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# Radiation Induced Defects Behavior in the Presence of Interfaces

**Enrique Martinez, Blas Uberuaga and Alfredo Caro**

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Alamos, New Mexico 87545, USA*

**CMIME IE Thrust**

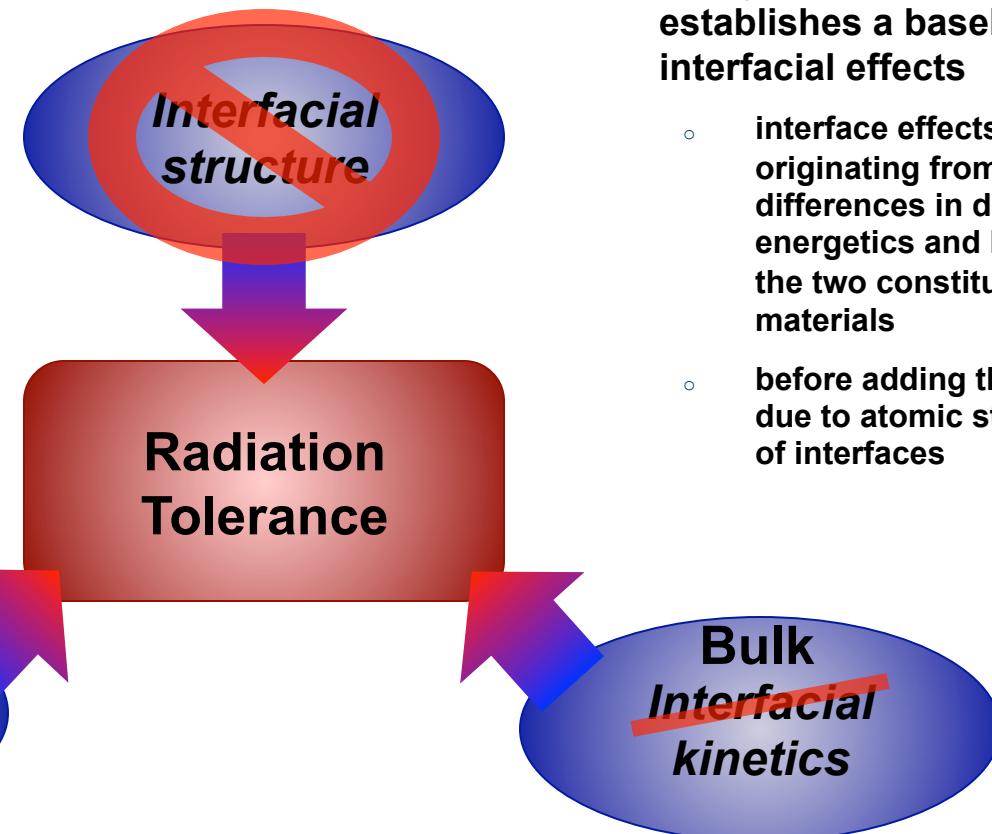
# Outline

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- **Motivation**
- **Coherent Interfaces**
  - Effect of  $\Delta\mu$  and  $\Delta E$  on defect accumulation and recombination
- **Semi-coherent Interfaces**
  - Development of an off-lattice kinetic Monte Carlo code
    - Why we need this kind of code
    - What we can do with it (so far)
  - Applications. Vacancy diffusion in the presence of Twist Boundaries
  - Error Estimation
- **Conclusions**

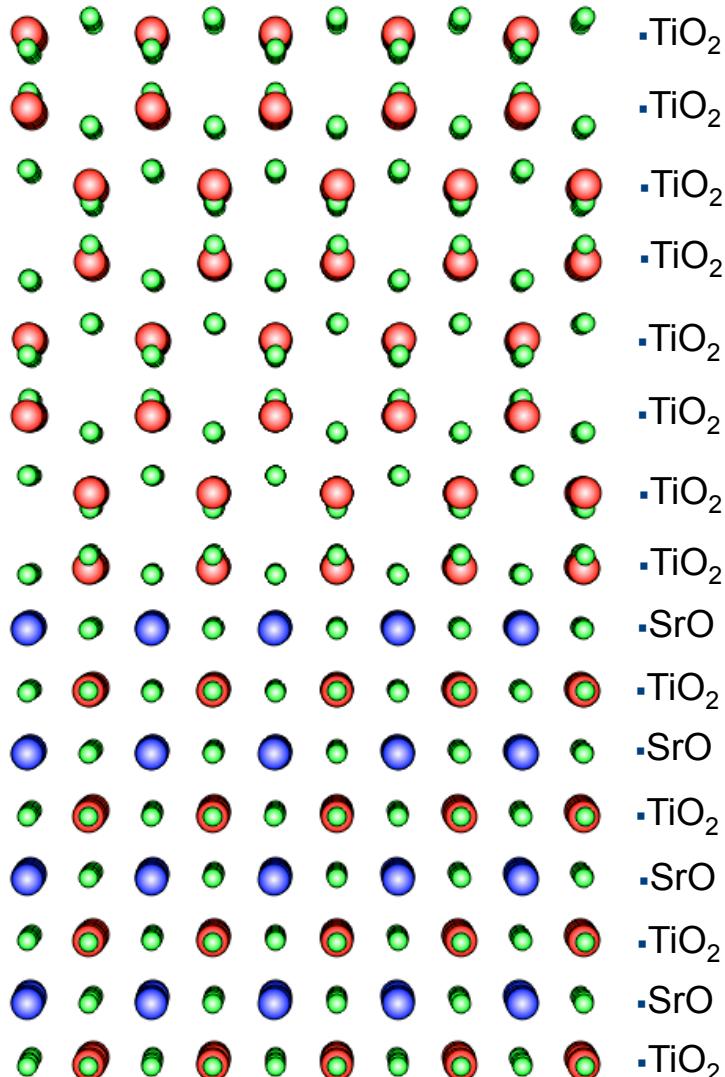
# Radiation damage at “structureless” interfaces

- “Structureless” in this context refers to interfaces where
  - the response is dominated by defect energetics and kinetics
  - with only a minor contribution from the interface atomic structure.
- Such interfaces are likely to be coherent or nearly-coherent with a small lattice mismatch.



- Study of such interfaces establishes a baseline for interfacial effects
  - interface effects originating from differences in defect energetics and kinetics in the two constituent materials
  - before adding the effects due to atomic structures of interfaces

# The atomic structure of the $\text{TiO}_2/\text{SrTiO}_3$ interface



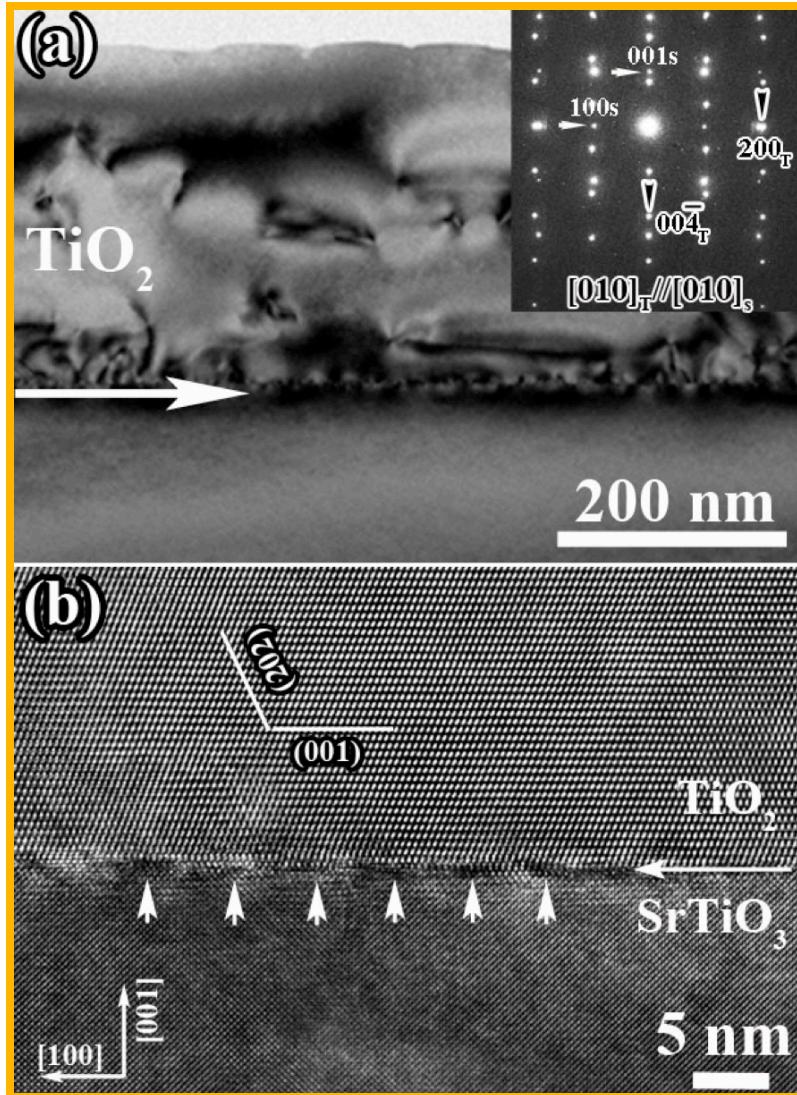
- (001) direction of  $\text{SrTiO}_3$  characterized by alternating layers of  $\text{SrO}$  and  $\text{TiO}_2$ 
  - Stoichiometric layers
- Layers in (001) direction of anatase  $\text{TiO}_2$  has the same structure as in  $\text{SrTiO}_3$
- Nearly coherent interface

J N C L A S S I F I E D

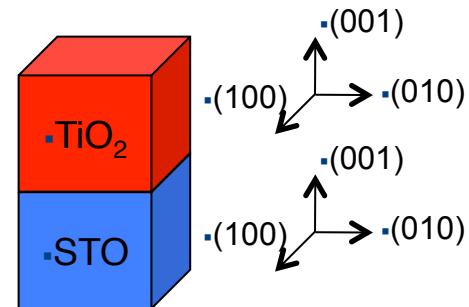
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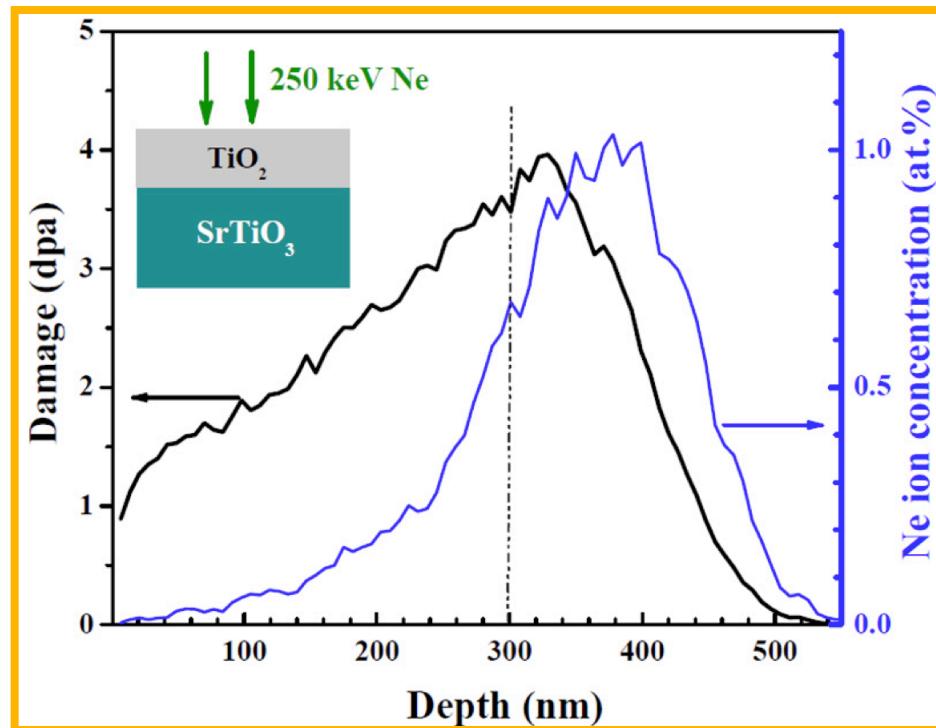
# Structure of as-grown $\text{TiO}_2/\text{SrTiO}_3$ interface



