## Anomalous Stress Response of Ultrahard $WB_n$ Compounds

Quan Li, <sup>1, 2, \*</sup> Dan Zhou, <sup>3, 2</sup> Weitao Zheng, <sup>1</sup> Yanming Ma, <sup>1, †</sup> and Changfeng Chen<sup>2, ‡</sup>

<sup>1</sup>College of Materials Science and Engineering and State Key Laboratory
of Superhard Materials, Jilin University, Changchun 130012, China

<sup>2</sup>Department of Physics and High Pressure Science and Engineering Center,
University of Nevada, Las Vegas, Nevada 89154, USA

<sup>3</sup>Laboratory of Clean Energy Technology, Changchun University of Science and Technology, Changchun 130022, China
(Dated: October 6, 2015)

Boron-rich tungsten borides are premier prototypes of a new class of ultrahard compounds. Here we show by first-principles calculations that their stress-strain relations display surprisingly diverse and anomalous behavior under a variety of loading conditions. Most remarkable is the dramatically changing bonding configurations and deformation modes with rising boron concentration in  $WB_n$  (n=2, 3, 4), resulting in significantly different stress responses and unexpected indentation strength variations. This novel phenomenon stems from the peculiar structural arrangements in tungsten borides driven by boron's ability to form unusually versatile bonding states. Our results elucidate the intriguing deformation mechanisms that define a distinct type of ultrahard materials. These new insights underscore the need to explore unconventional structure-property relations in a broad range of transition-metal light-element compounds.

PACS numbers: 62.20.-x,81.40.Jj,61.50.Ah,71.20.Be

Strong covalent solids, such as diamond and cubic boron nitride, contain rigid bonding networks that can sustain large stresses under diverse loading conditions. However, low thermal stability of diamond and high synthesis cost of traditional superhard materials have prompted intensive search for alternative materials that exhibit excellent stability and strength but are easier to synthesize. Recently a new design strategy introduces covalent-bond-forming light elements (LEs) to combine with heavy 4d or 5d transition metals (TMs) [1–11]. The key idea here is that such TM-LE compounds may benefit from the excellent properties of the constituent components where the heavy TM atoms introduce a high density of valence electrons into the compounds to resist compression while the LE atoms help form a strong covalent bonding network to enhance structural integrity. A high concentration of LE atoms is considered essential to achieving optimal high strength and hardness [7–13].

Outstanding among TM-LE compounds is a series of binary tungsten borides with the LE-to-TM ratio up to 4:1 obtained under relatively easy synthesis conditions [7–9]. These tungsten borides exhibit excellent mechanical properties, making them premier prototypes among this new class of ultrahard materials. Understanding these materials requires a full resolution of their crystal structures complicated by boron's ability to form versatile bonding states [14–17]. We recently identified the structures of  $WB_n$  compounds over a wide range of boron content using a global structural search method CALYPSO [18, 19] that has proven effective in solving complex crystal structures [20-23]. The detailed knowledge of the structural arrangements in tungsten borides [24] makes it possible to explore the structure-strength relations and unveil the underlying atomistic deformation mechanisms in these promising materials.

In this work, we determine the stress-strain relations from first-principles calculations for boron-rich tungsten borides WB<sub>n</sub> (n=2, 3, 4) under compressive, tensile, pure shear, and Vickers indentation shear strains. Of particular interest is their variation with changing boron content. A positive correlation between strength and boron content has been widely expected, driving extensive recent efforts on the synthesis and characterization of high-boron-content WB<sub>n</sub> ( $n \ge 3$ ). Surprisingly, our results show no such correlation; instead, we find that the  $WB_n$  compounds exhibit versatile stress response at different boron content, and WB<sub>2</sub> shows the best overall strength performance under various loading conditions. This anomalous phenomenon stems from the peculiar bonding arrangements in the  $WB_n$  compounds with the changing boron concentration n, which introduces different deformation mechanisms that produce distinct mechanical properties in these boron-rich tungsten borides. Our findings challenge the conventional wisdom and offer key insights into the fundamental structural deformation modes and strength of tungsten borides. These results suggest that unconventional structure-property relations may also exist in a broad range of TM-LE compounds such as other TM borides and TM nitrides and carbides that possess complex bonding configurations.

