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**DESIGN AND ANALYSIS OF RTGS  
FOR CRAF AND CASSINI MISSIONS**

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# DESIGN AND ANALYSIS OF RTGs FOR CRAF AND CASSINI MISSIONS

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## ABSTRACT

This report consists of two parts. Part I describes the development of novel analytical methods needed to predict the BOM performance and the subsequent performance degradation of the mutually obstructed RTGs for the CRAF and Cassini missions. Part II applies those methods to the two missions, presents the resultant predictions, and discusses their programmatic implications.

The results indicate that JPL's original power demand goals could have been met with two standard GPHS RTGs for each mission. But subsequently JPL significantly increased both the power level and the mission duration for both missions, so that they can no longer be met by two standard RTGs. The resultant power gap must be closed either by reducing JPL's power demand (e.g., by decreasing contingency reserves) and/or by increasing the power system's output. One way under active consideration which more than meets the system power goal would be the addition of a third RTG for each mission.

However, the author concluded that it may be possible to meet or closely approach the CRAF power demand goals with just two RTGs by relatively modest modification of their design and/or operating conditions. To explore that possibility, the effect of various modifications - either singly or in combination - was analyzed by Fairchild. The results indicate that modest modifications can meet or come very close to meeting the CRAF power goals with just two RTGs. Elimination of the third RTG would yield substantial cost and schedule savings.

## INTRODUCTION

The Jet Propulsion Laboratory (JPL) is in the process of designing spacecraft for NASA's upcoming CRAF and Cassini missions [1]. The CRAF (Comet Rendezvous and Asteroid Flyby) spacecraft, planned for a 1995 launch, will fly by at least one asteroid, orbit a comet, and rendezvous with a comet, where it will conduct scientific observations for more than two years. The Cassini spacecraft, scheduled for a 1996 launch, will also fly by at least one asteroid, fly by Jupiter, orbit the planet Saturn, repeatedly fly by a number of Saturn's moons, and send an instrumented probe, called the Huygens probe, into the atmosphere of Saturn's moon Titan.

Both missions will be launched by unmanned Titan-4/Centaur-G' boosters, and both are part of a series which will use a new generation of cost-effective modular spacecraft (Mariner/Mark 2) that can easily be modified to accomplish a variety of missions to comets and asteroids and to the outer solar system.

Because of their great distance from the sun, both missions will be powered by RTGs (Radioisotope Thermoelectric Generators). Since the planned schedules do not allow sufficient time for development and flight qualification of new RTG designs, both missions will use generators that are essentially identical to the GPHS-RTGs flown on the recently launched Galileo mission. JPL's original plans called for two such RTGs on each mission. Their construction will be under the direction of the U.S. Department of Energy's Office of Special Applications (DOE/OSA), which commissioned Fairchild Space Company to conduct RTG studies in support of JPL's design efforts.

The location and orientation of the two RTGs on the spacecraft are functions of numerous, often conflicting, design constraints. The CRAF/Cassini baseline design that JPL initially asked Fairchild to study is depicted in Figure 1.

Figure 1. CRAF/CASSINI Spacecraft, Baseline Design

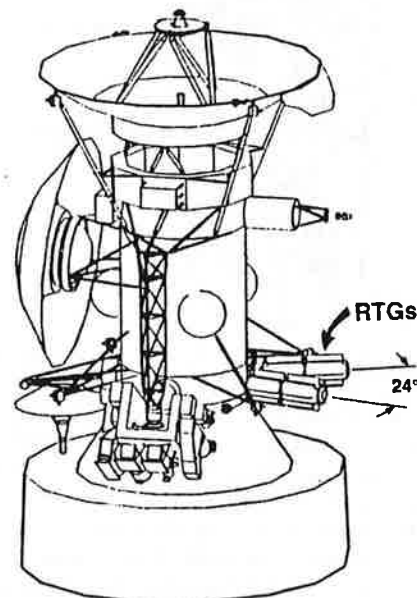


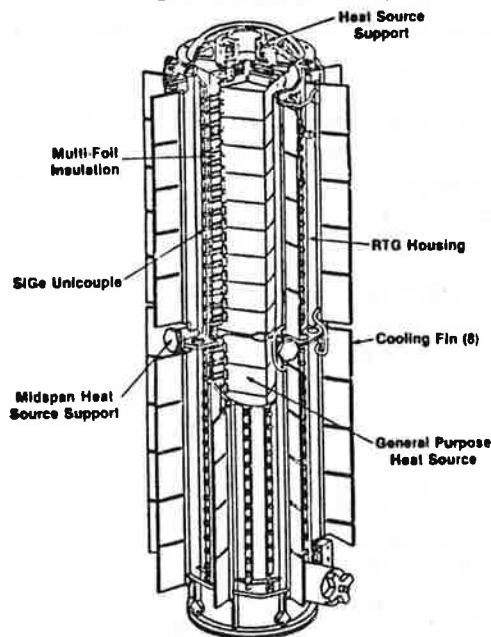
Figure 1 shows the two RTGs cantilevered from the cylindrical spacecraft like radial spokes, with a separation angle of 24 degrees between them. Such proximity leads to mutual obstruction of the two RTGs' heat rejection paths. This obstruction can result in significant axial and circumferential variation of the cold-junction temperatures and the electrical performance of the thermoelectric couples in each RTG.  $\approx$

To assess the effect of mutual obstructions on the RTGs' output power, Fairchild personnel were asked to analyze the baseline configuration shown in Figure 1, as well as some alternative configurations. Because of the unconventional problem, this required the development of novel analysis methods and computer codes, which are described in this paper. The analytical results reported here will serve as a guide for JPL's spacecraft design decisions.

### RTG DESIGN DESCRIPTION

Figure 2 shows a cutaway view of the GPHS-RTG [2, 3, 4] that will be used on the CRAF and Cassini missions. Each 1.15 m-long RTG contains an axial stack of 18 General Purpose Heat Source modules [5], which radiate their heat to a surrounding cylindrical array of 576 thermoelectric unicouples arranged in 36 layers of 16 couples.

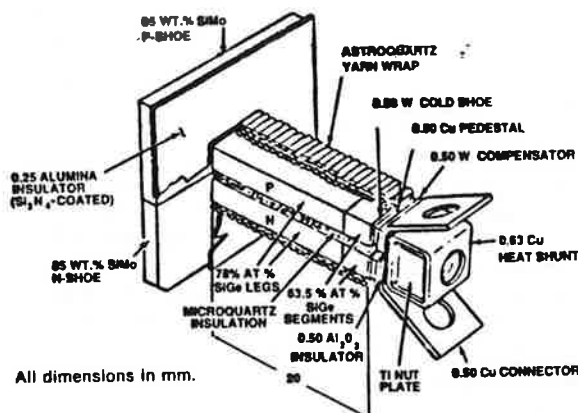
Figure 2. GPHS RTG



As shown in Figure 3, each unicouple contains a thermoelectric n- and p-leg. These are electrically connected at their hot ends by a hot-shoe, which serves to collect the heat radiated by the centrally located heat source stack and concentrate it at the thermoelectric legs. The cold end of each leg is series- and parallel-connected to adjoining couples to form the RTG's electrical network. The couples' cold ends are bolted to the RTG housing, to which they

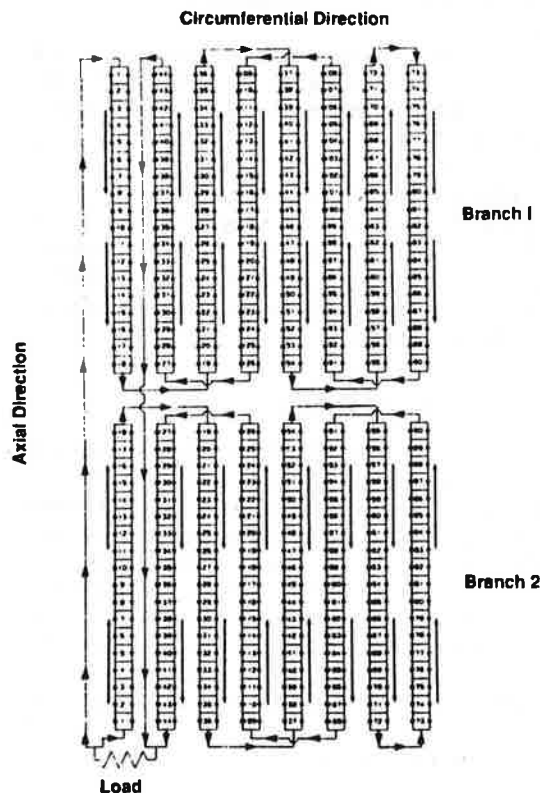
reject their waste heat. The RTG's waste heat is dissipated by radiation from its housing and its eight equispaced radiator fins, as shown in Figure 2.

Figure 3. Unicouple



The series-parallel network of the GPHS-RTG is shown in Figure 4, which depicts a rolled-out schematic of the cylindrical array. The electrical network consists of two parallel branches. Each branch contains 144 series-connected groups of two parallel couples. The rather tortuous current path shown is designed to minimize the RTG's self-induced magnetic field.

Figure 4. Current-Path Through GPHS RTG



## ANALYSIS

The close proximity of the two RTGs on the CRAF/Cassini spacecraft, as illustrated by the baseline design shown in Figure 1, can result in significant mutual blockage of their heat rejection paths. Such blockage would result in circumferential variation of the RTG's housing and cold-junction temperatures. Determining the effect of that temperature variation on the RTGs' power output requires a very detailed and careful analysis, because we are looking for relatively small differences between large numbers, and because even small differences can be quite significant if the mission is power-constrained (as are the CRAF and Cassini missions).

Previous RTG analyses usually made the simplifying assumption that all of the thermoelectric couples in a generator's series-parallel network operate at the same hot- and cold-junction temperatures and at the same current and voltage. For unobstructed RTGs, such a simplified analysis is a useful initial design tool, since it permits closed-form solutions for the optimum area ratio  $A_n/A_p$  of the thermoelectric n- and p-legs and for the optimum output voltage. For these optimized parameters, it yields simple expressions for the maximum material efficiency of the thermoelectric couples, and for the required RTG design parameters [6, 7, 8].

But the above simplifying assumptions can introduce significant errors even for unobstructed RTGs, because all RTGs have appreciable axial temperature variations due to unavoidable end losses by radiation and by conduction through the heat source support structure. A more exact analysis, which accounted for the axial temperature variations in a Martian RTG, was reported last year [7, 9]. But that analysis still assumed that the RTG has an axisymmetric view of space and the Martian ground, and therefore no circumferential temperature variation.

The present paper develops a Fairchild-generated methodology and generalized computer code for analyzing the performance of arbitrarily obstructed RTGs with both axial and circumferential temperature, voltage, and current variations, and applies that methodology to the specific example of the CRAF/Cassini baseline design depicted in Figure 1.

## COUPLED THERMAL AND ELECTRICAL ANALYSIS

Figure 5 presents the energy balance for a thermoelectric uncouple of leg length  $L$ , leg areas  $A_n$  and  $A_p$ , operating between cold- and hot-junction temperatures  $T_1$  and  $T_2$ . It gives the couple's thermal conductance  $K$ , electric resistance  $R$ , and open-circuit voltage  $V_o$  in terms of the temperature-averaged thermal conductivity  $\bar{k}$ , electrical resistivity  $\bar{\rho}$ , and Seebeck coefficient  $\bar{S}$  of the thermoelectric n- and p-

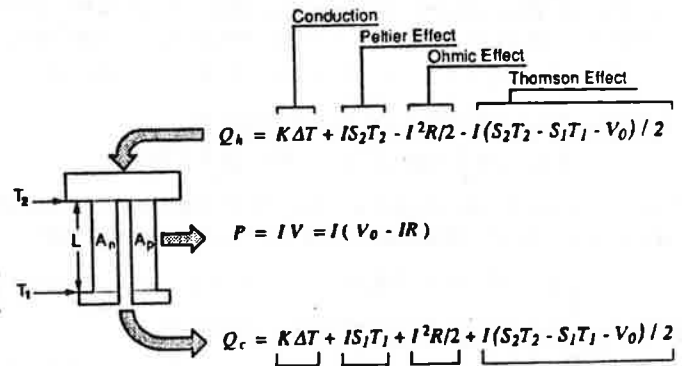
materials. As shown, the heat input rate  $Q_h$  at the couple's hot end and the heat rejection rate  $Q_c$  at its cold end each consists of four terms: normal heat conduction, Peltier effect, Ohmic dissipation, and Thomson effect. As can be seen, three of those four terms are current-dependent. Therefore, the thermal and electrical analyses cannot be performed separately, but must be conducted simultaneously and interactively.

