



LAWRENCE
LIVERMORE
NATIONAL
LABORATORY

Megahertz pulse trains enable multi-hit serial femtosecond crystallography experiments at X-ray free electron lasers

S. Holmes, H. J. Kirkwood, R. Bean, K. Giewekemeyer, A. V. Martin, M. Hadian-Jazi, M. O. Wiedorn, D. Oberthur, L. Adriano, M. A. Coleman, M. Frank, M. Shelby, H. Chapman, A. Mancuso, A. Barty, C. Darmanin, B. Abbey

January 30, 2023

Nature Communications

Disclaimer

This document was prepared as an account of work sponsored by an agency of the United States government. Neither the United States government nor Lawrence Livermore National Security, LLC, nor any of their employees makes any warranty, expressed or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States government or Lawrence Livermore National Security, LLC. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States government or Lawrence Livermore National Security, LLC, and shall not be used for advertising or product endorsement purposes.

1
2 **Title : Megahertz pulse trains enable multi-hit serial femtosecond crystallography**
3 **experiments at X-ray free electron lasers**
4

5
6 **Authors**

7 S. Holmes¹, H. J. Kirkwood², R. Bean², K. Giewekemeyer², A.V. Martin³, M. Hadian-Jazi^{1,2}, M. O.
8 Wiedorn⁴, D. Oberthur⁴, L. Adriano⁵, N. Al-Qudami², S. Bajt⁵, I. Barák⁶, S. Bari⁵, J. Bielecki², S.
9 Brockhauser², M. A. Coleman⁷, F. Cruz-Mazo⁸, C. Danilevski², K. Dörner², A. M. Gañán-Calvo⁸,
10 R. Graceffa², M. Heymann⁹, H. Fangohr², M. Frank⁷, S. Hauf², A. Kaukher², Y. Kim², B. Kobe¹⁰,
11 J. Knoška^{4,18}, T. Laurus⁵, R. Letrun², L. Maia², M. Messerschmidt²⁰, M. Metz⁴, T. Michelat², G.
12 Mills², S. Molodtsov², D. C. F. Monteiro^{11,12}, A. Münnich², G. Previtali², A. Round², T. Sato^{2,5}, R.
13 Schubert², J. Schulz², M. Shelby⁷, C. Seuring^{4,11}, J. A. Sellberg¹³, M. Sikorski², A. Silenzi², S.
14 Stern², J. Sztuk-Dambietz², J. Szuba², M. Trebbin^{14,15}, P. Vagovic², T. Ve¹⁶, B. Weinhausen², K.
15 Wrona², P. L. Xavier^{4,19}, C. Xu², O. Yefanov⁴, K. A. Nugent^{1,17}, H. Chapman^{4,11,18}, A. Mancuso^{2,1},
16 A. Barty⁴, C. Darmanin^{1*}, B. Abbey^{1*}.

17
18
19 **Affiliations**

20 ¹Australian Research Council (ARC) Centre of Excellence in Advanced Molecular Imaging,
21 Department of Chemistry and Physics, La Trobe Institute for Molecular Sciences, La Trobe
22 University, Bundoora, VIC 3086, Australia.

23 ²European XFEL, Holzkoppel 4, 22869 Schenefeld, Germany.

24 ³School of Science, RMIT University, Melbourne, VIC 3000, Australia.

25 ⁴Center for Free-Electron Laser Science, Deutsches Elektronen-Synchrotron DESY, Notkestrasse
26 85, 22607 Hamburg, Germany.

27 ⁵DESY, Notkestrasse 85, 22607 Hamburg, Germany.

28 ⁶Institute of Molecular Biology, SAS, Dubravská cesta 21, 845 51 Bratislava, Slovakia.

29 ⁷Lawrence Livermore National Laboratory, 7000 East Avenue, Livermore, CA 94550, USA.

30 ⁸Dept. de Ingeniería Aeroespacial y Mecánica de Fluidos, ETSI, Universidad de Sevilla, 41092
31 Sevilla, Spain.

32 ⁹Max Planck Institute of Biochemistry, Am Klopferspitz 18, D-82152 Am Klopferspitz 18, D-
33 82152 Martinsried.

34 ¹⁰School of Chemistry and Molecular Biosciences, Institute for Molecular Bioscience and
35 Australian Infectious Diseases Research Centre, University of Queensland, Brisbane, QLD 4072,
36 Australia.

37 ¹¹The Hamburg Centre for Ultrafast Imaging, Universität Hamburg, Luruper Chaussee 149,
38 Hamburg 22761, Germany.

39 ¹²Hauptman-Woodward Medical Research Institute, 700 Ellicott St., Buffalo, NY 14203, USA,

40 ¹³Biomedical and X-ray Physics, Department of Applied Physics, AlbaNova University Center.
41 KTH Royal Institute of Technology, SE-106 91 Stockholm, Sweden.

42 ¹⁴Department of Chemistry, State University of New York at Buffalo, 760 Natural Sciences
43 Complex, Buffalo, New York 14260, USA.

44 ¹⁵Hauptman-Woodward Medical Research Institute, 700 Ellicott Street, Buffalo, New York 14203,
45 USA.

46 ¹⁶Institute for Glycomics, Griffith University, Southport, QLD 4222, Australia.

47 ¹⁷Laser Physics Centre, Research School of Physics Australian National University, ACT, 2601
48 Australia. ¹⁸Department of Physics, Universität Hamburg, Luruper Chaussee 149, 22761 Hamburg,
49 Germany.

50 ¹⁹Max-Planck Institute for the Structure and Dynamics of Matter, Luruper Chaussee 175, 22761
51 Hamburg, Germany.

52 ²⁰School of Molecular Science, Arizona State University, Tempe, AZ, 85281, USA.

53

54 *Correspondence to: c.darmanin@latrobe.edu.au and b.abbey@latrobe.edu.au

55

56

57

58 The pulse structure of the European X-ray Free Electron Laser (EuXFEL) enables Serial
59 Femtosecond Crystallography (SFX) data to be collected at megahertz (MHz) repetition rates. The
60 future capabilities of MHz repetition rate are still to be explored. The combination of the high-
61 sensitivity, ultra-fast, AGIPD detector and the sub-microsecond inter-pulse spacing of the EuXFEL,
62 allows multiple diffraction images to be collected from the same crystal as it traverses the X-ray
63 beam in a liquid jet. Here we exploit the MHz pulse structure of the EuXFEL to measure diffraction
64 patterns (up to 2.1 Å resolution) from the same protein crystal at two time-points separated by less
65 than a microsecond. We term this approach ‘Multi-Hit SFX’. This is the first indication using the
66 MHz regime that a single crystal can be illuminated twice, and a structure can be generated from
67 the second hit data with no signs of radiation damage visible in the Fourier difference maps. This

68 technique takes advantage of the larger, faster to set-up beam sizes possible at the instrument. These
69 results pave the way towards a new technique for developing *in-situ* single crystal, sub-
70 microsecond, timing experiments at MHz repetition rate XFELs.

71

72 **Introduction**

Serial femtosecond crystallography (SFX) has changed the way in which we can probe macromolecular structures and investigate their dynamics. SFX is especially well suited for structure determination with small crystals and time-resolved systems exhibiting irreversible processes, which cannot be measured using conventional synchrotron or lab-based X-ray sources¹⁻⁵. One hurdle to implementing SFX is obtaining a large enough data set for high-resolution 3D structure determination, which historically has meant high sample consumption rates. At traditional XFEL facilities, repetition rates are low and obtaining a large data set takes time⁶. With the development of new sources like the EuXFEL⁷, capable of MHz repetition rates, both data collection times and sample consumption may be significantly reduced. Another avenue of research open to MHz XFEL facilities is the ability to use the pulse structure of the source to perform novel timing experiments⁷⁻⁹. Typically, X-ray pump/X-ray probe experiments at XFELs enable the study of dynamics such as bond breaking on sub-picosecond time-scales¹⁰⁻¹³. The spacing between consecutive pulses at the EuXFEL, is normally sub-microsecond. This creates an opportunity to perform a new class of time-resolved experiments in which molecular dynamics occurring on sub-microsecond time-scales can be studied on the same crystal. This timing regime is ideally suited to the study of fast (sub-millisecond) dynamics in proteins including transition states (i.e. active to inactive form)¹⁴⁻¹⁸, as well as helix motions¹⁹ and side-chain rotations²⁰. The unique pulse train structure of the EuXFEL offers the opportunity to collect femtosecond diffraction data with variable sub-microsecond time delays. An example problem that could benefit from multi-hit SFX is understanding the mechanism of the *Bacillus subtilis* response regulator, SpoOF, which is involved in sporulation. SpoOF can induce a shift in protein conformation on picosecond to millisecond time-scales¹⁶. Previous NMR studies have shown the secondary structure dynamics of the protein involves a complex series of movements occurring over a wide range of time-scales. These include rotation of bonds in methyl groups (thought to occur within nanoseconds) to side chain flipping of buried residues (which occurs on the scale of seconds); critical protein-protein contacts between SpoOF and its binding partners take place on millisecond time-scales¹⁶. Examples which maybe important for this application include studying protein-protein complexes or binding of co-factors such as ATP or NADP which induce conformational changes within proteins and result in the

transition between active to inactive forms of the enzyme which occur within millisecond time frames. Multi-hit SFX performed on sub-microsecond time-scales would be able to capture the intermediate helical/sheet conformational changes in the secondary structures of proteins prior to the important side chain flipping. The advantage of this new technique is to evolve the existing time-resolved SFX capabilities¹⁰⁻¹³ enabling capture of intermediate states ‘live’ based on the timing between the pulses rather than activation/mixing times in the jet. This work also questions whether a high number of photons and a good focus from XFEL is absolutely required if we have good crystal quality.

