

## Letter

# First polarisation measurement of coherently photoproduced $J/\psi$ in ultra-peripheral Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

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## ABSTRACT

The first measurement of the polarisation of coherently photoproduced  $J/\psi$  mesons in ultra-peripheral Pb–Pb collisions, using data at  $\sqrt{s_{NN}} = 5.02$  TeV, is presented. The  $J/\psi$  meson is measured via its dimuon decay channel in the forward rapidity interval  $-4.0 < y < -2.5$  using the ALICE detector at the CERN LHC. An event sample corresponding to an integrated luminosity of  $750 \mu\text{b}^{-1} \pm 5\%$  (syst) is analysed. Hadronic activity is highly suppressed since the interaction is mediated by a photon. The polar and azimuthal angle distributions of the decay muons are measured, and the polarisation parameters  $\lambda_\theta$ ,  $\lambda_\phi$ ,  $\lambda_{\theta\phi}$  are extracted. The analysis is carried out in the helicity frame. The results are found to be consistent with a transversely polarised  $J/\psi$ . These values are compared with previous measurements by the H1 and ZEUS experiments. The polarisation parameters of coherent  $J/\psi$  photoproduction in Pb–Pb collisions are found to be consistent with the  $s$ -channel helicity conservation hypothesis.

## 1. Introduction

The study of vector meson photoproduction at high energies has received great interest in recent years [1–3]. Ultra-peripheral collisions (UPCs) at the CERN Large Hadron Collider (LHC) allow access to the highest available photon energies for such photon-induced interactions. In UPCs, the impact parameter is larger than the sum of the radii of the two colliding nuclei, so they are beyond the range of the nuclear force and do not interact hadronically. A quasi-real photon is emitted from one of the charged nuclei and this in turn might fluctuate into a quark–antiquark pair, which can scatter elastically off the other nucleus. Vector meson photoproduction off nuclei can be either coherent, with the photon coupling coherently to the nucleus as a whole, or incoherent, where instead the photon couples to a single nucleon. Much emphasis has been put on how this process could give access to space and momentum features of the momentum distribution of gluons of the hadronic target. In particular, several model calculations indicate that vector meson photoproduction is sensitive to nuclear modifications of the partonic structure for gluons that carry a rather small fraction of the momentum of the nucleons [4,5]. Kinematics dictates that the momentum fraction carried by the gluons is given by  $x = \frac{M_{J/\psi}}{\sqrt{s_{NN}}} e^{\pm y}$ , where  $M_{J/\psi}$  is the mass of the  $J/\psi$ ,  $y$  is the rapidity, and  $\sqrt{s_{NN}}$  is the centre-of-mass rapidity per nucleon–nucleon collision. For rapidities  $y \neq 0$  one thus has a two-fold ambiguity in the value of  $x$ . In the middle of the acceptance of the muon spectrom-

eter  $y = -3.25$ . This corresponds to  $x$ -values of  $2 \times 10^{-3}$  and  $5 \times 10^{-5}$ . The largest contribution to the cross section comes from the higher value. Recent experimental data from ALICE, CMS, and LHCb [6–10] have found experimental evidence of gluon shadowing [11].

In previous analyses of exclusive vector meson photoproduction off protons at lower energies, using both real and virtual photons, the vector meson has generally been found to have transverse polarisation, as would be expected if the vector meson keeps the polarisation of the incoming photon ( $s$ -channel helicity conservation, SCHC) [12,13]. However, although early results from SLAC [14] and from the CERN SPS [15] with  $\rho^0$  mesons produced from real photons are consistent with SCHC, a much more recent study of  $\rho^0$ ,  $\omega$ , and  $\phi$  photoproduction from CLAS [16–18] reports violation of SCHC.

Exclusive photoproduction of vector mesons has also been studied at HERA in ep collisions. Both the H1 [19] and ZEUS [20] collaborations observed transverse polarisation in  $J/\psi$  photoproduction, thus favouring SCHC, in their photoproduction data. Note that at HERA, photoproduction, at very low photon  $Q^2$  ( $Q^2 \lesssim 0.05 \text{ GeV}^2$ ), is distinguished from electroproduction at somewhat higher  $Q^2$ , even though both are photon–hadron interactions. Here  $Q^2$  is the four-momentum transfer squared from the lepton to the nucleon vertex. In electroproduction the photon is more virtual, and can be partially longitudinally polarised. Accordingly, in  $J/\psi$  electroproduction the degree of longitudinal polarisation is seen to increase with  $Q^2$  [21]. Photoproduction can also be studied in heavy-ion UPCs. Polarisation measurements have been re-

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ported by STAR for  $\rho^0$  photoproduction [22] in collisions of gold nuclei at RHIC, and are consistent with SCHC.

For all previous vector meson photoproduction analyses at the LHC, when calculating the experimental acceptance, the vector meson has been assumed to have the same helicity as that of the initial photon that interacted with the target, i.e. SCHC. Since there is no fundamental physics principle that justifies SCHC, new polarisation and spin parameter measurements of photoproduced vector mesons at high energies could provide important insights.

In this Letter, the first polarisation measurement of coherent  $J/\psi$  photoproduction in ultra-peripheral Pb–Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV is presented. These events are defined by the reaction  $\gamma + \text{Pb} \rightarrow J/\psi + \text{Pb}$ , where the photon is emitted from one of the two colliding lead nuclei, and the  $J/\psi$  meson is reconstructed in the forward rapidity interval  $-4.0 < y < -2.5$ .

## 2. Detector description

The ALICE detector and its performance are described in [23,24]. Muons from  $J/\psi$  decays are measured in the single-arm muon spectrometer. For coherent  $J/\psi$  photoproduction, the hadron activity is required to be highly suppressed. For this reason, vetoes are applied using the V0 and ALICE Diffractive (AD) detectors. The muon spectrometer covers the pseudorapidity interval  $-4.0 < \eta < -2.5$ . It consists of a ten interaction length absorber followed by five tracking stations, the third of which is placed inside a dipole magnet with a 3 Tm integrated magnetic field, a 7.2 interaction length iron wall, and a trigger system located downstream of the iron wall. Each tracking station is made of two planes of cathode pad chambers, while the trigger system consists of four planes of resistive plate chambers arranged in two stations. Muon tracks are reconstructed using the tracking algorithm described in [25]. The V0 detector is composed of the V0A and V0C sub-detectors, consisting of 32 cells each and covering the pseudorapidity intervals  $2.8 < \eta < 5.1$  and  $-3.7 < \eta < -1.7$ , respectively. The AD detector [26] is composed of the ADC and ADA sub-detectors with 4 cells each, located at  $-19.5$  m and  $+16.9$  m along the beam direction from the interaction point and covering the pseudorapidity intervals  $-7.0 < \eta < -4.9$  and  $4.7 < \eta < 6.3$ , respectively [27]. The V0 and AD detectors are scintillator arrays with a time resolution better than 1 ns, allowing a distinction between beam–beam and beam–gas interactions.

## 3. Data selection and analysis

This analysis follows the same selections as those recently described in [28], which combines the ultra-peripheral Pb–Pb datasets at centre-of-mass energy per nucleon–nucleon collision  $\sqrt{s_{NN}} = 5.02$  TeV collected in 2015 and 2018. The resulting dataset corresponds to an integrated luminosity of about  $750 \mu\text{b}^{-1} \pm 5\%$  (syst). The UPC event sample was recorded using a hardware trigger that requires at least two unlike-sign tracks in the muon spectrometer, and no activity in the V0A, ADA, and ADC detectors. The use of these forward detectors ensures the exclusivity condition. The single muon threshold for the transverse momentum  $p_T$  for the muon trigger system has been set to  $p_T = 1$  GeV/c [29]. At the offline level further selections are applied. First, exactly two unlike-sign tracks are required in the muon spectrometer, where both tracks must be compatible with high-quality reconstructed muons as described in [28]. The tracks must fulfil the requirement on the radial coordinate of the track at the end of the absorber, and on the extrapolation to the nominal vertex [30]. Track segments in the tracking chambers are matched with corresponding segments in the trigger chambers. In addition, each muon and each dimuon are required to be in the pseudorapidity  $-4.0 < \eta < -2.5$  and rapidity  $-4.0 < y < -2.5$  intervals, respectively. In order to select events where coherent  $J/\psi$  production is strongly enhanced, the dimuon  $p_T$  is required to be less than 0.25 GeV/c, effectively rejecting most of the non-coherent  $J/\psi$  events as discussed in [28]. Finally, the V0A, ADA, and ADC vetoes are applied, and up to

two hits are allowed in V0C, since the V0C acceptance overlaps with the that of the muon spectrometer. These offline level selections have wider time windows than those used for the trigger electronics, together with a more refined algorithm, allowing for a more precise rejection. Possible V0 and AD detector veto inefficiencies could originate from independent hadronic or electromagnetic pile-up interactions. The mean efficiency correction factor is 95%, accounting for V0A and AD veto inefficiencies, while the V0C veto is found to be fully efficient. Since the angular distributions studied in this analysis are not affected by the forward veto selections, only the muon-related efficiencies are considered.

The detector acceptance and efficiency ( $A \times \epsilon$ ) for  $J/\psi$  and for  $\psi'$  are evaluated using a sample of events generated with STARlight 2.2.0 [31], which models the  $J/\psi$  as transversely polarised. The decay muons are then propagated using a model of the detector implemented in GEANT 3.21 [32] with realistic conditions. They are then reconstructed with the same algorithm as that used for the experimental data. This sample is also used to construct the response matrices which are needed for this measurement. Typical values of the  $A \times \epsilon$  range from 0.41 for  $\cos \theta \sim 0$ , and down to 0.024 at the limit of the acceptance i.e.  $\cos \theta \sim 0.56$  (see next section for definition of  $\theta$ ).

