

A Prototype System for Dynamically Polarized Neutron Protein Crystallography

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Abstract

The sensitivity of Neutron Macromolecular Crystallography to the presence of hydrogen makes it a powerful tool to complement X-ray crystallographic studies using protein crystals. The power of this technique is currently limited by the relative low neutron flux provided by even the most powerful neu-

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tron sources. The strong polarization dependence of the neutron scattering cross section of hydrogen will allow us to use Dynamic Nuclear Polarization to dramatically improve the signal to noise ratio of neutron diffraction data, delivering order of magnitude gains in performance, and enabling measurements of radically smaller crystals of larger protein systems than are possible today. We present a prototype frozen spin system, built at Oak Ridge National Laboratory to polarize single protein crystals on the IMAGINE beamline at the High Flux Isotope Reactor (HFIR). Details of the design and construction will be described, as will the performance of the system offline and during preliminary tests at HFIR.

Keywords: polarized target, dynamic nuclear polarization, protein crystallography, neutron scattering

1. Introduction and Motivation

Macromolecular crystallography is a powerful tool to measure the structure of proteins. This is usually done by measuring the diffraction from a protein crystal in a high intensity x-ray beam. The incredibly high flux of these beams has resulted in the determination of tens of thousands of protein structures. Unfortunately, since x-rays scatter primarily from the high Z atoms in a sample, these structures often lack direct information about the location of the lighter elements, especially hydrogen. This remains true even when the diffraction data extend to atomic resolution. Neutron Macromolecular Crystallography (NMC) presents a potential solution to this problem. The relatively high cross section for neutron scattering from hydrogen enables neutron diffraction to determine the hydrogen positions in proteins even at

13 medium resolutions (2.0\AA), information that is otherwise unattainable [1][2].
14 The greatest hurdle to the wide spread utilization of this technique is the rel-
15 atively low neutron flux available at even the most intense neutron sources
16 in the world [3][4]. This flux limitation requires data sets to be taken over
17 longer time, typically days or weeks, and more critically, using crystals that
18 are several orders of magnitude larger than those required at an x-ray facil-
19 ity. This crystal size limitation ends up the limiting factor for the technique,
20 requiring researchers to spend considerable amounts of effort to grow large
21 enough crystals, often without success. The spin dependence of the neutron-
22 hydrogen cross section presents a possible resolution to this problem [5].

23 The cross section for neutron scattering can be split into coherent and
24 incoherent parts. For a crystal sample, the coherent scattering contains in-
25 formation about the positions of atoms in the crystal. Incoherent scattering
26 enters as a background. Both the coherent and incoherent cross sections
27 display a large spin dependence, especially in the case of polarized neutrons
28 scattering from polarized hydrogen. The coherent cross section is given by
29 equation 1 and the incoherent is given by equation 2, where b_0 is the spin
30 independent part of the scattering amplitude, b is the spin dependent part of
31 the scattering amplitude, p is the neutron polarization, P is the nuclear po-
32 larization, and I is the spin of the nucleus [6]. A plot of the spin dependence
33 of the cross section is shown in Figure 1. The flux (or crystal size) needed
34 for a particular measurement, is reduced by the square of any increase in
35 the signal to noise ratio. This can result in order or magnitude or greater
36 reductions in crystal size without the need for deuteration.

$$\left(\frac{d\sigma}{d\Omega}\right)_{coh} = b_0^2 + 2b_0bIpP + b^2I^2P^2 \quad (1)$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{inc} = \frac{b^2}{4}[I(I+1) - pPI - P^2I^2] \quad (2)$$

37 The relatively low magnetic moment of the hydrogen nucleus means that
 38 even for a high magnetic field (5 T) and low temperature (1 K), the nuclear
 39 polarization for hydrogen is low (about 0.5%). Dynamic Nuclear Polarization
 40 (DNP) [7] uses a combination of high field, low temperature and microwave
 41 irradiation to produce very high (potentially 90% or higher) nuclear polar-
 42 ization in samples which have been prepared by the addition of unpaired
 43 electrons. This technique has been applied previously to Small Angle Neu-
 44 tron Scattering (SANS) measurements [8][9]. Tests with polarized single
 45 crystals of small molecule systems have also been conducted [10]. In order
 46 to test the effectiveness of DNP in improving NMC measurements of protein
 47 crystals, a prototype system was built to polarize these crystals, and mea-
 48 surements were made using a beam of polarized neutrons. This prototype
 49 will be described below.

