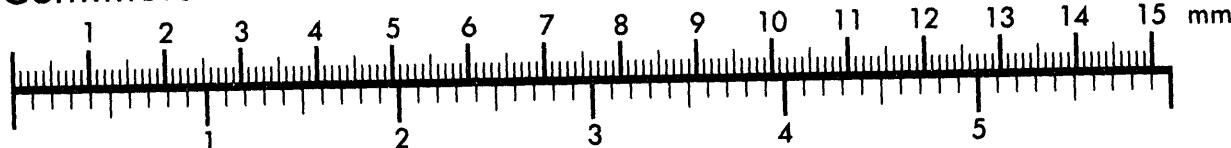


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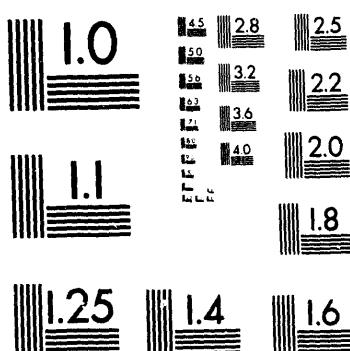
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# Thermo-Mechanical Analysis of the White-Beam Slits for an Undulator Beamline at the Advanced Photon Source

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## Abstract

A set of precision horizontal and vertical white-beam slits has been designed for an undulator beamline at the Advanced Photon Source. Due to the powerful x-ray heat flux emitted by the undulator, it is difficult to control the thermal distortion within the desired range of 1-2 microns. We analyzed many conceptual designs in an attempt to minimize the thermal distortion of the slits. Even with 1-mm-thick, low-Z material (graphite) coated on the heating surface of a traditional slit, the maximum thermal distortion is over 25 microns. A three-piece slit was then designed to satisfy the requirements. It consists of one large block, two tungsten knife edges, and an OFHC cooling tube (filled with copper mesh) brazed inside the large block. The thermal distortion at the knife edges of this three-piece slit has a relative displacement of less than 2 microns.

## 1 Introduction

A set of precision horizontal and vertical white-beam slits has been designed [1] for an undulator beamline at the Advanced Photon Source (APS). The slit, a knife-edge-type precision device, is required to have very small thermal distortion during operation with beam. The traditional slit consists of a large block and an OFHC cooling channel inside the block. Some users require that the thermal distortion be as small as about 1-2 microns at the knife edge. Due to the powerful x-ray heat flux coming from APS Undulator A [2], it is an exceedingly difficult problem to reduce the thermal distortion to less than 2 microns.

Dealing with the heat flux from Undulator A is an engineering challenge. At normal incidence, the peak power is about  $175 \frac{W}{mm^2}$  at 27.5 m, with a total power of 3.8 kW. The Gaussian deviation is about 1.22 mm, and the parabolic deviation is about 3.8 mm. The footprint of the x-ray beam is very small, which

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induces a very high thermal stress. A FORTRAN code allows the element (in the ANSYS finite element code used) subjected to the x-ray beam to get the power distribution automatically at its coordinate position.

Enhanced heat transfer technology [3] was used to increase the convective heat transfer coefficient with water from  $1 \frac{W}{cm^2 \cdot ^\circ C}$  to about  $3 \frac{W}{cm^2 \cdot ^\circ C}$  for the OFHC tube. However, we assumed a convective heat transfer coefficient of  $2 \frac{W}{cm^2 \cdot ^\circ C}$  in conservative calculations.

## 2 Material Considerations

It is very important that the slit body be made of a material having a very low thermal expansion coefficient and a high thermal conductivity. Both TZM (vacuum arc-cast molybdenum-0.5% titanium-0.1% zirconium alloy) and tungsten alloy (HD18) [4] were considered for the block material. OFHC is the material for the cooling tube. During operation, both the HD18 and TZM will reach temperature levels of about  $400^\circ C$ . Therefore, the material properties at such elevated temperatures were considered in the analysis. Table 1 shows the properties of the materials considered for the slit.

Table 1. Properties of the Materials Considered for the Slit

		Room Temperature			400°C	
Properties		TZM	HD18	OFHC	TZM	HD18
Conductivity	$K \frac{W}{cm \cdot ^\circ K}$	1.45	1.65	3.95	1.3	1.35
Young's Modulus	$E \text{ MPa}$	$3.15 \times 10^5$	$3.65 \times 10^5$	$1.19 \times 10^5$	$3.0 \times 10^5$	$3.1 \times 10^5$
Thermal Exp. Coef.	$\alpha \frac{1}{^\circ C}$	$5.4 \times 10^{-6}$	$4.5 \times 10^{-6}$	$17.7 \times 10^{-6}$	$5.6 \times 10^{-6}$	$4.6 \times 10^{-6}$
Poisson Ratio	$\nu$	0.293	0.28	0.307	.361	.28
Yield Strength	$\sigma_y \text{ MPa}$	1000	580	300	800	380
Ultimate Strength	$\sigma_u \text{ MPa}$	1000+	1500+	350	1000+	1500+

When the temperature rises from room temperature to about  $400^\circ C$ , the yield strength of HD18 will drop from 580 MPa to 380 MPa. Therefore, we decided to use TZM as the material for the large block because its yield strength (test at 0.05 inch per inch per minute, vacuum  $10^{-5}$  mm Hg) is above 1000 MPa at room temperature and is about 700 MPa at  $500^\circ C$ . (Without the vacuum test condition, the TZM also performs well; for a 0.06-inch stress-relieved sheet, the yield strength is above 1000 MPa at room temperature and is

about 690 MPa from 400°C to 650°C.)

