

Conf - 9110146-3

ANL/CP--74079

DE92 004130

Misalignment Sensitivity of an Inclined Crystal Monochromator*

A.T. Macrander, and W.K. Lee
Advanced Photon Source
Argonne National Laboratory
9700 South Cass Avenue
Argonne, IL 60439

DEC 13 1991

October, 1991

DISCLAIMER

CZ

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

*This work supported by the U.S. Department of Energy, BES-Materials Sciences,
under contract no. W-31-109-ENG-38

MASTER

DISTRIBUTION OF THIS DOCUMENT IS UNLIMITED

p2

MISALIGNMENT SENSITIVITY OF AN INCLINED CRYSTAL MONOCHROMATOR

A.T. Macrander and W-K Lee

Advanced Photon Source, Argonne National Laboratory, Argonne, IL 60439

abstract

The sensitivity of a novel inclined crystal monochromator design to misalignments has been calculated, and compared to data. Rocking curve line narrowing as well as broadening can occur because the asymmetry factor of dynamical diffraction given by $b = \mathbf{s}_0 \cdot \mathbf{n} / \mathbf{s}_H \cdot \mathbf{n}$ can have an absolute value larger or smaller than unity. Here $\mathbf{s}_0 \cdot \mathbf{n}$ and $\mathbf{s}_H \cdot \mathbf{n}$ are the direction cosines of the incident and diffracted beams, respectively, and \mathbf{n} is the inward surface normal. An inclined double crystal monochromator which is perfectly aligned would have $b = -1$ for both crystals, and only then would the diffraction be symmetric. We have computed b and rocking curve widths for inclination angles of 70.53° and 85.00° , and we compare the 70.53° case to data for silicon {111} reflections using 8 KeV ($\text{CuK}\alpha_1$) radiation. The 70.53° case applies to $(11\bar{1})$ reflection from a (111) oriented crystal. We report that rotations around the reciprocal lattice vector have the expected effect on b .

I. Introduction

The inclined crystal arrangement has recently been tested as a high heat load monochromator¹ and a sensitivity to rotations around $\bar{\mathbf{H}}$, the reciprocal lattice vector, was invoked to explain observed rocking curve widths. Rotation around $\bar{\mathbf{H}}$ constitutes a degree of freedom available with a symmetric reflection. This degree of freedom, which we have called ρ , is sacrificed in the inclined case, and a misalignment in ρ results in truly asymmetric diffraction, i.e., diffraction with b not equal to -1.

After deriving the relevant trigonometric relations, we present the supporting data which we have obtained.

II. Asymmetry factor as a function of azimuthal angle

The coordinate system we have used is shown in Fig. 1a. The inclination angle is denoted as β . The \hat{z} axis lies along \vec{H} , and the \hat{x} axis lies in the crystal surface along the intersection of a Bragg plane (the x-y plane) with the surface. \hat{n} is defined to be the inward surface normal. (Choosing the inward instead of the outward normal agrees with the convention of Zachariasen².) In our coordinate system \hat{n} is given by

$$\hat{n} = \hat{y} \sin \beta - \hat{z} \cos \beta \quad [1]$$

We consider an incident plane wave with a wavevector along the direction \hat{s}_0 which is shown in Fig. 1b. \hat{s}_0 is given by

$$\hat{s}_0 = \hat{x} \cos \theta \cos \rho + \hat{y} \cos \theta \sin \rho - \hat{z} \sin \theta \quad [2]$$

Here θ is the angle \hat{s}_0 makes with respect to the Bragg plane and ρ is the azimuthal rotation about \vec{H} . The direction of the Bragg diffracted beam, \hat{s}_H , is also shown in Fig. 1b and is given by

$$\hat{s}_H = \hat{x} \cos \theta \cos \rho + \hat{y} \cos \theta \sin \rho + \hat{z} \sin \theta \quad [3]$$

From Eqs. 1-3 it is straightforward to write down the asymmetry factor given by

$b = \hat{s}_0 \cdot \hat{n} / \hat{s}_H \cdot \hat{n}$. The result is

$$b = \frac{\sin \beta \cos \theta \sin \rho + \sin \theta \cos \beta}{\sin \beta \cos \theta \sin \rho - \sin \theta \cos \beta} \quad [4]$$

We note that for $\rho = 0$, we have $b = -1$ which corresponds to an aligned inclined crystal. (Also, when $\rho = -90^\circ$ we have the conventional asymmetric situation at low incidence angle with $b = -\sin(\theta - \beta) / \sin(\theta + \beta)$).

