- Strain rates from 0.0001/sec to 5000/sec.

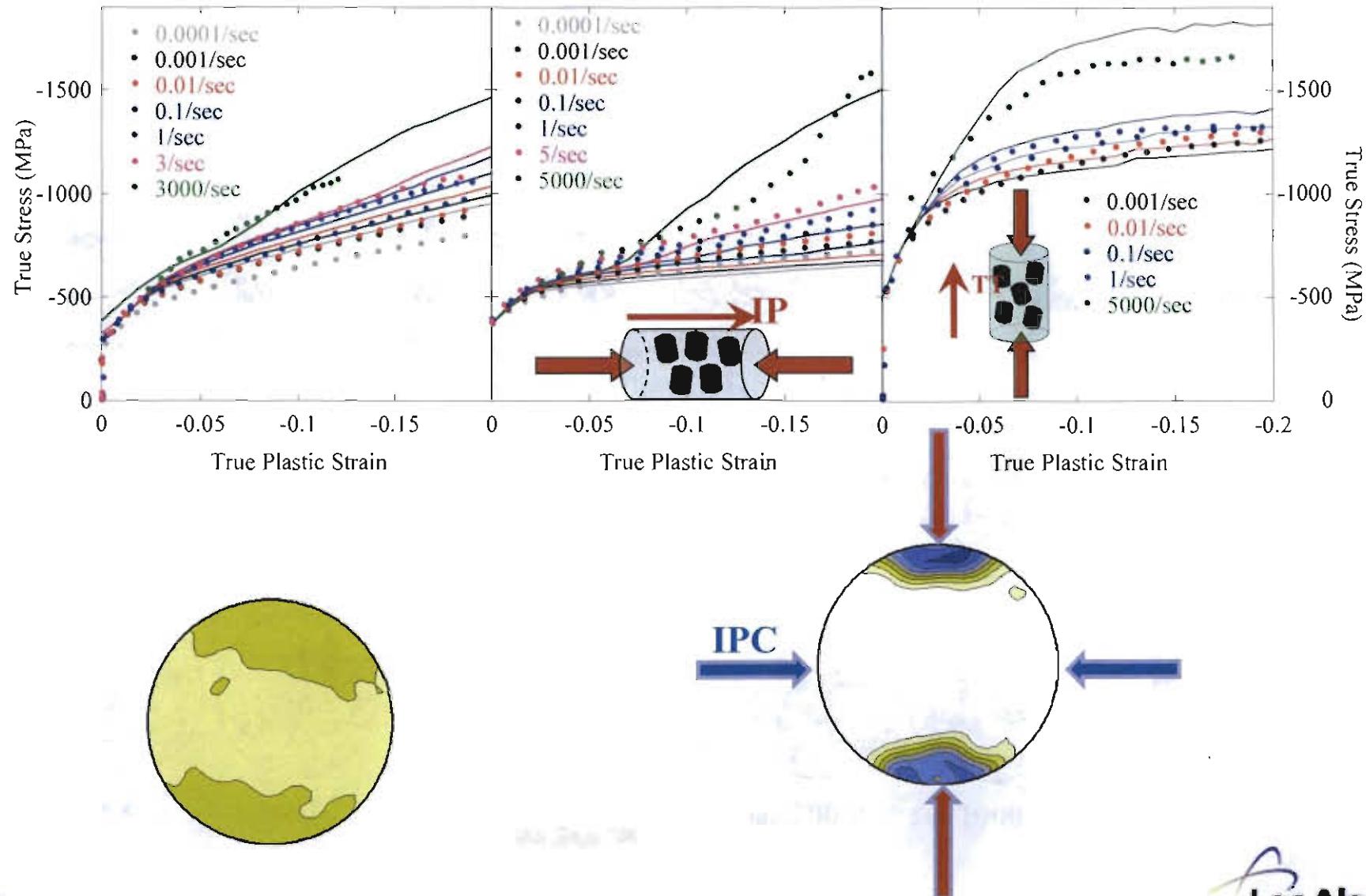
## Be: slip vs twinning competition

### Beryllium: In-Plane Compression