### 3 Design Progression

We analyzed many large block conceptual designs in order to

1. reduce the maximum temperature of the large block to under 400°C if tungsten is used,
2. reduce the maximum thermal distortion to the required 1-2 microns,
3. reduce the interfacial stress between the big block and the OFHC cooling tube,
4. reduce the localized stress concentration at the support area.

The slit design progressed in several stages due to severe distortion problems under the predicted heat load.

The analysis proves that a single piece of large block will always have a thermal distortion of about 90 microns (when subjected to Undulator A heat flux) with the best possible boundary conditions for cooling and supporting. Even with a 1-mm-thick, low-Z material (graphite) coated on the heating surface  $\triangle ADEF$  of the large block (see Fig. 1), the maximum thermal distortion is still over 25 microns. In order to reduce the thermal distortion to 2 microns, a three-piece slit was designed [1] that consists of one large block (as shown in Fig. 1), two tungsten knife-edge small blocks (as shown in Fig. 2), and an OFHC cooling tube brazed inside the large block. The major function of the large block is to act as an absorber that will take most of the heat flux. This results in a maximum deflection of about 100 microns in the large block (final design). The thermal distortion at the knife edges of this three-piece slit has a relative displacement of less than 2 microns.

Figure 1 depicts the large slit block geometry of the final design. It is about  $300 \times 26 \times 52$  mm (x, y, and z coordinates, respectively). Due to the limited space, the grazing incidence angle is fixed at  $2.5^\circ$  on both  $\triangle ADEF$  and  $\triangle ABGH$  surfaces (as shown in Fig. 1).

The critical surfaces [1] shown in Fig. 1 are :

1.  $\triangle ACD$ , which causes a clear cut of the x-ray beam in the horizontal plane.  $U_y$  (displacement in y-axis) is the only important displacement at this surface.
2.  $\triangle ABC$ , which creates a clear cut of the x-ray beam in the vertical plane of the x-ray beam.  $U_z$  is the only important displacement at this surface.

Both  $\triangle ACD$  and  $\triangle ABC$  will never be hit by the x-ray beam [1]. When the beam hits  $\diamond ADEF$  (in Fig. 1), the vertical beam is stretched out along the x-axis. When the beam hits  $\diamond ABGH$  (in Fig. 1), the horizontal beam is again stretched out along the x-axis. However, the beam-spreading length in the x direction on  $\diamond ABGH$  is two times more than that on  $\diamond ADEF$ , which implies that the temperature and resulting thermal stress on  $\diamond ADEF$  will be about twice those on  $\diamond ABGH$ .

A closed form solution (see [5]) is used for the design optimization. The heat equation for a thermal plate (as shown in Fig. 3) subjected to a Gaussian-distributed heat flux is

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0. \quad (1)$$

The boundary conditions are

$$\begin{aligned} -k \frac{\partial T}{\partial y}(x, 0) &= q_o \exp\left(-\frac{x^2}{r_o^2}\right) \quad 0 \leq x \leq a. \\ -k \frac{\partial T}{\partial y}(x, 0) &= 0 \quad a < x \leq b \\ -k \frac{\partial T}{\partial y}(x, t) &= h_c(T - T_\infty) \\ \frac{\partial T}{\partial x}(0, y) &= \frac{\partial T}{\partial x}(b, y) = 0. \end{aligned} \quad (2)$$

The solutions give

$$\theta = C_o \left[ \eta - 1 - \frac{1}{B_i} \right] + \sum_{m=1}^{\infty} C_m e^{-\lambda_m \eta} \cos(\lambda_m \xi) \left[ 1 + \frac{\lambda_m - B_i}{\lambda_m + B_i} e^{2\lambda_m(\eta-1)} \right], \quad (3)$$

where

- T: Temperature ( $^{\circ}\text{C}$ )
- k: Thermal conductivity ( $\text{W}/\text{m}\cdot^{\circ}\text{C}$ )
- XYZ: Fixed coordinate
- q: Heat flux ( $\text{W}/\text{m}^2$ )
- t: Thickness of plate (m)
- h: Convective heat transfer coefficient ( $\text{W}/\text{m}^2 \cdot ^{\circ}\text{C}$ )
- $T_\infty$ : Ambient temperature (assumed to be  $24^{\circ}\text{C}$ )
- b: Width of the plate (m)
- a: Width where q is applied to the large block (m)
- $r_o$ : Standard deviation (m),

and

$$\begin{aligned} \alpha &= \frac{a}{t}, \quad \beta = \frac{b}{t}, \quad \xi = \frac{x}{t} \\ \eta &= \frac{y}{t}, \quad \theta = \frac{k(T - T_\infty)}{qt}, \quad B_i = \frac{h_c t}{k}, \end{aligned} \quad (4)$$

$$C_o = -\frac{r\sqrt{\pi}}{2\alpha} \operatorname{erf}\left(\frac{\alpha}{r}\right); \quad (5)$$

$\operatorname{erf}(\dots)$  is known to be the *error function*, and

$$C_m = -\frac{\sqrt{\pi} r \exp\left(-\frac{r^2 \lambda_m^2}{4}\right) \operatorname{Re} \left[ \operatorname{erf}\left(\frac{2\alpha + ir^2 \lambda_m}{2r}\right) \right]}{\lambda_m \alpha \left( -1 + e^{-2\lambda_m} \frac{\lambda_m - Bi}{\lambda_m + Bi} \right)}, \quad \lambda_m = \frac{m\pi}{\beta}, \quad m \in N. \quad (6)$$

