

# Calcium and zinc differentially affect the structure of lipid membranes

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## ABSTRACT

Interactions of calcium ( $\text{Ca}^{2+}$ ) and zinc ( $\text{Zn}^{2+}$ ) cations with biomimetic membranes made of dipalmitoyl-phosphatidylcholine (DPPC) were studied by small angle neutron diffraction (SAND). Experiments show that the structure of these lipid bilayers is differentially affected by the two divalent cations. Initially, both  $\text{Ca}^{2+}$  and  $\text{Zn}^{2+}$  cause DPPC bilayers to thicken, while further increases in  $\text{Ca}^{2+}$  concentration result in the bilayer thinning, — eventually reverting to having the same thickness as pure DPPC. The binding of  $\text{Zn}^{2+}$ , on the other hand, causes the bilayers to swell to a maximum thickness, and the addition of more  $\text{Zn}^{2+}$  does not result in a further thickening of the membrane. Agreement between our results obtained using oriented planar membranes and those from vesicular samples implies that the effect of cations on bilayer thickness is the result of electrostatic interactions, rather than geometrical constraints due to bilayer curvature. This notion is further reinforced by MD simulations. Finally, the radial distribution functions reveal a strong interaction between  $\text{Ca}^{2+}$  and the phosphate oxygens, while  $\text{Zn}^{2+}$  shows a much weaker binding specificity.

## INTRODUCTION

Plasma membrane properties such as membrane fluidity, bending and compressibility moduli, electrostatics, and aggregation and fusion are associated with ions that are present in the cytosol and the extracellular fluid. Not surprisingly, divalent metal cations have been found to play a prominent role with regards to bilayer structure.<sup>1</sup> Metals such as calcium, magnesium, iron, manganese, copper, zinc, nickel and cobalt, although known to be toxic beyond a certain concentration, are essential to life. Among the large group of divalent metal cations, calcium and zinc attract a special interest due to their peculiar properties.

Calcium belongs to the alkaline earth metals and plays an important role in many cellular processes.<sup>2-4</sup> It has been shown to contribute to specific interactions within bacterial membranes that ~~resulted in the~~ limitation of water penetration, to mention but one example.<sup>5</sup> Zinc, amid the first-row of transition metals, is ~~also found likewise~~ attractive to study as it is second only to iron in terms of abundance and importance in biological systems.<sup>6</sup>  $\text{Zn}^{2+}$  is known to play a fundamental role in several critical cellular functions such as protein metabolism, gene expression, the structural and functional integrity of biomembranes, and in metabolic processes.<sup>7,8</sup> Although a complete understanding of the physicochemical processes taking place in biomembranes has yet to be established, their functionality is known to depend strongly on the type of ion, the chemical composition of the membrane's interface, its thermodynamical state, and degree of hydration. For example, both  $\text{Zn}^{2+}$  and  $\text{Ca}^{2+}$  have been found to deviate from an inversely linear correlation between the ion's ionic radius and its effect on shifting the lipid phase transition temperature.<sup>1</sup> What makes these cations different from others, and from each other, remains an open question.

In spite of the importance of divalent cations in biological membranes, their adsorption to phosphatidylcholine membranes and their influence on lipid bilayers is far from being understood.<sup>9,10</sup> While dehydrated planar stacks of multibilayers swelled upon the addition of different cations without any apparent structural changes to bilayers themselves<sup>11</sup> — similarly to the results of high monovalent salt concentration studies<sup>12,13</sup> suggesting ~~however~~ the possibility of phase separation<sup>14</sup> — a more complex behavior was displayed by fully hydrated, curved bilayers in the presence of  $\text{Ca}^{2+}$ .<sup>10,15,16</sup> Effects due to curvature, nevertheless, may be entangled with those due to electrostatic interactions, as has been suggested recently.<sup>17</sup>

It is important to decouple effects due to electrostatic interactions from those imposed by geometrical constraints found in curved vesicular bilayers, in order to associate the mode of

interaction of the individual cations with the lipid membranes. Small-angle neutron-diffraction (SAND) experiments were designed to utilize planar bilayers, extending previous studies on the effects of calcium and zinc on biomimetic membranes of dipalmitoyl-phosphatidylcholine (DPPC). SAND allows for the in situ manipulation of sample conditions, but more importantly, provides quantitative data on the distribution of structural moieties, their sizes, shapes and correlation lengths.<sup>18</sup> Variations in diffraction maxima can signal important changes to membrane structure, which in turn, may have biological implications. Our current observations agree well with the notion that ions affect the interface of membranes, results that can most likely be rationalized in terms of electrostatic interactions, rather than geometrical constraints due to bilayer curvature. Our results then highlight the importance of these cations with regards to membrane structure, and possibly function.

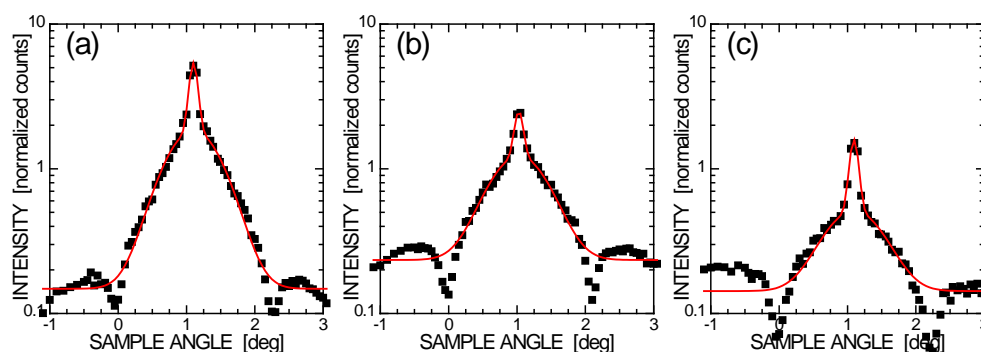
The refinement of the structural changes measured experimentally on the molecular scale was also performed on the atomistic level using results from molecular dynamics (MD) simulations. Although MD simulations offer structural resolution not attainable by experiment, prudence must nonetheless be exercised.<sup>13,19,20</sup> Based on a combination of experimentally obtained SAND data and MD simulations, we are able to observe structural changes in lipid bilayers due to interactions with  $\text{Ca}^{2+}$  and  $\text{Zn}^{2+}$  cations, and assign differences between the cations to differences in the specificity of their interactions with the membrane.

