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# A New Crater Near InSight: Implications for Seismic Impact Detectability on Mars

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## Key Points

- A new 1.5 m diameter impact crater formed on Mars ~40 km from the InSight lander between February and April 2019.
- Three candidate seismic events occurred during this time frame, but none of them can be definitively associated with the new crater.
- We revise our expectations for InSight impact detections above the background noise to be ~2 per Earth year, with large uncertainties.

**Index terms:**

5420 PLANETARY SCIENCES: SOLID SURFACE PLANETS - Impact phenomena, cratering (6022, 8136)  
 7299 SEISMOLOGY - General or miscellaneous  
 8136 TECTONOPHYSICS - Impact phenomena (5420, 6022)

**Abstract**

A new 1.5 meter diameter impact crater was discovered on Mars only ~40 km from the InSight lander. Context camera images constrained its formation between February 21 and April 6, 2019; follow-up HiRISE images resolved the crater. During this time period, three seismic events were identified in InSight data. We derive expected seismic signal characteristics and use them to evaluate each of the seismic events. However, none of them can definitively be associated with this source. Atmospheric perturbations are generally expected to be generated during impacts; however, in this case, no signal could be identified as related to the known impact. Using scaling relationships based on the terrestrial and lunar analogs and numerical modeling, we predict the amplitude, peak frequency, and duration of the seismic signal that would have emanated from this impact. The predicted amplitude falls near the lowest levels of the measured seismometer noise for the predicted frequency. Hence it is not surprising this impact event was not positively identified in the seismic data. Finding this crater was a lucky event as its formation this close to InSight has a probability of only ~0.2, and the odds of capturing it in before and after images is extremely low. We revisit impact-seismic discriminators in light of real experience with a seismometer on the martian surface. Using measured noise of the instrument, we revise our previous prediction of seismic impact detections downwards, from ~a few to tens, to just ~2 per Earth year, still with an order of magnitude uncertainty.

**Plain Language Summary**

A small new impact crater was discovered on Mars very close to the InSight lander. Photographs from a camera in orbit show it formed between February 21 and April 6, 2019. Three seismic events were detected by InSight during this time. We estimate what seismic data from the impact would have looked like and whether or not each of the seismic events was caused by the new impact, but none of them can be definitely linked. We predict the size, frequency, and length of time of the signal that would have come from this impact. Even though this impact is very close to InSight, it's small, so it was not a large seismic event. The signal would be near the quietest the instrument ever gets. There is only a 1 in 5 chance per Earth year that a crater would have formed this close to InSight, and a much lower chance that it would be imaged, thus we were very lucky to find this crater. Using what we know about the instrument on the ground, we update the number of impacts we expect to find with InSight to ~2 each Earth year, with a lot of uncertainty.

## 1 A new impact constrained by orbital images

On April 6, 2019, an image taken by the Context camera (CTX; Malin *et al.*, 2007) on the Mars Reconnaissance Orbiter (MRO) revealed a new dark spot that was not present in a previous image taken on February 21 (Fig. 1), only ~40 km from the newly-landed InSight mission (Interior Exploration using Seismic Investigations, Geodesy and Heat Transport; Smrekar *et al.*, 2019). Detecting an impact in both seismic data and orbital images would be an exciting development, leading to a number of scientific advances (Daubar *et al.*, 2018). This would be a seismic source with a known location, and thus a known distance and direction. A certain location and depth would allow modeling of seismic ray paths through the interior that could constrain seismic velocities and the physical properties of the material through which the rays traveled. This would improve models of interior structure and the seismic attenuation of Mars. An impact clearly observed in both orbital and seismic data would also provide a calibration of the seismic source parameters such as moment, cutoff frequency, and seismic efficiency (the ratio of impact energy to radiated seismic energy). The seismic efficiency, for example, is not well constrained, with values in the literature ranging from  $10^{-6}$  to  $10^{-2}$  (Daubar *et al.*, 2018 and references therein). High resolution images of newly formed craters would characterize crater sizes, leading to an empirical relationship between impact size and observed seismic amplitudes. Enough such observations would also result in an independent measurement of the current impact rate, anchoring absolute bombardment rates. Thus identifying an impact in seismic data that was also imaged from orbit would satisfy many important scientific goals. So naturally, this event was of immediate interest to the InSight team.

A high-resolution 25 cm pixel scale image from the High Resolution Imaging Science Experiment (HiRISE; McEwen *et al.*, 2007) was acquired shortly thereafter. The HiRISE image resolved a ~1.5 meter diameter impact crater at the location of the new dark spot (Fig. 1D), showing that an impact event occurred in the short period of time constrained by the before and after CTX images, between 21 February (03:56:17 UTC) and 6 April (08:19:17 UTC) in 2019. This occurrence is not especially rare; ~900 new dated impacts have been discovered in the last ~decade on Mars using similar techniques (Malin *et al.*, 2007; Daubar *et al.*, 2013), although the imaging date constraints are usually on the order of a few years rather than a month. This impact was also extraordinary in its location very close to the recently-landed InSight mission. At this distance, the prospect of detecting the impact event using the seismic and atmospheric instrumentation on InSight was an exciting possibility. This is the only impact we know to have formed this close to the lander during the time since InSight landed on Mars on 26 November 2018.

The new crater is located at 3.866°N planetocentric latitude, 135.613°E longitude, just 37.36 km from InSight, which landed at 4.502°N, 135.623°E (Parker *et al.*, 2019). It is located along an azimuth of 180.9°, almost directly south of the lander. The asymmetric low-albedo blast zone pattern around the crater (Fig. 1D), caused by the disturbance of light-toned dust during the impact, indicates a somewhat oblique impact coming from the southwest direction. Small dark spots to the southwest of the crater could be blast zones around secondary craters or multiple smaller primary craters in a clustered impact that formed when the impactor fragmented in the atmosphere (Daubar *et al.*, 2019). Craters within these smaller dark spots are not resolved. The pattern of dark spots is more consistent with a clustered impact than with secondary craters;

secondary craters would be concentrated downrange rather than uprange, and typically have more symmetric radial patterns. In either case, the contribution of the group of smaller craters to a combined seismic signal would be negligible compared with that of the main ~1.5 m diameter crater (Schmerr *et al.*, 2019).

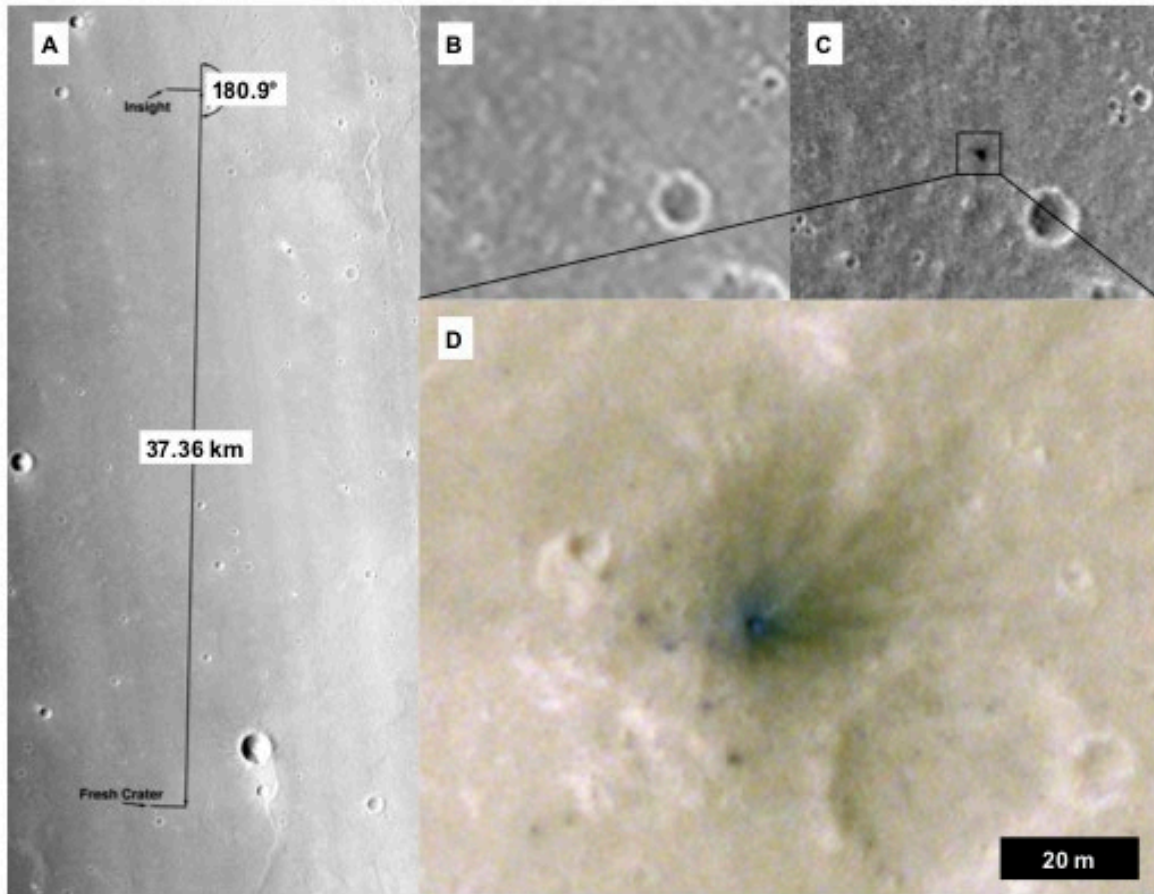


Figure 1.

*New crater observations. (A) CTX context image showing locations of InSight lander and new dated impact. (B) CTX image K14\_068929\_1845\_XN\_04N224W\_190221 taken February 21, 2019 (6 m/px). (C) CTX image K16\_059495\_1829\_XN\_02N224W\_190406 taken April 6, 2019, showing new dark spot that was not present in previous image. (D) Cutout from HiRISE image ESP\_060128\_1840 (COLOR RDR; 25 cm/px) showing new impact crater. North is up, and images have been stretched for contrast. Image credits: NASA/JPL/MSSS (CTX); NASA/JPL/University of Arizona (HiRISE).*

A second HiRISE image was acquired to obtain stereo data, but the crater is not resolved in the resulting Digital Terrain Model (DTM). (See anaglyph in Fig. S1.) A depth of a few tens

of centimeters is estimated for the new crater. Although this depth is not resolved in the DTM, an estimate was possible by scaling from larger, resolved, craters in the DTM.

In subsequent sections we derive the expected seismic and atmospheric signals that would have been produced by this known impact and have the potential to have been detected by InSight (Section 2). In Section 3, we describe the search of the seismic data during the time period constrained by the before and after CTX images, and the three candidate seismic events that were found. We then evaluate which of those seismic events, and any associated atmospheric signals, might be connected with the formation of the new crater. Finally, in Section 4, we use InSight mission experience thus far to re-evaluate the seismic impact discriminators we identified before landing, and we present updated expectations for impact detections with InSight in light of real data acquired since landing.

## 2 Predicted signals from the new impact crater

### 2.1 Predicted impact parameters from the observed crater

To assess the detectability of the observed  $\sim 1.5$  m diameter crater by InSight, we first estimate the impactor parameters. The geology of the impact target area is very similar to that in the immediate vicinity of the InSight lander, which has been characterized in detail (Golombek *et al.*, 2020). The material in which the crater formed is likely to be a loose, porous regolith with very low cohesive strength ( $\lesssim 50$  kPa). The diameter of meter-scale impact craters formed in such a material is expected to scale as a power of the vertical impactor momentum, with only minor additional dependence on other impactor parameters (Holsapple, 1993; Holsapple and Housen, 2007). For a  $1.5 \pm 0.25$  m diameter crater, the predicted vertical impactor momentum is 100-3000 Ns, depending on the cohesive strength of the regolith (Fig. S2). The lower limit applies if the martian regolith can be represented as cohesionless dry sand; a nominal upper limit applies if the martian regolith has an effective cohesive strength of 50 kPa. An even higher impactor momentum is possible, but that would require a cohesive strength of a well-cemented terrestrial soil, which is not compatible with observations of the martian regolith made in the vicinity of the InSight lander (Golombek *et al.*, 2020).

The seismic source of the impact can be expressed as an equivalent seismic moment, which scales approximately linearly with impactor momentum according to two independently derived, semi-empirical scaling relationships (Shishkin, 2007; Gudkova *et al.* 2011, 2015; reviewed in Daubar *et al.*, 2018). For an impactor momentum of 100-3000 Ns, these relationships predict an equivalent seismic moment of  $10^6$ - $10^7$  Nm (Fig. S3).

The estimated impactor momentum implies an impactor mass  $\sim 0.1$  to  $\sim 1$  kg, depending on impact speed. Meteoroids in this mass range are substantially decelerated by Mars' atmosphere (Fig. S4) and are predicted to lose approximately 90% of their initial kinetic energy, 75% of their initial speed, and 30% of their initial mass by ablation and drag before striking the ground (Table S1). Thus vertical impact speeds at the ground in the range of only 1-3 km/s are expected for typical pre-entry meteoroid encounter speeds of 5-15 km/s (Le Feuvre and Wieczorek, 2008; JeongAhn and Malhotra, 2015) and entry angles of 15-90°. At these relatively

slow impact speeds, and taking into account the uncertainty in impactor momentum, estimates of the impact energy range from approximately 0.1 MJ to 2 MJ (see supplemental section S1).

An independent test of these energy estimates is provided by the empirical relationship from Teanby and Wookey (2011) between crater diameter ( $D$ ) and impact energy ( $E$ ), based on laboratory and field impact experiments, explosive analogues, and the Apollo artificial lunar impacts:

$$D = 8.8_{-3.5}^{+2.6} \times 10^{-3} E^{0.32 \pm 0.01} \left( \frac{g_{\oplus}}{g} \right)^{3/16}, \quad (1)$$

where  $g_{\oplus}$  is Earth's gravity ( $9.81 \text{ ms}^{-2}$ ) and  $g$  is Mars' gravity ( $3.73 \text{ ms}^{-2}$ ). The error bars incorporate scatter in the source data and the uncertainties in impact conditions. Using this relationship gives an estimated ground impact energy of  $5.3 \pm 1.8 \text{ MJ}$ , which is somewhat larger than our previous estimate. We attribute this difference to the fact that most of the data used to construct Eq. (1) are from experiments in terrestrial soils and rocks that have a much higher cohesive strength than the strength we adopt for the martian regolith based on in situ and remote sensing of this region. Therefore, this scaling relationship provides an upper bound on the impact energy.

## 2.2 Predicted seismic signals based on energy and moment scaling

The estimated ground impact energy can be used to obtain a first order prediction of seismic P-wave amplitude  $v$  at source-receiver distance  $x$  using scaling relations developed for terrestrial impacts (Teanby, 2015):

$$v(x, E) = ax^b E^c \quad (2)$$

where scaling law constants  $a=5.6 \times 10^{-5}$ ,  $b=-1.6$ , and  $c=0.5$  under Mars conditions (Teanby, 2015). The overall uncertainty on  $v(x, E)$  is a factor of four. This relationship is strictly only valid over the range of energies and distances used by Teanby (2015), which cover  $\sim 400\text{-}10,000 \text{ kg TNT equivalent}$  ( $\sim 2 \times 10^3\text{-}4 \times 10^4 \text{ MJ}$ ) (excluding the very high energy buried nuclear explosions) and  $0.5\text{-}1200 \text{ km}$  ranges. These events had peak seismic frequencies in the range  $1\text{-}16 \text{ Hz}$ , with the Apollo lunar and Carancas Earth impacts peaking from  $1\text{-}10 \text{ Hz}$ . We can also estimate the longest timescale in the source function using crater excavation timescale,  $\tau = \sqrt{D/g}$   $\sim 0.6 \text{ s}$ , implying a frequency content of  $>1 \text{ Hz}$ . Therefore, the scaling relationship is a reasonable, although not ideal, match to conditions for the new martian crater, with the P-wave frequency content likely peaking at a few Hz or slightly higher.

