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**EXTENDED CALCULATIONS OF OECD/NEA PHASE II-C
BURNUP CREDIT CRITICALITY BENCHMARK PROBLEM
FOR PWR SPENT FUEL TRANSPORT CASK BY USING
MCNP-4B2 CODE AND JENDL-3.2 LIBRARY**

March 2003

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Extended Calculations of OECD/NEA Phase II-C Burnup Credit Criticality Benchmark Problem for
PWR Spent Fuel Transport Cask by Using MCNP-4B2 Code and JENDL-3.2 Library

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The reactivity effect of the asymmetry of axial burnup profile in burnup credit criticality safety is studied for a realistic PWR spent fuel transport cask proposed in the current OECD/NEA Phase II-C benchmark problem. The axial burnup profiles are simulated in 21 material zones based on in-core flux measurements varying from strong asymmetry to more or less no asymmetry. Criticality calculations in a 3-D model have been performed using the continuous energy Monte Carlo code MCNP-4B2 and the nuclear data library JENDL-3.2.

Calculation conditions are determined with consideration of the axial fission source convergence. Calculations are carried out not only for cases proposed in the benchmark but also for additional cases assuming symmetric burnup profile. The actinide-only approach supposed for first domestic introduction of burnup credit into criticality evaluation is also considered in addition to the actinide plus fission product approach adopted in the benchmark. The calculated results show that k_{eff} and the end effect increase almost linearly with increasing burnup axial offset that is defined as one of typical parameters showing the intensity of axial burnup asymmetry. The end effect is more sensitive to the asymmetry of burnup profile for the higher burnup. For an axially distributed burnup, the axial fission source distribution becomes strongly asymmetric as its peak shifts toward the top end of the fuel's active zone where the local burnup is less than that of the bottom end. The peak of fission source distribution becomes higher with the increase of either the asymmetry of burnup profile or the assembly-averaged burnup. The conservatism of the assumption of uniform axial burnup based on the actinide-only approach is estimated quantitatively in comparison with the k_{eff} result calculated with experiment-based strongest asymmetric axial burnup profile with the actinide plus fission product approach.

Keywords: Criticality Safety, Spent Fuel, Transport Cask, Burnup Credit, Axial Burnup Profile, End Effect, Asymmetry, Burnup Axial Offset, MCNP-4B2, JENDL-3.2

* MEXT Nuclear Researchers Exchange Program (Vietnam Atomic Energy Commission)

PWR 使用済み燃料輸送キャスクを対象とした OECD/NEA 燃焼度クレジット臨界ベンチマーク問題 II-C への MCNP-4B2 コード及び JENDL-3.2 ライブラリーを用いた拡張計算

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(2003 年 1 月 30 日受理)

現在進行中の OECD/NEA ベンチマーク問題 II-C において提案された、現実的な PWR 使用済み燃料輸送容器を対象とした、燃焼度クレジット臨界安全性における軸方向燃焼度分布の非対称性による反応度効果について研究した。炉内中性子束測定に基づいた強い非対称から、非対称が概ね消失するまでの変化を考慮すべく、軸方向燃焼度分布は 21 の組成領域により模擬されている。連続エネルギーモンテカルロコード MCNP-4B2 と核データライブラリー JENDL-3.2 を用いて、3 次元モデルの臨界計算を実施した。

軸方向核分裂源分布の収束性を考慮して計算条件を決定した。ベンチマーク問題で指定されたケースに加えて、対称な軸方向燃焼度分布を仮定した場合の計算も実施した。また、ベンチマーク問題におけるアクチニドと核分裂生成物を考慮する手法に加え、燃焼度クレジットを国内の臨界評価に導入する第一段階として、アクチニドのみを考慮する手法についても採用した。計算の結果、軸方向燃焼度非対称性の強度を示す典型的なパラメータとして定義された燃焼度軸方向オフセットの増加に伴って、実効増倍率及び端部効果はほぼ直線的に増加することが示された。また、より高い燃焼度に対して、燃焼度分布非対称性の端部効果への感度はより高い。軸方向分布を持つ燃焼度に対して、軸方向核分裂源分布は、燃料領域下端部よりも局所的燃焼度が低い上端部に向かってピークがシフトするという強い非対称となった。燃焼度分布の非対称性のみならず平均燃焼度の増加にも伴って、核分裂源分布のピーク値はより高くなった。実測値から得られた最も非対称性の強い軸方向燃焼度分布を用いてアクチニドと核分裂生成物を考慮する手法と軸方向燃焼度の一様分布を仮定したアクチニドのみを考慮する手法の実効増倍率計算結果を比較することにより、後者の保守性を定量的に評価することができる。

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1. Introduction

In the framework of the OECD/NEA burnup credit (BUC) criticality safety research activities, the effect of axial burnup profile on criticality of PWR spent fuel was studied for an infinite array of fuel rods with a finite axial height in the Phase II-A benchmark [1], and for a realistic spent fuel transport cask in the Phase II-B benchmark [2]. The end effect, which is defined as the reactivity difference between the neutron effective multiplication factor (k_{eff}) calculated with an explicit representation of the axial burnup profile, and k_{eff} calculated assuming a uniform axial burnup, was evaluated. It has been known from the Phases II-A and II-B as well as from other studies [3][4] that the end effect becomes positive for PWR spent fuel of an average burnup of more than about 30 MWd/kgU. This means that the assumption of uniform axial burnup is non-conservative for a high burnup, while it is conservative for a low burnup of below 30 MWd/kgU.

In the Phases II-A and II-B, the end effect was evaluated with an assumption of a symmetric axial burnup profile. In reality, however, the axial burnup profile of spent fuel is asymmetric due to asymmetry of the reactor's core configuration and operation history. In the late 2000, J. Ch. Neuber and Th. Lamprecht (Germany) have proposed the OECD/NEA Phase II-C burnup credit criticality benchmark for PWR spent fuels* loaded into the transport cask that is the same as in the Phase II-B benchmark in order to study the effect of the asymmetry of axial burnup profiles. The benchmark is based on a set of actual shapes of outstanding quality with the error in the nodal burnup being less than 5%. The axially asymmetric change with burnup is also taken into account. Thus, the OECD/NEA Phase II-C benchmark problem forms an experiment-based database for international studies of BUC criticality safety for PWR spent fuel of high burnup.

To study the effect of the asymmetry of axial burnup profiles proposed in the Phase II-C benchmark problem, we have performed three-dimensional criticality calculations by using the continuous energy Monte Carlo code MCNP-4B2 [5] and the nuclear data library JENDL-3.2 [6]. As additional cases to the benchmark problem, we propose symmetric burnup profiles of 32 MWd/kgU and 50 MWd/kgU in order to cover the range of burnup profile variation from pure symmetry to strong asymmetry [7]. Consequently, our calculations also include the cases of symmetric axial burnup profiles. Moreover, calculations based on the actinide-only approach for BUC criticality safety are also performed in addition to the actinide plus fission product approach. Such additional calculations are carried out as an extension of the original OECD/NEA Phase II-C benchmark problem. The obtained results contribute to the currently required international comparison in the OECD/NEA expert group on burnup credit and help to better understand the physics of PWR spent fuel criticality safety.

A description of the Phase II-C benchmark problem including the additional cases is presented

* Siemens Konvoi fuel assembly type FOCUSTM: Specifications can be downloaded from the website; <http://www.nea.fr/html/science/wpncs/buc/specifications/phase-2c/index.html>

in Chap. 2. The calculation method and model are presented in Chap. 3. In order to determine the calculation conditions concerning Monte Carlo simulation, convergence not only for k_{eff} but also for fission distribution should be carefully achieved. Hence, prior examinations for calculation conditions and convergence checks are performed and described with calculated results in Chap. 4. At the end of the report, a sample input of MCNP-4B2 is provided in Appendix I to show the calculation model in detail, and the calculated results for the fission source distributions illustrated in the section 4.4.2. are listed in the Tables in Appendix II.

2. Description of the Benchmark Problem

2.1. Configuration

The transport cask configuration of Phase II-C benchmark is similar to the one already used in Phase II-B (cf. Ref.2) in structure and arrangement of fuel assemblies. According to Phase II-C specification downloaded from the website, some modifications have been made due to the fact that the fuel assembly of the Konvoi type for Phase II-C is somewhat bigger than the 17x17 fuel assembly type for Phase II-B. The configuration is summarized below as shown in Fig. 1 and Fig. 2.

Cask shell

Inner diameter:	146cm,	Inner cavity height:	496cm,
Outer diameter:	206cm,	Outside height:	556cm,
Material: Stainless Steel of 30cm thickness, the same as for the upper and lower lids			

Cask interior

21 fuel assemblies are arranged in a 5x5 array without corner positions, separated by basket wall and positioned centrally within the basket regions. The cask is completely filled with water.

Assembly basket

Inner basket compartment dimensions: 23.5cm x 23.5cm x 496 cm per assembly position
 Basket wall thickness: 1cm (only one basket wall between two adjacent assembly positions)
 Material: Borated stainless steel (1wt% natural boron)

Fuel assembly (18x18-24 lattice)

Number of fuel rods: 300, Number of guide thimbles: 24, Rod pitch: 1.27cm
 Spacer grid: to be ignored
 Upper and lower hardware: 30.0 and 19.5cm high, respectively (width: 22.86cm),
 assumed as a smeared mixture of 50vol% water and 50vol% stainless steel.

Fuel rod

Pellet diameter:	0.822cm,	Cladding inner diameter:	0.822cm,
Cladding outer diameter:	0.95cm,	Cladding material:	Zircaloy,
Active fuel length:	390cm (fuel compositions will be described in the next section.),		
Full fuel length:	439.5cm,	Initial ²³⁵ U enrichment:	4.0 wt%,
Upper inactive zone:	17.1cm,	Lower inactive zone:	32.4cm

End plugs are ignored. The fuel pellet region is treated as void inside the inactive zone between top (or bottom) end of the active zone and upper (or lower) hardware, and the others inside the inactive zone are treated the same as the active zone.

Guide thimble

Inner diameter:	1.11cm,	Outer diameter:	1.232cm,
Material:	Zircaloy		

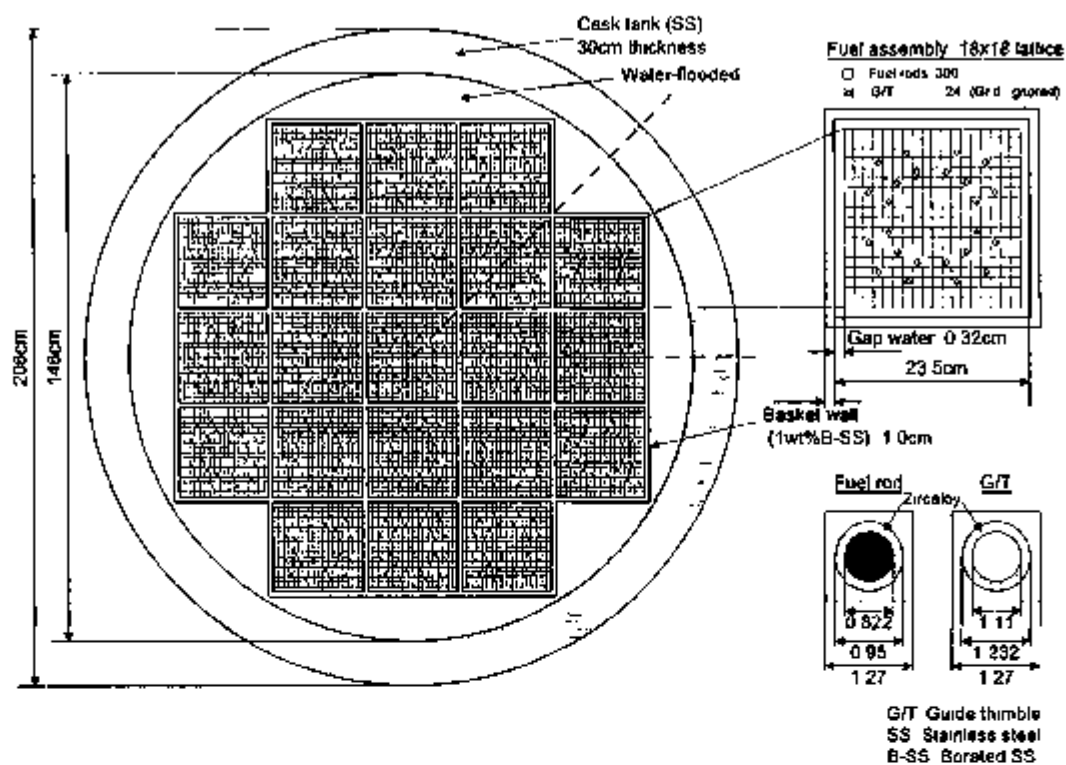


Fig. 1 Configuration of transport cask for Phase II-C benchmark

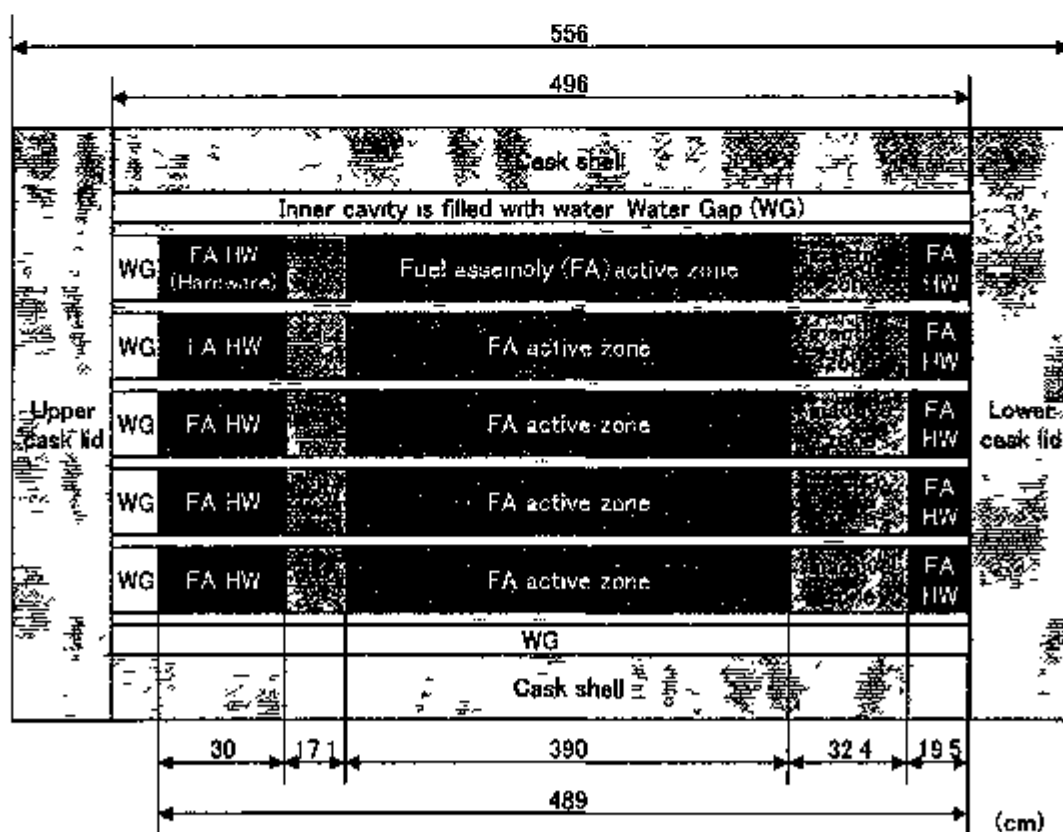


Fig. 2 Vertical cross section of transport cask for Phase II-C benchmark

2.2. Axial Burnup Profiles for Benchmark

2.2.1. Proposed Cases in Phase II-C

For the benchmark, initial ^{235}U enrichment of 4.0wt% and two average burnups of the fuel assemblies of 32MWd/kgU and 50MWd/kgU, and cooling time of five years are assumed.

Axial burnup profiles for Phase II-C benchmark problem are based on evaluations of 850 axial burnup shapes in fuel assemblies derived from in-core 3D power density (flux) distribution measurements. According to Phase II-C specifications, axial burnup shapes are evaluated by 850 various assembly-wise averaged burnups $\hat{B}_j$ (as shown by B(average) in Fig. 3); here each

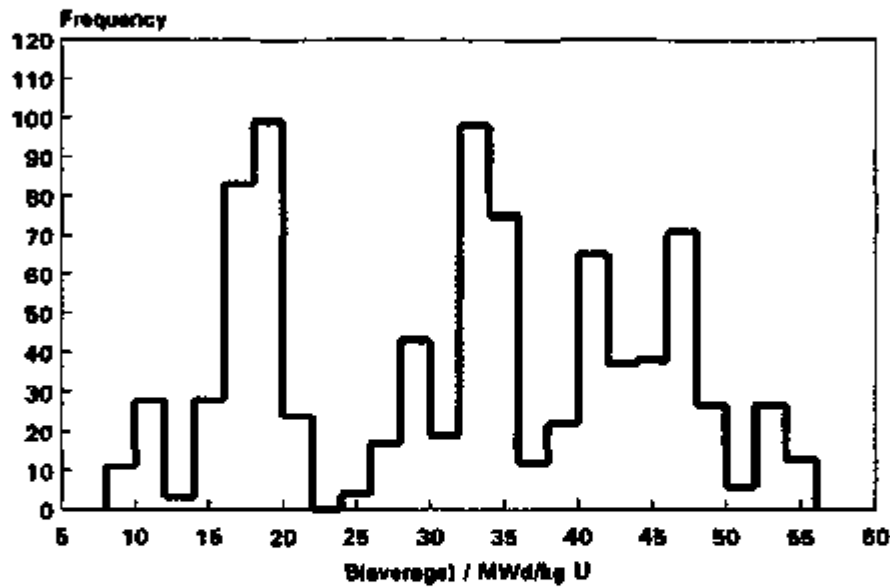


Fig. 3 Distribution of assembly-wise averaged burnups

assembly is referred to as j ($j=1,850$). $\hat{B}_j$ is given as the average of B_y of 850 axial burnup data monitored at 32 equidistant axial nodes, here, i means axial nodes ($i=1,32$).

$$\hat{B}_j = \frac{1}{32} \sum_{i=1}^{32} B_{yi} \quad (1)$$

The analyses are carried out by means of the relative burnup distribution $\alpha_y (= B_y / \hat{B}_j)$. The sample mean $\hat{\alpha}_i$ of the ratio α_y as shown by solid line in Fig. 4 is the basic case proposed in Phase II-C.

$$\hat{\alpha}_i = \frac{1}{N} \sum_{j=1}^N \alpha_{yi} = \sum_{j=1}^N \frac{B_{yi}}{\hat{B}_j}, \quad (i=1, \dots, 32, N=850) \quad (2)$$

The dashed line on a log scale in Fig.4 shows the standard deviation of axial burnup data,

simultaneously. It is found that the variance is significantly greater at the top end of fuel zone than that at the others.

$$\hat{\sigma}_i = \sqrt{\frac{1}{N-1} \sum_{j=1}^N (\alpha_{ij} - \hat{\alpha}_i)^2}, (i=1, \dots, 32, N=850) \quad (3)$$

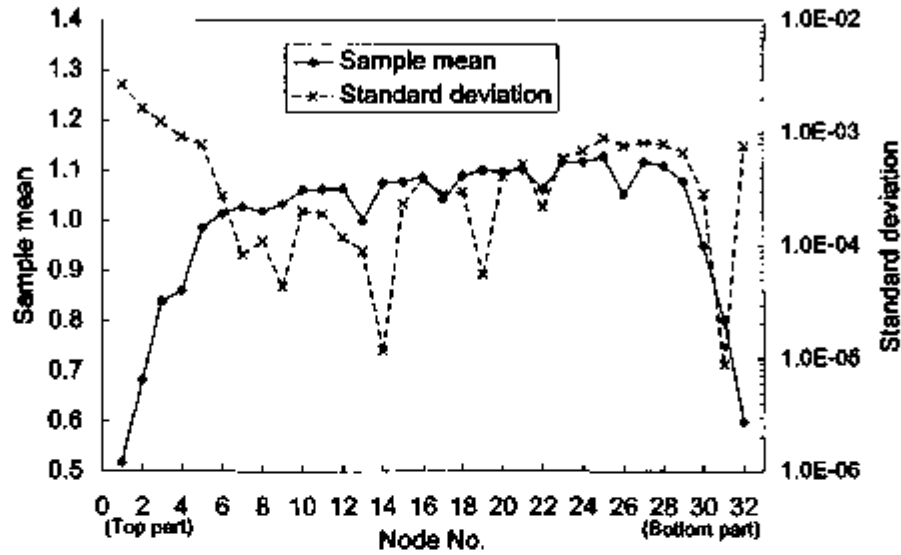


Fig. 4 Sample mean and standard deviation of axial burnup profiles

In consideration of the measured data whereby the decrease of the height of the burnup shape at the top nodes results in higher asymmetry of the shape and increase of the height of the shape at top nodes results in lower asymmetry, various burnup shapes are proposed in the benchmark by adding the mean values $\hat{\alpha}_i$ of the top six nodes by $\Delta\alpha_i$ ($i=1,6$). The variation of the top nodes is also counterbalanced by multiplying the mean values of the nodes $m=18$ through $m=29$ with f_m in order to preserve the sum of $\hat{\alpha}_i$ as described below.

$$f_m = s(m-17)+1, (m=18, \dots, 29) \quad (4)$$

$$s = \frac{\sum_{i=1}^6 \Delta\alpha_i}{\sum_{m=18}^{29} (m-17)\hat{\alpha}_m} \quad (5)$$

The variations $\Delta\alpha_i$ are bounded by the minimum and maximum values of the ratio α_{ij} observed in the sample of axial shapes under examination. It is taken into account that the minimum and maximum values depend on the average burnup. The observed minimum and maximum values are listed in Table 1.

Table 1 The observed minimum and maximum values of the ratios α_0

Node No. i	Average Burnup MWd/kg U	Min $j = 1, \dots, 850 (\alpha_{ij})$	Max $j = 1, \dots, 850 (\alpha_{ij})$
1	32	0.444	0.638
	50	0.492	0.665
2	32	0.621	0.778
	50	0.666	0.763
3	32	0.786	0.926
	50	0.818	0.917
4	32	0.808	0.936
	50	0.845	0.920
5	32	0.950	1.047
	50	0.978	1.036
6	32	0.977	1.062
	50	1.006	1.051

The cases proposed in the benchmark are given in Table 2. The cases with the serial numbers 2 through 7, 1, and 8 through 13 in this order represent the variation of the shape from strong asymmetry (case No.2) to more or less no asymmetry (case No.13). The cases No. 14 and 15 represent those often had when a fuel assembly is positioned at the edge of the core during its first

Table 2 Cases proposed by OECD/NEA BUC Phase II-C benchmark

Serial No.	Identification ¹⁾	Node					
		1	2	3	4	5	6
1	BbbCyyA22222***	Mean***	Mean	Mean	Mean	Mean	Mean
2	BbbCyyA11111	Min***	Min	Min	Min	Min	Min
3	BbbCyyA11112	Min	Min	Min	Min	Min	Mean
4	BbbCyyA11122	Min	Min	Min	Min	Mean	Mean
5	BbbCyyA11222	Min	Min	Min	Mean	Mean	Mean
6	BbbCyyA12222	Min	Min	Mean	Mean	Mean	Mean
7	BbbCyyA22222	Min	Mean	Mean	Mean	Mean	Mean
8	BbbCyyA22223	Mean	Mean	Mean	Mean	Mean	Max****
9	BbbCyyA22233	Mean	Mean	Mean	Mean	Max	Max
10	BbbCyyA22333	Mean	Mean	Mean	Max	Max	Max
11	BbbCyyA23333	Mean	Mean	Max	Max	Max	Max
12	BbbCyyA33333	Mean	Max	Max	Max	Max	Max
13	BbbCyyA33333	Max	Max	Max	Max	Max	Max
14	BbbCyyA12223	Min	Mean	Mean	Mean	Mean	Max
15	BbbCyyA12233	Min	Min	Mean	Mean	Max	Max
16	BbbCyyA32221	Max	Mean	Mean	Mean	Mean	Min
17	BbbCyyA32211	Max	Max	Mean	Mean	Min	Min

*: Bbb stands for the average burnup (B32: 32MWd/kgU, B50: 50MWd/kgU).

Cyy stands for the cooling time (C05: five years cooling time is just considered in Phase II-C).

Annunnn is for the identification of the variation of the top six nodes (1: Min, 2: Mean, 3: Max).

**: Reference case (cf. Solid line in Fig. 4)

***: Mean stands for $\hat{\alpha}_j$ given by Eq. (2), Min and Max are given in Table 1.

cycle. In the cases No. 16 and 17, the difference of the averaged burnup of 6 top nodes from that of the lower half is almost conserved. All cases are well suited to study the impact of the axial shape on the local reactivity within the fuel zone.

All burnup shapes are divided into 21 axial zones. For each of the 21 axial zones, the benchmark provides the number densities of twenty-eight nuclides such as actinides and fission products. The division of the axial shapes is shown in Table 3.

Table 3 Division of the axial burnup shape in 21 zones

Axial Zone No.	Nodes collapsed	Length of the zone cm	Height of the upper bound of the zone, cm (with respect to the bottom of the fuel zone)
1	32	12.19	12.19
2	31	12.19	24.38
3	30	12.18	36.56
4	29	12.19	48.75
5	27, 28	24.38	73.13
6	26	12.18	85.31
7	23, 24, 25	36.57	121.88
8	22	12.18	134.06
9	20, 21	24.38	158.44
10	18, 19	24.37	182.81
11	17	12.19	195.00
12	14, 15, 16	36.56	231.56
13	13	12.19	243.75
14	10, 11, 12	36.56	280.31
15	7, 8, 9	36.57	316.88
16	6	12.18	329.06
17	5	12.19	341.25
18	4	12.19	353.44
19	3	12.19	365.63
20	2	12.18	377.81
21	1	12.19	390.00

2.2.2. Additional Cases as an Extension of Phase II-C

The benchmark provides the axial burnup shapes varying from slight asymmetry to strong asymmetry. Naturally, one is interested in knowing the end effect for a symmetric axial burnup shape, although it is probably not a realistic one. We proposed symmetric shapes for two average burnups as additional cases to the benchmark problem in order to cover the range of burnup shape variation from a pure symmetry to a strong and realistic asymmetry.

The burnup shapes for symmetric cases, identified as BbbC05Symmetric, are assumed from those of the cases BbbC05A333333, which is the least asymmetric burnup shape, by averaging burnup data of two nodes that are located symmetrically about mid point in the axial direction. The

axial burnup shapes of BbbC05A333333 and BbbC05Symmetric for 32MWd/kgU and 50MWd/kgU are presented in Fig. 5 and Fig. 6, respectively. We can see that, the symmetric burnup shapes BbbC05Symmetric improve its symmetry compared with the least asymmetric burnup shapes BbbC05A333333.