- **Epitaxial [004]-oriented  $\text{TiO}_2$  film on (001)  $\text{SrTiO}_3$** 
  - Deposited via pulsed laser deposition
  - XRD confirms  $\text{TiO}_2$  is anatase polymorph
  - Film about 300 nm thick
- **Sharp  $\text{TiO}_2/\text{SrTiO}_3$  interface**
  - Misfit dislocation spacing: 5.6 nm
- **Orientation relationship:**

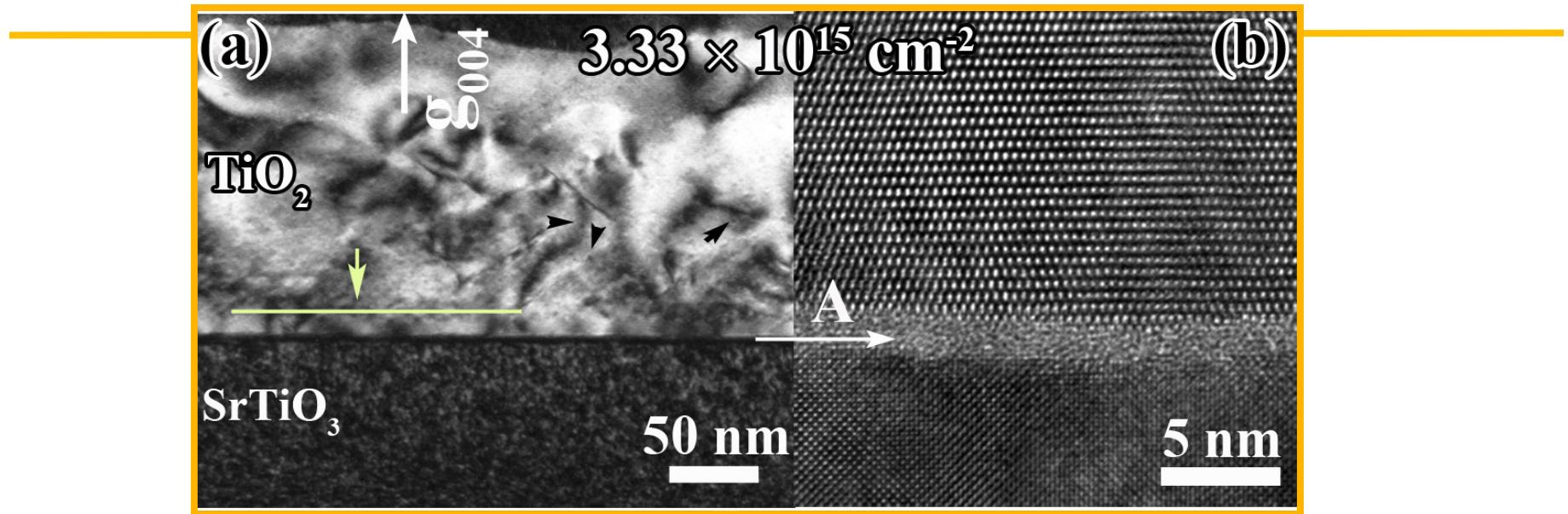


# Irradiation of $\text{TiO}_2/\text{SrTiO}_3$ interface



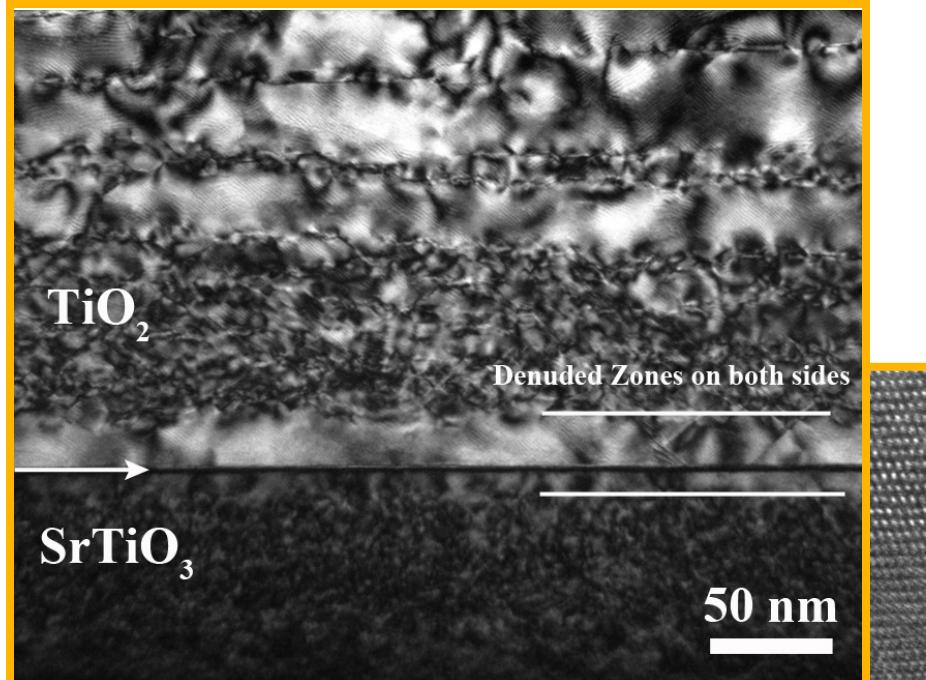
- 250 keV Ne irradiations
- Peak dpa is at about 330-340 nm
  - About 30-40 nm deeper than interface
- Films irradiated to  $3.3 \times 10^{15}$  to  $1.7 \times 10^{16}$  ions/cm<sup>2</sup>
  - Corresponds to 1-6 dpa at interface, based on SRIM

# Irradiation of $\text{TiO}_2/\text{SrTiO}_3$ interface: room temperature

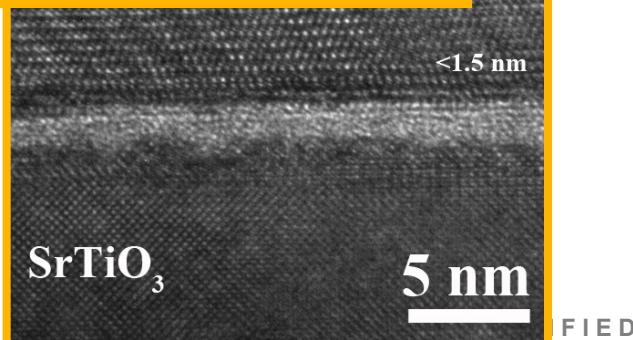


- Denuded zone forms on  $\text{TiO}_2$  side of interface
- Amorphous layer forms on  $\text{SrTiO}_3$  side of interface
- Denuded zone persists even as amorphous layer forms
- *Denuded zone independent of atomic structure of interface*

# Irradiation of $\text{TiO}_2/\text{SrTiO}_3$ interface: 500 C



- Denuded zone forms on **both** sides of the interface
- Very thin amorphous layer forms at the interface
- Dislocation walls form in  $\text{TiO}_2$  thin film



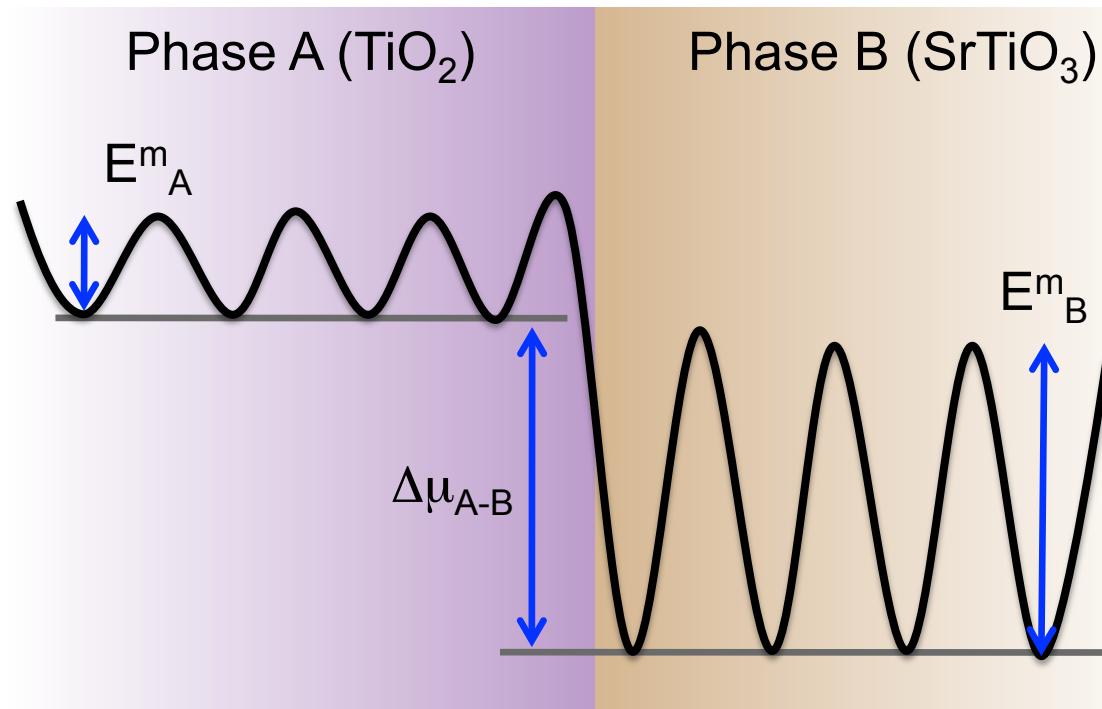
## Two determiners of defect properties near structureless interfaces

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- Defect transport driven by differences in chemical potential  $\mu$ 
  - Defects move from regions of high chemical potential to regions of low chemical potential
- Defect transport controlled by bulk migration energies
  - Whether defects can cross interface depends on migration barriers relative to temperature in both phases

