only in the $\sigma_{xx} \to \infty$ limit and for a particular kind of disordering potential [8,9]. The most important aspects of the diagram, its topology and exponents, are guesses. The paper presented by D. C. Tsui [10] at this conference shows experimental "flows" that are nearly perfect replicas of Fig. 1. While one could argue that this agreement is an accident, since the coulomb interactions in Prof. Tsui's samples should have made the non-interacting theory irrelevant, I think the correct interpretation is that "typical" coulomb interaction strengths always give rise to flows which look like the noninteracting theory when projected onto the $\sigma_{xx} - \sigma_{xy}$ plane, and that the topology of the guessed noninteracting flows is correct. I also think that the agreement between locations of the fixed points is meaningful. There is reason to believe that the fixed point should be near the maximum of the self-consistent born conductivity of Ando [1,9,11] when the correlation length of the impurity potential is much shorter than the localization length. When this is not the case, the position of the fixed points should depend on the functional form of the impurity potential.

If the transitions $0 \to 1/3$, $1/3 \to 2/3$ and $2/3 \to 1$, are all equivalent to the transition $0 \to 1$, then we must assign to each a fixed point like than in Fig. 1. One guesses the σ_{xx} coordinate of these points to be 1/9 of the $0 \to 1$ value, since the self-consistent born conductivity scales as the square of the particle's charge. The second important ramification of the exactness of these analogies is that the $0 \to 1/3$, ... etc., fixed points are topologically incompatible with the $0 \to 1$ fixed point without the addition of at least one new, totally repulsive fixed point to separate them. Assuming that there is only one, we obtain the flow diagram for the fractional quantum Hall effect shown in Fig. 2.

The new fixed point, denoted by a box in Fig. 2, corresponds physically to the phase transition in which increasing disorder destroys the fractional states. That this is the case is shown in more detail in Fig. 3. The