50 **2. DNP Apparatus**

51 *2.1. Requirements*

52 The DNP system used for NMC measurements shares aspects from the
 53 more commonly seen systems used in medium energy nuclear physics [11][12][13]
 54 or NMR [14]. In order to not block the acceptance required for diffraction
 55 measurements, the system has to run in frozen spin mode. The refrigera-
 56 tor, magnet, microwave, and NMR systems will be described in the following

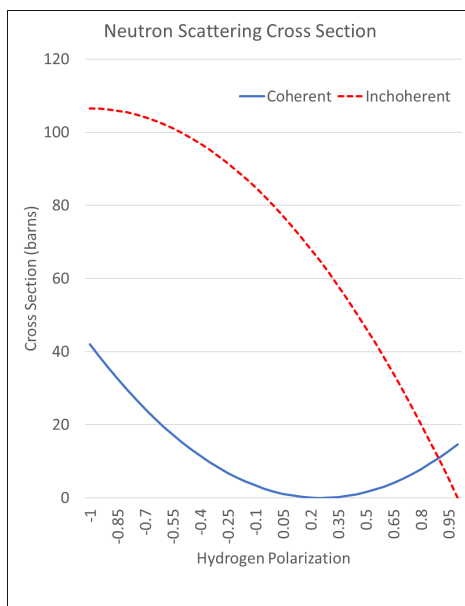


Figure 1: The coherent and incoherent cross sections for polarized neutrons scattering from a hydrogen nucleus, as a function of the nuclear polarization, as calculated in [6].

57 sections, as will the custom sample interface which was required for compat-
 58 ibility with the protein crystal samples to be polarized.

59 2.2. Refrigerator

60 The refrigerator used was a commercially available, cryogen-free, dilution
 61 refrigerator⁶. This was fitted with a custom tail and sample interface, in-
 62 cluding a sample loading system (see section 2.6). Before the addition of the
 63 sample interface, the refrigerator was capable of reaching a base temperature
 64 of 7 *mK*. The refrigerator has a listed cooling power of 400 μW at 100 *mK*.
 65 Tests at higher temperatures showed that it was capable of 12 *mW* of cooling

⁶Bluefors Oy, Arinatie 10, 00370, Helsinki, Finland

66 at 1 K .

67 *2.3. Magnet*

68 The magnet used was a 5 T , warm bore, cryogen-free, vertical magnet⁷.
69 The magnet has a 100 mm clear bore, and a better than one part in 10^4
70 uniformity over a 2 cm diameter spherical volume. In order to facilitate the
71 transition from the polarizing configuration to the data taking configuration,
72 the magnet was mounted on a custom designed lift table. This table was
73 capable of manually raising or lowering the magnet while it was at full field.
74 The minimum time to raise or lower the magnet was limited by the eddy
75 current heating in the refrigerator and sample space, but usually was less
76 than ten minutes. A diagram of the system in polarizing mode and data
77 taking mode is shown in Figure 2.

78 *2.4. Microwaves*

79 The small sample size ($< 1 \text{ mm}^3$) allowed the use of a relatively inex-
80 pensive solid state microwave source⁸. This source was capable of generating
81 300 mW of microwave power at 140 GHz . The frequency was extremely
82 stable, and set via USB communication with the device. A variable attenua-
83 tor and a power meter were used to control the amount of microwave power
84 sent into the system. For the cryogenic wave-guide, a 0.188 $in.$ O.D. copper
85 nickel tube was used. A sheet of brass at the end of the wave-guide was used

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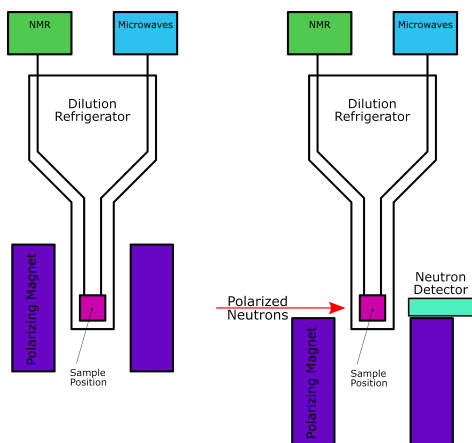


Figure 2: A diagram of the system in polarizing mode (left) and data taking mode (right). In order to move between the two configurations, the magnet was moved down, and a detector moved in close to the sample.

86 to reflect the microwaves at 90° so that they would intercept the sample
 87 position.