## MATERIALS AND METHODS

### Sample preparation

1,2-Dipalmitoyl-*sn*-glycero-3-phosphocholine (DPPC) was purchased from Avanti Polar Lipids (Alabaster, AL) and used without further purification.  $\text{CaCl}_2$  and  $\text{ZnCl}_2$  salts, and organic solvents were obtained from Fisher Scientific Company (Ontario). Approximately 15 mg of lipid

(thin film comprising of ~2,500 bilayers ~~with~~having a total thickness ~15  $\mu\text{m}$  when spread onto a 25 x 60  $\text{mm}^2$  silicon wafer) was co-solubilized with appropriate amounts of salt in a 1:1 (vol/vol) solution of chloroform:trifluoroethanol, or in deionized water (18.3  $\text{M}\Omega\text{ cm}$ ), and mixed thoroughly using a vortex mixer, or several freeze-thaw cycles in the case of the water solution samples. The dispersions prepared in organic solvent were deposited onto silicon wafers and rocked during evaporation in a glove box to form highly oriented multilayer stacks, while the remaining traces of solvent ~~traces~~ were removed by placing the samples under vacuum.<sup>21</sup> High quality oriented stacks were confirmed by observing 7-9 orders of diffraction peaks, and the narrow width of the central peak in the rocking curves shown in e.g., Fig. 1 and Figs. S1-S3.



**Figure 1.** Uncorrected rocking curves (i.e., normalized intensities measured as a function of sample angle) from aligned (a) DPPC multibilayers, and DPPC mixed with  $\text{Ca}^{2+}$  at 1:0.2 molar ratio in (b) solvent or (c) water solution. Data were collected with the detector located at the position of the first Bragg diffraction. The sample mosaicity, estimated from the Gaussian fit to the central narrow peak, is ~0.06 degree for all three samples. Note that two minima are the result of increased neutron absorption, which occurs when the sample is rotated so that either the incident beam or the diffracted beam is parallel to the substrate. These “wings” were not corrected as was done elsewhere.<sup>22</sup>

Rocking curves for samples with increasing concentration of cations (shown in Fig. S2 of ~~the~~ supporting information) revealed an ever-decreasing degree of sample alignment. This was more pronounced in the case of  $\text{Zn}^{2+}$  ions, where a visual inspection suggested samples having macroscopic domains of crystalized ions. These samples were excluded from further measurements. The situation was ameliorated by modifying the sample preparation procedure, where DPPC and salt were co-solubilized in water, as opposed to organic solvent. The  $\text{Zn}^{2+}$  containing bilayers mixed in water solution were deposited on a silicon wafer that was heated to  $50^{\circ}\text{C}$ , until the excess water evaporated. For both sample preparation protocols, care was taken to form lipid multilayers in the fluid phase, and to anneal the samples for several hours upon ~~the~~ rehydration.

Comparing the rocking curves of samples prepared through the deposition from solvent and water solution (Fig. 1b and 1c) confirms that the two samples are of comparable quality in the case of DPPC+ $\text{Ca}^{2+}$  (1:0.2 mol/mol), while a significant improvement was achieved in the case of ~~the~~  $\text{Zn}^{2+}$  loaded samples. All ~~the~~  $\text{Zn}^{2+}$  containing bilayers were then deposited from water solution. All samples considered in our study had rocking curves with mosaicity of less than 0.1 degree, as shown for the selected samples in figures S1-S3.

### Small-Angle Neutron Diffraction

Neutron diffraction data were collected at the Canadian Neutron Beam Centre's (CNBC) N5 beamline located at the National Research Universal (NRU) reactor (Chalk River, Ontario, Canada). 2.37 Å wavelength neutrons were selected by the (002) reflection of a pyrolytic graphite (PG) monochromator, and a PG filter was used to eliminate ~~its~~ higher order reflections. Despite reducing the intensity of the higher harmonics (i.e.,  $\lambda/2$ ,  $\lambda/3$ ,  $\lambda/4$  etc.) down to ~1%, diffraction peaks from the silicon monocrystal wafers had to be discarded on a case-by-case

basis. Si peaks were identified by the fact that they remained stationary as a function of temperature or changes in hydration, while peaks due to the membrane shifted in angle (see Fig. S7-S9 in the Supporting Information).

Incident neutrons were collimated (6x50 mm<sup>2</sup> collimator). At the sample, the neutron beam was ~8 mm wide (FWHM) with a horizontal angular divergence of 26'. A set of 7x40mm<sup>2</sup> slits — before and after the sample — were utilized to place the neutron beam at the sample center-of ~~the sample~~. Neutrons diffracted by the sample, and collimated by another rectangular collimator of 4x50 mm<sup>2</sup> opening with a horizontal angular divergence of ~17' were then collected by a multi-wire <sup>3</sup>He detector. We define the sample and detector angles as  $\Psi$  and  $\Phi$ , respectively.