Figure 5. Uncouple Energy Balance

$$\text{Open-Circuit Voltage: } V_o = \int (S_n + S_p) dT$$

$$\text{Couple Conductance: } K = [\bar{k}_n A_n + \bar{k}_p A_p] / L$$

$$\text{Couple Resistance: } R = [\bar{\rho}_n / A_n + \bar{\rho}_p / A_p] L$$



To analyze an RTG of a given design with an unsymmetrically obstructed heat rejection path, a detailed three-dimensional thermal model of the RTG and its environment must be constructed. The hot junction and cold junction of each thermoelectric element in the RTG are represented as discrete nodes. The model cannot be analyzed by means of a standard thermal analysis code, because the connectors between the couple's hot and cold junctions are not simple thermal conductors but include the current-dependent Peltier, Ohmic, and Thomson effects. The rate at which heat enters the connector's hot-end and leaves its cold-end are not equal, since part of the heat entering each couple is converted to electrical energy. That part must in effect be represented as a heat sink for each couple.

The electrical analysis is further complicated by the constraint that each RTG's thermoelectric couples are in general interconnected in a complex series-parallel network, and that all couples grouped in parallel must operate at the same output voltage, and that all couple groups in series must produce the same current.

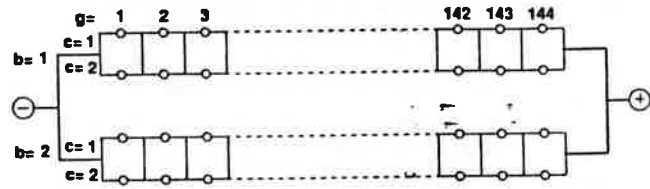
The RTG analysis methodology developed by Fairchild is generic, not just for the GPHS-RTG network shown in Figure 4. In general, the equivalent

circuit of an RTG network consists of  $B$  parallel branches, with each branch containing  $G$  series-connected groups of  $C$  parallel couples. Each thermoelectric element in the RTG is designated by a branch number  $b$ , group number  $g$ , and couple number  $c$ , where:

$$\begin{aligned} 1 &\leq b \leq B, \\ 1 &\leq g \leq G, \\ 1 &\leq c \leq C. \end{aligned}$$

For the case of the GPHS-RTG,  $B=2$ ,  $G=144$ ,  $C=2$ , and the equivalent circuit of the generator is shown in Figure 6.

Figure 6. Equivalent Circuit of GPHS RTG



## ELECTRICAL ANALYSIS

For each couple in the RTG, the difference between its open-circuit voltage  $V_o(b,g,c)$  and its internal voltage drop  $I(b,g,c) R(b,g,c)$  equals its output voltage  $V(b,g,c)$ . Since the couples in each group  $b,g$  are connected in parallel, their output voltage must equal the group voltage  $V(b,g)$ .

$$V_o(b,g,c) - I(b,g,c) R(b,g,c) = V(b,g,c) = V(b,g)$$

Therefore, the current through couple  $b,g,c$  is:

$$I(b,g,c) = [V_o(b,g,c) - V(b,g)] / R(b,g,c).$$

The sum of each group's  $C$  couple currents equals the group current  $I(b,g)$ . Since all groups in each branch  $b$  are connected in series, this must equal the branch current  $I(b)$ :

$$\sum_{c=1}^C [V_o(b,g,c) / R(b,g,c)] - V(b,g) \sum_{c=1}^C [1/R(b,g,c)] = I(b,g) = I(b)$$

Therefore, the voltage produced by group  $b,g$ , is:

$$V(b,g) = \left\{ \sum_{c=1}^C [V_o(b,g,c) / R(b,g,c)] - I(b) \right\} / \sum_{c=1}^C [1/R(b,g,c)]$$

The sum of the  $G$  group voltages in each branch equals the branch voltage  $V(b)$ . Since the  $B$  branches are connected in parallel, this also equals the RTG's output voltage  $V_{RTG}$ :

$$\sum_{g=1}^G \left\{ \sum_{c=1}^C [V_o(b,g,c) / R(b,g,c)] / \sum_{c=1}^C [1/R(b,g,c)] \right\} - I(b) \sum_{g=1}^G \left\{ 1 / \sum_{c=1}^C [1/R(b,g,c)] \right\} = V(b) = V_{RTG}.$$

Solving for the branch current  $I(b)$  for each branch and summing the currents for the  $B$  parallel branches gives the RTG output current:

$$I_{RTG} = \sum_{b=1}^B \left\{ \left[ \sum_{g=1}^G \left\{ \sum_{c=1}^C [V_o(b,g,c) / R(b,g,c)] / \sum_{c=1}^C [1/R(b,g,c)] \right\} - V_{RTG} \right] / \sum_{g=1}^G \left\{ 1 / \sum_{c=1}^C [1/R(b,g,c)] \right\} \right\}.$$

The above equation represents the current-voltage characteristic of the RTG network. It can be expressed in the condensed form

$$I_{RTG} = I_{SC} - V_{RTG}/R_{RTG},$$

where  $I_{SC}$  is the RTG's short-circuit current, defined by

$$I_{SC} = \sum_{b=1}^B \left[ \sum_{g=1}^G \left\{ \sum_{c=1}^C [V_o(b,g,c) / R(b,g,c)] / \sum_{c=1}^C [1/R(b,g,c)] \right\} / \sum_{g=1}^G \left\{ 1 / \sum_{c=1}^C [1/R(b,g,c)] \right\} \right],$$

$R_{RTG}$  is the RTG's internal resistance, defined by

$$R_{RTG} = 1 / \sum_{b=1}^B \left[ 1 / \sum_{g=1}^G \left\{ 1 / \sum_{c=1}^C [1/R(b,g,c)] \right\} \right],$$

and the product  $I_{SC} R_{RTG}$  is the RTG's open-circuit voltage

## ITERATIVE COMPUTATIONS

For each iteration in the analysis, the code uses each couple's cold- and hot-junction temperatures  $T_1$  and  $T_2$  (from the preceding iteration) to compute its temperature-averaged properties  $\bar{k}$  and  $\bar{\rho}$ , open-circuit voltage  $V_o$ , thermal conductance  $K$  and electrical resistance  $R$ :

$$\begin{aligned}\bar{k}_n(b, g, c) &= \int_{T_1(b, g, c)}^{T_2(b, g, c)} k_n(T) dT / [T_2(b, g, c) - T_1(b, g, c)] \\ \bar{k}_p(b, g, c) &= \int_{T_1(b, g, c)}^{T_2(b, g, c)} k_p(T) dT / [T_2(b, g, c) - T_1(b, g, c)] \\ \bar{\rho}_n(b, g, c) &= \int_{T_1(b, g, c)}^{T_2(b, g, c)} \rho_n(T) k_n(T) dT / \bar{k}_n(b, g, c) [T_2(b, g, c) - T_1(b, g, c)] \\ \bar{\rho}_p(b, g, c) &= \int_{T_1(b, g, c)}^{T_2(b, g, c)} \rho_p(T) k_p(T) dT / \bar{k}_p(b, g, c) [T_2(b, g, c) - T_1(b, g, c)] \\ V_o(b, g, c) &= \int_{T_1(b, g, c)}^{T_2(b, g, c)} [S_n(T) + S_p(T)] dT \\ K(b, g, c) &= [\bar{k}_n(b, g, c) A_n + \bar{k}_p(b, g, c) A_p] / L \\ R(b, g, c) &= [\bar{\rho}_n(b, g, c) / A_n + \bar{\rho}_p(b, g, c) / A_p] L\end{aligned}$$

Using these values of  $V_o$ ,  $K$ , and  $R$  for each couple and the prescribed RTG voltage  $V_{RTG}$ , the code computes the branch current  $I(b)$  for each of the  $B$  branches,

$$I(b) = \left[ \sum_{g=1}^G \left\{ \sum_{c=1}^C [V_o(b, g, c) / R(b, g, c)] \right\} / \sum_{c=1}^C [1/R(b, g, c)] \right] - V_{RTG} / \sum_{g=1}^G \left\{ 1 / \sum_{c=1}^C [1/R(b, g, c)] \right\},$$

the group voltage  $V(b, g)$  for each of the  $G$  groups in each branch,

$$V(b, g) = \left\{ \sum_{c=1}^C [V_o(b, g, c) / R(b, g, c)] - I(b) \right\} / \sum_{c=1}^C [1/R(b, g, c)],$$

and the couple current  $I(b, g, c)$  for each couple in the RTG,

$$I(b, g, c) = [V_o(b, g, c) - V(b, g)] / R(b, g, c).$$

The individual couple currents are then used to compute the hot-end heat input rate  $Q_H(b, g, c)$  and the cold-end heat rejection rate  $Q_C(b, g, c)$  for each couple:

$$\begin{aligned}Q_H(b, g, c) &= K(b, g, c) [T_2(b, g, c) - T_1(b, g, c)] - I^2(b, g, c) R(b, g, c) / 2 \\ &\quad + I(b, g, c) [S_2(b, g, c) T_2(b, g, c) + S_1(b, g, c) T_1(b, g, c) + V_o(b, g, c)] / 2, \\ Q_C(b, g, c) &= K(b, g, c) [T_2(b, g, c) - T_1(b, g, c)] + I^2(b, g, c) R(b, g, c) / 2 \\ &\quad + I(b, g, c) [S_2(b, g, c) T_2(b, g, c) + S_1(b, g, c) T_1(b, g, c) - V_o(b, g, c)] / 2,\end{aligned}$$

where  $S_1$  and  $S_2$  denote the Seebeck coefficients at the cold- and hot-junction of the couple.

The code inserts these heat flow rates for each couple in the RTG into the detailed thermal analysis model for the next iteration; and repeats the procedure until convergence is achieved.

## CODE VALIDATION

The code was first tested in the analysis of an unobstructed generator, with axial temperature variation but no circumferential variation. It was validated by using it to analyze the performance of the electrically heated thermoelectric generator (ETG) that had been employed as the engineering test unit [10], a prototype of the GPHS-RTGs used for the Galileo mission. The reason for using the ETG instead of the RTG as a validation check for the code is that the ETG test measurements include the thermocouples' hot-shoe temperatures, but - because of practical difficulties - the RTG measurements do not.

The analysis of the ETG's performance was based on thermoelectric properties of SiGe aged for one year to account for pre-test outgassing and processing. The temperature-dependent values of resistivity, conductivity, and Seebeck coefficient of the n- and p-material used in the analysis are summarized in Table 1.

Table 1. Thermoelectric Properties Model

Temp °C	Seebeck μV/K		Resistivity mΩ . cm		Conductivity mW/cm . K	
	N	P	N	P	N	P
0	90	114	0.79	0.91	51.4	58.4
50	113	128	0.87	0.95	50.2	57.2
100	135	140	0.95	1.01	49.0	56.3
150	153	152	1.04	1.08	47.9	55.0
200	168	163	1.13	1.15	46.8	54.1
250	185	173	1.28	1.24	45.8	52.9
300	215	183	1.72	1.32	44.7	51.7
350	265	192	2.77	1.42	43.9	51.0
400	304	202	3.81	1.53	43.1	49.9
450	317	211	4.17	1.64	42.5	49.1
500	317	220	4.02	1.74	41.8	48.2
550	312	229	3.75	1.87	41.5	47.5
600	306	240	3.46	2.06	41.0	46.8
650	298	256	3.13	2.38	40.8	46.2
700	289	287	2.82	2.98	40.6	45.7
750	280	320	2.53	4.49	40.6	45.3
800	273	331	2.28	4.76	40.7	45.2
850	268	330	2.09	4.54	40.9	45.4
900	264	323	1.95	4.17	41.4	46.0
950	263	318	1.86	3.84	42.1	47.1
1000	263	316	1.82	3.57	43.4	49.2
1050	262	314	1.77	3.22	45.6	52.2
1075	262	314	1.75	3.10	46.9	54.3

The analytical results were compared with the ETG test measurements. For the same RTG thermal power and the same average cold-junction temperature, the experimental measurements and the analytical results produced by the code were in very good agreement. The average hot-junction temperatures agreed within 5°C (986 versus 981), and the electrical power outputs agreed within 1 watt (296 versus 297).