## 4. Angular distributions

The angular distribution ( $W$ ) of the decay muons is described in terms of the three polarisation parameters  $\lambda_\theta$ ,  $\lambda_\phi$ , and  $\lambda_{\theta\phi}$  using the following expression as a function of the polar angle  $\theta$  and of the azimuthal angle  $\phi$ , which respects parity conservation [33]:

$$W(\cos \theta, \phi) \propto \frac{1}{3 + \lambda_\theta} \left[ 1 + \lambda_\theta \cos^2 \theta + \lambda_\phi \sin^2 \theta \cos 2\phi + \lambda_{\theta\phi} \sin 2\theta \cos \phi \right]. \quad (1)$$

It is also possible to study the data by using the following expressions

$$W(\cos \theta) \propto \frac{1}{3 + \lambda_\theta} \left[ 1 + \lambda_\theta \cos^2 \theta \right], \quad (2)$$

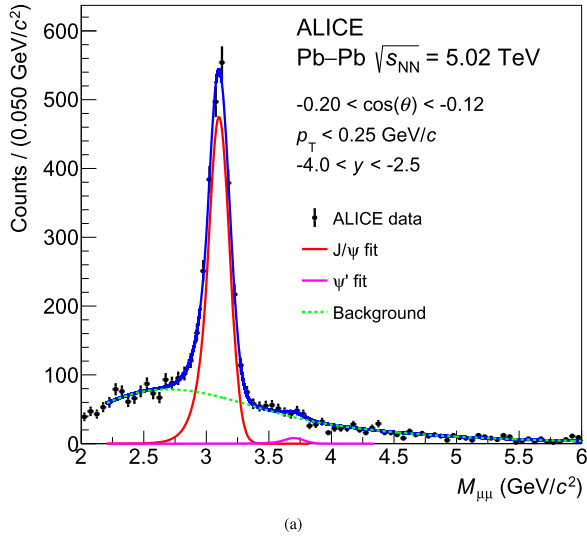
$$W(\phi) \propto 1 + \frac{2\lambda_\phi}{3 + \lambda_\theta} \cos 2\phi,$$

obtained by integrating Eq. (1) in the corresponding angular variable.

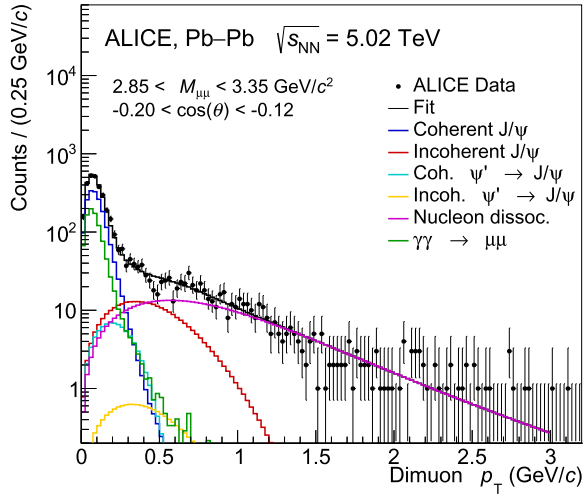
The polarisation parameters quantify the degree of polarisation. When all parameters are consistent with zero, an isotropic angular distribution is obtained. When  $(\lambda_\theta, \lambda_\phi, \lambda_{\theta\phi}) = (1, 0, 0)$  the vector meson is transversely polarised, while  $(\lambda_\theta, \lambda_\phi, \lambda_{\theta\phi}) = (-1, 0, 0)$  corresponds to a purely longitudinal polarisation.

The angular distribution of the decay muons is studied using the helicity frame. This frame has been used in recent studies of the polarisation of inclusive  $J/\psi$  production in pp and Pb–Pb collisions by the ALICE Collaboration [34–37], and by the PHENIX Collaboration in pp collisions [38]. In the helicity frame the z-axis coincides with the flight direction of the  $J/\psi$  in the Pb–Pb centre-of-mass frame. The y-axis is perpendicular to the plane formed by the collision axis and the direction of the  $J/\psi$  in the Pb–Pb centre-of-mass frame, while the x-axis is chosen so as to have a right-handed triplet [39]. The angles refer to the direction of the positive decay muon in this frame. Another commonly used definition of the axes is the Collins–Soper frame [33]. For  $J/\psi$  events at low  $p_T$  both frames are compatible, so results are only reported in terms of the helicity frame. The use of different reference frames has also been justified in terms of how the background could behave differently than the signal in the various frames. As discussed below, the background in this analysis is small which is expected for coherent  $J/\psi$  events.

To extract the  $J/\psi$  yield, the dimuon invariant mass distribution is fitted in the interval  $2.2 < M_{\mu\mu} < 6.0$  GeV/c<sup>2</sup>, requiring in addition that  $p_T < 0.25$  GeV/c to enhance the contribution from coherent production. The study is performed in twenty-four bins of  $\cos \theta$  over the interval  $-1 < \cos \theta < 1$ , and six bins in  $\phi$  in the interval between 0 and  $2\pi$ , to obtain yields in intervals of  $\cos \theta$  and  $\phi$ . As an example,



(a)



(b)

Fig. 1. A fit to the raw invariant mass (a) and  $p_T$  distributions for events in the  $-0.20 < \cos \theta < -0.12$  range evaluated in the helicity frame. The bottom panel (b) shows the dimuon  $p_T$  around the  $J/\psi$  mass range. The displayed uncertainties are statistical.

Fig. 1(a) shows the invariant mass distribution of coherent  $J/\psi$  candidates in the interval  $-0.20 < \cos \theta < -0.12$ . We follow the procedure described in [28] for fitting the invariant mass distributions. The data are fitted to a background model (dimuon continuum) and to a signal model ( $J/\psi$  and  $\psi'$  resonances). For the background, the fitting function is given by a fourth-order polynomial multiplied by an exponential. The exponential component becomes important for  $M_{\mu\mu} > 4.0 \text{ GeV}/c^2$ . The coefficients of the polynomial function are fixed to those obtained from Monte Carlo simulations using the STARlight [31] generator for the  $\gamma\gamma \rightarrow \mu^+\mu^-$  process. The exponential tail behaviour is expected for such an electromagnetic process. At low dimuon masses the detector acceptance decreases due to the trigger requirement, so the exponential part is not needed and a simple fourth-order polynomial is used. The  $J/\psi$  and  $\psi'$  signals are fitted with one-sided Crystal Ball functions [40]. The parameters for the Crystal Ball functions are fixed to those obtained from Monte Carlo templates. The mean and width of the  $J/\psi$  peak are free parameters in the fit. The mean of the  $\psi'$  peak is fixed to the world average value [41] and its width to that of the  $J/\psi$  multiplied by the ratio of the widths of the  $\psi'$  to the  $J/\psi$ , as measured from the analysis of reconstructed STARlight events. The ratio is approximately 1.09.

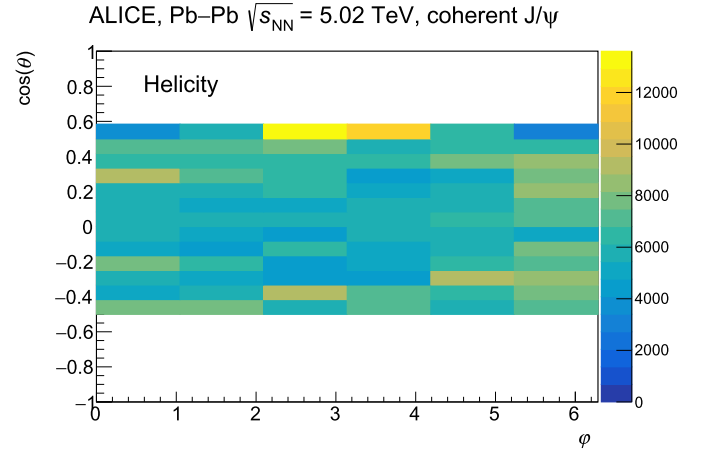


Fig. 2. Two-dimensional angular distribution of coherent  $J/\psi$  candidates in the helicity frame, after being unfolded in  $\varphi$  and corrected for  $A \times \epsilon$  in  $\cos \theta$ .

The  $J/\psi$  yield extracted from the fit to the invariant mass distributions comprises five contributions, namely (i) coherent production, (ii) incoherent production, and (iii) feed-down  $J/\psi$  production from coherent  $\psi'$  production, incoherent  $\psi'$  production, and (v) incoherent  $J/\psi$  production with associated nucleon production. To extract the coherent  $J/\psi$  yield, the residual background has been estimated from fits to the dimuon  $p_T$  distribution. This is possible since the various processes have different momentum distributions of the produced  $J/\psi$  mesons. Incoherent  $J/\psi$  photoproduction accompanied by nucleon dissociation was also taken into account in order to describe the high  $p_T$  tail, using a template based on the H1 parametrization of the dissociative  $J/\psi$  photoproduction [42]. Following this approach, which has been described in a recent publication [28], the background contribution is found to be smaller than 4.6% for the integrated  $\cos \theta$  and  $\varphi$  intervals used in this analysis. Note this fraction does not include the  $\gamma\gamma \rightarrow \mu^+\mu^-$  background which has been subtracted from fits to the invariant mass distribution, as described above. Since the residual background is small, and assuming that those  $J/\psi$  events have the same polarisation to that of coherent  $J/\psi$  events, their contributions are not subtracted from the sample. The  $J/\psi$  yield is estimated from the dimuon invariant mass distribution in a number of  $(\cos \theta, \varphi)$  cells assuming it all to be coherent. For illustration purposes, Fig. 1(b) shows the  $p_T$  distribution of candidates in the  $J/\psi$  invariant mass peak region ( $2.85 < M_{\mu\mu} < 3.35 \text{ GeV}/c^2$ ) and in the same  $\cos \theta$  interval of Fig. 1(a).

To correct for detector effects, the data are unfolded and corrected for  $A \times \epsilon$  in cells of  $\cos \theta$  and  $\varphi$ , and in the 1D distributions,  $A \times \epsilon$  corrections are applied to the data in  $\cos \theta$ . Since the STARlight Monte Carlo has been found to give a reasonable description of the data, in particular for coherent  $J/\psi$  production (low  $p_T$  events) at forward rapidity, we take this to be a natural starting point. The unfolding uses the Bayes' iterative method [43] as implemented in the RooUnfold framework [44]. The number of iterations in the unfolding procedure was set to fifteen. This number was found to be optimal at minimising the average uncertainty, following the prescription discussed in [45]. Since the generated  $\cos \theta$  is recovered in the reconstructed  $\cos \theta$  with a spread smaller than the bin width, bin migration effects are negligible and the data do not need unfolding in  $\cos \theta$ . Therefore, as in addition the sample size is small, the unfolding is performed only for  $\varphi$ , as described below. Fig. 2 shows the corrected  $J/\psi$  distribution in  $\cos \theta$  and  $\varphi$ , where the  $\varphi$  distribution in addition has been unfolded. The unfolding procedure increases the uncertainties on the estimated number of  $J/\psi$ s in each cell significantly. The first step is to fit the  $\cos \theta$  distribution as a 1D distribution, treating the uncertainties in each bin as independent of those of the other bins. The  $\cos \theta$  distribution is fitted using a  $\chi^2$  fitting method and the equivalent of Eq. (2) with the normalisation modified to take