~~Diffraction~~ data were collected following the standard symmetric diffraction scans, in which reciprocal space was probed along the  $z$  direction of the scattering vector  $Q$ . This was achieved by rotating the sample angle by half the detector angle ( $\Phi=2\Psi$ ), resulting in the in-plane  $Q$  component being zero. The  $Q_{in}$  component was obtained in separate measurements following the mode used for collecting rocking curve data. In this mode, the detector is positioned at the first Bragg reflection and the sample is scanned through a number of angles enabling one to interrogate reciprocal space along constant  $|Q|$  values (i.e.,  $Q_z^2 + Q_{in}^2 = Q^2$ ). The diffraction curve mode is used to obtain the one-dimensional neutron scattering length density (NSLD) profile, while rocking curves across the diffraction peak allow for the evaluation of sample quality, independent of the peak position.<sup>23</sup> Note, that measured intensities need to be corrected for incident flux, sample absorption, and the Lorentz correction according to well-established procedures for obtaining proper NSLD profiles.<sup>22,24</sup> However, this is not necessary in the case of rocking curves, where the evaluated parameter is only the peak width.

Samples were held vertically in an air-tight hydration chamber at 25°C. The chamber's bottom was filled with a saturated  $\text{K}_2\text{SO}_4$  (97% RH)<sup>25</sup> solution of different  $\text{D}_2\text{O}/\text{H}_2\text{O}$  mixtures. The well-controlled hydration conditions, which were slightly below the full hydration condition, restrict bilayer undulations important for determining bilayer material properties,<sup>26</sup> but assure that the sample is stable throughout the course of the experiment. Under these conditions, the amount of inter-bilayer hydration water is sufficient to avoid structural changes due to steric constraints.<sup>24</sup>

### Molecular Dynamics Simulations

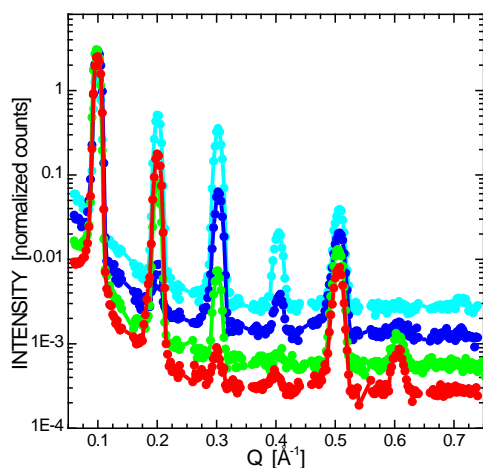
Molecular Dynamics simulations using GROMACS 5.0.4<sup>27</sup> [package](#) were performed using 128 DPPC and 3655 water molecules. 18 water molecules were replaced by cations in  $\text{Me}^{2+}$  loaded bilayers (additional 36 water molecules were replaced by  $\text{Cl}^-$  counter ions) to achieve a DPPC: $\text{Me}^{2+}$  ratio of 1:0.14 mol/mol, while 64  $\text{Me}^{2+}$  cations (and 128  $\text{Cl}^-$ ) were needed for the [ratio](#) 1:0.5 mol/mol [ratio](#). The initial topology was taken from literature.<sup>28</sup> Modeling was carried out in three stages: i) minimization of the energy; ii) NVT and NPT equilibration of the system; and iii) molecular dynamics (MD) calculations with the GROMOS96\_53A6 force field and parameters of expansion taken from Berger lipids.<sup>29</sup> Simulation convergence was assessed using [the area per lipid-parameter](#), which suggested an equilibration time of about 30ns (see Fig. S11). The production run continued for an additional 50ns.

GridMAT-MD<sup>30</sup> was used to calculate the average deviation of lipid membrane thickness ( $D_M$ ) and area per lipid ( $A_L$ ), while radial distribution functions were calculated using the GROMACS auxiliary program. SIMtoEXP<sup>19</sup> was used to verify the simulated results with those obtained from neutron diffraction experiments (see Fig. S12-S14).

## RESULTS AND DISCUSSION



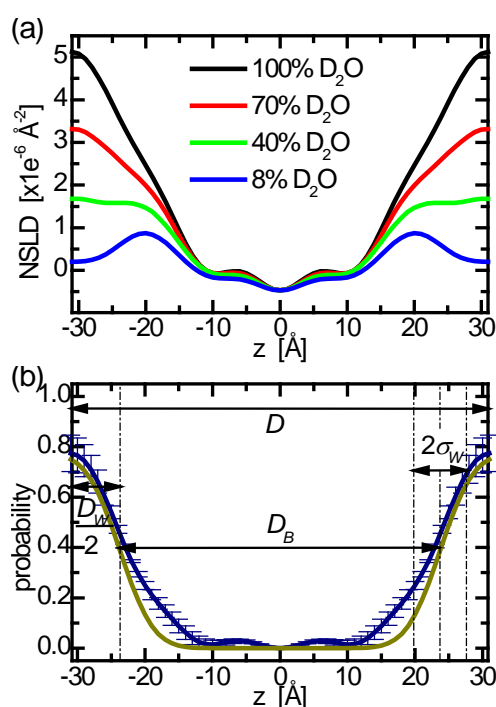
We screened multiple copies of samples using the above described rocking curve measurements. The best quality samples exhibited narrow central peaks and up to 7-9 orders of quasi-Bragg maxima (Fig. 2, Fig. S4-S6, S10). All diffraction experiments involved contrast variation measurements, in which different D<sub>2</sub>O/H<sub>2</sub>O ratio solutions were utilized to hydrate the samples. This approach allowed us to solve the scattering phase problem for our centrosymmetric system by requiring the scattering form factors to change linearly as a function of D<sub>2</sub>O content.<sup>31</sup> We also employed the procedure suggested for unambiguous phase determination as described in details elsewhere.<sup>24</sup>



**Figure 2.** Uncorrected diffraction curves (i.e., normalized intensities measured as a function of scattering vector) from aligned DPPC multibilayers containing Zn<sup>2+</sup> (DPPC:Zn<sup>2+</sup>=1:0.096 mol/mol), and measured in 100% (red), 70% (green), 40% (blue), and 8% (cyan) D<sub>2</sub>O contrast conditions. The curves were shifted vertically for clarity of viewing.