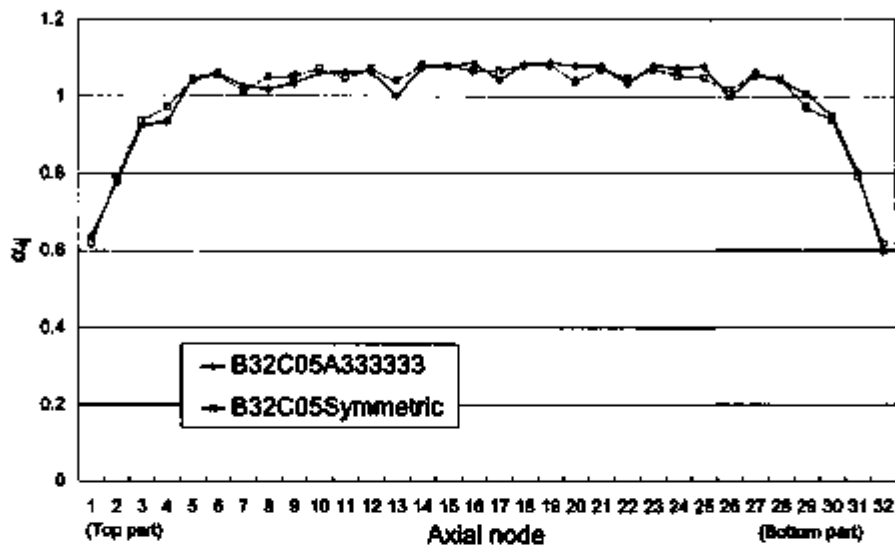


Fig. 5 Axial burnup shapes of B32C05A333333 and B32C05Symmetric

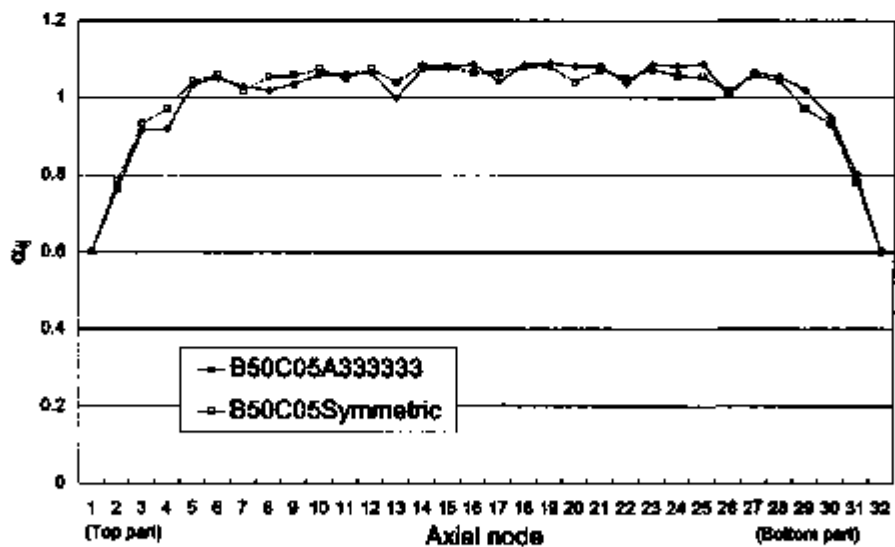


Fig. 6 Axial burnup shapes of B50C05A333333 and B50C05Symmetric

The number densities of each nuclide in the 32 equidistant nodes for BbbC05Symmetric are obtained by averaging the ones in the corresponding 32 equidistant nodes for BbbC05A333333 in the same manner as explained above for burnup shape.

The original OECD/NEA Phase II-C benchmark problem involves fifteen nuclides of fission product in the fuel composition. Here, the calculations for these cases are referred to as the actinide plus fission product (AC+FP) approach. In addition, we carried out calculations for the actinide-only (AC-only) approach, in which the fission products are neglected.

2.2.3. Introduction of Burnup Axial Offset

As a representative parameter of intensity of axial burnup asymmetry, we introduce burnup axial offset (burnup A. O.). The definition of the burnup A. O. is as follows. The values of burnup A. O. are listed in Table 4.

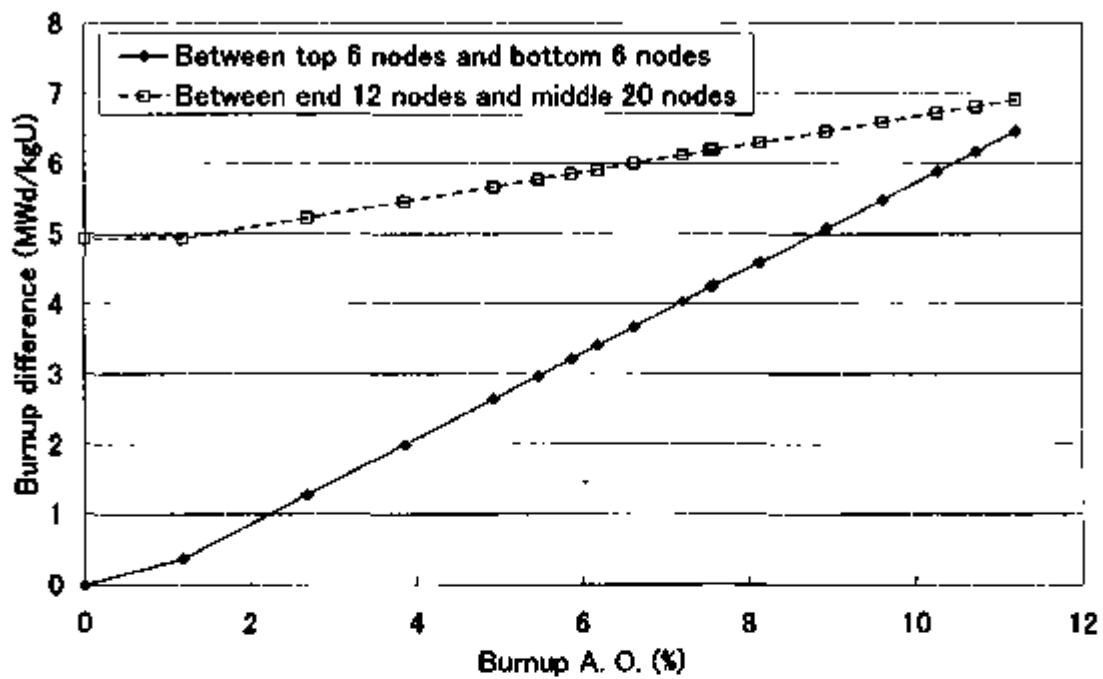
Burnup A. O. $\equiv$ (average burnup of lower 16 nodes – that of upper 16 nodes) / that for all nodes

Table 4 Burnup A. O.

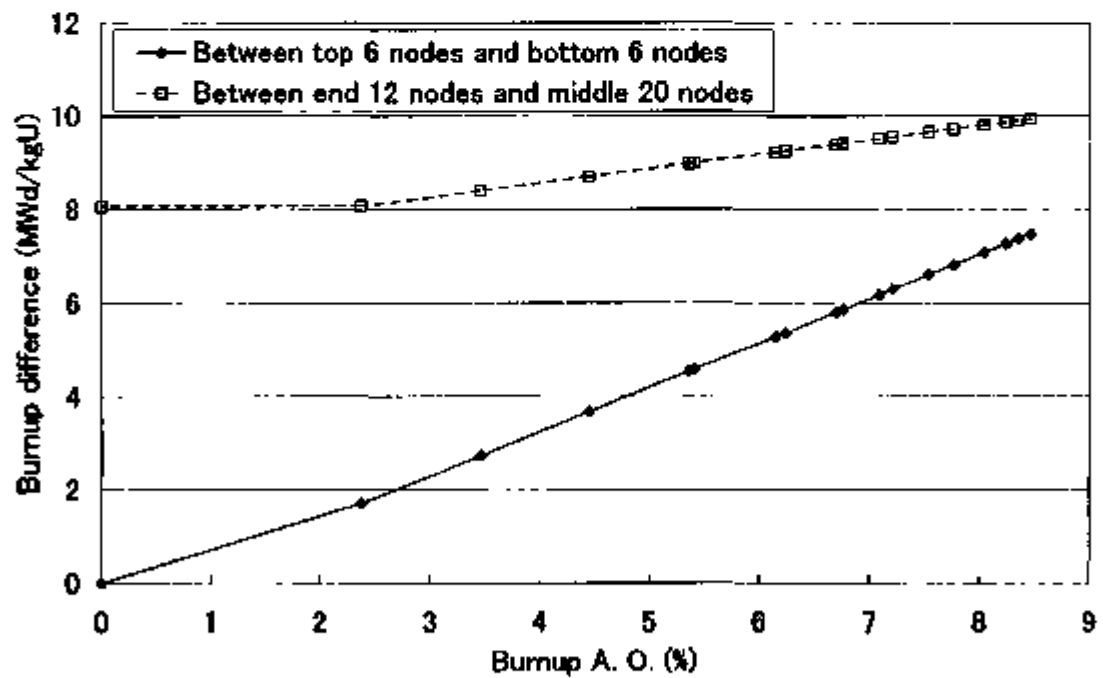
Case ID	Serial No.	Burnup A.O. (%)	Case ID	Serial No.	Burnup A.O. (%)
B32C05Symmetric	--	0.00	B50C05Symmetric	--	0.00
B32C05Uniform	--	0.00	B50C05Uniform	--	0.00
B32C05A222222	1	7.21	B50C05A222222	18	7.21
B32C05A111111	2	11.20	B50C05A111111	19	8.46
B32C05A111112	3	10.73	B50C05A111112	20	8.35
B32C05A111122	4	10.27	B50C05A111122	21	8.24
B32C05A111222	5	9.60	B50C05A111222	22	8.04
B32C05A112222	6	8.93	B50C05A112222	23	7.77
B32C05A122222	7	8.14	B50C05A122222	24	7.54
B32C05A222223	8	6.62	B50C05A222223	25	6.76
B32C05A222233	9	5.87	B50C05A222233	26	6.15
B32C05A222333	10	4.93	B50C05A222333	27	5.41
B32C05A223333	11	3.86	B50C05A223333	28	4.45
B32C05A233333	12	2.68	B50C05A233333	29	3.46
B32C05A333333	13	1.19	B50C05A333333	30	2.38
B32C05A122223	14	7.55	B50C05A122223	31	7.09
B32C05A112223	15	7.59	B50C05A112223	32	6.70
B32C05A322221	16	6.19	B50C05A322221	33	6.24
B32C05A332211	17	5.48	B50C05A332211	34	5.36

In Table 4, the uniform burnup case is identified by BbbC05Uniform.

Region-wise averaged burnup differences both between top 6 nodes and bottom 6 nodes, and between end part of 12 nodes and middle part of 20 nodes, increase linearly with increasing burnup A. O. as shown in Fig. 7. Consequently, the burnup A. O. shows the intensity of not only burnup profile asymmetry but also the neutron decoupling between upper and lower fuel parts.



(a) 32 MWd/kgU



(b) 50 MWd/kgU

Fig. 7 Burnup differences vs. burnup A. O.

3. Calculation Method, Model and Conditions

3.1. Calculation Method

The calculation method is a combination of the continuous energy Monte Carlo code MCNP-4B2 [5] and the nuclear data library JENDL-3.2 [6]. The KCODE option of MCNP has been employed in criticality calculations for 3-D model of the transport cask.

3.2. Calculation Model and Conditions

It was found that the input cards of the primary results as shown in the reference [7] had some mistakes. The lack of the $S(\alpha,\beta)$ identifier affects the k_{eff} results significantly, and the re-calculated results described in the next chapter show that the old k_{eff} results are underestimated by about 1%. In the re-calculations, the input cards were corrected and modified. The three major modifications are summarized below. A sample input is shown in Appendix I.

- (1) The radial 1/8 symmetry of the cask configuration as shown by dashed line in Fig. 1 is modeled explicitly in the input cards. It contributes to reduce the number of initial source points, namely, computer memory. All calculations are carried out on EWS: SUN/Blade100 (Solaris 8) where the MCNP is installed with the dynamic memory option. The dynamic storage of the sample input as shown in the appendix is 6810117 words, which is still more than the default value of the static memory parameter (MDAS = 4000000).
- (2) The 32 axial equidistant cells (not 21 non-equidistant cells) are explicitly specified in the input cards in order to show the fission source distribution smoothly. In addition, the compositions of 21 non-equidistant zones are redistributed to the corresponding 32 equidistant cells.
- (3) The initial guess for fission source is flat and distributed uniformly in all fissile fuel cells for all cases (symmetric, uniform, asymmetric cases).

In addition, calculation conditions are reexamined. At first, we assume 300 skipped cycles to check the axial fission source convergence for a pin cell configuration. We have to increase the skipped cycles to reach convergence in some cases of slow convergence. The number of active histories, namely, the product of the number of scored cycles and the number of histories per cycle, is taken to be 2,000,000 in order to estimate k_{eff} (1σ is about 0.0005) and the end effect (more or less 1%) except for the uniform burnup cases. It is necessary to increase the number of active histories in cases of uniform burnup to obtain a converged fission source distribution supposed to be a cosine-like curve in shape. The combination of the number of skipped cycles and the number of histories per cycle is determined as a result of convergence check for fission source distribution, which is described in section 4.2.

The numbers of neutron history for the final estimation are listed in Table 5.

Table 5 Numbers of neutron history

Numbers of history*	Cases	Comments
20000 x (300 + 100)	Symmetric 4 cases	At 32 & 50MWd/kgU for both AC-only and AC+FP approach
5000 x (300 + 4000)	Uniform 4 cases	
5000 x (300 + 400)	Asymmetric cases except below	1 σ for k_{eff} is nearly equal to 0.0005
5000 x (1500 + 400)	B32C05A233333, B32C05A333333 for AC+FP approach	Fission source slow convergence needs many skipped cycles.
5000 x (4500 + 400)	B32C05A233333 for AC-only app.	
5000 x (6000 + 400)	B32C05A333333 for AC-only app.	

*: Number of histories per cycle x (number of skipped cycles + number of scored cycles)

4. Calculated Results

4.1. Prior Examination of Skipped Cycles Based on the Pin Cell Configuration

It is important to check the slow convergence caused by the axial burnup distribution in the benchmark. The aim in this section is to determine the number of the skipped cycles enough to reach convergence of fission source distribution. Hence, we have examined the convergence of the fission source distribution for an infinite array system of pin cells that is modeled as full axial structure of the cask for the Phase II-C. We adopt the pin cell configuration prior to the realistic cask because it has the advantage that effects caused by the axial burnup distribution can be distinguished from effects of the radial structure (18x18-24 fuel lattice, basket wall of B-SS and water gap, 21 fuel assemblies arranged in a 5x5 array, cask wall, and so on). The simple configuration decreases the probability of units without neutron generation points arisen and reduces the calculation time to process the configuration when the Monte Carlo simulations are executed.

As representative cases, the symmetric case: B32C05Symmetric, the uniform case: B32C05Uniform and one of the asymmetric cases: B32C05A222222 are analyzed by the AC+FP approach. Criticality calculations are performed by varying the number of skipped cycles with other calculation conditions of 20,000 histories per cycle, 100 scored cycles, and flat and uniform initial guess of fission source. Figure 8 shows calculated k_{eff} per skipped cycles. Increase of skipped cycles of more than 100 affects slightly the k_{eff} trends for the three cases as shown in Fig. 8. In the following sections, we examine the number of skipped cycles to reach convergence of the fission source distribution for the three cases, respectively.

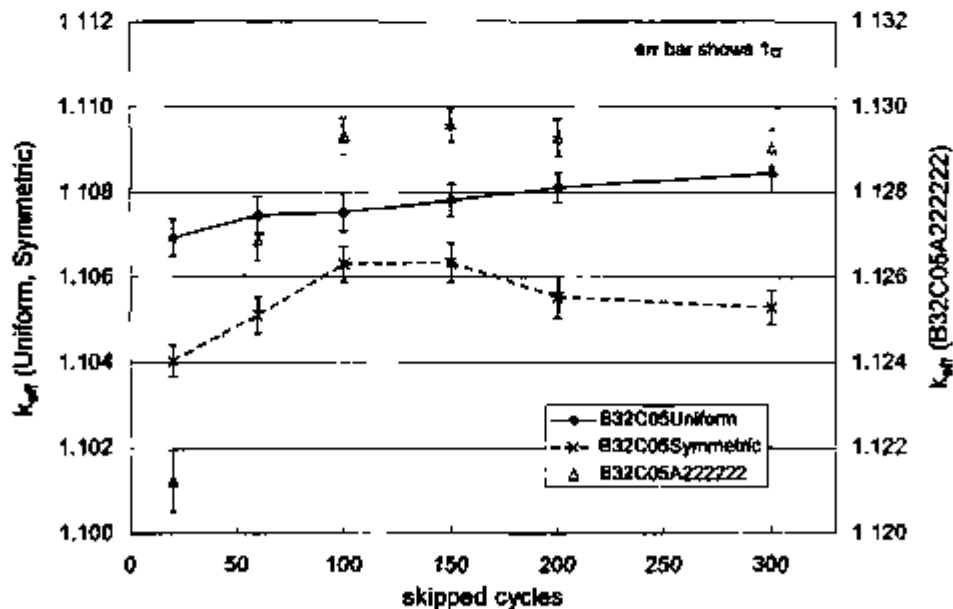


Fig. 8 Calculated k_{eff} vs. skipped cycles for the pin cell model based on Phase II-C
[20,000 histories/cycle, 100 cycles scored]

4.1.1. Symmetric Burnup Case

Figure 9 shows calculated results of fission source distributions with various numbers of skipped cycles for the symmetric burnup case identified by B32C05Symmetric.

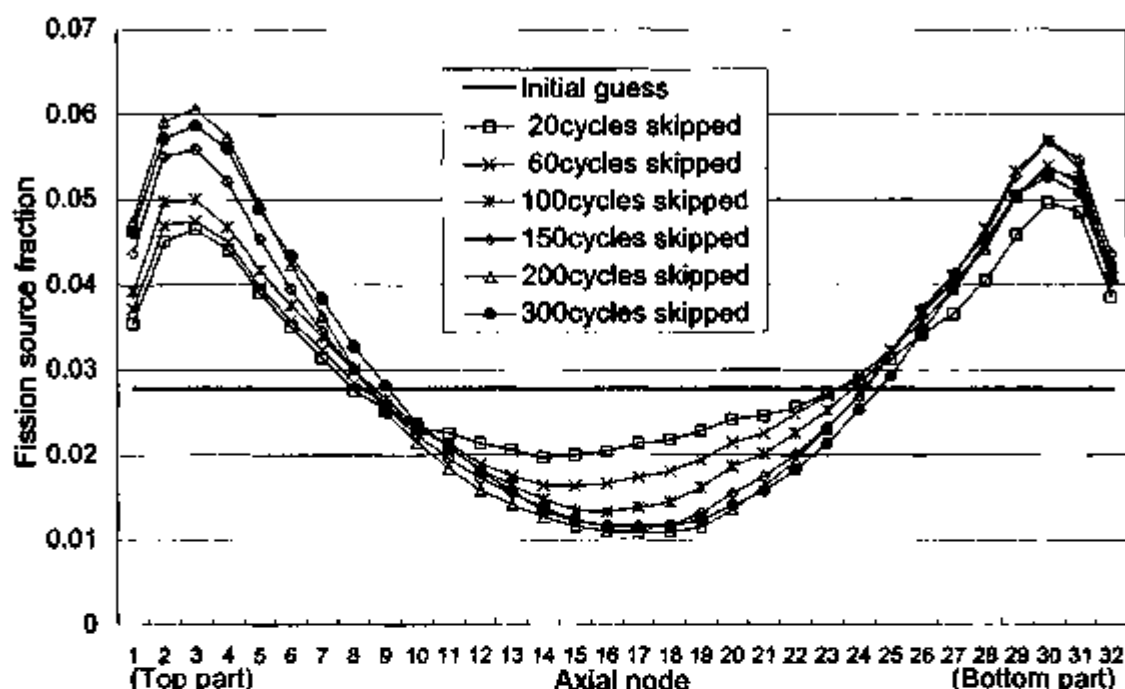


Fig. 9 Fission source distributions of B32C05Symmetric for various skipped cycles based on the pin cell model [20,000 histories/cycle, (20-300) skipped, 100 cycles scored]

As shown in Fig. 9, it is found that the initial guess affects the fission source in case of insufficient skipped cycles. The fission source at middle parts of the fuel decreases gradually with the increase of skipped cycles, and becomes almost constant for more than 150 skipped cycles. On the other hand, it seems that fission source at the top and bottom parts increases with the increase of skipped cycles, and still oscillates slightly even for 300 skipped cycles. Let's study the trends of fission fractions at the 3 parts (top, middle and bottom), which are obtained by summing those of six nodes at the top and bottom parts, and twelve nodes at the middle part shown in Fig. 9. As shown in Fig. 10, the fission fractions are almost constant for more than 200 scored cycles. Even though the burnup distribution is axially symmetric, the fission source is not symmetric such that the fission fraction at the top part is greater than that at the bottom part. It may be caused by a difference of the reflector effect between the upper and lower structures. The reactivity difference however, is found to be so small that it becomes very difficult to estimate accurately how it affects the fission source distribution as a result of some calculations.

Therefore, we conclude that 300 cycles suffice as the number of skipped cycles for the axially symmetric burnup case to estimate the fission source distribution. The combination of the number of histories per cycle and the number of scored cycles and its effect on the fission source distribution will be discussed in the next section based on the cask configuration.

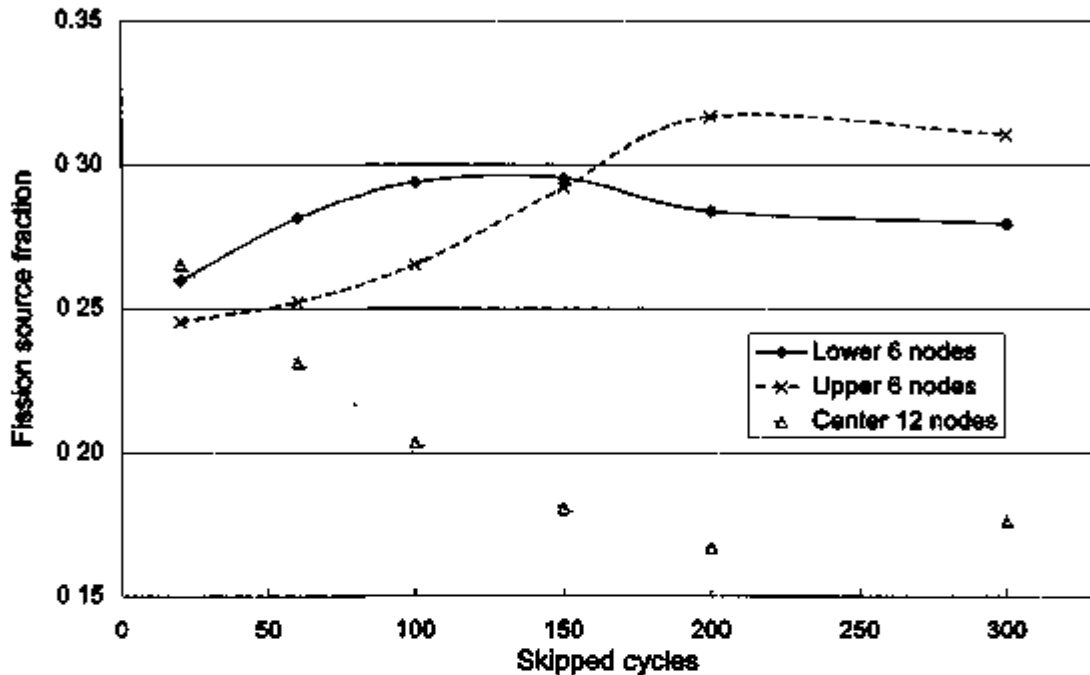


Fig. 10 Fission fractions of 6 nodes at the top and bottom parts and the center 12 nodes for various skipped cycles based on the pin cell model of B32C05Symmetric

4.1.2. Uniform Burnup Case

Figure 11 shows calculated results of fission source distributions with various numbers of skipped cycles for the uniform burnup case identified by B32C05Uniform.

As shown in Fig. 11, it is found that the fission source at the middle part is affected by the initial guess of flat distribution in case of insufficient skipped cycles. It is supposed to become near cosine-like in shape with increase of the skipped cycles because the reflector structure difference between top and bottom parts little affects the fission source distribution. But, no matter by how many the skipped cycles increase, the fission source distribution could not be represented by a smooth curve. It is necessary to increase the total scored histories, namely the product of the number of the histories per cycles and the number of the scored cycles, in order to obtain a smooth fission source distribution like a cosine curve.

Therefore, we conclude that 300 skipped cycles are enough for the uniform burnup case to estimate the fission source distribution. The results for the increase of the total scored histories will

be provided in the subsection 4.3.2 based on the cask configuration.

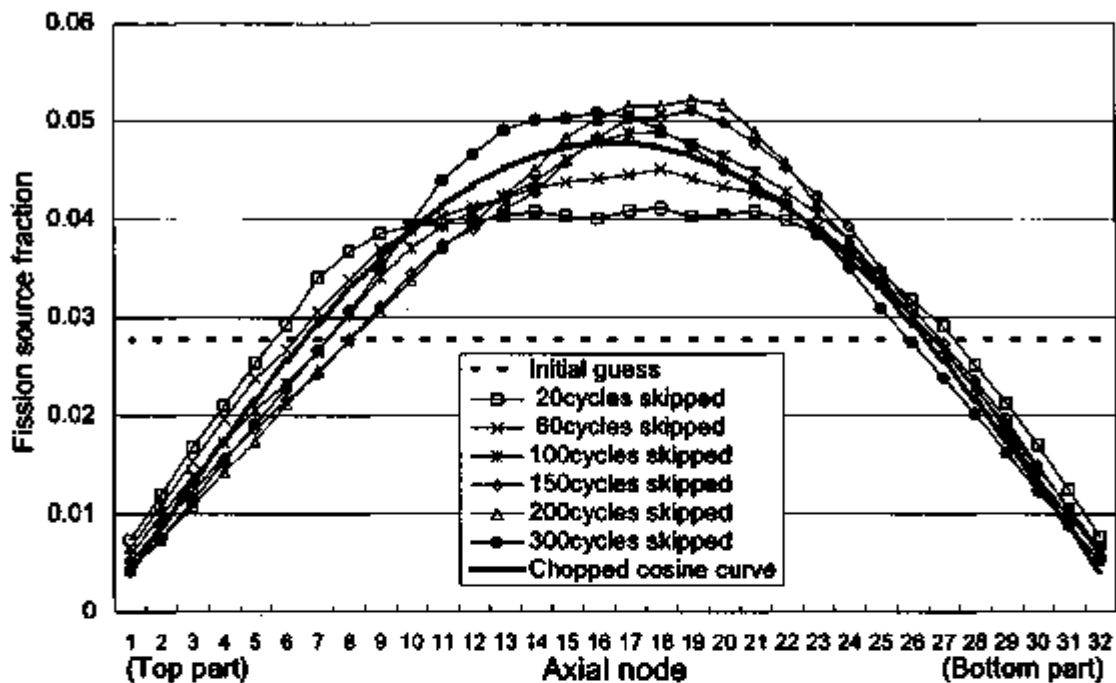


Fig. 11 Fission source distributions of B32C05Uniform for various skipped cycles based on the pin cell model (20,000 histories/cycle, (20-300) skipped, 100 cycles scored)

4.1.3. Asymmetric Burnup Case

Figure 12 shows calculated results of fission source distributions with various numbers of skipped cycles for the asymmetric burnup case identified by B32C05A222222.

Regarding the asymmetric burnup cases, the fission source distribution localizes on the top part because of its low burnup. As shown in Fig. 12, the peak value at the top part increases and the source level at the bottom part decreases with increase of the skipped cycles. Figure 13 shows the fission fractions of 4 nodes at the top and bottom parts. The fission fraction at the top 4 nodes becomes almost constant at more than 150 skipped cycles. The fission fraction at the bottom 4 nodes is shown by the logarithm axis on the right hand side in Fig. 13 because of its very low source level. It is still decreasing at 300 skipped cycles with increase of skipped cycles as shown in Fig. 13. But the change of the fission source at the bottom part is very small and does not affect the k_{eff} estimation. Therefore, we regard the fission source distribution as converged at the skipped cycles of 300 or more.