A first order prediction of seismic P-wave amplitude  $v$  for the new event is shown in Figure 2 compared to the range of measured InSight noise levels in the  $1\text{-}16 \text{ Hz}$  bandpass from Lognonné *et al.* (2019, 2020). The estimated P-wave amplitude at the observed range of  $37 \text{ km}$  is  $0.8\text{-}4 \times 10^{-9} \text{ ms}^{-1}$  for  $0.1\text{-}2.0 \text{ MJ}$  and  $6 \times 10^{-9} \text{ ms}^{-1}$  with a factor of four uncertainty for the  $5.6 \text{ MJ}$  upper bound. Furthermore, Wójcicka *et al.* (2020) use numerical impact simulations to propose a recasting of the amplitude scaling in terms of impact momentum instead of energy, which relates



to seismic moment more closely to linearly. When applied to impacts in relevant analog materials, this recasting results in a reduction in predicted seismic amplitudes by up to two orders of magnitude for small craters. Overall, these scaling laws have large uncertainties, and predictions span three orders of magnitude, but all imply a modest signal-to-noise ratio (SNR), with a likely SNR of only  $\sim 1$  on average. These amplitude estimates are also in reasonable agreement with peak ground velocities predicted from numerical waveform simulations of the impact event (see supplemental section S2). During the detection period, the continuous seismometer data coverage is limited to 10 sps (5 Hz Nyquist) sampling except during exceptional periods where 20 sps or 100 sps was collected (Fig. S10). Therefore, any seismic energy over 5 Hz is unlikely to have been recorded for the majority of the time in question. The combination of low SNR, high frequency content, and low sample rate implies this event would have been very difficult to detect seismically.

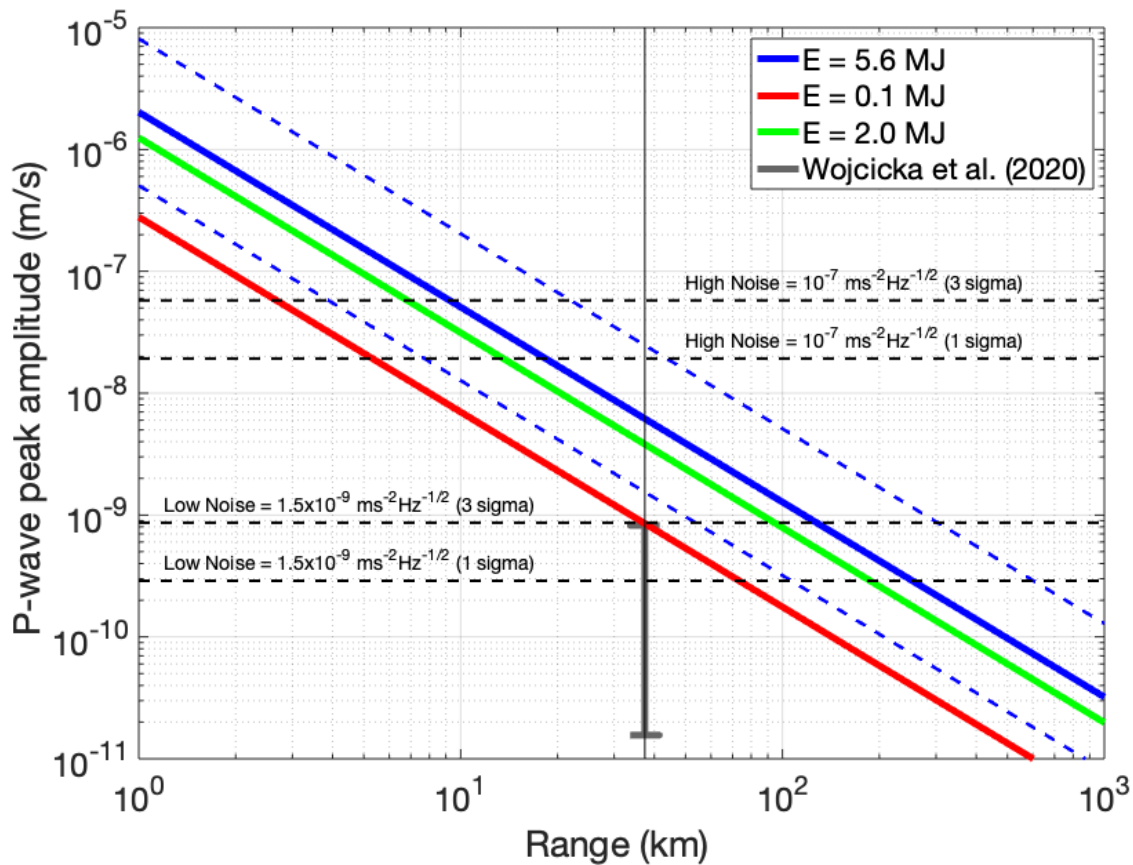


Figure 2.

Estimated amplitude of P-wave signal from the 1.5 m diameter new impact. The amplitude is estimated using the impact energy scaling relationship from Teanby (2015) in equation (2) as described in the text. Solid lines show nominal amplitude prediction from scaling relations and uncertainty for three potential impact energies, a nominal range (red to green) and an upper limit (blue). Dashed blue lines show uncertainty on the upper limit  $E=5.6$  MJ case. Horizontal

dashed lines show range of seismic noise measured at 4 Hz at the InSight landing site for 1 and 3 sigma (Lognonné et al., 2019, 2020), and vertical black line shows the distance between InSight and the new crater. Gray vertical bar shows range of predictions from Wójcicka et al. (2020) from numerically derived impact momentum scaling. Seismic noise amplitudes are converted to equivalent velocities by integrating the amplitude spectral density (ASD) noise for the 1-16 Hz bandwidth, using equation 14 in Teanby (2015). Amplitude estimates for the observed impact are at or below the noise levels.

### 2.3 Predicted seismic signals based on lunar impact analogies

The closest seismic analog for this impact is the Lunar Module of Apollo 14 (LM), which impacted 67 km from the Apollo Lunar Surface Experiments Package (ALSEP) station of Apollo 14. Its amplitude was about 40 data units (DU) on the vertical Long Period (LP) axis in peaked mode, corresponding to 2 nm of ground displacement at 2 sec. See Lognonné et al. (2009) for a detailed analysis of this and other lunar impacts.

At such a small epicentral distance, intrinsic attenuation can be neglected, and the seismic signal is mostly constrained by the elastic propagation properties, which are mostly diffusive on the Moon, and the source parameters. These parameters are summarized in Table 1.

| Impact                   | Distance from seismometer (km) | Velocity (km/s) | Angle (° from vertical)                  | Mass (kg) | Mv (Mv <sub>z</sub> ) (kg m/s)                                                                | Rim Diam (m) | Depth (m)        | Formation time |
|--------------------------|--------------------------------|-----------------|------------------------------------------|-----------|-----------------------------------------------------------------------------------------------|--------------|------------------|----------------|
| LM impact on the Moon    | 67                             | 1.68            | 86.4°                                    | 2383      | 4×10 <sup>6</sup><br>(2.5×10 <sup>5</sup> )                                                   | 6.5          | 1.37             | 0.94 sec       |
| New 1.5-m crater on Mars | 37.4                           | 1-3             | Not well constrained; moderately oblique | 0.1-1     | 1.4 ×10 <sup>2</sup> -<br>4.3 ×10 <sup>3</sup><br>(1×10 <sup>2</sup> -<br>3×10 <sup>3</sup> ) | 1.5          | A few tens of cm | ~0.30-0.35 sec |

Table 1.

Comparison between the source parameters of the Apollo 14 Lunar Module (LM) impact and the CTX-image-constrained impact that formed the new 1.5-m diameter crater discussed in this paper. Parameters from the LM impacts are from Lognonné et al. (2009) and references therein. A 45° impact angle is assumed for the martian impact, although this is only weakly constrained. Formation time is estimated from Holsapple (1993) and using  $0.5\sqrt{(D/g)}$  as an estimate of the

crater growth time (Schmidt and Housen, 1987). Known values are given in bold, other values are inferred.

According to Sato and Korn (2007), the maximum amplitude of a pulse propagating in the multiple forward scattering regime is proportional to  $\sqrt{\frac{1}{xT_m}}$ , where  $x$  is the hypocentral distance and  $T_m$  is a characteristic time scale.  $T_m$  depends on the heterogeneity of the medium as:

$$T_m = \sqrt{\pi} \frac{\langle \varepsilon^2 \rangle D^2}{2\alpha\beta} \quad (3)$$

Where  $\beta$  is the wave propagation speed,  $\alpha$  is the correlation length of the random fluctuations, and their variance is  $\langle \varepsilon^2 \rangle$ . This theory predicts that the typical maximum amplitude is proportional to  $\sqrt{\frac{\alpha\beta}{x^3 \langle \varepsilon^2 \rangle}}$ . Note these formulae are valid in media with velocity and density with gaussian fluctuations (Sato and Korn, 2007).

We do not expect the correlation distance to differ significantly between Mars and the Moon, but fluctuations are certainly stronger on the Moon because scattering is stronger. As the diffusivity is inversely proportional to  $\langle \varepsilon^2 \rangle$ , we expect the amplitude to be 5 to 10 times larger on Mars than on the Moon, for the same source and distance, following initial comparisons of the crustal diffusivity (Lognonné *et al.*, 2020).

With these assumptions, we can convert amplitudes of impacts detected on the Moon to the martian situation. Following previous work (Lognonné *et al.*, 2009; Gudkova *et al.*, 2011; 2015), we assume that the amplitude of the signal is linearly related to the vertical momentum, which implies a source for the martian impact smaller in moment than the LM source by a factor of ~83-2500. On the other hand, the difference in diffusion makes the maximum amplitude of the signal larger by a factor 5-10. Last but not least, the difference in distance for the LM impact at 67 km makes the signal larger by a factor of 2.37 for an -1.5 exponential decay, comparable to the -1.6 power law decay of local magnitudes on Earth at short distance (Richter, 1958). Combining these factors, this suggests a martian signal smaller than the lunar one by a factor of 8.3-500 without a geometrical spreading correction; with that correction, it would be smaller by a factor of 3.5-210.

The duration of the signal can also be addressed with similar analogies. Martian signals are expected to have much shorter durations than lunar ones due to the ratio of diffusivities. Rise times are found to be in the range of 600-800 sec for lunar impacts (Gillet *et al.*, 2017) and are expected to be reduced by a factor of 30-100 for Mars. Signals with SNR of 3 will have durations of about 2-3 times the rise time, leading to durations in the range of 20-60 seconds for each phase in this case.

In summary, based on early estimates of the diffusivity of Mars, we expect this impact on Mars to have a signal smaller in amplitude by a factor of 3.5 to 210 compared to the Apollo 14

LM impact recorded by the Apollo 12 vertical LP instrument. Martian impact signals are also expected to have much shorter durations of ~20-60 seconds.

In order to estimate propagation time differences between the main phases, we use the fact that an event at a distance of 37 km in a homogenous martian crust will propagate down to a depth of ~500 m. We can expect that most of the energy of this event will therefore be guided in the first half km of depth, for which seismic velocities are expected up to 2000 m/s for P waves and 1000 m/s for S waves. Most of the energy will be in surface waves, for which typical group velocities are computed with Mineos software (<https://geodynamics.org/cig/software/mineos/>) (Fig. 3a). These are shown for one possible model of the shallow subsurface structure based on constraints proposed by Lognonné *et al.* (2020) for the first five meters of depth, measurements of the seismic velocities of layers of volcanic material (Lesage *et al.*, 2018) down to 1 km depth, and the TAYAK reference model below that (Smrekar *et al.*, 2019). See supplemental section S3 for details of the model. Note that apart from the first five meters, this model is merely representative, constrained only by earth analog. Propagation times (Fig. 3b) range from 11-15 sec to 80 sec for the four first spheroidal/toroidal surface wave branches. We note that the ratio between the fundamental and the harmonic group velocities can be much larger than the standard  $\sqrt{3}$  ratio between the velocities of P and S body waves used by MQS (see Section 3). As an example, the 78 seconds between the two phases of event S0116a (discussed in section 3.1) are compatible with a slow packet propagating at 360 m/s (roughly the shear wave velocity at the base of the bedrock in our model) and a second packet propagating four times faster, which is roughly the P-wave velocity at a depth of ~100 meters, as proposed by Lesage *et al.* (2018). A difference of several minutes between the arrival of the first and second pulse is also found in event S0105a (see section 3.1). Second arrivals such as these might also be fundamental scattered Rayleigh waves, while the first arrivals could be overtones propagating in the deeper bedrock. The group velocity of the subsurface models also shows a clear variation of the group velocity just above 0.5 Hz, which might be the reason the signal has a cutoff frequency ~0.5 Hz.

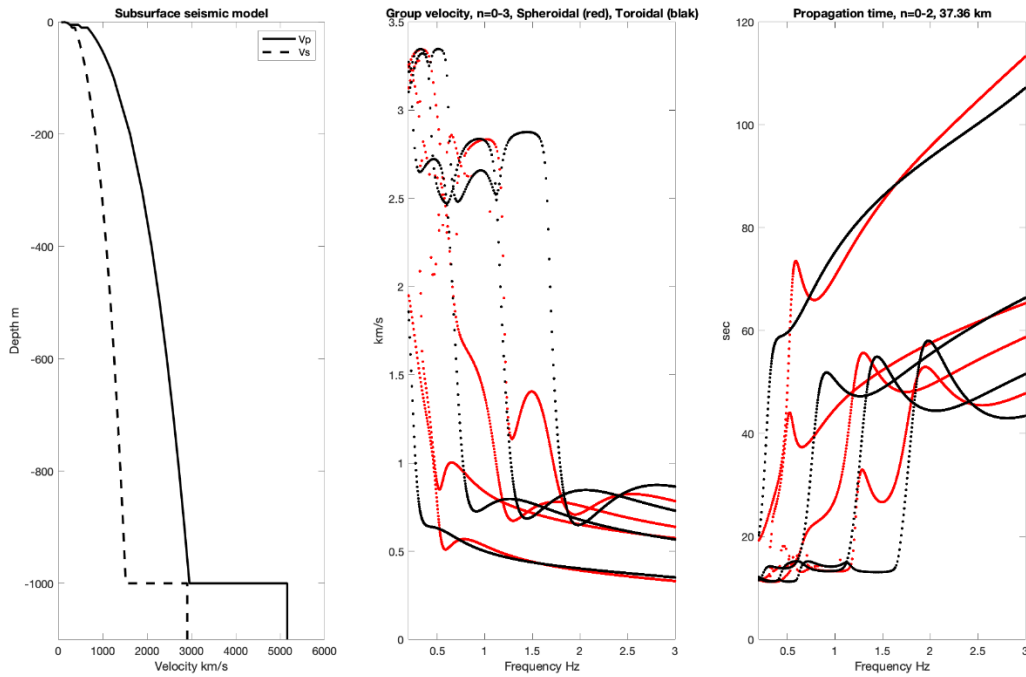


Figure 3. (Left) Subsurface seismic model. (Middle) Group velocities of the fundamental Rayleigh and Love waves and of the three first spheroidal (red) and toroidal (black) overtones. (Right) Propagation time of the surface wave packet to a distance of 37.36 km, as a function of frequency up to 3 Hz. Model data is provided in Table S2.

In summary, based on lunar data extrapolated to Mars, the shallow layers and diffusivity of Mars suggests that for an event at the distance of the new crater, we expect phase durations of 30 sec to 1 minute, with differences in phase arrivals up to about one minute.

## 2.4 Predicted atmospheric signals

A meteor entering the atmosphere and causing an impact crater would generate at various atmospheric levels in the entry path and at impact time both low-frequency gravity waves (typically 0.01-0.001 Hz) and high-frequency acoustic waves (frequencies above 0.01 Hz, typically 1-100 Hz) (Spiga *et al.*, 2018; Revelle 1976; Garcia *et al.*, 2017; Karakostas *et al.*, 2018). Those signals could be detected by a high-sensitivity pressure sensor operating continuously such as the pressure sensor in the Auxiliary Payload Sensor Suite (APSS) on board InSight (Daubar *et al.*, 2018).

An airburst signal would be characterized by two arrivals: first, the main seismic signal of surface waves excited at the location of the impact; and second, the blast wave through the atmosphere exciting the ground at the lander (Stevanović *et al.*, 2017). A differential travel time

of ~2 minutes is expected between two such signals due to the difference in wave propagation speeds of 230 m/s in the air and 1.5 km/s in the subsurface over the 37 km distance from the impact to the lander. Such a signal would be much smaller than InSight's pressure sensor limit of detectability, so SEIS would be the only way to detect such a phenomenon.

Atmospheric entry modeling demonstrates that for this scale of impact the majority of the meteoroid's kinetic energy is transferred to the atmosphere during deceleration and ablation, and only a small fraction is directly coupled to the ground by the surviving fragment(s). The relatively large blast zone surrounding the crater (Fig. 1D) is testament to this partitioning. However, previous work suggests that detection of the direct ground impact is more likely than detection of airburst-generated acoustic and gravity waves near the ground surface (Garcia *et al.*, 2017) as InSight's detection capability of acoustic and gravity waves produced by airbursts and surface explosions is negatively affected by atmospheric attenuation and propagation conditions less favorable than on Earth (Lognonné *et al.*, 2016). Moreover, numerical modeling (based on the methodology of Karakostas *et al.*, 2018) suggests that even in the endmember case of all the meteoroid kinetic energy being deposited in the atmosphere, the resulting air-coupled seismic waves would still not be detectable by the InSight instruments. Acoustic ray propagation models (Garcia *et al.*, 2017; Spiga *et al.*, 2018) show the trajectories of infrasound rays do not reach the InSight lander, which is in an unfortunate shadow zone at this distance from the impact (Fig. S11). Considering both atmospheric wave propagation conditions and meteor energy scaling, we therefore do not expect the acoustic and gravity waves generated by the meteoroid that formed the 1.5 m crater to be detected by InSight.