We have used Eq.4 to obtain values for b and have used them to calculate values for the FWHM of the Darwin-Prins curves. Values for the FWHM were calculated by dividing FWHM values for symmetric diffraction by $|b|^{1/2}$. Values for Si (11 $\bar{1}$) diffraction from a (111) surface at 5 KeV are shown in Fig. 2. For this situation $\beta = 70.529^\circ$ and $\theta = 23.292^\circ$. The same situation at 8 KeV is shown in Fig.3. Here $\theta = 14.221^\circ$. Finally, values for Si (11 $\bar{1}$) diffraction for $\beta = 85^\circ$ at 5 KeV are shown in Fig.4. For this latter situation the surface is cut 14.471° from (111). We see that the Darwin width is very sensitive to ρ at large inclination angles.

III. Experimental Verification

We obtained rocking curves in a double crystal arrangement as a function of the azimuthal rotation, ρ , of the second crystal. The first crystal was set for symmetric (111) diffraction on a tube source and remained stationary. The fixed anode was Cu and we used $\text{CuK}\alpha$ radiation at 8.04 KeV. The crystal that we used was float zone Si with a (111) oriented surface, and we made rocking curves for the (11 $\bar{1}$) reflection of the second crystal. The results are shown in Fig. 3. Since the zero for ρ was somewhat

arbitrary, we have adjusted the data by applying an arbitrary offset to the ρ values. The narrowest rocking curve which we obtained is shown in Fig.4 and had FWHM of 8.4 arcsec. This value is very close to the ideal Darwin-Prins value of 7.5 arcsec, and reveals that at large ρ the FWHM of the second crystal was only 3.8 arcsec since the quadrature addition of 7.5 arcsec for the first crystal and 3.8 arcsec for the second crystal yield the observed result of 8.4 arcsec. The other FWHM values plotted in Fig.3 were obtained in the same way, i.e., by subtracting in quadrature a value of 7.5 arcsec for the first crystal from the observed FWHM.

IV. Summary

Values both larger and smaller than the ideal Darwin-Prins FWHM were obtained by adjusting ρ in our experiments. This observation and the reasonable agreement between rocking curve data on the one hand and calculations on the other hand corroborate Eq. 4 .