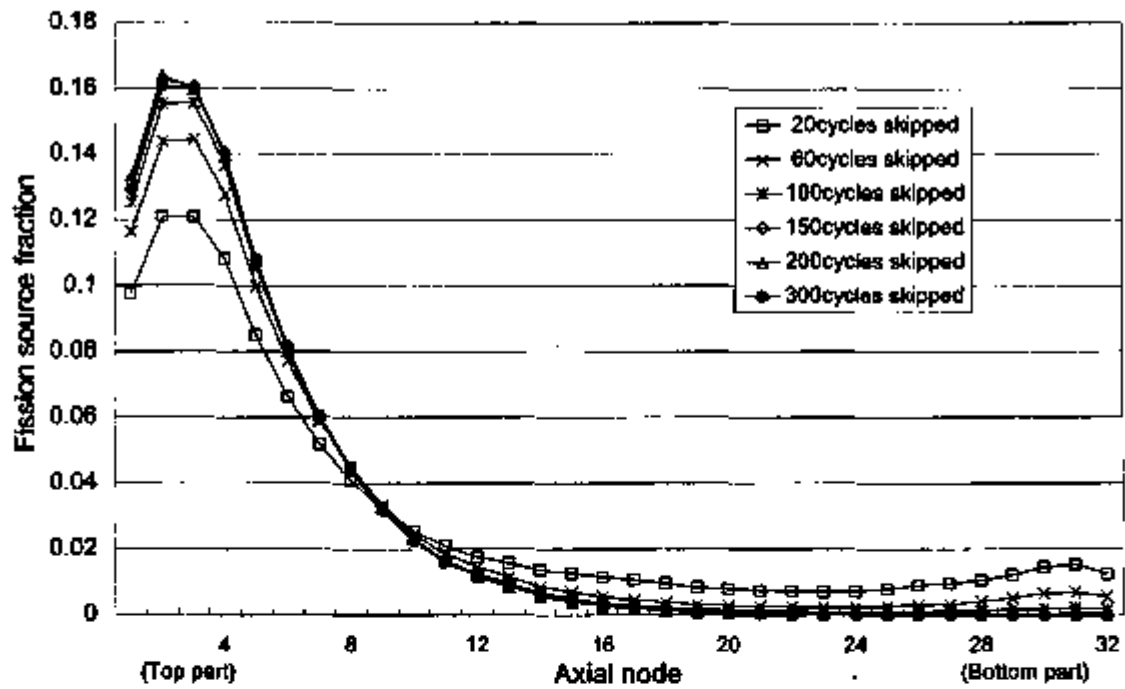


Fig. 12 Fission source distributions of B32C05A222222 for various skipped cycles based on the pin cell model [20,000 histories/cycle, (20-300) skipped, 100 cycles scored]

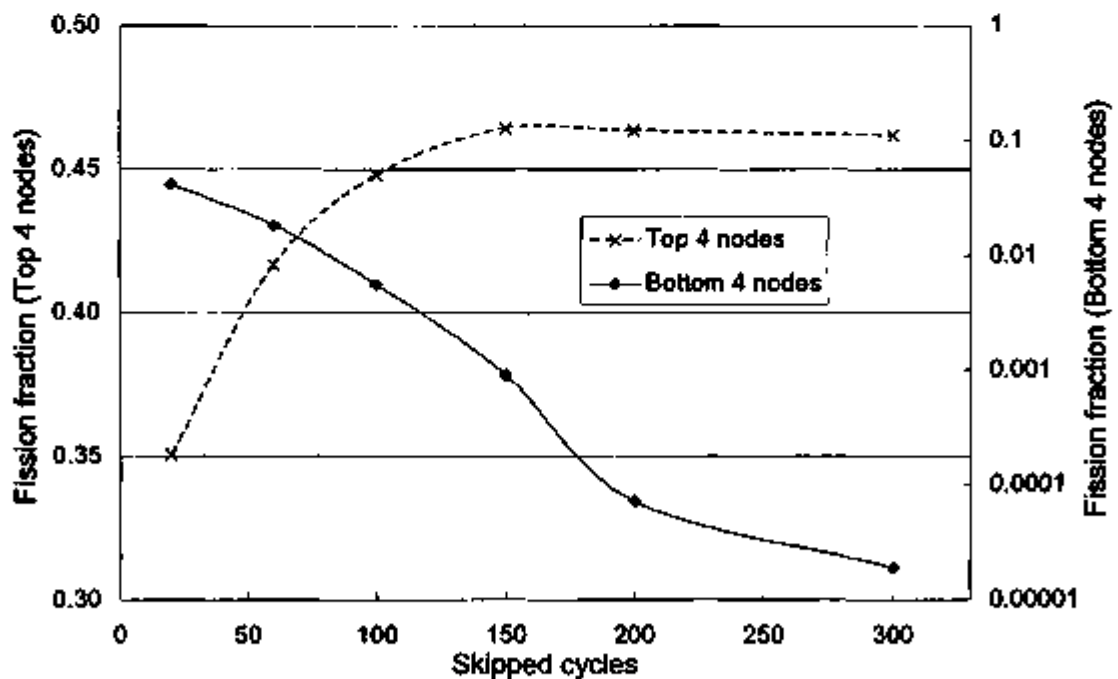


Fig. 13 Fission fractions of 4 nodes at the top and bottom parts for various skipped cycles based on the pin cell model of B32C05A222222

4.2. Radial Fission Distribution Based on Single Assembly and Color Set Configurations

In order to check the radial convergence of fission source distribution, we have also examined the single assembly configuration and 2x2 multi assembly configuration named color set configuration, here. In the color set configuration, three fuel assemblies and a water channel are considered as shown in Fig. 14. The radial boundary conditions are reflective, and the 1/8 radial symmetry is also considered for both configurations.

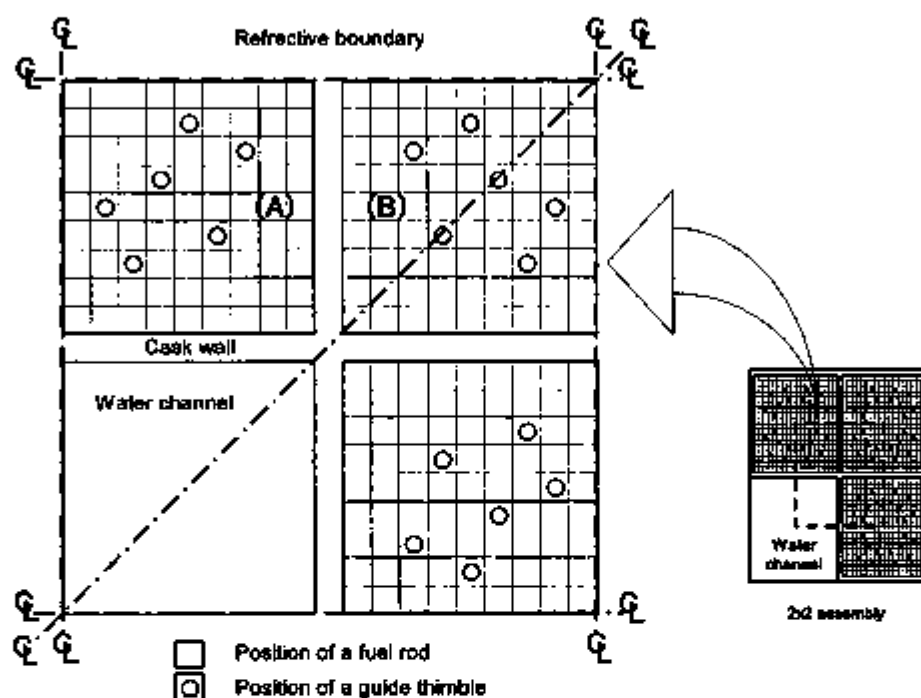


Fig. 14 Calculation model of the color set configuration

The calculated results of the radial fission source distribution for the single assembly configuration are shown in Fig. 15, 16 and 17 corresponding to the case ID of B32C05Symmetric, B32C05Uniform and B32C05A222222, respectively. In Fig. 15 to 17, the calculation conditions are 150 skipped cycles and 20000 histories per cycle, and the results for both 10 and 100 scored cycles are shown.

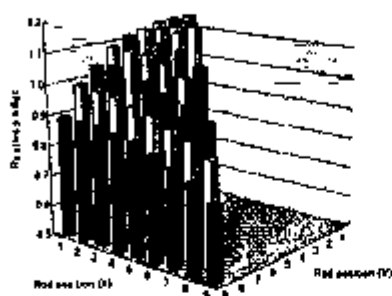
As shown in Fig. 15 to 17, the radial fission source distributions for the three cases inside the assembly are almost the same as a chopped cosine curve in shape in spite of the axial fission source distribution difference, and they have already converged with the number of 150 skipped cycles.

The calculated result of the radial fission source distribution for the color set configuration is shown in Fig. 18 corresponding to the case ID of B32C05A222222. The calculation conditions are 300 skipped cycles and 20000 histories per cycle, and the results for both 10 and 100 scored cycles are shown. As shown in Fig. 18, the averaged fission source inside the quarter assembly (B) shown in Fig. 14 is greater than that inside the eighth assembly (A) shown in Fig. 14. And the radial fission

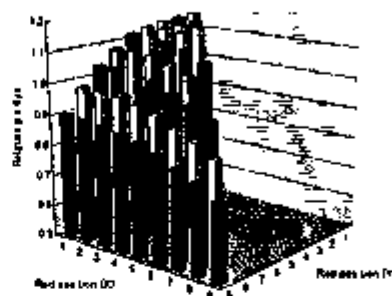
source distributions inside the assemblies are similar to a chopped cosine curve in shape

The radial fission source distribution for the color set configuration not only inside the assembly but also assembly-wise averages has already converged as shown in Fig. 18

These facts show that the radial source convergence is faster than the axial one. In addition, the calculation time increases very much in order to print out the pin-wise fission source. Hence, for the cask configuration, only the axial convergence is checked in the next section

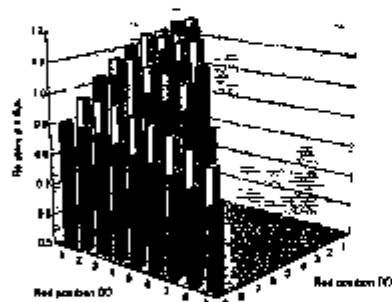


(a) 10 scored cycles

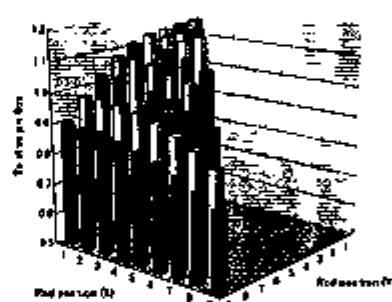


(b) 100 scored cycles

Fig. 15 Radial fission source distribution of single assembly configuration for B32C05Symmetric (difference of colors only means to distinguish between Y axis rod position)

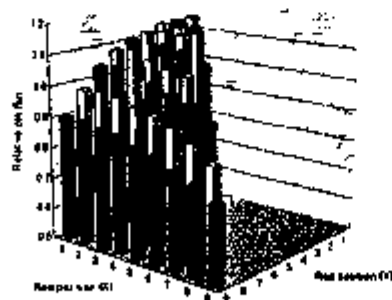


(a) 10 scored cycles

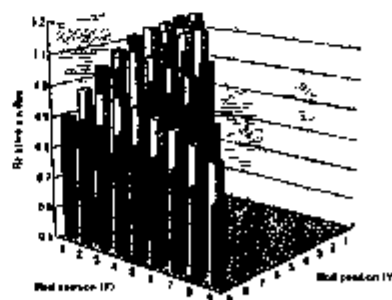


(b) 100 scored cycles

Fig. 16 Radial fission source distribution of single assembly configuration for B32C05Uniform (difference of colors only means the same as the other figures)



(a) 10 scored cycles



(b) 100 scored cycles

Fig. 17 Radial fission source distribution of single assembly configuration for B32C05A222222

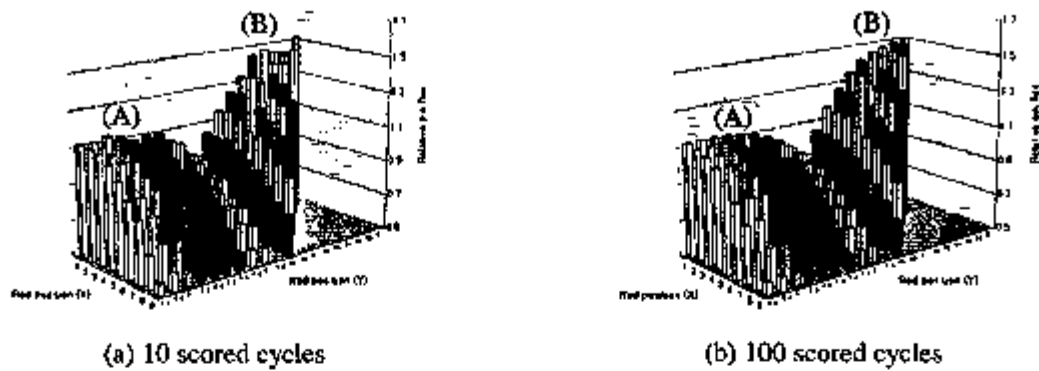


Fig. 18 Radial fission source distribution of color set configuration for B32C05A222222

4.3. Convergence Check of Axial Fission Source Distribution Based on Cask Configuration

As a result of the prior examinations described in the previous sections, we determine the number of skipped cycles to be set at 300 in all cases of the cask calculation. Calculations for the Phase II-C cask configuration are carried out. As calculation conditions, the combination of the number of histories per cycle and the number of scored cycles are examined. Then, we check the axial source convergence and reset the calculation conditions to improve convergence as described in the section 3.2.

4.3.1. Symmetric Burnup Case

The calculation conditions of 3-D cask model described in Chapter 2 for B32C05Symmetric based on the AC+FP approach are examined as listed in Table 6. The calculated results for the axial fission source are shown in Fig. 19

Table 6 Calculation conditions for B32C05Symmetric examined

No.	Skipped cycles	Scored cycles	Histories per cycle	Total scored histories
1	300	4,000	500	2,000,000
2	300	2,000	1,000	2,000,000
3	300	400	5,000	2,000,000
4	300	2,000	5,000	10,000,000
5	300	100	20,000	2,000,000

As shown in Fig. 19, it is found that the results for each condition are quite different. In order to check the fission source convergence, for instance, the fission source trends per scored cycle for No. 4 and 5 are shown in Fig. 20 and Fig. 21, respectively. As shown in Fig. 20, the fission source

distribution changes with the increase of scored cycles. The same trend is observed for the calculation conditions of No. 2 and 3. In these cases, the number of histories per cycle is so small that the fission source distribution trends become unstable. On the other hand, 20,000 histories per cycle in case No. 5 is enough for the fission source distribution trend to be stable as shown in Fig.21.

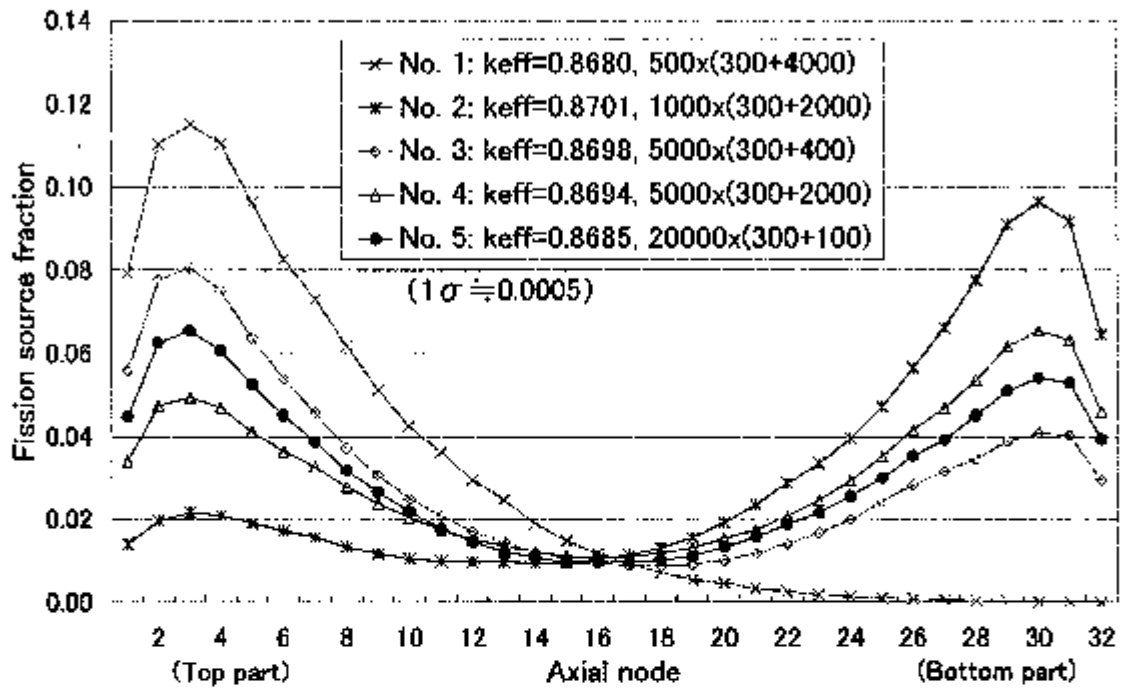


Fig. 19 Fission source distributions of B32C05Symmetric for various conditions

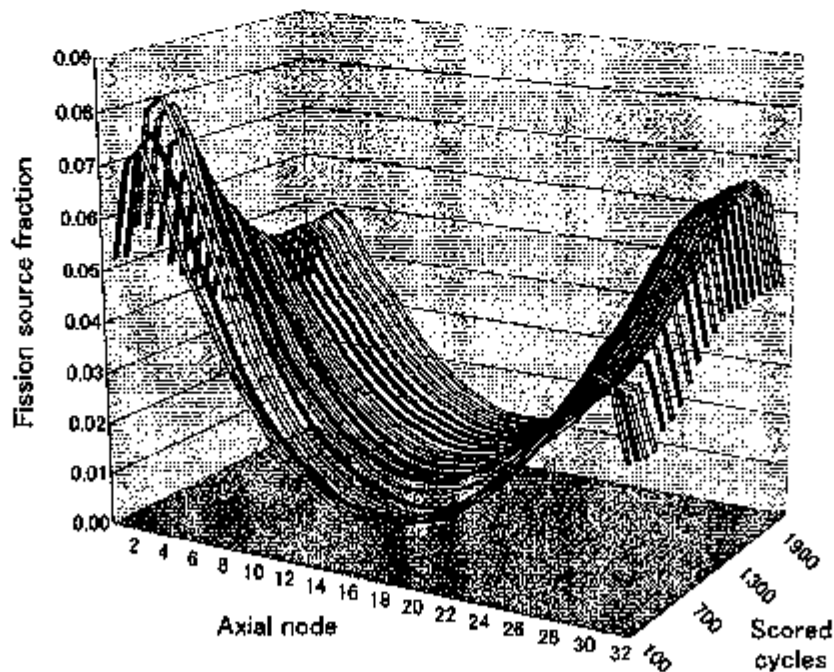


Fig. 20 Fission source distribution trend of B32C05Symmetric (calculation condition No. 4)

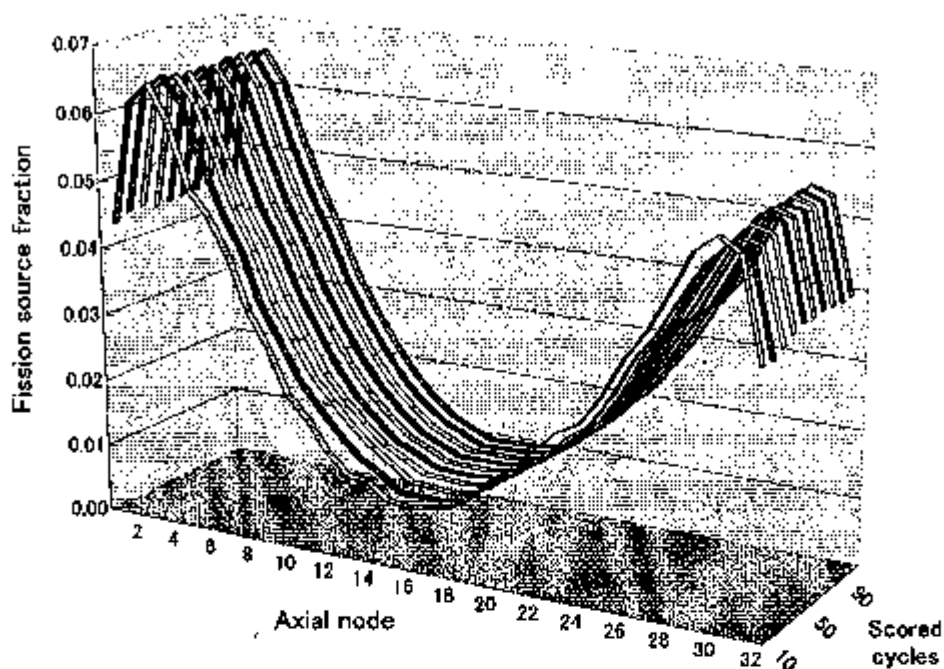


Fig. 21 Fission source distribution trend of B32C05Symmetric (calculation condition No. 5)

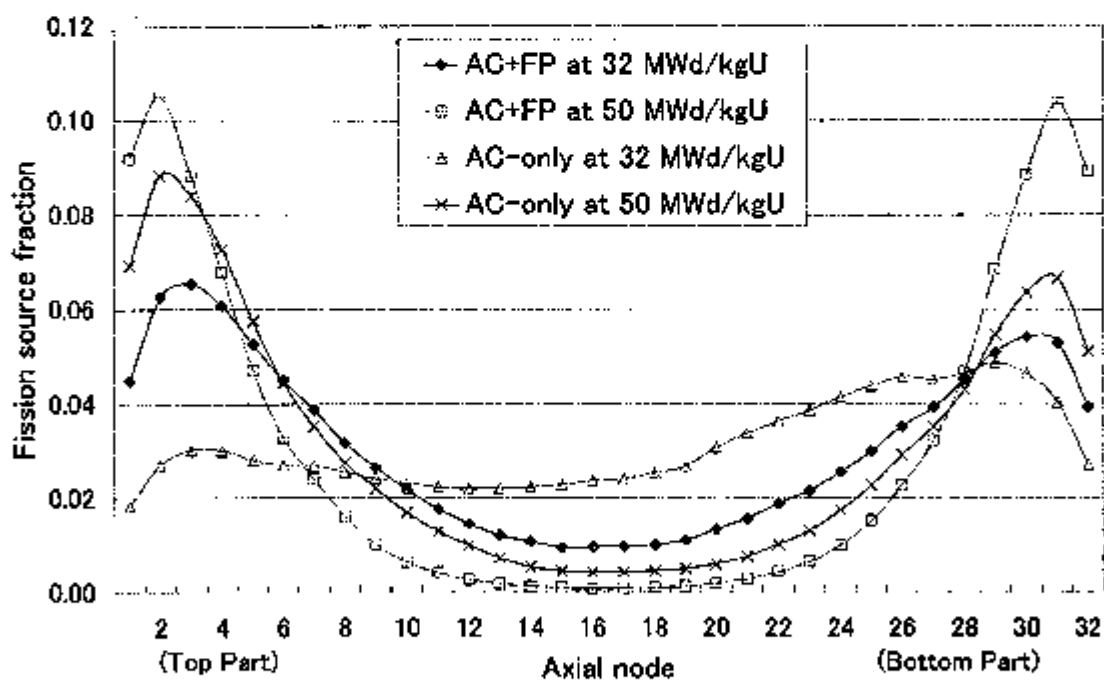


Fig. 22 Calculated results for fission source distribution in cases of the symmetric burnup

Therefore, we determine the calculation conditions for all symmetric cases to be set as No. 5 in Table 6. The calculated results for the fission source distributions are almost symmetric as shown in Fig. 22.

4.3.2. Uniform Burnup Case

The calculation conditions of 3-D cask model for B32C05Uniform based on the AC+FP approach are examined as listed in Table 7. The calculated results for the axial fission source are shown in Fig. 23

Table 7 The calculation conditions for B32C05Uniform examined

No.	Skipped cycles	Scored cycles	Histories per cycle	Total scored histories
1	300	1,000	2,000	2,000,000
2	300	400	5,000	2,000,000
3	300	100	20,000	2,000,000

The converged fission source distributions for the uniform burnup case are expected to be cosine-like in shape, and the results of the fission source distribution for each condition are a little bit different from cosine curve as shown in Fig. 23. It is necessary to increase the total scored histories in order to obtain smoother and more cosine-like shape. The calculated results of all uniform cases for condition No.2 in Table 7 and for the case of increased histories are shown in Fig. 24(a) and Fig. 24(b), respectively. The calculation condition of the latter, that is 300 skipped cycles, 4,000 scored cycles and 5,000 histories per cycle are considered appropriate to give the final results for the uniform burnup cases in this report.