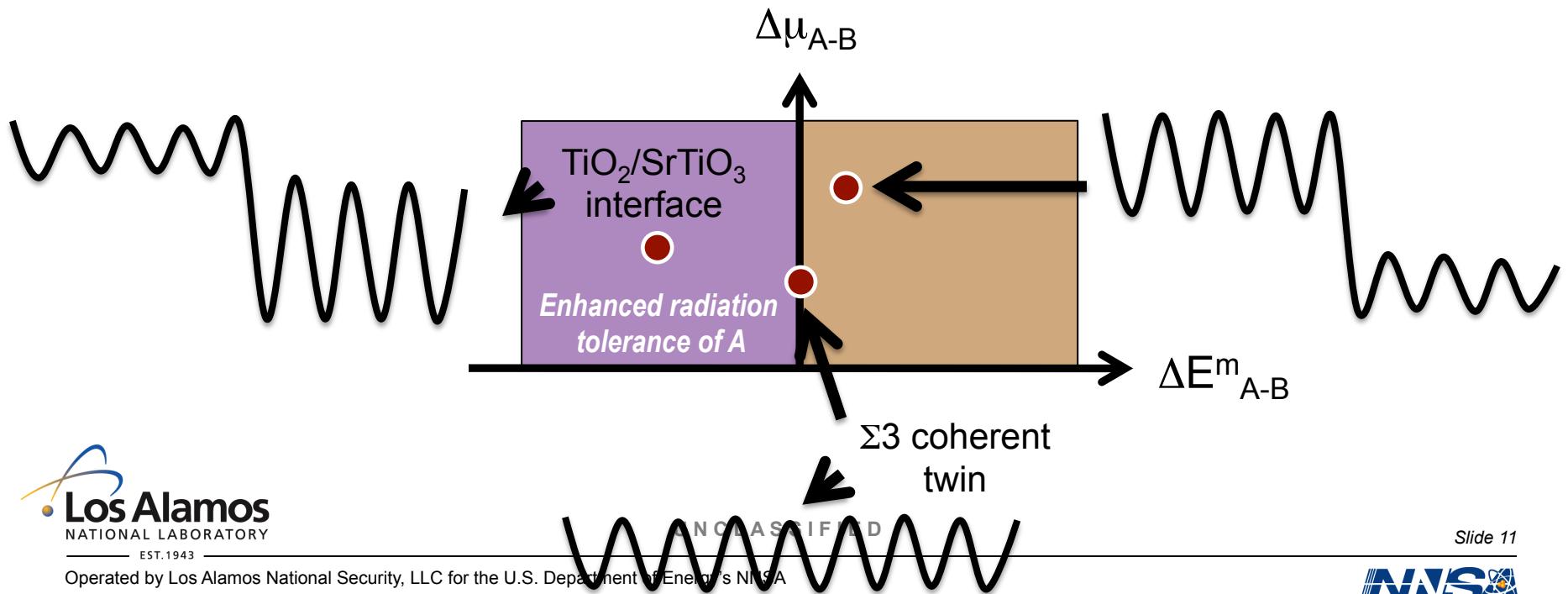
## Two determiners of defect properties near structureless interfaces

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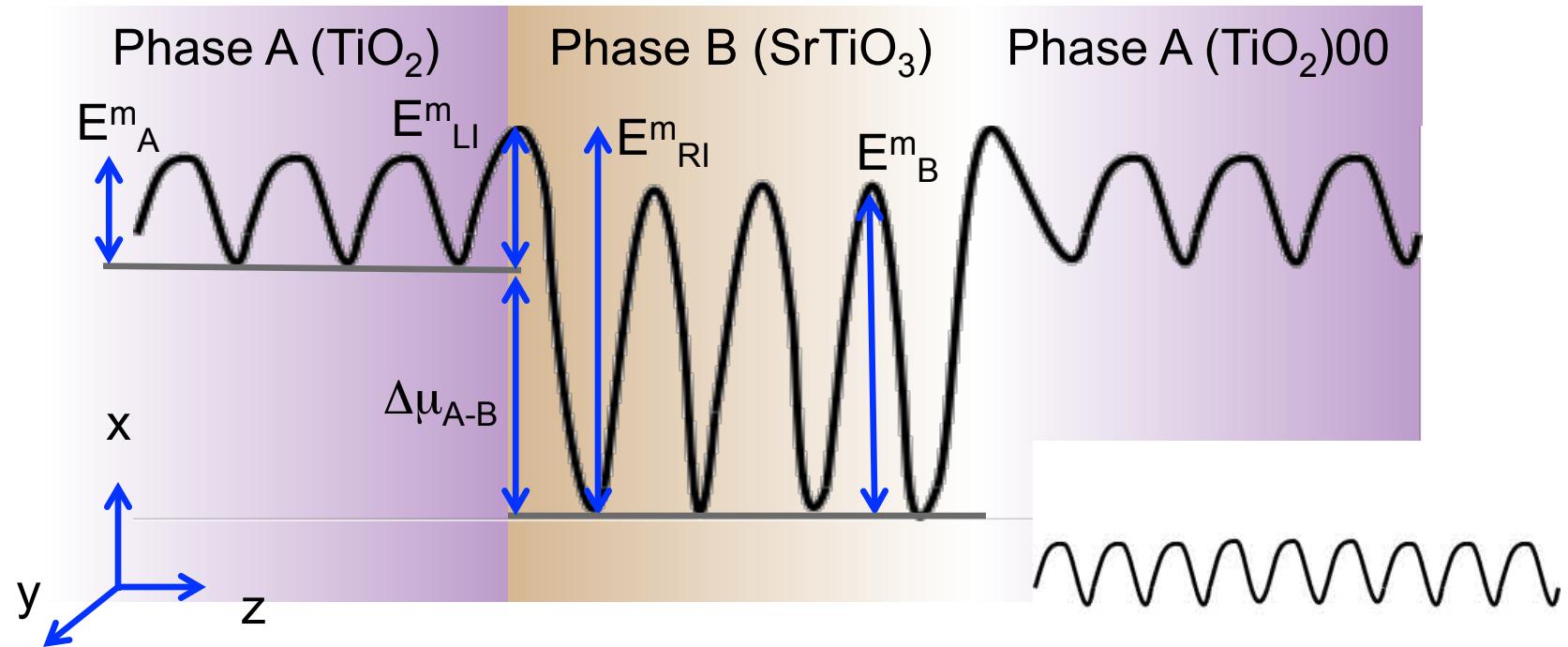


# Towards a Figure of Merit for “structureless” interfaces

- Differences in chemical potential  $\mu$  drive defect flow from one side to the other
  - Define phase A such that  $\mu_A > \mu_B$
- Differences in mobility  $E^m$  determine rate of defect buildup at interface
- *Kinetic sinks*



# Kinetic Monte Carlo model for structuraless interfaces



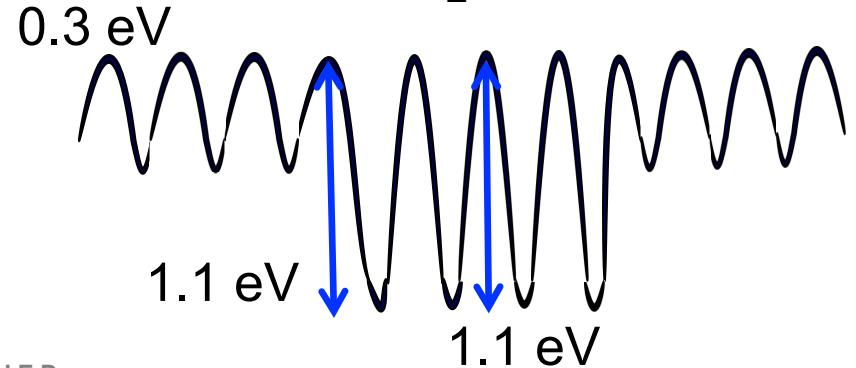
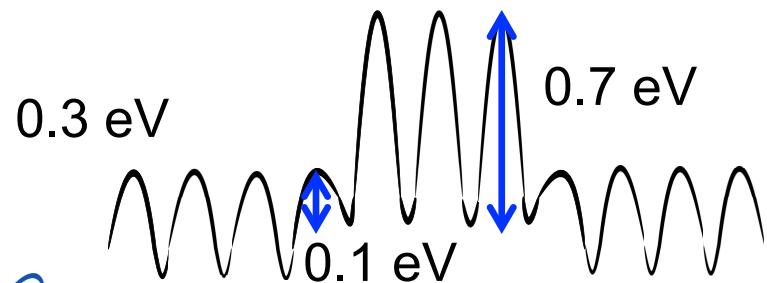
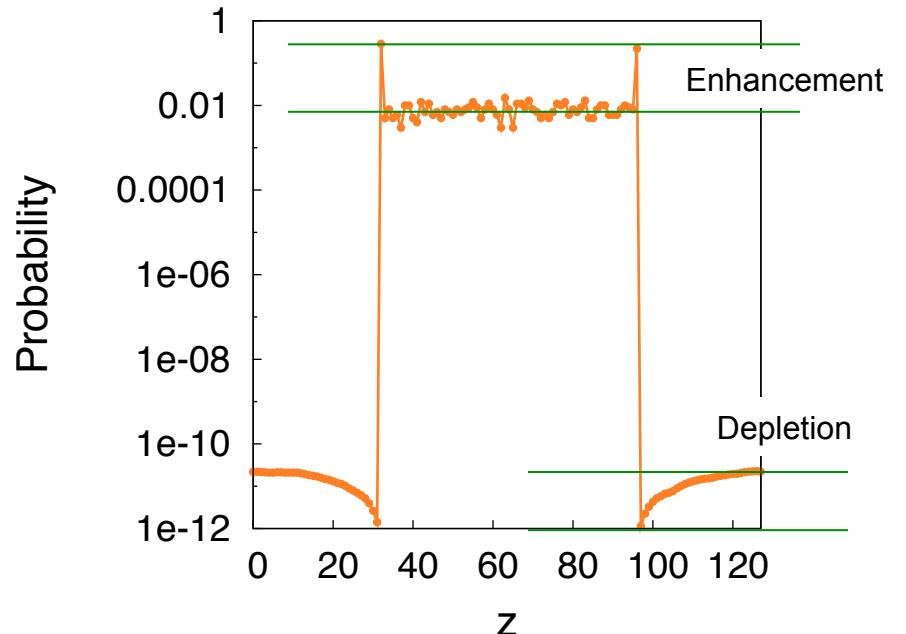
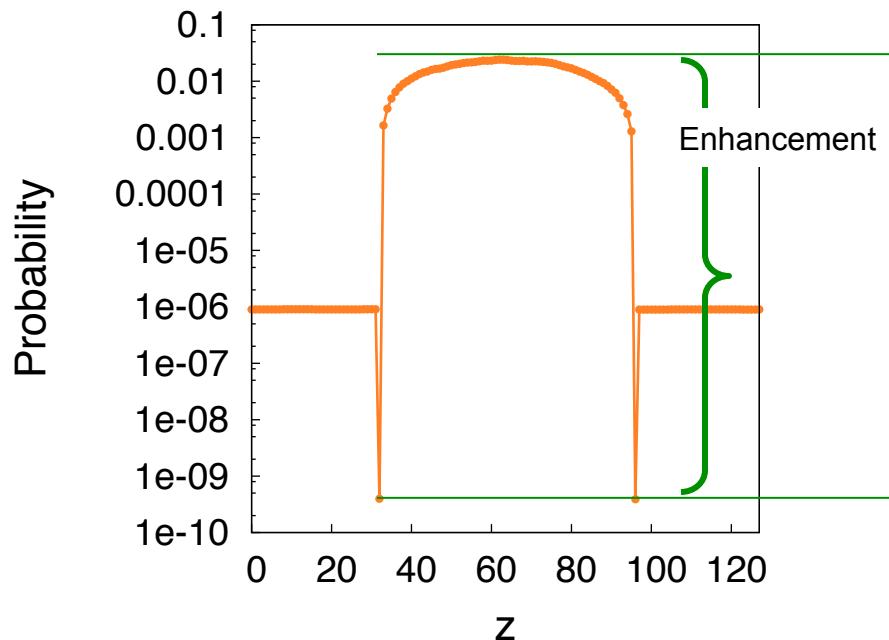
We have systematically change the migration parameters, run 1000 KMC independent calculations for each set and calculate the resulting defect profile.