73

74 The most common method for delivering samples to the XFEL beam is via a liquid jet, formed
75 using a flow focusing nozzle (commonly referred to as GDVN)²¹⁻²³, in which the sample is
76 continuously replenished with respect to the X-ray beam. Typically, no single crystal flowing in a
77 liquid jet can contribute to more than one diffraction pattern at an XFEL due to the crystal being
78 destroyed by each single beam interaction or the jet velocity being too fast with respect to the X-
79 ray pulse repetition rate²⁴⁻²⁵. The normally ultra-intense femtosecond duration pulses available at
80 XFELs, typically allow high-resolution diffraction data to be collected from suitably crystalline
81 samples during a single ‘hit’, prior to any significant nuclear motion occurring. The first
82 experiments at EuXFEL addressed the issue of jet explosion and its effect on the crystals. As a
83 result, multiple jet speeds (42, 78 and 102 m/s) were tested¹. However, we discovered using the
84 EuXFEL it is possible, with the appropriate beam intensity and flow rate, that a single crystal may
85 be hit multiple times as it passes through the beam without exhibiting degradation of measured
86 diffraction intensities. Hence one infers there is also no extensive crystal damage. This "multi-hit
87 diffraction" scenario is clearly distinct from "multi-crystal diffraction", in which two or more
88 crystals arrive at the X-ray beam at the same time, and which can be identified in the diffraction
89 pattern and by using multi-crystal indexing options²⁶⁻²⁸. Wiedorn *et. al*¹ also investigated the hit
90 rates as a function of train pulse number for each jet speed, and confirmed that there is no bias on
91 which train pulse number hits a crystal.

92

93 From our analysis, we can obtain high-resolution diffraction data from the first and second hits of
94 a single crystal in a MHz SFX experiment, hence the term ‘Multi-Hit SFX’. This is in stark contrast
95 to conventional liquid-jet SFX experiments in which each single crystal gives rise to just one
96 diffraction pattern prior to exiting the beam or being destroyed. Though it should be emphasised
97 that the caveat to performing multi-hit SFX is that the dose from the initial hit must be below the

98 threshold for significant radiation damage of the crystal to occur. Whilst this likely precludes the
99 use of nanocrystals, due to the weak diffraction from these samples, it does allow for multi-hit,
100 femtosecond diffraction experiments to be performed on microcrystals. Microcrystals are often
101 used for time-resolved SFX using either an optical pump or rapid mix-and-inject to initiate
102 structural changes within the sample¹¹.

103

104 Multi-hit SFX can be thought of as a bridge between SFX and traditional crystallography. SFX;
105 where thousands of crystals are hit once and destroyed, and traditional crystallography; where one
106 crystal is continuously rotated and illuminated multiple times before significant radiation damage
107 occurs. Previously, this phenomenon has only been detected using highly-viscous jet streams. This
108 is because, unlike liquid jets, highly-viscous jets can flow at speeds slow enough that even at lower
109 (e.g. non-MHz) repetition rate XFEL facilities multiple hits can be observed²⁹⁻³¹. A key benefit, in
110 addition to reducing sample consumption, is that multi-hit SFX at a MHz XFEL source can be used
111 to probe multiple time-points sub-microseconds apart and given good crystal quality, the need for
112 high photon counts and a good focus from the XFEL is no longer an essential requirement. The
113 important questions we address in this paper are:

114

- 115 • How often are multi-hits (in this case double hits) observed under the conditions probed
116 during the initial experiment at the SPB/SFX instrument^{1,32}?
- 117 • Why do multi-hits occur and what is the influence of the experimental geometry?
- 118 • How do multi-hits impact the data quality and resolution?
- 119 • Under what experimental conditions will more than one hit of the same sample occur and
120 how can this information be used to tailor future MHz XFEL experiments?

121

122 Based on the analysis presented by Wiedorn et al. (2018)¹ we have constructed a model that
123 demonstrates how crystals are hit either once or twice by the beam and demonstrate that under the
124 conditions used here the second hits do not negatively impact the quality of the data. This implies
125 that the data from multiple hits of the same crystal can be used for structure determination. This
126 result allow users of XFEL sources to select experimental parameters that can increase the
127 probability of multiple hits, or at least try not to minimise them, potentially increasing the amount
128 of usable data per crystal in an SFX experiment and paving the way to novel forms of microsecond,
129 time-resolved measurements.

130 **Results**

131

132 **EuXFEL experimental set-up and injector jet speeds**

133 SFX was carried out on the SPB/SFX instrument³² at the EuXFEL using compound refractive lenses
134 (CRL) to focus the beam as described in Wiedorn et al. (2018)¹. Lysozyme (HEWL) crystals were
135 delivered to a 1.125 MHz XFEL beam using a flow focused jet (also referred to as GDVN),²¹⁻²³ at
136 three different target jet speeds for the experiment: 42 m/s, 78 m/s and 102 m/s¹. The two slowest
137 jet speeds were checked experimentally¹, whilst the fastest jet speed (102 m/s) was theoretically
138 determined (see supplementary data). A summary of the liquid and gas flow rates and their
139 equivalent jet speeds from Wiedorn *et. al.* (2018)¹ are shown in Table 1, alongside the theoretically
140 calculated 102 m/s jet speed.

141

142 **Identifying consecutive hits for a single lysozyme crystal**

143 To assess this multi-hit phenomenon, lysozyme data from the first SPB/SFX commissioning round
144 of the EuXFEL experiment was analysed. The X-ray data statistics corresponding to the three
145 different jet speeds tested during this experiment are shown in Table 1. The number of images
146 collected is compared to the number of images with hits detected by Cheetah³³ and indexed by
147 *CrystFEL*^{34,35}, as well as the percentage of those patterns indexed where crystals were hit twice as
148 they passed through the beam. A crystal was determined a ‘multi-hit’ if it satisfied the following
149 criteria; i) the two hits were from consecutive pulses and ii) the patterns showed very similar
150 crystallographic orientations (the angle between each pair of basis vectors were less than 5 degrees
151 and the lengths of the basis vectors were in agreement to within 10%, determined using *CrystFEL*’s
152 *Whirligig* program^{34,35}).

153

154 The data shows that, as expected, when the jet speed increases the percentage of double hits
155 decreases. For the slowest (42 m/s), intermediate (78 m/s), and fastest (102 m/s) jet speeds the
156 percentage of double hits were 6.4%, 0.9% and 0.3%, respectively.

157

158 **Crystal rotation and Bragg peak analysis**

159 To confirm that our method for determining crystals hit twice was reliable, we calculated the change
160 in crystal rotation for all the diffraction data collected from consecutive hits in the pulse trains.
161 Rotation of the crystal as it ‘tumbles’ in the liquid jet between consecutive X-ray pulses results in
162 small changes in the position of the Bragg peak. A significant increase in the number of crystals
163 that had <5 degree change in orientation between consecutive images was observed for the 42 m/s

164 jet speed (Fig. 1 and supplementary Figs. S1 and S2). Fig 1 illustrates the change in crystal rotation
165 between consecutive pulses as the crystal passes through the X-ray beam. A significant increase in
166 the number of crystals with a <5 degree change in rotation can be seen, which is indicative of the
167 fact that many of the crystals in this region are hit twice by the XFEL beam before they either exit
168 the interaction region or are destroyed. By contrast, the number of crystals with change in rotation
169 of >5 degrees between consecutive hits does not vary significantly. This suggests that for the 42
170 m/s jet speed the crystal does not rotate more than 5 degrees between the first and second hit for
171 double hit crystals. These results verify the predictions of the *whirligig* program and show that a
172 significant fraction of the crystals were hit twice as they traversed the X-ray interaction region. The
173 *whirligig* program was therefore used to determine the number of multiple hits that occurred for the
174 two faster jet speeds. Due to the lower hit rates for these two jets speeds (78 and 102 m/s) the
175 number of diffractions patterns collected which contained Bragg peaks, is less than compared to
176 the 42 m/s data set (Table 1).

177

178 **Modelling the multi-hit scenario**

179 Now to the question of how if, as expected, the crystal is destroyed in the beam it is possible to get
180 more than one diffraction measurement from a single crystal. Irrespective of crystal concentration,
181 multiple hits on a single crystal are primarily determined by the jet speed, beam focus size, and
182 pulse repetition rate. To illustrate how multiple hits occur, a graphical representation of the crystal
183 path through the beam based on the experimental data is presented in Fig. 2. For the slowest jet
184 speed (42 m/s) the crystal is initially hit within the tail (lowest intensity) region of the beam
185 followed by a second hit within the FWHM of the beam. As exposure to the second hit is within
186 the most intense part of the X-ray beam, it leads to the destruction of the crystal and this is consistent
187 with the observation that no third hit of the crystal was observed. As the jet speed increases to 78
188 m/s, multiple crystal hits become possible only within the tail regions of the X-ray beam, where the
189 crystal is clipped by both ends of the beam (Fig. 2). Therefore, the intensities of both these hits are
190 less than the second hit for the 42 m/s jet speed. At the fastest jet speed (102 m/s), multiple crystal
191 hits can only occur in the extreme tail regions of the X-ray beam (Fig. 2). Therefore, the slowest jet
192 speed (42 m/s) resulted in the greatest number of double-hits.

193

194 **X-ray beam profile**

195 The X-ray beam size and profile determines the percentage of double hits observed during the
196 experiment and therefore is fundamental to achieving double-hits. To independently verify that our
197 model for the double-hit condition is correct, we calculated the beam size and profile using optical

198 measurements accounting for the point spread function of the YAG crystal and the optical
199 resolution of the system (see Methods). The beam profile was consistent with a Lorentzian model
200 as observed previously at the SPB/SFX beam line. The Lorentzian full-width-at-half-maximum
201 (FWHM) ($2\gamma = 50\%$) of each optical image was experimentally determined with a 7.5% image
202 threshold applied to eliminate background scatter. The full width (FW) ($7.04\gamma = 82.38\%$) was then
203 calculated and the variation in beam width was determined as shown in Fig. 3. The mean, minimum,
204 and maximum FWHMs were calculated to be 18.7 μm , 11.9 μm and, 23.3 μm , respectively. Whilst
205 the mean, minimum, and maximum, FWs were calculated to be 65.8 μm , 41.7 μm and 82.1 μm ,
206 respectively. The experimentally determined beam size was consistent with the optical measured
207 X-ray beam size, which shows double hits are possible.

208

209 **Lysozyme crystal movement through the X-ray beam**

210 The crystal trajectory through the X-ray beam can be calculated based on the experimental and
211 theoretical jet speeds provided in Table 1. The minimum distance a lysozyme crystal of length 8
212 μm (with respect to the flow direction) could clear between two X-ray pulses whilst having at least
213 1 μm of crystal within the beam window and with an inter-pulse spacing of 886 ns was calculated.
214 For the 42 m/s jet, a crystal will clear a minimum distance of 35.4 μm , whereas for the 78 m/s jet,
215 a crystal will clear a minimum distance of 56.7 μm and, at the fastest jet speed 102 m/s, the crystal
216 will clear a minimum distance of 79.9 μm . By minimum distance, we refer to the distance calculated
217 using the minimum jet speed including a 5% margin of error. This error is consistent with the
218 observation that hits are observed with the 102 m/s jet speed (Table 1). Fig. 2 shows a schematic of
219 how far the crystal has travelled at each jet speed. Based on this information it confirms that for
220 each of the three jet speeds, the X-ray beam diameter was sufficiently large that double hits occur,
221 which supports our experimental observations.