The analysis indicates that, while the tube I.D. (inside diameter) of the large block is fixed at 12.7 mm and  $h=2\frac{W}{cm^2 \cdot ^\circ C}$ , the lowest surface temperature is obtained if the wall thickness (the distance between the heating surface and the inside wall of the tube) equals 3-4 mm as shown in Fig. 4. The average wall thickness of the design is about 14.5 mm (11 mm at  $\overline{AF}$ , 18 mm at  $\overline{ED}$  as shown in Fig. 1). A decrease in the value of  $h$  from  $2\frac{W}{cm^2 \cdot ^\circ C}$  to  $1\frac{W}{cm^2 \cdot ^\circ C}$  will increase the maximum temperature from  $410^\circ C$  to  $500^\circ C$ . Increases in the I.D. also help to decrease the maximum temperature. Also, while the average wall thickness is fixed at 14.5 mm and  $h=2\frac{W}{cm^2 \cdot ^\circ C}$ , increasing the tube I.D. from 5 mm to 10 mm will decrease the maximum temperature from  $636^\circ C$  to  $450^\circ C$ . It was suggested that the cooling channel be designed so that it remains parallel to the heating surface in order to keep the wall thickness at a constant 5 to 6 mm. However, this idea compounds the design difficulty so we stayed with the design as shown in Fig. 1.

The most significant design changes after optimization were:

1. The outer diameter of the cooling channel was increased from 16 mm to 19 mm as shown in Fig. 1.
2. The cooling channel is closer to the heating surface, from 11 mm at  $\overline{AF}$  to 18 mm at  $\overline{ED}$  (as shown in Fig. 1), which is about an average of 2 mm less than in the original design.

## 4 Results and Discussion

### 4.1 Large Block

The worst possible case for the temperature field is when the beam center hits 3 mm away from  $\overline{AD}$  and 15 mm away from  $\overline{ED}$  on  $\diamond ADEF$  (in Fig. 1). The worst possible case for both thermal distortion and the stress field is when the beam center hits around  $\overline{AF}$  and 3 mm from  $\overline{AD}$  on  $\diamond ADEF$  (as shown in Fig.

1). Because the maximum temperature is not a concern, we will focus on the worst case for the thermal distortion and stress field in the following discussions.

Figure 5 shows the temperature contour of the large block. Figure 6 shows the displacement ( $U_z$ ) contour. The calculated results for the original and final designs are shown in Table 2.

Table 2. Calculated Results from ANSYS Finite Element Code for the L5 Slit

	Final Design			Early Design		
	(Tube O.D.=19-mm)			(Tube O.D.=16-mm)		
$h(\frac{W}{cm^2 \cdot ^\circ C})$	1	2	3	1	2	3
Max. Temp on Slit Surface ( $^\circ C$ )	471	424	406	600	540	520
Max. Eff. Stress on Surface (MPa)	384	374	372	380	374	370
Max. Temp at Interface ( $^\circ C$ )	191	138	120	180	150	140
Max. Eff. Stress at Interface (MPa)	273	177	141	300	230	165
Max. Deflection $U_y$ at $\Delta$ ACD (micron)	20	16	15	25	20	20
Max. Deflection $U_z$ at $\Delta$ ABC (micron)	110	90	85	150	125	115

The stress component  $\sigma_x$  on  $\diamond ADEF$  is much larger than the other two components,  $\sigma_y$  and  $\sigma_z$ , because the thermal gradient is much larger in that direction (the beam-spreading direction). The maximum displacement of the whole structure is in the  $x$  direction ( $U_x$ ), if free expansion is allowed. The slit design [1] allows the large block to expand freely along the  $x$ -axis; the minimum  $\sigma_x$  on  $\diamond ADEF$  is about 410 MPa (compressive stress). If the expansion along the  $x$ -axis is totally constrained, the minimum  $\sigma_x$  becomes about 460 MPa (compressive stress). Because the displacement component  $U_x$  in the large block is not a concern, to allow it to expand freely in the longitudinal ( $x$ ) direction will reduce the stress component  $\sigma_x$  by 12%; also, the effective stress will be reduced by about 10%. The other two displacement components,  $U_y$  and  $U_z$ , are relatively small at the pin-hole area as shown in Fig. 1. The thermal stress is small even when the block is fully constrained in both the  $y$  and  $z$  directions around the pin-hole area.

For the final design, increasing  $h$  from  $1 \frac{W}{cm^2 \cdot ^\circ C}$  to  $2 \frac{W}{cm^2 \cdot ^\circ C}$  will help to:

1. decrease the displacement component  $U_z$  from 110 microns to 90 microns,
2. decrease the maximum interfacial stress (between the large block and the OFHC cooling tube) from 273 MPa to 177 MPa, which helps the OFHC remain in elastic deformation, and
3. decrease the maximum temperature at the bonding surface from  $191^\circ C$  to  $138^\circ C$ , which will keep the OFHC from approaching the annealing temperature ( $175^\circ C$ ).

Increasing the depth in the z direction will increase  $I_y$ . This helps to reduce  $U_z$  on the  $\triangle ABC$ , but will decrease the natural frequency from the vibration point of view. The height in the original design was 36 mm compared to 52 mm in the final design (as shown in Fig. 1). To increase the height another 10 mm to 62 mm will help to reduce  $U_z$  at  $\triangle ABC$  by 8 microns with the same boundary conditions.

When the x-ray beam hits both  $\diamond ADEF$  and  $\diamond ABGH$ , the power deposited on them is complicated. Figure 7 (model for the original design) shows the temperature contour of such a case. It is easy to see that the x-ray beam extends much longer on  $\diamond ABGH$  than on  $\diamond ADEF$ . When more x-ray beam deposits on  $\diamond ABGH$  and less on  $\diamond ADEF$ , the resulting temperature, stress, and distortion are less severe. Besides, more x-rays will pass through  $\angle DCB$  instead of being deposited on the large block.

