Our next task, then, is to settle on the input information $F_2(\omega,q_o^2)$ and $F_3(\omega,q_o^2)$. To be safely asymptotic we would like to have these for "large enough" q_o^2. Detailed structure function information for the neutrino reactions is however still lacking. Apart from everything else, therefore, we cannot at present proceed in a really quantitativeway. However, in order to see qualitatively what kinds of effects are to be expected in the present framework, we adopt the following illustrative hypothesis. Let us suppose that $q_o^2 \approx 5(\mathrm{GeV})^2$ is already just sufficiently asymptotic so that, for electroproduction, we can employ the SLAC-MIT results for $F_2(\omega,q_o^2)$. We may then employ a simple parton model (from whose predictions we are expecting substantial departures only at much larger q^2) to translate this into the F_2 structure function for neutrino reactions at q_o^2. The details of one such approach, and fit, are discussed for example by Albright and Jarlskog[13]. We shall adopt a slightly modified version[14] of their Eq. (3.8c) to represent the neutrino structure function F_2, averaged over protons and neutrinos, at $q_o^2 \approx 5(\mathrm{GeV})^2$. Concerning $F_3(\omega,q_o^2)$ we make use of the fact that at CERN energies (where departures from scaling are presumably still small) the cross section ratio[2], $\sigma^\nu/\sigma^{\bar\nu} = 2.6 \pm 0.2$, is fairly close to its upper bound, $\sigma^\nu/\sigma^{\bar\nu} \leq 3$. The bound corresponds to $F_3=-\omega F_2$. It will simplify matters, and will perhaps not be too misleading for our qualitative purposes, if we accept this relation at the reference momentum transfer q_o^2. In any case, asymptotic freedom implies at very large q^2 that $|F_3|/\omega F_2 \to 0$, hence that $\sigma^\nu/\sigma^{\bar\nu} \to 1$ as $\mathcal{E} \to \infty$. Our input hypotheses merely helps us to