Of course, operators involving the right-handed fields become relevant only in models which contain right-handed charged currents. For the anomalous dimension coefficients ϕ_K one finds

$$\phi^{+}_{LL} = \frac{12}{33 - 2n}, \qquad \phi^{-}_{LL} = -\frac{6}{33 - 2n}$$

$$\phi^{+}_{LR} = \frac{24}{33 - 2n}, \quad \phi^{-}_{LR} = -\frac{3}{33 - 2n}, \quad \phi^{(0)}_{LR} = +\frac{4}{33 - 2n}$$

where n is the number of quark color triplets that enter in the strong interactions. The coefficient $\phi^{(0)}_{LR}$, which is relevant for the operator[18] $\theta^{(0)}_{LR}$, was not computed in the above papers. We have found that its Wilson coefficient C_K vanishes up to second order in the momentum dependent strong coupling constant. This becomes small at large momenta and it therefore seems reasonable to drop the operator $\theta^{(0)}_{LR}$ from consideration. We note that the operator θ^{+}_{LL} is symmetric in color indices; θ^{-}_{LL} antisymmetric. The former belongs to the 8 representation of strong SU(3), the latter also to the 27 representation. In $\theta^{(+)}_{LR}$ the color indices of the right-handed quarks are contracted with those of the left-handed quarks; θ^{-}_{LR} has a more complicated color structure. Both θ^{+}_{LR} and θ^{-}_{LR} belong to strong SU(3) octets in the models considered.

The question of octet enhancement[15] depends on the logarithmic factors in Eq. (10), on the Wilson coefficients C_K, and on the matrix elements of the operators θ_K. It is in the spirit of the present discussion to suppose