

where

$$G_m(q_0^2; \alpha_\beta) = \int_1^\infty \frac{d\omega}{\omega^{m+2}} \omega^{-\alpha_\beta+1} \frac{d}{d\omega} \omega^{\alpha_\beta} F(\omega, q_0^2) \quad (\text{A.4})$$

is known from the input data (which satisfies the requirement that $F(1, q_0^2) = 0$). In this way we are led to consider the moment problem

$$F_m(q^2) = G_m(q_0^2; \alpha_\beta) \lambda^{-b_m}$$

where

$$b_m = a_m + (\log \lambda)^{-1}.$$

The new problem has exactly the same structure as the original one, with $F_n(q_0^2) \rightarrow G_n(q_0^2; \alpha_\beta)$ and $a_n \rightarrow b_n$. By repeated use of this trick we can arrange (over some range of λ which is big enough for our needs) that the modified C_β are all positive. Indeed, with sufficient repetition we can arrange that the index ν encountered in (b) is always greater than unity. This last allows us to avoid modified Bessel functions of negative index. The latter are singular at the origin and would be a nuisance for numerical work.