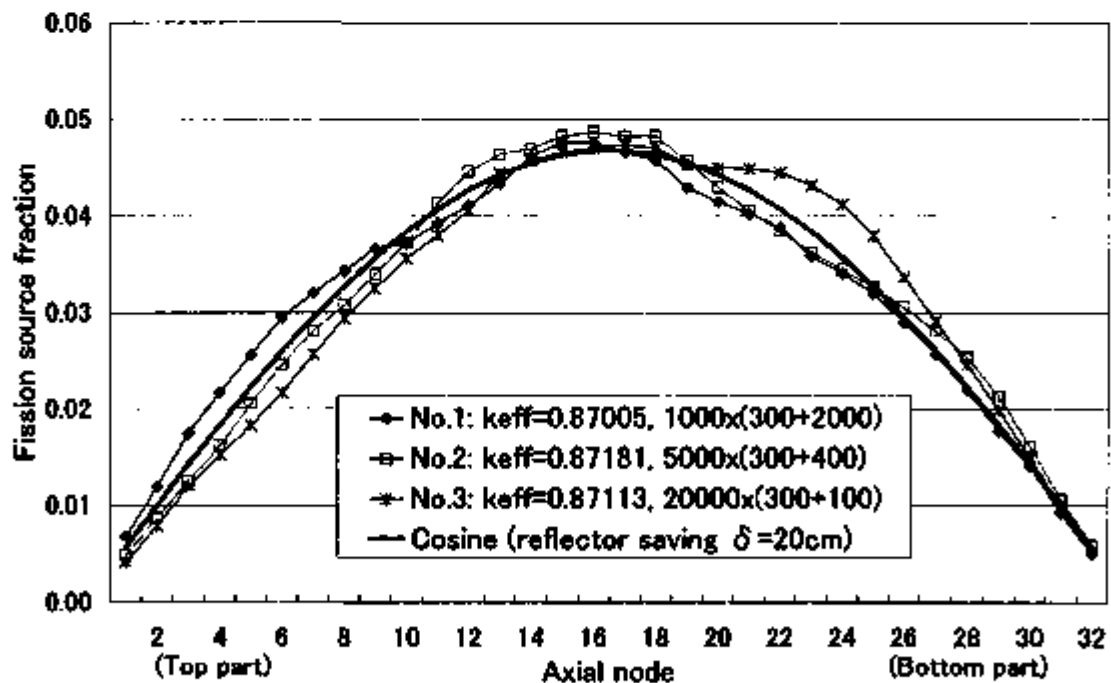
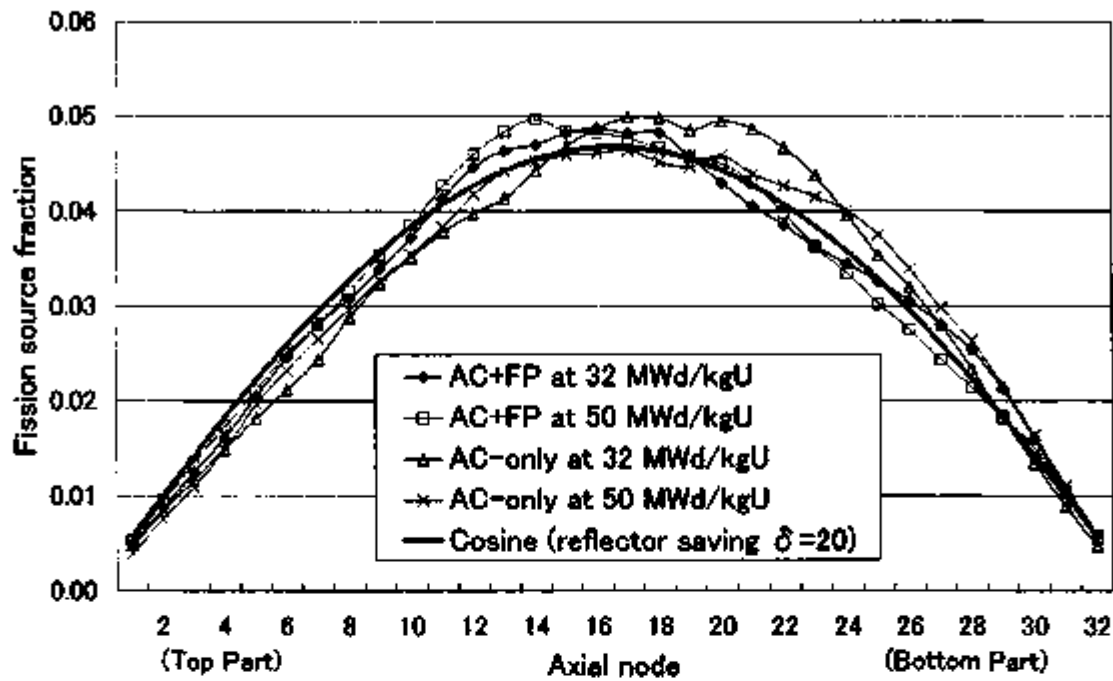
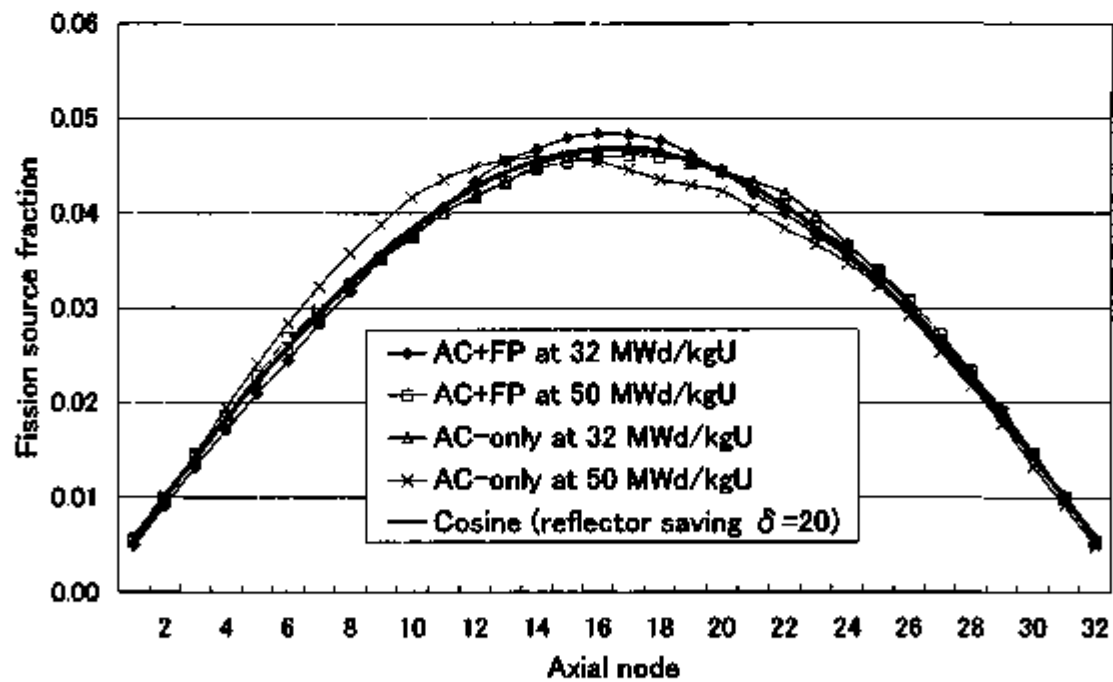


Fig. 23 Calculated results for fission source distribution in cases of the uniform burnup



(a) Not final results: 5000 histories per cycle x (300 skipped cycles + 400 scored cycles)



(b) Final results: 5000 histories per cycle x (300 skipped cycles + 4000 scored cycles)

Fig. 24 Calculated results for fission source distribution in cases of the uniform burnup

In order to check the fission source convergence of final results, the fission source trend per scored cycle of B32C05Uniform, for instance, is shown in Fig. 25. As shown in Fig. 25, the fission source distribution is observed to reach near convergence at 4000 scored cycles.

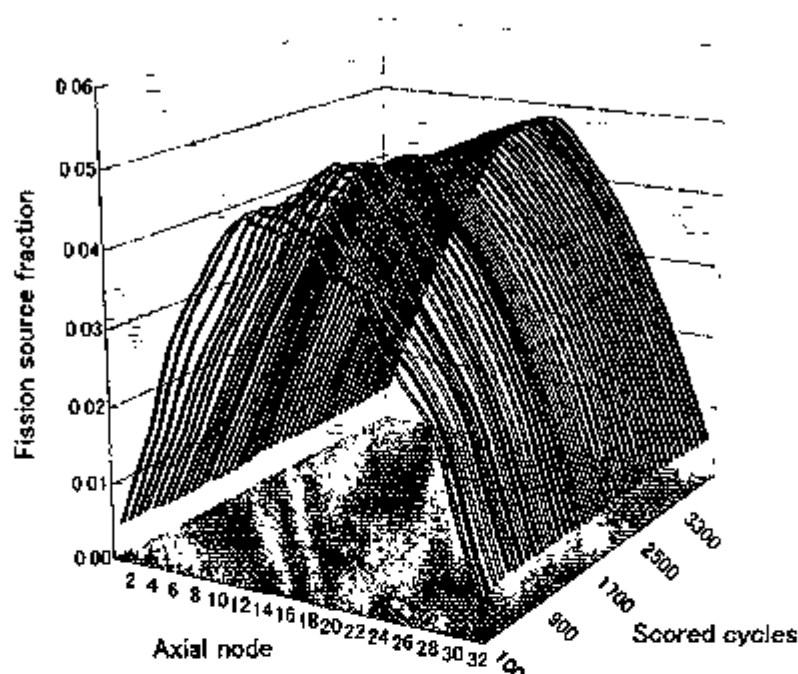


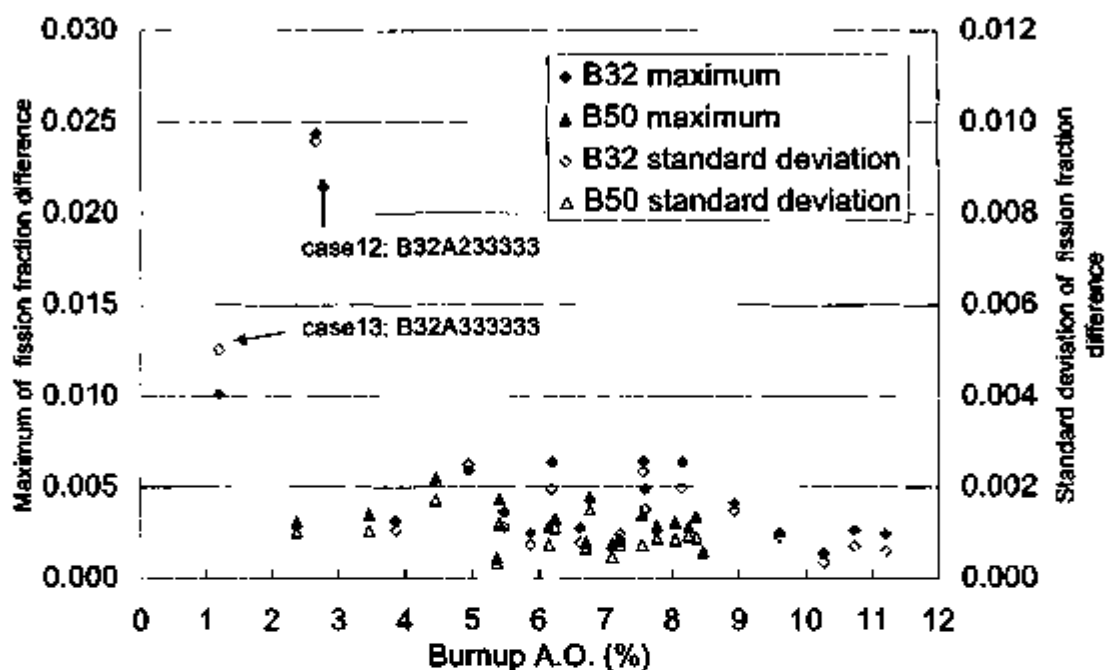
Fig. 25 Fission source distribution trend of B32C05Uniform in case of 5000 histories per cycle

4.3.3. Asymmetric Burnup Case

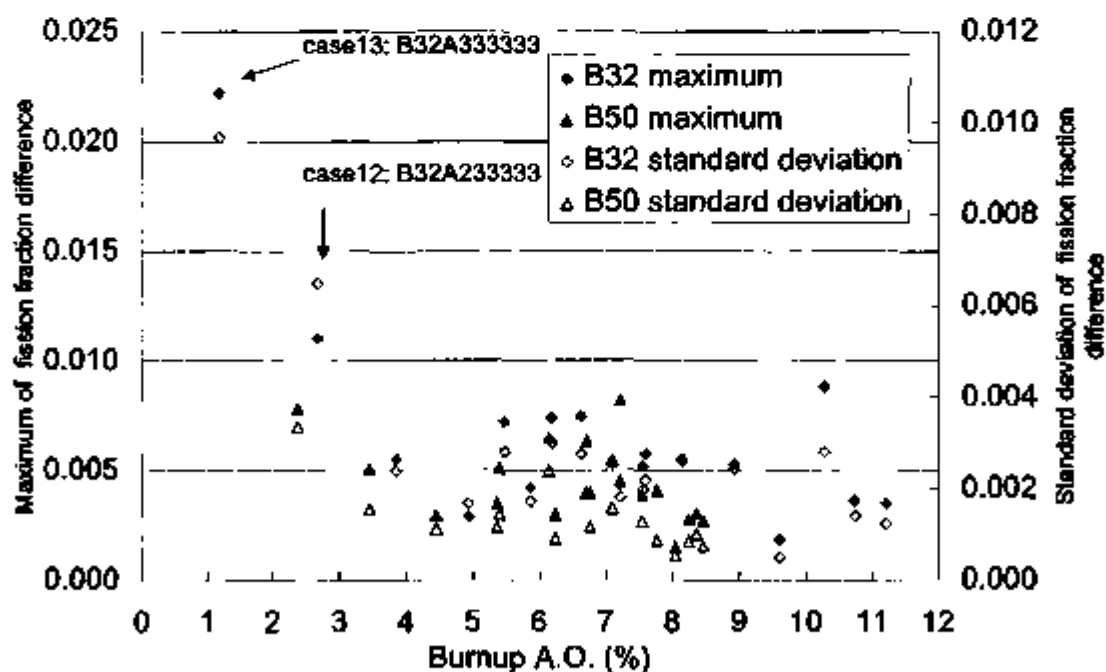
It is desirable to examine various calculation conditions for the asymmetric burnup cases in the same way as the previously mentioned cases of symmetric and uniform burnup. From the viewpoints of calculation time, we carry out calculations of only two different conditions such as 1,000 histories per cycle x 2000 scored cycles and 5,000 histories per cycle x 400 scored cycles. Both are with 300 skipped cycles. For most cases, the fission source distribution trends for the latter condition are slightly more stable than the former, and they reach near convergence. But it was found that there exist four cases of insufficiently converged fission source. As a function of the burnup A. O. (see subsection 2.2.3.), the maximum and the standard deviation for differences of fission source distribution between two conditions are shown in Fig. 26. In general, both conditions give almost the same result of fission source distribution except the insufficiently converged cases such as B32C05A233333 and B32C05A333333 for the both AC+FP and AC-only approaches as shown in Fig. 26. This is caused by their slightly asymmetric burnup. So, the skipped cycles are reexamined for them to reach convergence based on the cask configuration. The calculated results of the four cases for more than 300 skipped cycles are shown in Fig. 27 to Fig. 30, respectively. It is found that

the largest number of skipped cycles shown in each figure gives almost converged fission source distribution of the corresponding four cases. In order to check the fission source convergence of final results, the fission source trends per scored cycle of the four cases are shown in Fig. 31 to Fig. 34.

Therefore, the calculation conditions for all cases are finally determined as listed in Table 5.



(a) AC+FP approach



(b) AC-only approach

Fig. 26 Fission source differences between 2 conditions in cases of asymmetric burnup

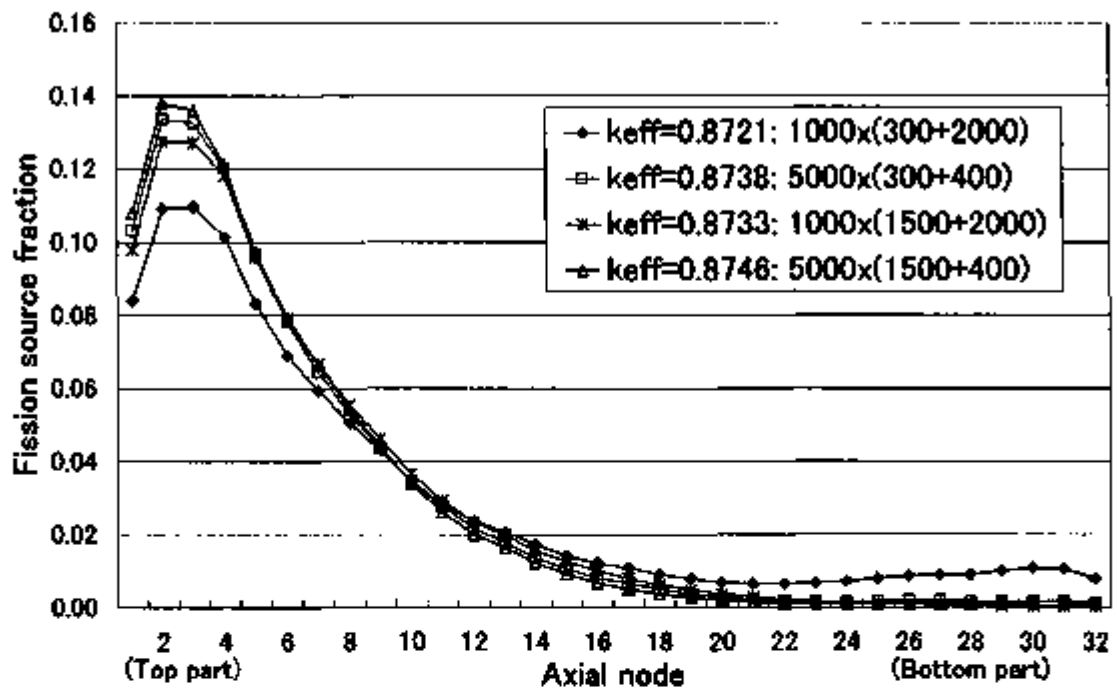


Fig. 27 Reexamination of calculation conditions for B32C05A233333 based on the AC+FP approach

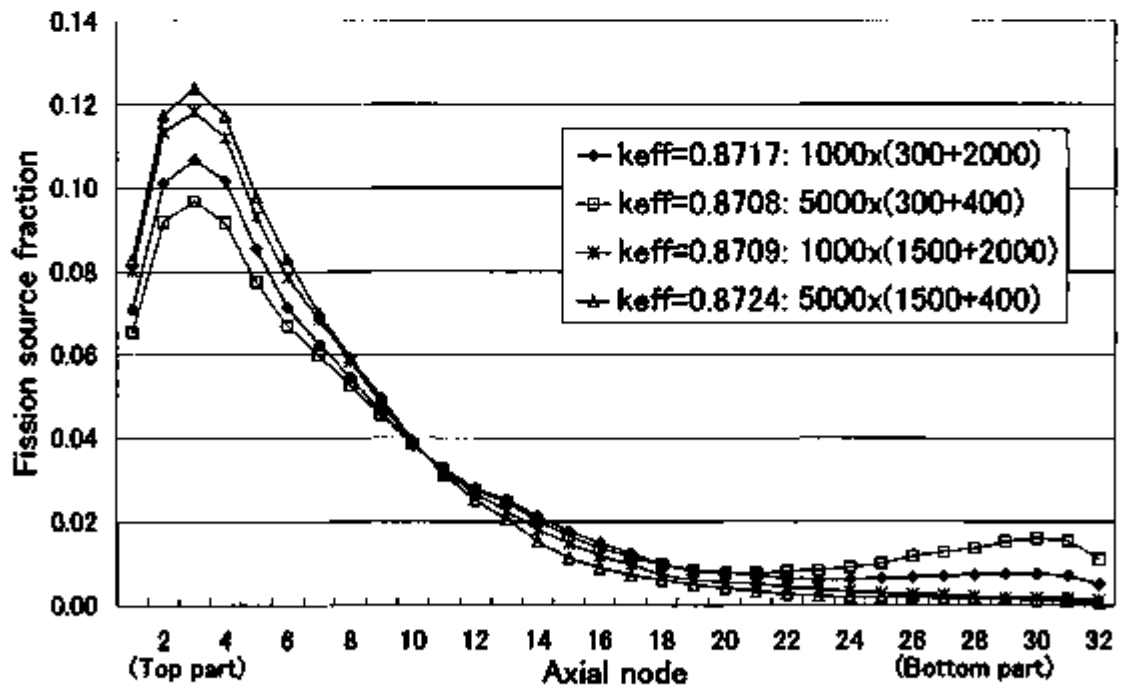


Fig. 28 Reexamination of calculation conditions for B32C05A333333 based on the AC+FP approach

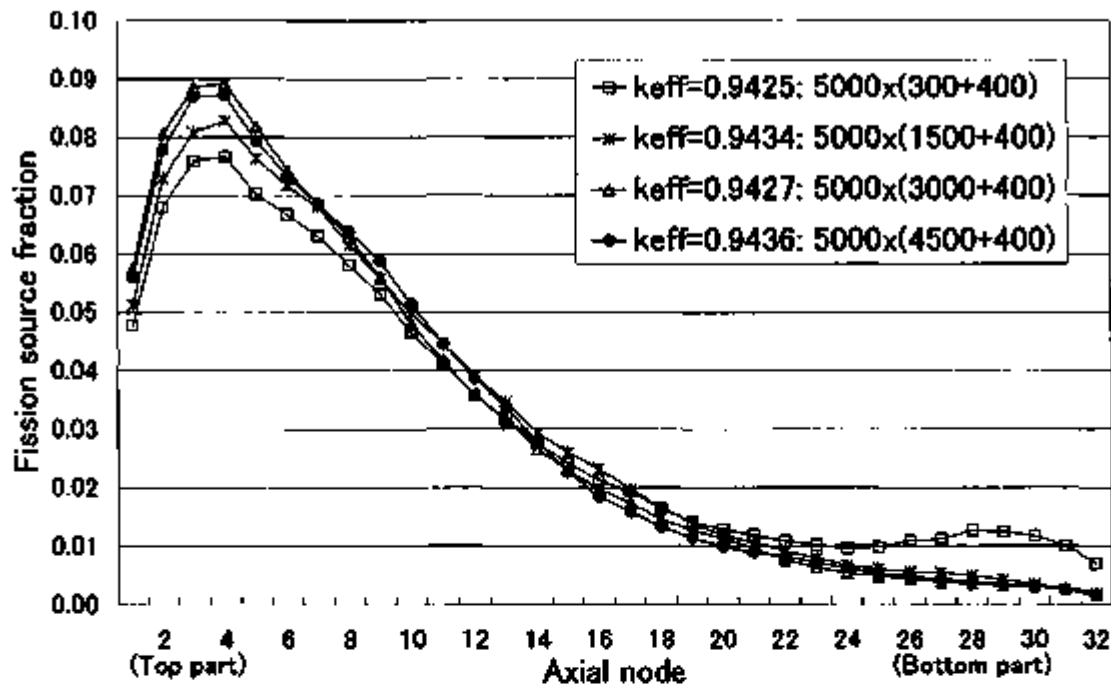


Fig. 29 Reexamination of calculation conditions for B32C05A233333
based on the AC-only approach

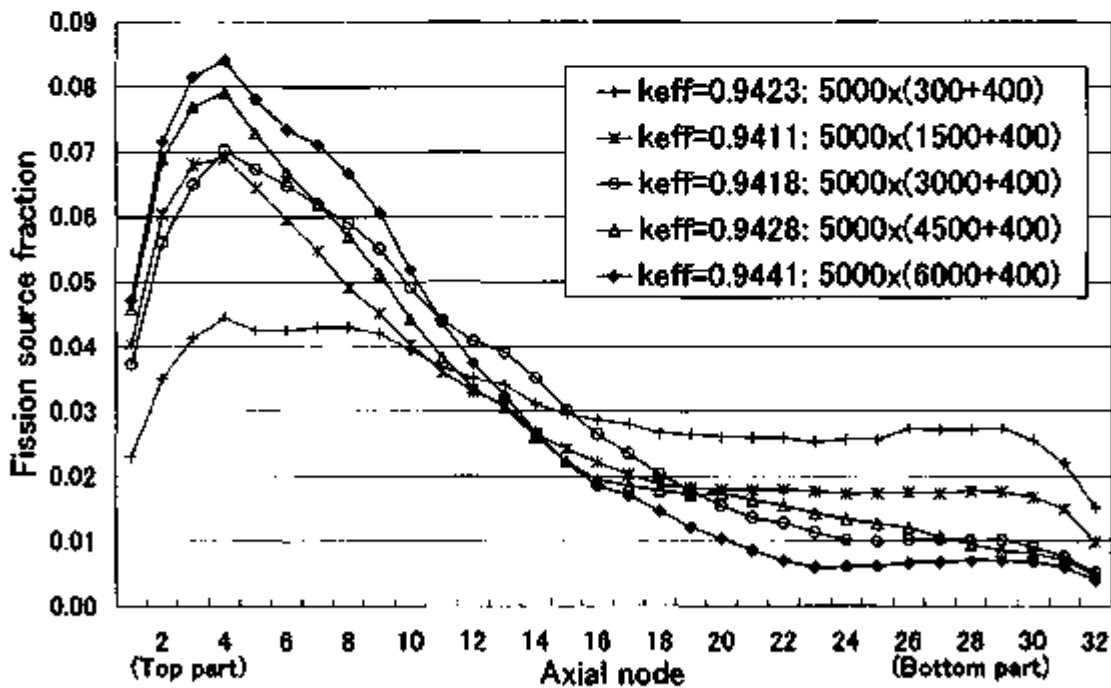


Fig. 30 Reexamination of calculation conditions for B32C05A333333
based on the AC-only approach

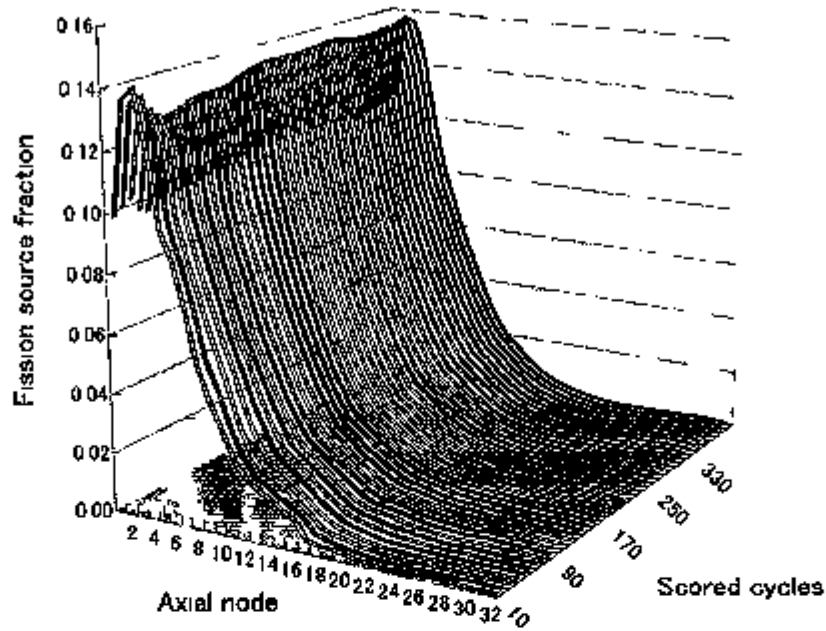


Fig. 31 Fission source distribution trend of B32C05A233333 in case of 1500 skipped cycles and 5000 histories per cycle, based on AC+FP approach

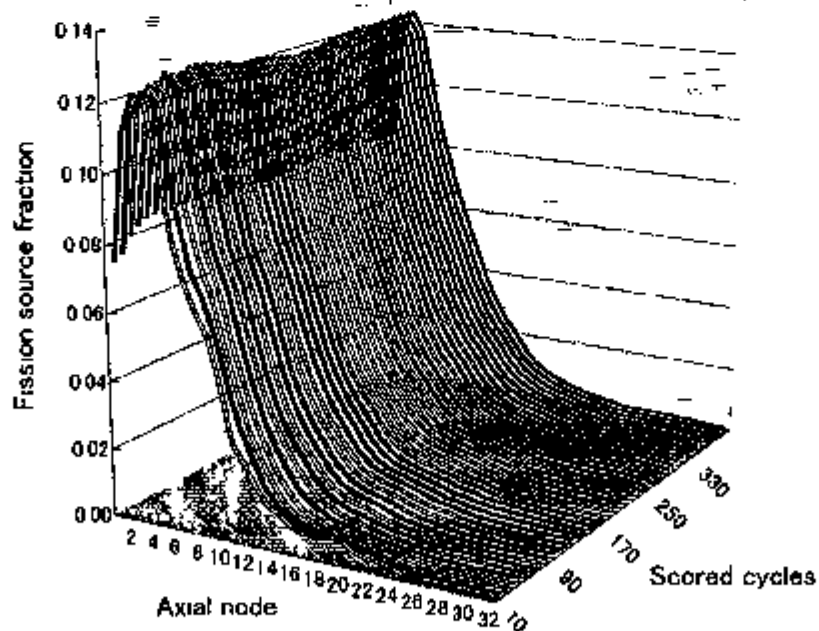


Fig. 32 Fission source distribution trend of B32C05A333333 in case of 1500 skipped cycles and 5000 histories per cycle, based on AC+FP approach

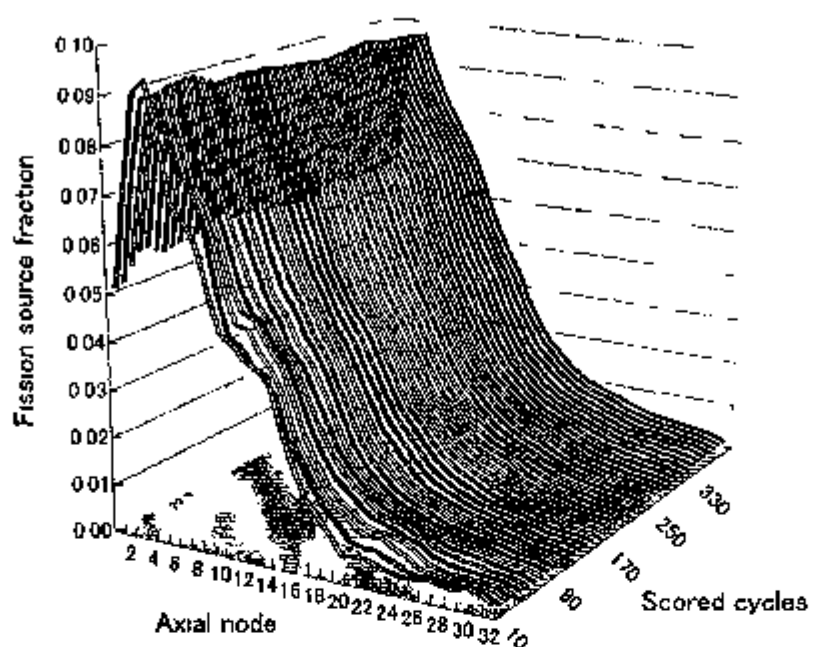


Fig 33 Fission source distribution trend of B32C05A233333 in case of 4500 skipped cycles and 5000 histories per cycle, based on AC-only approach