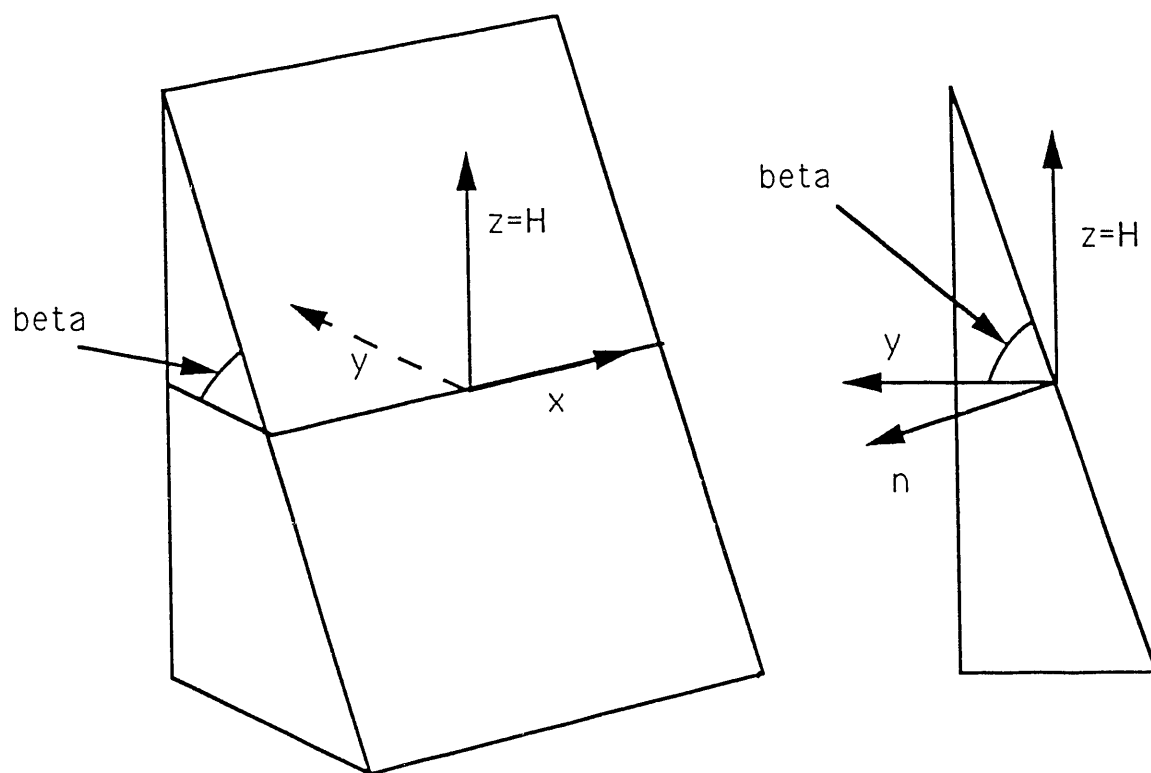
Acknowledgements

We would like to acknowledge the considerable input of D.M. Mills concerning alignment issues for an inclined crystal monochromator.

¹ A.T. Macrander, W.K. Lee, R.K. Smither, D.M. Mills, S. Rogers, and A. Khounsary, . these proceedings

² W.H. Zachariasen, "Theory of X-ray Diffraction in Crystals", (Dover , NY, 1945).

(a)



(b)