### 3 Candidate seismic events in the time period of interest

#### 3.1 Description of SEIS data and the candidate events

The time between the before and after CTX images was a period of immense interest in the data coming from InSight. The most relevant data was from the seismometer, Seismic Experiment for Interior Structure (SEIS; Lognonné *et al.*, 2019). This temporal search window occurred as SEIS commissioning was being finalized, only a few weeks after SEIS was placed on the ground (17 January 2019) and the Wind and Thermal Shield (WTS) had been placed over it (2 February 2019), allowing the lowest possible noise on the instrument. Fortunately, continuous data collection (InSight Mars SEIS data service, 2020) had already transitioned to being round-the-clock, and three-component Very-Broad-Band (VBB) and Short-Period (SP) data at 10 sps, sometimes also at 20 sps, was available throughout the time period (Fig. S10).

Figure 4 provides an overview of the completeness, the noise, and the occurrence of seismic signals in the data within the search window. Seismic noise on Mars clearly falls into a daily pattern, with low noise only occurring between ~16:00 LMST (Local Mean Solar Time) to ~02:00 LMST (Lognonné *et al.*, 2020, Giardini *et al.*, 2020). Outside of this time, there is a substantial increase in noise, with steady winds in the early morning followed by a gusty midday period. During these times, only very strong seismic signals can be detected. Furthermore, not all days include a significant quiet period. Thus, there are large daily and day-to-day variations in



our capacity to detect weak seismic events on Mars using SEIS data. During the search window, we estimate weak signals could be reliably detected only ~30% of the time.

Despite these limitations to the data, three potential seismic events were identified between the times of the constraining CTX images (Fig. 4). Although all three are weak signals, there are unique aspects of these events that deserve examination. We discuss the characteristics of each of them, and the likelihood that each is the signal resulting from the observed new crater.

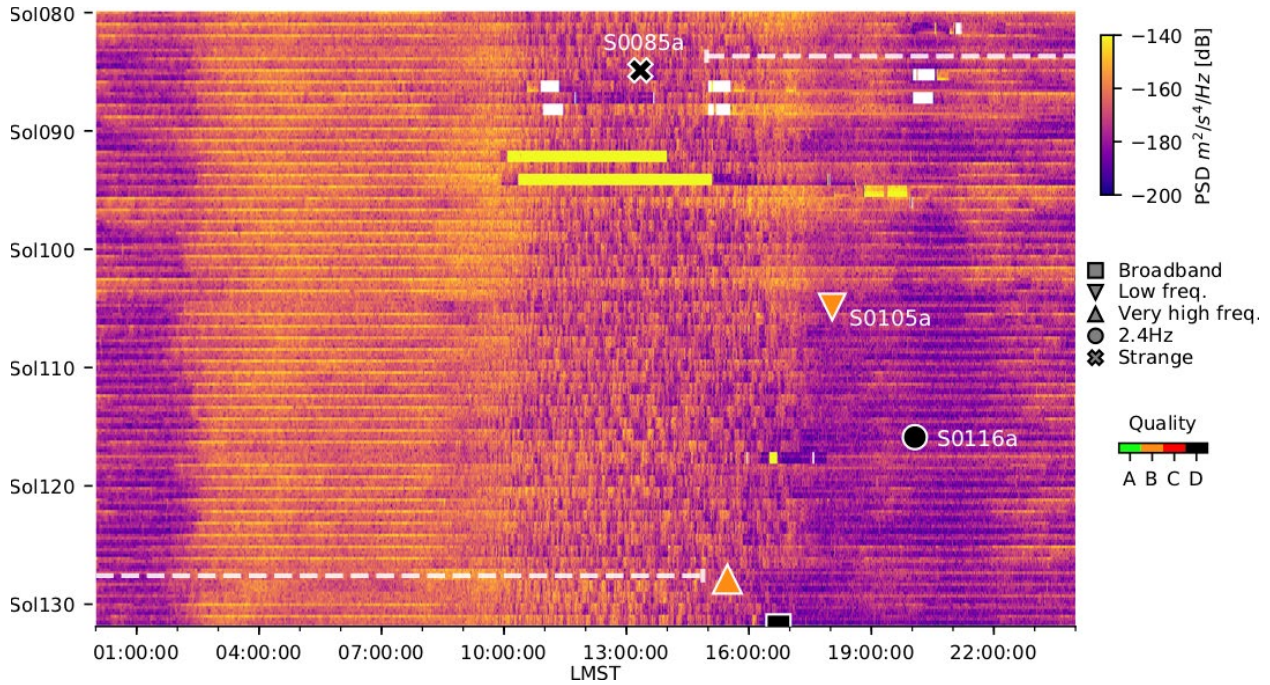


Figure 4.

*Spectrogram stack from InSight sol 80 (16 February 2019) to sol 132 (10 April 2019). This period bounds the impact search window from sol 84 14:55 LMST (21 February 2019 03:56 UTC) to sol 127 14:18 LMST (6 April 2019 08:20 UTC), indicated by start/end of the white dashed lines. Each horizontal line in this plot corresponds to a sol-long acceleration spectrogram from 20 s to 4 Hz for the vertical VBB component. White and yellow bars indicate data gaps and amplitude saturation, respectively, occurring during sensor calibration and hammering of the heat flow probe. The three events detected and discussed in this paper are marked with symbols corresponding to the event type, while event quality is indicated by symbol color (see legend). Two events that occurred just after the end of the search window are also indicated.*

The Marsquake Service (MQS, Clinton *et al.*, 2018) is tasked with reviewing all data from SEIS, detecting and characterizing seismic energy, and maintaining a catalogue of marsquakes. MQS detects events by careful manual review of all continuous data. Over the course of the mission so far, the most effective approach to identifying marsquakes has proven to be data review using spectrograms. Standard MQS operations produce daily spectrograms with a

window length of 50 seconds for frequencies below one Hz and 10 seconds for higher frequencies. In the first months, two major event families have been observed (Giardini *et al.*, 2020; InSight Marsquake Service, 2020). The first family is characterized by events with energy dominant at lower frequencies, visible as a 10-20 minute-long energy surplus between 0.1 and 3 Hz. This family comprises the two event types, Low Frequency (LF) and Broadband. The largest of these events (named S0173a and S0235b; Giardini *et al.*, 2020, Lognonné *et al.*, 2020) have clearly identifiable P- and S-waves, with clear polarization showing the direction as seen from the lander, followed by long codas of scattered energy. Smaller events of this type have polarities that are less clear or are not detectable, but the envelope of the waveforms and their spectral content supports the interpretation that they are smaller versions of the same type of event. The second major family includes High Frequency (HF) events, characterized by an energy content mainly above 1 Hz, an extended coda, and a lack of polarization. An additional curious feature of the InSight landing site is a local seismic resonance at 2.4 Hz. For larger HF events, the spectrum can be matched by a general earthquake spectrum, taking into account source size and attenuation, modulated by an amplification of 12 dB in spectral energy around 2.4 Hz. For smaller HF events, only this peak is visible, while the bulk of the energy is below the ambient noise level. Events in this family are classified as High Frequency, Very High Frequency, or 2.4 Hz. A handful of events have been documented as “Strange” if they do not fit into any of these standard event types.

During the time period of the impact search, one event was found during standard MQS operations. It has the label S0105a (the first seismic event to occur on sol 105 of the mission) and is a Low Frequency event. It was in fact the first seismic event detected during the whole mission. After the CTX discovery of the impact, a review of all data during this period was performed by the InSight team, both within and independently of the MQS team. This review took into account the improved understanding of marsquake character that had accumulated from other events in the meantime. During this review, two additional events in the time period were identified: a small High Frequency event on sol 116 and one unclassifiable seismic signal on sol 85. We first describe the three events in detail:

S0085a [2019/02/22 02:58:15 UTC, 13:35 LMST; MQS classification: Strange signal]

A summary of the S0085a event is shown in Figure 5. This event, which appears to be unique among events detected on InSight thus far, consists of a very narrow-banded energy surplus at 0.7 Hz, with a bandwidth of 0.05 Hz. There is a slight rise in frequency over the course of the event, from 0.57 Hz up to 0.7 Hz. The signal is visible dominantly on the north component, with weak traces in the Z (vertical) component. This indicates a clear N/S azimuth. The signal occurs only hours after the opening of the search window, during the part of the day with high atmospheric noise. In fact, it is interrupted by several wind bursts creating noise more than 10 dB above the signal itself. Because it can only be resolved during the intermittent quiet periods, the exact start and end times cannot be positively identified, but the event lasts at least 10 minutes. The very narrow bandwidth does not fit any expected seismic mechanism (including impacts). A similar signal has not been observed a second time during the mission, especially not during a quieter period, which would allow a better classification of its character. No particular lander activity was going on at the time of this event that could explain it. Given the high



atmospheric noise surrounding this time, it cannot be discounted that it could be of random origin.

This event was not detected using standard analysis, but extending a method that exploits the ratios of the average energy residing between 2.4 Hz  $\pm$  0.2 Hz, to different frequency bands of the SP's and VBB's North, East and Vertical (Z) components. The algorithm was implemented in steps of 0.4 Hz with 50% overlapping windows in frequency and avoiding injection of tick noise (cross-talk noise generated by the SEIS temperature signal on the VBB and SP seismic data). The resulting outliers were inspected against the average energy in the Energy Short-Term Average (ESTA) channel (defined as the root mean square of data filtered within a 0.5 second time window [Lognonné *et al.*, 2019]), to ensure they occurred during calmer atmospheric periods, and to allow for further investigation.

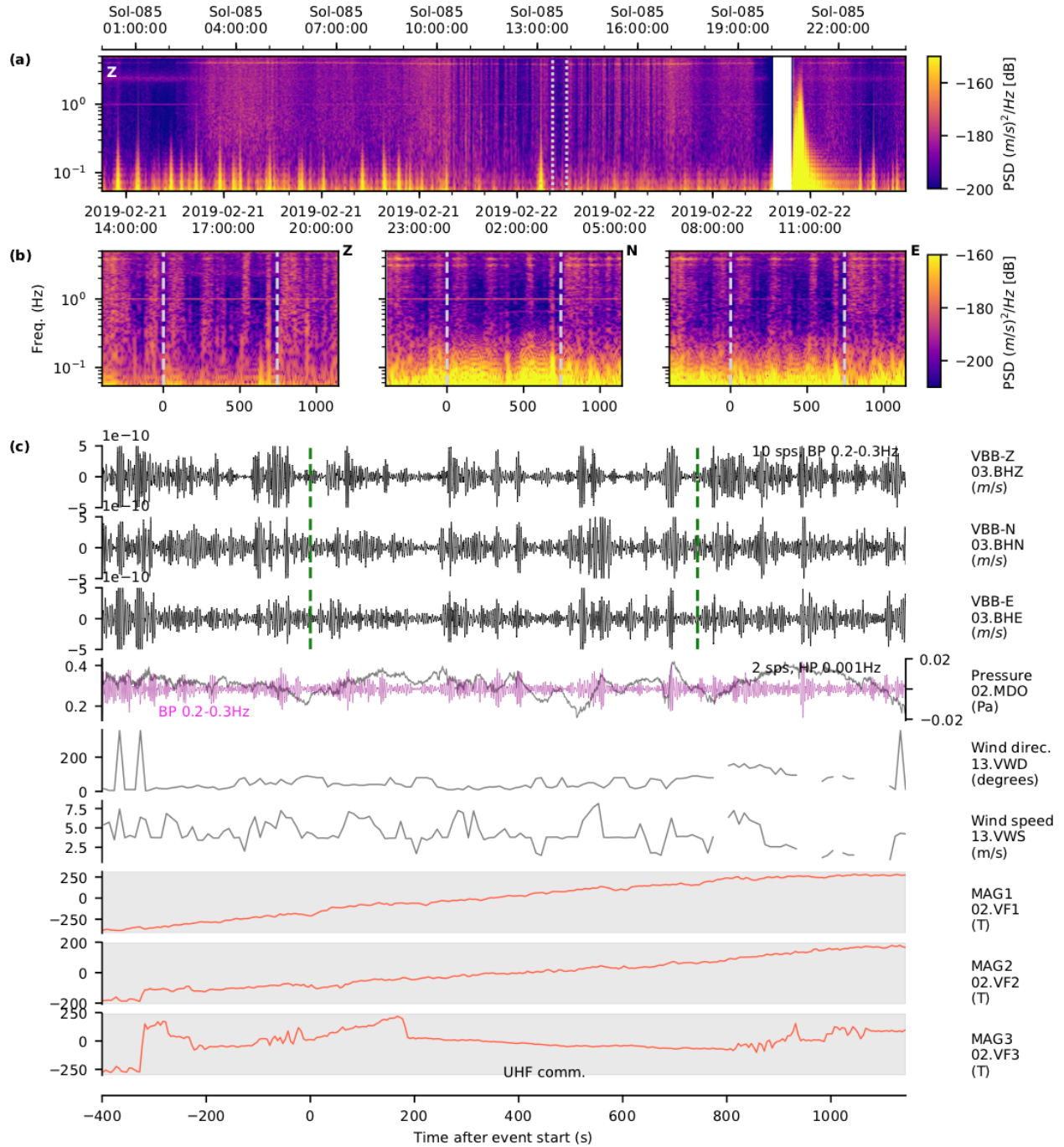


Figure 5.

Summary of S0085a event. (a) provides the context of the event in the full sol spectrogram on the VBB vertical (Z) component. (b) shows spectrograms for all 3 VBB components rotated into Z, north (N) and east (E) orientations. The start and end time are indicated by the vertical dashed white lines in (a) and (b). (c) shows timeseries from the 3 VBB velocity, pressure, wind direction, wind speed, and 3 magnetometer channels. The data are filtered as indicated to accentuate seismic and pressure signals. The vertical green dotted lines in (c) indicate the event start and end times. In general in these summary figures, additional phase picks in green and glitch

568 *windows in red are overlain on the seismic channels; and on the magnetometer channels*  
 569 *indications of any reported lander activity are shown in gray. For this event, however, the lander*  
 570 *has UHF communications, there are no major glitches, and the event is too weak for MQS to*  
 571 *identify phases. This event is extremely faint and not visible in the time series. The event is the*  
 572 *very narrow band of energy at 0.7 Hz visible on the N component spectrogram. As explained in*  
 573 *the text, this signal may not have a seismic source.*  
 574

575 S0105a [2019/03/14 21:03:31, 18:07 LMST; MQS classification: Low Frequency, Quality C]  
 576

577 This Low Frequency event consists of two energy pulses, each without clear polarization  
 578 (Fig. 6). It occurs around sunset, just after the transition from the high atmospheric situation of  
 579 the day into the very quiet early evening. The amplitude of this event is so low that it could only  
 580 have been reliably detected during ~25% of the time period of the impact. The total length of the  
 581 signal is ~15 minutes, with at least 5 minutes uncertainty, given the relatively high noise level.  
 582 The spectral energy is above the ambient noise between 0.3 and 0.5 Hz for the first pulse and  
 583 0.15 and 0.5 Hz for the second pulse. The spectrum of the two pulses is comparable to that of  
 584 event S0173a, currently the largest LF event in the MQS catalog, but 16 dB lower at 0.3 Hz. The  
 585 phases are emergent, and phase arrival picks for the two energy pulses cannot be made in the  
 586 time domain, and so are made using a spectrogram and accordingly assigned high uncertainties  
 587 of +/-20s. In the time domain, the separation of the two pulses is also similar to that of S0173a  
 588 (160 seconds for S0105a vs 155 seconds for S0173a). The similarity of the signal of this event  
 589 and other low frequency events is shown in Fig. 3 from Giardini *et al.*, 2020, and consistent with  
 590 other larger events of this type, we assign P and S phases to the onset of these pulses. It would be  
 591 difficult to convincingly assign these phase arrivals to P and S waves without the context of the  
 592 wider seismicity so far recorded by InSight. Other interpretations may also be plausible, as  
 593 discussed above in section 2.3, though this weak event is generally similar to stronger and more  
 594 well-understood events.  
 595