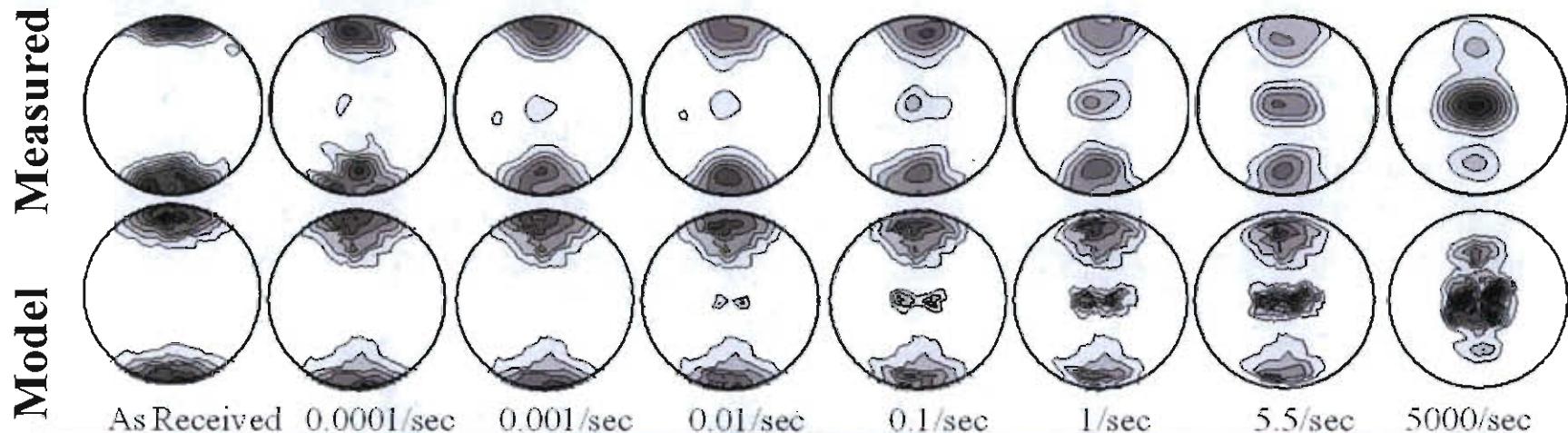
Rapid transition from dislocation-controlled to twin-controlled hardening as the temperature decreases (Zr) or as the strain rate increases (Be)