222

223 **Data quality check**

224 Quality control checking of the data is necessary to ensure we don't see any degradation in the
225 second hit data with respect to the first hit data. To check this, we compared the Bragg intensities
226 and resolution of the two data sets (first and second hit on the same crystal). The crystal diffraction
227 data was used to generate three independent powder plots (first hits, second hits, and single hits).
228 For the slowest jet speed used in this experiment (42 m/s), we observe that the integrated intensity
229 of the first of the double hits was lower than the second hit, with the integrated intensity for the
230 second hit more similar to that measured from crystals hit just once by the X-ray beam (Fig. 4a and
231 5a). This supports the interpretation that the second hit occurs within the FWHM of the X-ray

232 beam. As the jet speed increases the diffraction intensity profiles change generating a larger
233 difference in the integrated intensity between the double hit crystal and the single hit crystal. At a
234 jet speed of 78 m/s, the integrated intensity profiles for the first and second hit of the double hit
235 crystal is decreased further and found to be more similar in intensity to each other rather than the
236 single hit crystal intensity profile (Fig. 4b, 5a). This supports the fact that the crystal is hit within
237 both the tail regions of the X-ray beam, either side of the FWHM. At the fastest jet speed (102 m/s)
238 the difference in integrated intensities between the double hits and single hits is even more
239 pronounced. The integrated intensity profiles of all double hit crystals are very similar and much
240 less intense than the single hit crystals (Fig. 4c, 5a), which is again consistent with our model as
241 shown in Fig. 2. Irrespective of the integrated intensity data, if we specifically select for only high-
242 resolution reflections in the data (Fig. 5b), little difference ($\leq 5\%$) is detected in the data obtained
243 above 2.5 Å resolution, when comparing crystal hits for the same jet speed. When comparing the
244 data from the different jet speeds, an overall reduction in the percentage of high-resolution peaks
245 can be seen for the 42 m/s jet speed, however, this can be explained by the larger quantity of data
246 obtained for the slowest jet speed compared to the two higher jet speeds (Table 1). We would expect
247 the higher jet speeds to have been comparable to the 42 m/s jet if the equivalent amount of data was
248 collected. To further check if radiation damage has occurred within the crystal after the first hit, we
249 compared two data sets, single hit crystal data (where crystals were only hit once) versus the
250 second hit crystal data (where the same crystal was hit for the second time). Comparisons of the
251 two structures showed very little differences. Importantly, closer examination of the disulfide
252 bonds, which are highly prone to radiation damage, also showed no evidence of radiation damage
253 in the Fourier difference map (Fig. 6A). A statistical comparison between the single hit and second
254 hit data (in the double hit pair) show the data quality is similar (Fig. 6 and Table 2). The Wilson
255 plots for both data sets were included in the supplementary file (Supplementary Fig. S4). CC* is a
256 commonly used metric to assess data quality for structure determination. If we compare half the
257 data set of the first hits to half of the data sets of the second hits, there are double the reflections in
258 the second hit data set and the CC* are comparable (0.85 from a total of 6506 reflection pairs)
259 further indicating there is no significant change in quality between the two data sets (Fig. 6B and
260 Table 2).

261

262 Discussion

263

264 During this experiment, different jet speeds were tested to study the effect of jet explosion at the
265 EuXFEL. The jet explosion begins at the centre of the interaction region and produces an initial

266 gap that expands in the vertical direction aligned with the liquid jet. From our observations, the gap
267 initiates in most cases from the vaporization of a jet slice with a thickness about 1/5 to 1/4 of the
268 jet diameter, and initially expands according to a time power law that has been investigated.³⁶ That
269 power law yields initial expansion velocities of the order of 200 m/s for our experiments, which
270 decay approximately as $(t/t_0)^{-3/7}$, where t_0 is approximately to 1 ns in our experiments. Thus, if the
271 crystal is hit near the boundaries of the irradiated jet segment, the gap will not reach the crystal
272 before the next pulse arrives. It is, therefore, important to establish the pulse power, beam size, jet
273 diameter, and jet speed to avoid the interaction of the crystal with the expanding front of the opening
274 gap formed by a previous pulse. Nevertheless, for the jet speeds tested in this experiment we did
275 not see an effect from jet explosions. In these conditions it was indeed possible to collect high
276 quality data on the same crystal in a subsequent pulse. This is further supported by the results of
277 Wiedorn *et. al.*¹ which shows the jet recovery for each jet speed tested and the Yefanov *et. al.*⁹
278 results which compared structures from different pulses within the Megahertz train with no
279 observable difference between the first and any subsequent pulse seen.

280

281 The results from this experiment represent the first demonstration of multi-hit SFX using a flow-
282 focusing injector. The same experiment is not currently possible at lower repetition rate sources
283 with GDVNs, as the unique pulse structure of the MHz XFEL source is required for these
284 experiments. Furthermore, the new focusing system now installed at SPB/SFX allows tuning of the
285 horizontal and vertical focus size independently and could be used to maximise double hits. For
286 example, by using a vertical line focus with a similar width to the jet and a height, tuned to a length
287 as selected for by the distance the crystal will travel for the given jet speed between X-ray pulses.
288 Hence, multi-hit SFX could be a standard mode of operation for time-resolved experiments on
289 microcrystals with this configuration. Our data shows that depending on the jet speed and beam
290 profile the second hit in the double hit series can be more intense than the first with comparable
291 integrated intensities to that of the single hit crystals. Our observations and analysis show that the
292 initial interaction of the X-ray beam with the crystal must occur within the low intensity region of
293 the beam profile, i.e. within the ‘tail’ regions in order for a second interaction to occur. The
294 experimental value of the beam size using the YAG (accounting for the PSF and optical resolution)
295 is consistent with the beam size required to generate double hits.

296

297 The time interval between pulses at the EuXFEL is extremely short when compared to other non-
298 MHz XFEL sources, however sub-microsecond pulse spacing is still long on the time-scale of
299 radiation damage³⁷. Crystal dose estimates within the tail regions, calculated based on the

300 measured beam and crystal size^{38,39}, give an average absorbed dose of 0.3 Gy for the first hit
301 compared to 1.5 Gy within the FWHM of the beam. Typically, protein crystals can withstand a
302 radiation dose of 30 MGy^{40,41} under cryo-cooled conditions as it slows the effects of radiation
303 damage, whilst at room temperature the limit decreases by approximately two orders of magnitude,
304 i.e. ~ 0.2 MGy, as a result of the diffusion of free radicals⁴¹⁻⁴⁴. In addition, room temperature
305 radiation damage becomes dependent on dose rate⁴⁵. It has been reported that this dose limit can be
306 increased (30-150 MGy)^{3,46,47} by the introduction of highly intense femtosecond pulses at XFEL
307 facilities due to the pulse duration being short enough to significantly reduce the effects of radiation
308 damage under room temperature conditions^{37,48}. Radiation damage effects at XFELs is still an
309 active area of research⁴⁹⁻⁵¹ and is dependent on both dose rate, and the specific choice of sample.
310 However, in the present case, the dose estimates calculated for the double hit crystals measured
311 here fall well below published thresholds for radiation damage and is supported by the fact that
312 no significant degradation in the quality of diffraction was observed in the second hit structure.
313 With this in mind, we have developed an explanation for the observed integrated intensity data
314 considering the distance the crystal travels between X-ray pulses for the three different jet speeds.

315

316 Our interpretation is consistent with the observed changes in the measured integrated intensity
317 profile and it explains why for the 42 m/s jet speed, the second hit has a similar intensity profile to
318 the single hit crystal data (Fig. 4). This was initially contrary to expectations based on the
319 assumption that the crystal would be damaged during the first hit reducing the measured intensity
320 of the second hit. For the two higher jet speeds, the inter-pulse spacing and beam size are such that
321 to observe double hits the crystal must only interact with the tails of the beam for both the first and
322 second hit (Figures 2 & 3). In other words, for the faster jet speeds, it is not possible to observe a
323 double hit where one of the pair of hits is the result of an interaction within the central portion of
324 the beam. This is consistent with a weaker measured intensity for this data. However, it is important
325 to note that even in the tails of the beam, there is sufficient intensity at the XFEL to obtain high-
326 resolution diffraction from micron sized crystals, as has been observed in previous experiments³.

327

328 For the 78 m/s and 102 m/s jet speeds we demonstrate the estimated area of the beam 'seen' by the
329 crystal can be determined simply by considering the distance the crystal travels between pulses.
330 Whereas for the slowest jet speed (42 m/s) the second hit must occur within the high flux interaction
331 region to generate similar diffraction to the single hits and guarantee destruction of the crystal (Fig.
332 2). Similarly, for the faster jet speeds, the first hit occurs in the tail region of the beam where the
333 incident flux is low and does not significantly damage the crystal which reduces the chance of

334 double hits occurring. Using the beam size, jet speed and an understanding of how these parameters
335 relate to the region of crystal interaction, it is now possible to determine parameters, which either
336 minimise or maximise the chances of multi-hit SFX. This is critical, particularly for current and
337 future experiments at the EuXFEL, where we can investigate alternate time-resolved techniques
338 using the unique time structure of the source. To optimise for multiple hits it is possible to tailor
339 the jet speed and beam focus. More importantly, the spacing between the successive pulses in a
340 EuXFEL pulse train can be modified to allow for multi-hit SFX to probe different time regimes.
341 Multi-hit SFX will be a critical consideration for when the EuXFEL operates at planned 4.4 MHz
342 repetition rates, where multiple hits are significantly more frequent (and potentially unavoidable)
343 than at the 1.125 MHz that the current experiment was performed at. To predict how multi-hit SFX
344 at 1.125 MHz compares to 4.4 MHz, we have provided a figure to estimate optimal, beam size and
345 jet speed to achieve multiple hits. Fig. 7 demonstrates the criteria to obtain multi-hit regime for the
346 current rate (1.1 MHz, Fig. 7a) and future rate (4.4 MHz, Fig. 7b) of the EuXFEL.