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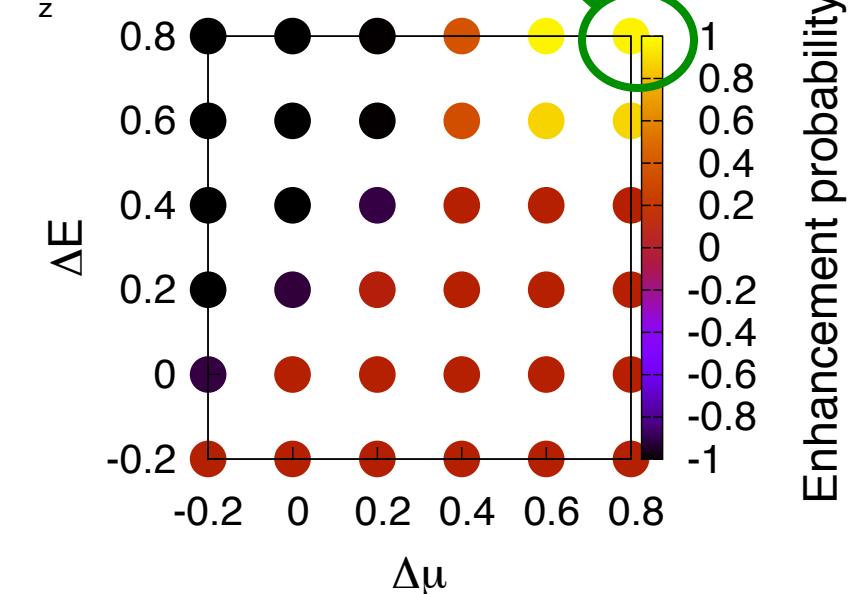
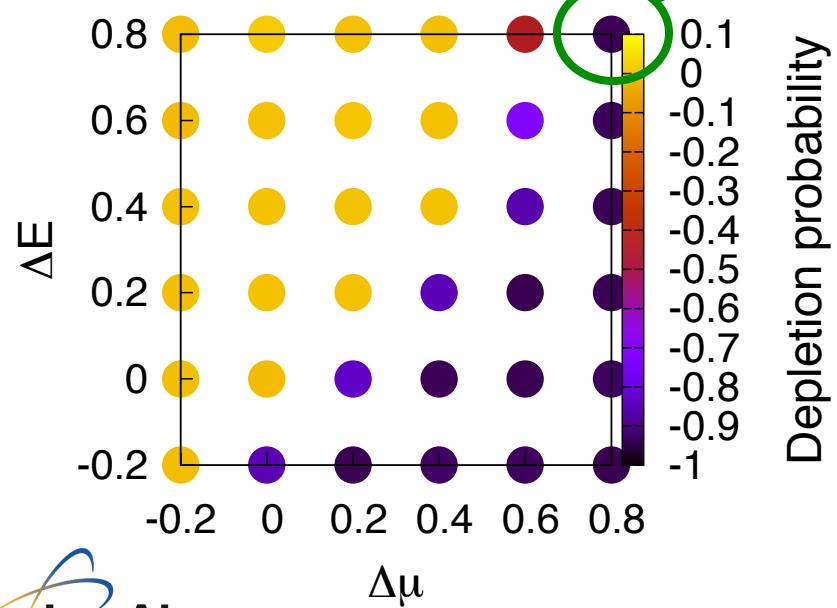


# Concentration profiles



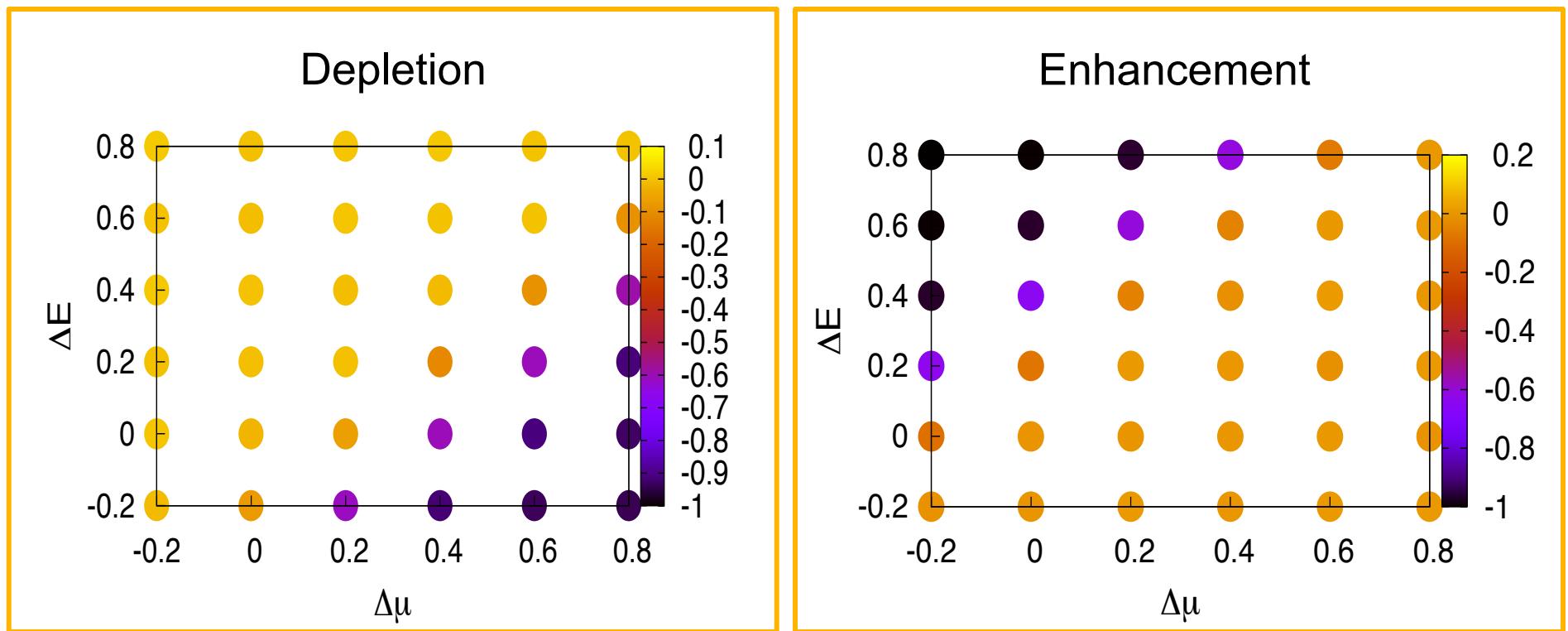
## $\Delta\mu - \Delta E$ maps at 300 K

The depletion probability is related to the length of the denuded zone



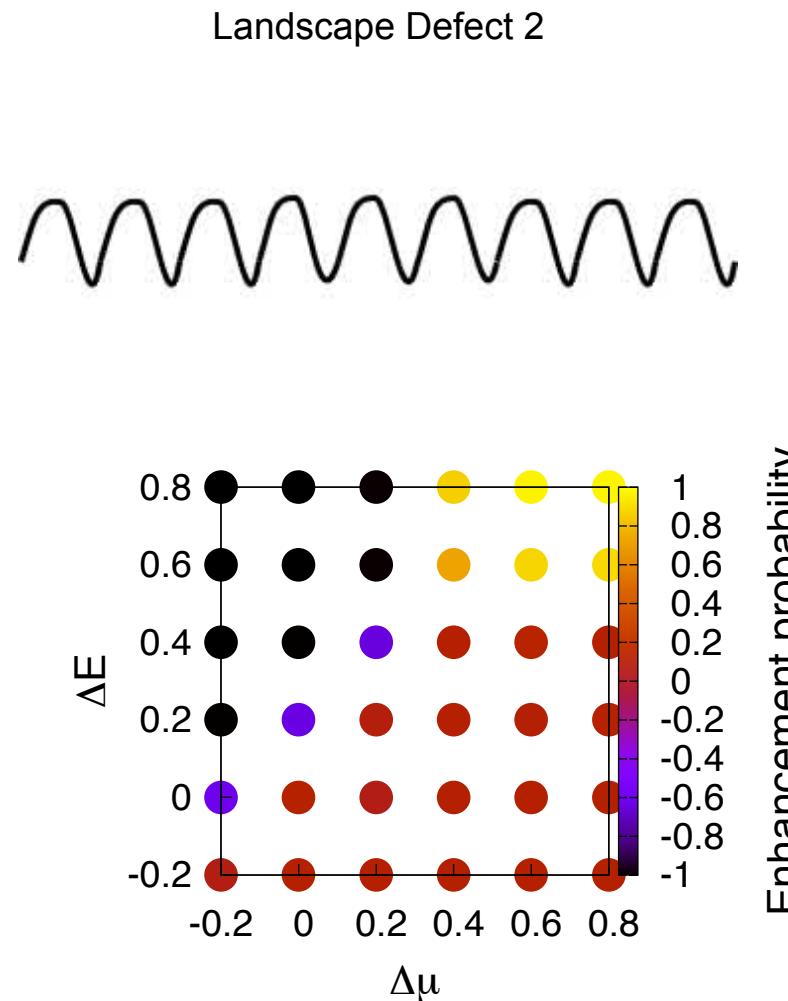
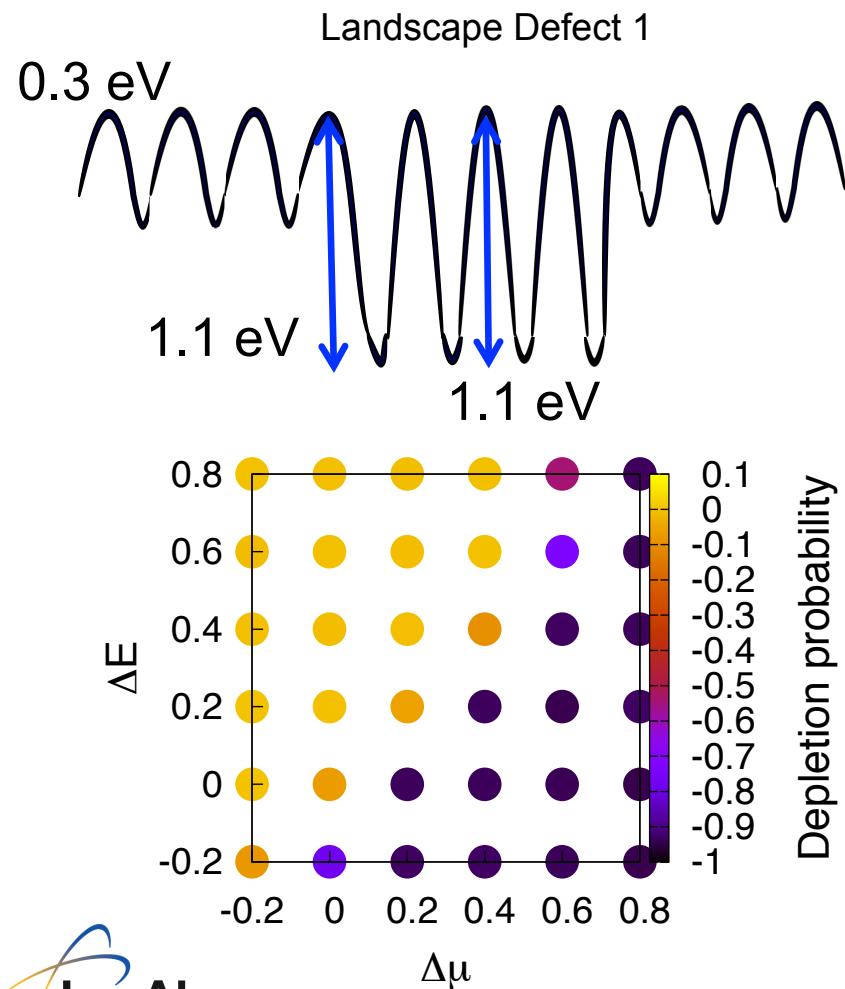
The enhancement probability is related to the amorphization