The one-dimensional NSLD bilayer profiles were obtained by Fourier transformation of the corrected contrast varied diffraction data.<sup>24</sup> While the profile corresponding to the 100% D<sub>2</sub>O

contrast condition provides the best estimate of the bilayer's steric thickness, lower contrast profiles reveal details in the region of the lipid head-group region. This is best exemplified by the 8% D<sub>2</sub>O data, where the net NSLD of water is zero, essentially yielding the lipid only NSLD profile of the bilayer (Fig. 3a). In addition, the contrast varied diffraction data can be subtracted from each other either in real space (e.g., NSLD profiles from Fig. 3a) or in reciprocal space (e.g., structure factors obtained from Fig. 2), resulting in the probability distributions of water and lipid molecules.<sup>24</sup>



**Figure 3.** (a) DPPC bilayers prepared with Zn<sup>2+</sup> (DPPC:Zn<sup>2+</sup>=1:0.096 mol/mol) shown through the one-dimensional NSLD profiles calculated from SAND measurements at different contrast conditions. (b) The water probability distribution obtained using the NSLD subtraction procedure (blue line), and by fitting the error function to the contrast varied diffraction data (green line).

The multilayer repeat distance  $D$  is the sum of the bilayer thickness  $D_B$  and interbilayer water space  $D_W$ .  $\sigma_W$ , defined by the width of error function, determines the roughness of water-bilayer interface.

Figure 3b shows water probability distributions obtained ~~from~~ by subtractions in real and reciprocal space. The former approach does not require any assumptions on the functional form of the NSLD profile, nor of the probability distribution resulting from a direct subtraction. The water distribution is obtained from averaging all of the separately subtracted pairs of contrast varied NSLD profiles, which also provides an estimate of the standard deviation error (blue line in Fig. 3b). On the other hand, the advantage of the latter approach lies in the simultaneous evaluation of all the different contrast variation diffraction data, while assuming that the water probability follows an error function, as described previously.<sup>24</sup> We adopted this approach for the data analysis used to obtain the inter-bilayer water layer thickness  $D_W$ , roughness of water-bilayer interface  $\sigma_W$ , and bilayer thickness  $D_B$  from the resulting error function ~~directly~~ (see Fig. 3b for the graphics of  $D=D_B+D_W$  definition).

The above-obtained parameters allow us to calculate the area per unit cell  $A_{UC}$  from the known DPPC molecular volume  $V_L$  (1147 Å<sup>3</sup>).<sup>32</sup> and volumes estimated for Ca<sup>2+</sup> (20.1 Å<sup>3</sup>) and Zn<sup>2+</sup> (12.5 Å<sup>3</sup>) ions.<sup>33</sup>  $A_{UC}$  is an important value of the lateral forces taking place within lipid bilayers, and also allows us to place the data on an absolute scale. The form factor  $F_0$  corresponding to the forward scattering that is not available from the scattering experiment, is calculated in a manner similar to what was done previously for X-ray scattering<sup>34</sup>

$$F_0^{abs} = (b_{UC} - \rho_W V_{UC}) / A_{UC} . \quad (1)$$

Here,  $b_{UC}$  represents the neutron scattering length<sup>35</sup> of the unit cell moieties whose volume  $V_{UC}$  displaces an equivalent volume of water molecules, with a neutron scattering length density  $\rho_w$ . All parameters in the brackets are well known, while  $A_{UC}$  is determined from the probability distributions discussed above. The calculation then consists of several iterations, converging on the final structural and normalization parameters.

It is worth noting that the calculated water probability distributions ~~of water~~ (see Fig. 3b) can assume values from 0 to less than 1. This suggests that the two component system consists of a region occupied completely by lipid ( $z=-15$  to  $15$  approximately) and a ~~region of~~ lipid/water mixture region (below  $z=-15$  and above  $z=15$ ). Interestingly, the water probability levels out before reaching a value of 1, suggesting that the system does not contain a pure water layer. It is also interesting to note a systematic difference between the reciprocal space and real space analyses, which varies by up to  $2 \text{ \AA}$ . While the first method averages over the four different contrast conditions in reciprocal space, the second method treats the diffraction data obtained at different contrast conditions individually, and then averages the results in real space. More important, however, is that the two methods show excellent agreement when it comes to relative structural changes.<sup>24</sup> Table 1 and 2 list the parameters evaluated according to the second method.

**Table 1.** Parameters calculated for DPPC multilayers with  $\text{Ca}^{2+}$  ions at  $25^\circ\text{C}$  and 97%RH.

DPPC: $\text{Ca}^{2+}$ [mol/mol]	D [ $\text{\AA}$ ]	$D_W$ [ $\text{\AA}$ ]	$\sigma_W$ [ $\text{\AA}$ ]	$D_B$ [ $\text{\AA}$ ]	$A_{UC}$ [ $\text{\AA}^2$ ]
<b>1:0</b>	60.3	13.6	4.0	46.7	49.1
<b>1:0.100</b>	63.5	15.6	3.9	47.8	48.0
<b>1:0.125</b>	63.9	15.3	4.3	48.6	47.3
<b>1:0.143</b>	63.5	12.5	3.0	51.0	45.0
<b>1:0.200</b>	63.5	13.0	3.9	50.5	45.5
<b>1:0.300</b>	64.0	15.6	4.7	48.4	47.6
<b>1:0.400</b>	63.9	16.6	4.0	47.4	48.8
<b>1:0.500</b>	64.4	17.0	4.2	47.4	48.9

