

Semi-orthogonal Wavelets for Elliptic Variational Problems

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Abstract

In this paper, we give a construction of wavelets which are (a) semi-orthogonal with respect to an arbitrary elliptic bilinear form $a(\cdot, \cdot)$ on the Sobolev space $H_0^1((0, L))$ and (b) continuous and piecewise linear on an arbitrary partition of $[0, L]$. We illustrate this construction using the model problem

$$-\epsilon^2 u'' + u = f$$

$$u(0) = u(L) = 0.$$

We also construct α -orthogonal Battle-Lemarié type wavelets which fully diagonalize the Galerkin discretized matrix for the model problem with domain \mathbb{R} .

Finally, we describe a hybrid basis consisting of a combination of elements from the semi-orthogonal wavelet basis and the hierarchical Schauder basis. Numerical experiments indicate that this basis leads to robust, scalable Galerkin discretizations of the model problem which remain well-conditioned independent of ϵ , L , and the refinement level K .

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1 Introduction

In this section, we review some basic theory about Galerkin discretizations of elliptic variational problems and their relationship to the Riesz bounds of the underlying basis (cf. [2, 5]).

Let \mathcal{H} be a Hilbert space with inner product $\langle \cdot, \cdot \rangle_{\mathcal{H}}$. Let $a(\cdot, \cdot)$ be a symmetric coercive continuous bilinear form on \mathcal{H} , that is a is a symmetric bilinear form such that

$$C\|v\|_{\mathcal{H}}^2 \leq a(v, v) \leq D\|v\|_{\mathcal{H}}^2$$

for some positive constants C and D . Define $\|\cdot\|_E := \sqrt{a(\cdot, \cdot)}$ to be the energy norm generated by a . The coercivity and continuity of a imply that the energy norm is equivalent to the norm associated with \mathcal{H} .

Let $\mathcal{H}' (\cong \mathcal{H})$ denote the dual of \mathcal{H} . Consider the elliptic variational problem:

$$\begin{aligned} \text{Given } F \in \mathcal{H}', \text{ find } u \in \mathcal{H} \text{ such that} \\ a(u, v) = F(v), \quad \forall v \in \mathcal{H}. \end{aligned} \quad (1)$$

Let \mathcal{V} be a finite dimensional subspace of \mathcal{H} . Then the Galerkin approximate solution $u_{\mathcal{V}}$ is the unique solution of (1) with \mathcal{H} replaced by \mathcal{V} . Let $\Phi = (\phi^1, \dots, \phi^N)^T$ be a basis for \mathcal{V} . (Throughout this paper, a basis will be arranged as a column vector.) Then $u_{\mathcal{V}} = c^T \Phi$ can be found by solving the linear system

$$a(\Phi, \Phi)c = F(\Phi), \quad (2)$$

where $a(\Phi, \Phi)$ is the $N \times N$ matrix $(a(\phi^i, \phi^j))$ and $F(\Phi)$ is the column vector $(F(\phi^1), \dots, F(\phi^N))^T$.

For large N , it is usually impractical to solve the linear system (2) using direct methods. When the matrix $A^{\Phi} := a(\Phi, \Phi)$ is well-conditioned, the system can be efficiently solved using iterative methods. We say that $\underline{\alpha}$ (respectively $\bar{\alpha}$) is a lower (upper) Riesz bound for the basis Φ with respect to $\|\cdot\|_E$ if

$$\underline{\alpha} c^T c \leq \|c^T \Phi\|_E^2 \leq \bar{\alpha} c^T c. \quad (3)$$

Define $\underline{\alpha}_{\Phi}$ ($\bar{\alpha}_{\Phi}$) to be the largest (smallest) lower (upper) Riesz bound for Φ with respect to $\|\cdot\|_E$. Observe that

$$\|c^T \Phi\|_E^2 = c^T A^{\Phi} c.$$

Since A^{Φ} is symmetric and positive definite we have

$$\begin{aligned} \|A^{\Phi}\|_2 &= \max_c \frac{c^T A^{\Phi} c}{c^T c} = \bar{\alpha}_{\Phi} \\ \|(A^{\Phi})^{-1}\|_2^{-1} &= \min_c \frac{c^T A^{\Phi} c}{c^T c} = \underline{\alpha}_{\Phi}. \end{aligned}$$