**Table 1**

Summary of the systematic uncertainty contributions, presented as absolute values. The  $\cos \theta$  range systematic uncertainty refers to the fitted range variation, the signal extraction to the choice of the description of the  $J/\psi$ , the unfolding systematic uncertainty is due to the choice of the number of iterations, the response matrix refers to the input distribution in generating the matrix, and the trigger systematic uncertainty is associated to the single muon  $p_T$  selection used for the trigger efficiency calculation.

| Systematics                 | 1D               | 2D               |                   |                           |
|-----------------------------|------------------|------------------|-------------------|---------------------------|
|                             | $\lambda_\theta$ | $\lambda_\theta$ | $\lambda_\varphi$ | $\lambda_{\theta\varphi}$ |
| $\cos \theta$ range         | 0.079            | 0.142            | 0.002             | 0.056                     |
| signal extraction           | 0.002            | 0.026            | 0.002             | 0.008                     |
| unfolding                   | 0.0              | 0.019            | 0.004             | 0.004                     |
| response matrix             | 0.010            | 0.009            | 0.008             | 0.004                     |
| single muon $p_T$ threshold | 0.203            | 0.196            | 0.022             | 0.019                     |
| Total                       | 0.218            | 0.244            | 0.023             | 0.060                     |

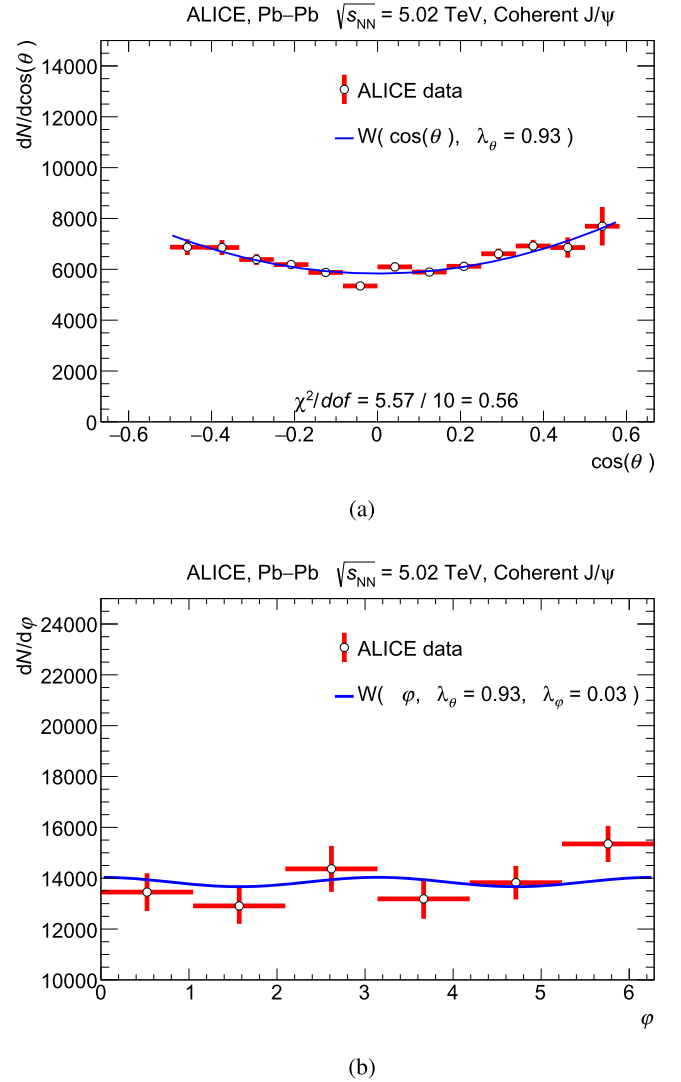
into account the actual limits of integration,  $-0.5 < \cos \theta < 0.583$ . This gives  $\lambda_\theta = 0.93 \pm 0.05(\text{stat}) \pm 0.22(\text{sys})$ .

For the  $\varphi$ -related parameters  $\lambda_\varphi$  and  $\lambda_{\theta\varphi}$ , the 2D distribution is required, giving the best estimates of the content of each  $(\cos \theta, \varphi)$  cell, where the  $\varphi$  distribution is unfolded independently in each  $\cos \theta$  band. As the contents of the unfolded cells for a given  $\cos \theta$  are correlated, an (input) covariance matrix was determined separately for each  $\cos \theta$  band. The contents of each cell and their uncertainty, as determined before any unfolding from the fit to the invariant mass distribution, are used to generate 10000 samples with fluctuating contents according to a normal distribution. The contents of all bins are varied simultaneously, and the unfolding algorithm applied. The covariance matrix for the cell contents after unfolding is evaluated using  $\text{cov}(x, y) = \langle xy \rangle - \langle x \rangle \langle y \rangle$ . Finally, the whole (2D) angular distribution is fitted with Eq. (1), using a  $\chi^2$  method, the unfolded contents for the central values in each cell, and the set of covariance matrices. A free fit for all three parameters yields  $\lambda_\theta = 0.82 \pm 0.44$ ,  $\lambda_\varphi = 0.03 \pm 0.04$ , and  $\lambda_{\theta\varphi} = 0.10 \pm 0.09$ . However, for the final fit the  $\lambda_\theta$  value has been fixed to the value obtained from the 1D fit, allowing only  $\lambda_\varphi$  and  $\lambda_{\theta\varphi}$  to vary. The  $\chi^2/NDF$  for this final fit is 114.2/175.

Table 1 presents a summary of the various sources of systematic uncertainty on the measurement of the polarisation parameters for coherent  $J/\psi$  mesons. To study how the polarisation parameters vary as the acceptance changes when  $\cos \theta$  increases, the polar angle distribution is fitted in the  $\cos \theta$  ranges  $[-0.5, 0.583]$  and  $[-0.5, 0.5]$ , as the analysis is carried out in a restricted  $\cos \theta$  interval. The observed variation is taken as a systematic uncertainty. The uncertainty from the signal extraction is estimated by changing the  $J/\psi$  line shape from a single-sided Crystal Ball function to a double-sided one. The background fraction is thus adjusted accordingly in the fit. The uncertainty from the unfolding is assessed by varying the number of iterations by  $\pm 1$  with respect to the number chosen for the central value. The impact of the response matrix as a source of systematic uncertainty in the results is also investigated. Three different matrices corresponding to transverse polarisation, flat distribution and longitudinal polarisation have been used, and the differences in the fitted parameters were evaluated to determine the corresponding systematic uncertainty. Another source of systematic uncertainty is the trigger and tracking efficiency. This is studied by varying the offline  $p_T$  selection from 0.85 to 1.15 GeV/c for single muons, corresponding to a  $\pm 5\%$  variation on the trigger efficiency. The choice of  $p_T$  selection and the  $\cos \theta$  range are correlated. The larger the  $p_T$ , the smaller the allowed  $\cos \theta$  range. Finally, the global systematic uncertainty on the polarisation parameters is obtained by adding the results of the various sources in quadrature.

## 5. Results and discussion

The projections of the angular distribution of the yield of coherently produced  $J/\psi$  in terms of  $\cos \theta$  and  $\varphi$  are shown in Figs. 3(a) and 3(b),



**Fig. 3.** Projections of the coherent  $J/\psi$  data points in (a)  $\cos \theta$  and (b)  $\varphi$ , for events in the  $-4 < y < -2.5$  rapidity interval. The curves are obtained by setting the lambda parameters in Eq. (2) to the results of the fit to the distribution given in Table 2. The uncertainties on the data points are statistical.

**Table 2**

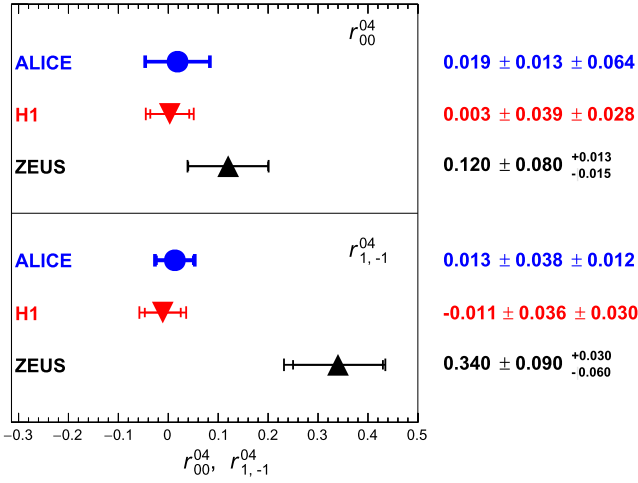
Measured polarisation parameters for coherent  $J/\psi$  in the helicity frame. The central values are given with statistical and total systematic uncertainties.

| $\lambda_\theta$         | $\lambda_\varphi$        | $\lambda_{\theta\varphi}$ |
|--------------------------|--------------------------|---------------------------|
| $0.93 \pm 0.05 \pm 0.22$ | $0.03 \pm 0.04 \pm 0.02$ | $0.10 \pm 0.09 \pm 0.06$  |

respectively. The numerical values of the polarisation parameters are given in Table 2. The numerical values of the polarisation parameters are given in Table 2, with the statistical uncertainty defined using  $\Delta\chi^2 = 2.3$ , since this is a two-dimensional distribution [46,47], and the systematic error calculated as described above. Note that the  $\lambda_\theta$  result is taken from the 1D fit. The results are found to be consistent with  $(\lambda_\theta, \lambda_\varphi, \lambda_{\theta\varphi}) = (1, 0, 0)$ , i.e. transverse polarisation, and they satisfy the required positivity constraints [48].

Alternative formalisms exist for describing the polarisation parameters [13,48,49]. In order to provide a comparison with H1 [19] and ZEUS [20] results, the measured  $\lambda_\theta$  and  $\lambda_\varphi$  values are converted to spin-density matrix elements  $r_{\lambda\lambda'}^{ik}$  [13]. The conversion from lambda parameters to spin-density matrix elements  $r_{\lambda\lambda'}^{ik}$  can be obtained from the following expressions [19]:





**Fig. 4.** ALICE results displayed in terms of spin density matrix elements, along with statistical and systematic uncertainties. A comparison with the available lowest  $Q^2$  ranges of the results from the H1 and ZEUS collaborations is also provided. H1 results [19] correspond to  $Q^2 \sim 0.05 \text{ GeV}^2$  consistent with photoproduction, while the shown ZEUS measurement [20] is for  $2 < Q^2 < 5 \text{ GeV}^2$ .

$$W(\cos \theta) \propto 1 + r_{00}^{04} + (1 - 3r_{00}^{04}) \cos^2 \theta$$

$$W(\varphi) \propto 1 + r_{1,-1}^{04} \cos 2\varphi.$$
(3)

By comparing Eq. (2) and Eq. (3), the spin-density matrix elements are expressed as:

$$r_{00}^{04} = \frac{1 - \lambda_\theta}{3 + \lambda_\theta}$$

$$r_{1,-1}^{04} = \frac{\lambda_\varphi}{2} (1 + r_{00}^{04}).$$
(4)

Fig. 4 shows the results of the comparison. In this analysis, the HERA measurements are given in the helicity frame. The ALICE result presented in Fig. 4 is compatible with the H1 measurement, which also uses a photoproduction sample ( $Q^2 \lesssim 0.05 \text{ GeV}^2$ ), and is lower than the ZEUS measurement, which uses a sample with  $2 < Q^2 < 5 \text{ GeV}^2$  [20]. This is the lowest  $Q^2$  point in the ZEUS range, and it can be seen to be tending towards the photoproduction value found by H1, but is already in the electroproduction regime.