As a further check on the validity of the analytical model, it was used to compute the BOM performance of one of the Galileo flight RTGs (F4). For the same thermal power, it yielded an electrical output of 285.6 watts, compared to the measured output power of 287.7 watts in Earth orbit. This agreement is better than the estimated 1.2% telemetry error.

The good agreement between the analytical and experimental results for both the ETG and the RTG lends confidence to our use of the same model and assumptions for subsequent analyses of the initial RTG output at various thermal powers and external environments.

## ANALYSIS OF OBSTRUCTED CRAFT/CASSINI RTGs

The application of the code to the analysis of the obstructed CRAFT/Cassini RTGs started with the construction of a 1912-node radiation-interchange analysis model. The model represented the housing and fins of the GPHS-RTG pictured in Figure 2 and the spacecraft shown in Figure 1. The ITAS (Integrated Thermal Analysis System) code [11], which accounts for the effect of mutual reflections, was used to compute over 102,000 radiation interchange factors between all surface nodes that are within each other's view. The computed radiation interchange factors were then inserted into a detailed thermal and electrical analysis model consisting of ~2900 node points.

The coupled thermal and electrical analysis was carried out by means of the previously discussed computer code. The code was based on the SINDA thermal analysis program [12], modified by us to incorporate the current-dependent Peltier, Ohmic, and Thomson effects on thermocouple conductance and to represent the electrical power generation in each couple as an effective heat sink. In each iteration, the modified code computed each couple's heat input rate and heat rejection rate and inserted them into the thermal analysis for the next iteration. After the solution had homed in to prescribed convergence criteria, it was used to calculate the RTG's electrical output and efficiency.

To illustrate typical results, the converged BOM solution for the baseline configuration and a thermal power of 245 watts from each of the 18 heat source modules is summarized in Tables 2 through 5. The tables display the results for the flattened-out cylindrical array of 576 thermocouples in the RTG.



As can be seen, the network's two branches have respective output currents 5.14 and 5.16 amp, for a total output of 10.30 amps per RTG. The 144 series-connected couple groups in each branch produce 30.0 volts. Subtracting 0.34 volts for ohmic losses in the RTG's series leads leaves a net output of 29.66 volts and 309 watts per RTG. The average material efficiency of the couples is 7.90%; the average couple efficiency (including the effect of contact resistances and electrode losses) is 7.55%; and the net system efficiency (including the effect of heat losses through the thermal insulation and through the heat source support structure) is 7.01%.

Equally detailed coupled thermal and electrical analyses were performed for thermal power levels ranging from 225 to 250 watts per heat source module, for both obstructed and unobstructed RTGs. The unobstructed units had axial temperature variations but no circumferential variation. The principal results for the mutually obstructed RTGs in the CRAF/Cassini baseline configuration (Figure 1) and the comparative results for an unobstructed RTG are displayed in Figures 7, 8, and 9.

For a separation angle of 24 degrees between the two RTGs, Figures 7 and 8 respectively show the effect of the thermal power  $Q$  per heat source module on the RTG power output and on the average hot-junction temperature. As can be seen, both relationships are essentially linear, and the figures present least-square curve fits for the respective curves. These curve fits are useful design tools, and will be used in predicting the effect of fuel decay on RTG power and temperature.

Figure 7. Effect of Thermal Power on BOM Electrical Power

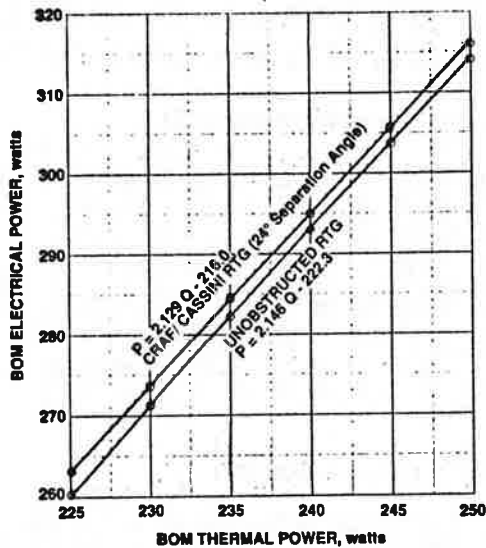
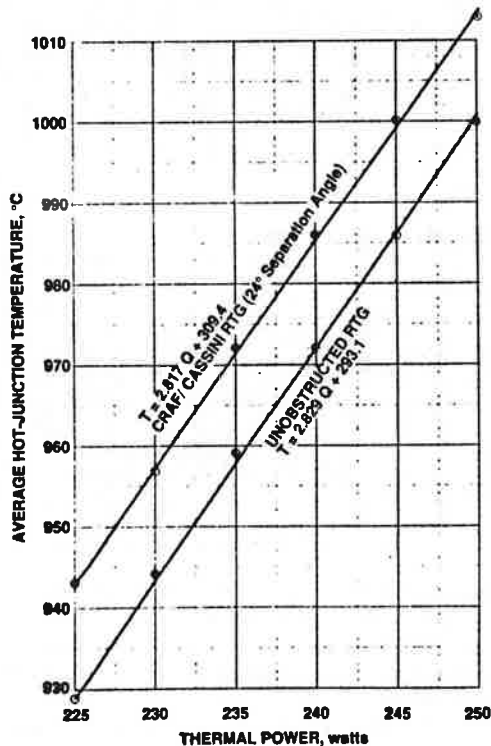
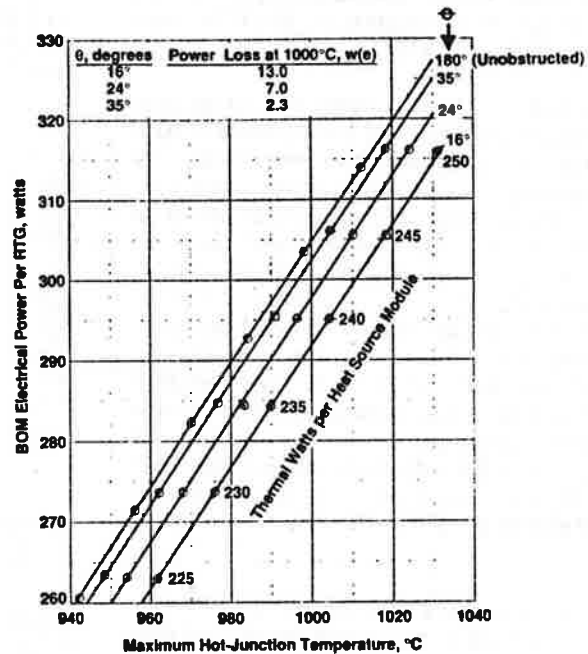


Figure 8. Effect of Thermal Power on Hot-Junction Temperature



The two figures show that, for a given thermal power, the obstruction by the neighboring RTG has very little effect on the RTG's power output, but has a significant effect on its hot-junction temperature, which affects the degradation rate. Figure 9 presents cross-plots showing the relationship between maximum hot-junction temperatures and BOM power outputs for unobstructed RTGs and for mutually obstructed RTGs with separation angles of 16°, 24°, and 35°.

Figure 9. Effect of Separation Angle  $\theta$  Between RTGs on Power-Versus-Temperature Characteristics



The performance of the obstructed and unobstructed RTGs should be compared for the same maximum hot-junction temperature, since that is what determines the RTG's degradation rate and lifetime. Previous SiGe flight units (LES 8/9, Voyager, Galileo) were designed for a maximum hot-junction temperature of 1000°C. For that temperature, Figure 9 shows that the blockage by its neighbor reduces the RTG's power output by 13.0 watts (4.3%) for a 16-degree separation angle, by 7.0 watts (2.5%) for the 24-degree separation of JPL's original baseline design and by 2.3 watts (0.8%) for the 35-degree separation of JPL's revised spacecraft design. This demonstrates the sensitivity of power output to separation angle.

The 7-watt performance penalty for the 24-degree separation angle may not seem large. But both missions are severely power-constrained, and the JPL designers were unwilling to give up 14 watts from the two RTGs. Subsequently, the spacecraft designs were changed to a 35-degree separation angle, which results in much less power loss.

## LONG-TERM PERFORMANCE DEGRADATION

The system designer is interested not only in the RTGs' BOM power, but also in its power history throughout the mission. The principal factors that diminish an RTG's output power  $P$  with time  $t$  are fuel decay and thermoelectric degradation. The combined effect of these factors can be represented by

$$\frac{P}{P_1} = \left( \frac{P_u}{P_1} \right) \left( \frac{P}{P_u} \right), \quad (15)$$

where  $P_1$  is the RTG's BOM power, the ratio  $P_u/P_1$  is the effect of fuel decay alone (i.e., just the effect of thermal power decrease, if there were no degradation), and the ratio  $P/P_u$  is the effect of thermoelectric degradation for the RTG's predicted temperature history.

As shown in Figure 7, in the range of practical interest the RTG's undegraded output power  $P_u$  is a linear function of its thermal power  $Q$ :

$$P_u = P^* + \eta^* Q. \quad (16)$$

The coefficients  $P^*$  and  $\eta^*$  must be determined by the previously described detailed thermal and electrical analyses, taking account of the location, orientation, and mutual blockage of the RTGs on the spacecraft. For example, for two standard GPHS-RTGs separated by an angle of 35-degrees (JPL's revised baseline design), least-square analysis of the analytic results yields the curve-fit coefficients

$$P^* = -218.0 \text{ w}, \text{ and } \eta^* = 0.1188.$$

Note that the coefficient  $\eta^*$  is a differential efficiency, which is the sum of the actual efficiency  $\eta$  and a term that reflects the fact that increasing  $Q$  increases the temperature drop  $\Delta T$  between hot- and cold-junctions, which in turn increases the efficiency

$$\eta^* = \frac{dP}{dQ} = \eta + \left[ \frac{\partial(\Delta T)}{\partial Q} \right] \left[ Q \frac{\partial \eta}{\partial(\Delta T)} \right] \quad (16a)$$

The effect of fuel decay on the thermal power  $Q$  at mission time  $t$  is given by

$$Q = Q_1 \exp(-\lambda t), \quad (17)$$

where  $Q_1$  is the BOM value of  $Q$ , and  $\lambda$  is the Pu-238 decay constant, which is related to its half-life  $\tau$  (87.74 years) by

$$\lambda = (\ln 2) / \tau. \quad (18)$$

Thus, the effect of fuel decay alone is given by

$$\frac{P_u}{P_1} = 1 - \frac{1 - e^{-\lambda t}}{(P^*/\eta^* Q_1) + 1} \quad (19)$$

Predicting the thermoelectric performance degradation with time is generally a complex problem, since many different degradation mechanisms are at play. A detailed method for making such predictions is the DEGRA code developed by V. Raag (1973), but

that code does not account for the detailed effects of mutual obstruction between neighboring RTGs. In the same period, A. Mowery (1975) performed statistical analyses of the results of various tests on thermoelectric converters that were similar to the converters in the GPHS-RTGs, but which were electrically heated and which operated at constant thermal powers and constant temperatures. These statistical analyses showed that the measured power output of a constant-temperature converter is a linear function of the square root of time  $t$ , and that the coefficient of  $\sqrt{t}$  is an Arrhenius function of the absolute hot-junction temperature  $T$ . In the range of practical interest, Mowery found that the measured effect of thermoelectric degradation could be quite accurately correlated by the simple empirical equation

$$P / P_u = 1 - \{ \exp [-T^* / T] (t / t^*) \}^{1/2}, \quad (20)$$

$$\text{where } T^* = 30,960^\circ K \text{ and } t^* = 0.648 \text{ sec}$$

This empirical expression was found to give excellent agreement with the experimental data from the electrically heated tests at various constant thermal powers and temperatures. However, Mowery's equations cannot be directly applied to predict the performance degradation of RTGs, because an RTG operates at a diminishing thermal power and temperature, and therefore at diminishing degradation rates. To account for that effect, the present author proposed the following simple modification of Mowery's original equation:

$$P / P_u = 1 - \left\{ \int_0^t \exp [-T^* / T(t')] dt' / t^* \right\}^{1/2} \quad (21)$$

Note that for the special case of constant  $T$ , the modified equation reduces to Mowery's validated Eq. 20. For variable  $T$ , the validity of Eq. 21 was subsequently confirmed by successfully predicting the results of long-term tests at diminishing temperatures.