Therefore, the spectral condition number of A^{Φ} , $\text{cond}(A^{\Phi})$, is related to the Riesz bounds for Φ in the following way:

$$\text{cond}(A^{\Phi}) = \bar{\alpha}_{\Phi} / \underline{\alpha}_{\Phi}. \quad (4)$$

Suppose Ψ is another basis for \mathcal{V} and suppose W is the nonsingular $N \times N$ matrix such that

$$\Psi = W^T \Phi.$$

Then $u_{\mathcal{V}} = d^T \Psi$ may also be found by solving

$$a(\Psi, \Psi)d = F(\Psi). \quad (5)$$

Note that

$$A^{\Psi} = a(\Psi, \Psi) = W^T a(\Phi, \Phi) W. \quad (6)$$

Thus the linear system (5) resulting from (2) by a change of basis can also be considered to arise from (2) by preconditioning with W .

2 Multiscale Transformations

Suppose

$$\mathcal{V}_0 \subset \mathcal{V}_1 \subset \dots \subset \mathcal{V}_k \subset \dots$$

is a one-sided sequence of nested finite-dimensional subspaces of \mathcal{H} such that $\bigcup \mathcal{V}_k = \mathcal{H}$. Define $\mathcal{W}_0 := \mathcal{V}_0$ and, for $k \geq 1$, choose \mathcal{W}_k in \mathcal{V}_k so that

$$\mathcal{V}_k = \mathcal{V}_{k-1} \dot{+} \mathcal{W}_k \quad (7)$$

where $\dot{+}$ denotes the direct sum. Let Φ_k be a basis for \mathcal{V}_k and let ψ_k be a basis for \mathcal{W}_k (we choose $\psi_0 = \Phi_0$). Then

$$\Psi_k := \begin{pmatrix} \psi_0 \\ \vdots \\ \psi_k \end{pmatrix}$$

is also a basis for \mathcal{V}_k . Let W_k be the multiscale transformation such that

$$\Psi_k = W_k^T \Phi_k,$$

and let T_k be the two-scale transformation such that

$$\begin{pmatrix} \Phi_{k-1} \\ \psi_k \end{pmatrix} = T_k^T \Phi_k.$$

Observe that

$$W_k = T_k \begin{pmatrix} T_{k-1} & 0 \\ 0 & I_{k-1} \end{pmatrix} \dots \begin{pmatrix} T_1 & 0 \\ 0 & I_1 \end{pmatrix}$$

where I_j is the $n \times n$ identity matrix with $n = \text{card}(\psi_j)$.

Fix K and let $\Psi = \Psi_K$, $\Phi = \Phi_K$, and $W = W_K$. We assume that (a) multiplication by W can be implemented with a fast algorithm (this is the case for compactly supported wavelet bases), (b) A^{Ψ} is well-conditioned, and (c) $F(\Phi)$ can be easily approximated. Algorithm A summarizes the solution of the discretized problem given in (2) using the multiscale transform W .

Algorithm A:

- Approximate F^Φ .
- Calculate $F^\Psi = W^T F^\Phi$.
- Solve $A^\Psi d = F^\Psi$.
- $c = Wd$.

It is interesting to note that this algorithm does not use the decomposition matrix W^{-1} .

3 Wavelet Construction

Let $(X_k)_{k \geq 0}$ be a given sequence of nested knot sequences on $[0, 1]$ satisfying

- $X_k = (x_k^j)_{0 \leq j \leq N_k}$
- $0 = x_k^0 < \dots < x_k^j < \dots < x_k^{N_k} = L$
- $x_{k+1}^j = x_k^j$

Let ϕ_k^j be the piecewise linear continuous function with knot sequence X_k such that $\phi_k^j(x_k^{j'}) = \delta_{j,j'}$. Let $\Phi_k = (\phi_k^1, \dots, \phi_k^{N_k-1})^T$, then Φ_k is a nodal basis for \mathcal{V}_k which is the usual finite element space of piecewise linear continuous functions on $[0, L]$ with knot sequence X_k .

