is where the dynamics will come in, and the result will depend in general on the choice.

The essential points are as follows. We first take the meson matrices $\mathcal{M}(\phi(x))$ and $\mathcal{M}^{+}(\phi(x))$ which behave as $(3_R, 3_L^*)$ and $(3_L, 3_R^*)$. We also have definite expressions for the vector and axial currents $v_\mu{}^i$ and $a_\mu{}^i$. These will be expanded and expressed in terms of $\phi^i$ and $\partial_\mu \phi^i$. In discussing the leptonic decays of mesons, we replace the basic hadronic current $j_\mu$ in $H_W$ by the above meson current. If we regard this as an effective Hamiltonian, we get automatically all possible processes like $K \to \mu\nu$, $K \to \mu\nu + \pi$, $K \to \mu\nu + 2\pi$. In the case of non-leptonic meson decays, one way may be to form a product of meson currents (probably with a phenomenological coupling constant rather than $G_0$).

Another way is to introduce a spurion matrix $S \sim \lambda_6$, which should belong to $(8_L, 1_R)$ according to the basic current-current Hamiltonian. We then form an invariant with respect to $SU(3)_L \times SU(3)_R$ out of $S$, $\mathcal{M}$ and $\mathcal{M}^+$. A simplest non-trivial form is

$$H_{eff} = \mathrm{Tr}(\partial_\mu \mathcal{M} \, S \, \partial_\mu \mathcal{M}^+)$$

The derivatives are necessary since otherwise $\mathcal{M}\mathcal{M}^+ = 1$. By expanding $\mathcal{M}$ and $\mathcal{M}^+$ we can reproduce the relations obtained earlier about $K \to \pi, 2\pi, 3\pi$ decays. In particular, the energy and mass dependence of the amplitudes arises from the derivatives in the above formula.