347

348 In addition to the integrated intensity, we also studied the resolution of the diffraction data from
349 both single and double hits. A direct comparison shows that whilst the integrated intensity varies,
350 there are no significant differences in the resolution of single hits compared to the first and second
351 hits that make up the double hit pair and this is supported by the structural results that no radiation
352 damage is present. This suggests that all the data collected during this experiment can be merged
353 into a single dataset for structural analysis without detriment to the quality, allowing us to increase
354 the amount of data without increasing crystal consumption. Furthermore, it also implies that under
355 the experimental configuration used here, the shock-waves generated in the liquid jet, as a result of
356 the first hit, do not appear to impact the crystal quality or packing. If the crystal quality was affected
357 by the first X-ray pulse, we would expect the higher resolution diffraction peak(s) to disappear;
358 indicative of crystal damage, however this was not observed which supports the fact that no
359 radiation damage occurred. A comparison of the single hit and second hit molecular structures
360 confirm this, showing no signs of radiation damage in the Fourier difference maps.

361

362

363 With reference to the key questions posed at the start of the paper, we find that in the experimental
364 configuration used here, double-hits occur between 0.3 and 6.3% of the time depending on beam-
365 size and jet speed, with the slowest jet speed being optimal for multi-hit SFX. This percentage is
366 expected to increase at 4.4 MHz. The quality and frequency of double-hits is dominated by the
367 experimental geometry, though for smaller beam sizes the effect of radiation damage will become

368 increasingly significant so larger beams may be more desirable for time-resolved multi-hit SFX
369 experiments. Under the present experimental conditions, we observed no evidence of radiation
370 damage, or damage from XFEL induced shock-waves when looking at the resolution and integrated
371 intensity data or the molecular structures.

372

373 In summary, multi-hit SFX will be developed as a new mode of sub-microsecond time-resolved
374 serial crystallography, taking advantage of the unique time structure of MHz XFEL sources and
375 large beam focuses. This method will become increasingly important as additional MHz XFEL
376 sources come on-line, and as facilities turn to faster and faster MHz repetition rates. The beam sizes
377 used for this experiment can be easily realised within the standard set-up at SFX/SPB and requires
378 no modification of beam-line or upstream optics. In addition, control over the initial sample state
379 (structure) and subsequent dynamics of the injected crystals can be achieved via a rapid mix-and-
380 inject approach using a chemical trigger. This provides an opportunity for measuring molecular
381 dynamics at intermediate time-scales between those currently achievable with SFX (femtosecond
382 regime) and serial millisecond crystallography (SMX) (millisecond regime). Through analysis of
383 the individual Bragg peaks, by comparing plots from the first and second hits as well as successfully
384 solving the structure based solely on second hit data, we demonstrate that for the system studied
385 here, radiation damage effects are not observed. Most importantly, the second hit structure did not
386 reveal any radiation damage surrounding the disulfide bonds. It should be emphasised however,
387 that for proteins that are particularly radiation sensitive, the effects of radiation damage (especially
388 specific damage) may not be easily separable from the molecular dynamics. Finally, we have
389 outlined the criteria to observe multi-hits in terms of jet speed and beam size (Fig. 7) and confirmed,
390 through discussions with beam-line staff, that these parameters may be readily achieved at the
391 European XFEL.

392

393

394 **Methods**

395 Further details on the sample preparation and delivery, as well as the experimental set-up and initial
396 data analysis can be found in ref [1]. The beam focusing at the SPB/SFX beam-line for this
397 experiment was achieved using compound refractive lenses (CRLs). The crystal size varied
398 between 6 x 6 x 6 μm and 8 x 8 x 8 μm as characterised via an optical microscope. A subset of the
399 data generated and analysed from the first EuXFEL experiment (run numbers r0069-r0082, r0145-
400 r0150, r0153) in the form of *CrystFEL* stream files were used to identify crystals which had been

401 hit consecutively, it is this subset that was used for the analysis presented here. The data used for
402 this paper is available from CXIDB-ID80 (www.cxidb.org/id-80.html).

403

404 **Identification of multi-hits in the data**

405 The stream files with the already indexed crystal data were used to identify the crystals classified
406 as multi-hits. The stream files were filtered for these multi-hits using the *whirligig* script from the
407 *CrystFEL* crystallography suite^{34,35}, which defines crystals hit twice as those that had similar
408 crystallographic orientations in consecutive frames. Based on our analysis of the relative change in
409 crystal orientation between consecutive hits we confirmed that for multi-hit crystals the change in
410 angle between each pair of basis vectors was less than 5 degrees. In addition, the lengths of the
411 basis vectors had to be within 10% agreement. Filtering using the *whirligig* program based on these
412 criteria alone does not account for the pulse/train format of the EuXFEL. Therefore, python code
413 was implemented to filter out any adjacent frames that were not from the same pulse train. This
414 data was further sorted into three separate categories: data from crystals hit just once, data from the
415 first hit of crystals hit twice, and data from the second hit of crystals hit twice. These three stream
416 files were then used as the input files for further analysis in Python 3, using multiple parameters
417 (integrated intensity, a^* , b^* , c^* , h , k , l) within the stream files to calculate $1/d$, normalised integrated
418 intensity plots.

419

420 Python 3 was also used to analyse the *CrystFEL* stream files, where the change in orientation
421 between all consecutive images was calculated in an identical manner to the *whirligig* program.
422 The output was plotted as a histogram, showing the change in angle (degrees) for consecutive
423 images to confirm those consecutive hits that were selected as our double hit data sets were truly
424 from the same crystal.

425

426 **Analysis of data quality**

427 Merging and scaling of the Bragg peaks were performed using Partialator in the *CrystFEL suite*^{34,35}.
428 Figures of merit were calculated using `compare_hkl` (Rsplit, $CC_{1/2}$, CC^*) and `check_hkl` (SNR,
429 multiplicity, completeness), also using the *CrystFEL suite*^{34,35}. To generate a complete second hit
430 data-set, experimental data from all jet speeds, were merged. The statistics for the single hits and
431 second hits data sets were generated (Table 2) and Wilson plots for both data sets are shown in the
432 supplementary (Supplementary, Fig. S4).

433

434 Structure refinement was performed in CCP4i2 using the MTZ output from CrystFEL. A solvent
435 free version of lysozyme (pdb code 6FTR) was used as the initial starting model for Molecular
436 Replacement in Phaser⁵² and the R_{free} flags were generated (utilizing 10% of the data) followed by
437 iterative cycles of Refmac5⁵³ refinement and rebuilding of the model in Coot⁵⁴. MolProbity⁵⁵ and
438 Xtriage (Phenix)⁵⁶ tool was used to validate the model.

439

440 To assess the quality of the data a Fourier difference ($F_o - F_c$) map was generated using Refmac5⁵³
441 (Fig. 6A).

442

443 **X-ray Beam Profile Analysis**

444 To find the optimum focus, a photon energy scan of the accelerator was carried out. The
445 nominal energy for the SFX data collection was 9.232 keV. The beam size and profile were
446 estimated based on 6773 individual YAG images collected using the in-line microscope positioned
447 within in the chamber. The optical images were generated from single shots using a 15 μm thick
448 Ce:YAG screen (refer to Supplementary Fig. S3). The point spread function (PSF) of the YAG was
449 determined based on published estimates⁵⁷ (see Supplementary Fig. S3) and estimated to be
450 approximately 2 μm . Using Edmund standards, the optical microscope resolution was determined
451 to be 8 μm . A Lorentzian distribution was fitted to the optical images using Python code. The actual
452 beam size was determined via a convolution taking into account the PSF for the YAG and optical
453 microscope. A Lorentzian FWHM ($\text{FWHM} = 2\gamma$) of each image was determined using a 3×3
454 median filter and a 7.5% noise threshold. Furthermore, the Full Width ($\text{FW} = 7.04\gamma$) of 82.38% was
455 determined based on the Lorentzian above the noise threshold; and for each image it was calculated
456 and plotted along with the overall mean, minimum and maximum FW.

457

458 **References**

459

- 460 1. Wiedorn, M. O. *et al.* Megahertz serial crystallography. *Nat. Commun.* **9**, 4025 (2018).
- 461 2. Chapman, H. N. *et al.* Femtosecond X-ray protein nanocrystallography. *Nature* **470**, 73–77
462 (2011).
- 463 3. Boutet, S. *et al.* High-Resolution Protein Structure Determination by Serial Femtosecond
464 Crystallography. *Science* **337**, 362–364 (2012).
- 465 4. Gati, C. *et al.* Atomic structure of granulin determined from native nanocrystalline
466 granulovirus using an X-ray free-electron laser. *Proc. Natl. Acad. Sci.* **114**, 2247–2252
467 (2017).
- 468 5. Suga, M. *et al.* Native structure of photosystem II at 1.95 Å resolution viewed by
469 femtosecond X-ray pulses. *Nature* **517**, 99–103 (2015).
- 470 6. Liu, H. & Lee, W. The XFEL Protein Crystallography: Developments and
471 Perspectives. *Int. J. Mol. Sci.* **20**, 3421 (2019)
- 472 7. Pandey, S. *et al.* Time-resolved serial femtosecond crystallography at the European XFEL.
473 *Nat. Methods* **17**, 73–78 (2020)
- 474 8. Gisriel, C. *et al.* Membrane protein megahertz crystallography at the European XFEL. *Nat.*
475 *Commun.* **10**, 1–11 (2019).
- 476 9. Yefanov, O. *et al.* Evaluation of serial crystallographic structure determination within
477 megahertz pulse trains. *Struct. Dyn.* **6**, 064702 (2019).
- 478 10. Aquila A, Hunter MS, Doak RB, Kirian RA, Fromme P, White TA, *et al.* Time-resolved
479 protein nanocrystallography using an X-ray free-electron laser. *Optics Express.*
480 **20**(3):2706–2716 (2012).
- 481 11. Tenboer J, Basu S, Zatsepin N, Pande K, Milathianaki D, Frank M, *et al.* Time-resolved
482 serial crystallography captures high-resolution intermediates of photoactive yellow
483 protein. *Science* **346**(6214):1242–1246 (2014).
- 484 12. Allen M. Orville Entering an era of dynamic structural biology... *BMC Biology*, **16**, 55
485 (2018)
- 486