We next describe two choices for \mathcal{W}_k .

3.1 Hierarchical Schauder Basis

One simple choice for \mathcal{W}_k satisfying (7) is the well known Schauder basis (cf. [7, 9, 10])

$$\psi_k^j := \phi_k^{2j-1}, \quad j = 1, \dots, N_k$$

illustrated in Figure 1.

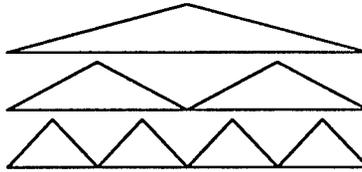


Figure 1: Schauder Basis $\Psi_2 = (\psi_0^T \ \psi_1^T \ \psi_2^T)^T$ with dimension 7 on a uniform partition.

Next we construct the two-scale transformation for the Schauder basis. Denote the length of the subinterval $[x_k^{j-1}, x_k^j]$ by

$$\Delta_k^j := x_k^j - x_k^{j-1}.$$

The function values $h_k^{j,j'}$ and $g_k^{j,j'}$ for ϕ_{k-1}^j and ψ_k^j respectively at the knot $x_k^{j'}$ are given by

$$h_k^{j,j'} = \begin{cases} \frac{\Delta_k^{j'}}{\Delta_{k-1}^j}, & j' = 2j - 1 \\ 1, & j' = 2j \\ \frac{\Delta_k^{j'+1}}{\Delta_{k-1}^j}, & j' = 2j + 1 \\ 0, & \text{otherwise} \end{cases}$$

$$g_k^{j,j'} = \delta_{j', 2j-1}.$$

Then

$$\phi_{k-1}^j = \sum_{j'} h_k^{j,j'} \phi_k^{j'}$$

$$\psi_k^j = \sum_{j'} g_k^{j,j'} \phi_k^{j'}$$

Now let H_k be the $(N_k - 1) \times (N_{k-1} - 1)$ matrix $H_k = (h_k^{j,j'})_{j',j}$ and let G_k be the $(N_k - 1) \times (N_{k-1})$ matrix $G_k = (g_k^{j,j'})_{j',j}$. Thus, the two-scale transformation for the Schauder basis is given by

$$T_k = (H_k \ G_k).$$

3.2 Semi-orthogonal Sombbrero Wavelets

Here we choose \mathcal{W}_k to be the orthogonal complement of \mathcal{V}_{k-1} in \mathcal{V}_k with respect to the scalar product $a(\cdot, \cdot)$, that is

$$\mathcal{W}_k := \mathcal{V}_k \cap \mathcal{V}_{k-1}^{\perp a}.$$

Regardless of the choice of basis ψ_k for \mathcal{W}_k , the matrix A^{Ψ_k} is then decoupled between levels so that it is a block diagonal matrix:

$$A^{\Psi_k} = \text{diag}(A^{\psi_0}, A^{\psi_1}, \dots, A^{\psi_k}).$$

We next give a procedure for constructing a local basis of wavelets for \mathcal{W}_k . Let

$$B := B(k) = a(\Phi_{k-1}, \Phi_k)$$

where we suppress the k dependence when the choice for k is unambiguous. Note that

$$\mathcal{W}_k = \{g^T \Phi_k \mid g \in \ker B\}.$$

We will use certain subblocks of B in our construction. To this end we define R_{i_1, i_2} to be the $(i_2 - i_1 + 1) \times (N_k - 1)$ matrix whose i -th row is the $(i - i_1 + 1)$ -th row of the $(N_k - 1) \times (N_k - 1)$ identity matrix. Then the $[i_1, i_2] \times [j_1, j_2]$ block of B is given by

$$B_{j_1, j_2}^{i_1, i_2} := R_{i_1, i_2} B R_{j_1, j_2}^T.$$

Let

$$C_n := \begin{cases} B_{1, 2n-1}^{1, n} & \text{for } n = 2, 3, \\ B_{2n-5, 2n-1}^{n-3, n} & \text{for } 3 \leq n \leq N_{k-1} - 1 \\ B_{2n-5, 2n-1}^{n-3, n-1} & \text{for } n = N_{k-1} \end{cases}$$

For $4 \leq n < N_{k-1}$, the matrix C_n is a 4×5 matrix which generically has a kernel of dimension one. This kernel then corresponds to a wavelet with support contained in $[x_k^{2n-6}, x_k^{2n}] = [x_{k-1}^{n-3}, x_{k-1}^n]$. More generally, we define the following procedure for constructing a local basis for \mathcal{W}_k .