88 2.5. NMR System

89 The polarization was monitored using a Liverpool [15] q-meter and a
 90 pre-amplifier card which provided a gain of 500. This amplification was
 91 necessary to see the thermal equilibrium (TE) polarization signal for the
 92 extremely small samples used. Even with this amplification, the TE signal
 93 area could not be measured reliably enough to calibrate the NMR system
 94 through the standard technique of comparing the TE signal to the known
 95 TE polarization (although increases in the signal size were clear evidence of
 96 polarization). In order to estimate the maximum polarization, larger (a few
 97 μL) frozen samples (of a solution which had the same chemical composition
 98 as the crystals) were polarized. These samples were large enough to use the

99 TE signal for calibration. A number of these samples were measured, and
100 with some variation, showed maximum polarizations of about 50%, usually
101 after 8 hours or more of polarizing. The error on these measurements was still
102 quite high (20% or higher relative error), as the low cooling and microwave
103 power of the system limits the maximum sample size significantly. This error,
104 combined with the unknown systematic error that comes from comparing the
105 large sample polarizations to that of the crystal, results in a system that is
106 only capable of qualitative estimates of the polarization. The NMR system
107 was, however, sufficient to monitor the relative decay of the polarization,
108 since the polarized signals are two orders of magnitude larger than the TE
109 signals needed for calibration.

110 *2.6. Sample Space*

111 The sample space and loading system has unique requirements, owing to
112 the nature of the samples used. Protein crystals contain about 50% solvent
113 by volume, and in order to retain their crystal structure, must be frozen very
114 quickly to vitrify this solvent. This is generally accomplished by immersion
115 in liquid nitrogen. Once frozen, the samples are mechanically very robust,
116 but must remain below around 140 *K* to prevent the formation of crystalline
117 ice, which would disrupt the protein lattice. Frozen crystals must therefore
118 be loaded with the refrigerator already cooled. Protocols for NMC at cryo-
119 genic temperatures have been developed [16] and demonstrated previously on
120 IMAGINE [17][18]. As an additional requirement, while the sample is being
121 polarized, it must have sufficient thermal contact with the mixing chamber
122 to remove the heat deposited by the microwaves. The very high neutron
123 capture cross section of ^3He makes it impossible to simply locate the sample

124 in the mixing chamber, so some other thermal link must be established. The
125 most effective thermal link from the mixing chamber to the sample is a high
126 conductivity tube, filled with super-fluid helium.

127 The sample was loaded though a thin walled stainless steel tube, which led
128 from the atmosphere side of the the refrigerator top plate, down to the mixing
129 chamber plate. This tube was thermally anchored at the 40 *K*, 4 *K*, and
130 still plates of the refrigerator using copper clamps and commercially produced
131 C101 copper braided connections. At the mixing chamber position, the tube
132 switched from stainless steel to copper, with a thick flange attaching the
133 copper to the mixing chamber plate. This thick walled copper tube extended
134 down the tail section of the refrigerator, into the magnet. The last section of
135 the tube was made of PCTFE (polychlorotrifluoroethylene, a material similar
136 to Teflon) and attached via an indium seal to a stainless steel flange at the end
137 of the copper tube. The PCTFE sample space had a thin walled (0.025 *in.*)
138 section to accommodate the beam and the diffracted neutrons. The bottom
139 of the PCTFE sample space was fitted with a second indium seal, where a
140 ceramic, SMA style, RF feed-through was connected in order to get an NMR
141 coil as close to the sample as possible. A picture of the PCTFE sample
142 space can be seen in Figure 3. The vacuum can and radiation shields had
143 thin aluminum windows along the beam path to minimize the background as
144 much as possible. The vacuum windows were 0.01 *in.* thick, while the three
145 radiation shields were less than 0.001 *in.* thick along the beam path.

146 The sample itself was located in a 0.125 *in.* O.D. PCTFE capillary,
147 mounted via a short aluminum transition piece, to the sample stick. The
148 stick consists of a long, fiberglass reinforced epoxy (G10) shaft, glued into

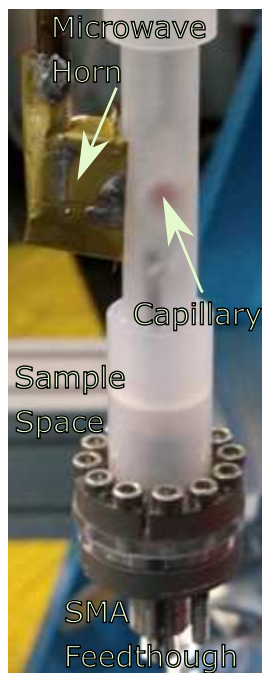


Figure 3: A photograph of the PCTFE section of the sample space. The microwave reflector is visible in the left side of the image. The SMA feed-through and lower indium seal are also visible. The sample has been indicated in this case by painting the capillary red (visible through the PCTFE).