## 6. Summary

The first measurement of the polarisation of coherent  $J/\psi$  photoproduction in ultra-peripheral Pb–Pb collisions, using data at  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ , has been carried out by the ALICE Collaboration. The  $\lambda_\theta$ ,  $\lambda_\varphi$ ,  $\lambda_{\theta\varphi}$  parameters have been evaluated in the helicity reference frame in the rapidity interval  $-4 < y < -2.5$ . The  $\lambda_\theta$  parameter has been found to be consistent with unity, and the other two parameters are found to be consistent with zero, indicating that coherently produced  $J/\psi$  mesons are transversely polarised as required for  $s$ -channel helicity conservation. The measured polarisation parameters are found to be compatible with those reported by the H1 Collaboration for  $J/\psi$  photoproduction in electron–proton interactions at lower energies, and evaluated in the helicity frame and using the spin density matrix formalism. This analysis represents the first experimental evidence that coherent  $J/\psi$  in ultra-peripheral Pb–Pb collisions at high energies is consistent with the photoproduction mechanism and the  $s$ -channel helicity conservation hypothesis.

## Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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## Data availability

Data will be made available on request.

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## ALICE Collaboration

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Caines <sup>137, [ib](#)</sup>, A. Caliva <sup>97, [ib](#)</sup>, E. Calvo Villar <sup>101, [ib](#)</sup>, J.M.M. Camacho <sup>109, [ib](#)</sup>, P. Camerini <sup>23, [ib](#)</sup>, F.D.M. Canedo <sup>110, [ib](#)</sup>, S.L. Cantway <sup>137, [ib](#)</sup>, M. Carabas <sup>113, [ib](#)</sup>, A.A. Carballo <sup>32, [ib](#)</sup>, F. Carnesecchi <sup>32, [ib](#)</sup>, R. Caron <sup>127, [ib](#)</sup>, L.A.D. Carvalho <sup>110, [ib](#)</sup>, J. Castillo Castellanos <sup>129, [ib](#)</sup>, F. Catalano <sup>24, [ib](#)</sup>, C. Ceballos Sanchez <sup>141, [ib](#)</sup>, I. Chakaberia <sup>74, [ib](#)</sup>, P. Chakraborty <sup>47, [ib](#)</sup>, S. Chandra <sup>134, [ib](#)</sup>, S. Chapeland <sup>32, [ib](#)</sup>, M. Chartier <sup>118, [ib](#)</sup>, S. Chattopadhyay <sup>134, [ib](#)</sup>, S. Chattopadhyay <sup>99, [ib](#)</sup>, T. Cheng <sup>97,6, [ib](#)</sup>, C. Cheshkov <sup>127, [ib](#)</sup>, B. Cheynis <sup>127, [ib](#)</sup>, V. 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Cosentino <sup>112, [ib](#)</sup>, F. Costa <sup>32, [ib](#)</sup>, S. Costanza <sup>21,55, [ib](#)</sup>, C. Cot <sup>130, [ib](#)</sup>, J. Crkovská <sup>94, [ib](#)</sup>, P. Crochet <sup>126, [ib](#)</sup>, R. Cruz-Torres <sup>74, [ib](#)</sup>, P. Cui <sup>6, [ib](#)</sup>, A. Dainese <sup>54, [ib](#)</sup>, M.C. Danisch <sup>94, [ib](#)</sup>, A. Danu <sup>63, [ib](#)</sup>, P. Das <sup>80, [ib](#)</sup>, P. Das <sup>4, [ib](#)</sup>, S. Das <sup>4, [ib](#)</sup>, A.R. Dash <sup>125, [ib](#)</sup>, S. Dash <sup>47, [ib](#)</sup>, A. De Caro <sup>28, [ib](#)</sup>, G. de Cataldo <sup>50, [ib](#)</sup>, J. de Cuveland <sup>38</sup>, A. De Falco <sup>22, [ib](#)</sup>, D. De Gruttola <sup>28, [ib](#)</sup>, N. De Marco <sup>56, [ib](#)</sup>, C. De Martin <sup>23, [ib](#)</sup>, S. De Pasquale <sup>28, [ib](#)</sup>, R. Deb <sup>133, [ib](#)</sup>, S. Deb <sup>48, [ib](#)</sup>, R.J. Debski <sup>2, [ib](#)</sup>, K.R. Deja <sup>135</sup>, R. Del Grande <sup>95, [ib](#)</sup>, L. Dello Stritto <sup>28, [ib](#)</sup>, W. Deng <sup>6, [ib](#)</sup>, P. Dhankher <sup>18, [ib](#)</sup>, D. Di Bari <sup>31, [ib](#)</sup>, A. Di Mauro <sup>32, [ib](#)</sup>, R.A. Diaz <sup>141,7, [ib](#)</sup>, T. Dietel <sup>114, [ib](#)</sup>, Y. Ding <sup>6, [ib](#)</sup>, R. Divià <sup>32, [ib](#)</sup>, D.U. Dixit <sup>18, [ib](#)</sup>, Ø. Djuvsland <sup>20</sup>, U. Dmitrieva <sup>140, [ib](#)</sup>, A. Dobrin <sup>63, [ib](#)</sup>, B. Dönigus <sup>64, [ib](#)</sup>, J.M. Dubinski <sup>135, [ib](#)</sup>, A. Dubla <sup>97, [ib](#)</sup>, S. Dudi <sup>90, [ib](#)</sup>, P. Dupieux <sup>126, [ib](#)</sup>, M. Durkac <sup>106</sup>, N. Dzalaiova <sup>12</sup>, T.M. Eder <sup>125, [ib](#)</sup>, R.J. Ehlers <sup>74, [ib](#)</sup>, V.N. Eikeland <sup>20</sup>, F. Eisenhut <sup>64, [ib](#)</sup>, D. Elia <sup>50, [ib](#)</sup>, B. Erasmus <sup>103, [ib](#)</sup>, F. Ercolessi <sup>25, [ib](#)</sup>, F. Erhardt <sup>89, [ib](#)</sup>,