## Dislocation Based Hardening Law Allows Us to Predict Rate Dependent Flow curve

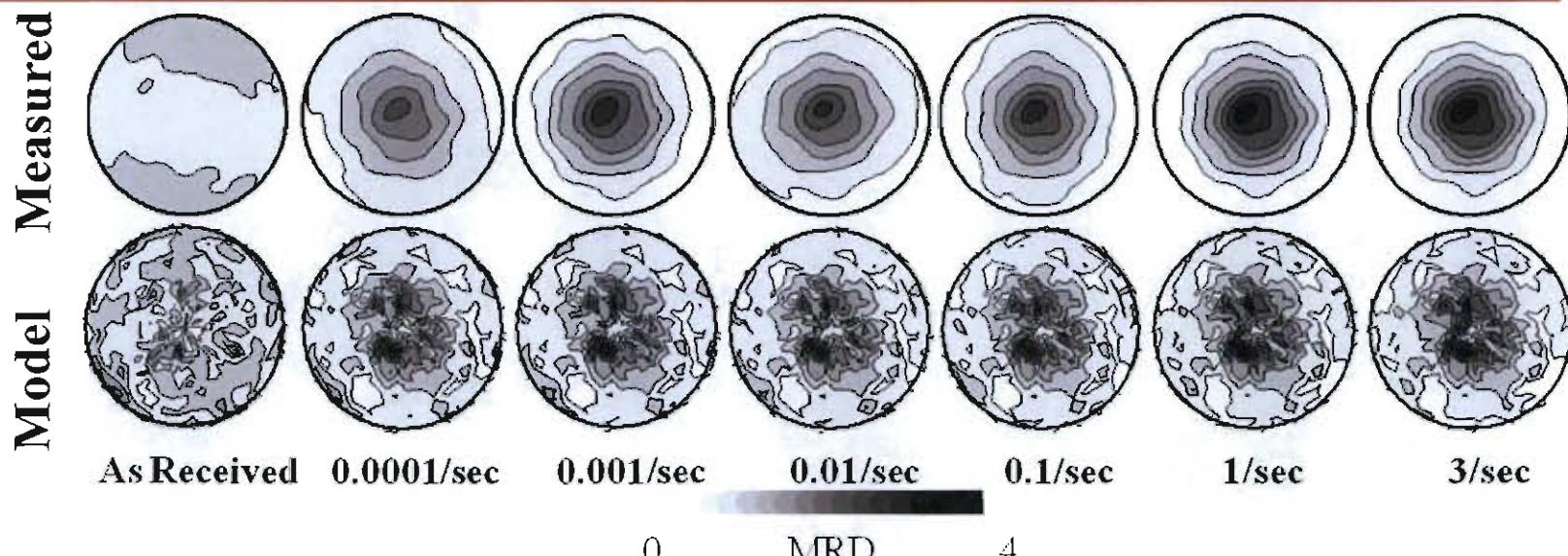


## Dislocation Based Hardening Law Allows Us to Predict Rate Dependent Texture

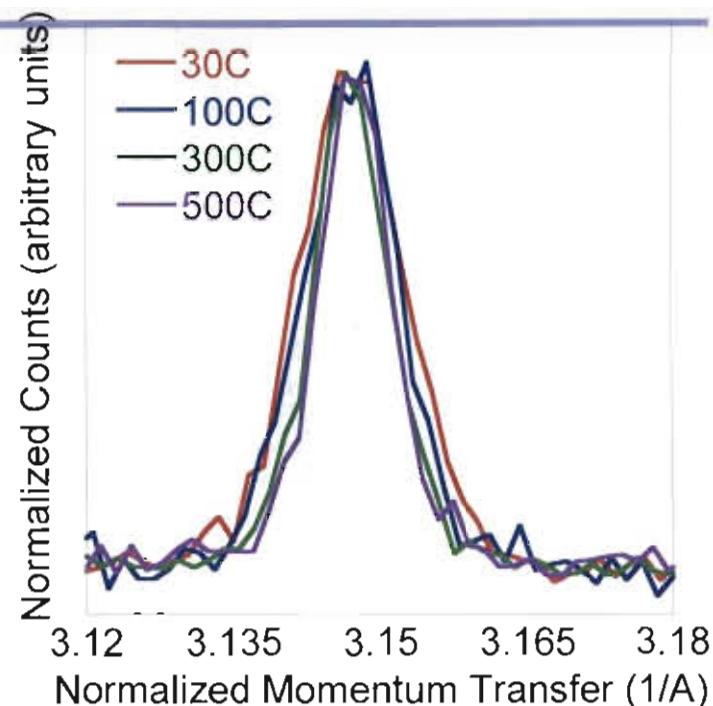
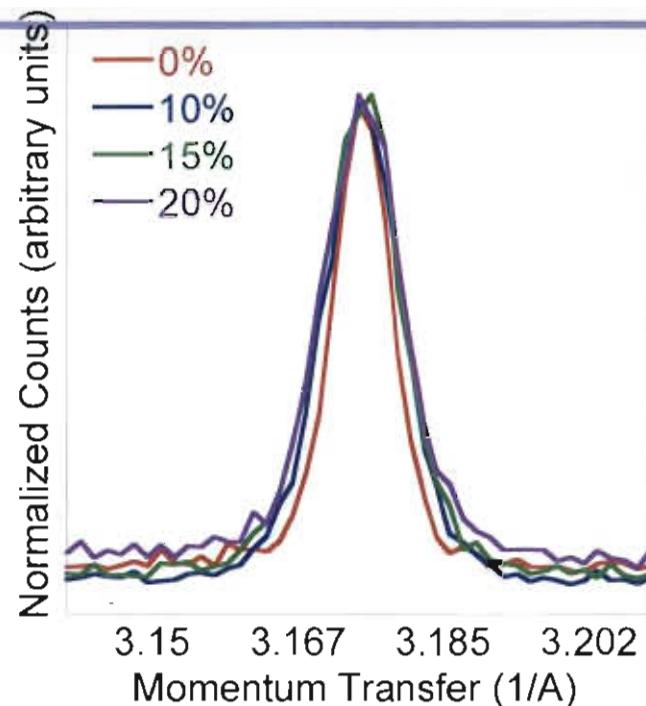
### Increased Strain Rate



In this case, the experiment drove the model development.