596 Based on the time elapsed between these pulses, this event is estimated to be located at a  
 597 distance of  $27 \pm 5^\circ$  ( $1600 \pm 300$  km). For S0173a, a polarization analysis was also possible,  
 598 resulting in a direction of the events as seen from the lander of  $91 \pm 5^\circ$ ; thus it has been concluded  
 599 this is the signal of a marsquake located in the Cerberus Fossae graben system (Giardini *et al.*,  
 600 2020). This fault system is the only place on Mars where more than one marsquake has been  
 601 located so far, in agreement with pre-mission hypotheses of seismic activity there. A possible  
 602 interpretation of the S0105a event is therefore that it is a smaller tectonic marsquake in a similar  
 603 location to S0173a. As no polarization could be determined for S0105a, this interpretation must  
 604 remain preliminary. The low signal-to-noise ratio also implies that no depth could be estimated  
 605 for this or any other event in the impact time period.  
 606

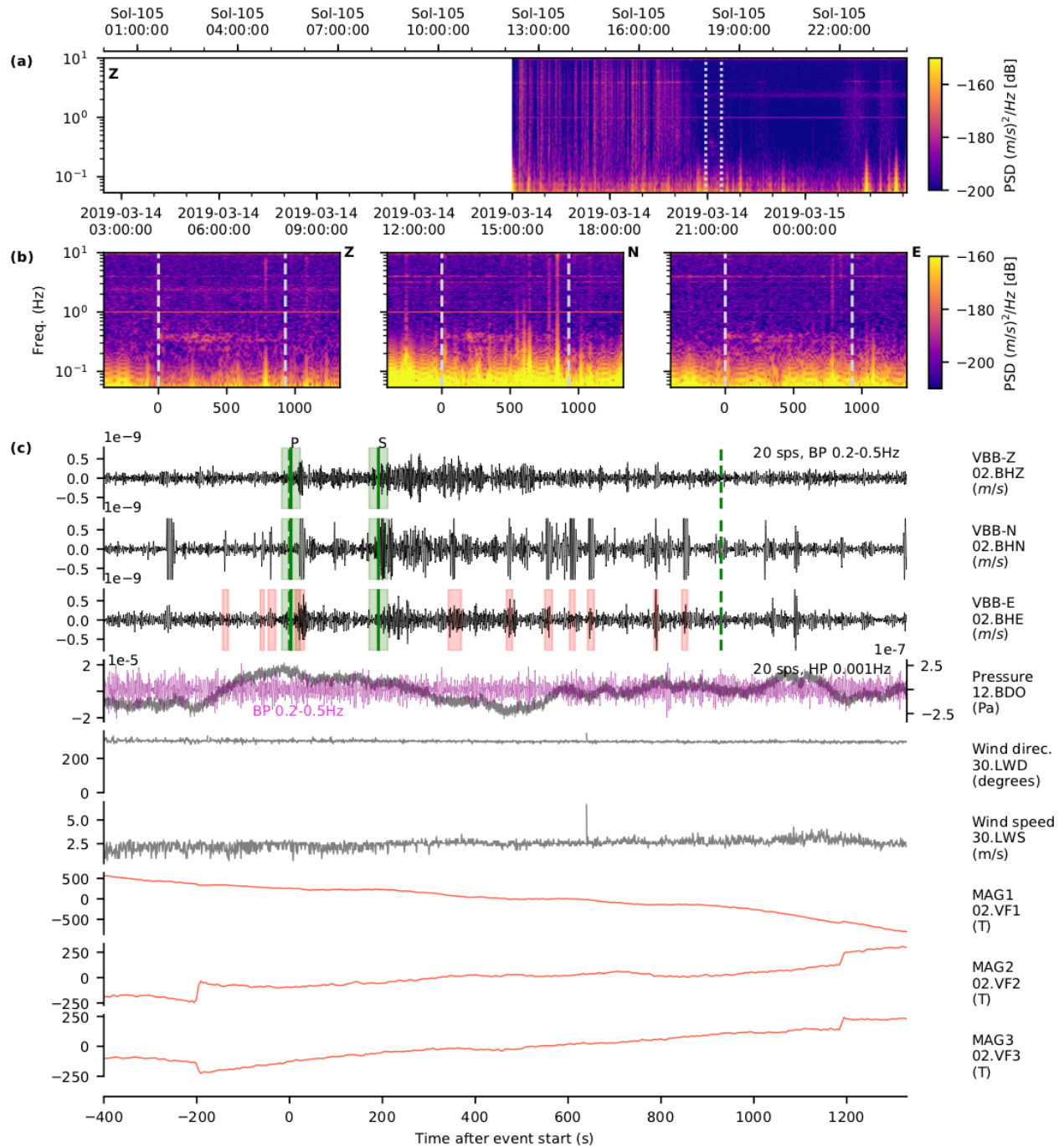


Figure 6.

Summary of S0105a event, following Figure 5. During this event, there are multiple glitches (red shaded windows), most clearly visible in the E component, and no lander activity. MQS also identifies P and S phases (green solid vertical lines). Event energy is visible on all 3 components in both time series and spectrograms.

S0116a [2019/03/26 06:27:19 UTC, 20:11 LMST; MQS classification: High Frequency 2.4 Hz, Quality D]

This High Frequency 2.4 Hz event, summarized in Figure 7, consists of an energy surplus around the 2.4 Hz mode of about 7 dB in displacement power, concentrated into two pulses separated by  $78 \pm 10$  seconds. At the time of detection, this event was unique, but as of the time of writing, we have come to realize that it was just the first occurrence of a general class of similar events, termed “2.4 Hz events”. These are currently understood as being small High-Frequency events. HF events are interpreted as shallow-source events occurring in a highly scattering layer in the upper crust, probably shallower than the source region of the LF events. The absolute distance of the HF events cannot be determined yet, as crustal seismic velocities are so far unknown. The convention for these HF events is to label the start of each pulse as Pg and Sg phases. From the separation of the two phases, a relative distance can be estimated. The S0116a event is about four times closer than the majority of the HF events occurring later in the mission, so it seems to have emanated from a difference source region. Only a handful of other events share a similarly short Sg-Pg interval. Nevertheless, for it to have occurred at the detected impact site, the shear wave velocity in the medium would have to be as low as 210 m/s (assuming a  $v_P/v_S$  ratio of  $\sqrt{3}$ ). Such a velocity is found in bedrock layers 5-10 m deep (Lognonné *et al.*, 2020), but is unlikely at these shallow depths. In the MQS catalogue, all HF events are given an estimated location using an assumed S wave velocity of  $v_S = 2.3$  km/s, and P wave velocity of  $v_P = \sqrt{3}v_S = 4.0$  km/s. Using those assumed velocities, this event has an estimated distance of  $11^\circ$ ,  $\sim 640$  km from the InSight lander. As the event is only visible as an excitation of the 2.4 Hz mode, its original source spectrum cannot be constrained.

The amplitude of this event is so low that it could have been detected during only  $\sim 20\%$  of the day during the time period of the known impact.

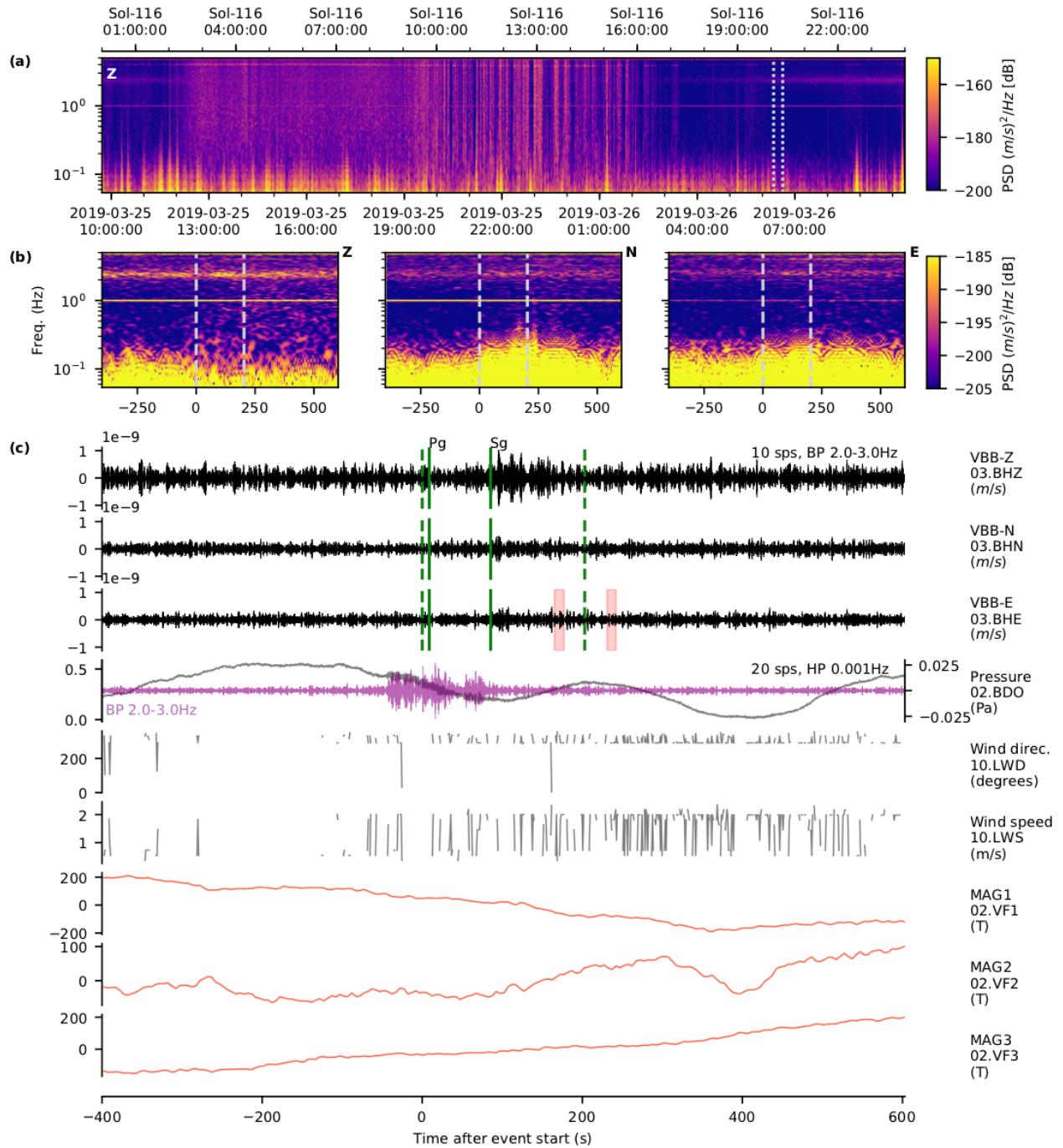


Figure 7.

Summary of S0116a event, following Figure 5. During this event, there are two minor glitches towards the end of the event (red shaded windows), and no lander activity. MQS also identifies tentative Pg and Sg phases (green vertical bars). Event energy at the 2.4 Hz resonance is weakly visible on all three components in both time series and spectrograms. An anomalous high frequency disturbance in the 2 sps pressure precedes the event, extending into the first minutes.



We note that nearly exactly 1 sol after the search period closed, the very high frequency event S0128a occurred. It can be seen in Figure 4, but outside the search period defined by the dashed white lines. This was one of the largest events so far recorded, and one of the events located closest to the InSight lander, although it is still estimated to be roughly  $\sim 8^\circ \pm 6^\circ$  ( $\sim 480$  km  $\pm 350$  km) away (Giardini *et al.*, 2020, Lognonné *et al.*, 2020). Although the uncertainties on this distance estimate are large, they still do not encompass the small distance to this known impact. Additionally, the timing of the CTX images has been closely compared to this event timing, and the event does not fall within the possible time period for the new impact.

### 3.2 Evaluating seismic data for the candidate events

With regards to the three events detected within the search period, how can we evaluate which, if any, is the recording of the known image-constrained impact? Aside from their occurrence within the search period between 2/21/19 and 4/6/19, there are few other positive indicators that each of the signals was caused by the impact. Scaling relationships and analog comparisons predict the observed impact would create a seismic signal with peak energy  $\sim$ a few Hz, with a peak amplitude of the P-wave  $\sim 0.8$ -4 nm/s. This range is also in good agreement with amplitudes from the numerical wave propagation simulations (supplemental section S2). However, none of the three candidate events includes energy above 2.4 Hz. The predicted duration of the event is  $\sim 30$  seconds to one minute, although this is difficult to compare directly due to scattering. However, all candidate events have durations of over several minutes. We know the impact occurred at a backazimuth of  $180.9^\circ$ , so any polarization present in the signal should be in the north-south direction. S0105a and S0116a have no indication of polarization, though S0085a does include energy only in the N-S component, which is a match. Here we detail how well each of the candidate signals matches these expected characteristics (Table 2).

Table 2.

*Expected characteristics of the seismic signal produced by the known impact, compared to the characteristics of each of the candidate seismic events. Matching characteristics are marked with a green check mark, non-matching characteristics are marked with a red “X”, and neutral or undetected characteristics are marked with a black “~”. Distance to source is measured from orbital images for the known impact and estimated for seismic events by MQS.*

|                                    | Unambiguous seismic event? | Amplitude (nm/s)                    | Peak Frequency (Hz)                                        | Polarization                                        | Duration (min)           | Distance to source (km)                  |
|------------------------------------|----------------------------|-------------------------------------|------------------------------------------------------------|-----------------------------------------------------|--------------------------|------------------------------------------|
| <i>Predicted for known impact:</i> | <i>Uncertain</i>           | <i><math>\sim 0.8</math>-4 nm/s</i> | <i><math>\sim 2</math>-3 Hz most likely for body waves</i> | <i><math>180.9^\circ</math> (approximately N/S)</i> | <i>30 sec - 1 minute</i> | <i>37.4 km (<math>0.65^\circ</math>)</i> |



|        |                                                                                |                                                                                                                                                             |                                       |                               |                                  |                                 |
|--------|--------------------------------------------------------------------------------|-------------------------------------------------------------------------------------------------------------------------------------------------------------|---------------------------------------|-------------------------------|----------------------------------|---------------------------------|
| S0085a | No,<br>very unusual<br>signal occurring<br>in noisiest time<br>period<br><br>~ | 0.3 nm/s<br>(North; Vertical<br>and East not<br>above noise)<br>(bandpass 0.5-1<br>Hz,<br>6 pole)<br><br>(approximately<br>right)<br><br>✓                  | 0.7 Hz<br><br>(too low)<br><br>✗      | N/S<br><br>✓                  | ~10 min<br><br>Too long<br><br>✗ | Unknown<br><br>~                |
| S0105a | Yes,<br>clear LF event<br><br>~                                                | 1.5 nm/s (East),<br>0.5 nm/s<br>(Vertical; North<br>affected by<br>glitches)<br>(bandpass 0.2-<br>0.67 Hz, 6 pole)<br><br>(approximately<br>right)<br><br>✓ | 0.15-0.5 Hz<br><br>(too low)<br><br>✗ | None<br>identifiable<br><br>~ | ~15 min<br><br>Too long<br><br>✗ | 1600±300<br>km (27±5°)<br><br>✗ |
| S0116a | No,<br>weak 2.4 Hz<br>resonance<br><br>~                                       | 0.7 nm/s<br>(Vertical)<br>0.5 nm/s (East)<br>0.5 nm/s (North)<br>(bandpass 2.2-2.8<br>Hz, 6 pole)<br><br>(approximately<br>right)<br><br>✓                  | 2.4 Hz<br><br>(reasonable)<br><br>✓   | None<br>identifiable<br><br>~ | ~3 mins<br><br>Too long<br><br>✗ | Unknown<br><br>~                |

684

685 S0085a:

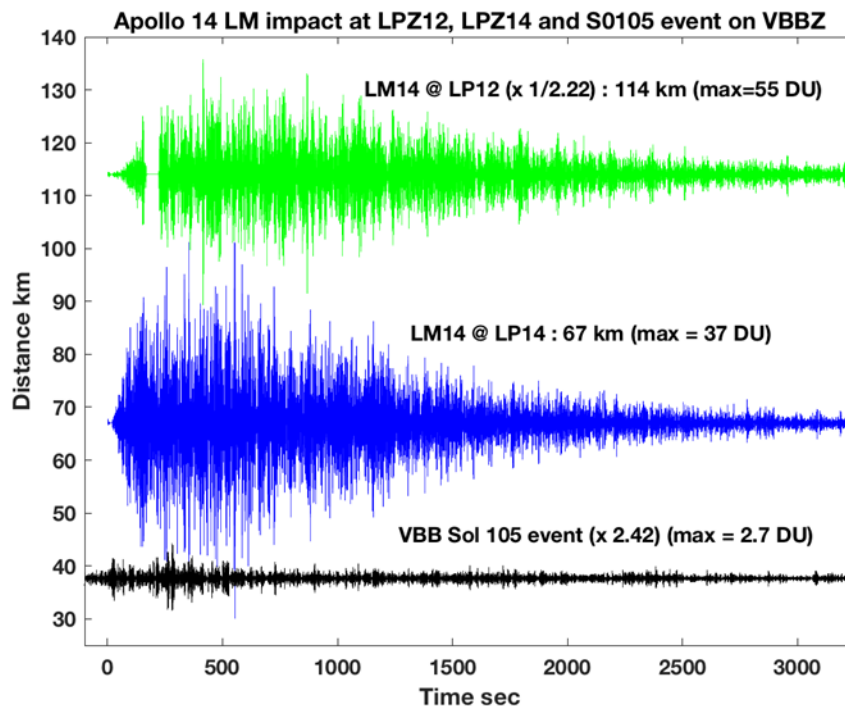
686 The event on sol 85 is the only one of the candidates with a measurable polarization, and  
687 it is in the correct direction relative to the impact. However, it is possible that this event may not  
688 be a seismic event at all – it could in fact be instrument-generated rather than a natural external  
689 source. Many spacecraft-induced signals will have a similar N/S polarization, as the InSight  
690 lander is towards the north of SEIS. Compared to other observed events, it has very narrow-band  
691 energy with an apparent dispersion, which is not expected for an impact. Even if this were a true  
692 seismic event, we cannot definitively identify it with the impact.