Tungsten borides adopt a rich variety of structural forms. Our global structure search [24] over a wide concentration range identified five stable tungsten borides; but the two compounds with lower boron concentrations,  $W_2B$  and WB, exhibit much lower strength due to a lack of direct boron-boron bonding network, which sets them apart from the three boron-rich  $WB_n$  (n=2, 3, 4), and are not regarded as ultrahard and thus not considered in the present work. The experimentally synthesized  $WB_2$  and  $WB_3$  have been assigned the structures with

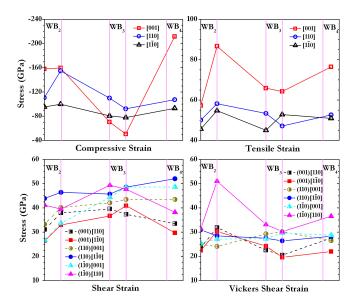


FIG. 1: (Color online) Calculated peak stress values under compressive, tensile, pure shear, and Vickers shear strains for (left to right)  $P6_3/mmc$ -4u WB<sub>2</sub>,  $P6_3/mmc$ -2u WB<sub>2</sub>,  $P6_3/mmc$ -4u WB<sub>3</sub>, R-3m-6u WB<sub>3</sub>, and  $P6_3/mmc$ -2u WB<sub>4</sub>.

 $P6_3/mmc$ -4u and  $P6_3/mmc$ -4u symmetry, respectively; interestingly, both are metastable according to the calculated convex hull [24]. Calculations also identify thermodynamically stable  $P6_3/mmc$ -2u WB<sub>2</sub>, R-3m-6u WB<sub>3</sub>, and  $P6_3/mmc$ -2u WB<sub>4</sub>. We examine these five tungsten boride structures and determine their stress response under various loading conditions. The objective here is to establish fundamental structure-property relations, especially the influence of boron content on mechanical strength and the underlying deformation mechanisms. The results may provide key benchmarks for understanding this new class of ultrahard materials. This is especially important considering boron-rich tungsten boride specimens usually contain multiple phases embedded in an amorphous boron matrix [8], which has impeded efforts to extract clear trends of mechanical properties of each individual phase.

 $6u~{\rm WB_3}$ , and  $P6_3/mmc\text{-}2u~{\rm WB_4}$  structures, respectively. The energy convergence is about 1 meV per atom, with residual stresses and forces in the fully relaxed structures less than 0.1 GPa and 0.001 eV/Å, respectively. The stress-strain relation is obtained by calculating stress response to structural deformation along specific loading paths using a quasistatic relaxation method that simulate various loading conditions and determine the corresponding ideal strength and deformation modes [31–34].

We have determined stress responses along different deformation paths under various types of strains, and the lowest peak stress in each case defines the corresponding ideal strength, which is the minimum stress needed to plastically deform a perfect crystal under that type of strain. These results establish key benchmarks for the intrinsic mechanical property of a material [35– 40]. Meanwhile, the stress responses along the paths with higher stress peaks provide additional insights into material behavior under different deformation modes (e.g., along another tensile direction or shear slip plane), and such information is crucial to understanding and optimizing material performance. In particular, to determine the indentation strengths we have carried out stress-strain calculations under indentation shear strains in the lowindex crystalline planes where crystal cleavage commonly occur, and the obtained results can be compared to actual indentation measurements that are performed on such cleavage planes.

We have computed full stress-strain curves and determined the peak stress and strain values under various strain conditions [41]. We present in Fig. 1 the calculated peak stresses that offer a comprehensive quantitative description of the value, range and trend of the stress response of  $WB_n$  with the changing boron content under each type of strain. Most conspicuous and surprising among these results is a total lack of correlation between the peak stress values and the boron content, which is in stark contrast to the common expectation that the strength of tungsten borides would increase with rising boron content [7–13]. In fact, the stress peaks largely decrease going from WB<sub>2</sub> to WB<sub>3</sub> and then recover, only partially in most cases, for WB<sub>4</sub>, making the relatively low-boron-content WB<sub>2</sub> the strongest among them with the best overall strength performance.