## 4.2 Knife Edge

Only HD 18 was considered as the material for the two knife-edge slits. The two knife edges slits are V1, which defines the vertical beam, and H1, which defines the horizontal beam. Edge V1 is sitting on  $\diamond EDMN$  in Fig. 1. Edge H1 is sitting behind the V1. Both V1 and H1 have very small surface areas, which will be subjected to the normal incident x-ray radiation from Undulator A. The small surface area of H1 is about 50 micron  $\times$  71 mm, located from  $\overline{CD}$  to 50 microns in the positive y direction of  $\overline{CD}$  (in Fig. 1). The dimension in the z direction is the same as  $\overline{CD}$  at 71 mm. The small surface area of V1 is about 100 micron  $\times$  65 mm, located from  $\overline{CB}$  to 100 microns in the positive z direction of  $\overline{CB}$  in Fig. 1. The dimension in the y direction is the same as  $\overline{CB}$  at 65 mm.

Edge V1 will absorb 114 W in the worst case. Edge H1 will be subjected to about 50 W in the worst case. V1 and H1 do not have a cooling system of their own. The cooling of V1 and H1 is conducted through  $\diamond EDMN$

to the large block. The thermal distortion on the two knife edges is very small, about 10 microns for V1 in the worst case as shown in Fig. 2. Because of the thermal distortion component  $U_z$  of the large block, V1 will never have a surface area larger than 100 micron  $\times$  65 mm that is subjected to the x-ray beam. Suppose  $U_z$  is 80 microns for the large block, the surface area subjected to the x-ray beam will decrease to 20 microns  $\times$  65 mm, which will induce only a 1-micron thermal distortion.

Edge V1 remains perfectly flat even at the worst case; there is only a 1-2 micron difference between the center and the two sides of the knife edge as shown in Fig. 2. This is because the cooling tube is far away (at least 30 mm) from the heating surface; the temperature distribution is more even throughout the knife edge, so that the displacement component  $U_z$  along the knife edge is nearly constant. The relative displacement  $\Delta_{2-1}$  (relative to point 1) is less than 1 micron as shown in Fig. 2.

## 5 Future Work

The idea of coating a low-Z material (for example, graphite) on the heating surface ( $\diamond ADEF$ ) of the large block was also evaluated in preliminary calculations. If it becomes technically possible to braze 1-mm-thick graphite ( $k=12 \frac{W}{cm \cdot K}$ ) on tungsten or TZM, the thermal-structural results become very promising. The maximum temperature on the tungsten is calculated to be about  $165^\circ C$ . The maximum displacement component  $U_z$  is about 28 microns.

For each low-Z material element (in the finite element model), the absorbed power can be written as

$$q_{\text{Absorb}} \left( \frac{W}{\text{Volume}} \right) = \int_{E_{\text{min}}}^{E_{\text{max}}} S(E) \cdot \exp^{-\alpha(E) \cdot z} \cdot \alpha(E) \cdot dE. \quad (7)$$

For high-Z materials, the transmitted power that hits on the surface of the tungsten (or TZM), after passing through the thickness of the low-Z material, can be written as:

$$q_{\text{transmited}} \left( \frac{W}{\text{Area}} \right) = \int_{E_{\text{min}}}^{E_{\text{max}}} S(E) \cdot \exp^{-\alpha(E) \cdot z} \cdot dE, \quad (8)$$

where

$$S(E) = \text{source spectrum}$$

$$\alpha(E) = \frac{1}{\text{thickness}} \text{ absorption coefficient of low-Z material (function of material and energy)}$$

$$z \text{ thickness.}$$

The program USP [6] was used to calculate the Undulator A spectrum, and the program TRANSMIT (part of SHADOW [7]) was used to calculate the function of  $\alpha(E)$ .

## 6 Acknowledgments

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Figure Captions

**Figure 1:** Computer Model of the Final Design for L1 Vertical Slit

**Figure 2:** The Displacement ( $U_z$ ) Contour of the Knife Edge Slit

**Figure 3:** Channel Model and Simplified Plate Model for Closed Form Solution

**Figure 4:** Parametric Study of the Maximum Temperature ( $^{\circ}C$ ), Various  $h$  and Wall Thicknesses

**Figure 5:** Temperature Contour of the Large Block when the Maximum  $U_z$  occurs

**Figure 6:** The Displacement ( $U_z$ ) Contour of the Large Block when the Maximum  $U_z$  Occurs

**Figure 7:** Temperature Contour of the Large Block when the X-ray Beam Hits Both  $\triangle ADEF$  and  $\triangle ABGH$