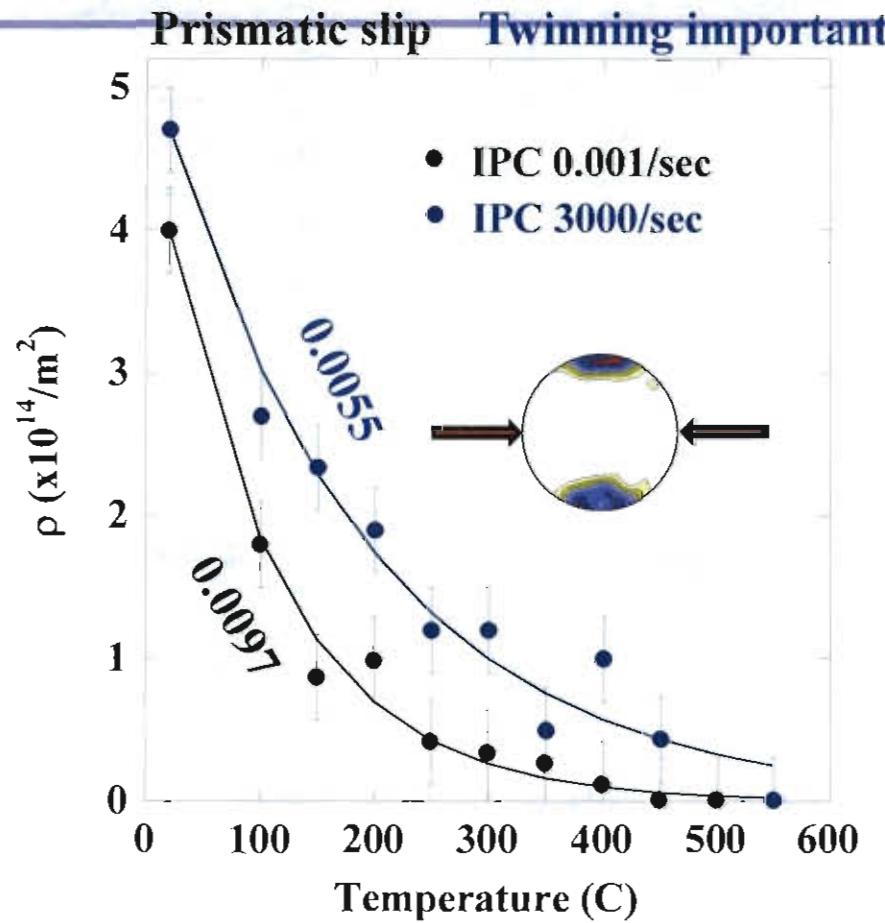


## Ability to Calculate Dislocations Drives Us To Measure Dislocations



- Diffraction peak breadth is related to the dislocation density.
- Completed in-situ measurements of peak breadth during deformation and subsequent annealing.
- Qualitatively, breadth increases with plastic deformation, decreases with annealing.

## Rate of Recovery Appears Dependent on Deformation Conditions



These experiments are now pushing the model development.

## Summary and Conclusions

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Polycrystal models are based on crystallinity.

Account for slip and twin systems and their interactions explicitly

Macroscopic anisotropic response is a combined result of: crystal anisotropy, microstructure inside crystal, and texture

Thermal, elastic and plastic regimes can be modeled

Stress-strain response (of grains and PX) can be predicted

Texture evolution (anisotropy evolution) can be predicted

Dislocation density evolution can be predicted

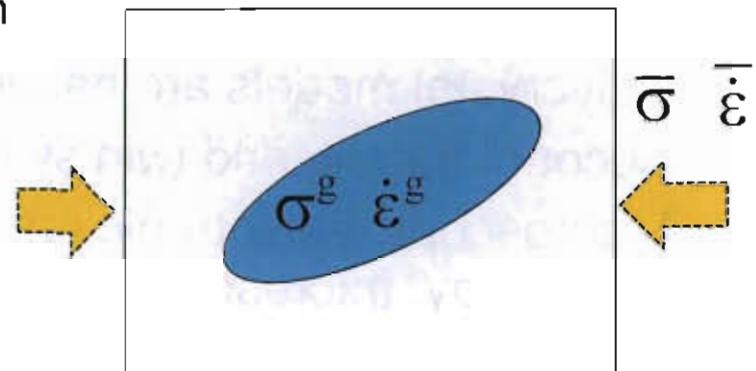
## VPSC model → interaction equation, shape effects