693

S0105a:

The event on sol 105, on the other hand, is a clear seismic event. Its amplitude (0.5 nm/s in the 0.2-0.67 Hz bandwidth on Vertical and 1.5 nm/s on East) is at the lower end of that predicted. The spectral peak is at a frequency lower than that predicted. If this is an impact at  $\sim 40$  km, it would have to be explained why tens of other seismic events detected so far look very similar to this one. It is exceedingly unlikely that multiple small impacts occurred in this region in the same time period and we do not have any images of them; although we do not have complete repeat image coverage of the region out to  $\sim 40$  km away from InSight in order to rule this out completely. Without definitive criteria for discriminating between impact and tectonic sources (see Section 4.1), we cannot exclude the possibility that one of these similar events is also an impact.

Figure 8 compares the S0105a signal with two Apollo impact records. All signals have been filtered with a 6<sup>th</sup>-order Butterworth bandpass (0.2-0.67 Hz), and SEIS data are expressed in Apollo Digital Units (DU). Amplitudes in Fig. 8 have been corrected with respect to distance using a -1.5 power law dependency with respect to the Apollo 14 LM impact recorded by the Apollo 14 LP seismometer, while the non-corrected amplitudes are given for each trace. The amplitude of the S0105 event is approximately 15 times smaller than that of the lunar LM impacts, which is within our estimate of a factor of 3.5-210 (see Section 2.3). The amplitude of the signal at 0.5 Hz can therefore be explained by the size of the known impact. However, neither the lack of high frequencies nor the duration of this event are compatible with what we expect for this impact.



## Figure 8.

*Comparison between the martian event S0105a (black) and the Apollo 14 LM impact as recorded at two Apollo stations (green for Apollo 12 and blue for Apollo 14). The S0105a event has been deglitched (Lognonné et al., 2020) and converted into Apollo data units (DU) by using the Apollo Transfer function of the LP instruments. All events have been filtered with a 6<sup>th</sup> order Butterworth bandpass between 0.2 and 0.67 Hz and corrected for the different distances by using a -1.5 power law with distance. Amplitude in DU as well as geometrical correction values are given on the figure. The very impulsive first arrival identified by MQS for S0105a is at time 0, followed by a second arrival 160 seconds later.*

## S0116a:

The event on sol 116 has an amplitude at the lower end of that predicted, though this is on top of the 2.4 Hz resonance, so is likely amplified. It has a higher frequency than the other events, which is reasonable for a small, local event. No polarization was detected, so no direction or distance can be estimated. In the months since this event was recognized, hundreds of other similar events have occurred, again making it unlikely this is due to an impact, which would occur relatively infrequently.

In summary, none of the three events can be unambiguously identified as the seismic signature of the new impact. The S0105a event can be explained as a relatively small tectonic marsquake in the Cerberus Fossae region. The S0116a could possibly be caused by the impact, but given its low amplitude, it cannot be further classified or analyzed. Both S0105a and S0116a are similar to numerous other events in the marsquake catalogue, suggesting they are not produced by a local impact signal, which we expect to be a rare occurrence. The S0085a signal is extremely weak, and its very narrow-band nature suggests it is not likely to have been caused by an external seismic event.

We note the extreme variation in diurnal noise means that significantly larger events than the three identified here may be hidden in the data. As noted, the amplitude we predict for this impact is quite close to the measured noise levels of SEIS during the least noisy time periods (Fig. 4). Given daily and seasonal variations in temperature and wind activity, the noise levels are lowest in the evening (Lognonné et al., 2020; Fig. 4). Signals on the order of the predicted amplitude would only be observable (at ~3 SNR) for ~20-30% of the time. Thus, the actual signal from this impact could very likely have occurred at a time when noise swamped the signal.

Another observational bias could occur due to the 2.4 Hz signals. These are a resonance seen in numerous other events (Giardini et al., 2020), and this is also near the peak frequency expected for impact event. Such a resonance could enhance smaller signals, allowing detection of signals that are otherwise ~ten times smaller if they are near the resonance. This might help our detection likelihood, but it is also a narrow band, making source discrimination more difficult.

### 3.3 Evaluating atmospheric data for the candidate events

Although no obvious atmospheric signals were associated with the three seismic candidates, each event was investigated to eliminate that possibility:

S0085a occurred in the local early afternoon (unlike the S0105a and S0116a events). Atmospheric variability for S0085a is thus mostly governed by convective turbulence (cells and vortices), usually found in daytime hours. There were no vortex signals close to the event that might have affected the seismic signal. Under normal conditions, gravity waves are not usually detected in daytime hours, based on the first 300 sols of InSight measurements (Banfield *et al.*, 2020), and no gravity waves were detected around the time of this event.

There was no notable atmospheric signal associated directly with the S0105a event. Pressure and temperature measurements were uneventful, and the wind was steady and low. Two hours after the seismic event, a gravity wave of strong amplitude ( $\pm 0.5$  Pa) with a period of 400-600 seconds was detected. However, between sols 100-150, similar signals were very frequently seen at these local times ( $\sim 20:00$  LMST). Furthermore, given the proximity of the impact to InSight, a propagation speed that would cause a two hour delay between the seismic signal and the atmospheric wave packet is far too low to be realistic. This is based on typical gravity wave phase speeds estimated by Banfield *et al.* (2020) of  $\sim 20$ -30 m/s. The gravity wave signal reached InSight two hours after the seismic event; even accounting for background wind, the gravity wave would be too fast to have been emitted by the atmospheric entry of a meteoroid at 37.4 km distance. It is thus not likely to be related to the seismic event.

For S0116a, a gravity wave signal was found in the pressure signal at a time near the seismic event. However, it started about a quarter of an hour before the seismic signal, which implies the two are unrelated. Wind and temperature measurements behaved as usual for evening conditions at the InSight landing site. Interestingly, around the start of the event, just before the seismic signal, the pressure signal also underwent high-frequency fluctuations in the infrasonic range (*i.e.*  $< 20$  Hz). Though an impulsive pressure signal could be expected from a close impact event, the long duration HF pressure wave packet we observed is similar to scattered pressure signals related to explosions seen in infrasound records on Earth (Green *et al.*, 2011). Such scatterings of acoustic energy can occur when small-scale gravity waves perturb the lower atmosphere wave guide (e. g. Green *et al.*, 2011; Damiens *et al.*, 2017). Nonetheless, other facts concur to put aside the impact hypothesis as a source of the observed pressure fluctuations: (1) high-frequency pressure fluctuations are recorded by InSight almost every sol in the evening, (Banfield *et al.*, 2020), (2) acoustic propagation models (Garcia *et al.*, 2017; Spiga *et al.*, 2018) show that the InSight lander is in a shadow zone for infrasound waves generated at the impact location (Fig. S11) and (3) owing to the noise levels of the respective instruments, if infrasound signals were seen in InSight pressure data, they would also be seen in the seismic data. They would also have to be at a significantly larger distance than this case (Martire *et al.*, 2020). This makes it very difficult to ascribe these particular pressure fluctuations to the sol 116 seismic event, or to any impact-induced phenomena at the distance of the known impact.

Regarding an airburst signal, none of the candidate events have two distinct arrivals with the expected temporal spacing of  $\sim 2$  minutes (Section 2.4), even if they were above the detection

threshold. So we do not believe an impact airburst was detected for this event. To summarize, while interesting atmospheric signals were noticed during the three events, they are not likely related to either the seismic events in question or to the impact event.

## 4 Discussion

### 4.1 Re-assessment of seismic impact discriminators

Over the first months of the InSight mission, we have learned that marsquakes (whether sources are impact or tectonic) differ from our previous experience with either terrestrial or lunar analogs. The impact discriminators we planned on using before arriving at Mars (Daubar *et al.*, 2018) have limited utility given the reality of martian seismic signals recorded thus far. The marsquakes observed so far are small in amplitude, with surprisingly long durations, and with apparently low attenuation / high Q. This makes many of these characteristics difficult to distinguish. We re-assess (*in italics*) each of those planned discriminators in light of real seismic data from InSight:

- 1) First motion: Impacts create positive pressure impulses, creating a positive first motion, in a direction away from the source.
  - *Despite the low noise recorded by InSight during periods of the day, marsquake signals have proven to be very small. In all but the largest signals seen so far, phase arrivals are emergent, so noise obscures the direction of first arrivals. Scattering in the regolith randomizes the energy.*
  - *Even if we had clear first motions, quakes with a double couple source would have a positive first motion 50% of the time anyway, assuming a random orientation of sources.*
- 2) S-wave energy: Impacts produce more P-waves than S-waves.
  - *A quake could also have low S energy for an unfavorable source orientation.*
  - *S-waves are obscured by scattered P energy, so this is hard to determine for small events.*
- 3) Magnitude ratio: Impacts produce fewer surface waves, so impacts should have a strong difference between magnitudes based on body waves and those based on surface waves.
  - *Surface waves are not being detected for any martian events (Giardini et al., 2020). The absence or diminished presence of surface wave energy, therefore, cannot be used as an impact discriminator, because all events lack surface waves.*
- 4) Frequency content: Different source mechanisms lead to a smaller cutoff frequency for impacts.
  - *Cutoff frequencies for the largest of the detected martian events, where they can be determined, are typically near 6 Hz but can rise up to 12 Hz (Giardini et al., 2020). This cutoff frequency is much higher than the ~1-3 Hz expected for impacts (Daubar et al., 2018).*
- 5) Depth phases: Impacts occur at the surface, implying no depth reflected phases.
  - *Additional phases beyond P and S arrivals have not been identified in any events thus far (Giardini et al., 2020) because of scattering and the resulting extended codas, so a lack of depth phases cannot be used to indicate an impact.*

## 4.2 Revised predictions of impact detections by InSight

As this is the only impact known to have occurred this close to InSight during its prime science monitoring phase thus far, we wish to evaluate how likely this particular impact event was. Using an estimated current cratering rate, we can estimate the probability of a ~1 m diameter crater forming within ~50 km of InSight in one Earth year. Unfortunately, the cratering rate for impacts of this scale is not well constrained. As an estimate, we use a production function based on an extrapolation of the fragmentation model of Williams *et al.* (2014) pinned to the production function based on observed dated craters from Daubar *et al.* (2013) (see Teanby 2015 for more details). The resulting rate is  $\sim 2 \times 10^{-5}$  impacts >1 m diameter/km<sup>2</sup>/Earth year. The uncertainty on this value is probably at least a factor of 10 in both directions. For this impact rate, the probability of one impact in any given circle of radius 50 km each Earth year is  $\sim 0.2$ . Thus this event is not completely unlikely, but we were quite lucky to catch it in the images, which have covered only a small fraction of that area multiple times since InSight landed.

Based on measured noise levels of SEIS on the ground at Mars, we can revise our pre-landing estimates of the number and size of impact detections to expect during the InSight mission. Teanby and Wookey (2011) and Teanby (2015) estimated seismic impact detection rates with predicted Mars seismic noise. We can now update these predictions using measured noise levels from the first few months of InSight operations. Teanby and Wookey (2011) model results for large impacts predict their peak seismic energy will be in the 1-2 Hz frequency range (where the SEIS-VBB instrument is most sensitive). Teanby (2015) compiled observations from small impacts and explosions to suggest that their peak seismic energy will be in the 1-8 Hz frequency range (where the SEIS-SP instrument is most sensitive). Typical SEIS noise levels are  $0.3\text{--}10 \times 10^{-9}$  m/s<sup>2</sup>/Hz<sup>1/2</sup> in the 1-8 Hz range (Lognonné *et al.*, 2020), although during much of the martian day SEIS sees considerably higher noise than these levels. Using scaling relationships developed in previous work (Teanby, 2015; Teanby and Wookey 2011), we can predict P-wave amplitudes for different size impacts at various distances (Fig. 2). To get the expected frequency of impacts of different sizes, we use the production function developed by Teanby (2015) that uses new dated craters from Daubar *et al.* (2013) extrapolated to smaller diameters to account for the observational rollover using the Williams *et al.* (2014) fragmentation model. As signals at the noise level are very difficult to detect (as demonstrated by this paper!), we use a more conservative restriction of SNR $\sim 3$  to be realistic. InSight's noise measurements show that the martian day can be roughly split into low-noise and high-noise time periods. Assuming a typical low noise level of  $1.5 \times 10^{-9}$  m/s<sup>2</sup>/Hz<sup>1/2</sup> at 4 Hz (Lognonné *et al.*, 2020) is appropriate approximately 50% of the time, and the remaining 50% of the time it is too noisy for any detections, we predict just  $\sim 2$  detections of impact events per Earth year, during times when higher continuous rate data are collected. Furthermore, seismic amplitudes of signals from small craters could be even lower (Wójcicka *et al.*, 2020) resulting in even fewer detections. There are still large uncertainties on the predicted detection rate, at least an order of magnitude. However, given that we have yet to unambiguously detect any impacts in the seismic data, either the large end of this range is increasingly unlikely, or – more likely – we have not yet learned enough about martian seismic signals to recognize impacts in the data.

Because of this revised expectation of seismic detections of impacts, we have reversed our operational approach to detecting impacts. Instead of examining the seismic data for possible impact-induced signals, then following up with orbital images appropriate for the expected size and location of the impact, as described in Daubar *et al.* (2018), we are instead examining orbital images for new impacts, as indicated by dark spots or albedo changes near InSight. When more of these are found, we will examine the seismic data during the image-constrained time periods in a manner similar to the analysis presented here.

## 5 Conclusions

The exciting and lucky observation of a new impact occurring very close to the InSight lander during its prime mission presented a first opportunity to test our understanding of the seismic detectability of small impacts on Mars. Three potential candidate events were identified in the seismic data during the time period constrained by the before and after orbital CTX images; however, we are not able to determine that any of those seismic or atmospheric signals were definitively associated with that impact event. This is mainly because although the impact was nearby, it was quite small, forming only a  $\sim 1.5$  m diameter crater, and likely was created by a significantly decelerated impactor. We predict that the signals produced by this impact were very close to the measured minimum noise amplitudes seen by the InSight seismometers, and for a good portion of the time, the observed noise levels are well above the predicted impact signal amplitude. Thus a lack of detection for an impact of this size and at this location is disappointing, but not surprising.

There are many uncertainties in our predictions of seismic signals from the known impact, for example in converting crater size to seismic moment. The attenuation and scattering properties of the martian crust are not yet completely understood, nor is the velocity structure of the subsurface. Given the uncertainties in our predictions, it is still possible that the known crater was indeed responsible for one of the three candidate seismic events, although we cannot support that conclusion with our current knowledge. As InSight reveals more about the properties of the martian interior, the uncertainties in our predictions will be reduced. Future efforts at numerical modeling of this specific impact and coupled seismic modeling of the resulting wave propagation may reveal additional things to look for in the seismic or atmospheric data that may allow us to identify future impacts, if not this particular event. As we did not positively detect this impact in the data, we can at least conclude that we are not grossly underestimating the seismic amplitudes from impact events. Likewise, we see no definitive signals associated with this impact in the atmospheric data, nor do we expect that would be likely in this specific case.

The process of searching within the continuous seismic data from InSight for evidence of an event associated with an image-constrained impact has refined our understanding of impact-generated seismic signals through forward modelling and allowed us to reevaluate our predictions of impact detectability. Using the now-known noise levels of the SEIS instrument on



Mars, we expect  $\sim 2$  impact detections with  $\text{SNR} > 3$  each Earth year. This is assuming continued high sampling rates able to detect higher frequency peaks, which have lately begun. Our continued efforts to search orbital images for new dated impacts near InSight will almost certainly result in more image-constrained impacts. This work has provided a template workflow to help us quickly identify future impact seismic signals associated with image-detected craters. We continue to listen for impacts on Mars.

