Evidently, $A_{(0)}$ does not commute with the components of charged fields. In this gauge, then, the dependence of the electric field upon the charged fields is made explicit through the decomposition of the electric field into transverse and longitudinal parts,

$$F_{(0)(k)} = -\partial_{(0)} A_{(k)} - \partial_{(k)} A_{(0)}$$
$$= F^{(T)}_{(0)(k)} + F^{(L)}_{(0)(k)} .$$

The inference that the transverse fields are the independent dynamical variables of the electromagnetic field in this gauge is confirmed on examining the generators $G_{\delta A}$ and $G_{\delta F}$. Indeed,

$$G_{\delta A} = -\int d\sigma \; F_{(0)(k)} \, \delta A_{(k)} = -\int d\sigma \; F^{(T)}_{(0)(k)} \, \delta A_{(k)} ,$$

and

$$G_{\delta F} = \int d\sigma \; \delta F_{(0)(k)} \, A_{(k)} = \int d\sigma \; \delta F^{(T)}_{(0)(k)} \, A_{(k)}$$

in view of the transverse nature of $A_{(k)}$, Eq. (62). We can now derive the commutation properties of these dynamical variables from

$$\left[A_{(k)}, G_{\delta A}\right] = i\,\delta A_{(k)} \quad , \quad \left[F^{(T)}_{(0)(k)}, G_{\delta A}\right] = 0 ,$$

$$\left[F^{(T)}_{(0)(k)}, G_{\delta F}\right] = i\,\delta F^{(T)}_{(0)(k)} , \quad \left[A_{(k)}, G_{\delta F}\right] = 0 ,$$

on taking into account the restrictions

$$\partial_{(k)} A_{(k)} = \partial_{(k)} \, \delta F^{(T)}_{(0)(k)} = 0$$

produced by the transverse nature of these quantities. The Lagrange multiplier device permits us to deduce that

$$i\left[A_{(k)}(x) , F^{(T)}_{(0)(\ell)}(x')\right] = \delta_{(k)(\ell)} \, \delta_\sigma (x - x') + \partial'_{(\ell)} \lambda_{(k)} .$$