M.R. Ersdal<sup>20</sup>, B. Espagnon<sup>130, <sup>id</sup></sup>, G. Eulisse<sup>32, <sup>id</sup></sup>, D. Evans<sup>100, <sup>id</sup></sup>, S. Evdokimov<sup>140, <sup>id</sup></sup>, L. Fabbietti<sup>95, <sup>id</sup></sup>,  
M. Faggin<sup>27, <sup>id</sup></sup>, J. Faivre<sup>73, <sup>id</sup></sup>, F. Fan<sup>6, <sup>id</sup></sup>, W. Fan<sup>74, <sup>id</sup></sup>, A. Fantoni<sup>49, <sup>id</sup></sup>, M. Fasel<sup>87, <sup>id</sup></sup>, P. Fedichio<sup>29</sup>,  
A. Feliciello<sup>56, <sup>id</sup></sup>, G. Feofilov<sup>140, <sup>id</sup></sup>, A. Fernández Téllez<sup>44, <sup>id</sup></sup>, L. Ferrandi<sup>110, <sup>id</sup></sup>, M.B. Ferrer<sup>32, <sup>id</sup></sup>,  
A. Ferrero<sup>129, <sup>id</sup></sup>, C. Ferrero<sup>56, <sup>id</sup>, <sup>IV</sup></sup>, A. Ferretti<sup>24, <sup>id</sup></sup>, V.J.G. Feuillard<sup>94, <sup>id</sup></sup>, V. Filova<sup>35, <sup>id</sup></sup>, D. Finogeev<sup>140, <sup>id</sup></sup>,  
F.M. Fionda<sup>52, <sup>id</sup></sup>, F. Flor<sup>115, <sup>id</sup></sup>, A.N. Flores<sup>108, <sup>id</sup></sup>, S. Foertsch<sup>68, <sup>id</sup></sup>, I. Fokin<sup>94, <sup>id</sup></sup>, S. Fokin<sup>140, <sup>id</sup></sup>,  
E. Fragiaco<sup>57, <sup>id</sup></sup>, E. Frajna<sup>46, <sup>id</sup></sup>, U. Fuchs<sup>32, <sup>id</sup></sup>, N. Funicello<sup>28, <sup>id</sup></sup>, C. Furget<sup>73, <sup>id</sup></sup>, A. Furs<sup>140, <sup>id</sup></sup>,  
T. Fusayasu<sup>98, <sup>id</sup></sup>, J.J. Gaardhøje<sup>83, <sup>id</sup></sup>, M. Gagliardi<sup>24, <sup>id</sup></sup>, A.M. Gago<sup>101, <sup>id</sup></sup>, C.D. Galvan<sup>109, <sup>id</sup></sup>,  
D.R. Gangadharan<sup>115, <sup>id</sup></sup>, P. Ganoti<sup>78, <sup>id</sup></sup>, C. Garabatos<sup>97, <sup>id</sup></sup>, T. García Chávez<sup>44, <sup>id</sup></sup>, E. Garcia-Solis<sup>9, <sup>id</sup></sup>,  
C. Gargiulo<sup>32, <sup>id</sup></sup>, K. Garner<sup>125</sup>, P. Gasik<sup>97, <sup>id</sup></sup>, A. Gautam<sup>117, <sup>id</sup></sup>, M.B. Gay Ducati<sup>66, <sup>id</sup></sup>, M. Germain<sup>103, <sup>id</sup></sup>,  
A. Ghimouz<sup>124</sup>, C. Ghosh<sup>134</sup>, M. Giacalone<sup>51, <sup>25</sup>, <sup>id</sup></sup>, P. Giubellino<sup>97, <sup>56</sup>, <sup>id</sup></sup>, P. Giubilato<sup>27, <sup>id</sup></sup>, A.M.C. Glaenger<sup>129, <sup>id</sup></sup>,  
P. Glässel<sup>94, <sup>id</sup></sup>, E. Glimos<sup>121, <sup>id</sup></sup>, D.J.Q. Goh<sup>76</sup>, V. Gonzalez<sup>136, <sup>id</sup></sup>, M. Gorgon<sup>2, <sup>id</sup></sup>, S. Gotovac<sup>33</sup>, V. Grabski<sup>67, <sup>id</sup></sup>,  
L.K. Graczykowski<sup>135, <sup>id</sup></sup>, E. Grecka<sup>86, <sup>id</sup></sup>, A. Grelli<sup>59, <sup>id</sup></sup>, C. Grigoras<sup>32, <sup>id</sup></sup>, V. Grigoriev<sup>140, <sup>id</sup></sup>, S. Grigoryan<sup>141, <sup>1</sup>, <sup>id</sup></sup>,  
F. Grosa<sup>32, <sup>id</sup></sup>, J.F. Grosse-Oetringhaus<sup>32, <sup>id</sup></sup>, R. Grosso<sup>97, <sup>id</sup></sup>, D. Grund<sup>35, <sup>id</sup></sup>, G.G. Guardiano<sup>111, <sup>id</sup></sup>,  
R. Guernane<sup>73, <sup>id</sup></sup>, M. Guilbaud<sup>103, <sup>id</sup></sup>, K. Gulbrandsen<sup>83, <sup>id</sup></sup>, T. Gündem<sup>64, <sup>id</sup></sup>, T. Gunji<sup>123, <sup>id</sup></sup>, W. Guo<sup>6, <sup>id</sup></sup>,  
A. Gupta<sup>91, <sup>id</sup></sup>, R. Gupta<sup>91, <sup>id</sup></sup>, R. Gupta<sup>48, <sup>id</sup></sup>, K. Gwizdzziel<sup>135, <sup>id</sup></sup>, L. Gyulai<sup>46, <sup>id</sup></sup>, M.K. Habib<sup>97</sup>,  
C. Hadjidakis<sup>130, <sup>id</sup></sup>, F.U. Haider<sup>91, <sup>id</sup></sup>, H. Hamagaki<sup>76, <sup>id</sup></sup>, A. Hamdi<sup>74, <sup>id</sup></sup>, M. Hamid<sup>6</sup>, Y. Han<sup>138, <sup>id</sup></sup>,  
R. Hannigan<sup>108, <sup>id</sup></sup>, M.R. Haque<sup>135, <sup>id</sup></sup>, J.W. Harris<sup>137, <sup>id</sup></sup>, A. Harton<sup>9, <sup>id</sup></sup>, H. Hassan<sup>87, <sup>id</sup></sup>, D. Hatzifotiadiou<sup>51, <sup>id</sup></sup>,  
P. Hauer<sup>42, <sup>id</sup></sup>, L.B. Havener<sup>137, <sup>id</sup></sup>, S.T. Heckel<sup>95, <sup>id</sup></sup>, E. Hellbär<sup>97, <sup>id</sup></sup>, H. Helstrup<sup>34, <sup>id</sup></sup>, M. Hemmer<sup>64, <sup>id</sup></sup>,  
T. Herman<sup>35, <sup>id</sup></sup>, G. Herrera Corral<sup>8, <sup>id</sup></sup>, F. Herrmann<sup>125</sup>, S. Herrmann<sup>127, <sup>id</sup></sup>, K.F. Hetland<sup>34, <sup>id</sup></sup>, B. Heybeck<sup>64, <sup>id</sup></sup>,  
H. Hillemanns<sup>32, <sup>id</sup></sup>, B. Hippolyte<sup>128, <sup>id</sup></sup>, F.W. Hoffmann<sup>70, <sup>id</sup></sup>, B. Hofman<sup>59, <sup>id</sup></sup>, B. Hohlweger<sup>84, <sup>id</sup></sup>,  
G.H. Hong<sup>138, <sup>id</sup></sup>, M. Horst<sup>95, <sup>id</sup></sup>, A. Horzyk<sup>2, <sup>id</sup></sup>, Y. Hou<sup>6, <sup>id</sup></sup>, P. Hristov<sup>32, <sup>id</sup></sup>, C. Hughes<sup>121, <sup>id</sup></sup>, P. Huhn<sup>64</sup>,  
L.M. Huhta<sup>116, <sup>id</sup></sup>, T.J. Humanic<sup>88, <sup>id</sup></sup>, A. Hutson<sup>115, <sup>id</sup></sup>, D. Hutter<sup>38, <sup>id</sup></sup>, J.P. Iddon<sup>118, <sup>id</sup></sup>, R. Ilkaev<sup>140</sup>,  
H. Ilyas<sup>13, <sup>id</sup></sup>, M. Inaba<sup>124, <sup>id</sup></sup>, G.M. Innocenti<sup>32, <sup>id</sup></sup>, M. Ippolitov<sup>140, <sup>id</sup></sup>, A. Isakov<sup>86, <sup>id</sup></sup>, T. Isidori<sup>117, <sup>id</sup></sup>,  
M.S. Islam<sup>99, <sup>id</sup></sup>, M. Ivanov<sup>97, <sup>id</sup></sup>, M. Ivanov<sup>12</sup>, V. Ivanov<sup>140, <sup>id</sup></sup>, M. Jablonski<sup>2, <sup>id</sup></sup>, B. Jacak<sup>74, <sup>id</sup></sup>, N. Jacazio<sup>32, <sup>id</sup></sup>,  
P.M. Jacobs<sup>74, <sup>id</sup></sup>, S. Jadlovská<sup>106</sup>, J. Jadlovsky<sup>106</sup>, S. Jaelani<sup>82, <sup>id</sup></sup>, L. Jaffe<sup>38</sup>, C. Jahnke<sup>111, <sup>id</sup></sup>,  
M.J. Jakubowska<sup>135, <sup>id</sup></sup>, M.A. Janik<sup>135, <sup>id</sup></sup>, T. Janson<sup>70</sup>, M. Jercic<sup>89</sup>, S. Jia<sup>10, <sup>id</sup></sup>, A.A.P. Jimenez<sup>65, <sup>id</sup></sup>,  
F. Jonas<sup>87, <sup>125</sup>, <sup>id</sup></sup>, J.M. Jowett<sup>32, <sup>97</sup>, <sup>id</sup></sup>, J. Jung<sup>64, <sup>id</sup></sup>, M. Jung<sup>64, <sup>id</sup></sup>, A. Junique<sup>32, <sup>id</sup></sup>, A. Jusko<sup>100, <sup>id</sup></sup>, J. Kaewjai<sup>105</sup>,  
P. Kalinak<sup>60, <sup>id</sup></sup>, A.S. Kalteyer<sup>97, <sup>id</sup></sup>, A. Kalweit<sup>32, <sup>id</sup></sup>, V. Kaplin<sup>140, <sup>id</sup></sup>, A. Karasu Uysal<sup>72, <sup>id</sup>, <sup>V</sup></sup>, D. Karatovic<sup>89, <sup>id</sup></sup>,  
O. Karavichev<sup>140, <sup>id</sup></sup>, T. Karavicheva<sup>140, <sup>id</sup></sup>, P. Karczmarczyk<sup>135, <sup>id</sup></sup>, E. Karpechev<sup>140, <sup>id</sup></sup>, M.J. Karwowska<sup>32, <sup>135</sup>, <sup>id</sup></sup>,  
U. Kebschull<sup>70, <sup>id</sup></sup>, R. Keidel<sup>139, <sup>id</sup></sup>, D.L.D. Keijdener<sup>59</sup>, M. Keil<sup>32, <sup>id</sup></sup>, B. Ketzer<sup>42, <sup>id</sup></sup>, S.S. Khade<sup>48, <sup>id</sup></sup>,  
A.M. Khan<sup>6, <sup>id</sup></sup>, S. Khan<sup>15, <sup>id</sup></sup>, A. Khanzadeev<sup>140, <sup>id</sup></sup>, Y. Kharlov<sup>140, <sup>id</sup></sup>, A. Khatun<sup>117, <sup>15</sup>, <sup>id</sup></sup>, A. Khuntia<sup>107, <sup>id</sup></sup>,  
M.B. Kidson<sup>114</sup>, B. Kileng<sup>34, <sup>id</sup></sup>, B. Kim<sup>104, <sup>id</sup></sup>, C. Kim<sup>16, <sup>id</sup></sup>, D.J. Kim<sup>116, <sup>id</sup></sup>, E.J. Kim<sup>69, <sup>id</sup></sup>, J. Kim<sup>138, <sup>id</sup></sup>,  
J.S. Kim<sup>40, <sup>id</sup></sup>, J. Kim<sup>69, <sup>id</sup></sup>, M. Kim<sup>18, <sup>94</sup>, <sup>id</sup></sup>, S. Kim<sup>17, <sup>id</sup></sup>, T. Kim<sup>138, <sup>id</sup></sup>, K. Kimura<sup>92, <sup>id</sup></sup>, S. Kirsch<sup>64, <sup>id</sup></sup>, I. Kisel<sup>38, <sup>id</sup></sup>,  
S. Kiselev<sup>140, <sup>id</sup></sup>, A. Kisiel<sup>135, <sup>id</sup></sup>, J.P. Kitowski<sup>2, <sup>id</sup></sup>, J.L. Klay<sup>5, <sup>id</sup></sup>, J. Klein<sup>32, <sup>id</sup></sup>, S. Klein<sup>74, <sup>id</sup></sup>, C. Klein-Bösing<sup>125, <sup>id</sup></sup>,  
M. Kleiner<sup>64, <sup>id</sup></sup>, T. Klemenz<sup>95, <sup>id</sup></sup>, A. Kluge<sup>32, <sup>id</sup></sup>, A.G. Knospe<sup>115, <sup>id</sup></sup>, C. Kobdaj<sup>105, <sup>id</sup></sup>, T. Kollegger<sup>97</sup>,  
A. Kondratyev<sup>141, <sup>id</sup></sup>, N. Kondratyeva<sup>140, <sup>id</sup></sup>, E. Kondratyuk<sup>140, <sup>id</sup></sup>, J. Konig<sup>64, <sup>id</sup></sup>, S.A. Konigstorfer<sup>95, <sup>id</sup></sup>,  
P.J. Konopka<sup>32, <sup>id</sup></sup>, G. Kornakov<sup>135, <sup>id</sup></sup>, M. Korwieser<sup>95, <sup>id</sup></sup>, S.D. Koryciak<sup>2, <sup>id</sup></sup>, A. Kotliarov<sup>86, <sup>id</sup></sup>,  
V. Kovalenko<sup>140, <sup>id</sup></sup>, M. Kowalski<sup>107, <sup>id</sup></sup>, V. Kozuharov<sup>36, <sup>id</sup></sup>, I. Králik<sup>60, <sup>id</sup></sup>, A. Kravčáková<sup>37, <sup>id</sup></sup>, L. Krcal<sup>32, <sup>38</sup>, <sup>id</sup></sup>,  
L. Kreis<sup>97</sup>, M. Krivda<sup>100, <sup>60</sup>, <sup>id</sup></sup>, F. Krizek<sup>86, <sup>id</sup></sup>, K. Krizkova Gajdosova<sup>32, <sup>id</sup></sup>, M. Kroesen<sup>94, <sup>id</sup></sup>, M. Krüger<sup>64, <sup>id</sup></sup>,  
D.M. Krupova<sup>35, <sup>id</sup></sup>, E. Kryshen<sup>140, <sup>id</sup></sup>, V. Kučera<sup>32, <sup>id</sup></sup>, C. Kuhn<sup>128, <sup>id</sup></sup>, P.G. Kuijper<sup>84, <sup>id</sup></sup>, T. Kumaoka<sup>124</sup>,  
D. Kumar<sup>134</sup>, L. Kumar<sup>90, <sup>id</sup></sup>, N. Kumar<sup>90</sup>, S. Kumar<sup>31, <sup>id</sup></sup>, S. Kundu<sup>32, <sup>id</sup></sup>, P. Kurashvili<sup>79, <sup>id</sup></sup>, A. Kurepin<sup>140, <sup>id</sup></sup>,  
A.B. Kurepin<sup>140, <sup>id</sup></sup>, A. Kuryakin<sup>140, <sup>id</sup></sup>, S. Kushpil<sup>86, <sup>id</sup></sup>, J. Kvapil<sup>100, <sup>id</sup></sup>, M.J. Kweon<sup>58, <sup>id</sup></sup>, J.Y. Kwon<sup>58, <sup>id</sup></sup>,