Solving stress equilibrium equation for the VP inclusion embedded in the VP medium leads to the interaction equation

$$(\dot{\varepsilon}_{ij}^g - \bar{\dot{\varepsilon}}_{ij}) = -\tilde{M}_{ijkl} (\sigma_{kl}^g - \bar{\sigma}_{kl})$$

where  $\tilde{M}_{ijkl} = (I - S)^{-1}_{ijmn} S_{mnpq} \bar{M}_{pqkl}$

and  $S(\bar{M})$  is the Eshelby tensor



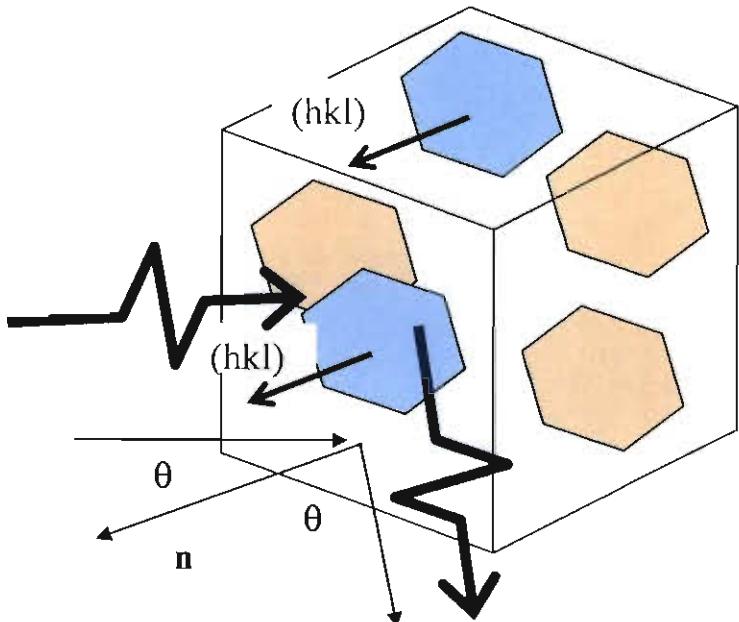
→  $\bar{M}$  large → medium very compliant →  $\sigma_{kl}^g = \bar{\sigma}_{kl}$  → lower bound Sachs

→  $\bar{M}$  small → medium very stiff →  $\dot{\varepsilon}_{ij}^g = \bar{\dot{\varepsilon}}_{ij}$  → upper bound Taylor

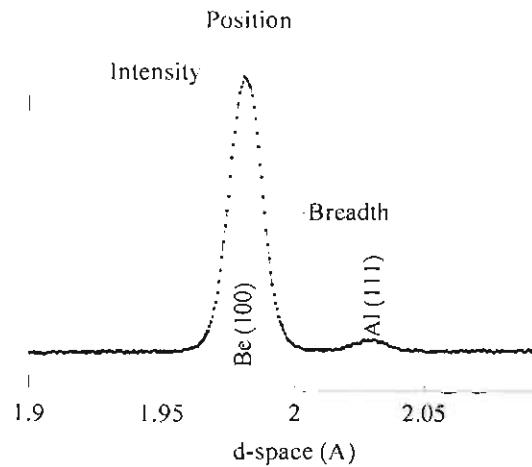
→  $\sim$   $\bar{M}$  components are a mix of small and large → relaxed constraints

→ **self-consistent cases** → intermediate → secant, affine, tangent

# Neutron & X-ray diffraction measurements of internal strains



Neutron Diffraction From Be-Al Composite



Subset of diffracting grains with plane  $(hkl)$  perpendicular to the diffraction vector  $\mathbf{n}$ .

The plane spacing  $d^{hkl}$  is given by the Bragg diffraction condition:  $d^{hkl} \sin \theta = n\lambda$

The difference between stressed and unstressed lattice parameters gives the elastic strain perpendicular to the plane  $(hkl)$

$$\varepsilon^{hkl} = \frac{d^{hkl} - d_o^{hkl}}{d_o^{hkl}}$$

Peak shift  $\rightarrow$  elastic strains  $\rightarrow$  internal stress

Peak intensity  $\rightarrow$  volume fraction of grains  $\rightarrow$  texture

Peak width  $\rightarrow$  elastic strain fluctuations  $\rightarrow$  dislocation density