As shown in Figure 8, the RTG's average hot-junction temperature  $T$  is a linear function of the RTG's thermal power  $Q$ :

$$T = T_0 + \alpha Q. \quad (22)$$

where the coefficients  $T_0$  and  $\alpha$  are obtained from detailed thermal and electrical analysis of the RTGs for various values of  $Q$ . For example, least-square analysis of the results for two mutually obstructing RTGs separated by an angle of 35 degrees yields the curve-fit coefficients

$$T_0 = 582.6^\circ K \text{ and } \alpha = 0.1565^\circ K / \text{w}.$$

Inserting Eqs. 22 and 17 into 20, we obtain

$$\frac{P}{P_u} = 1 - \left\{ \int_0^t \exp \left[ - \frac{T^*}{T_0 + \alpha Q_1 \exp(-\lambda t')} \right] dt' / t^* \right\}^{1/2} \quad (23)$$

Defining a dimensionless variable

$$x \equiv T^* / [T_o + \alpha Q_1 \exp(-\lambda t)], \quad (24)$$

Equation 23 reduces to

$$\frac{P}{P_u} = 1 - \left\{ \frac{1}{\lambda t^*} \int_{x_1}^{x_2} \frac{\exp(-x) dx}{x[1 - (T_o/T^*)x]} \right\}^{1/2}, \quad (25)$$

$$\text{where } x_1 \equiv T^* / (T_o + \alpha Q_1) \quad (25a)$$

$$\text{and } x_2 \equiv T^* / (T_o + \alpha Q_1 e^{-\lambda t}). \quad (25b)$$

Equation 25 can be partitioned into

$$\frac{P}{P_u} = 1 - \left\{ \frac{1}{(\lambda t^*)} \left[ \int_{x_1}^{x_2} \frac{e^{-x} dx}{x} + \int_{x_1}^{x_2} \frac{e^{-x} dx}{(T^*/T_o) - x} \right] \right\}^{1/2},$$

for which exact solutions in terms of exponential integrals are readily obtained. However, when we consider the magnitude of the x-values of interest, an almost exact but much simpler solution is available. Let us consider a typical problem with a BOM thermal power  $Q_1$  of 4410 watts and a mission life  $t$  of 12 years. Thus, from Eqs. 25a and 25b,

$$\begin{aligned} x_1 &= 24.325, \\ x_2 &= 25.590. \end{aligned} \quad (26)$$

For that range of x-values, the denominator in the integrand of Eq. 25 varies only from 13.18 to 13.26. Therefore, using a constant average value for that denominator is accurate to within  $\pm 0.3\%$ , and Eq. 25 can be quite accurately approximated by

$$\frac{P}{P_u} \approx 1 - \left\{ \frac{1}{(\lambda t^*) \bar{x} [1 - (T_o/T^*) \bar{x}]} \int_{x_1}^{x_2} e^{-x} dx \right\}^{1/2} \quad (27)$$

$$\text{where } \bar{x} \equiv T^* / (T_o + \alpha Q_1 e^{-\lambda t/2}). \quad (27a)$$

Integration of Equation 27 gives

$$\frac{P}{P_u} = 1 - \left\{ \frac{\exp(-x_1) - \exp(-x_2)}{(\lambda t^*) \bar{x} [1 - (T_o/T^*) \bar{x}]} \right\}^{1/2}, \quad (28)$$

where  $x_1$ ,  $x_2$ , and  $\bar{x}$  are given by Eqs 25a, 25b, and 27a. The above solution accounts for diminishing temperatures and degradation rates with time. Finally, Eqs. 19 and 28 can be inserted into 15 to obtain the combined effect of fuel decay and thermoelectric degradation for a given BOM thermal power  $Q_1$  and mission time  $t$ . (29)

$$P = (P^* + \eta^* Q_1 e^{-\lambda t}) \left\{ 1 - \left[ \frac{\exp(-T^*/T_1) - \exp(-T^*/T)}{(\lambda t^*)(T^*/\bar{T})[1 - (T_o/\bar{T})]} \right]^{1/2} \right\}$$

$$\text{where } T_1 = T_o + \alpha Q_1 \quad (29a)$$

$$T = T_o + \alpha Q_1 e^{-\lambda t} \quad (29b)$$

$$\bar{T} = T_o + \alpha Q_1 e^{-\lambda t/2} \quad (29c)$$

## VALIDATION OF DEGRADATION MODEL

The above equations were applied to the Q-1 RTG, which was the qualification unit for the GPHS-RTGs used on the Galileo mission. The computed results are shown in Figure 10. The curve labeled  $Q/Q_1$  shows the thermal power decrease due to isotope decay (Eq. 17), and the curve labeled  $T$  shows the resultant decrease in average hot-junction temperature (Eq. 29b). The curve labeled  $P_u/P_1$  shows the power loss due to fuel decay for undegraded thermoelectric performance (Eq. 19), and the curve labeled  $P/P_u$  shows the power loss due to thermoelectric degradation (Eq. 28). The combined effect of fuel decay and thermoelectric degradation is shown by the curve labeled  $P/P_1$  (Eq. 15). As can be seen, the combined effect for that case is predicted to result in a 23.6% power loss in 12 years.

Figure 10. Effect of Fuel Decay and Thermoelectric Degradation on Performance of Galileo Qualification RTG (Q-1)

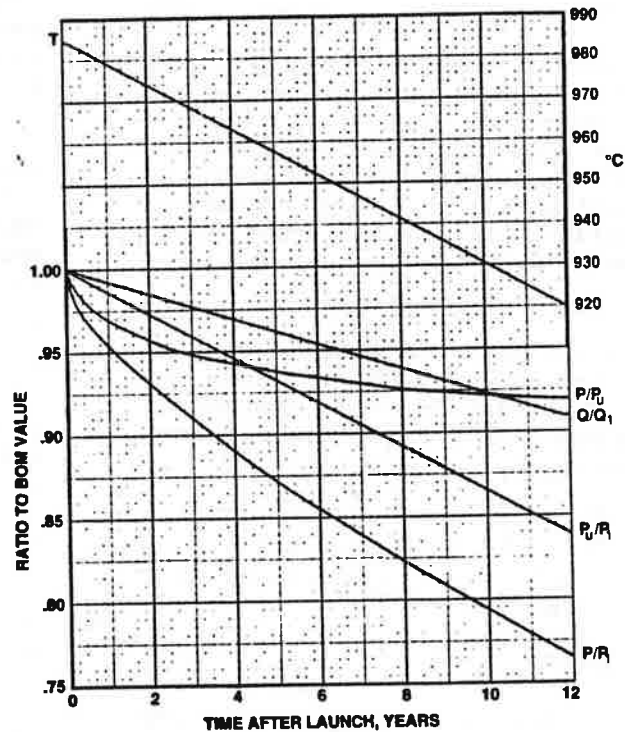
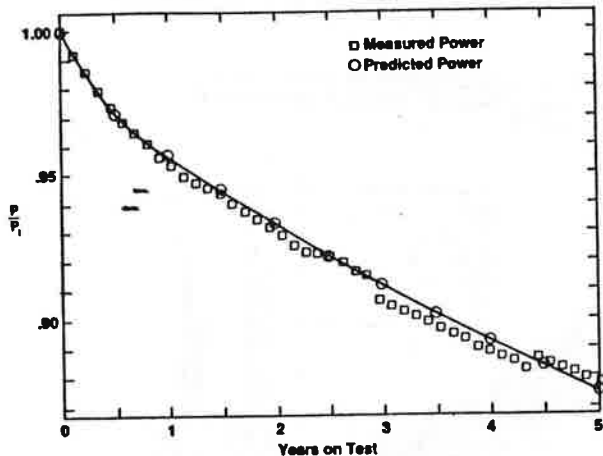


Figure 11 shows a comparison of these predictions with the measured results of a five-year test of the Q-1 RTG at diminishing temperatures. As can be seen, for the range of times and temperatures tested, the rather simple prediction model used showed surprisingly good agreement with the experimental results.

Figure 11. Comparison of Predicted and Measured Power History of Q-1 RTG



### PARAMETRIC PERFORMANCE PREDICTIONS FOR CRAFT AND CASSINI RTGS

The analytical methodology described above was used to predict the long-term power degradation of the two CRAFT and Cassini RTGs integrated with their respective spacecraft. (See Figure 1.) Parametric analytical results for two RTGs with a 35 degree separation angle are shown in Figure 12. For various fuel loadings, the solid curves show the variation of RTG power with time. For each solid curve, the figure shows the thermal power (in each of the 18 heat source modules) and the corresponding plutonium enrichment (for the same fuel density as that used in the Galileo RTGs). The numbers listed are for two points in time: at the chemical separation of the plutonium from the irradiated neptunium-237 targets (BOL), and three years later at launch (BOM). Thus, the figure can be employed as a useful design tool by the CRAFT and Cassini mission planners.

For each solid curve in Figure 12, its intersections with the dashed curves indicate the decrease of the RTG's average hot-junction temperature with time for that fuel loading. (Each constant-temperature curve was obtained by inserting  $T$  into Eq. 29b, solving for  $Q_1$  at each mission time, and using that value of  $Q_1$  in Eqs. 29, 29a, and 29c to compute the corresponding output power  $P$ .)

The bold solid curve in the figure represents the fuel loading in the Galileo RTGs, and the bold dashed curve represents the 1000°C hot-junction limit adopted in previous missions. Note that this limit is exceeded by the BOM temperatures at higher fuel loadings. But modest excesses may be allowed, since the results displayed in Figure 12 already include proper allowances for increased degradation at the higher temperatures, based on results of (accelerated) tests conducted above 1000°C.

Figure 12. Effect of Fuel Loading on RTGs' Power and Temperature Histories (35° Separation, 18 Heat Source Modules, 36 Couple Rings, 3" Fins, 30 Volts/RTG)

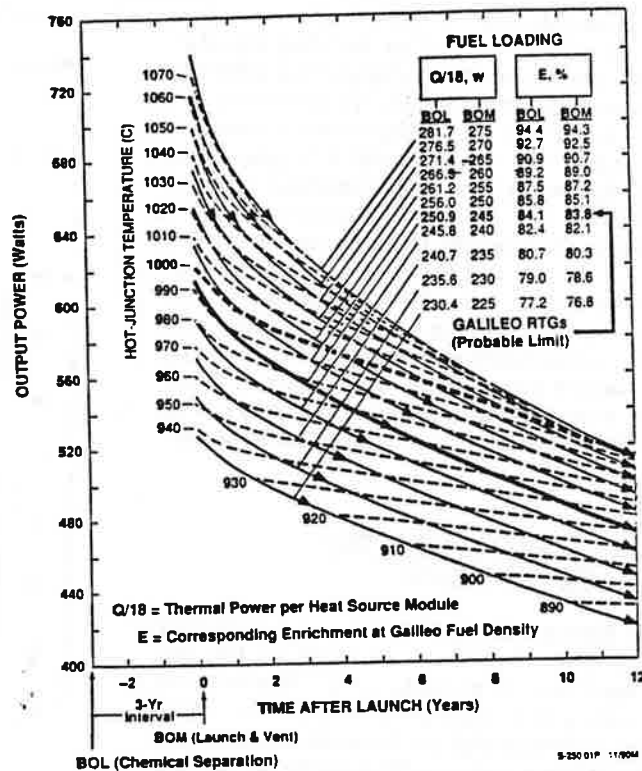


Figure 12 shows that increasing the BOM thermal power (by increasing the fuel loading) leads to a substantial increase of the RTG's electrical power at the beginning of mission, but that this benefit diminishes towards the end of mission, particularly for high fuel loadings and long mission times. This comes about because higher thermal powers result in higher temperatures and therefore higher degradation rates. In fact, at unrealistically high hot-junction temperatures increasing the thermal power can actually decrease the EOM electrical power.