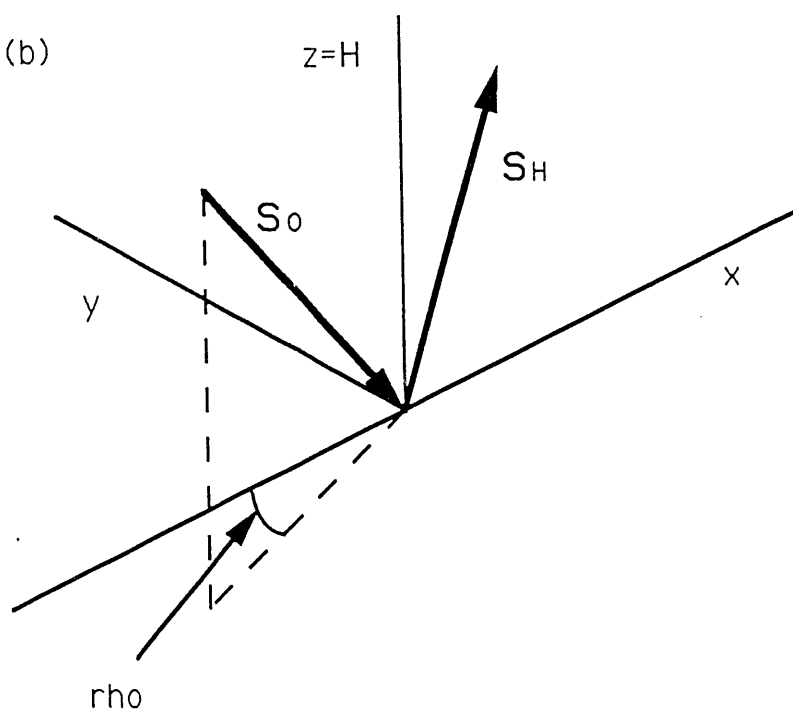
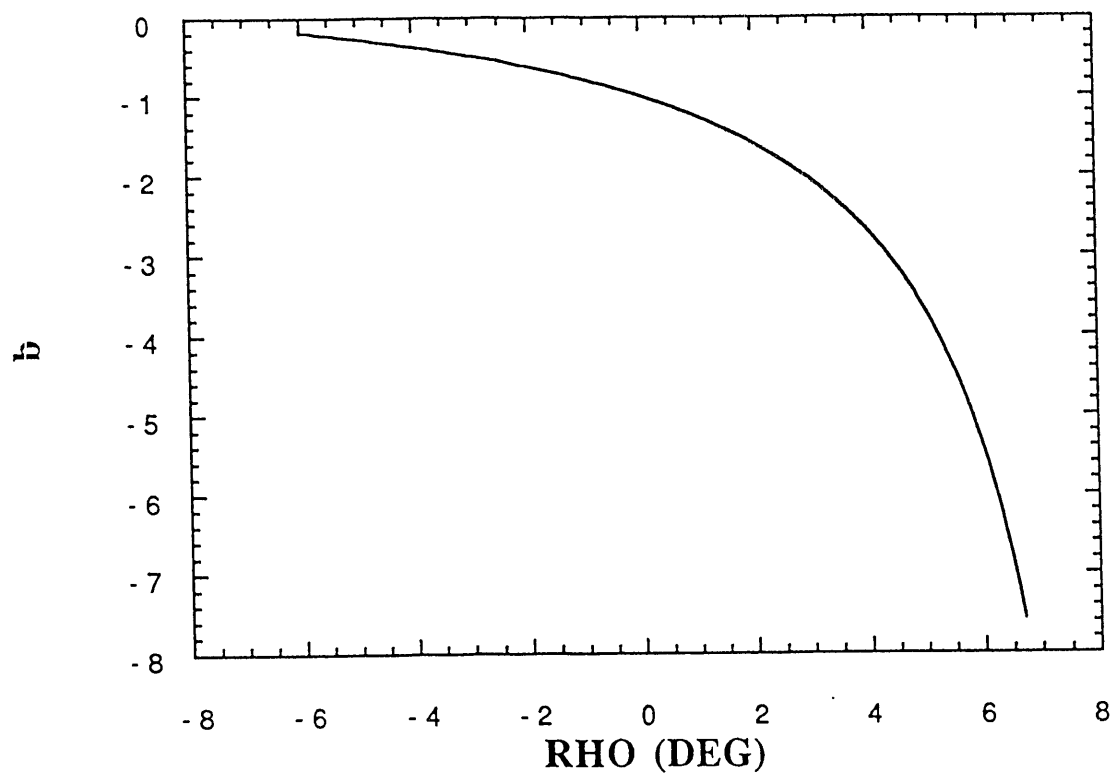