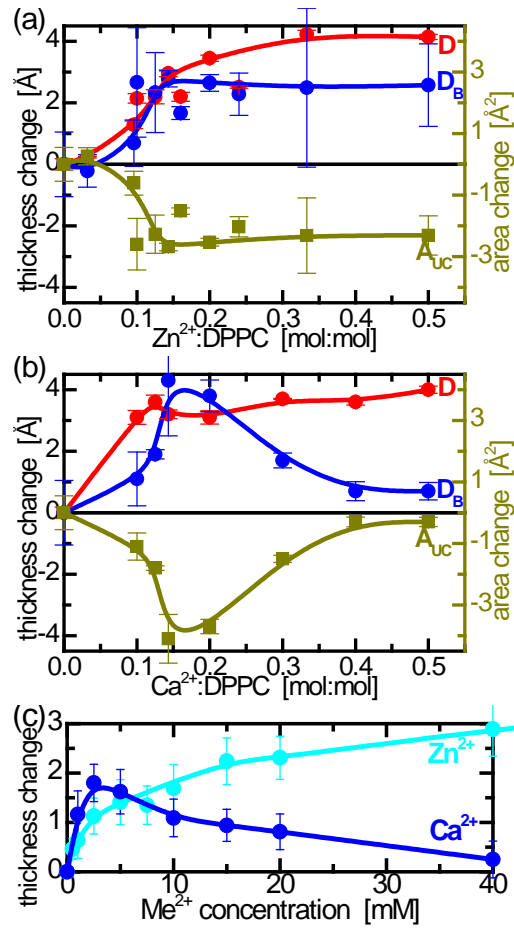
**Table 2.** Parameters calculated for DPPC multilayers with  $\text{Zn}^{2+}$  ions at 25°C and 97%RH.

DPPC: $\text{Zn}^{2+}$ [mol/mol]	$D$ [Å]	$D_W$ [Å]	$\sigma_W$ [Å]	$D_B$ [Å]	$A_{UC}$ [Å <sup>2</sup> ]
<b>1:0</b>	60.3	13.6	4.0	46.7	49.1
<b>1:0.032</b>	60.6	14.1	4.5	46.5	49.4
<b>1:0.096</b>	61.6	14.2	3.9	47.4	48.5
<b>1:0.100</b>	62.5	13.1	4.5	49.4	46.5
<b>1:0.125</b>	62.5	13.5	5.0	49.0	46.8
<b>1:0.143</b>	63.3	13.8	4.2	49.5	46.4
<b>1:0.160</b>	62.5	14.1	4.1	48.4	47.6
<b>1:0.200</b>	63.8	14.4	3.9	49.4	46.6
<b>1:0.240</b>	62.8	13.8	3.3	49.0	47.1
<b>1:0.333</b>	64.6	15.4	4.1	49.2	46.8
<b>1:0.500</b>	64.5	15.2	3.9	49.3	46.8

The ~~obtained~~ results confirm that both cations bind to DPPC bilayers. This fact is reflected in the increased lamellar repeat distance  $D$  that is most likely the result of charge induced repulsion between bilayers. This observation is consistent with previously published results, in which zwitterionic lipid multilayers swelled in the presence of different types of salts.<sup>11,12,15,36</sup> For example,  $\text{CaCl}_2$  at 1-50 millimolar concentrations induced fluid multibilayers to swell to the point that they became unbound (i.e., infinite  $D$ ), ~~and~~ — not unlike our current observations — and  $D$  reached ~~had~~ an upper limit in the gel phase.<sup>15,36</sup> Inoko et al.<sup>37</sup> explained that the initial swelling was due to an increase in inter-bilayer electrostatic repulsion that was screened at high salt concentrations. Supposedly membrane structure was not affected at high salt concentrations.<sup>12</sup> Interestingly, our low salt concentration results reveal that both  $D_B$  and  $D_W$  contribute to the increases in  $D$ . Moreover,  $\text{Ca}^{2+}$  and  $\text{Zn}^{2+}$  do not have the same effect on  $D$ .

Fig. 4a shows a monotonic increase of  $D$  and  $D_B$  up to a  $\text{Zn}^{2+}$ :DPPC ratio of 0.14 mol/mol, after which they both seem to plateau, with ~~the bilayer thickness~~  $D_B$  experiencing almost twice

the change of  $D_W$  (see Table 2). Calculated area per unit cell that comprises of a DPPC molecule and the appropriate amount of cations follows the thickness changes, decreasing with an increasing  $\text{Zn}^{2+}$ :DPPC ratio. This is due to the fact that lateral area and transverse thickness are inversely related (i.e.,  $A_{UC}=V_{UC}/D_B$ ).



**Figure 4.** (a)  $\text{Zn}^{2+}$  and (b)  $\text{Ca}^{2+}$  cation-induced changes to bilayer structure expressed through the lamellar repeat distance  $D$  (red), bilayer thickness  $D_B$  (blue), and area per unit cell  $A_{UC}$  (dark yellow). (c) Previously published results of cation-induced changes to the bilayer structure of unilamellar vesicles dispersed in water solution.<sup>10,16</sup> Changes are calculated with respect to neat

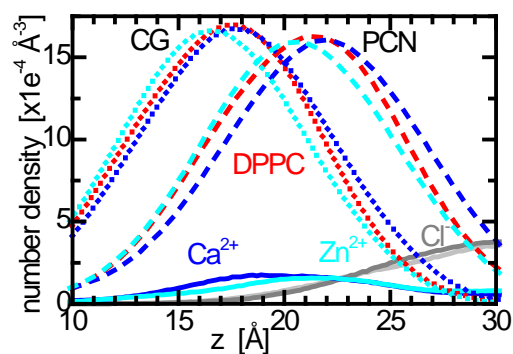
DPPC bilayers, and uncertainties are estimated from the partial derivatives of the chi-square function.