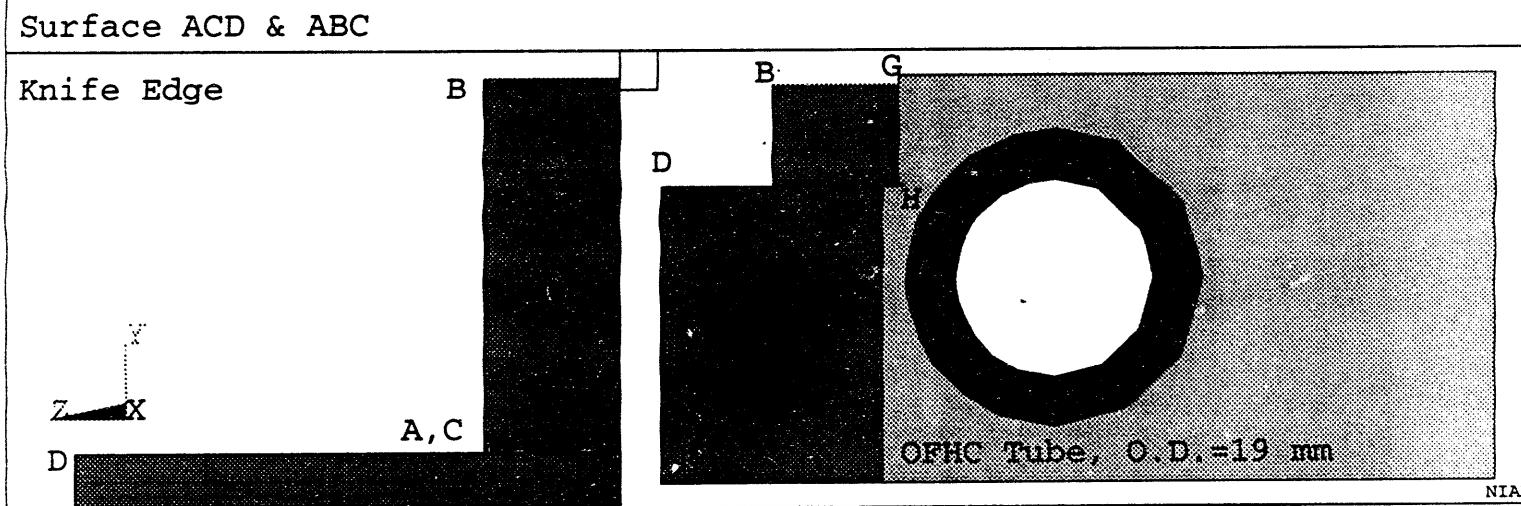
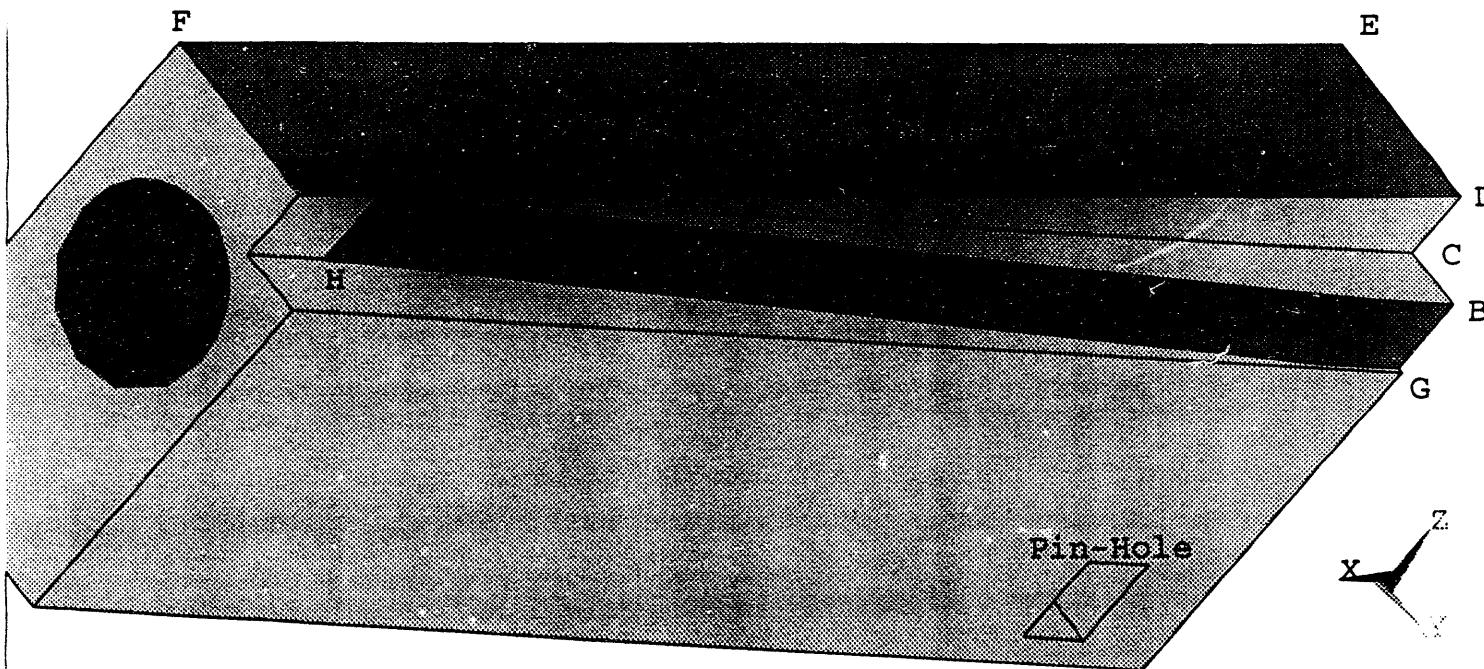