Y. Kwon <sup>138, [ib](#)</sup>, S.L. La Pointe <sup>38, [ib](#)</sup>, P. La Rocca <sup>26, [ib](#)</sup>, A. Lakrathok <sup>105</sup>, M. Lamanna <sup>32, [ib](#)</sup>, R. Langoy <sup>120, [ib](#)</sup>, P. Larionov <sup>32, [ib](#)</sup>, E. Laudi <sup>32, [ib](#)</sup>, L. Lautner <sup>32,95, [ib](#)</sup>, R. Lavicka <sup>102, [ib](#)</sup>, T. Lazareva <sup>140, [ib](#)</sup>, R. Lea <sup>133,55, [ib](#)</sup>, H. Lee <sup>104, [ib](#)</sup>, G. Legras <sup>125, [ib](#)</sup>, J. Lehrbach <sup>38, [ib](#)</sup>, T.M. Lelek <sup>2</sup>, R.C. Lemmon <sup>85, [ib](#)</sup>, I. León Monzón <sup>109, [ib](#)</sup>, M.M. Lesch <sup>95, [ib](#)</sup>, E.D. Lesser <sup>18, [ib](#)</sup>, P. Lévai <sup>46, [ib](#)</sup>, X. Li <sup>10</sup>, X.L. Li <sup>6</sup>, J. Lien <sup>120, [ib](#)</sup>, R. Lietava <sup>100, [ib](#)</sup>, I. Likmeta <sup>115, [ib](#)</sup>, B. Lim <sup>24, [ib](#)</sup>, S.H. Lim <sup>16, [ib](#)</sup>, V. Lindenstruth <sup>38, [ib](#)</sup>, A. Lindner <sup>45</sup>, C. Lippmann <sup>97, [ib](#)</sup>, A. Liu <sup>18, [ib](#)</sup>, D.H. Liu <sup>6, [ib](#)</sup>, J. Liu <sup>118, [ib](#)</sup>, I.M. Lofnes <sup>20, [ib](#)</sup>, C. Loizides <sup>87, [ib](#)</sup>, S. Lokos <sup>107, [ib](#)</sup>, J. Lömker <sup>59, [ib](#)</sup>, P. Loncar <sup>33, [ib](#)</sup>, J.A. Lopez <sup>94, [ib](#)</sup>, X. Lopez <sup>126, [ib](#)</sup>, E. López Torres <sup>7, [ib](#)</sup>, P. Lu <sup>97,119, [ib](#)</sup>, J.R. Luhder <sup>125, [ib](#)</sup>, M. Lunardon <sup>27, [ib](#)</sup>, G. Luparello <sup>57, [ib](#)</sup>, Y.G. Ma <sup>39, [ib](#)</sup>, A. Maevskaya <sup>140</sup>, M. Mager <sup>32, [ib](#)</sup>, A. Maire <sup>128, [ib](#)</sup>, M.V. Makariev <sup>36, [ib](#)</sup>, M. Malaev <sup>140, [ib](#)</sup>, G. Malfattore <sup>25, [ib](#)</sup>, N.M. Malik <sup>91, [ib](#)</sup>, Q.W. Malik <sup>19</sup>, S.K. Malik <sup>91, [ib](#)</sup>, L. Malinina <sup>141, [ib](#)</sup>, I.VIII, D. Mal'Kevich <sup>140, [ib](#)</sup>, D. Mallick <sup>80, [ib](#)</sup>, N. Mallick <sup>48, [ib](#)</sup>, G. Mandaglio <sup>30,53, [ib](#)</sup>, S.K. Mandal <sup>79, [ib](#)</sup>, V. Manko <sup>140, [ib](#)</sup>, F. Manso <sup>126, [ib](#)</sup>, V. Manzari <sup>50, [ib](#)</sup>, Y. Mao <sup>6, [ib](#)</sup>, G.V. Margagliotti <sup>23, [ib](#)</sup>, A. Margotti <sup>51, [ib](#)</sup>, A. Marín <sup>97, [ib](#)</sup>, C. Markert <sup>108, [ib](#)</sup>, P. Martinengo <sup>32, [ib](#)</sup>, J.L. Martinez <sup>115</sup>, M.I. Martínez <sup>44, [ib](#)</sup>, G. Martínez García <sup>103, [ib](#)</sup>, S. Masciocchi <sup>97, [ib](#)</sup>, M. Masera <sup>24, [ib](#)</sup>, A. Masoni <sup>52, [ib](#)</sup>, L. Massacrier <sup>130, [ib](#)</sup>, A. Mastroserio <sup>131,50, [ib](#)</sup>, O. Matonoha <sup>75, [ib](#)</sup>, P.F.T. Matuoka <sup>110</sup>, A. Matyja <sup>107, [ib](#)</sup>, C. Mayer <sup>107, [ib](#)</sup>, A.L. Mazuecos <sup>32, [ib](#)</sup>, F. Mazzaschi <sup>24, [ib](#)</sup>, M. Mazzilli <sup>32, [ib](#)</sup>, J.E. Mdhuli <sup>122, [ib](#)</sup>, A.F. Mechler <sup>64</sup>, Y. Melikyan <sup>43,140, [ib](#)</sup>, A. Menchaca-Rocha <sup>67, [ib](#)</sup>, E. Meninno <sup>102, [ib](#)</sup>, A.S. Menon <sup>115, [ib](#)</sup>, M. Meres <sup>12, [ib](#)</sup>, S. Mhlanga <sup>114,68</sup>, Y. Miake <sup>124</sup>, L. Micheletti <sup>56, [ib](#)</sup>, L.C. Migliorin <sup>127</sup>, D.L. Mihaylov <sup>95, [ib](#)</sup>, K. Mikhaylov <sup>141,140, [ib](#)</sup>, A.N. Mishra <sup>46, [ib](#)</sup>, D. Miśkowiec <sup>97, [ib](#)</sup>, A. Modak <sup>4, [ib](#)</sup>, A.P. Mohanty <sup>59, [ib](#)</sup>, B. Mohanty <sup>80, [ib](#)</sup>, M. Mohisin Khan <sup>15, [ib](#)</sup>, VI, M.A. Molander <sup>43, [ib](#)</sup>, Z. Moravcova <sup>83, [ib](#)</sup>, C. Mordasini <sup>95, [ib](#)</sup>, D.A. Moreira De Godoy <sup>125, [ib](#)</sup>, I. Morozov <sup>140, [ib](#)</sup>, A. Morsch <sup>32, [ib](#)</sup>, T. Mrnjavac <sup>32, [ib](#)</sup>, V. Muccifora <sup>49, [ib](#)</sup>, S. Muhuri <sup>134, [ib](#)</sup>, J.D. Mulligan <sup>74, [ib](#)</sup>, A. Mulliri <sup>22, [ib](#)</sup>, M.G. Munhoz <sup>110, [ib](#)</sup>, R.H. Munzer <sup>64, [ib](#)</sup>, H. Murakami <sup>123, [ib](#)</sup>, S. Murray <sup>114, [ib](#)</sup>, L. Musa <sup>32, [ib](#)</sup>, J. Musinsky <sup>60, [ib](#)</sup>, J.W. Myrcha <sup>135, [ib](#)</sup>, B. Naik <sup>122, [ib](#)</sup>, A.I. Nambrath <sup>18, [ib](#)</sup>, B.K. Nandi <sup>47, [ib](#)</sup>, R. Nania <sup>51, [ib](#)</sup>, E. Nappi <sup>50, [ib](#)</sup>, A.F. Nassirpour <sup>17,75, [ib](#)</sup>, A. Nath <sup>94, [ib](#)</sup>, C. Nattrass <sup>121, [ib](#)</sup>, M.N. Naydenov <sup>36, [ib](#)</sup>, A. Neagu <sup>19</sup>, A. Negru <sup>113</sup>, L. Nellen <sup>65, [ib](#)</sup>, S.V. Nesbo <sup>34</sup>, G. Neskovic <sup>38, [ib](#)</sup>, D. Nesterov <sup>140, [ib](#)</sup>, B.S. Nielsen <sup>83, [ib](#)</sup>, E.G. Nielsen <sup>83, [ib](#)</sup>, S. Nikolaev <sup>140, [ib](#)</sup>, S. Nikulin <sup>140, [ib](#)</sup>, V. Nikulin <sup>140, [ib](#)</sup>, F. Noferini <sup>51, [ib](#)</sup>, S. Noh <sup>11, [ib](#)</sup>, P. Nomokonov <sup>141, [ib](#)</sup>, J. Norman <sup>118, [ib](#)</sup>, N. Novitzky <sup>124, [ib](#)</sup>, P. Nowakowski <sup>135, [ib](#)</sup>, A. Nyanin <sup>140, [ib](#)</sup>, J. Nystrand <sup>20, [ib](#)</sup>, M. Ogino <sup>76, [ib](#)</sup>, A. Ohlson <sup>75, [ib](#)</sup>, V.A. Okorokov <sup>140, [ib](#)</sup>, J. Oleniacz <sup>135, [ib](#)</sup>, A.C. Oliveira Da Silva <sup>121, [ib](#)</sup>, M.H. Oliver <sup>137, [ib](#)</sup>, A. Onnerstad <sup>116, [ib](#)</sup>, C. Oppedisano <sup>56, [ib](#)</sup>, A. Ortiz Velasquez <sup>65, [ib](#)</sup>, J. Otwinowski <sup>107, [ib](#)</sup>, M. Oya <sup>92</sup>, K. Oyama <sup>76, [ib](#)</sup>, Y. Pachmayer <sup>94, [ib](#)</sup>, S. Padhan <sup>47, [ib](#)</sup>, D. Pagano <sup>133,55, [ib](#)</sup>, G. Paić <sup>65, [ib](#)</sup>, S. Paisano-Guzmán <sup>44, [ib](#)</sup>, A. Palasciano <sup>50, [ib](#)</sup>, S. Panebianco <sup>129, [ib](#)</sup>, H. Park <sup>124, [ib](#)</sup>, H. Park <sup>104, [ib](#)</sup>, J. Park <sup>58, [ib](#)</sup>, J.E. Parkkila <sup>32, [ib](#)</sup>, R.N. Patra <sup>91</sup>, B. Paul <sup>22, [ib](#)</sup>, H. Pei <sup>6, [ib](#)</sup>, T. Peitzmann <sup>59, [ib](#)</sup>, X. Peng <sup>6, [ib](#)</sup>, M. Pennisi <sup>24, [ib](#)</sup>, L.G. Pereira <sup>66, [ib](#)</sup>, D. Peresunko <sup>140, [ib](#)</sup>, G.M. Perez <sup>7, [ib](#)</sup>, S. Perrin <sup>129, [ib](#)</sup>, Y. Pestov <sup>140</sup>, V. Petráček <sup>35, [ib](#)</sup>, V. Petrov <sup>140, [ib](#)</sup>, M. Petrovici <sup>45, [ib](#)</sup>, R.P. Pezzi <sup>103,66, [ib](#)</sup>, S. Piano <sup>57, [ib](#)</sup>, M. Pikna <sup>12, [ib](#)</sup>, P. Pillot <sup>103, [ib](#)</sup>, O. Pinazza <sup>51,32, [ib](#)</sup>, L. Pinsky <sup>115</sup>, C. Pinto <sup>95, [ib](#)</sup>, S. Pisano <sup>49, [ib](#)</sup>, M. Płoskoń <sup>74, [ib](#)</sup>, M. Planinic <sup>89</sup>, F. Pliquet <sup>64</sup>, M.G. Poghosyan <sup>87, [ib](#)</sup>, B. Polichtchouk <sup>140, [ib](#)</sup>, S. Politano <sup>29, [ib](#)</sup>, N. Poljak <sup>89, [ib](#)</sup>, A. Pop <sup>45, [ib](#)</sup>, S. Porteboeuf-Houssais <sup>126, [ib](#)</sup>, V. Pozdniakov <sup>141, [ib](#)</sup>, I.Y. Pozos <sup>44, [ib](#)</sup>, K.K. Pradhan <sup>48, [ib](#)</sup>, S.K. Prasad <sup>4, [ib](#)</sup>, S. Prasad <sup>48, [ib](#)</sup>, R. Preghenella <sup>51, [ib](#)</sup>, F. Prino <sup>56, [ib](#)</sup>, C.A. Pruneau <sup>136, [ib](#)</sup>, I. Pshenichnov <sup>140, [ib](#)</sup>, M. Puccio <sup>32, [ib](#)</sup>, S. Pucillo <sup>24, [ib](#)</sup>, Z. Pugelova <sup>106</sup>, S. Qiu <sup>84, [ib](#)</sup>, L. Quaglia <sup>24, [ib](#)</sup>, R.E. Quishpe <sup>115</sup>, S. Ragoni <sup>14, [ib](#)</sup>, A. Rakotozafindrabe <sup>129, [ib](#)</sup>, L. Ramello <sup>132,56, [ib](#)</sup>, F. Rami <sup>128, [ib](#)</sup>, T.A. Rancien <sup>73</sup>, M. Rasa <sup>26, [ib](#)</sup>, S.S. Räsänen <sup>43, [ib](#)</sup>, R. Rath <sup>51, [ib](#)</sup>, M.P. Rauch <sup>20, [ib](#)</sup>, I. Ravasenga <sup>84, [ib](#)</sup>, K.F. Read <sup>87,121, [ib](#)</sup>, C. Reckziegel <sup>112, [ib](#)</sup>, A.R. Redelbach <sup>38, [ib](#)</sup>, K. Redlich <sup>79, [ib](#)</sup>, VII, C.A. Reetz <sup>97, [ib](#)</sup>, H.D. Regules-Medel <sup>44</sup>, A. Rehman <sup>20</sup>, F. Reidt <sup>32, [ib](#)</sup>, H.A. Reme-Ness <sup>34, [ib](#)</sup>, Z. Rescakova <sup>37</sup>, K. Reygers <sup>94, [ib](#)</sup>, A. Riabov <sup>140, [ib](#)</sup>, V. Riabov <sup>140, [ib](#)</sup>, R. Ricci <sup>28, [ib](#)</sup>, M. Richter <sup>19, [ib](#)</sup>, A.A. Riedel <sup>95, [ib](#)</sup>, W. Riegler <sup>32, [ib](#)</sup>, C. Ristea <sup>63, [ib](#)</sup>, M. Rodríguez Cahuantzi <sup>44, [ib](#)</sup>, S.A. Rodríguez Ramírez <sup>44, [ib](#)</sup>,