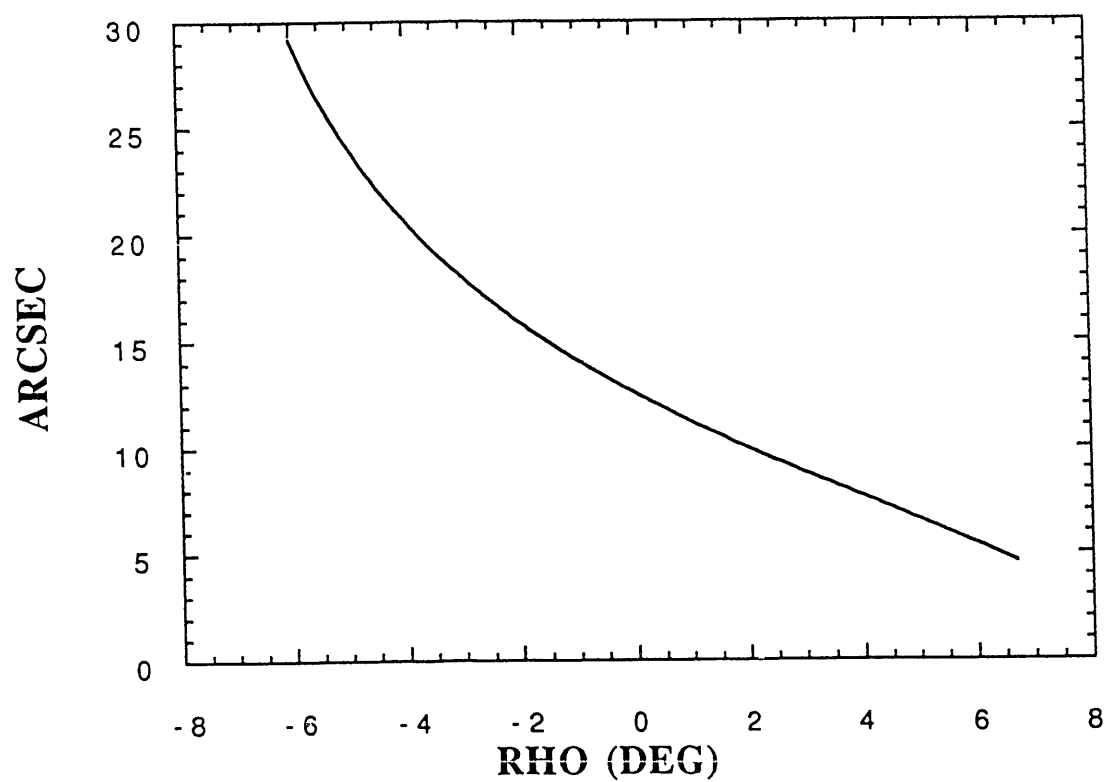