Let $K_n := \ker(C_n)$ and, for $n \geq 3$, let K_n^0 denote the subspace of K_n consisting of the elements in $w \in K_n$ whose last two components are both zero. In the generic case, K_n^0 is the trivial subspace. Let

$$\Phi_k^n := \begin{cases} (\phi_k^1, \phi_k^2, \phi_k^3)^T, & n = 2 \\ (\phi_k^{2n-5}, \dots, \phi_k^{2n-1})^T, & 3 \leq n \leq N_{k-1}. \end{cases} \quad (8)$$

Algorithm B:

- Let w_3 denote a basis for C_3 and set $\psi_k^3 = \{w^T \Phi_k^3 \mid w \in w_3\}$.
- For $n = 3, \dots, N_{k-1}$, do
 - Choose K_n^1 so that $K_n = K_n^0 \dot{+} K_n^1$ and choose a basis w_n for K_n^1 .
 - Set $\psi_k^n = \{w^T \Phi_k^n \mid w \in w_n\}$.
- $\psi_k = \bigcup_{n=3}^{N_{k-1}} \psi_k^n$.

We next give a sufficient condition that the above procedure produces a basis for \mathcal{W}_k . For $4 \leq n < N_{k-1}$, we note that C_n has the following block form

$$C_n = \begin{pmatrix} D_n & E_n \\ 0 & F_n \end{pmatrix} \quad (9)$$

where D_n is 3×3 , E_n is 3×2 , and F_n is 1×2 .

Lemma 1 *Let ψ_k be the set produced by Algorithm B. Suppose*

$$\text{range } D_n \supset E_n(\ker F_n) \quad (4 \leq n \leq N_{k-1}). \quad (10)$$

Then ψ_k is a basis for $\mathcal{W}_k := \mathcal{V}_k \cap \mathcal{V}_{k-1}^\perp$.

Proof:

Let $S^{n_1, n_2} := \{f \in W_k \mid \text{supp}(f) \subset [x_k^{2n_1}, x_k^{2n_2}]\}$. Let $B_n := B_{1, 2n-1}^{1, n}$. Note that $f \in S^{0, n}$ if and only if $f = y^T(\phi_k^1, \dots, \phi_k^{2n-1})^T$ for some $y \in \ker B_n$ and that $f \in S^{n-3, n}$ if and only if $f = y^T \Phi_k^n$ for some $y \in \ker C_n$. Hence, the proof will be complete if we can show that

$$S^{0, n} = S^{0, n-1} + S^{n-3, n}, \quad (4 \leq n \leq N_{k-1}). \quad (11)$$

Observe that

$$B_n = \begin{pmatrix} B_{n-1} & E_n \\ 0 & F_n \end{pmatrix} \quad (12)$$

for $n \geq 4$.

Suppose $v \in \ker F_n$, then by (10) there is some $u \in \ker C_n$ such that $(u^4, u^5)^T = v$. Suppose $y \in \ker B_n$. From (12) it is clear that $w := (y_{2n-2}, y_{2n-1})^T \in \ker F_n$ and hence there is some $u \in \ker C_n$ such that the last two components of y agree with the last two components of u . We then obtain

$$\ker B_n = P_1(\ker B_{n-1}) + P_2(\ker C_n) \quad (13)$$

where P_1 is the padding operator that takes a vector v of length $2n-3$ to one of length $2n-1$ by appending two zeros to v and P_2 is the padding operator that takes a vector v of length 5 to one of length $2n-1$ by prepending $2n-6$ zeros to v .