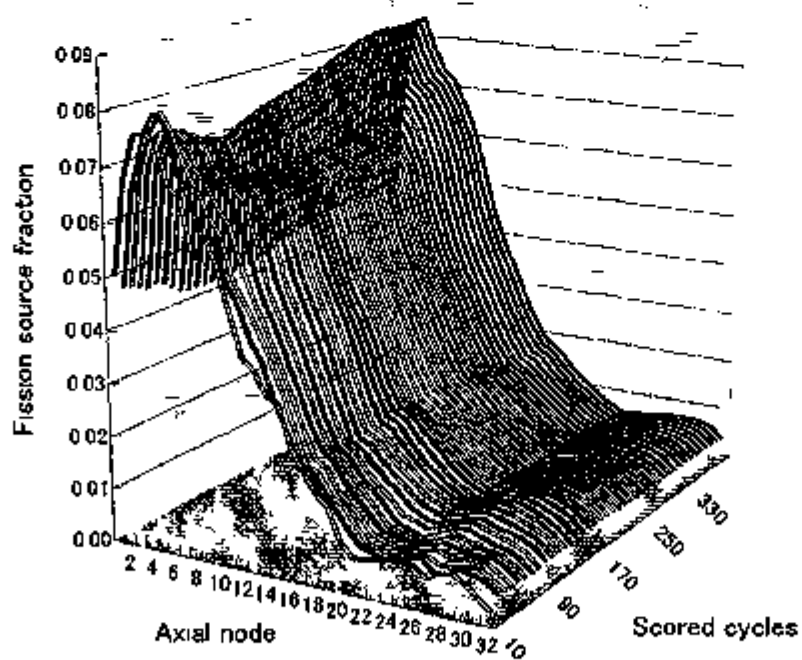


Fig 34 Fission source distribution trend of B32C05A333333 in case of 6000 skipped cycles and 5000 histories per cycle, based on AC-only approach

4.4. Final Results

4.4.1. k_{eff} and End Effect

Calculated k_{eff} results, one standard deviations and end effects for the final estimations are listed in Table 8. Here, the end effect is defined as a reactivity difference between the uniform burnup case and the others in consideration of the axial burnup profiles as formulated below.

$$\text{End effect } (\% \Delta k/k) = 100 \times \frac{k_{eff}(\text{axial burnup profile}) - k_{eff}(\text{uniform burnup})}{k_{eff}(\text{uniform burnup})}$$

Table 8 Calculated k_{eff} results for both AC+FP and AC-only approaches

Case ID	Serial No.	Burnup A.O. (%)	AC+FP approach			AC-only approach		
			k_{eff}	1σ	End effect (% $\Delta k/k$)	k_{eff}	1σ	End effect (% $\Delta k/k$)
B32C05Symmetric	--	0.00	0.86848	0.00051	-0.36	0.94071	0.00045	-0.53
B32C05Uniform	--	0.00	0.87166	0.00015	---	0.94571	0.00016	---
B32C05A222222	1	7.21	0.88971	0.00050	2.07	0.95218	0.00048	0.68
B32C05A111111	2	11.20	0.90013	0.00050	3.27	0.96120	0.00050	1.64
B32C05A111112	3	10.73	0.90042	0.00051	3.30	0.96001	0.00055	1.51
B32C05A111122	4	10.27	0.89959	0.00049	3.20	0.96077	0.00051	1.59
B32C05A111222	5	9.60	0.89709	0.00047	2.92	0.95874	0.00050	1.38
B32C05A112222	6	8.93	0.89390	0.00050	2.55	0.95695	0.00048	1.19
B32C05A122222	7	8.14	0.89102	0.00049	2.22	0.95332	0.00048	0.80
B32C05A222223	8	6.62	0.88674	0.00049	1.73	0.95265	0.00052	0.73
B32C05A222233	9	5.87	0.88630	0.00049	1.68	0.95158	0.00049	0.62
B32C05A222333	10	4.93	0.88410	0.00047	1.43	0.94881	0.00048	0.33
B32C05A223333	11	3.86	0.87746	0.00049	0.67	0.94411	0.00047	-0.17
B32C05A233333	12	2.68	0.87327	0.00047	0.18	0.94357	0.00052	-0.23
B32C05A333333	13	1.19	0.87093	0.00047	-0.08	0.94408	0.00049	-0.17
B32C05A122223	14	7.55	0.88986	0.00049	2.09	0.95247	0.00050	0.71
B32C05A112223	15	7.59	0.89203	0.00049	2.34	0.95273	0.00052	0.74
B32C05A322221	16	6.19	0.88671	0.00049	1.73	0.95300	0.00049	0.77
B32C05A332211	17	5.48	0.88283	0.00048	1.28	0.95254	0.00050	0.72
B50C05Symmetric	--	0.00	0.75590	0.00047	2.03	0.84242	0.00047	0.76
B50C05Uniform	--	0.00	0.74084	0.00013	---	0.83610	0.00014	---
B50C05A222222	18	7.21	0.78772	0.00046	6.33	0.86562	0.00047	3.53
B50C05A111111	19	8.46	0.79403	0.00046	7.18	0.87013	0.00049	4.07
B50C05A111112	20	8.35	0.79362	0.00046	7.12	0.87012	0.00050	4.07
B50C05A111122	21	8.24	0.79372	0.00045	7.14	0.87093	0.00049	4.17
B50C05A111222	22	8.04	0.79299	0.00044	7.04	0.86943	0.00047	3.99
B50C05A112222	23	7.77	0.79037	0.00047	6.69	0.86867	0.00048	3.90
B50C05A122222	24	7.54	0.78843	0.00047	6.42	0.86630	0.00045	3.61
B50C05A222223	25	6.76	0.78679	0.00047	6.20	0.86569	0.00048	3.54
B50C05A222233	26	6.15	0.78565	0.00043	6.05	0.86415	0.00047	3.35
B50C05A222333	27	5.41	0.78221	0.00045	5.58	0.86155	0.00050	3.04
B50C05A223333	28	4.45	0.77597	0.00043	4.74	0.85658	0.00050	2.45
B50C05A233333	29	3.46	0.76778	0.00046	3.64	0.85038	0.00050	1.71
B50C05A333333	30	2.38	0.76141	0.00046	2.78	0.84740	0.00046	1.35
B50C05A122223	31	7.09	0.78910	0.00046	6.51	0.86599	0.00046	3.57
B50C05A112223	32	6.70	0.79017	0.00044	6.66	0.86763	0.00049	3.77
B50C05A322221	33	6.24	0.78211	0.00044	5.57	0.86290	0.00048	3.21
B50C05A332211	34	5.36	0.77461	0.00047	4.56	0.85828	0.00044	2.65

The relation between k_{eff} and burnup A. O. is shown in Fig. 35 to Fig. 38. Here, the error bars show 3σ of k_{eff} . The fitting lines are obtained by the following equations as shown in Fig. 35 to Fig. 38.

$$k_{eff} [32\text{MWd/kgU, AC+FP approach}] = 0.0031 \times \text{Burnup A. O. (\%)} + 0.8672$$

$$k_{eff} [50\text{MWd/kgU, AC+FP approach}] = 0.0049 \times \text{Burnup A. O. (\%)} + 0.7531$$

$$k_{eff} [32\text{MWd/kgU, AC-only approach}] = 0.0019 \times \text{Burnup A. O. (\%)} + 0.9399$$

$$k_{eff} [50\text{MWd/kgU, AC-only approach}] = 0.0036 \times \text{Burnup A. O. (\%)} + 0.8406$$

As shown in Fig. 35 to Fig. 38, it is found that k_{eff} in case of 32 MWd/kgU decreases by considering the axial burnup profile at 0 burnup A. O. both for the AC+FP and AC-only approaches. On the other hand, in case of 50 MWd/kgU, k_{eff} increases at 0 burnup A. O. This is explained as follows. The consideration of axial burnup profile makes the fission source distribution to be localized in the top and the bottom parts. In case of low burnup, the localization of fission source distribution increases the negative reactivity of neutron leakage that is supposed to be major effect. At 32 MWd/kgU for symmetric burnup profile, total reactivity is somewhat negative. With increase of average burnup, the burnup difference between the end part of low burnup and the middle part of high burnup increases, and the localization of fission source distribution to the low burnup is also enhanced. Therefore, the increase of average burnup contributes to the positive reactivity effect. In a certain point of average burnup more than 32 MWd/kgU, the positive reactivity effect seems to be greater than the negative one. Hence, the total reactivity at 50 MWd/kgU for symmetric burnup profile becomes positive.

The burnup difference between the top part of highest burnup and the other part increases with the increase of burnup A. O., namely, with increase of the axial asymmetry of burnup profile, and the localization of fission source distribution is also enhanced. So, k_{eff} increases almost linearly with the increase of burnup A. O. at both 32 and 50 MWd/kgU. The discrepancy between k_{eff} results and the fitting line seems to be large in comparison with 3σ of k_{eff} . One of the reasons is supposed to be that the burnup shape inside the upper part affects the k_{eff} results because the burnup A.O. represents the region averaged burnup difference, but does not reflect the fine burnup shape inside the top part.

It is obviously seen that the AC-only approach is conservative in comparison with the AC+FP approach. The k_{eff} margin at both 32 and 50 MWd/kgU is at least 4 % $\Delta k/k$.

The relation between the end effect and the burnup A. O. is shown in Fig. 39. The fitting lines are also shown in Fig. 39. The same tendency as k_{eff} is observed in Fig. 39. Both the increase of average burnup and the consideration of FP contribute to the increase of not only the end effect itself but also the tendency of dependence on the burnup A. O.

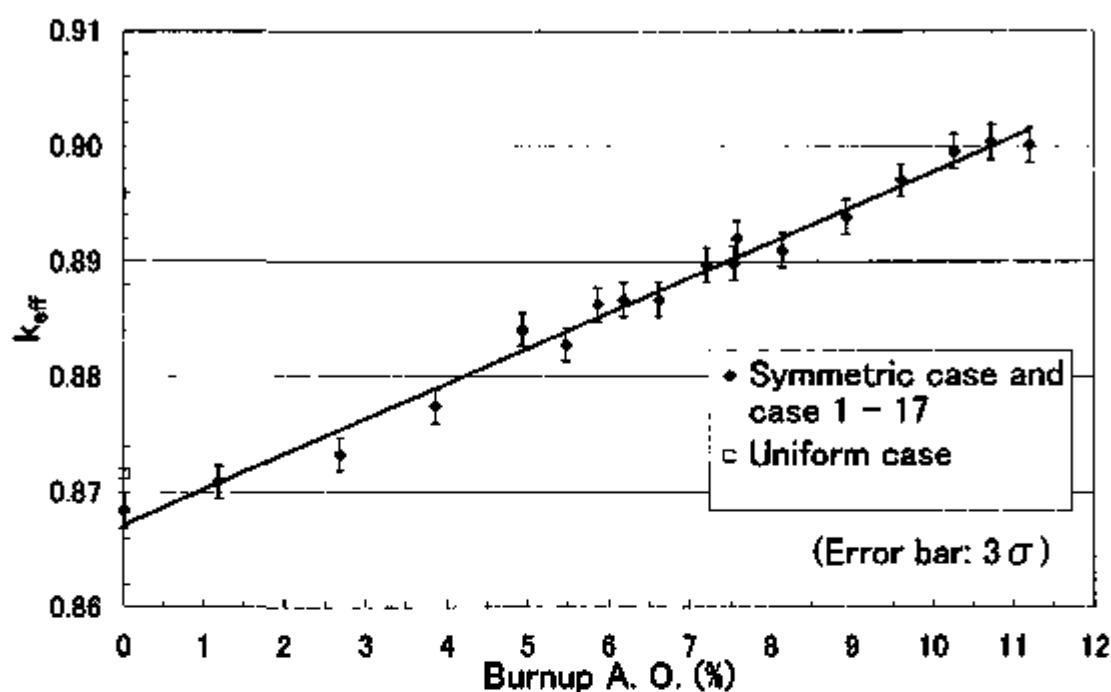


Fig. 35 k_{eff} vs. burnup A. O. at 32 MWd/kgU based on AC+FP approach

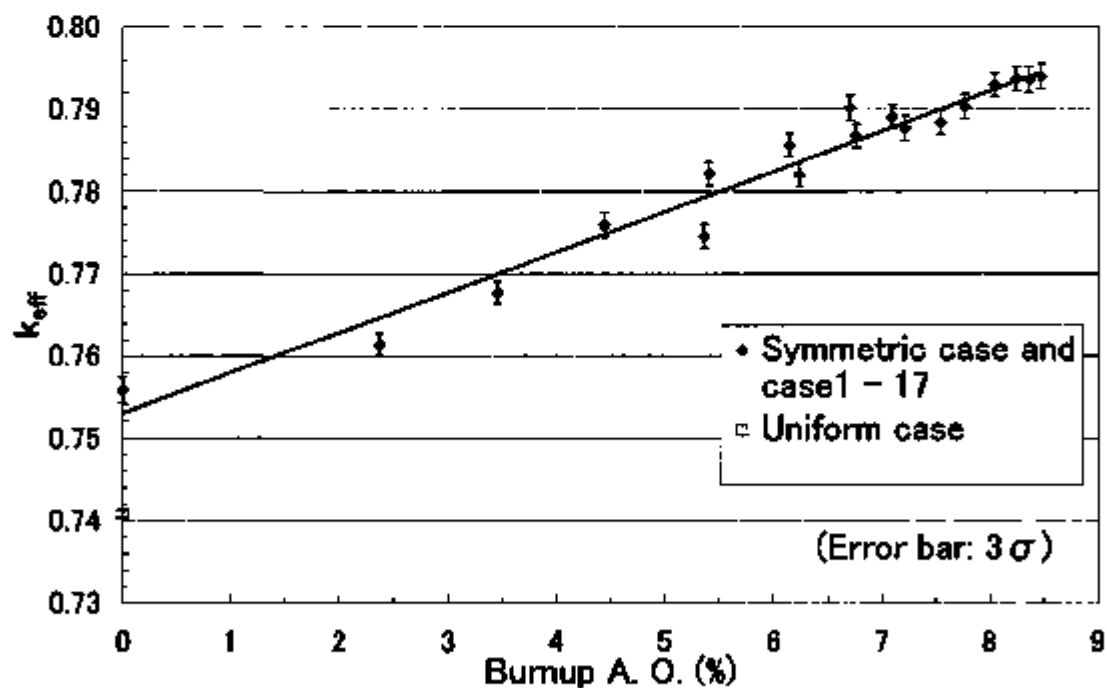


Fig. 36 k_{eff} vs. burnup A. O. at 50 MWd/kgU based on AC+FP approach

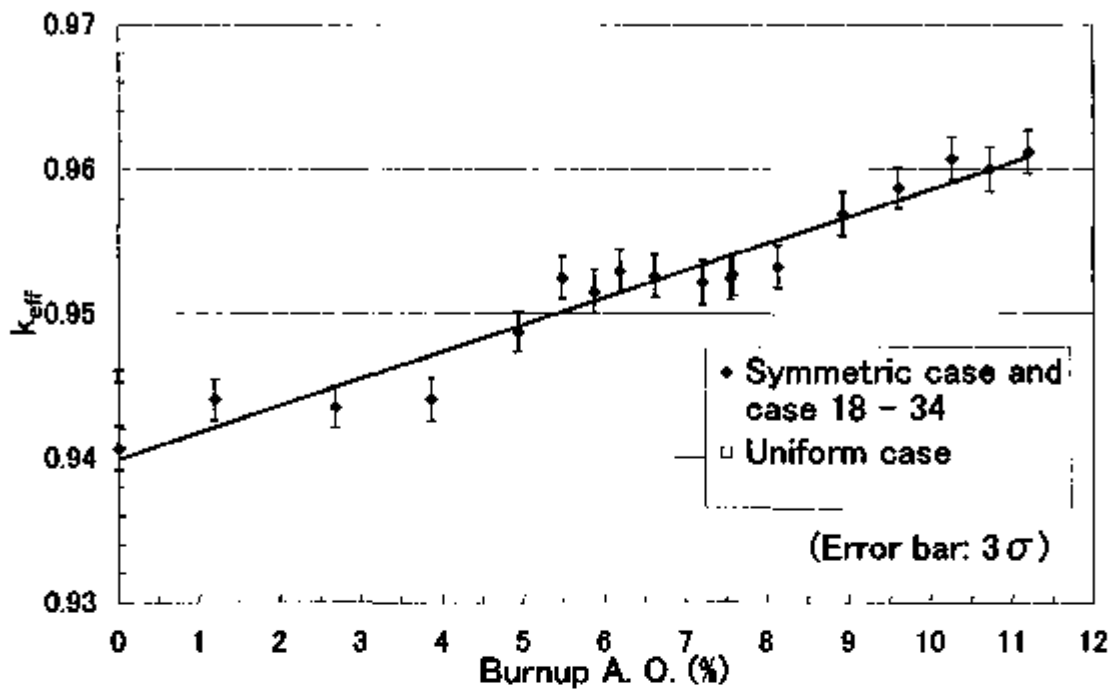


Fig. 37 k_{eff} vs. burnup A. O. at 32 MWd/kgU based on AC-only approach

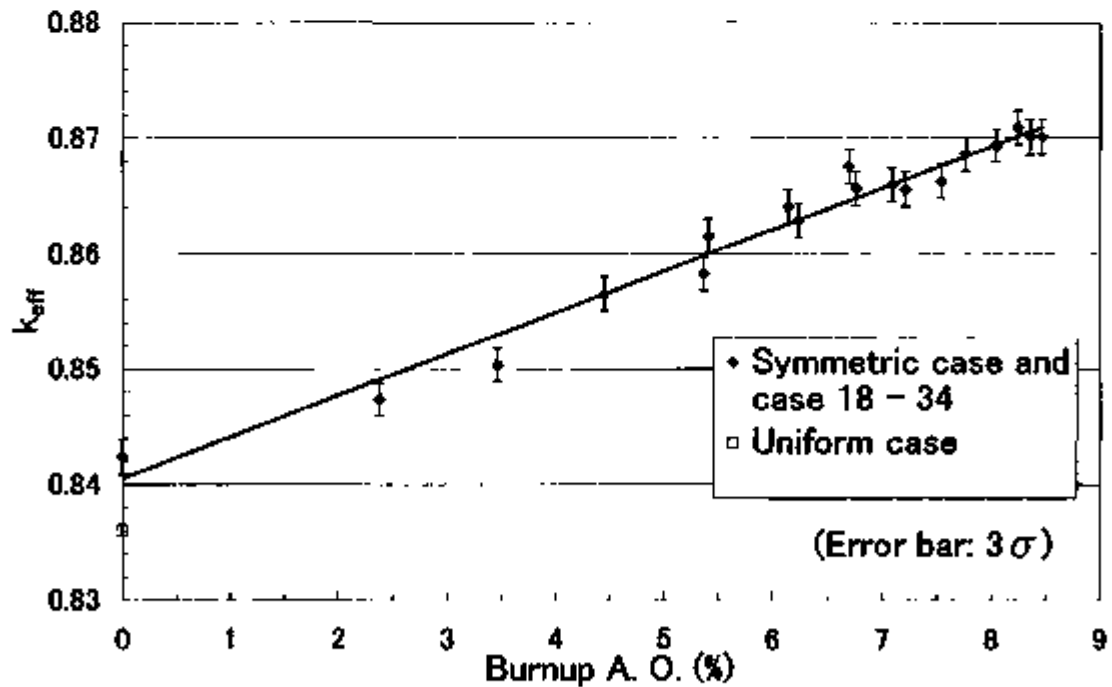


Fig. 38 k_{eff} vs. burnup A. O. at 50 MWd/kgU based on AC-only approach

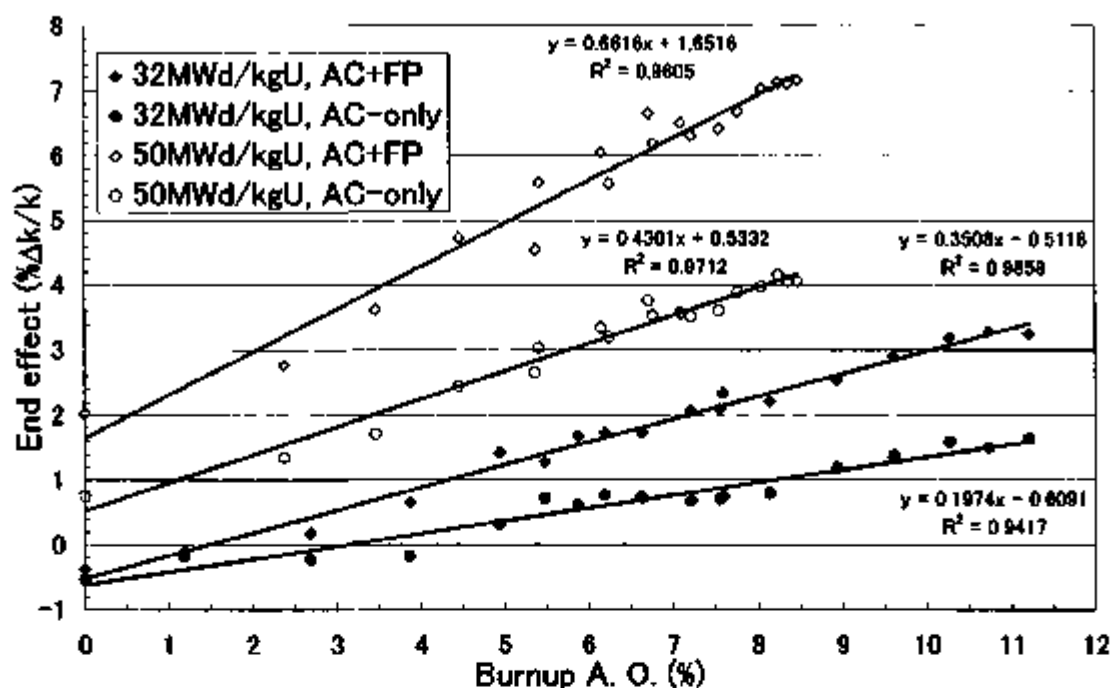


Fig. 39 End effect vs. burnup A. O.

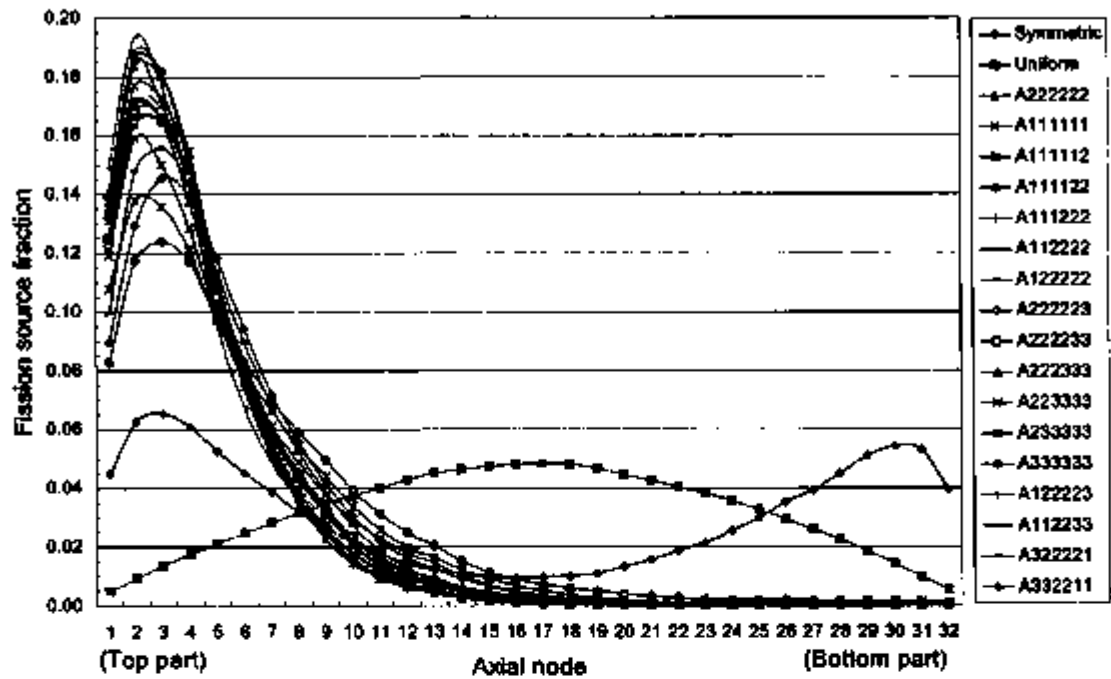
4.4.2. Fission Source Distribution

The calculated results of the fission source distribution for all cases are shown in Fig. 40 to Fig. 43 with both linear and semi-log scales. We cannot predict the very low source levels at the bottom end with enough reliability. Nonetheless, the precision of estimated k_{eff} 's and fission source distributions are acceptable. So, we neglect the fission source fluctuation at the bottom end in judging convergences of calculations. It is seen that the increase of average burnup and the burnup A. O., and the consideration of FP enhance the localization of the fission source distribution at the upper part and increase the maximum value. In order to illustrate this tendency, we introduce the fission A. O. defined as follows.

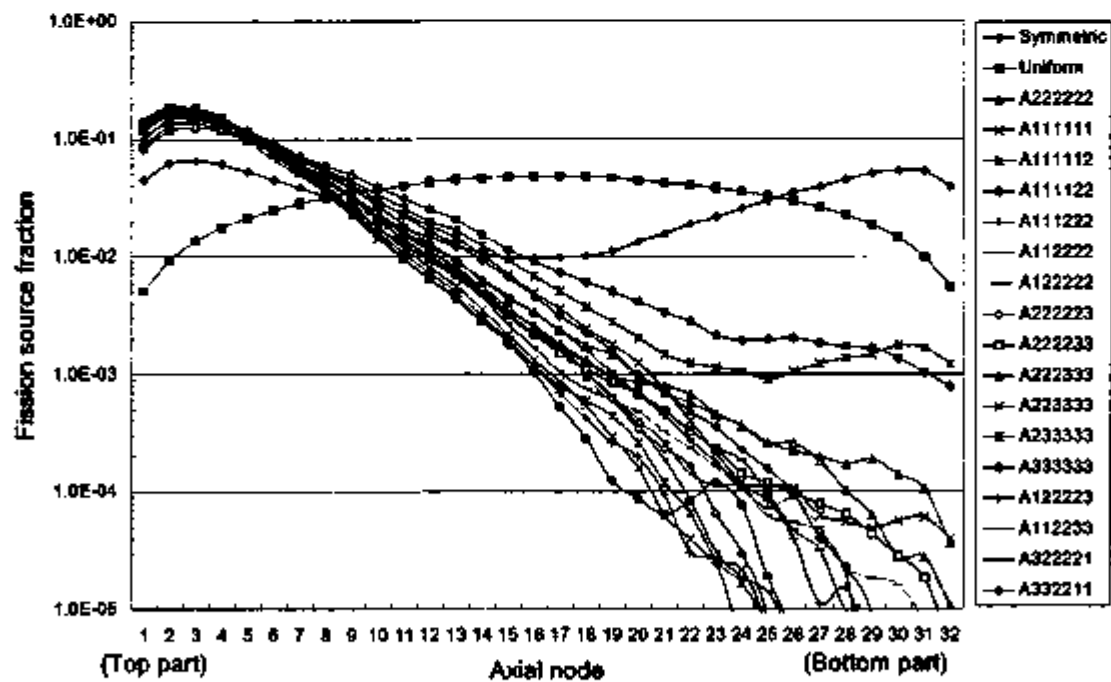
$$\text{Fission A. O.} = \frac{(\text{Fission source average for upper 16 nodes} - \text{that for lower 16 nodes})}{\text{fission source average for all 32 nodes}}$$

The calculated results of the fission A. O. as a function of burnup A. O. are shown in Fig. 44 and Fig. 45. The fission A. O. increases with the increase of the burnup A. O. as shown in Fig. 44 and Fig. 45. The tendency of the fission source distribution is different from that of k_{eff} . In case of strong asymmetry of burnup profile, the fission A. O. is very large, and it is supposed to change drastically toward zero, namely as approaching to symmetric shape with a low value of burnup A. O.

The above observation shows that the localization of fission source distribution could happen drastically as the asymmetry degree increases in burnup profile.