## $\Delta\mu - \Delta E$ maps at 800 K



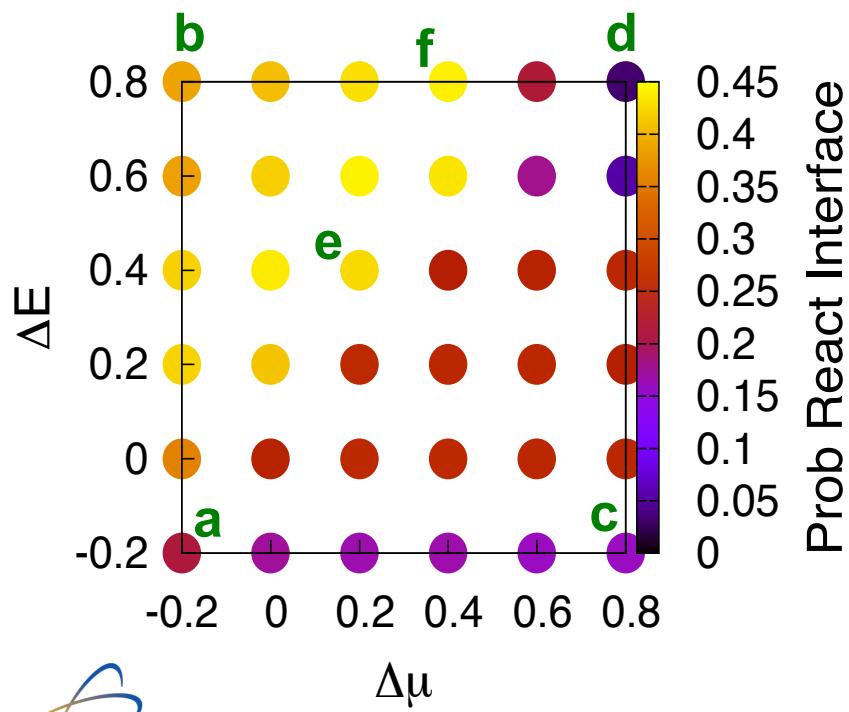
Increasing the temperature reduces the effect. It is a steady state effect

# Two Reacting Defects at 300 K

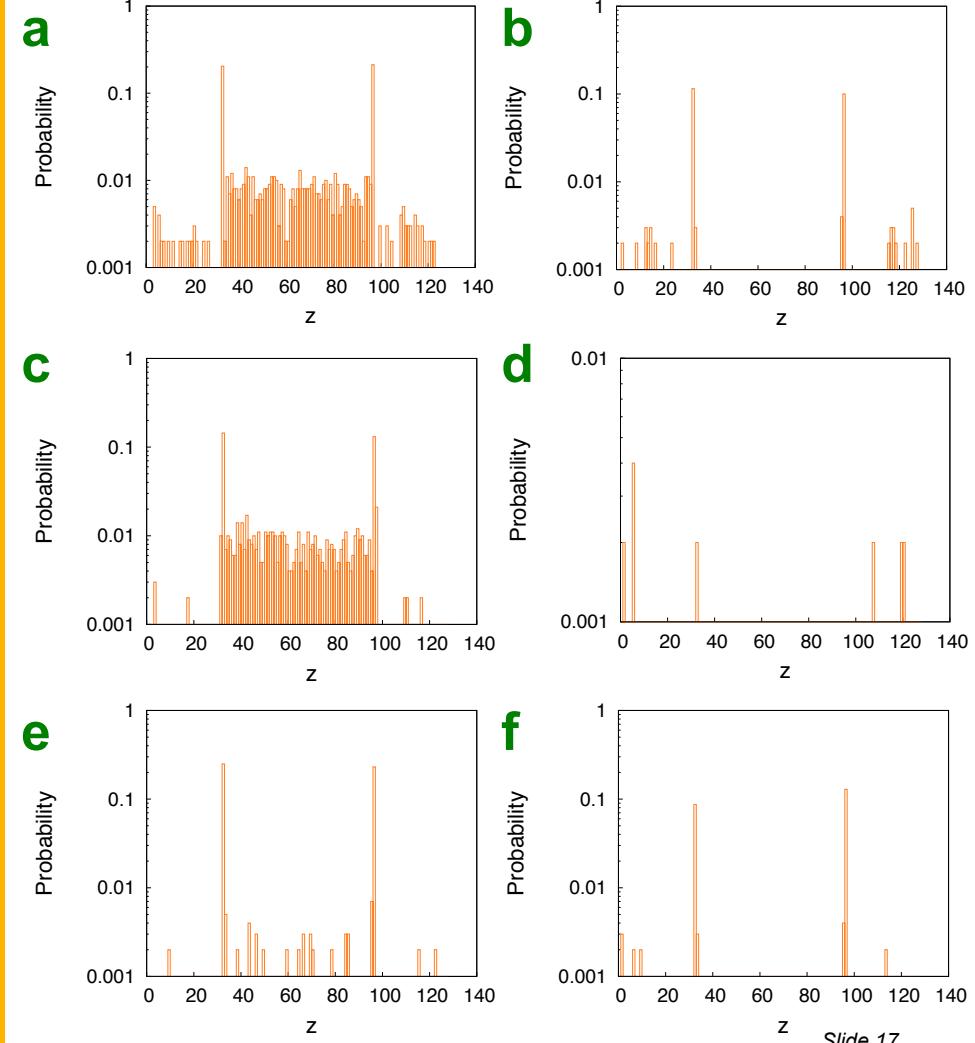


# Reaction Probability at the Interface

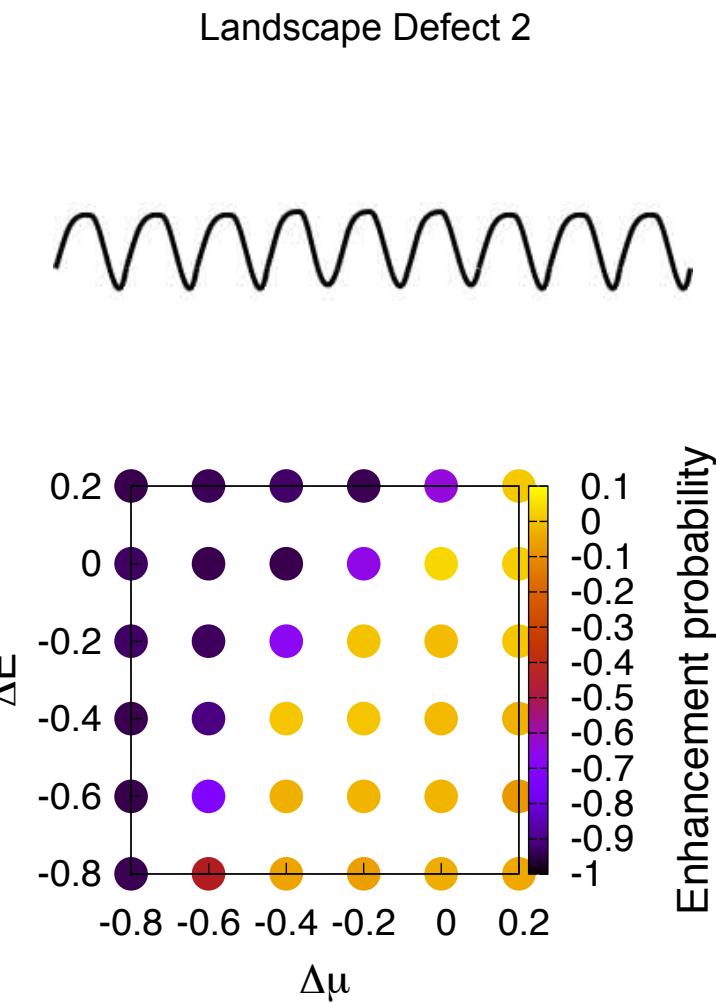
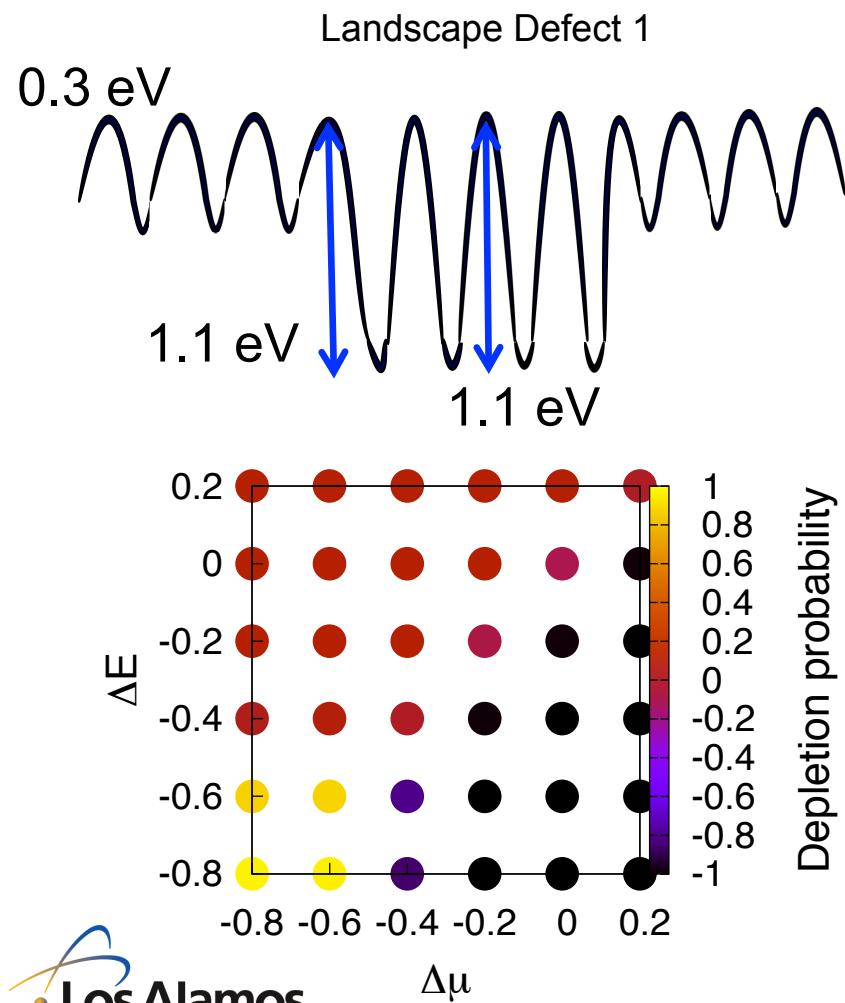
Probability of reaction at the interface provided that a reaction has taken place