Bilayers with  $\text{Ca}^{2+}$  behave differently from those with  $\text{Zn}^{2+}$ . While they both experience similar changes for ion:DPPC ratios up to 0.14 mol/mol, they differ at higher ratios. The most pronounced difference is exhibited by  $D_B$ , which experiences a decrease at high concentrations of  $\text{Ca}^{2+}$ , displaying-with a maximum at DPPC: $\text{Ca}^{2+}$ =1:0.14 mol/mol.  $D_B$  decreases almost to the level of pure DPPC bilayers at high  $\text{Ca}^{2+}$  DPPC ~~the high~~ ratios, while the changes are compensated by the changes in  $D_W$ , resulting in a  $D$  that is practically unaffected by changes in  $\text{Ca}^{2+}$  concentration. We can therefore disregard the effect that adjacent bilayers have on each other in our stacked system. This is supported by our previously published results<sup>10,16</sup> (reproduced in Fig. 4c), whereby unilamellar vesicles dispersed in the water solution loaded with either  $\text{Zn}^{2+}$  or  $\text{Ca}^{2+}$  ions displayed trends similar to those in our present data. The data also imply that a small curvature present in the unilamellar vesicles is not an important factor affecting the bilayer structure of ion containing DPPC membranes.

It is interesting to note that recently published studies indicate a behavior of cation-containing lipid bilayers to be different at low and high hydration conditions. Alsop et al. reported that lipid headgroups hydrated by less than 7 water molecules (50% RH) are not affected by the addition of salt cations, and any changes are manifested by an increase of inter-bilayer water layer thickness.<sup>11</sup> Changes to lipid bilayer structure apparently require the presence of a sufficient number of water molecules. The critical number of water molecules for neat POPC bilayers has been suggested to be 3, while in the case of POPC +  $\text{Ca}^{2+}$  and POPC +  $\text{Zn}^{2+}$  complexes the critical number is  $\sim 7$ .<sup>1</sup> Our observation of changes to  $D_B$  and  $D_W$  at 97% RH support the high

hydration scenario. These changes — in the form of headgroup-salt complexes — in turn induce changes to the ordering of the lipid hydrocarbon chains, as noted previously.<sup>15</sup>

We used MD simulations to determine the atomistic structural details of the most interesting concentration of cations, namely the 1:0.14 molar ratio of DPPC:Me<sup>2+</sup>. Simulation results were first validated by the experimentally obtained diffraction data using the simulation-to-experiment (SIMtoEXP) approach.<sup>19</sup> The SIMtoEXP procedure calculates scattering form factors from simulations without model assumptions using a Fourier Transform and comparing them directly to experiment. Doing this reveals local deviations. Differences between experiment and simulations are not unexpected due to the less than perfect force fields used,<sup>20</sup> and the small differences in hydration levels. Simulations give rise to smoother profiles, while those from experiment contain more features (compare e.g., Fig. 3a and S12-S14). Simulations, however, are capable of deconstructing the bilayer into smaller structural subunits that allow the determination of important structural information. Figure 5 focuses on the head-group's phosphate and carbonyl groups, which are anticipated to be the primary targets for cations binding due to their charges.



**Figure 5.** Simulated number density distributions along the bilayer normal (bilayer is centered at  $z=0$  Å) focusing on the carbonyl-glycerol (CG - dotted lines) and phosphate-choline groups

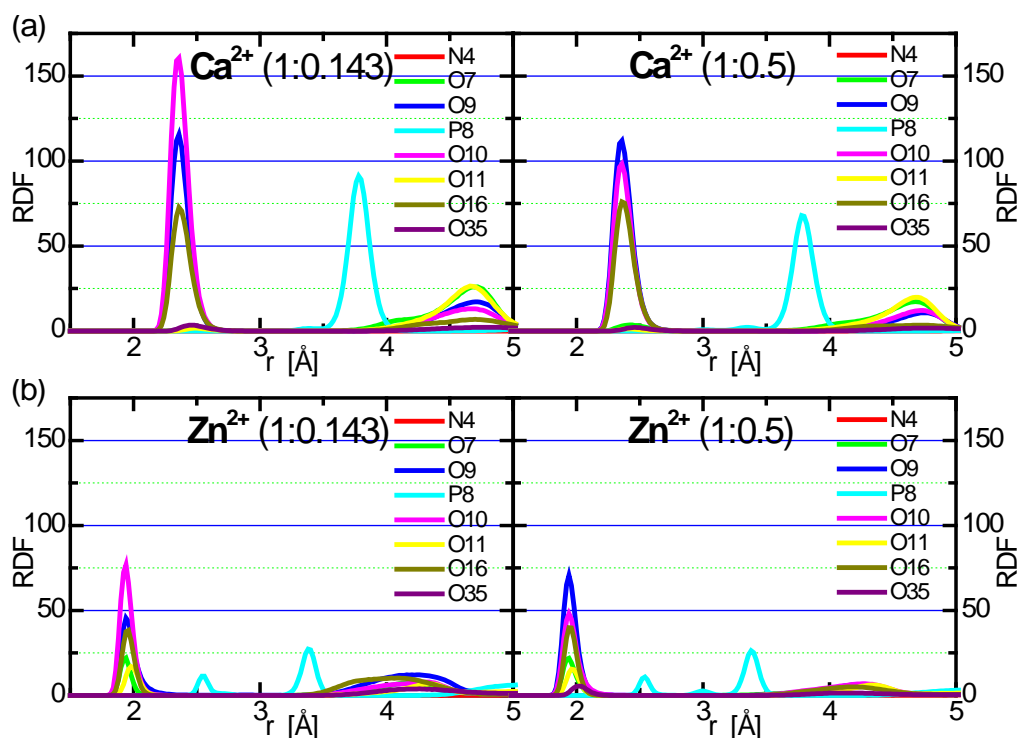


(PCN - dashed lines) for a pure DPPC bilayer (red), and bilayers containing  $\text{Ca}^{2+}$  (blue) or  $\text{Zn}^{2+}$  (cyan) cations. Note the nearly identical distributions of  $\text{Cl}^-$  counter ions (gray), mostly outside the bilayer, for both salts.