The distinct stress responses of  $WB_n$  are sensitive to the type of strains applied. Under compressive strains, the stress peaks drop significantly from  $WB_2$  to  $WB_3$ and then largely recover for  $WB_4$  with the highest peak achieved along its [001] direction. Under tensile strains, the peak values initially trend downward or remain essentially unchanged as the boron content rises, and then partially recover. Under pure shear strains, all the stress peaks trend upward from  $WB_2$  to  $WB_3$ , but then they move in diverging directions for  $WB_4$ , with the lowest peak (i.e., the ideal shear strength) reduced essentially to the same value as that for  $WB_2$ . Of particular interest are

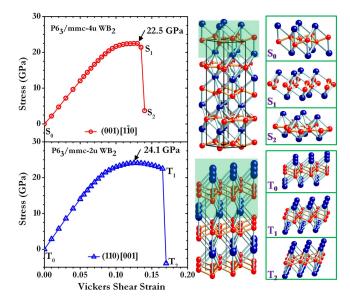


FIG. 2: (Color online) The calculated stress-strain relations and key structural snapshots for the two WB<sub>2</sub> phases in the weakest Vickers shear directions. The small (red) and large (blue) balls represent B and W atoms, respectively, which is also the case in the following figures.

the stress responses under indentation shear strains since indentation strength has been the primary focus of recent extensive studies of tungsten borides and other LE-TM compounds. Results in Fig. 1(d) show that the indentation stress peaks trend lower from WB<sub>2</sub> to WB<sub>3</sub> and then rise only slightly for WB<sub>4</sub>, but still remain at or below the values for WB<sub>2</sub>. Below we analyze the atomistic deformation modes in these WB<sub>n</sub> compounds with the lowest stress peaks (i.e., the ideal strength) under indentation strains, and the results also shed light on the anomalous stress variations under other strain conditions.

We first examine the deformation modes in  $P6_3/mmc$ - $4u \text{ WB}_2 \text{ under the indentation strain in the } (001)[1\overline{1}0]$ shear direction where the lowest stress peak (i.e., the ideal indentation strength) is obtained (Fig. 2). Structural snapshots show that the slightly buckled boron honeycomb sheet (see the highlighted region in Fig. 2) is the main load-bearing component. Near the peak stress at the indentation shear strain  $\epsilon$ =0.135, several boron bonds break, turning the boron sheet into zigzag chains. The normal compressive pressure beneath the indentor in the [001] direction squeezes these zigzag chains, resulting in B-B rebonding that leads to the formation of a graphitelike flat honeycomb boron layer, accompanied by a stress release. The situation in  $P6_3/mmc$ -2u WB<sub>2</sub> is similar, but it involves a different mechanism; instead of a bond breaking and rebonding process, the indentation strains cause a stress concentration that drives a collective bond flipping in the buckled boron sheet (see Fig. 2) to release the stress past the peak value. In general, the boron sublattices in these WB<sub>2</sub> structures form

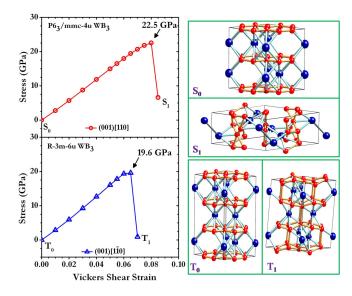


FIG. 3: (Color online)The calculated stress-strain relations and key structural snapshots for the two WB<sub>3</sub> phase in the weakest Vickers shear directions.

a strongly bonded framework that with stand large deformation strains, while the tungsten atoms provide a high electronic density to resist uniform compression. This is the essence of new design principle for ultrahard TM-LE compounds [1–11]. The strong covalent boron network in WB<sub>2</sub> can produce very high indentation strength in certain crystal orientations; for example, the peak stress under the  $(1\bar{1}0)[110]$  indentation shear strains reaches 51.0 GPa for  $P6_3/mmc$ -2u WB<sub>2</sub>, which is close to the value for cubic boron nitride. This offers a promising avenue for enhancing the indentation strength through optimal design of crystal orientation relative to the direction of external loading.

