$$\exp - \int_1^{\lambda} (d\lambda'/\lambda') \gamma_A \, g(\lambda') \rightarrow \eta_A (\log \lambda)^{C_A/2b} \quad , \lambda \rightarrow \infty , \quad (9)$$

according to Eq. (3). Here C_A is defined by $\gamma_A = C_A g^2 + 0(g^4)$ and η_A is an unknown constant reflecting the contribution of the low momentum region.

In summary, the effect of an asymptotically free strong interaction is to modify the free quark effective interaction as follows: replace m_c by

$$m_c(M_W^2) \sim \eta \, m_c \, (\log \frac{M_W^2}{\mu^2})^{-12/(33-2n)} ,$$

where η is an unknown constant and the operators 4S + T by

$$\sum_{A=1}^{2} f_A \, \eta_A \, (\log \frac{M_W^2}{\mu^2})^{C_A/2b} \, \theta_A .$$

Thus, the operator that dominate as $M_W \rightarrow \infty$ is θ_1 with the dominance factor

$$(\log \frac{M_W^2}{\mu^2})^{3(7.51-8/3)/(33-2n)} .$$

For a four-quark model n = 4 and the dominance factor is $(\log M_W^2/\mu^2)^{0.6}$. It is amusing to note that the free quark effective operator 4S + T has been transformed almost completely into the scalar operator:

$$4S + T \rightarrow 6.6 \, \eta_1 \theta_1 = 6.6 \, \eta_1 \, [S + 0.008 \, T] .$$

Hopefully, the unknown constant η_1 is of order one and may be absorbed