Then (11) follows from (13) and the proof is complete. \blacksquare

4 Uniform Partition

In this section we give the construction of piecewise-linear wavelets on a uniform partition which are semi-orthogonal with respect to the bilinear form associated with the following model problem:

$$\begin{aligned} -\epsilon^2 u'' + u &= f \\ u(0) = u(L) &= 0. \end{aligned} \quad (14)$$

Let $\langle \cdot, \cdot \rangle$ be the usual inner product on $L_2([0, L])$ and let \mathcal{H} be the Sobolev space $\mathcal{H}_0^1([0, L])$ for some $L \in \mathbb{R}^+$. We assume f is such that $F := \langle f, \cdot \rangle$ is in \mathcal{H}' . Then the weak formulation becomes: Find $u \in \mathcal{H}$ so that

$$a(u, v) := \epsilon^2 \langle u', v' \rangle + \langle u, v \rangle = \langle f, v \rangle \quad \forall v \in \mathcal{H}. \quad (15)$$

Moreover, suppose (X_k) is a uniform partition of $[0, L]$. In order to avoid special cases associated with boundary wavelets, we let $L = 4$ and $N_0 = 4$:

$$x_k^j = \frac{j}{2^k}, \quad j = 0, \dots, N_k = 2^k L.$$

Let $\phi(x) = (1 - |x - 1|)_+$ and define

$$\phi_k^j := 2^{k/2} \phi(2^k \cdot - j).$$

(Here we have chosen a different normalization of ϕ_k^j than the normalization used in the non-uniform case.) Then

$$\phi_k^j = \sum_{l=-1}^1 h_l \phi_{k+1}^{j-l}$$

where $h_{-1} = h_1 = \frac{1}{2\sqrt{2}}$, and $h_0 = \frac{1}{\sqrt{2}}$.

As in the previous section, let Φ_k^n be defined by (8). Because of the differentiation in the scalar product a , the ϵ in the model problem is scaled differently at each level resulting in a level dependent parameter ϵ_k given by

$$\epsilon_k := 2^k \epsilon.$$

In this case, C_n is independent of $4 \leq n \leq N_{k-1} - 1$ and its kernel is the space spanned by the vector

$$w = (24\epsilon_k^2 - 1, 6, 48\epsilon_k^2 - 10, 6, 24\epsilon_k^2 - 1)^T \quad (16)$$

The kernel of C_3 (respectively, $C_{N_{k-1}}$) contains w plus an additional vector w_L (respectively, w_R) given below:

$$w_L = (9 + 72\epsilon_k^2, -6, 1 - 24\epsilon_k^2, 0, 0)^T$$

and

$$w_R = (0, 0, 1 - 24\epsilon_k^2, -6, 9 + 72\epsilon_k^2)^T.$$

Then we let $\psi_k^1 = w_L^T \Phi_k^3$, $\psi_k^j = w^T \Phi_k^{j+1}$ for $2 \leq j \leq N_{k-1} - 1$ and $\psi_k^{N_{k-1}} = w_R^T \Phi_k^{N_{k-1}}$.

The wavelet ψ is shown in Figure 2 for selected ϵ and for a larger set in Figure 3. Another more general construction of semi-orthogonal wavelets on a uniform grid was given in [3, 4].

4.1 Unbounded Domain: Riesz bounds and Battle-Lemarié type wavelets

We next consider the simpler choice of domain \mathbb{R} . In this case we can calculate the Riesz bounds for the wavelet bases ψ_k for \mathcal{W}_k using Fourier transform techniques. For $\theta \in L_2(\mathbb{R})$ we define the Gramian symbol E_θ (with respect to the scalar product $a(\cdot, \cdot)$) by

$$E_\theta(\omega) = \frac{1}{2\pi} \sum_{n \in \mathbb{Z}} a(\theta, \theta(\cdot - n)) e^{in\omega} \quad (17)$$