149 an aluminum top section. The aluminum top section is keyed, allowing it
150 to be inserted or removed using a tool designed specifically for this purpose.
151 The top of the stick is located at the level of the dilution refrigerator mixing
152 chamber. The top of the stick also has a spring energized Teflon seal, which
153 causes it to make tight contact with the inner wall of the sample space at the
154 position of the mixing chamber interface. This seal reduces the amount of
155 gas and superfluid that entered the section of the loading tube which leads
156 to room temperature. A thin walled 0.125 *in.* stainless steel tube is used to

157 supply helium gas below this seal. This tube was thermally anchored at the
158 40 K , 4 K , and still plates of the refrigerator.

159 **3. Operation**

160 Crystals were soaked for 30 to 60 minutes in a solution containing 100 mM
161 hydroxy-TEMPO and flash frozen (with some of samples first being passed
162 through a glycerol cryo-protectant solution) in liquid nitrogen before being
163 placed in the capillary and stored under liquid nitrogen. In order to ensure
164 that the crystals didn't warm during loading, samples were loaded into the re-
165 frigerator while it was in standby mode (no helium circulating in the dilution
166 unit, pulse tube at base temperature). After the refrigerator recovered from
167 the heat associated with inserting the sample (which was at liquid nitrogen
168 temperature), helium was condensed in the dilution unit. Once the dilution
169 unit was running, helium was flowed through the small stainless steel tube
170 into the sample space, through a mass flow controller. After approximately
171 16 mL of helium had been condensed in the sample space, the mixing cham-
172 ber and sample temperature (as measured by a pair of resistors mounted
173 to the SMA feed-through) were within about 35 mK of each other, corre-
174 sponding to the lowest achievable sample temperature. This temperature
175 varied from sample to sample, but generally the lowest sample temperature
176 attainable was between 200 and 250 mK . This variation is presumed to be
177 primarily related to the effectiveness of the energized Teflon seal, which can
178 change due to the presence of ice or other contaminants, either on the seal
179 itself or on the mating surface in the sample space.

180 With the magnet at full field, and raised till the sample was in the uniform

181 region, microwaves would be transmitted to sample. As measured at the top
182 of the cryostat, the optimum microwave power to produce high polarization
183 was 10-20 *mW*, corresponding to about a *mW* at the sample position. Polar-
184 ization was generally quite slow, requiring eight or more hours of microwave
185 irradiation, so the samples were usually polarized overnight in order to op-
186 timize operational efficiency. The microwaves were then turned off, and the
187 sample allowed to cool. Re-cooling was relatively quick, requiring about 30
188 minutes. The magnet was then lowered, and neutron diffraction data was
189 collected for a time equivalent to the spin lattice relaxation time T_1 . This
190 means that the effective average in-beam polarization was about half of the
191 starting polarization. The magnet could then be moved back up and the
192 process repeated. Optionally, the sample temperature could be increased
193 while in the data taking position by means of a heater located on the mixing
194 chamber plate. Raising the temperature shortens T_1 greatly, allowing for
195 unpolarized data to be taken for comparison to the polarized results.

196 **4. Performance and Results**

197 Over a dozen protein crystal samples were tested at the IMAGINE beam-
198 line over two separate HFIR operating cycles under different conditons (in-
199 cluding room temperature tests, cold tests, and polarized tests) in order to
200 first commission the apparatus and the detector, and then measure the ef-
201 fect of the polarization on the diffraction. The samples were hen-egg-white
202 lysozyme and T4-lysozyme, well characterized standard proteins that are
203 typical of the medium unit cell (around 100Å) systems that are amenable to
204 neutron analysis. The crystals were of varying size, with most of them being



Figure 4: Apparatus in the data taking position. Detector shown with a 3-D printed shield to reduce background.