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All data used in this work are publicly available via the Planetary Data System (PDS; <https://pds.nasa.gov/>). Specifically, CTX images can be found at <https://pds-imaging.jpl.nasa.gov/volumes/mro.html>, HiRISE images can be found at <https://www.uahirise.org/>, and InSight APSS/TWINS/PS data can be found at [https://atmos.nmsu.edu/data\\_and\\_services/atmospheres\\_data/INSIGHT/insight.html](https://atmos.nmsu.edu/data_and_services/atmospheres_data/INSIGHT/insight.html). InSight SEIS data is available in the form of a seismic event catalogue (DOI: 10.12686/a6) and waveform data (DOI: 10.18715/SEIS.INSIGHT.XB\_2016) that are publicly available from the IPGP Datacenter and IRIS-DMC, as well as raw data available in the PDS at <https://pds-geosciences.wustl.edu/missions/insight/seis.htm>. Apollo seismic data are available in raw form at <https://darts.isas.jaxa.jp/planet/seismology/apollo/index.html>, and the data are available in SEED format from the IPGP Data Center for lunar data (Code XA, <http://datacenter.ipgp.fr/data.php>).

986 Seismic modeling results and parameters are available on the IGP data center at  
987 [https://doi.10.18715/JGR\\_NewCraterMod\\_2020](https://doi.10.18715/JGR_NewCraterMod_2020).  
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1137



## Supplemental material

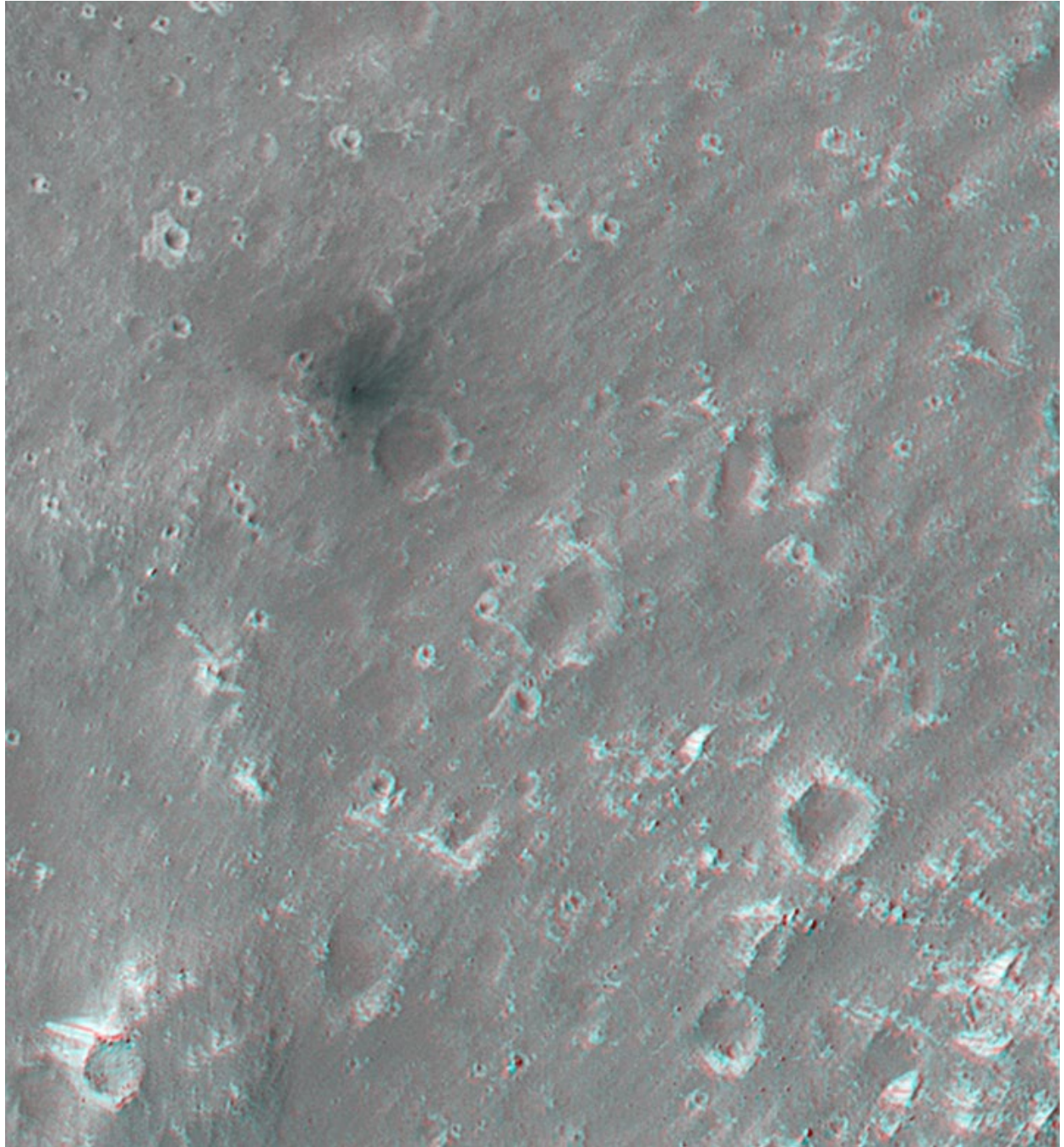
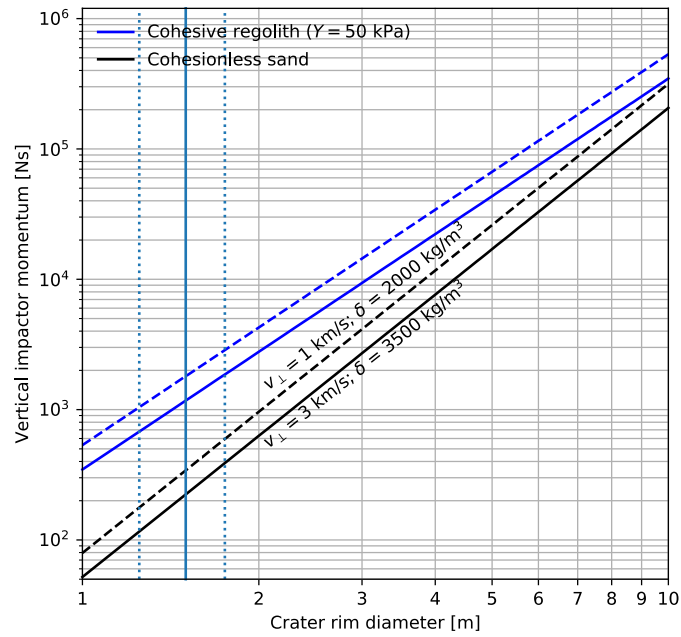


Figure S1.

*Red-blue anaglyph combining HiRISE images ESP\_060062\_1840 and ESP\_060128\_1840. The depth of the crater is subtle, but detectable by visual inspection of this anaglyph. The surrounding area is relatively flat and smooth. Image credit: NASA/JPL/University of Arizona.*



1146



1147

1148 Figure S2.

1149 *Impactor momentum (vertical component) as a function of crater diameter for meter-scale*  
 1150 *craters on Mars. Lines show estimates based on empirical scaling relationships, for two assumed*  
 1151 *target models: cohesionless sand (Eq. S3; black lines); cohesive regolith (Eq. S4; blue lines).*  
 1152 *Pairs of lines of the same color illustrate the small sensitivity to the combination of vertical*  
 1153 *component of the impact velocity and impactor density. Vertical blue lines indicate the observed*  
 1154 *crater diameter of 1.5 m (solid line) and associated uncertainty (dotted lines).*

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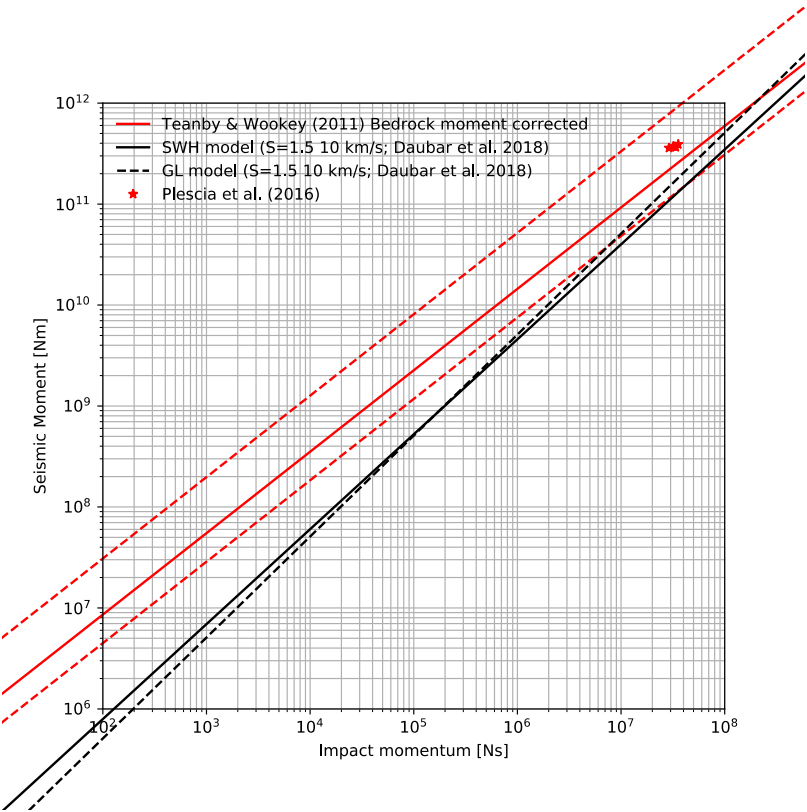
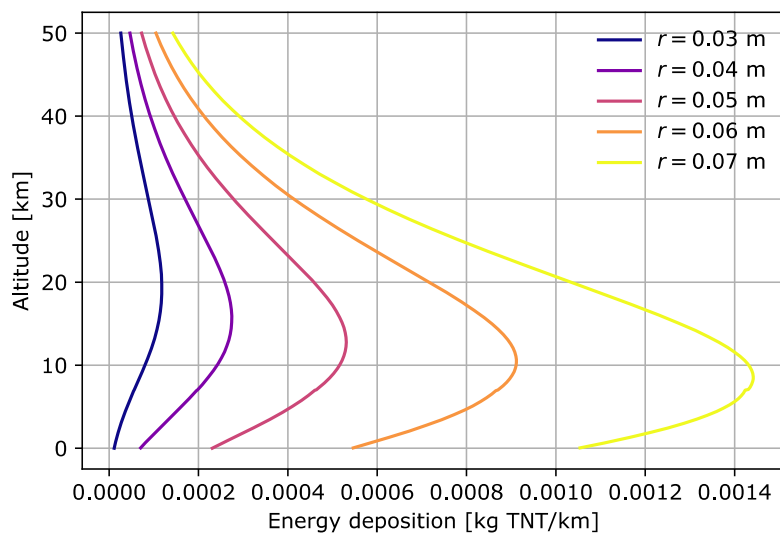
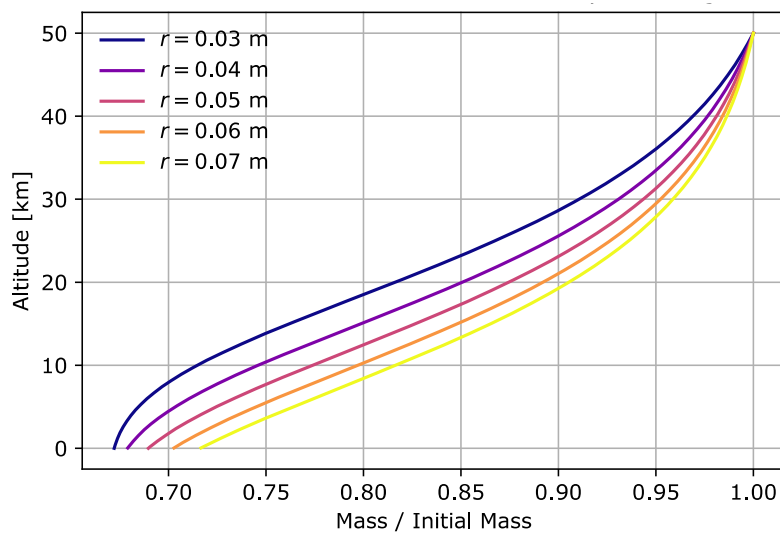


Figure S3.

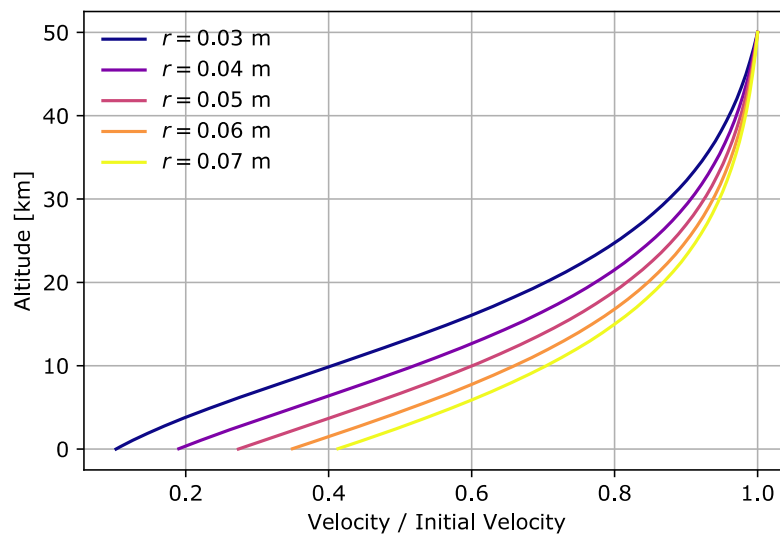
*Seismic moment as a function of impact momentum according to three semi-empirical models described in Daubar et al. (2018).*



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1164

1165 Figure S4.

1166 *Kinetic energy, mass and velocity loss during martian atmospheric entry of kg-size stony*  
1167 *meteoroids. Pre-entry meteoroid radius,  $r$ , is shown in the legend. Other meteoroid parameters:*  
1168 *pre-entry speed, 10 km/s; pre-entry angle to horizontal, 45 degrees; density, 2860 kg/m<sup>3</sup>;*  
1169 *strength, 1 MPa.*

1170

Table S1.

Atmospheric entry outcomes for several cases of meteoroids on Mars

| Scenario<br>Num. | Pre-entry <sup>1</sup> |              |                  |                | At ground      |              |                 |                      |                  |                | Crater<br>diameter |
|------------------|------------------------|--------------|------------------|----------------|----------------|--------------|-----------------|----------------------|------------------|----------------|--------------------|
|                  | Radius<br>[cm]         | Mass<br>[kg] | Momentum<br>[Ns] | Energy<br>[MJ] | Radius<br>[cm] | Mass<br>[kg] | Speed<br>[km/s] | Vert. vel.<br>[km/s] | Momentum<br>[Ns] | Energy<br>[MJ] | Sand<br>[m]        |
| 1                | 3                      | 0.3          | 2,290            | 16.2           | 2.6            | 0.2          | 1.03            | 0.73                 | 157              | 0.057          | 1.2                |
| 2                | 4                      | 0.8          | 5,420            | 38.3           | 3.5            | 0.5          | 1.90            | 1.35                 | 694              | 0.46           | 1.9                |
| 3                | 5                      | 1.5          | 10,600           | 74.9           | 4.4            | 1.0          | 2.74            | 1.93                 | 1,980            | 1.89           | 2.6                |
| 4                | 6                      | 2.6          | 18,300           | 129.3          | 5.3            | 1.8          | 3.49            | 2.47                 | 4,440            | 5.40           | 3.3                |
| 5                | 7                      | 4.1          | 29,100           | 200.5          | 6.3            | 2.9          | 4.13            | 2.92                 | 8,500            | 12.3           | 4.1                |

<sup>1</sup>Pre-entry parameters: speed, 10 km/s; angle, 45 degrees; density, 2860 kg/m<sup>3</sup><sup>2</sup>Crater rim diameter in sand estimated using Eq. S1; crater rim diameter in weakly cohesive regolith estimated using Eq. S2, with  $Y = 50$  kPa.Atmospheric entry model coefficients (see McMullan and Collins, 2019): Drag coefficient:  $C_D = 1.5$ ; Ablation parameter  $C_{ab} = 1.2 \times 10^{-8}$  kg/J.