### APPLICATION TO TWO UNEQUAL RTGS

JPL's baseline CRAFT and Cassini mission plans called for the use of two GPHS RTGs on each spacecraft. Under the corresponding DOE plan, two freshly fueled RTGs (F6 and F7) were to be used for the Cassini mission, which has the higher power demand. In that case, Figure 12 could be applied directly to predict the RTGs' power output. However, the CRAFT mission under that plan was slated to use only one freshly fueled RTG (F2) and one aged spare unit (F5) left over from the Galileo program. The fuel of that unit was encapsulated in 1982, and will have decayed for 13 years by the time of the CRAFT spacecraft's planned 1995 launch. As a result, by the beginning of the mission its thermal power will have dropped from 252 watts to 227 watts per heat source module.



During almost all of the 13-year storage interval, the old RTG will have been filled with an argon cover gas, to spoil its multifoil thermal insulation, lower its hot-junction temperature, and virtually eliminate its thermoelectric degradation. Thus, the spacecraft will be launched with two RTGs having similar thermoelectric properties but substantially different thermal powers and temperatures. To determine the effect of the two unequal RTGs on each other, a detailed BOM thermal and electrical analysis of the baseline-configured RTGs with respective thermal powers of 245 and 227 watts per heat source module was carried out. The results fell right on the curves shown in Figures 7 and 8 for obstructed but equal RTGs. In other words, each RTG is affected by the obstruction of its neighbor, but its resultant power loss is essentially independent of modest deviations in its neighbor's thermal power. Therefore, Figure 12 can also be used to predict the performance history of two unequal RTGs by using an appropriate BOM thermal power for each RTG.

### PROGRAMMATIC IMPLICATIONS FOR CRAF AND CASSINI MISSIONS

Prior to September 1990, only fragmentary information about the power demand for the two missions had been issued by JPL. Preliminary analyses using the above-described Fairchild methodology indicated that JPL's baseline plan would comfortably meet the then-specified power demand for the CRAF mission, and would come close to meeting it for the Cassini mission. The then-existing status is summarized in Table 6.

Table 6. Power Margins in CRAF/CASSINI RTGs (Aug. 1990)

MISSION	CRAF				CASSINI			
	Rendezvous		Perihelion		Arrival		EOM	
Years After Launch	5.01	7.61	6.50	10.30				
Mission Power Requirement (watts)	458	432	524	478				
RTG #	1	2	1	2	1	2	1	2
Age of Fuel at Launch (years)	3	13	3	13	3	3	3	3
Thermal Power per H.S. Module (watts)								
At Fuel Encapsulation	252	252	252	252	252	252	252	252
At Launch (BOM)	245	227	245	227	245	245	245	245
Hot-Junction Temperature (°C)								
At Launch	1000	952	1000	952	1000	1000	1000	1000
At Event	973	924	960	910	966	966	946	946
Electrical Power (watts)								
At Launch	305	267	305	267	305	305	305	305
At Event	265	237	253	226	258	258	241	241
Total at Event	502	479	516	482				
Power Margin (watts)	44	47	-8	4				

In late September 1990, as the result of ongoing mission studies, JPL issued much more detailed (and significantly higher) power demand goals to meet the missions' engineering and science requirements. Figures 13 and 14 display these power demand goals for the CRAF and Cassini missions in the form of histograms that have been superimposed on the power supply capability curves for two RTGs. As can be

seen, for achievable thermal powers (e.g., 245 watts BOM per module), the original JPL plan falls significantly short of meeting the revised power demand goals, particularly in the case of the Cassini mission.

Figure 13. Power Demand Versus Supply, for CRAF Mission (35° Separation, 18 Heat Source Modules, 36 Couple Rings, 3" Fins, 30 Volts/RTG)

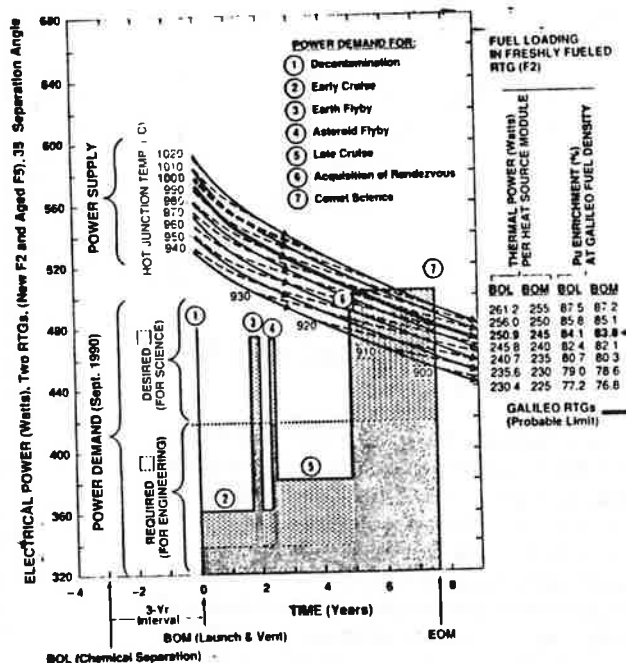
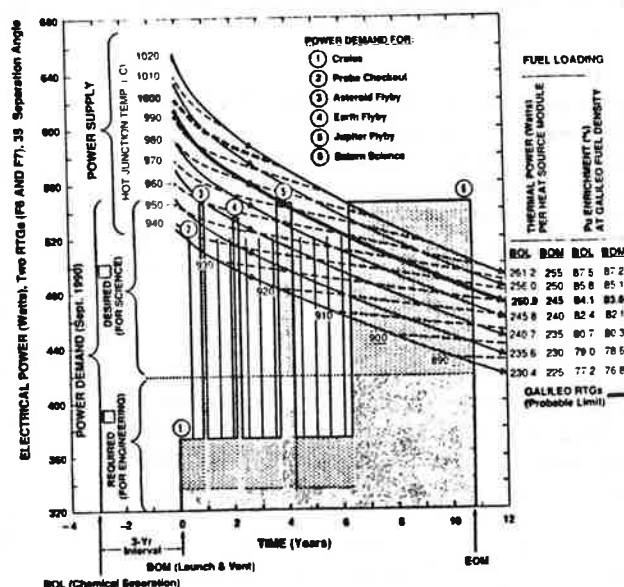


Figure 14. Power Demand Versus Supply, for CASSINI Mission (35° Separation, 18 Heat Source Modules, 36 Couple Rings, 3" Fins, 30 Volts/RTG)

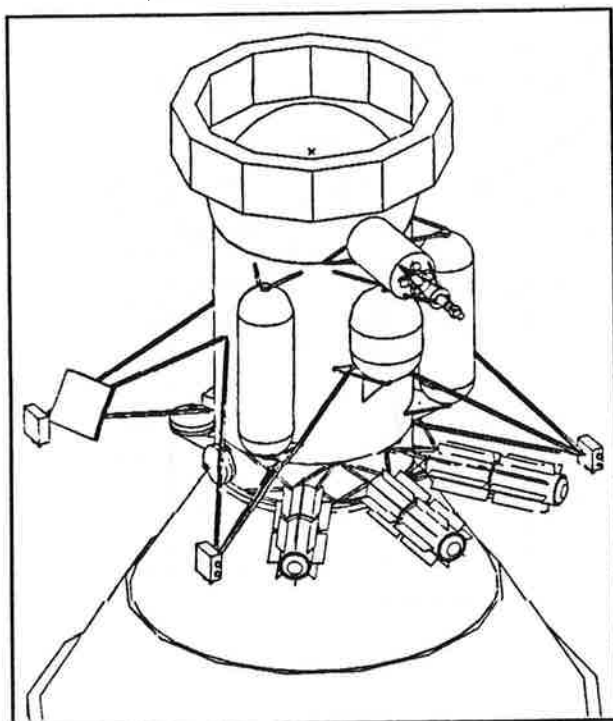


Subsequently, in late October, the power shortfall was further exacerbated when JPL specified additional increases in power demand and mission duration, for both missions. The mission duration for CRAF was increased from 7.6 to 9.6 years, and the duration of Cassini was increased from 10.7 to 12.0 years. These time increases were introduced to allow additional gravity-assist maneuvers, in order to provide substantially higher mass margins for the spacecraft.

Because the twice-revised power demand goals for both missions significantly exceed the capability of two standard RTGs, JPL is studying various modifications of their mission plans. To support those studies, Fairchild personnel is conducting studies to assess the programmatic impact of adding a third RTG to each mission and to assess the technical effort of possible modest modifications of the RTG design and/or operating procedures to raise the power output of two RTGs to the desired level. The Fairchild programmatic assessment was discussed in a report by Mr. Eck (FSC-ESD-217/90/478B, 1990). The technical assessment is the subject of the present report.

In the case of the Cassini mission, the latest power demand and mission duration are so high that there appears to be no chance of meeting them with two RTGs. Therefore, it was decided to focus on the use of three RTGs for the Cassini mission, two freshly fueled ones (F2 and F6) and one aged, lower-power unit (F5). The additional mass of the third RTG can be accommodated because of the increased mass margin made possible by the added gravity-assist maneuver.

Figure 15. Mounting of Three RTGs on CASSINI Spacecraft (34-Degree Separation Between Neighbors)



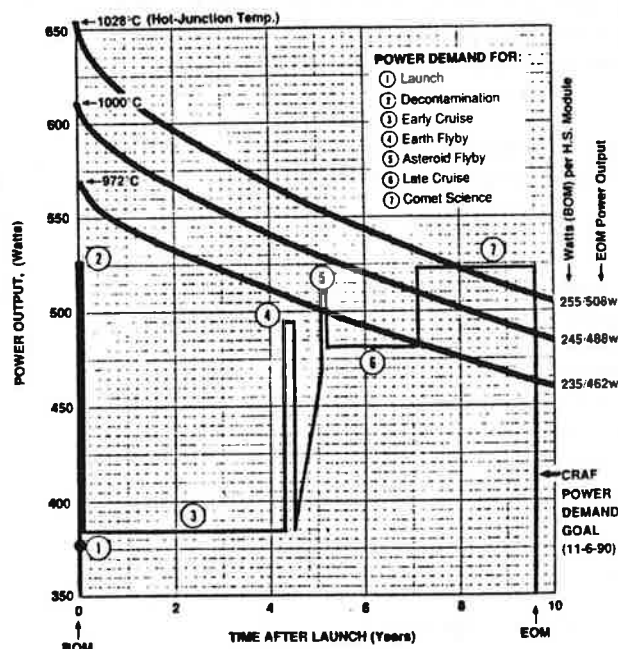
Present plans envisage placing the three Cassini RTGs near each other, with a separation angle of 34 degrees between neighboring RTGs. (See Figure 15). The middle one of the three RTGs is the most obstructed and would normally run at the highest temperature, but this problem can be alleviated by placing the aged RTG (F2) with its lower thermal power in the middle. Detailed thermal and electrical analyses of the three-RTG configuration have been conducted by Fairchild. The three RTGs can provide ample margin to meet the current Cassini power demand, even for the extended mission duration. The detailed results will be presented after discussing the results for the more complicated CRAF problem.

The latest power demands for the CRAF mission could of course be similarly met by the addition of a third RTG. But building an additional RTG (F9) for the CRAF mission would be costly and time-consuming. Thus both the manufacturing cost of a third RTG and the programmatic cost of the extra time required to build it would adversely impact the cost of the CRAF mission. Therefore, there is a great incentive to try to meet that mission's power demand with just two RTGs.

Figure 16 compares the predicted power output of two standard GPHS RTGs (for various fuel loadings) with JPL's latest power demand goals for the CRAF mission (furnished by R. Campbell on 11/6/90).

Figure 16. Revised CRAF Power Demand (11-6-90) Versus Output on Two Standard RTGs

(35° Separation, 18 Heat Source Modules, 36 Couple Rings, 3" Fins, 30 Volts/RTG)



The CRAF power demand profile is shown in the form of a histogram which covers the cruise phase ( $0 < t < 7.1$  yrs) and the comet science phase ( $7.1 < t < 9.42$  yrs). As can be seen, the two standard RTGs can meet cruise power goals, but with an achievable fuel loading (240 to 245 thermal watts BOM per heat source module) they fall short of meeting the JPL-specified 524-watt goal for the comet science phase. For 245-watt heat source modules, the shortfall is seen to be about 13 watts at the beginning and 35 watts at the end of that phase. Thus, unless JPL can reduce its power requirements by that amount (e.g., by reducing the system-level and subsystem-level contingency reserves built into their estimates), the two RTGs would have to be modified to increase their power output at EOM while still meeting the cruise power demand.

