the four-momentum squared t carried by the nonleptonic weak interaction and evaluate the dispersion integral at $t = 0$. We may furthermore suppose that the integral is dominated by low-lying meson states with suitable quantum numbers. Now if scalar and pseudoscalar meson octets of low mass exist, while 27-plets do not, then we have an explanation of octet predominance over the $\underline{27}$. Moreover, this explanation can be generalized, more or less, to the electromagnetic as well as strong mass differences.[19)]

An important feature[8)] of the "tadpole" mechanism is that the charge-conjugation properties of the parity-violating spurion octet and of the known pseudoscalar meson octet are opposite. For the $|\Delta Y| = 1$ part of the nonleptonic interaction, this means that the "tadpole" mechanism of octet enhancement is forbidden by SU(3) symmetry; of course, the enhancement effect could still be considerable despite the SU(3)-forbiddenness, like the ratio of $K_1^0 \to 2\pi$ and $K^+ \to 2\pi$ rates.[12,22)] If we turn to the $\Delta Y = 0$ nonleptonic interaction, however, we see that the disagreement in charge conjugation properties between the parity-violating spurion and the pseudoscalar mesons completely forbids the tadpole mechanism of octet enhancement. For $\Delta Y = 0$, then, tadpoles are powerless to make $|\Delta \underline{I}| = 0$ more important than $|\Delta \underline{I}| = 2$ in parity-violating weak forces.

A somewhat different mechanism for octet enhancement has been proposed by Dashen and Frautschi,[23)] extending earlier suggestions by Cutkosky and Tarjanne.[24)] For purposes of describing the mechanism, we consider the specific case of electromagnetic mass corrections.