- 487 13. Schmidt, M. Time-Resolved Macromolecular Crystallography at Pulsed X-ray Sources.
488 *Int. J. Mol. Sci.* **20**, 1401 (2019) doi:10.3390/ijms20061401
- 489 14. Grünbein, M. L. *et al.* Megahertz data collection from protein microcrystals at an X-ray
490 free-electron laser. *Nat. Commun.* **9**, 3487 (2018).
- 491 15. Volkman, B. F., Lipson, D., Wemmer, D. E. & Kern, D. Two-State Allosteric Behavior in a
492 Single-Domain Signaling Protein. *Science* **291**, 2429–2433 (2001).
- 493 16. Moy, F. J. *et al.* Assignments, Secondary Structure, Global Fold, and Dynamics of
494 Chemotaxis Y Protein Using Three- and Four-Dimensional Heteronuclear (¹³C,¹⁵N)
495 NMR Spectroscopy. *Biochemistry*, **33**, 10731-10742 (1994).
- 496 17. Feher, V. A. & Cavanagh, J. Millisecond-timescale motions contribute to the function of
497 the bacterial response regulator protein Spo0F. *Nature* **400**, 289–293 (1999).
- 498 18. Nicolaï, A., Delarue, P. & Senet, P. Decipher the Mechanisms of Protein Conformational
499 Changes Induced by Nucleotide Binding through Free-Energy Landscape Analysis: ATP
500 Binding to Hsp70. *PLOS Comput. Biol.* **9**, e1003379 (2013).
- 501 19. Xie, M. *et al.* Functional protein dynamics on uncharted time scales detected by
502 nanoparticle-assisted NMR spin relaxation. *Sci. Adv.* **5**, eaax5560 (2019).
- 503 20. Hansen, A. L. & Kay, L. E. Quantifying millisecond time-scale exchange in proteins by
504 CPMG relaxation dispersion NMR spectroscopy of side-chain carbonyl groups. *J. Biomol.*
505 *NMR* **50**, 347–355 (2011).
- 506 21. Oberthuer, D. *et al.* Double-flow focused liquid injector for efficient serial femtosecond
507 crystallography. *Sci. Rep.* **7**, 44628 (2017).
- 508 22. Gañán-Calvo, A. M. Generation of Steady Liquid Microthreads and Micron-Sized
509 Monodisperse Sprays in Gas Streams. *Phys. Rev. Lett.* **80**, 285–288 (1998).
- 510 23. DePonte, D. P. *et al.* Gas dynamic virtual nozzle for generation of microscopic droplet
511 streams. *J. Phys. Appl. Phys.* **41**, 195505 (2008).
- 512 24. Neutze, R., Wouts, R., van der Spoel, D., Weckert, E. & Hajdu, J. Potential for
513 biomolecular imaging with femtosecond X-ray pulses. *Nature* **406**, 752–757 (2000).
- 514 25. Stan, C. A. *et al.* Liquid explosions induced by X-ray laser pulses. *Nat. Phys.* **12**, 966–971
515 (2016)
- 516 26. Beyerlein, K. R. *et al.* FELIX: an algorithm for indexing multiple crystallites in X-ray
517 free-electron laser snapshot diffraction images. *J. Appl. Crystallogr.* **50**, 1075–1083
518 (2017).
- 519 27. Ginn, H. M. *et al.* TakeTwo: an indexing algorithm suited to still images with known
520 crystal parameters. *Acta Crystallogr. D. Struct. Biol.* **72**, 956–965 (2016).

- 521 28. Gildea, R. J. *et al.* New methods for indexing multi-lattice diffraction data. *Acta*
522 *Crystallogr. D Struct. Biol.* **70**, 2652–2666 (2014).
- 523 29. Liu, W., Ishchenko, A. & Cherezov, V. Preparation of microcrystals in lipidic cubic phase
524 for serial femtosecond crystallography. *Nat. Protoc.* **9**, 2123–2134 (2014).
- 525 30. Nogly, P. *et al.* Lipidic cubic phase serial millisecond crystallography using synchrotron
526 radiation. *IUCrJ*, **2**, 168–176 (2015).
- 527 31. Weierstall Uwe. Liquid sample delivery techniques for serial femtosecond crystallography.
528 *Philos. Trans. R. Soc. B Biol. Sci.* **369**, 20130337 (2014).
- 529 32. Mancuso, A. P., *et al.* The Single Particles, Clusters and Biomolecules and Serial
530 Femtosecond crystallography instrument of the European XFEL: initial installation. *J.*
531 *Synchrotron Radiat.* **26**, (2019).
- 532 33. Barty, A. *et al.* Cheetah: software for high-throughput reduction and analysis of serial
533 femtosecond X-ray diffraction data. *J. Appl. Crystallogr.* **47**, 1118–1131 (2014).
- 534 34. White, T. A. *et al.* CrystFEL: a software suite for snapshot serial crystallography. *J. Appl.*
535 *Crystallogr.* **45**, 335–341 (2012).
- 536 35. White, T. A. *et al.* Recent developments in CrystFEL. *J. Appl. Crystallogr.* **49**, 680–689
537 (2016).
- 538 36. Ganan-Calvo, A. M. The scaling of exploding liquid jets under intense X-ray pulses. *PRL*
539 **123**, 064501-5 (2019).
- 540 37. Holton, J. M. A beginner’s guide to radiation damage. *J. Synchrotron Radiat.* **16**, 133–142
541 (2009).
- 542 38. Coughlan, H. D. *et al.* Bragg coherent diffraction imaging and metrics for radiation
543 damage in protein micro-crystallography. *J. Synchrotron Radiat.* **24**, 83–94 (2017).
- 544 39. Marman, H., Darmanin, C. & Abbey, B. The Influence of Photoelectron Escape in
545 Radiation Damage Simulations of Protein Micro-Crystallography. *Crystals* **8**, 267 (2018).
- 546 40. Owen, R. L., Rudiño-Piñera, E. & Garman, E. F. Experimental determination of the
547 radiation dose limit for cryocooled protein crystals. *Proc. Natl. Acad. Sci.* **103**, 4912–4917
548 (2006).
- 549 41. Owen, R. L. *et al.* Outrunning free radicals in room-temperature macromolecular
550 crystallography. *Acta Crystallogr. D.* **68**, 810–818 (2012).
- 551 42. Gotthard, G., Aumonier, S., De Sanctis, D., Leonard, G., von Stettena, D., Royant, A.
552 Specific radiation damage is a lesser concern at room temperature. *IUCrJ*, **6**, 665–680,
553 (2019). <https://doi.org/10.1107/S205225251900616X>

- 554 43. de la Mora, E., Coquelle, N., Bury, C.S., Rosenthal, M., Holton, J.M., Carmichael, I,
555 Garman, E.F., Burghammer, M., Colletier, J.F., Weik, M. Radiation damage and dose limits
556 in serial synchrotron crystallography at cryo- and room temperatures *PNAS*, **117**, 4142-
557 4151, (2020) <https://doi.org/10.1073/pnas.1821522117>
- 558 44. Ebrahim, A., Moreno-Chicano, T., Appleby, M.V., Chaplin, A.K., Beale, J.H., Sherrell, D.A.,
559 Duyvesteyn, H.M.E., Owada, S., Tono, K., Sugimoto, H., Strange, R.W., Worrall, J.A.R.,
560 Axford, D., Owenb, R.L., Hougha, M.A. Dose-resolved serial synchrotron and XFEL
561 structures of radiation-sensitive metalloproteins. *IUCrJ*, **6**, 543–551, (2019).
562 <https://doi.org/10.1107/S2052252519003956>
- 563 45. Kern, J. *et al.* Simultaneous Femtosecond X-ray Spectroscopy and Diffraction of
564 Photosystem II at Room Temperature. *Science* **340**, 491–495 (2013).
- 565 46. Liu, W. *et al.* Serial Femtosecond Crystallography of G Protein–Coupled Receptors. *Science*
566 **342**, 1521–1524 (2013).
- 567 47. Lomb, L. *et al.* Radiation damage in protein serial femtosecond crystallography using an x-
568 ray free-electron laser. *Phys. Rev. B Condens. Matter Mater. Phys.* **84**, 214111 (2011).
- 569 48. Nass, K. Radiation damage in protein crystallography at X-ray free-electron lasers. *Acta*
570 *Crystallogr. Acta D.* **75**, 211–218 (2019).
- 571 49. Abbey, B. *et al.* X-ray laser–induced electron dynamics observed by femtosecond
572 diffraction from nanocrystals of Buckminsterfullerene. *Sci. Adv.* **2**, e1601186 (2016).
- 573 50. Amin, M., Badawi, A. & Obayya, S. S. Radiation Damage in XFEL: Case study from the
574 oxygen-evolving complex of Photosystem II. *Sci. Rep.* **6**, 36492 (2016).
- 575 51. Standfuss, J. Membrane protein dynamics studied by X-ray lasers—or why only time will
576 tell. *Current Opin. Struc. Biol.* **57**, 63–71, (2019) <https://doi.org/10.1016/j.sbi.2019.02.001>
- 577 52. McCoy, A. J. *et al.* Phaser crystallographic software. *J. Appl. Crystallogr.* **40**, 658–674
578 (2007)
- 579 53. G. N. Murshudov, P. Skubák, A. A. Lebedev, N. S. Pannu, R. A. Steiner, R. A. Nicholls, M.
580 D. Winn, F. Long, and A. A. Vagin, *REFMAC5* for the refinement of macromolecular crystal
581 structures. *Acta Crystallogr., Sect. D: Biol. Crystallogr.* **67**, 355 (2011).
- 582 54. Emsley, P., Lohkamp, B., Scott, W. G. & Cowtan, K. Features and development of Coot.
583 *Acta Crystallogr. D Biol. Crystallogr.* **66**, 486–501 (2010)
- 584 55. Chen, V. B. *et al.* MolProbity: all-atom structure validation for macromolecular
585 crystallography. *Acta Crystallogr. D Biol. Crystallogr.* **66**, 12–21 (2010).
- 586 56. Afonine, P., Grosse-Kunstleve, R.W., Chen, V.B., Headd, J.J. Moriarty, N.W. Richardson,
587 J.S. Richardson, D.C., Urzhumtsev, A., Zwart, P.H., Adams, P.D. Phenix.model_vs_data: a

high-level tool for the calculation of crystallographic model and data statistics, *J. App. Cryst.* **43**, 669-676 (2010) <https://doi.org/10.1107/S0021889810015608>

57. Koch, A., Raven, C., Spanne, P., Snigirev, A. X-ray imaging with submicrometer resolution employing transparent luminescent screens. *J. Opt. Soci. America A.* **115**, 1940-1951, (1998)

Acknowledgments

General: We acknowledge European XFEL in Schenefeld, Germany, for provision of X-ray free electron laser beam-time at Scientific Instrument SPB/SFX and use of the XBI biological sample preparation laboratory, enabled by the XBI User Consortium and would like to thank the instrument group and facility staff for their assistance and everyone who attended the first experiment.