#### OPTIONS FOR INCREASING THE POWER AT ENCOUNTER

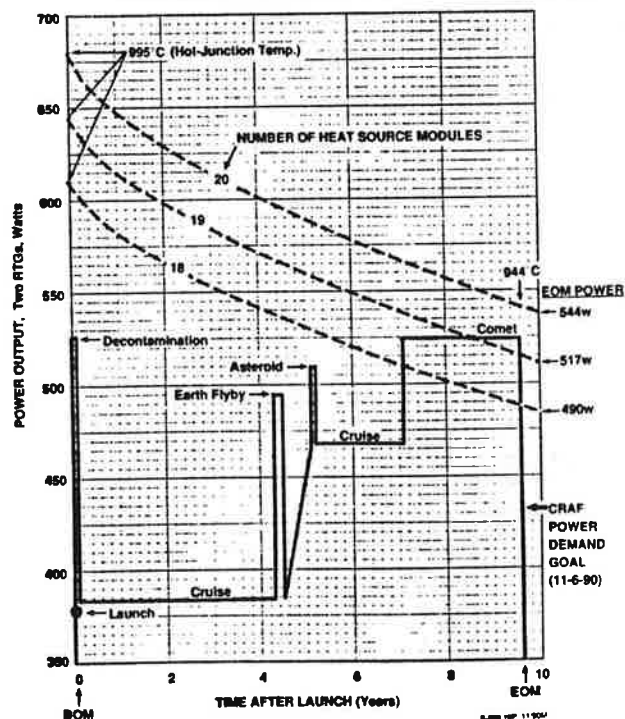
There are several design or operational modifications that could be used singly or in combination to increase the RTG's power output:

1. "Stretch" the standard RTG, to raise the number of heat source modules from 18 to 19 or 20.
2. Increase the fuel loading by increasing its density and/or its isotopic enrichment.
3. Delay venting the RTG's cover gas, to reduce the thermoelectric degradation during the coast phase.
4. Reduce the number of thermoelectric couples and/or the leg area in each couple (to increase the temperature drop between junctions and to raise the efficiency).
5. Increase the size of the radiator fins, to decrease the RTGs' cold-junction temperatures.
6. Provide external switching to supply the desired 30-volt output with the two RTGs in series during the power-rich cruise phase and in parallel during the power-lean encounter phases. (Operating each RTG at 15 instead of 30 volts increases the currents and the Peltier cooling rates, which lowers the hot-junction temperatures and the thermoelectric degradation rates.)
7. Use controlled-rate venting of the helium generated by alpha decay to spoil the thermal insulation and lower the hot-junction temperatures and power output just enough to match the power demand profile. This would minimize the thermoelectric degradation.

The above options, used either singly or in various combinations, were subjected to detailed analyses, to assess their effectiveness in closing the gap between power supply and demand, and to determine which - if any - warrant further investigation. The computed results are presented below. In each case, the results are compared with a histogram of the power demand goals furnished by JPL on 11/6/90. It is recognized that these goals are likely to change as the result of changed science demand and as the result of reducing the assumed multi-level contingency reserves. But it is useful to employ the same histogram as a common yardstick, to facilitate comparison of the various options.

**"Stretching" the RTG length by 2 or 4 inches to allow raising the number of heat source modules from 18 to 19 or 20 would be a straight-forward method of increasing the RTG's power output. As shown in Figure 17, even the 19-module RTG would virtually meet all of JPL's power demand goals. However, such a design change would constitute a significant deviation from the standard RTG design, and would therefore require some new qualification tests. (Note that the curves for 19 and 20 heat source modules were obtained by simple scaling, which implicitly assumes the same efficiency as the 18-module case. This is conservative, since efficiency generally increases with power output, because end losses consume a smaller fraction of the thermal power in larger RTGs.)**

Figure 17. Effect of "Stretching" the Standard RTG (35° Separation, 245 Watts BOM per Heat Source Module, 36 Couple Rings, 3" Fins, 30 Volts/RTG)



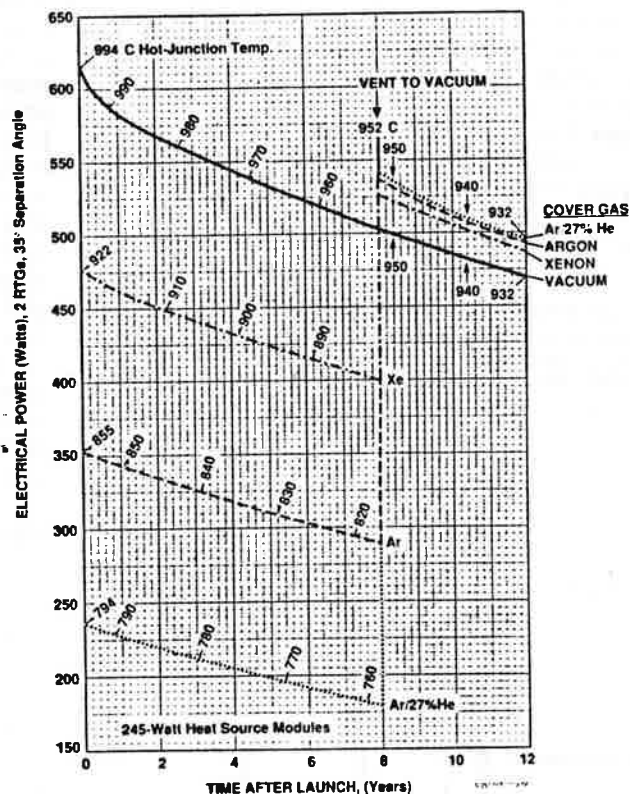


Moreover, increasing the fuel enrichment may not be possible without new fuel production, which may not be available in time in the required quantities. Even if higher-enrichment fuel did become available, beyond a certain enrichment the resultant higher fuel temperature could also lead to problems in pellet formation requiring costly and time-consuming new fabrication development and safety tests. In any case, until new fuel is available we are limited to the present inventory. A recent study (Eck 1990) indicates that the present inventory, by proper selection, would yield an isotopic enrichment of 82.9% at the time of chemical separation of the plutonium. This corresponds to a BOL thermal power of 247.3 watts per heat source module, which would decay to a BOM power of 241 watts at launch, 3.25 years later.

The third stratagem for raising the encounter power, delayed venting of the RTG's cover gas, could be used to spoil the converter's multifoil thermal insulation until the mission's comet phase, when the maximum output power is needed. Spoiling the thermal insulation would significantly decrease the RTG's hot-junction temperatures and therefore its degradation rate during the mission's cruise phase, making more power available during the comet phase. Unfortunately, the same cover gas that increases the comet science power also diminishes the cruise power. Thus, delayed venting could be quite useful for missions in which the required ratio of cruise power to encounter power is low, but much less useful for a mission like CRAF where that ratio has quite a high value (481 w/521 w = 92%).

are the only cover gases for which RTG test data are available. (Additional tests to expand the data base are under way at Mound.) This figure assumes venting of the cover gas after eight years, and confirms that delayed venting has the anticipated effect of increasing the power at EOM. The eight-year venting delay would increase the power of the two RTGs at 12 years by about 30 watts for the argon-helium mixture, almost as much for pure argon, and about 21 watts for the least conductive gas, xenon.

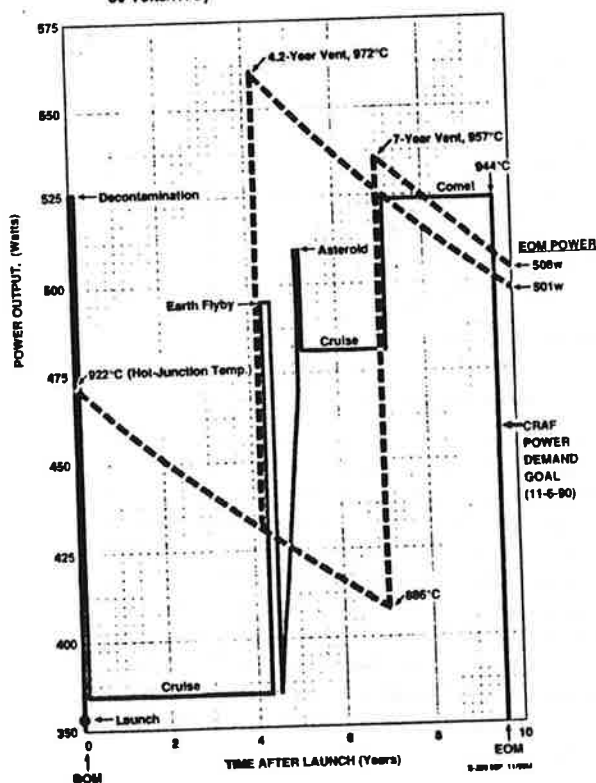
**Figure 18. Effect of Delayed Venting for Various Cover Gases**



These EOM power increases are respectable improvements, but unfortunately all of these cover gases lead to excessive losses during the mission's cruise phase. The least excessive power loss is obtained with xenon, the inert gas with the lowest thermal conductivity. However, to be effective this would require a means of preferentially discharging the alpha-generated helium while retaining the xenon, e.g. by use of a selective vent. Selective vents made of Viton were used in the Viking RTGs, but the adequacy of those vents for the present application has not been demonstrated.

The effect of delayed xenon venting is illustrated in Figure 19. Two output profiles are shown, one for venting at the beginning of the comet science phase (7.1 yrs), and one for venting just prior to the Earth flyby (4.2 yrs). As can be seen, the power demands for Earth flyby, asteroid flyby, and late cruise cannot be met if venting is delayed until comet rendezvous, but can be met if the RTG is vented just before Earth flyby. Neither option meets the power demand peak for decontamination. Both options meet the demand goal at the beginning of the comet science phase, but neither meets it at the end of that phase. The 7-year vent case misses the required EOM goal by 16 watts, and the 4.2-year vent case misses by 22 watts. Both of these power deficits are smaller than for the vent-at-launch case (35 watts), but they are still not good enough. Does that mean that the use of delayed venting is of no interest? Not necessarily. It may not be adequate when used by itself, but may be adequate in conjunction with one of the other stratagems (e.g., reduced number of couples).

Figure 19. Effect of Delayed Xenon Vent on Output of Two Standard RTGs  
(35° Separation, 18 Heat Source Modules, 36 Couple Rings, 3" Fins, 30 Volts/RTG)



The fourth stratagem involves reduction of the thermoelectric area, either by reducing the number of couples in the RTG, or by reducing the leg area in each couple. The former option is easier to implement, because it avoids the need of developing a new couple. In either case, reducing the thermoelectric

cross-section of the RTG increases the temperature drops between hot and cold junctions, which increases the conversion efficiency and therefore the RTG's power output for a given thermal power. The deleted couple rings were removed from the two ends of the RTG. This has the added benefit of improving the uniformity of the hot-junction temperatures, by compensating for the drop-off due to end losses (see Table 3).

Figure 20 compares the performance of the standard (vented-at-launch) 36-ring RTGs with those of the 34-ring and 32-ring designs; i.e., designs in which one or two rings have been deleted from each end of the RTG. As seen, these deletions lead to a substantial increase in the BOM output power. But, as shown, they do so only at the cost of substantially higher hot-junction temperatures. The consequent rise in degradation rates largely defeats the purpose of deleting the couple rings. As can be seen, for the 34-ring case increased degradation reduces the 27-watt gain in BOM power to only 4 watts by the end of mission; and for the 32-ring design the BOM power gain of 53 watts is reduced to a 13-watt power loss at EOM. All three cases fail to meet the EOM power demand goal. Thus, deletion of couple rings is not useful by itself, and would only be useful in conjunction with another method, to reduce the hot-junction temperature. Either delayed venting and/or longer fins may be used for that purpose.