Table S2

Upper crust model: Density,  $v_p$  and  $v_s$  models used for group velocity computations at depths of 5 meters to 1 km. At depths less than 5 meters, model B9 of Lognonné *et al.* (2020) is used, while at depths below 1 km, the TAYAK model is used (Smrekar *et al.*, 2019). This model is used for the computation of the surface waves group velocity shown in Fig. 3.

| Depth (m) | Density (kg/m <sup>3</sup> ) | $v_p$ (m/s) | $v_s$ (m/s) |
|-----------|------------------------------|-------------|-------------|
| 5         | 1399.5                       | 375.5       | 223.5       |
| 6         | 1408.6                       | 466.7       | 277.8       |
| 7         | 1416.3                       | 481.3       | 286.5       |
| 8         | 1422.7                       | 494.4       | 294.3       |
| 9         | 1428.2                       | 506.1       | 301.3       |
| 10        | 1433.0                       | 516.9       | 307.7       |
| 10        | 2600.0                       | 658.6       | 370.3       |
| 20        | 2600.0                       | 756.5       | 414.0       |
| 30        | 2600.0                       | 840.9       | 452.7       |
| 40        | 2600.0                       | 915.6       | 487.8       |
| 50        | 2600.0                       | 982.9       | 519.9       |
| 60        | 2600.0                       | 1044.4      | 549.6       |
| 80        | 2600.0                       | 1154.2      | 603.2       |
| 90        | 2600.0                       | 1203.8      | 627.7       |
| 100       | 2600.0                       | 1250.6      | 650.9       |
| 200       | 2600.0                       | 1618.4      | 835.9       |
| 300       | 2600.0                       | 1886.3      | 972.3       |
| 400       | 2600.0                       | 2102.5      | 1082.8      |
| 500       | 2600.0                       | 2286.2      | 1176.9      |
| 600       | 2600.0                       | 2447.3      | 1259.5      |
| 700       | 2600.0                       | 2591.7      | 1333.4      |
| 800       | 2600.0                       | 2723.0      | 1400.7      |
| 900       | 2600.0                       | 2843.8      | 1462.5      |
| 1000      | 2600.0                       | 2956.0      | 1519.9      |

## S1. Supplemental text for Section 2

Initial estimates of impactor parameters for the 1.5-m diameter crater on Mars were determined using a combination of empirical scaling relationships (Daubar *et al.*, 2018; Holsapple, 1993; Holsapple and Housen, 2007; Teanby and Wookey, 2011) and semi-analytical atmospheric entry modeling (McMullan and Collins, 2019). Empirical crater scaling relationships for m-scale craters formed under martian gravity in porous, cohesionless or weakly cohesive target materials indicate that crater diameter scales approximately as a power of the vertical component of the impactor momentum at the ground, with only a weak additional dependence on impactor density and speed (Holsapple, 1993; Holsapple and Housen, 2007) (Fig. S2).

For a cohesionless sand-like regolith, the scaling relationship is

$$D = 2.06r^{0.83} \left(\frac{g}{v_{\perp}^2}\right)^{-0.17} \left(\frac{\delta}{\rho}\right)^{0.33}, \quad (\text{S1})$$

where  $D$  is crater diameter,  $r$  is impactor radius,  $g$  is surface gravity,  $\delta$  and  $\rho$  are the impactor and target density, respectively, and  $v_{\perp}$  is the vertical component of the impact velocity. For a weakly cohesive, porous granular soil, the scaling relationship is:

$$D = 2.06r \left(\frac{Y}{\rho v_{\perp}^2}\right)^{-0.205} \left(\frac{\delta}{\rho}\right)^{0.4} \quad (\text{S2})$$

where  $Y$  is a poorly-defined measure of the cohesive strength of the soil. Note that in these expressions  $D$  is the crater diameter at the preimpact level; the rim-to-rim diameter  $D_r$  is typically 30% larger. The most comparable data for small craters on Mars are the artificial impacts on the Moon (the Ranger spacecraft and Apollo S-IVB boosters; Plescia *et al.*, 2016) and Earth (White Sands Missile impacts; Moore, 1976) that produced m-to-decameter-scale craters. These data suggest that Eq. (S1) provides an upper bound on crater size, while a lower bound is provided by Eq. (S2) for a regolith strength parameter  $Y \approx 50$  kPa.

Rearranging to find the vertical impact momentum  $p_{\perp}$  as a function of the rim-to-rim crater diameter, gives for a cohesionless sand:

$$p_{\perp} = 79.4 \left(\frac{D_r}{1 \text{ m}}\right)^{3.6} \left(\frac{g}{3.71 \text{ m s}^{-2}}\right)^{0.61} \left(\frac{\rho}{1500 \text{ kg m}^{-3}}\right)^{1.19} \left(\frac{\delta}{2000 \text{ kg m}^{-3}}\right)^{-0.19} \left(\frac{v_{\perp}}{1000 \text{ ms}^{-1}}\right)^{-0.23}, \quad (\text{S3})$$

and for a weakly cohesive regolith:

$$p_{\perp} = 821 \left(\frac{D_r}{1 \text{ m}}\right)^3 \left(\frac{Y}{100 \text{ kPa}}\right)^{0.62} \left(\frac{\rho}{1500 \text{ kg m}^{-3}}\right)^{0.59} \left(\frac{\delta}{2000 \text{ kg m}^{-3}}\right)^{-0.2} \left(\frac{v_{\perp}}{1000 \text{ ms}^{-1}}\right)^{-0.23}. \quad (\text{S4})$$

For a  $1.5 \pm 0.25$  m diameter crater, and using 50 kPa as an upper limit for  $Y$ , the predicted range of impactor momentum is  $\sim 100$ -3000 Ns ( $\text{kg m s}^{-1}$ ), depending most sensitively on the cohesion (or effective strength) of the regolith target material (Fig. S2).

The approximate energy of this impact is estimated to be between 0.1-2.0 MJ. This range is bounded by a cohesionless sand-like target or regolith scaling to get a 1.5 m diameter crater. The smallest energy (0.1 MJ) is between scenario 1 and 2 in Table S1, based on the crater



diameter lying between those results. The upper bound of 2 MJ comes from scenario 3, where the regolith crater diameter gives 1.5 m exactly. Because this is based on these selected individual scenarios, this is a rough approximation.

To further constrain the likely impactor properties, we calculated the expected deceleration and mass-loss of kg-scale meteoroids in the martian atmosphere by numerically solving the standard conservation equations of meteor dynamics (*e.g.*, McMullan and Collins, 2019). We neglect meteoroid fragmentation for several reasons. First, inspection of the newly-formed crater suggests that any fragments that separated from the main meteoroid were less than 0.1% in relative mass and therefore made no significant contribution to the impact-generated seismic signal. Second, for meteoroids in the relevant mass range and typical meteoroid strengths greater than 0.1 MPa, substantial break-up induced by aerodynamic stresses is not expected (except for extreme pre-entry speeds).

Our atmospheric entry modeling shows that ~0.1 to ~1-kg impactors undergo substantial deceleration and partial ablation in the martian atmosphere (Fig. S4, Table S1). Energy transfer to the atmosphere of >90% initial kinetic energy and deceleration to 30% of initial impactor speed are typical for such impacts, for a wide range of impactor masses, densities, strengths and trajectory angles (Table S1). Deceleration is greater when the pre-entry impactor speed is higher, implying that impact speeds at the ground are relatively insensitive to pre-entry speed. A vertical impact velocity at the ground of 1-3 km/s is predicted for a wide range of impactor trajectories, densities and masses. Only very shallow (<20 degrees) trajectory or high meteoroid density (>5000 kg/m<sup>3</sup>) scenarios result in a substantially different outcome. While a moderately oblique impact trajectory is most likely and consistent with observations, a very shallow trajectory impact scenario is not supported by the shape of the crater or blast pattern.

## S2. Supplemental material on seismic modeling.

### S2.1 Seismic signal recorded at InSight

We model the signal of the impact recorded at InSight using the seismic wave propagation code AxiSEM (Nissen-Meyer *et al.*, 2014), which solves the seismic wave equation in an axisymmetric geometry using the spectral element method. We simulate seismic frequencies up to 2 Hz, which is relevant to the frequencies of the 3 candidate events. The shallow subsurface structure significantly affects wave propagation because low velocity layers near the surface (*i.e.*, regolith) trap high amplitude surface waves (*e.g.*, Daubar *et al.*, 2018). Here, we consider two different models of the shallow subsurface of Mars. First, we use a simple pre-InSight model consisting of an 80 m layer of regolith overlying a layer of fractured bedrock that extends to 1 km depth, which Daubar *et al.* (2018) used to benchmark impact-generated seismic wave propagation on Mars. Next we use the recent shallow subsurface model of Lognonné *et al.* (2020), which is constrained in part by new seismic data from InSight. Below each of the shallow subsurface models, we assume the seismic velocity model EH45TcoldCrust1 (Rivoldini *et al.*, 2011). Figure S5 shows a comparison of the two subsurface velocity models. In the upper 100 m, both structures include a very seismically slow regolith layer, with shear velocity  $v_s < 1$  km/s, which makes the seismic waveform simulations computationally challenging at high frequency.

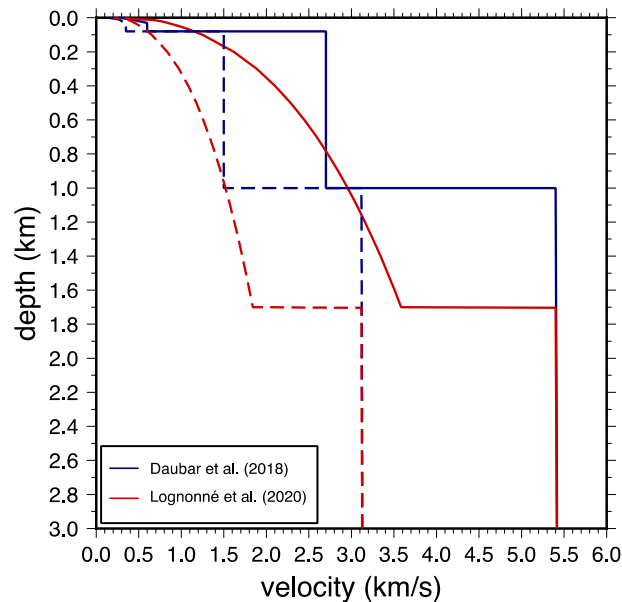


Figure S5.

Comparison of the two shallow subsurface seismic velocity structures used in the impact modeling. The blue line shows the pre-InSight estimate of structure.

To approximate the impact, we use an explosive source with a scalar moment of  $1 \times 10^7$  Nm (*i.e.*, the upper estimate of seismic moment from Section 2.2 in the main document) placed at the surface. We use a Gaussian moment rate function with a pulse width equal to the minimum mesh period of 0.5 s. The seismic source is located at  $3.866^\circ\text{N}$  latitude,  $135.613^\circ\text{E}$  longitude, and the receiver is located at  $4.502^\circ\text{N}$  latitude  $135.623^\circ\text{E}$  longitude. Figures S6 and S7 show the vertical and radial component synthetics computed with the shallow structural models of Daubar *et al.* (2018) and Lognonné *et al.* (2020) respectively. In both cases, the amplitudes are similar, although the duration and polarizations differ significantly. In the case of the pre-InSight shallow subsurface model (Fig. S6), the signal is strongly polarized on the radial component and lasts upwards of 100 s. These characteristics arise from the sharp discontinuity between the very slow regolith layer and the underlying bedrock near 80 m, which traps short period seismic energy in the shallowest subsurface. The strong scattering and regolith layers on Mars would likely generate a signal of extended duration (*i.e.* one with a surface wave coda), but the length of the wavetrain in these synthetics is due to the homogeneity of the regolith layer and the treatment of the regolith-bedrock boundary, and hence does not arise through the correct physical mechanism. In reality, the gradations in the actual regolith down to the bedrock boundary likely serve to dissipate some of this energy away from the peak. The structural model of Lognonné *et al.* (2020) includes a smooth gradient in shallow subsurface to a depth of 1.7 km. In this case, as seen in Figure S7, both the vertical and radial seismograms have similar amplitudes (*i.e.*, no strong radial polarization), and the signal duration is shorter ( $< 100$  s long). The waveforms are characterized by two wave pulses, both of which exhibit strong dispersion. The second pulse exhibits the largest peak amplitude of about  $1 \times 10^{-8}$  m/s, although this is likely to be an overestimate of the true signal amplitude as the effects of 3D structure are ignored. The modeled signal amplitudes are comparable to the seismic noise level during each of the candidate events, thus making them difficult to clearly identify in SEIS data.

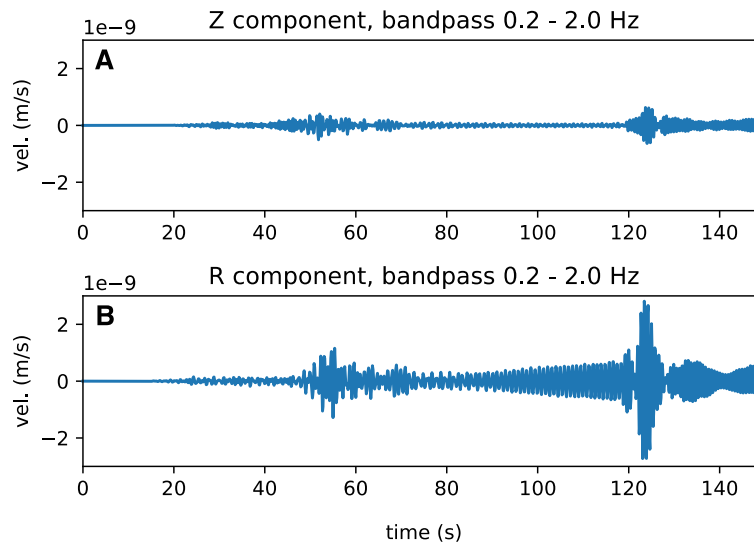


Figure S6.

*AxiSEM synthetics calculated for CTX-constrained impact event using the shallow structural model of Daubar et al. (2018). An isotropic source with a moment magnitude of  $1 \times 10^7$  Nm was used. Panels A and B show the vertical and radial component waveforms respectively. Synthetics are bandpass filtered between 0.2 – 2.0 Hz.*

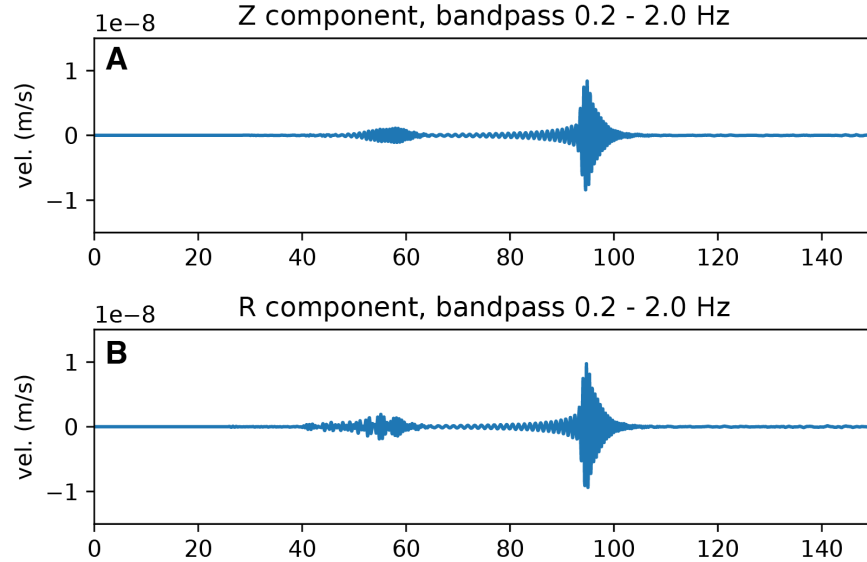


Figure S7.

*Same as Figure S6, but using the recent shallow subsurface model of Lognonné et al. (2020)*

## S2.2 Comparison with previous model benchmarking

To validate our modeling approach, we compare our simulation results to the martian impact modeling benchmark presented in Daubar *et al.* (2018), specifically the one shown in Figure 25 of Daubar *et al.* (2018). The velocity model is based on the structural model EH45TcoldCrust1 modified in the upper 1 km (*i.e.*, blue lines in Fig. S5). While the previous benchmark only modeled seismic waves at periods 5 s or longer, we extend the modeling so that it is accurate at frequencies up to 2 Hz. We use an isotropic seismic source with each diagonal moment tensor component equal to  $5 \times 10^{10}$  Nm and calculate the vertical velocity seismogram recorded at 50 km distance. Figure S8A shows the vertical velocity synthetic bandpass filtered between 0.05 – 0.2 Hz, which is in excellent agreement with the results shown in Figure 25 of Daubar *et al.* (2018). At periods of 5 s or longer the maximum amplitude is  $1 \times 10^9$  m/s. However, bandpass filtering between 0.05 – 1.0 Hz (Fig. S8B), the peak amplitudes are increased by 2 orders of magnitude. Additionally, the higher frequency waves reverberate in the shallow subsurface layer which generates a long coda of up  $> 100$  s (Fig. S8C).