Funding: We would also like to acknowledge the funding support for B.A., C.D, K.A.N and S.H from Australian Research Council Center of Excellence in Advanced Molecular Imaging (CE140100011), www.imagingcoe.org and the Australian Nuclear Science and Technology Organisation (ANSTO); the International Synchrotron Access Program (ISAP) managed by the Australian Synchrotron, part of ANSTO, and funded by the Australian Government. M.L.S., M.A.C. and M.F. performed, in part, under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344. M.L.S., M.A.C. and M.F. were supported by the NIH grant 1R01GM117342-01. I.B acknowledges support by Ministry of Education, Science, Research and Sport of the Slovak Republic and by grant from the Slovak Research and Development Agency under contract APVV-14-0181. M.L.S., M.A.C. and M.F. performed, in part, under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344. M.L.S., M.A.C. and M.F. were supported by the NIH grant 1R01GM117342-01. T.V is a recipient of ARC DECRA (DE170100783). B.K work is funded by National Health and Medical Research Council (NHMRC) grants 1107804 and 1160570 to BK and TV; BK is Australian Research Council (ARC) Laureate Fellow (FL180100109). F.C-M and A.M.G-C acknowledge Ministerio de Economía y Competitividad (Spain), Plan Estatal 2013-2016 Retos, project DPI2016-78887-C3-1-R. S.B.

620 acknowledged the Helmholtz Initiative and Networking Fund through the Young Investigators
621 Program and the Deutsche Forschungsgemeinschaft SFB755/B03.

622

623 **Author contributions:** The experiment was conceived by H.C., A.B. and A.P.M. as an open
624 experiment for the EuXFEL user community. A.B. coordinated and led the collaboration. B.A. and
625 C.D. conceived the idea of studying the multi-hit phenomena and edited the manuscript. S.H.
626 developed the model for double hit scenario and wrote the manuscript. H.K., A.V.M., M.H-J and
627 K.G. contributed to the beam profile data collection and analysis. M.O.W., A.B., A.P.M., K.A.N.
628 were involved in discussions. R.B. was the instrument scientist responsible for this experiment.
629 A.P.M. led development of the SPB/SFX instrument. M.O.W. and S.B. led the sample injection
630 team. C.D., D.O., R.S., I.B., M.A.C., M.F., M.S., T.V. B.K., K.D., R.L., G.P., S.H., M.M., G.P.,
631 P.L.X., H.F. and C.S. contributed to sample preparation. S.B., A.B., D.C.FM., M.T., M.O.W., J.S.,
632 J.B., R.G., H.F., and L.A. contributed to sample injector delivery. A.B. led the data processing team
633 and H.C., M.S., O.Y., A.M., K.A., A.T., and P.V., contributed to data processing. A.P.M., R.B.,
634 J.S., S.S., K.G., Y.K., M.M., G.M., A.R., P.V., B.W., and J.S., developed and operated the
635 SPB/SFX instrument at EuXFEL. F.C-M., A.M.G-C., S.M., J.A.S., B.A. M.A.C., M.F., and M.S.,
636 participated in the XFEL 2012 experiment. T.L., K.W., L.M., G.P., N.A-Q., S.B., T.M., C. Dan.,
637 A.S., C.X., S.Ha., developed XFEL DAQ and controls. J.S-D., T.L., A.K. and A.M., developed and
638 responsible for AGIPD. The manuscript was written by S.H., C.D. and B.A. with input from all
639 authors.

640

641 **Competing interests:** The authors declare no competing interest.

642

643

644

645

646

647

648

649

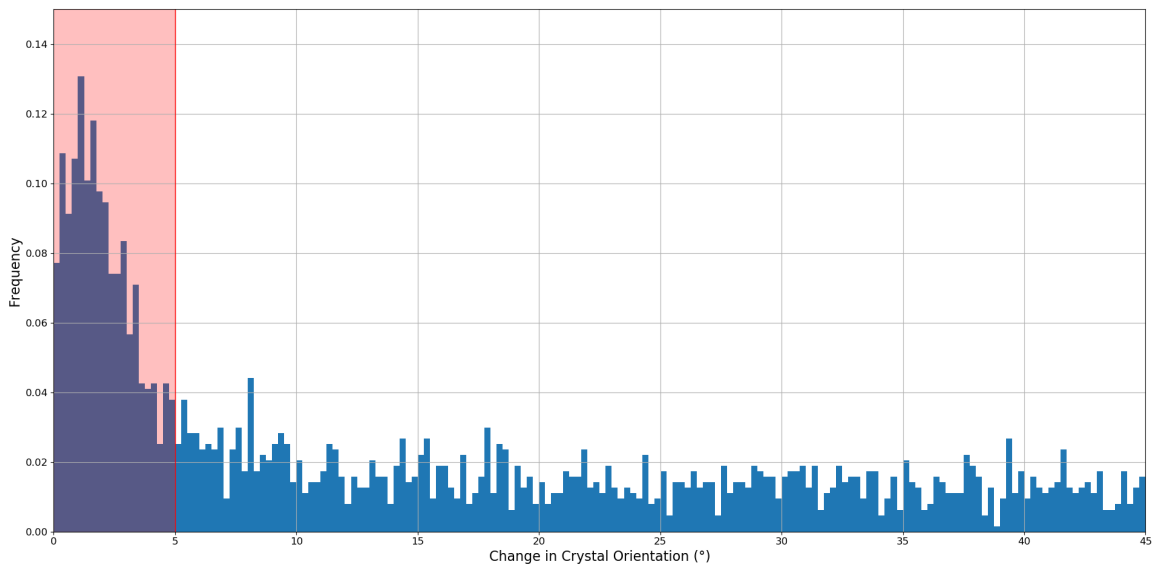
650

651

652

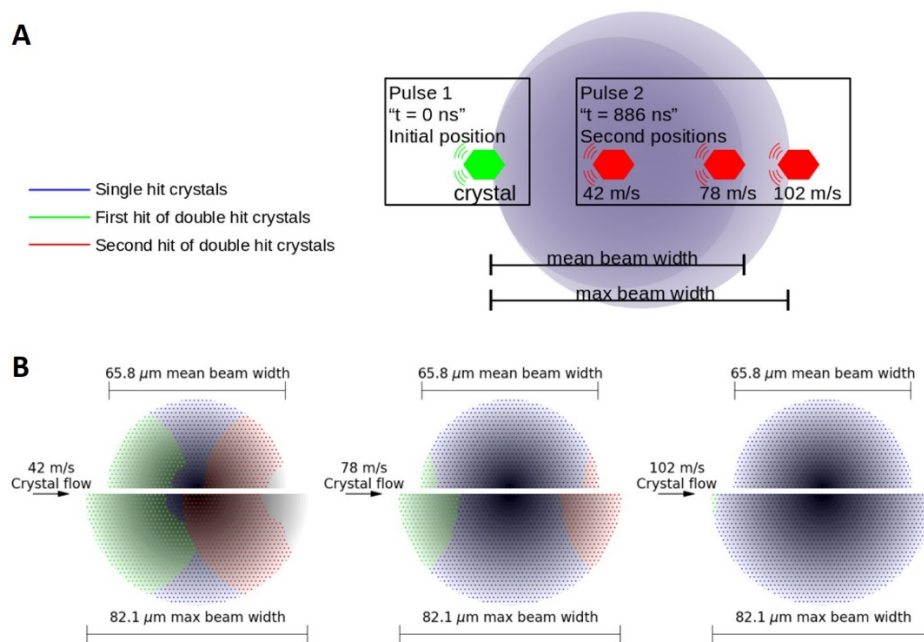
653

654
655
656
657
658
659
660
661 **Figures and Tables**



662
663
664
665
666
667
668
669
670
671
672

Fig. 1. Frequency of diffraction patterns as a function of the relative change in crystal rotation. This is characterized by reciprocal space vector \vec{a} , between consecutive diffraction measurements (separated by 886 ns) within the X-ray pulse train. The liquid jet flow rate was 42 m/s. An increase in the frequency of images with a change in orientation of less than 5 degrees can be observed due to the occurrence of double hits.