Fig. 1

Si(111) 70.529° INCLINED AT 5 KeV



FWHM OF THE DARWIN-PRINS CURVE



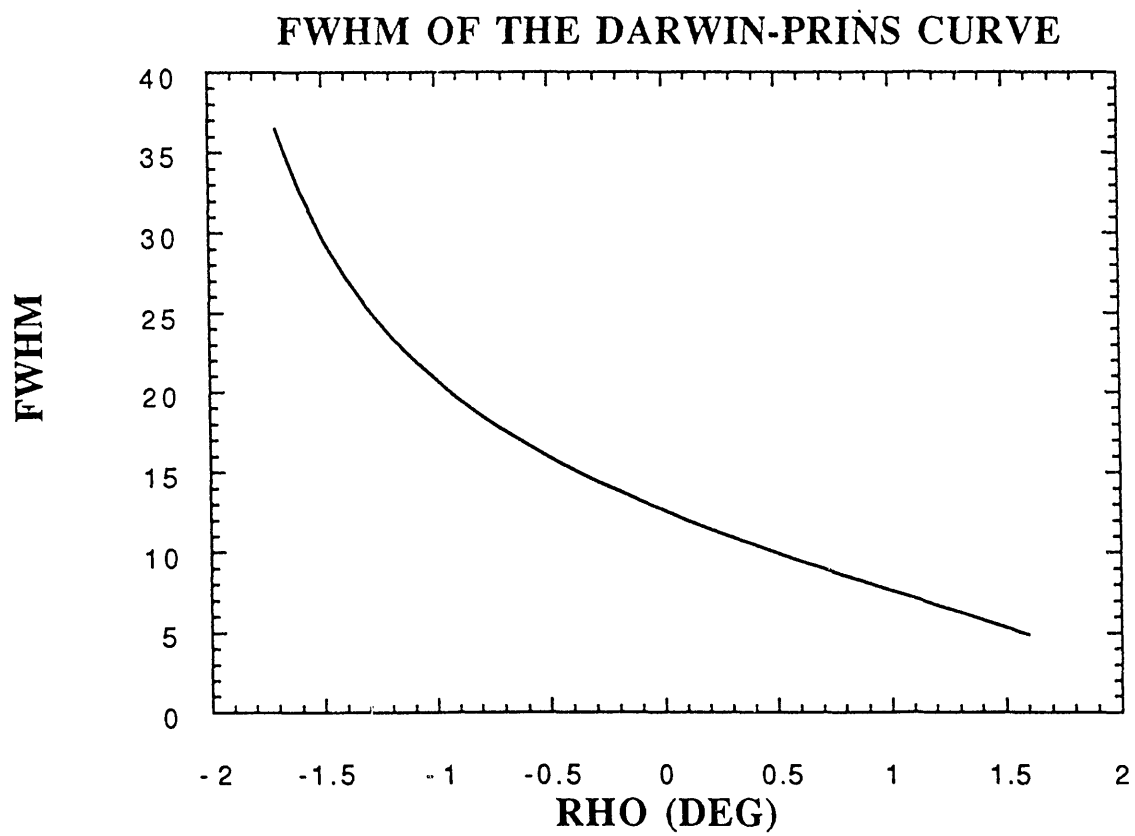
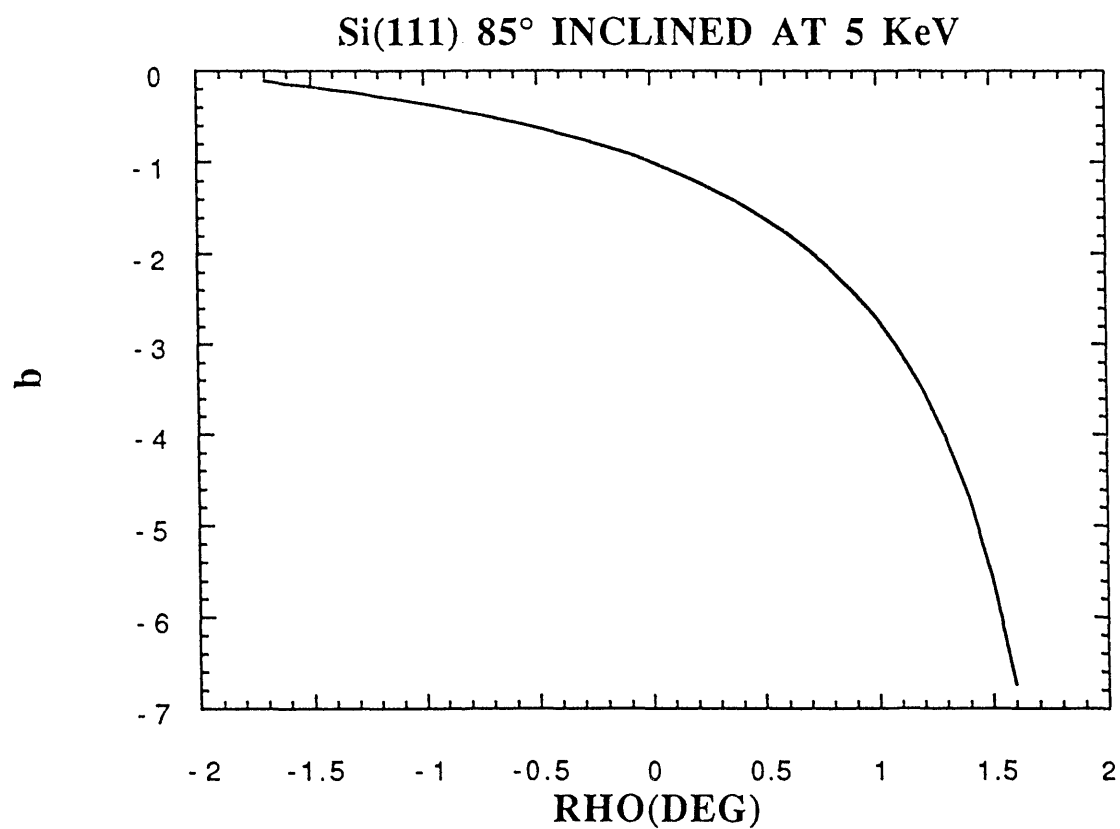