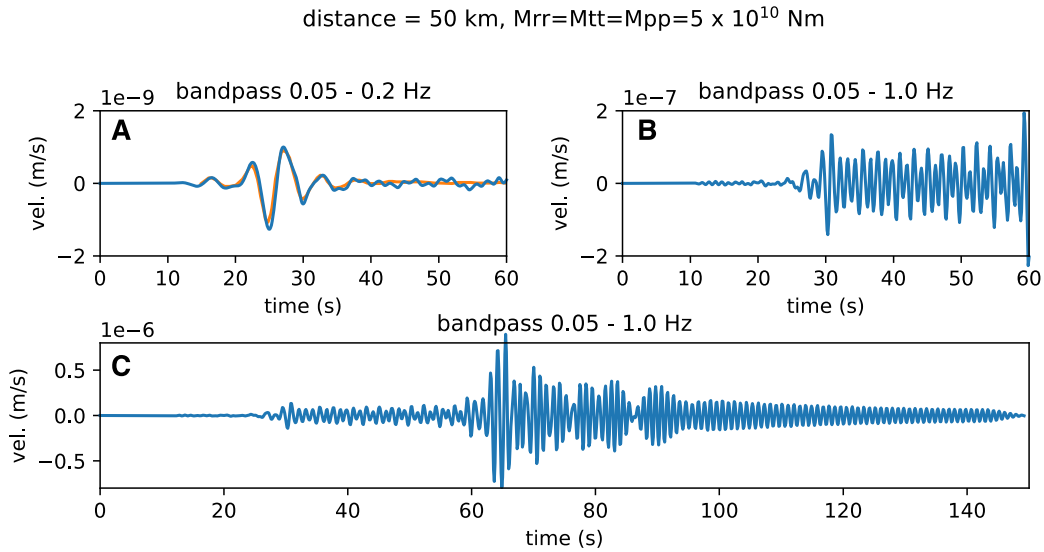


Figure S8.

Comparison with Daubar et al. (2018) benchmark. Panel A shows vertical velocity seismograms bandpass filtered between 0.05 and 0.2 Hz. The blue line shows the results from the present study and the orange line shows the AxiSEM results from Daubar et al. (2018), which were run at lower frequency. Panels B and C show the vertical velocity synthetics bandpass filtered between 0.05 and 1 Hz.

### S2.3 Future considerations on 3D effects and their modeling

Three-dimensional effects such as surface topography and subsurface velocity heterogeneities can scatter and distort seismic waves and thus influence seismic amplitudes, phases, and signal durations. However, modeling full 3D seismic wavefields at high frequency is computationally challenging, particularly for impact sources in a low-velocity planetary regolith. While a comprehensive modeling study of the effects of seismic scattering in the martian subsurface is beyond the scope of this work, we take this opportunity to compare the results and performance of commonly used seismic modeling codes to simulate a local impact event. Our intention is to explore the feasibility of current 3D codes to model high-frequency impact signals on Mars. Based on the parameters described in Section S2.2, we compare synthetic waveforms calculated using AxiSEM against two different commonly used seismic wave propagation codes: SPECFEM3D (Komatitsch and Tromp, 2002), and OpenSWPC (Maeda et al., 2017). SPECFEM3D and OpenSWPC are computationally intensive codes because they solve the elastic wave propagation problem on a fully 3D mesh, using the spectral element method and finite difference method respectively.

In order to simulate high frequency waves while still remaining computationally feasible, we limit the size of our domain in the 3D codes such that the maximum source-receiver distance

is 10 km. Another reason to limit the distances of the simulations is that SPECSEM3D and OpenSWPC modeling are both performed on a Cartesian mesh, whereas the mesh of AxiSEM conforms to the curvature of the planet's surface. All of the simulations use the recent shallow subsurface structure of Lognonné *et al.* (2020), but they are re-parameterized to limit parameterization issues introduced by the different formats used by the codes. The source is an explosive source with a moment magnitude of  $1 \times 10^7$  Nm, the upper bound estimate of section 2.2. The source time function is a Küpper wavelet with a pulse width of 1 s. The parameters used for this modeling benchmark are available in the IPGP data center at [https://doi.10.18715/JGR\\_NewCraterMod\\_2020](https://doi.10.18715/JGR_NewCraterMod_2020). Figure S9 shows a comparison of vertical component velocity synthetics for distances between 1 and 10 km modeled by each of the codes, filtered between 0.2 – 2.0 Hz. The modeled waveforms show good agreement, despite using different numerical methods and slightly different resolution due to different meshing strategies.

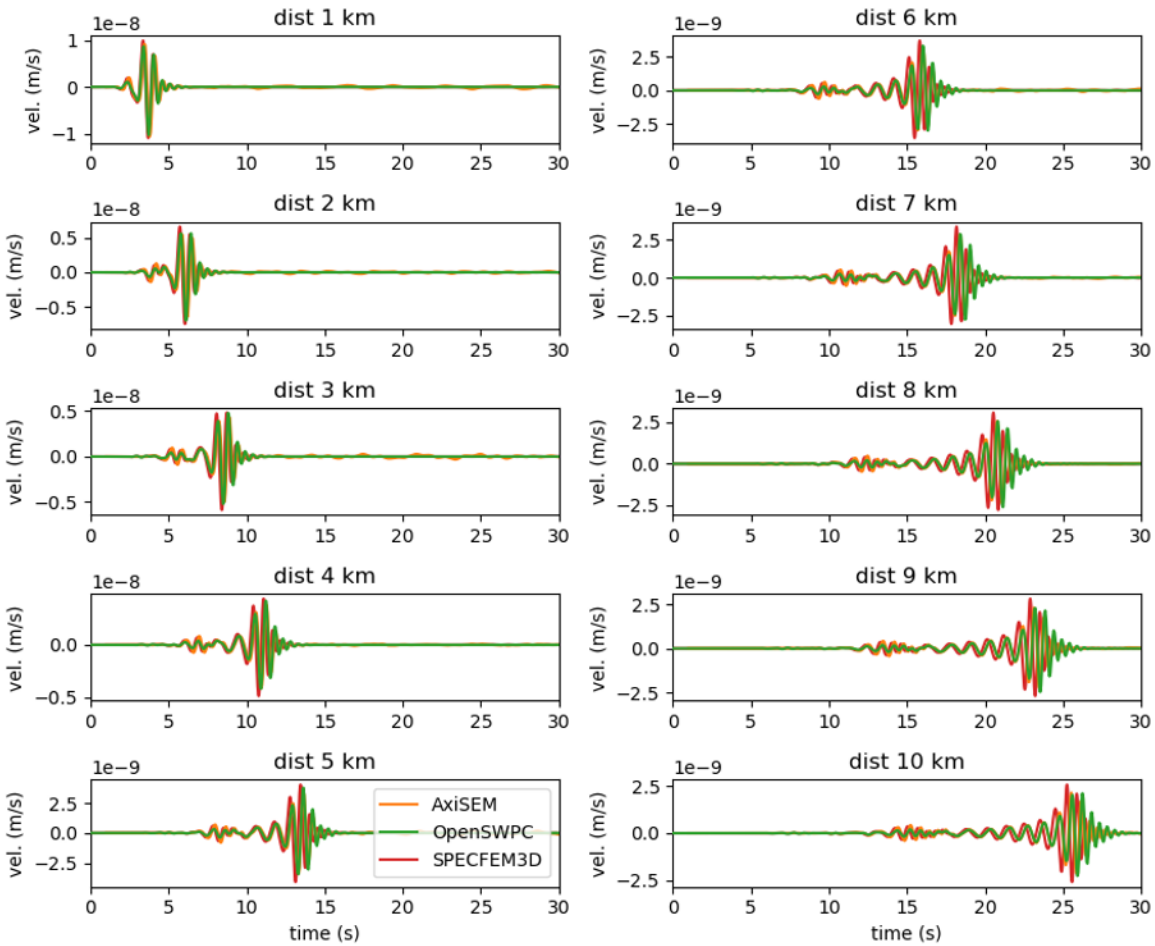


Figure S9. Comparison between 3D seismic waveform modeling codes and AxiSEM at distances between 1 – 10 km. All seismograms show vertical velocity. The source is an impact modeled as an explosion with a moment magnitude of  $1 \times 10^7$  Nm.

The tests shown here illustrate the potential of modeling the local signal of impact events on Mars in 3D at high frequency, and they also illustrate the modeling limitations. This work is in preparation for future modeling that will address how subsurface scattering and surface topography affect seismic signals in general, especially impact signals whose source is on the surface. Such work can be facilitated by limiting 3D modeling to local impact signals, which can be input into global scale simulation results to estimate signals at larger distances (*e.g.*, Schmerr *et al.*, 2019).

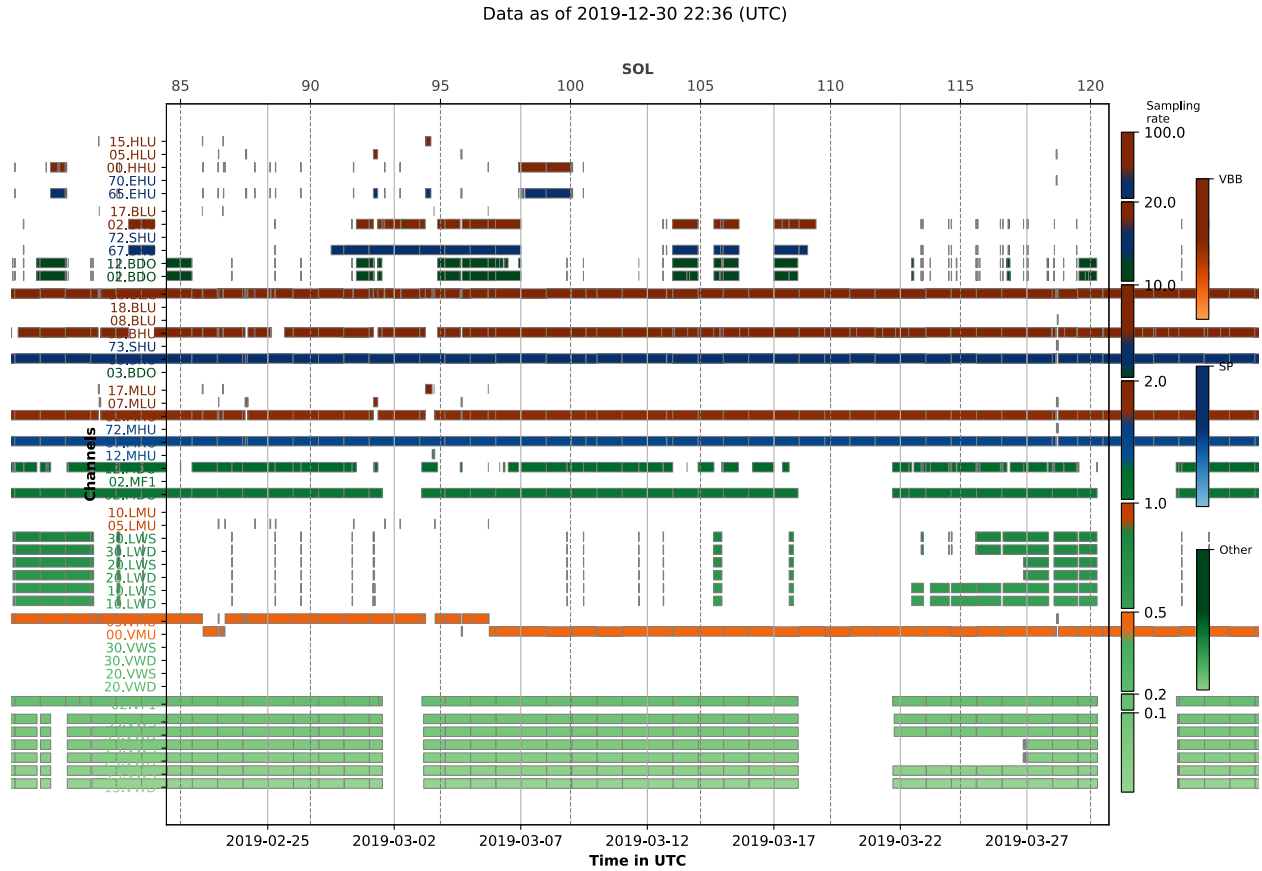


Figure S10.

Summary of available channels over the time period covering the impact search. Channels are ordered in terms of sampling rate (samples per second; sps) as indicated on the bar on the right. Colors indicate the sensor type. Channel naming conventions follow SEED standards as described in Lognonné et al. (2019). During the search period, 10 sps channels from VBB in high gain science mode (03.BH[UVW]) and SP (68.SH[UVW]) are generally always available. Occasionally, also 20 sps VBB and SP, as well as 100 sps SP, are available.



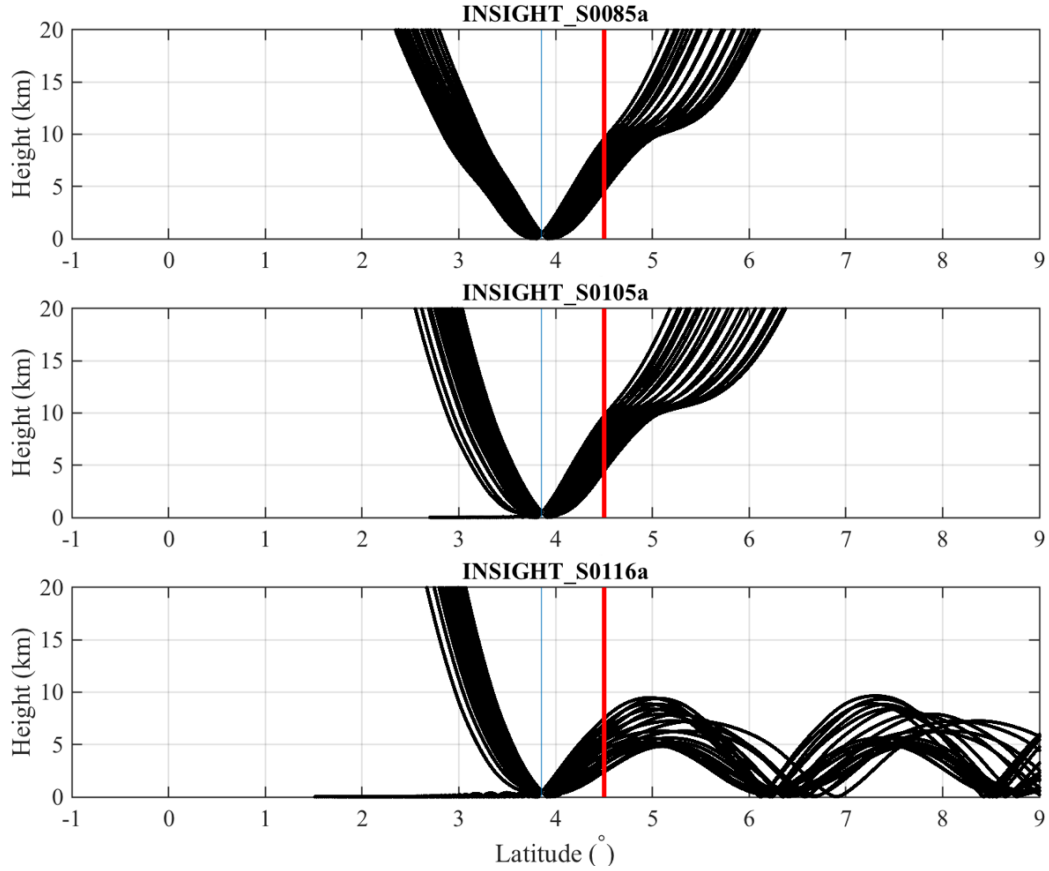


Figure S11.

Acoustic ray trajectories from a source located at the new impact crater at  $3.866^{\circ}\text{N}$  latitude (blue line). A standard climatological model is used to simulate the propagation medium at the time of each of the InSight events S0085a, S0105a, and S0116a. The InSight landing site to the north of the impact source is located in a shadow zone for such nearby acoustic source: none of the acoustic rays reaches the ground at InSight's location at latitude  $4.502^{\circ}\text{N}$  (red line). Differences in local times of the events imply changes in the atmospheric structure, causing significant differences in the morphology of acoustic propagation, with rays either refracted upward (S0085a and S0105a) or back downward (S0116a). In all cases, however, they only reach the ground far from the InSight lander (see Garcia et al., 2017 and references therein).

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