## Total Probability of Reaction

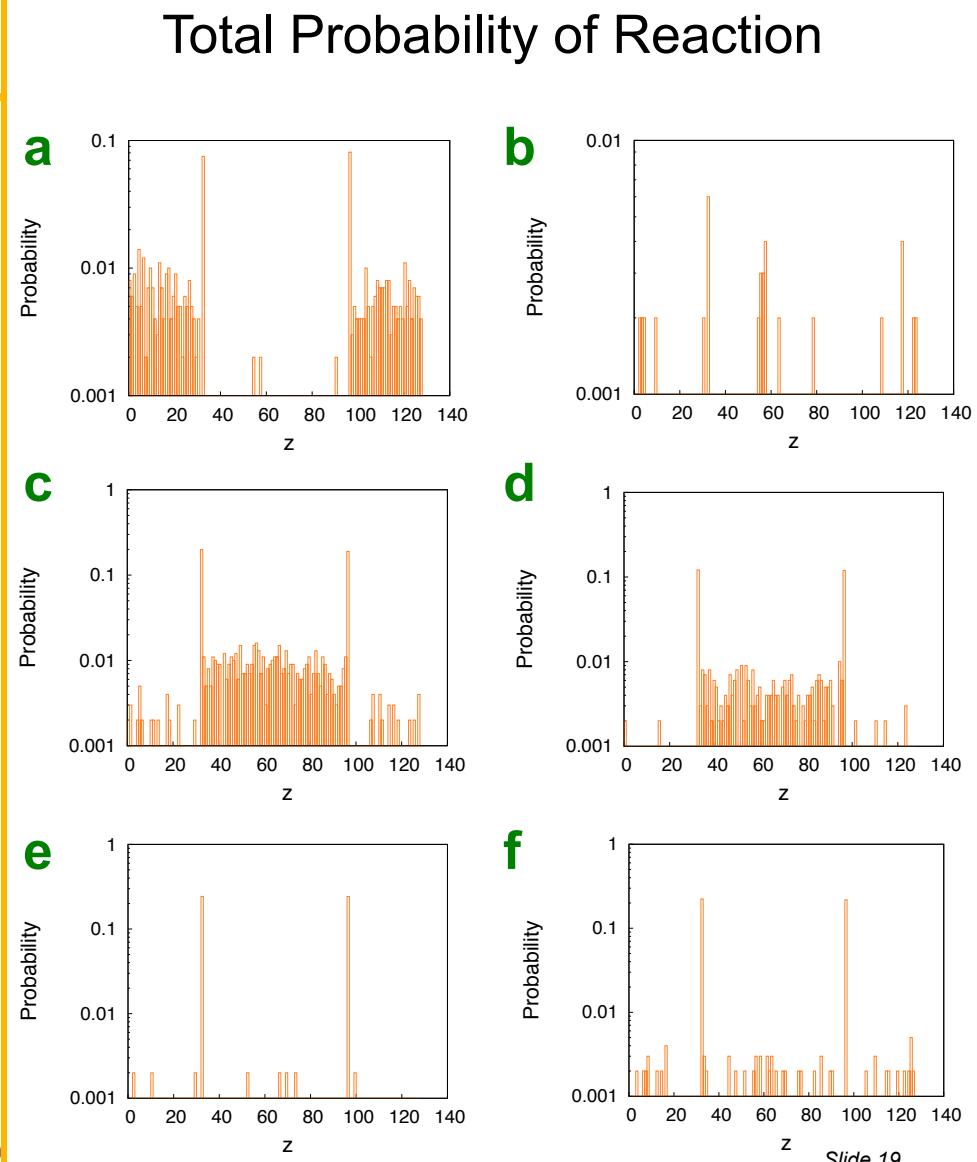
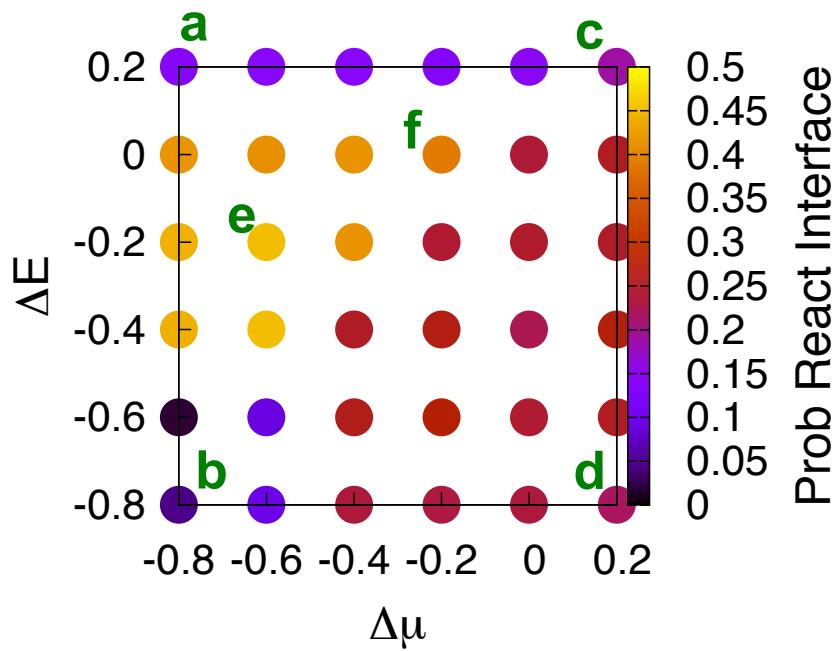


# Two Reacting Defects at 300 K



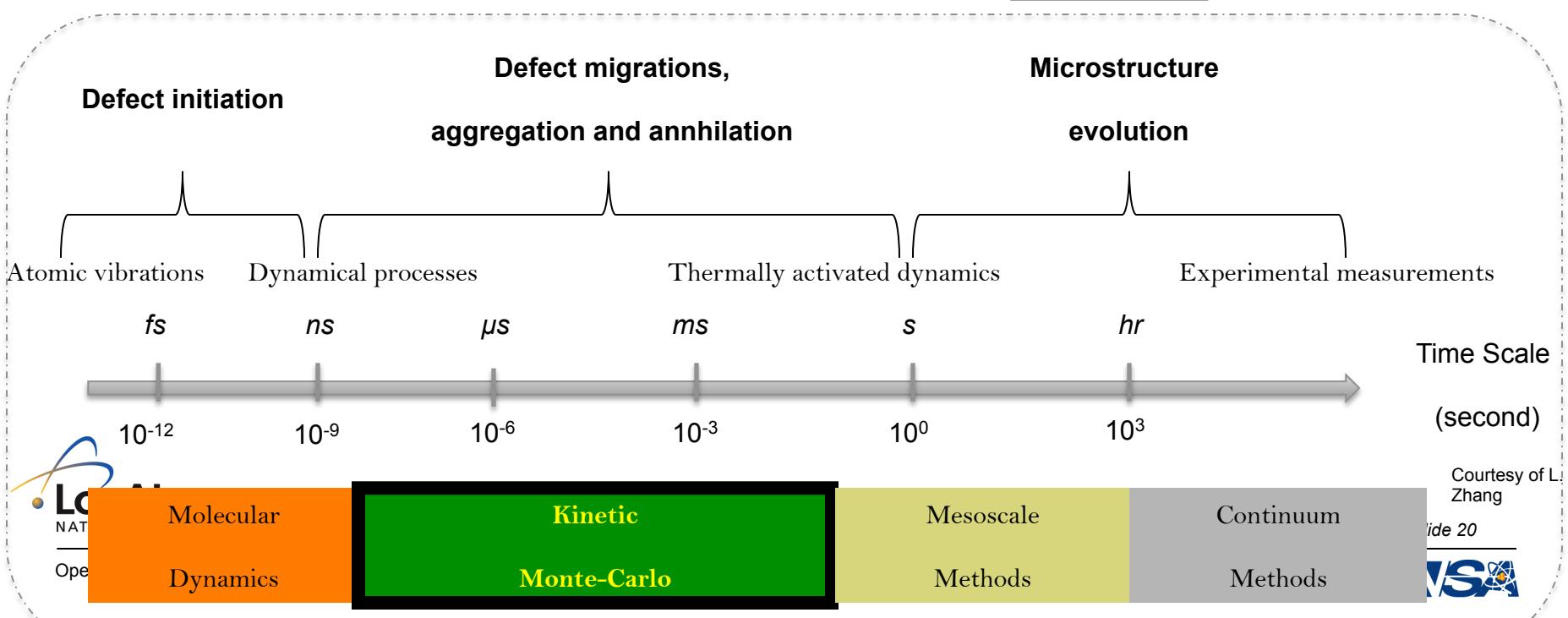
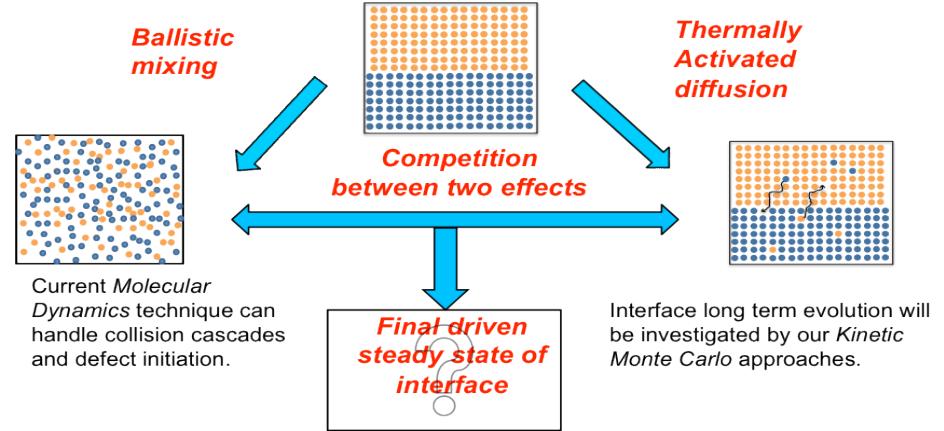
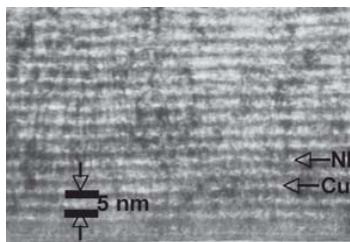
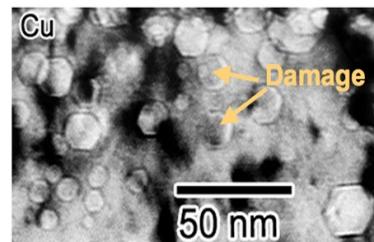
# Reaction Probability at the Interface

Probability of reaction at the interface provided that a reaction has taken place



# Motivation

## Sink behavior of Interfaces



# Our Kinetic Monte Carlo Approaches

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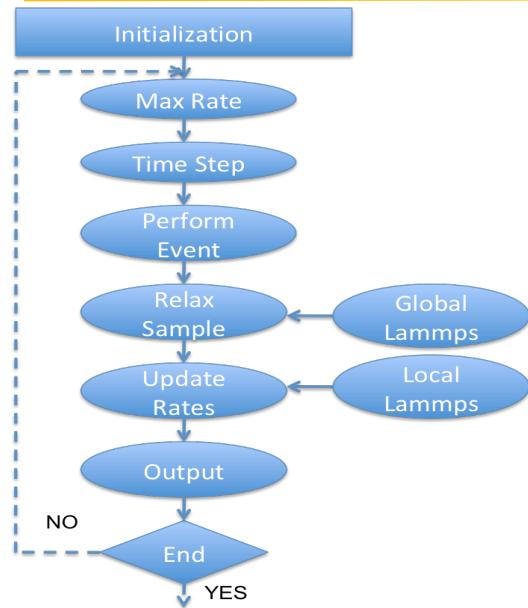
- **Self-learning KMC. (L. Vernon, B. Uberuaga, A. Voter)**

- Dynamically explores the potential energy surface to discover all processes.
- Calculate the rates accurately.
- Might be computationally demanding.

