Eq. (3) involves a combination of scalar and tensor coupling terms (we subsume the $1 - \gamma_5$ factors in this characterization). The mass factor $m_p^2$, also appears in Eq. (3). It is suggestive on the ideas of asymptotic freedom to treat this mass factor as a function of loop momentum and similarly to multiply the four quark operators by momentum dependent factors which reflect their anomalous dimensions (actually there are two different linear combinations of the scalar and tensor operators which constitute the objects having definite anomalous dimension). This approach ignores other operators that might be important at small distances, e.g., operators involving the strong interaction gluon fields, but it does incorporate gluon effects for the operators that are retained. The question, then, is whether the momentum dependent factors introduced by these considerations are substantial enough to modify Eq. (3) seriously. A discussion along these lines is carried out in Appendix I. It turns out that the anomalous dimensions of the mass operator and of the four quark operators combine to produce what is only a rather modest momentum dependence. There is therefore no obvious reason to expect a substantial change in the estimate of Eq. (3) -- what is at stake, recall, is two orders of magnitude.

We turn next to the analog of $K^0$-$\bar{K}^0$ mixing for mesons composed of a heavy and light quark pair, i.e., mixing of $D^0 \equiv (p'\bar{p})$ and $\bar{D}^0 \equiv (\bar{p}'p)$, where $p'$ is some new heavy quark. In the present absence of experimental information on such states we are no longer dealing with constraints, but rather with possibilities. For present purposes we ignore effects which