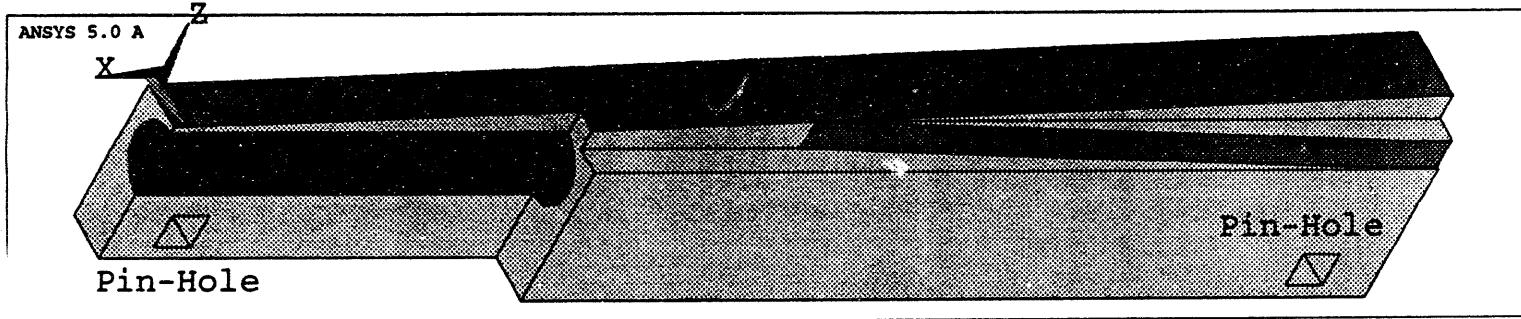
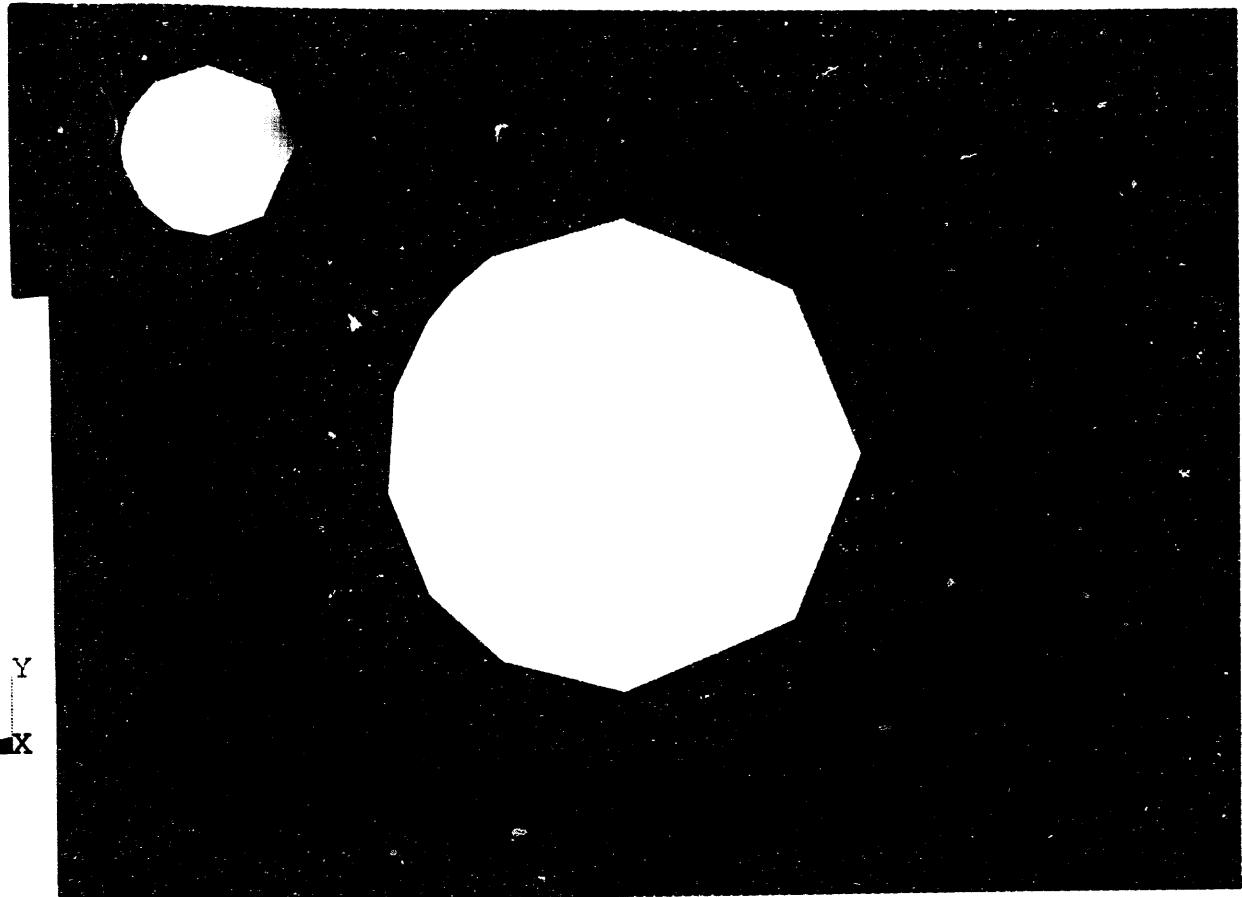