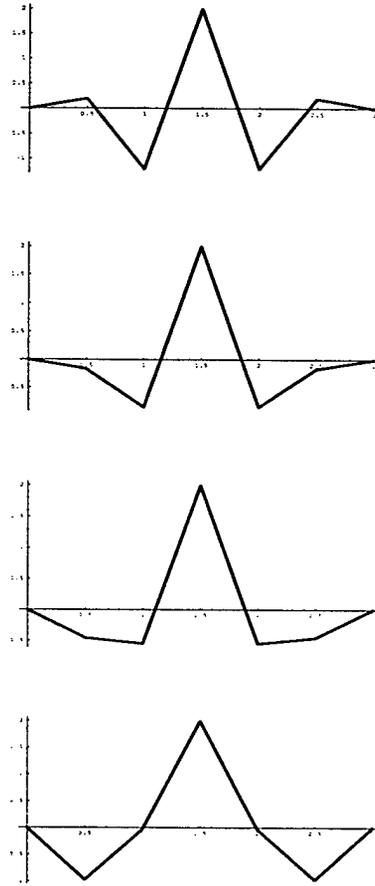


Figure 2: Semi-orthogonal Sombbrero for $\epsilon = 0, 0.3, 0.5, 3$ respectively on a uniform grid

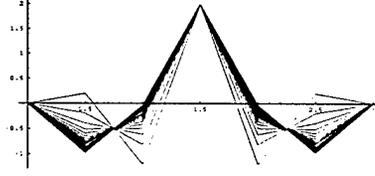
It is a standard result (see [8], for example) that the Riesz bounds $\underline{\alpha}_\Theta$ and $\bar{\alpha}_\Theta$ for the infinite basis $\Theta = (\theta(\cdot - n))_{n \in \mathbb{Z}}$ with respect to $a(\cdot, \cdot)$ are equal to the essential infimum and essential supremum of E_Θ , respectively. The L_2 -condition number of the infinite matrix (A^Θ) is then the ratio $\bar{\alpha}_\Theta / \underline{\alpha}_\Theta$.

In the case of our model problem with the sombrero wavelets ψ_k we get

$$E_{\psi_k}(\omega) = \alpha_0 + 2\alpha_1 \cos(\omega) + 2\alpha_2 \cos(2\omega)$$

where

$$\begin{aligned} \alpha_0 &= 12(3 + 122\epsilon_k^2 + 480\epsilon_k^4 + 1152\epsilon_k^6) \\ \alpha_1 &= 20/3 + 384\epsilon_k^2 - 2304\epsilon_k^4 - 9216\epsilon_k^6 \\ \alpha_2 &= (2/3)(1 - 24\epsilon_k^2)^2(-1 + 6\epsilon_k^2). \end{aligned}$$

Figure 3: Semi-orthogonal Sombbrero for ϵ between 0 and 3

It is an elementary, but tedious, exercise to verify that

$$\frac{\max_{\omega} E_{\psi_k}(\omega)}{\min_{\omega} E_{\psi_k}(\omega)} = \begin{cases} 4 \frac{(1+12\epsilon_k^2)^3}{9+432\epsilon_k^2} & \text{for } 0 \leq \epsilon_k \leq 0.33 \\ 4 \frac{9+432\epsilon_k^2}{(1+12\epsilon_k^2)^3} & \text{for } 0.36 \leq \epsilon_k < \infty \end{cases}$$

and that

$$\frac{\max_{\omega} E_{\psi_k}(\omega)}{\min_{\omega} E_{\psi_k}(\omega)} \leq 1.2 \quad \text{for } 0.33 \leq \epsilon_k \leq 0.36. \quad (18)$$

Since A^{Ψ_k} is block diagonal, A^{Ψ_k} can be preconditioned with a simple diagonal preconditioner so that the resulting A^{Ψ_k} satisfies

$$\text{cond}(A^{\Psi_k}) = \max_{j \leq k} \text{cond}(A^{\psi_j}).$$

Then (18) shows that $\text{cond}(A^{\Psi_k})$ is uniformly bounded for $0 \leq \epsilon_k \leq \epsilon^*$ for any fixed ϵ^* . For instance, we get the following:

$$\text{cond}(A^{\Psi_k}) < \begin{cases} 2.4 & \text{for } \epsilon_k < .5 \\ 271 & \text{for } \epsilon_k < 2 \\ 1330 & \text{for } \epsilon_k < 3 \end{cases}$$