K. Røed<sup>19, [id](#)</sup>, R. Rogalev<sup>140, [id](#)</sup>, E. Rogochaya<sup>141, [id](#)</sup>, T.S. Rogoschinski<sup>64, [id](#)</sup>, D. Rohr<sup>32, [id](#)</sup>, D. Röhrich<sup>20, [id](#)</sup>, P.F. Rojas<sup>44</sup>, S. Rojas Torres<sup>35, [id](#)</sup>, P.S. Rokita<sup>135, [id](#)</sup>, G. Romanenko<sup>141, [id](#)</sup>, F. Ronchetti<sup>49, [id](#)</sup>, A. Rosano<sup>30,53, [id](#)</sup>, E.D. Rosas<sup>65</sup>, K. Roslon<sup>135, [id](#)</sup>, A. Rossi<sup>54, [id](#)</sup>, A. Roy<sup>48, [id](#)</sup>, S. Roy<sup>47, [id](#)</sup>, N. Rubini<sup>25, [id](#)</sup>, D. Ruggiano<sup>135, [id](#)</sup>, R. Rui<sup>23, [id](#)</sup>, B. Rumyantsev<sup>141</sup>, P.G. Russek<sup>2, [id](#)</sup>, R. Russo<sup>84, [id](#)</sup>, A. Rustamov<sup>81, [id](#)</sup>, E. Ryabinkin<sup>140, [id](#)</sup>, Y. Ryabov<sup>140, [id](#)</sup>, A. Rybicki<sup>107, [id](#)</sup>, H. Ryttonen<sup>116, [id](#)</sup>, W. Rzesza<sup>135, [id](#)</sup>, O.A.M. Saarimaki<sup>43, [id](#)</sup>, R. Sadek<sup>103, [id](#)</sup>, S. Sadhu<sup>31, [id](#)</sup>, S. Sadovsky<sup>140, [id](#)</sup>, J. Saetre<sup>20, [id](#)</sup>, K. Šafařík<sup>35, [id](#)</sup>, S.K. Saha<sup>4, [id](#)</sup>, S. Saha<sup>80, [id](#)</sup>, B. Sahoo<sup>47, [id](#)</sup>, B. Sahoo<sup>48, [id](#)</sup>, R. Sahoo<sup>48, [id](#)</sup>, S. Sahoo<sup>61</sup>, D. Sahu<sup>48, [id](#)</sup>, P.K. Sahu<sup>61, [id](#)</sup>, J. Saini<sup>134, [id](#)</sup>, K. Sajdakova<sup>37</sup>, S. Sakai<sup>124, [id](#)</sup>, M.P. Salvan<sup>97, [id](#)</sup>, S. Sambyal<sup>91, [id](#)</sup>, I. Sanna<sup>32,95, [id](#)</sup>, T.B. Saramela<sup>110</sup>, D. Sarkar<sup>136, [id](#)</sup>, N. Sarkar<sup>134</sup>, P. Sarma<sup>41, [id](#)</sup>, V. Sarritzu<sup>22, [id](#)</sup>, V.M. Sarti<sup>95, [id](#)</sup>, M.H.P. Sas<sup>137, [id](#)</sup>, J. Schambach<sup>87, [id](#)</sup>, H.S. Scheid<sup>64, [id](#)</sup>, C. Schiaua<sup>45, [id](#)</sup>, R. Schicker<sup>94, [id](#)</sup>, A. Schmah<sup>94</sup>, C. Schmidt<sup>97, [id](#)</sup>, H.R. Schmidt<sup>93</sup>, M.O. Schmidt<sup>32, [id](#)</sup>, M. Schmidt<sup>93</sup>, N.V. Schmidt<sup>87, [id](#)</sup>, A.R. Schmier<sup>121, [id](#)</sup>, R. Schotter<sup>128, [id](#)</sup>, A. Schröter<sup>38, [id](#)</sup>, J. Schukraft<sup>32, [id](#)</sup>, K. Schwarz<sup>97</sup>, K. Schweda<sup>97, [id](#)</sup>, G. Scioli<sup>25, [id](#)</sup>, E. Scomparin<sup>56, [id](#)</sup>, J.E. Seger<sup>14, [id](#)</sup>, Y. Sekiguchi<sup>123</sup>, D. Sekihata<sup>123, [id](#)</sup>, I. Selyuzhenkov<sup>97,140, [id](#)</sup>, S. Senyukov<sup>128, [id](#)</sup>, J.J. Seo<sup>58, [id](#)</sup>, D. Serebryakov<sup>140, [id](#)</sup>, L. Šerkšnytė<sup>95, [id](#)</sup>, A. Sevcenco<sup>63, [id](#)</sup>, T.J. Shaba<sup>68, [id](#)</sup>, A. Shabetai<sup>103, [id](#)</sup>, R. Shahoyan<sup>32</sup>, A. Shangaraev<sup>140, [id](#)</sup>, A. Sharma<sup>90</sup>, B. Sharma<sup>91, [id](#)</sup>, D. Sharma<sup>47, [id](#)</sup>, H. Sharma<sup>107, [id](#)</sup>, M. Sharma<sup>91, [id](#)</sup>, S. Sharma<sup>76, [id](#)</sup>, S. Sharma<sup>91, [id](#)</sup>, U. Sharma<sup>91, [id](#)</sup>, A. Shatat<sup>130, [id](#)</sup>, O. Sheibani<sup>115</sup>, K. Shigaki<sup>92, [id](#)</sup>, M. Shimomura<sup>77</sup>, J. Shin<sup>11</sup>, S. Shirinkin<sup>140, [id](#)</sup>, Q. Shou<sup>39, [id](#)</sup>, Y. Sibiriak<sup>140, [id](#)</sup>, S. Siddhanta<sup>52, [id](#)</sup>, T. Siemiarczuk<sup>79, [id](#)</sup>, T.F. Silva<sup>110, [id](#)</sup>, D. Silvermyr<sup>75, [id](#)</sup>, T. Simantathammakul<sup>105</sup>, R. Simeonov<sup>36, [id](#)</sup>, B. Singh<sup>91</sup>, B. Singh<sup>95, [id](#)</sup>, R. Singh<sup>80, [id](#)</sup>, R. Singh<sup>91, [id](#)</sup>, R. Singh<sup>48, [id](#)</sup>, S. Singh<sup>15, [id](#)</sup>, V.K. Singh<sup>134, [id](#)</sup>, V. Singhal<sup>134, [id](#)</sup>, T. Sinha<sup>99, [id](#)</sup>, B. Sitar<sup>12, [id](#)</sup>, M. Sitta<sup>132,56, [id](#)</sup>, T.B. Skaali<sup>19</sup>, G. Skorodumovs<sup>94, [id](#)</sup>, M. Slupecki<sup>43, [id](#)</sup>, N. Smirnov<sup>137, [id](#)</sup>, R.J.M. Snellings<sup>59, [id](#)</sup>, E.H. Solheim<sup>19, [id](#)</sup>, J. Song<sup>115, [id](#)</sup>, A. Songmoolnak<sup>105</sup>, F. Soramel<sup>27, [id](#)</sup>, A.B. Soto-hernandez<sup>88, [id](#)</sup>, R. Spijkers<sup>84, [id](#)</sup>, I. Sputowska<sup>107, [id](#)</sup>, J. Staa<sup>75, [id](#)</sup>, J. Stachel<sup>94, [id](#)</sup>, I. Stan<sup>63, [id](#)</sup>, P.J. Steffanic<sup>121, [id](#)</sup>, S.F. Stiefelmaier<sup>94, [id](#)</sup>, D. Stocco<sup>103, [id](#)</sup>, I. Storehaug<sup>19, [id](#)</sup>, P. Stratmann<sup>125, [id](#)</sup>, S. Strazzi<sup>25, [id](#)</sup>, C.P. Stylianidis<sup>84</sup>, A.A.P. Suaide<sup>110, [id](#)</sup>, C. Suire<sup>130, [id](#)</sup>, M. Sukhanov<sup>140, [id](#)</sup>, M. Suljic<sup>32, [id](#)</sup>, R. Sultanov<sup>140, [id](#)</sup>, V. Sumberia<sup>91, [id](#)</sup>, S. Sumowidagdo<sup>82, [id](#)</sup>, S. Swain<sup>61</sup>, I. Szarka<sup>12, [id](#)</sup>, M. Szymkowski<sup>135, [id](#)</sup>, S.F. Taghavi<sup>95, [id](#)</sup>, G. TAILLEPIED<sup>97, [id](#)</sup>, J. Takahashi<sup>111, [id](#)</sup>, G.J. Tambave<sup>20, [id](#)</sup>, S. Tang<sup>126,6, [id](#)</sup>, Z. Tang<sup>119, [id](#)</sup>, J.D. Tapia Takaki<sup>117, [id](#)</sup>, N. Tapus<sup>113</sup>, L.A. Tarasovicova<sup>125, [id](#)</sup>, M.G. Tarzila<sup>45, [id](#)</sup>, G.F. Tassielli<sup>31, [id](#)</sup>, A. Tauro<sup>32, [id](#)</sup>, G. Tejeda Muñoz<sup>44, [id](#)</sup>, A. Telesca<sup>32, [id](#)</sup>, L. Terlizzi<sup>24, [id](#)</sup>, C. Terrevoli<sup>115, [id](#)</sup>, S. Thakur<sup>4, [id](#)</sup>, D. Thomas<sup>108, [id](#)</sup>, A. Tikhonov<sup>140, [id](#)</sup>, A.R. Timmins<sup>115, [id](#)</sup>, M. Tkacik<sup>106</sup>, T. Tkacik<sup>106, [id](#)</sup>, A. Toia<sup>64, [id](#)</sup>, R. Tokumoto<sup>92</sup>, N. Topilskaya<sup>140, [id](#)</sup>, M. Toppi<sup>49, [id](#)</sup>, F. Torales-Acosta<sup>18</sup>, T. Tork<sup>130, [id](#)</sup>, A.G. Torres Ramos<sup>31, [id](#)</sup>, A. Trifiró<sup>30,53, [id](#)</sup>, A.S. Triolo<sup>32,30,53, [id](#)</sup>, S. Tripathy<sup>51, [id](#)</sup>, T. Tripathy<sup>47, [id](#)</sup>, S. Trogolo<sup>32, [id](#)</sup>, V. Trubnikov<sup>3, [id](#)</sup>, W.H. Trzaska<sup>116, [id](#)</sup>, T.P. Trzcinski<sup>135, [id](#)</sup>, A. Tumkin<sup>140, [id](#)</sup>, R. Turrisi<sup>54, [id](#)</sup>, T.S. Tveter<sup>19, [id](#)</sup>, K. Ullaland<sup>20, [id](#)</sup>, B. Ulukutlu<sup>95, [id](#)</sup>, A. Uras<sup>127, [id](#)</sup>, M. Urioni<sup>55,133, [id](#)</sup>, G.L. Usai<sup>22, [id](#)</sup>, M. Vala<sup>37</sup>, N. Valle<sup>21, [id](#)</sup>, L.V.R. van Doremalen<sup>59</sup>, C. Van Hulse<sup>130, [id](#)</sup>, M. van Leeuwen<sup>84, [id](#)</sup>, C.A. van Veen<sup>94, [id](#)</sup>, R.J.G. van Weelden<sup>84, [id](#)</sup>, P. Vande Vyvre<sup>32, [id](#)</sup>, D. Varga<sup>46, [id](#)</sup>, Z. Varga<sup>46, [id](#)</sup>, M. Vasileiou<sup>78, [id](#)</sup>, A. Vasiliev<sup>140, [id](#)</sup>, O. Vázquez Doce<sup>49, [id](#)</sup>, O. Vazquez Rueda<sup>115, [id](#)</sup>, V. Vechernin<sup>140, [id](#)</sup>, E. Vercellin<sup>24, [id](#)</sup>, S. Vergara Limón<sup>44</sup>, L. Vermunt<sup>97, [id](#)</sup>, R. Vértesi<sup>46, [id](#)</sup>, M. Verweij<sup>59, [id](#)</sup>, L. Vickovic<sup>33</sup>, Z. Vilakazi<sup>122</sup>, O. Villalobos Baillie<sup>100, [id](#)</sup>, A. Villani<sup>23, [id](#)</sup>, G. VINO<sup>50, [id](#)</sup>, A. Vinogradov<sup>140, [id](#)</sup>, T. Virgili<sup>28, [id](#)</sup>, M.M.O. Virta<sup>116, [id](#)</sup>, V. Vislavicius<sup>75</sup>, A. Vodopyanov<sup>141, [id](#)</sup>, B. Volkel<sup>32, [id](#)</sup>, M.A. Völkl<sup>94, [id](#)</sup>, K. Voloshin<sup>140</sup>, S.A. Voloshin<sup>136, [id](#)</sup>, G. Volpe<sup>31, [id](#)</sup>, B. von Haller<sup>32, [id](#)</sup>, I. Vorobyev<sup>95, [id](#)</sup>, N. Vozniuk<sup>140, [id](#)</sup>, J. Vrláková<sup>37, [id](#)</sup>, C. Wang<sup>39, [id](#)</sup>, D. Wang<sup>39</sup>, Y. Wang<sup>39, [id](#)</sup>, A. Wegrzynek<sup>32, [id](#)</sup>, F.T. Weiglhofer<sup>38</sup>, S.C. Wenzel<sup>32, [id](#)</sup>, J.P. Wessels<sup>125, [id](#)</sup>, J. Wiechula<sup>64, [id](#)</sup>, J. Wikne<sup>19, [id](#)</sup>, G. Wilk<sup>79, [id](#)</sup>, J. Wilkinson<sup>97, [id](#)</sup>, G.A. Willems<sup>125, [id](#)</sup>, B. Windelband<sup>94, [id](#)</sup>, M. Winn<sup>129, [id](#)</sup>, J.R. Wright<sup>108, [id](#)</sup>, W. Wu<sup>39</sup>, Y. Wu<sup>119, [id](#)</sup>, R. Xu<sup>6, [id](#)</sup>, A. Yadav<sup>42, [id](#)</sup>,