Figure 1

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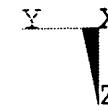
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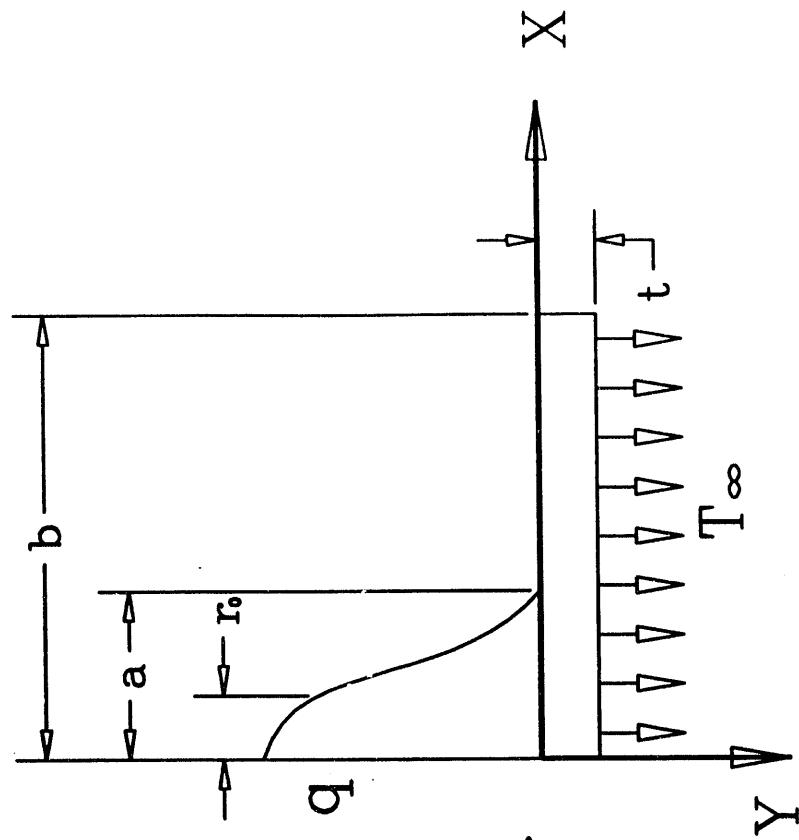
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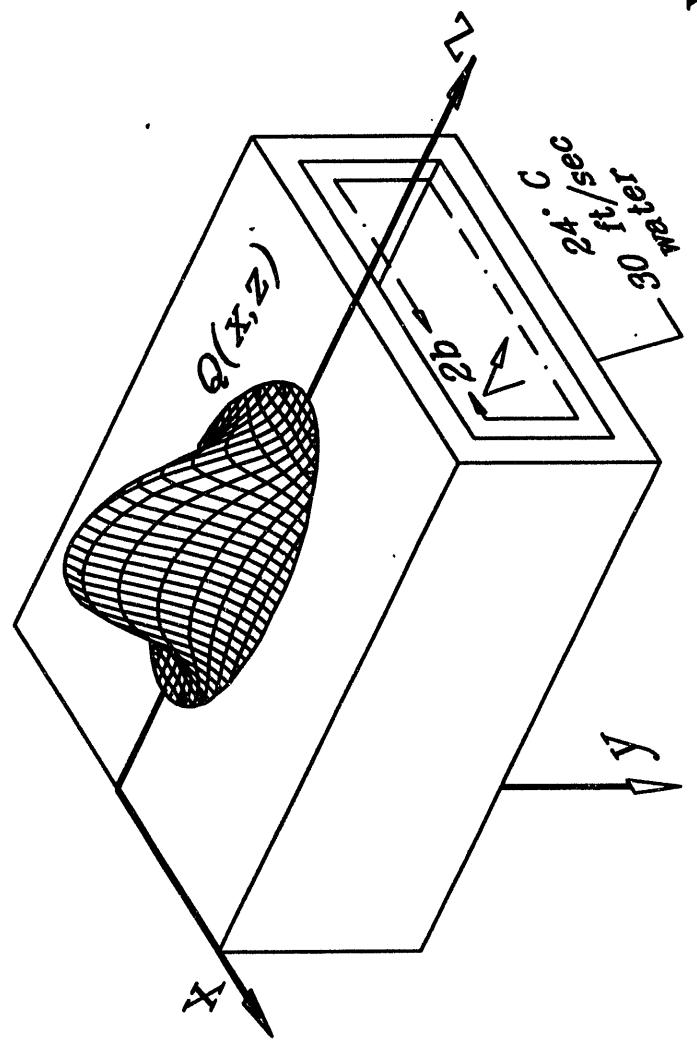