We now turn to the two  $WB_3$  phases (Fig. 3), both of which contain large voids in their structures [24]. This structural feature makes them less resistent to compression, thus explaining the large reduction in their peak stresses under compressive strains. This mechanism also underlies the reduction in the indentation strength in WB<sub>3</sub> phases since indentation produces a large compressive strain beneath the indenter, which crushes the boron layers in  $P6_3/mmc$ -4u WB<sub>3</sub> and significantly reduces the distance between the boron layers in R-3m-6u WB<sub>3</sub>. Moreover, the resulting large lateral volume expansion perpendicular to the [001] direction induced by the compressive pressure breaks up the honeycomb sheet at a low indentation shear strain ( $\epsilon < 0.085$ ). On the other hand, the strong boron networks in both  $P6_3/mmc$ -4u WB<sub>3</sub> and R-3m-6u WB<sub>3</sub> produce large pure shear strength (36.7 and 37.3 GPa), which suggests that they can resist large deformation in scratch hardness tests where no substantial normal compression is present, making them suitable for applications in hard wear-resistant coatings.

Most intriguing among the tungsten borides is WB<sub>4</sub>.

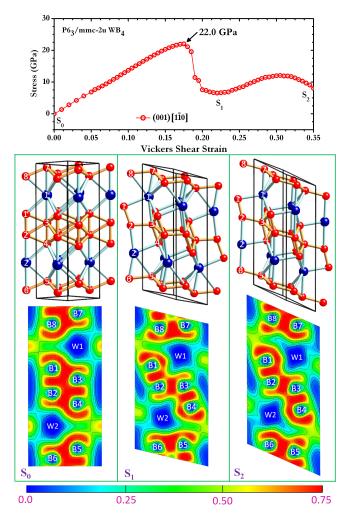


FIG. 4: (Color online) The calculated stress-strain relations and key structural snapshots for the WB<sub>4</sub> phase in the weakest Vickers shear directions. Also shown are the two-dimensional ELF plot in the  $(1\overline{1}0)$  crystalline plane.

It possesses the highest boron content in a dense bonding network, raising the expectation on superior strength and hardness. Surprisingly, our calculations (Fig. 1d) reveal that the lowest stress peak (i.e., ideal strength) under indentation strains for WB<sub>4</sub> is even slightly lower than the corresponding value for WB<sub>2</sub>; moreover, the other stress peaks of WB<sub>4</sub> are not much higher either. Equally striking is the observation that the stress-strain curve (Fig. 4) exhibits a creep-like behavior typically seen in metals but almost never in strong covalent solids. These results point to highly unusual bonding characters in WB<sub>4</sub>. Structural snapshots at several key strains (Fig. 4) show that the boron bilayer in the middle of the cell is the main load-bearing component where significant B-B bond stretch and bond breaking occur under large strains. To examine the nature of these bonding states, we plot the electron localization function (ELF) isosurfaces for  $P6_3/mmc$ -2u WB<sub>4</sub>. Three-center bonding is clearly visible in the buckled boron bilayers at equi-

librium (Fig. 4), which is made possible by the similar bond lengths of triangular  $B_1$ - $B_2$  (1.725 Å),  $B_1$ - $B_3$  (1.851 Å), and  $B_2$ - $B_3$  (1.982 Å) bonds in  $P6_3/mmc$ -2u WB<sub>4</sub>. As the indentation shear deformation in the  $(001)[1\overline{1}0]$ direction increases, the B<sub>1</sub>-B<sub>3</sub> and B<sub>2</sub>-B<sub>4</sub> bonds are continuously stretched, making a smooth transition from the three- to two-center covalent bonding states accompanied by a continuous charge transfer between the threeand two-center bonds. This peculiar transition produces a creep-like stress response with much reduced energy cost, resulting in the low peak stress, without the usual hard bond breaking seen in strong covalent solids such as diamond, cubic BN and BC<sub>2</sub>N [32-34, 37, 39] or even lower-boron-content tungsten borides, whose bonding environments do not support the three-center bonds. The multi-center bonding transition in WB<sub>4</sub> considerably constraints its indentation strength, and the same mechanism underlies its low ideal strength under pure shear strains (Fig. 1c). Meanwhile, the smooth transitions between different bonding states produce a simultaneously high-strength and high-ductility state in WB<sub>4</sub>, making it suitable for applications where both qualities are desired.

The anomalous stress responses of  $WB_n$  have important implications for understanding such a different type of ultrahard materials and guiding their optimization in synthesis and application. In particular, our results suggest that recent efforts almost exclusively focused on high-boron-content  $WB_3$  and  $WB_4$  may be misguided, and that emphasis should shift to the  $WB_2$  phases that, according to our calculations, exhibit the best overall strength performance under various loading conditions.