Fig. 3

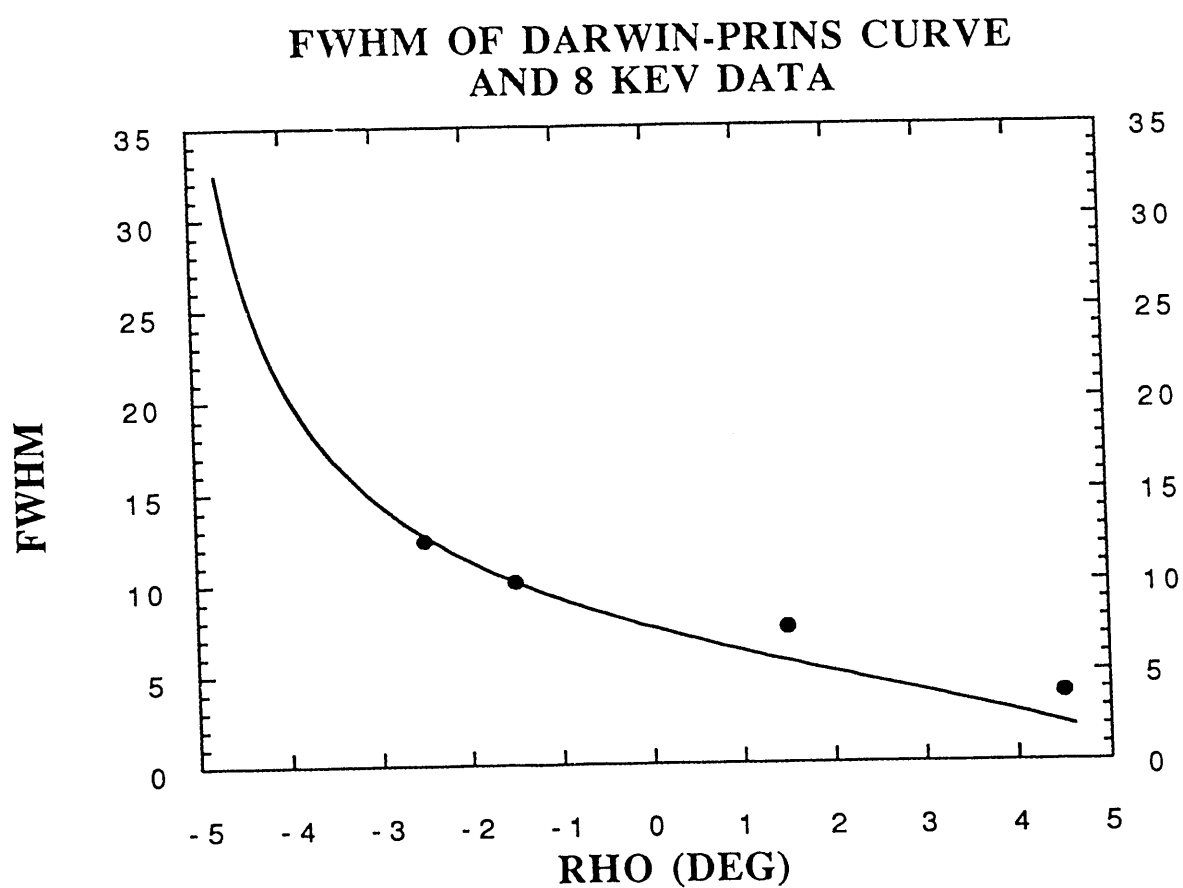
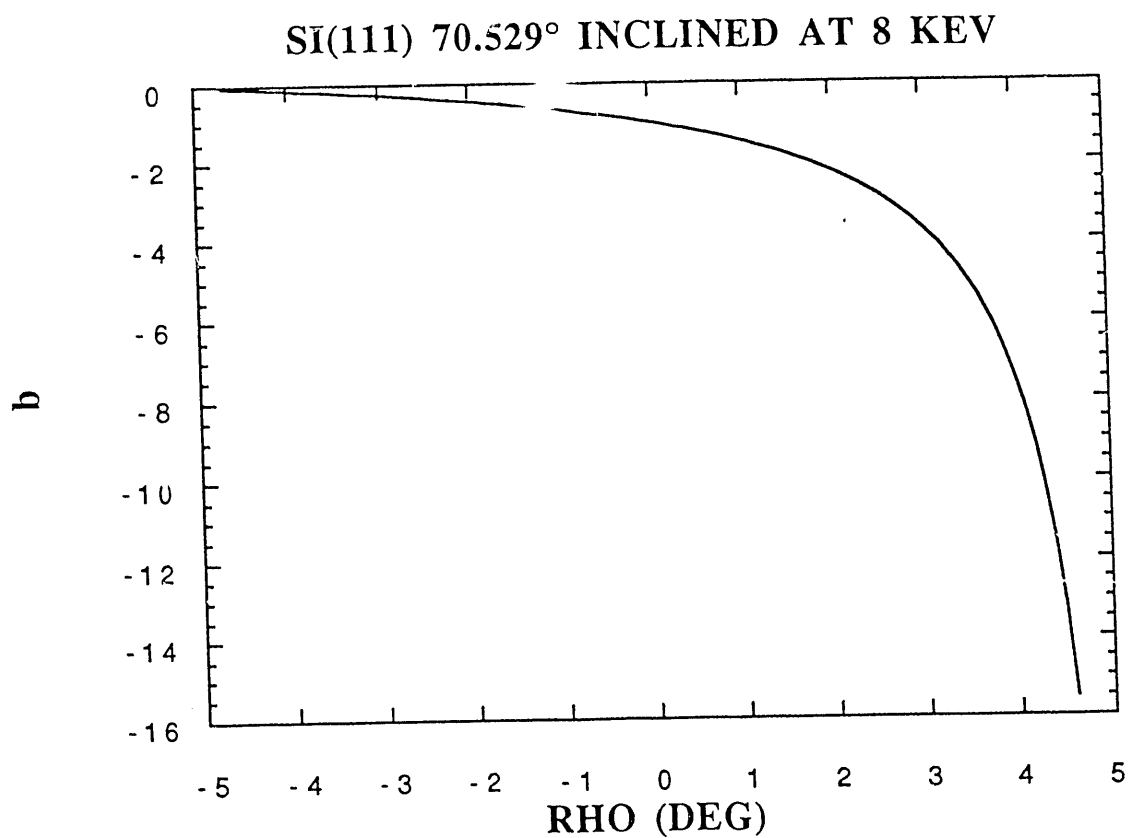