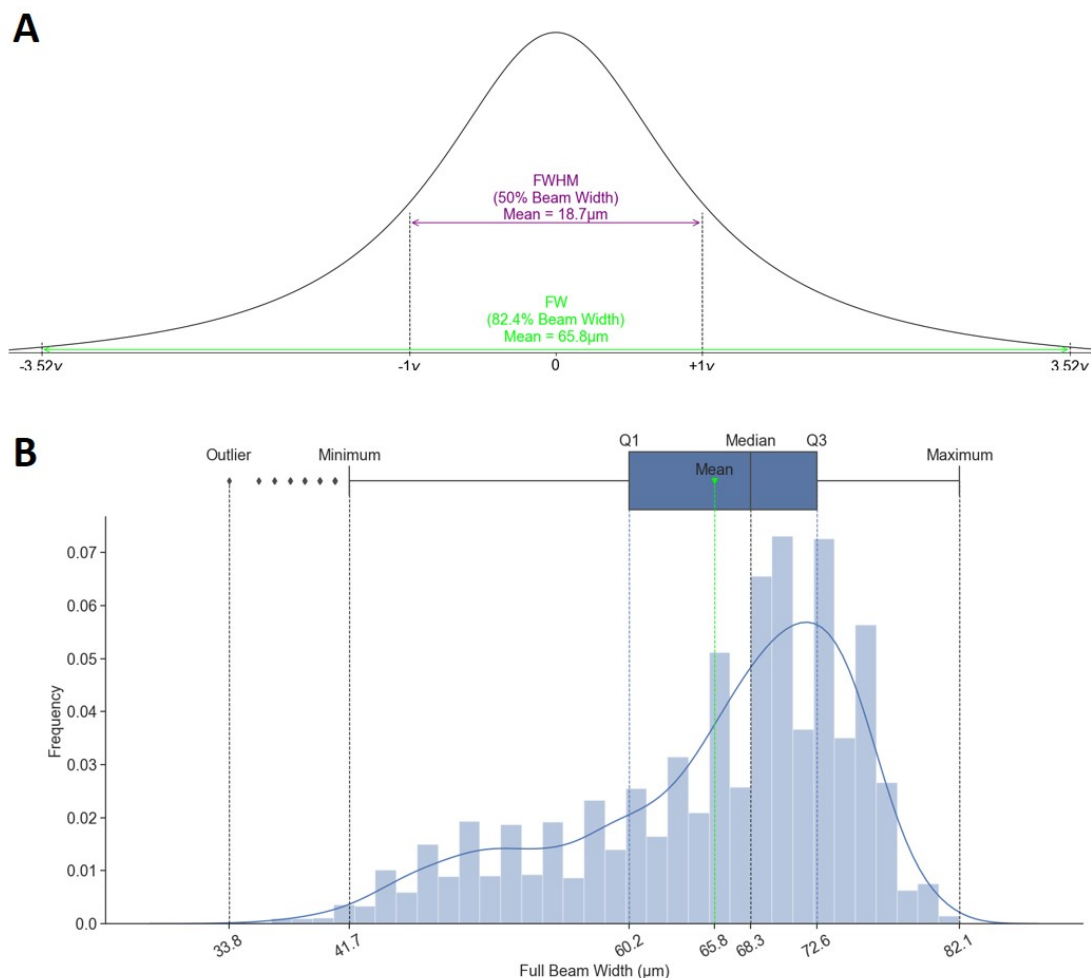
673

674

675 **Fig. 2. Distance travelled by the crystal.** **A.** Schematic (not to scale) illustrating the minimum
 676 distances travelled by an $8\mu\text{m}$ crystal for the different jet speeds overlaid with the average beam
 677 FW (dark purple shaded region) and the max beam FW (light purple shaded region). All three jet
 678 speeds are represented, the green crystal depicts the initial position, and the red crystal illustrates
 679 how far the crystal travels after the first hit for 42 m/s, 78 m/s, and 102 m/s jet speeds. **B.** Schematic
 680 representation of the crystal path through the X-ray beam for each of the three jet speeds for the
 681 mean beam FW (upper half) and maximum beam FW (lower half). The beam profile (shaded grey),
 682 is overlaid with the regions that the crystal travels through for the single hits (blue) as well as the
 683 first (green) and second (red) hits of the double hit crystal. Note for 42 m/s and 78 m/s, regions
 684 where no hits occur are possible.

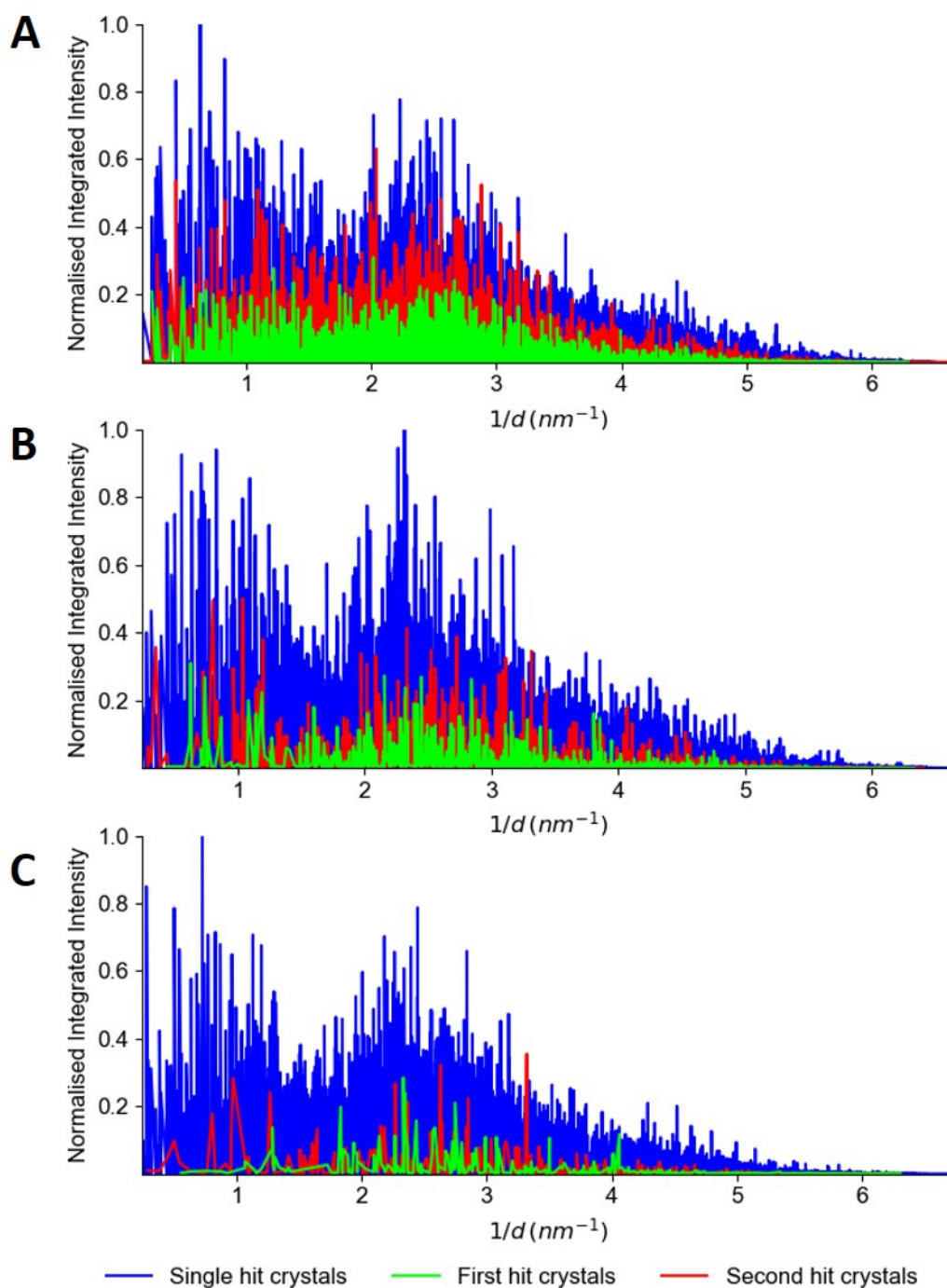
685

686



687

688 **Fig. 3. Characteristic beam profile. A.** A histogram showing the X-ray Beam profile. The beam
 689 profile was modelled using a Lorentzian distribution with a Full Width Half Maximum (FWHM)
 690 = 2γ (50% of beam) and Full Width (FW) = 7.04γ (82.4% of beam). **B.** The Lorentzian
 691 distribution used to determine the FWHM and FW is shown for all 6773 YAG images obtained.
 692 The furthest outlier, minimum, Q1 (25th percentile), mean, median, Q3 (75th percentile), and
 693 maximum have been indicated. To obtain the FWHM and FW for each YAG image, a 7.5% noise
 694 threshold was applied to the image combined with a 3×3 median filter to account for the noise.



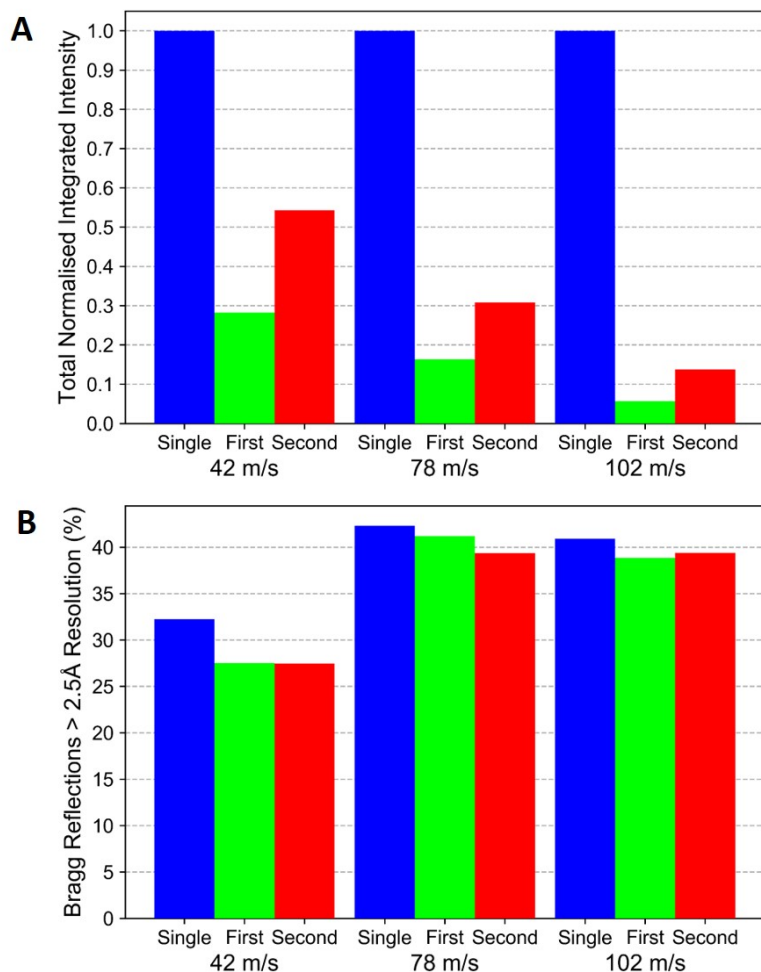
695

696

697 **Fig. 4. Normalized integrated intensity plots.** Integrated intensities were extracted from the data,
 698 normalized and plotted against $1/d$ (where d is the lattice spacing) for the **a.** 42 m/s jet data, **b.** 78
 699 m/s jet data, and **c.** 102 m/s jet data. Blue represents data for single hit crystals only; green represents
 700 the first hit of the double hit crystal; red represents the second hit of the double hit crystal. A
 701 threshold of $I/\text{sig}(I) > 2$ was applied to the analysis.