205 roughly cubic and approximately $600 \mu\text{m}$ on edge. IMAGINE is a quasi-Laue
206 diffractometer which utilizes a pair of elliptical neutron mirrors to produce
207 a focused neutron beam spot [19]. The standard detector was not utilized
208 for these proof of principle tests, since it was not possible to integrate it
209 with the polarizing magnet. Instead, a single arm detector was built, using a
210 prototype Anger camera which has been developed at ORNL. This detector
211 uses silicon photo diodes instead of the usual photo-multiplier tubes. This
212 allowed the detector to operate in the fringe field of the polarizing magnet.
213 Data was taken with the detector approximately 20 cm from the sample
214 position, over a 116 mm by 116 mm area. The detector was positioned at

215 different angles, between 0° and 60° . The neutron beam was polarized using
216 a V-cavity supermirror polarizer⁹.

217 About half of the crystals both polarized and produced sufficient diffrac-
218 tion to test the effect of polarization on the diffraction pattern (the crys-
219 tals that both polarized and diffracted were all T4-lysozyme crystals, which
220 is the sample material that was used for early offline tests). The spin re-
221 laxation time for the four best diffracting samples varied from 0.7 to 3.0
222 hours. All of the T4-lysozyme samples polarized to some extent. Measure-
223 ment of the diffraction showed a substantial enhancement in the intensity
224 for the case where the sample was polarized anti-parallel with the beam,
225 and a reduction in intensity for polarization parallel with the beam (relative
226 to the unpolarized intensity). Figure 5 shows a diffraction image for both
227 the polarized and unpolarized data and an intensity profile for an individual
228 reflection. Due to the limited acceptance of the detector and the unknown
229 orientation of the crystal, diffraction patterns were not indexed. Instead,
230 the integrated intensity for sub-sets of reflections was compared to average
231 enhancements calculated using estimates of the sample polarization (deter-
232 mined from NMR measurements). For the anti-parallel case, the integrated
233 intensity of the diffraction pattern typically increased by factors of 2 to 3,
234 with enhancements of individual reflections within the diffraction pattern
235 varying, depending on the relative contribution of hydrogen to that particu-
236 lar reflection. Comparing individual diffraction measurements this way has
237 the potential for large systematic errors, but the results of these averaged

⁹SwissNeutronics AG, Bruehlstrasse 28, CH-5313 Klingnau, Switzerland

238 measurements were consistent with samples starting with about 50% polar-
239 ization in the highest cases and decaying over the course of the data taking
240 as expected. This qualitative consistency between NMR and diffraction data
241 provides some confidence as to the level of polarization being produced in
242 the samples. A detailed description of the sample preparation and the ex-
243 perimental results will be discussed in a future publication.

244 **5. Conclusion and Outlook**

245 Tests using the system described in this paper have shown that it is
246 possible to increase the intensity of diffraction from a protein crystal by
247 manipulating the nuclear polarization. The test apparatus built is sufficient
248 to polarize these samples and allow neutron scattering data to be taken,
249 although the data collection time is limited by the relatively high sample
250 temperature in frozen spin mode. Since the tests described were completed,
251 the sample interface has been redesigned to lower the sample temperature and
252 correspondingly increase the spin relaxation time, with the goal of increasing
253 the average polarization during data collection. The upgraded apparatus is
254 scheduled to be tested in the IMAGINE beam-line in 2019.

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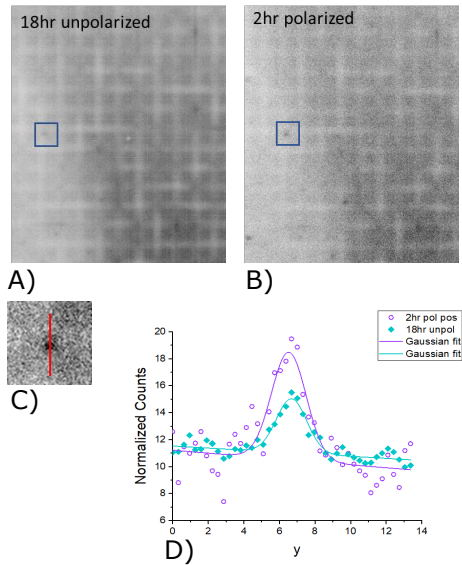


Figure 5: A) Eighteen hours of diffraction data with an unpolarized sample, showing a large fraction of the detector coverage. B) Two hours of diffraction data in the same configuration, but with the sample polarized anti-parallel to the polarization of the neutron beam. C) The time normalized profile of the detector counts, along with a Gaussian fit, in a single reflection from two different images, showing an enhancement in the polarized case. D) A detail showing the profile that was plotted. Both of these detector images were acquired using a T4-lysozyme crystal, a quasi-Laue polarized neutron beam, and the detector positioned at approximately 45 degrees. The contrast has been scaled to the data collection time.