A.K. Yadav<sup>134, </sup>, S. Yalcin<sup>72, </sup>, Y. Yamaguchi<sup>92, </sup>, S. Yang<sup>20</sup>, S. Yano<sup>92, </sup>, Z. Yin<sup>6, </sup>, I.-K. Yoo<sup>16, </sup>,  
 J.H. Yoon<sup>58, </sup>, S. Yuan<sup>20</sup>, A. Yuncu<sup>94, </sup>, V. Zaccolo<sup>23, </sup>, C. Zampolli<sup>32, </sup>, F. Zanone<sup>94, </sup>, N. Zardoshti<sup>32, </sup>,  
 A. Zarochentsev<sup>140, </sup>, P. Závada<sup>62, </sup>, N. Zaviyalov<sup>140</sup>, M. Zhalov<sup>140, </sup>, B. Zhang<sup>6, </sup>, L. Zhang<sup>39, </sup>,  
 M. Zhang<sup>6</sup>, S. Zhang<sup>39, </sup>, X. Zhang<sup>6, </sup>, Y. Zhang<sup>119</sup>, Z. Zhang<sup>6, </sup>, M. Zhao<sup>10, </sup>, V. Zhrebchevskii<sup>140, </sup>,  
 Y. Zhi<sup>10</sup>, D. Zhou<sup>6, </sup>, Y. Zhou<sup>83, </sup>, J. Zhu<sup>97,6, </sup>, Y. Zhu<sup>6</sup>, S.C. Zugravel<sup>56, </sup>, N. Zurlo<sup>133,55, </sup>

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