Knife Edge

NIAN



## Simplified Model with Curve-Fit Gaussian Profile (Closed Form)



# Channel Model with Real Beam Profile (ANSYS)

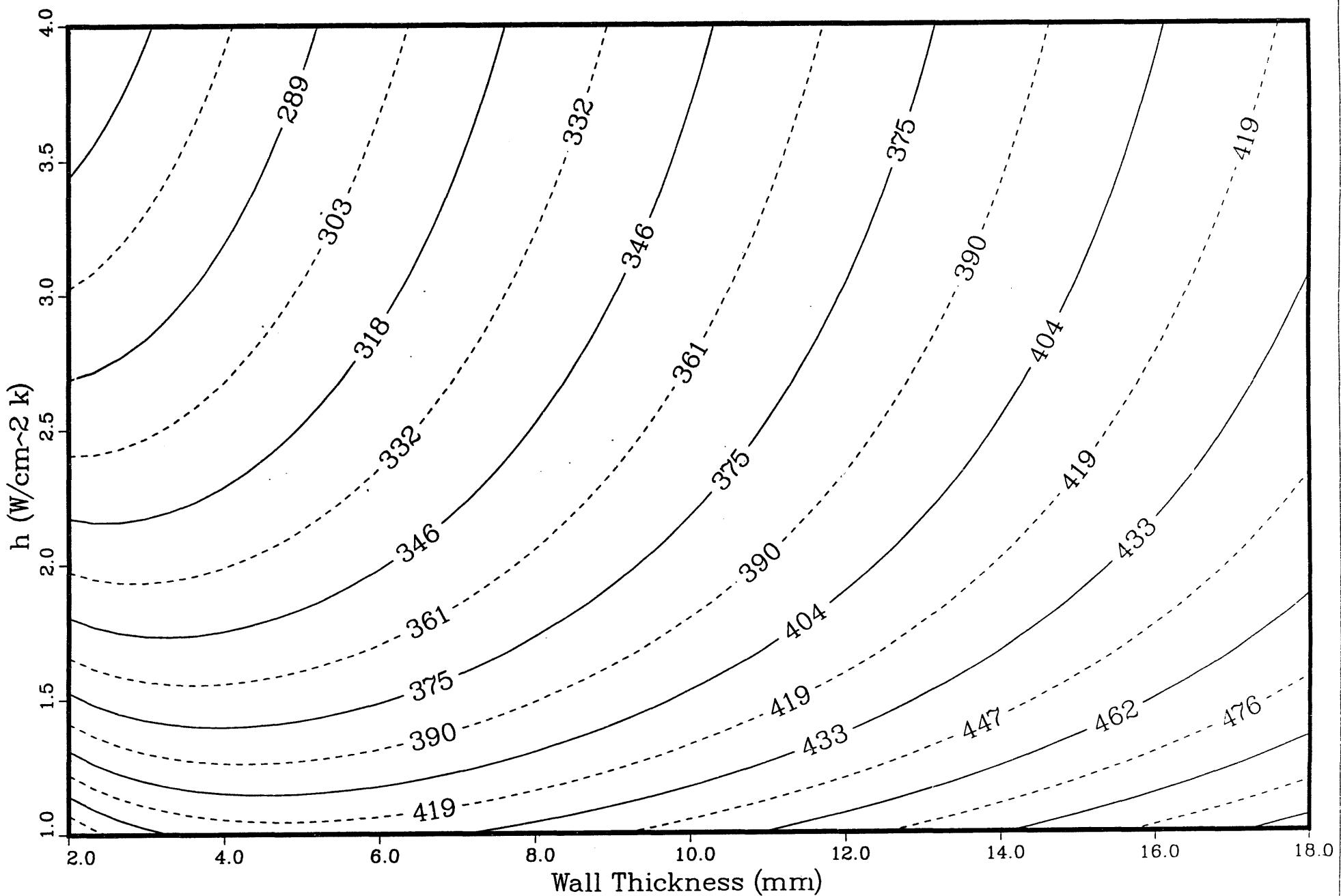


Fig. 4



ANSYS 5.0 A

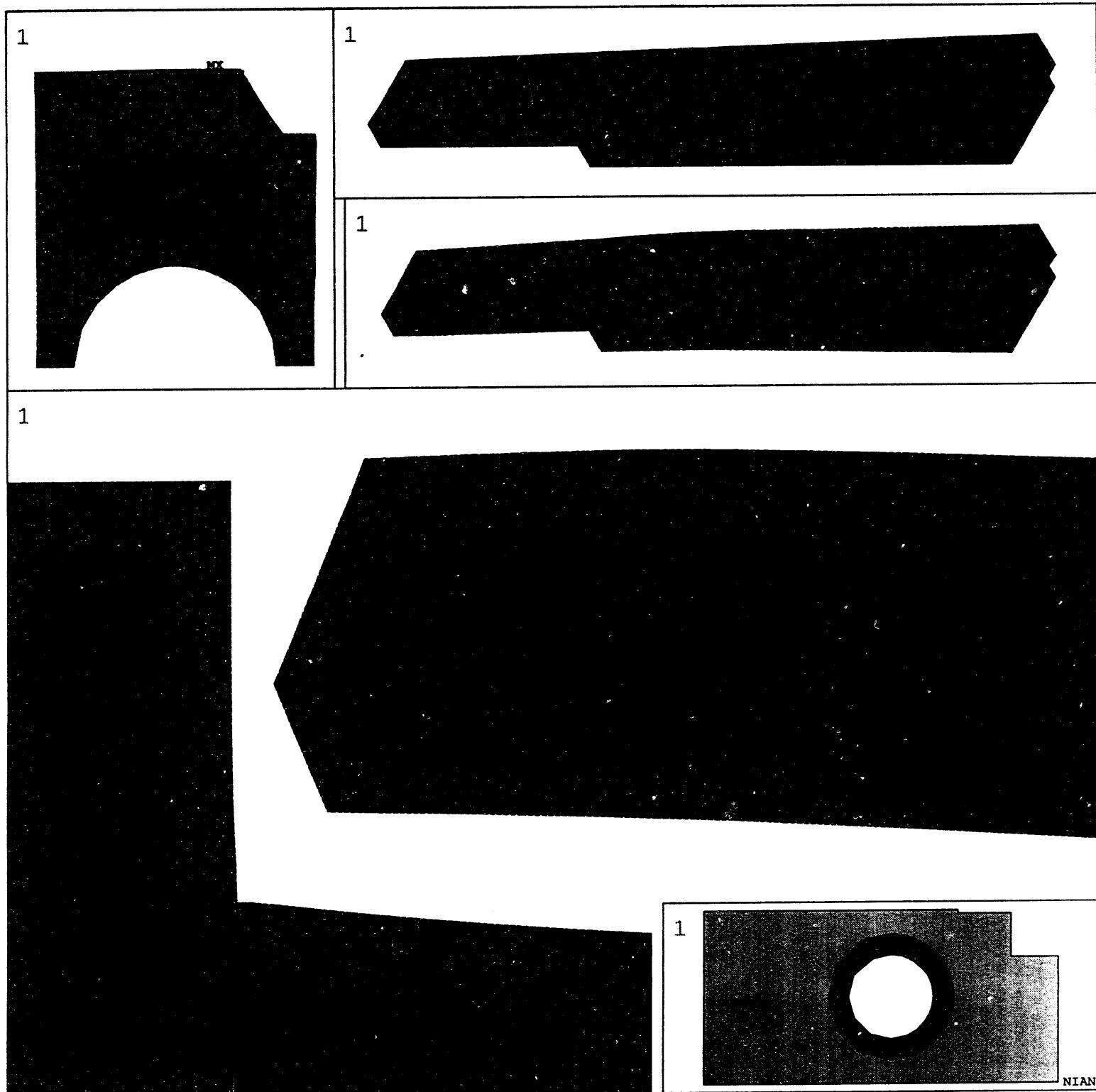


Figure 6

UZ Contour (mm)  
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