Figure 20. Effect of Deleting Unicouple Rings from Ends of 36-Ring RTGs  
(35° Separation, 18 Heat Source Modules, 245-Watt BOM, 36 Couple Rings, 3" Fins, 30 Volts/RTG)

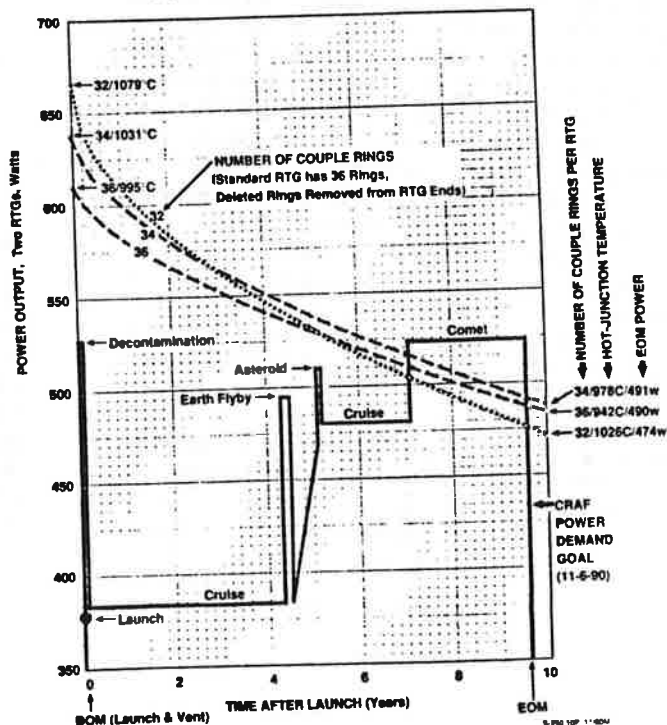


Figure 21 displays the effect of delayed venting on an RTG with 32 couple rings. Results are shown for venting at launch and at 4.2, 5.0, and 7.0 years after launch. As expected, the later the venting the lower the degradation and the higher the EOM power. For venting at 7.0 years, the EOM power is increased by 45 watts, which essentially meets the EOM power demand goal. But as shown, it does not meet the power goal for Earth and asteroid flyby or for late cruise. Venting at 5.0 years meets the asteroid and late-cruise goals, but not the Earth-flyby goal. The only strategy that meets the Earth, asteroid, and late-cruise goals requires venting at 4.2 years, but this yields an EOM power gain of only 28 watts, which is 22 watts short of the EOM power goal.

Figure 21. Effect of Delayed Xenon Vent on Output of Two 32-Ring RTGs  
(18 Heat Source Modules, 245-watt BOM, 3" Fins, 30 Volts/RTG)

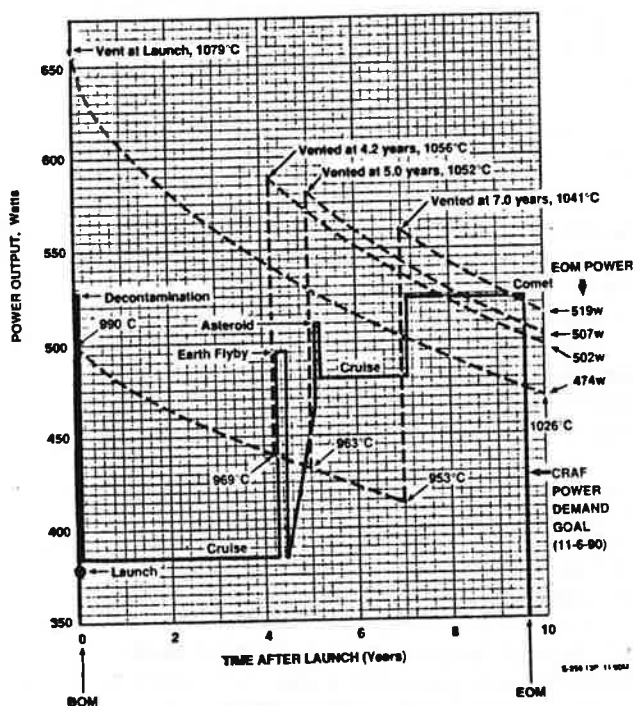
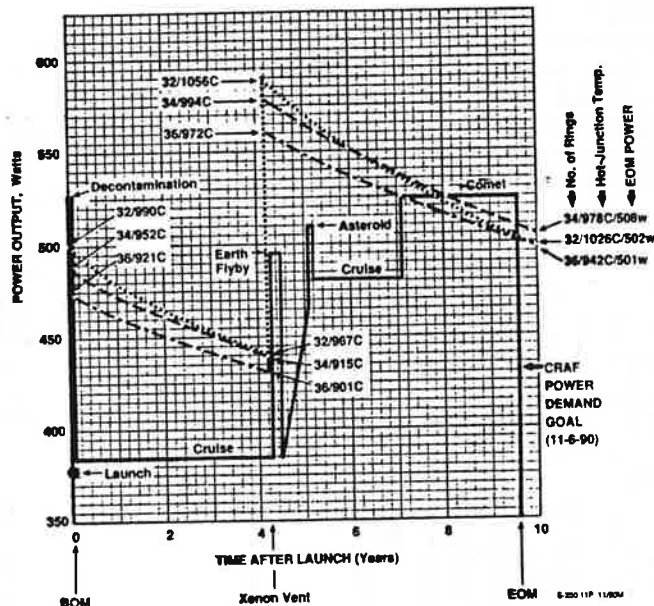


Figure 22 examines the effect of varying the number of couple rings on the output of RTGs vented at 4.2 years. As can be seen, the 34-ring design results in reasonable temperatures and yields the highest EOM power (508 watts, which is still 16 watts short of the EOM goal).

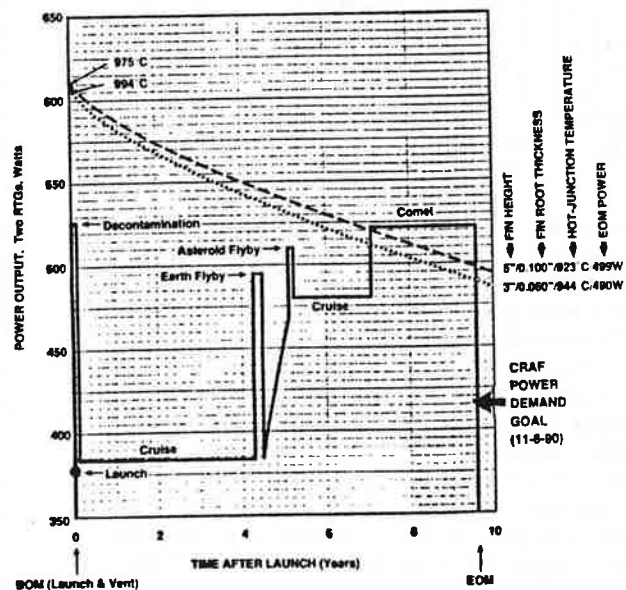
The effect of increased fin size on performance is depicted in Figure 23. The fins have a trapezoidal cross section with a 0.015" fin tip thickness. The standard design has a 0.060" fin root thickness with a 3" root-to-tip height. This is compared with a 5" fin having a 0.100" root thickness. The larger fins lower the RTGs's housing temperatures and cold-junction temperatures. But, as shown, they also lower the hot-

Figure 22. Effect of Number of Couple Rings on Output on Two RTGs Vented at 4.2 Years  
(18 Heat Source Modules, 245-Watt BOM, 3" Fins, 30 Volts/RTG)



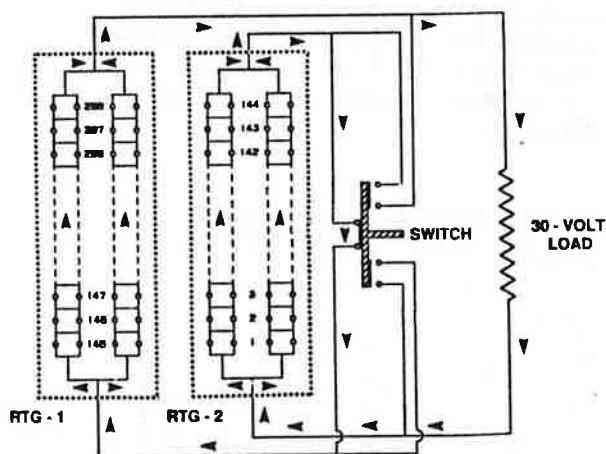
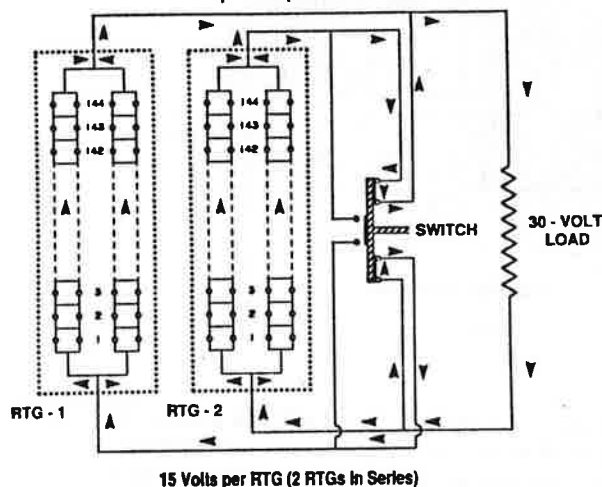
junction temperature by 19°C, so that there is virtually no change in  $\Delta T$  or in BOM power. However, the lower hot-junction temperature does reduce the degradation rate, resulting in a 9-watt increase in EOM power. Thus, while larger fins are helpful, their benefit is not nearly enough to meet the EOM power goal. But it may be a useful design change in conjunction with other modifications.

Figure 23. Effect of Fin Size on Output of Standard RTGs  
(35° Separation, 18 Heat Source Modules, 245-Watt BOM, 36 Rings, 30 Volts/RTG)



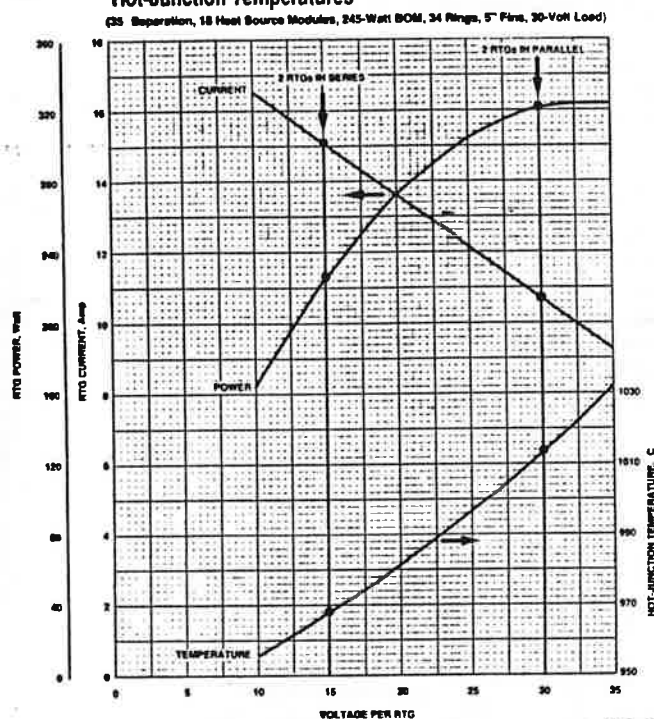
The next stratagem investigated is voltage switching. This is achieved by the addition of an external two-position switch which allows the two RTGs to be connected either in parallel or in series. The alternative circuit configurations are depicted in Figure 24. In either position, the RTGs are connected to a 30-volt load. As usual, the load voltage is regulated by a shunt dissipator. When the two RTGs are in parallel, each has an output of 30 volts. When they are in series, each has an output of 15 volts.

Figure 24. Voltage Switching Schematic  
30 Volts per RTG (2 RTGs in Parallel)



The resultant effects for a 34-ring RTG with 5" fins are shown by the curves of Figure 25, which were obtained by detailed analyses of the obstructed RTG at various voltages. The curve labeled "current" depicts the current-voltage characteristic of each RTG, and the other two curves show the corresponding variations of the output power and the hot-junction temperature. As can be seen, an output of 30 volts is close to the maximum-efficiency point (35 volts) while an output of 15 volts is way off-optimum. Lowering the RTG voltage to 15 volt drops the BOM power from 320 watts to 226 watts. But that would be acceptable during mission phases when the power demand is low.