For the unbounded domain case we use the following well known Fourier trick (cf. [8]) to construct an a -orthogonal basis for \mathcal{W}_k . Let $(\xi_\ell)_{\ell \in \mathbf{Z}}$ denote the Fourier coefficients of $\sqrt{1/E_{\psi_k}}$ and define

$${}^{BL}\psi_k^j := \sum_{\ell} \xi_\ell \psi_k^{j+\ell}.$$

In the case $\epsilon = 0$, we get the usual Battle-Lemarié wavelets. In this case, A^{Ψ_k} is the identity matrix. It is interesting to observe that ${}^{BL}\psi_0^0$ appears to converge pointwise to the Schauder wavelet ${}^h\psi_0^0 = \phi_1^1$ as ϵ goes to infinity. The wavelet ${}^{BL}\psi$ is shown in Figure 4 for selected ϵ .

4.2 Hybrid Basis

Our goal is to achieve a robust, fully scalable algorithm which is uniformly $\mathcal{O}(N_K)$ independent of the size of the problem L , the maximum refinement

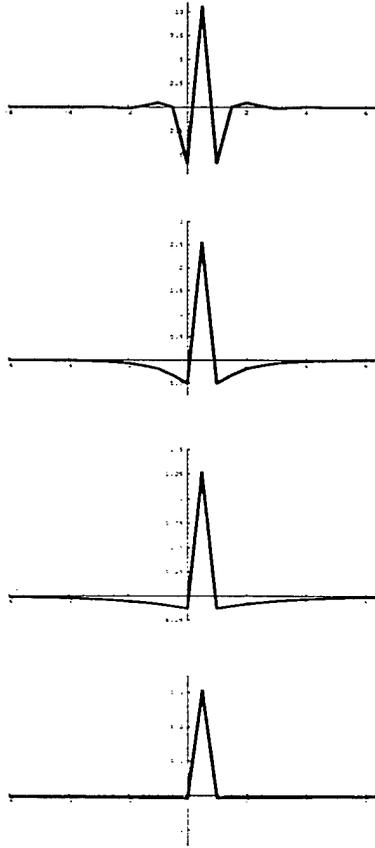


Figure 4: Battle-Lemarié type wavelets for $\epsilon = 0, 1, \sqrt{5}, 10$ respectively

level K , and the parameter ϵ . In this section we assume that our bases are normalized in the a -norm. This corresponds to a preconditioning of the form $D^{-1/2}AD^{-1/2}$ where D is the diagonal of A . We let ${}^h\Psi_k$ denote the normalized Schauder basis described in Section 3.1 and ${}^s\Psi_k$ the normalized Sombrero basis described in Section 4.

For the model problem, the semi-orthogonal basis is ill-conditioned for large ϵ and well-conditioned for small ϵ . One approach we have explored numerically is to use the hybrid basis

$${}^{sh}\Psi_K := \begin{pmatrix} {}^s\Psi_{\tilde{k}} \\ {}^h\psi_{\tilde{k}+1} \\ \vdots \\ {}^h\psi_K \end{pmatrix}$$

where \tilde{k} is chosen so that $\epsilon_{\tilde{k}} = \mathcal{O}(1)$. The resulting discretized matrix $A^{sh\Psi_K}$

is illustrated in Figure 5. Our numerical experiments indicate that the hybrid basis achieves the above mentioned goals for the model problem.

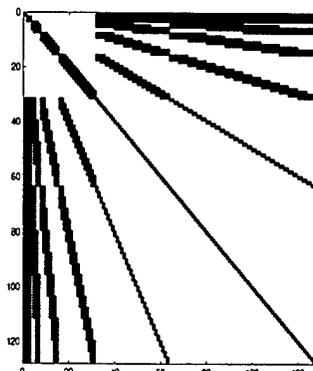


Figure 5: $A^{sh}\Psi_{\kappa}$ consisting of 4 levels with the semi-orthogonal basis combined with 2 additional levels with the Schauder basis.

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