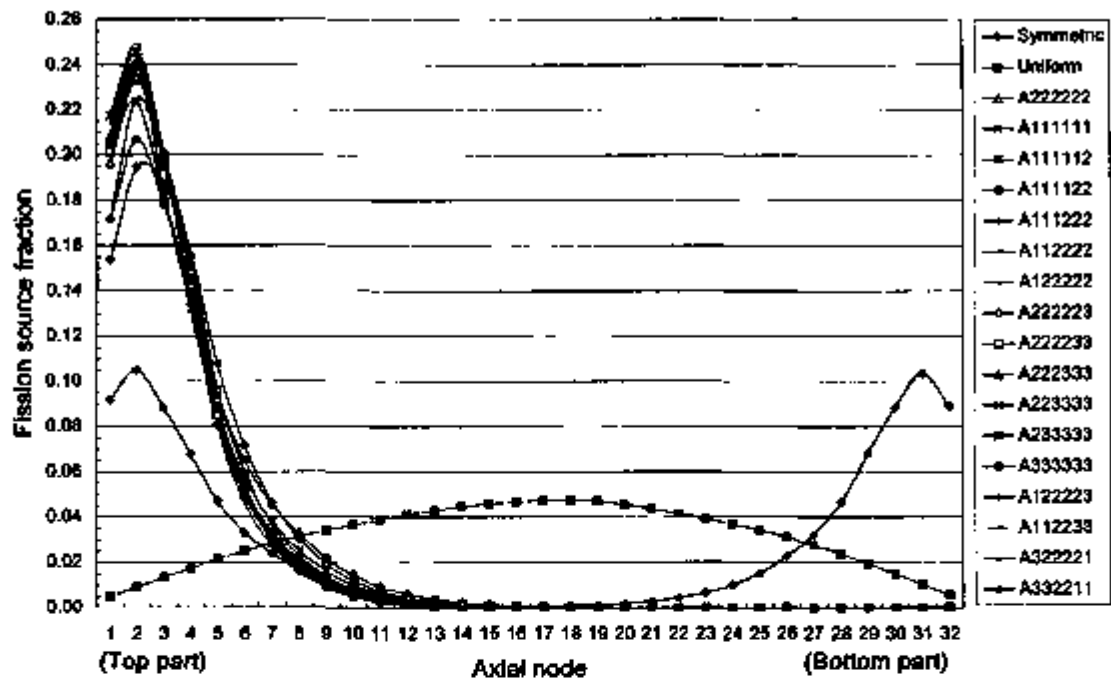


(a) Linear

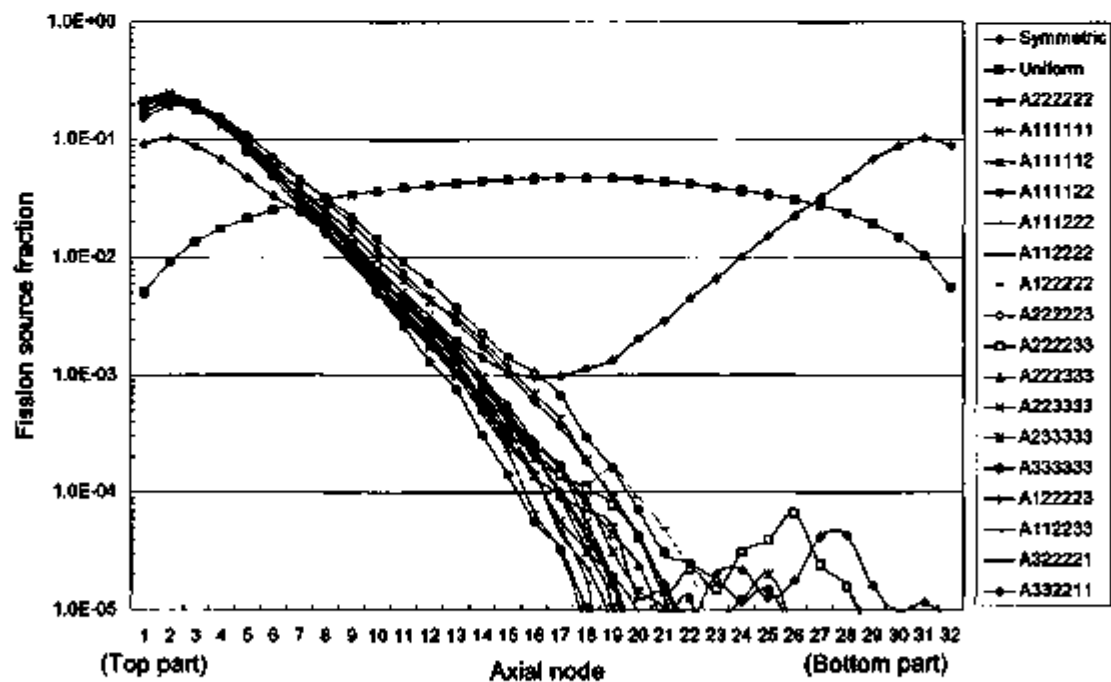


(b) Semi-log

Fig. 40 Calculated results of the fission source distribution at 32 MWd/kgU
based on AC+FP approach

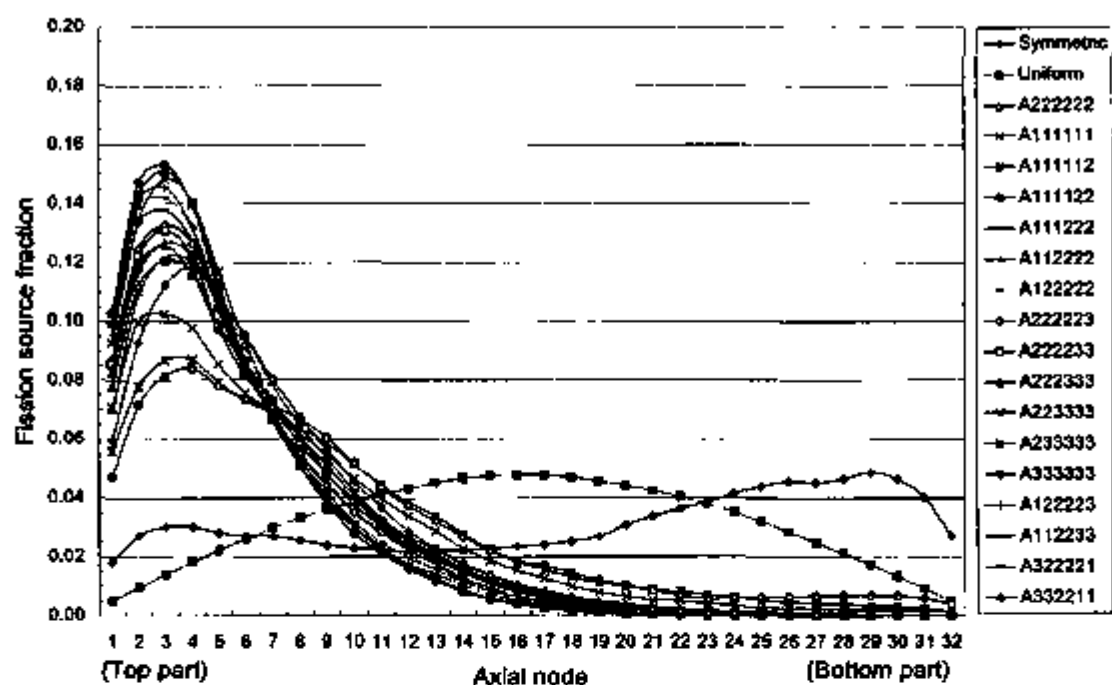


(a) Linear

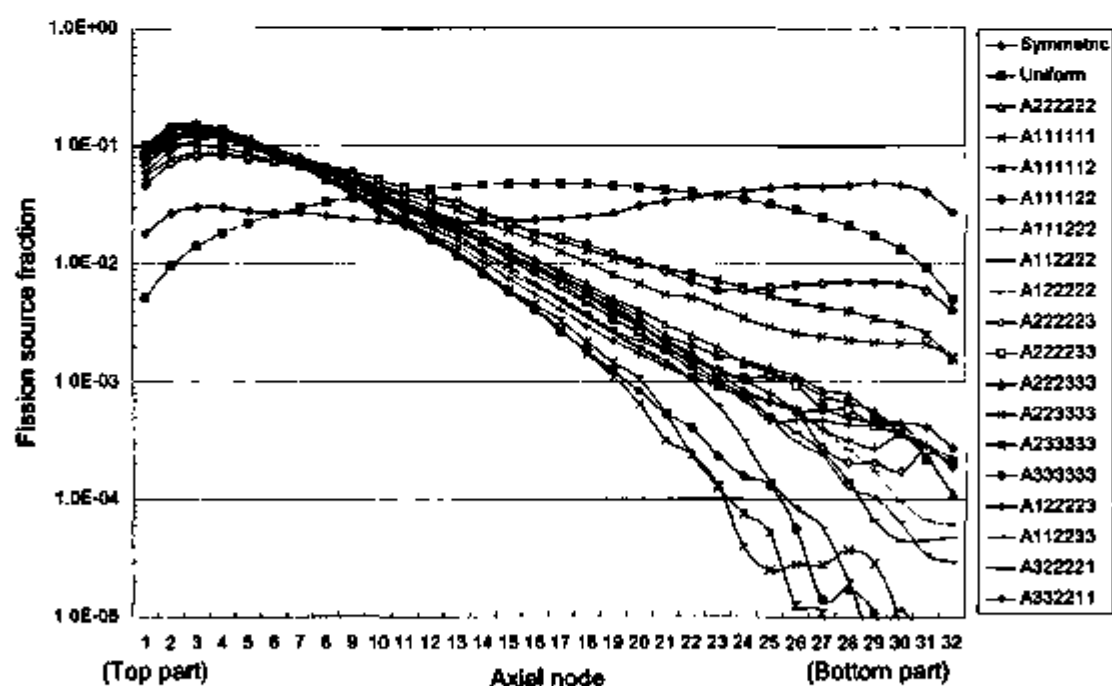


(b) Semi-log

Fig. 41 Calculated results of the fission source distribution at 50 MWd/kgU
based on AC+FP approach

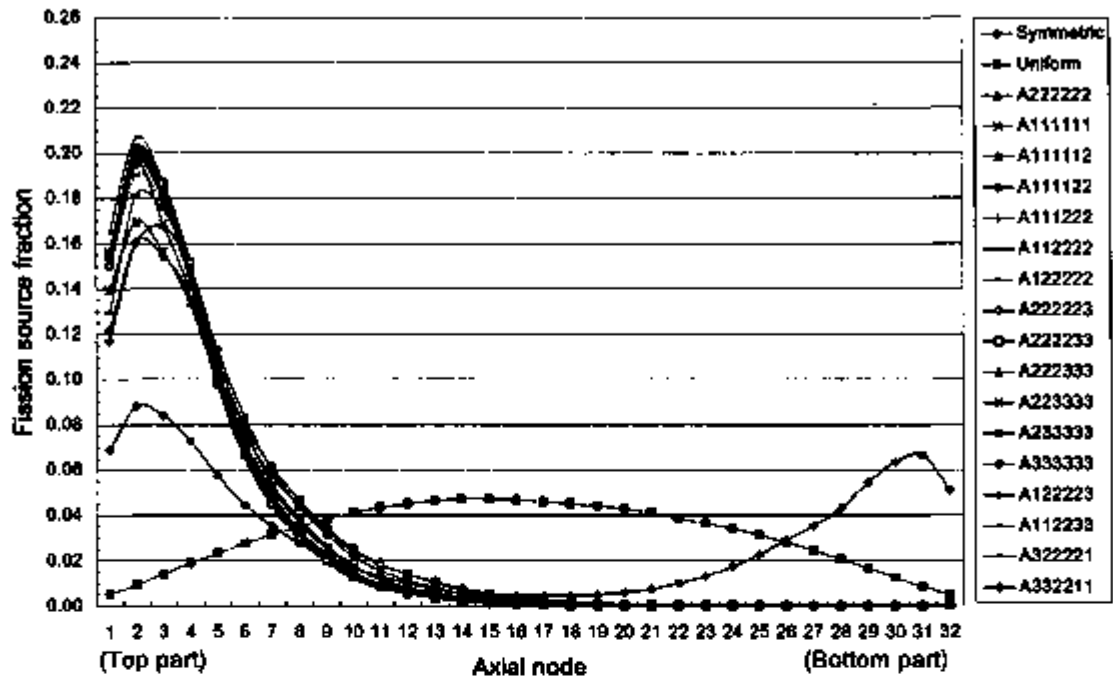


(a) Linear

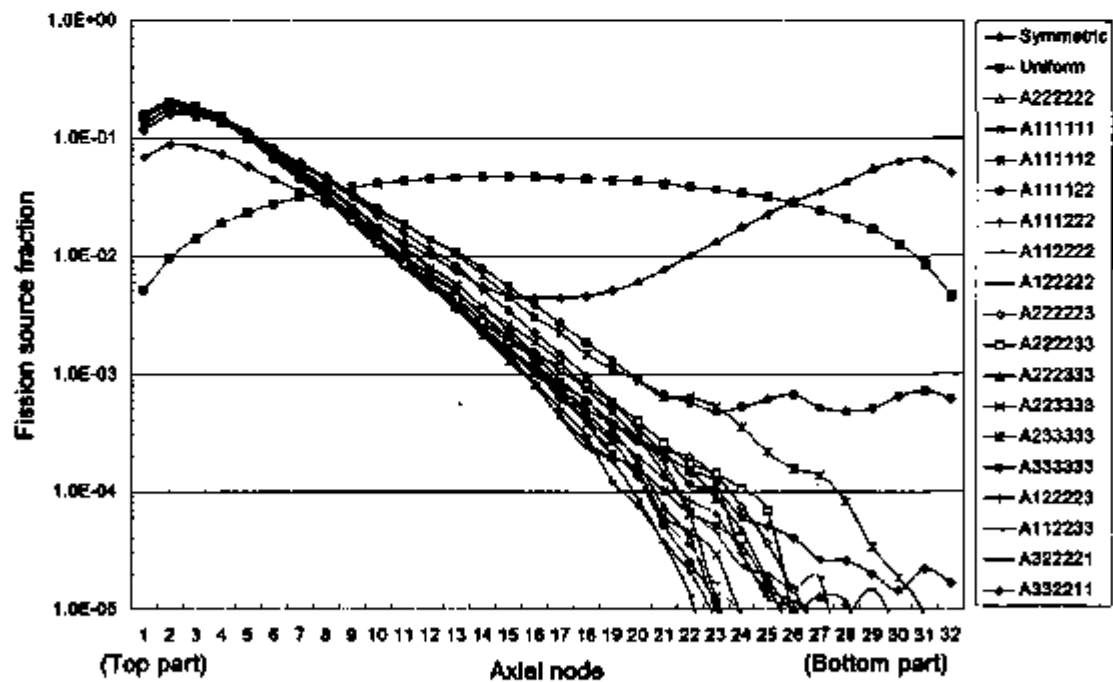


(b) Semi-log

Fig. 42 Calculated results of the fission source distribution at 32 MWd/kgU
based on AC-only approach



(a) Linear



(b) Semi-log

Fig. 43 Calculated results of the fission source distribution at 50 MWd/kgU
based on AC-only approach

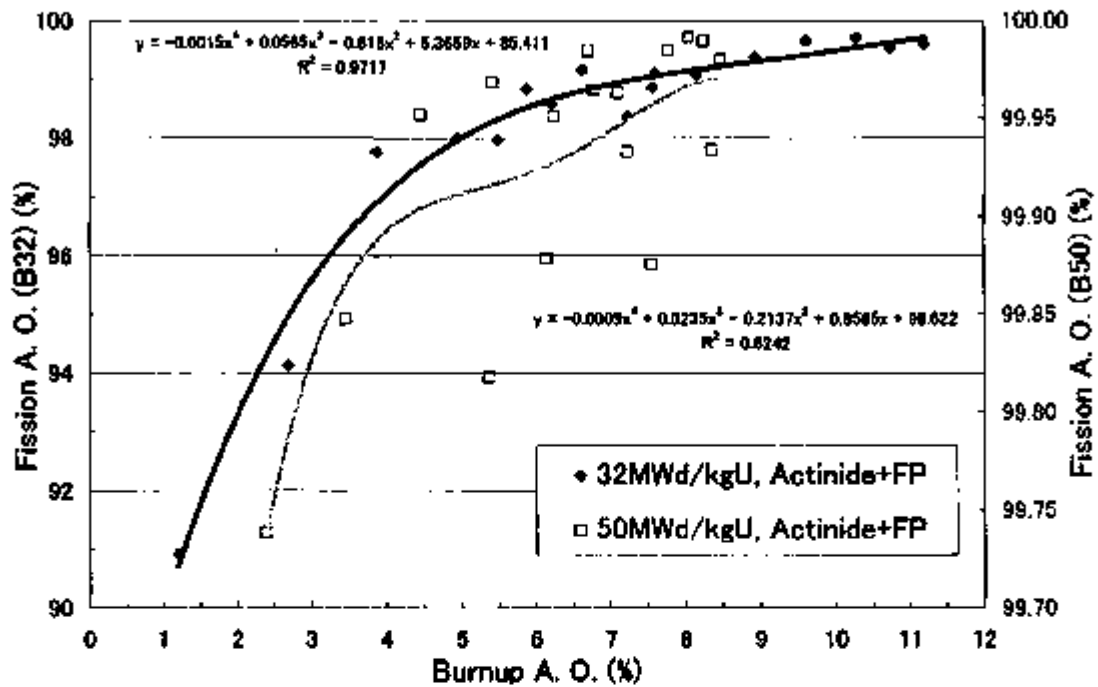


Fig. 44 Fission A. O. vs. burnup A. O. based on AC+FP approach

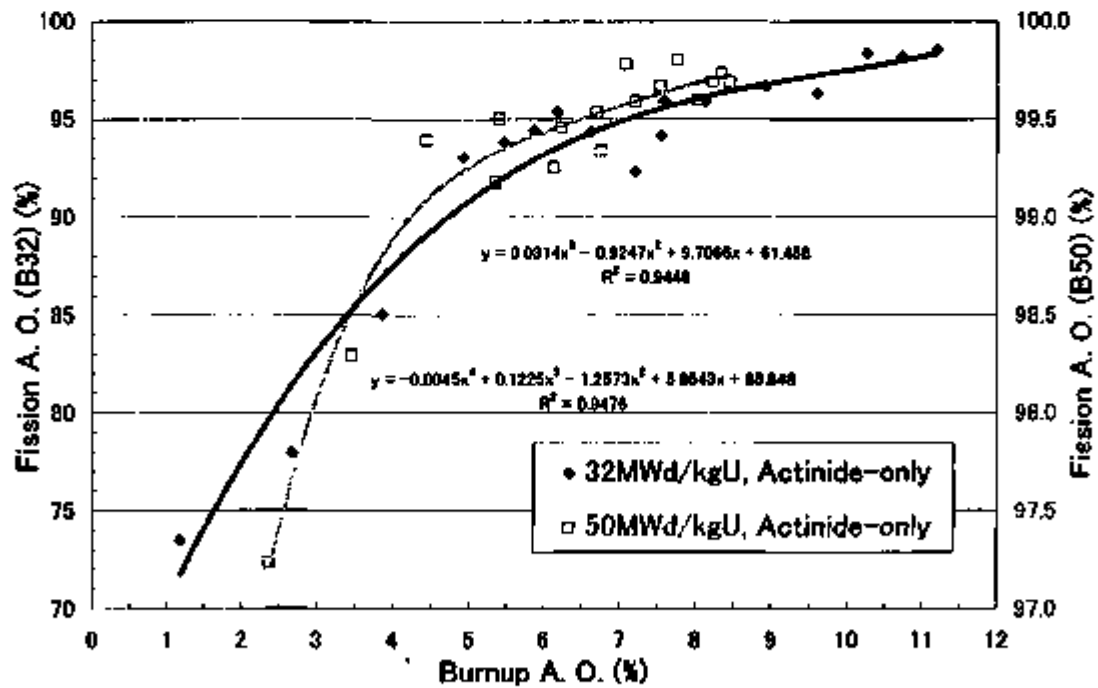


Fig. 45 Fission A. O. vs. burnup A. O. based on AC-only approach

4.4.3. Comparison of Calculation Time

The calculation time depends on not only the number of histories but also the calculation model and printout parameters to be estimated. The representative cases for the AC+FP approach adopted in this report are listed in Table 9. All calculation times listed in Table 9 are actual results for one case performed on the SUN/BLADE100 (Solaris8) machine. In case of the AC-only approach, it isn't necessary to read 15 FP nuclides data from the cross section library. So, the calculation time is reduced to about 80% of the AC+FP approach.

Table 9 Calculation time per execution for representative cases

Configuration	Numbers of history*	Calculation time
Pin cell	20000 x (300 + 100)	7.7 hr
Single assembly	20000 x (150 + 100)	6.5 hr
Color set	10000 x (300 + 200)	9.5 hr
Cask	20000 x (300 + 100)	10.0 hr
	5000 x (300 + 4000)	29.6 hr
	5000 x (300 + 400)	4.7 hr
	5000 x (1500 + 400)	12.1 hr
	5000 x (4500 + 400)	25.1 hr**
	5000 x (6000 + 400)	32.2 hr**
	1000 x (300 + 2000)	3.2 hr

*: Number of histories per cycle x (number of skipped cycles + number of scored cycles)

** : Actual results for the AC-only approach (calculations for AC+FP is not carried out)

5. Conclusion

Our study has been based on the current OECD/NEA Phase II-C benchmark problem for PWR spent fuel transport cask, which consists of a realistic set of axial burnup profiles for average burnups of 32 MWd/kgU and 50 MWd/kgU. The material compositions in each of the 21 non-equidistant material zones are provided in the benchmark in order to simulate the axial burnup profiles. Criticality calculations of 3-D model are carried out based not only on the actinide plus fission product approach specified in the benchmark, but also actinide-only approach as an extension of the Phase II-C, by using the continuous energy Monte Carlo code MCNP-4B2 with the cross section library JENDL-3.2. And the sensitivities of k_{eff} , end effect and fission source distribution to the asymmetry of the axial burnup profile are evaluated with introducing the burnup axial offset.

The adequacy of Monte Carlo calculation conditions is examined by the convergence check for the fission source distribution. In some cases recognized as very slow fission source convergence, the calculation conditions are reexamined. As a result of the final estimation, a general tendency is observed such that the end effect increases linearly with the increase of axial burnup profile asymmetry. The end effect is found to be more sensitive to the asymmetry of axial burnup profile as averaged burnup increases.

Even for slightly asymmetric axial burnup profiles, the axial fission source distribution becomes strongly asymmetric such that the peak shifts toward the top end of the fuel active zone where the local burnup is less than that of the bottom end. The peak of fission source distribution becomes higher with the increase of the asymmetry of burnup profile as well as the assembly-averaged burnup.

As far as conservatism of burnup credit criticality is concerned, the following conclusions can be drawn. The actinide plus fission product approach should be adopted carefully because an axially distributed burnup profile strongly affects the calculated k_{eff} so that the upper limit of axial asymmetry should be assumed from the real situation of burnup profile. On the other hand, the actinide-only approach assuming uniform burnup profile is still conservative with a reactivity margin of at least 4 % $\Delta k/k$ as a result of the Phase II-C benchmark problem. The trade-off between simplicity by assuming the above approach and its less incentive to an introduction of burnup credit, regarding the reactivity of fission products as a safety margin, should be further considered from the economical point of view.

References

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JENDL-3.2 can be referred to the website: <http://www.ndc.tokai.jaeri.go.jp/jendl/j32/j32.html>
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Appendix I

Sample Input of MCNP-4B2 Code for the OECD/NEA Phase II-C Burnup Credit Criticality Benchmark

This is a blank page.