The addition of  $\text{CaCl}_2$  or  $\text{ZnCl}_2$  salt affects gel phase DPPC bilayers, as is evident by the positional changes within their headgroup regions. The effect of  $\text{Cl}^-$  anions ~~is can be negligible, eted,~~ based on the fact that ~~their influence is they reside~~ mostly outside the bilayer. This is in agreement with the observation of weak interactions between various anions and the lipid liquid-condensed phase.<sup>38</sup> The ~~distinct~~ effects ~~as a result of of~~ cations ~~are~~, however, expected based on their ~~different~~ locations. While  $\text{Ca}^{2+}$  causes both the carbonyl and phosphate groups to shift slightly outwards, the shifting effect of  $\text{Zn}^{2+}$  is towards the bilayer center. Moreover, the cations themselves seem to be localized in opposite directions.  $\text{Ca}^{2+}$  penetrates deeper into the lipid bilayer, positioning itself about  $1.5\text{\AA}$  below the maximum of the phosphate distribution, while  $\text{Zn}^{2+}$  is positioned about  $1.5\text{\AA}$  above the phosphate maximum. Nevertheless, in both cases the distribution functions in Fig. 5 suggest the primary binding site for both cations to be the phosphate group, in contrast to what was observed in dehydrated bilayers.<sup>11</sup>

We further analyzed our simulation data in order to characterize the extent of the specific binding of the two cations. As confirmed by the radial distribution functions (RDF) calculated from our simulations, both cations interact strongly with the negatively charged oxygens. Previous simulation results suggested carbonyl oxygens to be their main target,<sup>39</sup> although other studies questioned the reliability of the force-fields utilized in the different simulations.<sup>13</sup> Our data (left hand graphs of Figure 6) clearly show that both  $\text{Ca}^{2+}$  and  $\text{Zn}^{2+}$  (at  $\text{DPPC}:\text{Me}^{2+}=1:0.143$  mol/mol) interact with carbonyl oxygens only a third as much as they interact with the phosphate

oxygens (O16 and O35 curves vs. O9 and O10 in Fig. 6 – see Fig. S15 for ~~the the schematics of~~ naming convention used). The graphs also show that  $\text{Ca}^{2+}$  forms a contact pair with phosphate oxygens (blue and magenta curves in Fig. 6a) that is about two times stronger than that formed by  $\text{Zn}^{2+}$  (blue and magenta curves in Fig. 6b). A similar ratio is also observed in the case of the RDF calculated for carbonyl oxygen (dark yellow curves in Fig. 6). In addition to confirming the preferential interactions of both cations with phosphate, the data show that  $\text{Ca}^{2+}$  binds with the phosphate group much more readily than  $\text{Zn}^{2+}$  at ~~a the~~ DPPC: $\text{Me}^{2+}$  molar ratio of 1:0.143. In fact, without exception  $\text{Ca}^{2+}$  binds to any head-group atom much more favorably than  $\text{Zn}^{2+}$ .



**Figure 6:** Radial distribution functions (RDF) for (a)  $\text{Ca}^{2+}$ - and (b)  $\text{Zn}^{2+}$ -lipid headgroup atom pairs from simulations performed at ion:DPPC ratios of 0.143 (left-hand) and 0.5 (right-hand).

The RDF is shown for different headgroup atoms, ~~starting from -in the order from~~ the bilayer exterior, inwards.

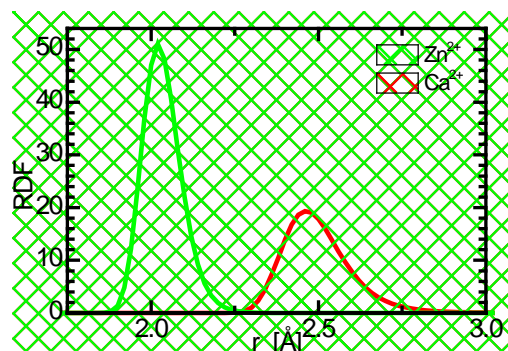
~~Interestingly, the~~ The situation changed ~~intriguingly~~ when an  $\text{Me}^{2+}$ :DPPC ratio of 0.5 mol/mol was used in our simulation – the ratio at which experimental results revealed differences in the effects between the two cations. Firstly, the specific binding with the phosphate O10 oxygen decreases markedly for both cations (see the right-hand graphs of Figure 6). However, this decrease is, ~~for the most part,~~ offset by an increased binding of  $\text{Zn}^{2+}$  with O9 oxygen. In the case of  $\text{Ca}^{2+}$  there is an overall decrease of phosphate group specific binding for  $\text{Ca}^{2+}$  by about 20%, despite a 3.5-fold increase in ~~ion~~ concentration. This decrease in specific binding then explains the decreased effect that  $\text{Ca}^{2+}$  cations have on ~~the experimentally observed~~ bilayer thickness and area per lipid ~~observed experimentally~~. The strong, initial effect on the lipid headgroup by  $\text{Ca}^{2+}$ , which also affects the hydrocarbon chain region, is likely screened at higher cation concentrations.<sup>15</sup> Moreover, the constant level of  $\text{Zn}^{2+}$  specific binding agrees well with what we saw experimentally regarding DPPC bilayer structure.

The differences between the structural changes experienced by DPPC bilayers when  $\text{Ca}^{2+}$  or  $\text{Zn}^{2+}$  are introduced may also be rationalized by the physical differences between the two cations.<sup>1</sup> Calcium ions have a larger ionic, or Pauling radius (1 Å), than zinc ions (0.75 Å). The preferred coordination number for bulk hydrated calcium has been reported to be between 6 and 8,<sup>1</sup> while for zinc ~~ions~~ this number is between 4 and 6.<sup>7</sup> In addition, the hydrogen bonding of water molecules extends beyond the ion's<sup>2</sup> primary hydration shell — their crystal arrangements in bulk water have been proposed to be  $\text{Zn}[\text{H}_2\text{O}]_6^{2+} \bullet [\text{H}_2\text{O}]_{12}^{40}$  and  $\text{Ca}[\text{H}_2\text{O}]_{6,1}^{2+} \bullet [\text{H}_2\text{O}]_{5,29}^{41}$ .