- **Event-driven KMC. (Me and Alfredo)**

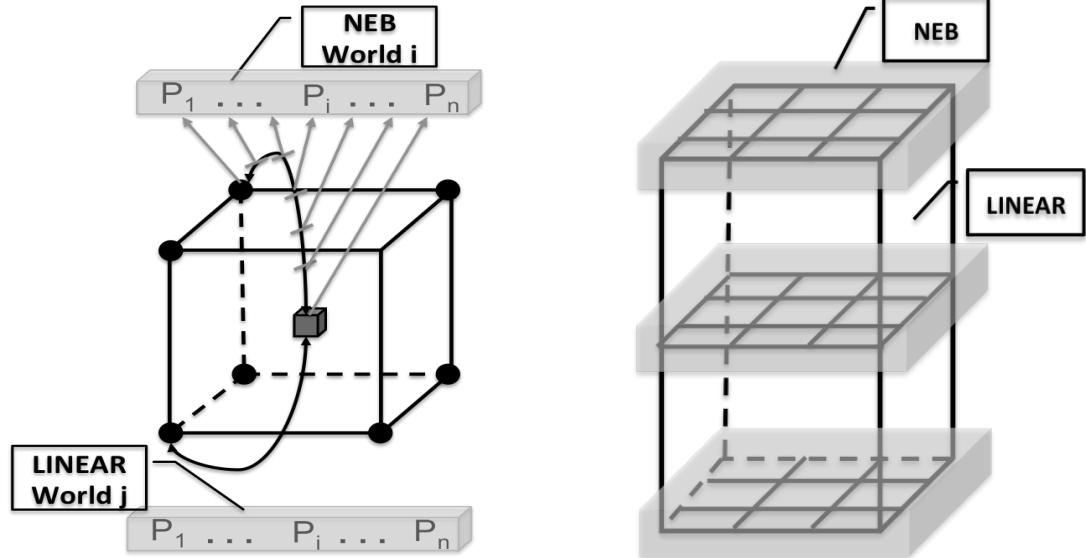
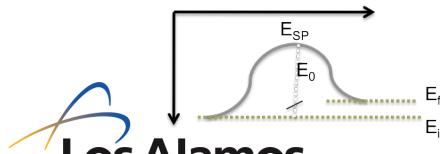
- Uses the local microstructure to guess potential processes.
- The accuracy in the rate calculation can be tuned.
- Computationally less demanding.

# Event-Driven KMC: Coupling MD LAMMPS – KMC



Set of possible events  
 Harmonic Transition State Theory  
 Particle rate:  $\Gamma = v \exp(-E_{SP}/kT)$ ; Two options to calculate  $E_{SP}$ :

1. Linear approximation  $E_{SP} = E_0 + (E_f - E_i)/2$
2. Nudged-elastic band (NEB)



Each vacancy exchange is handled by a set of processors. The global relaxation is performed using all processors

Hybrid models can be used depending on the complexity of the sample to obtain more accurate results

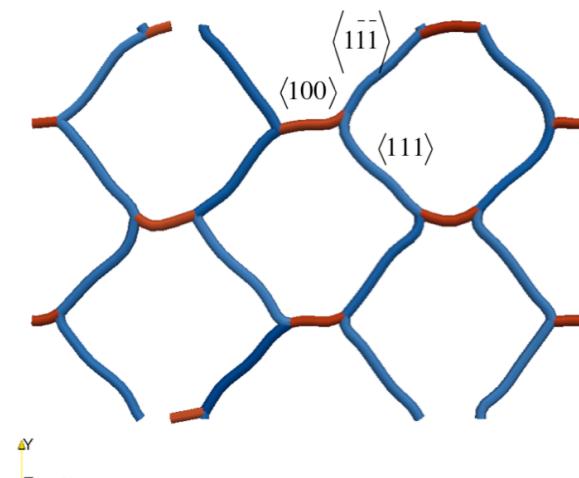
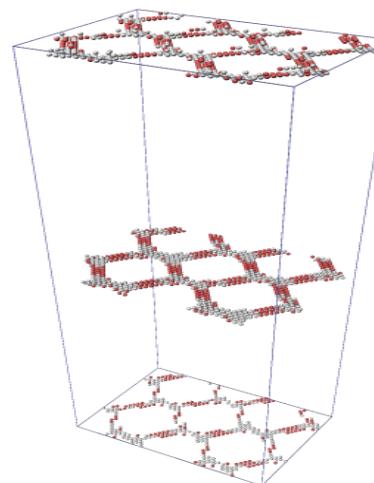
# Applications

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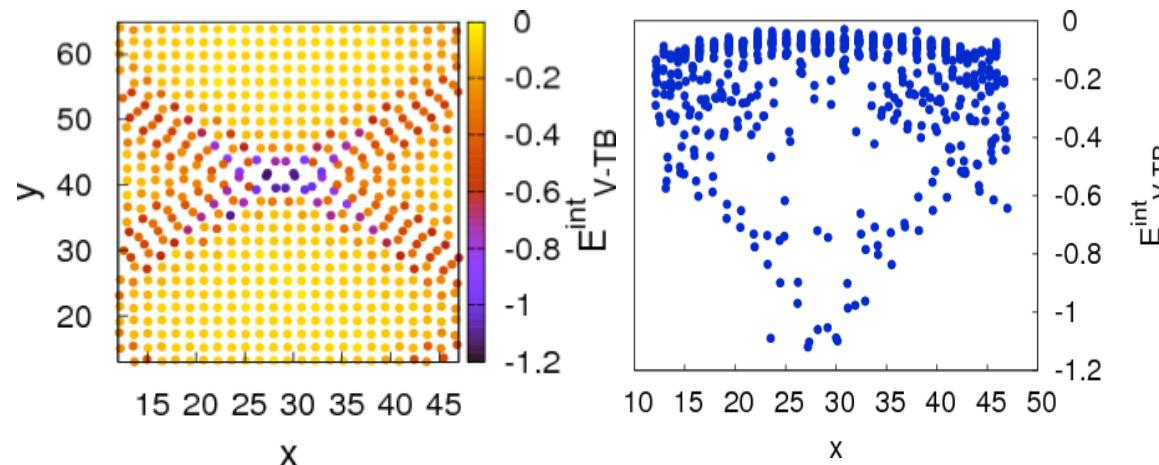
- **(110) 4-degrees Twist Boundary in bcc Fe.**
  - i. Interface structure
  - ii. Vacancy formation energies
  - iii. Vacancy accumulation using hybrid linear approx-NEB
  
- **(111) 2-degrees Twist Boundary in fcc Cu.**
  - i. Interface structure
  - ii. Vacancy formation energies
  - iii. Vacancy accumulation using hybrid linear approx-NEB

# Vacancy-Twist Boundary Interaction Energy in a (110) Fe Interface

Dislocation structure at the interface: 2 sets of  $a_0/2\langle 111 \rangle$  dislocations and 1 set of  $a_0\langle 100 \rangle$  screw dislocations



Interaction energy maps. Vacancies are attracted to dislocations. Thermodynamic driving force for the vacancies to accumulate at the interface.

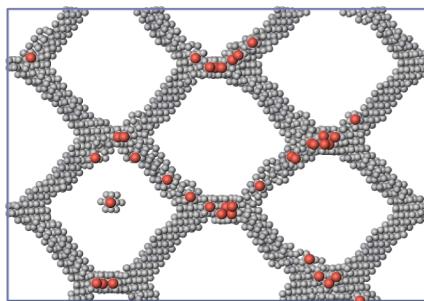


The interaction energy is most attractive at  $\langle 100 \rangle$  segments forming the MDIs

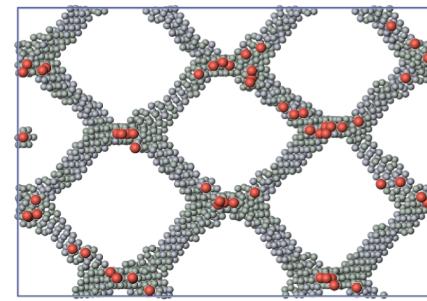
# Vacancy accumulation at the (110) Fe Interface

- **vacancies**
- Non-bcc atoms

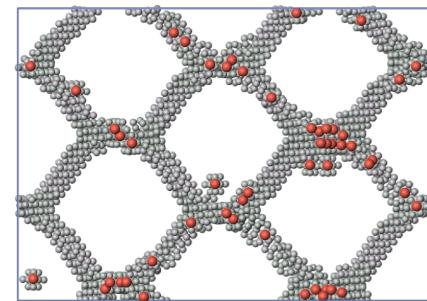
$$\Phi=10^{-4} \text{ V}/\text{A}^3 \cdot \text{s} = 1.16 \cdot 10^{-3} \text{ dpa}$$



$$\Phi=10^{-3} \text{ V}/\text{A}^3 \cdot \text{s} = 1.16 \cdot 10^{-2} \text{ dpa}$$



$$\Phi=10^{-2} \text{ V}/\text{A}^3 \cdot \text{s} = 1.16 \cdot 10^{-1} \text{ dpa}$$