41	cz	105.00	\$	Cash potter diameter	200 cm
4					
c					
c Equidistant 32 nodes in fuel active length					
101	pa	0.0	\$	bottom of active length	
101	pa	12.1875	\$		
102	pa	24.3750	\$		
103	pa	36.5625	\$		
104	pa	48.7500	\$		
105	pa	60.9375	\$		
106	pa	73.1250	\$		
107	pa	85.3125	\$		
108	pa	97.5000	\$		
109	pa	109.6875	\$		
110	pa	121.8750	\$		
111	pa	134.0625	\$		
112	pa	146.2500	\$		
113	pa	158.4375	\$		
114	pa	170.6250	\$		
115	pa	182.8125	\$		
116	pa	195.0000	\$		
117	pa	207.1875	\$		
118	pa	219.3750	\$		
119	pa	231.5625	\$		
120	pa	243.7500	\$		
121	pa	255.9375	\$		
122	pa	268.1250	\$		
123	pa	280.3125	\$		
124	pa	292.5000	\$		
125	pa	304.6875	\$		
126	pa	316.8750	\$		
127	pa	329.0625	\$		
128	pa	341.2500	\$		
129	pa	353.4375	\$		
130	pa	365.6250	\$		
131	pa	377.8125	\$		
132	pa	390.0	\$	top of active length	
u					
u Lower structure					
201	pa	-51.0	\$	bottom of core	
202	pa	-51.0	\$	bottom of water gap (Dcm)	
203	pa	-51.0	\$	bottom of lower hardware	
204	pa	-32.4	\$	bottom of fuel pin length	
u					
u Upper structure					
205	pa	407.1	\$	top of fuel pin length	
206	pa	437.1	\$	top of upper hardware	
207	pa	464.1	\$	top of water gap	
208	pa	474.1	\$	top of core	
h					
c					
d					
d Data cards					
h					
h Isotopic data for case 1 032C05A222222					
032235.37c		5.06830E-04			
032234.37c		7.30473E-04			
032236.37c		1.51404E-03			
032237.37c		4.18860E-03			
032238.37c		2.16017E-02			
042234.37c		1.04042E-07			
042236.37c		1.08028E-04			
042240.37c		2.67140E-05			
042417.37c		1.04105E-05			
042418.37c		1.04105E-05			
052417.37c		3.04320E-05			
052418.37c		1.04105E-07			
052419.37c		2.75000E-05			
052420.37c		2.75000E-05			
052421.37c		2.75000E-05			
052422.37c		2.75000E-05			
052423.37c		2.75000E-05			
052424.37c		2.75000E-05			
052425.37c		2.75000E-05			
052426.37c		2.75000E-05			
052427.37c		2.75000E-05			
052428.37c		2.75000E-05			
052429.37c		2.75000E-05			
052430.37c		2.75000E-05			
052431.37c		2.75000E-05			
052432.37c		2.75000E-05			
052433.37c		2.75000E-05			
052434.37c		2.75000E-05			
052435.37c		2.75000E-05			
052436.37c		2.75000E-05			
052437.37c		2.75000E-05			
052438.37c		2.75000E-05			
052439.37c		2.75000E-05			
052440.37c		2.75000E-05			
052441.37c		2.75000E-05			
052442.37c		2.75000E-05			
052443.37c		2.75000E-05			
052444.37c		2.75000E-05			
052445.37c		2.75000E-05			
052446.37c		2.75000E-05			
052447.37c		2.75000E-05			
052448.37c		2.75000E-05			
052449.37c		2.75000E-05			
052450.37c		2.75000E-05			
052451.37c		2.75000E-05			
052452.37c		2.75000E-05			
052453.37c		2.75000E-05			
052454.37c		2.75000E-05			
052455.37c		2.75000E-05			
052456.37c		2.75000E-05			
052457.37c		2.75000E-05			
052458.37c		2.75000E-05			
052459.37c		2.75000E-05			
052460.37c		2.75000E-05			
052461.37c		2.75000E-05			
052462.37c		2.75000E-05			
052463.37c		2.75000E-05			
052464.37c		2.75000E-05			
052465.37c		2.75000E-05			
052466.37c		2.75000E-05			
052467.37c		2.75000E-05			
052468.37c		2.75000E-05			
052469.37c		2.75000E-05			
052470.37c		2.75000E-05			
052471.37c		2.75000E-05			
052472.37c		2.75000E-05			
052473.37c		2.75000E-05			
052474.37c		2.75000E-05			
052475.37c		2.75000E-05			
052476.37c		2.75000E-05			
052477.37c		2.75000E-05			
052478.37c		2.75000E-05			
052479.37c		2.75000E-05			
052480.37c		2.75000E-05			
052481.37c		2.75000E-05			
052482.37c		2.75000E-05			
052483.37c		2.75000E-05			
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052506.37c		2.75000E-05			
052507.37c		2.75000E-05			
052508.37c		2.75000E-05			
052509.37c		2.75000E-05			
052510.37c		2.75000E-05			
052511.37c		2.75000E-05			
052512.37c		2.75000E-05			
052513.37c		2.75000E-05			
052514.37c		2.75000E-05			
052515.37c		2.75000E-05			
052516.37c		2.75000E-05			
052517.37c		2.75000E-05			
052518.37c		2.75000E-05			
052519.37c		2.75000E-05			
052520.37c		2.75000E-05			
052521.37c		2.75000E-05			
052522.37c		2.75000E-05			
052523.37c		2.75000E-05			
052524.37c		2.75000E-05			
052525.37c		2.75000E-05			
052526.37c		2.75000E-05			
052527.37c		2.75000E-05			
052528.37c		2.75000E-05			
052529.37c		2.75000E-05			
052530.37c		2.75000E-05			
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052565.37c		2.75000E-05			
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052567.37c		2.75000E-05			
052568.37c		2.75000E-05			
052569.37c		2.75000E-05			
052570.37c		2.75000E-05			
052571.37c		2.75000E-05			
052572.37c		2.75000E-05			
052573.37c		2.75000E-05			
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m8	§	Neologus data for case 1 B32C08A222222	92234 37c	5 00610E-00
			92235 37c	1 11400E-04
			92237 37c	1 43701E-02
			92238 37c	2 14071E-02
			92239 37c	3 43600E-05
			92240 37c	1 40600E-04
			92241 37c	4 11700E-00
			92242 37c	2 35160E-05
			92243 37c	3 72440E-00
			92244 37c	7 04622E-00
			92245 37c	1 70900E-00
			92246 37c	4 80957E-05
			92247 37c	4 81600E-05
			92248 37c	4 49940E-00
			92249 37c	2 80600E-05
			92250 37c	3 14600E-04
			92251 37c	4 98130E-05
			92252 37c	3 44800E-00
			92253 37c	2 84400E-00
			92254 37c	4 63210E-05
			92255 37c	1 62440E-07
			92256 37c	1 14177E-00
			92257 37c	4 64211E-07
			92258 37c	3 98220E-00
			92259 37c	3 99947E-00
			92260 37c	1 52200E-07
			92261 37c	4 58600E-02
m9	§	Neologus data for case 1 B32C08A222222	92262 37c	2 94300E-04
			92263 37c	6 51700E-00
			92264 37c	1 08400E-04
			92265 37c	1 05712E-05
			92266 37c	2 14300E-02
			92267 37c	3 18400E-00
			92268 37c	1 29010E-04
			92269 37c	4 48674E-05
			92270 37c	2 21600E-05
			92271 37c	1 54147E-00
			92272 37c	5 71040E-00
			92273 37c	1 47140E-00
			92274 37c	4 08447E-00
			92275 37c	4 50211E-05
			92276 37c	4 22040E-00
			92277 37c	2 40730E-00
			92278 37c	3 80400E-00
			92279 37c	4 70400E-00
			92280 37c	3 54940E-05
			92281 37c	2 72100E-05
			92282 37c	3 14600E-00
			92283 37c	1 51147E-07
			92284 37c	1 07300E-05
			92285 37c	4 17040E-07
			92286 37c	3 44600E-04
			92287 37c	3 73501E-00
			92288 37c	1 30400E-07
			92289 37c	4 50600E-02
m10	§	Neologus data for case 1 B32C08A222222	92290 37c	2 94400E-04
			92291 37c	3 91700E-00
			92292 37c	1 10440E-04
			92293 37c	1 50480E-05
			92294 37c	3 14700E-02
			92295 37c	3 47640E-00
			92296 37c	1 28471E-04
			92297 37c	4 83600E-00
			92298 37c	2 20180E-05
			92299 37c	1 29114E-00
			92300 37c	6 95310E-00
			92301 37c	1 46400E-00
			92302 37c	4 75120E-05
			92303 37c	4 73800E-05
			92304 37c	4 36504E-05
			92305 37c	2 56900E-05
			92306 37c	3 21700E-00
			92307 37c	4 80700E-00
			92308 37c	5 49040E-05
			92309 37c	2 80007E-00
			92310 37c	5 56900E-00
			92311 37c	1 23140E-07
			92312 37c	1 14310E-05
			92313 37c	4 70750E-07
			92314 37c	3 63630E-00
			92315 37c	3 88040E-04
			92316 37c	1 47810E-07
			92317 37c	4 50600E-02
m11	§	Neologus data for case 1 B32C08A222222	92318 37c	2 83300E-04
			92319 37c	6 02874E-00
			92320 37c	1 18220E-04
			92321 37c	1 09000E-05
			92322 37c	2 14220E-02
			92323 37c	3 43600E-00
			92324 37c	1 28422E-04
			92325 37c	4 61882E-00
			92326 37c	2 28313E-02
			92327 37c	1 20110E-00
			92328 37c	1 92804E-00
			92329 37c	1 54347E-00
			92330 37c	4 39600E-05
			92331 37c	1 14700E-00
			92332 37c	4 72200E-00
			92333 37c	4 14820E-00
			92334 37c	2 345180E-05
			92335 37c	3 19741E-00
			92336 37c	4 88810E-00
			92337 37c	3 40000E-05
			92338 37c	2 79040E-00
			92339 37c	5 63870E-00
			92340 37c	1 67000E-07
			92341 37c	1 17500E-05
			92342 37c	4 18521E-07
			92343 37c	3 92470E-00
			92344 37c	3 44001E-00
			92345 37c	1 45500E-07
			92346 37c	4 58600E-02
m12	§	Neologus data for case 1 B32C08A222222	92347 37c	2 83400E-04
			92348 37c	4 98800E-00
			92349 37c	1 09400E-04
			92350 37c	1 07470E-05
			92351 37c	3 14307E-02
			92352 37c	3 52704E-00
			92353 37c	1 28440E-04
			92354 37c	4 56220E-00
			92355 37c	2 78200E-05
			92356 37c	3 90407E-00
			92357 37c	6 82755E-00
			92358 37c	1 93700E-00
			92359 37c	4 46040E-00
			92360 37c	4 68571E-05
			92361 37c	4 28800E-00
			92362 37c	2 52870E-05
			92363 37c	3 13210E-00
			92364 37c	4 82147E-00
			92365 37c	3 37200E-05
			92366 37c	2 75000E-05
			92367 37c	4 48400E-00
			92368 37c	1 51830E-07
			92369 37c	1 08710E-05
			92370 37c	4 78700E-07
			92371 37c	3 62750E-04
			92372 37c	2 41804E-00
			92373 37c	1 43204E-07
			92374 37c	4 08500E-02
m13	§	Neologus data for case 1 B32C08A222222	92375 37c	3 18670E-04
			92376 37c	6 88671E-00
			92377 37c	1 06130E-04
			92378 37c	1 07740E-00
			92379 37c	2 14730E-02
			92380 37c	2 70870E-00
			92381 37c	1 25040E-04
			92382 37c	4 24750E-00
			92383 37c	2 06170E-00
			92384 37c	7 38870E-00
			92385 37c	6 28804E-00
			92386 37c	1 19120E-00
			92387 37c	4 30194E-05
			92388 37c	4 36101E-05
			92389 37c	3 87640E-00
			92390 37c	2 38321E-00
			92391 37c	2 79000E-00
			92392 37c	4 83044E-00
			92393 37c	3 52300E-00
			92394 37c	2 50800E-05
			92395 37c	1 18220E-00
			92396 37c	1 508120E-07
			92397 37c	1 00804E-00
			92398 37c	4 80137E-07
			92399 37c	3 87547E-00
			92400 37c	3 45400E-00
			92401 37c	1 28370E-07
			92402 37c	4 00600E-02
m14	§	Neologus data for case 1 B32C08A222222	92403 37c	2 85710E-04
			92404 37c	1 01640E-00
			92405 37c	1 08571E-04
			92406 37c	1 05700E-05
			92407 37c	2 14300E-02
			92408 37c	3 99940E-00
			92409 37c	1 28800E-04
			92410 37c	4 49800E-05
			92411 37c	2 21670E-00
			92412 37c	8 57200E-00
			92413 37c	6 71440E-00
			92414 37c	1 47430E-00
			92415 37c	4 80600E-05
			92416 37c	4 80100E-05
			92417 37c	4 22300E-00
			92418 37c	2 40850E-05
			92419 37c	3 02807E-00
			92420 37c	4 18500E-00
			92421 37c	4 72100E-07
			92422 37c	3 84200E-00
			92423 37c	3 73700E-00

m15	84155 37c	1 39618E-07	82162 37c	3 27360E-06
	8016 37c	4 5850E-02	83153 37c	2 82138E-06
	\$	Isotopic data for case 1 B32C05A222222	84155 37c	9 87156E-06
	82255 37c	3 08078E-04	8996 37c	4 3450E-02
	82256 37c	0 1136E-06	m16	\$
	82256 37c	1 06822E-04	\$	Isotopic data for case 1 B32C05A222222
	82257 37c	1 01206E-05	82255 37c	3 84048E-04
	82258 37c	2 14507E-02	82256 37c	6 85230E-04
	84238 37c	2 94449E-06	82257 37c	9 48098E-05
	84239 37c	1 24500E-04	83227 37c	7 3357E-04
m16	84240 37c	4 38810E-05	83228 37c	2 16418E-02
	84241 37c	2 14027E-05	84153 37c	1 75024E-06
	84242 37c	7 89877E-06	84239 37c	1 26826E-04
	85241 37c	8 47345E-06	84240 37c	3 56336E-05
	85242 37c	1 30813E-06	84241 37c	1 88608E-05
	42085 37c	4 45403E-05	84242 37c	4 75895E-06
	43089 37c	4 44807E-06	85241 37c	8 66832E-06
	44101 37c	4 08324E-05	85242 37c	4 38848E-07
	44102 37c	2 40338E-05	42095 37c	3 74038E-05
	47108 37c	2 90902E-06	43096 37c	3 72577E-05
m17	56135 37c	4 63359E-06	44101 37c	3 3284E-05
	59143 37c	3 27424E-05	45103 37c	2 08553E-06
	60145 37c	2 64710E-06	47108 37c	2 13402E-06
	62147 37c	8 22499E-06	86133 37c	3 89438E-05
	62148 37c	1 61218E-07	80143 37c	2 87787E-05
	62149 37c	1 03949E-06	80145 37c	2 23907E-05
	62150 37c	4 84785E-06	82147 37c	7 39520E-05
	62151 37c	3 74837E-06	82148 37c	1 47263E-07
	62152 37c	3 67548E-06	82150 37c	8 3328E-06
	64155 37c	1 32008E-07	82151 37c	4 25027E-07
m18	8016 37c	4 56500E-02	82152 37c	3 21811E-06
	\$	Isotopic data for case 1 B32C05A222222	83153 37c	2 78871E-06
	82255 37c	3 13335E-04	84154 37c	9 48172E-06
	82256 37c	6 15127E-04	9016 37c	4 36000E-02
	82257 37c	1 06914E-04	m17	\$
	82258 37c	9 06222E-04	\$	Isotopic data for case 1 B32C05A222222
	82259 37c	2 44634E-02	82256 37c	4 83518E-04
	84238 37c	2 86614E-06	82257 37c	1 26541E-05
	84239 37c	1 52879E-04	82258 37c	2 88878E-04
	84240 37c	4 30458E-04	82259 37c	5 84589E-06
m19	84241 37c	2 11218E-05	82259 37c	2 18408E-02
	84242 37c	7 85493E-06	84238 37c	1 08438E-06
	85241 37c	8 36270E-06	84239 37c	1 13645E-04
	85242 37c	1 25910E-06	84240 37c	2 83857E-05
	42085 37c	4 41633E-06	84241 37c	1 27913E-04
	43089 37c	4 41581E-05	84242 37c	2 74311E-06
	44101 37c	4 07186E-05	85241 37c	3 74723E-06
	44102 37c	2 40342E-05	85242 37c	2 86815E-07
	47108 37c	2 84807E-06	42086 37c	3 40638E-05
	56135 37c	4 60369E-05	43090 37c	3 04713E-05
m20	59143 37c	3 24634E-05	44181 37c	2 72344E-04
	60145 37c	2 64710E-06	86180 37c	1 72885E-05
	62147 37c	8 22499E-06	87108 37c	1 54833E-06
	62148 37c	1 61218E-07	80153 37c	5 24174E-05
	62149 37c	1 03949E-06	80143 37c	2 47651E-05
	62150 37c	4 84785E-07	80145 37c	1 86407E-06
	62151 37c	3 74837E-06	82147 37c	6 48037E-06
	62152 37c	3 67548E-06	82148 37c	1 34046E-07
	64155 37c	1 32008E-07	82150 37c	8 07044E-06
	8016 37c	4 56500E-02	82151 37c	3 80750E-07
m21	\$	Isotopic data for case 1 B32C05A222222	82152 37c	2 71832E-06
	82255 37c	3 13335E-04	83153 37c	2 78871E-06
	82256 37c	6 15127E-04	84154 37c	9 48158E-06
	82257 37c	1 06914E-04	84155 37c	8 87568E-06
	82258 37c	9 06222E-04	8996 37c	4 3450E-02
	82259 37c	2 44634E-02	m21	\$
	84238 37c	2 86614E-06	\$	Isotopic data for case 1 B32C05A222222
	84239 37c	1 52879E-04	82256 37c	5 51813E-04
	84240 37c	4 30458E-04	82257 37c	7 85193E-06
	84241 37c	2 11218E-05	82258 37c	6 74148E-05
m22	84242 37c	7 85493E-06	82259 37c	3 92843E-06
	85241 37c	8 36270E-06	82259 37c	2 17288E-02
	85242 37c	1 25910E-06	84238 37c	5 26368E-07
	42085 37c	4 41633E-06	84239 37c	1 01346E-04
	43089 37c	4 41581E-05	84240 37c	2 82383E-06
	44101 37c	4 07186E-05	84241 37c	8 19882E-06
	44102 37c	2 40342E-05	84242 37c	1 23788E-06
	47108 37c	2 84807E-06	85241 37c	2 37114E-06
	56135 37c	4 60369E-05	85242 37c	7 38294E-06
	59143 37c	3 24634E-05	42086 37c	2 40780E-05
m23	60145 37c	2 64710E-06	43090 37c	2 84090E-06
	62147 37c	8 22499E-06	44101 37c	2 08388E-05
	62148 37c	1 61218E-07	45103 37c	1 39830E-05
	62149 37c	1 03949E-06	47108 37c	8 93322E-07
	62150 37c	4 84785E-07	56135 37c	2 61255E-05
	62151 37c	3 74837E-06	59143 37c	3 88120E-06
	62152 37c	3 67548E-06	60145 37c	1 44868E-06
	64155 37c	1 32008E-07	62147 37c	5 34875E-06
	8016 37c	4 56500E-02	62148 37c	1 37121E-07
	\$	Isotopic data for case 1 B32C05A222222	62149 37c	4 83828E-06
m24	82255 37c	3 13335E-04	62150 37c	3 84723E-07
	82256 37c	6 15127E-04	62151 37c	2 71988E-06
	82257 37c	1 06914E-04	62152 37c	2 13689E-06
	82258 37c	9 06222E-04	64155 37c	4 35823E-06
	82259 37c	2 44634E-02	8996 37c	4 3450E-02
	84238 37c	2 86614E-06	m102	\$
	84239 37c	1 52879E-04	\$	Isotopic data for case 1 B32C05A222222
	84240 37c	4 30458E-04	24900 37c	7 589E-05
	84241 37c	2 11218E-05	26000 37c	1 484E-04
	84242 37c	7 85493E-06	40000 37c	4 288E-02
m25	85241 37c	8 36270E-06	m103	\$
	85242 37c	1 25910E-06	\$	Isotopic data for case 1 B32C05A222222
	42085 37c	4 41633E-06	1001 37c	5 875E-02
	43089 37c	4 41581E-05	8016 37c	2 327E-02
	44101 37c	4 07186E-05	m103	\$
	44102 37c	2 40342E-05	\$	Isotopic data for case 1 B32C05A222222
	47108 37c	2 84807E-06	24000 37c	1 743E-02
	56135 37c	4 60369E-05	26000 37c	1 735E-03
	59143 37c	3 24634E-05	26000 37c	5 935E-02
	60145 37c	2 64710E-06	26008 37c	7 727E-03
m26	62147 37c	8 22499E-06	m104	\$
	62148 37c	1 61218E-07	\$	Isotopic data for case 1 B32C05A222222
	62149 37c	1 03949E-06	24000 37c	1 743E-02
	62150 37c	4 84785E-07	26000 37c	1 735E-03
	62151 37c	3 74837E-06	26000 37c	5 935E-02
	62152 37c	3 67548E-06	26008 37c	7 727E-03
	64155 37c	1 32008E-07	m105	\$
	8016 37c	4 56500E-02	\$	Isotopic data for case 1 B32C05A222222
	\$	Isotopic data for case 1 B32C05A222222	24000 37c	1 881E-02
	\$	Isotopic data for case 1 B32C05A222222	26000 37c	1 884E-06
\$	Isotopic data for case 1 B32C05A222222	26000 37c	5 738E-03	
\$	Isotopic data for case 1 B32C05A222222	26000 37c	7 488E-03	
\$	Isotopic data for case 1 B32C05A222222	26000 37c	7 828E-04	
\$	Isotopic data for case 1 B32C05A222222	26000 37c	9 189E-03	
\$	Isotopic data for case 1 B32C05A222222	26000 37c	9 189E-03	
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72.7100 0.01 60 57.1	1.00	102	72.7100 0.01 60 57.1	2.00	2
72.7100 0.01 50 57.2	2.00	102	72.7100 0.01 50 57.1	3.00	2

*IBM source print specification calls for allowing 1/8" for the printing space.
*Form 511 cards are recycled 32 times for the job using.