Figure 25. Effect of RTG Voltage on BOM Power, Current, and Hot-Junction Temperatures



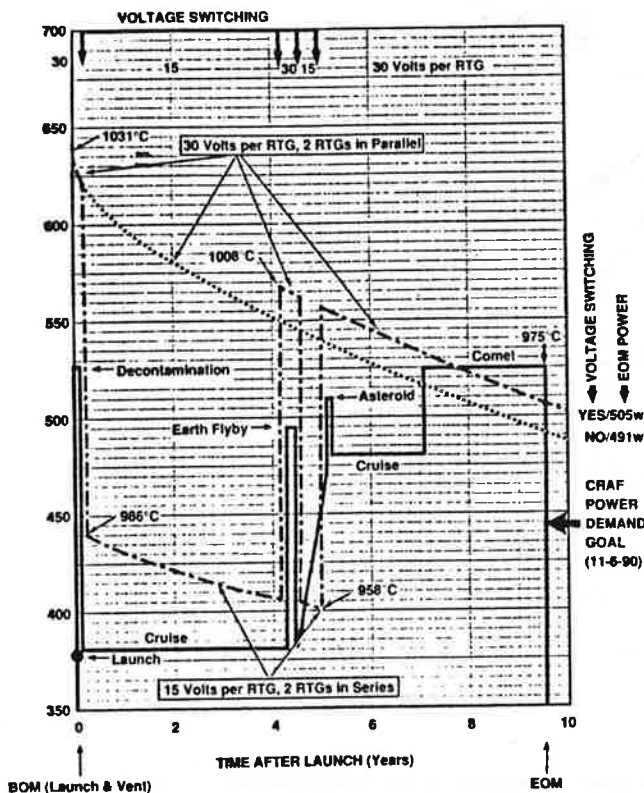
As shown in Figure 25, lowering the RTG's voltage from 30 to 15 volts raises its current from 10.7 to 15.1 amps. This increases the Peltier cooling rate, and - as shown - lowers the hot-junction temperature from 1014°C to 968°C. The 46°C reduction in hot-junction temperature would greatly decrease the thermoelectric degradation rates during periods when the lower power output is sufficient to meet the demand.

This is illustrated in Figure 26, which shows the effect of voltage switching on the output of a 34-ring RTG with standard 3" fins, vented at launch. Voltage switching has a similar effect as delayed venting. But one of its advantages is that - unlike delayed venting - the process is reversible, and can be invoked repeatedly to meet temporary power demand peaks. Another of its advantages is that it does not require a selective vent for preferential release of the helium.

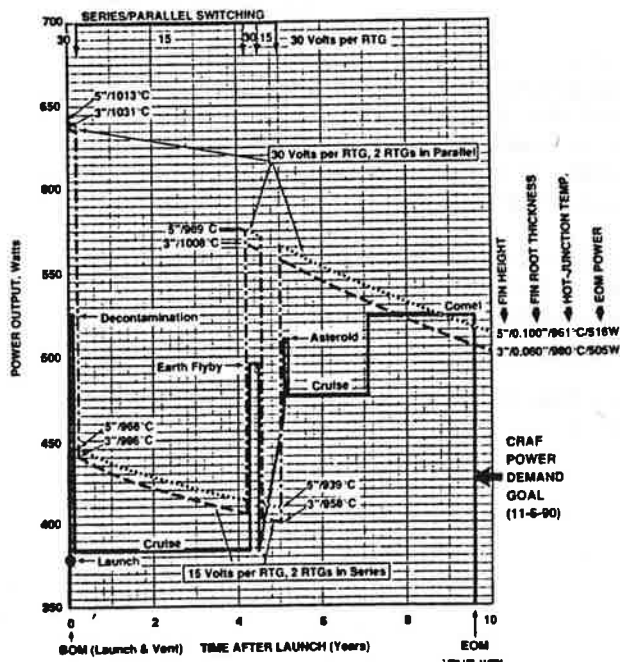
In the case illustrated by the dash-dot curve of Figure 26, the two RTGs operate in series (at 15 volts each) during the brief decontamination cycle at the beginning of the mission, during the Earth flyby, and during the asteroid flyby and the remainder of the mission. During the early cruise phases the two RTGs are in parallel, at 30 volts each, to reduce the thermoelectric degradation. As can be seen, this strategy meets all of the power demand goals except at the latter part of the comet science phase. Although the 505-watt EOM output for this case is seen to be 14 watts higher than what would be obtained without voltage switching, it is still 19 watts below the EOM power goal.



**Figure 26. Effect of Voltage Switching on Output of 34-Ring RTGs, Vented at Launch**  
(35° Separation, 18 Heat Source Modules, 245-Watt BOM, 3" Fins, 15/30 Volts/RTG)

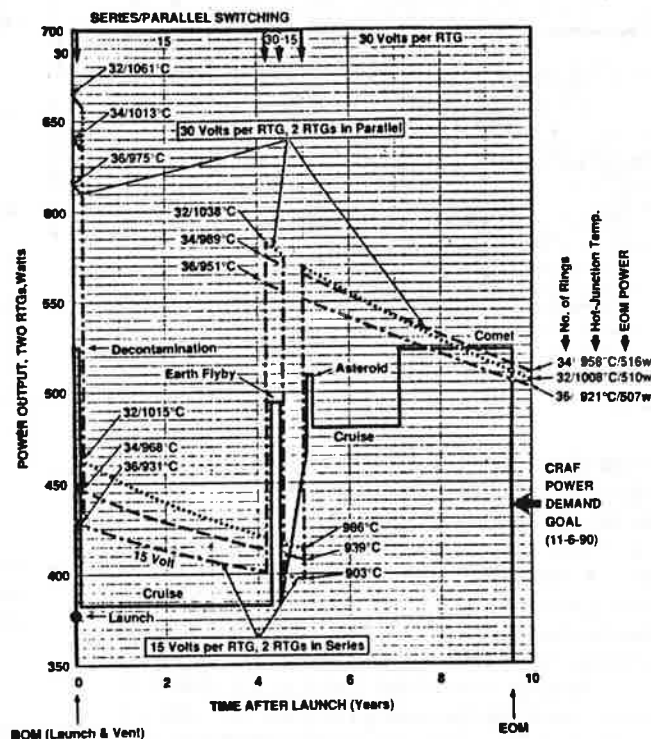


**Figure 27. Effect of Fin Size on Output of Dual-Voltage RTGs**  
(35° Separation, 18 Heat Source Modules, 245-Watt BOM, 34 Couple Rings)



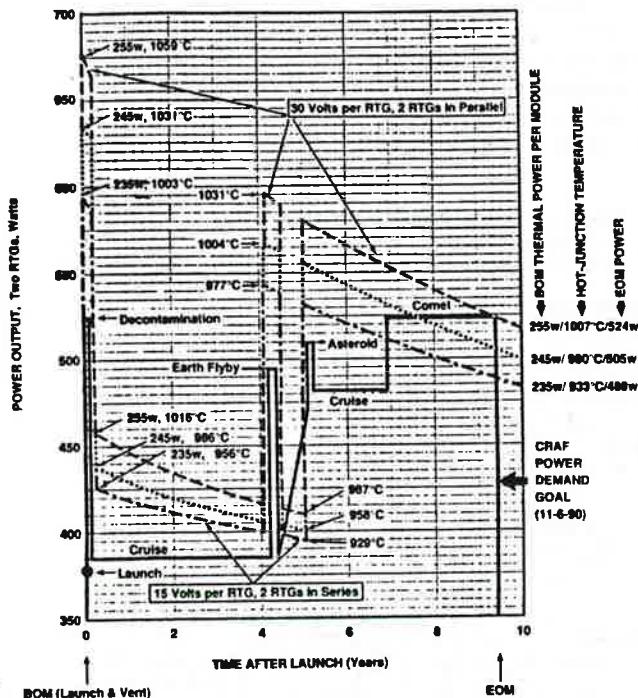
One way of further increasing the RTGs' power output is to combine voltage switching with larger radiator fins. As shown in Figure 27, this increases the EOM power by 11 watts (from 505 to 516 watts), and meets all of the power demand goals except those for the comet science phase. It exceeds that goal by 17 watts at the beginning of that phase, but falls 8 watts short of meeting it at the end of that phase.

**Figure 28. Effect of Number of Couple Rings on Output of Dual-Voltage RTGs**  
(35° Separation, 18 Heat Source Modules, 245-Watt BOM, 5" Fins)



Two more figures are shown to complete the picture for this option. Figure 28 depicts the effect of the number of couple rings on the output of the voltage-switched RTGs with 5" fins, and again shows 34 rings to be optimum. Finally, Figure 29 examines the effect of fuel loading on the output of 34-ring RTGs with voltage switching and 3" fins. It shows that if the BOM thermal power could be raised to 255 watts per heat source module, the EOM power goal could be met without changing the fins from their standard size.

Figure 29. Effect of Fuel Loading on Output of Dual-Voltage 34-Ring RTGs with 3" Fins (35° Separation, 18 Heat Source Modules, 15/30 Volts per RTG)

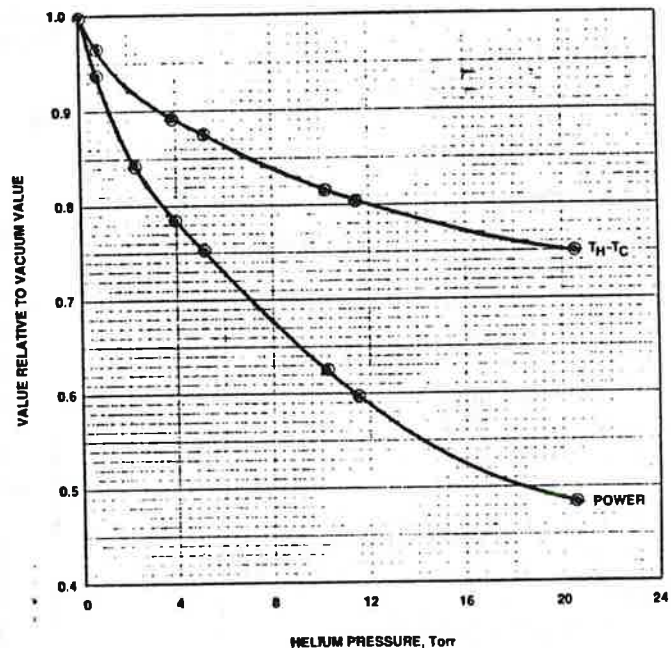


The final stratagem analyzed requires controlled-rate venting of the helium generated by the isotope's alpha decay. As indicated earlier in Figure 18, helium is such a good conductor and so effective in spoiling the thermal insulation that its retention above a continuum pressure would reduce the RTG's power output below the level required for the cruise phase. But if the helium pressure were low enough to be in the molecular-flow or transition regime, its thermal conductance could be low enough to yield the required power output. In principle, by careful control of the RTGs' helium pressure their power output could be regulated to just match the power demand histogram. Since that would minimize the hot-junction temperature history and thermoelectric degradation for the prescribed output profile, it would maximize the EOM power output.

The previously described analysis methods for vacuum RTGs must be modified to account for the effect of the lower pressure helium cover gas. Figure 30 shows the experimentally measured effect of helium pressure on a GPHS/RTG's power output  $P$  and thermoelectric temperature drop  $\Delta T$  relative to their respective vacuum values  $P_v$  and  $\Delta T_v$ . The data plotted represent recent measurements by Mound [ ] on the Galileo Qualification Unit (Q-1) after a five-year test. The measurements covered a pressure range from vacuum to 23 torr. As can be seen, the effect of increasing helium pressure on power and  $\Delta T$  is diminishing but has not quite leveled off yet. In other words, helium at 23 torr is still somewhat below the

continuum regime for this geometry. As expected, the fractional decrease in power output is approximately twice as great as that in  $\Delta T$ .

Figure 30. Effect of Helium Pressure on RTG Power and Temperature (Measured at Mound on Q-1 RTG)



The experimental data are cross-plotted in Figure 31, which shows that the relative temperature drop is essentially a linear function of the relative power output:

$$(30)$$

where  $\beta =$  from the slope of Figure 31. In the absence of further experimental data, our analysis assumed that Eq. (30), which gives the relative reduction of power and  $\Delta T$  for given helium pressures, applies not only to the Q-1 RTG but also to any similar RTG, throughout its mission.

From Eqs. 22 and 17 the hot-junction temperature at time  $t$  of the vacuum RTG with a BOM thermal power  $Q_1$  is given by

$$(31)$$

and from Eqs. 16 and 17 its undegraded power output is given by

$$(32)$$

Figure 31. Effectiveness of Controlled Helium Pressure In Lowering Temperature at Partial Power