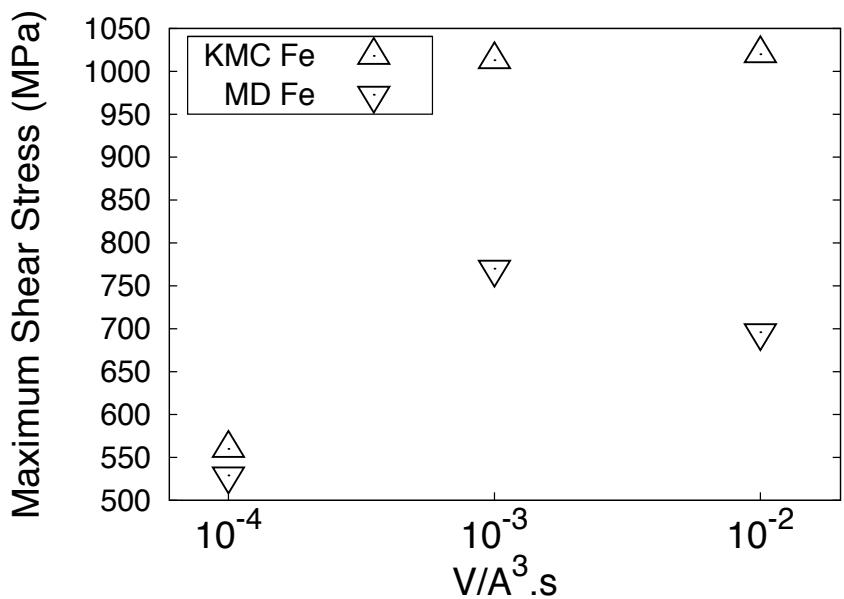
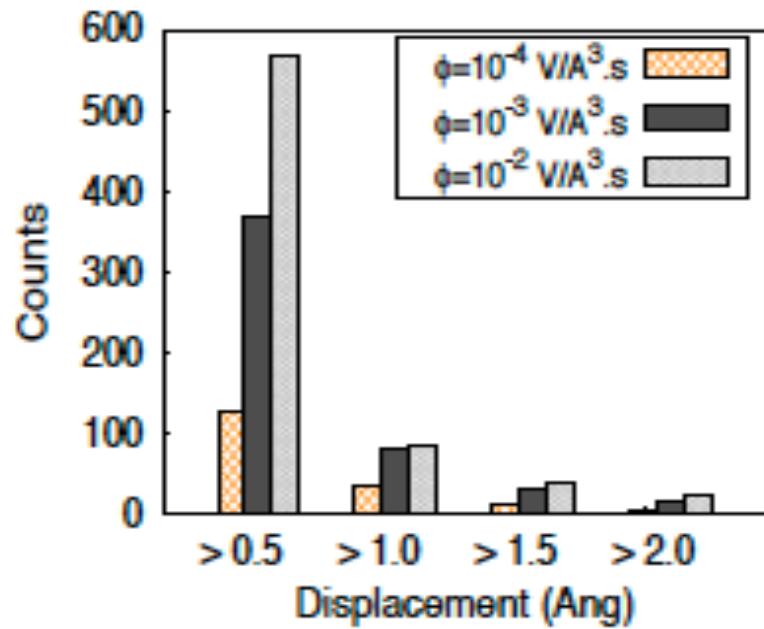
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Vacancies decorate the more energetic <100> segments forming the MDIs

25

# Error Estimation

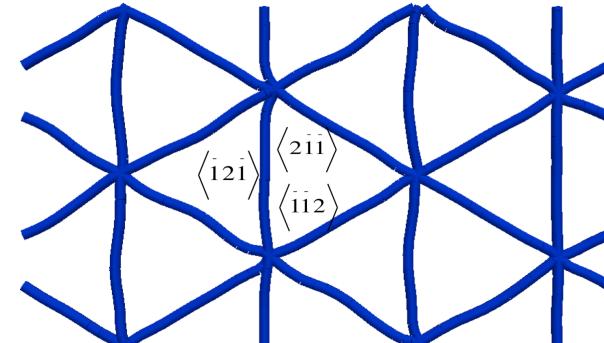
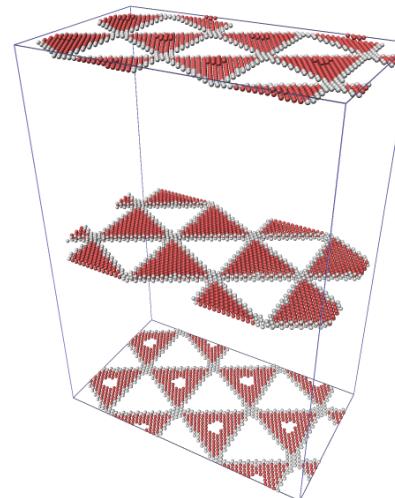
MD annealing for 1 ns at 500 K of KMC samples computing atomic displacements and maximum shear strength



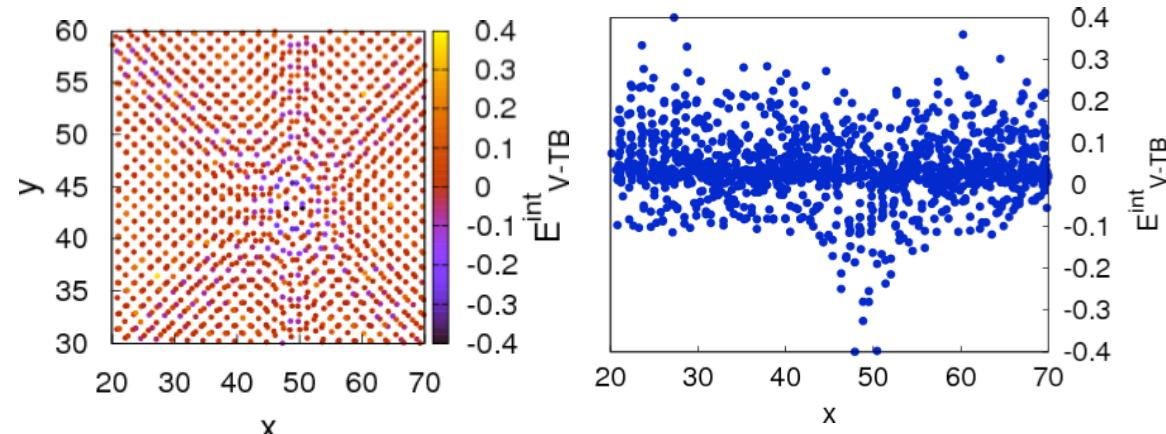
For the lowest dose rate investigated the number of events does not change the observable of interest. On the other hand, for larger dose rates, the maximum shear stress is  $\sim 30\%$  lower after annealing.

# Vacancy-Twist Boundary Interaction Energy in a (111) Cu Interface

Dislocation structure at the interface: 3 sets of  $a_0/6<112>$  Shockley partial screw dislocations



Interaction energy maps. Vacancies are attracted to dislocations. Thermodynamic driving force for the vacancies to accumulate at the interface.

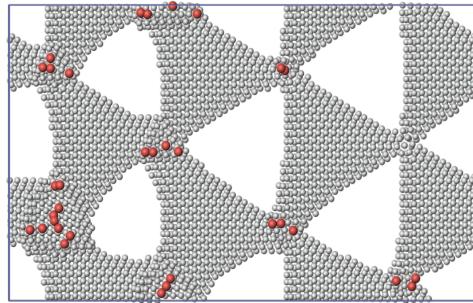


The interaction energy is most attractive at the MDIs

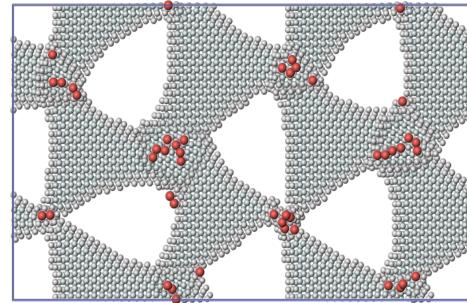
# Vacancy accumulation at the (111) Cu Interface

- **vacancies**
- Non-fcc atoms

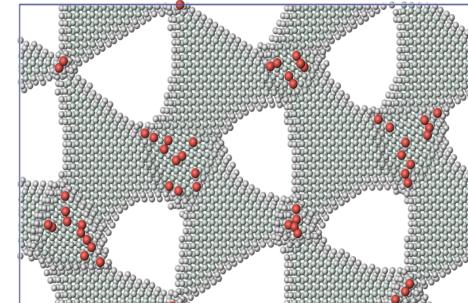
$\Phi=10^{-4} \text{ V}/\text{A}^3 \cdot \text{s} = 1.17 \cdot 10^{-3} \text{ dpa}$



$\Phi=10^{-3} \text{ V}/\text{A}^3 \cdot \text{s} = 1.17 \cdot 10^{-2} \text{ dpa}$



$\Phi=10^{-2} \text{ V}/\text{A}^3 \cdot \text{s} = 1.17 \cdot 10^{-1} \text{ dpa}$



• **Los Alamos**

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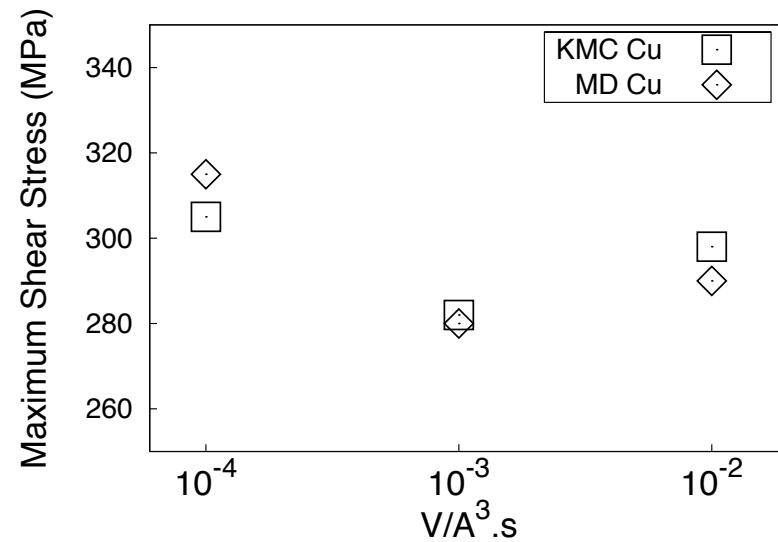
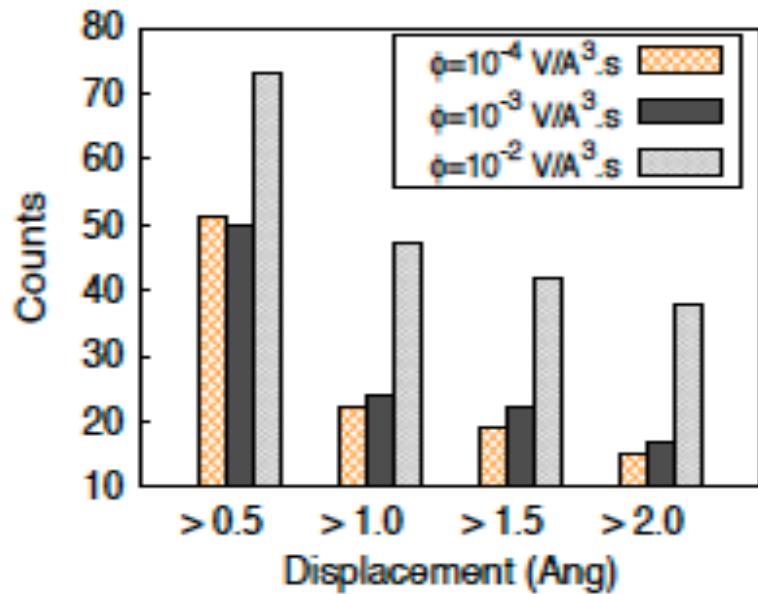
Vacancies agglomerate at the MDIs changing the dislocation structure

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# Error Estimation

MD annealing for 1 ns at 500 K of KMC samples computing atomic displacements and maximum shear strength



In this case the events do not change the observable of interest at any dose rate. The maximum difference in the shear stress is  $\sim 3\%$ .

# Conclusions

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- We have developed an hybrid KMC-MD off-lattice algorithm
- It is able to describe diffusion with all the fields (elastic, thermodynamic)
- We can reach real times far beyond MD capabilities.
- Application to CuNb interfaces (in collaboration with Liang Zhang-MIT).
- Improvement of the physics:
  - Implementation of interstitial diffusion.

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