In summary, our first-principles calculations reveal that, contrary to common expectations, increasing boron content does not lead to higher mechanical strength in boron-rich tungsten borides; instead, their ideal strength values show little improvement or even trend downward with rising boron content in most cases. This intriguing behavior is fundamentally rooted in boron's ability to form versatile bonding states, which produces distinct structural configurations, including the alternating boron and tungsten single-layer arrangement in WB<sub>2</sub>, the large voids in WB<sub>3</sub>, and the three-center bonds in the boron bilayers in WB<sub>4</sub>. While WB<sub>3</sub> exhibits reduced strength due to the conventional structural weakness associated with the large voids, the strength of WB<sub>4</sub> is limited by a more subtle quantum effect, which introduces three-center boron bonds that can transform to two-center bonds via a continuous charge transfer under shear strains. These constraints make WB<sub>2</sub> the best overall choice for high-strength performance; meanwhile, WB<sub>4</sub> shows remarkable simultaneous high strength and enhanced ductility under shear strains. This work represents a comprehensive study of boron-rich tungsten borides based on a global structure search and the results have important implications for a large class of transition-metal borides that share the same boron concentration range and likely similar bonding configurations and deformation mechanisms as revealed in the present study. The diverse atomistic mechanisms underlying the striking stress-strain relations in  $WB_n$  introduce new knowledge about structural deformation in strong solids, which calls for further studies of a broad range of additional TM-LE compounds to explore their structure-property relations that may also exhibit rich and unexpected phenomena.

This work was supported in part by DOE under Cooperative Agreement DE-NA0001982 at UNLV and NSFC (Nos. 11474125, 51202084, and 11274136) at JLU.

- \* Electronic address: liquan777@jlu.edu.cn
- † Electronic address: mym@jlu.edu.cn
- ‡ Electronic address: chen@physics.unlv.edu
- J. C. Crowhurst, A. F. Goncharov, B. Sadigh, C. L. Evans, P. G. Morrall, J. L. Ferreira, and A. J. Nelson, Science 311, 1275 (2006).
- [2] E. Gregoryanz, C. Sanloup, M. Somayazulu, J. Badro, G. Fiquet, H. Mao, and R. J. Hemley, Nat. Mater. 3, 294 (2004).
- [3] A. F. Young, C. Sanloup, E. Gregoryanz, S. Scandolo, R. J. Hemley, and H. Mao, Phys. Rev. Lett. 96, 155501 (2006).
- [4] S. H. Tolbert, R. W. Cumberland, S. M. Clark, J. J. Gilman, M. B. Weinberger, and R. B. Kaner, J. Am. Chem. Soc. 127, 7264 (2005).
- [5] H. Y. Chung, M. B. Weinberger, J. B. Levine, A. Kavner, J. M. Yang, S. H. Tolbert, and R. B. Kaner, Science 316, 436 (2007).
- [6] J. Q. Qin, D. W. He, J. H. Wang, L. M. Fang, L. Lei, Y. J. Li, J. Hu, Z. L. Kou, and Y. Bi, Adv. Mater. 20, 4780 (2008).
- [7] R. Mohammadi, A. T. Lech, M. Xie, B. E. Weaver, M. T. Yeung, S. H. Tolbert, and R. B. Kaner, Proc. Natl. Acad. Sci. U.S.A. 108, 10958 (2011).
- [8] M. Xie, R. Mohammadi, Z. Mao, M. M. Armentrout, A. Kavner, R. B. Kaner, and S. H. Tolbert, Phys. Rev. B 85, 064118 (2012).
- [9] Q. Gu, F. Krauss, and W. Steurer, Adv. Mater. 20, 3620 (2008).
- [10] X. Cheng, W. Zhang, X. Chen, H. Niu, P. Liu, K. Du, G. Liu, D. Li, H. Cheng, H. Ye and Y Li, Appl. Phys. Lett. 103, 171903 (2013).
- [11] M. Wang, Y. Li, T. Cui, Y. Ma, and G. Zou, Appl. Phys. Lett. 93, 101905 (2008).
- [12] Q. Tao, D. Zheng, X. Zhao, Y. Chen, Q. Li, Q. Li, C. Wang, T. Cui, Y. Ma, X. Wang, and P. Zhu, Chem. Mater. 26, 5297 (2014).
- [13] M. Xie, R. Mohammadi, C. L. Turner, R. B. Kaner, A. Kavner, and S. H. Tolbert, Phys. Rev. B 90, 104104 (2014).
- [14] H. Tang and S. Ismail-Beigi, Phys. Rev. Lett. 99, 115501