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72 71(1 2 0) 64 50(8 0 0) 32
72 71(1 2 0) 64 50(8 0 0) 32
#
sp1 0 1 200367
c
ad2 6 0 0 0 479
sp2 -21 1
c
ad3 0 0 250 0
sp3 0 1
c
sp4 -3
#
p=mp 70 100 1 3
c=beta 175

```

Appendix II

Table of Fission Source Distributions

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Table A-1 Fission source fraction at 32 MWd/kgU based on AC+FP approach

Asst Node	Case - Symmetric	Case - Uniform	Case 01	Case 02	Case 03	Case 04	Case 05	Case 06	Case 07	Case 08	Case 09	Case 10	Case 11	Case 12	Case 13	Case 14	Case 15	Case 16	Case 17
32	0.03843	0.00354	1.1E-05	0	0	2.8E-08	0	0	7.1E-08	0	4.9E-06	3.8E-06	4E-06	0.00123	0.00079	0	0	0	0
31	0.05304	0.00896	2.8E-05	0	0	4.8E-06	0	0	8.4E-06	0	1.8E-05	0.00011	8E-05	0.00167	0.00104	0	0	0	4.2E-07
30	0.05423	0.01439	2.8E-05	0	0	9.2E-06	0	0	1.6E-05	0	2.6E-05	0.00014	5.7E-05	0.00176	0.00138	3.2E-06	4E-07	6.2E-07	6.1E-07
29	0.05107	0.01985	6.3E-05	0	0	1.3E-06	0	0	1.8E-05	0	4.4E-05	0.00019	4.9E-05	0.00148	0.00167	4.8E-06	1.1E-06	8E-06	9.1E-06
28	0.04515	0.02255	0.0001	8.2E-07	0	4.5E-06	0	1.2E-06	2.3E-05	0	6.4E-06	0.00017	5.6E-06	0.00138	0.00172	1.5E-06	2E-05	9.8E-06	2.3E-05
27	0.03831	0.02805	0.00018	4.1E-07	2E-06	4.1E-06	0	6.4E-06	4.9E-06	0	7.8E-05	0.0002	8.3E-05	0.00124	0.00185	1.1E-06	4.4E-06	3.2E-05	4E-06
26	0.03638	0.02848	0.00026	1.4E-06	3.1E-06	6.7E-06	7.3E-07	8.7E-06	8.7E-05	2.5E-09	9.5E-05	0.00022	0.00011	0.00106	0.00202	4E-05	5.5E-06	4.7E-05	9.4E-05
25	0.03006	0.03271	0.00028	7.4E-06	7.7E-06	1.9E-05	1.1E-06	1.4E-05	7.3E-05	7.7E-06	0.00012	0.00028	0.00011	0.00081	0.00186	8.4E-06	6.1E-06	9.6E-05	0.00016
24	0.02563	0.03578	0.00036	2.2E-05	1.7E-05	7.7E-05	5.1E-06	1.9E-05	0.00011	3E-05	0.00014	0.00036	0.00012	0.00106	0.00184	0.00011	0.00011	0.00018	0.00023
23	0.02193	0.03818	0.00049	2.8E-05	2.8E-05	0.00012	2.4E-05	3.1E-05	0.00018	6.4E-05	0.00023	0.00044	0.00021	0.00118	0.00214	0.00018	0.00011	0.00023	0.00035
22	0.0189	0.04088	0.00068	3.9E-06	6.7E-05	8.4E-06	3E-06	6.5E-05	0.00029	0.00018	0.00033	0.00058	0.0004	0.00125	0.00283	0.00028	0.00014	0.00033	0.00048
21	0.01574	0.04265	0.0008	6.3E-05	0.00012	6.4E-05	9.4E-05	0.00019	0.00033	0.00023	0.00049	0.0007	0.00078	0.00149	0.00337	0.00045	0.00028	0.00048	0.00073
20	0.01345	0.04473	0.00086	0.00018	0.00029	8.7E-05	0.0002	0.00036	0.00048	0.00034	0.00068	0.00087	0.00128	0.00205	0.00415	0.00067	0.0004	0.00068	0.00099
19	0.01113	0.04571	0.00097	0.00028	0.00044	0.00012	0.00028	0.0006	0.0006	0.00082	0.00086	0.00145	0.00179	0.0028	0.00504	0.00093	0.00088	0.00104	0.00161
18	0.01018	0.04813	0.00132	0.00084	0.0006	0.00026	0.00041	0.00075	0.00091	0.00104	0.0011	0.00178	0.00245	0.00372	0.0061	0.00126	0.00104	0.0016	0.00226
17	0.00883	0.04836	0.0018	0.00087	0.0008	0.00053	0.0007	0.00105	0.00154	0.00168	0.00163	0.00236	0.00381	0.00506	0.00737	0.00168	0.00158	0.00233	0.00313
16	0.00676	0.04817	0.00259	0.00133	0.00118	0.00103	0.00118	0.00168	0.00211	0.00247	0.00215	0.00337	0.00487	0.00688	0.00908	0.00224	0.00214	0.00331	0.00458
15	0.00558	0.04749	0.00342	0.00217	0.00198	0.00183	0.00198	0.00266	0.00312	0.0037	0.00322	0.00436	0.00709	0.00936	0.01143	0.00325	0.00307	0.00484	0.00667
14	0.01089	0.04539	0.005	0.0033	0.00263	0.00294	0.00307	0.00452	0.00488	0.00546	0.0052	0.00625	0.0103	0.01234	0.01528	0.00483	0.00463	0.00647	0.00928
13	0.01226	0.04516	0.00721	0.0056	0.0043	0.00456	0.00602	0.00683	0.00744	0.00804	0.00776	0.00823	0.01445	0.0165	0.02062	0.0074	0.00668	0.00927	0.01289
12	0.01486	0.04308	0.00976	0.00807	0.0065	0.00641	0.00727	0.00815	0.01016	0.01117	0.01127	0.01211	0.01813	0.02008	0.02524	0.01008	0.00828	0.01284	0.0164
11	0.01779	0.04036	0.01429	0.0119	0.00868	0.00882	0.01072	0.01271	0.01375	0.01584	0.01683	0.01861	0.02309	0.02617	0.03148	0.01423	0.01302	0.01881	0.02237
10	0.02187	0.03767	0.02103	0.01783	0.01461	0.01445	0.01628	0.01783	0.01968	0.02315	0.02362	0.02289	0.02928	0.03408	0.03995	0.02151	0.01788	0.02729	0.03057
9	0.02684	0.03477	0.03123	0.02866	0.02382	0.02486	0.02444	0.02581	0.02879	0.03356	0.03229	0.03289	0.03787	0.04416	0.04874	0.0316	0.02582	0.0387	0.0416
8	0.0319	0.03172	0.04414	0.03884	0.03497	0.03621	0.03547	0.03688	0.04062	0.04579	0.04377	0.04455	0.04905	0.05439	0.05928	0.04301	0.03808	0.0524	0.0583
7	0.0389	0.0284	0.06111	0.05562	0.0525	0.05202	0.05109	0.05426	0.0586	0.06008	0.05815	0.05863	0.06115	0.06659	0.07014	0.06806	0.06889	0.06616	0.07221
6	0.04511	0.02489	0.06297	0.06084	0.07881	0.07366	0.07384	0.07881	0.07893	0.07887	0.07735	0.07843	0.07889	0.07946	0.08315	0.07895	0.0874	0.08888	0.09485
5	0.0527	0.02147	0.11188	0.11276	0.11289	0.10776	0.10721	0.1063	0.10816	0.10913	0.10822	0.10824	0.09831	0.09731	0.08993	0.10725	0.09895	0.11424	0.11859
4	0.06084	0.01777	0.14707	0.15189	0.18482	0.18286	0.18881	0.18479	0.18817	0.18646	0.18623	0.18774	0.12877	0.1213	0.11739	0.14449	0.14325	0.14354	0.13959
3	0.06547	0.01576	0.18513	0.17296	0.18081	0.1818	0.18217	0.17127	0.1885	0.18476	0.18624	0.18542	0.15025	0.13812	0.12402	0.18506	0.17788	0.15587	0.14885
2	0.06259	0.00938	0.18454	0.17713	0.18388	0.18661	0.18539	0.1856	0.17309	0.18514	0.18617	0.17037	0.15812	0.13805	0.11746	0.19088	0.14784	0.1284	0.1284
1	0.04486	0.00511	0.12127	0.13113	0.13682	0.13906	0.14053	0.14082	0.13344	0.12281	0.12489	0.12881	0.11826	0.10785	0.08288	0.1319	0.14814	0.09848	0.08992

Table A-2 Fission source fraction at 50 MWd/kgU based on AC+FP approach

Asel Node	Case - Symmetric	Case - Uniform	Case 16 A222222	Case 19 A111111	Case 20 A111112	Case 21 A111122	Case 22 A111222	Case 23 A112222	Case 24 A122222	Case 25 A222223	Case 26 A222233	Case 27 A222333	Case 28 A233333	Case 30 A333333	Case 31 A122223	Case 32 A112233	Case 33 A322221	Case 34 A332211
32	0.0893	0.00558	0	0	0	0	0	0	0	0	5.8E-09	0	0	0	0	0	0	6E-08
31	0.10374	0.01022	0	0	0	0	0	0	0	0	1.7E-06	0	0	0	0	0	0	1.2E-05
30	0.09856	0.01481	0	0	0	0	0	0	0	0	9.2E-07	0	0	0	0	0	0	9.7E-06
29	0.08847	0.01947	0	0	0	0	0	0	0	0	5E-08	0	0	0	0	0	0	1.6E-05
28	0.04677	0.02383	0	0	2.6E-07	0	0	0	0	0	1.6E-05	0	0	0	0	0	0	4.3E-05
27	0.03232	0.02779	0	0	1.6E-06	0	0	0	0	0	2.4E-05	0	0	0	1.8E-07	0	0	4.2E-05
26	0.02276	0.03125	0	0	8.6E-06	0	0	0	0	0	8.9E-05	0	0	0	4.9E-06	0	0	1.8E-05
25	0.01521	0.03429	0	0	2E-05	0	0	0	0	0	3.9E-05	0	0	4.8E-09	1.8E-05	6.7E-08	0	1.3E-05
24	0.01011	0.03696	0	0	1.2E-05	0	0	0	1.7E-08	9E-08	3.1E-05	0	0	7.4E-06	1.2E-05	4.8E-06	1.2E-06	2.2E-05
23	0.00655	0.03831	1.1E-08	0	7.8E-06	0	0	0	4.7E-08	1.3E-08	1.8E-05	0	0	1.7E-06	1.8E-05	7.8E-06	2.2E-06	2E-05
22	0.00447	0.04162	1.3E-05	0	3.9E-08	0	0	0	2.1E-05	3.8E-06	2.2E-05	2.8E-07	0	5E-06	2.4E-05	9.8E-06	2.2E-07	8.6E-06
21	0.00268	0.04382	9.6E-06	2.6E-09	1E-05	0	0	2.5E-08	4.9E-05	5.1E-08	1.4E-05	2.7E-06	2.3E-06	1.8E-05	3.1E-05	8.3E-06	1.8E-06	1.3E-05
20	0.00204	0.04673	2.4E-05	2.8E-07	1.4E-08	1.6E-06	5.2E-07	1E-06	9.9E-05	2.8E-08	4.2E-05	9.8E-06	4.9E-06	4.1E-05	7E-05	4E-06	4.2E-06	4.1E-05
19	0.00135	0.04704	4.4E-05	1E-05	3.1E-05	5.5E-06	3.2E-08	7.5E-06	0.00016	1.7E-05	7.9E-05	2E-05	1.9E-05	9.8E-05	0.00016	1.6E-05	1.7E-05	5.2E-05
18	0.00114	0.04745	7.7E-05	3.1E-05	5.6E-05	1E-05	7.6E-06	2.3E-05	0.00012	4.8E-05	0.00011	3.3E-05	5.8E-05	0.00019	0.0009	4.1E-05	1.7E-05	5.9E-05
17	0.00098	0.04748	0.00017	5.7E-05	0.00013	3.3E-05	3.1E-05	4.5E-05	0.00018	0.0001	0.00014	9.3E-05	0.00016	0.00042	0.00067	8.3E-05	4.9E-05	0.00038
16	0.00087	0.04852	0.00028	0.00014	0.0002	5.8E-05	5.8E-05	0.00016	0.00027	0.00021	0.00023	0.00023	0.00026	0.00069	0.00108	0.00014	0.00015	0.00021
15	0.00107	0.04598	0.00056	0.00026	0.0004	0.00014	0.00022	0.00028	0.00038	0.00032	0.00048	0.00045	0.00046	0.00117	0.00144	0.00029	0.00034	0.00042
14	0.00142	0.04482	0.00091	0.00052	0.00068	0.00031	0.00047	0.00062	0.00088	0.00058	0.00085	0.00079	0.00086	0.00187	0.0023	0.0008	0.00062	0.00083
13	0.00196	0.043	0.0018	0.001	0.00112	0.00078	0.00105	0.00112	0.00132	0.00116	0.00145	0.00154	0.00182	0.00317	0.00382	0.00133	0.00101	0.00168
12	0.00285	0.04112	0.00258	0.00176	0.0019	0.00131	0.00179	0.00189	0.00216	0.00201	0.00242	0.00243	0.00321	0.00427	0.00606	0.00215	0.00176	0.00276
11	0.00431	0.03907	0.00404	0.00313	0.00368	0.00263	0.00304	0.00331	0.00381	0.00348	0.00417	0.00421	0.00609	0.00643	0.00825	0.00368	0.00292	0.00493
10	0.00653	0.03664	0.00674	0.00564	0.00606	0.00504	0.00543	0.00619	0.00655	0.00639	0.00752	0.00707	0.00815	0.00883	0.0144	0.00603	0.00501	0.00807
9	0.01027	0.03438	0.01189	0.01005	0.01048	0.00969	0.00977	0.0105	0.01149	0.01163	0.01249	0.01202	0.01323	0.0162	0.02189	0.01026	0.00696	0.01376
8	0.0162	0.0313	0.02	0.01788	0.01802	0.01683	0.01739	0.01794	0.01944	0.01979	0.02029	0.01947	0.02085	0.02502	0.03253	0.01815	0.01562	0.02289
7	0.02412	0.02841	0.03261	0.031	0.03084	0.02935	0.02955	0.03082	0.03276	0.03281	0.03329	0.03114	0.03306	0.03817	0.04674	0.03098	0.02723	0.03664
6	0.03303	0.02521	0.05432	0.03411	0.03236	0.03137	0.03065	0.03185	0.03347	0.0334	0.03661	0.03502	0.03162	0.03732	0.05533	0.03135	0.04849	0.0604
5	0.04723	0.02138	0.09065	0.06067	0.06787	0.06632	0.06639	0.0656	0.06759	0.06128	0.06512	0.06155	0.06273	0.06975	0.09589	0.06785	0.06136	0.06707
4	0.06799	0.01758	0.14652	0.14518	0.1483	0.14651	0.14511	0.14321	0.14386	0.14516	0.13579	0.1319	0.13669	0.14197	0.14474	0.14381	0.15477	0.15597
3	0.08847	0.01351	0.19828	0.19625	0.19932	0.20084	0.20143	0.19987	0.1938	0.18845	0.19796	0.19884	0.18282	0.16144	0.17836	0.197	0.20106	0.19872
2	0.1051	0.00916	0.23396	0.23795	0.23772	0.23987	0.24195	0.24158	0.23706	0.23438	0.23686	0.24558	0.24812	0.22371	0.20709	0.23843	0.24925	0.22542
1	0.06178	0.00504	0.19701	0.20148	0.20384	0.20481	0.20575	0.20702	0.20621	0.1998	0.20143	0.20873	0.21731	0.20658	0.17175	0.20694	0.21414	0.17128

Table A-3 Fission source fraction at 32 MWd/kgU based on AC-only approach

Adnl node	Case -- Symmetry	Case -- Uniform	Case 01 A222222	Case 02 A111111	Case 03 A111112	Case 04 A111122	Case 05 A111222	Case 06 A112222	Case 07 A122222	Case 08 A222223	Case 09 A222233	Case 10 A222333	Case 11 A223333	Case 12 A233333	Case 13 A333333	Case 14 A122233	Case 15 A112233	Case 16 A322221	Case 17 A332221
32	0 02726	0 06809	0 00011	1E-06	0	0	0 00023	1 3E-06	6E-06	0 00016	0 0002	0 00011	0 00166	0 00156	0 00464	0 0002	2 9E-05	4 7E-05	0 00027
31	0 04037	0 00829	0 00024	3E-07	0	2 8E-07	0 00029	5 8E-06	6 7E-06	0 00028	0 00026	0 00022	0 00208	0 00254	0 00595	0 00026	3 4E-06	4 5E-05	0 00041
30	0 04651	0 01349	0 00036	6 4E-06	0	3 2E-06	0 00042	1 2E-06	9 5E-05	0 00017	0 00036	0 00037	0 00213	0 00312	0 00876	0 00035	6 3E-06	4 6E-06	0 00044
29	0 04676	0 0174	0 00056	2 9E-05	1 6E-07	1 1E-05	0 00042	6 5E-06	0 00017	0 0002	0 00042	0 00052	0 00217	0 00345	0 00899	0 00027	6 0001	6 7E-05	0 00045
28	0 0486	0 02123	0 00067	3 7E-05	1 4E-06	1 7E-05	0 00043	2 1E-05	0 00026	0 00021	0 00059	0 00077	0 00226	0 00401	0 00704	0 00031	0 00012	0 00014	0 0005
27	0 04823	0 02491	0 00078	2 6E-05	1 1E-06	1 4E-05	0 00047	5 7E-06	0 00039	0 00028	0 0006	0 00086	0 00244	0 00431	0 0068	0 00039	0 00023	0 00026	0 00056
26	0 04681	0 02868	0 00103	2 6E-05	1 3E-05	5 7E-06	0 00047	8 5E-05	0 00061	0 00054	0 00063	0 00119	0 00259	0 00477	0 00668	0 00058	0 00031	0 00037	0 00057
25	0 04396	0 03212	0 00121	2 5E-05	5 2E-05	0 00013	0 00047	0 00014	0 00069	0 00069	0 00112	0 0013	0 00294	0 00541	0 00629	0 0006	0 0005	0 00052	0 00066
24	0 04166	0 03626	0 00144	3 9E-05	7 5E-05	0 00018	0 00071	0 00032	0 00077	0 00084	0 00106	0 00151	0 00345	0 00624	0 00808	0 00103	0 00061	0 00037	0 00078
23	0 03671	0 03612	0 00183	0 00013	0 00013	0 00023	0 00062	0 00062	0 00068	0 0012	0 00121	0 00165	0 00429	0 00719	0 00996	0 00127	0 00118	0 00096	0 00102
22	0 03646	0 04054	0 00244	0 00024	0 00024	0 0004	0 00106	0 00101	0 00115	0 00172	0 00153	0 00206	0 00516	0 00826	0 00704	0 00164	0 00154	0 00127	0 00143
21	0 03383	0 04244	0 00306	0 00032	0 00054	0 00054	0 00134	0 0015	0 00139	0 00225	0 00191	0 00247	0 00554	0 0089	0 00859	0 00217	0 00184	0 00163	0 00203
20	0 03067	0 04438	0 00402	0 00065	0 00106	0 00063	0 00173	0 00192	0 00185	0 00297	0 00262	0 00354	0 00677	0 00995	0 01039	0 00301	0 00222	0 00252	0 00316
19	0 02694	0 04578	0 00509	0 00111	0 00148	0 00122	0 00223	0 00257	0 00277	0 00367	0 00341	0 00476	0 00809	0 01141	0 01226	0 00417	0 00286	0 00348	0 00445
18	0 02643	0 04722	0 00668	0 00173	0 00216	0 00181	0 00263	0 00351	0 00393	0 00535	0 00482	0 006	0 01023	0 01322	0 01463	0 00649	0 00368	0 00468	0 00605
17	0 02426	0 04753	0 00837	0 00262	0 00326	0 00267	0 00406	0 00494	0 00517	0 00748	0 00673	0 00754	0 0126	0 01597	0 01708	0 00712	0 00484	0 00621	0 00606
16	0 02372	0 04762	0 01072	0 00411	0 0045	0 00406	0 00552	0 00676	0 00694	0 00969	0 00921	0 00943	0 01567	0 01845	0 01851	0 00911	0 00645	0 00619	0 01085
15	0 02295	0 04757	0 0137	0 00568	0 00602	0 0057	0 00749	0 00867	0 00931	0 01236	0 01231	0 01203	0 0189	0 02286	0 02211	0 01136	0 00652	0 01103	0 01406
14	0 02242	0 0468	0 01738	0 00691	0 00671	0 00618	0 01043	0 0126	0 01227	0 01611	0 01654	0 01569	0 02296	0 02763	0 0266	0 01502	0 01161	0 01506	0 01792
13	0 02226	0 04532	0 0226	0 01271	0 01253	0 01191	0 01425	0 01775	0 01689	0 02036	0 02024	0 02129	0 02864	0 03395	0 03241	0 01894	0 0158	0 02083	0 02313
12	0 02204	0 04331	0 02662	0 01718	0 01688	0 01601	0 01765	0 02247	0 02063	0 0236	0 02537	0 02616	0 03369	0 0386	0 0373	0 02382	0 01977	0 02625	0 02885
11	0 02247	0 04124	0 03293	0 023	0 02236	0 02162	0 02369	0 02854	0 02716	0 02864	0 0317	0 03272	0 04	0 04456	0 04374	0 03038	0 02504	0 03291	0 03675
10	0 02318	0 03872	0 04091	0 03007	0 02978	0 02787	0 03061	0 03648	0 03539	0 03514	0 03937	0 04082	0 0467	0 05147	0 05167	0 03792	0 03186	0 04267	0 04583
9	0 02412	0 03617	0 05088	0 04009	0 03926	0 03617	0 04061	0 04439	0 04623	0 04458	0 04822	0 05145	0 05488	0 0588	0 06068	0 04864	0 04205	0 05371	0 05647
8	0 02564	0 03332	0 06064	0 05344	0 06111	0 05103	0 05291	0 05474	0 05918	0 05529	0 0586	0 06094	0 06185	0 06383	0 06679	0 06071	0 05332	0 06483	0 0676
7	0 02714	0 02989	0 07243	0 07036	0 06656	0 06713	0 06756	0 06799	0 07425	0 06765	0 07041	0 07129	0 0691	0 06973	0 07117	0 07337	0 06712	0 07661	0 08062
6	0 02718	0 02825	0 08498	0 0833	0 08666	0 08674	0 0854	0 08444	0 09065	0 08283	0 08165	0 08193	0 07566	0 0734	0 07346	0 06549	0 06319	0 09501	0 09524
5	0 02617	0 02234	0 09666	0 11863	0 11302	0 11034	0 10726	0 10416	0 10929	0 10433	0 09832	0 09762	0 08553	0 07938	0 07912	0 10518	0 10429	0 10753	0 1083
4	0 03028	0 01839	0 11725	0 14066	0 13696	0 13992	0 13246	0 12761	0 12796	0 12642	0 12227	0 11908	0 096	0 08714	0 08411	0 12236	0 13168	0 12022	0 11627
3	0 03032	0 01619	0 12073	0 14847	0 14123	0 13932	0 14666	0 13762	0 1324	0 13094	0 13102	0 12605	0 10253	0 08691	0 08167	0 12676	0 14202	0 12174	0 11242
2	0 02721	0 0067	0 11229	0 13658	0 143	0 14723	0 14109	0 13361	0 12336	0 12463	0 12235	0 11891	0 10002	0 07791	0 07161	0 11683	0 13679	0 10911	0 09291
1	0 01834	0 00531	0 07794	0 09267	0 09832	0 10261	0 09821	0 09554	0 08765	0 08664	0 08585	0 08267	0 07097	0 06619	0 0472	0 0839	0 0963	0 08628	0 05969

Table A-4 Fission source fraction at 50 MWd/kgU based on AC-only approach

Aztl node	Case - Symmetric	Case - Uniform	Case 18	Case 19	Case 20	Case 21	Case 22	Case 23	Case 24	Case 25	Case 26	Case 27	Case 28	Case 29	Case 30	Case 31	Case 32	Case 33	Case 34
			A222222	A111111	A111112	A111122	A111222	A112222	A122222	A222223	A222233	A222333	A223333	A233333	A333333	A122223	A112223	A322221	A332211
32	0.06117	0.00464	0	0	0	0	0	0	0	0	0	0	0	4.3E-07	0.00062	0	0	0	1.7E-05
31	0.06062	0.00861	0	0	0	0	0	0	0	0	0	4.1E-06	0	8.5E-06	0.00073	0	0	6.3E-07	2.2E-05
30	0.06368	0.01268	0	0	0	0	0	0	0	0	0	7.9E-06	0	1.9E-05	0.00085	0	0	5.7E-06	1.5E-05
29	0.0548	0.01684	0	0	0	0	0	0	0	6.9E-07	0	4E-06	0	3.4E-05	0.00051	0	0	1.5E-05	2E-05
28	0.0431	0.02074	0	0	0	0	5.9E-06	0	0	1.6E-06	0	1.2E-05	0	8.2E-06	0.00048	0	1.8E-06	8.8E-06	2.6E-05
27	0.03838	0.02456	0	3.5E-06	7.6E-10	0	1.9E-06	0	0	8.1E-06	3.2E-06	1.3E-05	0	0.00014	0.00052	0	1.8E-05	7.8E-06	2.7E-05
26	0.02926	0.02826	0	8.7E-06	4.4E-06	0	9.5E-06	0	0	1.5E-06	6.5E-06	1E-05	8.5E-07	0.00016	0.00068	5.5E-06	1.5E-05	1.2E-05	4E-05
25	0.02271	0.03157	1.1E-07	6.7E-06	1.6E-05	1.4E-06	1.3E-05	0	9.6E-07	3.7E-05	6.9E-05	1.9E-05	1.1E-06	0.00022	0.00081	5E-06	2E-05	1.6E-05	5.1E-05
24	0.01784	0.03424	4E-06	8.3E-06	3.5E-05	3.3E-06	3E-05	9.6E-07	7.9E-06	7.4E-05	0.0001	4.7E-05	4.5E-06	0.00035	0.00052	3.8E-06	2.3E-05	3E-06	6.1E-05
23	0.01324	0.03661	1.1E-05	2.9E-05	5.1E-05	1.2E-05	8.4E-05	2.2E-06	1.7E-05	0.00014	0.00014	8.9E-05	1.3E-05	0.00053	0.00048	8.4E-06	4.5E-05	0.00012	0.0001
22	0.01022	0.03879	4E-05	4.5E-05	8.6E-05	2.6E-05	8.3E-05	1.3E-05	3.4E-05	0.0002	0.00017	0.00015	8.2E-05	0.00042	0.00058	2.1E-05	7.1E-05	0.00015	0.00012
21	0.00767	0.04115	7.5E-05	6.2E-05	0.0001	5.2E-05	0.00013	3.6E-05	5.7E-05	0.00023	0.00026	0.0002	0.00015	0.00084	0.00067	3.7E-05	0.00017	0.00021	0.00021
20	0.0061	0.04303	0.00017	0.00016	0.00018	0.00013	0.0002	8.6E-05	0.00013	0.00029	0.00039	0.00028	0.00028	0.00069	0.0006	7.6E-05	0.00027	0.00026	0.00034
19	0.00613	0.04423	0.00033	0.0002	0.00019	0.00027	0.0003	0.00018	0.00022	0.00048	0.00058	0.00038	0.00047	0.00111	0.0013	0.00012	0.00037	0.00038	0.00058
18	0.00463	0.0452	0.0006	0.00038	0.00027	0.00041	0.00048	0.00023	0.00041	0.00075	0.0008	0.0006	0.00076	0.00149	0.00186	0.0003	0.00051	0.0006	0.00097
17	0.00447	0.04817	0.0008	0.00068	0.00044	0.00083	0.00072	0.00043	0.00078	0.00109	0.00119	0.00076	0.00129	0.00225	0.00271	0.00051	0.00061	0.00089	0.0016
16	0.00445	0.0468	0.00108	0.001	0.0008	0.00112	0.00109	0.00081	0.0014	0.00154	0.00142	0.00131	0.00189	0.00305	0.0038	0.00085	0.00116	0.00152	0.00226
15	0.00463	0.04732	0.00158	0.00168	0.0013	0.00164	0.00162	0.00138	0.00208	0.00208	0.0019	0.0021	0.00267	0.00457	0.00553	0.00148	0.00158	0.00241	0.00347
14	0.00553	0.04727	0.00237	0.00233	0.00214	0.00254	0.00223	0.00233	0.00285	0.00316	0.00276	0.00318	0.00373	0.00682	0.00786	0.00237	0.0023	0.0037	0.00518
13	0.00743	0.04847	0.00383	0.00403	0.00375	0.00388	0.00354	0.00382	0.00427	0.00467	0.00436	0.00514	0.00571	0.01037	0.01067	0.00407	0.00364	0.0068	0.00832
12	0.01008	0.04632	0.00681	0.00686	0.00677	0.00683	0.0065	0.00688	0.00821	0.00878	0.00839	0.00725	0.00788	0.01363	0.01386	0.00692	0.00637	0.00844	0.01131
11	0.01313	0.04351	0.00871	0.00824	0.00894	0.00828	0.00868	0.00892	0.00938	0.0097	0.0096	0.01003	0.01148	0.01841	0.01873	0.00837	0.00788	0.01212	0.01569
10	0.01702	0.04119	0.0142	0.0137	0.01347	0.01288	0.01328	0.01398	0.01638	0.01483	0.01308	0.01508	0.01748	0.02411	0.02322	0.0144	0.01286	0.01781	0.02201
9	0.02214	0.03847	0.02188	0.0208	0.02171	0.02053	0.02068	0.02193	0.02451	0.0217	0.02004	0.02313	0.02677	0.03334	0.03466	0.02236	0.01911	0.02883	0.03176
8	0.02768	0.03522	0.03378	0.0307	0.03252	0.0308	0.03226	0.03247	0.03823	0.03223	0.03053	0.03293	0.03703	0.04414	0.04681	0.03358	0.02834	0.03882	0.04363
7	0.0364	0.03166	0.0495	0.04584	0.04798	0.04752	0.0468	0.04812	0.05182	0.04723	0.04528	0.04737	0.05189	0.0577	0.06146	0.04881	0.04392	0.05429	0.06037
6	0.04459	0.02773	0.07235	0.07003	0.07088	0.07008	0.06938	0.07138	0.07908	0.06979	0.06695	0.06783	0.07131	0.07574	0.07906	0.07084	0.06533	0.0788	0.08351
5	0.0578	0.02345	0.10414	0.10497	0.10489	0.1036	0.10201	0.10338	0.10614	0.10421	0.09949	0.09846	0.09988	0.09975	0.10349	0.1042	0.09814	0.10877	0.11386
4	0.07291	0.01806	0.14826	0.15174	0.15105	0.15198	0.14827	0.14802	0.14788	0.15082	0.15171	0.14317	0.13894	0.13408	0.13801	0.1485	0.149	0.15086	0.14941
3	0.08416	0.01421	0.18182	0.18413	0.18404	0.18388	0.18495	0.18098	0.17708	0.18289	0.18648	0.18222	0.18891	0.18827	0.18455	0.17888	0.18704	0.17528	0.18775
2	0.0883	0.00864	0.19588	0.19788	0.19712	0.19841	0.20075	0.19967	0.19018	0.19499	0.20181	0.20212	0.19832	0.18962	0.18059	0.1802	0.20706	0.18144	0.18084
1	0.08925	0.00624	0.18138	0.18441	0.18253	0.18567	0.18702	0.18696	0.14817	0.15034	0.15583	0.15838	0.16634	0.13868	0.12171	0.15828	0.16478	0.1286	0.11722

国際単位系 (SI) と換算表

表1 SI基本単位および補助単位

量	名称	記号
長さ	メートル	m
質量	キログラム	kg
時間	秒	s
電流	アンペア	A
熱力学温度	ケルビン	K
物質の量	モル	mol
光強度	カンデラ	cd
平面角	ラジアン	rad
立体角	ステラジアン	sr

表3 固有の名称をもつ SI 組立単位

量	名称	記号	他の SI 単位による表現
周波数	ヘルツ	Hz	s ⁻¹
力	ニュートン	N	m kg/s ²
圧力、圧力	パスカル	Pa	N/m ²
エネルギー、仕事、熱電	ジュール	J	N m
工率、放射束	ワット	W	J/s
電気量、電荷	クーロム	C	A s
電位、電圧、起電力	ボルト	V	W/A
静電容量	ファラド	F	C/V
電気抵抗	オーム	Ω	V/A
コンダクタンス	ジーメン	S	A/V
磁束	ウェーバ	Wb	V s
磁束密度	テスラ	T	Wb/m ²
インダクタンス	ヘンリー	H	Wb/A
セルシウス温度	セルシウス度	°C	
光強度	カンデラ	cd	sr
放射線量	グレイ	Gy	J/kg
放射線当量	シーベルト	Sv	J/kg

表2 SIと併用される単位

名称	記号
分、時、日	min, h, d
度、分、秒	°, ', "
リットル	L, l
トン	t
電子ボルト	eV
原子質量単位	u

1 eV = 1.60218 × 10⁻¹⁹ J
1 u = 1.66054 × 10⁻²⁷ kg

表4 SIと共に暫定的に維持される単位

名称	記号
オングストローム	Å
バール	bar
ガリ	Gal
キュリー	Ci
レントゲン	R
ラジアン	rad
レム	rem

1 Å = 0.1 nm = 10⁻¹⁰ m
1 b = 100 fm² = 10⁻²⁸ m²
1 bar = 0.1 MPa = 10⁵ Pa
1 Gal = 1 cm/s² = 10⁻² m/s²
1 Ci = 3.7 × 10¹⁰ Bq
1 R = 2.58 × 10⁻⁴ C/kg
1 rad = 0.01 Gy = 10⁻² Gy
1 rem = 0.01 Sv = 10⁻² Sv

表5 SI接頭語

指数	接頭語	記号
10 ¹⁸	エクサ	E
10 ¹⁵	ペタ	P
10 ¹²	テラ	T
10 ⁹	ギガ	G
10 ⁶	メガ	M
10 ³	キロ	k
10 ²	ヘクト	h
10 ¹	デカ	da
10 ⁻¹	デシ	d
10 ⁻²	センチ	c
10 ⁻³	ミリ	m
10 ⁻⁴	マイクロ	μ
10 ⁻⁵	ナノ	n
10 ⁻⁶	ピコ	p
10 ⁻⁷	フェムト	f
10 ⁻⁸	アト	a

(注)

- 表1～5は「国際単位系」第5版、国際度量衡局 1985年刊行による。ただし、1 eV および 1 u の値は CODATA の 1986 年推奨値によった。
- 表4には商型、リットル、アール、ヘクタールも含まれているが日常の単位なのでここでは省略した。
- bar は、JIS では液体の圧力を表す場合に限り表2のカチコリーに分類されている。
- EC 閣僚理事会指令では bar、barn および「血圧の単位」mmHg を表2のカチコリーに入れている。

換 算 表

力	N (= 10 ⁵ dyn)	kgf	lbf
1	0.101972	0.224809	
9.80665	1	2.20462	
4.44822	0.453592	1	

圧 度 1 Pa = 1 N/m² = 10⁵ dyn/cm² (g/cm s²)

動粘度 1 m²/s = 10⁴ St (ストークス) (cm²/s)

圧	MPa (= 10 bar)	kgf/cm ²	atm	mmHg (Torr)	lbf/in ² (psi)
1	10.1972	9.80665	7.50062 × 10 ⁻¹	145.038	
0.00980665	1	0.967841	735.559	14.2233	
0.101325	1.01325	1	760	14.6959	
1.33322 × 10 ⁻⁴	1.35951 × 10 ⁻³	1.31579 × 10 ⁻²	1	1.93358 × 10 ⁻²	
6.89476 × 10 ⁻²	7.03070 × 10 ⁻²	6.80460 × 10 ⁻²	51.7149	1	

エネルギー	J (= 10 ⁷ erg)	kgf·m	kW·h	cal (計量法)	Btu	ft·lbf	eV
1	0.101972	2.77778 × 10 ⁻²	0.238849	9.47813 × 10 ⁻⁴	0.737562	5.24150 × 10 ⁻¹¹	
9.80665	1	2.72407 × 10 ⁻¹	2.34270	9.29487 × 10 ⁻²	7.23301	6.12082 × 10 ⁻¹⁰	
3.6 × 10 ⁵	3.67098 × 10 ⁴	1	8.59899 × 10 ⁻¹	3412.13	2.65522 × 10 ³	2.24694 × 10 ⁻⁷	
4.18605	0.426858	1.16279 × 10 ⁻¹	1	3.96759 × 10 ⁻¹	3.08747	2.61272 × 10 ⁻¹¹	
1055.06	107.506	2.93072 × 10 ⁻¹	252.042	1	778.172	6.58515 × 10 ⁻¹¹	
1.35582	0.138255	3.76616 × 10 ⁻¹	0.323890	1.28506 × 10 ⁻¹	1	8.46233 × 10 ⁻¹⁴	
1.60218 × 10 ⁻¹⁹	1.63377 × 10 ⁻²⁰	4.45059 × 10 ⁻¹⁴	3.82743 × 10 ⁻¹⁸	1.51857 × 10 ⁻²²	1.18171 × 10 ⁻¹¹	1	

1 cal = 4.18605 J (計量法)

= 4.184 J (熱化学)

= 4.1855 J (15 °C)

= 4.1868 J (国際蒸気表)

仕事率 1 PS (仏馬力)

= 75 kgf·m/s

= 735.499 W

放射線	Bq	Ci
1	2.70270 × 10 ¹⁰	
3.7 × 10 ¹⁰	1	

吸収線量	Gy	rad
1	100	
0.01	1	

照射線量	C/kg	R
1	3876	
2.58 × 10 ⁻⁴	1	

線量当量	Sv	rem
1	100	
0.01	1	