702

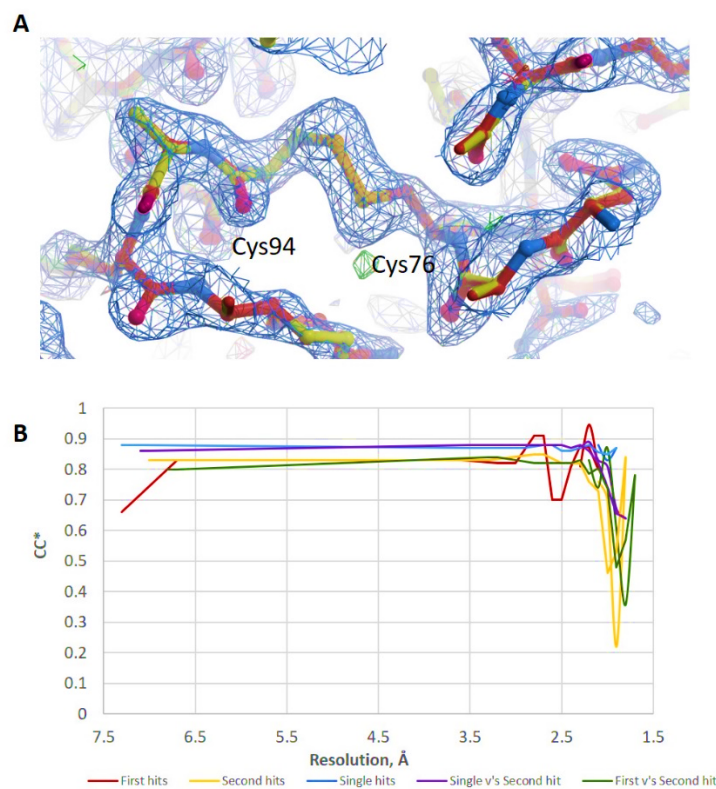
703



705

706 **Fig. 5. Comparison of the data between the single and double hit crystals. A.** The area under
 707 the curves in Fig. 4 were obtained and are seen here in order to more clearly show the differences.
 708 For the slowest jet speed, 42 m/s, the second hit had the most similar results compared to the single
 709 hit data as both crystals are hit within the FWHM region of the X-ray beam. For the faster jet speeds
 710 (78 and 102 m/s) the first and second hit data are similar in intensity as the crystal is hit within the
 711 tail region of the x-ray beam and not within the FWHM as occurs for the single hit crystal. **B.** This
 712 shows the percentage of the reflections with a better than 2.5 Å resolution for each hit type and jet
 713 speed. Note: The resolution for the fast jets is perceived to be higher in this figure but this is not
 714 statistically significant as the amount of data for each jets speed varies greatly, as denoted in Table
 715 1. The data-sets were also collected on different days which resulted in varying beam intensities. A
 716 threshold of $I/\text{sig}(I) > 2$ was applied to the analysis.

717

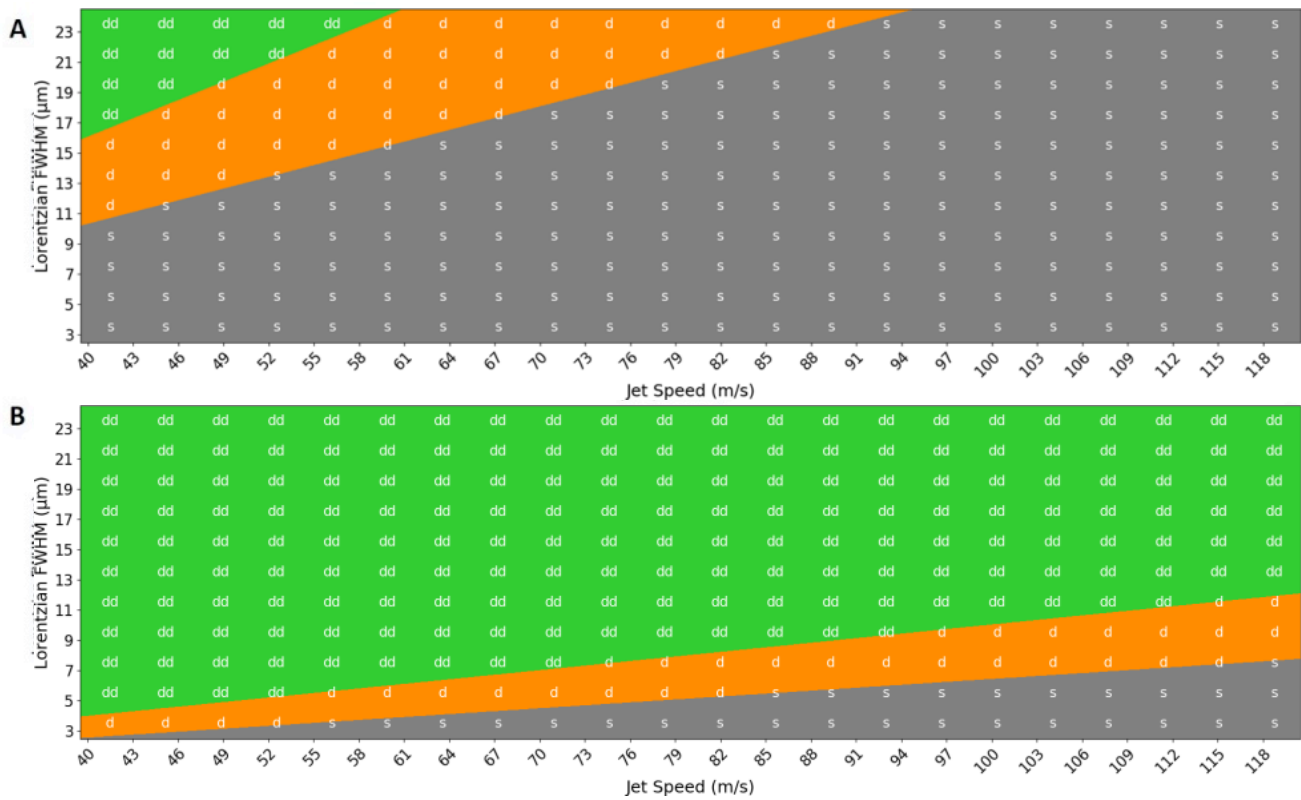


718

719 **Fig. 6. Lysozyme structural analysis.** A. Electron density map of the Lysozyme double hit
 720 structure showing no radiation damage is present surrounding the disulfide bond between Cys94
 721 and Cys76, for the second hit data. The $2F_o-F_c$ map at 1 sigma is shown in blue and overlaid with
 722 the Fourier difference (F_o-F_c) map at 3 sigma (green). B. CC* for data separated into single hits
 723 (blue), first hits (red) and second hits (yellow) as well a correlation of merged data of the second
 724 hit relative to the first hit (green) and single hit (purple) data sets indicating they are very similar.

725

726



727

728

729 **Fig. 7. Parameters for optimizing the collection of double-hit data at the EuXFEL.** A. EuXFEL
 730 repetition rate of 1.1 MHz during this experiment and B. A Future capability of EuXFEL where a
 731 4.4 MHz repetition rate is possible. The green shaded area (dd) indicates parameter combinations
 732 that will result in double-hits that allow the second hit to occur within the FWHM of the beam, the
 733 orange shaded area (d) indicates parameter combinations that will result in double-hits that allow a
 734 second hit to occur within in the tail region of the beam, and the grey shaded area (s) indicates
 735 parameter combinations that will result in only single hits. This was modeled independent of crystal
 736 size (i.e. crystal center-to-crystal center hits).

737

738

739

740

741

742

743

744

745

746

747

748 **Table 1. Summary of jet speeds, experimental conditions and statistics for lysozyme crystals.**

749 Based on Wiedorn et al. (2018)¹, this table shows the flow rates of the gas and liquid used as well
 750 the experimentally and theoretically determined jet speeds for the current analysis. The data
 751 statistics were calculated utilising the *CrystFEL* software suite^{34,35}. As can be seen below, as the jet
 752 speed increases the number of double hits decrease.

753

Target jet speed	50 m/s	75 m/s	110 m/s*
Liquid Flow ($\mu\text{l}/\text{min}$)	15	13	13
Gas Flow (mg/min)	23	50	85
Experimental jet speed (m/s)	42 ± 2.1	78 ± 3.9	-
Theoretical jet speed (m/s)	-	-	102 ± 5.1
Total no. Images	440,000	60,000	240,000
No. hits	10,726 (2.4%)	1,638 (2.7%)	3,733 (1.6%)
No. Indexed Frames	9,970 (93%)	1,509 (92.1%)	3,474 (93.1%)
No. double hits	595 (6.3% crystals)	14 (0.9% crystals)	10 (0.3% crystals)
No. single hits	8,780 (93.7% crystals)	1,481 (99.1% crystals)	3,454 (99.7% crystals)

754

755 * A '100 m/s' jet speed was determined experimentally for a 13 $\mu\text{l}/\text{min}$ liquid and 80 mg/min gas
 756 flow rate giving a speed of 105 m/s, which is similar to the theoretically calculated 102 m/s jet
 757 speed.

758

759

760

761

762

763

764

765

766

767 **Table 2:** SFX and refinement statistics for lysozyme single hit crystal data and lysozyme second
 768 hit crystal data.

Parameter	Single hit	Second hit
Photon energy (mean value)	9232 eV	9232 eV
X-ray focus		
Pulse energy at sample (assuming 50% beamline transmission)	290 μ J	290 μ J
Pulse length	50 fs	50 fs
Space group	P4 ₃ 2 ₁ 2	P4 ₃ 2 ₁ 2
Unit cell		
<i>a</i> , <i>b</i> , <i>c</i> (Å)	79.6, 79.6, 37.73	79.6, 79.6, 37.73
α , β , γ (°)	90, 90, 90	90, 90, 90
No. of hits/indexed lattices	10,106 / 10,725	962 / 1,073
No. of unique reflections	7,418 (535)	7,263 (494)
CC _{1/2}	0.906 (0.796)	0.615 (0.413)
CC*	0.975 (0.942)	0.873 (0.764)
Data Refinement		
Resolution range (Å)	21.66-2.10 (2.155-2.10)	21.66-2.10 (2.155-2.10)
Completeness (%)	99.84% (100%)	97.75% (92.34%)
Redundancy	47.24 (28.34)	6.66 (3.98)
<i>I</i> / σ (<i>I</i>)	5.1 (4.2)	2.9 (3.1)
<i>R</i> _{work}	0.152 (0.114)	0.249 (0.244)
<i>R</i> _{free}	0.216 (0.183)	0.299 (0.355)
Rmsd bonds (Å) /Rmsd angles (°)	0.0007 / 1.453	0.004 /1.254
Ramachandran (%)		
favored	98.43%	96.85%
allowed	1.57%	3.15%
outliers	0.00%	0.00%
Wilson B-factor (Å²)		
Average	19.58	20.4
Protein	20.61	21.44
Ligands	37.73	37.24
Ions	22.99	40.32
Waters	29.96	28.01
PDB code		
CXIDB data deposition		

769 Note: Statistics for the highest-resolution shell are shown in parentheses

770

771

772
773
774
775
776
777
778
779
780
781
782
783
784