Fig.4

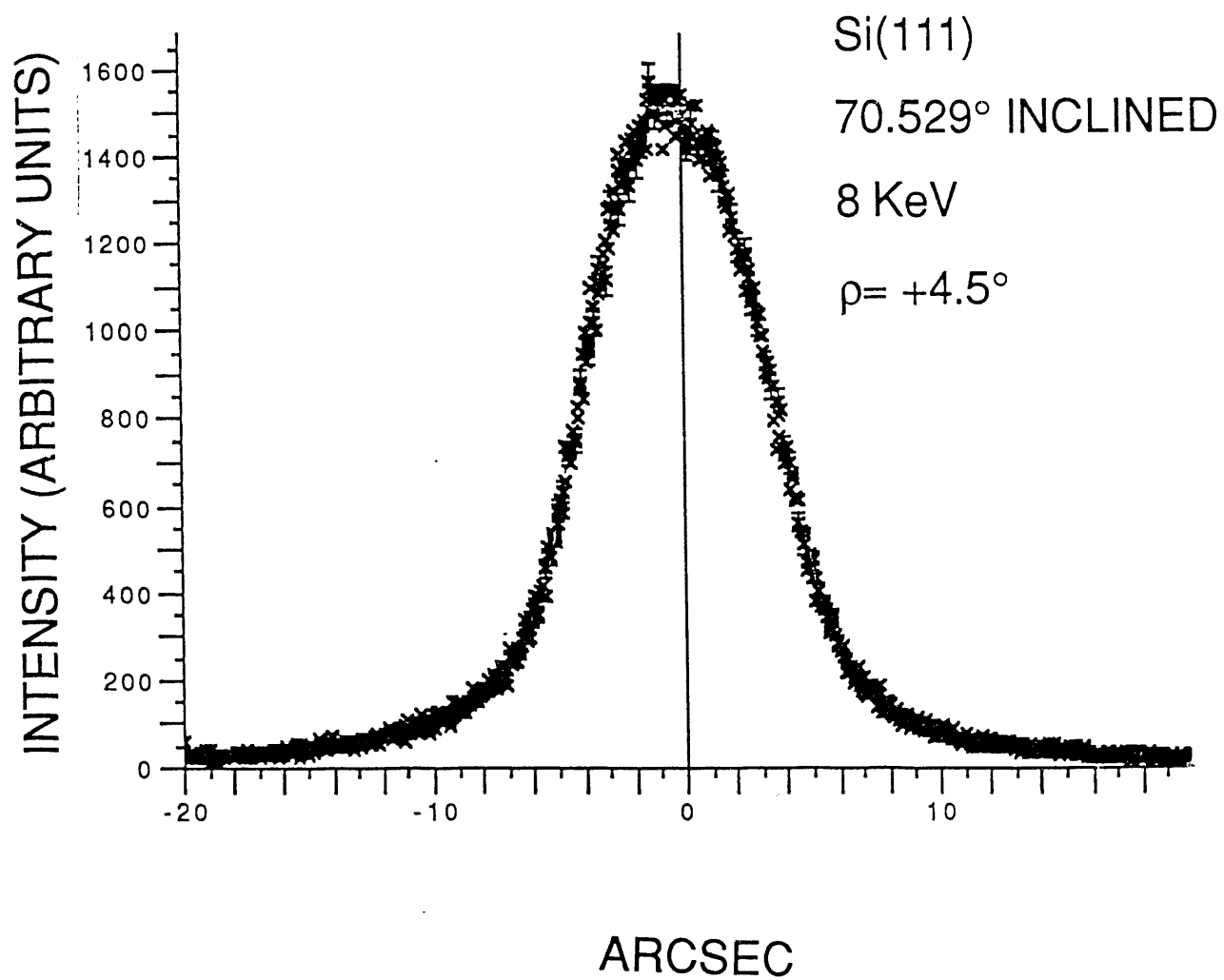


Fig. 5

END

DATE
FILMED
01/30/92

I