Calcium ions as a consequence of their larger size, require smaller amounts of energy for the removal of hydration shell water.<sup>42</sup>

Up to this point we have discussed experimental and simulation data. It is in order to have a closer look at interactions of ion-containing lipids and ions with water—one of the tiniest and often forgotten biomolecules.<sup>43</sup> The most of the data discussed above have been obtained utilizing various experimental methods. Theoretical approaches, however, on the other hand, can provide additional information regarding in both the structure and dynamics of ion hydration. The hydration energies that are routinely calculated regularly are unfortunately often poor or require the input of unrealistic ion characteristics,<sup>41</sup> — although there are recent MD simulations that are in good provide quantitative agreement with experimental data.<sup>39</sup> The difficulties lie in the fact that water molecules surrounding the ions are strongly polarized, something not well addressed by—that is in turn a weak spot of—MD force fields.<sup>44</sup> In addition, the water binding capacity of ions is nevertheless distinctly reduced in the case of ion-lipid complex formation, suggesting that both ion and lipid sacrifice-give up to 50% of their hydration shell.<sup>6</sup> Ions have been found to disrupt water structure, reducing its ability to solvate lipid head-groups, and whereas a monotonic dependence on the size of monovalent cations has ave-been reported.<sup>45</sup>

According to our MD simulation results, divalent  $\text{Ca}^{2+}$  ions create about 1.6 times fewer pairs with surrounding water molecules than do  $\text{Zn}^{2+}$  ions (Fig. 7) when interacting with DPPC bilayers. This smaller lesser- $\text{Ca}^{2+}$  hydration shell most likely allows for its-more proximal closer and stronger contacts with the lipid phosphate.



**Figure 7:** Radial distribution functions (RDF) for  $\text{Zn}^{2+}$ - and  $\text{Ca}^{2+}$ -water pairs determined from MD simulations. The peak position provides a mean radius ~~for the of~~ primary hydration shell (larger for  $\text{Ca}^{2+}$  when compared to  $\text{Zn}^{2+}$ ). The area under the peaks corresponds to the number of water molecules present in the hydration shell of each cation interacting with ~~the~~ lipid bilayer (~~this value is~~ smaller for  $\text{Ca}^{2+}$  ~~- when~~ compared to  $\text{Zn}^{2+}$ ).

Our simulation data corroborate that the electronic structure of calcium results ~~for in~~ its high ~~affinity for~~  $\text{PO}_4$ ; ~~affinity in particular~~ (~~also partly~~ for  $\text{CO}_2$  also).<sup>12</sup> The electronic structure of zinc, on the other hand, differs from that of other divalent metal ions. Calculations have suggested significant variability in the modes of binding of hydrated  $\text{Zn}^{2+}$ , while their binding energies vary only by small amounts.<sup>46</sup> This can then explain the low specificity of zinc binding to lipid bilayers and the non-specific changes to their structure. In our case, this characteristic may have manifested itself in the monotonic increase of ~~the experimentally determined~~ structural parameters, which are unaffected at higher concentration of the cation. ~~The other divalent metals may, however behave differently. For example,  $\text{Mg}^{2+}$  and  $\text{Fe}^{2+}$  are for example known to have ionic size and hydration shells similar to those that of  $\text{Zn}^{2+}$ ,<sup>40,47</sup> yet they have been reported to cause even higher repulsion between the lipid bilayers, compared to the than in the case of larger~~

Ca<sup>2+</sup>.<sup>11</sup> It is clear that such findings may propose the high specificities of those cations that can likely prevail other structural properties in their overall effect. After all, how to better explain special physiological roles these cations play in the functionality of membrane? The further studies examining the lipid-ion interactions in the case of Mg<sup>2+</sup>, Fe<sup>2+</sup> are needed to clearly in order to answer this perplexing intriguing question.

## CONCLUSIONS

Using SAND and MD simulations we observed differences in how Ca<sup>2+</sup> and Zn<sup>2+</sup> affect DPPC bilayers – i.e., changes in lipid bilayer thickness. For both cations, bilayer thickness increased, reaching a maximum at ~1:0.14 mol/mol DPPC:Me<sup>2+</sup>. However, while further increases in Ca<sup>2+</sup> concentration resulted in the bilayer thinning to a level not dissimilar from that of pure DPPC bilayers, Zn<sup>2+</sup> binding exhibited the behavior of a typical binding isotherm, reaching a level where further addition of the cation resulted in no change to the bilayer. The radial distribution functions of cation-lipid pairs calculated from MD simulations for several lipid head-group atoms revealed Ca<sup>2+</sup> binding strongly to the phosphate group, while Zn<sup>2+</sup> had a much lower binding affinity for all head-group atoms. Our results also confirmed that the observed structural changes could be rationalized in terms of electrostatic interactions, rather than of geometrical constraints due to bilayer curvature, thus reinforcing the chemical importance of the two cations. In addition, our observations imply that lipid-ion interactions do not only depend on which cation is present, but also on the lipid's thermodynamic state and its level of hydration. It will be interesting to expand these studies by examining the effect of other cations, and/or the role of lipid structural properties. An understanding of lipid bilayer structural changes as a result of different salt environments provides a foundation for better insights into the structure-function relationships that most certainly take place in more complicated biomembrane systems.

## ASSOCIATED CONTENT

~~Complete~~ Additional experimental and simulation data ~~neutron rocking curves and diffraction data, and graphics of a direct validation of MD simulations by experimental data~~ are included in the as-Supporting Information. This material is available free of charge via the Internet at <http://pubs.acs.org>.

## ACKNOWLEDGMENT

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