- (2007).
- [15] A. R. Oganov and V. L. Solozhenko, J. Superhard Mater. 31, 285 (2009).
- [16] A. R. Oganov, J. Chen, C. Gatti, Y. Ma, Y. Ma, C. W. Glass, Z. Liu, T. Yu, O. O. Kurakevych, and V. L. Solozhenko, Nature 457, 863 (2009).
- [17] W. Zhou, H. Sun, and C. F. Chen, Phys. Rev. Lett. 105, 215503 (2010).
- [18] Y. Wang, J. Lv, L. Zhu, and Y. Ma, Comput. Phys. Commun. 183, 2063 (2012); the CALYPSO code is free for academic use, see details at http://www.calypso.cn.
- [19] Y. Wang, J. Lv and L. Zhu, and Y. Ma, Phys. Rev. B 82, 094116 (2010).
- [20] J. Lv, Y. Wang, L. Zhu, and Y. Ma, Phys. Rev. Lett. 106, 015503 (2011).
- [21] L. Zhu, H. Wang, Y. Wang, J. Lv, Y. Ma, Q. Cui, Y. Ma, and G. Zou, Phys. Rev. Lett. 106, 145501 (2011).
- [22] H. Liu and Y. Ma, Phys. Rev. Lett. 110, 025903 (2013).
- [23] M. Zhang, H. Liu, Q. Li, B. Gao, Y. Wang, H. Li, C. Chen, and Y. Ma, Phys. Rev. Lett. 114, 015502 (2015).
- [24] Q. Li, D. Zhou, W. Zheng, Y. Ma, and C. F. Chen, Phys. Rev. Lett. 110, 136403 (2013).
- [25] J. P. Perdew, K. Burke, and M. Ernzerhof, Phys. Rev. Lett. 77, 3865 (1996).
- [26] G. Kresse, and J. Furthmüller, Phys. Rev. B 54, 11169 (1996).
- [27] G. Kresse, and D. Joubert, Phys. Rev. B 59, 1758 (1999).
- [28] M. P. Teter, M. C. Payne, and D. C. Allan, Phys. Rev. B 40, 12255 (1989).
- [29] D. M. Bylander, L. Kleinman, and S. Lee, Phys Rev. B 42, 1394 (1990).
- [30] H. J. Monkhorst and J. D. Pack, Phys. Rev. B 13, 5188 (1976).
- [31] Y. Zhang, H. Sun and C. F. Chen, Phys. Rev. Lett. 93, 195504 (2004).
- [32] Z. C. Pan, H. Sun, and C. F. Chen, Phys. Rev. Lett. 98, 135505 (2007).
- [33] Z. C. Pan, H. Sun and C. F. Chen, Phys. Rev. B 79, 104102 (2009).
- [34] Z. C. Pan, H. Sun, Y. Zhang, and C. F. Chen, Phys. Rev. Lett. 102, 055503 (2009).
- [35] D. Roundy, C.R. Krenn, M. L. Cohen, and J.W. Morris Jr., Phys. Rev. Lett. 82, 2713 (1999).
- [36] S. Ogata, J. Li, and S. Yip, Science 298, 807 (2002).
- [37] Y. Zhang, H. Sun, and C. F. Chen, Phys. Rev. Lett. 94, 145505 (2005).
- [38] Y. Zhang, H. Sun, and C. F. Chen, Phys. Rev. B 73, 144115 (2006).
- [39] Y. Zhang, H. Sun, and C. F. Chen, Phys. Rev. B 76, 144101 (2007).
- [40] J. Yang, H. Sun, and C. F. Chen, J. Am. Chem. Soc. 130, 7200 (2008).
- [41] See Supplemental Material for all the calculated full stress-strain curves and the detailed information on the peak stress and strain values under the compressive, tensile, pure shear and Vickers indentation shear conditions.